

Superfluidity and coherence of Bose gases in reduced dimensions: from harmonic traps to uniform fluids



Lauriane Chomaz – PhD defense November 10th, 2014







Macroscopic quantum phenomena.

Superfluidity



Superconductivity



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Superconductivity



In 2D: special meaning. Important for understanding new material :



Semiconductors heterostructures

Graphene



Complex oxide heterostructures



YBa2Cu3O7 (.3) lattice

 $High-T_c$ superconductors

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Infinite uniform gas of non/weakly interacting Bosons at temperature T:

Classical world



describe by trajectories
No coherence

Infinite uniform gas of non/weakly interacting Bosons at temperature T:

Classical world Quantum world Quantum field: $\hat{\psi}(r) \equiv \text{annihilates an atom in } r$ $\langle \hat{\psi}^{\dagger}(r)\hat{\psi}(r) \rangle = \text{probability of finding an at. in } r$ $\langle \hat{\psi}^{\dagger}(r)\hat{\psi}(r) \rangle = \text{probability of finding an at. in } r$ $\langle . \rangle \equiv \text{average over the thermal state of the gas}$ describe by trajectories No coherence Coherence $G_1(r) = \langle \hat{\psi}^{\dagger}(r) \hat{\psi}(0) \rangle$

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Bose statistic in 3D.

Bose-Einstein condensation (at $\lambda_T > C \stackrel{\frown}{d}$). $C = c^{st} \approx 1.38$ for 3D ∞ id.

State of matter stemming from Bose stat.

 $\hat{H} = \sum_{i \text{ atoms}} \hat{h}_i | \text{ single-part. Hamiltonian } \hat{h} = \frac{\hat{p}^2}{2m}.$ eigenstates = $|\varphi_k\rangle, \ p = \hbar k.$ ideal

Saturation of the population in $k \neq 0$ at $\lambda_T = C d$.

 \Rightarrow Macroscopic occupation of a single-particle state $|\varphi_0\rangle$.

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long-range phase coherence

$$G_1(r) \xrightarrow[r \to \infty]{} n_0 \neq 0$$

 $n_0 = \text{density in } |\varphi_0\rangle$

MIT 1997

 k_y

JILA, NIST.

1995

3D ∞

Outline.

Motivation: specificity of 2D and 3D-to-2D crossover.

Motivation: from harmonic to uniform traps.

+ Experimental setup: creating 2D uniform gases.

 Emergence of an extended phase coherence in 2D uniform gases at thermal equilibrium.

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$2D \otimes$ ideal gas, $\lambda_T > d$

Effect of bidimensionality.

no BEC in an infinite 2D gas at $T \neq 0$

(Peierls - 1935, Mermin-Wagner-Hohenberg - 1966-67)

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no BEC in an infinite 2D gas at $T \neq 0$ (Peierls - 1935, Mermin-Wagner-Hohenberg - 1966-67)

 $G_1(r) \stackrel{}{\underset{r o \infty}{\longrightarrow}} 0$

but Bose statistic => broader component in G_1 arises.

 $G_1(r) = G_1^{(ext{th})} + G_1^{(0)} + \dots \leftarrow ext{intermediate energy-states}$

high energy-states

$$G_1^{(\mathrm{th})}(r) \propto \exp\left(-\frac{\pi r^2}{\lambda_T^2}\right)$$

ideal gas

2D ∞

 \Rightarrow Fast increase of ℓ_c/λ_T when $\mathcal{D} \gtrsim 1$

Finite size effect

Fast increase of ℓ_c with \mathcal{D} (when ℓ dominates, *i.e.* $\mathcal{D} \gtrsim 1$).

At $\mathcal{D} \geq \mathcal{D}_{BEC}, \ \ell_c \gtrsim L$ \Rightarrow recover BEC : L $G_1(r) > 0 \ \forall r \in \text{gas.}$

 $G_1(L) \Leftrightarrow$ condensed density

 $\Rightarrow \mathcal{D}_{BEC} \sim 7\text{--}10 \text{ (for typical } L \text{ and } T)$

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Boltzmann	«Pre-superfluid» /«Quasi-condensate»	BEC/ Superfluid ${\cal D}$
$G_1(r) \propto e^{-\pi r^2/\lambda_T^2} \approx$	$G_1(r) \propto \rho_0 e^{-r/\ell} + \approx 7$	10 $G_1(r) > 0, \forall r \in \text{gas}$

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	NIST Gaithersburg, Boulder	Chicago, ENS, Seoul
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3rd dim. in 2D experiments.

