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## **<sup>35</sup>Cl NQR and calorimetric reinvestigation of the incommensurate phase of bis(4-chlorophenyl) sulfone : evidence for no lock-in transition**

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**Abstract.** — <sup>35</sup>Cl NQR and calorimetric measurements in BCPS have been reinvestigated between 77 K and room temperature (NQR) and between 100 K and 180 K ( $C_p$ ). It is shown that, in contradiction with conclusions found in literature, there is no lock-in transition at 115 K. These new results are in agreement with those obtained with other experimental techniques : Raman scattering, X-ray diffraction, neutron diffraction and proton NMR ; the latter indicate that the incommensurate phase, appearing on cooling at 150 K, persists down to 4.2 K, at least. The shape of the <sup>35</sup>Cl NQR spectrum in the BCPS incommensurate phase shows that the frequency of a chlorine nucleus is a function of the order parameter, through linear and quadratic terms.

### **1. Introduction.**

For about fifteen years, a great amount of experimental and theoretical studies concerning the structure and dynamics of incommensurate systems have been reported [1-4]. The collective aspect of the dynamics has been studied by coherent neutron scattering and light scattering and is now rather well understood, even in quite complex systems [1]. Resonance methods, such as NMR, NQR or EPR, have also been very fruitful for the analysis of the different aspects involved in these structures [1, 5, 6]. The resonance lineshape gives specific informations on the local nature of the modulation wave whereas spin-lattice relaxation measurements yield the spectral density of the corresponding excitation modes. Some fundamental questions still remain open concerning especially the influence of the characteristic excitations on local properties : modification of the state density, Debye-Waller factor and heat capacity.. Another point is the analysis of the evolution of the lattice dynamics in the incommensurate phase when going from a sinusoidal regime to a soliton-like one. The study of these questions requires a compound where the incommensurate phase exists over a large temperature region and, even

better, persists down to lowest temperatures. There are few examples of structurally incommensurate crystals where this is indeed the case and they always concern purely displacive phase transitions ( $\text{ThBr}_4$ ,  $\text{ThCl}_4$ , biphenyl, ...). In biphenyl and  $\text{ThBr}_4$ , this property allows the direct observation of the phason branch by coherent neutron scattering because of the low damping of these excitations at low temperature [7, 8]. In biphenyl, the appearance of satellites up to the third order in phase III below 20 K was explained by higher-order components in the modulation wave. Compounds from the  $\text{A}_2\text{BX}_4$  family, on the opposite, do not remain incommensurate down to the lowest temperatures but several results, often obtained by resonance techniques [1, 4], clearly indicate an evolution to non-sinusoidal regimes with soliton-like character in these compounds : unfortunately the large damping of the excitations does not permit direct analysis in these cases.

In this paper, we study a molecular compound in which an incommensurate phase has been found rather recently upon cooling below 150 K : bis(4-chlorophenyl) sulfone or BCPS ; we were led to this study because contradictory results are found in literature concerning the incommensurate phase in this compound. Let us first consider the major features of BCPS and then the contradictory results.

At room temperature, BCPS crystallizes in a monoclinic structure of space group  $C2/c$  with two molecules per unit cell. The molecule,  $(\text{ClC}_6\text{H}_4)_2\text{SO}_2$ , is made of two chlorophenyl groups connected to a sulfur atom located on the twofold symmetry axis of the molecule (Fig. 1) ; the dihedral angle between each benzene ring and the Cl-S-Cl plane is about  $84^\circ$ . Following the intensities of the first-order satellites reflections obtained by X-ray [9] and neutron [10] diffraction, it seems that the incommensurate phase exists from 150 K down to the lowest studied temperature (13 K). Direct observation of rather strong higher-order satellites by neutron diffraction below about 110 K [10] is explained by a departure from a sinusoidal regime in the incommensurate phase. Raman scattering clearly shows a soft mode that the authors assumed to be related to a twisting of the phenyl groups [11] and has an exceptionally low damping [12]. Some analogies with biphenyl were then considered as biphenyl is also a molecular compound which presents a second order displacive incommensurate phase with an instability associated to a soft internal degree of freedom. A crystallographic study of the BCPS incommensurate phase is now being made [13] and will inform us on the nature of the

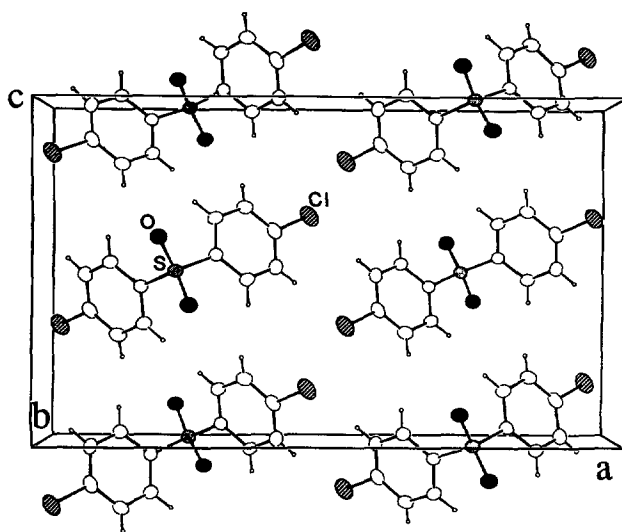


Fig. 1. — Crystalline structure of bis(4-chlorophenyl) sulfone.

involved motion. As in biphenyl [14], proton NMR measurements also led to the existence of a gapless phason in the incommensurate phase of BCPS [15].

