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Classification

Physics Abstracts

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The infrared emission spectrum of SiN between 2.2 and 4.4 μm

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Abstract. — The emission spectrum of a plasma of silane and nitrogen, excited by a radio frequency discharge, has been recorded with a high resolution Fourier transform spectrometer. We report here the first observation by IR spectroscopy of the $X^2\Sigma^+$ ($v = 1, 2$) and $A^2\Pi_1$ ($v = 2, 3, 4$) vibrational levels of $^{28}\text{Si}^{14}\text{N}$, through the observation of 724 unperturbed transitions of the (1-0), (2-0), (2-1), (3-1), and (4-2) vibronic bands of the $A^2\Pi_1 - X^2\Sigma^+$ system of this radical. The data were analyzed in a least squares fit procedure, with a standard deviation equal to $6 \times 10^{-4} \text{ cm}^{-1}$

Introduction.

Since silicon is chemically similar to carbon, the astrophysicists expect the silicon-containing molecules, analogous to the carbon-containing species already detected, to be present in space where until now only SiO [1], SiS [2], SiC₂ [3], SiC [4], SiC₄ [5], SiH₄ [6] and SiN [7] have been detected. Although the existence of this latter species was predicted for long by theoretical studies [8-11], it was likely observed only very recently. As the pure rotational lines of SiN have proved very difficult to detect in space, the knowledge of the A-X system, which occurs in the atmospheric infrared windows, can be of great use for the observation of this species.

The two lower electronic states of SiN, $X^2\Sigma^+$ and $A^2\Pi_1$, and the following $B^2\Sigma^+$ and $D^2\Pi$ states, are known since the first studies by Jevons [12] and Mulliken [13]. The identification of the next state $L^2\Pi$ is due to Bredohl *et al.* [14]. This paper contains references to all the previous works which will not be repeated here. A further step in the knowledge of the SiN radical was made by Yamada *et al.* [15] who observed the (0-0) band of the isotopic species of SiN by millimeter- and submillimeter-wave spectroscopy and fixed the location of the $A^2\Pi_1$ state with respect to the ground electronic state at $T_e(A^2\Pi_1) = 2032.4 \text{ cm}^{-1}$, allowing the determination of the position of the other $^2\Pi$ electronic states. All these results have been summarized by Foster [16], who fitted all the available band origin data, namely Bredohl's results [14] in the visible and near UV spectral region (3 100-5 600 Å, concerning the D-A, L-A, B-A, B-X systems) complemented with Foster's paper [17] where the B-A system and the results of the last three studies on the A-X

system in the $3.3 \mu\text{m}$ (1-0 band) region [18] and $5 \mu\text{m}$ (0-0 band) region [15, 19] are reanalyzed. The microwave spectrum of SiN in the $X^2\Sigma^+$ ($v = 0$) and ($v = 1$) levels was studied by Saito *et al.* [20] and Yamada *et al.* [15], leading to the most accurate measurements.

Experimental results.

The SiN radical is produced in a flowing mixture of N_2 and SiH_4 excited in a radio frequency discharge. The experimental setup consists of a plasma reactor [21] connected to a stepping-mode Fourier transform spectrometer [22]. The experimental conditions are the same as those which enabled the detection of the HNSi molecule [23]. The optical path is adjusted in the reactor through a White-type multireflection system to about 100 m, corresponding to a plasma column 30-50 m long. The plasma pressure and the rf power are set to respectively 70 mtorr and 720 W. The infrared emission is collected and sent to the interferometer equipped with a CaF_2 beamsplitter and liquid nitrogen cooled InSb detectors. The spectral range is limited by the optical filters and the detectors to the domain $2\,265\text{--}4\,530 \text{ cm}^{-1}$.

Two scans were recorded during seven hours each, with the maximal apodized resolution of 0.0054 cm^{-1} . They differ by the N_2 to SiH_4 ratio, that was set at 16 for the first spectrum, and at 1.6 for the second spectrum.

