

A characterization of Generalized functions of Bounded Deformation

Antonin Chambolle, Vito Crismale

▶ To cite this version:

Antonin Chambolle, Vito Crismale. A characterization of Generalized functions of Bounded Deformation. 2025. hal-04946846v1

HAL Id: hal-04946846 https://hal.science/hal-04946846v1

Preprint submitted on 13 Feb 2025 (v1), last revised 15 Feb 2025 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons Attribution 4.0 International License

A CHARACTERIZATION OF G(S)BD

ANTONIN CHAMBOLLE AND VITO CRISMALE

ABSTRACT. We show that Dal Maso's GBD space, introduced for tackling crack growth in linearized elasticity, can be defined by simple conditions in a finite number of directions of slicing.

1. INTRODUCTION

In [21] G. Dal Maso introduced the space of Generalized (Special) functions with Bounded Deformation (G(S)BD) to characterize the domain of a class of functionals associated to free discontinuity problems [22, 7] in linearized elasticity or elastoplasticity, in particular for the modeling of crack growth [23]. In these contexts the energy functionals depend on a variable $u: \Omega \subset \mathbb{R}^d \to \mathbb{R}^d$, representing the displacement, through its symmetrized gradient $e(u) = (Du+Du^T)/2$ (the deformation) and its (d-1)-dimensional jump set J_u (the crack set), with a constraint imposing - in a weak sense - Dirichlet boundary conditions. Yet an integrability bound on u in terms of a bound for such functionals is in general not available, so that it is not possible to control the deformation in the space of Radon measures.

An analogous issue arises both in the simplified *antiplane shear* setting, where the displacement is only in the vertical direction so it can be assumed scalar, and in the framework of finite elasticity or elastoplasticity, contexts in which the role of e(u) is played by the full gradient ∇u . This motivated the introduction of the space G(S)BV [4, 5] as the space of functions which are in (S)BV after composition $\psi \circ u$ with C^1 functions ψ having $\nabla \psi$ compactly supported.

However, $e(\psi \circ u)$ depends on the whole gradient of u, even for u smooth, which prevents from adopting a similar definition of G(S)BD. The solution found in [21] was to define this space by slicing: a measurable function $u: \Omega \to \mathbb{R}^d$ is in $GBD(\Omega)$ if there exists a bounded Radon measure Λ_u such that for every $\xi \in \mathbb{S}^{d-1}$ and for \mathcal{H}^{d-1} -a.e. $z \in \xi^{\perp}$ the slice $s \mapsto u_z^{\xi}(s) := \xi \cdot u(z + s\xi)$ has bounded variation in $\Omega_z^{\xi} := \{s \in \mathbb{R} : z + s\xi \in \Omega\}$ and

$$(\widehat{\mu}_{u})^{\xi}(B) := \int_{\xi^{\perp}} \left((|D^{a}u_{z}^{\xi}| + |D^{c}u_{z}^{\xi}|)(B_{z}^{\xi}) + \sum_{s \in J_{u}\xi \cap B_{z}^{\xi}} |[u_{z}^{\xi}](s)| \wedge 1 \right) d\mathcal{H}^{d-1}(z) \le \Lambda_{u}(B)$$

for every $B \subset \Omega$ Borel, with $B_z^{\xi} := \{s \in \mathbb{R} : z + s\xi \in B\}$. Still in [21], it is shown that it is enough to have the bound for all ξ in a dense subset of \mathbb{S}^{d-1} . However, it is well known that a much simpler characterization exists for the space BD of functions with *Bounded Deformation*: a measurable u is in $BD(\Omega)$ as soon as there is a basis $(e_i)_{i=1}^d$ of \mathbb{R}^d such as the following holds: for any $\xi \in \tilde{V} := \{e_i : i = 1, \ldots, d\} \cup \{e_i + e_j : 1 \le i < j \le d\}$ and for \mathcal{H}^{d-1} -a.e. $z \in \xi^{\perp}$, $u_z^{\xi} \in BV(\Omega_z^{\xi})$ and

$$|\mathrm{E} u\xi \cdot \xi|(\Omega) = |\mathrm{D}_{\xi}(u \cdot \xi)|(\Omega) = \int_{\xi^{\perp}} |Du_{z}^{\xi}|(\Omega_{z}^{\xi})d\mathcal{H}^{d-1}(z) < +\infty,$$

cf. [6, Section 3] and [30, Chap. II, Section 2.2] (see also [31, 30, 8] for general properties of BD functions). It is therefore natural to ask whether GBD functions can, similarly, be described by conditions in a finite number of appropriate directions. Our main result answers this affirmatively.

We show that a measurable u is in $GBD(\Omega)$ as soon as there exists a basis $(e_i)_{i=1}^d$ of \mathbb{R}^d and d(d-1)/2 vectors $\xi_{i,j} \in \text{span}\{e_i \pm e_j\}, 1 \le i < j \le d$ such that for any $\xi \in V = \{e_i : i = 1, \ldots, d\} \cup \{\xi_{i,j} : 1 \le i < j \le d\}$ and for \mathcal{H}^{d-1} -a.e. $z \in \xi^{\perp}, u_z^{\xi} \in BV(\Omega_z^{\xi})$ and

$$\Lambda_u^{\xi} := (\widehat{\mu}_u)^{\xi}(\Omega) < +\infty$$

In other words, we show that if we have a control on the d(d+1)/2 directions $\xi \in V$ through Λ_u^{ξ} , then the same control holds in fact for all $\xi \in \mathbb{S}^{d-1}$, and can be localized on Borel subsets $B \subset \Omega$ and controlled by a global finite measure Λ_u .

Our strategy is to approximate any measurable u with Λ_u^{ξ} bounded for all $\xi \in V$, by a family of functions, each in a suitable finite dimensional space, satisfying a uniform bound in *GBD*, according to the definition in [21] (Theorem 1). More precisely, for any $\varepsilon > 0$, we choose a family of discretizations of the domain into hypercubes of sidelength $\varepsilon > 0$ so that u is approximated in measure by the functions defined in any hypercube as the d-linear interpolation over the vertices (see Appendix B). Then we modify such functions, setting them to 0 on the "bad hypercubes", which are such that there are $\xi \in V$ and a couple of vertices joined by a segment parallel to ξ that intersect the set where the jump of u is larger than 1. Thanks to an averaging argument, we show that we can choose the discretizations in such a way that (i) the bound on Λ_u gives a uniform control on the perimeter of the union of the "bad hypercubes" while (ii) outside this set, the symmetric gradient of the d-linear approximation is controlled in L^1 by the finite differences along directions in V, thus again by the bound on Λ_u (see Appendix A).

From this approximation we derive (Corollary 1) that $u \in GBD$ by employing a suitable compactness result for bounded sequences in GBD that converge in measure [21, 25]. This strategy of discretization/reinterpolation for (G)BD functions was first used in [11, 12, 29] to show approximation results in (G)SBD.

As a further application (Corollary 2), we prove that any sequence of measurable functions satisfying uniformly the control on $(\Lambda_{u_k}^{\xi})_{\xi \in V_k}$ is precompact in *GBD*, in the sense of the general compactness and closedness result in [19] (we refer also to [3] for a related compactness result in *GBD*).

Eventually, we address an open question raised in the recent work [1] (see also [2]) concerning a nonlocal approximation of Griffith-type energies for brittle cracks (see e.g. [23, 13, 24, 28, 15, 20, 14, 16, 10, 27, 9, 26]). In some approximation regimes any limiting admissible deformation unaturally belong to the space therein called $GBV^{\mathcal{E}}(\Omega; \mathbb{R}^d)$, namely it satisfies the bound

$$\int_{\mathbb{S}^{d-1}} \widehat{\mu}_u^{\xi}(B) \, d\mathcal{H}^{d-1}(\xi) \le \Lambda_u(B) \quad \text{for } B \subset \Omega \text{ Borel},$$

for a suitable bounded Radon measure Λ_u . We prove that $GBV^{\mathcal{E}}$ coincide with GBD (Corollary 3), in particular we give positive answer to a question in [1] concerning the rectifiability of the measure obtained by integrating over directions ξ the jump parts of the measures $(\hat{\mu}_u)^{\xi}$.