Concerning the extension of the incommensurate phase, results from  $^{35}\text{Cl}$  NQR spectroscopy and reported in two independently published papers [16, 17] would indicate a lock-in transition at 115 K. Such a transition was also suggested by AC-calorimetry measurements showing a  $C_p$  anomaly at this temperature [18]. In order to clarify the contradiction between the mentioned lock-in transition at 115 K and the persistence of the incommensurate phase down to the lowest studied temperature, we have reinvestigated the thermodynamic behaviour of this transition. The second part of this paper deals with an experimental reanalysis of the  $^{35}\text{Cl}$  NQR results in this compound.

## 2. Differential scanning calorimetry study.

**2.1 EXPERIMENTAL.** — Bis(4-chlorophenyl) sulfone was purchased from Aldrich-Chemicals and next purified by recrystallization from benzene solution and zone refining ; single crystals were grown using two techniques, the Bridgman method and crystallization from a benzene solution. The calorimetric measurements were performed in the temperature range extending from 100 K to 180 K with a Perkin-Elmer DSC 7 differential scanning calorimeter controlled by a Digital DEC Station 425C micro-computer for data acquisition and processing. The measurement head was immersed in liquid nitrogen and flushed with dry helium gas. The temperature calibration of the instrument was made against the transition temperature (186.1 K) and melting point (279.9 K) of cyclohexane [19]. Energy was calibrated from the enthalpy of the phase transition of cyclohexane ( $79.52 \text{ J.g}^{-1}$ ). The calibration exactness was then checked against a sapphire single crystal plate of about 130 mg using the aluminium oxide heat capacity standards at 1 atmosphere [20]. Single crystal samples of BCPS cut from a large crystal and weighting from 60 mg to 80 mg, were carefully encapsuled in  $50 \mu\text{l}$  aluminium pans in order to ensure good heat transfer. The specific heat measurements were carried out at  $5 \text{ K.min}^{-1}$  and  $10 \text{ K.min}^{-1}$  rates on crystals grown from the melt and benzene solution. Only results obtained with the  $10 \text{ K.min}^{-1}$  scan rate are shown in figures 2 and 3, because with  $5 \text{ K.min}^{-1}$  results are quite similar though a little more noisy.

**2.2 RESULTS AND DISCUSSION.** — The specific heat curves recorded at increasing temperature in the temperature range 100 K-180 K for the Bridgman and the benzene solution grown crystals are shown in figure 2. Apart a small difference discussed below, the two samples present qualitatively similar behaviours : considering decreasing temperature, a single small endothermic anomaly with a rapid increase around 150 K and a smooth decay spreading over about 40 K. The fundamental difference between this result that appears in literature [18] is the absence, in the present study, of any anomaly around 115 K. As further discussed below, this is a strong argument for the nonexistence of a lock-in transition in BCPS at this temperature. The  $C_p$  anomaly beginning at 150 K clearly results from the second order incommensurate phase transition. It is worth mentioning that the shape of the anomaly depends on the method of preparation of the sample. Whereas there is a nice typical signature for the Bridgman grown single crystal giving  $T_1 = 150 \text{ K}$ , the signature is less sharp for the solution grown crystal ; the transition temperature is then more difficult to define and is around 147 K. This actually shows the important role which impurities or crystalline defects may have on the incommensurate phase transition and perhaps also explains the difference with the previous thermodynamic study. The excess of  $C_p$  is obtained from subtraction of the normal non transitional heat capacity which is described by an extrapolated seventh order polynomial curve following usual procedure. The result is shown in figure 3, for the Bridgman grown single crystal. The values of the enthalpy and the entropy of the transition are respectively  $52 \text{ J.mole}^{-1}$  and

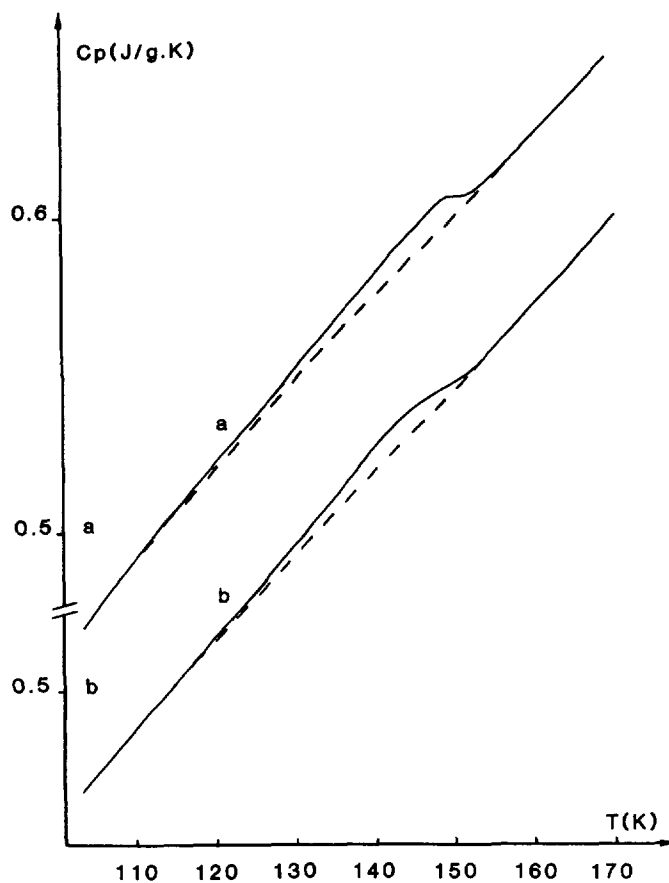


Fig. 2. — Specific heats measured on BCPS single crystals grown a) by the Bridgman method and b) from benzene solution.

