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An energy dependent phase shift analysis of proton-proton scattering between 700 and 1 300 MeV

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Résumé. — Une analyse en déphasages avec dépendance en énergie est présentée pour la réaction élastique pp \Rightarrow pp entre 700 et 1 300 MeV. La plupart des déphasages se connectent parfaitement avec ceux trouvés dans l'analyse de Saclay-Genève en dessous de 800 MeV. Grâce aux nouveaux résultats de Saturne II, la base de données est maintenant suffisante pour faire une telle analyse. Par contre nos résultats sont en désaccord avec ceux trouvés par Arndt *et al.*, et ce pour la plupart des déphasages. Le comportement des diagrammes d'Argand suggère un comportement résonnant pour l'onde partielle ${}^{1}G_{4}$ mais pas pour l'onde ${}^{3}F_{3}$.

Abstract. — An energy dependent pp \Rightarrow pp phase-shift analysis for elastic proton-proton scattering is reported in the energy range between 700 to 1 300 MeV. Almost all phase shifts smoothly connect with the previous results of the Saclay-Geneva phase shift analysis below 800 MeV. Due to recent Saturne II results the data base is now sufficiently abundant to perform such an analysis. Present results strongly disagree with the predictions of Arndt *et al.*, for most phase shifts. Argand diagrams suggest a non-resonant behaviour of the ³F₃ partial wave and a resonance-like behaviour for the ¹G₄ partial wave.

1. Introduction.

We present an energy dependent phase shift analysis (PSA) of elastic pp scattering in the energy range from 700 to 1 300 MeV. It makes use of a reasonable number of different spin observables measured with a good precision in a large angular domain. The PSA was made possible by the recent results obtained at Saturne II which have considerably enlarged the existing data base. This PSA is an extension and confirmation of the Saclay-Geneva PSA [1] in the energy range from 10 to 800 MeV. With the exception of the narrow energy interval between 970 and 1 040 MeV, the data base above 800 MeV up to now has been too sparse to perform even fixed energy analysis.

The Arndt *et al.* [3, 4] energy dependent PSA from 0 to 1.1 GeV differs considerably from the present PSA results for almost all phase shifts. Note that it makes use of an insufficient data base above 800 MeV, except around 1 GeV.

There exist also several other fixed energy PSA

above 800 MeV. All of them have quite insufficient data base and their results are very dispersed.

The present PSA uses the same amplitude formalism, one pion exchange (OPE) contributions, electromagnetic contributions and a phase shift energy dependence as described in reference [1]. These subjects are briefly reviewed in section 2, while section 3 is devoted to the data base. Results, their discussion and comparison with other representative analyses [3-8] and with the Paris potential predictions [9, 10] are presented in section 4. Conclusions are summarized in section 5.

Throughout this article we will use a four-index notation and the formalism presented in reference [11]. For data published before April 1981 we refer to reference [2] (reference to data already included in Ref. [2] are quoted as in Ref. [2] e.g. 78/B-200).

2. Formalism.

The observables are expressed in terms of five complex scattering amplitudes, the invariant scattering amplitudes a, b, c, d and e, as discussed in

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amplitudes is expressed by Legendre polynomial series and partial wave amplitudes α_{LJ} (Ref. [12] and formulae (2), (3) in Ref. [1]). Partial wave amplitudes, parametrized by nuclear bar phase shifts δ_{LI} are real, while at higher energies the δ_{LI} can be complex with Im $\delta \ge 0$. As in reference [1] we use the « five parameter representation » for the partial wave amplitudes assuming all ε_I to be real. This is justified by the fact that the fit to the data is reasonably good and that the number of imaginary phase shifts increases little compared to the lower energy interval. Indeed, only $\text{Im} {}^{1}\text{S}_{0}$, $\text{Im} {}^{3}\text{P}_{1}$ and $Im^{3}H_{5}$ were added at high energies to the imaginary parts of the ³P₂, ¹D₂, ³F₂, ³F₃ and ¹G₄ phase shifts used below 800 MeV in reference [1]. Other phase shift analyses in this energy range also use five parameters.

The expansion of the amplitudes was truncated at the total angular momentum J_{max} and orbital angular momentum L_{max} . The higher angular momentum states were replaced, as usual, by the Born approximation of the OPE contribution (Refs. [13, 14] and formulae (4) in Ref. [1]).

The Coulomb amplitudes and Coulomb-nuclear interferences were introduced as in references [12, 14]. For angular momenta L higher than L_{max} , we have added the magnetic moment correction [15] to the amplitude e.

The energy dependence of the phase shifts were fitted by a polynomial expansion of the form

$$\sum_{n=0}^{3} \frac{a_{LJn}}{n!} (T - T_0)^n , \qquad (1)$$

where T_0 is the central point of the interval (1 000 MeV for the present PSA) and a_{IJn} are variable parameters. In all cases it turned out that n = 3 was sufficient (for seven real phase shifts), and in some cases we adopted n = 2, n = 1 or n = 0. Proper threshold behaviour was assured by multiplying equation (1) by the OPE contributions, obtained from the appropriately calculated OPE elements of the K-matrix. The higher waves, as mentioned above, were taken to be pure OPE ones. It is well known that low-L OPE phase shifts do not correspond to the nuclear-bar phase shifts. We have therefore used the scattering length and effective range energy dependence for ${}^{1}S_{0}$ phase shift.

The energy dependence of the imaginary parts of phase shifts are parametrized as

$$\operatorname{Im} \delta_{LJ}(T) = \begin{cases} 0 & \text{for } T \leq T_{LJ} \\ \sum_{n=0}^{2} \frac{a_{LJn}}{n!} (T - T_{LJ})^{\left(L + n + \frac{1}{2}\right)} & \text{for } T_{LJ} < T \end{cases}$$
(2)

where the threshold energy T_{LI} is proper to each phase shift. It was fixed by the value found in the previous PSA if the imaginary part of a given phase shift was non-zero. For Im ${}^{1}S_{0}$ and Im ${}^{3}P_{1}$ phase shifts the corresponding T_{LJ} were treated as free parameters. The coefficients a_{LJ0} , a_{LJ1} and a_{LJ2} were varied. The parametrization, equation (2), of the threshold behaviour of the imaginary parts of the phase shifts is in agreement with general analyticity and unitarity requirements. Inelastic unitarity was imposed by constraining the Im δ_{LI} to be nonnegative.

