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# Single crystal growth by the Bridgman-Stockbarger technique of CaF<sub>2</sub>-rich - TmF<sub>3</sub> solid solutions: Evidence of mixed valence Tm<sup>2+</sup> and Tm<sup>3+</sup> cations

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## ABSTRACT

Crystal growth and numerical simulation of CaF<sub>2</sub>-rich-based solid solution with TmF<sub>3</sub> have been performed by the Bridgman-Stockbarger technique with initial Tm content 0.1 mol%, 1 mol% and 5 mol%. Centimeter-sized single crystals were obtained. Growth in the vicinity of (110) and (210) crystallographic directions have been identified as the natural crystal habit for the CaF<sub>2</sub>-rich- TmF<sub>3</sub> solid solution with an axial thermal gradient about G=12 K.cm<sup>-1</sup> and a pulling rate of v=4 mm.h<sup>-1</sup>. Optical absorption spectroscopy evidenced the presence of both Tm<sup>2+</sup> and Tm<sup>3+</sup> cations within the matrix. Their respective effective partition coefficients  $k_{\text{eff}}(\text{Tm}^{2+})$  and  $k_{\text{eff}}(\text{Tm}^{3+})$ , close to unity, have been calculated. For low TmF<sub>3</sub> initial content, with 0.1 mol.% and 1 mol.%,  $k_{\text{eff}}(\text{Tm}^{2+}) > 1$  and  $k_{\text{eff}}(\text{Tm}^{3+}) > 1$  whereas  $k_{\text{eff}}(\text{Tm}^{2+}) < 1$  and  $k_{\text{eff}}(\text{Tm}^{3+}) < 1$  for the attempt with 5 mol% TmF<sub>3</sub> initial content. For this latter, a global effective partition coefficient, including both Tm<sup>3+</sup> and Tm<sup>2+</sup> contents, is logically deduced to be  $k_{\text{eff}}([\text{Tm}]) < 1$ . Chemical analysis carried out by ICP-AES and LIBS confirmed this trend with an effective partition coefficient of the global [Tm] content in the crystal that can be lower than 1. The global estimated partition coefficient of [Tm] in CaF<sub>2</sub> matrix is thus ranging in between 0.906 and below 1 for 5 mol.% TmF<sub>3</sub> molar fraction. The expected congruent melting point in the CaF<sub>2</sub>-rich-region of the CaF<sub>2</sub>-TmF<sub>3</sub> phase diagram is evidenced and is higher than 1 mol.% and below 5 mol.% TmF<sub>3</sub> molar fraction. Finally, carbon environment and TmF<sub>3</sub> synproportionation reaction are pointed out as the main likely causes of Tm<sup>3+</sup> cation reduction into Tm<sup>2+</sup> during the growth process.

## 1. Introduction

Fluoride-based materials are well known [1] due to their intrinsic properties, such as a high optical transparency (VUV–VIS–IR range), low phonon energy, small refractive index, remarkable physicochemical properties (non-hygroscopy in most cases) and ease of machining and polishing. In addition, they have great ability to host various doping-cations with 3d or 4 f electronic configuration. Owing to their specific features, fluoride-based single crystals exhibit enhanced properties, which are of particular interest in numerous physical phenomena ranging from luminescence, scintillation and up-conversion to