$0.38 \text{ J.mole}^{-1}.\text{K}^{-1}$ ; they are both a little larger (20 %) than previously reported [18]. The entropy of the transition remains much smaller than  $R \cdot \ln(2)$  as this is the case for all displacive structural phase transitions.

### 3. $^{35}\text{Cl}$ -nuclear quadrupole resonance study.

3.1 EXPERIMENTAL. — Investigations were performed at the I.E.F. on a laboratory made pulsed NQR spectrometer [21]. The power stage of this spectrometer is now a commercial 1 kW peak wideband rf power amplifier (Kalmus 166LP), and an usual value for a  $90^\circ$  pulse width is 16 microseconds with a 14.5 mm inner diameter, 5 turn coil. As the free induction decay signal is mixed with the receiver recovery signal, the echo following a standard  $90^\circ$ - $\tau$ - $180^\circ$  pulse sequence was recorded. When the spectrum is broad, its shape is obtained from the plot, on the same diagram, of the echo signals recorded at different frequencies [22]: the horizontal position of each echo signal on the diagram is proportional to the spectrometer frequency and the line joining the points of maximum amplitude of each echo is an

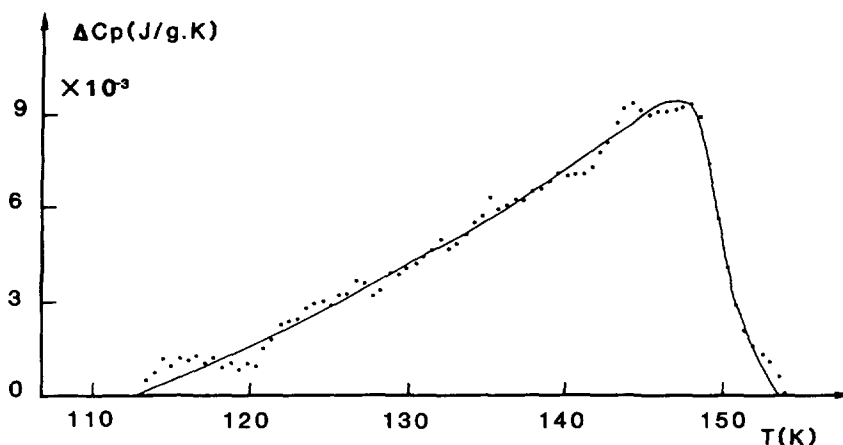


Fig. 3. — Specific heat excess observed in the Bridgman grown BCPS single crystal.

approximation of the shape of the spectrum, or spectral function  $f(\nu)$ . Most of the NQR work was done on a large polycrystalline sample of BCPS (Aldrich), recrystallized from the melt, about 14 mm in diameter and 3 g in weight. Some spectra were also recorded from the large and high-quality single crystal (diameter: 13.5 mm, weight: 2.6 g) which gave the  $\Delta C_p$  shown in figure 3.

**3.2 NQR SPECTRUM.** — Above 150 K, in the high temperature phase of BCPS, there is only one crystallographical site for the chlorine atoms resulting from the molecule and crystal symmetries; thus one half of the molecule suffices to generate the four molecules of the unit cell by symmetry operations. Accordingly, a single NQR resonance line is observed.

Below 150 K, the spectrum is broad and recorded as described in 3.1. Spectra recorded at different temperatures below 150 K are shown in figure 4; upon cooling, the spectrum broadens steadily and its shape changes from that of a bell to a more complicated one with three resolved maxima separated by two valleys. This structure persists down to liquid nitrogen temperature. The spectra recorded with the polycrystalline and single crystal samples at 77 K are presented in figure 5: they are quite similar which shows that the quality of the polycrystalline sample used for most of this work is quite satisfactory. At this temperature, the width of the spectrum is about 300 kHz. Such a spectral distribution is typical of incommensurate phases where no atom occupies the same position as any homologous atom in another crystal cell of the crystal. The spectra presented in figure 4 show a continuous evolution from 150 K to 90 K and no clear difference appears between the spectra at 110 K and 120 K which could suggest a transition occurring between these two temperatures, in particular at 115 K; moreover, no narrowing of the resonance peaks leading to a discrete spectrum, as expected for a commensurate phase, is observed.

In order to understand the discrepancy between the present results and those already published [16, 17], we have plotted in figure 6 the experimental maxima recorded by us (crosses) and Nakayama *et al.* [17] (closed circles). It seems that for temperatures above 115 K, these authors recorded the low frequency and central peaks of the spectrum, whereas for temperatures below 115 K they have recorded the low and high frequency peaks; the change in the peaks actually recorded, above and below 115 K, leads to the jump in the  $\nu(T)$  curve they report. This may result from a low signal-to-noise ratio or from some