3. Data base.

3.1 ANALYSED DATA. — To denote the experimental quantities, we use a four subscript notation [11] X_{srbt} , where the subscripts, s, r, b, t refer to the scattered, recoil, beam and target particle spin orientation, respectively.

In the present analysis we have used all the relevant data available in the compilation of reference [2] as well as recently published data. These latter are listed in table I (Refs. [16-48]). Some data from different experiments at Saturne II, included in this PSA, have only recently been submitted for publication.

A total of 4 356 independent pp data points were analysed. A more detailed repartition of the data is shown in table II. For each type of experiment the total number of data points as well as their occurrences is indicated.

3.2 INELASTIC TOTAL CROSS-SECTIONS. — The inelastic total cross-sections were introduced in the PSA as described in reference [1]. We have fitted the energy dependence of the total cross-section of the reaction j by an expression of the form [49]:

$$\sigma_{j}(T) = (T - T_{oj}) \exp\left[\sum_{i=0}^{n} c_{ij} P_{i}(T)\right],$$
 (3)

where P_i are polynomials of the energy T and the parameters T_{oj} , c_{oj} , ..., c_{nj} are fitted for each reaction *j*. The total inelastic cross-section was then calculated as a sum of all the reaction cross-sections, fitted together with the independent measurements of the σ_{tot} (inelastic). The σ_{tot} (inelastic) were introduced into the PSA in 5 MeV steps with the calculated errors. Our fitted pp reaction cross-section is in good agreement with the one obtained by VerWest and Arndt [50] up to 750 MeV. From there on, the 2π -channel contributions, not taken into account in reference [50], become significant, representing ~ 3 mb at 1 300 MeV. Our treatment of total inelastic cross sections is discussed in a separate article (see Ref. [49b]).

| Observable | Kin. Energy Range (MeV) | Angular Range CM deg. | Total Points | Accelerator | Ref. |
|---------------------------------------|---|--|--|---|---|
| $\sigma_{ m tot}$ | 767-1 260 | 0 | 5 | KEK | [16] |
| $\sigma_{\rm tot}$ (el) | 767-1 260 | 0 | 5 | КЕК | [16] |
| $\Delta \sigma_T$ | 702-1 220 791 741-791 730-1 271 | 0 0 0 0 | 12 1 2 4 | Saturne II LAMPF LAMPF ZGS-ANL | [17] [18] [19] [19] |
| $\Delta \sigma_L$ | 739-790 752-1 293 | 0 0 | 3 5 | LAMPF Saturne II | [20] [21] |
| . $\mathrm{d}\sigma/\mathrm{d}\Omega$ | 746 795 843 892 942 992 992 699-1 203 797 797 800 857 1 078 1 168 1 168 1 168 1 262 | $\begin{array}{c} 7.1-17.3\\ 8.0-17.5\\ 6.8-17.2\\ 7.2-17.3\\ 6.9-16.6\\ 6.5-16.7\\ 3.7-8.0\\ 90\\ 2.9-10.3\\ 1.2-11.4\\ 5.9-91\\ 22.3-88.6\\ 22.3-88.6\\ 22.3-88.6\\ 22.3-88.6\\ 22.3-88.6\\ 22.3-88.6\\ 22.3-88.6\\ \end{array}$ | 25 23 25 24 23 24 39 86 45 41 70 19 19 19 19 19 | Gatchina Gatchina Gatchina Gatchina Gatchina Gatchina Saturne II LAMPF LAMPF LAMPF KEK KEK KEK KEK KEK | [22] [22] [22] [22] [22] [23] [24] [25] [25] [26] [27] [27] [27] [27] [27] |
| A _{oono} | $\begin{array}{c} 699\\ 750\\ 699-800\\ 797\\ 800\\ 800\\ 800\\ 699-1\ 015\\ 702\\ 720\\ 720\\ 720\\ 720\\ 790\\ 834\\ 874\\ 934\\ 995\\ 1\ 095\\ 1\ 095\\ 719\\ 719\\ 894\\ 1\ 145\\ 874\\ 870\\ 1\ 095\\ 1\ 091\\ 1\ 295\\ \end{array}$ | $\begin{array}{c} 34.6-123.0\\ 34.6-122.8\\ \sim 40\\ 1.2-14.0\\ 10.4-47.0\\ 40.1-47.0\\ 14.3-75.4\\ 35-88\\ 42.2-77.7\\ 42.3-77.7\\ 42.2-81.7\\ 40.2-77.6\\ 40.2-77.6\\ 40.2-77.6\\ 40.2-77.6\\ 43.0-86.6\\ 42.8-86.8\\ 42.8-86.9\\ 42.9-86.9\\ 43.4-87.0\\ 46.6-86.1\\ 53.6-88.7\\ 57.2-88.9\\ 59.3-91.7\\ 28.6-47.9\\ 30.0-47.9\\ 28.3-53.5\\ 26.8-52.5\\ 28.6-55.6\end{array}$ | 9 8 3 35 9 4 14 1 050 19 19 20 20 20 20 20 20 20 20 20 20 20 20 23 23 23 23 23 23 23 23 23 23 23 21 19 17 17 17 17 17 19 | LAMPF LAMPF LAMPF LAMPF LAMPF LAMPF LAMPF Saturne II Saturne II | [28] [28] [29] [25] [26] [26] [30] [31] [31] [31] [31] [31] [31] [32] [32] [32] [32] [32] [32] [32] [32 |

Table I. — Recent pp measurements from 1981 to 1986.

