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R. Świetlik, A. Lapiński, L. Kushch, E. Yagubskii. Electronic and Vibrational Spectra of Et4N[Au(dmit)4TCNQ] Crystals. Journal de Physique I, 1996, 6 (12), pp.1643-1653. 10.1051/jp1:1996180 . jpa-00247271

HAL Id: jpa-00247271 https://hal.science/jpa-00247271v1

Submitted on 4 Feb 2008

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Electronic and Vibrational Spectra of $Et_4N[Au(dmit)_2TCNQ]$ Crystals

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(Received 7 February 1996, revised 23 April 1996, accepted 3 June 1996)

PACS.78.30.Jw – Organic solids, polymers PACS.78.40.Me – Organic solids and polymers

Abstract. — We present the results of spectral studies of $Et_4N[Au(dmit)_2TCNQ]$ salt containing alternating stacks of $Au(dmit)_2$ and TCNQ molecules. The polarized transmission spectra of single crystals are measured in the IR region (650 to 6500 cm⁻¹) as a function of temperature (90 to 300 K). Moreover, the powder absorption spectra in the frequency range (400 to 40000 cm⁻¹) are reported. From spectral evidence results that transferred electrons are localized on $Au(dmit)_2$ molecules in the mixed stacks. The intrastack charge transfer between $Au(dmit)_2$ and TCNQ molecules (~ 8000 cm⁻¹) and the interstack charge transfer transition between $Au(dmit)_2$ molecules (~ 5500 cm⁻¹) are found. Interactions of electrons with intramolecular vibrations of TCNQ (ag modes) and $Au(dmit)_2$ (C = C and C = S stretching) are suggested and discussed.

1. Introduction

The metal complexes $M(dmit)_2$ (M = Ni. Pd, Pt, Au and dmit = 4,5 dimercapto-1,3-dithiole-2-thione; Fig. 1) are known to form many highly conducting salts with different donor molecules or closed-shell cations [1]. In this series three Ni(dmit)₂ salts and three Pt(dmit)₂ salts exhibit superconductivity under pressure and one Ni(dmit)₂ salt becomes superconductor at ambient pressure [1,2]. Close S · S contacts, shorter than the sum of the corresponding van der Waals radii, promote effective molecular overlap inside and/or between the stacks. Up to now the most widely studied are compounds based on the dmit complexes of Ni, Pd and Pt but an interesting series of highly conducting salts based on Au(dmit)₂ has been also synthetized [3].

The TCNQ (tetracyano-p-chinodimethane) is another important organic acceptor yielding a great number of conducting charge transfer salts with different cations [4]. The TCNQ molecules stack face-to-face to form linear chains and the overlap of π -molecular orbitals along the stacking direction greatly exceeds interstack coupling. The consequence of such specific crystal structure is a high anisotropy of the physical properties, *i.e.* quasi-one-dimensionality.

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Fig. 1. — $M(dmit)_2$ and TCNQ molecules (M = Ni, Pd, Pt, Au).



Fig. 2. — Crystal structure of $Et_4N[Au(dmit)_2TCNQ]$ projected along the *c*-axis (above) and *b*-axis (below) [3].

Both Au(dmit)₂ and TCNQ form conducting salts with Et₄N⁺ (tetraethylammonium) cation. The Et₄N(TCNQ)₂ crystal is a quasi-one-dimensional semiconductor [5], whereas the Et₄N [Au(dmit)₂]₂ compound is known to exist in two modifications α and β , exhibiting metallic and semiconducting properties, respectively [3]. Moreover, an interesting crystal Et₄N [Au(dmit)₂TCNQ], containing mixed [Au(dmit)₂TCNQ]⁻ anion, has been also synthetized [3]. The crystal structure data of the Et₄N[Au(dmit)₂TCNQ]⁻ anion, has been also synthetized [3]. The crystal structure data of the Et₄N[Au(dmit)₂TCNQ] are following: triclinic unit cell, space group P1. a = 6.490(1), b = 7.330(1), c = 17.802(4) Å, $\alpha = 86.51(3)$, $\beta = 81.75(3)$, $\gamma = 78.24(3)^{\circ}$, V = 820.08 Å³, Z = 1 [3]. The Au(dmit)₂ and TCNQ molecules alternate each other forming stacks along b direction (Fig. 2). The stacks are arranged in layers parallel to ab plane alternating with the layers of Et₄N cations in c direction. Slightly shortened S · S contacts (3.637 Å) are observed between the stacks in the anion layers. The low room temperature conductivity of this crystal (~ 10⁻⁵ S/cm) gives evidence of strong charge localization in the mixed [Au(dmit)₂TCNQ] stack.