How to realize 2D systems in cold atoms experiments...

 $\mathcal{D}\equiv\lambda_T^2
ho$ –2D density

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* $k_BT \ll \hbar \omega_z$ thermally 2D, only j = 0 populated:

$$\mathcal{D}_{\mathrm{tot}} \approx \mathcal{D}_0.$$

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* $k_BT \ll \hbar \omega_z$ thermally 2D, only j = 0 populated:

$$\mathcal{D}_{\mathrm{tot}} \approx \mathcal{D}_0.$$

* $k_BT\gtrsim \hbar\omega_z$ all js must be accounted for.

$$\mathcal{D}_{\text{tot}} = \mathcal{D}_0 + \mathcal{D}_{\text{exc}}$$
$$G_1(r) = \sum_{j=0}^{\infty} G_1^{(j)}(r)$$

 $G_1^{(0)}(r) \equiv \text{broadest contribution}$

N. J. van Druten & W. Ketterle, PRL (1997)

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 \rightarrow When \mathcal{D}_{tot} large enough, atoms accumulate in j = 0: $\mathcal{D}_0 \sim \mathcal{D}_{tot}$ even if $k_B T \gg \hbar \omega_z$.

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Recover relevance for 2D physics!

 $\mathcal{D} \rightarrow \mathcal{D}_{0}$ $G_{1}^{(0)}(r)$ becomes dominant => enhanced coherence in plane!

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Harmonic confinement and local probing

Harmonic confinement and local probing

• Useful for exploring local properties : smooth variation = locally uniform. **Local Density Approximation**

EoS = relation between dimensionless parameters.

=> dictates all thermodynamics properties. can deduce EoS for any variables. In weakly interacting 2D gases: only 2 variables (approx. scale invariance).

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=> fit-free method for measuring the EoS by derivations/integrations of $E(V) = \frac{\hbar^2 n(V)}{m}$:

EoS = relation between dimensionless parameters.

=> dictates all thermodynamics properties. can deduce EoS for any variables. In weakly interacting 2D gases: only 2 variables (approx. scale invariance).





$$ightarrow (\mu/k_BT, \mathcal{D})$$



T. Yefsah & al., PRL (2011), R. Desbuquois & al., PRL (2014)

Superfluidity <=> frictionless response to a moving obstacle.

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Local probing : rotation of a tightly focused laser beam ($w = 2 \mu m / R_{TF} = 20 \mu m$) on "equi-potential" $r = c^{st}$ ($\mu_{loc} = \mu_0 - V(r)$) at velocity $v = c^{st}$.



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Measure heating





Superfluidity <=> frictionless response to a moving obstacle.



Limitations of harmonic confinement.

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Limiting for non-local measurements:

- Iong-range properties
- coherence and correlation functions

In particular at vicinity of a transition point:

- transitions become local phenomena
- divergences are smoothen

Entangles effects of interactions / Bose stat / dimensionality:

- strong influence of interactions (even if weak) / overall state of the gas.
- transverse condensation disappears : not a transition anymore

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Toward uniform traps...

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3D cold gas of ⁸⁷Rb







3D: Hadzibabic group, 2013





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Quasi-2D gases in boxes.



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Coherence in time-of-flight.



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Bimodality of ToF density distribution.



fit by a sum of 2 Gaussians

 $N = N_1 + N_2$

Bimodality of ToF density distribution.



fit by a sum of 2 Gaussians

 $N = N_1 + N_2$

bimodality parameter : $\Delta = N_1/N$

Bimodality of ToF density distribution.



Coherence via matter-wave interference.



 $\rho \,(\mu \mathrm{m}^{-2})$

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0

Coherence via matter-wave interference.



Release box confinement and let expand in-plane for τ such that the gases overlap **Visualize interferences in a unique plane.**









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Kibble Zurek mechanism.



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P. C. Hohenberg, B. I. Halperin, Rev. Mod. Phys. (1977).

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transition point, e.g. BEC



t = 0

at domains boundaries :

phase windings => topological defects

 $N_{
m d} \propto au_Q^{-2}$

transition point, e.g. BEC



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In situ images, no excitation: - Phonons = phase defects => not visible - Vortices = too small (charact. size $\approx 0.6 \ \mu m/$ imaging res. $\approx 1.2 \ \mu m$)





In situ images, no excitation: - Phonons = phase defects => not visible - Vortices = too small (charact. size $\approx 0.6 \ \mu m/$ imaging res. $\approx 1.2 \ \mu m$)

After 4 ms Release

«short» 3D TOF

reveals defects as density holes...