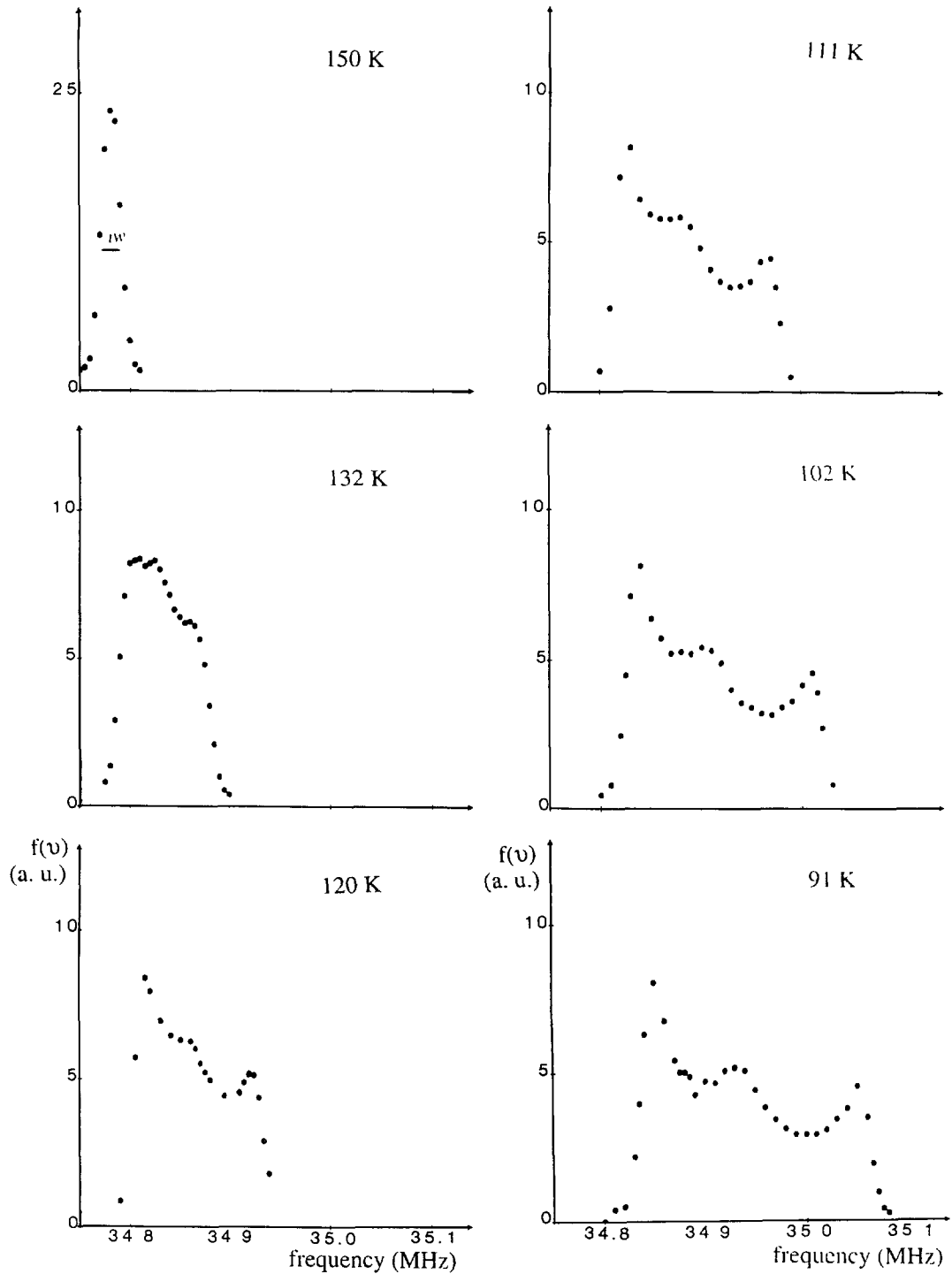


Fig. 4. — Temperature evolution of the  $^{35}\text{Cl}$  NQR recorded spectrum in the BCPS incommensurate phase. The horizontal bar  $i_w$  in the 150 K diagram shows the 18 kHz instrumental width of the spectrometer response. It is the same for all the spectra.

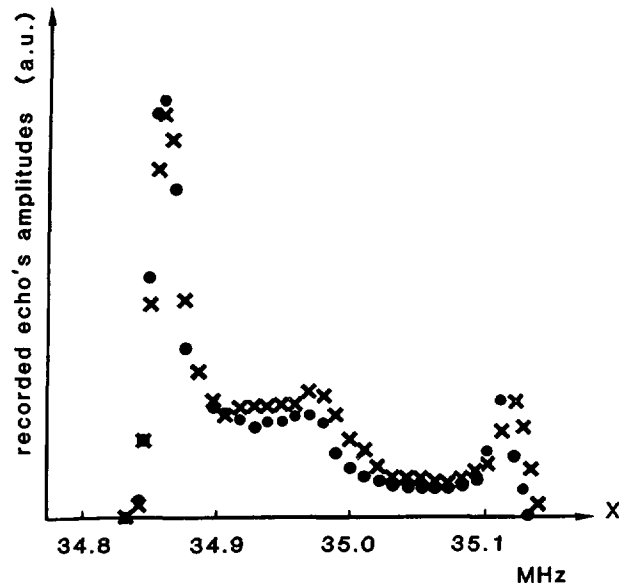


Fig. 5. — Comparison of the  $^{35}\text{Cl}$  NQR spectra in polycrystalline Bridgman grown ( $\times$ ) single crystal ( $\bullet$ ) samples of BCPS at 77 K. For technical reasons, phase detection was not used here and the receiver response is nonlinear. Though the single crystal sample gives a stronger signal, the two spectra are drawn with similar amplitude for the sake of convenient comparison.