The extraordinary crystal structure and the presence of mixed anion prompted us to perform the spectral investigations of $Et_4N[Au(dmit)_2TCNQ]$ crystals. Optical properties of the alternating stack and first of all the charge distribution are the most interesting problems. It is well known that infrared spectra of crystalline organic conducting systems are dominated by vibrational bands which are a consequence of electron-molecular vibration (EMV) coupling phenomena [6]. The eventual presence of EMV interactions in $Et_4N[Au(dmit)_2TCNQ]$ crystal is another interesting question.

2. Experimental

The plate-shaped Et₄N[Au(dmit)₂TCNQ] crystals of typical dimensions $1 \times 1 \times 0.03 \text{ mm}^3$ were grown from acetonitrile. The compound Et₄N[Au(dmit)₂] in the amount of 0.0216 g $(3 \times 10^{-5} \text{ mol})$ was dissolved in acetonitrile (30 ml) at room temperature. After heating of the solution up to t = 72 °C the TCNQ (0.02448 g, 1.2×10^{-4} mol) was added. Then the vessel with solution was placed into a thermostat (72 °C) and cooled down slowly to room temperature at the rate of 1 °C/hour. The obtained crystals were filtered off, washed up with acetonitrile and acetone and dried at t = 40 °C.

Polarized transmission spectra of single crystals, for electrical vector of polarized light parallel and perpendicular to the stacking axis, were measured within the frequency range (650 to 6500 cm^{-1}) by using the FT-IR Perkin Elmer 1725 X spectrometer equipped with a FT-IR microscope. We studied also the spectra of powders dispersed in KBr pellets in the range 400 to 40000 cm⁻¹; three different spectrometers were used: FT-IR Perkin Elmer 1725 X (400 to 7400 cm⁻¹), Specord NIR (4000 to 13500 cm⁻¹) and Specord UV-VIS (13000 to 40000 cm⁻¹). Moreover, temperature dependence of the middle IR transmission of both single crystals and powders in KBr pellets was studied down to T = 90 K. For comparison we measured also room temperature absorption spectra of the compounds $\text{Et}_4\text{N}[\text{Au}(\text{dmit})_2]$ and Et_4NI_3 using KBr pellet technique. Detailed analysis of infrared spectra and their temperature dependencies was performed after decomposition of the selected spectral features. For this purpose a standard PEAKFIT program was used and the bands were fitted by Lorentzian and/or Gaussian functions.

The value of average charge residing on TCNQ molecule can be estimated from the position of Raman band assigned to the totally symmetric TCNQ mode $a_g\nu_4$ [7]. For this reason we performed the Raman measurements of $Et_4N[Au(dmit)_2TCNQ]$ powders at room temperature using FT-IR Bruker IFS 66 spectrometer equipped with FRA 106 Raman accessory and Nd:YAG laser ($\lambda = 1.064 \ \mu m$).