In the same spirit, we prove (Proposition 2) a characterization of the space $GSBD^p$ (i.e. the GBD functions with slices in SBV, symmetrized gradient in L^p and jump set of finite \mathcal{H}^{d-1} -measure, see e.g. [24]), with a slight adaptation of the arguments in the proof of Theorem 1, and the corresponding compactness (Corollary 4), employing the general compactness result [18], see also [17].

2. Main results

Here Ω is a bounded¹ open set of \mathbb{R}^d , $d \geq 2$. Given $u : \Omega \to \mathbb{R}^d$, for any $\xi \in \mathbb{R}^d \setminus \{0\}$ and $z \in \xi^{\perp}$, we define:

$$u_z^{\xi}(s) := \xi \cdot u(z + s\xi) \quad \text{for } s \in \Omega_z^{\xi} := \{s \in \mathbb{R} : z + s\xi \in \Omega\}$$

We denote by $L^0(\Omega; \mathbb{R}^m)$, $m \ge 1$, the measurable functions from Ω into \mathbb{R}^m and by $v_k \to v$ in $L^0(\Omega; \mathbb{R}^m)$ we mean that v_k converge to v in \mathcal{L}^d -measure in Ω , for \mathcal{L}^d the Lebesgue measure on \mathbb{R}^d . We also let $Q := [0, 1)^d$ and, for any $\varepsilon > 0$,

$$\Omega_{\varepsilon} := \{ x \in \Omega \colon \operatorname{dist}(x, \partial^* \Omega) > \sqrt{d\varepsilon} \}.$$

We refer to [7] for general theory of BV functions. If $u \in BV$ we denote by $D^a u$ and $D^c u$ the absolutely continuous and the Cantor part of the bounded Radon measure Du.

Theorem 1. Let $u \in L^0(\Omega; \mathbb{R}^d)$ and assume there is an orthonormal basis $(e_i)_{i=1}^d$ of \mathbb{R}^d such that for any $\varepsilon > 0$ it holds that $\xi \in V = (e_i)_{i=1}^d \cup (e_i + e_j)_{1 \le i < j \le d}$, and that for \mathcal{H}^{d-1} -a.e. $z \in \xi^{\perp}$, $u_z^{\xi} \in BV(\Omega_z^{\xi})$ and

$$\Lambda_{u}^{\xi} := \int_{\xi^{\perp}} \left((|D^{a}u_{z}^{\xi}| + |D^{c}u_{z}^{\xi}|)(\Omega_{z}^{\xi}) + \sum_{s \in J_{u_{z}^{\xi}}} |[u_{z}^{\xi}(x)]| \wedge 1 \right) d\mathcal{H}^{d-1}(z) < +\infty.$$
(1)

Then there exist a dimensional constant C > 0 and a family $(u^{\varepsilon})_{\varepsilon > 0} \subset SBV^{\infty}(\Omega; \mathbb{R}^d)$ such that, for any $\varepsilon > 0$, $J_{u_{\varepsilon}} \subset \partial \mathcal{B}^{\varepsilon}$ with $\mathcal{B}^{\varepsilon}$ union of finitely many cubes of sidelength ε included in Ω ,

$$\int_{\Omega_{\varepsilon}} |e(u^{\varepsilon})| dx + \mathcal{H}^{d-1}(\partial^* \mathcal{B}^{\varepsilon} \cap \Omega_{\varepsilon}) \le C \Lambda_u^V, \quad \Lambda_u^V := \sum_{\xi \in V} \Lambda_u^{\xi}, \tag{2}$$

and $u^{\varepsilon} \to u \in L^0(\Omega; \mathbb{R}^d)$.

Remark 1. In Theorem 1, we assumed $(e_i)_{i=1}^d$ orthonormal for simplicity. The result would hold also for an arbitrary basis: this is easily checked by considering a matrix A which sends $(e_i)_{i=1}^d$ to an orthonormal basis and replacing u with $A^{-T}u(A^{-1}\cdot)$. One could of course also consider the integrals in (1) over arbitrary spaces E_{ξ} with $E_{\xi} \oplus \mathbb{R}\xi = \mathbb{R}^d$ rather than just ξ^{\perp} , see [21, Remark 4.11]. Additionally, as in [21, Remark 4.3], one could replace in (1) the threshold 1 (in $|[u_z^{\xi}]| \wedge 1$) by an arbitrary (ξ -dependent) threshold level $\beta_{\xi} > 0$, since $|x| \wedge \alpha \leq (1 \vee \alpha/\beta)(|x| \wedge \beta)$ for any $\alpha, \beta > 0$. Finally one could consider, for some pairs (i, j), slicing in the direction $e_i - e_j$ rather than $e_i + e_j$.

From Theorem 1 we deduce the following three consequences. We refer to [1, Definition 2.3] for the definition of the space $GBV^{\mathcal{E}}(\Omega; \mathbb{R}^d)$, characterized in Corollary 3.

Corollary 1. For any fixed basis $(e_i)_{i=1}^d$ of \mathbb{R}^d and $V := (e_i)_{i=1}^d \cup (e_i + e_j)_{1 \le i < j \le d}$ it holds that $GBD(\Omega) = \{ u \in L^0(\Omega; \mathbb{R}^d) : \Lambda_v^V < +\infty \},$ (3)

where, for any $u \in L^0(\Omega; \mathbb{R}^d)$, Λ_u^V is defined in (1), (2).

Corollary 2. Let $(u_k)_k \subset GBD(\Omega)$ such that for every $k \in \mathbb{N} \setminus \{0\}$ there exists $(e^k)_k$, $e^k := (e_i^k)_{i=1}^d$ orthonormal basis of \mathbb{R}^d with (recalling (1), (2))

$$\sup_{k\in\mathbb{N}}\Lambda_{u_k}^{V_k}<+\infty,$$

¹for simplicity, since the result could easily be localized.

for $V_k := (e_i^k)_{i=1}^d \cup (e_i^k + e_j^k)_{1 \le i < j \le d}$. Then there exist a Caccioppoli partition $\mathcal{P} = (P_n)_n$ of Ω , a sequence of piecewise rigid motions $(a_k)_k$ with

$$a_k = \sum_{n \in \mathbb{N}} a_k^n \chi_{P_n},\tag{4a}$$

 $|a_k^n(x) - a_k^{n'}(x)| \to +\infty \quad \text{for a.e. } x \in \Omega, \text{ for all } n \neq n', \tag{4b}$

and $u \in GBD(\Omega)$ such that, up to a (not relabelled) subsequence,

$$u_k - a_k \to u \quad a.e. \text{ in } \Omega,$$
 (5a)

$$\mathcal{H}^{d-1}(\partial^* \mathcal{P} \cap \Omega) \le \lim_{\sigma \to +\infty} \liminf_{k \to \infty} \mathcal{H}^{d-1}(J_{u_k}^{\sigma}).$$
(5b)

where $J_{u_k}^{\sigma} := \{ x \in J_{u_k} : |[u_k]|(x) \ge \sigma \}.$

Corollary 3. It holds that

$$GBV^{\mathcal{E}}(\Omega; \mathbb{R}^d) = GBD(\Omega)$$

We now provide analogous results for the space $GSBD^p$ of GBD functions with symmetrized gradient in L^p and jump set of finite \mathcal{H}^{d-1} -measure, see e.g. [15].