Analysis of holes properties (size, contrast)=> single vortices (\neq phonons, \neq pairs) Highly degenerate gas (PSD>100) => no thermally activated vortices.

BEC_{\perp} similar to standard 3D BEC



BEC in an ∞ 3D gas.

 $\bigstar Z$

0









BEC₁ similar to standard **3D BEC**



Quantitative test of KZM scaling.



Quantitative test of KZM scaling.



Quantitative test of KZM scaling.







Similar quench through ${\rm BEC}_{\perp}$

At freezing out:

unique patch in the central disk
ID chain of phase patches in the annulus





Similar quench through $\operatorname{BEC}_{\perp}$

At freezing out:

unique patch in the central disk ✤ID chain of phase patches in the annulus Merging of patches: $\Rightarrow 2\pi n$ phase winding = Supercurrents

 $n_{\rm w} = \langle |n| \rangle \propto t_{\rm evap}^{-d/2}$



Similar quench through ${\rm BEC}_{\perp}$

2D TOF



matter-wave interference



Spiral pattern : $2n\pi$ phase winding in the annulus (+access to *n*)

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 $\begin{array}{l} \label{eq:main_states} \hline \textbf{Merging of patches:} \\ = > 2\pi n \ \textbf{phase winding} \\ \hline \textbf{supercurrents} \end{array}$

 $n_{\rm w} = \langle |n| \rangle \propto t_{\rm evap}^{-d/2}$



Similar quench through ${\rm BEC}_{\perp}$

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Experiments	d
Holes	0.35 (8)
Spirals	0.38 (12)

interference

$$n_{\rm w} = \langle |n| \rangle \propto t_{\rm evap}^{-d/2}$$

Spiral pattern : $2n\pi$ phase winding in the annulus (+access to *n*)

Theory, Model
$$F: d = 0.33$$











Outlook.

Toward characterizing correlations and fluctuations:

- Characterization of G_1 : exponential vs. algebraic decay.
- Higher order correlation functions, e.g. G_2 (density fluctuations).

Toward strongly interacting 2D gas in uniform geometry:

- Increase of v_z : the accordion setup
- Regime where coherence established at BKT transition.
- New scaling expected from KZ mechanism.

Toward other geometries, e.g. transport measurement:

- Two uniform reservoirs linked by a small channel.
- Study of superfluidity via transport measurement along the channel

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Esslinger group

Thanks.

The Team:

Jean Dalibard Jérôme Beugnon Sylvain Nascimbène Kenneth J. Günter Christof Weitenberg Tom Bienaimé Tarik Yefsah Rémi Desbuquois Laura Corman David Perconte Benno Rem Julian Leonard Katharina Kleinlein Jean-Loup Ville

The Group:

BEC: Na,Yb, Dy Gaz de Fermi: Li, K-Li + The theory

The Lab

specially, the directors: Paul Indelicato Antoine Heidmann

+ its tutelles

L'X (teaching)

Linda Vary-Guevel

COLLÈGE DE FRANCE

aboratoire Kastler Brossel.

ECOLE NORMALE SUPERIES

hysique quantique et applications

The technicians & administratives:

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My parents Pierrot & Lucile Tonton Jean-Marc & Co. Les Revels Les Chomaz, Borel, Lopez

The Friends:

The atomic physicists The Master «LKB» The «Climbers» & co. The «Caennais» The unclassified

> To all the people present today



Temperature Measurement (I)





Temperature Measurement (II)

Renormalize all *T* and *U* by $\exp(2d_{\max}/w)$, $d_{\max} = \max$ and distance to trap center: w = waist of the 2D trap beam (Hermite-Gauss and box trap) $\exp(-2d_{\max}/w) = \text{reduction}$ of the heights of both traps compare to the center.



	U_0 / k_B	T ₀	η
$v_z = 365 \text{ Hz}$	969 (75) nK	275 (8) nK	3.49 (25)
$v_z = 1460 \text{ Hz}$	712 [115] nK ₃₈	203 (18) nK	3.50[25]



Kibble Zurek in an annulus



Stochastic origin:

random number of windings between -2 and 2 no bias on winding direction $\langle n \rangle = 0.001(12)$

Dynamic origin:

 $n_{\rm w}$ not compatible with thermal predictions
Kibble Zurek in an annulus



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$$n_{\rm w} \propto t_{\rm evap}^{-d/2}$$
$$d = 0.38 \,(12)$$