3. Results and Discussion

3.1. ELECTRONIC TRANSITIONS AND CHARGE DISTRIBUTION. — The electronic absorption spectra of $Et_4N[Au(dmit)_2TCNQ]$ and $Et_4N[Au(dmit)_2]$ powders in KBr pellets are shown in Figure 3. In both compounds three strong, similar bands D. E, F are observed and the salt $Et_4N[Au(dmit)_2TCNQ]$ exhibits additional, weak electronic absorptions A, B, C. The electronic transitions fall into two classes: those at higher energies result from intramolecular electronic excitations and the lower-energy bands are due to the interactions among the molecules and are called charge transfer bands. The assignment of electronic bands has been performed by comparison with the interpreted electronic spectra of other $Au(dmit)_2$ [8] and TCNQ [9,10] salts.

The strong band F at the frequency 33000 to 34000 cm⁻¹ with a shoulder at 35500 cm⁻¹ is related to a π - π^* transition of the dmit ligand. Another π - π^* band of the ligand is observed at 21000 cm⁻¹ for Et₄N[Au(dmit)₂TCNQ] and at slightly lower frequency 20500 cm⁻¹ for Et₄N[Au(dmit)₂] (band D in Fig. 3). The electronic absorptions F and D correspond to the bands at about 34200 cm⁻¹ and 21400 cm⁻¹, respectively, for Bu₄N[Au(dmit)₂]₂ in acetonitrile solution [8]. The band E at 27000 cm⁻¹ with a weak shoulder at 25000 cm⁻¹ can be assigned



Fig. 3. — Electronic absorption spectra of $Et_4N[Au(dmit)_2TCNQ]$ and $Et_4N[Au(dmit)_2]$ powders in KBr pellets (weight concentration 1:4000).

to the Au^{III} \leftarrow S charge transfer excitation but in this frequency region the intramolecular excitation of neutral TCNQ⁰ usually appears (27000 to 28000 cm⁻¹ for different TCNQ salts [9,10]). Therefore, in Et₄N[Au(dmit)₂TCNQ] the band E is a superposition of the Au^{III} \leftarrow S charge transfer and TCNQ⁰ intramolecular excitation. Analogous electronic absorption was found at 28600 cm⁻¹ with a shoulder at 26300 cm⁻¹ for Bu₄N[Au(dmit)₂]₂ in acetonitrile solution [8].

A striking feature of the electronic spectrum of $\text{Et}_4\text{N}[\text{Au}(\text{dmit})_2\text{TCNQ}]$ is the absence of a band at about 16000 cm⁻¹ which could be expected from comparison with the electronic spectra of other TCNQ salts — at this frequency the intramolecular excitation of TCNQ⁻ anion is usually observed [9,10]. This fact suggests that the average charge density on TCNQ molecules is close or equal to zero, *i.e.* the transferred electron is mostly localized on Au(dmit)₂ molecule. This conclusion is confirmed by our Raman measurements of $\text{Et}_4\text{N}[\text{Au}(\text{dmit})_2\text{TCNQ}]$ since the $a_g\nu_4$ Raman band is situated at 1450 cm⁻¹, *i.e.* nearly at the frequency corresponding to neutral TCNQ⁰ molecule (1454 cm⁻¹ [7]).

The relatively weak doublet C (12000 and 13000 cm⁻¹) is without doubt related to TCNQ. Analogous electronic feature is always found in spectra of TCNQ salts [9,10], though usually at slightly lower frequency. Since the publication of the paper [9] by Torrance *et al.*, such electronic absorption is generally assigned to the charge transfer process (TCNQ⁻+TCNQ⁻ \rightarrow TCNQ⁰ + TCNQ²⁻). The specific structure of the stacks in Et₄N[Au(dmit)₂TCNQ] and the estimated charge distribution prove that such charge transfer is impossible. Yakushi *et al.* [11] attributed band C to the lowest π - π * band of TCNQ⁻ ion but in Et₄N[Au(dmit)₂TCNQ] there are only neutral TCNQ⁰ molecules in the stacks. In conclusion the nature of TCNQ electronic absorption C remains unclear.