| r | ······································ | T | | | |
|-------------------|--|--------------------------|-----------------|-------------|---------|
| Observable | Kin. Energy Range (MeV) | Angular Range CM deg. | Total Points | Accelerator | Ref. |
| | 1 292 | 29.9-52.7 | 7 | Saturne II | [24] |
| A _{oono} | 1 205 | 52 5 00 7 | | Saturne II | [34] |
| | 1 293 | 50.95 | 20 | Saturne II | [34] |
| | 8/4 | 50-85 | 8 | Saturne II | [35] |
| | 890 | 27.6-88.0 | 13 | Gatchina | [36] |
| | 950 | 35-70 | 6 | Gatchina | [37] |
| | 761-950 | 50 | 8 | Gatchina | [38] |
| | 804 | 50, 85 | 3 | Gatchina | [38] |
| | 741-790 | 90 | 2 | LAMPF | [39] |
| | 800 | 4.8-23.8 | 4 | LAMPF | [40] |
| | 800 | 27.3-133.6 | 11 | LAMPF | i411 |
| $D = D_{max}$ | 800 | 23.8-69.2 | 4 | LAMPF | [28] |
| nono | 800 | 70-110 | 3 | Gatchina | |
| | 850 | 50-130 | 0 | Gatchina | |
| | 900 | 50 1 20 | 9 | Catchina | |
| | 900 | 50-150 | 9 | Gatchina | [42] |
| | 699 | 41.9-123.3 | 7 | LAMPF | [28] |
| R = D | 750 | 42.3-134.0 | 7 | LAMPF | [28] |
| | 800 | 4.8-23.8 | 4 | LAMPF | [40] |
| | 800 | 21.2-133.6 | 12 | LAMPF | [41] |
| | 699 | 41.1-123.3 | 8 | LAMPF | [28] |
| $A = D_{s'oko}$ | 750 | 42.3-134.0 | 7 | LAMPF | [28] |
| | 800 | 4.8-23.8 | 3 | LAMPF | 401 |
| | 800 | 21.2-133.6 | 12 | LAMPE | [41] |
| | 874 | 95-130 | 8 | Saturne II | [35] |
| | 800 | 4.8-23.8 | 4 | LAMPE | [40] |
| $R' = D_{k'oso}$ | 800 | 18.8-133.1 | 11 | LAMPF | [28] |
| | 699 | 34.6-123.3 | 8 | LAMPF | [28] |
| <i></i> | 750 | 34.6-134.0 | 8 | LAMPE | [28] |
| $A' = D_{k'oko}$ | 800 | 7 2-23 8 | 3 | LAMPE | |
| | 800 | 18 8-133 1 | 11 | LAMPE | [20] |
| | | 10.0-155.1 | 11 | | [20] |
| | 760-950 | 50, 90 | 10 | Gatchina | [38] |
| | 890 | 27.6-88.0 | 13 | Gatchina | [36] |
| | 950 | 35.0-98.0 | 12 | Gatchina | [37] |
| | 700-800 | 90 | 3 | LAMPF | [43] |
| | 800 | 5.7-37.0 | 23 | LAMPF | [44] |
| | 702 | 42.2-77.7 | 19 | Saturne II | [31] |
| | 720 | 42.3-77.7 | 19 | Saturne II | [31] |
| | 790 | 40.2-77.6 | 20 | Saturne II | [31] |
| A | 834 | 43 0-86 6 | 20 | Saturna II | [37] |
| ^ ~ 00nn | 87 <i>1</i> | 19.0-00.0 19.6 A7 0 | 11 | Saturne II | [32] |
| | 077 071 | 20.0-41.9 17 0 02 0 | 22 | Saturne II | [43] |
| | 0/4 | 42.0-00.0 | 23 | Saturne II | [32] |
| | 934 00 <i>5</i> | 42.8-80.9 | 23 | Saturne II | [32] |
| | 995 | 42.9-86.9 | 23 | Saturne II | [32] |
| | 1 095 | 28.3-52.0 | 13 | Saturne II | [45] |
| | 1 095 | 43.4-87.0 | 23 | Saturne II | [32] |
| | 1 295 | 28.6-55.6 | 15 | Saturne II | [45] |
| | 1 295 | 52.5-90.7 | 20 | Saturne II | [45] |
| | | | | · | L : * J |

Table I (continued).

| Observable | Kin. Energy Range (MeV) | Angular Range CM deg. | Total Points | Accelerator | Ref. |
|-------------------|----------------------------|--------------------------|-----------------|-------------|------|
| | 706 | 26.5-50.5 | 9 | ANL-ZGS | [46] |
| | 740 | 80-98 | 10 | LAMPF | [20] |
| | 790 | 80-98 | 10 | LAMPF | [20] |
| | 800 | 4.5-36.3 | 25 | LAMPF | [44] |
| | 806 | 29.5-50.5 | 8 | ANL-ZGS | [46] |
| | 823 | 29.5-50.5 | 8 | ANL-ZGS | [46] |
| | 1 012 | 29.5-53.5 | 9 | ANL-ZGS | [46] |
| | 1 271 | 29.5-50.5 | 8 | ANL-ZGS | [46] |
| A _{ookk} | 719 | 49.6-87.6 | 11 | Saturne II | [47] |
| | 834 | 47.5-88.5 | 22 | Saturne II | [47] |
| | 874 | 34.5-51.9 | 10 | Saturne II | [48] |
| | 874 | 47.5-88.6 | 22 | Saturne II | [47] |
| | 934 | 46.7-89.5 | 23 | Saturne II | [47] |
| | 995 | 46.8-89.6 | 23 | Saturne II | [47] |
| | 1 095 | 30.5-55.9 | 14 | Saturne II | [48] |
| | 1 095 | 44.7-89.7 | 23 | Saturne II | [47] |
| | 1 295 | 28.6-93.0 | 35 | Saturne II | [48] |
| | 706 | 29.5-47.5 | 7 | ANL-ZGS | [46] |
| | 806 | 29.5-50.5 | 8 | ANL-ZGS | [46] |
| | 823 | 29.5-50.5 | 8 | ANL-ZGS | [46] |
| Aoosk | 874 | 50.4-89.4 | 21 | Saturne II | [47] |
| | 1 012 | 29.5-50.5 | 8 | ANL-ZGS | [46] |
| | 1 271 | 29.5-50.5 | 8 | ANL-ZGS | [46] |
| Nonkk | 874 | 50-85 | 8 | Saturne II | [35] |

Table I (continued).