Proposition 2. Let $u \in L^0(\Omega; \mathbb{R}^d)$, p > 1, and assume there is an orthonormal basis $(e_i)_{i=1}^d$ of \mathbb{R}^d such that for any $\varepsilon > 0$ it holds that $\xi \in V = (e_i)_{i=1}^d \cup (e_i + e_j)_{1 \le i < j \le d}$, and almost $z \in \xi^{\perp}$, $u_z^{\xi} \in SBV^p(\Omega_z^{\xi})$ with $D^a u_z^{\xi} = (u_z^{\xi})' \mathcal{L}^1$ and

$$\Lambda_u^{p,\xi} := \int_{\xi^\perp} \left(\int_{\Omega_z^\xi} |(u_z^\xi)'|^p \, dt + \# J_{u_z^\xi} \right) d\mathcal{H}^{d-1}(z) < +\infty.$$
(6)

Then there exist a constant C > 0 depending only on d, p, and a family $(u^{\varepsilon})_{\varepsilon>0} \subset SBV^{\infty}(\Omega; \mathbb{R}^d)$ such that, for any $\varepsilon > 0$, $J_{u_{\varepsilon}} \subset \partial \mathcal{B}^{\varepsilon}$ with $\mathcal{B}^{\varepsilon}$ union of finitely many cubes of sidelength ε included in Ω ,

$$\int_{\Omega_{\varepsilon}} |e(u^{\varepsilon})|^p dx + \mathcal{H}^{d-1}(\partial^* \mathcal{B}^{\varepsilon} \cap \Omega_{\varepsilon}) \le C \Lambda_u^{p,V}, \quad \Lambda_u^{p,V} := \sum_{\xi \in V} \Lambda_u^{p,\xi}, \tag{7}$$

and $u^{\varepsilon} \to u \in L^0(\Omega; \mathbb{R}^d)$. In particular $u \in GSBD^p(\Omega)$.

Corollary 4. With the notation of Proposition 2, let $(u_k)_k \subset GSBD^p(\Omega)$, $(e^k)_k$, $(V_k)_k$ with $e^k := (e_i^k)_{i=1}^d$ orthonormal basis of \mathbb{R}^d , $V_k := (e_i^k)_{i=1}^d \cup (e_i^k + e_j^k)_{1 \le i < j \le d}$ be sequences such that

$$\sup_{k\in\mathbb{N}}\Lambda^{p,V_k}_{u_k}<+\infty$$

Then there exist a Caccioppoli partition $\mathcal{P} = (P_n)_n$ of Ω , a sequence of piecewise rigid motions $(a_k)_k$ with

$$a_k = \sum_{n \in \mathbb{N}} a_k^n \chi_{P_n} \,, \tag{8a}$$

$$a_k^n(x) - a_k^{n'}(x)| \to +\infty \quad \text{for a.e. } x \in \Omega, \text{ for all } n \neq n',$$
(8b)

and $u \in GSBD^{p}(\Omega)$ such that, up to a (not relabelled) subsequence,

$$u_k - a_k \to u \quad a.e. \text{ in } \Omega,$$
 (9a)

$$e(u_k) \rightharpoonup e(u) \quad in \ L^p(\Omega; \mathbb{M}^{d \times d}_{sum}),$$
(9b)

$$\mathcal{H}^{d-1}(J_u \cup \partial^* \mathcal{P}) \le \liminf_{k \to \infty} \mathcal{H}^{d-1}(J_{u_k}).$$
(9c)

From now on we prove the announced results, following the order of presentation.

Proof of Theorem 1. For $\xi \in V$ and $z \in \xi^{\perp}$ such that $u_z^{\xi} \in BV(\Omega_z^{\xi})$, we denote by μ_z^{ξ} the measure (on Ω_z^{ξ}):

$$\mu^{\xi}_z = |D^a u^{\xi}_z| + |D^c u^{\xi}_z| + \sum_{s \in J_{u^{\xi}_z}} (|[u^{\xi}_z](s)| \wedge 1) \mathcal{H}^0 \, {\sqsubseteq} \, \{s\}.$$

Fixed $\varepsilon > 0$, we evaluate the integral

$$\varepsilon^{d-1} \int_Q \sum_{\xi \in V} \sum_{i \in \varepsilon \mathbb{Z}^d} \left(\left| \xi \cdot (u(\varepsilon y + i + \varepsilon \xi) - u(\varepsilon y + i)) \right| \wedge 1 \right) dy$$

where here the sum is over $i \in \varepsilon \mathbb{Z}^d$ such that $\varepsilon y + i \in \Omega \cap (\Omega - \varepsilon \xi)$. Given $\xi \in V$, then the change of variable $x = i + \varepsilon y$ shows that:

$$\begin{split} \varepsilon^{d-1} \int_{Q} \sum_{i \in \varepsilon \mathbb{Z}^{d}} \left(|\xi \cdot (u(\varepsilon y + i + \varepsilon \xi) - u(\varepsilon y + i))| \wedge 1 \right) dy \\ &= \varepsilon^{-1} \int_{\Omega \cap (\Omega - \varepsilon \xi)} \left(|\xi \cdot (u(x + \varepsilon \xi) - u(x))| \wedge 1 \right) dx, \end{split}$$

and writing $x = z + s\xi$, $z \in \xi^{\perp}$, $s \in \Omega_z^{\xi} \cap (\Omega_z^{\xi} - \varepsilon)$ we find that this is also expressed as

$$\varepsilon^{-1} \int_{\xi^{\perp}} \int_{\Omega_z^{\xi} \cap (\Omega_z^{\xi} - \varepsilon)} \left(|u_z^{\xi}(s + \varepsilon) - u_z^{\xi}(s)| \wedge 1 \right) |\xi| ds \, d\mathcal{H}^{d-1}(z).$$

Now, since (for \mathcal{H}^{d-1} -a.e. z) $u_z^{\xi} \in BV(\Omega_z^{\xi})$, then for a.e. s: either there is $s' \in (s, s + \varepsilon) \cap J_{u_z^{\xi}}$ with $|[u_z^{\xi}](s)| \ge 1$, and then clearly $|u_z^{\xi}(s + \varepsilon) - u_z^{\xi}(s)| \land 1 \le \mu_z^{\xi}((s, s + \varepsilon))$, or there isn't, in which case, $\mu_z^{\xi} \bigsqcup (s, s + \varepsilon) = |Du_z^{\xi}| \bigsqcup (s, s + \varepsilon)$, which implies that $|u_z^{\xi}(s + \varepsilon) - u_z^{\xi}(s)| \le \mu_z^{\xi}((s, s + \varepsilon))$. It follows that

$$\begin{split} \int_{\Omega_z^{\xi} \cap (\Omega_z^{\xi} - \varepsilon)} \Big(|u_z^{\xi}(s + \varepsilon) - u_z^{\xi}(s)| \wedge 1 \Big) |\xi| \, ds \\ & \leq |\xi| \int_{\Omega_z^{\xi} \cap (\Omega_z^{\xi} - \varepsilon)} \int \chi_{(s, s + \varepsilon)}(t) d\mu_z^{\xi}(t) \, ds \leq \varepsilon |\xi| \int_{\Omega_z^{\xi}} d\mu_z^{\xi}(t) \, ds \leq \varepsilon$$

where we have used Fubini's theorem and that $\chi_{(s,s+\varepsilon)}(t) = \chi_{(t-\varepsilon,t)}(s)$. Combining the previous estimates, we end up with

$$\begin{split} \varepsilon^{d-1} \int_Q \sum_{\xi \in V} \sum_{i \in \varepsilon \mathbb{Z}^d} \left(|\xi \cdot (u(\varepsilon y + i + \varepsilon \xi) - u(\varepsilon y + i))| \wedge 1 \right) dy \\ & \leq \sum_{\xi \in V} |\xi| \int_{\xi^\perp} \int_{\Omega_z^{\xi}} d\mu_z^{\xi} d\mathcal{H}^{d-1}(z) \leq \sum_{\xi \in V} |\xi| \Lambda_u^{\xi} =: M. \end{split}$$

Hence, there is a set $Q^{\varepsilon} \subset Q$ of measure at least 1/2 such that for any $y \in Q^{\varepsilon}$:

$$\varepsilon^{d-1} \sum_{\xi \in V} \sum_{i \in \varepsilon \mathbb{Z}^d} \left(|\xi \cdot (u(\varepsilon y + i + \varepsilon \xi) - u(\varepsilon y + i))| \wedge 1 \right) \le 2M.$$
⁽¹⁰⁾

Let us choose $y^{\varepsilon} \in Q^{\varepsilon} \cap Q_{\varepsilon}$, where $Q_{\varepsilon} \subset Q$ is the set given by Proposition 7 in Appendix B. For $\varepsilon > 0$ small enough, we define the function u^{ε} for every $x \in \Omega$ as

• the multilinear interpolation of the values of u at the vertices of the cube $\varepsilon y^{\varepsilon} + i + \varepsilon [0, 1)^d$, $(i \in \varepsilon \mathbb{Z}^d)$ which contains x if these vertices are all in Ω and if for all $\xi \in V$ and any pair of vertices $(j, j + \xi)$ of the cube,

$$|\xi \cdot (u(\varepsilon y^{\varepsilon} + j + \varepsilon \xi) - u(\varepsilon y^{\varepsilon} + j))| \le 1;$$

• 0, else.

