# Superfluidity and coherence of Bose gases in reduced dimensions: from harmonic traps to uniform fluids 



> Lauriane Chomaz - PhD defense November 10 ${ }^{\text {th }}, 2014$


## Macroscopic quantum phenomena.

Superfluidity


Superconductivity


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## Superconductivity



In 2D: special meaning. Important for understanding new material :


Semiconductors heterostructures


Graphene

$\mathrm{YBa}_{2} \mathrm{Cu}_{3} \mathrm{O}_{7}$ (.3) lattice

## Quantum coherence.

Infinite uniform gas of non/weakly interacting Bosons at temperature T:
Classical world

describe by trajectories
No coherence

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Classical world Quantum world


Quantum field: $\hat{\psi}(\boldsymbol{r}) \equiv$ annihilates an atom in $\boldsymbol{r}$ $\left\langle\hat{\psi}^{\dagger}(\boldsymbol{r}) \hat{\psi}(\boldsymbol{r})\right\rangle=$ probability of finding an at. in $\boldsymbol{r}$ $\langle.\rangle \equiv$ average over the thermal state of the gas
describe by waves Coherence

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G_{1}(r)=\left\langle\hat{\psi}^{\dagger}(r) \hat{\psi}(0)\right\rangle
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## Characteristic lengths in ideal quantum gases :

* Coherence length:

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\ell_{c} \equiv \text { charact. of decay of } G_{1}(r)
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※ Thermal wavelength:

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\lambda_{T}=\sqrt{2 \pi \hbar^{2} / m k_{B} T}
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\% Mean distance between particles: $d$

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\% Mean distance between particles: $d$
~ Boltzmann statistics

$$
\lambda_{T}<d
$$


$G_{1}(r) \propto \exp \left(-\pi r^{2} / \lambda_{T}^{2}\right)$
Coherence over $\sim \lambda_{T}$

$$
G_{1}(r)=\left\langle\hat{\psi}^{\dagger}(r) \hat{\psi}(0)\right\rangle
$$

describe by trajectories describe by waves
Coherence

No coherence


Bose statistics

$$
\lambda_{T}>d
$$



Modif. $G_{1}$ and $\uparrow \ell_{c}$. depends on the dimension
$3 \mathrm{D} \infty \quad$ Bose statistic in 3D.

## ideal gas, $\lambda_{T}>d$ <br> Bose-Einstein condensation (at $\left.\lambda_{T}>C \overleftarrow{d}\right) . \quad \begin{aligned} & \mathrm{C}=\mathrm{c}^{\mathrm{ct}} \approx 1.38 \\ & \text { for } 3 \mathrm{D} \\ & \text { id }\end{aligned}$



Bose-Einstein condensation (at $\left.\lambda_{T}>C \overleftarrow{d}\right)$. for $3 \mathrm{C}=\mathrm{c}=$
$\Leftrightarrow$ State of matter stemming from Bose stat.

$$
\hat{H}=\sum_{i \text { atoms }} \hat{h}_{i} \left\lvert\, \begin{aligned}
& \text { single-part. Hamiltonian } \hat{h}=\frac{\hat{\boldsymbol{p}}^{2}}{2 m} . \\
& \text { eigenstates }=\left|\varphi_{\boldsymbol{k}}\right\rangle, \boldsymbol{p}=\hbar \boldsymbol{k} .
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$\Rightarrow$ Saturation of the population in $k \neq 0$ at $\lambda_{T}=\mathrm{C} d$.
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$\leftrightarrow$
long-range phase coherence

$$
\begin{aligned}
& \boldsymbol{G}_{\mathbf{1}}(r) \underset{r \rightarrow \infty}{\longrightarrow} n_{\mathbf{0}} \neq \mathbf{0} \\
& n_{0}=\text { density in }\left|\varphi_{0}\right\rangle
\end{aligned}
$$

## Outline.

Motivation: specificity of 2D and 3D-to-2D crossover.

Motivation: from harmonic to uniform traps.

