

Tm-169 MÖSSBAUER STUDY OF THULIUM HYDRIDES (II): TmH3

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Introduction

The aptitude of the 8.4 keV resonance of ¹⁶⁹Tm for the study of thulium hydrides has been demonstrated in a preceeding contribution to this conference [1]. The reason for this lies, beside the low gamma energy, in the resolution power of the ¹⁶⁹Tm resonance for quadrupole interactions. ¹⁶⁹Tm, a 100% isotope, is in its trivalent state a non-Kramers ion with a ³H₆ configuration of the 4f¹² electron shell, which responds very sensitively to crystalline electric field (CEF) interactions. When the Tm³⁺ site symmetry is lower than cubic, rather large quadrupole splittings can be observed with the 8.4 keV $3/2 \rightarrow 1/2$ transition of 169 Tm. They may reach values up to 150 mm/s at liquid helium temperatures, where only the lowest CEF levels of the J=6 4f-configuration are populated. At higher temperatures, eq4f, the 4f contribution to the observed electric field gradient eq, is temperature averaged and the lattice contribution, eq_{latt}, becomes dominant. This behaviour has been observed and analysed in the early days of Mössbauer spectroscopy [2].

The information obtained from the present study of TmH_3 will be used to discuss the controversial results of two recent studies of Er^{3+} in YH₃ and ErH_3 , respectively [3,4].

Sample Preparation and Experiments

The TmH₃ sample was prepared by exposing

about 100 mg of Tm metal at 870 K to a H_2 pressure of 1400 Torr for one hour. During the the cooling of the sample, the H_2 pressure was raised to about 3000 Torr. The X-ray diffraction analysis showed only the hexagonal rare earth (R.E.) trihydride phase [5]. The gravimetric determination of the H content of the sample yielded a formula TmH_{3.1(2)}. The Mössbauer measurements were performed similar as described in [1].



TmH₃ at various temperatures

Results and Discussion

Fig. 1 shows two spectra of TmH₃ at 300 K and 4.2 K. Both spectra reveal a typical quadrupole doublet, and, as in the 169 Tm study of the Tm dihydrides [1], no indication of any magnetic hyperfine interactions. The temperature dependence of the quadrupole splitting, $\Delta E_Q = e^2 qQ/2$, is given in Tab. 1 for a series of measurements between 4.2 K and 350 K. Measurements at T > 300 K were followed by studies at lower temperatures to ensure that no hydrogen loss of the sample had occured. A plot of the data of Tab. 1 shows that ΔE_0 approaches a value of 14(1) mm/s at high temperatures.

From the quadrupole splittings observed in TmH₂, information on the CEF groundstate of the Tm³⁺ ion can be extracted by arguments similar to those applied recently to a Mössbauer study of Er³⁺ in ErH₃ and YH₃ [3] and to a susceptibility and magnetization study of ErH₃ [4], respectively. It was assumed that in the isostructural R.E. trihydrides the CEF interaction of the R.E. ions can be described by hexagonal CEF parameters, although the H lattice actually lowers the point symmetry at the R.E. ions to C_2 . This assumption is supported by a 155 Gd Mössbauer study [6], where an axially symmetric eq_{latt} was observed, indicating that the hexagonal lattice of the R.E. ions is mainly responsible for equatt.

Using the CEF-level diagrams of Segal and Wallace [7] for hexagonal fields (which, however, use only fourth and sixth order CEF operators), Flood [4] derived from the magnetic data of ErH₂ a negative value of x, the ratio of fourth order to sixth order CEF parameters. In contrast to this, Suits et al. [3] found from their Mössbauer data of Er^{3+} in YH₃ a positive value of x. Transforming the Er³⁺ CEF parameters to the Tm^{3+} case, the negative sign (and the magnitude) of x from Ref. 4 would lead to two singlets with predominantly $m_{T} = \pm 6$ character for the lowest CEF states in TmH2, while the positive sign of x from Ref. 3

results in a Tm³⁺ singlet groundstate with $m_{J} = \pm 3$ character [7]. Using a value of $eq_{latt} = 5 \cdot 10^{17} \text{ V/cm}^2$ from Ref. 6 and the above CEF states (including the shielding parameters of Ref. 2), a rough estimate yields $\Delta E_Q \cong$ 150 mm/s at 4.2 K for the former and $\Delta \tilde{E}_{0} \cong 50$ mm/s for the latter case.

Together with our experimental results, this clearly favours the CEF parameters obtained from the ¹⁶⁶Er Mössbauer study [3]. In principle, the temperature dependence of ΔE_0 can be used to further elucidate the CEF levels of Tm^{3+} in TmH_3 . However, since ΔE_0 decreases quite fast in the low temperature region, additional data are needed for a more detailed analysis.

Finally, it should be noted that the high temperature value of eq (ΔE_Q (T \rightarrow 400 K) = 14 mm/s corresponds to eq = 6 \cdot 10¹⁷ V/cm², which most likely can be regarded as an approach to eq_{latt}) is about 20% larger than the value obtained for GdH₃ [6]. Such a small difference can easily be attributed to the smaller volume of the unit cell of TmH2. This work will be continued.

Table 1: Temperature dependence of the ¹⁶⁹Tm quadrupole splitting in TmH₃

Temperature	⊅ E _O
<u>(K)</u>	<u>(mm/s)</u>
4.2	65.5(2)
55	56.2(2)
77	48.7(2)
260	18.1(6)
300	15.9(4)
330	15.4(6)
350	14.6(6)

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