Figure 1 shows a part of the first spectrum. This spectrum contains a number of lines that are absent from the second one. Some of them correspond to the $A^2\Pi_1(v = 1) - X^2\Sigma^+(v = 0)$ transition of SiN given by Foster *et al.* [18]. Through an iterative procedure with direct approach, we assigned the new lines from this band and from the previously unknown (2-0), (2-1), (3-1) and (4-2) bands of this system. 724 unperturbed lines of the spectrum were observed with a signal-to-noise ratio of about 15. The distribution

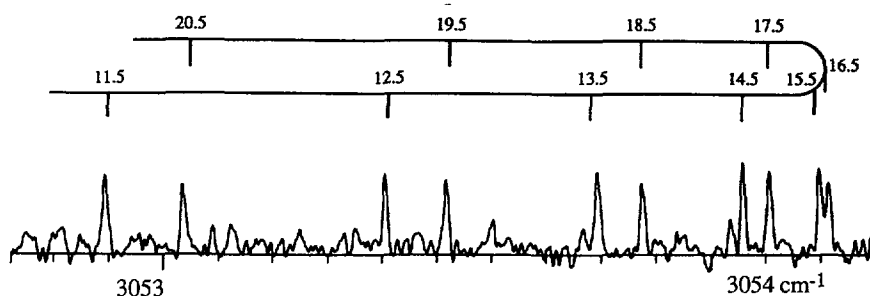


Fig. 1. — High resolution spectrum of the emission from the $(\text{N}_2 + \text{SiH}_4)$ rf discharge ($R_{21}(J)$ branch of the (1-0) band of SiN). With Resolution : 0.0054 cm^{-1} , rf Power : 720 W, P : 70 mtorr, $\text{N}_2/\text{SiH}_4 = 16$.

Table I. — *Distribution of the observed lines in the different bands.*

Band	Maximum observed J value	Number of lines
(1-0)	51.5	313
(2-0)	47.5	131
(2-1)	32.5	84
(3-1)	48.5	121
(4-2)	30.5	75

of the lines in the different bands is given in table I. We failed to detect any other band, which is not surprising since the Franck-Condon factors and the concentration are not high enough.

We observed transitions for J values up to 51.5 whereas in the former infrared studies [15, 18, 19] the highest observed J value was 31.5. This allowed us to include in the Hamiltonian high order distortion terms that have not been taken into account previously.

Theory and analysis.

The effective Hamiltonian used in the present analysis has been developed by Brown *et al.* [24] :

$$H = H^0 + H^{\text{rot}} + H^{\text{fs}}$$

H^0 is the vibrational-electronic part of the Hamiltonian.

H^{rot} is the rotational Hamiltonian and has the form :

$$H^{\text{rot}} = B(r) \mathbf{R}^2 = B(r)(\mathbf{J} - \mathbf{L} - \mathbf{S})^2.$$

As the molecule rotates, it swells in size, and we take into account the rotational stretching by adding to the rotational Hamiltonian a quartic centrifugal distortion term : $-D(r) \mathbf{R}^4$ [25].

H^{fs} is the fine-structure Hamiltonian. It takes into account two contributions, the spin-orbit and the spin-rotation interactions :

$$H^{\text{fs}} = H^{\text{so}} + H^{\text{sr}}$$

These two parts have a complex microscopic form and are described with the usual phenomenological expressions :

$$H^{\text{so}} = A(r) \mathbf{L} \cdot \mathbf{S}$$

and

$$H^{\text{sr}} = \gamma(r) \mathbf{R} \cdot \mathbf{S} = \gamma(r)(\mathbf{J} - \mathbf{L} - \mathbf{S}) \cdot \mathbf{S}.$$

We disregard here in the fine-structure Hamiltonian the spin-spin interaction H^{ss} which vanishes in the case of doublet states, and the hyperfine-structure Hamiltonian H^{hfs} (which includes the effects of the nuclear moments in interaction with the electronic and magnetic fields of the diatomic) since these effects are not observed at our resolution.

The Hamiltonian actually used is constructed by applying the perturbation theory with the Van Vleck transformation [26], and consists of two matrices representing respectively the $^2\Sigma^+$ and the $^2\Pi_1$ Born-Oppenheimer states. The parameters of the Λ -doubling interaction resulting from the perturbation of the A $^2\Pi_1$ state by both the X and B $^2\Sigma$ states, but mainly by the ground state, are included as corrections elements inside the Born-Oppenheimer states matrices. Each matrix is splitted into two symmetry blocks corresponding to e and f levels.

