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Measurement of cross sections for ²²Na, ²⁰⁻²²Ne and ³⁶⁻⁴²Ar in the spallation of Mg, Al, Si, Ca and Fe. Production ratios of some cosmogenic nuclides in meteorites

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Résumé. — Nous avons mesuré, par spectrométrie de masse des gaz rares et par comptage γ , les sections efficaces σ des réactions nucléaires conduisant à ²²Na, ²⁰⁻²²Ne et ³⁶⁻⁴²Ar dans des cibles de Mg, Al, Si, Ca et Fe bombardées par des protons de 0,059 à 24 GeV. Ces nouvelles mesures et celles existant dans la littérature nous ont permis d'évaluer les rapports des sections efficaces $\sigma_{\rm R}/\sigma_{\rm S}$ en fonction de l'énergie de quatre paires d'isotopes, l'un stable S et l'autre radioactif $R : {}^{20}\text{Na}/{}^{20}\text{Ne}, {}^{26}\text{Al}/{}^{21}\text{Ne}, {}^{39}\text{Ar}/{}^{38}\text{Ar}$ et ${}^{36}\text{Cl}/{}^{36}\text{Ar}$. Ces paires d'isotopes permettent de calculer l'âge d'exposition des météorites, si on connaît le rapport de leur taux de production $P_{\rm R}/P_{\rm S}$. Notre travail établit des limites aux rapports $P_{\rm R}/P_{\rm S}$ dans les météorites de fer à partir de la seule considération des rapports $\sigma_{\rm R}/\sigma_{\rm S}$. Dans les météorites de pierre on peut seulement montrer, par cette approche, l'importance des effets de composition chimique et de blindage.

Abstract. — We measured the cross sections of ²²Na, ^{20,21,22}Ne and ^{36,38,39,42}Ar in nuclear reactions induced by the proton bombardment of Mg, Al, Si, Ca and Fe at several energies between 0.059 and 24 GeV using highsensitivy mass-spectrometry for noble gases and γ -counting for ²²Na. We then evaluated the cross section ratio *versus* energy for the ²²Na/²²Ne, ²⁶Al/²¹Ne, ³⁹Ar/³⁸Ar and ³⁶Cl/³⁶Ar radioactive-stable pairs of isotopes, which are useful for measuring the exposure age of meteorites. From the results alone, we can set some limits to the corresponding production ratios in iron meteorites, but not in stony objects because a knowledge of chemical abundances of Na, Mg, Al, Si, S, K, Ca, Fe, and shielding conditions are required.

1. Introduction.

Cosmic-ray-produced nuclides provide sensitive tools for studying exposure history of extra-terrestrial materials such as meteorites and lunar rocks. These nuclides are produced through spallation reactions induced by primary and secondary cosmic-ray particles in the near surface (up to a few meters) of bodies in the solar system that are not radiation-shielded by an atmosphere or strong magnetic fields. A detailed study of the exposure history of a meteoritic or lunar sample should comprise several chapters dealing with : the duration of exposure (or exposure age T); the flux of galactic and solar cosmic rays and its variations in time and space; the shielding history of the sample, involving possibly multiple breakups; the space erosion by micrometeorites. The measurement of as many cosmogenic nuclides as possible is required to decipher the relative importance of so many variables.

In the present paper we investigate a group of methods for measuring T, which involve a pair of one stable S and one radioactive R cosmogenic nuclides. The production rate P of any cosmogenic nuclide may be defined as the number of atoms produced per time unit through nuclear reactions. It depends on the intensity and energy spectrum of cosmic radiations and also on the chemical composition and the shielding depth of the sample. The growth curve for a radioactive nuclide is such as $R = P_{R}(1 - e^{-\lambda T})$. After several half-lives the activity R is in secular equilibrium and $R = P_{R}$. So a given isotope R integrates the cosmic ray flux over several half-lives. The concentration of a stable nuclide is given by $S = S_0 + P_S T$ where S_0 is the contribution of sources other than spallation reactions. In favorable cases (mainly isotopes of noble gases), S_0 can be neglected or can be subtracted. Ps is generally unknown. Assuming cons-

^(*) E.R.A. nº 144 du C.N.R.S.

