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Liquid crystals from poly(4,4′-methoxybiphenyl methacrylate)

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1. Introduction.

Side-chain liquid crystalline polymers have generally been synthesized with a flexible spacer between the pendant mesogenic groups and the polymer backbone [1]. The flexible spacer is currently thought to mechanically decouple the pendant groups from the polymer backbone, allowing for the former to organize in a mesomorphic way and for the latter to adopt a random conformation [2]. A side-chain polymer without flexible spacers has been reported in the literature, which however showed a smectic phase [3]. In this note, another example of such a polymer will be discussed.

The present work is part of a systematic study of side-chain liquid crystalline poly(methacrylates) having 4′-methoxybiphenyl mesogenic pendant groups and oligo(ethylene oxide) spacers (abbreviated as PM-n in the following):

\[
\text{CO}_2\text{O} - \left(\text{CH}_2\text{CH}_2\text{O}\right)_n - \Phi - \Phi - \text{OCH}_3
\]

The thermal properties [4] and the structure [5,6] of the polymers with n = 2, 3 and 4 have been described recently. Their mesomorphic structure was found to be smectic E. The pendant groups are arranged in single layers according to a two-dimensional pseudo-centred rectangular lattice, and the polymer backbones run parallel to the smectic layers along the [110] direction. Using stereochemical and structural arguments, the polymer chains were described to adopt a ribbon-like geometry with the pendant groups hanging
on the same side of the backbone as shown in figure 1.

The liquid crystalline behaviour of the polymers with \( n = 0 \) and 1 is still under study. As shown in reference [4], thermodynamic and structural equilibrium for these polymers is, indeed, difficult to reach owing, presumably, to their high viscosity. With low-molecular weight PM-O polymers, whose viscosity was not exceedingly high, several interesting structural results have, however, been possible to obtain and will be presented below.

2. Results and discussion.

The PM-O polymer used in the present work (referred as RP38) was free-radically polymerized [7]. Its weight-average molecular weight was \( 5.3 \times 10^3 \) and its index of polydispersity was 1.6 [8]. Its thermal behaviour was studied by thermogravimetry, differential scanning calorimetry, and optical microscopy [4] : 

\[
\begin{align*}
S_{OB} \xrightarrow{124^\circ C} S_A \xleftarrow{185^\circ C} & S_P \xrightarrow{235^\circ C} I 
\end{align*}
\]

Transition temperatures reported are only approximate due to the equilibrium difficulties stated above.

2.1 The smectic P phase (\( S_P \)).— The \( S_P \) phase is found to occur at high temperatures, between 185°C and 235°C. Its structure was described previously as “palisades of ribbons stacked in one dimension” [9]. The ribbon-like polymer molecules (Fig.1) lie on their flat sides in layers (Fig.2a) and the polymer backbones are oriented alternately to the right and left. The layer spacing is equal to the thickness of the ribbons (8.4 Å), and not to the length of the fully extended monomeric unit (about 16 Å) as is usually the case for smectic phases with the mesogenic groups normal to the layers (Fig.2b).
The texture observed (Fig.3) under the polarizing microscope for this phase is of the classical Schlieren type with two and four-brush point singularities. This is characteristic of a nematic phase, and it is now useful to understand why this texture is observed with a layered structure. It is worth recalling the Schlieren texture is also observed with smectic C phases; but here only four-brush point singularities can be seen. In smectic C mesophases, the \( \hat{\varepsilon} \) director (which is the direction of the projection of the tilted molecules on the smectic planes) plays a rôle similar to that of the classical nematic \( \hat{n} \) director. The absence of two-brush point singularities is due merely to reasons of symmetry: while in nematics rod-like molecules adopt a random up to down orientation, making \( \hat{n} \) equivalent to \( -\hat{n} \), in smectic C phases the molecules are all pointing in the same direction, resulting in a non-equivalence of \( \hat{\varepsilon} \) and \( -\hat{\varepsilon} \) [10,11]. Because the pendant groups lie parallel to the layers, the smectic P phase described in this paper can be considered as a "special smectic C phase" with a 90-degree tilt angle. The \( \hat{\varepsilon} \) director of this "special smectic C" is, of course, equivalent to the \( \hat{\nu} \) director defined in figure 1. As the pendant groups point equally to the right and left (Fig.2a), the \( \hat{\varepsilon} \) director (here \( \hat{\nu} \)) becomes equivalent to a true nematic director, \( \hat{n} \) (\( \hat{\varepsilon} \) becomes equivalent to \( -\hat{\varepsilon} \)).