solid-state laser effect. A wide range of fluoride matrix have been investigated through time such as CaF<sub>2</sub>, BaF<sub>2</sub> or CeF<sub>3</sub> [1–4] and many others such as: TbF<sub>3</sub>, LiYF<sub>4</sub> (LYF), LiTbF<sub>4</sub>, KYF<sub>4</sub>, LiLuF<sub>4</sub> (LuLiF), KY<sub>3</sub>F<sub>10</sub> (KYF), KTb<sub>3</sub>F<sub>10</sub> (KTF), BaY<sub>2</sub>F<sub>8</sub>, LiCaAlF<sub>6</sub> (LiCAF), LiSrAlF<sub>6</sub> (LISAF), KMgF<sub>3</sub> (KMF), BaLiF<sub>3</sub>, KAlF<sub>4</sub>, Na<sub>5</sub>Al<sub>3</sub>F<sub>14</sub> [1,5–13]. When doped with suitable cations, fluoride compositions open up opportunities in many application fields such as window materials, tunable lasers, defence, telecommunication, optical lenses or remote sensing. Crystal growth of bulk fluorides are mainly carried out by the Czochralski technique [1,12, 14], Bridgman technique in a slightly lower extent [9,15], and seldom by Top Seeded Solution Growth [6,16] or micro-pulling down technique

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[17]. In particular, Bridgman technique features particular advantages by comparison with the other techniques. Indeed, it enables the setting of an accurate thermal profile in both liquid and solid phases for a steady controlled growth and homogeneous cooling-down of as-grown boules. It also enables the control of the atmosphere in order to perform single crystal growth with stoichiometric compositions, low amount of ionic vacancies, low oxygen pollution and no hydroxyl substitution.

Because rare-Earth ions are attractive for achieving optical applications in the visible, near and mid-infrared spectral range, in the following, Tm doping cation is considered in order to perform high quality Tm<sup>3+</sup>-doped CaF<sub>2</sub> single crystal growth. Tm<sup>3+</sup> is a very interesting doping cation [12,18–20] because of its electronic configuration, which allows several optical emission at 1.5 μm, 1.9 μm and 2.3 μm. Hence, Tm<sup>3+</sup>-doped fluoride single crystals are promising for the realization of eye-safe laser, remote sensing of atmospheric species such as CH<sub>4</sub>, infrared counter-measurements or for medicine purposes. However, as previously reported [21], CaF<sub>2</sub> doped with TmF<sub>3</sub> exhibit both Tm<sup>3+</sup> and Tm<sup>2+</sup> cations. In the present work, we report the crystal growth, the numerical modelling and the characterization on the chemical behaviour of TmF<sub>3</sub> in single crystalline CaF<sub>2</sub> solid solutions obtained by the Bridgman-Stockbarger method. That aims at clarifying the origin of the Tm<sup>3+</sup> reduction into Tm<sup>2+</sup> through their respective chemical distribution and segregation. Three initial TmF<sub>3</sub> contents were assessed: 0.1 mol.%, 1 mol.% and 5 mol.% in order to screen the CaF<sub>2</sub>-rich region of CaF<sub>2</sub>-TmF<sub>3</sub> phase diagram where the most interesting compositions for the above-mentioned optical applications are found. Structural study and crystal habit characterized by Laue experiments enable to identify preferential growth directions in the as-grown crystals. Optical absorption characterization of the crystals is displayed leading to the evidence of mixed valence of Tm<sup>3+</sup> and Tm<sup>2+</sup> cations within the CaF<sub>2</sub> matrix. LIBS coupled to ICP-AES analysis and optical measurements enable to characterize, respectively, the chemical homogeneity and optical absorption of the boules at different stages of the growth. Effective partition coefficients of thulium are calculated by these two methods and are discussed with respect to the CaF<sub>2</sub>-TmF<sub>3</sub> phase diagram.

