

Caractérisation expérimentale et modélisation de l'OAPS en tant que précurseur de membranes polyOAPS-imides

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THÈSE

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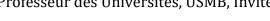
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Caractérisation expérimentale et modélisation de l'OAPS en tant que précurseur de membranes polyOAPS-imides

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Experimental characterization and modelling of OAPS as a precursor for polyOAPS-imide membranes

Table of contents

I.	Introduction			
II.	Thermoresistant polymers	9		
III.	 POSS and OAPS in the literature 1) Infrared spectroscopy (IR) 2) Nuclear magnetic resonance (NMR) 3) X-ray diffraction (XRD) 4) Gel permeation chromatography (GPC) 	20 31 32 38 39		
IV.	 Characterization of a commercial and a controlled OAPS 1) Density of the commercial OAPS 2) IR spectroscopy of OAPS 3) NMR spectroscopy of OAPS a) The experimental NMR spectra b) Predictions of the NMR spectra b.1) Description of the methods b.2) Case study: the meta-OAPS isomer b.2.1) The ²⁹Si-NMR predicted spectra for the 1-ring meta-OAPS b.2.2) The ¹³C-NMR predicted spectra for the 1-ring meta-OAPS b.2.3) The ¹H-NMR predicted spectra for the 1-ring meta-OAPS b.2.4) The ²⁹Si-NMR predicted spectra for the 8-ring meta-OAPS b.2.5) The ¹³C-NMR predicted spectra for the 8-ring meta-OAPS b.2.6) The ¹⁴C-NMR predicted spectra for the 8-ring meta-OAPS b.2.7) Reproducibility: ²⁹Si-NMR, ¹³C-NMR and ¹H-NMR for meta-OAPS b.2.7) Reproducibility: ²⁹Si-NMR, ¹³C-NMR and ¹H-NMR for meta-OAPS b.2.7) Reproducibility: ²⁹Si-NMR predicted spectra for meta-OAPS b.2.7) Reproducibility: ²⁹Si-NMR, ¹³C-NMR and ¹H-NMR for meta-OAPS b.2.7) Reproducibility: ²⁹Si-NMR predicted spectra for meta-OAPS b.3) The ortho and para isomers b.4) Comparison of the three isomers b.4.1) The ²⁹Si-NMR predicted spectra for meta-OAPS, ortho-OAPS and para-OAPS b.4.2) The ¹³C-NMR predicted spectra for meta-OAPS, ortho-OAPS and para-OAPS b.4.3) The ¹⁴-NMR predicted spectra for meta-OAPS, ortho-OAPS and para-OAPS b.4.3) The ¹⁴-NMR predicted spectra for meta-OAPS, ortho-OAPS and para-OAPS 	42 43 49 57 57 65 65 68 69 70 73 76 77 80 82 85 89 90 92 98 109		
V.	 Hybrid Membranes based on OAPS 1) Fabrication of polyOAPS-imide membrane 2) TGA analysis of the membranes 3) IR spectroscopy characterization of the membranes 4) Permeance and gas selectivity of the new membrane 5) Morphology of the polyOAPS-imide membranes 	114 114 117 118 119 123		

VI.	Force-field development for OAPS	128
VII.	Conclusions	137
VIII.	Annexes and References	139
	Annexe 1	139
	Annexe 2	142
	References	151

Abbreviation list

2D-NMR	Two-dimensional nuclear magnetic resonance spectroscopy
6FDA	4,4'-(Hexafluoroisopropylidene)diphthalic anhydride
6FPPr	2,6-bis(trifluoromethylphenylene) pyridine
6FPT	2,5-bis(3-trifluoromethylphenylene)
APS	aminophenylsilsesquioxane
BMI	Bismaliemide
BPADA	4,4'-(4,4'-isopropylidenediphenoxy)bis(phtalic anhydride)
BTDA	Benzophenone-3,3',4,4'-tetracarboxylic dianhydride
BTFDA-DAB	2,2'-bis (p -trimethoxyphenyl) hexafluoropropane dianhydride-3,3'- diaminobenzidine
CUT	Continuous use temperature
DFT	Density functional theory
DMAC	Dimethylacetamide
DMF	Dimethylformamide
EDS	Energy dispersive spectroscopy
FBP-6FPT	2,5-bis(3-trifluoromethylphenylene) thiophene
FE-SEM	Field emission scanning electron microscope
FFV	Fractional free volume
EB	Electron beam
GC	Gradient corrected
GIAO	Gauge-including atomic orbital
GPC	Gel permeation chromatography
GS	Gas separation
HDT	Heat distortion temperature
HETCOR	Heteronuclear correlation spectroscopy
HOSE	Hierarchical Organisation of Spherical Environments
IGA	Isothermal gravimetric analysis
IE	Ion exchange
LDA	Local density approximation
MAE	Mean absolute error
MBR	Membrane bioreactors
MD	Molecular dynamics
meta-APS	meta-aminophenylsilsesquioxane
MF	Micro-filtration
MMMs	Mixed matrix membranes
MWD	Molecular weight distribution

NF	Nano-filtration
NMP	N-Methyl-2-pyrrolidone
NMR	Nuclear magnetic resonance
OAPS	Octa(aminophenyl)silsesquioxane
ODA	4,4'-Oxydianiline
ODPA	4,4'-Oxydiphthalic anhydride
ONPS	Octa(nitrophenyl)silsesquioxane
OPS	Octa(phenyl)silsesquioxane
Ortho-OAPS	ortho-aminophenylsilsesquioxane
PA	Polyamide
Para-OAPS	para-aminophenylsilsesquioxane
PBI	Polybenzimidazole
PC	Polycarbonate
PCM	Polarizable Continuum Model
PEEK	Polyetheretherketone
PEI	Polyetherimide
PI	Polyimide
PMDA	Pyromellitic dianhydride
POPs	Poly[organophasphazenes]
POSS	Polyhedral oligomeric silsesquixanes
PSA	Pressure swing adsorption
PSF	Polysulfone
PTFE	Polytetrafluoroethylene
QM	Quantum mechanical
RMSE	Root mean square error
RO	Reverse osmosis
RTI	Relative Temperature Index
SCRF	Self-Consistent Reaction Field
TAP	2,4,6-triaminopyrimidine
Tg	Glass temperature
TGA	Thermal gravimetric analysis
T _m	Melting temperature
VDW	Van der Waals
VOL	Volatile Organic Liquids
VSA	Vacuum swing adsorption
XRD	X-ray diffraction

ZIF Zeolitic imidazolate frameworks

Abstract

Gas separation by dense polymer membranes is a promising alternative to the cryogenic distillation or adsorption separation processes due to its much lower energy costs. Unfortunately, polymers tend to lose their structural integrity at high temperatures. To improve the thermomechanical resistance while maintaining the gas-sieving properties, hybrid ultrathin membranes based on inorganic POSS (polyhedral oligosilsesquioxanes) cross-linked with organic imides have recently been developed using interfacial polymerization. These polyPOSS-imide membranes were indeed found to perform under tougher operating conditions than conventional polymers. However, the aliphatic arms of the inorganic POSS precursors used so far were too flexible and prone to thermal degradation.

This work aims to attain more stable hybrid membranes using a new inorganic POSS precursor, the octa(aminophenyl)silsesquioxane (OAPS). This siloxane-based cage has three different isomers depending on the *meta*, ortho and para position of the amine with respect to the phenyl group. An extensive literature search showed that the isomer ratios were not clearly defined and that several synthesis routes gave different results. This is problematic as the nature of the isomer will strongly affect the resistance and gas selectivities of the resulting membranes. Experimental characterizations including picnometry, infrared spectroscopy (IR), 1-dimensional and 2-dimensional nuclear magnetic resonance (NMR) have been performed for a commercial OAPS containing all three isomers and a controlled OAPS containing only the para and meta isomers. To better identify the isomers, the experimental results were compared to predictions of the IR and NMR spectra by quantum mechanical methods such as Density Functional Theory and by machine-learning methods such as MestreNova. OAPS was then used as a precursor to make new hybrid membranes with increased thermoresistance. These polyOAPS-imide membranes were characterized using IR, thermal gravimetric analysis (TGA), scanning electron microscopy (SEM), energy dispersive spectroscopy (EDS) and permeability measurements. They were found to degrade well over 400°C and to preserve gas selectivities of ~10 for H₂/CH₄ and H₂/N₂ at 250°C. In parallel, a force-field was developed to carry out molecular dynamics simulations of the three OAPS isomers in order to characterize their differences at the atomic level.

Résumé en français

La séparation des gaz par les membranes polymères denses est une alternative prometteuse aux procédés de séparation par distillation cryogénique ou par adsorption en raison de ses coûts énergétiques beaucoup plus faibles. Malheureusement, les polymères ont tendance à perdre leur intégrité structurale à hautes températures. Afin d'améliorer leur résistance thermomécanique tout en conservant leurs propriétés de séparation des gaz, des membranes hybrides ultrafines à base de POSS (polyhedral oligosilsesquioxanes) inorganiques réticulés avec des imides organiques ont récemment été développées par polymérisation interfaciale. Ces membranes polyPOSS-imides se sont effectivement montrées plus performantes dans des conditions difficiles que les polymères conventionnels. Par contre, les motifs aliphatiques des précurseurs POSS utilisés initialement étaient trop flexibles et sujets à la dégradation thermique.

L'objectif de ce travail est de développer des membranes hybrides plus stables à partir d'un nouveau précurseur POSS inorganique, l'octa(aminophenyl)silsesquioxane (OAPS). Cette cage siloxane possède trois isomères différents selon la position en meta, en ortho ou en para du groupement amine par rapport au groupement phényle. Une recherche de littérature a montré que les proportions des isomères étaient mal définies et que plusieurs voies de synthèse donnaient des résultats différents. Cela pose problème car la résistance et la sélectivité des membranes vont fortement dépendre de la nature des isomères. Un OAPS commercial contenant les trois isomères et un OAPS contrôlé contenant uniquement les isomères para et meta ont été caractérisés expérimentalement par picnométrie, spectroscopie infrarouge (IR), ainsi que par résonance magnétique nucléaire (RMN) unidimensionnelle et bidimensionnelle. Afin de mieux identifier les isomères, les résultats expérimentaux ont été comparés aux prédictions des spectres IR et RMN par des méthodes de mécanique quantique telles que la Density Functional Theory et par des méthodes d'apprentissage automatique telles que MestreNova. Le précurseur OAPS a ensuite été utilisé pour fabriquer de nouvelles membranes hybrides à thermorésistance élevée. Ces membranes polyOAPS-imides ont été caractérisées par IR, analyse thermogravimétrique (ATG), microscopie électronique à balayage (MEB), spectroscopie à dispersion d'énergie (EDS) et mesures de perméabilités. Elles ont pu conserver leur stabilité bien au-delà de 400°C et une sélectivité de ~10 pour les séparations H₂/CH₄ et H₂/N₂ à 250°C. En parallèle, un champ de force a été développé pour effectuer des simulations de dynamique moléculaire des trois isomères OAPS afin de caractériser leurs différences au niveau atomique.

I. Introduction

Due to the fast increase in population and the growth in energy consumption and industry, it is predicted that energy-related CO₂ emissions will increase by ~50% from 36 Gigatons in 2019 to 55 Gigatons in 2030 with the current and planned policies.¹⁻² The combustion of fossil fuels causes environmental issues and global problems, including greenhouse gases and air toxics, with the main contributor being coal.³ Over the past decades, there has been a growing interest in the development of membrane-based gas separation technologies to overcome the aforementioned environmental challenges and for specific applications such as O₂ enrichment, H₂ recovery and sweetening of natural gas. A growing interest in inorganic and polymer membranes for use in various industrial sectors has emerged. New fairly advanced polymers and copolymers such as polybenzimidazoles, thermally rearranged poly(benzoxazole)s and poly(benzoxazole-co-imides) have been produced.⁴⁻⁸

A membrane is generally defined as a selective barrier between two phases.⁹ The separation of gases can refer to any technique that may have multiple products or may be used to purify only one product. As shown in Figure 1, the feed is divided into two streams by the membrane, including the permeate, which is the part of the feed passing through the membrane, and the retentate, which is the part of the feed that cannot pass through the membrane. Depending on the purpose of the separation, the product will either be the permeate or the retentate.

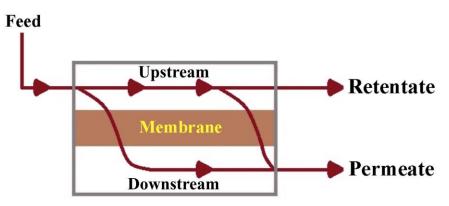


Figure 1. Separation of the components in a mixture by a membrane.

A gas takes the shape of its container and reaches a uniform density inside the container even in the presence of other objects. If gases are not confined to the container, they

are known as vapours that circulate in space. Atoms or molecules of matter move freely in the gaseous state and are less compact than in the solid and liquid states.⁹

The most important properties of gases include:

- 1. High compressibility: external forces can compress the gas and reduce its volume; when removing the external forces, the gas volume increases again.
- 2. Thermal expansion: when a gas sample is heated, its volume increases, and when it is cooled, its volume decreases.
- 3. Low viscosity: gases flow more easily than liquids.
- 4. Low density: the density of gases is in grams per cubic decimeter, while the density of liquids and solids is in grams per cubic centimetre.
- 5. Infinite miscibility: gases can be mixed in any proportion.

Important variables that describe a gas are its pressure P, temperature T and volume V. In a gas mixture, the behaviour of the gas depends on its number of molecules,⁹ unless the interactions of the gas-gas molecules are strong such as in the case of water vapour.

For the separation of gases, the most important membrane parameters are permeability and selectivity. Permeability is a volume current that passes through the unit of the membrane surface area for a known membrane thickness and at a given pressure difference (Equation 1):⁹⁻¹⁰

$$P = J \frac{l}{\Delta p} \tag{1}$$

where J is the gas flux through a membrane of thickness l, and Δp indicates the pressure difference across the membrane. J is defined by Fick's first law (Equation 2)¹⁰ as:

$$J = -D\frac{\partial C}{\partial x} \tag{2}$$

where *D* is the diffusion coefficient and $\left(\frac{\partial C}{\partial x}\right)$ is the penetrant concentration as a function of the distance *x* across the membrane.

2

By combining Equations 1 and 2 while assuming that the downstream pressure is negligible, one obtains Equation 3:

$$P = DS \tag{3}$$

where *S* is the solubility coefficient.

The separation factor $\alpha_{A/B}$, which refers to a mixture of gases, is defined as the ability of a membrane to separate gas *A* from gas *B* (Equation 4):¹⁰

$$\alpha_{A/B}^* = \frac{y_A / y_B}{x_A / x_B} \tag{4}$$

where x and y are the upstream and downstream mole fractions of A and B respectively. Equation 5 shows the ideal permselectivity $\alpha_{A/B}$ of the membrane, which is the ratio of the permeabilities or permeances of the individual gases:

$$\alpha_{A/B} = \frac{P_A}{P_B} \tag{5}$$

In other words, the difference in the permeation of the penetrants measured individually determines the selectivity of the membrane. Equation 5 is often a good approximation for the true separation factor (Equation 4), providing that there are no specific effects due to the mixture. For example, in the case of oxygen purification, $\alpha_{02/N2}$ is the extent to which the membrane can separate oxygen from nitrogen with oxygen being the desired penetrant.¹¹

The membrane separation technology selectively separates mixtures through pores and/or fine gaps in a continuous molecular structure. It is based on the physical or chemical interactions of the mixture's components with the membrane material. Membrane separations are often categorized by the size of the pores and the separation driving force. This includes microfiltration (MF), Ultrafiltration (UF), Nanofiltration (NF), Ion exchange (IE), Reverse osmosis (RO), and gas separation membranes, which have either the smallest pores or no pores.⁹ According to the IUPAC notation,¹²⁻¹⁴ microporous membranes have pore diameters of less than 2 nm (20 Å) and macroporous membranes have pore diameters greater than 50 nm (500 Å); the mesoporous category thus lies in the middle with pore diameters between 2

and 50 nm (20-500 Å). Nonporous membranes do not have any clearly-defined pores. Figure 2 shows different mechanisms of separation depending on the pore diameters, from the smallest to the largest:^{9, 14-15}

- Solution-Diffusion: this mechanism is generally considered to have a threestep process. Firstly, the gas molecules are adsorbed at the membrane surface on the upstream side. This is followed by the diffusion of the gas molecules through the matrix. Finally, the gas molecules desorb on the downstream side. This mechanism occurs for nonporous membranes, i.e. those based on polymers.
- Molecular Sieving: this mechanism happens when the pore sizes are sufficiently small (3.0–5.2 Å), leading to the separation of gas molecules that differ in kinetic diameter. The pore sizes of the membrane are so small that only the smaller gas molecules can permeate.
- Knudsen Diffusion: Knudsen separation can be obtained with mesoporous membranes having pore sizes smaller than 50 nm. It gives relatively low separation selectivities. In this mechanism, the gas molecules are separated according to their mean free path.
- Convection flow: when the pore diameters are longer than the mean free path of the gas molecules, the convection flow occurs, and no efficient separation can be achieved.

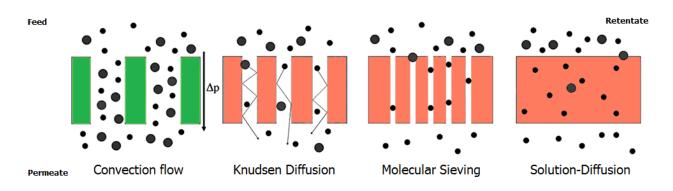


Figure 2. Convection flow and the three other possible mechanisms for membrane gas diffusion depending on the pore sizes.

Compared to alternative separation methods, the advantages of the membrane processes are the reduced power consumption, the low volume and the lack of need for a large space, the various possibilities in shapes, sizes and design, the low-pressure drop and the high mass transfer. Its disadvantages are the short membrane lifespan, the relatively low selectivities and permeabilities, the plasticization phenomena, the limited mechanical properties, the defects, the ageing and the high manufacturing costs.¹⁶ Within this context, inorganic nanomaterials have been used to provide higher permeabilities, selectivities and resistances in the new generation of mixed-matrix membranes (MMMs) and hybrid membranes. MMMs are polymer matrices in which inorganic fillers are dispersed at the nanometer level. They combine the resistance benefits of the inorganic particles and the polymeric and the inorganic phases being connected by covalent bonds.¹⁸

Membrane processes have now reached the first stage of maturity in many important processes such as:⁹

- Nitrogen or oxygen separation from the air;
- Separation of hydrogen from other gases such as nitrogen and methane;
- Hydrogen recovery from vapours in ammonia plants;
- Hydrogen recovery in oil refining processes;
- Methane separation from the other biogas compounds;
- Water vapour extraction from natural gas and other gases;
- CO₂ separation from natural gas;
- H₂S withdrawal from natural gas;
- Volatile Organic Liquids (VOL) removal from vapour atmospheres.

An example is the use of membranes for partial air separation, which was developed as a promising alternative to adsorption and cryogenics.¹⁹ Cryogenic distillation is a largescale process producing very pure gas and liquid products.²⁰ The basic choice between cryogeny or non-cryogeny mostly depends on the number of products that must be prepared (e.g. oxygen or nitrogen or both), the required production rates of each gas and/or liquid product and the required purities of the products.⁹ Compared to a membrane process, cryogenic air separation produces highly-purified gases or liquids. However, cryogenic distillation may be complicated in practice. Initially, pre-cooled air is compressed to 650 kPa in a multistage process. It is then throttled to cool to very low temperatures (below around -180°C) and liquefy. Nitrogen and oxygen are subsequently separated in a fractional distillation column. While liquid oxygen settles to the bottom of the column, nitrogen boils and can be removed through the top of the column. The remaining contents then move to a low-pressure distillation column to enhance the oxygen purity. The nitrogen stream is condensed to generate the reflux stream and liquid nitrogen is produced as well.²¹ The final result is oxygen with a purity of more than 99%. The cryogenic distillation process is illustrated in Figure 3.²²

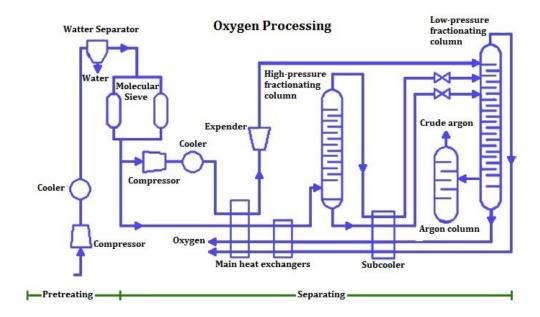


Figure 3. A cryogenic distillation system.²²

Adsorption separation (pressure swing adsorption or PSA) is another process in which gases tend to be adsorbed to solid surfaces under pressure. A gas is adsorbed under high pressures to a surface such as a carbon molecular sieve, and reducing the pressure then desorbs the gas. PSA can be used to separate a mixture of gases because different gases are adsorbed more or less strongly to different solid surfaces. Vacuum swing adsorption (VSA) is another technique that operates at near-ambient temperatures and pressures.⁹

The current polymeric membranes can compete with these traditional technologies in specific energy consumptions when the required product mole fraction is low. In air separation, an optimized membrane module with a selectivity of more than 8 can replace a large-scale cryogenic distillation if the required molar fraction of oxygen is 50%.²³ Table 1 presents a cost comparison of the three technologies for the removal of 8% N₂ from natural gas. The total capital cost and the processing cost of the membrane technology are much lower than for the other conventional techniques.

Technology	Total capital cost (US\$/Mscfd)	Processing cost (US\$/Mscf)
Cryogenic distillation	11200	1.30
Pressure swing adsorption	1300	1.65
Membranes	280	0.30

Table 1. Cost comparison of the technologies for the removal of 8% N₂ from natural gas.²⁴⁻²⁵

Mscfd: million standard cubic feet per day; Mscf: million standard cubic feet

Unfortunately, dense polymers are not very adapted to harsh conditions. Indeed, the gas separation performance of highly-thermostable polymers such as e.g. polyimides declines above 200°C, and by now, only a few high-temperature applications have been reported for such membranes.⁴ Table 2 illustrates the highest testing temperatures for gas separations by thermoresistant polymer-based membranes. T_g is the glass transition temperature, above which the polymer becomes too flexible and its mechanical performances severely decline. In addition, selectivity can decrease at high pressures of certain gases such as CO₂ because of plasticization effects.²⁶⁻²⁷

Heat-resistant polymeric membranes	Membrane material	Maximum GS test temperature (°C)	Maximum reported T _g (°C)
Fluorinated polymeric membranes	FBP-6FPT	75	253
Inorganic polymeric membranes	POPs	100	91
Carbocyclic polymeric membranes	PEEK	260	230
Heterocyclic polymeric membranes	PBI	400	436
Ladder polymeric membranes	BTFDA-DAB	90	> 400
Polymer blends membranes	BMI/PEI	35	> 400
Mixed matrix membranes	PBI/ZIF-90	400	> 400

Table 2. The highest testing temperatures for gas separations (GS) in thermoresistant dense membranes.²⁸

Within this context, the Films in Fluids group at the University of Twente (The Netherlands) recently prepared a series of novel hyper-crosslinked inorganic-organic ultrathin membranes using interfacial polycondensation. The idea was to combine the high thermoresistance of inorganic materials with the efficient separation properties of organic materials. These so-called polyPOSS-imides consist of a large molecular network of alternating polyhedral oligomeric silsesquioxanes (POSS), i.e. silicon oxide cages with the basic formula $R_nSi_nO_{1.5n}$ (n = 6, 8, 10, 12), and organic precursors at the basis of aromatic

imides.⁴ These hybrid membranes were indeed found to preserve their gas separations characteristics at elevated temperatures and pressures, i.e. up to 300°C and 60 bar^{4, 29} and as such, it can become a key-enabling technology for industrial-scale gas separations under harsh conditions.⁴

In collaboration with the University of Twente, molecular modelling studies were performed for these initial hybrid membranes at University Savoie Mont Blanc (USMB, France).³⁰⁻³¹ They pointed to a weak link in the structures, which is the aliphatic chain linker between the inorganic and the organic moieties. The present work aims to attain more stable hybrid ultrathin membranes by using a new POSS precursor. the octa(aminophenyl)silsesquioxane (OAPS), which has three isomers depending on the meta, para and ortho position of the amines on the phenyl groups.

- In the first stage, an extensive literature search was carried out on thermoresistant polymers and the POSS family.

- In a second stage, experimental investigations including density, infrared (IR), and nuclear magnetic resonance (NMR) characterizations were performed for both a commercial OAPS containing all three isomers and for a controlled OAPS containing only the *para* and *meta* isomers. To better identify the nature of the isomers, the experimental results were compared to predictions of the IR and NMR spectra by quantum mechanical and machine-learning methods.

- In a third stage, the commercial OAPS was used as a precursor to make novel hybrid polyOAPS-imide membranes with increased thermoresistance. These were characterized using IR, thermal gravimetric analysis (TGA), scanning electron microscopy (SEM), energy dispersive spectroscopy (EDS) and permeability measurements.

- In a fourth stage, a force-field was developed for OAPS to be used in subsequent molecular dynamics (MD) simulations, which provide direct structural and dynamical information at the atomic level as a function of time³² by integrating the classical equations of motion.³³

II. Thermoresistant polymers

The main impetus for the development of thermoresistant polymers initially came from the aerospace field.³⁴ During the late fifties and early sixties, synthetic activity was at its pinnacle. But it declined because of the decrease in aerospace expenses, the difficulties in manufacturing the required materials, and the cost of these materials.³⁴ Indeed, thermal stability and processability are two opposite requirements in the production of thermoresistant polymers. As a result, few polymers have achieved commercial exploitation.

The strength of chemical bonds affects the thermal stability,³⁴ and the primary bond energy between the atoms in the polymer chain is considered as being the source of thermal resistance.^{35,28} Table 3 shows a list of typical bond strengths.

C-S 273	C-N	307	C-H	416	P-O	528	C=C	609
	Si-H	319	C-F	428	P-C	580	C=N	617
	Si-C	328	Si-N	437				
	C-Cl	340	Si-O	445				
	C-C	349						
	C-O	361						

Table 3. Bond strengths (kJ/mole).^{34, 36}

The emphasis on bond strength resulted in the development of inorganic polymers. For example, in Table 3, the Si-O bond energy in polysiloxanes is almost 450 kJ/mole. The vibrational energy increases with heat, so the energy of bond dissociation is decisive in the thermostability.³⁷ To achieve thermoresistant polymers:³⁴

- (a) The most powerful chemical bonds should be used
- (b) Rearrangement reactions should not be made easy
- (c) Maximum resonance stabilization should be used
- (d) All the bond angles in the ring structures should be under the same tensions
- (e) There should be polybinding as much as possible.

The condition (a) is not often the limiting one, because most bonds are strong enough to be stable. Almost all the systems degrade because of condition (b) because the rearrangement reactions tend to increase with high temperatures. Since more energy is needed to break the bonds that have resonance stabilization, condition (c) is an advantage. Condition (d) refers to the angles keeping the atoms in close proximity when a bond is broken, so bond healing is possible upon dissipating the excess energy through the molecule. To meet condition (e), there should be more than one link for a skeletal atom in the polymer chain. These conditions relate to an ideal case, but in practice, a polymer has rarely a perfect structure.³⁴

Table 4 shows the relative ranking of the linking groups in polymers according to the temperature at which a 25% weight loss happens in a period of 2 hours in the presence or the absence of oxygen. The amide and carbonyl links show the highest thermal stability.

Table 4. Thermal and thermal-oxidative stability of some simple flexible linking groups. These are thetemperatures in °C for a 25% weight loss in 2h without and with oxygen, respectively.

Linking group	Thermal stability (°C)	Thermal-oxidative stability (°C)
c	500	389
о н с	500	431
F F F 	469	N/A
c	457	447
s	436	418
н н сс н н	429	383
н с н	408	N/A
0	N/A	368

Secondary forces such as hydrogen bonding and dipole-dipole interactions are another important factor since the sum of these forces determines the cohesion between neighbouring chains. These forces have a crucial influence on the T_g and the melting point T_m . As such, it is

usually possible to improve the polymer resistance by increasing the number of polar side groups or by making the arrangements of the chains more regular.

The most common polymers are non-crosslinked and are called thermoplastics. The chains can also be chemically crosslinked, thus leading to the thermosets, which are produced by the curing of prepolymers or polymers.²⁸ Curing is a chemical process by which the prepolymer or polymer in a viscous or solid-state is converted to a material with a higher molecular mass. This can be performed by chemical reactions induced by thermal curing, electron-beam irradiation (EB curing), photo-irradiation (photo-curing) or via mixing with a chemical crosslinking agent.^{35, 40} Molding or casting can be performed during curing by heating usually up to more than 200°C. Thermosets are often more brittle, more rigid, less soluble and harder than thermoplastics.²⁸ Low-crosslinked thermosets with long chains (between the crosslinks) are flexible while highly-crosslinked ones with shorter chains are rigid.²⁸

Degradation and thermal decomposition happen through pyrolysis or thermolysis of the polymers. As explained above, to overcome decomposition and degradation, the polymer should preferably include aromatic and/or heterocyclic backbones (not aliphatic backbones), resonance structures and strong chemical bonds with high bond enthalpies. As such, thermoresistant polymers are often made up of alternate moieties of aromatic rings and heterocyclic units, which is the case for e.g. polyimides (PIs) (Figure 4) and their derivatives.

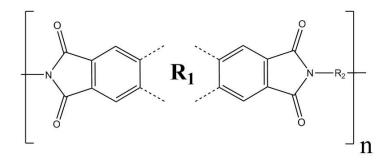


Figure 4. The general structure of a polyimide, R_1 and R_2 are usually aromatic and/or heterocyclic groups.

Durability in the range of 250-350°C is now well established for the PIs.⁴¹ Hydrogen bonding has also an important role in other highly polar polymers such as polyurethanes and polyamides.^{28, 37} As for the PIs, the resonance stabilization of the aromatic structures supplies further bond strength and thermoresistance.³⁵ Indeed, thermosets such as phenolic resins

present a high thermostability due to their high aromatic content, making them suitable for use under harsh conditions.⁴²

When polymers are exposed to heat, thermoplastic and thermoset polymers show different behaviours. Thermoplastic polymers are melted or softened and return to their solidstate by heating and cooling. In this repeatable process, the covalent structure of the thermoplastic polymer does not change,⁴³ and weak forces such as Van der Waals forces, dipole-dipole interactions or hydrogen bonding allow for these materials to be recycled.⁴⁰ On the other hand, due to the existence of crosslinks between the molecules, thermosets do not melt or soften but break down at elevated temperatures.⁴⁰ However, it should be noted that for highly-thermoresistant thermoplastic polymers, the molten temperature can be quite close to the degradation temperature too.

There is no standard definition for the thermal stability of polymers. For the user, the criterion is that the features of interest should be kept under the desired range of temperatures.³⁴ However, P. E. Cassidy ²⁰ has defined the thermoresistance of a polymer as:

- Maintenance of the polymer chemical features above 300°C in air measured by Thermogravimetric Analysis (TGA).
 or
 - ſ
- Maintenance of the polymer physical characteristics for a long time at 280°C in air evaluated by Isothermal Gravimetric Analysis (IGA).

TGA is a popular technique to study the deterioration of the polymers because of its wide range of applications and reliability:³⁴

- 1- Easy measurements
- 2- Wide choice of commercial instrumentation
- 3- Small samples (milligrams or less)
- 4- Can be accomplished under any atmosphere
- 5- No specific fabrication procedure is needed. Infusible and insoluble materials can also be tested
- 6- No limitations in heat and mass transfer at low heating rates.⁴⁴

IGA is a similar thermogravimetric technique, but it measures the weight change with time at a specified temperature, giving information on the material stability and composition. It is routinely used in industry.⁴⁵⁻⁴⁶ IGA is closer to the real conditions, but TGA is easier to use in an academic laboratory since it analyses the samples in a shorter time.⁴⁷

Since the mechanisms for loss of the thermal properties are not exactly addressed, other researchers have interpreted differently the definition of thermoresistant polymers. Indeed, terms such as "high" or "good" have been frequently used to describe thermoresistance without addressing quantification.²⁸ Some questions have been raised, such as the environment for studying thermoresistance, the stability of the material at a specific temperature, etc.³⁷

There are six categories of thermoresistant polymers:²⁸

- 1- Fluorinated polymers such as polytetrafluoroethylene (PTFE), which have high C-F bond energies and shielding of the carbon chains
- 2- Ladder polymers such as polypyrrone, which have polybondings, stiff chains and high $T_{\rm g}$
- 3- Inorganic polymers such as polyphosphazenes and siloxanes, which have strong bonds
- 4- Carbocyclic polymers such as polysulfones (PSFs) and polyamides (PAs), which have stiff chains, resonance stability and high T_g
- 5- Heterocyclic polymers such as PIs and polybenzimidazoles (PBIs), which have the same features as carbocyclic polymers.

There are also three other ways to boost the chain stiffness by incorporating inorganic or polymeric additives or by introducing crosslinks into the target polymer:

- 6- Hybrid inorganic/polymers;
- 7- Polymer blends;
- 8- Thermosets.

Figure 5 shows a spectrum of low to high-performance thermoplastics polymers in comparison to their cost. In this Figure, thermoplastics are divided into three main types according to their thermal, chemical, and mechanical stabilities. Some highly thermoresistant

thermoplastics are listed in Table 5.²⁸ The Continuous Use Temperature (CUT), also called the Relative Temperature Index (RTI), is the temperature at which 50% of the features (such as mechanical strength) are maintained after 100000 hours. It characterizes durability. The Heat Distortion Temperature (HDT) is the temperature at which a material distorts under a certain stress. The typical stress is 1.82 MPa but values between 0.45 and 8 MPa have also been used.⁴⁸

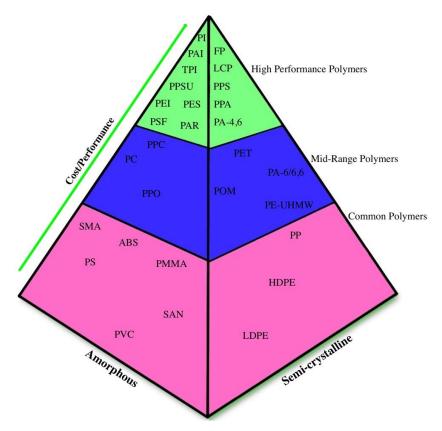


Figure 5. The pyramidal scheme of typical thermoplastic materials.²⁸

Table 5. Some highly thermoresistant thermoplastic polymers. ^{28, 49-}	Table 5.	Some highly	y thermoresistant	thermoplastic	polymers. ^{28, 49-3}
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Abbreviation	Polymer	T_{g} (°C)	CUT ^a (°C)	HDT ^b (°C)/Stress (MPa)
PAI	Polyamide-Imide	275	250	279/1.82
PI	Polyimide	250	288	246/1.8
PEE	Polyethersulfone	224	177-204	204-238/1.82
PEI	Polyetherimide	213	177-204	199-216/1.8
PSF	Polysulfone	190	149-171	171-182/1.82
PEEK	Polyetheretherketone	143	204-232	177-321/1.8
PPA	Polyphthalamide	134	204-232	117-260/1.8
PPS	Polyphenylene Sulfide	92	204-232	149-288/1.8

^a Continuous Use Temperature.

^b Heat Deflection Temperature.

From a practical point of view, the processability of the membrane is another factor that affects thermoresistance in addition to stability and durability. Figure 6 shows that these three factors have to be unified.



Figure 6. The three important factors to prepare thermoresistant membranes.²⁸

The processing difficulties (green in Figure 6) mainly result from the low solubility of thermoresistant polymers in the common solvents.²⁸ Some approaches to improve the situation are:

- 1- The use of the aprotic polar solvents such as DMAC, DMF, and NMP.⁵⁶⁻⁵⁷
- 2- Adding inorganic salts such as LiCl or CaCl₂ to the aforementioned solvents.⁵⁸⁻⁵⁹
- 3- The use of acids, such as HF, CHF₃, H₂SO₄ and sulfonic acid.⁶⁰
- 4- Thermal phase inversion, in which the polymer is dissolved at a high temperature in a solvent and the solution is cooled to precipitate the polymer.⁶⁰

As mentioned above, polyimides (PIs) are well suited to thermoresistant applications due to their remarkable thermo-oxidative durability and mechanical features.⁶¹⁻⁶² Besides, they exhibit high CO₂/CH₄, H₂/CO₂, H₂/N₂, and CO₂/N₂ selectivities. They are thus very good candidates for gas separation applications.⁶³ PIs are heterocyclic and can be aliphatic or aromatic. To overcome the difficulties of synthesis and processing caused by their limited solubility and infusibility, noticeable efforts have been done to incorporate bulky substituents, alkyl side chains, biphenylene, and twisty units into the polymer backbones.^{41, 62, 64-67, 68-71} Besides, it is known that polymers having –CF₃ exhibit boosted solubility and that the strong C-F bond further improves the chemical and thermal stability. Consequently, a lot

of studies have been carried out to explore the structural properties of $-CF_3$ substituted monomers in PIs. The $-CF_3$ group improves the fractional free volume (FFV) and disrupts the chain packing, which leads to increased permeability.²⁸ The PIs based on the 4,4'- (hexafluoroisopropylidene)diphthalic anhydride (6FDA) monomer (Figure 7) present indeed very high CO₂ permeabilities and CO₂/CH₄ selectivities:^{28, 72}

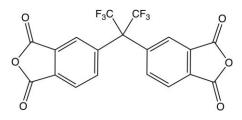


Figure 7. The 6FDA dianhydride.