Table II. — Summary of pp data points between 700 and 1 300 MeV used in PSA.

| - | 697-1 3 | 00 MeV |
|--|----------------------------|---------|
| Experimental Quantity | Number of Points | Percent |
| $\sigma_{tot} \ \sigma_{tot}(el) \ \sigma_{tot}(inel)$ | 108 9 32 | 3.4 |
| $\begin{array}{c} \Delta \sigma_T \\ \Delta \sigma_L \end{array}$ | 21 11 | 0.7 |
| ${ m d}\sigma/{ m d}\Omega$ | 1 292 | 29.7 |
| A _{oono} | 1 906 | 43.8 |
| $D = D_{nono}$ $R = D_{s'oso}$ $A = D_{s'oko}$ $R' = D_{k'oso}$ $A' = D_{k'oko}$ | 58 40 48 15 30 | 4.4 |

| | 697-1 300 MeV | | | | |
|---|---------------------|---------|--|--|--|
| Quantity | Number of Points | Percent | | | |
| A _{oonn} A _{ookk} A _{oosk} | 393 325 60 | 18.0 | | | |
| $N_{_{onkk}}$ | 8 | | | | |
| Total | 4 356 | 100.0 | | | |
| Contributed Points | 4 260 | 97.8 | | | |
| Omitted Points | 96 | 2.2 | | | |

3.3 FITTING PROCEDURE. — All experimental data were fitted according to the standard χ^2 -method, including the error matrix calculation (see e.g. Ref.

[51]) with the statistical errors quoted by the authors. The systematic errors were taken into account by introducing a variable normalization factor multiplying each data set. These factors were kept only if their values differed from one by more than their errors. We found that most of the remaining normalization factors are applied to differential cross-sections. Our χ^2 sum thus correctly reflects the systematic experimental uncertainties. Experimental points for which the χ^2 -contribution was larger than 10 were omitted from the analysis. These represent 2.2% of the data but in about one half of the omitted data (1%) the quoted statistical errors were very small, probably underestimated in original references.

Several incompatible sets which exist below 800 MeV were discussed in reference [1]. Above 800 MeV no important discrepancies between different data sets were found.

The phase shifts varied in the present PSA and the number of free coefficients necessary to describe their energy dependences are listed in the bottom row of table III. The influence of higher *n* terms in equations (1) and (2) on the χ^2 value were carefully tested as well as the influence of all possible imaginary parts of phase shifts. It was found necessary to fit the J = 8 real parts of phase shifts in order to reproduce the spin dependent parameters up to 1.3 GeV. The phase shifts with $J_{max} = 6$ are sufficient to fit the data up to 1.1 GeV. The imaginary parts of all phase shifts not used in our analysis presented a negligible influence on the χ^2 value.

3.4 AMBIGUITIES. — The direct reconstruction of

the scattering matrix and conditions for an unique solution of the PSA have been discussed in detail in references [13, 52-59]. We present here a brief summary relevant to this PSA.

A direct reconstruction of the five complex $pp \Rightarrow pp$ scattering amplitudes at one energy and one scattering angle θ_{CM} requires determining 9 real values from the existing data (one value is arbitrary due to the overall common phase). The set of measured observables allowing the reconstruction of the scattering matrix at one energy and one angle is called « the complete set of observables » [52, 53]. Since the amplitudes are bilinear functions of the observables, the number of observables for a complete set vary from 9 to 18 depending on the choice of experiments [54-57]. An optimal set would contain at least 11 different observables. The complete set of observables in a large angular region was measured for the first time at SIN at few energies and scattering amplitudes were reconstructed [58, 59].

The expansion of amplitudes as a Legendre polynomial series truncated at a given L_{max} and the replacement of the remaining terms by OPE are powerful constraints [60]. For fixed energy PSA these constraints decrease not only the required number of experimental points for each observable, but could, in principle decrease even the number of necessary observables. Indeed, statistical criteria often allow one to reject solutions having a small probability and the PSA can be unambiguous even if the number of independent observables is considerably lower than eleven. In connection with the phase shift analysis, Klepikov [61-63] introduced the concept of a « sufficient set of observables » (called also

| T _{kin} | | | Real Parts of pp Phase Shifts [deg] | | | | | |
|---|--|--|--|---|---|---|--|--|
| (MeV) | ¹ S ₀ | ³ P ₀ | ³ P ₁ | ³ P ₂ | ¹ D ₂ | £2 | ³ F ₂ | ³ F ₃ |
| 700 750 800 850 900 950 1 000 1 050 1 100 1 150 1 200 1 250 1 300 | $\begin{array}{r} - 33.98 \\ - 37.71 \\ - 38.73 \\ - 38.01 \\ - 36.46 \\ - 34.83 \\ - 33.84 \\ - 34.13 \\ - 36.26 \\ - 40.78 \\ - 48.18 \\ - 58.92 \\ - 73.45 \end{array}$ | $\begin{array}{r} - 41.48 \\ - 42.90 \\ - 44.01 \\ - 45.42 \\ - 45.85 \\ - 45.85 \\ - 45.88 \\ - 45.82 \\ - 45.60 \\ - 45.25 \\ - 44.77 \\ - 44.20 \\ - 43.57 \end{array}$ | $\begin{array}{r} - 46.50 \\ - 45.31 \\ - 44.52 \\ - 44.09 \\ - 44.00 \\ - 44.21 \\ - 44.66 \\ - 45.30 \\ - 46.95 \\ - 46.95 \\ - 47.83 \\ - 48.67 \\ - 49.39 \end{array}$ | 20.85 20.26 19.70 19.19 18.73 18.33 18.02 17.78 17.63 17.57 17.60 17.72 17.93 | $\begin{array}{r} 6.69\\ 4.50\\ 2.51\\ 0.71\\ -\ 0.88\\ -\ 2.27\\ -\ 3.45\\ -\ 4.43\\ -\ 5.19\\ -\ 5.74\\ -\ 6.08\\ -\ 6.21\\ -\ 6.12\end{array}$ | $\begin{array}{r} - \ 0.05 \\ - \ 1.40 \\ - \ 2.00 \\ - \ 1.93 \\ - \ 1.31 \\ - \ 0.28 \\ 1.04 \\ 2.49 \\ 3.91 \\ 5.14 \\ 5.99 \\ 6.30 \\ 5.87 \end{array}$ | $\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$ | $\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$ |
| Free Coef. | . 4 | 4 | 4 | 4 | 4 | 4 | 3 | 4 |

Table III. — Real parts of pp phase shifts in degrees.