Experimental setup: creating 2D uniform gases.

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ideal gas, $\lambda_{T}>d$

no BEC in an infinite 2D gas at $T \neq 0$
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G_{1}(r)=\underset{\uparrow}{G_{1}^{(\mathrm{th})}}+G_{1}^{(0)}+\ldots \leftarrow \text { intermediate energy-states }
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$G_{1}^{(\text {th })}(r) \propto \exp \left(-\frac{\pi r^{2}}{\lambda_{T}^{2}}\right)$

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low energy-states

$$
\begin{aligned}
G_{1}^{(\mathrm{th})}(r) \propto \exp \left(-\frac{\pi r^{2}}{\lambda_{T}^{2}}\right) \quad & G_{1}^{(0)}(r) \propto \exp \left(-\frac{r}{\ell}\right) \\
\ell & =\frac{\lambda_{T}}{\sqrt{4 \pi}} \exp \left(\frac{\mathcal{D}}{2}\right) \\
\mathcal{D} & =\left(\frac{\lambda_{T}}{d}\right)^{2}
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& \text { At } \mathcal{D} \geq \mathcal{D}_{\mathrm{BEC}}, \ell_{c} \gtrsim L \\
& \quad \Rightarrow \text { recover BEC }: L \\
& G_{1}(r)>0 \forall r \in \text { gas. }
\end{aligned} \downarrow \square \square \square
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## Inter-particle interactions

A distinct type of phase transition possible:
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Quasi-long
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\xrightarrow[\rho_{\mathrm{s}}=0]{G_{1}(r) \propto e^{-\pi r^{2} / \lambda_{T}^{2}}}
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& \alpha=1 /\left(\rho_{\mathrm{s}} \lambda_{T}^{2}\right) \\
& \rho_{\mathrm{s}}=4 / \lambda_{T}^{2}
\end{aligned}
$$

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| Boltzmann | «Pre-superfluid»//«Quasi-condensate» | BEC/ Superfluid | $\mathcal{D}$ |  |
| :---: | :---: | :---: | :---: | :---: |
| $G_{1}(r) \propto e^{-\pi r^{2} / \lambda_{T}^{2}}$ | $\approx 1$ | $G_{1}(r) \propto \rho_{0} e^{-r / \ell}+\ldots$ | $\approx 7-10$ | $G_{1}(r)>0, \forall r \in$ gas |

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NIST Gaithersburg, Boulder

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## $3^{\text {rd }}$ dim. in 2D experiments.

How to realize 2D systems in cold atoms experiments...


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$$
\mathcal{D} \equiv \lambda_{T}^{2} \rho \leftarrow 2 \mathrm{D} \text { density }
$$

\& $k_{B} T \ll \hbar \omega_{z}$ thermally 2D, only $j=0$ populated:

$$
\mathcal{D}_{\mathrm{tot}} \approx \mathcal{D}_{0}
$$

$\& k_{B} T \gtrsim \hbar \omega_{z}$ all $j$ 's must be accounted for.

$$
\begin{aligned}
& \mathcal{D}_{\mathrm{tot}}=\mathcal{D}_{0}+\mathcal{D}_{\mathrm{exc}} \\
& G_{1}(r)=\sum_{j=0}^{\infty} G_{1}^{(j)}(r)
\end{aligned}
$$

$G_{1}^{(0)}(r) \equiv$ broadest contribution
$3^{\text {rd dim.: Transverse condensation. }}$

# $3^{\text {rd }}$ dim.: Transverse condensation. 

 quantum states in Bose statistics.