Table 6 reports studies performed on the gas separation applications of 6FDA-based PIs:

Membrane	Full Name	P (atm)	H_2	O_2	N_2	СО	CH ₄
6FDA-ODA ⁷³	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and oxydianiline	1-10	-	4.3	0.83	23	0.38
6FDA-MDA ⁷³	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and	1-20	-	4.6	0.81	19	0.42
	methylenedianiline						
6FDA-IPDA ⁷³	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and	1-10	-	7.5	1.34	30	0.70
	isopropylidenedianiline						
6FDA-6FpDA74	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and 4.4'	10	-	16	3.40	64	1.60
	(hexafluoroisopropylidene)-dianiline						
6FDA-m-PDA ⁷⁵	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and meta-	1.4-8.2	20.3	2.6	0.36	8.2	0.14
	phenylenediamine						
6FDA-2,4-DATr ⁷⁵	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and 2,4-	1.4-8.2	87.2	7.4	1.31	28.6	0.71
	diaminotoluene						
6FDA-3,5-DBTF ⁷⁵	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and 3.5-	1.4-8.2	58.6	6.4	1.17	21.6	0.45
	diaminobenzotrifluoride						
6FDA-4BDAF ⁷⁶	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and 4.4'-	3-8.8	46	5.4	0.98	19	0.51
	bis(aminophenoxyphenyl) hexafluoropropane						
6FDA-3,3'-ODA ⁷⁶	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and 3,3' -oxydianiline	3-10	14	0.68	0.10	2.1	0.03
6FDA-3BDAF76	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and 3,3'-bis-	3-8.8	21	1.3	0.24	6.3	0.13
	(aminophenoxyphenyl) hexafluoropropane						
6FDA-p-PDA ⁷⁶	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and	2.8-6.8	23	2.1	0.38	11.8	0.18
1	paraphenylenediamine						
6FDA-mp'ODA77	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and 3,4'-oxydianiline	10	23.7	1.6	0.26	6.1	0.12
6FDA-APAP ⁷⁷	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and 2-(3-amino-(4-	10	38.2	2.89	0.47	10.7	0.22
	phenyl)-2-amino-phenyl)propane						

6FDA-BATPHF77	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and 2,2-bis(4-(4-	10	55.4	6.5	1.30	22.8	0.70
	amino-2-trifluoropropane dianhydride and 2,2-bis(4,4amino-2-						
	trifluoromethylphenoxy) phenyl) hezafluoropropane						
6FDA-DAF ⁷⁸	2,2-bis(3.4-dicarboxyphenyl) hexafluoropropane dianhydride and 2,7-	10	-	7.85	1.27	32.2	0.63
	diaminofluorene						

The studies of Sen and Banerjee⁷⁹ on other types of dianhydrides such as 4,4'oxydiphthalic anhydride (ODPA), 2,2-Bis[4-(3,4-dicarboxyphenoxy) phenyl]propane dianhydride (BPADA), 3,3',4,4'-benzophenone tetracarboxylic dianhydride (BTDA) and pyromellitic dianhydride (PMDA) (Figure 8) also showed dramatically higher permeabilities and selectivities than the typical PSF, cellulose acetate (CA) and polycarbonate (PC) membranes.

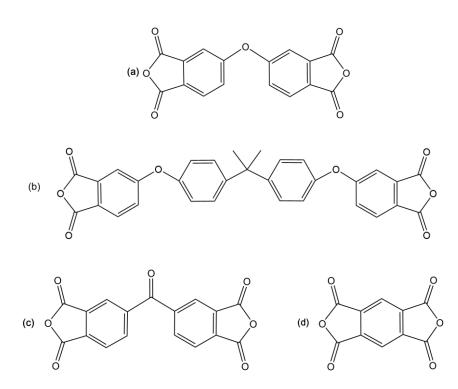


Figure 8. (a) ODPA, (b) BPADA, (c) BTDA and (d) PMDA.

Al-Masri et al.⁸⁰ synthesized dense 6FDA-based PIs membranes with methyl substituents at high yields (Figure 9), which were easily soluble in typical high boiling point aprotic solvents such as NMP, DMF, and DMSO, and also in low boiling point solvents such as THF, CH₂Cl₂, and CHCl₃. Increasing the ratio of the methyl to imide nitrogen content resulted in higher permeabilities, diffusivities, and solubilities, as well as T_g being higher than 300°C (Table 7).⁸⁰

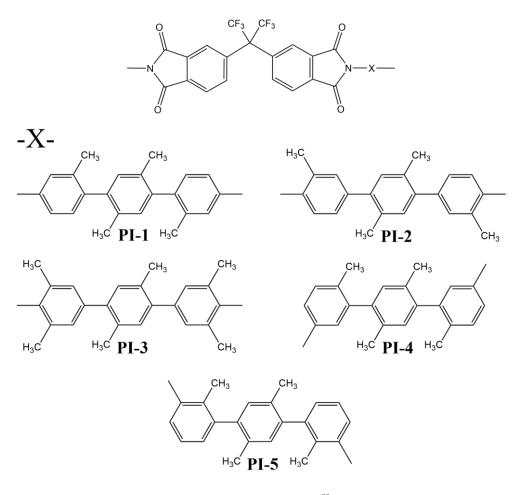


Figure 9. Different 6FDA-based PIs synthesized by Al-Masri et al.⁸⁰ Note that PI-3 has two methyl substituents on the first and third rings.

 Table 7. Single-gas permeabilities in Barrer for the different PIs of Al-Masri presented in descending order as compared to a polysulfone (PSF-1) membrane.⁸⁰

PI	P _{He}	P _{H2}	<i>P</i> _{CO2}	P ₀₂	P _{N2}	P _{CH4}
PI-3	230	350	360	67.0	16.5	15.0
PI-5	160	210	190	32.0	7.3	5.6
PI-4	84	100	62	12.0	2.4	1.9
PI-2	47	51	not measured	8.6	1.6	1.1
PSF-1	71	72	32	7.2	1.4	0.9

It is clear that isomerism makes a difference. A report on different PIs isomers⁸¹ indicated that the PMDA-4,4'-ODA polyimide has higher permeability and a lower selectivity toward He, CO₂, CH₄, O₂, and N₂ than its PMDA-3,3'-ODA isomer. The authors attributed the differences to the intrasegmental (rotational) mobility of the *para* PMDA-4,4'-ODA as compared to the *meta* PMDA-3,3'-ODA isomer. The phenyl rings in the *para*-isomer can rotate around their principal axis whereas no rotation is possible in the *meta*-isomer. Dynamic

mechanical analyses, as well as measurements of CO₂ diffusivity in both isomers, support this hypothesis. Similarly, a series of hyperbranched copolyimides (HBPI)s based on the commercially-available monomers ODPA (Figure 8), ODA, and 2,4,6-triaminopyrimidine (TAP) (Figure 10)⁸² show a considerably lower permeability for all gases than a linear ODPA-ODA polyimide. The permeabilities improved with increasing content of the ODA amine. However, low permeabilities are not always only a drawback. Compared to other polyimides, the HBPIs presented very high selectivities, which indicates that these materials can be used in gas separation applications such as CO₂/N₂, O₂/N₂, CO₂/CH₄, H₂/CH₄, etc.

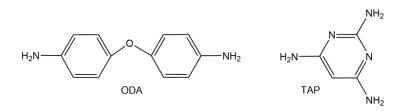


Figure 10. The ODA and TAP monomers.

In addition to the influence of the chemical structures and the high-temperature issues, the performance of the PIs membranes decreases when they are exposed to high pressures of penetrants such as CO₂. This phenomenon, which is called plasticization, reduces the T_g and the thermal stability while the permeability increases along with a loss in selectivity.⁸³

To further improve the thermostability of PI membranes, many methods such as polymer-polymer blending,⁸⁴ inorganic-polymer blendings⁸⁵, and crosslinking⁸⁶ have been employed. The latter is at the basis of the hypercrosslinked polyPOSS-imides networks designed by the University of Twente. For the current project, the 6FDA, PMDA, and ODPA dianhydrides (Figure 8) have been selected as possible organic precursors for the imides.

III. POSS and OAPS in the literature

Compounds containing Si-O have been studied for many years.⁸⁷ In the last 2 decades, the new field of silsesquioxane chemistry, based on the RSiO_{3/2} repeat units with R being an organic fragment, has gained interest. These materials between organic silicon polymers and inorganic compounds display versatile hybrid features, which are due both to their potentially-reactive and easily-modified R-Si part and to their thermally-stable and chemically-inert inorganic Si-O-Si part. Various structures from ladder polymers to highly ordered species with the common formula (RSiO_{3/2})_n, where n is usually 6, 8 or 10, have been synthesized.⁸⁸

POSS (Figure 11)⁸⁹ is the abbreviation of polyhedral oligosilsesquioxanes, i.e. n = 8. These chemically-cubic siloxane cages have nanoscale dimensions in the range of 1-3 nm depending on the organic substituent R along with a very robust inorganic framework.⁸⁷

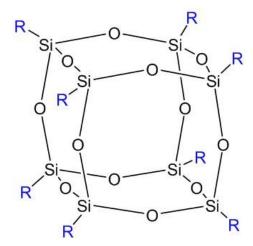


Figure 11. Polyhedral oligosilsesquioxanes (POSS)⁸⁹ - T₈R₈ structures.

The full nomenclature of POSS is complicated, so the nomenclature used for siloxane polymers is employed. A silicon atom connected to three oxygen atoms then connected to another silicon is defined by "T", and if R = H as in Figure 11, the abbreviation of the structure is T₈H₈. In T₈R₈, R can also be alkyl, aryl, alkoxy, siloxy, phenyl, aminophenyl, etc.⁸⁷ Many routes have been reported to synthesize POSS, but all of them can be listed in two reaction types:

a) Preparing POSS derivatives by making a T_n core from a precursor having fewer than n silicon atoms,

b) Chemically modifying the functional groups of an existing T_n core to get a new POSS derivative.⁸⁷

In recent years, the incorporation of POSS into polymer matrices has attracted attention in order to create nanocomposites with enhanced mechanical and thermal properties.⁹⁰ Their rigid cage-like structure, nanoscale size and ability to tailor properties by introducing different functional groups make POSS effective candidates for high-performance membranes.⁹¹ They can be introduced by physical blending, chemical polymerization, and crosslinking.⁹²⁻⁹⁴ The thermal stability of the POSS-based membranes depends on the amount of POSS in the polymer matrix, the POSS-polymer interface interactions and the amount of crosslinking. In addition, the local polymer topology, interactions and mobility can be changed and as a result, macroscopic features such as the T_g , thermal and mechanical stability can be improved. Indeed, the preparation of hybrid membranes by incorporating inorganic fillers into high-performance polymers is aimed at overcoming the weaknesses of both types of materials, e.g. the poor processability of the inorganic fillers and the low thermal and mechanical stabilities of the polymers. However, the preparation of defect-free and robust nanocomposite membranes still remains a challenge.²⁸

As far as gas separation properties are concerned, POSS has been shown to affect the permeability characteristics in mixed-matrix membranes (MMM). Incorporation of a commercial poly(ethylene glycol)-functionalized POSS (PEG-POSS) in two grades of poly(ether block-amide), PEBAX[®] MH 1657 and PEBAX[®] 2533, showed a higher permeability for both MMMs and a higher selectivity for PEG-POSS/PEBAX[®] 2533.⁹⁵ A MMM comprising of octa(aminophenyl)silsesquioxane (OAPS) in 6FDA-methylenedianiline (6FDA-MDA) up to 20% wt displayed a decrease in the permeabilities for all the gases tested, except for CO₂ with increasing OAPS loading.⁹⁶ Another report showed improved selectivities for CO₂ over H₂ in POSS/PEBAX[®] MMMs.⁹⁷ The authors attributed this observation to stronger interactions between the polymer and CO₂ and reduced transport channels for H₂. As far as hybrid membranes are concerned, Dasgupta et al.⁹⁸ synthetized polyimide-POSS (PI-POSS) membranes with the POSS being attached at both ends of the polyimide chains. The permeability of CO₂, N₂, O₂ and CH₄ increased, while the selectivities for gas pairs such as CO₂/CH₄ and O₂/N₂ were comparable to those in the pure PI membranes.

Nieck E. Benes, Michiel J. Raaijmakers and colleagues at the University of Twente have recently synthesized hybrid hypercrosslinked polyPOSS-imide membranes, where the inorganic POSS is linked through an extensive periodic network by organic aromatic imide bridges.^{4, 99-100} As shown in Figure 12, an organic solvent containing the organic monomer and an aqueous phase comprising an octa(aminopropyl)silsesquioxane POSS are contacted. A polycondensation reaction occurs at the interface and this results in a polyPOSS-(amic acid) intermediate. The subsequent step is a heat treatment to convert the amic acids to cyclic imides (imidization). The chemical structures involved in the polycondensation and imidization steps are provided in Figure 13. In the following part of this work, the name be "POSS" will used either for the entire family or for this specific octa(aminopropyl)silsesquioxane when referring to the work of the University of Twente.

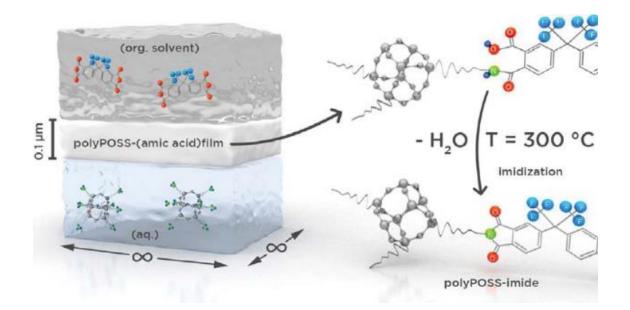


Figure 12. Interfacial polycondensation reaction of 6FDA in toluene and octa(aminopropyl)silsesquioxane POSS in water and the subsequent conversion of the amic acid to the cyclic imide (imidization) by a heat treatment up to 300°C.⁴

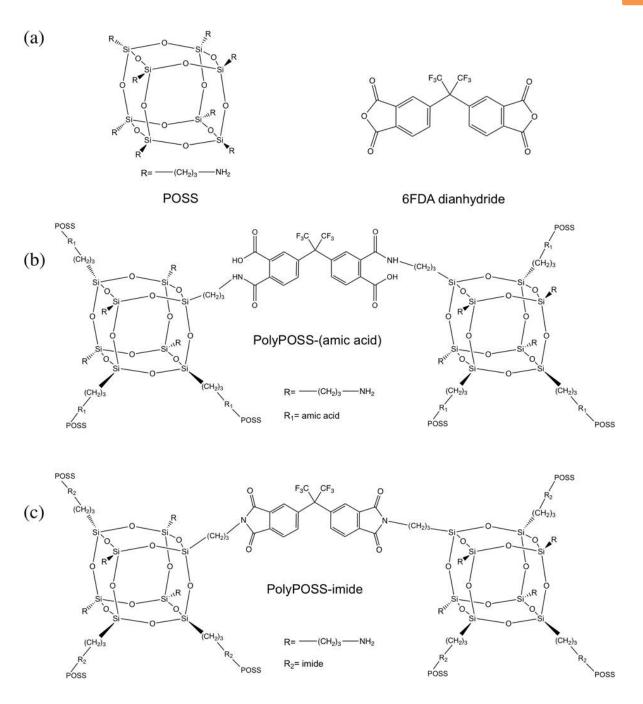


Figure 13. Chemical structures of (a) the POSS and 6FDA dianhydride monomers, (b) the interfacial polycondensation leading to the polyPOSS-(amic acid), and (c) the imidization leading to the 6FDA polyPOSS-imide network.³¹

As can be seen in Figure 14, the polyPOSS-imide thin films present various gas separation and permeabilities at a trans-membrane pressure of 2 bar and over a large range of temperatures, which arise from the rigid POSS and the nature of the mobile organic moieties (here, BPDA and BPADA, chemical structures provided on the right side of Figure 14).⁹⁹⁻¹⁰⁰

Chapter 3

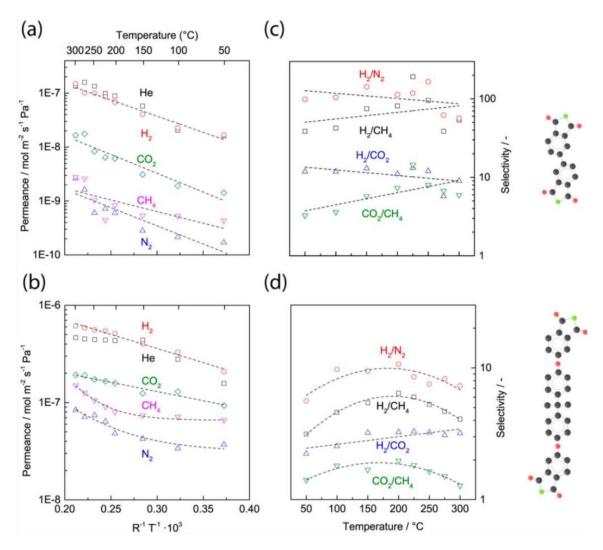


Figure 14. Left side: Arrhenius plot of the logarithm of the single-gas permeance as a function of 1/T for polyPOSS-imides derived from (a) BPDA and (b) BPADA. Right side: the corresponding ideal gas selectivities for H₂/N₂, H₂/CH₄, H₂/CO₂ and CO₂/CH₄ as a function of T.¹⁰⁰

Raaijmakers et al. reported the polyPOSS-imide activation energies for gas permeation as being 5–7 higher than those in conventional polyimides.⁴ They attributed these relatively high activation energies to the high-energy barriers for gas diffusion in the rigid networks. Some of these membranes could maintain the selectivity for H₂/CH₄ and H₂/N₂ and even increase the selectivity for H₂/CO₂ up to 300°C. A TGA analysis confirmed that no degradation happened below 300°C.¹⁰⁰ The impact of elevated pressures of plasticizing gases on the sorption, dilation and swelling of the polyPOSS-imide membranes was also investigated. Compared with conventional polyimides such as Matrimid[®] (3,3',4,4'-benzophenone tetracarboxylic dianhydride and diaminophenylindane), the sorption and swelling of polyPOSS-imide membranes were found to be much higher after exposure to CO₂

and CH₄ under pressures up to 50 bar. It was attributed to the large quantity of open space in the hybrid networks and the flexible responses of the networks.²⁹

Studying polyPOSS-imide networks via molecular dynamics (MD) simulations confirmed their remarkable abilities for gas separations up to 300°C.²⁶ A comparison of uncrosslinked PMDA and POSS mixtures, the amic-acid intermediates, and the final PMDA polyPOSS-imide networks indicated that the volume shrinkage during imidization was less than the quantity of water removal, thus leading to a large void space and an improved gas solubility in the polyPOSS-imides. The model membranes also showed solid-like elastic moduli at room temperature and at 300°C when subjected to uniaxial tension.³⁰ In another study, new molecular models were obtained for the flexible 6FDA precursor and compared to those with the shorter and rigid PMDA. Crosslinking 6FDA with POSS led to a larger thermally-induced dilation of the networks and more volume loss per H₂O over imidization than for PMDA.³¹ Figure 15 shows the isotropic dilation for the 6FDA and PMDA polyPOSS-imide is more resistant but only at high dilations. The -(CH₂)₃- linker was found to play the main role in the first part of the deformation for both samples and as such to be a "weak link".

POSSs with stronger R substituents such as phenyl groups should offer a higher thermostability than the aliphatic groups. However, the new POSS precursor should also have NH₂ functional groups, as the polycondensation requires them to react with the dianhydrides and produce the links between organic and inorganic fragments. Since the imide part proved its reliability in thermoresistant membranes,^{4, 99} the dianhydride precursors can be conserved.

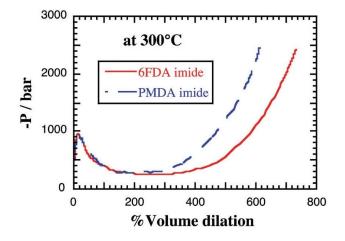


Figure 15. Isotropic dilation for model 6FDA and PMDA polyPOSS-imide networks at 300°C.³¹

Although the phenyl derivatives of POSS offer high thermal stability (> 400°C/air), they have not attracted as much attention as the other types of POSS, because the most typical phenyl POSS, i.e. [PhSiO_{1.5}]₈ referred to as octa(phenyl)silsesquioxane (OPS), presents a low solubility (less than 200 mg in 1 litre of boiling CH₂Cl₂). Therefore, the isolation of this compound is difficult and it can only be prepared in amounts of hundred milligrams. A better yield can be obtained via base-catalyzed synthesis.¹⁰¹⁻¹⁰⁴ Figure 16 shows the reaction pathway for the synthesis of OPS by Brown et al.¹⁰²

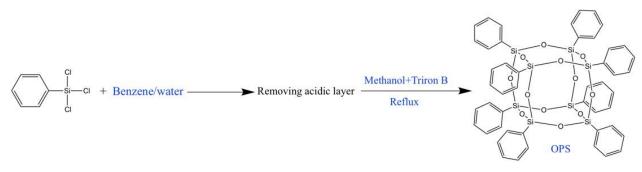


Figure 16. Octa(phenyl)silsesquioxane (OPS).¹⁰¹

However, as noted above, our new POSS precursor will require -NH₂ groups to carry out polycondensation. Nitration of OPS was the first electrophilic substitution reported by Olssen and Gronwall in 1961.¹⁰⁵ The authors described the NO₂ nitro groups of the resulting octa(nitrophenyl)silsesquioxane (ONPS) as inert since they could not reduce them to amines, and consequently, no new effort was reported for 38 years.¹⁰¹ Tamaki, Laine et al.¹⁰⁶ revisited the reaction and realized that these nitro groups can be easily reduced using formic acid/Et₃N reducing Pd/C the agent with as the catalyst. This produced as octa(aminophenyl)silsesquioxane (OAPS).

Figure 17 shows the three possible pure isomers for OAPS depending on the position of the amine substituents, i.e. the *meta*, *ortho* and *para* forms. However, in practice, these are ideal structures. In reality, the synthesis route summarized in Figure 18 leads to a mixture of OAPS cages with the different groups on the same cage being probably in different forms. For example, a cage can have five -NH₂ groups in the *meta* position, two in the *ortho* position, one in the *para* position, and so on.

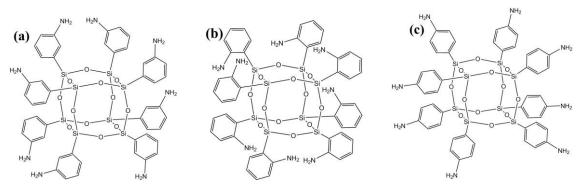


Figure 17. The pure (a) meta (b) ortho and (c) para isomers of OAPS.

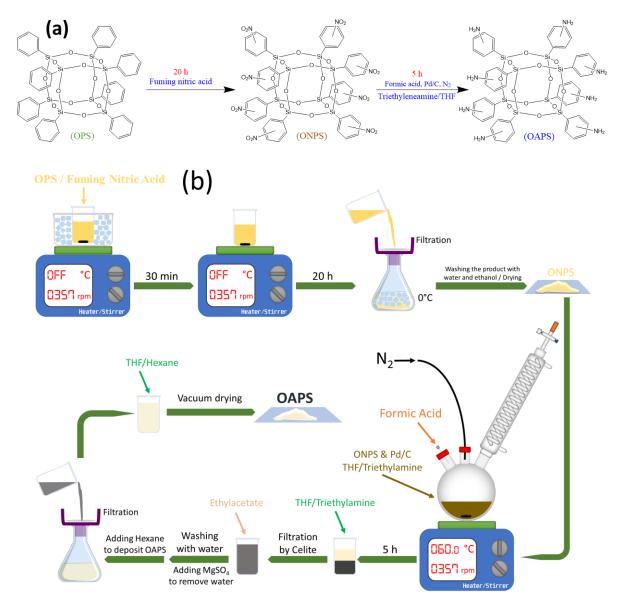


Figure 18. (a) The summary and (b) the details of the synthesis of ONPS (octa(nitrophenyl)silsesquioxane) and OAPS (octa(aminophenyl)silsesquioxane) from OPS (octa(phenyl)silsesquioxane) by Laine's method.¹⁰⁶

Rather than modifying the OPS structure, another less-used route is to start directly from aminophenyltrimethoxysilane (APTMS) precursors for the aminophenylsilsesquioxane

(APS) structures. This was done by Lee et al.,¹⁰⁷ by using the *meta* and *para* isomers of APTMS shown in Figures 19(a-b). The *ortho* isomer was not included. Both *meta*-aminophenylsilsesquioxane (*meta*-APS) and *para*-aminophenylsilsesquioxane (*para*-APS) isomers were synthesized in the same way (Figure 19(c)): their *meta*-APTMS and *para*-APTMS precursors reacted individually with tetrabutylammonium hydroxide and were then quenched with glacial acetic acid. The separate products (pure *meta*-APS and pure *para*-APS) each led to a mixture of cage sizes ranging from 8 to 14 silicon units with the dodecamer being predominant. This alternative approach allows for better control of the isomer proportions, but its downside is that the cage sizes are less well-controlled than with Laine's nitration/reduction of OPS (Figure 18).

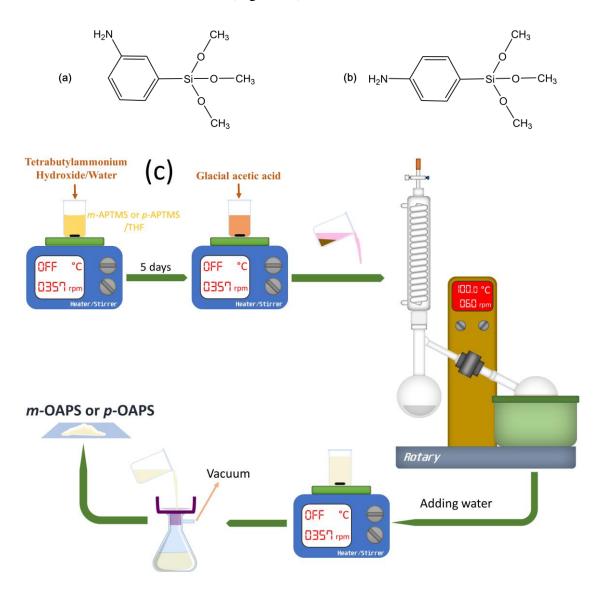


Figure 19. The (a) *meta* and (b) *para* isomers of APTMS (small precursors for isomer-specific OAPS) along with (c) the details of the synthesis of OAPS from APTMS by Lee's method.¹⁰⁷

More recently, the Norwegian SINTEF research organization,¹⁰⁸ which is a partner of the University of Twente, used a controlled sol-gel conversion upon hydrolysis and condensation of the *meta* (30-40%) and *para* (60-70%) APTMS precursors bought from Gelest.¹⁰⁹ The final product is a 21% w/w solution of *para* and *meta* isomers of OAPS in DMSO. Although Gelest can produce the APTMS isomers individually, the mixture of *meta* (30-40%) and *para* (60-70%) isomers of APTMS is less expensive.¹⁰⁹

As a precursor for hybrid structures, OAPS has already been used in blends^{94, 96, 110} or in covalent-bond reactions.^{90, 98} An example is the polyOAPS-PMDA structure (Figure 20),¹¹¹ which was built from both monomers reacting directly in NMP (and not through interfacial polymerization). It was able to show a thermal stability up to 500°C with a 5% mass loss.

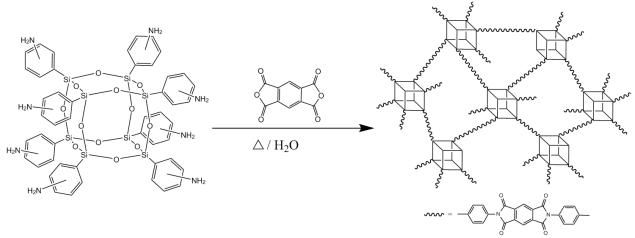


Figure 20. PolyOAPS-PMDA synthesized in NMP.111

Unfortunately, this polyOAPS-PMDA was also found to be brittle and efforts were made to employ more flexible and longer dianhydrides such as ODPA (Figure 21).¹⁰¹ Networks built from OAPS/ODPA/ODA (4,4'-oxydianiline)¹¹² and OAPS/ODA/BMI (bismaleimide)¹¹³ have been reported as well.

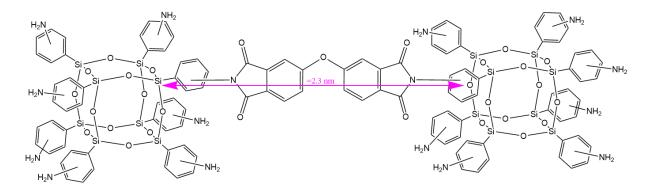


Figure 21. PolyOAPS-ODPA synthesized in NMP.¹⁰¹

The problem with mixing the monomers in the same solvent is that a shapeless product appears, whereas it would be preferable to produce a defect-free, uniform and thin layer to operate as a gas separation membrane. This is the main reason for attempting to use the interfacial polymerization route, as was done by the University of Twente (Figure 12). The difficulty of this approach is that it requires two immiscible solvents, one for the OAPS precursor and one for the dianhydride precursor, in order to create an interface where the thin layer can be generated.

The off-white OAPS powder is soluble in THF, DMF, 1,4-dioxane, acetone, DMSO and ethyl acetate but it is insoluble in water, benzene, ethanol and petroleum ether.¹¹³ The dianhydrides are soluble in DMF, NMP, DMAc, THF, toluene, and so on.¹¹⁴ For making polyOAPS-imide thin layers, two immiscible solvents should be chosen. Table 8 shows different solvents and their miscibility.¹¹⁵

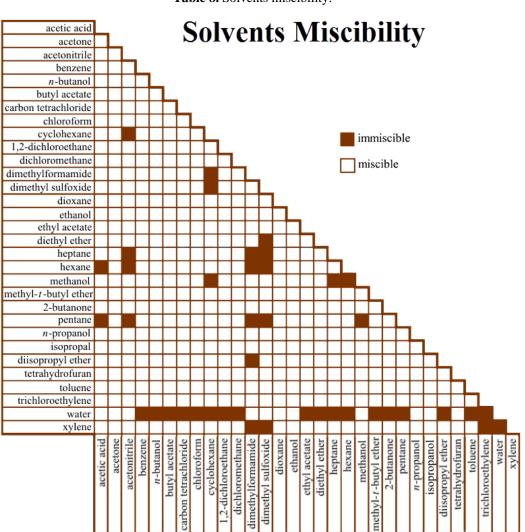


 Table 8. Solvents miscibility.¹¹⁵

Although there are different possibilities, the typical solvents used at the University Twente are toluene for the dianhydride, water for the octa(aminopropyl)silsesquioxane POSS and DMSO/water for OAPS. Since DMSO (the solvent for OAPS) and toluene (the solvent for the dianhydride) are miscible, a certain amount of water (which is immiscible with toluene) must be added to the DMSO in order to have two immiscible phases.

As far as the structure of the new inorganic precursor is concerned, several experimental techniques were used to characterize OPS, ONPS and OAPS, i.e. the intermediate and final products in Laine's nitration/reduction approach (Figure 18) as well as those obtained by Lee et al. from the APTMS precursors (Figure 19). Many efforts focused on elucidating the ratios of the different *para*, *meta* and *ortho* isomers of OAPS.

1) Infrared spectroscopy (IR)

Almost any compound containing covalent bonds absorbs radiation frequencies in the infrared region of the electromagnetic spectrum. This region lies at approximately 700 nm to 1 mm in terms of wavelengths. An important use of IR spectroscopy is to obtain structural information since the absorption of each bond type (N-H, C-H, O-H, C-X, C=O, C-O, C-C, C=C, C-N, and so on) is only found in specific portions of the vibrational infrared region.¹¹⁶ This is typically 400 to 4000 cm⁻¹ in terms of wavenumbers.

Figure 22 displays the comparative IR spectra of OPS, ONPS and OAPS, i.e. the route towards the functionalization of amine groups in the nitration/reduction approach (Figure 18).¹¹⁷ The peak at 1115 cm⁻¹ corresponds to the Si-O-Si stretching in the cages,¹¹⁷ and is present in all three compounds. The two peaks at 1350 and 1530 cm⁻¹ in the IR spectra of ONPS disappear after reduction to OAPS. The two peaks at 3220 and 3380 cm⁻¹ can be attributed to the -NH bond in OAPS. These comparative IR spectra show that the OPS phenyl rings have indeed been functionalized by -NH₂ in the final OAPS product.

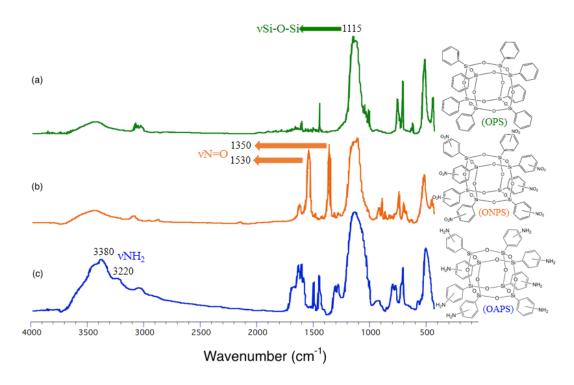


Figure 22. IR spectra of (a) OPS, (b) ONPS and (c) OAPS.¹¹⁷

2) Nuclear magnetic resonance (NMR)

Nuclear magnetic resonance (NMR) provides information on magnetically distinct atoms. For instance, when hydrogen nuclei are studied, one can determine the number of each distinct type of H and information about their immediate environment. Resonance of the ¹H hydrogen in a molecule can be detected since this isotope is magnetically active, has a spin of ½ and can be found in most organic compounds. Similar information can be achieved for the carbon nuclei (the ¹³C nucleus also has a spin 1/2 but has a natural abundance of only 1%)¹¹⁸ and for some other atoms such as silicon, nitrogen, oxygen, fluorine, etc.¹¹⁶ For the OAPS and related structures, hydrogen nuclear magnetic resonance (¹H-NMR), carbon-13 nuclear magnetic resonance (¹³C-NMR) and silicon-29 nuclear magnetic resonance (²⁹Si-NMR) characterizations were all attempted on ONPS and OAPS obtained by the nitration/reduction route (Figure 18). Only the ²⁹Si-NMR spectra for OAPS have been reported following the synthesis route by Lee et al. (Figure 19).¹⁰⁷

Since the ¹H-NMR spectrum of ONPS is sharper and clearer than that of OAPS, Tamaki et al.¹⁰⁶ measured the isomer ratios by using both ¹H-NMR and ¹³C-NMR on ONPS (Figure 23). The assumption was that the isomer ratios of OAPS (which has amines) should be the same as those of ONPS (which has nitro groups) since, based on the reduction reaction, each -NO₂ should be displaced by one -NH₂.¹⁰⁶ Steric hindrance and electronwithdrawing effects of the POSS core were expected to prevent the formation of the *ortho* isomer. As such, the authors interpreted the ten peaks of the ONPS ¹³C-NMR spectrum (Figure 23(a)) as ten different carbon environments for the *meta* and *para* isomers. There were three small extra peaks at 123, 137, and 159 ppm, but they were described as "insignificant". The triplet peak appearing at 8.73 ppm in the ¹H-NMR spectra (Figure 23(b)) was assigned to the most positively charged hydrogen placed between the nitro group and the siloxy group for the *meta* isomer, which the authors used to calculate the *meta*-isomer ratio. The surface area of the shift at 8.73 ppm was assumed to be equal to the unit, and the total surface area of the shift at 8.73 ppm multiplied by the 4 *meta* phenyl hydrogens with respect to the total surface area of all phenyl hydrogens is (4 /7.76) \approx 0.52. The ratio of the *meta* isomer was thus found to be 52%.