2.2 The smectic A phase (S\textsubscript{A}).— The S\textsubscript{A} phase, characterized by typical fan textures under the microscope [4], was found to occur in the temperature range between 124 and 185°C. To confirm the smectic A nature of the material, X-ray patterns of magnetically (1.2 T) oriented samples were registered (Fig.4). These consist of three sharp, equidistant, and very slightly arced Bragg reflections in the small-angle region, indicative of a lamellar structure with the layers oriented perpendicular to the magnetic field. The layer thickness of 15.8 Å corresponds to the length of one fully extended monomeric unit. The X-ray patterns also exhibit diffuse crescent-shaped scattering in the wide-angle region at about 4.5 Å. This scattering was oriented parallel to the direction of the magnetic field indicating that the pendant groups are consistently parallel to the field and perpendicular to the smectic layers.
high-temperature $S_p$ phase, the ribbons in $S_A$ are expected to be parallel to one another within the layers, presumably with little positional correlation between adjacent ribbons. However, no Schlieren textures can be detected with $S_A$ as $\hat{p}$ is not an optical director. X-ray diagrams taken with fibers drawn from the melt and annealed in the $S_A$ temperature range confirmed the structure just described and showed that the smectic layers are oriented parallel to the draw direction.

It is worth noting in the small-angle region of the X-ray patterns (Fig.4) the presence of four weak and very diffuse spots paired to form two diffuse streaks parallel to the smectic layers. The Bragg spacing of the streaks roughly corresponds to twice the smectic-layer thickness, indicating a tendency towards a modulated $S_A$ structure [12]. In each smectic layer there is a tendency for "ferroelectric" clustering of the ribbons, and between smectic layers, there is a tendency for "antiferroelectric" stacking of the cluster. The wavelength of the in-plane modulation (roughly 30 Å), related to the correlation between clusters, is of the order of three to four ribbon thicknesses.

2.3 The smectic OB phase ($S_{OB}$).— The $S_{OB}$ phase was found to occur at temperatures below 124°C. Its optical texture was similar to that of $S_A$ [4]. The X-ray patterns registered with magnetically oriented samples shown in figure 5 consisted of three to four sharp, equidistant, and rather slightly arced (although more so than in $S_A$) Bragg reflections in the small-angle region, indicative of a lamellar structure with the layers oriented perpendicularly to the magnetic field. The layer spacing measured equal to that of $S_A$, is consistent with a single-layered arrangement of upright ribbons as shown in figure 2b. In the wide-angle region and in a direction perpendicular to that of the magnetic field, the X-ray patterns showed three rather sharp and rather slightly arced Bragg reflections at 4.89, 4.38, and 3.92 Å coming from the lateral ordered packing of the mesogenic pendant groups. Considering the packing of the pendant groups alone, these reflections can be readily indexed as the 010, 100, and 110 reflections from a two-dimensional oblique lattice with the cell parameters $a = 4.61$ Å, $b = 5.15$ Å, and $\gamma \cong 108^\circ$ (Fig.6). The molecular area per pendant group, $S = ab \times \sin \gamma = 22.6$ Å$^2$, is in agreement with the usual values obtained with ordered rod-like smectics.

The first neighbours around each pendant group are in the [100], [010], and [110] directions and located at distances of 4.61, 5.15, 5.76 Å respectively. Knowing from CPK molecular models that the distance between every other pendant group along the polymer chain is about 5 Å, it is clear that the polymer backbones in the layers must run along the [010] direction, as illustrated in figure 6. The lateral distance between adjacent polymer chains is 8.8 Å, in agreement with
the known thickness of the ribbons. It is interesting to note that the packing of the polymer chains in the oblique lattice here is very close to that of the smectic E phase seen for the higher homologs of PM-O where the polymer chains run along the [110] direction of the pseudo-centre rectangular lattice [5,6]. The oblique lattice can indeed be considered as a slightly deformed pseudo-centred rectangular lattice with the cell parameters $A = |a| = 4.61 \text{ Å}$, $B = |a + 2b| = 9.91 \text{ Å}$, and $\Gamma = 82^\circ$. The slight difference between the two lattices is obviously due to the difference in the constraints caused by the lack of flexible spacers in PM-O.

3. Experimental.
The custom-made X-ray camera used was equipped with two bent gold-plated glass mirrors crossed to give a point-“focused” direct beam. It used Ni-filtered copper radiation from an Elliott GX20 rotating anode X-ray generator. With a sample to film distance of 100 mm, it was aligned to allow registration of Bragg reflections in the spacing range from 3 to 100 Å with a resolution of about 1000 Å. Samples were oriented magnetically in the X-ray beam using a custom-made magnetic sample holder. This was composed of a 1.2T permanent magnet with X-ray beam passages allowing for the study of the orientation parallel and perpendicular to the magnetic field. The sample holder was electronically thermostatted between room temperature and $250^\circ \text{C}$ within $1^\circ \text{C}$; the variation of the magnetic field was reversible and less than 3 % over this temperature range.

References