Fig. 4. — Polarized infrared absorption spectra at room temperature of $Et_4N[Au(dmit)_2TCNQ]$ single crystal in the region 6500 - 2100 cm⁻¹. The electrical vector of polarized light is parallel ($\mathbf{E} \parallel \mathbf{b}$) and perpendicular ($\mathbf{E} \perp \mathbf{b}$) to the stacking axis (note the change of frequency scale at 3200 cm⁻¹)

The electronic bands A (5500 cm⁻¹) and B (8000 cm⁻¹) in powder spectra (Fig. 3) were also studied using the transmission technique on thin single crystals (Fig. 4). The band A is polarized perpendicular and the band B parallel to the stacking direction b and we relate them to inter- and intrastack charge transfer, respectively. The band A is very weak and corresponds to the side-by-side charge transfer between Au(dmit)₂ and TCNQ inside the stack. In single crystal spectrum the band A is centered at lower frequency (~ 5100 cm⁻¹) than in powder spectrum (~ 5500 cm⁻¹); the band B seems to lie also at lower frequency but it cannot be positioned more precisely since the measurement was made up to 6500 cm⁻¹. In powder spectrum band A is weaker than band B but this difference is not so large as in the single crystal spectrum. This could be an indication that band A is polarized perpendicularly to the (001) crystal face and in the single crystal spectrum its component in the b direction is only observed. Since such electronic transition is rather unexpected, it seems that the difference of relative intensities between powder and single crystal spectrum is due to an influence of powdering effect.

3.2. VIBRATIONAL MODES. — The specific feature of the organic conductors is the existence of strong electron-molecular vibration (EMV) interactions which manifest themselves by strong bands in the infrared spectra. Obviously, the EMV coupling takes place when the frequency of vibrational mode lies within the region of a suitable, strong charge transfer band. In the $Et_4N[Au(dmit)_2TCNQ]$ salt the high frequency of intrastack charge transfer (~ 8000 cm⁻¹) and the very low intensity of interstack charge transfer (~ 5100 cm⁻¹) give rise to the weakness of expected infrared bands related to EMV coupling phenomena. The IR spectrum of



Fig. 5. — Infrared absorption spectra at room temperature of $\text{Et}_4N[\text{Au}(\text{dmit})_2]$ and $\text{Et}_4N[\text{Au}(\text{dmit})_2]$ TCNQ] powders in KBr pellets in the region $1650 - 400 \text{ cm}^{-1}$ (weight concentration 1:1000).

 $Et_4N[Au(dmit)_2TCNQ]$ should be dominated by IR active vibrations of $Au(dmit)_2$, TCNQ as well as cation molecules, nevertheless the presence of weak EMV features is also expected.

The IR absorption spectra of $Et_4N[Au(dmit)_2]$ and $Et_4N[Au(dmit)_2TCNQ]$ powders are illustrated in Figure 5. The spectra of both compounds are similar and several additional bands In $Et_4N[Au(dmit)_2TCNQ]$ are due to TCNQ molecule. In $Et_4N[Au(dmit)_2]$ spectrum the bands are of higher intensity, especially those at 1477 cm^{-1} and 1061 cm^{-1} related to C = C and C = Sstretching vibrations of Au(dmit)₂ molecule, respectively. The structure of $Et_4N[Au(dmit)_2]$ is not known but it may be supposed that Au(dmit)₂ molecules interact stronger (probably faceto-face interaction) than in $Et_4N[Au(dmit)_2TCNQ]$, therefore the IR effect of EMV coupling is stronger. The IR transmission spectra on $Et_4N[Au(dmit)_2TCNQ]$ single crystals were investigated in two polarizations: the electrical vector of polarized light was parallel or perpendicular to the stacking direction b (Figs. 4 and 6). The transmission method on partially transparent single crystal is very useful for measurements of very weak modes, therefore, the single crystal spectra are very reach and in comparison to powder spectra the relative intensities of many bands are different. The FT-NIR Raman spectrum of Et₄N[Au(dmit)₂TCNQ] is dominated by the strong TCNQ bands: 1609 cm⁻¹ ($a_g \nu_3$), 1450 cm⁻¹ ($a_g \nu_4$), 1196 cm⁻¹ ($a_g \nu_5$), 952 cm⁻¹ $(a_g \nu_6)$, 708 cm⁻¹ $(a_g \nu_7)$ and 345 cm⁻¹ $(a_g \nu_9)$; the other strong bands at 906 cm⁻¹, 497 cm⁻¹ and 325 cm⁻¹ may be related to Au(dmit)₂ molecule.

