

ON THE PERTURBATION OF THE ATOMIC STRUCTURE OF LIQUID POTASSIUM BY THE ADDITION OF KCl

J .F. Jal, J. Dupuy, P. Chieux

To cite this version:

J .F. Jal, J. Dupuy, P. Chieux. ON THE PERTURBATION OF THE ATOMIC STRUCTURE OF LIQUID POTASSIUM BY THE ADDITION OF KCl. Journal de Physique Colloques, 1980, 41 (C8), pp.C8-257-C8-261. $10.1051/jphyscol:1980866$. jpa-00220375

HAL Id: jpa-00220375 <https://hal.science/jpa-00220375v1>

Submitted on 4 Feb 2008

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

ON THE PERTURBATION OF THE ATOMIC STRUCTURE OF LIQUID POTASSIUM BY THE ADDITION OF KC1

J.F. JAL, J. DUPUY and P. CHIEUX^{*}

Département de Physique des Matériaux, Université Claude Bernard, Lyon Villeurbanne, France.
* *Institut Laue-Langevin, Grenoble, France*.

Abstract: We present and analyse the effect of the addition of up to 25% of KC1 on the structure of molten potassium. At these metallic concentrations salient features of the salt structure are becoming apparent. The small angle scattering behavior and preliminary excess volume data point to strong deviation from regularity at very high metal concentration.

The solutions of the alkali metals in their molten salts have been long known [I1 for their considerable change in electrical transport proper-
ties versus addition of metal They offer an inter 10^{4} . ties versus addition of metal. They offer an interesting example of non metal to metal transition in disordered systems. A detailed study of the concen-
tration fluctuations near the liquid-liquid mis-
cibility gap in the system K in KBr has already
 $\overrightarrow{10}$ \overrightarrow{E} tration fluctuations near the liquid-liquid miscibility gap in the systen K in KBr has already been made by small angle neutron scattering [2], [3]. On the other hand, it is quite interesting to $\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$ $\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$ investigate the short range order behavior of this disordered system where ionicity sets in progressively.

In the present study we have obtained the structure factors of several solutions of potassium in KC1 in a concentration range where the systen is metallic [4] $10^{-4} \Omega^{-1}$ cm⁻¹ < σ < $10^{-3} \Omega^{-1}$ cm⁻¹ and where the concentration fluctuations are definitively very small [2] (see figure 1). The experiments have been performed on the D4 machine at the Institut Laue Langevin (Grenoble) [51.

I. The experimental results - The structuye factors

The structure factors of five solutions of potassium chloride in potassium from 75% to 95% in potassium content are displayed on figure 2. The smoothed structure factor of potassium given for comparison has not been extended to the whole **Q** range since it was found identical to the literature

Fig. 1: Phase diagramme and electrical conductivi ty of $K - KC1$

value [61. The data have been corrected for background, furnace, vanadium container and multiple scattering in the usual manner [7]. The incoherent scattering, assumed to he isotropic, has been obtained at large scattering angles fron the ratio of incoherent to total scattering cross section. This is however still a first approach of the data acquisition and treatment. As it is visible on the graph not all the runs have been performed to the same accuracy, the 90% concentration being especially poor. A close look at the low angles reveals also some difficulties left in the background correction, as noticeable on the 952 sample near 0.75 Å^{-1} . On the other hand, detection problems have sometimes contaminated the high angle values (see the 75% sample). We believe however that the

Fig. 2: The structure factors for K - KC1. Temperatures 867°,830°,800°,725°,780°C from 75 to 95%.

present state of data analysis is sufficient to assert the main characteristics of those structure patterns. Several features are of interest. First, a significant small angle scattering signal is observed at all those concentrations, although it is very small at 95%. The main peak of the potassium structure factor is considerably lowered upon the addition of small amounts of KC1 and has practically disappeared at a concentration of 85% in potassium. On the-other hand it is at about this concentration that we begin to detect a characteristic peak of the salt structure. It is a pity that the main peak in the partial structure factor S_{CC} expressing the correlation of charges in KC1 just happens to be at the main peak position of the structure factor of the pure metal preventing any analysis at this stage. However we see at the position $Q_N \approx 2.33 \text{ A}^{-1}$ the main peak of the partial structure factor S_{NN} , [8], expressing the correlations of density or mass in the salt. Finally, one notices also the almost structureless behavior of the 90% concentration where the disappearance of the metal peak matches the appearance of the salt peak.

11. The Small Angle Scattering

A high value of the thermodynamic limit S(o) accompanied by a significant small angle scattering is not uncommon in liquid metallic alloys [9] [lo]. It has been attributed to concentration fluctuations or to clustering effects. An accurate analysis of the **D4** low angle data is not appropriate considering the poor experimental resolution at those angles and the difficulty to untangle the small angle signal from the foot of the structure factor main peak. It is however of interest to evaluate approximatively the correlation length or the size of the corresponding fluctuations. A crude fit.

a value of about 3×5 , the same value being obtained in a Guinier analysis. Such a small fluctuation would not have been detected on conventional small angle machines such as DII or 017 at the I.L.L. It is of course difficult to give a definite physical picture of so short fluctuations or to assess clearly if they are related to the phase separation which appears at much lower concentration in metal. Accurate thermodynamic values and nore resolution in the low angle data would be quite useful. One may however make a comment on the significant feature which appear on the figure 2, i.e. the very low small angle signal at 953.