| T _{kin} | | | Real Parts of pp Phase Shifts [deg] | | | | | |
|---|--|--|--|---|---|---|---|---|
| (MeV) | ³ F ₄ | ¹ G ₄ | ê ₄ | ³ H ₄ | ³ H ₅ | ³ H ₆ | ¹ I ₆ | З |
| 700 750 800 850 900 950 1 000 1 050 1 100 1 150 1 200 1 250 1 300 | 6.46 6.51 6.53 6.51 6.45 6.37 6.26 6.12 5.96 5.77 5.57 5.34 5.10 | 4.25 4.03 3.79 3.54 3.28 3.03 2.77 2.50 2.24 1.98 1.71 1.45 1.19 | $\begin{array}{r} -1.95\\ -2.06\\ -2.10\\ -2.08\\ -2.00\\ -1.84\\ -1.62\\ -1.32\\ -0.94\\ -0.49\\ 0.05\\ 0.67\\ 1.37\end{array}$ | $\begin{array}{r} - 0.35 \\ - 0.43 \\ - 0.52 \\ - 0.61 \\ - 0.70 \\ - 0.81 \\ - 0.91 \\ - 1.02 \\ - 1.13 \\ - 1.25 \\ - 1.37 \\ - 1.49 \\ - 1.62 \end{array}$ | $\begin{array}{r} -1.39\\ -1.40\\ -1.41\\ -1.41\\ -1.40\\ -1.30\\ -1.38\\ -1.36\\ -1.33\\ -1.30\\ -1.27\\ -1.23\\ -1.19\end{array}$ | $ \begin{array}{c} 1.68\\ 1.69\\ 1.68\\ 1.65\\ 1.61\\ 1.55\\ 1.48\\ 1.41\\ 1.32\\ 1.21\\ 1.10\\ 0.98\\ 0.85 \end{array} $ | $ \begin{array}{r} 1.16\\ 1.13\\ 1.08\\ 1.03\\ 0.99\\ 0.93\\ 0.87\\ 0.81\\ 0.75\\ 0.68\\ 0.62\\ 0.55\\ 0.48 \end{array} $ | $\begin{array}{r} -1.19\\ -1.24\\ -1.29\\ -1.33\\ -1.37\\ -1.41\\ -1.45\\ -1.49\\ -1.52\\ -1.55\\ -1.58\\ -1.61\\ -1.64\end{array}$ |
| Free Coef. | 2 | 2 | 3 | 2 | 2 | 2 | 2 | 1 |

Table III (continued).

Table III (continued).

| T_{kin} | Real Parts of pp Phase Shifts [deg] | | | | | | | | |
|--|--|---|--|---|---|--|--|--|--|
| (MeV) | ³ J ₆ | ³ J ₇ | ³ J ₈ | ¹ K ₈ | 88 | ³ L ₈ | | | |
| 700 750 800 850 900 950 1 000 1 050 1 100 1 150 1 200 1 250 | $\begin{array}{r} 0.20\\ 0.13\\ 0.04\\ -\ 0.05\\ -\ 0.15\\ -\ 0.27\\ -\ 0.39\\ -\ 0.52\\ -\ 0.66\\ -\ 0.80\\ -\ 0.96\\ -\ 1.12\end{array}$ | $\begin{array}{r} -1.14\\ -1.31\\ -1.49\\ -1.67\\ -1.86\\ -2.05\\ -2.25\\ -2.46\\ -2.67\\ -2.88\\ -3.10\\ -3.32\end{array}$ | 0.56 0.61 0.65 0.69 0.73 0.77 0.81 0.85 0.89 0.92 0.96 0.99 | $\begin{array}{r} 0.53\\ 0.39\\ 0.23\\ 0.06\\ -\ 0.12\\ -\ 0.31\\ -\ 0.50\\ -\ 0.70\\ -\ 0.91\\ -\ 1.12\\ -\ 1.34\\ -\ 1.56\end{array}$ | $\begin{array}{r} - \ 0.30 \\ - \ 0.32 \\ - \ 0.34 \\ - \ 0.35 \\ - \ 0.37 \\ - \ 0.38 \\ - \ 0.39 \\ - \ 0.41 \\ - \ 0.42 \\ - \ 0.43 \\ - \ 0.44 \\ - \ 0.45 \end{array}$ | 0.09 0.11 0.14 0.16 0.19 0.22 0.25 0.28 0.31 0.35 0.39 0.42 | | | |
| 1 300 | - 1.29 | - 3.55 | 1.02 | - 1.78 | - 0.46 | 0.46 | | | |
| Free Coef. | 2 | 2 | 1 | 2 | 1 | 2 | | | |

« necessary experiments ») allowing to obtain a unique (or more precisely two-valued above the pion production threshold) PSA solution. Klepikov has determined the necessary number of experimental points as well as the number of observables which must be measured. His method can minimize the data taking time at a given energy but necessitates an interactive procedure between measurements and the PSA [64]. Here we restrict ourself to the case when all measured observables are determined with a good precision in a large angular range. In fact, a precise measured angular dependence of one observable at one energy must be fitted by a polynomial with degree L_{max} and, consequently, allows one to determine about L_{max} free parameters. The total cross sections (σ_{tot} , $\Delta \sigma_T$ and $\Delta \sigma_L$) determine one extra free parameter. For L_{max} equal 8 five independent experiments can determine about 40 free parameters at one energy (we use 22 real parts and 8 imaginary parts of phase shifts). Of course, our sufficient set of observables does not assure a unique solution of the PSA.

The number of free parameters for an energy dependent PSA increases by a factor two or three with respect to an energy fixed PSA. On the other hand the number of parameters that can be to determined increases by the same factor (we use 73 free parameters for the present PSA). Moreover, phase shift energy dependence provides additional powerful constraints. For example the continuity of phase shifts within an energy interval automatically decreases the possible multiplicity of solutions.