The frequencies of IR bands at room temperature and their assignment are listed in Table I. The attributions of $Au(dmit)_2$ vibrational bands are based upon those previously established for other molecules containing dmit ligand [12–15] or analogous molecules [16] since the normal

TCNQ ⁰ [17]		Et₄N[Au(dmit)₂TC		CNQ]	Et₄N[Au(dmit) ₂]	Assignment
obs	calc	crystal		powder	powder	
		Elb	E⊥ b			
3065 3053	3041 3050	3059 3043 3006 2986 2976 2946 2910	3059 3043 3006 2986 2946	3058 3043 3005 2986 2973 2945 2922 2852	3000 2976 2940	$\begin{array}{c} TCNQ \ b_{1u} v_{18} \\ TCNQ \ b_{2u} v_{32} \\ Et_4 N \end{array}$
2229 2228 2228	2230 2226 2230	2217 2214	2217 2213	2218		$\begin{array}{c} TCNQ \ a_{g}v_{2} \\ TCNQ \ b_{2u}v_{33} \\ TCNQ \ b_{1u}v_{19} \end{array}$
1602 1545 1540	1602 1542 1530	1598 1543 1537 1483 1474	1921 1614 1543 1535	1543 1537	1483	TCNQ $a_{g}v_{3}$ TCNQ $b_{1u}v_{20}$ TCNQ $b_{2u}v_{34}$ Et ₁ N
1454	1454	1461 1457 1450 1446 1432	1457 1454 1446 1431	1456 1450 1446 1432	1477 1457 1446 1435	TCNQ $a_{s}v_{4}$ Au(dmt) ₂ C=C Au(dmt) ₂ C=C Et ₄ N Et ₄ N
1405	1404	1389 1385	1394 1385	1391	1397 1391	$ \begin{array}{c} \operatorname{Et}_4 \mathrm{N} \\ \operatorname{TCNQ} \mathrm{b}_{1\mathrm{u}} \mathrm{v}_{21} \\ \operatorname{Et}_4 \mathrm{N} \\ \operatorname{Et}_4 \mathrm{N} \end{array} $
1354	1360	1363 1353 1331	1363 1353 1331 1268 1246	1363	1365	Et ₄ N TCNQ b _{2u} v ₃₅
1207	1207	1195 1185 1174	1185 1168	1181 1172	1182 1171	TCNQ agv5 Et₄N Et₄N
1125	1119	1054	1128 1118 1110 1094			TCNQ $b_{2u}v_{37}$
998	1013	1034 1061 1052 1035 1028 1006 996 991	1059 1051 1035 1030 1006 998 992	1056 1050 1033 1026 1005 998	1061 1027 1001	$ \begin{array}{l} Et_4N\\ Et_4N\\ Au(dmt)_2\ C=S\\ Au(dmt)_2\ C=S\\ Au(dmt)_2\ C=S\\ TCNQ\ b_{1u}v_{22}\\ Et_4N \end{array} $
962 859	961	905 831 783	961 907 787 776	905 831 784 778 520 514 468	909 788 514 468	$\begin{array}{l} TCNQ \ b_{1u} \mathbf{v}_{23} \\ Au(dmit)_2 \ C-S \\ TCNQ \ b_{3u} \mathbf{v}_{\varsigma_0} \\ Et_4N \\ Et_4N \\ Au(dmit)_2 \ Au-dmit \\ Au(dmit)_2 \ Au-dmit \\ Au(dmit)_2 \ Au-S \end{array}$

Table I. — Frequencies (cm^{-1}) and assignments of the infrared spectra of $Et_4N[Au(dmit)_2]$ and $Et_4N[Au(dmit)_2TCNQ]$ at room temperature (T = 300 K).