## 2. Experimental procedure

### 2.1. Crystal growth

CaF<sub>2</sub>-rich solid-solution single crystals doped with initial 0.1 mol.%, 1 mol.% and 5 mol.% TmF<sub>3</sub> contents were grown by using the vertical Bridgman-Stockbarger method [3]. Crushed and sieved pure CaF<sub>2</sub> optical UV-VIS crystals from Crystran Ltd. UK and commercial powder from Alfa Aesar of TmF<sub>3</sub> with 99,999 % (5 N) purity were chosen as raw materials. Raw materials were mixed and inserted thoroughly in a 10-cm long and 1-cm wide graphite crucible without prior synthesis. Graphite is used because carbon cannot react with CaF<sub>2</sub> and TmF<sub>3</sub> and its wettability is virtually zero with respect to liquid fluorides [22], making easy to extract boules without them sticking to the inner surface of the crucible. The crucible, after closing with a graphite screw-on cap at its top, is set in the furnace. The solid solution is synthesized and melted when the CaF<sub>2</sub>-TmF<sub>3</sub> mixture is heated up in 5 hours, under a  $5 \times 10^{-5}$  torr vacuum atmosphere, above the melting point of CaF<sub>2</sub> with a maximum temperature of 1785 K. The melt is soaked during 3 hours then pulled down with  $v=4 \text{ mm.h}^{-1}$  translation speed. Prior to the growth, a longitudinal temperature distribution along the axis of the furnace was measured with a S-type thermocouple and a thermal gradient about  $G=12 \text{ K.cm}^{-1}$  was achieved in the adiabatic zone where the solid-liquid interface stands. Once the crucible was translated over all its length within the furnace, the translation is stopped and the furnace is cooled down within 24 h. The as-grown crystals are then cleaved [23] with a hammer and a metallic blade into 10–30 ellipsoidal-like-shaped slices of ~1.5–3 mm thickness along the well-known (111) directions [23,24]. As-cleaved slices are hereinafter

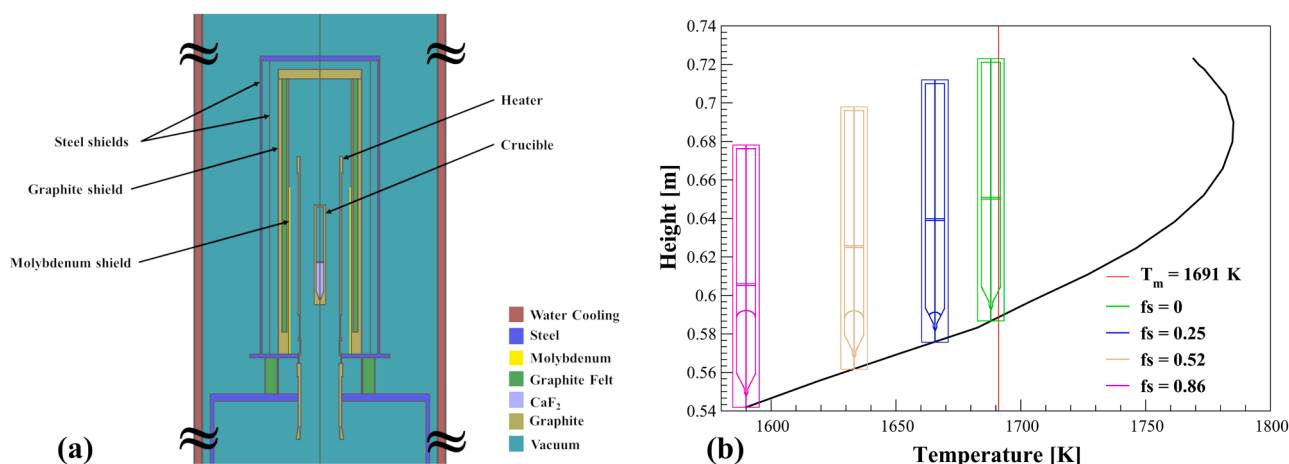
used for optical, chemical and structure characterization.

### 2.2. Numerical modelling

CrysMAS code is a finite-volume numerical tool developed by the Crystal Growth Laboratory within the Fraunhofer Institute of Integrated Systems and Device Technology (IISB) at Erlangen in Germany [25–27] and was used for the numerical modelling of the crystal growth experiment. It allows simulating heat transfer, fluid flow and solid-liquid interface shape during crystallization with respect to latent heat, radiative transport, and conduction and convection heat phenomena. The features of the furnace are described in our previous work [3] and were schematized via a 2D-Computer-Aided-Design-like tool integrated to the software. The global geometry of the furnace is presented in Fig. 1a. The temperature gradient in the ampoule is generated by a shaped graphite heater. The thermokinetic parameters of the materials constituting the Bridgman device are listed in Table 1.