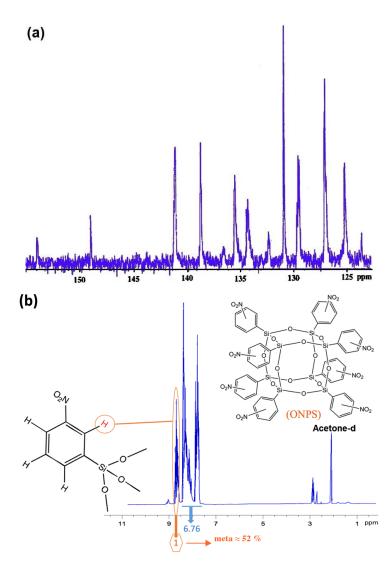


Figure 23. The experimental (a) ¹³C-NMR and (b) ¹H-NMR spectra of ONPS in acetone-d6.¹⁰⁶

By assuming that the *ortho* isomer was negligible, the ratio of *meta* to *para* could then be calculated as being 52% to 48%.¹⁰⁶ By using TGA analysis, the authors revealed that the decomposition of both ONPS and OAPS begins at the earliest at 350°C.

Kim et al.¹¹⁹ produced OAPS in a similar way to Tamaki et al.¹⁰⁶ (Figure 18) but with some modifications on the conversion of ONPS to OAPS. They achieved a higher yield of OAPS by using a lower concentration of the reagents over a longer time (20 h) and at a lower temperature (40°C). As can be seen in Figure 24, the authors used the cleavage of the Si-C bond in two different ways (peroxidation and D₂O/F⁻) to detach the rings from the ONPS corners and isolate the organic substituents. The ¹H-NMR spectra of nitrophenol and deuterated nitrobenzene were then used to re-measure the isomer ratio of ONPS. Although the ¹H-NMR spectra were not provided in the paper, the ratios were reported as being 70%, 25% and 5% for *meta, ortho* and *para* respectively for both cleavage methods. Takahashi et al.¹²⁰ redid the Si-C cleavage using peroxidation and obtained 76%, 21% and 3%, whereas Sulaiman et al.¹²¹ reported 65%, 10% and 25% for *meta, ortho* and *para* respectively, which likewise indicated the presence of the three isomers. It should be noted that the corresponding author of the papers by Tamaki et al.,¹⁰⁶ Kim et al.,¹¹⁹ Takahashi et al.,¹²⁰ and Sulaiman et al.¹²¹ is Prof. Richard Laine, so these results originate from the same group using slightly modified routes.

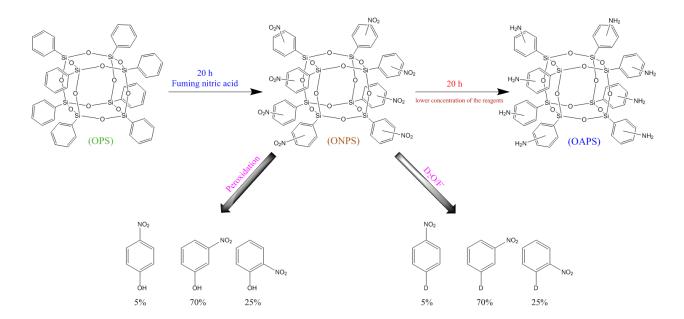


Figure 24. Laine's modified method at the ONPS step, Si-C cleavage of ONPS and the resulting isomer ratios.¹¹⁹

In another modification of Laine's method, Zhang et al. shortened the time of nitrifying OPS to 6 hours and used hydrazine as well as FeCl₃ instead of the expensive Pd/C catalyst to reduce ONPS to OAPS.¹¹³ The yields of the yellow ONPS and the off-white OAPS were 92% and 86% respectively. But there were no attempts to calculate the isomer ratios.

Although Tamaki et al.¹⁰⁶ had written that the formation of *ortho* isomer was unlikely, Kim et al.¹¹⁹ claimed that 25% of the isomers were *ortho* and that there was very little *para*.¹¹⁹ Due to the electron-withdrawing nature of the cage and the steric hindrance, Tamaki et al.¹⁰⁶ expected to have a majority of *meta*, then *para* and then *ortho* at last (or no *ortho* at all). These conflicting observations, which have not been resolved yet, are definitely a problem when OAPS is obtained by the nitration/reduction approach.

However, the assumption that the isomer ratios are the same in ONPS and OAPS was confirmed by the ¹H-NMR spectrum of OAPS (Figure 25).¹²² Tamaki et al.¹⁰⁶ and Huang et al.¹¹⁷ calculated the integration ratio of the peaks for the amine groups with respect to the aromatic groups to obtain the amount of quantitative reduction of the nitro to amine groups. This ratio is 2:1 for aromatic to amine H, which indicates that there is one amine per phenyl ring and that the nitro groups are completely displaced with amine groups.

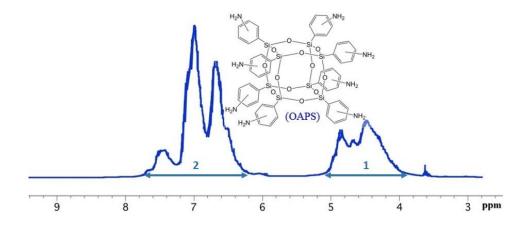


Figure 25. The ¹H-NMR spectrum of OAPS in acetone- d_6 and the ratio of aromatic H (6.2 to 7.8 ppm) to amine H (3.9 to 5 ppm).¹⁰⁶

The ¹³C-NMR spectrum of OAPS in Figure 26 apparently shows ten peaks. Tamaki et al.¹⁰⁶ interpreted the same number of the peaks in the ONPS (Figure 23(a)) and the OAPS ¹³C-NMR spectra as the presence of ten different carbons for the *meta* and *para* isomers. But by assuming that the results of Kim et al.¹¹⁹ are correct, there should be 16 peaks: 4 for *para*,

6 for *meta*, and 6 for *ortho*. Since the peaks overlap and some peaks are too small, it is not clear how many individual peaks appear. That may be the reason why Tamaki et al.¹⁰⁶ used the description "insignificant peaks" for the three smallest peaks in the ¹³C-NMR spectrum of ONPS, without attempting to provide an explanation for their presence.

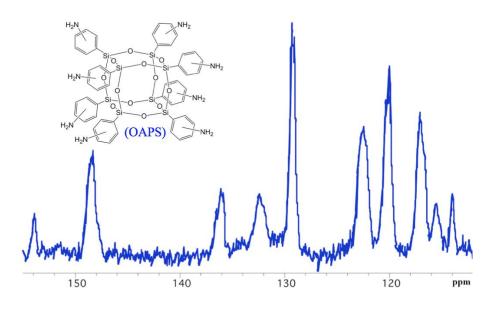


Figure 26. ¹³C-NMR spectra of OAPS in acetone-*d*₆.¹²²

The solid-state ²⁹Si-NMR spectra of ONPS and OAPS show two overlapping peaks (Figure 27).^{113, 122-123} These studies were published before the work of Kim et al.¹¹⁹ determining the isomer ratios via the Si-C cleavage of ONPS. Zhang et al.¹¹³ concluded that these peaks were related to the *meta* and *para* isomers, but as explained before, it was not known at the time that the *ortho* should also be present. This remains again unresolved. The peak is much more narrow in the ²⁹Si-NMR spectra of the simpler OPS, which suggests that all Si atoms feel the same environment there.

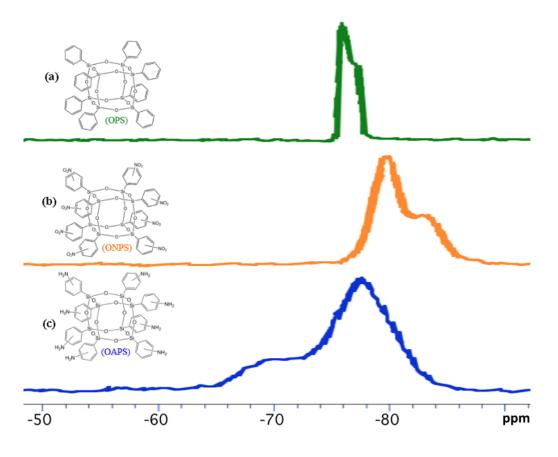


Figure 27. Solid-state ²⁹Si-NMR spectra of OPS (a), ONPS (b) and OAPS (c).^{113, 123}

To assess their alternative synthesis approach based on aminophenyltrimethoxysilanes (Figure 19), Lee et al.¹⁰⁷ purchased a commercial OAPS prepared through nitration/reduction by Mayaterials Inc., and compared its ²⁹Si-NMR spectra with that of the pure *para*-APS and *meta*-APS obtained from the APTMS precursors. We refer to APS (aminophenyl silsesquioxane) when the cages can have different sizes, i.e. they are not necessarily octamers as in OAPS. As can be seen in Figure 28, the commercial OAPS in (c) shows broad resonances in the ²⁹Si-NMR spectrum while the *para*-APS (a) and *meta*-APS (b) exhibit much more narrow spectra. The authors suspected that the nitration/reduction of OPS did not fully preserve the integrity of the Si-O cage structures and that OAPS was converted into an ill-defined resin when using Laine's approach. They attributed the many narrow peaks in Figures 28 (a) and (b) to the presence of cages with different numbers of Si and identified the dodecamer as being probably predominant. Indeed, the main orange peak of *meta*-APS matches the sharpest peak of the commercial OAPS, but the main purple peak of *para*-APS corresponds to a region where there are smaller peaks.

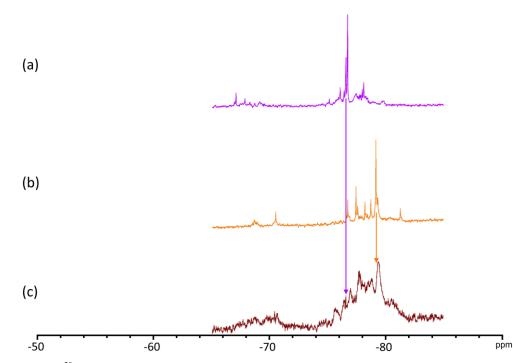


Figure 28. The ²⁹Si-NMR spectra of APS (with the dodecamer being supposed to be predominant) obtained either from (a) *para*-APTMS, (b) *meta*-APTMS along with that of (c) OAPS produced through nitration/reduction of OPS.¹⁰⁷ All spectra were measured in DMSO-*d6*.

3) X-ray diffraction (XRD)

XRD crystallography is widely used to obtain information on the atomic-scale structure.¹²⁴ Figure 29 presents the XRD pattern of OAPS obtained by nitration/reduction.¹¹³ The sharper peak at 2θ =8° shows the long-range order in OAPS (d = 1.1 nm), which corresponds to the diameter of the OAPS cage. The broad halo centred at 2θ =19°, which usually appears in the XRD pattern of amorphous materials, is the result of random packing at the molecular level caused by the amine-substituted isomers.¹¹³

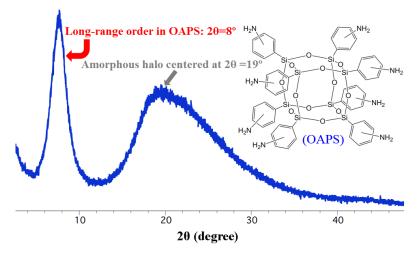


Figure 29. The XRD pattern of OAPS.¹¹³

Lee et al.¹⁰⁷ also used XRD to assess the differences between their synthesis from aminophenyltrimethoxysilanes and the nitration/reduction approach (Figure 30). All three OAPS showed the same reflection peak at $2\theta = 7.8^{\circ}$ (attributed to a d-spacing of 1.126 Å) which corresponds to the peak at 8° reported in Figure 29. However, the intensity varied in the order pure *para*-APS > pure *meta*-APS > OAPS by nitration/reduction. The authors concluded that the lower intensity of the latter was probably due to its ill-defined resin-like structure. They also suggested that since *para*-APS contains only one possible site for its NH₂ substituent, it displays more order than *meta*-OAPS having two possible sites. However, it remains unclear why the XRD pattern for the pure *meta*-OAPS (intact cages according to Lee et al.) looks closer to that of the nitration/reduction OAPS (ill-defined cages according to Lee et al.), whereas their NMR spectra are much more different (Figure 28). It should also be remembered that there is a mixture of cage sizes in the Lee approach, whereas all of them should be eight-cornered in the Laine approach.

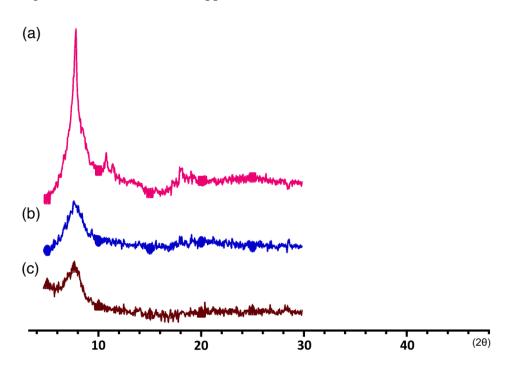


Figure 30. XRD patterns of (a) *para*-APS cage mixtures (p-NH₂C₆H₅[SiO1.5])_n, (b) *meta*-APS cage mixtures (m-NH₂C₆H₅[SiO1.5])_n where n=8,10,12; and (c) OAPS prepared by nitration/reduction. Curves are vertically shifted for clarity.¹⁰⁷

4) Gel permeation chromatography (GPC)

Gel permeation chromatography or GPC is a technique used for the determination of molecular weight distributions (MWD).¹²⁵ The GPC curves of the nitration/reduction OAPS

prepared using three different temperatures for the reduction of ONPS (Figure 31) were accompanied by shoulders at 50°C and 60°C probably representing some oligomeric species (10-20%). On the other hand, the oligomerization did not appear at 40 °C. It was suggested that a lower temperature could thus result in a more uniform product. However, the authors did not provide the corresponding molecular weights.

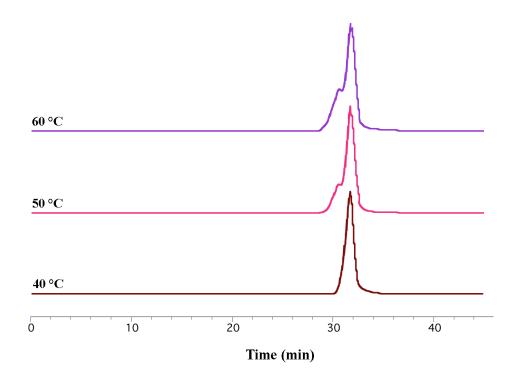
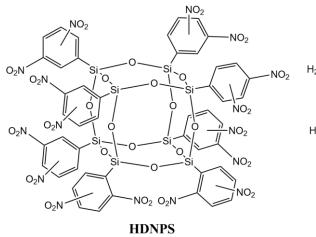
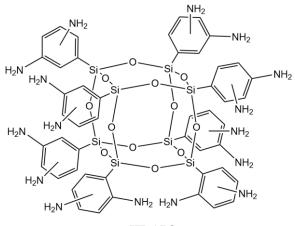


Figure 31. GPC curves of OAPS produced by the reduction of ONPS at three different temperatures.¹¹⁹

It should be noted that there are other phenyl-derived silsesquioxanes than OPS, ONPS and OAPS. These include for example the structures shown in Figure 32, i.e. hexadeca(aminophenyl)silsesquioxane (HDAPS) and its hexadeca(nitrophenyl)silsesquioxane (HDNPS) well precursor as as even complex structures such more as octa(maleimidophenyl)silsesquioxane (OMIPS).^{101, 111, 113, 120} Nagendiran et al.¹²⁶ synthesized crosslinked polyaspartimides by polymerizing OMIPS with three different diamines. However, all these structures introduce a very high degree of complexity and they are difficult to get commercially. As such, it was decided to keep OAPS as the most interesting candidate.





HDAPS

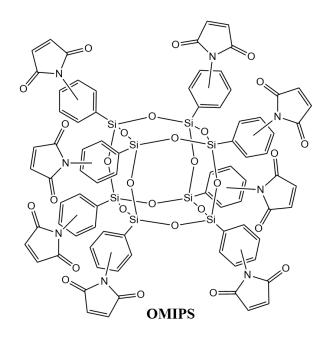


Figure 32. Other examples of phenyl-derived silsesquioxanes.

IV. Characterization of a commercial and a controlled OAPS

The experimental part of the research was mostly performed in the *Films in Fluids* research group directed by Prof. Nieck Benes at the University of Twente. After selecting OAPS as a suitable candidate precursor to make new hybrid polyOAPS-imide membranes, it was first purchased from Gelest Inc.¹⁰⁹ Based on the information presented by Gelest, this "commercial OAPS", which came as a powder, was synthesized using the Laine nitration/reduction route. In agreement with Choi et al.¹¹² and Kim et al.,¹¹⁹ Gelest¹⁰⁹ claimed that their product is a mixture of *meta:ortho:para* with a 60:30:10 ratio and typically contains 20% of dimer.

Later in the work, a (60%*para*/40%*meta*) OAPS was synthesized using the Lee aminophenyltrimethoxysilane route. As mentioned before, this synthesis was carried out by SINTEF, a Norwegian partner of the University of Twente. They used an inexpensive mixture of the *meta* (30-40%) and *para* (60-70%) APTMS isomers also provided by Gelest as the initial reactants. This "controlled OAPS" was dissolved in DMSO before being provided to the University of Twente.

The main structural differences between the so-called commercial and controlled OAPS can be summarized as follows:

- the commercial OAPS should contain all three isomers while the controlled OAPS does not have any *ortho* isomer.

- the controlled OAPS contains a lot more of the *para* isomer than the commercial OAPS.

- the commercial OAPS is based solely on eight-cornered Si siloxane cages, while there could be larger siloxane cage structures too in the controlled OAPS.

Both types of OAPS were characterized experimentally using picnometry as well as IR and NMR spectroscopies. In addition, IR and NMR predictions were performed by different softwares (machine-learning approaches such as MestreNova and quantum mechanical approaches such as DFT with Gaussian) in order to better understand the experimental results.¹²⁷⁻¹²⁸

1) Density of the commercial OAPS

Figure 33 shows (a) the dark yellow commercial OAPS powder from Gelest¹⁰⁹ and (b) the controlled OAPS solution in DMSO. The latter can be obtained as a powder following the evaporation of the solvent.

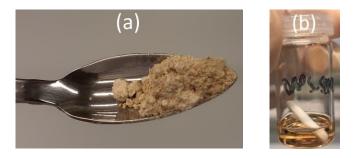


Figure 33. The OAPS samples: (a) commercial powder and (b) controlled solution.

Since the OAPS are in the form of powders, the AccuPyc II 1340 pycnometer (Figure 34) was used to measure their densities. However, the technique first had to be tested on NaCl and on a spherical metallic standard. As can be observed in Figure 34, an inert gas (such as helium or nitrogen) flows into a chamber of known volume (valve "a" opens then closes), and the sample is sealed in the instrument chamber. Then the gas expands into another precise internal volume (valve "b" opens). The pressure drops upon filling the sample chamber and then discharging it into the second empty compartment allows for the measure of the sample volume. The gas molecules rapidly fill the pores and only the solid phase of the sample displaces the introduced gas. Dividing the measured volume by the weight of the sample gives the gas displacement density. The pressure then vents off by opening valve "c".¹²⁹



Figure 34. A photo (left) and a schematic diagram (right) of the AccuPyc II 1340 pycnometer; a, b and c are the valves.¹²⁹

This technique is non-destructive as it uses the gas displacement method to measure the volume. Furthermore, it generates much more precise and reproducible densities than the traditional Archimedes water displacement method,¹²⁹ since the small gas molecules can diffuse in the pores, cavities, and any open space. The reproducibility (precision) is guaranteed within +/- 0.02% of the nominal full-scale cell chamber volume.⁸⁹ Since the full-scale cell chamber volume of the AccuPyc II 1340 pycnometer used is 74.051 cm³, the nominally guaranteed precision should be ± 0.015 cm³.

The pycnometer calculates the volume of the sample by using Equation 6:⁸⁹

$$V_{sample} = V_{cell} - \frac{V_{expansion}}{\frac{P_1}{P_2} - 1}$$
(6)

where P_1 is the gauge pressure after filling by the inert gas (helium) and P_2 is the gauge pressure after expansion. V_{cell} and $V_{expansion}$ are already known via a calibration procedure, whereas P_1 and P_2 are measured by the gauge included inside the pycnometer. The weight of the sample should be measured by a precise balance and inserted into the memory of the pycnometer.

Before beginning each test, the pycnometer performs purges to sweep the compartments. The default number of the purges and P_1 were set to 10 and 1.5 bar respectively. Although P_1 fluctuates cycle-by-cycle, it is recorded at each cycle to calculate the volume in the most precise way. The number of cycles for each analysis was set to 30, with the steady-state being reached for each cycle. The pycnometer calculated the average volume and density by averaging 30 measurements for each test. The average temperature was also obtained by averaging 30 recorded temperatures for the 30 measurements. The standard deviations were measured for the volumes, densities and temperatures.

To test the accuracy of the pycnometer, a NaCl sample with a specific density of 2,17 g/cm^3 was first analyzed. Although the purge process sweeps the compartments, the water absorption for a material such as NaCl should also be considered. It was thus put in a vacuum oven to ensure that there was no water absorbed. The pycnometer applied 10 purges before starting each test to remove the moisture probably absorbed by the NaCl sample during its displacement from the oven into the cell (sample holder) of the pycnometer. The volumes and

densities of NaCl were found to fluctuate with the changes in ambient temperature during the day. Figure 35(a) shows the changes in volume and temperature as a function of the test number. Figure 35(b) shows the inverse correlation between temperature and density.

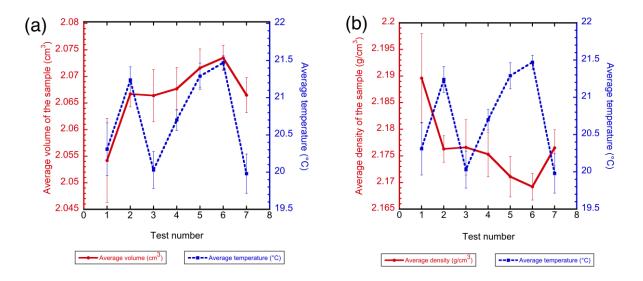


Figure 35. The average temperature along with (a) the average volume and (b) the average density of NaCl as a function of the test number.

The differences between the highest and lowest volumes and densities are 0.0193 cm³ and 0.0204 g/cm³ respectively, while the difference between the temperatures is 1.16° C. The volume and density of NaCl change by 0.0166 cm³ and 0.0176 g/cm³ per °C, respectively. Obviously, the density is inversely proportional to the temperature.

The compartments and internal components of the pycnometer are made of a metal alloy, so the fluctuations in temperature can also affect the volumes of the compartments (Equation 6), in addition to the sample volume. To study the effect of each factor, two different experiments were performed. Initially, the empty cell was tested and then a spherical metallic standard in steel with an accurately known volume (accessory of the pycnometer) was placed inside the cell to compare its nominal volume with the measured volume.

Figure 36 illustrates the impact of temperature on the volumes measured by the pycnometer when using a blank cell. Ideally, all measured volumes (V_{sample}) should be zero. But the V_{sample} values are positive, even though they remain in the ± 0.015 cm³ guaranteed error range. This is because the V_{cell} parameter is recorded with the last calibration and remains constant until the next calibration. The calibration process is done with standard

metal samples. All measured volumes are positive, and the average temperature at which the calibration was performed is probably different from the temperature at which the volume of the blank cell is measured. Indeed, the calibration process of the pycnometer takes almost 2 hours (the same as a complete test), but the average temperature of the calibration over the 2 hours cannot be obtained and only the temperature of the final cycle is presented. On the other hand, in a complete test, the pycnometer averages 30 recorded temperatures for the 30 measurements. In the present case, the recorded temperature for the calibration process was 21.2°C. According to Figure 36, the difference between the highest and lowest volumes is 0.0084 cm³ while the difference between the temperatures is 1.162°C. The volume of the sample in the blank cell changes by 0.0072 cm³ per °C. This was 0.0166 cm³/°C for NaCl, which is more than twice the volume variation of the blank cell. It can be a result of the cooperative impact of NaCl and the cell volume variations resulting from the changes in the ambient temperature.

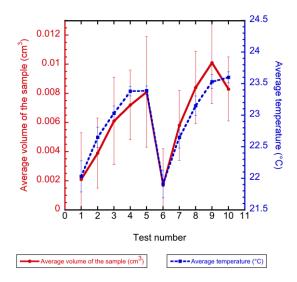


Figure 36. The average volume of the sample and the average temperature of the blank cell as a function of the test number.

To further investigate, a spherical metallic standard with a nominal volume of 6.3714 cm³ and a mass of 95.1566 g was tested. Figure 37 shows how the volumes and densities change with respect to the average temperatures. As before, there is a direct correlation between volume and temperature (Figure 37(a)) while Figure 37(b) shows the inverse correlation with the density. The differences between the highest and lowest volumes and densities are 0.0097 cm³ and 0.0227 g/cm³, respectively, while the difference between the temperatures is 1.99°C. The volume and density of the standard change by 0.0049 cm³ and

0.0114 g/cm³ per °C, respectively. The volume change is less than that of NaCl but the limitation of the precision should also be considered.

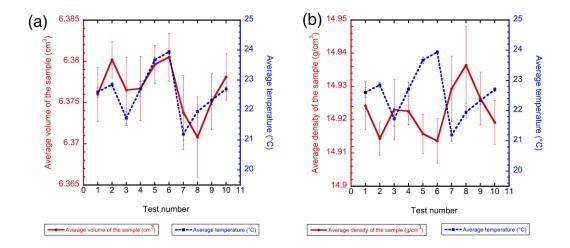


Figure 37. The average temperature along with (a) the average volume and (b) the average density of the spherical metallic standard, as a function of the test number.

Following these preliminary tests, 0.9087 g of the commercial OAPS was put in a vacuum oven to evaporate any moisture. It was then immediately weighed and transferred rapidly into the pycnometer to limit as much as possible the absorption of humidity. The same number of purges (10) and cycles (30) were set for each test. Figures 38(a) and 38(b) show the volumes and densities measured for that commercial OAPS. The differences between the highest and lowest volumes and densities are 0.0084 cm³ and 0.0179 g/cm³, respectively, while the difference between the temperatures is 2.58°C. The volume and density of the commercial OAPS change by 0.0032 cm³ and 0.0069 g/cm³ per °C, respectively. The density of the commercial OAPS sample averaged over the ten tests and was found to be 1.39 g/cm³ with a standard deviation of 0.01 g/cm³.

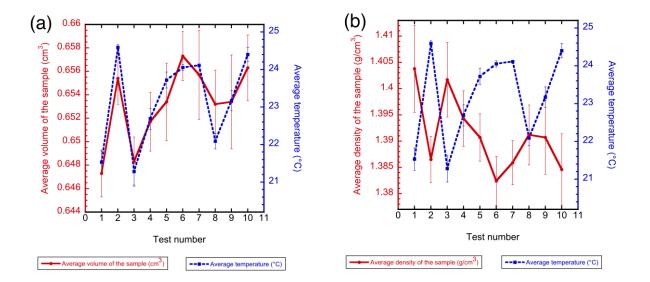


Figure 38. The average temperature along with (a) the average volume and (b) the average density of the commercial OAPS, as a function of the test number.

According to the results presented above for NaCl, the blank cell, the sphere and the commercial OAPS, it is clear that the temperature affects the measurements. To remove the impact of this factor, the pycnometer does a pressure calibration before each cycle. Thirty pressure calibrations are automatically performed with each test, but during each cycle (2 to 4 minutes), the ambient pressure is assumed to remain constant. As mentioned before, the nominally guaranteed precision is ± 0.015 cm³ and it is compatible with the error bars presented in the graphs. But it would be better to actually maintain the temperature constant during the calibrations and the tests. Micromeritics Instrument Corpcompany⁸⁹ proposes a thermoelectric device for the accurate control of the temperature from 15–36 °C (± 0.1 °C), which is adjustable in 0.1 °C increments. This could be an option for use with temperature-sensitive samples when the ambient temperature cannot be adequately controlled.

Following the evaporation of its solvent, the density of the controlled OAPS was measured with the pycnometer in the same way as for the commercial OAPS. It was found to be 1.37 g/cm^3 with a standard deviation of 0.01 g/cm^3 . Both values are reasonably close and thus provide confidence in the results presented above.

2) IR spectroscopy of OAPS

As noted before, almost all inorganic or organic compounds containing covalent bonds absorb frequencies in the infrared region of the electromagnetic spectrum. Its wavelengths l are shorter than those associated with microwaves but longer than those associated with visible light.¹¹⁶ Table 9 shows the energy transitions in the different regions of the electromagnetic spectrum. The vibrational transitions of the infrared region are especially interesting for chemical purposes and are referred to in terms of wavenumbers \bar{v} expressed in cm⁻¹ by taking the reciprocal of the wavelengths. IR wavenumbers typically cover the 400 to 4000 cm⁻¹ range.¹¹⁶

Region of Spectrum	Energy Transitions
X-rays	Bond breaking
Ultraviolet/visible	Electronic
Infrared	Vibrational
Microwave	Rotational
Radio frequencies	Nuclear spin (nuclear magnetic resonance)
	Electron spin (electron spin resonance)

In order to further compare the commercial and controlled OAPS samples, experimental IR spectra were obtained for both of them at the University of Twente using a Perkin Elmer ATR-IR spectrometer. Only a small amount of OAPS powder was used directly and the samples needed no specific preparation process. As for the density, since the controlled OAPS was initially dissolved in DMSO, its solvent was first evaporated by heating before performing the IR analysis.

In parallel, Density Functional Theory (DFT)¹³⁰ calculations were carried out on separate model structures of the pure *para*, *meta* and *ortho* isomers of OAPS (Figure 17) with the Gaussian software.¹²⁸ The Avogadro software¹³¹⁻¹³² was used to analyze the Gaussian outputs as described in Annexe 1. There are several quantum mechanical methods available to calculate the vibrational frequencies starting from model molecular structures. The *ab initio* methods, such as Hartree-Fock (HF), which investigate all the details of the electron movements are very computationally expensive. However it was shown that many properties of a molecule can be derived from its electron density distribution only,¹³³⁻¹³⁴ and this gave

rise to DFT approaches. Depending on the level of theory used, this allows for lower computational requirements than the *ab initio* methods.¹³⁴

The basis of quantum mechanics is the Schrödinger Equation.¹³⁴ Due to complex interactions between the electrons, an analytical solution for the Schrödinger equation is only possible for a few nuclei and electrons. Thus, it needs to be solved numerically by different approximations.¹³⁴ HF is an approximation named after its developers D. R. Hartree and W. A. Fock. It is based on the so-called Slater-determinant, which is an asymmetrized product of n one-electron wave functions, instead of using one wave function depending on nelectrons.¹³⁴ On the other hand, as mentioned above, DFT exploits the properties of the target molecule based on its electron density. There are three categories of Density Functional methods: (i) Local Density Approximation (LDA), which assumes the uniformity of the electron density throughout the molecule, ¹³⁵⁻¹³⁶ (ii) Gradient Corrected (GC) methods which account for the non-uniformity of the electron density,¹³⁷ (iii) Hybrid methods which use interesting features of the *ab initio* methods, particularly the Hartree-Fock theory, to improve the functionals used in DFT.¹³⁸⁻¹⁴⁰ Electronic correlation means that the electrons cannot be considered "independent" in a many-electron system such as a molecule. They are surrounded by an effective potential that accounts for all interactions with the other electrons, and which influences the excited-state features.¹⁴¹ In addition to the correlations, there are exchanges, which means that, due to their quantum nature, different electrons cannot be discerned from each other when their wave functions overlap. This has a direct consequence on the symmetry of the many-body wavefunctions and the Hartree-Fock method is able to discern correlation effects from exchanges.

The B3LYP (Becke, 3-parameter, Lee-Yang-Parr) level of theory is a hybrid exchange-correlation functional that linearly combines the Hartree-Fock exact exchange functional and the GGA (Generalized Gradient Approximation) corrections with the LDA electron-electron and electron-nuclei energies. GGA uses the density gradient to correct for the deviation from uniform densities.¹⁴⁰

A very important factor that must be considered in the DFT calculations is the basis set for expanding the molecular orbitals.¹⁴⁰ A basis set is a set of functions (called the basis functions) used to provide the electronic wave functions in the Hartree-Fock method or DFT theory in order to turn the partial differential equations of the model into algebraic equations suitable for efficient implementation on a computer. A large basis set is expected to present

accurate results but also requires a large computational cost. On the other hand, a small basis set is computationally efficient but can produce inaccurate results. It is therefore desirable to have the accuracy and at the same time use basis sets that are as compact as possible.¹⁴⁰

The minimal basis set is the minimum number of basis functions required to describe the ground state of atoms in a molecule. For example, the minimal basis for O $(1s^2, 2s^2, 2p^4)$ contains five basis functions:1s, 2s, 2p_x, 2p_y, 2p_z. In addition to the minimum number of basis functions, extra basis functions can be added to define the expanded basis sets. There are various kinds of extended basis sets including multiple-zeta (double, triple, etc), splitvalence, polarization and/or diffuse functions:¹⁴⁰

- In double-zeta basis sets, each orbital in the minimal basis set is replaced by two basis functions, for instance, the double-zeta basis set for O contains ten functions instead of five. The same principle applies to other multiples.
- Split-valence basis sets simplify the computational efforts by using the double (or more)-zeta description only for the valence orbitals, while the inner-shell electrons are described by the minimal basis assumption.
- Polarization functions are shown between two parentheses and are characterized by an orbital quantum number *l* larger than the largest *l* in the minimal basis set approximation, *i.e.* they provide more room space for the electrons to get away from each other and minimize electron-electron repulsions. As an example, for O: $1s^2$, $2s^2$, $2p^4$, the "d, f ..." orbitals (l = 2, 3...) are polarization functions, whereas for H: $1s^1$, the "p" (l = 1) orbitals are polarization functions.
- Diffuse functions such as the very shallow Gaussian basis functions more accurately display the "tail" parts of the atomic orbitals, which are far from the atomic nuclei.¹⁴² Diffuse functions are indicated by + and can be important for describing anions, large soft molecular structures and long-range interactions such as Van der Waals forces, as well as for accomplishing electronic excited-state and electric field property calculations.^{128, 140}

In the present work, we use split-valence basis sets. Their notation comes from the group of John Pople and is typically expressed as X-YZWG.¹⁴³⁻¹⁴⁴ Here, X shows the number of primitive Gaussians including each core atomic orbital basis function. The Y, Z and W indicate that the valence orbitals consist of three basis functions each, the first one

comprising a linear combination of Y primitive Gaussian functions, the second one comprising Z primitive Gaussian functions and the third one comprising W primitive Gaussian functions. G is an initial for Gaussian. For example, 6-311G means that the inner shell atomic orbital (1s orbital) is a linear combination of six primitives and each valence shell atomic orbital is split into three parts (triple zeta) using three, one and one primitive Gaussians respectively.^{140, 144-145} As explained above, "G" means Gaussian.¹²⁸ Another example, which is less expensive but still accurate enough in some cases, is the 6-31G(d,p) basis set, which means that, as in the previous example, the inner shell atomic orbital (1s orbital) is a linear combination of six primitive Gaussians, but in that case, each valence shell atomic orbital is split into an inner and outer part (double zeta) using three and one primitive Gaussians, respectively. In addition, (d,p) polarization functions are used for heavy atoms and hydrogens respectively.¹⁴⁴ It was chosen here to perform the geometry optimizations and predict the vibrational frequencies for the three isomers of OAPS along with B3LYP, which is the most widely-used functional due to its superior quality in the energy evaluation of small molecules.¹⁴² B3LYP also reproduces very well the geometries of molecules.¹⁴⁶⁻¹⁴⁹ Hence, all calculations in this Section were carried out using B3LYP as the hybrid density functional with the polarization functions added to a double zeta basis set 6-31G(d,p).¹⁴⁷⁻¹⁵¹

Since the IR spectra arise from changes in dipole moments during the analyses, different isomers with different amounts of symmetry may present different spectra, i.e. the change in the dipole moment may diminish in a more symmetrical structure such as the *para* isomer. Indeed, Figure 39 shows that there are differences between the experimental bands of (a) the commercial OAPS and those of (b) the controlled OAPS, whereas they only differ by the nature and proportions of the isomers. Interestingly, there are no signs of large peaks around 3500 cm⁻¹, which would be characteristic of -OH groups and thus of open siloxane cages. This suggests that most or all of the cages are intact.

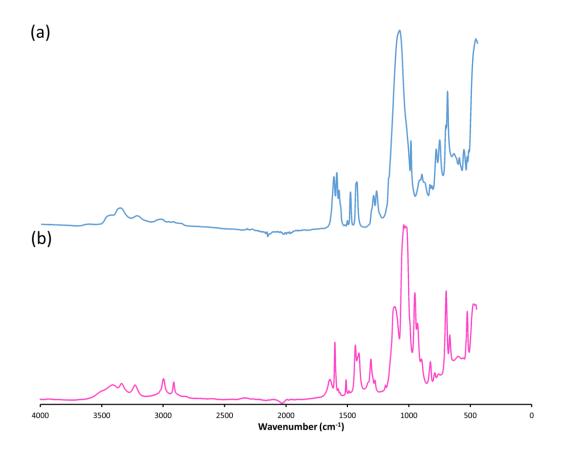


Figure 39. A comparison of the experimental IR spectrum for (a) the commercial and (b) the controlled OAPS.