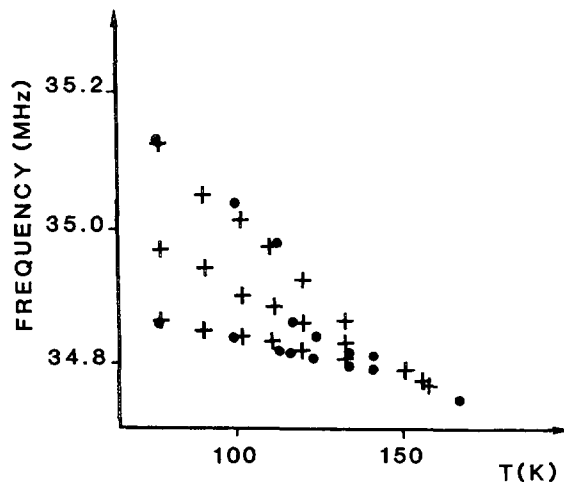


Fig. 6. — Temperature evolution of the different maxima in the BCPS spectra recorded in the present work ( $+$ ) compared to the results in reference [16] ( $\bullet$ ).

difference in the spectrum shape related to the sample preparation. Indeed, in the temperature range where the BCPS incommensurate phase is studied here, the frequencies of the maxima in the resonance curve vary smoothly with temperature and Nakayama's data can be reinterpreted consistently with the present investigation. The present reinterpretation of the data thus excludes any lock-in transition in this compound down to 77 K. In this study, the shape of the

recorded spectra confirms the incommensurate character of the BCPS phase appearing below 150 K. It is possible to discuss the results obtained further and propose a tentative model of the temperature dependence of the NQR spectrum. Such an analysis would be useful because the electric field gradient, and also the NQR frequencies, are related to the crystalline structure and constitute parameters to monitor the incommensurate distortion in the crystal.

In the case of a sinusoidal one-dimensional static modulation, the special shape of the resonance line with three maxima, indicates that the NQR frequency of a chlorine nucleus is function of its displacement  $u$  with linear and quadratic terms [23] :

$$\nu = \nu_0 + a_1 u + a_2 u^2 \quad (1)$$

where  $\nu_0$  is the resonance frequency in the high temperature phase. The line shape can be derived from this equation by introducing the modulation :

$$u = A \cos [\Phi(x)] \quad (2)$$

where  $\Phi(x) = \Phi_0 + q_s \cdot x$  and  $q_s$  is the modulation wave vector. Therefore, the resonance frequency becomes :

$$\nu = \nu_0 + \nu_1 \cos [\Phi(x)] + \nu_2 \cos^2 [\Phi(x)] , \quad (3)$$

with  $\nu_1 = a_1 A$  and  $\nu_2 = a_2 A^2$  ;  $A$  is expected to vary as  $(T_1 - T)^\beta$ , where  $\beta$  is the order parameter exponent. The spectral function, which is defined as  $f(\nu) = \text{constant}/(|d\nu/dx|)$  has two or three singularities according to the respective values of  $\nu_1$  and  $\nu_2$ , functions of the temperature [23] :

When  $|\nu_2| \leq |\nu_1/2|$ , there are two singularities :

$$\nu_a = \nu_0 + \nu_1 + \nu_2 \quad \text{and} \quad \nu_b = \nu_0 - \nu_1 + \nu_2 .$$

When  $|\nu_2| > |\nu_1/2|$ , a third singularity appears at :

$$\nu_c = \nu_0 - \nu_1^2 / (4 \nu_2) .$$

These singularities correspond to the extrema of the function (3). When three singularities are present, one gets used to decompose the calculation of the spectral function between each extremum and called respectively  $g(\varphi_+)$  and  $g(\varphi_-)$  [23] :

$$f(\nu) = \begin{cases} g(\varphi_+) & \text{for } \nu_b < \nu \leq \nu_a \\ g(\varphi_-) + g(\varphi_+) & \text{for } \nu_c \leq \nu \leq \nu_b \end{cases} \quad (4)$$

where

$$g(\varphi) = \frac{\text{constant}}{|\sin(\varphi) \cdot [\nu_1 + 2 \nu_2 \cdot \cos(\varphi)]|} ,$$

with

$$\cos(\varphi_\pm) = \frac{-\nu_1 \pm \sqrt{\nu_1^2 + 4(\nu - \nu_0)\nu_2}}{2\nu_2}$$

Parameters  $\nu_0$ ,  $\nu_1$ ,  $\nu_2$ ,  $T_1$  and  $\beta$  are determined by fitting the frequencies of the peaks in the experimental spectra. However, the response of the spectrometer, acting as a convolution of the physical spectrum by the spectrometer response function, brings some perturbations to the recorded spectrum : i) the top of asymmetrical peaks is slightly shifted towards the smallest slanted side of the peak ; ii) the intensity of the peaks with respect to the valleys is reduced ; iii) the noise is reduced. To obtain more accurate values of the frequencies of the peaks, a



deconvolution [24] of the spectra was performed before determining these frequencies. The Bayes algorithm was used taking an 18 kHz half-height wide Gaussian for the deconvolution function. This instrumental function was determined by taking the Fourier transform of the echo given by a sample which has a flat spectrum over a frequency range much greater than the instrumental function, as it is the case in BCPS between the spectrum maxima at the lowest temperature studied (77 K) ; its width was then confirmed from the plot of a sharp NQR line, the one of BCPS at room temperature, recorded as done for broad spectra. A least square fit using the frequencies of the maxima in the deconvoluted spectra leads to the following parameters values, the transition temperature being  $T_1 = 150$  K :

$$\begin{aligned}\beta &= 0.5 \\ \nu_1 &= 7.25 \cdot (T_1 - T)^\beta \text{ kHz} \\ \nu_2 &= 2.14 \cdot (T_1 - T)^{2\beta} \text{ kHz} .\end{aligned}$$

