

Introduction of an equation of state for nirogen based described by the free Helmholz energy in the model of a spherical cavitation bubble (type Rayleigh-Plesset)

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Introduction

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Mécanique

In the last instants of the contraction of an inertial cavitation bubble, the gases pressure value of hundred of MPa and up to several GPa. Van der Waals equation of state which have replace the ideal gas law long time ago starts to diverge from exepimental value of the gaz behaviour typically as pressure overdrives 1 GPa [1,2]; for example, for fixed temperature and gas concentration, the real pressure is not so important than the calculated prevure with the Van der Waals pressure. Others equations of state or correlations exist to describe gases at such hight pressure among which : the 'cubic' said equation of state, extension of Van der Waals formulation, as the Peng-Richardson (PR) state or correlation founded on the free energy of Helmholz [3,4].

According to Storey and Szeri [5] who modelised the chemical activity of a bubble, using equation of state for very high pressure do not change the nature of the physics and chemical phenoma in the bubble, but modify they quantification.

The target that drives this work is to precise temperature and pressure close to the bubble in an undercooled water to forecast the number of crystal ice a bubble can produce (a continuation of [6]); it can be approximated that undercooled water is free of vapor and chosen to work with only nitrogen in the bubble, as it is experimentally rather easily practicable.

The dynamic equation

bubble dynamics (Rayleight-Plesset, Gilmore, Keller-Miksis), it is Whatever the expressed with the second derivative of the radius R which contains both the gaz pressure $\mathbf{P_g}$ and its derivative $\dot{\mathbf{P}_g}$ (below, Keller-Miksis' one from [7])

$$\ddot{\mathbf{R}} = \frac{1}{\mathbf{R}\left(1 - \frac{\dot{\mathbf{R}}}{c_{l}}\right)} \left[-\frac{3}{2} \dot{\mathbf{R}}^{2} \left(1 - \frac{\dot{\mathbf{R}}}{c_{l}}\right) + \frac{1}{\rho} \left(1 + \frac{\dot{\mathbf{R}}}{3 c_{l}}\right) \left(\mathbf{P_{g}} - \mathbf{P_{a}} - \mathbf{P_{o}}\right) + \frac{\mathbf{R}}{\rho_{l} c_{l}} \dot{\mathbf{P}}_{g} - 4\mu \frac{\dot{\mathbf{R}}}{\mathbf{R}} - \frac{2\sigma}{\rho \mathbf{R}} \right]$$

With R, \dot{R} and \ddot{R} the bubble radius and its two first derivatives according to time; c_1 the wave celerity, ρ the liquid volumic mass and μ the dynamic viscosty.

'Cubic' equations of state

This expression contains both the gas pressure Pg and its derivatives against time, which depends of the chosen equation of state and the gas-liquid thermal transfert :

Van der Waals gas and a polytropic gas transformation (coef. κ):

 $P_{g} = \left(P_{o} + \frac{2\sigma}{R_{o}}\right) \left(\frac{R^{3} - 8,54^{3}}{R^{3} - 8,54^{3}}\right)^{\kappa}$ T : the gaz temperature, 8,54 coefficient for N₂ T is directly deducted from VdW law

Van der Waals gas and explicit model of gas-liquid heat transfert :

$$P_g(t) = \frac{N_{tot} k T}{\frac{4\pi}{3} [R(t)^3 - Ro/8, 54]}$$
 N_{tot}: total number of gas molecules
R_o the initial bubble radius

For explicit gas-liquid transfert, a 'cubic equation of state' as Peng-Robinson (PR) is possible :

$$P_g = \frac{RT}{V_m - b} - \frac{a}{V_m(V_m + b) + b(V_m - b)}$$

 V_m : molar volume ; 'a' and 'b', coefficients.

When using an explicit heat transfert, T is determined through internal energy U and results from statistic thermodynamics.

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Formulation for the Helmholz free energy correlation

Free Helmholtz energy based corretation have been introduced in bubble model by way of tabulated value []; for nitrogen, the correlation used was limited to 2 GPa and 1000 K [] but other correlation is given up to 25 GPa and 5000K [2]. The correlation is build from the ideal α^{o} and the residual α^{r} free energy, given by the following kind of expressions; coefficients are given in references.

$$\alpha^{o}(\delta,\tau) = \ln(\delta) + \alpha_{1}\ln(\tau) + \alpha_{2} + \alpha_{3}\ln(\tau) + \alpha_{4}\ln(\tau^{-1}) + \alpha_{5}\ln(\tau^{-2}) + \alpha_{6}\ln(\tau^{-3}) + \alpha_{7}\ln(1 - e^{-\alpha_{8}\tau})$$

$$\alpha^{r}(\delta,\tau) = \sum_{1}^{6} N_{k}\delta^{i_{k}}\tau^{j_{k}} + \sum_{7}^{32} N_{k}\delta^{i_{k}}\tau^{j_{k}}e^{-\delta^{l_{k}}} + \sum_{33}^{36} N_{k}\delta^{i_{k}}\tau^{j_{k}}e^{(-\varphi_{k}(\delta-1)^{2}-\beta_{k}(\tau-\gamma_{k})^{2})}$$

With : $\delta = \frac{\rho}{\rho_c}$, $\tau = \frac{T_c}{T}$ and 'c' the index for critical values.