Figure 40 compares separately the predicted IR spectra for each of the three isomers, i.e. in (a) *meta*, in (b) *ortho* and in (c) *para*, with the experimental IR spectra of the commercial OAPS (Figure 39(a)). The different vibrational modes were identified using the Avogadro software (Annexe 1). The predicted spectra are obviously restricted to specific optimized structures while the experimental spectrum takes into account a very large number of possible structures, hence its much broader peaks. However, Figure 40(a) for the *meta* isomer still indicates a good compatibility between the experimental and the predicted spectra for wavenumbers below 1700 cm⁻¹. The predicted bands for the *ortho* isomer (Figure 40(b)) are broader, especially for the Si-O-Si stretching vibrations. Figure 40(c) for the *para* isomer indicates lower intensities for most of the predicted bands due to its more symmetrical structure. This may prove that there are indeed higher proportions of the *ortho* and *meta* isomers in the commercial OAPS. All three predicted spectra show slightly overestimated bands for wavenumbers beyond 3000 cm⁻¹, but such overestimations for the N-H vibrations have already been observed in other reports,¹⁵²⁻¹⁵⁶ and this is a part of the spectrum which has very low intensities.

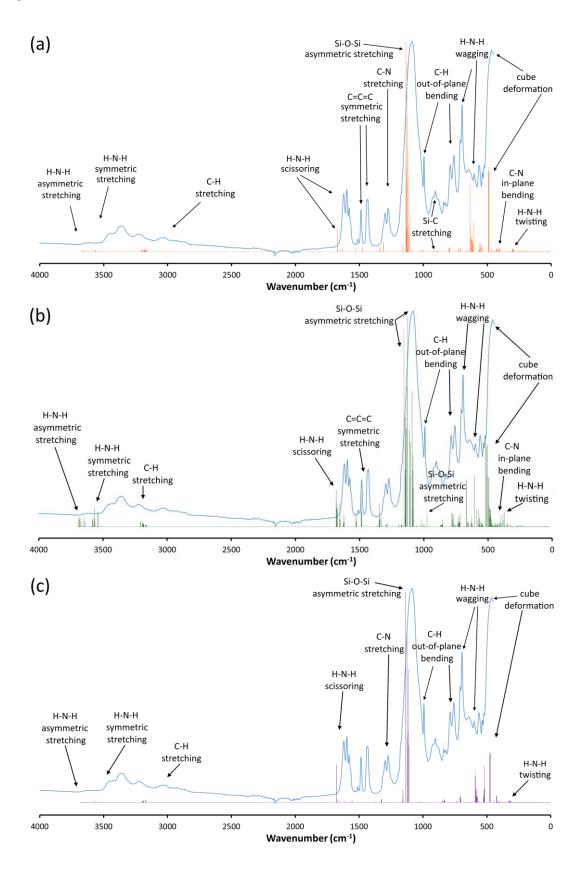


Figure 40. A comparison of the experimental IR spectrum for the commercial OAPS (blue) with the predicted IR spectra for model (a) *meta* (orange), (b) *ortho* (green), and (c) *para* (purple) pure isomers of OAPS obtained from DFT calculations at the B3LYP/6-31G(d,p) level.

Figure 41 compares the experimental spectra of both (a) the commercial and (b) the controlled OAPS with the three predicted IR spectra assembled. A lot of information can be obtained:

- As shown in Figure 40, the compatibility between the experimental and the three predicted spectra assembled is rather good for the commercial OAPS, which seems to agree with the presence of the *ortho* isomer, despite what has been reported by Tamaki et al.¹¹⁸

- A sharp peak of H-N-H wagging at 520 cm⁻¹ (yellow rectangle) is observed for the controlled OAPS while there is only a small trace of it for the commercial one. Unlike for the other isomers, this rather strong peak was predicted for the *para*-OAPS, which could confirm its high proportion in the controlled OAPS.

- Two strong peaks for the out-of-plane vibrations of C-H appear at 756 cm⁻¹ and 784 cm⁻¹ (red rectangle) for the commercial OAPS, whereas they almost disappear for the controlled OAPS. Both of them correspond to predicted bands for the *ortho* isomer.

- The strong peak at 1115 cm⁻¹ (dark blue rectangle) in the commercial OAPS illustrates the Si-O-Si stretching vibrations which are common to the three isomers because of the silicon oxide cages.¹¹⁷ In the controlled OAPS, this broad peak is split into two peaks (light blue rectangles), with only one overlapping with the predicted peaks. It is not easy to assign these peaks to the specific *para* and *meta* isomers but it is clear that the *ortho* isomer leads to broader predicted bands in this region.

- The sharp peak at 1485 cm⁻¹ for the commercial OAPS (brown rectangle) corresponds to a predicted band for the *ortho*-OAPS and does not appear in the controlled OAPS.

- The predicted bands for the *ortho*-OAPS from 1280 cm⁻¹ to 1690 cm⁻¹ are stronger than for the other two isomers (see Figure 40(b)). This agrees with the experimental spectrum of the commercial OAPS showing stronger peaks in the same region than the controlled one.

- Other peaks could be specific to one isomer, although many of them overlap in the experimental spectra. For example, the peak appearing at 670 cm⁻¹ (pink rectangle) in the commercial OAPS seems to correspond to the *ortho* isomer. The peak at 710 cm⁻¹ (green rectangle), which is sharp in the controlled OAPS matches with a predicted peak for the *para*-OAPS.

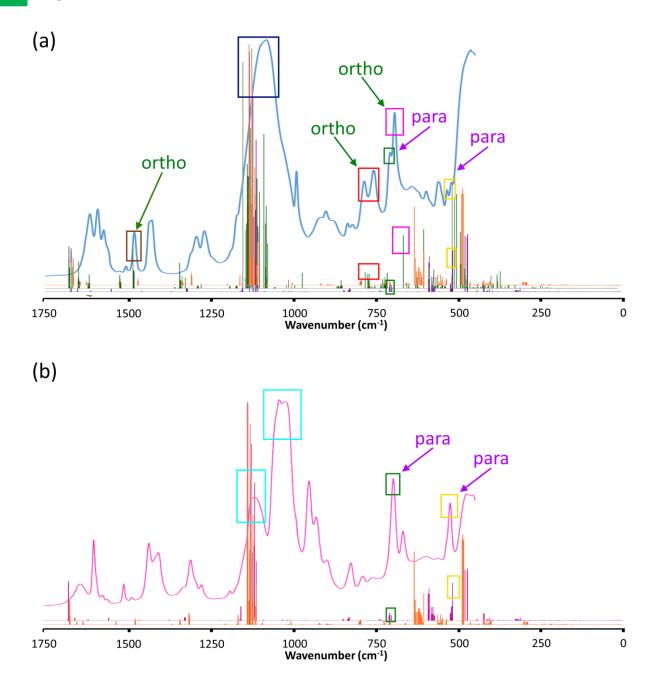


Figure 41. A comparison of (a) the commercial (blue) and (b) the controlled (pink) OAPS experimental spectra with the predicted IR spectra for the *meta* (orange), *ortho* (green), and *para* (purple) OAPS isomers up to 1750 cm⁻¹.

These compared analyses seem to confirm the different natures of both OAPS samples in terms of isomers and show how DFT predictions on a limited number of model structures can be a useful tool to help the interpretation of the experimental spectra.

3) NMR spectroscopy of OAPS

a) The experimental NMR spectra

The NMR spectrum is associated with the emission of electromagnetic radiation by a nucleus at a given frequency of radiation. The ²⁹Si-NMR, ¹³C-NMR, and ¹H-NMR spectra of the commercial and controlled OAPS were recorded at the University of Twente using the Bruker AVANCE III 600 MHz NMR spectrometer. DMSO-d6 was used as the solvent of the samples.

Figure 42 shows the ²⁹Si-NMR spectra of the commercial (blue line) and controlled (pink line) OAPS compared to the OAPS purchased by Lee et al. (brown line).¹⁰⁷ As for our commercial OAPS, Lee's purchased OAPS was obtained by nitration/reduction of OPS. A good compatibility can be observed between the three spectra in Figure 42. However, there are differences between the controlled OAPS spectrum in (b) and both OAPS spectra obtained from nitration/reduction in (a)-(c). Indeed, the intensities of the main shifts at -77.8 ppm and -79.4 ppm (violet and orange rectangles) in the controlled OAPS are in reverse order compared to those of the nitration/reduction OAPS. The increase in intensity at -77.8 ppm in Figure 42(b) can be attributed to the higher portion of the *para* isomer while the decrease at -79.4 ppm is probably due to the lower portion of the *meta* isomer in the controlled OAPS. On the other hand, there are no clearly defined signatures of the *ortho* isomer in the ²⁹Si-NMR spectra, whereas it should be present in the nitration/reduction OAPS samples.

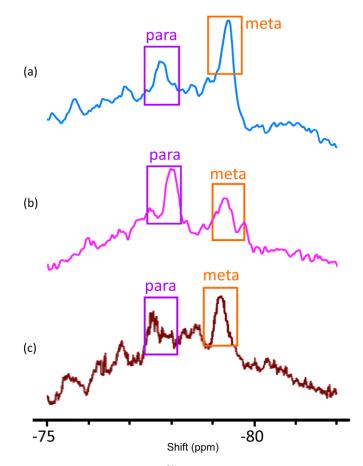


Figure 42. A comparison of the experimental ²⁹Si-NMR spectra for (a) the commercial OAPS, (b) the controlled OAPS, and (c) the OAPS purchased by Lee et al.,¹⁰⁷ all of them in DMSO-d6, along with a proposed isomer assignment for the main peaks.

Similarly, Figure 43 shows the good compatibility of the ¹³C-NMR spectrum for the commercial OAPS (light blue line) with that of OAPS produced through nitration/reduction by Krishnan et al. (dark blue line).¹⁵⁷ There are at least two peaks at 132 ppm and 154 ppm in the nitration/reduction samples (green rectangles in (a)-(c)), that do not have any equivalent peaks in the (b) controlled OAPS spectrum. These two peaks should again be related to the presence of the *ortho* isomer. The controlled OAPS spectrum also shows weaker peaks in comparison to the other two spectra. Surprisingly, there is a peak in (b) (red rectangle) that shows such a small trace that it cannot be easily revealed in the nitration/reduction samples. That may be why Tamaki et al.¹⁰⁶ reported only ten observable peaks for their ¹³C-NMR spectrum and concluded that this reflected the presence of the *meta* and *para* isomers only. However, as explained in this work, there are definitely indications of the presence of the *ortho* isomer too in the nitration/reduction samples.

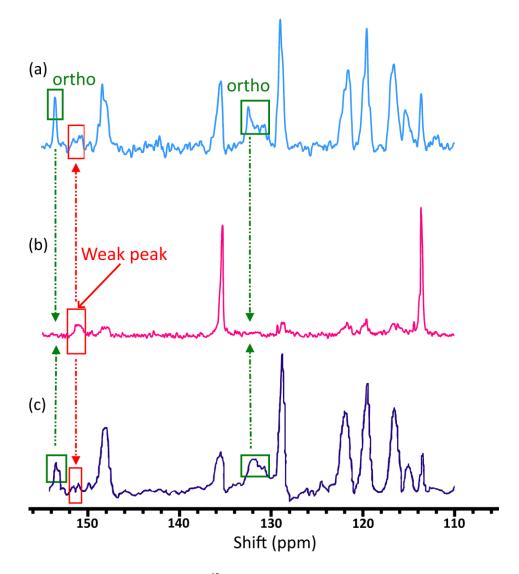


Figure 43. A comparison of the experimental ¹³C-NMR spectra for (a) the commercial OAPS, (b) the controlled OAPS and (c) the OAPS prepared by Krishnan et al.,¹⁵⁷ all of them in DMSO-d6, along with a proposed assignment for *ortho*-specific peaks.

The ¹H-NMR spectra of the commercial and controlled OAPS are compared to that obtained by nitration/reduction by Zhang et al.¹¹³ in Figure 44. Both commercial (light blue line) and Zhang (dark blue line) OAPS spectra are again quite similar but different from the controlled OAPS spectrum (pink line). They are also compatible with the ¹H-NMR spectra published by Tamaki et al.¹⁰⁶ and Huang et al.¹¹⁷ Two broad regions can be distinguished: one from 4.2 ppm to 5.4 ppm, which corresponds to the amine hydrogens, and the other one from 6.2 ppm to 7.8 ppm, which is related to the four aromatic hydrogens of the OAPS phenyl arms. The latter appears in a slightly more narrow range, i.e. 6 ppm to 7.3 ppm, for the controlled OAPS.

The main differences between the three spectra are:

- The yellow rectangles at 5.3 ppm, red rectangles at 6.6 ppm and green rectangles at 7.3 ppm show peaks that do not appear in the (b) controlled OAPS spectrum and can thus be linked to the *ortho* isomer.

- The orange rectangles at 5.15 ppm indicate small peaks in both (a) the commercial and (c) Zhang OAPS, which become stronger for the (b) controlled OAPS with a higher proportion of *para* isomer. Similarly, the sharp peak at 6.25 ppm in the controlled OAPS (brown rectangle) is probably also related to the aromatic *para* hydrogens.

As mentioned before (Figure 25), the number of amine functional groups on each arm or ring connected to the core can be obtained by dividing the integration of the peaks. This ratio is equal to 2 for all OAPS understudy, which means that there are two amine hydrogens along with four aromatic hydrogens, i.e. one -NH₂ per ring for all three samples. However, it is difficult to distinguish the aromatic hydrogens of the different OAPS isomers. This is the reason why Tamaki et al.¹⁰⁶ tried to use the ¹H-NMR spectrum of ONPS with its sharper peaks (Figure 23) to calculate the isomer ratio.

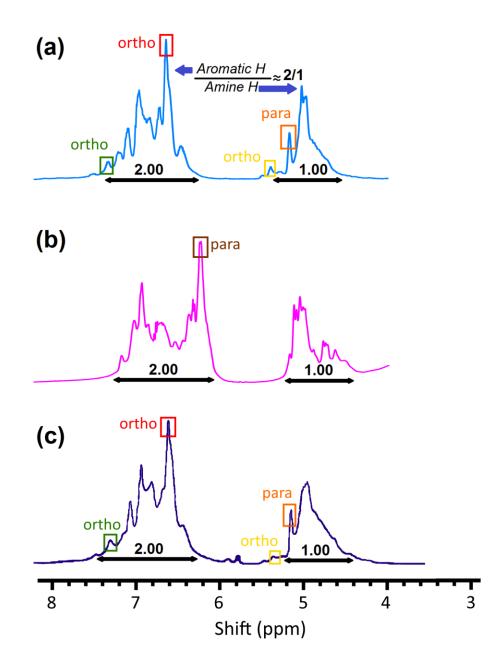


Figure 44. A comparison of the experimental ¹H-NMR spectra and surface areas of (a) the commercial OAPS,
(b) the controlled OAPS, and (c) the OAPS prepared by Zhang et al.,¹¹³ all of them in DMSO-d6, along with a proposed isomer assignment for the main peaks.

In addition to conventional one-dimensional NMR (1D-NMR) methods, twodimensional NMR spectroscopy (2D-NMR), which has been developed since the 1980s, analyses the samples with two frequency axes rather than with one.¹⁵⁸⁻¹⁶⁰ It is aimed at facilitating the interpretation of NMR spectra for complex molecules.¹⁶⁰ 2D-NMR techniques can be divided into three main categories, each of which is based on several techniques (Table 10).¹⁶¹ In homonuclear correlation spectroscopy, both axes of a correlation spectrum display shifts for a single isotope, commonly hydrogen (¹H). Heteronuclear correlation spectroscopy (HETCOR) is linked to two different isotopes such as carbons and hydrogens. Through-space correlation spectroscopy studies nuclei that are physically close to each other, irrespective of the presence of a bond between them.

Homonuclear 2D-NMR	Heteronuclear 2D-NMR	Through-space correlation 2D-NMR
Correlation spectroscopy (COSY) Exclusive correlation spectroscopy (ECOSY)	Heteronuclear single-quantum correlation spectroscopy (HSQC)	Nuclear Overhauser effect spectroscopy (NOESY)
Total correlation spectroscopy (TOCSY) Incredible natural-abundance double-quantum transfer experiment (INADEQUATE)	Heteronuclear multiple-bond correlation spectroscopy (HMBC)	Rotating-frame nuclear Overhauser effect spectroscopy (ROESY)

 Table 10. 2D-NMR techniques.¹⁶¹

In a heteronuclear single-quantum correlation (HSQC) spectrum, the 1D-NMR of one nucleus is plotted versus the 1D-NMR of another type of nucleus.¹⁶¹ In our particular case, the HSQC spectrum correlates the carbon nuclei with the hydrogens directly connected to them, i.e. one-bond couplings. The ¹³C spectrum is presented on the *y*-axis and the ¹H spectrum on the *x*-axis. To interpret such a spectrum, each ¹³C shift is selected on the *y*-axis and a line is drawn horizontally, i.e. parallel to the ¹H-NMR *x*-axis, until a cross-peak is encountered. Similarly, each ¹H shift is selected on the *x*-axis and a line is drawn vertically until a cross-peak is encountered. These cross-peaks identify which hydrogens are connected to which carbons and vice versa. When a line encounters no cross-peak, it implies that this carbon has no connected hydrogen.

Figure 45(a) shows the HSQC spectrum obtained for the commercial OAPS. Ten cross-peaks can be observed, indicating ten different carbons connected each to one hydrogen. It confirms that the ¹H-NMR shifts from 6.2 ppm to 7.8 ppm refer to the aromatic hydrogens. The three isomers of the commercial OAPS possess sixteen different environments for carbon atoms, ten of which are connected to the hydrogens and the six others to either silicon or nitrogen. The black rectangle shows the ¹H-NMR signals presenting no connections to the carbon atoms. These peaks between ~4.5 ppm to ~5.5 ppm are the amine hydrogens, as reported by Tamaki et al.¹⁰⁶ The brown rectangle indicates a ¹H-NMR shift specific of the *ortho* amine hydrogens (Figure 44). The orange and purple rectangles show the ¹³C-NMR signals displaying no connections to the hydrogens. The yellow

rectangles for ¹H and red rectangles for ¹³C give examples of overlapping peaks. The blue rectangles indicate points, which are probably related to incomplete or open cages as their intensities are low. The green rectangles indicate the ¹³C-NMR shifts related to the *ortho* isomer (Figure 43). Such an HSQC spectrum clearly confirms the presence of the three isomers and the report of Kim et al.¹¹⁹ is thus more accurate than that of Tamaki et al.¹⁰⁶

Figure 45(b) shows the HSQC spectrum of the controlled OAPS. Six cross-peaks corresponding to six different carbons connected each to one hydrogen can be observed, as expected from the presence of the *meta* and *para* isomers only. The four remaining carbons are those connected to either silicon or nitrogen. The purple rectangle shows again the ¹³C-NMR signals having no connections to the hydrogens. The ¹H-NMR peaks between ~4.5 ppm to ~5.5 ppm are the amine hydrogens as indicated by the black rectangle. The *ortho* amine shifts between 5.2 ppm and 5.5 ppm (yellow rectangles in Figure 44 and brown rectangle in Figure 45(a)) are absent from Figure 45(b). The pink rectangles display expanded cross-points, which could be related to the different sizes of the OAPS cages in the controlled sample. Indeed, different cages sizes possess different numbers of oxygens and as such, they probably have slightly different electron-withdrawing tendencies. The ¹³C-NMR peaks specific to the *ortho* isomer (green rectangles in Figure 45(a)) disappear in Figure 45(b), thus confirming its absence. Such an HSQC spectrum shows that there are indeed only two isomers out of the possible three in the controlled OAPS sample.

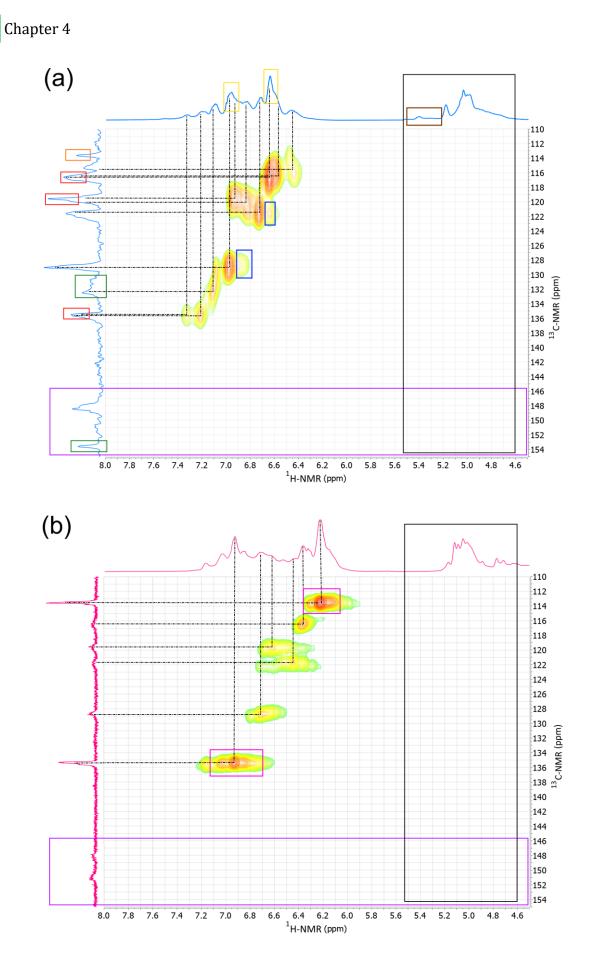


Figure 45. The HSQC spectra of (a) the commercial OAPS and (b) the controlled OAPS.

b) Predictions of the NMR spectra

b.1) Description of the methods

There are both fast and slow computational methods to predict NMR spectra. All fast NMR prediction methods use databases of assigned data. In a nutshell, these approaches can be summarized as:¹²⁷

- Increments methods (from tables of fragments)
- Machine Learning methods
- HOSE-code methods

The most common algorithms for NMR prediction use reference databases where a large number of spectra and their assigned 2D or 3D chemical structures are the source knowledge. The MestreNova software¹²⁷ uses both machine learning and the HOSE (Hierarchical Organization of Spherical Environments) code prediction, which was first described by Bremser.^{127, 162} The HOSE code starts at the carbon atom whose shift is to be predicted, looks one bond away from the carbon, and tries to find this environment in its database. If it is successful, it moves two bonds away and tries again and so on until it either comes across something which is not represented in the database or it reaches the boundary of the molecule (Figure 46). The HOSE code approach works very well for structures, which are represented in the reference collection.

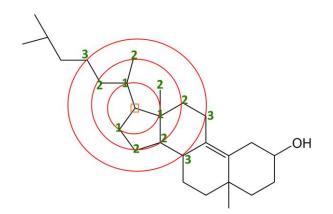


Figure 46. The prediction path for ¹³C-NMR by HOSE.¹²⁷

"NMRPredict" is the name of the prediction module in MestreNova. If the environment of an atom can be predicted to be three neighbouring atoms or more (NMRPredict goes to a maximum of five neighbouring atoms), the prediction can be

Chapter 4

considered very reliable. However, if the query structure is not well represented in the database and the atom environment can only be predicted to one or two neighbouring atoms, the prediction cannot be relied upon. To improve the situation, NMRPredict also uses a unique Neural Network algorithm,¹²⁷ which is more general and error-tolerant than the HOSE code and is more accurate at predicting shifts for atoms whose environments are not in its database. It should be noted that even a quality reference database of 10000 compounds and the best software may deliver poor prediction accuracy due to specific conformational features.¹²⁷ MestreNova includes a module for minimization of the input structures, which tries to find the optimized geometries according to specific parameters fed to the software by the developer (no details on the techniques and parameters used for minimization are provided). Another factor is the chemical environment. MestreNova calculates the NMR spectrum of a single-point structure at the moment, which is not the case for a molecule dissolved in a solvent. Although there is an option to select a common solvent in the software in order to carry out more accurate calculations, the prediction depends on the database and the machine learning feature ability. In the present case, DMSO was chosen as the solvent for the calculations as it was used for the experimental NMR spectrometry.

In addition to the MestreNova approach, it is also possible to predict an NMR spectrum with much slower quantum mechanical (QM) calculations such as the DFT calculations carried out by Gaussian (Annexe 2). As explained in the previous section, it did give good results at the B3LYP/6-31G(d,p) level for the IR spectra. However, in the case of the NMR spectra, results can be very dependent on the level of theory used. As such, different levels of theory with a solvent correction, i.e. here SCRF=(solvent=DMSO) and the GIAO formulation first had to be tested,¹⁶³⁻¹⁶⁵ in order to perform both the optimization and calculation of the NMR chemical shifts for the different OAPS isomers. The SCRF (Self-Consistent Reaction Field) keyword¹⁶⁶ accounts for the presence of the solvent (here DMSO) by locating the solute in a cavity within the solvent reaction field using the Polarizable Continuum Model (PCM).¹⁶⁷ This method, which represents the solute cavity through a set of overlapping spheres was initially introduced by Tomasi et al.¹⁶⁸⁻¹⁶⁹ and Pascual-Ahuir et al.,¹⁷⁰ and it has been further improved by Tomasi, Barone, Mennucci et al.¹⁷¹⁻¹⁷⁹ and other researchers.¹²⁸ The GIAO (Gauge-Including Atomic Orbital) formulation results in all matrix elements involving the basis functions being independent of the origin.¹¹⁸ Since Gaussian calculates the absolute shieldings, the same level of theory has to be used for tetramethylsilane (TMS) as the reference for the chemical shifts.

To do an in-depth research, different DFT functionals used in the literature for NMR predictions,¹⁸⁰⁻¹⁸³ *i.e.* B3LYP, B3PW91, and HCTH407 were used along with 6-311+G(2d,p) as the main basis set and PCM as the default SCRF method in order to predict the chemical shielding values. Here, the 2d term in the basis set adds two sets of d functions on the heavy atoms, and p adds one set of p functions on the hydrogens.¹⁸⁴ The molecules were previously all optimized at the B3LYP/6-31G(d,p) level as it is a good compromise between time and accurate energy minimisation.¹⁴⁹ As far as the functionals are concerned, the first one, B3LYP, has already been presented in the previous Section about the IR predictions. The second one, B3PW91, is a hybrid functional representing Becke's three parameters and combining Perdew and Wang's gradient-corrected exchange (published in 1991) with the correlation functions and Hartree-Fock exchanges.¹⁸⁵⁻¹⁸⁶ The third one, HCTH407, which is a GGA approach,¹⁸⁷ refers to Hamprecht, Cohen, Tozer, and Handy,¹⁸⁸ who fitted their parameters by using 407 atomic and molecular training systems. Contrary to the hybrid B3LYP and B3PW91, there are no HF functions for HCTH407.¹⁸⁸ In addition, the B3LYP/6-31G(d,p) level of theory used for the IR spectra was attempted for NMR prediction and its results will also be presented. These different combinations of DFT functionals and basis sets, which were tested for all three types of NMR spectra (²⁹Si, ¹³C and ¹H), are summarized in Table 11.

While they have been shown to be appropriate for ¹³C-NMR and ¹H-NMR, conventional DFT approaches are not always able to produce as good predictions for the less common ²⁹Si-NMR spectra.^{183, 189-198} As such, in that specific case, the B3PW91 and HCTH407 functionals were also tested with a more extended basis set, i.e. 6-311++G(2df,2pd). The 2df term adds two d along with one f function for the heavy atoms, and the 2pd term adds two p and one d functions for the hydrogens. In addition, even more powerful Hartree-Fock calculations were attempted with the main 6-311+G(2d,p) and the extended 6-311++G(2df,2pd) basis sets. In the HF model, each electron is described by an orbital and the total wave function is presented as a product of orbitals,¹⁹⁹ thus making it a lot more computationally expensive than DFT. These approaches are thus also mentioned in Table 11, but as noted above, the predictions with the extended 6-311++G(2df,2pd) basis set and the heavy HF approach were restricted to ²⁹Si-NMR. Indeed, it will be shown that for ¹³C-NMR and ¹H-NMR, a sufficient degree of accuracy can be obtained using more manageable (i.e. lower) basis sets and DFT methods.

	Guussian sortware.					
	B3LYP	B3PW91	HCTH407	HF		
6-31G(d,p)	*	-	-	-		
6-311+G(2d,p)	*	*	*	*		
6-311++G(2df,2pd)	-	*	*	*		

 Table 11. The different combinations of methods and basis sets used to calculate the chemical shifts using the Gaussian software.

The procedure to get chemical shifts from the calculations by Gaussian is detailed in Annexe 2. As noted above, Gaussian is in general very heavy in terms of calculations (for a specific OAPS isomer, DFT and HF calculations take typically 10 and 20 days respectively, to complete both optimization and NMR prediction), whereas MestreNova produces the results in a few minutes. As such, it is interesting to attempt and compare both approaches.

b.2.) Case study: the meta-OAPS isomer

Due to the highly time-consuming DFT calculations, a siloxane cage with one *meta* ring and seven hydrogens linked to the Si (Figure 47(a)) was first fed to both MestRenova and Gaussian in order to calculate the chemical shifts for the various combinations of Table 11. In a second stage, one method was discarded and the other calculations were repeated on an 8-ring *meta*-OAPS (Figure 47(b)) since it is expected that the presence of the rings connected to the silicons will somewhat displace the shifts. Note that, in the following Figures, the 8-ring *meta*-OAPS is simply shortened to *meta*-OAPS.

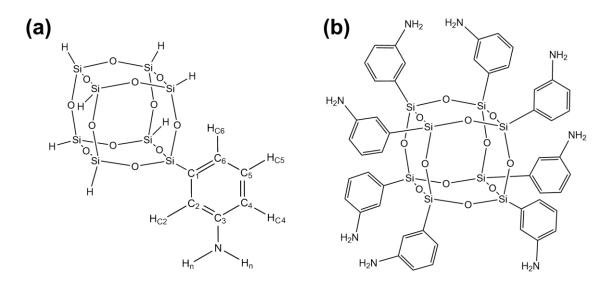


Figure 47. The molecular structures of a) the 1-ring *meta*-OAPS and b) the 8-ring *meta*-OAPS. The phenyl ring carbons are named C₁, C₂, C₃, C₄, C₅ and C₆. H_{C2}, H_{C4}, H_{C5} and H_{C6} are the aromatic hydrogens connected to C₂, C₄, C₅ and C₆ respectively. H_n indicates both amine hydrogens.

b.2.1.) The ²⁹Si-NMR predicted spectra for the 1-ring meta-OAPS

Figure 48 shows the ²⁹Si-NMR shifts predicted by Gaussian¹²⁸ using the combinations presented in Table 11 (dashed lines) and MestreNova¹²⁷ (line) for the 1-ring *meta*-OAPS along with two experimental spectra in the range from -50 to -90 ppm. In the predicted spectra, only the shift of the silicon connected to the aminophenyl ring was taken into account. The blue experimental spectrum corresponds as before to the commercial OAPS obtained by nitration/reduction (all three isomers present, same as in Figure 42) and displays clear chemical shifts at -79.4 ppm, -77.8 ppm, -76.9 ppm, and -75.7 ppm along with a smaller one at -69 ppm. The orange experimental spectrum is that of the pure *meta* isomer obtained from the condensation of APTMS precursors by Lee et al.¹⁰⁷ It shows a distinct peak at -79.1 ppm, which is close to the sharpest peak of the blue spectrum at -79.4 ppm and could indicate that this belongs to the *meta* isomer.

As far as MestreNova is concerned, the predicted chemical shift for the Si connected to the single *meta*-ring appears at -59.3 ppm which differs from the experimental spectra by 20 ppm. The shift predicted by Gaussian at the B3LYP/6-31G(d,p) level of theory is at -61.9 ppm, which is also very distant from the main experimental peaks around -79.4 ppm. On the other hand, the Gaussian predictions with the 6-311+G(2d,p) and 6-311++G(2df,2pd) basis sets are closer to the experiment. For the former, the Si shift is underestimated by ~10 ppm for B3LYP and HCTH407 and by ~3 ppm for B3PW91, while it is overestimated by ~2 ppm by HF. The situation is further improved for HF by using the more extended 6-311+G(2d,p) basis set since it predicts the Si shift at -80 ppm, i.e. very close to the sharpest experimental peak. This is not the case for both B3PW91 and HCTH407 whose predicted shifts with 6-311++G(2df,2pd) are even further from the experiment than when associated with the lower 6-311+G(2d,p) basis set. It clearly shows that both the methods and the basis sets significantly affect the final accuracy of the prediction.

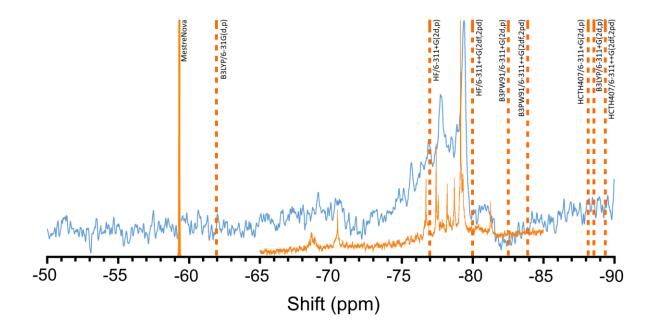


Figure 48. The experimental ²⁹Si-NMR spectra of the commercial OAPS (blue line) and the pure *meta* isomer of Lee et al. (orange line)¹⁰⁷ compared with the predicted chemical shifts by Gaussian (dashed lines) and MestreNova (lines) for the 1-ring *meta*-OAPS. Note that in the following NMR Figures, the same colour code and representations will systematically be used: Gaussian will be represented by dashed lines, MestreNova by lines, the experimental commercial OAPS by a blue line and the experimental pure *meta* and *para* spectra of Lee et al. by an orange and a purple line respectively.

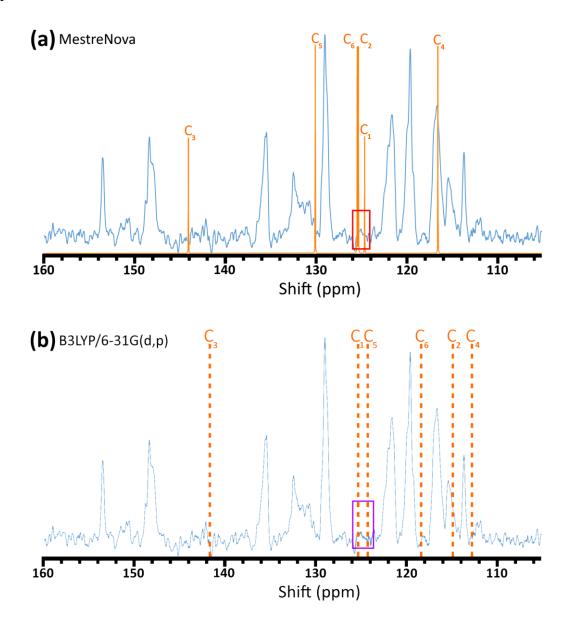
b.2.2.) The ¹³C-NMR predicted spectra for the 1-ring meta-OAPS

Figure 49 shows separately several predicted ¹³C-NMR spectra for the 1-ring *meta*isomer along with the experimental ¹³C-NMR spectrum for the commercial OAPS (blue line). Lee et al.¹⁰⁷ did not report any ¹³C-NMR spectra for their pure *meta*-OAPS and *para*-OAPS structures.

There again, the predicted peaks appear at different chemical shifts to the experimental ones, although they are in general much closer for the lower basis sets than for the ²⁹Si-NMR spectra. One overlap is predicted at ~125 ppm by MestreNova (red rectangle in (a)). In the Gaussian predictions, some shifts are found within less than 2 ppm and could also be possible overlaps. This is the case for B3LYP when associated with 6-31G(d,p) at ~125 ppm (purple rectangle in (b)) and when associated with 6-311+G(2d,p) at ~136 ppm (green rectangle in (c)). Both B3PW91 and HCTH407 indicate overlaps at 134.5 ppm (pink rectangle in (d)) and 127.9 ppm (brown rectangle in (e)), respectively. As noted before, the existence of overlaps (which do not allow to see the peaks separately) can explain why

Tamaki et al.¹⁰⁶ observed ten peaks and excluded the *ortho* isomer in their nitration/reduction OAPS.

Each predicted shift in Figure 49 can be assigned to one specific aromatic carbon according to the definition of Figure 47(a). The carbon shifts are in the order $C_4 < C_2 < C_6 < C_5 < C_1 < C_3$ for the Gaussian predictions. The order is similar for MestreNova, except for the C₁ shift which is predicted between C₄ and C₆ rather than close to C₅. However, it should be remembered that the MestreNova predictions are carried out on directly-inbuilt and optimized structures within that code, and as such, they can slightly differ from the Gaussian-inbuilt and optimized structures.