tancy for $P_{\rm R}$ and $P_{\rm S}$, a simple expression for exposure age is thus :

$$T = \frac{P_{\rm R}}{P_{\rm S}} \frac{S}{R} \tag{1}$$

if secular equilibrium has been reached for the radioisotope and if the meteorite did'nt have a complex history. Convenient R-S pairs involve cosmogenic nuclides that are produced by similar nuclear reactions and will generally have very similar atomic masses. The key parameter in such a method is the production ratio $P_{\rm R}/P_{\rm S}$, because S and R are measured quantities in the sample. At first [1-5], many production ratios were calibrated against some adopted exposure age. This approach does not provide an independent method and ignores likely effects from target chemistry and shielding. Also, absolute and independent calibrations are necessary to compare ages from different methods and obtain additional information about the exposure history of a sample or about the temporal variations in the fluxes of cosmic-ray particles. However some production ratios have been tentitavely inferred from cross section data [6-9]. Independent calibrations were used for the ⁸¹Kr-⁸³Kr method [10] and for the ⁴⁰K-⁴¹K method which applies only to iron meteorites [11]. We will focus here on four *R-S* pairs : 22 Na- 22 Ne, 26 Al- 21 Ne, 39 Ar- 38 Ar and ³⁶Cl-³⁶Ar and determine the limits of their production ratios $P_{\rm R}/P_{\rm S}$ starting from cross sections ratios $\sigma_{\rm R}/\sigma_{\rm S}$.

We measured new cross sections for ²²Na, ²⁰⁻²²Ne and ³⁶⁻⁴²Ar in Mg, Al, Si, Ca and Fe targets bombarded with 59 MeV to 24 GeV protons. From these new measurements and from literature data, we deduced excitation functions $\sigma(E)$ for ²²Na, ²²Ne, ²⁶Al, ²¹Ne, ³⁹Ar, ³⁸Ar, ³⁶Cl and ³⁶Ar in targets that are abundant in meteorites. Neutron-induced reactions were mostly estimated. We then examined the cross section ratios $\sigma_{\rm R}/\sigma_{\rm S}$ and tried to deduce limits on the corresponding $P_{\rm R}/P_{\rm S}$ ratios and their possible variations with chemistry and shielding.

2. Experimental.

Pure target foils of Mg (99.9 % or 3N), A1 (5N) and Fe (5N), and also crystals of Si (4N) and CaF₂ (4N) were stacked, 2 or 3 at a time, with monitor and guard A1 (5N) foils. Exact thickness was determined by weighing and was (20 to 50) 10^{-3} mm for foils, and 0.6 to 1 mm for crystals. The targets were irradiated between 1969 and 1980 in the accelerators of Louvain-La-Neuve (0.059 and 0.075 GeV), Orsay (0.150, 0.168 and 0.200 GeV), CERN-SC (0.600 GeV), Saturne (1.0, 1.05 and 2.5 GeV) and CERN-PS (24 GeV). Both internal and external proton beams were used at 1 and 24 GeV, internal beam only at 0.150 and 0.600 GeV, and external beam only at 0.059, 0.075, 0.168, 0.200 and 2.5 GeV. Special care was taken to minimize secondary particle and recoil effects.

Proton fluences were measured according to the 27 Al(p, X) 22 Na monitor reaction whose cross sections are 28.1, 23.6, 17, 16.7, 16.1, 15.3, 11.7 and 10.0 mb, respectively at 0.059, 0.075, 0.150, 0.168, 0.200, 1, 2.5 and 24 GeV [12]. Between 2 × 10¹⁵ and 4 × 10¹⁷ protons struck the targets. By comparing results from either internal and external irradiation, or young and old targets, no diffusion loss for gaseous products neon and argon was observed. More details about targets and irradiations are given in references [13] and [14].

Some Mg, Si and Ca targets were counted for the 1.27-MeV γ peak of ²²Na in a 4-keV resolution Ge(Li) gamma ray spectrometer, calibrated with ²²Na standards.