Thermodynamics magnitudes are defined from α^{o} , α^{r} and derivatives, for example P, U, Z (and also Cv, Cp, S and sound velocity not displayed here).

$$P_{g} = \rho R_{g} T \left[1 + \delta \left(\frac{\partial \alpha^{r}}{\partial \delta} \right)_{\tau} \right] \quad Z = \frac{P_{g}}{\rho R_{g} T} = 1 + \delta \left(\frac{\partial \alpha^{r}}{\partial \delta} \right) \quad U = R_{g} T \tau \left[\left(\frac{\partial \alpha^{o}}{\partial \tau} \right)_{\delta} + \left(\frac{\partial \alpha^{r}}{\partial \tau} \right)_{\delta} \right]$$

Introduction in a cavitation bubble

The formulation of free Helmholtz energy equation of state does not permit to express directly $\mathbf{P}_{\mathbf{g}}$ nor $\dot{\mathbf{P}}_{\mathbf{g}}$.

The system of differential equations of the bubble model (dynamic, heat transfert, molecule transfert) enable to get a each time of their integration R, U and the total number N_{tot} of gas molecules because their derivatives can be analytically expressed : from R and N_{tot} them, it is possible to calculate the volumic mass (ρ) or molecule conentration : with given U and ρ , it is possible thanks to the equation of state, to research the corresponding temperature by resolving a first order equation in T then to calculate the pressure.

It still misses $\dot{\mathbf{p}}_{\mathbf{g}}$ to calculate $\ddot{\mathbf{R}}$ first it is calculated R(t+dtt), N(t+dtt) and U(t+dtt) with dtt small compared to the dt use to integrate the system of differential equation of the bubble model : then the same procedure used at t gives : T(t+dtt) and Pg(t+dtt); the same calculation backward to t-dtt gives T(t-dtt) and Pg(t-dtt); so are estimated :

$$\dot{T} = \frac{T(t - dtt) - T(t + dtt)}{2 dtt} \qquad \dot{P}_g = \frac{P_g(t - dtt) - P_g(t + dtt)}{2 dtt}$$

is obtained with no element external to the equation of state

References :

[1]: S. Bidi et al: "Numerical Study of real gaz effects during bubble collapse using a desequilibrium multiphose model "; Ultrasonics Sonochemistry, 90, 2022. [2] D. Qin et al. : "Numerical investigation on acousite cavitation characteristics on an air-vopor bubble : effect of equation of state for interior gases", Ultrasonics sonochemistry, 97 (2023). [3] R.T. Jacobsen et al. : "A thermodynamic property formulation for nitr reference quality equation of state for nitrogen"; Int. J. of thermophysics, 7(3), 1986. C [4] R. Span et al. : "A reference quality equation of state for nitrogen"; Int. J. of thermophysics, 19 (4), 1998. [5] B.D. Storey, A.J. Szeri : "Water vapor, sonoluminescence and sonochemistry", Proc. Royal Soc. London, Serie A456, p1685-1709, 1999-2000. [6] Cogne *et al.*: Theorical model of ice nucleation induced by acoustic nucleation, Part I: pressure and temperature profile around a single bubble. Ultrasonic Sonochemistry, 29, 2016

[7] Toegel *et al.* : "Does water vapor prevent upscaling sonoluminescence", Phys. Rev. Let., 85 (15), 2000,

Calculation for a 1 mm initial radius bubble submitted to a constant pressure, as does Bidi et al. [1]The model of these authors is more more sophiticated than ours; we compare the radius and tenyperature evolution for our model with Helmholtz free energy and a VdW law; the minimum radii and maximum averaged température given by Bidi et al., with o for the HelmHoltz model and + for Peng-Robinson model.

The scheme of calculation we use gives similar result than the one employed by Bidi et al. [1], for the radius, rather similar for temperature for PR and Helmholtz. Considering pressure (no result VdW Helmholtz displayed in [1]), Helmholtz and VdWs are close in this case, but for smaller bubble [2], first result have shown larger





Results and comparison with [1]



differences.c



Conclusion This calculation scheme have to be introduced in the ice nucleation model to evaluate its quantitative influence.

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