## 3rd dim.:Transverse condensation.



Saturation of the population of the quantum states in Bose statistics.

$$
\mathcal{D}^{(D)}(\mu)=\left(\frac{\lambda_{T}}{L}\right)^{D} \int f(E-\mu) D(E) \mathrm{d} E
$$

$$
\mathcal{D}^{(D)}(\mu) \underset{\mu \rightarrow 0}{\longrightarrow} \begin{cases}\mathcal{D}_{\mathrm{c}} & \Leftrightarrow \text { BEC } \\ \infty & \Leftrightarrow \text { no BEC }\end{cases}
$$

## $3^{\text {rd }}$ dim.: Transverse condensation.



2D+ configuration


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When $\mathcal{D}_{\text {tot }}$ large enough, atoms accumulate in $j=0: \mathcal{D}_{0} \sim \mathcal{D}_{\text {tot }}$ even if $k_{B} T \gg \hbar \omega_{z}$.


Recover relevance for 2D physics!

$$
\mathcal{D} \rightarrow \mathcal{D}_{0}
$$

$G_{1}^{(0)}(r)$ becomes dominant
$=>$ enhanced coherence in plane!
N. J. van Druten \& W. Ketterle, PRL (1997)

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## Harmonic confinement and local probing



## Trapping region :

extrema of magnetic field / light intensity. $=>$ typical dependance $V(r) \sim r^{2}$

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## Trapping region :

extrema of magnetic field / light intensity. => typical dependance $V(r) \sim r^{2}$
\& Useful for exploring local properties : smooth variation => locally uniform.
Local Density Approximation
Trapped density profile
Homogeneous equivalent


## Example of the Equation of State.

EoS $=$ relation between dimensionless parameters.
=> dictates all thermodynamics properties. can deduce EoS for any variables.
In weakly interacting 2D gases: only 2 variables (approx. scale invariance).

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15
$=>$ fit-free method for measuring the EoS by derivations/integrations of $E(V)=\frac{\hbar^{2} n(V)}{m}$ : check scale invariance

$$
\rightarrow\left(\mu / k_{B} T, \mathcal{D}\right)
$$


T. Yefsah \& al. , PRL (20II), R. Desbuquois \& al. , PRL (20|4)

## Example of the superfluid behavior.

Superfluidity $<=>$ frictionless response to a moving obstacle.

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## Local probing :

rotation of a tightly focused laser beam
$\left(w=2 \mu \mathrm{~m} / R_{\mathrm{TF}}=20 \mu \mathrm{~m}\right)$
on "equi-potential" $r=\mathrm{C}^{\text {st }}$

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\left(\mu_{\mathrm{loc}}=\mu_{0}-V(r)\right)
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Measure heating

$T(v)=T_{v=0}+\kappa \tau \max \left[v^{2}-v_{c}^{2}, 0\right]$

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\left(\mu_{\mathrm{loc}}=\mu_{0}-V(r)\right)
$$


discrepancy (partly) due to $w \neq 0$ at velocity $v=\mathrm{c}^{\text {st }}$.
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## Limitations of harmonic confinement.

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\% Limiting for non-local measurements:

- long-range properties
- coherence and correlation functions
\% In particular at vicinity of a transition point:
- transitions become local phenomena
- divergences are smoothen
$\%$ Entangles effects of interactions / Bose stat / dimensionality:
- strong influence of interactions (even if weak) / overall state of the gas.
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$\Rightarrow$ Toward uniform traps...


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## Creating uniform 2D gases in experiments.

$x \xrightarrow{2} y$
3D cold gas
of ${ }^{87} \mathrm{Rb}$

## Creating uniform 2D gases in experiments.



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$x_{4}^{z} y$

$+$
Hermite-Gauss beam $=2 \mathrm{D}$ conf.
$+$
Hollow beam
= Uniform conf.

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3D: Hadzibabic group, 2013

## Creating uniform 2D gases in experiments.



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## Quasi-2D gases in boxes.



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## Coherence in time-of-flight.



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## 3D ToF



Release all trapping potential and let expand for long $\tau$ : access velocity distribution.