Noble gases Ne and Ar were analysed by means of a Micromass 60° sector, 12 cm radius mass-spectrometer equipped with a gas extraction and purification line, two activated charcoals and two calibration systems. After at least two years following the irradiation (2 to 12 years) targets and blank foils were stored, up to 10 at a time, in the extraction-purification line in which an ultrahigh vacuum (around 10^{-9} torr) was progressively attained by pumping and baking. Targets were baked two days between 60 and 80 °C. We checked that this procedure did not outgas any significant amount of reaction-produced neon or argon. Each foil was then melted inside a molybdenum crucible heated under vacuum by an electron-bombardment oven, between 700 and 1 600 °C depending on the metal or crystal. Attempts to separate trapped gases of atmospheric isotopic composition and nuclear reaction-produced gases by step-wise heating were unsuccessful. The noble gases were purified in several steps, using Ti, CuO-Pd and Al-Zr getters between 200 and 800° for eliminating chemically active species. Argon was adsorbed on a first activated charcoal trap at -196 °C. About half of the neon was introduced into the mass spectrometer. During isotopic analysis. most memory argon was adsorbed on a second charcoal trap at -196 °C. Interferences at m/e = 20 from $H_2^{18}O^+$ or ${}^{40}Ar^{2+}$ ions, and at m/e = 22 from CO_2^{2+} ions were monitored by measuring also m/e = 18, 40 and 44 peaks. No interference was detected at m/e = 21. In most cases, the second half of neon was used for isotopic dilution and was measured independently after mixing with a known amount of atmospheric neon. In Ca and Fe targets, argon was also analysed. Almost all ⁴⁰Ar was atmospheric, except a small part (always less than 1%) deduced in an iterative procedure from the charge-dispersion curves of products. Corresponding corrections were applied to ³⁶Ar and ³⁸Ar. Background peaks due to hydrocarbon ions had to be corrected at m/e = 39 and m/e = 42. In this work, flat top peaks were obtained for a resolution of 250 (1 % valley). Calibration was made by frequently introducing air standards from a pipette systems or from an independent system composed of several known volumes and measured pressure using a capacitance manometer. The isotope

ratios were corrected for memory or pumping effects, isotopic discrimination, blank, hydrocarbons (for 38,39,42 Ar) and interfering ions. Complete outgassing of melted targets was checked in reextraction runs, and a small correction (0 to 1 % in most cases) was applied. Depending on cross sections, target thicknesses and fraction of gases that was analysed, 10⁷ to 10¹¹ atoms of neon 20, 21, 22 or argon 36, 38, 39, 40 and 42 were measured.

3. Results.

The cross sections measured for 22 Na in the targets Mg, Si and Ca, of $^{20-22}$ Ne in Mg, Al, Si, Ca and Fe, and of $^{36-42}$ Ar in Ca are reported in tables I, II and III.

The results are averages of N independent measurements weighed by individual uncertainties. The root mean square errors included the various error sources for each measurement : target preparation and impurities; monitoring and alignment; secondary particles; blanks, interfering ions or hydrocarbons (the associated error was 50 % of the correction); zero time corrections to mass-spectrometer analysis; re-extractions; calibration. Recoil corrections for spallation (0 to 3 % in most cases) were calculated according to Sauvageon [15]. Correction for radioactive decay was made for ³⁹Ar and ⁴²Ar. Also ²²Ne was corrected for decay of ²²Na in table II. Many errors cancel out for the isotope ratios.

3.1 ²²Na. — Our measurements in Mg (Table I) agree within errors with those of Raisbeck and Yiou, 1975 [16]. Additional cross sections at other energies can be found in Korteling and Caretto, 1970 [17] and the compilation of de Lassus and Tobailem, 1972 [18]. Excitation function $\sigma(E)$ in Si was already well documented from Raisbeck and Yiou [16] and Korteling and Caretto [17], and also from the compilation of Tobailem, 1981 [19]. In Ca, there was only one

Table I. — Cross sections σ (in mb) for ²²Na in Mg, Si and Ca, as functions of the incident energy E_p (in GeV). N is the number of independent measurements.

Target	E p (GeV)	N	σ (mb)
Mg	1.0	1	33.1 ± 3.3
	2.5	1	22.2 ± 2.2
	24.0	1	21.7 ± 2.2
Si	0.150	3	18.5 ± 1.9
	24.0	2	15.2 ± 1.5
Ca	0.168	1	0.44 <u>+</u> 0.04
	1.0	1	6.7 ± 0.7
	2.5	1	7.1 ± 0.7

measurement at 400 MeV from Korteling and Caretto, 1967 [20].