The anomalies of the thermodynamic limit S(o) in a binary system depend essentially on two terms, the value of the partial structure factor $S_{cc}^{\dagger}(\sigma)$ related to the derivative of the chemical potential versus concentration [11] and a contrast coefficient C which involves the difference of the neutron scattering length per unit volume between the two components [I21

$$
S(o) = A + C Scc(o)
$$
 (1)

the A-term depends on the isothermal compressibility of the system. The contrast is

$$
C = \frac{1}{\Sigma_1 c_1 b_1^2} \left| \frac{\overline{v}_1 b_2 - \overline{v}_2 b_1}{c_1 \overline{v}_1 + c_2 \overline{v}_2} \right|^2
$$
 (2)

where c_i , b_i and \overline{v}_i are the concentration in mole fraction, the scattering length and partial molar volume of the two constituants, K and KC1 in our case **(X).We** have displayed on figure 3 the estimated volume of the two constituants, K and KCl in our

case (x). We have displayed on figure 3 the estimated

variation of S_{CC}(o) and of the contrast C versus $\frac{Fig. 3a}{\text{volume}}$. The contrast effect on S(o). 3b) Excess

conce concentration for the system K in KC1. The $S_{cc}(0)$ values have been obtained from neutron small angle $\frac{1}{x}$ The choice of K and KCl as the two basic
constituants for writing the equations (1) and (2) constituants for writing the equations (I) and (2)
does not imply that these are really the two terms of the fluctuation process. Complementary experi- KRr and developed in reduced coordinates around ments are required for such a determination. the liquid-liquid critical point of K in KC1. On

of an Ornstein Zernike law on the 85% sample gives the basis of vapour pressure measurements [13], they decrease smoothly on the potassium rich side. The contrast coefficient has been computed first on a regular solution basis (curve R) and shows a maximum near 95Z in K. However, preliminary excess volume neasurements as obtained from neutron transmission data display a positive excess volume in both K in KCl and K in KBr systems at concentrations rich in metal (see figure 3b). On this basis, the maximum in the contrast C would be around 85X (see curve ΔV_{FX} in figure 3). The very low small angle scattering at 95% supports therefore the idea of strong deviations from regularity in the very concentrated metallic system.

volume and deviation from regularity.

111. The Fourier Transforms - Comparison with Au-Cs

The effect of addition of small quantities of KC1 to potassium is well illustrated again on figure 4 by the Fourier transforms of the total structure factors. The concentrations80% and 85% having intermediate values between those displayed have been omitted. Difficulties in Fourier transforming such as on the 909; sample which was of poor statistics have still to be overcomed but the main trends are obvious. On the 90% sample and probably already on *95X* appears the characteristic distance d_N of the KCl structure which is quite different from the first and second nearest neighbor d_1 and d_2 distances in potassium. We see on the total distribution function $g_N(r)$ of KCl^{nat} at 818°C [8] that the peak position d_M is typical of the first nearest neighbors i.e. the K⁺- Cl⁻ distance. One notices however that the distance of closest

Fig. 4: The Fourier transforms of K - KC1.

approach which is nearly the same for all concentrations studied is slightly larger than for the pure salt (2.25 A). The first neighbor metal peak considerably diminishes with the addition of KC1. At 95X the diminution is approximatively described however by simply taking into account the temperature and yolume expansion effectilk].

The macroscopic properties of the M in MX and the M in MAu systems [15] present strong analogies. This is also true if one compares their structure factors [161. In the Cs-Au case, the $characteristic Cs⁺-Au⁻ distance of the salt was$ observed at a concentration of 20% Au in Cs and if the analogy with K in KC1 is correct, it should be detectable at even much higher concentrations of Cs .

These first results are stimulating since they point to a significant salt-like structure even at high metal concentration. It is essential however to obtain accurate values of the partial structure factors and especially the charge charge partial S_{cc} at one concentration at least, in order to throw some light on the onset of ionicity in these metallic systems.

References:

- [1] see BREDIG, M.A., in "molten salt chemistry" Ed. BLANDER, M., Wiley, N.Y. (1964)
- [2] J.F.JAL,P.CHIEUX,J.DUPUY, J.Physique 41 (1980)
- *C33* **P.CFTIEUX,P.DP~dAV,J.DUPV',J.?.JAL,Colloque WEYL** 17(137?)to be published in S,Phys.Chen. (1980) [4] We thank H.BRONSTEIN for kindly supplying his
- original data
- [5] J.F.JAL,P.CHIEUX, J.DUPUY, to be published
- [6] M.JOHNSON, to be published
- [7] S.EISENBERG, thèse de spécialité, LYON 1977
- [8] **J.DEP.I?IETJ,J.DUPU??,J.Physique** 36,191 (1975)
- [9] W.KNOLL,S.STEEB,Z. Naturforsch.33a,472 (1978)
- [10] M.J.HUIGBEN, W. VAN DER LUGT, W.A.M.REIMERT, J.TH.M. DE HOSSON, C. VAN DIGK, Physica, 97B, 338-364 (1 979)
- [11] A.B.BHATIA, D.E.THO?NTON, Phys. Pev. B2, 3004 (1970)
- [12] P.DAMAY, P. CHIEUX, Z. Naturforsch. 34a, 804 (1979)
- [13] P.KOZOULIA, thèse de spécialité, Univ. de Provence **(1** 980)
- [14] J.BLETRY, J.F.JAL, J.DUPUY, to be published
- **1152** P.HENSEL,Adv.in Physics *23<4),* **555** (1979)
- [I61 S.STEEB,I?.MARTIN, ILL annuel report 1979,p.240
- [17] J. BLETRY, 2. Naturforsch. 32a, 445 (1977).

are 0.16\AA and 0.08\AA in the amorphous and crystalline alloys, respectively. The cause of this discrepancy is still under investigation (one possibility is a double-peak feature in the pair distribution function of the Mo-Ni distances). However, several general features can bc seen from Table **11.** Flrst. the distances are about O.IOA shorter in the amorphous than In the crystalline phase. Second, coordination number of the first shell at the Ni atoms in $a-MoNi$ is high \sim 12, the value found for densely close packed structure. In comparison, the crystalline phases (Tables I and **11)** show low coordination number \sim 7. Third, the root mean square deviations (Debye-Waller factors) of \sim 0.10 Å in the a-MoNi are large (in comparison to the corresponding values of $\sim 0.05\text{\AA}$ in the crystalline matenals

> Table 2. The distance r. Debye-Waller factor σ . and the coordination number N of the first shell of $a-Mo₅₀Ni₅₀$ and $c-Mo₅₀Ni₅₀$.

Discussion

The low coordination number of $\overline{N} \sim 7$ at the first shell for both Mo and Ni edges and the appearance of the second shell at \sim 3.0A for the Mo edge in c-MoNi suggest that the bcc-type local short range order prevails in the c-MoNi alloy. In contrast, the absence of the second shell at 3 OA together with the large coordination number for the Ni edge in a-MoNi suggest that the structure of a-MoNi alloy differs from that of the crystalline alloys and may be descnbed by a more or less dense random packing of binary spheres. This is in harmony with the results obtained by X -ray scattering measurements on many glassy metal-metal alloys, such as $Zr_{70}(Pd, Ni, Co, Fe)_{10}^{(6)}$ and $Zr-Cu$ and $Nb-Ni^{(7)}$.

The Dcbye-Waller factor decreases with lowering temperatures but remains high of \sim 0.10Å at 77 K for the amorphous alloys. This

indicates a substantial structural disorder in a-MoNi which is comparable to that reported for a- $Pd_{80}Ge_{20}^{(8)}$ ($\sigma \sim 0.1$). It is premature at this point to discuss whether the unlike pair distances $r_{M_0-N_1}$ show any dccreasc in the amorphous state as commonly observed for many glassy alloys because factors such as large anharmonicity in σ or asymmetry in the pair distribution function can cause error in the distance determination. For the amorphous systems with large structural disorders a gencrallzed EXAFS formulation must be used.'9) The results will be published elsewhere.

The EXAFS data of Zr-Cu glassy alloys are in many features similar to those of the MoNi. Preliminary results indicate that the relative ratio of the coordination number of $Zr-Zr$ pairs to that of $Zr-Cu$ pairs decreases with increasing content of Cu.⁽¹⁰⁾

We are grateful for the experimental opportunity and personnel help at Stanford Synchrotron Radiation Laboratory.

References

- [I] H. S. Chen. Rep. Prog. Phys 43. 353 (1980).
- [2] B. K. Teo, P. A. Lec. A. L. Simons, P. Eisenberger, and B. M. Kincaid, J Am Chcm. Soc. 99, 3854 (1977)
- [3] P A. Lec, B. K. Teo, and A. L. Simons, J. Am. Chem. Soc. 99 3856 (1977)
- [41 B. K. Teo and P. A. Lee. J. Am. Chem. Sac. **101,** 2815 (1979).
- 151 P. Eisenberger and B. M. Kincaid, Science. **200,** 1441 (1978).
- I61 Y. Waseda and H. S. Chen. Rapidly Quenched Metals **111,** Val. 2 (London: Metals Society) pp. 415 (1978).
- [7] H. S. Chen and Y. Waseda, Phys. Stat. Sol~di (a) **51,** 593 (1979).
- [8] T. M. Hayes, J. W. Allen, J Tauc, B. C. Giessen and J. J. Hauscr. Phys. Rev Lett 40, 1282 (1978).
- 191 (a) G. Brown and P. Eisenberger. Solid State Comm.. 29, 481 (1979); (b) T. M. Hayes. J. B Boycc, and J. L. Beeby, J Phys., **C11,** 2931 (1978).
- [10] H. S. Chen, B. K. Teo and P. Eisenberger, unpublished.