Chapter 4

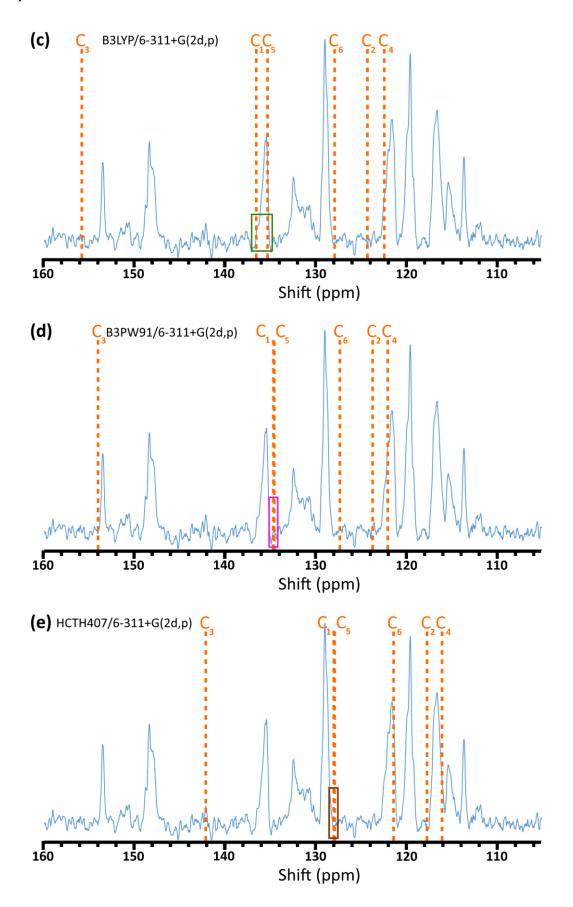


Figure 49. A comparison of the OAPS ¹³C-NMR experimental spectrum (blue line) with the chemical shifts predicted for the aromatic carbons of the 1-ring *meta*-OAPS by Gaussian (dashed lines) and MestreNova (lines).

As the five predicted spectra are not easy to distinguish, the quality of the prediction was assessed by calculating the differences between all predicted peaks and their closest experimental peaks according to Equation 7. To eliminate the background, only the experimental peaks whose intensity was at least 25% of the main peak at 129 ppm were selected. Although this is somewhat subjective, this allows for the comparison of all predicted spectra under the same conditions.

$$Difference = \sqrt{\sum (\text{predicted shifts - closest experimental shifts})^2}$$
(7)

The differences in ppm are reported in Table 12 and show that Gaussian at the B3PW91/6-311+G(2d,p) level of theory gives the best prediction for the 1-ring *meta*-OAPS.

 Table 12. The differences between the experimental and the predicted ¹³C-NMR peaks for the 1-ring *meta*-OAPS according to Equation 7.

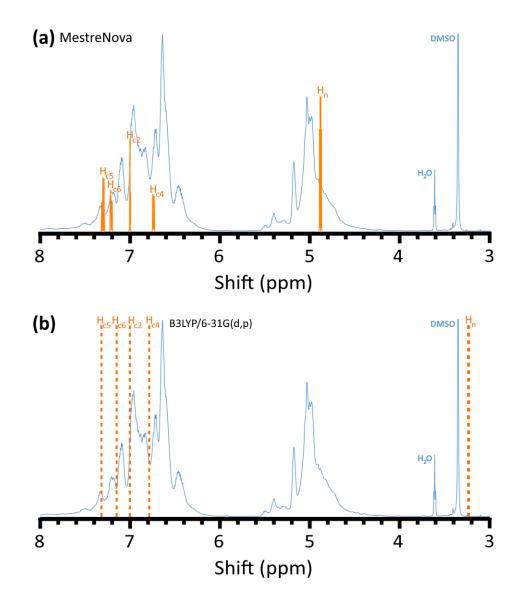
Method	MestRenova	B3LYP 6-31G(d,p)	B3LYP 6-311+G(2d,p)	B3PW91 6-311+G(2d,p)	HCTH407 6-311+G(2d,p)
Difference / ppm	7.4	7.8	4.0	3.3	6.5

b.2.3.) The ¹H-NMR predicted spectra for the 1-ring meta-OAPS

Figure 50 shows the predicted ¹H-NMR spectra for the 1-ring *meta*-isomer along with the experimental ¹H-NMR spectrum for the commercial OAPS (blue line). Lee et al.¹⁰⁷ did not report any ¹H-NMR spectra for their pure *meta*-OAPS and *para*-OAPS structures. In the predictions, the ¹H-NMR shifts of the hydrogens directly connected to the Si were omitted. As for the aromatic carbons in Figure 49, each predicted shift can be assigned to a specific hydrogen as defined in Figure 47(a).

The (a) MestreNova spectrum matches rather well the experimental spectrum for both aromatic (left part of the spectrum) and amine (right part of the spectrum) hydrogens. The Gaussian-predicted shifts for the aromatic hydrogens at the (b) lower B3LYP/6-31G(d,p) level of theory are also fairly good consistent with the experimental spectrum, while the shifts obtained from the larger (c-e) 6-311+G(2d,p) basis set appear to be slightly overestimated. In

addition, they predict possible overlaps as indicated by the pink rectangles. The hydrogen shifts are in the order $H_{C4} < H_{C2} < H_{C6} < H_{C5}$ for all predictions, i.e. they follow the same order as the Gaussian predictions of their attached carbons. However, unlike MestreNova, all Gaussian-predicted peaks appear to be strongly underestimated for the amine hydrogens. The same problem was reported in other studies,²⁰⁰⁻²⁰¹ and was attributed to the NH₂ groups forming hydrogen bonds with the solvent, which is not properly taken into account by the PCM formalism.²⁰¹⁻²⁰²



Chapter 4

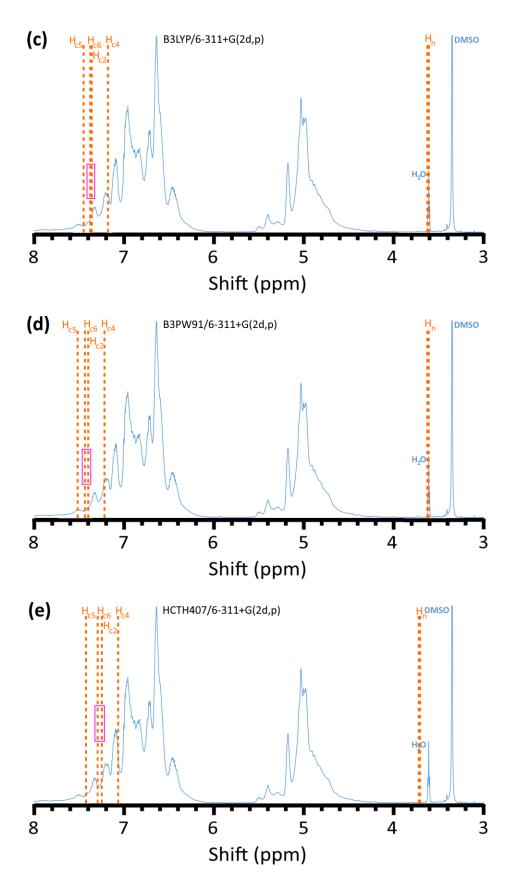


Figure 50. A comparison of the OAPS ¹H-NMR experimental spectrum (blue line) with the chemical shifts predicted for the aromatic and amine hydrogens of the 1-ring *meta*-OAPS by Gaussian (dashed lines) and MestreNova (lines).

Chapter 4

Since it is more difficult to assign each ¹H-NMR predicted shift to a specific experimental one than for the ¹³C-NMR spectra (Equation 7 and Table 12), the methods were here compared based on the distance between the lowest shift for the aromatic H and the sharpest experimental peak at 6.64 ppm (Table 13).

 Table 13. The differences between the main experimental and the closest predicted ¹H-NMR peaks for the aromatic hydrogens of the 1-ring *meta*-OAPS.

Method	MestRenova	B3LYP 6-31G(d,p)	B3LYP 6-311+G(2d,p)	B3PW91 6-311+G(2d,p)	HCTH407 6-311+G(2d,p)
Difference / ppm	0.10	0.10	0.60	0.66	0.49

Although MestreNova could not produce results close to the experimental ²⁹Si-NMR spectrum (Figure 48), it shows matchable shifts for the ¹³C-NMR and especially for the ¹H-NMR spectra (Table 13). So it was used again to evaluate the NMR shifts of the 8-ring isomers. DFT at the lower 6-31G(d,p) level of theory also performed quite well for the aromatic hydrogens in the ¹H-NMR spectrum (Table 13), but it was much further from the experiment for both the ²⁹Si-NMR and the ¹³C-NMR spectra. Consequently, it was discarded and only B3LYP, HCTH407, and B3PW91 at the main 6-311+G(2d,p) level of theory were further tested with Gaussian for the 8-ring structures.

b.2.4.) The ²⁹Si-NMR predicted spectra for the 8-ring meta-OAPS

Figure 51 shows the predicted ²⁹Si-NMR shifts for the 8-ring *meta*-OAPS along with the same two experimental spectra as in Figure 48. All Gaussian-predicted shifts are moved to the left in comparison to the 1-ring *meta*-OAPS, which confirms that the presence of the eight aminophenyl rings displaces the shifts. On the other hand, the MestreNova results do not depend on the structure being a 1-ring or an 8-ring *meta*-OAPS, which is a drawback of this approach. They remain also clearly too far from the experimental ²⁹Si-NMR peaks.

At the main 6-311+G(2d,p) level of theory, the DFT shifts predicted by B3LYP and HCTH407 move to -86.8 ppm and -85.7 ppm, respectively. The B3PW91-predicted shift is displaced by 3.7 ppm to -78.8 ppm, which now falls within the experimental range, whereas the HF-predicted peak at -74.9 ppm is further away. With the extended 6-311++G(2df,2dp)

basis set, B3PW91 and HF predict peaks at -80.3 ppm and -78.0 ppm, which are also within the experimental range, unlike the HCTH407-predicted shift at -85 ppm.

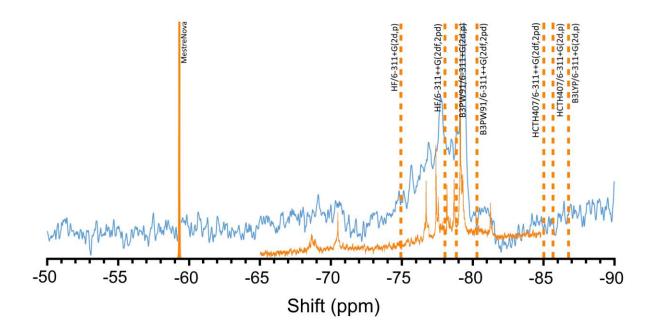


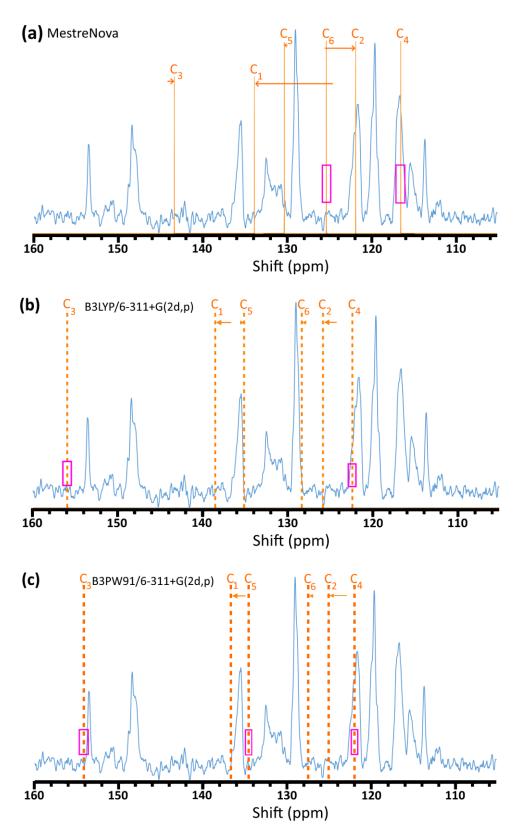
Figure 51. The experimental ²⁹Si-NMR spectra of the commercial OAPS (blue line) and the pure *meta* isomer of Lee et al. (orange line)¹⁰⁷ compared with the predicted chemical shifts by Gaussian (dashed lines) and MestreNova (lines) for the 8-ring *meta*-OAPS.

As for the one-ring OAPS (Figure 48), the Gaussian approaches with B3PW91/6-311+G(2d,p) and HF/6-311++G(2df,2pd) predict the closest ²⁹Si-NMR shifts to the experimental peaks of the pure *meta* isomer. Consequently, these will be the methods tested to predict the ²⁹Si-NMR shifts for the *ortho* and *para* isomers.

b.2.5.) The ¹³C-NMR predicted spectra for the 8-ring meta-OAPS

Figure 52 shows the predicted ¹³C-NMR shifts for the 8-ring *meta*-OAPS along with the commercial OAPS experimental spectrum. Unlike for the 1-ring *meta*-OAPS (Figure 49), no overlaps are predicted by neither MestRenova nor Gaussian. The pink rectangles in Figure 52 indicate the few predicted shifts which are the same as for the 1-ring *meta*-OAPS. All other shifts were found to differ, indicating how the presence of the aminophenyl rings affects the results. The displacements of the predicted chemical shifts from the 1-ring to the 8-ring *meta*-OAPS are shown by arrows (for MestreNova and for Gaussian), whose lengths are proportional to the changes. As shown in (a), the largest change is found

for MestreNova with the 1-ring shift at ~124 ppm being displaced by ~10 ppm to the left in the 8-ring structure. This results in the ¹³C-NMR MestreNova predicted shifts associated with specific carbons being now in the same order as for Gaussian, which is also identical to that in the one-ring *meta* structure.



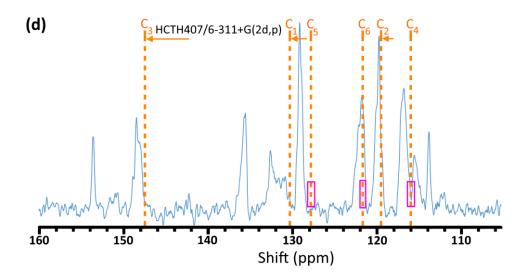


Figure 52. A comparison of the commercial OAPS ¹³C-NMR experimental spectrum (blue line) with the chemical shifts predicted for the aromatic carbons of the 8-ring *meta*-OAPS by Gaussian (dashed lines) and MestreNova (lines). The arrows — and — indicate the differences in shifts with respect to the 1-ring *meta*-OAPS generated by Gaussian and MestRenova respectively.

There again, the quality of the prediction was assessed according to the difference between the predicted peaks and their closest experimental peaks (Equation 7). As for the 1-ring analyses (Table 12), only the experimental peaks whose intensity was at least 25% of the main peak at 129 ppm were selected to eliminate the background. These differences in ppm are reported in Table 14 for the various methods tested and show that HCTH407/6-311+G(2d,p) now gives the best prediction.

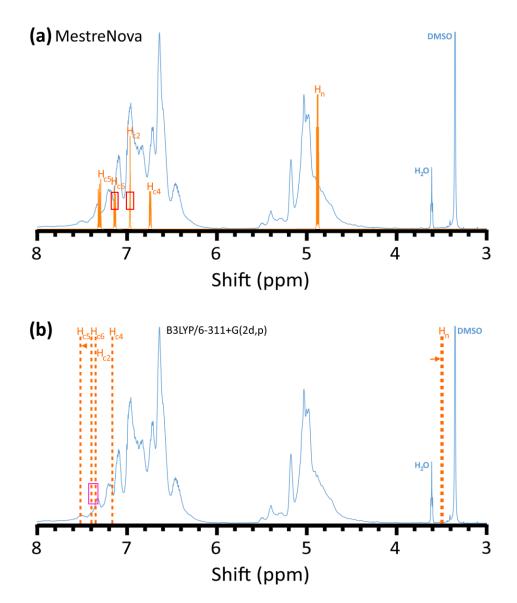
 Table 14. The differences between the experimental and the predicted ¹³C-NMR peaks for the 8-ring *meta*-OAPS according to Equation 7.

Method	MestRenova	B3LYP 6-311+G(2d,p)	B3PW91 6-311+G(2d,p)	HCTH407 6-311+G(2d,p)
Difference / ppm	2.5	2.3	2.1	1.4

Compared to the 1-ring *meta*-OAPS, the HCTH407 method at the 6-311+G(2d,p) level of theory seems to be better than B3PW91, although the latter provides the second-best prediction. This change in order is due to the three HCTH407-predicted shifts at 116.2 ppm (C4), 119.7 ppm (C2), and 121.8 ppm (C6) being very close to their respective experimental peaks. As such, Gaussian with HCTH407/6-311+G(2d,p) will be the method used to predict the ¹³C-NMR shifts for the *ortho* and *para* isomers.

b.2.6) The ¹H-NMR predicted spectra for the 8-ring meta-OAPS

Figure 53 shows the predicted ¹H-NMR spectra for the 8-ring *meta*-isomer along with the commercial OAPS experimental spectrum. The red rectangles indicate the chemical shifts predicted by (a) MestreNova, which are somewhat displaced to the right in comparison to the 1-ring *meta*-OAPS (Figure 50). As a whole, they match even better with the experimental data, thus suggesting that MestreNova is well adapted to ¹H-NMR prediction. On the contrary, the three Gaussian methods (b-d) still overestimate the chemical shifts for the aromatic hydrogens (with the pink rectangles showing possible overlaps) and underestimate the chemical shifts for the amine hydrogens. The displacements of the predicted chemical shifts from the 1-ring to the 8-ring *meta*-OAPS are again shown by arrows, but they remain very small. Consequently, the relative order for the specific ¹H-NMR shifts remains the same as for the 1-ring structure.



Chapter 4

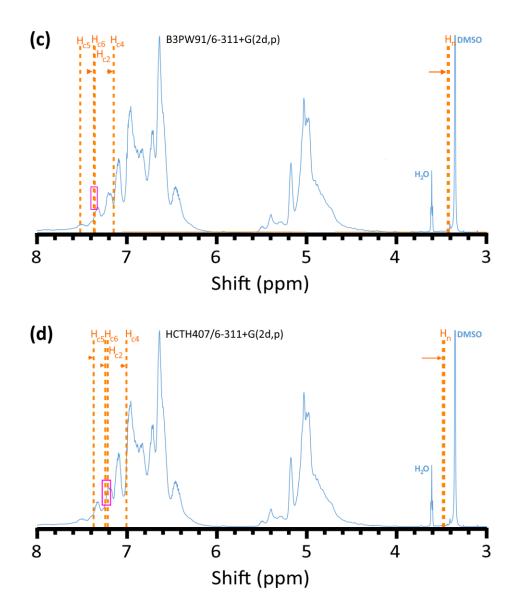


Figure 53. A comparison of the commercial OAPS ¹H-NMR experimental spectrum (blue line) with the chemical shifts predicted for the aromatic and amine hydrogens of the 8-ring *meta*-OAPS by Gaussian (dashed lines) and MestreNova (lines). The arrows findicate the differences in predicted shifts for the 8-ring *meta*-OAPS with respect to the 1-ring *meta*-OAPS by Gaussian.

As for the 1-ring *meta*-OAPS (Table 13), the methods were compared by calculating the distance between the lowest shift of the aromatic H and the sharpest experimental peak at 6.64 ppm. They are reported in Table 15 for the 8-ring *meta*-OAPS. If MestreNova still gives the best prediction, the HCTH407 method at the 6-311+G(2d,p) level of theory is now the most appropriate Gaussian approach. This is interesting as this method was selected as the optimal one for ¹³C-NMR.

Method	MestRenova	B3LYP 6-311+G(2d,p)	B3PW91 6-311+G(2d,p)	HCTH407 6-311+G(2d,p)
Difference / ppm	0.01	0.51	0.50	0.17

 Table 15. The differences between the main experimental and the closest predicted ¹H-NMR peaks for the aromatic hydrogens of the 8-ring *meta*-OAPS.

MestreNova clearly produces more compatible ¹H-NMR chemical shifts for *meta*-OAPS than the time-consuming Gaussian DFT calculations but its major downside is that it performs very poorly for both ¹³C-NMR and ²⁹Si-NMR spectra. As it would be better to be able to predict all types of spectra with similar approaches, it was decided to further prospect the ¹H-NMR OAPS spectra using both MestreNova and the best DFT approach, *i.e.* HCTH407/6-311+G(2d,p). As such, both methods were tested in parallel to predict the *ortho* and *para* ¹H-NMR spectra in addition to that of the *meta* isomer.

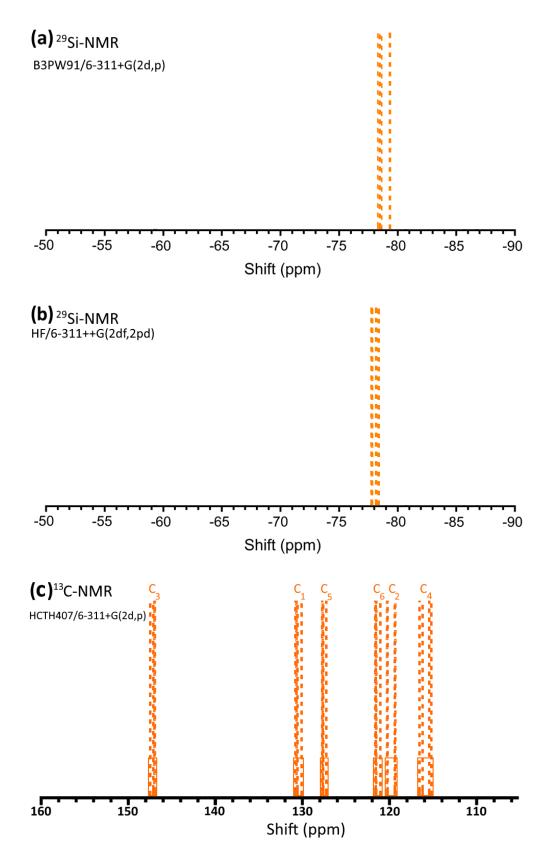
b.2.7) Reproducibility: ²⁹Si-NMR, ¹³C-NMR, and ¹H-NMR for meta-OAPS

Before studying both other isomers, Figure 54 assesses the reproducibility of the optimized methods for the prediction of the *meta*-OAPS ²⁹Si-NMR, ¹³C-NMR, and ¹H-NMR spectra. For this purpose, three more structures of the 8-ring *meta*-OAPS were prepared from scratch as before, *i.e.* first by Hyperchem²⁰³ and then optimized by Gaussian at the B3LYP/6-31G(d,p) level. Three new 8-ring *meta*-OAPS structures were also drawn and optimized by MestRenova. The NMR-predicted shifts were obtained using:

* both B3PW91/6-311+G(2d,p) and HF/6-311++G(2df,2pd) for ²⁹Si-NMR

- * HCTH407/6-311+G(2d,p) for ¹³C-NMR,
- * both MestreNova and HCTH407/6-311+G(2d,p) for ¹H-NMR

in order to test the reproducibility of the predictions initially carried out on a single 8-ring *meta*-OAPS structure (Figures 51-53).



Chapter 4

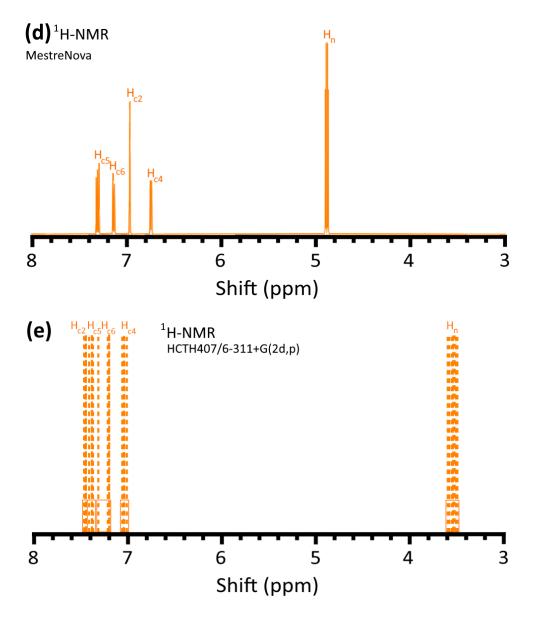


Figure 54. Reproducibility of the predicted spectra for (a-b) ²⁹Si-NMR (c) ¹³C-NMR and (d-e) ¹H-NMR using the optimized methods described above for four different 8-ring *meta*-OAPS structures

For the ²⁹Si-NMR shifts, four different structures predict results within a difference of less than 1 ppm at the same level of theory. For the ¹³C-NMR shifts, the differences are less than 2 ppm as indicated by the width of the orange rectangles. For the ¹H-NMR shifts, no differences are observed between the four different structures with MestreNova. It should be mentioned that this code calculates the spin-spin coupling along with the shifts, thus splitting each peak into a few smaller peaks. For example, the peak at 7.3 ppm for H_{c5} is split into three smaller peaks because of its adjacents H_{c4} and H_{c6} on the aromatic ring. For the ¹H-NMR shifts calculated with Gaussian, the differences are less than 0.2 ppm as indicated by the width of the orange rectangles. The reproducibility is thus rather good for all cases tested.

b.3) The ortho and para isomers

Figures 55(a) and 55(b) show the different chemical environments for the carbons and hydrogens of the *ortho* and *para* isomers of OAPS, respectively. Due to the symmetry in the *para* isomer, there are only four carbon types and three hydrogen types with different chemical environments. On the other hand, as for the *meta* isomer, each carbon and each hydrogen in the *ortho* isomer presents a different chemical environment.

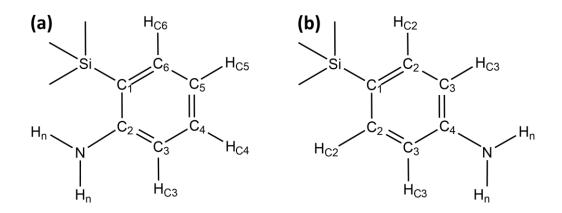
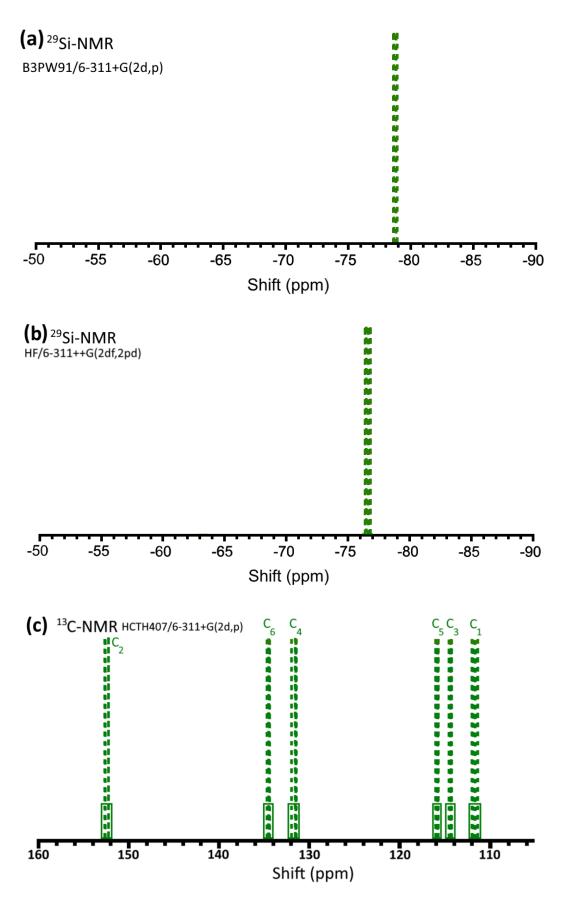


Figure 55. The molecular structures of a) the *ortho* and b) the *para* isomers connected to OAPS.

Figures 56 and 57 display the spectra predicted by the optimized methods for the ²⁹Si-NMR, ¹³C-NMR, and ¹H-NMR shifts of the 8-ring *ortho*-OAPS and *para*-OAPS isomers, respectively. As for Figure 54, they were calculated on four different structures each in order to directly assess their reproducibility. The predicted shifts associated with specific aromatic carbons and hydrogens are indicated in the ¹³C-NMR and ¹H-NMR analyses.

The reproducibility of the predictions for the *ortho* isomer (Figure 56) is very good. The predicted ²⁹Si-NMR shifts by (a) B3PW91/6-311+G(2d,p) are within less than 0.5 ppm for the four different structures whereas those predicted by (b) the HF methods are within less than 1 ppm. The ¹³C-NMR predictions all fall within less than 1 ppm (c). The ¹H-NMR shifts predicted by MestreNova are identical for the four structures (d), while those predicted by HCTH407/6-311+G(2d,p) (e) also show a very good reproducibility within a range smaller than 0.1 ppm for each specific hydrogen.



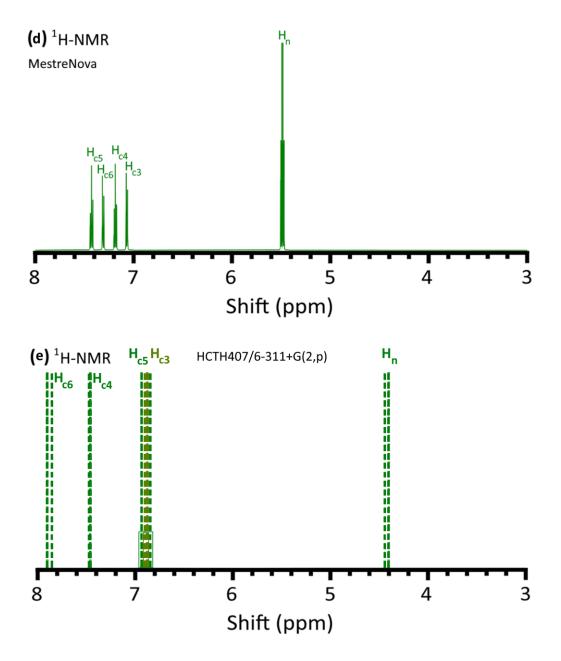
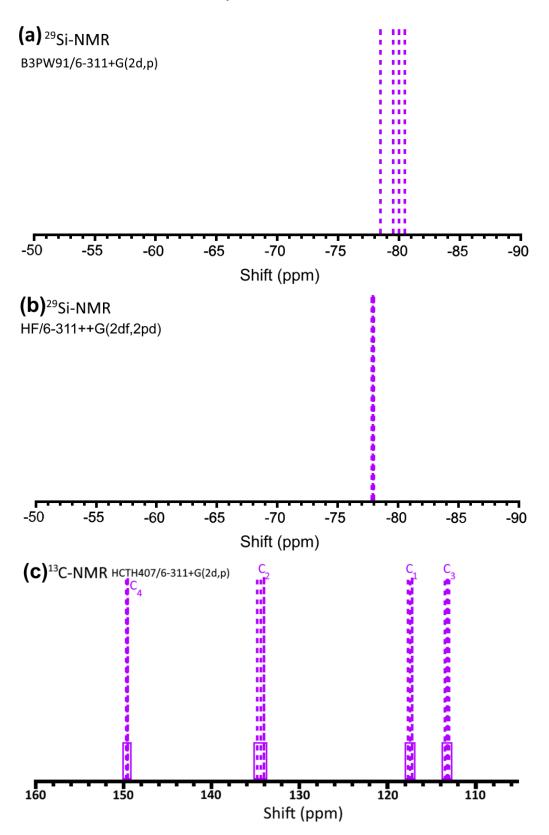


Figure 56. Reproducibility of the predicted spectra for (a-b) ²⁹Si-NMR (c) ¹³C-NMR and (d-e) ¹H-NMR using the optimized methods for four different 8-ring *ortho*-OAPS structures

For the *para* isomer (Figure 57), the predicted ²⁹Si-NMR shifts at (a) B3PW91/6-311+G(2d,p) cover a wider range for the four different structures than for the *ortho* isomer, but they still fall within only 2 ppm. On the other hand, HF/6-311++G(2df,2pd) predicts more centralized data (b). The range for the predicted specific ¹³C-NMR shifts (c) is less than 1 ppm and as expected, there are only four different carbon types in the *para* isomer. Similarly, MestreNova (d) predicts three ¹H-NMR shifts and the HCTH407/6-311+G(2d,p) level of theory (e) displays a range of less than 0.2 ppm for the amine shifts and less than 0.1 ppm for the aromatic shifts (e). The surface areas of the three peaks in MestreNova (d) are equal

Chapter 4

because there are two hydrogens in each chemical environment (one set of two amine hydrogens and two sets of two aromatic hydrogens). Gaussian only provides the shifts and no data about the intensities can be directly extracted from its results.



88

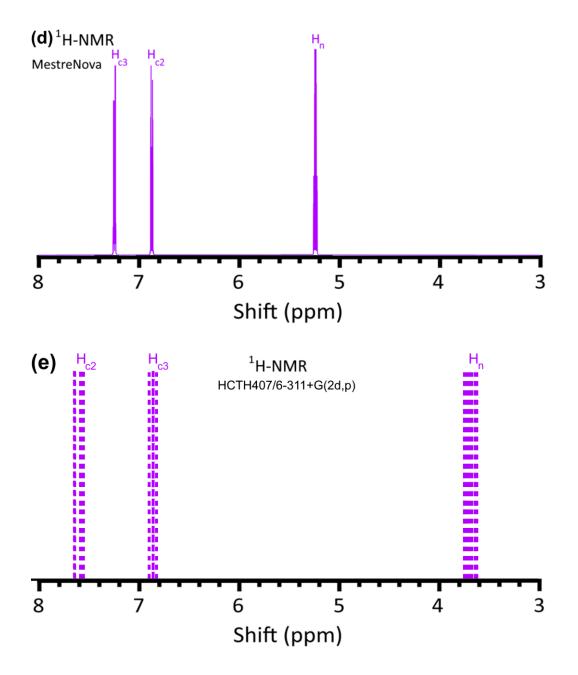


Figure 57. Reproducibility of the predicted spectra for (a-b) ²⁹Si-NMR (c) ¹³C-NMR and (d-e) ¹H-NMR using the optimized methods for four different 8-ring *para*-OAPS structures

b.4) Comparison of the three isomers

To compare all three isomers with the experimental spectra and account for the slight differences in geometries, the predicted shifts were averaged on the four different structures studied previously for each isomer.

b.4.1) The ²⁹Si-NMR predicted spectra for meta-OAPS, ortho-OAPS, and para-OAPS

Following the *meta*-OAPS tests (Section b.2.), the ²⁹Si-NMR spectra were predicted for all three isomers using both DFT at the B3PW91/6-311+G(2d,p) level of theory and the heavier HF/6-311++G(2df,2pd) approach. The average predicted spectra for *meta*-OAPS (orange), *para*-OAPS (purple) and *ortho*-OAPS (green) are compared in Figure 58 to the experimental (a) commercial, (b) controlled and (c) pure *meta*-OAPS and *para*-OAPS of Lee et al.¹⁰⁷ The *x*-axis scale was zoomed to the -75 to -82 ppm range to better distinguish the differences.

For the B3PW91/6-311+G(2d,p) method, the average predicted shits of the *ortho* and *meta* isomer differ by less than 0.2 ppm, which suggests an overlap. This could explain why there are similar experimental peaks in (a) and (b), whereas there should be no *ortho* isomer in the latter. However, the predicted shift of the *para* isomer is on the right side of both other isomers, unlike the order of the intensities in the commercial OAPS (which should have less *para* and more *meta* isomer), the controlled OAPS (which should have more *para* and less *meta* isomer), as well the pure samples of Lee et al. Although, there is a much better agreement for B3PW91/6-311+G(2d,p) than for the other DFT predictions tested (Figures 48 and 51), it does not seem to be able to properly reproduce the isomer differences.

Using a different method can produce different results.²⁰⁴ Indeed, HF/6-311++G(2df,2pd) shows not only slightly different shifts but also a different order for the shifts. The *para* peak is now on the left side of the *meta*, which is more consistent with both experimental spectra. The *ortho* peak is displaced towards higher values and it is difficult to assign it to a specific part of the commercial *ortho*-containing spectra. However, there is not any clear *ortho* signature in any of the experimental spectra (Figure 42), and as such, HF/6-311++G(2df,2pd) probably provides a better prediction than B3PW91/6-311+G(2d,p).

It is difficult to get fully satisfactory results from the ²⁹Si-NMR prediction methods tested here, which could be due to the fact that the peaks all fall within a very restricted ppm range. It is also difficult to explain why the main experimental peak of Lee et al. for the pure *para* isomer (c) does not clearly appear neither in the commercial (a) nor in the controlled (b) samples, especially considering that the latter should have a significant amount of *para* isomer. So there are still issues in both the predicted and experimental shifts.