3.2 20,21,22 Ne. — Excitation functions for stable Ne isotopes were quite well documented in Mg, Al and Si targets in the low energy part $E_p < 70$ MeV (see Ref. [21, 22], and our results notably extend our knowledge to higher energies. No neon cross section in Ca was measured until now. Our measurements (Table II) disagree in a systematic way with those of Goebel *et al.*, 1964 [8] as is evident in figure 1 where we plotted excitation functions for $^{20-22}$ Ne in Fe. However we agree with Bieri and Rutsch, 1962 [23].

3.3 ${}^{36-42}$ Ar. — Ar measurements in Ca are reported (Table III) at 0.168, 1.0 and 2.5 GeV. 39 Ar cross sections at 1.0 and 2.5 GeV are mostly upper limits, due to strong hydrocarbon interference at m/e = 39. Previous results were available at 600 MeV (Stoenner *et al.*, 1970 [24]) and 12 to 45 MeV (Walton *et al.*, 1976 [25]). We also measured argon isotopes in iron targets, and our results were in agreement with those of Regnier, 1979 [14]. This agreement shows that there is no significant loss of spallation neon or argon in « old » iron targets.

Table II. — Cross sections $\dot{\sigma}$ (in mb) for neon isotopes in Mg, Al, Si, Ca and Fe, as functions of the incident energy E_p (in GeV). Values for ²²Ne are either corrected for ²²Na, or cumulative (²²Ne + ²²Na). The ²²Na cross section adopted for correction took into account all published values and not only that of table I. The cumulative cross sections in the last column are more reliable because they imply much less correction to the measured ²²Ne. N is the number of independent measurements.

Target	E _p (GeV)	N	20 _{Ne} or(mb)	21 _{Ne} or(mb)	22 _{Ne} or(mb)	22 _{Ne +} 22 _{Na} or(mb)
Mg	0.059	2	43.3 ± 5.2	40.9 ± 4.9	15.2 ± 2.0	77.2 ± 10.2
	0.075	2	45.2 ± 5.3	43.5 ± 5.1	18.3 ± 2.4	69.3 ± 9.1
	0.168	1	25.4 ± 3.3	27.1 ± 3.6	10.6 ± 1.6	41.6 ± 6.3
	0.200	2	28.2 ± 3.4	30.0 ± 3.7	-12.6 ± 1.7	42.6 ± 5.7
	1.0	1	35.8 ± 5.9	33.6 ± 4.9	. 22.2 ± 3.7	51.8 ± 8.0
	2.5	1	14.6 ± 4.3	18.4 ± 2.8	8.2 ± 1.4	31.7 ± 5.4
	24.0	1	20.6 ± 2.7	20.9 ± 2.8	9.0 ± 1.3	30.7 ± 4.4
Al	0.059	2	14.3 ± 1.8	32.1 ± 4.1	9.9 ± 1.5	38.0 ± 5.8
	0.075	2	26.1 ± 3.1	32.9 ± 3.9	11.5 ± 1.5	35.1 ± 4.6
	0.150	2	22.6 ± 2.5	24.3 ± 2.7	10.4 ± 1.2	27.4 ± 3.2
	0.168	2	22.0 ± 2.5	23.6 ± 2.7	10.1 ± 1.2	26.8 ± 3.2
	0.200	2	26.3 ± 2.9	27.7 ± 3.1	12.9 ± 1.5	29.0 ± 3.4
	1.0	з	22.2 ± 2.8	22.4 ± 2.8	11.3 ± 1.6	26.6 ± 3.8
	1.05	2	28.5 ± 3.4	29.4 ± 3.4	16.5 ± 2.0	31.8 ± 3.9
	2.5	3	21.2 ± 2.4	21.9 ± 2.5	12.0 ± 1.3	23.7 ± 2.6
	24.0	2	18.1 ± 2.2	18.8 ± 2.4	10.4 ± 1.7	20.4 ± 3.3
Si	0.150	3	17.0 ± 1.8	15.8 ± 1.7	4.1 ± 0.5	21.3 ± 2.6
	24.0	2	17.8 ± 2.0	15.9 ± 1.8	4.9 ± 0.6	18.6 ± 2.3
Ca	0.168	1	0.83 ± 0.09	0.53 ± 0.06	0.21 ± 0.04	0.65 ± 0.12
	1.0	1	11.9 ± 1.4	11.3 ± 1.3	5.3 ± 0.9	12.0 ± 2.0
	2.5	1	12.4 ± 1.4	11.4 ± 1.3	5.1 ± 0.8	12.2 ± 1.9
Fe	0.600	2	1.04 ± 0.20	1.04 ± 0.20	0.80 ± 0.15	1.07 ± 0.20
	2.5	1	6.5 ± 0.8	6.5 ± 0.8	4.3 ± 0.6	6.5 ± 0.9
	24.0	2	8.0 ± 1.2	7.8 ± 0.9	5.6 ± 0.9	7.9 ± -1.3