Observe that u^{ε} is a polynomial in the cubes of the first type, so in particular it does not jump therein, that it is continuous accross a facet between two such cubes, and that the number of cubes of the second type is bounded by $(2/\varepsilon)^{d-1}2M$ so that, denoting by $\mathcal{B}^{\varepsilon}$ the union of cubes of the second type, it holds that

$$J_{u^{\varepsilon}} \subset \partial \mathcal{B}^{\varepsilon}, \quad \mathcal{H}^{d-1}(\partial \mathcal{B}^{\varepsilon}) \le CM, \quad |\mathcal{B}^{\varepsilon}| \le CM\varepsilon$$
 (11)

for a dimensional C > 0. In particular, the convergence properties of the multilinear interpolations stated in Proposition 7 also hold for u^{ε} , since they differ, in Ω , on a set of vanishing measure. In addition, if $\overline{Q}_{\varepsilon} = \varepsilon y^{\varepsilon} + i + \varepsilon [0, 1)^d$ is a cube of the first type, we have that

$$\int_{\overline{Q}_{\varepsilon}} |e(u^{\varepsilon})| dx \leq C \varepsilon^{d-1} \sum_{\xi \in V} |\xi \cdot (u(\varepsilon y^{\varepsilon} + i + \varepsilon \xi) - u(\varepsilon y^{\varepsilon} + i))|$$

by Proposition 4 (in Appendix A) for p = 1, and a change of variables.

Therefore, summing the above two estimates over the cubes all included in Ω (which cover Ω_{ε}), we find that, for ε small enough, $u^{\varepsilon} \in SBV^{\infty}(\Omega)$ with

$$\int_{\Omega_{\varepsilon}} |e(u^{\varepsilon})| dx + \mathcal{H}^{d-1}(J_{u^{\varepsilon}} \cap \Omega_{\varepsilon}) \le 2CM,$$

for a dimensional constant C > 0, and this gives (2) recalling that $M = \sum_{\xi \in V} |\xi| \Lambda_u^{\xi}$.

Proof of Corollary 1. Let $(u^{\varepsilon})_{\varepsilon>0}$ be the sequence provided by Theorem 1 and fix $\Omega' \subset \Omega$. By the slicing properties of BV functions and the bound (2) it follows that defining $\Lambda^{\varepsilon} := |e(u^{\varepsilon})|dx + \mathcal{H}^{d-1} \sqcup \partial \mathcal{B}^{\varepsilon}$, then $\Lambda^{\varepsilon}(\Omega') \leq C \Lambda_u$ and one has, for all $\xi \in \mathbb{S}^{d-1}$ and all Borel $B \subset \Omega'$,

$$\int_{\xi^{\perp}} \left(\int_{B_z^{\xi} \setminus J_{(u^{\varepsilon})_z^{\xi}}} |e((u^{\varepsilon})_z^{\xi})| dt + \#(J_{(u^{\varepsilon})_z^{\xi}} \cap B_z^{\xi}) \right) d\mathcal{H}^{d-1}(z) \le \Lambda^{\varepsilon}(B)$$
(12)

so that $(u^{\varepsilon})_{\varepsilon}$ are equibounded in $GBD(\Omega')$, in the sense of [21].

Moreover, as $u^{\varepsilon} \to u$ in $L^{0}(\Omega; \mathbb{R}^{d})$, by [19, Theorem 1.1] we conclude that $u \in GBD(\Omega')$. In fact, the convergence of u^{ε} to u implies that the infinitesimal rigid motions can be taken as equal to 0 (alternatively, one could use [21, Corollary 11.2 and Theorem 11.3] together with the fact that since $u^{\varepsilon} \to u$ in $L^{0}(\Omega; \mathbb{R}^{d})$, there is an increasing function $\psi_{0} \colon \mathbb{R}^{+} \to \mathbb{R}^{+}$ with $\lim_{s \to +\infty} \psi_{0}(s) = +\infty$ such that $\|\psi_{0}(u^{\varepsilon})\|_{L^{1}(\Omega)}$ is uniformly bounded w.r.t. $\varepsilon > 0$, see e.g. [25, Lemma 2.1]). Then we conclude that $u \in GBD(\Omega)$ by the arbitrariness of $\Omega' \subset \subset \Omega$.

Proof of Corollary 2. For any $u_k \operatorname{let} (u_k^{\varepsilon})_{\varepsilon>0}$ be the approximating family provided by Theorem 1. Then for any $k \in \mathbb{N} \setminus \{0\}$ we can find $\varepsilon_k > 0$ such that, for $\tilde{u}_k := u_k^{\varepsilon_k}$ it holds that $\tilde{u}_k \in SBV^{\infty}(\Omega; \mathbb{R}^d)$ with $J_{\tilde{u}_k} \subset \partial^* \mathcal{B}^{\varepsilon_k}$ and $\mathcal{B}^{\varepsilon_k}$ union of finitely many cubes of sidelength ε_k included in Ω , such that

$$\|\arctan(\tilde{u}_k - u_k)\|_{L^1} \le \frac{1}{k},$$

$$\int_{\Omega_{\varepsilon_k}} |e(\tilde{u}_k)| dx + \mathcal{H}^{d-1}(\partial^* \mathcal{B}^{\varepsilon_k} \cap \Omega_{\varepsilon_k}) \le C \Lambda_{u_k}^{V_k}.$$
(13)

Recalling (12), it follows that the sequence $(\tilde{u}_k)_k$ satisfies the assumptions of [19, Theorem 1.1] on any $\Omega_{\varepsilon_{\overline{k}}}$ for $k \geq \overline{k}$, and therefore by that compactness result and by a diagonal argument on Ω_{ε_k} , up to a (not relabelled) subsequence, there are a Caccioppoli partition $\mathcal{P} = (P_n)_n$ of Ω , a sequence of piecewise rigid motions $(a_k)_k$ satisfying (4), and $u \in GBD(\Omega)$ such that

$$\tilde{u}_k - a_k \to u \quad \text{a.e. in } \Omega.$$
 (14a)

By the first in (13) we deduce (5a). At this stage, (5b) follows arguing exactly as in [19, Section 3.3: Lower semicontinuity]. \Box

Proof of Corollary 3. We prove that $GBV^{\mathcal{E}}(\Omega; \mathbb{R}^d) \subset GBD(\Omega)$, the opposite inclusion being true by definition of $GBD(\Omega)$ in [21]. The proof follows by averaging. We fix an orthonormal basis $(e_i)_{i=1}^d$, and let $\widehat{V} = \{e_i, i = 1, \dots, d; \frac{e_i + e_j}{\sqrt{2}}, 1 \leq i < j \leq d\}$. Then, we have that:

$$\int_{SO(d)} \sum_{\xi \in \widehat{V}} \Lambda_u^{R\xi} d\mu(R) = \frac{d(d+1)}{2} \int_{\mathbb{S}^{d-1}} \Lambda_u^{\xi} d\mathcal{H}^{d-1}(\xi)$$

(for μ the Haar measure over SO(d)), which is finite whenever $u \in GBV^{\mathcal{E}}(\Omega; \mathbb{R}^d)$. We then pick any rotation for which $\sum_{\xi \in R \widehat{V}} \Lambda_u^{\xi}$ is less than average and consider the basis $(Re_i)_{i=1}^d$ in Theorem 1. We conclude since $\Lambda_u^{R(e_i+e_j)} = \sqrt{2} \Lambda_u^{R(e_i+e_j)/\sqrt{2}}$.