Concerning the present PSA we know that the solution is unambiguous at least at two energies namely 800 and 1000 MeV. In the vicinity of 800 MeV the number of independent observables corresponds to conditions close to a complete set of experiments in a large angular range. At $\sim 1 \text{ GeV}$ there exist 10 to 11 independent observables in the angular range $40^{\circ} \le \theta_{\rm CM} \le 90^{\circ}$ but some of these have relatively large experimental errors. From 800 to 970 MeV there exist eight independent observables and above 1.04 GeV their number is reduced to five but covering almost the entire angular range. Imaginary parts of the amplitudes at $\theta_{CM} = 0^{\circ}$ are well determined by existing measurements of σ_{tot} , $\Delta \sigma_T$ and $\Delta \sigma_L$. The Gatchina $d\sigma/d\Omega$ data at small angles [22, 23] determine the real parts of amplitudes at $\theta = 0^{\circ}$ below 1 GeV. Energy dependences of the phase shifts are strongly constrained by the $d\sigma/d\Omega$ points at $\theta_{CM} \sim 90^\circ$ measured in a jet-target experiment at Saturne II [24] up to 1.2 GeV. The results of A_{0000} measured in the angular range $35^{\circ} \le \theta_{\rm CM} \le 88^{\circ}$ in small energy steps [30] give

another excellent constraints for phase shift energy dependence below 1 GeV. The extension of the PSA up to 1.3 GeV was possible only by adding the recent analysing power (A_{oono}) and spin correlation $(A_{oonn}$ and $A_{ookk})$ data [31-34, 45, 47, 48] measured at Saturne II with a good precision in a large angular region $(20^{\circ} \le \theta_{CM} \le 100^{\circ})$.

Some of the phase shift energy dependences, mainly above 1 GeV, may change when all results from Saturne II are available.

4. Results and discussion.

Phase shift values found in the present PSA are given in tables III and IV. Both real and imaginary phase shifts are given in degrees. The χ^2 value of the fit per degree of freedom is 1.41.

The starting values for the analysis were taken from our fixed energy solution at 1 GeV [1]. The thresholds of the imaginary parts of S, P, D, F, G and H waves were studied. The thresholds for Im ${}^{1}S_{0}$, Im ${}^{3}P_{1}$ and Im ${}^{3}H_{5}$ are found to be 750, 920 and 900 MeV, respectively, the threshold for other imaginary parts were taken from reference [1] and fixed here. Note, that the contribution of the Im ${}^{3}P_{0}$, found to be negligibly small in our analysis, is hard to separate from that of the other P-wave imaginary parts.

The energy dependences of the S, P, D, F, G, H, I, J, K and L phase shifts and mixing parameters are shown in figures 1-14 together with our previous results below 800 MeV (full lines) and the starting points at 1 GeV (black dots). One can see that

| T (MeV) | Imaginary Parts of pp Phase shifts [deg] | | | | | | | | |
|---------------|--|--------------------------------|--------------------------------|--------------------------------|--------------------------------|--------------------------------|--------------------------------|--------------------------------|--|
| | Im ¹ S ₀ | Im ³ P ₁ | Im ³ P ₂ | Im ¹ D ₂ | Im ³ F ₂ | Im ³ F ₃ | Im ¹ G ₄ | Im ³ H ₅ | |
| 700 | | | 10.03 | 12.09 | 1.44 | 10.33 | 0.29 | | |
| 750 | | | 12.17 | 14.44 | 2.41 | 13.85 | 0.50 | | |
| 800 | 0.04 | | 14.26 | 16.74 | 3.70 | 17.47 | 0.81 | | |
| 850 | 0.20 | | 16.26 | 18.91 | 5.32 | 20.92 | 1.22 | 0.00 | |
| 900 | 0.55 | | 18.16 | 20.89 | 7.27 | 23.96 | 1.73 | 0.00 | |
| 950 | 1.14 | 0.01 | 19.96 | 22.58 | 9.50 | 26.41 | 2.33 | 0.00 | |
| 1 000 | 1.99 | 0.12 | 21.63 | 23.93 | 11.96 | 28.20 | 3.01 | 0.00 | |
| 1 050 | 3.13 | 0.39 | 23.19 | 24.87 | 14.56 | 29.42 | 3.74 | 0.00 | |
| 1 100 | 4.61 | 0.87 | 24.63 | 25.34 | 17.19 | 30.37 | 4.78 | 0.02 | |
| 1 1 50 | 6.43 | 1.61 | 25.96 | 25.30 | 19.68 | 31.62 | 5.16 | 0.06 | |
| 1 200 | 8.63 | 2.64 | 27.20 | 24.70 | 21.86 | 34.09 | 5.73 | 0.17 | |
| 1 250 | 11.24 | 3.98 | 28.35 | 23.51 | 23.50 | 39.12 | 6.09 | 0.40 | |
| 1 300 | 14.26 | 5.66 | 29.43 | 21.72 | 24.34 | 48.48 | 6.15 | 0.88 | |
| Free Coef. | 1 | 1 | 3 | 3 | 2 | 3 | 3 | 1 | |

Table IV. — Imaginary parts of pp phase shifts in degrees.

almost all phase shifts connect smoothly with our solution at lower energies. The apparent disconnection of the Im ${}^{1}G_{4}$ is probably due to the fact that in the lower energy interval (500-800 MeV) we have not allowed ${}^{1}S_{0}$ to be imaginary. The « corridor of errors », shown only for the fixed energy analysis at 1 GeV, was calculated as the square root of the corresponding diagonal element of the error matrix. It is much narrower than the error corresponding to the « confidence level 1 σ » due to the large contribution of the non-diagonal elements in the error matrix, as was discussed in reference [1].

Our results in figures 1-13 are compared with the energy dependent PSA of Arndt et al. [3] (dotdashed line) and with the results of representative fixed-energy phase shift analyses above 800 MeV (triangles [5], open circles [6], diamonds [7]). Below 800 MeV a comparison with the locally energydependent PSA of Dubois et al. [8] is also shown (black squares). The imaginary parts of phase shifts are also compared with the Paris potential prediction [10] (dashed line). All phase shifts have been converted into our representation; this transformation mainly changes the imaginary parts. Only the phase shifts left free are shown in the figures. The comparison of the results below 800 MeV was discussed in reference [1]. The energy dependences above 800 MeV show a general agreement with



Fig. 1. — Energy dependences of ${}^{1}S_{0}$ and ${}^{3}P_{0}$ phase shifts. — Present analysis and PSA of reference [1]; • Previous fixed energy solution at 1 GeV; — Analysis of reference [3]; \bigtriangledown Analysis of reference [5]; \bigcirc Analysis of reference [6]; \diamondsuit Analysis of reference [7]; \blacksquare Analysis of reference [8]; OPE and electromagnetic correction.