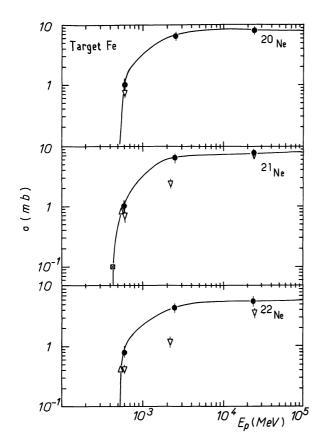


Fig. 1. — Measured cross sections σ in mb for the production of ²⁰Ne, ²¹Ne and ²²Ne in iron bombarded with protons of energy *E* in MeV. Filled symbols : this work ; others : ∇ ref. [8], \triangle ref. [23], \boxtimes ref. [40]. Solid lines are adopted excitation functions (see text).

Table III. — Cross sections σ (in mb) for ${}^{36-42}$ Ar in Ca, versus energy E_p (in GeV). Large error bars for 39 Ar are due to hydrocarbon interference at 1.0 and 2.5 GeV.

	σ ₃₆ (mb)	о ⁷ 38 (mb)	σ ₃₉ (mb)	σ ₄₂ (mb)
1	39.9 ± 4.1	82.6 ± 8.3	1.12 ± 0.12	0.032 ± 0.006
1	29.0 ± 3.0	58.4 ± 5.9	1.1 ± 0.9	
1	20.6 ± 2.1	42.1 ± 4.2	0.9 ± 0.7	
	1 1 1	1 39.9 ± 4.1 1 29.0 ± 3.0	1 39.9 ± 4.1 82.6 ± 8.3 1 29.0 ± 3.0 58.4 ± 5.9	$\begin{array}{c} 1 \\ 39.9 \pm 4.1 \\ 29.0 \pm 3.0 \\ 58.4 \pm 5.9 \\ 1.1 \pm 0.7 \\ 1.1 \\ 1.1 \\ 0.7 $

4. Discussion.

From our results and others found in literature, we examined the σ_R/σ_S cross section ratios, versus incident energy, of some radioactive-stable pairs of isotopes that are in common use for calculating the exposure age T of meteorites. Our general approach was to obtain excitation functions for all appropriate nuclear reactions, either from experimental cross sections or from systematics. We largely used the Silberberg and Tsao [26] formula for spallation of Mg to Fe targets at more than 100 MeV. When possible we calibrated the formula using experimental cross sections. In the lower energy part we used nuclear

systematics from analog reactions in neighbouring targets, applying appropriate corrections for different thresholds. We also used the Keller *et al.* [27] method as well as N/Z ratio corrections. Secondary neutrons are very important in meteorites, especially when size and geometry conditions favour their growth. At energies lower than 100 MeV, they are responsible for most inelastic collisions. More details are discussed in Baros [13]. From the excitation functions, we computed the cross section ratios versus energy curves.