Proof of Proposition 2. We argue as in the proof of Theorem 1, with the measures μ_z^{ξ} replaced by the measures

$$(\mu^p)_z^{\xi} := |(u_z^{\xi})'|^p d\mathcal{L}^1 + \mathcal{H}^0 \bigsqcup J_{u_z^{\xi}}$$

thus evaluating

$$\varepsilon^{d-1} \int_Q \sum_{\xi \in V} \sum_{i \in \varepsilon \mathbb{Z}^d} \left(|\xi \cdot (u(\varepsilon y + i + \varepsilon \xi) - u(\varepsilon y + i))|^p \wedge 1 \right) dy \leq \sum_{\xi \in V} \int_{\xi^\perp} \int_{\Omega_z^{\xi}} d\mu_{\xi, z}^p d\mathcal{H}^{d-1}(z) \leq \Lambda_u^{p, V}.$$

Therefore the proof proceeds as done for Theorem 1, first choosing suitably y^{ε} from the previous inequality and Proposition 7 in Appendix B, and then replacing $|\xi \cdot (u(\varepsilon y^{\varepsilon} + i + \varepsilon \xi) - u(\varepsilon y^{\varepsilon} + i))|$ by $|\xi \cdot (u(\varepsilon y^{\varepsilon} + i + \varepsilon \xi) - u(\varepsilon y^{\varepsilon} + i))|^p$, to define the approximating functions u^{ε} converging in measure to u. We still control $\mathcal{B}^{\varepsilon}$, the union of cubes of the second type, as in (11); on any $\overline{Q}_{\varepsilon} = \varepsilon y^{\varepsilon} + i + \varepsilon [0, 1)^d$ of the first type, we have that

$$\int_{\overline{Q}_{\varepsilon}} |e(u^{\varepsilon})|^p dx \leq C \varepsilon^{d-1} \sum_{\xi \in V} |\xi \cdot (u(\varepsilon y^{\varepsilon} + i + \varepsilon \xi) - u(\varepsilon y^{\varepsilon} + i))|^p$$

by Proposition 4 for p > 1 and a change of variables. Summing up we obtain the bound (7). \Box

Proof of Corollary 4. For any u_k let $(u_k^{\varepsilon})_{\varepsilon>0}$ be the approximating family provided by Proposition 2. Then for any $k \in \mathbb{N} \setminus \{0\}$ we can find $\varepsilon_k > 0$ such that, for $\tilde{u}_k := u_k^{\varepsilon_k}$ it holds that $\tilde{u}_k \in SBV^{\infty}(\Omega; \mathbb{R}^d)$ with $J_{\tilde{u}_k} \subset \partial^* \mathcal{B}^{\varepsilon_k}$ and $\mathcal{B}^{\varepsilon_k}$ union of finitely many cubes of sidelength ε_k included in Ω , such that

$$\|\arctan(\tilde{u}_k - u_k)\|_{L^1} \leq \frac{1}{k},$$

$$\int_{\Omega_{\varepsilon_k}} |e(\tilde{u}_k)|^p dx + \mathcal{H}^{d-1}(\partial^* \mathcal{B}^{\varepsilon_k} \cap \Omega_{\varepsilon_k}) \leq C \Lambda_{u_k}^{p, V_k}.$$
(15)

Then $(\tilde{u}_k)_k$ satisfies the assumptions of [18, Theorem 1.1] on any $\Omega_{\varepsilon_{\overline{k}}}$ for $k \geq \overline{k}$, and therefore by that compactness result and by a diagonal argument on Ω_{ε_k} , up to a (not relabelled) subsequence, there are a Caccioppoli partition $\mathcal{P} = (P_n)_n$ of Ω , a sequence of piecewise rigid motions $(a_k)_k$ satisfying (4), and $u \in GSBD^p(\Omega)$ such that

$$\tilde{u}_k - a_k \to u$$
 a.e. in Ω . (16a)

By the first in (15) we deduce (9a). Eventually, we get (9b) and (9c) by [21, Theorem 11.3] (recall [25, Lemma 2.1]) applied to $(u_k - a_k)_k$.

APPENDIX A. DISCRETE ENERGY ESTIMATE

In what follows, $(e_i)_{i=1}^d$ denotes the canonical basis of \mathbb{R}^d .

Lemma 3. Consider the unit cube $Q = [0,1]^d \subset \mathbb{R}^d$. Let $v \in (\mathbb{R}^d)^{\{0,1\}^d}$ be given at all vertices of Q such that $v_i(x+e_i) = v_i(x)$ for any $x \in \{0,1\}^d$ with $x_i = 0$ and $v_i(x+e_i+e_j)+v_j(x+e_i+e_j) = v_i(x) + v_j(x)$ for any $x \in \{0,1\}^d$ with $x_i = x_j = 0$. For $x \in Q$, we also denote by v(x) the multilinear interpolation of the values v at the vertices (affine on each $[x, x+e_i]$ for any $x \in Q$ with $x_i = 0$). Then e(v) = 0 in Q (so that, in fact, v is affine with skew-symmetric gradient).

Proof. We first assume d = 2. Then,

 $v(x) = (1 - x_1)(1 - x_2)v(0, 0) + (1 - x_1)x_2v(0, 1) + x_1(1 - x_2)v(1, 0) + x_1x_2v(1, 1).$ One easily sees that $\partial_1 v_1 = \partial_2 v_2 = 0$ everywhere. One has

$$\partial_1 v_2(x) + \partial_2 v_1(x) = (1 - x_2)(v_2(1, 0) - v_2(0, 0)) + x_2(v_2(1, 1) - v_2(0, 1)) + (1 - x_1)(v_1(0, 1) - v_1(0, 0)) + x_1(v_1(1, 1) - v_1(1, 0))$$

Since $v_2(1,0) = v_2(1,1)$, $v_2(0,1) = v_2(0,0)$, etc, this is also:

/ 1

$$\begin{split} \partial_1 v_2(x) + \partial_2 v_1(x) &= (1 - x_2)(v_2(1, 1) - v_2(0, 0)) + x_2(v_2(1, 1) - v_2(0, 0)) \\ &+ (1 - x_1)(v_1(1, 1) - v_1(0, 0)) + x_1(v_1(1, 1) - v_1(0, 0)) \\ &= v_2(1, 1) - v_2(0, 0) + v_1(1, 1) - v_1(0, 0) = 0. \end{split}$$

Hence e(v) = 0.

Now, consider $d \ge 3$: then for any $\{i, j\} \subset \{1, \ldots, d\}$, the case d = 2 shows that $e(v)_{i,j}(x) = 0$ for any x with $(x_i, x_j) \in [0, 1]^2$ and $x_k \in \{0, 1\}, k \notin \{i, j\}$. Since at any other $x \in Q, e(v)_{i,j}(x)$ is a convex interpolation of those values, we find that $e(v)_{i,j} = 0$ everywhere. Similarly, $e(v)_{i,i}(x) = \partial_i v_i(x) = 0$ for any $x \in Q$ and $i \in \{1, \ldots, d\}$. Hence e(v) = 0.

