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RAMAN STUDY OF GEOMETRICAL DISORDER AT SEMICONDUCTORS SURFACES

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<u>Résumé</u>.- Un désordre de nature géométrique est provoqué à la surface d'échantillons par un pollisage mécanique grossier. Son influence sur les propriétés vibroniques et électroniques de InSb, InAs et GaSb est discutée.

Abstract.- Geometrical disorder is induced at the surface of samples by rough mechanical polishing. Its influence on both vibrational and electronic properties are reported in InSb, InAs and GaSb.

1.- Introduction.- It is well known that optical properties of crystals (reflectivity, absorption,luminescence) depend strongly on surface preparation. It is prejudicial when studying bulk properties but provides, nevertheless, a mean of investigating damages at the surfaces. Among other optical techniques Raman spectroscopy is quite suitable to look at these damages since the penetration depth of the light can be small enough to study the disorder in the vicinity of the surface. Evans and Ushioda (1) were the first to realize this possibility. They took advantage of the broadening which characterizes the first order Raman lines of polished surfaces to analyse the disorder induced by the polishing process. Recently, Carles et al (2) found that the Raman spectra of pure InAs samples alterated by mechanical polishing display two new bands labelled D.A.T.A. and D.A.L.A. which were interpreted as first order disorder this so-called "geometrical" disorder.

We present here results relative to InSb, InAs and GaSb. The influence of the geometrical disorder on the electronic properties of these compounds is also discussed.

2.- Experiment.- The samples analyzed were single crystals which (100) surfaces were mechanically polished with 0.3  $\mu$  alumina powder and then rinced in methanol. To avoid scattering from air the spectra were always recorded with the samples inside the vacuum chamber of the cryostat. They were not analyzed.

3.- Results and discussion.- We display in figures 1 and 2 respectively, the direct traces for InSb and InAs over a wide range of temperature. Arrows indicate phonons frequencies from ref. 3 and 4. We present only the frequency range below the  $TO(\Gamma)$  and  $LO(\Gamma)$  which are still present in the spectra and well polarized (2).

In figure 1, besides the 2TA scattering characteristic of crystalline InSb two new bands (D.A.T.A. and D.A.L.A.) appear confirming the results already found in InAs (2). At room temperature the D.A.L.A. is somewhat obscured by the end of the 2TA scattering (2IA(K) at 121 cm<sup>-1</sup>). As temperature decreases the second order scattering gets strongly attenuated (see the 2TA(X) for example) and the D.A.L.A. band merges around 137 cm<sup>-1</sup> at 25 K. The apparent broadening of the low energy tail of the TO( $\Gamma$ ), as temperature decreases, exemplifies the presence of disorder activated optical processes.

In figures 1 and 2, the intensities of the D.A.T.A. and the D.A.L.A. become stronger, compared to that of the second order scattering, as one goes from 300 K to 25 K. This

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Figure 1.- Evolution with temperature of the Raman spectrum of InSb alterated mechanically.

Figure 2.- Same as fig.1, for InAs.

λ=5145Å

214 (%)

150

200

behavior shows that they correspond to first order processes. In figure 2, the intensity of the TO-TA band around 170 cm-1 follows the expected temperature dependance for a substractive combination.

The frequencies of the maxima of the D.A.T.A. and the D.A.L.A. are respectively half those of the 2TA(X) and 2LA(L) indicating that they correspond to transverse and longitudinal acoustical modes respectively. The D.A.T.A. and the D.A.L.A. appearing in both InSb and InAs near phonon frequencies characteristic of each compound rule out the eventuality of In-In modes. Also As-As or Sb-Sb modes would lie outside the frequency range explored (5).

At the lowest temperatures the shape of the D.A.T.A. mirrors that of the 2TA scattering : same kinds of singularities characteristic of L, X and K points, same relative intensities of the bands at these points. The first assertion already proves that phonons of the whole Brillouin zone contribute to the D.A.T.A. in accordance with the breakdown of the  $\aleph = 0$  selection rule. The second one, together with the fact that, in the III-V compounds, the 2TA scattering reflects the transverse accoustical one phonon density of states (3-4) with the scale doubled, implies that the D.A.T.A. too images this density of states.

From these facts, Carles et al. (2) proposed to fit, over the whole acoustical region, the disorder activated scattering in InAs by the expression :

$$I_1 = C_1(\omega) g'(\omega) [\bar{n}(\omega,T) + 1]$$

 $g'(\omega)$  is the one phonon density of states of InAs convoluted by a gaussian to account for the broadening of the structures ;  $\bar{n}$  is the Bose Einstein factor ;  $C_1(\omega)$  is the coupling constant of the light with the vibrational modes. The fit was excellent ta-