For the temperature dependence of  $\nu_0$  in (3), resulting from the averaging effect of the molecular angular oscillations and described by the Bayer's theory, we used the approximation [25] :

$$\nu_0 = aT + b$$

with  $a = -1.0 \text{ kHz.K}^{-1}$  and  $b = 34.946 \text{ MHz}$ .

This set of values is different from that found above  $T_1$ , where the temperature dependence of the frequency is quite linear :

$$a = -2.92 \text{ kHz.K}^{-1} \quad \text{and} \quad b = 35.22 \text{ MHz} .$$

The difference between the  $a$  values, above and below the transition, is an indication of a corresponding difference in the crystal dynamics. This is consistent with the change observed in the intensity of the Bragg (0, 2, 0) reflection and the slopes of the temperature dependence of the crystal parameter  $a$  as shown in figures 1 and 2 of reference [10].

Concerning the determination of  $\beta$ , the results of the fitting procedure appear to be very sensitive to the extension of the temperature range considered below  $T_1$ . Changes of more than 0.15 were found on the values of  $\beta$  according to the temperature range considered in this fitting procedure. A value of  $\beta$  close to 0.5 was found, using the whole temperature range, and is used henceforth. This value, which corresponds to the mean field approximation is consistent with the results of others methods of investigation : X-ray diffraction [9], neutron diffraction [10] and Raman scattering [11, 12].

With the determined parameters and the 18 kHz wide Gaussian instrumental response as convoluting function, spectra at 132 K, 120 K, 111 K, 102 K and 91 K have been calculated (Fig. 7). They are in relatively good agreement with experimental recorded spectra shown in figure 4.

A possible difference between theoretical and experimental curves may be the non validity of the sinusoidal model for the static modulation when going further below  $T_1$ , especially below 100 K. The evolution of the modulation towards a nonsinusoidal regime is suggested by the analysis of high order satellites in neutron diffraction experiments [10], and the discussion of X-ray diffraction measurements taking these satellites into account [13]. In this last case, better crystallographic description is obtained by introducing second harmonic in the modulation. Calculation of the resulting NQR spectra would now be quite complex because much more parameters have to be introduced (weight and phase of the second harmonic relative to the first one). The effect of this so-called soliton regime on NQR spectra has been qualitatively discussed [23]. The modulation is described as commensurate quasimacroscopic

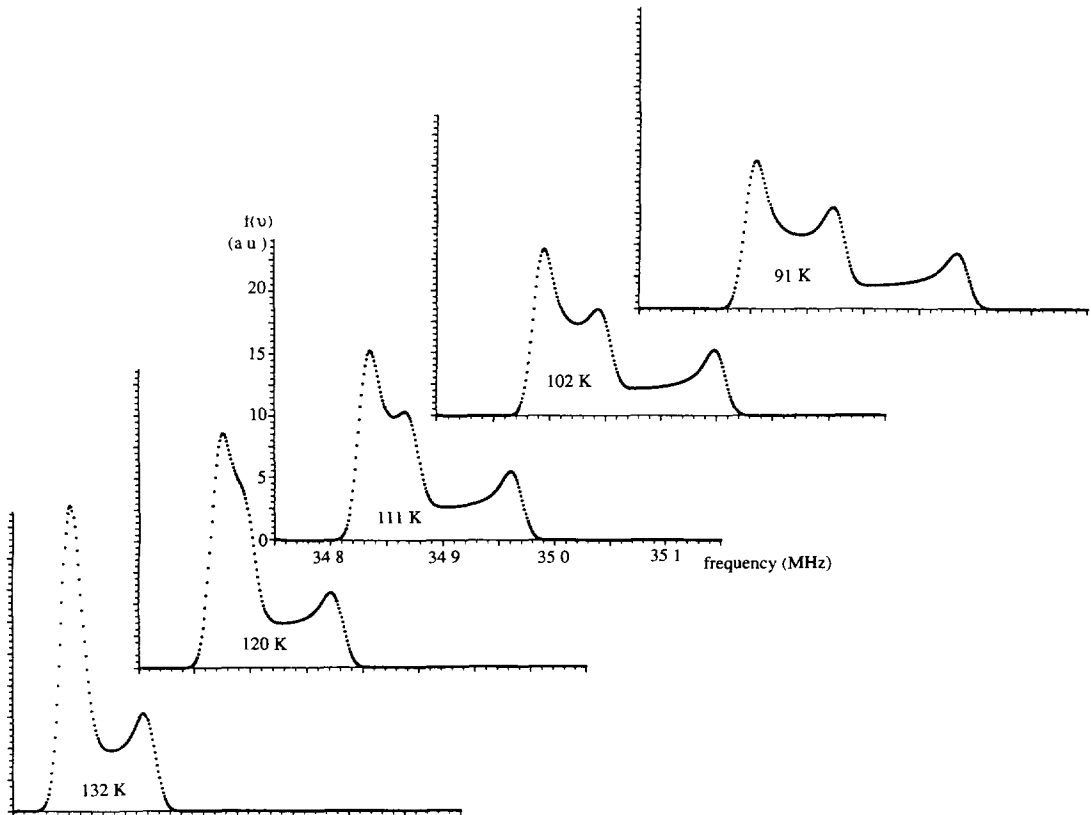


Fig. 7. — Spectral function curves calculated with the parameters derived from the experimental data (see text), including convolution with an 18 kHz half-height wide Gaussian simulating the instrumental response.