Proposition 4. Consider the unit cube Q, $v \in (\mathbb{R}^d)^{\{0,1\}^d}$ given at all vertices of Q, and the *d*-linear interpolation of v inside Q. Then for every $p \in [1, +\infty)$ there is a constant C > 0 depending on d and p such that:

$$\int_{Q} |e(v)|^{p} dx \leq C \left(\sum_{i=1}^{a} \sum_{\substack{x \in \{0,1\}^{d} \\ x_{i}=0}} |v_{i}(x+e_{i}) - v_{i}(x)|^{p} + \sum_{\substack{i,j=1 \\ x_{i}=x_{j}=0}}^{d} \sum_{\substack{x \in \{0,1\}^{d} \\ x_{i}=x_{j}=0}} |v_{i}(x+e_{i}+e_{j}) + v_{j}(x+e_{i}+e_{j}) - v_{i}(x) - v_{j}(x)|^{p} \right) \quad (17)$$

Proof. If not, there is a sequence $(v_n)_{n\geq 1}$ which satisfies the reverse inequality, with C replaced with n, arbitrarily large. Renormalizing, we may assume that $||e(v_n)||_{L^p(Q)} = 1$ and thus that the right-hand side of (17), evaluated for v_n , goes to zero as $n \to \infty$. By Poincaré-Korn inequality, there is a sequence rigid motions (affine functions $a_n(x) = A_n x + b_n$ with skew-symmetric gradient A_n) such that:

$$\int_{Q} |v_n(x) - a_n(x)|^p dx \le c \int_{Q} |e(v_n)|^p dx = c$$

where c depends only on the dimension and on p. Letting $v'_n = v_n - a_n$, it is bounded in $L^p(Q)$. Since it is, in fact, finite-dimensional (it is a polynomial of degree at most d), it is relatively compact and, up to a subsequence, converges to a limit v' with $||e(v')||_{L^p(Q)} = 1$. Yet the righthand side of (17), for v', vanishes, so that v' satisfies the assumptions of Lemma 3 and we deduce e(v') = 0. This is a contradiction, hence the corollary is true. An alternative, and maybe simpler proof, consists in arguing that both left-hand and right-hand sides of (17) define *p*th-powers of norms on the finite dimensional space $(\mathbb{R}^d)^{\{0,1\}^d} / \sim$ where $u \sim v$ if and only if the *d*-linear interpolation of u - v is an infinitesimal rigid motion.

Appendix B. Convergence of Discretizations

Let $\mathcal{Q}(x) = \chi_{[-1/2,1/2)^d}(x)$, $\Delta(x) = \prod_{i=1}^d (1-|x_i|)^+$, for $x = (x_1, \ldots, x_d) \in \mathbb{R}^d$, and for every $k \in \mathbb{N}$ let $T_k : \mathbb{R}^m \to \mathbb{R}^m$ be defined by

$$T_k(t) = t\left(1 \wedge \frac{k}{|t|}\right).$$

We remark that for $u \in L^0(\mathbb{R}^d; \mathbb{R}^m)$ it holds that for any R > 0,

$$\lim_{\ell \to \infty} \left| \{ x \in B_R(0) : |u(x)| > \ell \} \right| = \left| \bigcap_{\ell > 0} \{ x \in B_R(0) : |u(x)| > \ell \} \right|$$
$$= \left| \{ x \in B_R(0) : |u(x)| = +\infty \} \right| = 0.$$
(18)

Moreover, for any $\varepsilon > 0$ and $y \in Q$ we denote by

$$\mathcal{D}_y^{\varepsilon} := \{ \varepsilon y + z \colon z \in \varepsilon \mathbb{Z}^d \},\$$

the discretization of \mathbb{R}^d formed by cubes (with the same orientation) of sidelength ε and anchor point y and by $u_y^{\varepsilon} \colon \mathbb{R}^d \to \mathbb{R}^m$ the corresponding discretized function defined by

$$u_y^{\varepsilon}(x) := \sum_{\overline{x} \in \mathcal{D}_y^{\varepsilon}} u(\overline{x}) \Delta(\frac{x - \overline{x}}{\varepsilon}) \quad \text{for all } x \in \mathbb{R}^d.$$
(19)

In the following notation for the norms, we denote by $L^1(B_k)$ the space $L^1(B_k(0); \mathbb{R}^m)$, with $B_k \equiv B_k(0) \subset \mathbb{R}^d$ the ball of center 0 and radius k.

Lemma 5. Let $u \in L^0(\mathbb{R}^d; \mathbb{R}^m)$ and assume that for all $\varepsilon > 0$ there exists $y_{\varepsilon} \in Q$ such that for any $k \in \mathbb{N} \setminus \{0\}$

$$\lim_{\varepsilon \to 0} \left\| \sum_{\overline{x} \in \mathcal{D}_{y_{\varepsilon}}^{\varepsilon}} T_k(u(\overline{x})) \mathcal{Q}(\frac{-\overline{x}}{\varepsilon}) - T_k(u) \right\|_{L^1(B_k)} + \left\| \sum_{\overline{x} \in \mathcal{D}_{y_{\varepsilon}}^{\varepsilon}} T_k(u(\overline{x})) \Delta(\frac{-\overline{x}}{\varepsilon}) - T_k(u) \right\|_{L^1(B_k)} = 0.$$
(20)

Then

$$\lim_{\varepsilon \to 0} \|T_k(u_{y_\varepsilon}^\varepsilon) - T_k(u)\|_{L^1(B_k)} = 0 \quad \text{for any } k \in \mathbb{N} \setminus \{0\}.$$
(21)

Proof. Let us fix $k \in \mathbb{N} \setminus \{0\}$. Observe that if $\ell > k$, since $T_k = T_k \circ T_\ell$ and T_k is 1-Lipschitz,

$$\left\| T_k \left(\sum_{\overline{x}} T_\ell(u(\overline{x})) \Delta(\frac{\cdot - \overline{x}}{\varepsilon}) \right) - T_k(u) \right\|_{L^1(B_k)} \le \left\| \sum_{\overline{x}} T_\ell(u(\overline{x})) \Delta(\frac{\cdot - \overline{x}}{\varepsilon}) - T_\ell(u) \right\|_{L^1(B_k)} \to 0$$
(22)

as $\varepsilon \to 0$, thanks to (20). This means we just need to bound

$$\left\|T_k(u_{y_{\varepsilon}}^{\varepsilon}) - T_k\left(\sum_{\overline{x}} T_\ell(u(\overline{x}))\Delta(\frac{-\overline{x}}{\varepsilon})\right)\right\|_{L^1(B_k)}.$$

Observe that this quantity is zero, except possibly in the cubes with center a point \overline{x} such that $|u(\overline{x})| > \ell$, in which case it is bounded (in norm) by 2k. Hence we get the bound

$$\left\| T_k(u_{y_{\varepsilon}}^{\varepsilon}) - T_k\left(\sum_{\overline{x}} T_\ell(u(\overline{x}))\Delta(\frac{-\overline{x}}{\varepsilon})\right) \right\|_{L^1(B_k)} \leq 2k(2\varepsilon)^d \#\{\overline{x} \in \mathcal{D}_{y_{\varepsilon}}^{\varepsilon} \cap B_{k+\varepsilon\sqrt{d}} \colon |u(\overline{x})| > \ell\}.$$
(23)

For any $x \in \overline{x} + [-\varepsilon/2, \varepsilon/2)^d$ with $|u(\overline{x})| > \ell$, either $|u(x)| > \ell/2$, or

$$|T_{\ell}(u(\overline{x})) - T_{\ell}(u(x))| \ge |T_{\ell}(u(\overline{x}))| - |T_{\ell}(u(x))| = \ell - |u(x)| \ge \frac{\ell}{2}.$$

Hence (for ℓ large enough),

$$\begin{split} \varepsilon^{d} \# \{ \overline{x} \in \mathcal{D}_{y_{\varepsilon}}^{\varepsilon} \cap B_{k+\varepsilon\sqrt{d}} \colon |u(\overline{x})| > \ell \} \\ &\leq \left| \left\{ x \in B_{k+\varepsilon\sqrt{d}} \colon |u(x)| > \frac{\ell}{2} \right\} \right| + \left| \left\{ x \in B_{\ell} \colon \left| \sum_{\overline{x}} T_{\ell}(u(\overline{x}))\mathcal{Q}(\frac{x-\overline{x}}{\varepsilon}) - T_{\ell}(u(x)) \right| \ge \frac{\ell}{2} \right\} \right| \\ &\leq \left| \left\{ x \in B_{k+\varepsilon\sqrt{d}} \colon |u(x)| > \frac{\ell}{2} \right\} \right| + \frac{2}{\ell} \int_{B_{\ell}} \left| \sum_{\overline{x}} T_{\ell}(u(\overline{x}))\mathcal{Q}(\frac{x-\overline{x}}{\varepsilon}) - T_{\ell}(u(x)) \right| dx. \end{split}$$

We deduce from (20), (22), and (23) that for any $\ell > k$,

$$\limsup_{\varepsilon \to 0} \left\| T_k(u_{y_\varepsilon}^\varepsilon) - T_k(u) \right\|_{L^1(B_k)} \le 2^{d+1} k \left| \left\{ x \in B_{k+\varepsilon\sqrt{d}} : |u(x)| > \frac{\ell}{2} \right\} \right|.$$

Sending $\ell \to \infty$, by (18) we obtain (21).