The effective matrix elements for the X $^2\Sigma^+$ state are :

$$\langle \Sigma | H | \Sigma \rangle = T_v - B_v Z(-Z \pm 1) - D_v Z^2(-1 \pm Z)^2 - \frac{1}{2}(-Z \pm 1)[\gamma_v - \gamma_{vJ} Z(-Z \pm 1)]$$

with $Z = J + 1/2$.

For the A $^2\Pi_1$ state, we used the effective Hamiltonian of Brown *et al.* [24], and a basis set of Hund's case (a) wavefunctions. The second order expansion of the Hamiltonians has been found to be necessary to fit our data.

The effective Hamiltonian matrix elements for the $A^2\Pi_1$ state are :

$$\begin{aligned} \langle \Pi_{1/2} | H | \Pi_{1/2} \rangle &= T_v - \frac{A_v}{2} + \left(B_v - \frac{A_{vD}}{2} \right) (X+2) - D_v (X+1)(X+4) - \\ &\quad - \frac{1}{2} [p_v + p_{vJ} J(J+1)] (-1 \pm \sqrt{X+1}) - \frac{1}{2} [q_v + q_{vJ} J(J+1)] (-X-2 \pm 2\sqrt{X+1}) \\ \langle \Pi_{3/2} | H | \Pi_{3/2} \rangle &= T_v + \frac{A_v}{2} + \left(B_v + \frac{A_{vD}}{2} \right) X - D_v X(X+1) + \frac{1}{2} X [q_v + q_{vJ} J(J+1)] \\ \langle \Pi_{3/2} | H | \Pi_{1/2} \rangle &= \left[2D_v(X+1) - B_v + \frac{1}{4} [p_v + p_{vJ} J(J+1)] + \right. \\ &\quad \left. + \frac{1}{2} [q_v + q_{vJ} J(J+1)] (-1 \pm \sqrt{X+1}) \right] \sqrt{X} \end{aligned}$$

where $X = (J + 1/2)^2 - 1$.

In all these formulae, the upper sign refers to the e-levels, and the lower sign to the f-levels. The e and f labelling is done according to the recommendations of reference [27].

The energy levels are obtained by diagonalizing the blocks of the matrix constructed for each J -value. The transitions are then calculated by applying the selection rules :

$$\begin{aligned} \Delta J &= \pm 1 : e \leftrightarrow e, f \leftrightarrow f \\ \Delta J &= 0 : e \leftrightarrow f. \end{aligned}$$

Results and discussion.

Two fits were performed for the analysis of this system. In the first one we treated simultaneously all the unperturbed line wavenumbers of the (1-0), (2-0), (2-1), and (3-1) bands, together with the data of the (0-0) band given in the references [15] and [19]. We included also the lines of the (1-0) band measured by Foster *et al.* [18] but not observed in the present work. The second calculation treated the (4-2) band alone. This two fit procedure could not be avoided since the vibronic levels involved in this band are different from the levels of the other bands. The RMS deviation for each fit is about $6 \times 10^{-4} \text{ cm}^{-1}$. The molecular constants are summarized in table II for the $X^2\Sigma^+$ state and in table III for the

Table II. — *Molecular parameters of ^{28}SiN in the $X^2\Sigma^+$ state (in cm^{-1}). The numbers in parentheses represent one standard deviation in units of the last digit.*

	v=0	v=1	v=2 ^a
T	0 ^b	1138.37420(26)	2263.83945(30) ^c
B	0.7280952(32)	0.7224925(32)	0.716972(14)
$10^5 D$	0.11760(21)	0.11808(22)	0.1226(15)
$10^2 \gamma$	1.6853(12)	1.9604(11)	1.3430(41)
$10^5 \gamma_J$	0.0116(15)	0.0505(18)	0.1310(91)

^a The results in this column come from the fit of the (4-2) band alone.

^b Fixed

^c Relative to T₄ ($A^2\Pi$). See text and Table III.