regions, corresponding to locked-in phase, separated by a periodic array of narrow « discommensuration » regions. The phase  $\Phi(x)$  of the order parameter is then a nonlinear function of the coordinate  $x$  which leads to the complexity of the expression (4). The evolution of the spectrum, as temperature is lowered, is expected to be characterized by a progressive increase of the intensity of the maxima associated to the nearly commensurate regions and a simultaneous decrease of the intensity between these maxima [23]. Both models, the crystallographic one, characterized by the presence of high order harmonics and a sinewave modulation, and the discommensuration model, predict the appearance of new lines in the spectrum. NQR measurements at lower temperature would show these differences. Other authors [17] have carried on the  $^{35}\text{Cl}$  NQR study in BCPS down to liquid helium temperature. They observed new lines that they interpreted as characteristic of a low temperature commensurate phase. As it is clear now that there is no lock-in transition in this compound at least down to 4.2 K, a new NQR investigation should be done to continuously follow the spectrum evolution in this whole temperature range. This would precise the deformation of the static incommensurate modulation when going towards a nonsinusoidal regime with the apparition of narrow solitons.

**3.3 SPIN-LATTICE RELAXATION.** — In the temperature dependence of the  $^{35}\text{Cl}$  spin-lattice relaxation time, Pusiol *et al.* [16] found an argument in favour of a lock-in phase transition at

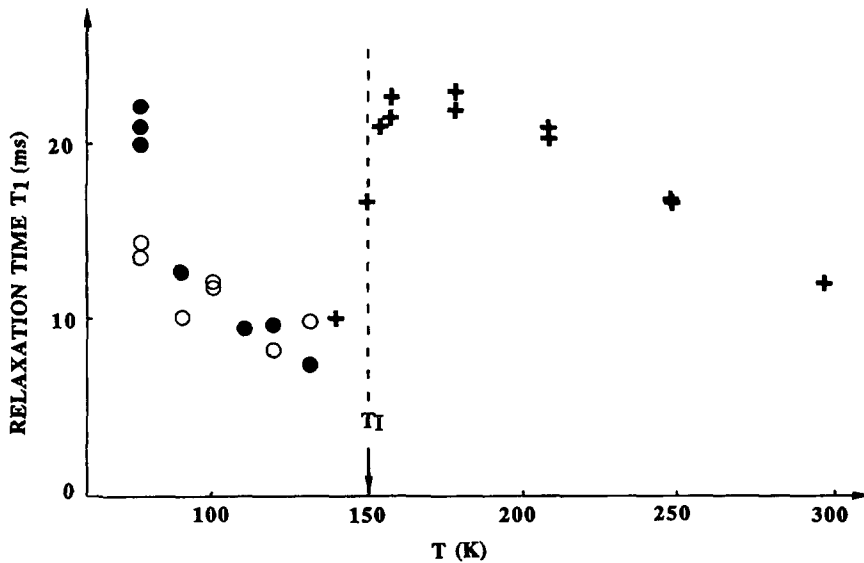


Fig. 8. — Temperature evolution of the spin-lattice relaxation times measured in the high temperature phase (+) on the center of the line, and, in the incommensurate phase, on the low-frequency (O) and high-frequency (●) peaks. In the incommensurate phase, the relaxation is not exponential and the value that can be proposed as characteristic time depends upon the method employed to calculate it; the values shown in the figure are half the time necessary for the signal to recover to  $(1 - 1/e^2)$  of its full value after a saturation pulse.

115 K : they observed a sharp increase in  $T_1$  below 110 K. This is usually observed when the phason spectrum vanishes at the transition from an incommensurate to a commensurate phase [23]. We made the same kind of measurements as Pusiol and the results are shown in figure 8. In the normal phase, above 150 K, where the spectrum consists of a single line, narrow as compared to the frequency spectrum covered by a saturation pulse,  $T_1$  values are quite similar to those of Pusiol. Below 150 K, in the incommensurate phase, the spectrum becomes wider and wider; below 140 K,  $T_1$  was measured at different frequencies in the spectrum and the values at both ends of the spectrum maxima were recorded. Between 140 K and 100 K,  $T_1$  does not much change through the spectrum and remains close to 8 ms; the fact that, between 140 K and 100 K, the relaxation time  $T_1$  does not vary appreciably through the spectrum is unexplained in the frame of the usual model of spin-lattice relaxation for quadrupolar nuclei in incommensurate phase [5, 23]. Below 100 K,  $T_1$  increases sharply and its value, at 77 K, varies from about 20 ms on the high frequency peak to 15 ms on the low frequency peak and 10 ms in the high frequency valley. These values are only indicative as the relaxation is not exactly exponential in this temperature region and the signal-to-noise ratio is poor in the valley. On the whole, the overall  $T_1$  behaviour is similar to that reported by Pusiol.