Fig. 6. — Polarized infrared absorption spectra (1670 – 650 cm⁻¹) at room temperature of Et₄N [Au(dmit)₂TCNQ] single crystal for the electrical vector of polarized light parallel ($\mathbf{E} \parallel \mathbf{b}$) and perpendicular ($\mathbf{E} \perp \mathbf{b}$) to the stacking direction.

mode analysis of $M(dmit)_2$ molecule has not been done as yet. The TCNQ bands are assigned according to the paper [17] and the bands of Et_4N cation on basis of our IR data on Et_4NI_3 salt. In order to make the best use of experimental evidence the band attribution has been carried out very carefully considering orientation of the molecules with respect to the crystal face (001), the powder and single crystal spectra, the relative intensities for different polarizations as well as the temperature dependencies.

From the crystallographic data results that molecular planes of Au(dmit)₂ and TCNQ are nearly perpendicular to the (001) crystal face and the long molecular axes are tilted in this way that they have components in the direction perpendicular to the stacks (Fig. 2). Such orientation of the molecules with respect to the (001) crystal face causes that in-plane, IR active b_{1u} and b_{2u} modes should be mainly registered in the direction perpendicular to the stacks whereas the out-of-plane b_{3u} modes parallel to the stacks. This agrees with the polarization of the bands related to IR active TCNQ modes of b_{1u} and b_{2u} symmetry, and their frequencies are close to those for neutral molecule (Fig. 6 and Tab. I). The band 831 cm⁻¹ polarized parallel to the stacks, corresponds to the out-of-plane $b_{3u}\nu_{50}$ TCNQ mode and exhibits larger shift from the TCNQ⁰ frequency; it seems to us that it can be due to an influence of the neighbouring Au(dmit)₂ anions. In the direction parallel to the stacks **E** || **b** we find several bands which can be attributed to the totally symmetric a_g modes of TCNQ molecule. The best example is the mode $a_g\nu_3 = 1598 \text{ cm}^{-1}$ but also $a_g\nu_2 = 2217 \text{ cm}^{-1}$ and $a_g\nu_5 = 1195 \text{ cm}^{-1}$ are well separated by the decomposition procedure. We observe also $a_g\nu_4 = 1461 \text{ cm}^{-1}$ but this band is difficult to analyse since it strongly overlaps with Au(dmit)₂ and Et₄N bands. The a_g



Fig. 7. — Temperature dependence of the relative integral intensity of some IR bands in the polarized absorption spectra of $Et_4N[Au(dmit)_2TCNQ]$ single crystal. Note the bands assigned to the $Au(dmit)_2$ and TCNQ molecules and different polarizations (**E** || **b** and **E** \perp **b**).

modes are normally IR-inactive but their weak IR activity may appear due to the influence of lower symmetry crystal field or to the effect of coupling to charge transfer. To answer this question we studied temperature evolution of the IR spectra assuming that the bands being a consequence of EMV coupling phenomenon strongly increase the intensity when temperature decreases. The temperature dependence of the relative integral intensity for different TCNQ and Au(dmit)₂ bands is shown in Figure 7. The intensity of TCNQ a_g bands grows strongly when the temperature decreases, what suggests that the a_g modes are activated by coupling with charge transfer transition between Au(dmit)₂ and TCNQ in stacks. It is interesting that the intensities of the greater part of b_{1u} and b_{2u} TCNQ bands grow weaker when temperature decreases (Fig. 7) with the exception of b_{1u} $\nu_{20} = 1543$ cm⁻¹ and b_{2u} $\nu_{33} = 2217$ cm⁻¹ (b_{2u} ν_{33} coincides with a_g $\nu_2 = 2217$ cm⁻¹ and it can be the reason of different behaviour). We do not understand the temperature dependence of TCNQ IR active modes — may be it is a sign of continuous structural modifications and/or changes of charge distribution.