Table III. — *Molecular parameters of ^{28}SiN in the $A^2\Pi_i$ state (in cm^{-1}). The numbers in parentheses represent one standard deviation in units of the last digit.*

	v=0	v=1	v=2	v=3	v=4 ^a
T	1971.774771(90)	2991.353782(63)	3998.54016(15)	4993.30588(30)	5975.623 ^b
A	-89.09357(17)	-88.87314(12)	88.67265(21)	88.48635(28)	-88.30789(44)
$10^4 A_J$	0.549(10)	0.5451(14)	0.5887(48)	0.6688(70)	0.719(14)
B	0.6673586(24)	0.66191761(33)	0.6564594(32)	0.6509795(33)	0.645508(15)
$10^5 D$	0.11325(26)	0.113133(17)	0.11348(22)	0.11367(23)	0.1174(17)
$10^3 p$	6.302(28)	1.946(11)	0.628(18)	1.508(22)	1.291(68)
$10^4 q$	0.111(45)	0.2265(65)	0.442(12)	0.547(14)	-0.507(78)
$10^5 p_J$	0.1200(62)	0.13933(96)	0.1215(21)	0.0742(24)	0.049(15)
$10^7 q_J$	0.214(52)	0.1399(38)	0.084(11)	0.051(10)	not significant

^a The results in this column come from the fit of the (4-2) band alone.

^b Fixed at the value estimated after expanding the first four $T_v(A^2\Pi)$ in power of $(v+0.5)$.

$A^2\Pi_i$ state. All parameters were free, with the exception of the vibronic terms $T_0(X^2\Sigma^+)$ and $T_4(A^2\Pi_i)$, that were set respectively at 0 and 5 975.623 cm^{-1} . This extrapolated value of $T_4(A^2\Pi_i)$ has been estimated by expanding the first four vibrational terms of the excited state using the usual development :

$$T_v = T_e + \omega_e(v + 1/2) - \omega_e x_e(v + 1/2)^2 + \omega_e y_e(v + 1/2)^3$$

As a result of this procedure, the (4-2) band origin (3 711.78655(30) cm^{-1}) computed in the second fit made possible to fix up $T_2(X^2\Sigma^+)$ at 2 263.839(42) cm^{-1} . The equilibrium vibrational constants of the ground state were derived up to the first anharmonic correction $\omega_e x_e$, since only three vibrational terms are available.

The equilibrium rotational constants were obtained through a least squares fit of the set of B_v constants for the vibrational levels, with the usual sign conventions :

$$B_v = B_e - \alpha_e(v + 1/2) - \beta_e(v + 1/2)^2$$

Table IV shows the equilibrium constants of the X state. Our results are in good agreement with the published ones. Since the position of the $X^2\Sigma^+$ ($v = 2$) level is not known with a sufficient accuracy, we preferred not to take into account the value of $\omega_e y_e = -0.69 \times 10^{-2} \text{cm}^{-1}$ given by Foster [16]. Introducing this value leads to values for ω_e and $\omega_e x_e$ slightly lower than reported, but remaining still in the uncertainty limits of Foster.

For the $A^2\Pi_i$ state, the comparison of the vibrational constants cannot be made straightforwardly because Foster used Zare's « unique perturber approximation » [28]. Instead of applying transformation relationships between the two formulations [24], what can lead to biased results if the parameters are correlated, a second analysis was performed using Zare's Hamiltonian. This led to another set of molecular and equilibrium parameters. The results of both analysis are given in table V. The difference between the two values of $T_e(A^2\Pi)$ comes from a different definition of this constant in Zare's Hamiltonian, where it includes also the rotational constant B_e [29]. Nevertheless, our vibrational development shows quite large discrepancies with Foster's one. We ascribe these discrepancies to the fact

Table IV. — *Equilibrium parameters of ^{28}SiN in $X^2\Sigma^+$ state (in cm^{-1}). The numbers in parentheses represent one standard deviation in units of the last digit.*

	This work		Previous works	Refs
B_e	0.730927(15)		0.7311(1)	Bredohl et al.(14)
$10^3 \alpha_e$	5.685(30)		5.65(3)	Bredohl et al.(14)
$10^4 \beta_e$	4.1(12)			
T_e	0 ^a		0 ^a	
ω_e	1151.284(43)		1151.30(11)	Foster (16)
$\omega_e x_e$	6.455(21)		6.463(28)	Foster (16)
$10^2 \omega_e y_e$			-0.69(20)	Foster (16)

^aFixed.