Recalling the notation (19), the following lemma holds.

Lemma 6. Let $u \in L^0(\mathbb{R}^d; \mathbb{R}^m)$ and $(y_{\varepsilon})_{\varepsilon>0}$ be such that (21) holds. Then

$$u_{y_{\varepsilon}}^{\varepsilon} \to u \quad locally \ in \ L^0(\mathbb{R}^d; \mathbb{R}^m).$$
 (24)

0

Proof. Fix $\eta > 0$, R > 0. For any x and any $k \in \mathbb{N} \setminus \{0\}, k > 2\eta$, it holds that if $|u_{y_{\varepsilon}}^{\varepsilon}(x) - u(x)| > \eta$ then either |u(x)| > k/2, or $|T_k(u_{y_{\varepsilon}}^{\varepsilon}(x)) - T_k(u(x))| > \eta$. Hence, for $k \ge R$,

$$\left| \left\{ x \in B_R(0) : |u_{y_{\varepsilon}}^{\varepsilon}(x) - u(x)| > \eta \right\} \right| \le \left| \left\{ x \in B_R(0) : |u(x)| > \frac{k}{2} \right\} \right| + \frac{1}{\eta} \|T_k(u_{y_{\varepsilon}}^{\varepsilon}) - T_k(u)\|_{L^1(B_k)}.$$

Hence, for any fixed k,

$$\limsup_{\varepsilon \to 0} \left| \left\{ x \in B_R(0) : |u_{y_\varepsilon}^\varepsilon(x) - u(x)| > \eta \right\} \right| \le \left| \left\{ x \in B_R(0) : |u(x)| > \frac{k}{2} \right\} \right|$$

The conclusion follows letting $k \to \infty$.

Proposition 7. For any $u \in L^0(\mathbb{R}^d; \mathbb{R}^m)$ and for all $\varepsilon > 0$ there exists a measurable set $Q_{\varepsilon} \subset Q = [0,1)^d$ such that both $\lim_{\varepsilon \to 0} |Q_{\varepsilon}| = 1$ and $u_{y_{\varepsilon}}^{\varepsilon} \to u$ locally in $L^0(\mathbb{R}^d; \mathbb{R}^m)$ if $y_{\varepsilon} \in Q_{\varepsilon}$ for every $\varepsilon > 0$.

Proof. By Lemma 5 and Lemma 6, it is enough to find a set Q_{ε} such that (20) holds for any choice of $y_{\varepsilon} \in Q_{\varepsilon}$. For $y \in Q$, we let:

$$\Phi_{\varepsilon}(y) = \sum_{k \ge 1} 2^{-k} \int_{B_k} \left| \sum_{\overline{x} \in \mathcal{D}_y^{\varepsilon}} T_k(u(\overline{x})) \mathcal{Q}(\frac{-\overline{x}}{\varepsilon}) - T_k(u) \right| + \left| \sum_{\overline{x} \in \mathcal{D}_y^{\varepsilon}} T_k(u(\overline{x})) \Delta(\frac{-\overline{x}}{\varepsilon}) - T_k(u) \right| dx.$$

For any bounded function with compact support $v \in L^{\infty}(\mathbb{R}^d)$, and a continuous function ψ : $\mathbb{R}^d \to [0,1]$ with compact support such that $\sum_{z \in \mathbb{Z}^d} \psi(x-z) = 1$ for any $x \in \mathbb{R}^d$, one has

$$\begin{split} \int_{Q} \int \Big| \sum_{\overline{x} \in \mathcal{D}_{y}^{\varepsilon}} v(\overline{x}) \psi(\frac{x-\overline{x}}{\varepsilon}) - v(x) \Big| dx \, dy &= \int \Big| \int_{Q} \sum_{\overline{x} \in \mathcal{D}_{y}^{\varepsilon}} (v(\overline{x}) - v(x)) \psi(\frac{x-\overline{x}}{\varepsilon}) dy \Big| dx \\ y' &= \varepsilon y \\ &\leq \int \int_{\varepsilon Q} \sum_{z \in \varepsilon \mathbb{Z}^{d}} |v(y'+z) - v(x)| \psi(\frac{x-z-y'}{\varepsilon}) \varepsilon^{-d} dy' \, dx = \int \int \varepsilon^{-d} \psi(\frac{x-y'}{\varepsilon}) |v(y') - v(x)| dy' dx \\ &= \frac{x-y' - \varepsilon \xi}{\varepsilon} \int \psi(\xi) \|v(\cdot - \varepsilon \xi) - v\|_{L^{1}} d\xi \end{split}$$

where we have set $\psi_{\varepsilon} := \psi(\cdot/\varepsilon)\varepsilon^{-d}$.

Hence, for $v = T_k(u)\chi_{B_k}$ and $\psi = \mathcal{Q}, \Delta$,

$$\int_{Q} \Phi_{\varepsilon}(y) dy \leq \sum_{k \geq 1} 2^{-k} \int (\mathcal{Q}(\xi) + \Delta(\xi)) \|T_k(u)(\cdot - \varepsilon \xi) - T_k(u)\|_{L^1(B_k)}$$

Since the terms in the sum above are uniformly bounded (by $4k2^{-k}(1+|B_k|)$ for ε small enough) and go to zero, we deduce that

$$\lim_{\varepsilon \to 0} \int_Q \Phi_\varepsilon(y) dy = 0$$

We can then choose $Q_{\varepsilon} := \left\{ y \in Q : \Phi_{\varepsilon}(y) \leq \sqrt{\int_{Q} \Phi_{\varepsilon}} \right\}$, which is such that $|Q \setminus Q_{\varepsilon}| \leq \sqrt{\int_{Q} \Phi_{\varepsilon}}$. If one chooses $y_{\varepsilon} \in Q_{\varepsilon}$, then for each k, both L^{1} norms in (20) are bounded by $2^{k} \sqrt{\int_{Q} \Phi_{\varepsilon}}$, which goes to zero. Then we conclude by Lemma 6.

Remark 2. Proposition 7 implies that $u_{y_{\varepsilon}}^{\varepsilon}$ converge to u a.e. in \mathbb{R}^d , up to a subsequence $\varepsilon_n \downarrow 0$. We remark that such convergence may be recovered directly, without using Lemma 6, by modifying the argument in the proof of Proposition 7 as follows. With the notation in the proof of Proposition 7, first let ε_n such that $\sqrt{\int_Q \Phi_{\varepsilon_n}} \leq 2^{-n}$. Then, letting $Q_n := \bigcap_{h\geq n} Q_{\varepsilon_h}$, we find that $|Q \setminus Q_n| \leq 2^{-n+1}$ and for any $y \in Q_n$ and $h \geq n$, $\Phi_{\varepsilon_h}(y) \leq 2^{-h}$. Therefore, for $\tilde{Q} := \bigcup_{n\geq 1} Q_n$, we find that $|\tilde{Q}| = 1$ and for any $y \in \tilde{Q}$, $\Phi_{\varepsilon_h}(y) \leq 2^{-h}$ for h large enough (depending on y).