4.1 22 Na/ 22 Ne. — In figure 2 we plotted the cross section ratio $\sigma(^{22}Na)/\sigma(^{22}Ne)$ versus energy for incident protons (solid lines) or neutrons (broken lines) in several target elements. In stony meteorites such as chondrites (about 90 % of observed falls), the main target elements for both 22 Na and 22 Ne are Mg and Si; then there is some contribution from Al, S, Ca. Fe and other elements. Iron meteorites are composed of Fe-Ni alloys with microscopic to centimeter-sized inclusions. In chondrites there is a clear difference between the cross section ratio inferred from : i) proton-induced reactions, mostly from measured values, giving $\sigma(^{22}Na)/\sigma(^{22}Ne)$ ratios between 0.5 and 0.8, ii) neutron induced reactions in Mg, corresponding to $\Delta A \lesssim 4$. Neutron channels are mostly evaluated from systematics. The $Mg(n, x)^{22}Ne$ excitation function was taken from Reedy et al., 1979 [28], for energies below 20 MeV. The cross-section ratio $\sigma(^{22}Na)/$ σ ⁽²²Ne) was evaluated by Nishiizumi *et al.*, 1980 [7] to be between 0.50 and 0.65. We agree with their upper limit but the present results tend to slightly decrease the lower limit in cases where low energy neutrons are important. The situation is simpler in iron meteorites where mostly particles with energy higher than 500 MeV are effective, specifically for samples reasonably free of sulfur or phosphorus-rich inclusions, From cross sections only, the production ratio $P(^{22}\text{Na})/P(^{22}\text{Ne})$ is 0.27 \pm 0.03 in iron meteorites.

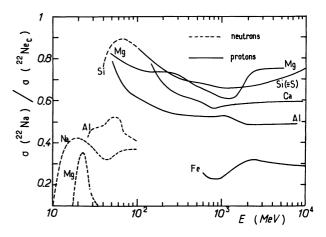


Fig. 2. — Cross section ratio $\sigma(^{22}Na)/(\sigma(^{22}Na) + \sigma(^{22}Ne))$ versus incident energy *E* for protons (solid lines) and neutrons (dashed lines) in several targets. Low energy neutrons favour ²²Ne over ²²Na in the very important magnesium target.

4.2 26 Al/ 21 Ne. — A priori this pair of radioactivestable isotopes is not a good candidate for calculating exposure ages, because ²⁶Al and ²¹Ne have quite different atomic masses. Thus they are expected to arise from different reactions in targets such as Al or Si, and of course in Mg where ²⁶Al is only produced through the ${}^{26}Mg(p, n)$ channel, while ${}^{21}Ne$ is supplied by the important $^{24}Mg(n, \alpha)$ reaction among others. Also ²⁴Mg is the most abundant isotope of Mg. Neutrons are not able to produce ²⁶Al from Mg, except for very rare reactions like ${}^{26}Mg(n, n\pi^{-}){}^{26}Al$. However there exists a great deal of both ²¹Ne and ²⁶Al data in meteorites and we investigate in what extent the ²⁶Al/²¹Ne method is reliable in considerations about nuclear reactions. In figure 3, we show the cross section ratio $\sigma({}^{26}\text{Al})/\sigma({}^{21}\text{Ne})$ versus energy in Mg to Fe targets. The ordinate covers more than 6 orders of magnitude although the neutron channels in Mg are not reported (they give near 0 values). ²⁶Al excitation functions were largely deduced from the Saclay compilations [12, 18, 19] and some unpublished results from this laboratory [29]. The 24 Mg(n, α)²¹Ne excitation function deduced by Reedy et al., 1979 [28], from their measurements at 14.1 and 14.7 MeV was adopted, but the shape of the peak is quite uncertain. The ²⁷Al(n, 2n)²⁶Al channel was estimated at low energy from the compilation of Stehn et al., 1964 [30] about the ²⁶Al^m data. According to Norman, 1979 [31], the relative yields of ²⁶Al^g and ²⁶Al^m deduced from the ²⁵Mg(p, γ) reaction in the 0.3-1.72 MeV energy range are in the ratio $Y(m)/Y(g) \simeq$ 1/3. However, Reedy [32], examining the first 52 levels of ²⁶Al, found that the ratio of levels decaying to the ground state is about the same as those going to the meta state when the excitation energies are less than 6 MeV, so that the Y(m)/Y(g) ratio would be near unity. The last point of view was adopted in figure 3 because it applies to (n, 2n) reactions specifically.

When stony meteorites are considered, Al and Si target elements (high ratio) compete with Mg (low ratio) whatever the energy range. A knowledge of the balance between chemical and energy effects requires the use of models of interaction of cosmic rays and meteorites, but that is not the purpose of the present paper.