We studied also the temperature evolution of the Au(dmit)₂ bands in the IR spectra of Et₄N[Au(dmit)₂TCNQ]. For polarization parallel to the stacks $\mathbf{E} \parallel \mathbf{b}$ the Au(dmit)₂ band at 1450 cm⁻¹ (C = C stretching) exhibits temperature changes analogous to those of TCNQ a_g modes. We suggest that it is also a consequence of the coupling of Au(dmit)₂ vibration to charge transfer transition within the stacks. For polarization $\mathbf{E} \perp \mathbf{b}$ similar strong temperature dependence has not been found for any band, *i.e.* the Au(dmit)₂ vibrational modes do not interact with interstack charge transfer.



Fig. 8. — Temperature dependence of the relative integral intensity of some IR bands in the absorption spectra of $Et_4N[Au(dmit)_2TCNQ]$ powders in KBr pellet (note the Au(dmit)_2 and TCNQ bands).

Another interesting fragment of $Et_4N[Au(dmit)_2TCNQ]$ spectrum is the region of C = Sstretching vibrations $(1070 - 1020 \text{ cm}^{-1})$. The comparison of powder and single crystal spectra shows that the C = S bands are weak in single crystal spectra for both polarizations but are very strong in powder spectrum (Figs. 5 and 6). It means that these bands are polarized nearly perpendicularly to the (001) crystal face, probably parallel to the long molecular axis. For this reason the C = S bands were mainly analysed in the spectra of $Et_4N[Au(dmit)_2TCNQ]$ powders. In Figure 8 we show temperature dependence of the relative integral intensity of C = S, C-S and Au-S bands of Au(dmit)₂ molecule, and $b_{3u}\nu_{50}$ band of TCNQ molecule as found from the powder spectrum. The C = S band at 1033 cm⁻¹ grows strongly when temperature decreases, whereas the bands 1050 cm^{-1} and 1026 cm^{-1} decrease their intensity (similar behaviour was found for analogous bands in single crystal spectra). The high intensity and strong temperature dependence of the band 1033 cm^{-1} could be due to the strong interaction of C = S vibration with the electrons. It is important to emphasize that the C = S mode seems to be polarized along the long axis of $Au(dmit)_2$ molecule, *i.e.* it can be b_{1u} mode. Analogous suggestion about coupling of electrons with b_{1u} modes was recently made for crystals based on BEDT-TTF molecule [18, 19].

4. Conclusions

We have performed the studies of the electronic and vibrational spectra of $Et_4N[Au(dmit)_2 TCNQ]$ crystal with the extraordinary stacks consisting of alternating TCNQ and $Au(dmit)_2$ molecules. In the analysis of the spectra we mainly focused our attention on the problem

of charge distribution and the electron-molecular vibration coupling phenomena. From the electronic and Raman spectra results that the transferred electrons are localized on Au(dmit)₂ molecule. In the low-energy region of electronic spectrum two charge transfer bands are observed: the intrastack transition between Au(dmit)₂ and TCNQ molecules at about 8000 cm⁻¹ and the weak interstack transition between neighboring Au(dmit)₂ molecules at about 5000 cm⁻¹. We suggest that the C = C vibrations of Au(dmit)₂ and the a_g modes of TCNQ couple with the intrastack charge transfer excitation. Any signs of the electron-molecular vibration coupling with the interstack charge transfer are not found. On the other hand, it is supposed the interaction of electrons with C = S vibrations polarized along the long axis of Au(dmit)₂ molecule. The intensities of the bands related to IR active vibrations of Au(dmit)₂ and TCNQ molecules decrease when the crystal is cooled down.

Acknowledgments

We are indebted to Dr. M. Połomska for help in measurements of the FT-NIR Raman spectra. L.A.K and E.B.Y acknowledge Russian Foundation for a financial support (Grant RF93-030).

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