Inverse 2D-axisymmetric global simulations were performed in order to determine the thermal gradient in the furnace. For this, a control point placed at the symmetry axis at a height of 18.5 cm from the heater bottom, where the temperature of 1688 K was measured in experiment. This, in turn was adapted for local simulation of the crucible. The heat transfer mechanisms that are considered during the simulations are conduction, convection and radiative heat transfer. For this study, the main goal of the numerical simulation was to evaluate the level of convection and consequently the melt mixing in different stages of the crystal growth. Therefore, in order to save computational time, a local numerical model was considered consisting of an axisymmetric graphite crucible with an inner diameter of 10 mm and an outer graphite wall with a thickness of 2 mm placed in the temperature gradient obtained from global simulation (Fig. 1b). At the bottom of the crucible, the geometry contains a tip with a maximum diameter of 2 mm and a height of 4 mm. From this point, the crystal thickness gradually increases until it reaches a final crystal radius of 5 mm. A 1 mm thick graphite floater is placed 58 mm from the tip of the crucible similarly with the experimental set-up. The floater is designed to prevent excessive volatilization of the liquid solution in the upper part of the crucible where the temperature is the highest. The floater is always in contact with the liquid phase at any stage of the growth and it is almost liquid-tight near the inner wall of the graphite crucible. Moreover, previous simulations have shown that the addition of the floater has no significant influence at the beginning of the growth. Its influence is efficient when the solidified fraction (*fs*) approaches *fs*=1, so that the floater acts as a thermal shield and prevents, at the end of the growth, a too high cooling rate at the free surface level. The solid-liquid interface is assumed to be present on the melting point isotherm, where latent heat is released due to the solidification process. CrysMAS uses the well-established Gebhart view factors to describe the ratio of radiation absorbed by any other surface versus the total emitted radiation from given surface and the n-Band model for treatment of the spectral behaviours. which was proven to produce accurate results with relatively low computational cost [36, 37]. Note that the material is totally transparent only in the defined spectral bands and is totally opaque in the rest of the spectrum.

Four stages of crystal growth were considered, as follows: C0 – *fs* = 0; C1 – *fs* = 0.25; C2 – *fs* = 0.52; C3 – *fs* = 0.86 (where *fs* indicates the solidified fraction) and the temperature distribution is studied. The temperature profile obtained from global simulation was used as boundary conditions for the lateral side of the crucible. For the top and bottom boundaries of the crucible, we used fixed temperature values. At the melt-crucible interface a zero-velocity boundary condition was imposed. The physical properties for the materials used in the numerical modelling are listed in Table 1 and are assumed to be constant with temperature (if not specified otherwise). Grid refinement and time steps used in these computations were tested to ensure accurate results.



**Fig. 1.** (a) Furnace and ampoule geometry. (b) Temperature gradient and four crucible positions for different solidification fraction. The black line represents the longitudinal distribution of the temperature with respect to the position (height) of the crucible in the furnace.

**Table 1**

Thermo-physical parameters used for modelling with CrysMAS code.