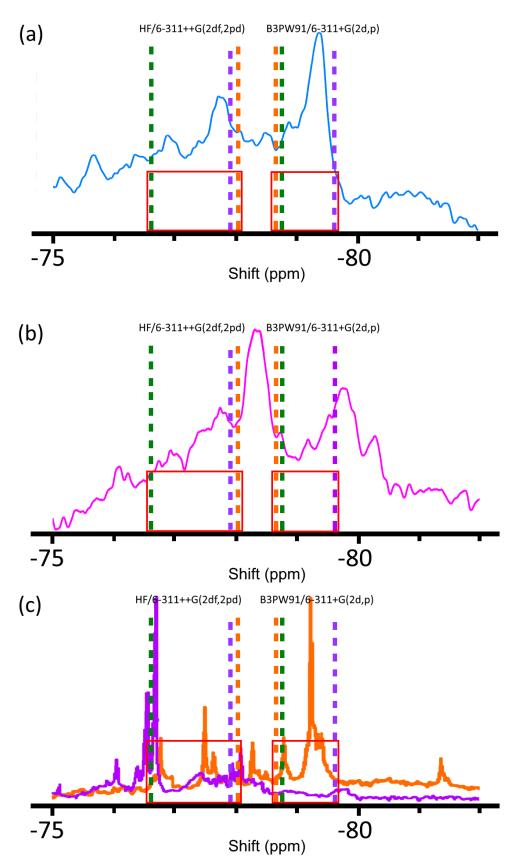


Figure 58. The experimental ²⁹Si-NMR spectra of (a) the commercial OAPS (blue line) (b) the controlled OAPS (pink line) and (c) the pure *meta* (orange line) and *para* (purple line) isomers of Lee et al.¹⁰⁷ compared to the predicted shifts by Gaussian (dashed lines) with B3PW91/6-311+G(2d,p) and HF/6-311++G(2df,2pd) for the there OAPS isomers.

b.4.2) The ¹³C-NMR predicted spectra for meta-OAPS, ortho-OAPS, and para-OAPS

The average ¹³C-NMR spectra were predicted for all three isomers using DFT at the HCTH407/6-311+G(2d,p) level of theory. They are compared in Figure 59 to those of the experimental (a) commercial and (b) controlled samples. As before, the predicted peaks are identified by their colours with orange representing the *meta* isomer, green representing the *ortho* isomer and purple representing the *para* isomer.

In Figure 59(a), the predicted *ortho* green shift at 111.5 ppm (P) matches a weak experimental peak for the commercial OAPS sample. A similar match can be observed for two *para* purple shifts at 117.5 ppm (K) and 149.5 ppm (B). The *meta* orange shift at 130.5 ppm (G) and the *ortho* green shift at 131.5 ppm (F) are close to two medium-intensity experimental peaks. Both *meta* orange shifts at 120.5 ppm (J) and 121 ppm (I) are adjacent to stronger experimental peaks. The overlap at 134.5 ppm (*ortho* green in D and *para* purple in E) is close to the experimental peak at 136 ppm. The *meta* orange shift at 127.8 ppm (H) is within less than 1 ppm from the sharpest experimental peak at 129 ppm. This good compatibility between the various peaks indicates that the HCTH407/6-311+G(2d,p) level of theory seems to produce rather accurate ¹³C-NMR predictions for the commercial OAPS. In addition, the same number of experimental and predicted shifts confirms that all three isomers are present in this sample.

Figure 59(b) shows that each predicted *meta* or *para* shift is also very close to an adjacent experimental peak in the controlled OAPS sample. The predicted B and C shifts are next to weak experimental signals at 151 ppm and 148 ppm, respectively. The same compatibility can be observed for the G, H, I, J, K and L predicted shifts with adjacent experimental signals. Two *para* shifts, E and O, are compatible with the main experimental peaks at 135 ppm and 113 ppm respectively. As expected, there are only ten peaks in the experimental spectrum. This confirms the absence of the *ortho* isomer in the controlled sample, whereas there are sixteen peaks in (a) for the commercial OAPS.

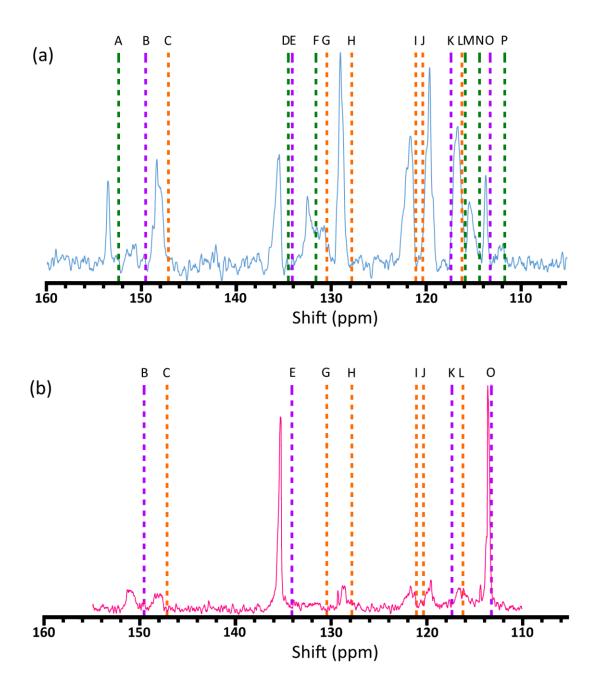


Figure 59. The experimental ¹³C-NMR spectra of (a) the commercial OAPS (blue line) and (b) the controlled OAPS (pink line) compared with the predicted shifts by Gaussian (dashed lines) at the HCTH407/6-311+G(2d,p) level of theory for the three OAPS isomers.

Tables 16 and 17 give the precise chemical shift values for each peak attributed to each specific carbon of each isomer (Figures 47 and 55) in the commercial and controlled OAPS samples, respectively. They also report in each case the absolute errors e_i between the experimental and predicted shifts. The mean absolute error (MAE) and root mean rectangle error (RMSE) can be calculated from:²⁰⁵

$$MAE = \frac{1}{n} \sum_{i=1}^{n} |e_i|$$
(8)

$$\text{RMSE} = \sqrt{\frac{1}{n} \sum_{i=1}^{n} e_i^2}$$
(9)

where *n* is the number of shifts.

For the commercial sample (Table 16), the MAE and RMSE are 0.80 ppm and 0.90 ppm, respectively. For the controlled sample (Table 17), the MAE and RMSE are 0.80 ppm and 0.95 ppm. These RMSE values are much smaller than the values of 1.89 ppm reported by Abil et al.¹⁸³ and 4.46 ppm reported by Zhang et al.²⁰⁶ for a large variety of small molecules without any silicon atoms. Such small errors indicate that HCTH407/6-311+G(2d,p) is indeed pretty accurate in predicting the experimental ¹³C-NMR chemical shifts for OAPS. In addition, the number of peaks confirms the presence (Table 16) and absence (Table 17) of the *ortho* isomer in the respective experimental samples.

C-Type	δ_{calc}	δ_{exp}	$ e_i $	Identification	Isomer type
C 1	130.45	130.70	0.25	G	meta
C ₂	120.36	119.55	0.81	J	meta
C ₃	147.13	148.34	1.21	С	meta
C ₄	116.30	116.64	0.34	L	meta
C 5	127.82	128.97	1.15	н	meta
C ₆	121.11	121.60	0.49	I	meta
C 1	111.75	111.78	0.03	Р	ortho
C ₂	152.41	153.53	1.12	Α	ortho
C ₃	114.40	115.28	0.88	N	ortho
C ₄	131.59	132.45	0.86	F	ortho
C 5	115.90	115.28	0.62	м	ortho
C ₆	134.53	135.54	1.01	D	ortho
C 1	117.42	116.64	0.78	к	para
C ₂	134.09	135.54	1.45	E	para
C ₃	113.28	113.60	0.32	0	para
C ₄	149.57	151.00	1.43	В	para

Table 16. Comparison between the predicted δ_{calc} and experimental δ_{\exp} ¹³C NMR shifts in ppm along with their absolute error e_i in ppm for each peak identified in the commercial OAPS sample (Figure 59(a)).

C-Type	δ_{calc}	δ_{exp}	$ e_i $	Identification	Isomer type
C ₁	130.45	129.28	1.17	G	meta
C ₂	120.36	119.74	0.62	J	meta
C ₃	147.13	147.71	0.58	С	meta
C ₄	116.30	116.13	0.17	L	meta
C 5	127.82	128.20	0.38	Н	meta
C ₆	121.11	121.51	0.40	I	meta
C 1	117.42	116.13	1.29	К	para
C ₂	134.09	135.35	1.26	E	para
C ₃	113.28	113.59	0.31	0	para
C 4	149.57	151.36	1.79	В	para

Table 17. Same as Table 16 but for in the controlled OAPS sample (Figure 59(b)).

The calculated shielding tensors and experimental chemical shifts can be related by using a linear relationship described with a linear equation $\delta_{calc-scaled} = a\delta_{calc} + b$ which cancels systemic errors (a and b are slope and intercept, $\delta_{calc-scaled}$ is the scaling calculated chemical shifts close to the experimental chemical shifts, and δ_{calc} is calculated shielding tensors).^{183, 198, 207} Figure 60 shows the linear correlation¹⁸³ between the δ_{calc} and δ_{exp} ¹³C-NMR shifts for the commercial OAPS (Table 16), which can be fitted by:

$$\delta_{\text{calc-scaled}} = 1.04 \ \delta_{\text{calc}} - 4.36 \tag{10}$$

The slope is close to 1 and the correlation is very good (R = 0.9992).

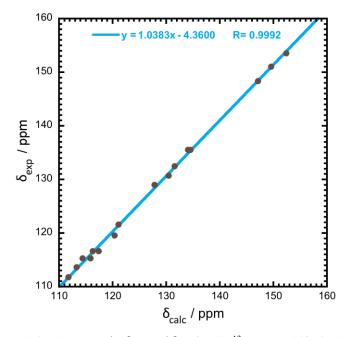


Figure 60. Linear correlation between the δ_{calc} and δ_{exp} for the ¹³C-NMR shifts in the commercial OAPS

Equation 10 can be used to scale the δ_{calc} originating from Gaussian in order for them to be even closer to the experimental δ_{exp} . These scaled predicted shifts $\delta_{calc-scaled}$ are reported in Table 18 for the commercial OAPS. They decrease the MAE and RMSE to 0.44 ppm and 0.54 ppm, respectively. Such a correlation can thus further improve the chemical shifts predicted at the HCTH407/6-311+G(2d,p) level of theory for OAPS.

Table 18. Comparison between the predicted scaled $\delta_{\text{calc-scaled}}$ and experimental δ_{exp} ¹³C NMR shifts in ppm along with their absolute error e_i in ppm for each peak identified in the commercial OAPS sample (Figure

С Туре	$\delta_{\text{calc-scaled}}$	δ_{exp}	$ e_i $	Identification	Isomer type
C 1	131.09	130.70	0.39	G	meta
C ₂	120.61	119.55	1.06	J	meta
C ₃	148.41	148.34	0.07	С	meta
C 4	116.39	116.64	0.25	L	meta
C 5	128.36	128.97	0.61	н	meta
C ₆	121.39	121.60	0.21	I	meta
C 1	111.67	111.78	0.11	Р	ortho
C ₂	153.89	153.53	0.36	Α	ortho
C ₃	114.42	115.28	0.86	Ν	ortho
C 4	132.27	132.45	0.18	F	ortho
C 5	115.98	115.28	0.70	М	ortho
C ₆	135.32	135.54	0.22	D	ortho
C 1	117.56	116.64	0.92	к	para
C ₂	134.87	135.54	0.67	Е	para
C ₃	113.26	113.60	0.34	Ο	para
C 4	150.94	151.00	0.06	В	para

59(a)).

Figure 61 shows likewise the very good linear correlation (R = 0.9983) between the δ_{exp} and δ_{calc} ¹³C-NMR shifts for the controlled OAPS (Table 17), which can be fitted by:

$$\delta_{\text{calc-scaled}} = 1.05 \ \delta_{\text{calc}} - 5.85 \tag{11}$$

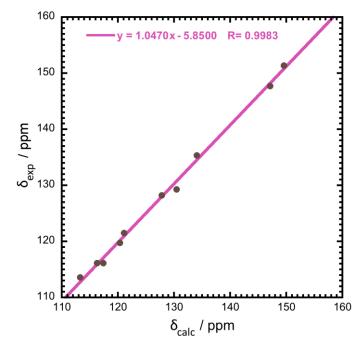


Figure 61. Linear correlation between the δ_{calc} and δ_{exp} for the ¹³C-NMR shifts in the controlled OAPS

As for Figure 60, the slope of the linear correlation in Figure 61 is close to 1, which confirms the accuracy of the predicted ¹³C-NMR shifts for both types of OAPS. The $\delta_{calc-scaled}$ obtained from applying Equation 11 are reported in Table 19 for the controlled OAPS. The MAE and RMSE decrease to 0.66 ppm and 0.75 ppm, respectively.

C-Type	$\delta_{\text{calc-scaled}}$	δexp	$ e_i $	Identification	Isomer type
C 1	130.73	129.28	1.45	G	meta
C ₂	120.17	119.74	0.43	J	meta
C 3	148.19	147.71	0.48	С	meta
C 4	115.91	116.13	0.22	L	meta
C 5	127.98	128.20	0.22	н	meta
C ₆	120.95	121.51	0.56	I	meta
C 1	117.09	116.13	0.96	к	para
C ₂	134.54	135.35	0.81	Е	para
C ₃	112.76	113.59	0.83	Ο	para
C 4	150.75	151.36	0.61	В	para

Table 19. Comparison between the predicted scaled $\delta_{\text{calc-scaled}}$ and experimental δ_{exp} ¹³C NMR shifts in ppm along with their absolute error e_i in ppm for each peak identified in the controlled OAPS sample (Figure 59(b)).

These analyses confirm that Gaussian at the HCTH407/6-311+G(2d,p) level of theory is able to predict very accurately the ¹³C-NMR shifts for all OAPS isomers, regardless of

whether there are three isomers (commercial sample) or only two isomers (controlled sample).

b.4.3) The ¹H-NMR predicted spectra for meta-OAPS, ortho-OAPS, and para-OAPS

The ¹H-NMR spectra were predicted for all three isomers using both MestreNova and Gaussian at the HCTH407/6-311+G(2d,p) level of theory. They are compared to those of the experimental (a) commercial and (b) controlled samples in Figure 62 for MestreNova and in Figure 63 for Gaussian. The letters used to identify the predicted peaks for the aromatic hydrogens are those of their adjacent aromatic carbons (Tables 18-19). The amine hydrogens connected to the nitrogen are referred to by the name of their isomer.

For the aromatic hydrogens of the commercial sample (left regions of Figures 62(a) and 63(a)), ten slightly overestimated shifts are predicted by both MestreNova and Gaussian, with four of them overlapping. As a result, only eight peaks are observable from 6.7 ppm to 7.9 ppm, in agreement with the eight observable experimental shifts (red rectangles). However, it is worth noting that the relative order of these peaks is somewhat different between MestreNova and Gaussian, which means that both codes do not assign the same peaks to the same isomer. For the amine hydrogens (right regions of Figures 62(a) and 63(a)), three MestreNova shifts are predicted in the range 4.9-5.5 ppm, *i.e.* close to the three experimental peaks. This is compatible with the presence of the *ortho* isomer in addition to the *meta* and *para* isomers in the commercial sample. Gaussian also predicts three separate shifts albeit in an underestimated range, i.e. 3.5-4.5 ppm. Unlike the aromatic hydrogens, both predicted sets of amine hydrogen shifts are in the same relative order.

The situation is similar for the predicted amine hydrogen shifts (right regions of Figures 62(b) and 63(b)), where the experimental peak related to the *ortho* isomer disappears. For the aromatic hydrogens (left regions of Figures 62(b) and 63(b)), the predicted shifts are again slightly displaced towards higher values and their identification differs according to the code used. There are six of them, which relate to the six observable experimental shifts for the *meta* and *para* isomers (red rectangles).

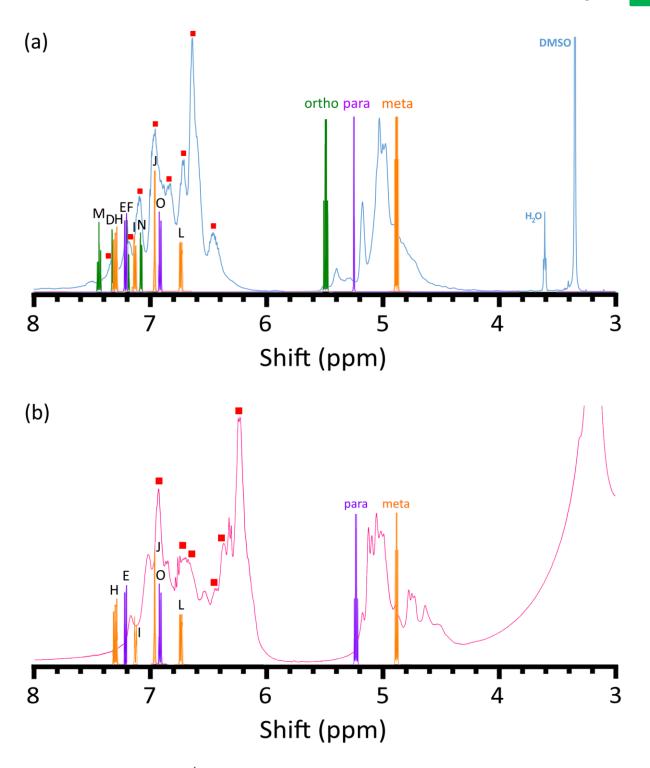


Figure 62. The experimental ¹H-NMR spectra of (a) the commercial OAPS (blue line) and (b) the controlled OAPS (pink line) compared with the predicted shifts by MestreNova (orange, green and purple lines) for the three isomers.

Chapter 4

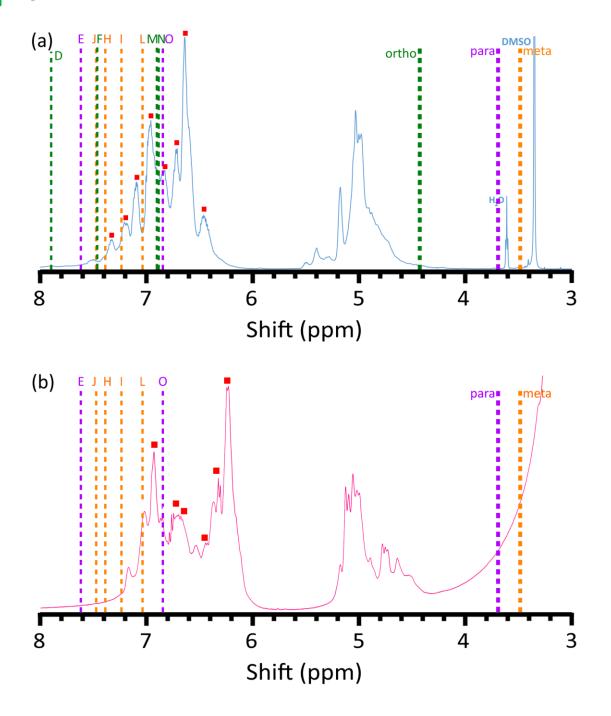


Figure 63. Same as Figure 62 but for the Gaussian-predicted shifts at the HCTH407/6-311+G(2d,p) level of theory.

Tables 20 and 21 report the absolute errors e_i between the experimental δ_{exp} and the MestreNova-predicted δ_{calc} ¹H-NMR shifts in both the commercial and controlled OAPS samples. The same analyses carried out with the Gaussian predictions at the HCTH407/6-311+G(2d,p) level of theory are provided in Tables 22 and 23.

The MestreNova-predicted MAE and RMSE for the ¹H-NMR shifts of the commercial OAPS (Table 20) are 0.16 ppm and 0.18 ppm, respectively. If the amine

hydrogens are removed from the calculations, the aromatic hydrogens MAE and RMSE are 0.18 ppm and 0.19 ppm. The MAE and RSME for the controlled OAPS (Table 21) are 0.39 ppm and 0.42 ppm when all hydrogens are taken into account, and 0.47 ppm for both parameters when only the aromatic hydrogens are considered.

The corresponding Gaussian-predicted MAE and RMSE for all hydrogens in the commercial OAPS (Table 22) are 0.60 ppm and 0.74 ppm, respectively. Those related to the aromatic hydrogens are 0.37 ppm and 0.38 ppm. The MAE and RMSE for the controlled OAPS (Table 23) are 0.85 ppm and 0.90 ppm, whereas they decrease to 0.69 and 0.70 after removing the amine shifts. The fact that the MAE and RMSE are close to each other indicates that the errors are not scattered.

Table 20. The predicted δ_{calc} by MestreNova and experimental δ_{exp} ¹H NMR shifts in ppm along with their absolute error e_i in ppm for each peak identified in the commercial OAPS sample (Figure 62(a))

Н-Туре	δ_{calc}	δexp	<i>e</i> _i	Identification	Isomer type
H _{C2}	6.96	6.71	0.25	J	meta
H _{C4}	6.74	6.46	0.28	L	meta
H _{C5}	7.30	7.21	0.09	н	meta
H _{C6}	7.13	6.95	0.18	I	meta
Hn	4.88	5.09	0.21	meta	meta
H _{C3}	7.08	6.84	0.24	Ν	ortho
H _{C4}	7.19	7.09	0.10	F	ortho
H _{C5}	7.43	7.32	0.11	Μ	ortho
H _{C6}	7.31	7.21	0.10	D	ortho
Hn	5.49	5.40	0.09	ortho	ortho
H _{C2}	7.25	7.09	0.16	E	para
H _{C3}	6.87	6.63	0.24	0	para
Hn	5.24	5.18	0.06	para	para

Table 21. Same as Table 20 but for the controlled OAPS sample (Figure 62(b)).

Н-Туре	δ_{calc}	δ _{exp}	$ e_i $	Identification	Isomer type
H _{C2}	6.96	6.45	0.51	J	meta
H _{C4}	6.74	6.23	0.51	L	meta
H _{C5}	7.30	6.93	0.37	н	meta
H _{C6}	7.13	6.70	0.43	I	meta
Hn	4.88	4.78	0.10	meta	meta

H _{C2}	7.25	6.75	0.50	E	para
Нсз	6.87	6.37	0.50	0	para
Hn	5.25	5.06	0.19	para	para

Table 22. The δ_{calc} by Gaussian at the HCTH407/6-311+G(2d,p) level of theory and δ_{exp} ¹H NMR shifts in ppm along with their e_i in ppm for each peak identified in the commercial OAPS sample (Figure 63(a)).

Н-Туре	δ_{calc}	δexp	<i>e</i> _i	Identification	Isomer type
H _{C2}	7.44	7.09	0.35	J	meta
H _{C4}	7.04	6.72	0.32	L	meta
H _{C5}	7.39	6.96	0.43	н	meta
H _{C6}	7.25	6.84	0.41	I	meta
Hn	3.48	5.03	1.55	meta	meta
Нсз	6.88	6.64	0.24	Ν	ortho
H _{C4}	7.46	7.09	0.37	F	ortho
H _{C5}	6.90	6.64	0.26	Μ	ortho
H _{C6}	7.89	7.32	0.57	D	ortho
Hn	4.43	5.40	0.97	ortho	ortho
H _{C2}	7.62	7.21	0.41	E	para
H _{C3}	6.90	6.46	0.44	0	para
Hn	3.69	5.18	1.49	para	para

Table 23. Same as Table 22 but for the controlled OAPS sample (Figure 63(b))

Н-Туре	δ_{calc}	δ_{exp}	$ e_i $	Identification	Isomer type
H _{C2}	7.44	6.75	0.69	J	meta
H _{C4}	7.04	6.37	0.67	L	meta
H _{C5}	7.39	6.70	0.69	н	meta
H _{C6}	7.25	6.45	0.80	I	meta
Hn	3.48	4.78	1.30	meta	meta
H _{C2}	7.62	6.93	0.69	E	para
Нсз	6.86	6.23	0.63	Ο	para
Hn	3.69	5.06	1.37	para	para

Figures 64 and 65 illustrate the linear correlations for the commercial OAPS between the δ_{exp} and the δ_{calc} for the ¹H-NMR shifts predicted by both MestRenova (Table 20) and Gaussian (Table 22), respectively while taking into account either (a) all hydrogens or (b) only the aromatic hydrogens. All cases are well correlated (R > 0.97) with the slopes being slightly different depending on the prediction method used and on whether the amine ¹H-NMR shifts are included or excluded.

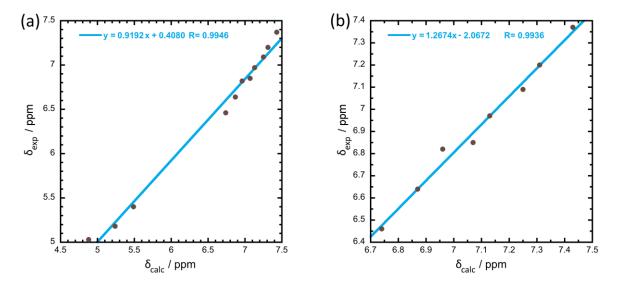


Figure 64. Linear correlation between the MestreNova-predicted δ_{calc} and the δ_{exp} ¹H-NMR shifts taking into account either (a) all hydrogens or (b) only the aromatic hydrogens in the commercial OAPS.

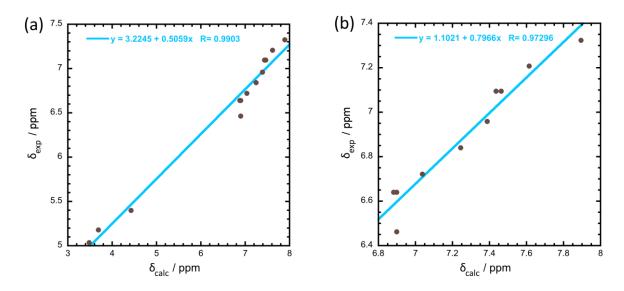


Figure 65. Same as Figure 64 but for the Gaussian-predicted shifts at the HCTH407/6-311+G(2d,p) level of theory.

As before, the scaled predicted shifts $\delta_{calc-scaled}$ were obtained for the commercial OAPS from the correlations. However, in this case, both types of approaches were investigated, i.e. either while including all hydrogens (Figures 64(a) and 65(a)) or the aromatic hydrogens only (Figures 64(b) and 65(b)). The latter is expected to be more precise than the former due to the problems related to the amine predictions.

Chapter 4

The MestRenova-predicted $\delta_{calc-scaled}$ are reported in Tables 24-25. Compared to the unscaled predictions, the MAE and RMSE fall by 50% to 0.08 ppm and 0.09 ppm, respectively, for all hydrogens (Table 24). Ignoring the amine shifts (Table 25) further reduces the MAE and RMSE to 0.03 ppm.

Table 24. The scaled MestreNova-predicted $\delta_{\text{calc-scaled}}$ and experimental δ_{exp} ¹H NMR shifts in ppm along with their absolute error e_i in ppm for each hydrogen peak identified in the commercial OAPS (Figure 64(a)).

Н-Туре	$\delta_{calc-scaled}$	δexp	$ e_i $	Identification	Isomer type
H _{C2}	6.81	6.71	0.10	J	meta
H _{C4}	6.60	6.46	0.14	L	meta
H _{C5}	7.12	7.21	0.09	н	meta
H _{C6}	6.96	6.95	0.01	I	meta
Hn	4.89	5.09	0.20	meta	meta
Нсз	6.92	6.84	0.08	Ν	ortho
H _{C4}	7.02	7.09	0.07	F	ortho
H _{C5}	7.24	7.32	0.08	Μ	ortho
H _{C6}	7.13	7.21	0.08	D	ortho
Hn	5.45	5.40	0.05	ortho	ortho
Hc2	7.07	7.09	0.02	E	para
Нсз	6.72	6.63	0.09	0	para
Hn	5.22	5.18	0.04	para	para

 Table 25. Same as Table 24 but for each aromatic hydrogen peak identified in the commercial OAPS sample

 (Figure 64(b)).

Н-Туре	$\delta_{\text{calc-scaled}}$	δ_{exp}	$ e_i $	Identification	Isomer type
H _{C2}	6.75	6.71	0.04	J	meta
H _{C4}	6.48	6.46	0.02	L	meta
H _{C5}	7.18	7.21	0.03	н	meta
H _{C6}	6.97	6.95	0.02	I	meta
H _{C3}	6.91	6.84	0.08	Ν	ortho
H _{C4}	7.05	7.09	0.02	F	ortho
H _{C5}	7.35	7.32	0.03	Μ	ortho
H _{C6}	7.20	7.21	0.01	D	ortho
H _{C2}	7.12	7.09	0.03	E	para
Нсз	6.64	6.63	0.01	ο	para

The corresponding Gaussian-predicted $\delta_{calc-scaled}$ are reported in Tables 26-27. The full MAE and RMSE fall to 0.09 ppm and 0.10 ppm (Table 26), whereas ignoring the amine hydrogen shifts reduces them to 0.05 ppm and 0.06 ppm, respectively (Table 27). Such correlations can thus be used to achieve more accurate ¹H-NMR shifts by either prediction method.

Н-Туре	$\delta_{\text{calc-scaled}}$	δexp	$ e_i $	Identification	Isomer type
Hc2	6.99	7.09	0.10	J	meta
Hc4	6.78	6.72	0.06	L	meta
H _{C5}	6.96	6.96	0.00	н	meta
H _{C6}	6.89	6.84	0.05	I	meta
Hn	4.99	5.03	0.04	meta	meta
H _{C3}	6.71	6.64	0.07	Ν	ortho
Hc4	7.00	7.09	0.09	F	ortho
H _{C5}	6.71	6.64	0.07	м	ortho
H _{C6}	7.22	7.32	0.10	D	ortho
Hn	5.46	5.40	0.06	ortho	ortho
H _{C2}	7.08	7.21	0.13	E	para
Нсз	6.71	6.46	0.25	Ο	para
Hn	5.09	5.18	0.09	para	para

Table 26. The scaled Gaussian-predicted $\delta_{\text{calc-scaled}}$ and experimental δ_{exp} ¹H NMR shifts in ppm along with their absolute error e_i in ppm for each hydrogen peak identified in the commercial OAPS (Figure 65(a)).

 Table 27. Same as Table 26 but for each aromatic hydrogen peak identified in the commercial OAPS sample

 (Figure 65(b)).

Н-Туре	δ_{calc} -scaled	δexp	$ e_i $	Identification	Isomer type
H _{C2}	7.03	7.09	0.06	J	meta
H _{C4}	6.71	6.72	0.01	L	meta
H _{C5}	6.99	6.96	0.03	н	meta
H _{C6}	6.87	6.84	0.03	I	meta
H _{C3}	6.58	6.64	0.06	Ν	ortho
H _{C4}	7.05	7.09	0.04	F	ortho
H _{C5}	6.60	6.64	0.04	м	ortho
H _{C6}	7.39	7.32	0.07	D	ortho
H _{C2}	7.17	7.21	0.04	Е	para
H _{C3}	6.60	6.46	0.14	0	para

Figures 66 for MestreNova (Table 21) and 67 for Gaussian (Table 23) both show that the amine hydrogen ¹H-NMR shifts also change the slope of the linear correlations in the controlled OAPS sample with (a) including for all hydrogens and (b) including only for the aromatic hydrogens. The R is more than 0.97 in all cases, which is still good.

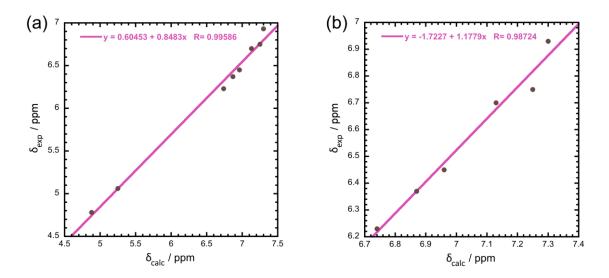


Figure 66. Linear correlation between the MestreNova-predicted δ_{calc} and the δ_{exp} ¹H-NMR shifts taking into account either (a) all hydrogens or (b) only the aromatic hydrogens in the controlled OAPS.

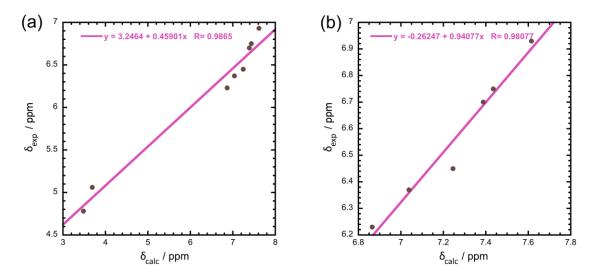


Figure 67. Same as Figure 65 but for the Gaussian-predicted shifts at the HCTH407/6-311+G(2d,p) level of theory.

The scaled predicted shifts $\delta_{calc-scaled}$ for the controlled OAPS were again calculated by using the correlations for either all hydrogens (Figures 66(a) and 67(a)) or only for the aromatic hydrogens (Figures 66(b) and 67(b)).

The MestRenova-predicted $\delta_{calc-scaled}$ are reported in Tables 28-29. Compared to the unscaled predictions, the full MAE and RMSE are reduced to 0.05 ppm and 0.07 ppm, respectively (Table 28), and to 0.03 ppm and 0.04 ppm after subtracting the amine shifts (Table 29).

Table 28. The scaled MestreNova-predicted $\delta_{\text{calc-scaled}}$ and experimental δ_{exp} ¹H NMR shifts in ppm along with their absolute error e_i in ppm for each hydrogen peak identified in the controlled OAPS (Figure 66(a)).

Н-Туре	$\delta_{\text{calc-scaled}}$	δ_{exp}	$ e_i $	Identification	Isomer type
H _{C2}	6.51	6.45	0.06	J	meta
H _{C4}	6.32	6.23	0.09	L	meta
Hcs	6.80	6.93	0.13	н	meta
H _{C6}	6.65	6.7	0.05	I	meta
Hn	4.74	4.78	0.04	meta	meta
H _{C2}	6.75	6.75	0.00	E	para
Нсз	6.43	6.37	0.06	Ο	para
Hn	5.06	5.06	0.00	para	para

 Table 29. Same as Table 28 but for each aromatic hydrogen peak identified in the controlled OAPS sample

 (Figure 66(b)).

Н-Туре	δ_{calc} -scaled	δ_{exp}	$ e_i $	Identification	Isomer type
H _{C2}	6.48	6.45	0.03	J	meta
H _{C4}	6.22	6.23	0.01	L	meta
H _{C5}	6.88	6.93	0.05	н	meta
H _{C6}	6.68	6.7	0.02	I	meta
H _{C2}	6.82	6.75	0.07	E	para
H _{C3}	6.37	6.37	0.00	Ο	para

The corresponding Gaussian-predicted $\delta_{calc-scaled}$ are reported in Tables 30-31. The full MAE and RMSE both fall to 0.12 ppm (Table 30), while ignoring the amine hydrogens further reduce the MAE to 0.08 ppm (Table 31). This improvement in accuracy when the amine hydrogens are omitted shows that precise values for the OAPS aromatic hydrogens shifts can be obtained from both prediction methods when associated with the appropriate scaling.

Н-Туре	$\delta_{ ext{calc-scaled}}$	δ_{exp}	$ e_i $	Identification	Isomer type
H _{C2}	6.66	6.75	0.09	J	meta
H _{C4}	6.48	6.37	0.11	L	meta
H _{C5}	6.64	6.70	0.06	н	meta
H _{C6}	6.57	6.45	0.12	I	meta
Hn	4.84	4.78	0.06	meta	meta
H _{C2}	6.74	6.93	0.19	E	para
Нсз	6.40	6.23	0.17	0	para
Hn	4.94	5.06	0.12	para	para

Table 30. The scaled Gaussian-predicted $\delta_{\text{calc-scaled}}$ and experimental δ_{exp} ¹H NMR shifts in ppm along with their absolute error e_i in ppm for each hydrogen peak identified in the controlled OAPS (Figure 67(a)).

 Table 31. Same as Table 30 but for each aromatic hydrogen peak identified in the controlled OAPS sample

 (Figure 67(b)).

Н-Туре	$\delta_{\text{calc-scaled}}$	δ_{exp}	$ e_i $	Identification	Isomer type
H _{C2}	6.73	6.75	0.02	J	meta
H _{C4}	6.36	6.37	0.01	L	meta
H _{C5}	6.55	6.70	0.15	н	meta
H _{C6}	6.69	6.45	0.24	I	meta
H _{C2}	6.90	6.93	0.03	Е	para
Нсз	6.20	6.23	0.03	0	para

Despite the slight displacements between predicted and experimental peaks, both MestreNova and Gaussian at the HCTH407/6-311+G(2d,p) level of the theory seem to be able to provide results that are compatible with the experimental ¹H-NMR OAPS shifts. Such results can even be quantitatively improved using the scaling factors displayed in Figures 64-67, especially when the amine hydrogen predictions are omitted.