 $\begin{array}{c} 0 \\ -20 \\ -40 \\ -60 \\ -80 \\ -$

Fig. 2. — Energy dependence of ${}^{3}P_{1}$ and ${}^{3}P_{2}$ phase shifts. Symbols are defined in figure 1.



Fig. 3. — Energy dependence of the phase shift ${}^{1}D_{2}$. and of the mixing parameter ε_{2} . Symbols are defined in figure 1.



Fig. 4. — Energy dependence of ${}^{3}F_{2}$, ${}^{3}F_{3}$ and ${}^{3}F_{4}$. Symbols are defined in figure 1.



Fig. 6. — Energy dependence of ${}^{3}H_{4}$, ${}^{3}H_{5}$ and ${}^{3}H_{6}$. Symbols are defined in figure 1.





Fig. 5. — Energy dependence of the phase shift ${}^{1}G_{4}$ and of the mixing parameter ε_{4} . Symbols are defined in figure 1.

Fig. 7. — Energy dependences of the phase shift ${}^{1}I_{6}$ and of the mixing parameter ε_{6} . Symbols are defined in figure 1.



Fig. 8. — Energy dependence of the phase shifts $^3J_6,\,^3J_7$ and $^3J_8.$ Symbols are defined in figure 1.



Fig. 9. — Energy dependence of ${}^{1}K_{8}$, ε_{8} and ${}^{3}L_{8}$. Symbols are defined in figure 1.



Fig. 10. — Energy dependence of the imaginary parts of the phase shifts ${}^{1}S_{0}$ and ${}^{3}P_{1}$. The dashed line is prediction from the Paris potential [10]. Other symbols have been defined in figure 1.

Fig. 12. — Energy dependence of the imaginary parts of \triangleright ${}^{3}F_{2}$, and ${}^{3}F_{3}$ phase shifts. The symbols are as in figures 1 and 10.



Fig. 11. — Energy dependence of the imaginary parts of ${}^{3}P_{2}$ and ${}^{1}D_{2}$ phase shifts. Symbols are defined in figures 1 and 10.





Fig. 13. — Energy dependence of the imaginary parts of ${}^1\mathrm{G}_4$ and ${}^3\mathrm{H}_5$ phase shifts.

results of Arndt *et al.* [3] only for the real parts of ${}^{3}P_{1}$, ${}^{1}D_{2}$, ${}^{3}F_{2}$, ${}^{3}F_{4}$, ${}^{3}H_{4}$, ${}^{3}H_{5}$ and ${}^{3}H_{6}$ phase shifts and for Im ${}^{3}P_{2}$ and Im ${}^{1}G_{4}$. All other phase shifts and mixing parameters are in strong disagreement. The results of their very recent version [4] differ only little from reference [3]. A most important disagreement can be seen in figure 2 for the Re ${}^{3}P_{2}$ phase shift from 400 MeV on. The present analysis confirms the behaviour observed in reference [1].

This analysis, as well as that of reference [1] below 800 MeV is in good agreement with the results of reference [8]. It is hard to compare our results with those of other analyses at fixed energy above 800 MeV, since fixed-energy PSA results are very dispersed. For the Re ${}^{1}S_{0}$ phase shift there is very good agreement between the present PSA and the results of references [5, 7]. For the Re ${}^{3}P_{0}$ and Re ${}^{3}P_{1}$ the fixed energy PSA results are between the present solution and that of reference [3]. We found the Re ${}^{3}P_{1}$ to be very close to the OPE contribution. The energy dependence of Re ${}^{3}F_{3}$ agree with the results from references [5, 6]. The imaginary parts differ considerably between different analyses for the ${}^{1}P_{0}$, ${}^{3}P_{1}$, ${}^{1}D_{2}$, ${}^{3}F_{2}$ and ${}^{3}F_{3}$ phase shifts. The solution, given by the present analysis will probably change above 1 GeV, but below this energy only small changes should be expected. Finally, we have compared our results for imaginary parts of phase shifts with the Paris potential group latest predictions [10], calculated from a dispersion relation approach [9]. The global energy trend for the ${}^{1}D_{2}$, ${}^{3}F_{3}$ and ${}^{1}G_{4}$ phase shifts is consistent with ours (see Figs. 11 to 13). A disagreement for S and P phase shifts is not surprizing, since the authors of references [9, 10]claim reasonable predictions starting from J = 3 on. Considering all their imaginary phase shifts, their total inelastic cross section reproduces the energy dependence of the data.

Our earlier single energy solution at 1 GeV [1] differs significantly from the present energy-dependent solution. In the previous analysis, in order to avoid energy dependent developpement of phase shifts, data were considered only in a narrow energy range (0.97-1.04 GeV). This provide us with an uncomplet set of observables between 45° CM and 90° CM with relative large experimental errors.

No new data have appeared in this energy range, but a significant amount (~ 210 data points equivalent to 1/3 of all data used at «1 GeV ») becomes available mainly from KEK, Saturne II and Gatchina in the interval ± 100 MeV around 1 GeV (see Tab. I). Moreover, these new data cover the smallangle region ($\geq 23^{\circ}$ CM). Under these circonstances it is not surprizing that the solutions differ, since the energy dependence represent a powerful constraint as discussed in section 3.4.

The Argand diagrams shown in figures 14a-h, also connect smoothly with our previous results. The Argand diagrams for ${}^{3}P_{2}$, ${}^{1}D_{2}$, ${}^{1}G_{4}$ and ${}^{3}H_{5}$ show an







Fig. 14a-h. — Argand diagrams for - Previous analysis every 40 MeV [1]; • Previous fixed-energy solution at 1 GeV [1]; - Present analysis every 100 MeV; - Analysis of reference [3]; \Box Analysis of reference [8] at 325, 425, 515, 580, 650 and 800 MeV. For each energy only the phase shifts with model-independent imaginary parts are represented.

anticlockwise behaviour in the entire energy range. The change of the initial anticlockwise behaviour of the ${}^{3}F_{3}$ phase shift into a slow clockwise movement above 800 MeV bring into question the existence of a resonance at the total energy of 2.22 GeV, as claimed by many authors. The shape of the Argand diagram for ${}^{3}P_{2}$ (Fig. 14c) may be consistent with a possible resonance structure in one of the P waves, suggested by Lomon [65]. The structure in Re ${}^{3}P_{2}$ above 400 MeV (Fig. 2) observed in reference [1] and confirmed here supports this suggestion. The singlet-spin state partial waves ${}^{1}D_{2}$ and ${}^{1}G_{4}$ show interesting behaviour between 1.1 and 1.3 GeV. To establish a possible resonance-like structure in the ${}^{1}G_{4}$ wave, which may occur at ~1.3 GeV, an extension of the PSA to higher energies is necessary. Our results are compared with those of references [3] (dot-dashed line) and [8] (open squares).