Description	Value	Reference
<i>CaF<sub>2</sub> crystal</i>		
Density ( $\text{kg}\cdot\text{m}^{-3}$ )	3200	[28]
Thermal Conductivity ( $\text{W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$ )	9.71	[29]
Specific heat capacity ( $\text{J}\cdot\text{kg}^{-1}\cdot\text{K}^{-1}$ )	$766.3 + 0.39\cdot T + 2.52\times 10^{-5}\cdot T^2$	[30]
Visibility bands ( $\mu\text{m}$ )	0.15 – 9	[31]
Melting temperature (K)	1691	[31]
Refraction index	1.43	[28]
Thermal expansion ( $\text{K}^{-1}$ )	$1.87\times 10^{-5}$	[28]
<i>CaF<sub>2</sub> melt</i>		
Density ( $\text{kg}\cdot\text{m}^{-3}$ )	$3767 - 0.694\cdot T$	[32]
Thermal Conductivity ( $\text{W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$ )	0.6	[28]
Specific heat capacity ( $\text{J}\cdot\text{kg}^{-1}\cdot\text{K}^{-1}$ )	1280	[33]
Latent heat ( $\text{J}\cdot\text{kg}^{-1}$ )	320000	[28]
Visibility bands ( $\mu\text{m}$ )	0.15–9	[31]
Refraction index	1.43	[28]
Viscosity ( $\text{kg}\cdot\text{m}^{-1}\cdot\text{s}^{-1}$ )	$0.000371205\cdot\exp(2470\cdot T^{-1})$	[34]
Volumetric expansion ( $\text{K}^{-1}$ )	0.0002682	[35]
<i>Graphite</i>		
Thermal Conductivity ( $\text{W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$ )	$115.74 - 0.0982\cdot T + 3.61\times 10^{-5}\cdot T^2 - 4.29\times 10^{-9}\cdot T^3$	[28]
Emissivity	0.8	[28]
<i>Molybdenum</i>		
Thermal Conductivity ( $\text{W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$ )	$163.7 - 0.062\cdot T + 9.72\times 10^{-6}\cdot T^2$	[28]
Emissivity	0.4	[28]

### 2.3. Chemical and physical analysis

Powders resulting from crushed single crystals and a 63- $\mu\text{m}$  sieving were analysed by X-Ray Diffraction (XRD). XRD patterns were collected on a Bragg-Brentano  $\theta$ - $\theta$  geometry diffractometer Bruker D8 Advance working with Cu anticathode (40 kV, 40 mA), 0.3° slits and Ni filter. The data were collected over an angular range of  $2\theta = 5^\circ$ -90° with a 0.006° step. XRD patterns were studied using the known space groups Fm-3m of CaF<sub>2</sub> [38]. Lattice parameters were refined by Rietveld method using the FullProf [39] software. Laue back-scattering patterns were recorded using a Photonic Science CCD camera by averaging of 6 images snapped for 1 minute. A stationary crystal is irradiated with a 0.7-mm collimated polychromatic X-ray beam supplied by a tungsten anticathode working at 35 kV and 40 mA (Seifert ID3003 generator). The polished samples were held on a goniometric head and set perpendicular to the X-ray beam using an autocollimator. OrientExpress software [40] was used for

pattern indexing, with an orientation accuracy is of 0.5°.

Thulium elemental content measurements were performed by using Inductively Coupled Plasma - Atomic Emission Spectroscopy (ICP-AES) using the Vista MPX (Varian). For the analysis, samples containing a mixture of calcium fluoride (CaF<sub>2</sub>) and thulium (Tm) were quantified at three designated concentration levels. Samples were weighed accordingly: ranging from 50.6 mg to 98.6 mg, to align with the predetermined concentrations. These samples underwent a controlled digestion process in a 1:1 hydrochloric acid solution within 50 mL digestion tubes, which involved sequential heating at 383 K for one hour, 393 K for 1.5 hours, and 413 K for four hours. After cooling, the digested samples were transferred to volumetric flasks and diluted to a final volume of 60 microliters with filtered water on the following day. The standard addition method was applied to counteract matrix effects, utilizing a thulium standard solution (1000 mg.l<sup>-1</sup> in 2–3 % nitric acid, Supelco reference 1.70361.0100, Merck/Aldrich) that was appropriately diluted. The calibration involved systematically increasing the volume of the standard solution added to a constant volume of the sample mixture, with adjustments in the volume of water to maintain consistent total volumes across all test solutions. This meticulous approach enabled precise calibration essential for accurate quantification.