However, it remains unclear why the relative order of the predicted aromatic hydrogens peaks differ, *i.e.* the identification of each peak as belonging to a specific isomer depends on the prediction method used. Since it was not possible to differentiate both methods using the 1D-NMR experimental spectra, it was attempted to predict the 2D-NMR spectra with each of them.

Chapter 4

b.4.4) 2D-NMR predictions

The ¹³C-NMR $\delta_{calc-scaled}$ shifts (Section b.4.2.) for both types of OAPS were combined with their respective aromatic ¹H-NMR $\delta_{calc-scaled}$ shifts (Section b.4.3.) in Figure 68 to predict 2D-NMR spectra comparable to the available experimental HSQC spectra (Figure 45). (a) and (b) show the predictions by MestreNova and Gaussian, respectively, for the commercial OAPS, while (c) and (d) are the corresponding predictions for the controlled OAPS.

In the predicted 2D-NMR spectrum of the commercial OAPS with MestreNova (Figure 68(a)), some peaks can easily be assigned. The G, P and K shifts on the ¹³C-NMR axis are related to the *meta*, *ortho*, and *para* aromatic carbons connected to the silicons in the OAPS cages. Similarly, the C, A and B shifts correspond to the carbons connected to the amine nitrogens. The other ¹³C-NMR shifts are connected to the aromatic hydrogens and their cross-points can be compared to the experimental HSQC spectrum. However, it is clear that the predicted cross-points are scattered and far from the experimental spectrum (for example M and N).

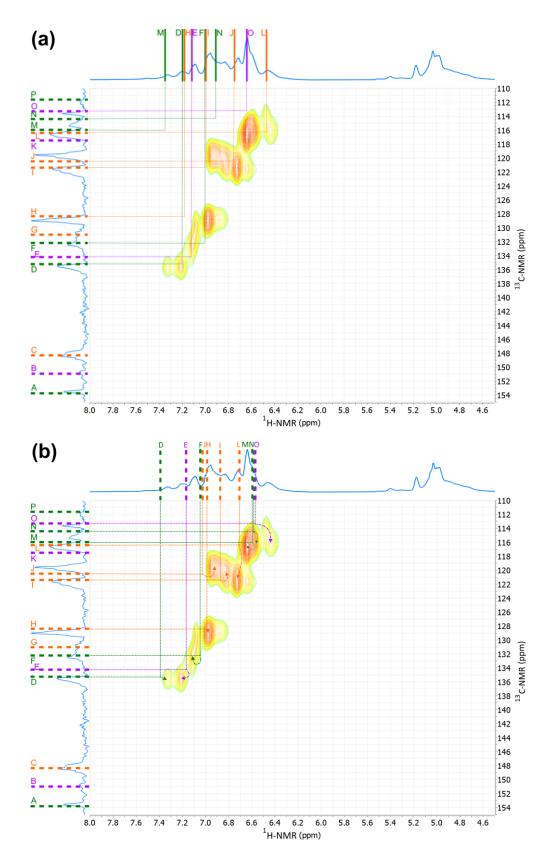
The situation improves significantly for the predicted 2D-NMR spectrum of the commercial OAPS with Gaussian (Figure 68(b)). The G, P, K, C, A and B shifts on the ¹³C-NMR axis can again be assigned to the carbons linked either to the OAPS cage silicons or to the amine nitrogens. As far as the cross-points are concerned, the *ortho* D, *para* E, *ortho* F and *meta* H shifts correspond very well to experimental contour lines. The other cross-points are slightly more displaced because of the differences between predicted and experimental peaks. The *meta* L peak seems to be the furthest from its corresponding contour line, which could be due to an underestimation in the ¹³C-NMR spectrum. But the Gaussian-predicted cross-points still reproduce rather well the ten contour points of the experimental HSQC spectrum, i.e. the prediction is much better than for MestreNova in (a).

In the predicted 2D-NMR spectrum of the controlled OAPS with MestreNova (Figure 68(c)), the G, K, B and C shifts on the ¹³C-NMR axis are related to the carbons connected either to the cage silicons or to the amine nitrogens. For the carbons connected to aromatic hydrogens, the black rectangles suggest that the two sharpest ¹³C-NMR peaks, which correspond to large contour cross-points, can be assigned to the *para* isomer. On the other hand, the H, I, J and L shifts can be assigned to the *meta* isomer. However, as for (a), the predicted shifts for the aromatic hydrogens lead to cross-points that are displaced with respect

Chapter 4

to the experimental signals, even if there is a better agreement than for the commercial OAPS.

The situation for the predicted 2D-NMR spectrum of the controlled OAPS improves once again with Gaussian (Figure 68(d)). The assignment of each peak to a specific isomer is similar to MestreNova but the cross-points are much closer to the experimental adjacent contour point. The black rectangles confirm that the sharp ¹³C-NMR peaks correspond to the *para* isomer.



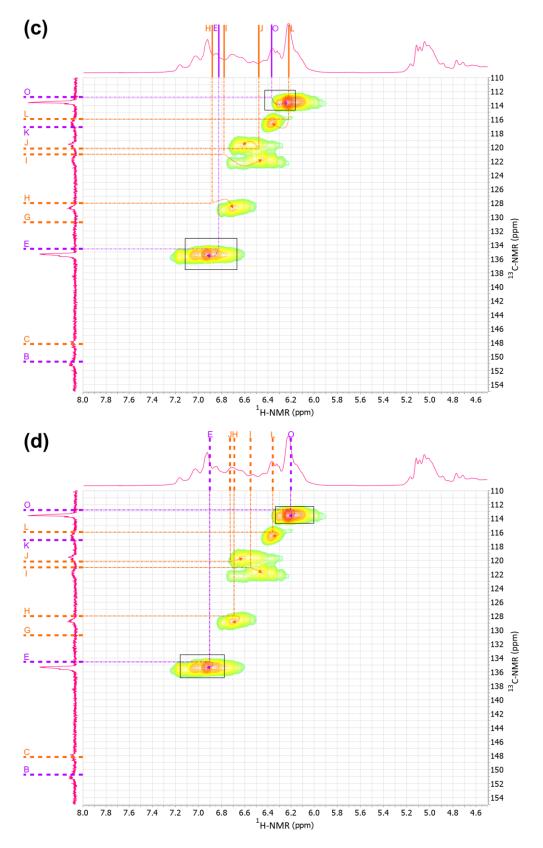


Figure 68. The experimental HSQC spectrum for the commercial (blue) and controlled (pink) OAPS compared to the $\delta_{calc-scaled}$ for the ¹³C-NMR shifts (y-axis) and the ¹H-NMR shifts (x-axis) predicted by (a and c) MestreNova and (b and d) Gaussian at the HCTH407/6-311+G(2d,p) level of theory. The colours for the isomers are orange for *meta*, green for *ortho* and purple for *para*.

While Section b.4.3. showed that it was difficult to unambiguously assess the quality of the MestreNova and Gaussian predictions for ¹H-NMR using the 1D spectra only, the 2D-NMR predictions prove that the DFT calculations are much better at predicting accurate chemical shifts for ¹H-NMR than MestreNova.

Consequently, the final isomer assignment for each peak is that predicted by Gaussian at the HCTH407/6-311+G(2d,p) level of theory. It is summarized in Table 32.

Table 32. The Gaussian-predicted $\delta_{calc-scaled}$ compared to the δ_{exp} for the ¹H NMR and ¹³C-NMR chemical shifts in ppm for each aromatic carbon and hydrogen peaks identified in the commercial and controlled OAPS samples and the assignment of each signal to a specific isomer (Tables 18, 19, 27, 31).

H and C Type	Commercial δ _{calc-scaled}	Commercial ð _{exp}	Controlled δ _{calc-scaled}	Controlled δ _{exp}	Identification	Isomer type
C ₁	131.09	130.70	130.73	129.28	G	meta
H _{C2} - C ₂	7.03-120.61	7.09-119.55	6.73-120.17	6.75-119.74	J	meta
C ₃	148.41	148.34	148.19	147.71	С	meta
H _{C4} - C ₄	6.71-116.39	6.72-116.64	6.36-115.91	6.37-116.13	L	meta
H _{C5} - C ₅	6.99-128.36	6.96-128.97	6.55-127.98	6.70-128.20	н	meta
Hc6 - C6	6.87-121.39	6.84-121.60	6.69-120.95	6.45-121.51	I	meta
C 1	111.67	111.78	-	-	Р	ortho
C ₂	153.89	153.53	-	-	Α	ortho
H _{C3} - C ₃	6.58-114.42	6.64-115.28	-	-	Ν	ortho
H _{C4} - C ₄	7.05-132.27	7.09-132.45	-	-	F	ortho
H _{C5} - C ₅	6.60-115.98	6.64-115.28	-	-	Μ	ortho
Hc6 - C6	7.39-135.32	7.32-135.54	-	-	D	ortho
C 1	117.56	116.64	117.09	116.13	К	para
H _{C2} - C ₂	7.17-134.87	7.21-135.54	6.90-134.54	6.93-135.35	Е	para
H _{C3} - C ₃	6.60-113.26	6.46-113.60	6.20-112.76	6.23-113.59	0	para
C 4	150.94	151.00	150.75	151.36	В	para

V. Hybrid membranes based on OAPS

1) Fabrication of polyOAPS-imide membranes

Following the extensive characterization of OAPS on its own, the commercial OAPS produced by Gelest and the 6FDA dianhydride (4,4'-(hexafluoroisopropylidene) diphthalic anhydride, Figure 7) produced by Sigma-Aldrich were selected as the monomers to fabricate new hybrid polyOAPS-imide membranes aimed at being more resistant to high temperatures and pressures than the initial polyPOSS-imides.

The fabrication procedure was based on the previous polyPOSS-imide work (Figure 12),⁴ and its three steps are schematically shown for OAPS in Figure 69. It is expected that the hypercrosslinked structures will be somewhat dependent on the OAPS isomer.²⁰⁸⁻²⁰⁹ We note that an OAPS cage could potentially carry several different isomers for the aminophenyl arms but as these are not able to be distinguished experimentally, the isomers are not specified in the schemes.

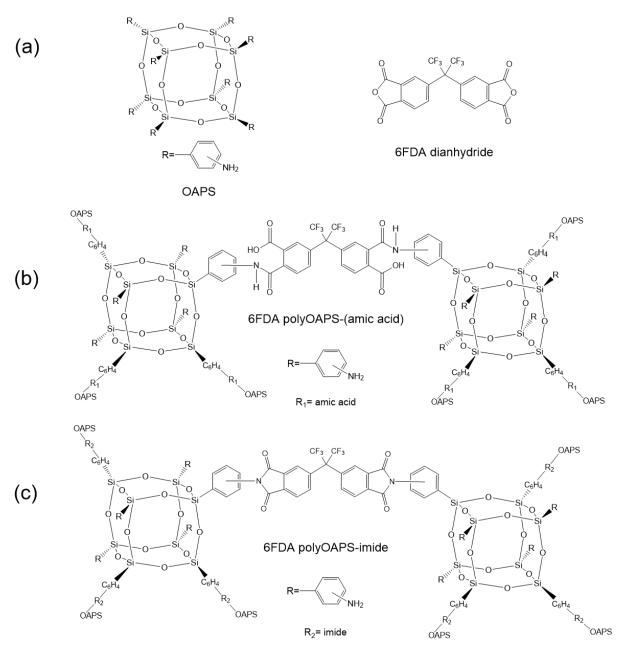


Figure 69. Chemical structures of (a) the OAPS and the 6FDA dianhydride monomers, (b) the interfacial polycondensation step leading to a 6FDA polyOAPS-(amic acid) network with an average connectivity of 4, and (c) the imidization step leading to the final 6FDA polyOAPS-imide network.

As for the polyPOSS-imides, 6FDA-based polyOAPS-imide membranes were produced through interfacial polycondensation between an organic phase containing the 6FDA dianhydride in toluene, and an aqueous phase containing OAPS in a mixture of DMSO/water with a ratio of 3/2 v/v. Solutions of 0.1% w/v 6FDA in toluene and 2.5% w/v OAPS in DMSO were first prepared separately. OAPS cannot be dissolved in water but, as explained in Section III, DMSO is soluble with toluene. As such, water was added in a second step to the OAPS/DMSO solution in order to reach a 3/2 DMSO/water ratio and create two separate phases between the 6FDA-containing and OAPS-containing solutions.

 α -alumina disks were chosen as the main mechanical supports for the membranes. In order to add a layer with smaller pores, a thin layer of γ -alumina (3 µm) was deposited on each α -alumina support by dip coating (Figure 70). Since the formation of the γ -alumina layer is very sensitive to dust, the coating procedure was performed in a cleanroom. A solution of Al(OH)Ac₂/polyvinyl alcohol (PVA) was first prepared. Then, the support was contacted for a few seconds to get completely coated with the alumina solution. The coated support was later displaced into a high-temperature (HT) oven and heated up to 650°C to undergo calcination.

Figure 70 summarizes the whole membrane fabrication procedure. Following the modification of the alumina support:

- The support was fixed in a Buchner funnel and the OAPS solution in DMSO/water was poured on it. After 20 minutes, the excess solution was poured out, and a rubber roller was rolled on the surface of the support to remove the droplets. A very thin layer of the OAPS solution remained on the surface of the support.
- The 6FDA solution in toluene was then added to start the polycondensation reaction. A new thin layer of hypercrosslinked polyOAPS-(amic acid) network (Figure 69(b)) was immediately formed. In this reaction, the limiting reactant is OAPS and the excessive one is 6FDA. After 13 minutes, the 6FDA solution was poured out and the support was taken out of the Buchner funnel. It was left in a vacuum oven for 48 hours at 50°C in order to evaporate the solvents.
- The polyOAPS-(amic acid) membrane on its support was subsequently transferred into the HT oven to perform the imidization, i.e. the transition from the amic acid to the imide form for 6FDA (Figure 69(c)). The step-by-step annealing program was as follows with a ramp rate of 5 °C/min to prevent any shrinkage of the membrane: 3 hours up to 100 °C, 3 hours up to 150 °C, 3 hours up to 200 °C, 2 hours up to 250 °C, 2 hours up to 300 °C. It was finally cooled for 3 hours to get back to ambient temperature. As can be seen in Figure 70, the final 6FDA polyOAPS-imide membrane is brown.

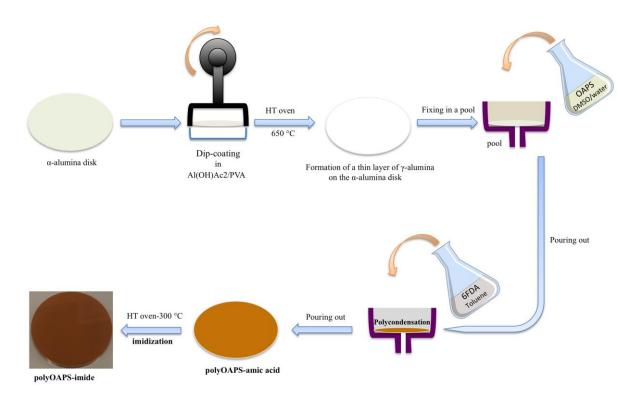


Figure 70. The fabrication procedure of an hypercrosslinked 6FDA-based polyOAPS-imide membrane.

2) TGA analysis of the membranes

Thermal gravimetric analyses (TGA) were carried out using an STA 449 F3 Jupiter (NETZSCH). Following the interfacial polycondensation, 15 mg of free-standing polyOAPS- (amic acid) membrane was transferred inside an aluminium pan and was warmed up under a nitrogen atmosphere (stream of 50 ml min⁻¹) with a heating rate of 10 °C/min. As shown in Figure 71, a first weight loss occurs at 150-300°C, indicating the imidization step and the production of water. The second weight loss at ~400°C shows the start of the degradation process for the organic part. The sample reaches an almost constant mass at ~650 °C, and the weight loss continues up to 1100 °C at a gentle slope related to the degradation of the inorganic part.

Simulation studies on the thermomechanical characteristics of the new polyOAPSimide membranes at 400°C show that the *ortho* isomer provides a better resistance for the 6FDA polyOAPS-imide membranes than the *meta* and *para* isomers.²⁰⁸ However, in all cases, the polyOAPS-imide membranes remain thermally superior to the polyPOSS-imides membranes, which degrade experimentally before 350°C.³¹ Following imidization, polyOAPS-imide networks can thus tolerate very high temperatures, which is crucial for the separation of hot gases.

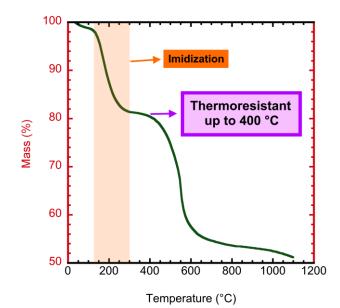


Figure 71. TGA analysis of a 6FDA polyOAPS-(amic acid) membrane

3) IR spectroscopy characterization of the membranes

The ATR-IR Perkin Elmer Paragon 1000 PC was used to compare the IR spectra for OAPS, 6FDA polyOAPS-(amic acid), and 6FDA polyOAPS-imide (Figure 72) in order to follow the various stages of the preparation procedure. (a) shows the whole IR spectra while (b) specifically displays the 1000-2000 cm⁻¹ range for clarity.

The Si-O-Si band, which can be observed at 1100 cm⁻¹ in all three spectra, is due to the OAPS cages. The polyOAPS-(amic acid) presents two typical amide bands at 1590 and 1650 cm⁻¹, corresponding to the N-H bending and the C=O stretching, respectively as well as a small shoulder at 1715 cm⁻¹ indicating the presence of the carboxylic acid functional group. Following the thermal treatment, the aforementioned peaks should disappear due to the imidization process. Indeed, the amic acid bands are no anymore clearly observable in the polyOAPS-imide spectrum, thus showing that imidization is efficient. However, small bands corresponding to the N-H bending can still be observed for the polyOAPS-imide. This is probably due to unreacted NH₂ groups on the OAPS cage. This is not unexpected as it is difficult for steric reasons to have all eight arms of the inorganic precursor reacting with a dianhydride (Figure 69), and as such, some of the initial amine groups remain present in the final network. The two new bands at 1720 and 1780 cm⁻¹ in the polyOAPS-imide spectrum can be attributed to the imide C=O symmetric and asymmetric stretching modes, respectively. The band at 1270 cm⁻¹ is the C-F stretching of the 6FDA organic bridge, which is present in both the polyOAPS-(amic acid) and polyOAPS-imide but is absent in the inorganic OAPS precursor.

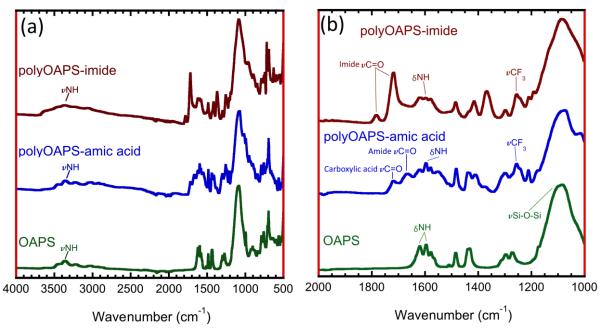


Figure 72. IR spectra of the OAPS, 6FDA polyOAPS-(amic acid), and 6FDA polyOAPS-imide.

Figure 72 thus confirms that the final product is indeed basically a polyOAPS-imide membrane.

4) Permeance and gas selectivity of the new membranes

An automatic single-gas permeation set-up was used in a dead-end mode at a transmembrane pressure of 2 bar and at five different temperatures (50, 100, 150, 200, and 250 °C) in order to measure the permeance of the new membranes. He, H₂, N₂, CO₂ and CH₄ were tested as pure feeds, and the permeance results were averaged over three membranes each. In the following Figures (Figs. 73-75), results related to the new 6FDA polyOAPSimide membranes will be presented on (a) the left while those related to the initial 6FDA polyPOSS-imides for the same temperature range will be displayed on (b) the right for comparison.

Figure 73 shows the permeances as a function of the kinetic diameter for the different gases. The polyOAPS-imide membranes based on 6FDA are able to maintain their

permeation characteristics at temperatures at least up to 250 °C (a). This is much higher than for conventional polymers which suffer from a performance loss at elevated temperatures. As expected for glassy polymers, the permeance tends to decrease with the increase in the gas kinetic diameter.^{4, 210} Compared to the polyPOSS-imides (b), the gas permeances in polyOAPS-imide membranes are a bit lower for He and H₂, similar for CO₂ and generally larger for N₂ and CH₄. The latter is probably related to the larger amount of free volume in the OAPS-based membranes, which allows for larger penetrants to permeate better. However, it should also be noted that permeances are not normalized to the membrane thicknesses. These are expected to be close but slight differences could affect the results.

For all gases except for CO₂, the permeance increases with the temperature in the order $CH_4 > N_2 > H_2 \approx$ He for both polyOAPS-imide and polyPOSS-imide matrices, indicating that the diffusion of the larger penetrants (CH₄ and N₂) is thermally improved to a greater extent than that of the smaller ones. The differences in free volumes between both types of membranes, i.e. polyOAPS-imides >> polyPOSS-imides,²⁰⁸⁻²⁰⁹ lead to more differences for the latter as a function of the temperature. Indeed, polyPOSS-imides are able to provide significantly more diffuse jumps for the larger penetrants at high temperatures with respect to those at ambient temperatures.²¹⁰ On the other hand, diffusion in polyOAPS-imides is easier both at ambient and high temperatures.

The differences as a function of temperature are less marked for CO₂ because of its higher solubility at low temperatures. The affinity toward CO₂ results from the trifluoromethyl and the imide groups, which confirms that the gas molecules diffuse through the organic part of the network.^{4, 209, 211} At higher temperatures, the CO₂ diffusivity increases while its solubility decreases, thus counteracting the influence of temperature.⁴ The selectivity of CO₂ over the other gases thus decreases as a function of temperature. Both types of membranes show close permeances for CO₂, which confirms that the 6FDA organic linker is more important than the inorganic precursor for the transport of this penetrant.

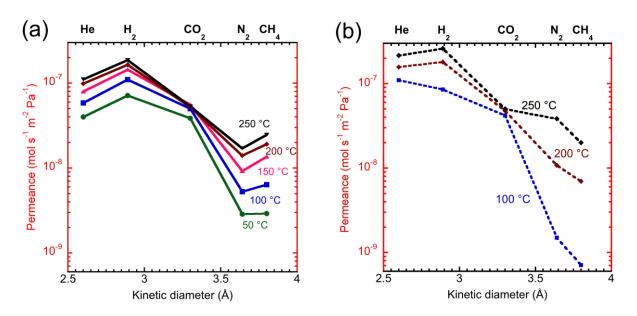


Figure 73. Gas permeances of (a) the new polyOAPS-imides and (b) the polyPOSS-imides at 2 bar.⁴

Figure 74 compares the Arrhenius plots of the permeances in both types of membranes on a logarithmic scale as a function of $R^{-1}T^{-1}$ (lower *x*-axis). The temperature is also provided in °C on the top *x*-axis for clarity. They show that the gas permeance is thermally activated. The activation energies were obtained from nonlinear fits of the temperature-dependent permeation data. They decrease in the order $CH_4 > N_2 > He \approx H_2 > CO_2$ for the polyOAPS-imides (a). It is the same as for the initial polyPOSS-imides,⁴ except for H₂ which is slightly larger than for He (b). The linear trends in the Arrhenius plots indicate that the activation energies of the sorption and diffusion are not severely affected by the temperature.²¹⁰ This is unlike the chain-rearrangement-induced permeability loss in conventional polymeric membranes at high temperatures.^{4, 212-213}

Activation energies are associated both with the penetrant size and with the molecular nature of the membrane. The larger penetrants generally exhibit larger activation energies in polymer membranes (except for CO₂ due to its high solubility).²¹⁰ A high activation energy results from a great resistance within the polymer matrix for the penetrant to perform a diffusive jump.²¹⁰ On the other hand, lower activation energies are associated with fewer energy barriers for gas permeation. As noted above, molecular motions generated at higher temperatures create larger gaps between the chains, which have more impact on the larger penetrants since they encounter more severe diffusive resistance at ambient temperatures.²¹⁰ The effect of temperature on the CO₂ permeance is low, and as a result, its activation energy is low in both types of membranes. For all other gases, the activation energies are higher in the polyPOSS-imide membranes (b). This is especially striking for CH₄ and N₂, where they

are more than twice the activation energies in the polyOAPS-imides (a). According to molecular modelling screening investigations,²⁰⁹ polyOAPS-imides have open structures resulting from the high free volume generated by the intrinsic planarity and rigidity of the phenyl rings, and as such a high sorption capacity. PolyPOSS-imides are much more compact, thus providing more efficient barriers at low temperatures.

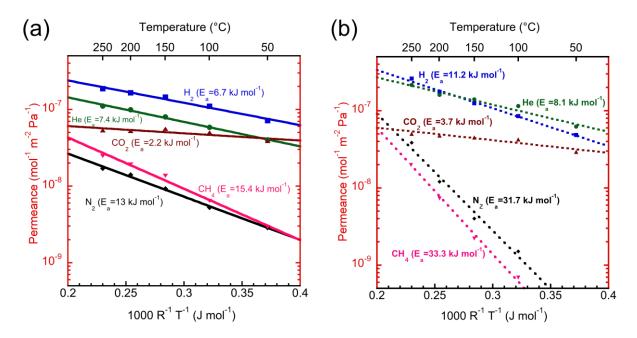


Figure 74. Arrhenius plots of gas permeances for He, H₂, CO₂, N₂ and CH₄ for (a) the new polyOAPS-imides and (b) the polyPOSS-imides at 2 bar.⁴

Figure 75 compares the ideal permselectivity of the gas pairs under study, i.e. H₂/CO₂, CO₂/CH₄, H₂/CH₄, CO₂/N₂, He/CO₂, He/CH₄ and H₂/N₂ as a function of temperature in both types of membranes. Regardless of the more open structures in the new polyOAPS-imides (a), which decrease the selectivities with respect to the more gas-sieving polyPOSS-imides (b), the gas selectivities are still satisfactorily maintained at high temperatures due to the rigidity of the hypercrosslinked polyOAPS-imide networks. Indeed, the molecular simulations²⁰⁹ confirm the superior mechanical stabilities of polyOAPS-imides over polyPOSS-imides, while molecular motions still allow for the diffusion of the gas molecules.⁴

Most of the selectivities over both membranes decrease as a function of the temperature due to the different gas activation energies. They vary more dramatically for the polyPOSS-imide membranes (b), because of their higher activation energies. On the other hand, the He/CO₂ and H₂/CO₂ selectivities increase with temperature due to the higher activation energies of He and H₂ over CO₂. The polyOAPS-imide membranes are able to

preserve gas selectivities of ~10 for H_2/CH_4 and H_2/N_2 up to 250°C which shows their interest in gas separations at high temperatures.

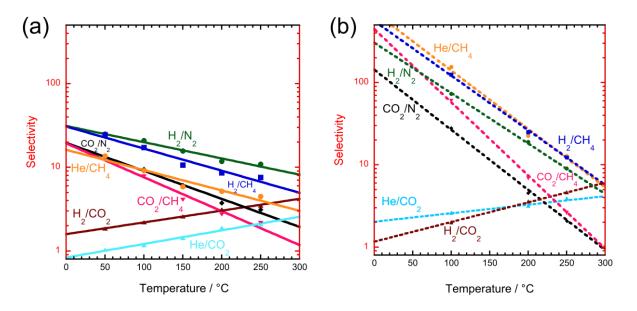


Figure 75. Permselectivities for (a) the new polyOAPS-imides and (b) the polyPOSS-imides at 2 bar.⁴

The permeation studies show that the polyPOSS-imide membranes have higher selectivities at temperatures below 300°C (b). But the trendlines suggest that the scenario can be reversed at higher temperatures than 300°C (a). While polyPOSS-imides will eventually degrade below 350°C, polyOAPS-imides will remain active under such harsh conditions, particularly for H₂/CO₂, He/CO₂ as well as H₂/N₂ to some extent.

This can be interesting as H₂ separation is a critical step in hydrocarbon fuel processing.²¹⁴ Helium is also important as it is used as an inert gas in welding, a carrier gas in scientific equipment and so on.²¹⁵⁻²¹⁶

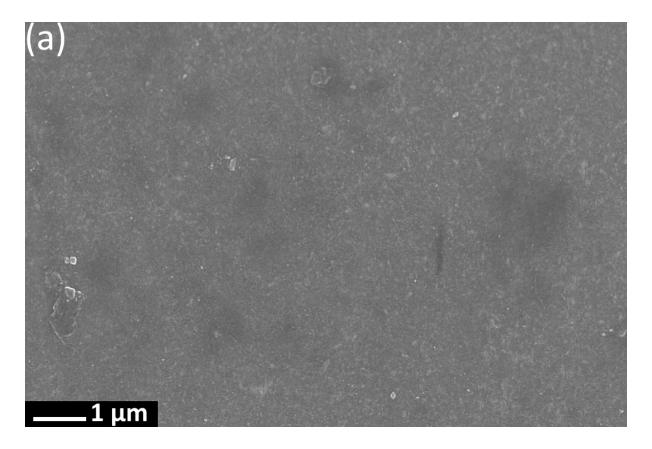
5) Morphology of the new polyOAPS-imide membrane

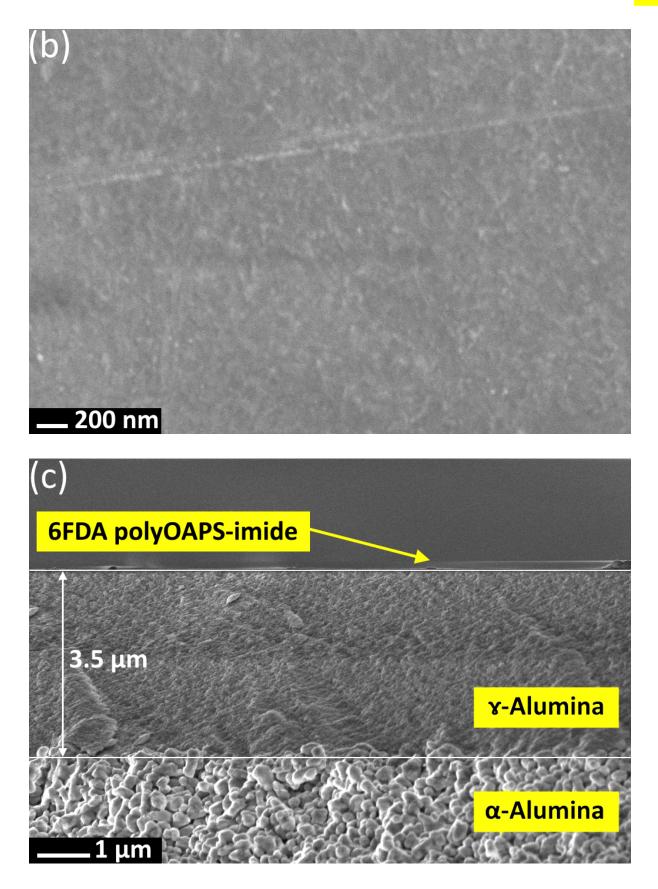
Figure 76 shows the surfaces on the scales of (a) 1 μ m and (b) 200 nm as well as the cross-sections on the scales of (c) 1 μ m and (d) 200 nm of a 6FDA polyOAPS-imide membrane studied by Field Emission Scanning Electron Microscopy (FE-SEM) with a ZEISS Gemini FE-SEM 500 apparatus.

Figures 76(a-b) confirm that the polyOAPS-imide surfaces are intact without any cracks a the corresponding scales. The polycondensation reaction thus results in dense layers

Chapter 5

with no pores and defects. Figures 76(c-d) show the cross-sections of the polyOAPS-imide membranes with their three layers. The lowest layer is the α -alumina with sintered grains. The middle layer is the γ -alumina with a thickness of ~3.5 µm (c). γ -alumina has a uniform morphology, indicating that the sol-gel process was properly performed. The top layer is the polyOAPS-imide membrane and its thickness can be evaluated as being ~45 nm (d). Such low thicknesses along with intact surfaces confirm the high efficiency of the membrane preparation procedure via interfacial polymerization followed by imidization.





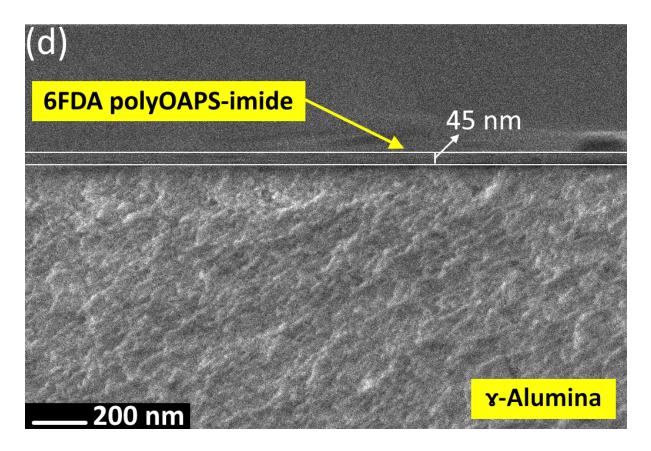


Figure 76. The surface (a and b) and cross-section (c and d) of a 6FDA polyOAPS-imide membrane at different scales.

FE-SEM can be combined with Energy Dispersive Spectroscopy (EDS) to probe the elements contained in the sample under study. Figure 77 displays the EDS spectrum of the same 6FDA polyOAPS-imide membrane characterized by FE-SEM in Figure 76. The presence of Si and N results from the OAPS inorganic precursor while F and C are related to the organic 6FDA links. O has the largest peak as it is present in the membrane and both two alumina layers. The Al peak also results from the alumina layers. There is a weak Na peak which probably corresponds to impurities.

The EDS spectrum thus confirms the presence of the appropriate atoms in the polyOAPS-imide membranes. However, it was not possible to extract an unambiguous quantitative assignment because of the small size of the C peak, which is probably related to technical problems.

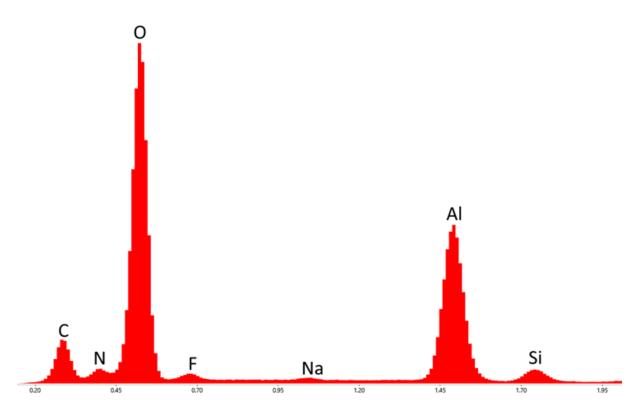


Figure 77. EDS spectrum of a 6FDA polyOAPS-imide membrane.

The morphology analyses confirm the intact structures of the polyOAPS-imide membranes, i.e. they are thin and can be selective. The intermediate γ -alumina layer forms a uniform structure, which preserves the hybrid membrane very well.

VI. Force-field development for OAPS

To investigate the molecular scale of polyOAPS-imide networks, the *gmq* program²¹⁷ was used by researchers at LEPMI, USMB, to carry out molecular modelling simulations of OAPS, imides and polyOAPS-imide networks.²⁰⁸⁻²⁰⁹ *gmq* is principally designed for molecular dynamics (MD) simulations of dense materials with periodic boundary conditions in three dimensions.²¹⁷ This program has displayed its performance in numerous papers published in the literature previously.^{30-31, 218-224}

Before generating the model networks, a specific force-field had to be developed for OAPS on its own. MD simulations are based on the forces applied to every atom, and the force-field describes the potential energy contribution of all other atoms. In *gmq*, a basic force-field includes bends, torsions, out-of-plane, van der Waals and electrostatic interactions as well as special -CH₃, -CH₂-, X-Y-H₂, -CH and linear triatomic constraints.²²⁵ Each atom interacts with the others via either bonded or non-bonded potentials and the total potential energy U_{pot} is as in Equation 12.^{10,217}

$$U_{pot} = \sum U_{bonded} + \sum U_{non-bonded}$$
(12)

The bending, torsional and out-of-plane potentials contribute to the bonded potentials, *i.e.* those involving closely-connected atoms (Equation 13):

$$\sum U_{bonded} = \sum_{\theta} U_{bend}(\theta) + \sum_{\tau} U_{tors}(\tau) + \sum_{i-sp2} U_{oop}(i)$$
(13)

where θ , τ , and *i*-sp² are the bending angles, the dihedral angles, and i the atoms with sp² hybridization respectively (Figure 78).