Reminder: KZM for 3D BEC

	d
MF	0.25
F Model	0.33

no evidence for a breakdown of KZM for 3D BEC (but no counter-evidence either...)₄₀

Direct Observation of vortices (II)



Direct Observation of vortices (II)



Direct Observation of vortices (II)



Dynamical orgin:

- Equilibrium expectation at final PSD (>100) = vanishingly small mean vortex number N_v . Experimentally $N_v \approx 0.6$
- BKT theory at final PSD = vortices must be tightly paired.
- Dissipative dynamic (variation of N_v) with a varying hold time \neq equilibrium. 41











Active Research

Steady state 2D Physics/ BKT Physics:

- Cheng Chin, Chicago, USA.
- ✦ Zoran Hadzibabic, Cambridge, UK.
- ✦ Yong-il Shin, Seoul, Korea.
- ✦ Bill Phillips, NIST, Maryland, USA.
- ◆ E.A. Cornell, JILA, Colorado, USA.
- ✦ Thomas Bourdel, Palaiseau, France.

Kibble Zurek mechanism with cold atoms

- ♦ B. P. Anderson, Tuscon, Arizona, USA.
- ✦ D. Stamper Kurn, Berkeley, California, USA.
- ♦ B. L. DeMarco, Urbana-Ch., Illinois, USA.
- ✦ Gabriele Ferrari, Trento, Italy.
- Immanuel Bloch, Munich, Germany.
- ✦ Zoran Hadzibabic, Cambridge, UK.







Symmetry Breaking in 2D

Peierls 1935

Crystalline order in 2D?



How do the fluctuations of the distance between two atoms vary with the distance itself?

$$\langle (\vec{u_a} - \vec{u_b})^2 \rangle \propto \log(d_{ab})$$

Crystalline order destroyed by long wavelength phonons (T≠0)

Mermin-Wagner-Hohenberg theorem (1966-67)

No breaking of continuous symmetry in dimensions < 3 (+ short range interactions)

No BEC in an infinite, uniform 2D gas

 $g_1(r) = \langle \psi^{\dagger}(r) \psi(0) \rangle ; \lim_{r \to \infty} g_1(r) = 0$

The superfluid transition in 2D

Not all phase transitions are excluded in 2D, only those due to symmetry breaking! universal jump of superfluid density



Coherence + microscopic nature of the transition : studied with cold atoms.

Previous studies with atomic gases and trapped ions:

- Kinks in ion chains: PTB (Mehlstäubler), Simon Fraser U. (Haljan), Mainz U. (Schmidt-Kaler), U. Freiburg (Schaetz)
- Quenched ferromagnetic spinor Bose-Einstein condensates: U. Berkeley (Stamper-Kurn)
- Formation of BEC in harmonic traps: U. Arizona (Anderson) & U. Queensland (Davis)
- Solitons in elongated BECs: U. Trento (Ferrari)
- Superfluid to Mott insulator: Munich (Bloch-Schneider)
- Formation of a 3D BEC in a box-like trap: U. Cambridge (Hadzibabic)



 $\boldsymbol{\boldsymbol{\zeta}}$

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Kibble Zurek mechanism for a 3D Bose gas

The static picture

Above the critical temperature, coherent patches of size *d*

For large $T : d \sim \lambda$ (thermal wavelength) Approaching $T_c : d$ diverges as $1/(T - T_c)^{\nu}$

Below T_c , full coherence: d = size of the box



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Dynamical aspects

For a finite cooling rate, this picture fails: coherence cannot be established fast enough

- → \dot{T} fixes the maximal size of phase-coherent patches: topological defects (vortex lines) may appear at the border of the patches and survive to the subsequent → thermalization $N_t \propto (\dot{T}/T)^{\alpha}$
 - KZ theory relates the number of topological defects N_t to the cooling rate:

One can also look for other signatures than topological defects: Hadzibabic's talk







N. J. van Druten & W. Ketterle, PRL (1997)





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 \rightarrow When \mathcal{D}_{tot} large enough, atoms accumulate in j = 0: $\mathcal{D}_0 \sim \mathcal{D}_{tot}$ even if $k_B T \gg \hbar \omega_z$.



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We cross two successive transitions:

T₁: transverse BEC which occurs when the 3D phase-space density reaches 2.6

Atoms accumulate in the ground state j = 0 of the z motion, coherence $\sim a_{ho}$ length T₂: superfluid BKT // full BEC corresponding to quasi-full coherence over the sample



Tentative scenario:

- Formation of independent coherent phase domains during transition T₁
- Going up to the superfluid regime allows us to visualize the vortices