The important increase in  $T_1$ , below 100 K, reveals a modification of the spectral correlation function of the phason branch. As the absence of major change in the spectrum around this temperature precludes the consideration of a lock-in phase transition, another mechanism should be found to explain this modification. It may result from an evolution of the incommensurate structure from a sinusoidal to a soliton regime, as discussed above. Then, theory predicts the opening of gaps in the phason branch in the reciprocal space for the wave vector corresponding to the ordered locked commensurate phase [26]. As the commensurate

phase has not been observed, this wave vector is not known in BCPS. This makes difficult a quantitative estimation of the effect of the modification of the dynamics on the NMR relaxation processes in this compound. To a first approximation, one may consider that going from a sinusoidal incommensurate dynamics to a soliton-like type yields a decrease of the spectral density of the correlation function. Since this is the efficient process for the nuclear spin-lattice relaxation, it could explain the increase of  $T_1$  below 100 K in BCPS.

#### 4. Conclusion.

The present  $C_p$  investigation in BCPS, from 300 K down to 90 K, reveals only one transition, displacive and second order at  $T_1 = 150$  K. The chlorine-35 NQR spectra recorded from room temperature down to 77 K, confirm the existence of an incommensurate phase below 150 K and, also, the absence of a lock-in transition in this temperature range. The three peak structure observed in the incommensurate phase and below 140 K, shows the presence of a quadratic term in addition to the linear one in the expression of the resonance frequency in terms of the atomic, or molecular, displacements. These new results, from calorimetry and NQR investigations, are in agreement with those obtained from neutron and X-ray diffraction, the latter methods showing the incommensurate phase persists down to the lowest temperature studied. BCPS provides an exceptional opportunity for studying the dynamics of a simple incommensurate system with very weakly damped excitations when passing from a sinusoidal incommensurate phase to a soliton-like phase. Further NQR investigation, on the temperature dependence of the frequencies and relaxation times, down to liquid helium temperature, would bring forth fundamental informations on static and dynamical aspects of the incommensurate phase in BCPS.

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#### References

- [1] Incommensurate Phases in Dielectrics, in modern problems in condensed matter sciences, vol. 1 and 2, R. Blinc and A. P. Levanyuk Eds. (North Holland, 1986).
- [2] Janssen T. and Janner A., Incommensurability in crystals, *Adv. Phys.* **36** (1987) 519-624.
- [3] Currat R. and Janssen T., Excitations in incommensurate crystal phases, *Solid State Phys.* **41** (1988) 201-302.
- [4] Cummins H. Z., Experimental studies of structurally incommensurate crystal phases, *Phys. Rep.* **185** (1990) 211-411.
- [5] Blinc R., Magnetic resonance and relaxation in structurally incommensurate systems, *Phys. Rep.* **79** (1981) 331-398.
- [6] Borsa F. and Rigamonti A., Structural Phase Transition II, K. A. Muller and H. Thomas Eds. (Springer-Verlag, Berlin, 1991) pp. 83-183.
- [7] Cailleau H., in [1], vol. 2, chap. 12.
- [8] Currat R., Bernard L. and Delamoye P., in [1], vol. 2, chap. 15.
- [9] Kasano H., Koshiba T., Kasatani H. and Terauchi H., *J Phys. Soc Jpn* **59** (1990) 408-411.
- [10] Etrillard J., Even J., Sougoti M., Launois P., Longeville S. and Toudic B., *Solid. State Commun.* **87** (1993) 47-51.
- [11] Ishii K., Nakayama H., Sakato T. and Kano H., *Phys. Soc. Jpn* **61** (1992) 2317-2326.

- [12] Sougoti M. and Ecolivet C., private communication.
- [13] Zuniga F. J., Perez-Mato J. M. and Breczenski T., to be published.
- [14] Liu S. B. and Conradi S., *Phys. Rev. Lett.* **54** (1985) 1287-1290.
- [15] De Souza R. E., Engelsberg M. and Pusiol D. J., *Phys. Rev. Lett.* **66** (1991) 1505-1508.
- [16] Pusiol D. J., Wolfenson A. E. and Brunetti A. H., *Phys. Rev.* **B 40** (1989) 2523-2528.
- [17] Nakayama H., Eguchi T. and Kishita M., Proceedings of the XIth International Symposium on Nuclear Quadrupole Resonance Spectroscopy, *Z. Naturforsch.*, London **47a** (1991) 232-236.
- [18] Saito K., Kamio H., Kikuchi K. and Ikemoto I., *Thermochim. Acta* **163** (1990) 241-248.
- [19] Aston J. G., Szasz G. J. and Fink H. L., *J. Am. Chem. Soc.* **65** (1943) 1135-1143.
- [20] Ginnings D. C. and Furakawa G. T., *J. Am. Chem. Soc.* **75** (1953) 522-527.
- [21] Christidès A., Thèse de 3<sup>e</sup> cycle, Orsay, France (1986).
- [22] Brown R. J. C., Christidès A., Gourdjji M. and Guibé L., *J. Mol. Struct.* **192** (1989) 355-367.
- [23] Blinc R., Prelovsek P., Rutar V., Seliger J. and Zumer S., in [1], vol. **1**, chap. 4.
- [24] Bourges P., Thèse de l'Université de Rennes, France (1989).
- [25] Kusida T., Benedek G. B. and Bloembergen N., *Phys. Rev.* **B 104** (1956) 1364-1377.
- [26] Bruce A. D. and Cowley R. A., *J. Phys.* **C 11** (1978) 3609-3630.