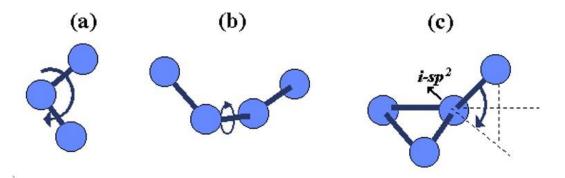


Figure 78. The bonding interactions: (a) bending, (b) torsion, and (c) out-of-plane

The bond distances b_0 are kept constant by rigid constraints in order to use a time step of 10^{-15} s and ensure equipartition of the energy (Equation 14):

$$\left| r_{ij} \right|^2 - b_0^2 = 0 \tag{14}$$

The bending potential keeps the bending angles θ defined by atoms *i-j-k* close to their equilibrium values (Equations 15 and 16):

$$U_{bend}(\theta) = \frac{1}{2}k_{\theta}(\cos\theta - \cos\theta_0)^2$$
(15)

$$\cos\theta = \frac{(r_{ij} \bullet r_{kj})}{|r_{ij}||r_{kj}|} \tag{16}$$

where k_{θ} is a constant indicating the flexibility of the angle and θ_0 is the equilibrium bending angle.

The torsional potential energy around the *j*-*k* bond in an (i, j, k, l) atom quadruplet is described as in Equations 17 and 18:

$$U_{tors}(\tau) = \sum_{n=0}^{6} a_n \cos^n \tau \tag{17}$$

$$\cos \tau = \frac{(r_{ij} \times r_{jk}) \bullet (r_{jk} \times r_{kl})}{\left| r_{ij} \times r_{jk} \right| \bullet \left| r_{jk} \times r_{kl} \right|}$$
(18)

where the a_n are the coefficients of a polynomial of order *n*. For this term, two different possibilities may occur. If the atoms separated by more than three bonds are defined as being non-bonded, Equation 17 represents the entire torsional energy. But if the atoms separated by more than two bonds are defined as being non-bonded, a part of torsional energy is described by Equation 17 while the non-bonded interactions between the end atoms of the torsion quadruplet, *i* and *l*, form the rest.

The out-of-plane potential is used to restrict the movement of a central trivalent atom i connected directly to three atoms j, k and l on a plane (Equation 19):

Chapter 6

$$U_{oop}(i) = \frac{1}{2}k_{oop}s^2\tag{19}$$

where *s* is the perpendicular distance of atom *i* from the plane of atoms *j*, *k*, and *l*, and k_{oop} is the force constant. *s* can be calculated through Equation 20:

$$s = r_{ji} \bullet \frac{(r_{jk} \times r_{jl})}{\left| r_{jk} \times r_{jl} \right|}$$
(20)

The non-bonded potentials include the Van der Waals interactions and the electrostatic interactions resulting from the charges on the atoms (Equation 21):

$$\sum U_{non-bonded} = \sum_{(i,j)nb} U_{vdw}(r) + \sum_{(i,j)nb} U_{coul}(r)$$
⁽²¹⁾

Different forms of Van der Waals (*vdw*) potentials are available in *gmq*, but the Lennard-Jones (LJ) 12-6 potential is mostly used (Equation 22):

$$U_{LJ}(|r_{ij}|) = 4\varepsilon \left(\left(\frac{\sigma}{r_{ij}}\right)^{12} - \left(\frac{\sigma}{r_{ij}}\right)^6 \right)$$
(22)

where ε is the well depth of the *LJ* potential and σ is the distance at which the potential is zero. The $(1/r_{ij})^{12}$ term describes the Pauli repulsions at short distances while $(1/r_{ij})^6$ accounting for the attractions between the atoms at longer distances.

For unlike-atom pairs, the cross-terms are calculated via standard combination rules²²⁶⁻²²⁸ (Equations 23 and 24):

$$\sigma_{(i,j)} = \frac{\sigma_{(i,i)} + \sigma_{(j,j)}}{2} \tag{23}$$

$$\varepsilon_{(i,j)} = \sqrt{(\varepsilon_{(i,i)} \times \varepsilon_{(j,j)})}$$
(24)

The electrostatic interactions between charged particles are described by the Coulombic potential (Equation 25):

$$U_{coul}\left(\left|r_{ij}\right|\right) = \frac{q_i q_j}{\left(4\pi\varepsilon_0 \left|r_{ij}\right|\right)}$$
(25)

where q_i and q_j represent the partial charges on the atoms and ε_0 is the vacuum permittivity. The Ewald summation approach is used to calculate electrostatic interactions in periodic systems.²²⁹

The Hyperchem program²⁰³ was used to build the primary OAPS molecular structures. The output of Hyperchem is a *.*ent* file containing the coordinates and the connectivities of the atoms. In each *ent* file, the atom names were replaced with the atom-type numbers defined for *gmq*. Table 33 shows the atom labels and atom-types needed to describe OAPS. The atom-types that are not given have been used elsewhere for MD simulations of aliphatic-based polyPOSS-imides.^{26, 30-31} Since OAPS had not been modelled in the former works, some specific atom-types (32 for Car3 and 33 for Hcar3) were added.

Atom label	Atom type	Description	
Si	1	Silicon in cage	
0	2	Oxygen in cage	
Npri	4	The nitrogen of the primary amine	
Hnpri	24	The hydrogen of the primary amine	
Car3	32	Aromatic carbon of the aminophenyl ring	
Hcar3	33	The hydrogen of the aminophenyl ring	

Table 33. Atom labels and atom types used for modelling OAPS with gmq

To obtain accurate force-field parameters, i.e. the equilibrium bond distances, the equilibrium bending angles, the partial charges as well as the optimized coordinates for OAPS starting structures, the *.*ent* files were converted to input files to carry out DFT optimizations with Gaussian at the B3LYP/6-31G(d,p) level of theory. As mentioned before, this basis set is a good compromise between time and accurate energy minimisation.¹⁴⁹ The VMD software²³⁰ was used for the visualization of the atomic structures. Figure 79 shows optimized *meta*-OAPS, *ortho*-OAPS and *para*-OAPS molecular structures.

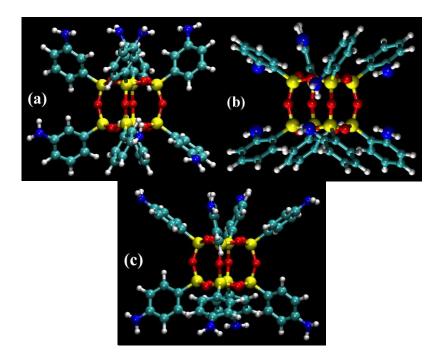
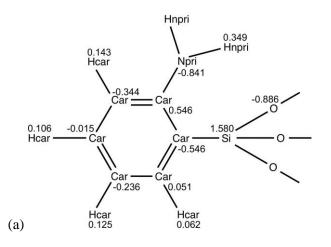


Figure 79. Optimized structures for the (a) meta, (b) ortho, and (c) para isomers of OAPS shown with VMD.²³⁰

The optimized partial charges (Equation 25) for each separate OAPS isomer extracted from the Gaussian outputs are given in Figure 80. The sum of the charges must be equal to zero for each structure.



Chapter 6

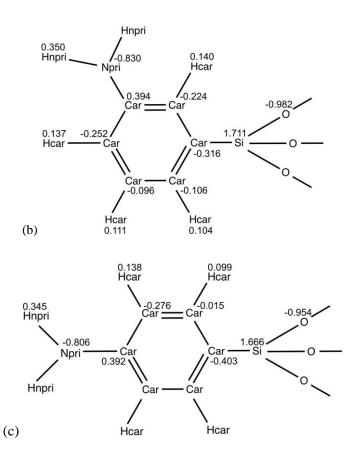
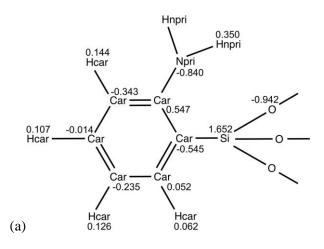


Figure 80. Individual charge distributions in the (a) ortho, (b) meta, and (c) para isomers of OAPS

In Figure 80, the charges of each isomer were calculated individually due to the different charge distributions. But to be able to simulate both individual isomers and mixtures of *ortho*, *meta*, and *para* isomers, it is crucial to have the same charges on the cage Si and O atoms. So after averaging the charges of Si and O, the other charges were adapted again to ensure total charge neutrality. The resulting final charge distributions are given in Figure 81.



Chapter 6

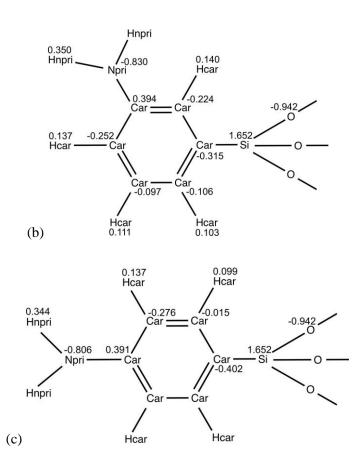


Figure 81. Final charge distributions in the (a) ortho, (b) meta, and (c) para isomers of OAPS

The equilibrium bond lengths (Equation 14) and bending angles (Equation 15) are other parameters, which can be extracted from the Gaussian-optimized structures over all three isomers. Table 34 shows the equilibrium bond lengths for OAPS.

Bond type	b_0 / Å
Si-O	1.65
Si-Car	1.85
Car-Car	1.40
Car-Hcar	1.08
Car-Npri	1.39
Npri-Hnpri	1.01

Table 34. Bond types in OAPS and their equilibrium lengths

The equilibrium bending angles extracted from Gaussian and their respective constants k_{θ} adapted from the TRIPOS force-field are given in Table 35.²³¹⁻²³⁴ Table 36 indicates the dihedral types with the coefficients of the polynomial a_0 , a_1 , a_2 and a_3 (Equation

17). Table 37 shows the out-of-plane parameters (Equation 19) and Table 38 the van der Waals parameters (Equation 22). All parameters that were not extracted from the optimized Gaussian outputs were adapted from the literature.²³¹

Bond-angle type	$ heta_0$ / deg	$k_{ heta}/\mathrm{kJ}\mathrm{mol}^{-1}$
O-Si-O	109.0	-
Si-O-Si	148.8	-
O-Si-Car	109.9	477.8
Si-Car-Car	120.9	879.1
Car-Car-Car	120.0	879.1
Car-Car-Hcar	120.1	879.1
Car-Car-Npri	120.6	1465.1
Car-Npri-Hnpri	115.4	constrained
Hnpri-Npri-Hnpri	112.2	constrained

Table 35. Bond-angle types in OAPS, equilibrium angles and their constants

Table 36. Dihedral types in OAPS and their polynomial constants

Dihedral type	a_0 / kJ mol ⁻¹	a_1 / kJ mol ⁻¹	a_2 / kJ mol ⁻¹	<i>a3 /</i> kJ mol ⁻¹
-Si-Car-	-0.135	0	0	0.480 (<i>a</i> ⁶ used instead of <i>a</i> ³)
-Car-Npri-	11.380	0	-11.380	0
-Car-Car-	16.736	0	-16.736	0

Table 37. Out-of-plane types in OAPS and their force constants

Out-of-plane atom	k_{oop} / kg.s ⁻²
Car-Si-Car-Car	667
Car-Car-Car-Hcar	667
Car-Npri-Car-Car	667

Like-atom pair ii	σii / Å	εij∕kβ∕ K
SiSi	3.385	294.38
00	2.955	102.15
NpriNpri	2.762	47.81
CarCar	3.029	53.84
HnpriHnpri	2.673	21.14
HcarHcar	2.673	21.14

Table 38. Non-bonded atom pairs in OAPS and Van der Waals parameters

This force-field was used by the molecular modelling team at LEPMI, USMB, to carry out separate MD simulations of the three OAPS isomers in the bulk state at room temperature. These OAPS simulation boxes were combined with different dianhydride-containing boxes (6FDA, PMDA and ODPA) in order to create model networks for a large number of different polyOAPS-imide hyper-cross-linked networks both at room and at high temperatures.²⁰⁸⁻²⁰⁹

Conclusions

VII. Conclusions

The gas separation process by dense polymer-based membranes is a promising alternative to the cryogenic distillation or adsorption separation processes due to its much lower energy costs. But this is difficult under harsh conditions since most polymers tend to lose their structural integrity at high temperatures. Hybrid ultrathin membranes based on inorganic POSS cross-linked with organic imides are able to improve the thermomechanical resistance while preserving the gas separation properties. These can perform under tougher operating conditions than conventional polymers, but the aliphatic arms of the POSS precursors have been shown to be prone to thermal degradation above 300°C.

This work aims to attain more thermally-stable hybrid membranes based on a related inorganic siloxane-based cage, octa(aminophenyl)silsesquioxane (OAPS), which replaces the aliphatic arm of the initial POSS with a phenyl arm. The difficulty with using OAPS is that it has three different isomers depending on the *meta*, *ortho*, and *para* position of the amine group with respect to the phenyl group. An extensive literature search showed that the isomer ratios were not clearly defined and that several synthesis routes gave different results. This is problematic as molecular modelling suggests that the nature of the isomer will strongly affect the resistance and selectivities of the resulting membranes.

Experimental characterizations were carried out both for a commercial OAPS containing all three isomers and for a controlled OAPS containing only the *para* and *meta* isomers. This included picnometry, which showed that the densities of the commercial and controlled OAPS samples were close, i.e. respectively 1.39 g/cm³ and 1.37 g/cm³. IR, as well as ²⁹Si, ¹³C and ¹H-NMR spectra, were also obtained for both types of OAPS. To better identify the isomers, the experimental results were compared to predictions of the IR and NMR spectra by quantum mechanical methods such as DFT with Gaussian and by machine-learning methods such as MestreNova.

IR predictions allowed for the identification of the various modes of vibrations and assigned some of the experimental peaks to specific isomers. The 1-D and 2-D NMR predictions were more difficult to carry out as the results were found to be highly dependent on the method and the level of theory used. Despite not being able to reproduce the experimental ²⁹Si-NMR shifts to a fully satisfactory level, the Gaussian-predicted shifts at the HCTH407/6-311+(2d,p) level of theory were consistent with the experimental ¹³C-NMR and

Conclusions

¹H-NMR shifts. MestreNova was also tested for the prediction of the ¹H-NMR shifts but did not give accurate results when compared with the 2-D HSQC OAPS experimental spectra. On the other hand, the Gaussian-predicted shifts were found to be well compatible with the contour clouds of the 2-D NMR spectra for both the commercial OAPS (3 isomers) and the controlled OAPS (2 isomers). Within this context, quantum mechanical methods are clearly superior to machine-learning methods, despite being computationally much more expensive. As a result, each peak in the ¹³C-NMR and ¹H-NMR OAPS spectra could be assigned to a specific isomer using the Gaussian predictions.

OAPS was then used as the inorganic precursor along with the organic 6FDA dianhydride to make new hybrid membranes with increased thermoresistance. The fabrication procedure was similar to that developed for the initial polyPOSS-imides, i.e. it combined interfacial polycondensation with imidization. The IR analyses assessed the extent of imidization and the TGA analyses proved that the new polyOAPS-imide membranes could preserve their structures up to 400°C. FE-SEM showed that these membranes had intact surfaces and thicknesses of ~45 nm, while the thickness of the intermediate γ -alumina was ~3.5 µm. Additional EDS analyses confirmed the presence of the appropriate elements. The new polyOAPS-imide membranes were found to preserve gas selectivities of ~10 for H₂/CH₄ and H₂/N₂ at 250°C while the selectivities increased as a function of temperature for He/CO₂ and H₂/CO₂. The apparatus could only test the permeation characteristics of the new membranes up to 250°C, but since the new membranes preserve their structure up to 400°C, it is likely that the permeation trends are valid at higher temperatures.

In parallel, a force-field was developed to carry out molecular dynamics simulations of the three OAPS isomers in order to characterize their differences at the atomic level. Different force-field parameters such as partial charges, bond length, and bond angles were extracted from the DFT calculations. They were used in another work to create model polyOAPS-imides in order to complement the experimental characterizations.

PolyOAPS-imide membranes thus appear as promising materials for gas separation under harsh conditions. The IR and NMR spectra for OAPS can help identify which isomers are present in the inorganic precursor. However, the membrane fabrication procedure such as the reaction time and the thermal treatment should be tested to a better extent in order to ensure reproducibility.

VIII. Annexes and References:

Annexe 1

How to extract IR results from Gaussian

This tutorial describes how to get the IR data from Gaussian for an optimized structure. *metaOAPS-IR* is used here as an example.

1. An input file of the optimized structure should be fed to Gaussian with the following IR (frequency) command: freq b3lyp/6-31G(d,p)

The geometry optimization and IR commands can be carried out in one single run using: opt freq b3lyp/6-31G(d,p)

The level of theory for both the optimization and frequency calculations must be the same and the basis sets (here b3lyp/6-31G(d,p)) can be chosen based on the accuracy needed. The Gaussian output is a file called *metaOAPS-IR.log*.

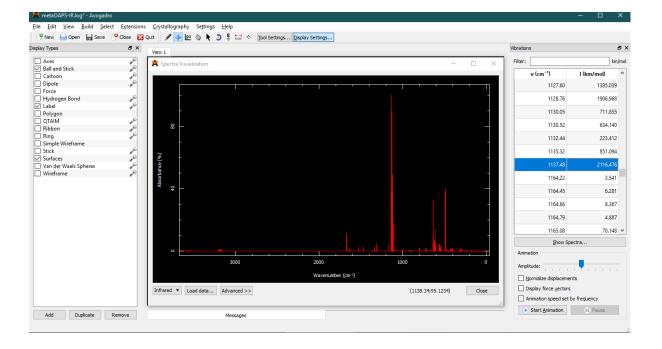
2. Avogadro is a freely-available software very useful to analyze the IR frequencies calculated by Gaussian. It works as described below:

* First open Avogadro

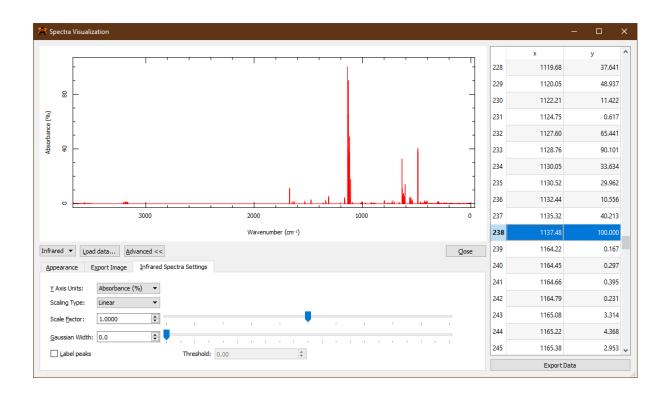
* Open the *log* file by clicking on the *metaOAPS-IR.log* file.

* A table on the right side shows the frequencies and the absolute intensities. Click on *"Show Spectra..."*. This displays the IR spectrum with the same wavenumbers as the Gaussian *log* file but with normalized intensities. Avogadro sets the highest intensity peak (here, 1137.48 cm⁻¹) to 100 and renormalizes the other intensities.

* "Animation" under the right table allows the user to identify the vibrations corresponding to each wavenumber. Close the "Spectra visualization" and click on one of the wavenumbers in the right table, then click on "Start Animation" to observe the type of vibration corresponding to the selected wavenumber.



* Click on "Show Spectra..." then on "Advanced" to see the possible functions. In the "Appearance" tab, the background and spectra colour can be changed. The spectra can be exported with the desired size using the "Export image" tab. The width of the peaks is adjustable with the "Gaussian width" in the "Infrared Spectra Setting" tab. A *.tsv* version of the data can be exported with "Export data" to plot them with Excel.



3. An alternative way to extract the frequencies and their intensities is to open the Gaussian *log* file using text editors. Open *metaOAPS-IR.log* and search "Frequencies" and "IR Inten", which indicate the wavenumber and their corresponding intensities, respectively. It should be noted that the intensity is not normalized with respect to the highest one as in Avogadro.

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120	1	0.18	-0.06	-0.02	0.31	-0.10	-0.03	-0.13	0.04	0.01			
121	1	0.02	0.03	0.01	0.07	0.08	0.03	-0.02	-0.03	-0.01			
122	7	0.00	0.02	0.01	0.00	0.03	0.01	0.00	-0.01	0.00			
123	1	-0.11	-0.04	-0.05	-0.16	-0.06	-0.08	0.04	0.02	0.02			
124	1	0.14	-0.02	0.03	0.20	-0.02	0.04	-0.05	0.01	-0.01			
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IR Int	en	21	16.4759			3.5411			6.2806	5			
Atom	AN	Х	Y	Z	х	Υ	Z	х	Y	Z			
1	14	0.06	-0.08	0.01	0.00	0.00	0.00	0.00	0.00	0.00			
2	8	-0.05	-0.03	0.02	0.00	0.00	0.00	0.00	0.00	0.00			
3	14	0.04	-0.02	-0.02	0.00	-0.01	0.00	0.00	0.00	0.00			
4	8	0.01	0.01	0.07	0.00	0.00	0.00	0.00	0.00	0.01			
5	14	0.05	-0.01	-0.03	0.00	0.00	0.00	0.00	0.01	-0.01			
6	8	-0.10	-0.06	0.03	0.01	0.00	0.00	-0.01	-0.01	0.00			
7	14	0.06	-0.07	-0.03	-0.01	0.00	-0.01	0.01	0.00	0.00			
8	8	0.00	0.01	0.04	0.00	0.00	0.01	0.00	0.00	0.00			
9	8	-0.19	0.30	-0.02	0.00	0.01	0.00	0.00	0.01	0.00			
<	5	5.15	0.50	5.02	5.00	0.01	0.00	5.00	0.01	0.00			>
									L	.n 11648, Col 13	100% Unix (LF)	UTF-8	

* The frequencies and their corresponding absolute intensities can then be found one by one by continuing to search for "Frequencies".

📗 metaOAPS-IR.log - Notepad				– 🗆 X
File Edit Format View Help				
Frequencies	1137.4818	1164.2193		1164.4484 ^
Red. masses	7.2584	1.4571		1.4561
Frc consts	5.5332	1.1636		1.1633
IR Inten	2116.4759	3.5411		6.2806 🗸
<				>
		Ln 11652, Col 30	100% Unix (LF)	UTF-8

Annexe 2

How to extract NMR shifts from Gaussian

This tutorial first describes how to get the Gaussian NMR data for the TMS (trimethylsilane) reference, and then for an arbitrary structure such as *meta*-ONPS.

1. An input file of the optimized TMS is fed to Gaussian using the following command: *nmr=(giao) B3LYP/6-31G(d,p) scrf=(solvent=dmso)*

The geometry optimization and NMR commands can also be carried out in one single run: Write at the start of the file: B3LYP/6-31G(d,p) Opt Density=Current Pop=MK, then: the atomic coordinates at the end: nmr=(giao) B3LYP/6-31G(d,p) scrf=(solvent=dmso)

The basis set (here, B3LYP/6-31G(d,p)) can be chosen depending on the accuracy needed. The Gaussian output is a file called *TMS.log*

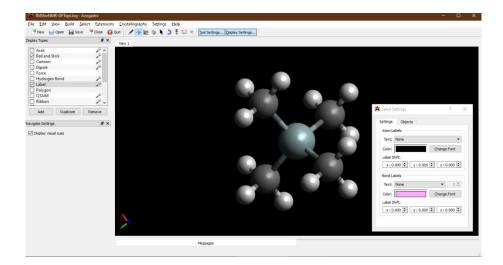
2. Avogadro is a freely-available software very useful to analyze the NMR shifts calculated by Gaussian. It works as described below:

* First open Avogadro

* Open the log file by clicking on the TMS.log file

* Make sure that the options such as "Tool settings..." and "Display settings..." are activated in the top bar. Then their menus can appear in the window.

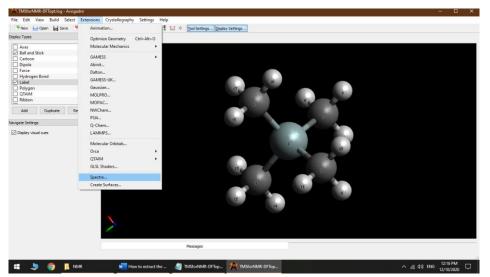
* If the atom numbers do not appear on the atoms, click on the size icon in the "Label" row. A small window appears, as in the following Figure.



* The "Atom number" can be selected from the "Text" menu:

🔀 Label Set	ttings				×
Settings	Objects				
Atom Lab	oels				
Text: 1	None				-
Color:	None Atom number				^
Label Si g	Element syml Symbol & Nu	mber in	Group		
	Symbol & Ato Formal charg Partial charg		er		
Bond La B	Element nam Residue nam	e			11
Text: F	Residue num	ber		_	~
Color:			Change	e Font	
Label Shi	ft:				
x:0.0	00 🗘 y	: 0.000	‡ z:	0.000	\$

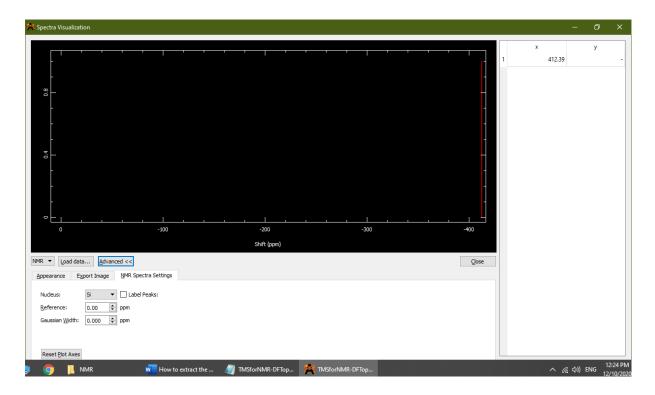
* Click on the "Extensions" menu and "Spectra" to get the chemical shielding values in ppm.



* A new window opens. Click on "Advanced" and "NMR spectra settings" respectively. Since TMS itself is the reference, make sure that the values of "Reference" and "Gaussian width" are zero. As can be seen below, the default "Nucleus" is Si and its chemical shielding is given

Annex II

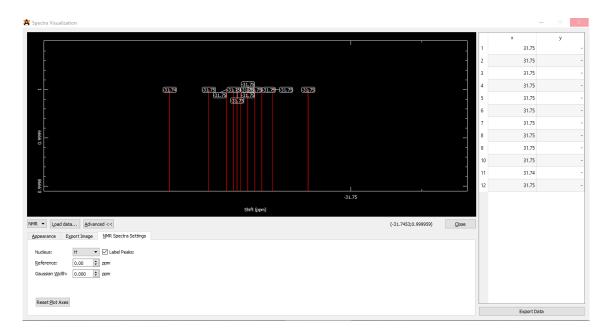
on the right side, i.e. 412.39 ppm. This is the reference for ²⁹Si-NMR at the B3LYP/6-31G(d,p) level of theory. It is also possible to see the shielding value on the red peak by checking the box "Label Peaks", which is on the same row as "Nucleus".



* By changing the "Nucleus" value to C, the chemical shielding values of the TMS carbons will appear on the right side. It is possible to zoom in on the peaks by using the mouse. Since in a single-point structure, the charge distribution and bond lengths can be slightly different (even for atoms with the same chemical environments), the chemical shielding values for the same atoms must be averaged. The average of the four values that appear in the Figure is 192.29 ppm and it can be used as the reference for ¹³C-NMR at the B3LYP/6-31G(d,p) level of theory.

🔀 Spectra Visualization						- 0	×
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(<u>192.28)</u> - -		1992 29 1992 29			4 192.29		
-		, , , , , , , , , , , , , , , , , , ,	-192.29				
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Audeus: C ▼ ∑ Lale Peaks: Beference: 0.00 € ppm Gaussian Width: 0.000 € ppm Reset Blot Axes							
					Export D	ata	

* All hydrogens in TMS have the same chemical environments and their average chemical shielding is 31.72 ppm. It is necessary to mention that the numbers in the table on the right side are not the atom numbers but the order of the hydrogens' chemical shieldings based on the atom numbers. For example, number 1 shows the chemical shielding for the hydrogen with the lowest atom number, which is the hydrogen with atom number 6. Number 6 shows the chemical shielding for the hydrogen with the sixth-highest atom number, which is the hydrogen with the order of the atom number, which is the hydrogen with the order of the atom number for the hydrogen with the sixth-highest atom number, which is the hydrogen with atom number 10. It is possible to find the order of the atoms by using either Avogadro or the input structure.



3. An alternative way to extract the chemical shielding is to open the log files using text editors. Open the TMS log file and search "Isotropic", as indicated in the Figure below:

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54 vectors produced by pass 2 Test12= 5.28D-15 1.85D-09 XBig12		
54 vectors produced by pass 3 Test12= 5.28D-15 1.85D-09 XBig12		
54 vectors produced by pass 4 Test12= 5.28D-15 1.85D-09 XBig12		
9 vectors produced by pass 5 Test12= 5.28D-15 1.85D-09 XBig12	2= 1.97D-15 8.39	D-09.
InvSVY: IOpt=1 It= 1 EMax= 5.33D-15		
Solved reduced A of dimension 279 with 54 vectors.		
Calculating GIAO nuclear magnetic shielding tensors.		
SCF GIAO Magnetic shielding tensor (ppm):		
1 Si <mark>Isotropic</mark> = 412.3930 Anisotropy = 0.0365		
XX= 412.4128 YX= 0.0298 ZX= 0.0087		
XY= -0.0201 YY= 412.3666 ZY= -0.0055		
XZ= 0.0093 YZ= -0.0147 ZZ= 412.3995		
Eigenvalues: 412.3627 412.3988 412.4173		
2 C Isotropic = 192.2856 Anisotropy = 7.8613		
XX= 189.6833 YX= -0.2563 ZX= -0.2092		
XY= -0.2562 YY= 194.3302 ZY= 3.8560		
XZ= -0.2053 YZ= 3.8536 ZZ= 192.8433		
Eigenvalues: 189.6597 189.6706 197.5265		
3 C Isotropic = 192.2855 Anisotropy = 7.8520		
XY= 0.9053 YY= 189.8570 ZY= -0.8267		
XZ= -3.8052 YZ= -0.8293 ZZ= 193.1672		
Eigenvalues: 189.6612 189.6751 197.5202		~
	Ln 854, Col 24	> 100% Unix (LF) UTF-8
	21.05.9, COLLA	

The value for atom number 1, 412.3930 ppm, is the chemical shielding for the Si atom in TMS at the B3LYP/6-31G(d,p) level of theory. The atom numbers 2, 3, 4, and 5 give the chemical shielding values for the four carbons of TMS and should be averaged to calculate the chemical shielding for C in TMS. In that example, the reference chemical shielding values for Si, C, and H were 412.39 ppm, 192.29 ppm, and 31.72 ppm respectively.

After extracting the chemical shielding values for Si, C, and H, the chemical shifts can be obtained from:

$$\delta = \sigma_{iso} - \sigma$$

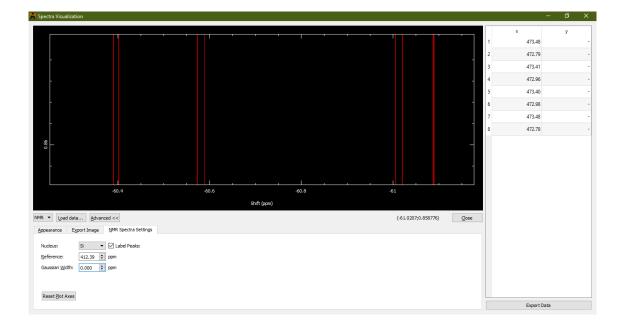
where δ is the chemical shift, σ is the chemical shielding value of the target, and σ_{iso} is the isotropic value or chemical shielding for the standard reference used in the NMR.

4. The level of theory for both the reference (TMS) and the structure under study must be the same. For example, if the structure is studied with the HF/6-31G* level of theory, the same level (HF/6-31G*) must be applied for TMS.

5. To calculate the chemical shifts of a non-reference structure such as *meta*-ONPS:

- Open *meta*-ONPS.log with Avogadro and display the atom numbers using "Label" and *P* Then, click on "Extensions"/"Spectra".
- Click on "Advanced", then NMR Spectra Settings.
- Select Si as the nucleus.

- Now type 412.39 in the Reference box. This is the reference obtained from TMS.
- Upon typing the reference, the chemical shifts appear in the graphical window. The table on the right still shows the chemical shielding values (ex: 473.79 ppm), while the plot displays the corresponding chemical shifts (-60.4 ppm), i.e. the difference in shielding values between the sample and reference atoms (see Equation above). It would be preferable if Avogadro could display directly the chemical shifts in the table, but this is not the case right now.
- By zooming in on the spectrum, it is possible to see the exact values of the chemical shifts.
- Since all the Si atoms have the same chemical environment in *meta*-ONPS, the average ²⁹Si-NMR chemical shifts for the eight Si with the level of theory understudy is -60.77 ppm.



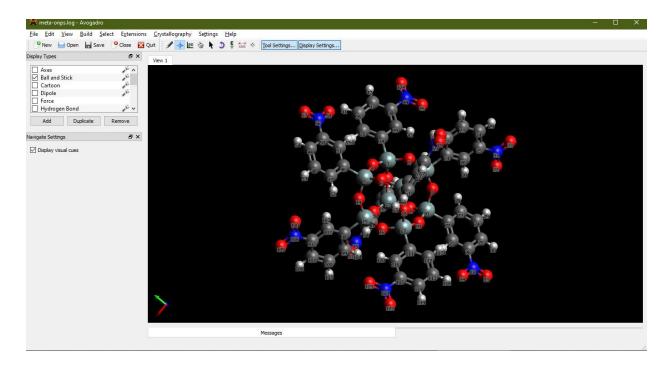
• In order to observe the chemical shifts for Carbon, the Nucleus should be changed to C and the reference shielding for C, i.e. 192.29 ppm, should be inserted into the Reference box.

Annex II

Spectra Visualization						- 0 ×
। । । । । । ।		[127.7 [127.7	72 [26.30 [26.43 [23.29 \$772][26.83][[25[8][125[8]]	1	x 64.33	у ^ -
144, 13 143,95 144, 1 4 4 3,95 143,63 144,29 143,64	137.48 186.90 136.84 137.42 137.13 136.84	127.72 127.72	29.6727493.05123129,7625.48 1224,24123.25 99126.83126.436.30 124.99 123.20123.20122.83	2	55.10	-
<u>[144.28]</u>	136.91		125.80 125.48	3	66.21	
				4	68.76	-
				5	48.05	
				6	64.95	-
-			-	7	64.28	
<u>6</u> –				8	54.57	-
				9	66.52	
				10	68.75	-
-				11	47.71	
				12	65.57	-
145 140	135	130	125	13	64.01	-
	Shift (ppm)			14	54.87	-
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Appearance Export Image NMR Spectra Settings				16	68.80	
Nucleus: C 💌 🗹 Label Peaks:				17	47.86	-
Reference: 192.00 🛊 ppm				18	65.17	
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				20	55.16	
Reset Plot Axes				21	67.01	- ~
					Export Da	a

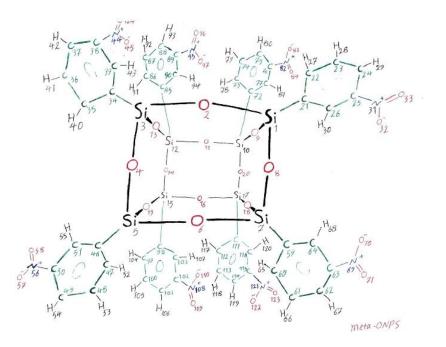
As can be seen, the shifts are very close to each other and the table on the right side does not include the atom numbers. It is possible to see better by zooming on the spectrum.

6. An alternative is to study the Gaussian *log* file directly along with the structure opened by Avogadro, which provides the atom numbers:



or by using directly a drawn Figure:

Annex II



For example, there are eight carbons connected to the Si atoms, all of which have the same chemical environment. Their atom numbers are: 21, 34, 46, 59, 72, 85, 98 and 111. It is then possible to find their corresponding chemical shielding values in the Gaussian *log* file:

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File	Edit	Format	View Help							
	21	С	Isotropic =	64.	3254	Anisotropy	=	181	.0098	0
<										>
			Ln 123666, Col 24		100%	Unix (LF)		UTF-8	3	
/ //	meta-o	onps.log	- Notepad					_		×
File	Edit	Format	View Help							
	34	С	Isotropic =	64.	2773	Anisotropy	=	180	.9549	0
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			Ln 123666, Col 24		100%	Unix (LF)		UTF-8	3	
/ //	meta-c	nps.log	- Notepad					—		×
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				64.	0124	Anisotropy	=	 180	.7151	×
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<										>
			Ln 123666, Col 24		100%	Unix (LF)		UTF-8	3	
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The average shielding value of the carbons connected to the silicons is 64.23 ppm. According to the reference chemical shielding which is 192.29 ppm, the average chemical **shift** of the carbons connected to the silicons is thus: 192.29 - 64.23 = **128.06 ppm**. The chemical shifts of the other atoms can be calculated according to the same procedure.

GaussView is another software prepared to visualize the chemical shielding values calculated by Gaussian. It is not free, but it is probably improved with respect to Avogadro.

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