Next, once the main result is proven, we can choose a (further) subsequence $\varepsilon_k \downarrow 0$ such that for each k,

$$||T_k(u_{y_{\varepsilon_k}}^{\varepsilon_k}) - T_k(u)||_{L^1(B_k)} \le 2^{-k}.$$

Then, $\sum_{k} |(T_k(u_{y_{\varepsilon_k}}^{\varepsilon_k}(x)) - T_k(u(x)))\chi_{B_k}(x)| < +\infty$ for a.e. $x \in \mathbb{R}^d$, which shows that

$$u_{u_{\alpha}}^{\varepsilon_k} \to u$$
 a.e. in \mathbb{R}^d as $k \to +\infty$.

As a consequence, if one considers instead a *subsequence* in Proposition 7, one can replace Q_{ε} by the set \tilde{Q} of full measure, independent on ε , and assume that the convergence is almost everywhere.

Acknowledgements. V.C. acknowledges the support of the MUR - PRIN project 2022J4FYNJ CUP B53D23009320006 "Variational methods for stationary and evolution problems with singularities and interfaces", PNRR Italia Domani, funded by the European Union via the program NextGenerationEU. Most of this work was done while A.C. was visiting the Dipartimento di Matematica, Università di Roma-La Sapienza, thanks to the "Dipartimento di Eccellenza Project CUP B83C23001390001".

References

- S. ALMI, E. DAVOLI, A. KUBIN, AND E. TASSO, On De Giorgi's conjecture of nonlocal approximations for free-discontinuity problems: The symmetric gradient case, Preprint 2024, https://arxiv.org/abs/2410.23908.
- S. ALMI AND E. TASSO, A general criterion for jump set slicing and applications, Preprint 2022, https://arxiv.org/abs/2212.09822.
- [3] S. ALMI AND E. TASSO, A new proof of compactness in G(S)BD, Advances in Calculus of Variations, (2022).
- [4] L. AMBROSIO, Existence theory for a new class of variational problems, Arch. Rational Mech. Anal., 111 (1990), pp. 291–322.
- [5] L. AMBROSIO, On the lower semicontinuity of quasiconvex integrals in SBV(Ω, \mathbf{R}^k), Nonlinear Anal., 23 (1994), pp. 405–425.

- [6] L. AMBROSIO, A. COSCIA, AND G. DAL MASO, Fine properties of functions with bounded deformation, Arch. Rational Mech. Anal., 139 (1997), pp. 201–238.
- [7] L. AMBROSIO, N. FUSCO, AND D. PALLARA, Functions of bounded variation and free discontinuity problems, Oxford Mathematical Monographs, The Clarendon Press, Oxford University Press, New York, 2000.
- [8] J.-F. BABADJIAN, Traces of functions of bounded deformation, Indiana Univ. Math. J., 64 (2015), pp. 1271– 1290.
- J.-F. BABADJIAN, F. IURLANO, AND A. LEMENANT, Partial regularity for the crack set minimizing the twodimensional griffith energy, J. Eur. Math. Soc. (JEMS), 24 (2022), pp. 2443–2492.
- [10] F. CAGNETTI, A. CHAMBOLLE, AND L. SCARDIA, Korn and Poincaré-Korn inequalities for functions with a small jump set, Math. Ann., 383 (2022), pp. 1179–1216.
- [11] A. CHAMBOLLE, An approximation result for special functions with bounded deformation, J. Math. Pures Appl. (9), 83 (2004), pp. 929–954.
- [12] —, Addendum to: "An approximation result for special functions with bounded deformation" [J. Math. Pures Appl. (9) 83 (2004), no. 7, 929–954; mr2074682], J. Math. Pures Appl. (9), 84 (2005), pp. 137–145.
- [13] A. CHAMBOLLE, S. CONTI, AND G. FRANCFORT, Korn-Poincaré inequalities for functions with a small jump set, Indiana Univ. Math. J., 65 (2016), pp. 1373–1399.
- [14] A. CHAMBOLLE, S. CONTI, AND F. IURLANO, Approximation of functions with small jump sets and existence of strong minimizers of Griffith's energy, J. Math. Pures Appl. (9), 128 (2019), pp. 119–139.
- [15] A. CHAMBOLLE AND V. CRISMALE, A density result in GSBD^p with applications to the approximation of brittle fracture energies, Arch. Rational Mech. Anal., 232 (2019), pp. 1329–1378.
- [16] —, Existence of strong solutions to the Dirichlet problem for the Griffith energy, Calc. Var. Partial Differential Equations, 58 (2019), pp. 27, No. 136.
- [17] ——, Compactness and lower semicontinuity in GSBD, J. Eur. Math. Soc. (JEMS), 23 (2021), pp. 701–719.
- [18] ——, Equilibrium configurations for nonhomogeneous linearly elastic materials with surface discontinuities, Ann. SNS Pisa Cl. Sci., 24 (2023), pp. 1575–1610.
- [19] —, A general compactness theorem in G(S)BD, Indiana Univ. Math. J., Preprint 2022, in press (2025).
- [20] S. CONTI, M. FOCARDI, AND F. IURLANO, Existence of strong minimizers for the Griffith static fracture model in dimension two, Ann. Inst. H. Poincaré C Anal. Non Linéaire, 36 (2019), pp. 455–474.
- [21] G. DAL MASO, Generalised functions of bounded deformation, J. Eur. Math. Soc. (JEMS), 15 (2013), pp. 1943–1997.
- [22] E. DE GIORGI AND L. AMBROSIO, New functionals in the calculus of variations, Atti Accad. Naz. Lincei Rend. Cl. Sci. Fis. Mat. Nat. (8), 82 (1988), pp. 199–210 (1989).
- [23] G. A. FRANCFORT AND J.-J. MARIGO, Revisiting brittle fracture as an energy minimization problem, J. Mech. Phys. Solids, 46 (1998), pp. 1319–1342.
- [24] M. FRIEDRICH, A piecewise Korn inequality in SBD and applications to embedding and density results, SIAM J. Math. Anal., 50 (2018), pp. 3842–3918.
- [25] —, A compactness result in GSBV^p and applications to Γ-convergence for free discontinuity problems, Calc. Var. Partial Differential Equations, 58 (2019), pp. 31, No. 86.
- [26] M. FRIEDRICH, C. LABOURIE, AND K. STINSON, On regularity for griffith almost-minimizers in the plane, Preprint 2023, https://arxiv.org/abs/2310.07670.
- [27] M. FRIEDRICH, M. PERUGINI, AND F. SOLOBRINO, Lower semicontinuity for functionals defined on piecewise rigid functions and on GSBD, J. Funct. Anal., 280 (2021), pp. 45, No. 108929.
- [28] M. FRIEDRICH AND F. SOLOMBRINO, Quasistatic crack growth in 2d-linearized elasticity, Ann. Inst. H. Poincaré Anal. Non Linéaire, 35 (2018), pp. 27–64.
- [29] F. IURLANO, A density result for GSBD and its application to the approximation of brittle fracture energies, Calc. Var. Partial Differential Equations, 51 (2014), pp. 315–342.
- [30] R. TEMAM, Mathematical problems in plasticity, Gauthier-Villars, Paris, 1985. Translation of Problèmes mathématiques en plasticité. Gauthier-Villars, Paris, 1983.
- [31] R. TEMAM AND G. STRANG, Duality and relaxation in the variational problem of plasticity, J. Mécanique, 19 (1980), pp. 493–527.

CEREMADE, CNRS, UNIVERSITÉ PARIS-DAUPHINE, UNIV. PSL, FRANCE AND MOKAPLAN, INRIA PARIS. *Email address*, Antonin Chambolle: antonin.chambolle@ceremade.dauphine.fr

DIPARTIMENTO DI MATEMATICA GUIDO CASTELNUOVO, PIAZZALE ALDO MORO 5, 00185 ROMA, ITALY. *Email address*, Vito Crismale: vito.crismale@mat.uniroma1.it

12