Table V. — *Equilibrium parameters of ^{28}SiN in the $A^2\Pi_i$ state (in cm^{-1}). The numbers in parentheses represent one standard deviation in units of the last digit.*

	This work		Previous works	Refs
	Brown's hamiltonian	Zare's hamiltonian	Zare's hamiltonian	
B_e	0.6700793(93)	0.6700771(95)	0.6700904(11)	Yamada et al.(15)
$10^3 \alpha_e$	5.4344(88)	5.4345(89)	5.4546(11)	Yamada et al.(15)
$10^6 \beta_e$	5.9(17)	5.9(17)		
T_e	2031.375(16)	2032.045(49)	2032.15(6)	Foster (16)
ω_e	1031.9448(15)	1031.9393(17)	1031.65(5)	Foster (16)
$\omega_e x_e$	6.17531(95)	6.1753(11)	5.989(29)	Foster (16)
$10^2 \omega_e y_e$	-0.467(17)	-0.467(20)	-3.02(35)	Foster (16)

that the weight attributed to the A state in his calculation was too small compared to the weight of the higher states.

Conclusion.

The interpretation of the behavior of the A -doubling and spin-orbit coupling parameters in the A state and of the spin-rotation constants in the X state in term of interaction between the vibrational levels of these two states has been discussed by Foster *et al.* [18] and Yamada *et al.* [15], who ascribe the drastic and irregular changes of these constants *versus* v , specially of p , q , and γ , to the fact that the A-X separation is comparable with the vibrational energy. Table VI shows the results of the present work together with the values reported by Yamada *et al.* [15]. The two experimental sets of parameters are in good agreement. The values computed in reference [15] are obtained by expanding the theoretical expressions of the A -doubling constants as a product of electronic and vibrational integrals, and fitting the electronic integrals, while the vibrational integrals are evaluated numerically from an RKR

Table VI. — Comparison between the experimental and calculated values of ^{28}SiN for the p, q, and A constants of the $A^2\Pi$ state, and the γ constants of the $X^2\Sigma^+$ state (cm^{-1}).

	This work	Previous work ⁽¹⁵⁾	Computed value ⁽¹⁵⁾
$10^3 p_v$			
v=0	6.302(28)	6.261(18)	6.244
v=1	1.946(11)	1.892(31)	1.940
v=2	-0.628(18)		-0.723
v=3	-1.508(22)		-1.831
v=4	-1.291(68)		-1.642
$10^4 q_v$			
v=0	0.111(45)	0.121(37)	0.185
v=1	-0.2265(65)	-0.159(17)	-0.174
v=2	-0.442(12)		-0.396
v=3	-0.547(14)		-0.488
v=4	-0.507(78)		-0.471
A_v			
v=0	-89.09357(17)	-89.09657(18)	-89.097
v=1	88.87314(12)	-88.87320(22)	-88.873
v=2	-88.67265(21)	88.65(2) ^a	-88.665
v=3	-88.48635(28)		-88.469
v=4	-88.30789(44)	-88.32(2) ^b	-88.284
$10^2 \gamma_v$			
v=0	1.6853(12)	1.684862(57)	1.685
v=1	1.9604(11)	1.97317(17)	1.966
v=2	1.3430(41)		1.384

^a Ref. 14

^b Ref. 17

calculation. They do not match closely the experimental results, but they show the same behavior, and can be useful for the detection of the hot bands, since the difficulties in observing the lines do not come from the approximated values of the constants but from the weakness of the transitions. For this reason, we did not perform similar calculations, even though our experimental results should improve the prediction of the parameters. Since the perturbation is expected to be maximal [15] in the $X^2\Sigma^+ v = 7, 8$ and $A^2\Pi v = 8, 9$ levels, the knowledge of some more vibrational levels, in the ground state as well as in the excited state, is necessary to obtain a good description of the crossing between the $A^2\Pi$ and $X^2\Sigma^+$ states of SiN.

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Note : The spectral positions of the transitions have been deposited in the Editorial Office of the Journal de Physique. These data are also available upon request to M. Elhanine.

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