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## velocity distribution $=$

Fourier transform of $G_{1}$


## Bimodality of ToF density distribution.



fit by a sum of 2 Gaussians

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N=N_{1}+N_{2}
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bimodality parameter : $\Delta=N_{1} / N$

## Bimodality of To density distribution.



For each $T$, fit $\Delta(N)$ by «heuristic» threshold function:
$\Delta=\left(1-\left(N_{\mathrm{c}} / N\right)^{0.6}\right) \Theta\left(N>N_{c}\right)$

fit by a sum of 2 Gaussian

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## Coherence via matter-wave interference.



## Coherence via matter-wave interference.



Release box confinement and let expand in-plane for $\tau$ such that the gases overlap
Visualize interferences in a unique plane.

## Universal behavior of the onset of coherence.



$$
\zeta \equiv k_{B} T / h \nu_{z}
$$

full symbols : bimodality open symbols : fringe contrast colors:
$\neq$ shapes,sizes

$$
\neq v_{z}
$$

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## Topological defects from KZM.



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## Direct observation of vortices.

$T \uparrow$ thermal cloud



$$
\underline{v_{\mathrm{z}}}=365 \mathrm{~Hz}:
$$

Coherence emerges
via $\mathrm{BEC}_{\perp}$.

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In situ images, no excitation:

- Phonons = phase defects $=>$ not visible -Vortices $=$ too small (charact. size $\approx 0.6 \mu \mathrm{~m} /$ imaging res. $\approx 1-2 \mu \mathrm{~m}$ )


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After 4 ms Release


## «short» 3D TOF

## reveals defects as

 density holes...

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Analysis of holes properties (size, contrast) $=>$ single vortices ( $\neq$ phonons, $\neq$ pairs) Highly degenerate gas (PSD $>100$ ) => no thermally activated vortices.

## Transverse BEC as a candidate for KZM.

$\mathrm{BEC}_{\perp}$ similar to standard 3D BEC

$\mathrm{BEC}_{\perp}$ occurs at
$<n(z=0)=2.612 \lambda_{T}^{-3}$

$B E C$ in an $\infty$ 3D gas.

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$\ell_{c}$
Same universality class but quasi-2D geometry cut divergences...
$\leftarrow$ transition point

3D BEC

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$\ell_{c} \leqslant a_{z} \longleftarrow$ size of the ground state of the harmonic oscillator along $z$

$$
a_{z}=\sqrt{\hbar / m \omega_{z}}
$$

$\ell_{c} \uparrow$| Same universality class but |
| ---: |
| quasi-2D geometry cut divergences... |
| $\ell_{c} \leqslant a_{z} \longleftarrow$ size of the ground state |
| harmonic oscillator alo |
| $a_{z}=\sqrt{\hbar / m \omega_{z}}$ |

$T_{\mathrm{c}}$
$\longleftarrow$ transition point

3D BEC
$\mathrm{BEC}_{\perp}$

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| harmonic oscillator alo |
| $a_{z}=\sqrt{\hbar / m \omega_{z}}$ |

Standard BEC description of $B E C_{\perp}$ not valid when it predicts $\ell_{c} \gtrsim \boldsymbol{a}_{\boldsymbol{z}}$

## Quantitative test of KZM scaling.



## Quantitative test of KZM scaling.



## Quantitative test of KZM scaling.



## KZ supercurrents in an annulus.

$10 \mu \mathrm{~m}$

$$
\begin{aligned}
& \text { In situ } \\
& \text { small central disk } \\
& \text { = phase reference } \\
& R_{0}<\hat{\xi} \\
& \text { Annulus } \\
& \text { ~1D for KZM } \\
& \delta R<\hat{\xi}
\end{aligned}
$$

## KZ supercurrents in an annulus.

$10 \mu \mathrm{~m}$


## Similar quench through $\mathrm{BEC}_{\perp}$

At freezing out:
\% unique patch in the central disk
$\%$ ID chain of phase patches in the annulus

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## $10 \mu \mathrm{~m}$

In situ

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Merging of patches:
$=>2 \pi n$ phase winding = Supercurrents
$n_{\mathrm{w}}=\langle | n| \rangle \propto t_{\mathrm{evap}}^{-d / 2}$