The energy dependences of the real to imaginary part ratio of the spin independent forward amplitude [Re (a + b)/Im(a + b)] and of the total cross section differences $\sigma_{1 \text{ tot}} = -\Delta \sigma_T/2$ and $-\Delta \sigma_L$ are shown in figures 15 to 17 (together with all existing data). We have not used dispersion relation predictions [66] as input into our analysis. Nevertheless, these are in excellent agreement with our solution as can be seen in figure 15. The $\sigma_{1 \text{ tot}}$ (Fig. 16) does not



Fig. 15. — Energy dependence of $\operatorname{Re}(a+b)/\operatorname{Im}(a+b)$ at $\theta = 0$. Dispersion relation prediction [64] is shown as dotted line.



Fig. 16. — Energy dependence of $\sigma_{1 \text{ tot}}(pp) = -\Delta \sigma_T(pp)/2$. — this PSA and reference [1]; — · · · · reference [3]; - · · · · · reference [3]; - · · · · · $\sigma_{1 \text{ tot}}(el)$, this analysis and reference [1]; • Saturne II, reference [17]; \Box TRIUMF, reference [67]; \diamond LAMPF, reference [18]; \circ Rice-LAMPF, reference [19]; \bigtriangledown Rice-ZGS, reference [19]; \bigtriangleup ZGS-ANL, reference 78/B-248/ in [2]; \times ZGS-ANL, reference 75/D-80/ in [2].

connect smoothly with the energy dependence established in reference [1] even if the experimental points are fairly well fitted. There is a lack of constraints below 700 MeV where the absolute value of $\sigma_{1 \text{ tot}}$ increases rapidly. Anyway the difference in the overlapping region is ≤ 0.5 mb which is comparable to the experimental errors. The $-\Delta\sigma_L$ total cross section difference (Fig. 17) connects smoothly with our previous results in the overlapping region. A lack of data around 1.2 GeV allows the decrease in $-\Delta\sigma_L$ values. Our predictions for $\sigma_{1 \text{ tot}}$ and $\Delta\sigma_L$ are compared with the predictions of reference [3]. In figures. 16 and 17 are also plotted predictions for the elastic parts of $\sigma_{1 \text{ tot}}$ and $-\Delta\sigma_L$ as dashed line. These results confirm that the elastic part of $\sigma_{1 \text{ tot}}$ is small and that for $\Delta\sigma_L$ (el) is large.

Our analysis fits well the observed crossing of zero analysing power A_{0000} at $\theta_{\rm CM} = 90^{\circ}$ from negative to positive sign above 1.15 GeV. This behaviour, definitely established at Saturne II [32, 45], can be seen in figure 18 where the angular dependence $A_{0000} = f(\theta_{\rm CM})$ is plotted for three energies.



Fig. 17. — Energy dependence of $-\Delta\sigma_L(pp)$. — this PSA and reference [1]; — reference [3]; — $-\Delta\sigma_L(el)$, this PSA and reference [1]; \blacksquare SIN, reference [68]; • Saturne II, reference [21]; \Box TRIUMF, reference [67]; O LAMPF, reference [20]; ∇ ZGS-ANL, reference 78/A-164/ in [2].

In figure 19 are shown the energy dependences of the spin correlation parameters A_{oonn} , A_{ookk} and A_{ooss} at $\theta_{CM} = 90^{\circ}$ as calculated from this PSA and from reference [1]. Also are plotted all data points above 700 MeV measured close to 90° CM. A local maximum at ~ 1.15 GeV in the A_{oonn} (as well as in A_{ooss}) energy dependence may be due to a lack of experimental data and needs confirmation. On the other hand the decrease observed at 1.3 GeV for all three parameters, is confirmed by existing data above this energy [45, 48]. This supports a possible singlet spin state resonance in ${}^{1}G_{4}$ partial wave around 1.3 GeV.

5. Conclusions.

The present PSA fits well all existing experimental data points including the new Saclay results. The



Fig. 18. — Angular dependences of $A_{\text{cono}} = f(\theta_{\text{CM}})$ at 1 100, 1 200 and 1 300 MeV.

slope of the analysing power near 90° CM is also well described. The energy dependence of the phase shifts smoothly connects with Saclay-Geneva PSA previous results below 800 MeV. An important disagreement is observed with the results of Arndt *et al.* for almost all phase shifts. The energy dependence of the phase shifts above 1.1 GeV may change when the 2 and 3 spin-index data measured at Saturne II are available.

The present PSA shows the general behaviour of phase shifts and provides good predictions for the design of pp elastic and inelastic scattering experiments as well as for values of nucleon-nucleon amplitudes used as input in nucleon-nucleus scattering models. It also provides useful information about a possible resonance-like behaviour in several partial waves.

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Fig. 19. — Energy dependences of the spin correlation parameters A_{oonn} , A_{ookk} and A_{ooss} at $\theta_{CM} = 90^{\circ}$. — this PSA and reference [1]; Symbols for A_{oonn} : ⊟ ZGS-ANL, reference 80/B-304/ in [2]; ⊙ ZGS-ANL, reference 78/L-52/ in [2]; ⊡ ZGS-ANL, reference 77/M-97/ in [2]; □ Saturne I, reference 67/C-26/ in [2]; ⊽ LAMPF, reference 80/M-111/ in [2]; * LAMPF, reference [43]; + Gatchina, reference [36]; ◇ Gatchina, reference [37]; △ Gatchina, reference [38]; • Saturne II, references [32, 45]; Symbols for A_{ookk} : ○ LAMPF, reference [20]; ⊕ ZGS-ANL, reference [69] and reference 78/A-165/ in [2] (numerical values); ■ Saturne II, references [47, 48].

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