In iron meteorites there are two important sources :

i) Fe and Ni targets give ratios between 0.3 and 0.4, and only more than 500 MeV particles are effective.

ii) S and P targets are only present at the 0.2-0.5 % concentrations (except in big inclusions) but they yield a cross section ratio 1 to 3. Particles with energies as low as 100 MeV are effective and they favour ²⁶Al over ²¹Ne. In iron meteorites, we thus expect a cross section ratio ²⁶Al/²¹Ne that depends on both chemical abundances of phosphorus and sulfur and shielding conditions (i.e. the number of 100 to 500 MeV secondary particles). However, high abundances of phosphorus or sulfur are exceptional [33] and from simple assumptions a cross section ratio $\sigma(^{26}Al)/\sigma(^{21}Ne) =$

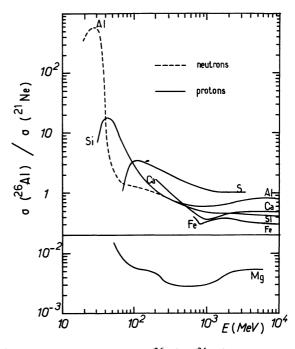


Fig. 3. — Cross section ratio $\sigma({}^{26}\text{Al})/\sigma({}^{21}\text{Ne})$ versus incident energy *E* for protons (solid lines) and neutrons (dashed lines) in several targets,

 0.37 ± 0.05 was deduced, in good agreement with the 0.38 value of Hampel and Schaeffer [6].

4.3 ³⁹Ar/³⁸Ar. — The most important targets for regulating the $\sigma({}^{39}\text{Ar})/\sigma({}^{38}\text{Ar})$ cross section ratio are K, Ca and Fe. The ratio is represented in figure 4 *versus* incident energy. Note the very different behaviour of neutron- and proton-induced reactions in Ca or K targets. The importance of the ${}^{39}\text{K}(n, p){}^{39}\text{Ar}$ reaction channel was first recognized in chondrites

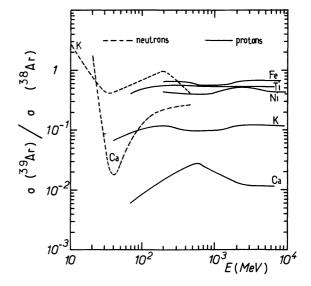


Fig. 4. — Cross section ratio $\sigma({}^{39}\text{Ar})/\sigma({}^{38}\text{Ar})$ versus incident energy *E* for protons and neutrons in several targets. The neutron induced reactions are mostly evaluated and should be checked by some measurements.

by Begemann *et al.*, 1967 [34]. It has a zero threshold and its cross section is 15 to 400 mb between 1 and 20 MeV. However its excitation function is not very well known between 4 and 12 MeV, that is the region of the maximum. The situation in chondrites is quite complicated since the cross section ratio depends on both the nature of the target (K, Ca or Fe) and the energy range of interest. In iron meteorites, disregarding Ca- or K-rich inclusions, only Fe and Ni produce significant amounts of Ar. The ³⁹Ar/³⁸Ar ratio is well known [14] in both targets and is fixed between 0.42 and 0.58 at all energies. From excitation functions only, a cross section ratio of 0.50 ± 0.05 is deduced in iron meteorites, in agreement with the 0.52 value in use for years in literature [9].

4.4 36 Cl/ 36 Ar. — This pair of isotopes is not very well suited in chondrites because the cross-section ratio 36 Cl/ 36 Ar surely depends on both chemistry and shielding conditions through the 39 K(n, α) 36 Cl, 40 K(n, $p\alpha$) 36 Cl reactions, and the 35 Cl(n, γ) 36 Cl reaction in big meteorites. Also it is very difficult to correct for trapped 36 Ar in chondrites. This method is well suited for iron meteorites and Schaeffer and Heymann, 1965 [9] proposed a cross section ratio of 0.82. From new measurements in Fe and Ni [14, 29, 35, 36] the cross section ratio $\sigma({}^{36}$ Cl)/ $\sigma({}^{36}$ Ar) is fixed between 0.78 and 0.81 at all energies in iron meteorites.