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## Similar quench through $\mathrm{BEC}_{\perp}$

2D TOF

matter-wave interference

At freezing out: \% unique patch in the central disk
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$$
\begin{gathered}
\begin{array}{c}
\frac{\text { Merging of patches: }}{=>2 \pi n \text { phase winding }} \\
=\text { Supercurrents }
\end{array} \\
n_{\mathrm{w}}=\langle | n| \rangle \propto t_{\mathrm{evap}}^{-d / 2}
\end{gathered}
$$



Spiral pattern: $2 n \pi$ phase winding in the annulus (+access to $n$ )

## KZ supercurrents in an annulus.



## Similar quench through $\mathrm{BEC}_{\perp}$

## 2D TOF <br> 

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\end{array}
\end{gathered}
$$

| Experiments | $d$ |
| :---: | :---: |
| Holes | $0.35(8)$ |
| Spirals | $\mathbf{0 . 3 8 ( 1 2 )}$ |

Theory, Model F : $d=0.33$


Spiral pattern : $2 n \pi$ phase winding in the annulus (+access to $n$ )

## Summary



## Summary



## Summary



## Summary



## Summary



## Outlook.

+ Toward characterizing correlations and fluctuations:
- Characterization of $G_{1}$ : exponential vs. algebraic decay.
- Higher order correlation functions, e.g. $G_{2}$ (density fluctuations).

+ Toward strongly interacting 2D gas in uniform geometry:
- Increase of $v_{z}$ : the accordion setup
- Regime where coherence established at BKT transition.
- New scaling expected from KZ mechanism.

+ Toward other geometries, e.g. transport measurement:
- Two uniform reservoirs linked by a small channel.
= Study of superfluidity via transport measurement along the channel



## The Team:

Jean Dalibard Jérôme Beugnon Sylvain Nascimbène Kenneth J. Günter ChristofWeitenberg

Tom Bienaimé Tarik Yefsah
Rémi Desbuquois
Laura Corman
David Perconte Benno Rem Julian Leonard Katharina Kleinlein Jean-Loup Ville

The Group: BEC: $\mathrm{Na}, \mathrm{Yb}, \mathrm{Dy}$ Gaz de Fermi: Li, K-Li

+ The theory


## The Lab

 specially, the directors:Paul Indelicato
Antoine Heidmann

+ its tutelles


## L'X (teaching)

Linda Vary-Guevel


COLLE GE
DE FRANCE
UPMC
Inel PARISUNIVERSITAS

## The Family:

My parents
Pierrot \& Lucile
Tonton Jean-Marc \& Co.
Les Revels
Les Chomaz, Borel, Lopez

## The Friends:

The atomic physicists
The Master «LKB»
The «Climbers» \& co.
The «Caennais»
The unclassified

## To all the people present today

The Jury:
Gabriele Ferrari
Bruno Laburthe-Tolra
Patrizia Vignolo
Chris Westbrook JD \& JB
special thanks to
Carmen Toderasc for help in organizing

## Temperature Measurement (I)

$\left.\begin{array}{ll}\text { ToF measurement for very } & N \\ \text { weakly gas configurations: } & \Delta v\end{array}\right\} \underset{\text { Bose law ideal gas }}{\text { All quantum, 3D }}\left\{\begin{array}{l}Z=e^{\beta \mu} \\ T\end{array}\right.$


## Temperature Measurement (II)

Renormalize all $T$ and $U$ by $\exp \left(2 d_{\max } / w\right), d_{\max }=$ maximal distance to trap center: $w=$ waist of the 2D trap beam (Hermite-Gauss and box trap) 350 $\exp \left(-2 d_{\max } / w\right)=$ reduction of the heights of both traps compare to the center.


|  | $U_{0} / k_{\mathrm{B}}$ | $T_{0}$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| $v_{\mathrm{z}}=365 \mathrm{~Hz}$ | $969(75) \mathrm{nK}$ | $275(8) \mathrm{nK}$ | $3.49(25)$ |
| $v_{\mathrm{z}}=1460 \mathrm{~Hz}$ | $712[115] \mathrm{nK}_{38}$ | $203(18) \mathrm{nK}$ | $3.50[25]$ |