5. Conclusions.

We measured new cross sections for ²²Na, ²⁰⁻²²Ne and ³⁶⁻⁴²Ar in Mg, Al, Si, Ca and Fe. Incident proton energies were 0.059 to 24 GeV. From our measurements, others from this laboratory, and a survey of literature we built excitation functions for ²²Na, ²²Ne, ²⁶Al, ²¹Ne, ³⁹Ar, ³⁸Ar, ³⁶Cl and ³⁶Ar in targets that produce significant amounts of these isotopes through cosmic ray induced nuclear reactions. We then deduced the σ_R/σ_S cross section ratios versus energy for the radioactive-stable pairs of isotopes ²²Na-²²Ne, ²⁶Al-²¹Ne, ³⁹Ar-³⁸Ar and ³⁶Cl-³⁶Ar.

With respect to nuclear reactions only, the corresponding production ratios $P_{\rm R}/P_{\rm S}$ commonly used to calculate the exposure age of meteorites are expected to depend on both the concentrations of the main target elements and the shielding conditions (thus, the distribution of interacting protons and neutrons *versus* energy) in a particular sample. However the cases of stony and iron meteorites are quite different.

Iron meteorites are chemically well differentiated. They are mostly composed of Fe-Ni alloys with some inclusions. Isotopes with A = 21 to 39 are all highenergy products except in some specific elements where they could be produced with high yields in inclusions at energies lower than 300 to 500 MeV, such as ²⁶Al in phosphorus and sulfur. Keeping this in mind, the limits for iron meteorites deduced from this work are $0.24 \leq P({}^{22}Na)/P({}^{22}Ne) \leq 0.30$; $0.32 \leq P(^{26}\text{Al})/P(^{21}\text{Ne}) \leq 0.42; \ 0.45 \leq P(^{39}\text{Ar})/$ $P(^{38}\text{Ar}) \leq 0.55 \text{ and } 0.78 \leq P(^{36}\text{Cl})/P(^{36}\text{Ar}) \leq 0.81.$ The main target elements for producing ²²Na, $^{20-22}$ Ne and 26 Al in stony meteorites are Mg and Si. Magnesium plays a prominent part in regulating the 22 Na/ 22 Ne and 26 Al/ 21 Ne production ratios, and the low energy neutron induced reactions in this element are still insufficiently known. From nuclear reactions only, no good limits can be deduced for the ${}^{26}Al/{}^{21}Ne$ production ratio. It is unlikely that this ratio is constant in the different classes of chondrites and for varying shielding conditions. In fact, there is no way to find a single $P({}^{26}\text{Al})/P({}^{21}\text{Ne})$ value that fits exposure ages obtained from others methods [37], in a wide selection of chondrites. We offer some reservations regarding the ²²Na/²²Ne ratio, since the importance of neutron induced reactions on Mg has still to be clarified. Also, 22 Na (half-life = 2.6 years) integrates the cosmic radiation over a very short time and appears to vary over a solar cycle as measurements by Evans et al., 1982 [38] in meteorite falls have shown. The ³⁹Ar/³⁸Ar cross section ratio appears also very sensitive to chemical abundances and energy.

More information about the above production ratios in meteorites could be obtained from models of interaction of cosmic rays with spherical or nearly spherical objects. We are presently working on such models, inspired by the Reedy-Arnold model of interaction between cosmic rays and the Moon surface [39]. Another useful result of the present work was to identify the reaction channels for which experimental information is most needed, such as the lowenergy neutron-induced reactions in K and Mg, and also in Na, Si and Ca. Such measurements are planned in the near future, and will provide a better knowledge of neutron-induced reactions in meteorites.

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Appendix. — Production of ³⁶Cl by high energy spallation. — Due to a mistake in the calculation of the proton fluences, the ³⁶Cl cross section reported from this laboratory [36] are incorrect. Corrected values, together with some new measurements [29], also from this laboratory, are given in table IV.

Table IV. — Measured production cross sections σ of ³⁶Cl in Fe, Ti and Ca as functions of the incident energy E_p . N is the number of independent measurements. The CaF₂ target irradiated at 0.15 GeV showed possible evidence for overheating in the accelerator.

E _p (GeV)	Target	N	$ \begin{array}{c} \sigma(^{36}\text{Cl}) \\ \text{(mb)} \end{array} $
0.15 1.05 24	Fe Ti Ca Fe Fe Ti Ca	1 2 1 3 3 3 2	$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$

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