## Characterizing the fringe contrast



$$
\Gamma=b+a\left(1-\left(N_{\mathrm{c}} / N\right)^{0.6}\right) \Theta\left(N>N_{c}\right)
$$



© 1-Body corr. on complex fringe contrast:

$$
\gamma(d)=\left|\left\langle\tilde{\rho}\left[k_{p}(y), y\right] \tilde{\rho}^{*}\left[k_{p}(y+d), y+d\right]\right\rangle\right|
$$

$$
\infty 1 \mathrm{D} \text { gases: } \gamma(d)=\left|G_{1}(d)\right|^{2}
$$

\% Look for extended coherence:

$$
\begin{gathered}
\Gamma=\langle\gamma(d)\rangle 2 \mu \mathrm{~m}<d<5 \mu \mathrm{~m} \\
{\left[\sup \left(\lambda_{T}\right)<2 \mu \mathrm{~m}\right]}
\end{gathered}
$$

## Kibble Zurek in an annulus



## Stochastic origin:

random number of windings between -2 and 2 no bias on winding direction $\langle n>=0.001$ (12)

## Dynamic origin:

$n_{\mathrm{w}}$ not compatible with thermal predictions

## Kibble Zurek in an annulus



## Stochastic origin:

random number of windings between -2 and 2 no bias on winding direction $\langle n\rangle=0.001$ (12)

## Dynamic origin:

$n_{\mathrm{w}}$ not compatible with thermal predictions


Reminder:
KZM for 3D BEC

|  | $d$ |
| :---: | :---: |
| MF | 0.25 |
| F Model | 0.33 |

no evidence for a breakdown of KZM for 3D BEC (but no counter-evidence either...) 40

## Direct Observation of vortices (II)


similar properties at a given $\tau$

+ high contrast


3 ms
4.5 ms

constant contrast and increasing size with $\tau$

## Direct Observation of vortices (II)



## Direct Observation of vortices (II)



similar properties at a given $\tau$ + high contrast


Hole Nature:
= single vortices
$\neq$ phonons
$\neq$ pairs of vortices
constant contrast and increasing size with $\tau$

## Dynamical orgin:

\& Equilibrium expectation at final PSD $(>100)=$ vanishingly small mean vortex number $N_{\mathrm{v}}$. Experimentally $N_{\mathrm{v}} \approx 0.6$
$\%$ BKT theory at final PSD $=$ vortices must be tightly paired.
$\because$ Dissipative dynamic (variation of $N_{\mathrm{v}}$ ) with a varying hold time $\neq$ equilibrium.






## Active Research

## Steady state 2D Physics/ BKT Physics:

$\downarrow$ Cheng Chin, Chicago, USA.
$\downarrow$ Zoran Hadzibabic, Cambridge, UK.
$\uparrow$ Yong-il Shin, Seoul, Korea.

- Bill Phillips, NIST, Maryland, USA.
$\downarrow$ E.A. Cornell, JILA, Colorado, USA.
$\downarrow$ Thomas Bourdel, Palaiseau, France.


## Kibble Zurek mechanism with cold atoms

- B. P. Anderson, Tuscon, Arizona, USA.
$\downarrow$ D. Stamper Kurn, Berkeley, California, USA.
- B. L. DeMarco, Urbana-Ch., Illinois, USA.
$\downarrow$ Gabriele Ferrari, Trento, Italy.
$\uparrow$ Immanuel Bloch, Munich, Germany.
$\uparrow$ Zoran Hadzibabic, Cambridge, UK.
$\downarrow$...





## Symmetry Breaking in 2D

## Peierls |935



## Crystalline order in 2D?

How do the fluctuations of the distance between two atoms vary with the distance itself?

$$
\begin{aligned}
& \left\langle\left(\overrightarrow{u_{a}}-\overrightarrow{u_{b}}\right)^{2}\right\rangle \propto \log \left(d_{a b}\right) \\
& \text { Crystalline order destroyed by long } \\
& \text { wavelength phonons }(T \neq 0)
\end{aligned}
$$

Mermin-Wagner-Hohenberg theorem (1966-67)
No breaking of continuous symmetry in dimensions < 3 ( + short range interactions)
No BEC in an infinite, uniform 2D gas

$$
g_{1}(r)=\left\langle\psi^{\dagger}(r) \psi(0)\right\rangle ; \lim _{r \rightarrow \infty} g_{1}(r)=0
$$

## The superfluid transition in 2D

Not all phase transitions are excluded in 2D, only those due to symmetry breaking! universal jump of superfluid density

$g_{1}\left(\left|\vec{r}-\vec{r}^{\prime}\right|\right)=\left\langle\psi^{\dagger}(\vec{r}) \psi\left(\vec{r}^{\prime}\right)\right\rangle$

Bound vortexantivortex pairs


Coherence + microscopic nature of the transition : studied with cold atoms.

## Previous studies with atomic gases and trapped ions:

- Kinks in ion chains: PTB (Mehlstäubler), Simon Fraser U. (Haljan), Mainz U. (Schmidt-Kaler), U. Freiburg (Schaetz)
- Quenched ferromagnetic spinor Bose-Einstein condensates: U. Berkeley (Stamper-Kurn)
- Formation of BEC in harmonic traps: U. Arizona (Anderson) \& U.

Queensland (Davis)

- Solitons in elongated BECs: U. Trento (Ferrari)
- Superfluid to Mott insulator: Munich (Bloch-Schneider)
- Formation of a 3D BEC in a box-like trap: U. Cambridge (Hadzibabic)

$$
\Delta=N_{1} / N
$$





## Kibble Zurek mechanism for a 3D Bose gas

The static picture
Above the critical temperature, coherent patches of size $d$
For large $T: d \sim \lambda$ (thermal wavelength)
Approaching $T_{c}: d$ diverges as $1 /\left(T-T_{c}\right)^{\nu}$
Below $T_{c}$, full coherence: $d=$ size of the box


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\end{aligned}
$$

Below $T_{c}$, full coherence: $d=$ size of the box


## Dynamical aspects

For a finite cooling rate, this picture fails: coherence cannot be established fast enough
$\rightarrow \dot{T}$ fixes the maximal size of phase-coherent patches: topological defects (vortex lines)
may appear at the border of the patches and survive to the subsequent
$\rightarrow$ thermalization
KZ theory relates the number of topological defects $N_{t}$ to the cooling rate:
One can also look for other signatures than topological defects: Hadzibabic's talk

## $3^{\text {rd }}$ dim.: Transverse condensation.



2D+ configuration


## $3^{\text {rd }}$ dim.: Transverse condensation.



## $3^{\text {rd }}$ dim.: Transverse condensation.



## 3rd dim.:Transverse condensation.


$>$ When $\mathcal{D}_{\text {tot }}$ large enough, atoms accumulate in $j=0: \mathcal{D}_{0} \sim \mathcal{D}_{\text {tot }}$ even if $k_{B} T \gg \hbar \omega_{z}$.


## Quench cooling of a thick 2D gi <br> $$
\frac{k_{\mathrm{B}} T}{\hbar \omega_{z}}=10
$$

We cross two successive transitions:
$T_{1}$ : transverse BEC which occurs when the 3D phase-space density reaches 2.6
Atoms accumulate in the ground state $j=0$ of the $z$ motion, coherence $\sim a_{\text {ho }}$ length
$\mathrm{T}_{2}$ : superfluid BKT // full BEC corresponding to quasi-full coherence over the sample


Tentative scenario:

- Formation of independent coherent phase domains during transition $T_{1}$
- Going up to the superfluid regime allows us to visualize the vortices