The Varian Vista MPX ICP-AES was employed to measure thulium concentrations, focusing on its spectral lines at 313.125 nm and 336.261 nm. The instrument settings were optimized to achieve accurate detection and quantification limits specific to thulium's spectral characteristics. The results of the analysis are listed in Table 2. Additionally, a recovery test on the standard solution showed no significant loss, with a recovery rate of 104.54 %. This confirms the high accuracy and reliability of the method, demonstrating effective compensation for potential analytical deviations and ensuring the integrity of the thulium quantification in these complex sample matrices.

LIBS-based imaging setup [41–44] was used to image the distribution of Ca, F and Tm in the different sections of as-grown crystal growth slices. A Nd:YAG laser (Centurion GRM, Lumibird) was used with 7 ns

**Table 2**

Tm content measured by ICP-AES for the different growth attempts.

Attempt	Wgh.% Tm	Wgh.% Tm accuracy	Mol.% Tm	Mol.% Tm accuracy
5 mol.% Tm-doped CaF <sub>2</sub>	9.74 %	1.09 %	4.92 %	0.60 %
1 mol.% Tm-doped CaF <sub>2</sub>	1.97 %	0.07 %	0.93 %	0.03 %
0.1 mol.% Tm-doped CaF <sub>2</sub>	0.243 %	0.012 %	0.1126 %	0.0060 %



pulse duration and 100 Hz repetition rate, emitting at the fundamental wavelength (1064 nm). A x15 magnification objective microscope (LMM-15X-P01, Thorlabs) was used to focus the laser beam onto the sample surface). An energy of 1 mJ was set for all the acquisition, with an argon flow directed to the plasma region of  $0.8 \text{ l}\cdot\text{min}^{-1}$  to prevent surface contamination by ablated material redeposition of the previous laser shots and to improve the sensitivity due to better plasma confinement. The experiment was done at room temperature and ambient pressure. During the analysis, the sample surface was scanned, through single laser pulses, in a pixel-by-pixel manner inducing the breakdown and sparks of the material. The light radiation emitted by the plasma was collected by an optical system and analysed using a spectrometer (Shamrock 500, Andor technology) equipped with an intensified charge-coupled device (ICCD) camera (istar, Andor technology). The spectrometer was configured in the spectral range from 345 nm to 385 nm ( $1200 \text{ l}\cdot\text{mm}^{-1}$  grating) in order to detect both Ca and Tm lines (see Fig. 1 in supplementary material). In this configuration, the spectral resolution was 0.08 nm. The ICCD camera was synchronized to the Q-switch of the laser. A gain of 500 was applied, with a delay and a gate width of  $0.5 \mu\text{s}$  and  $5 \mu\text{s}$ , respectively. For the shown LIBS imaging results, the step size between each laser shot was fixed to  $40 \mu\text{m}$ .

The intensity calibration was performed with respect to the 3 above mentioned reference samples analysed with ICP-AES. For these 3 samples, LIBS imaging was obviously carried out before sample dissolution required for ICP-AES analysis. For these 3 analyses, a mask was defined by removing the influence of the edges of the crystal sections which could be affected by de sample preparation. These masks were produced from the calcium images. An average LIBS spectrum was then determined for each section (i.e. average of all spectra contained in the mask). Assuming that the concentration of Tm was homogeneous in volume (i.e. correspondence between LIBS surface analysis and ICP-AES volume analysis), we built a quantitative model by plotting the LIBS intensity of Tm as a function of the concentration obtained with ICP-AES. A special attention was paid to the choice of the Tm line. Indeed, given the concentration range of interest (up to 10 % molar), the intense Tm lines observed on the spectrum suffered from self-absorption, resulting in a non-linear quantitative model. We, therefore, turned our attention to the

line at 383.8 nm, which is weak but unaffected by these saturation phenomena (see resulting model on Fig. 2 in supplementary material). A very good linearity is obtained, with a coefficient of determination of 0.9993. With the evaluation of the confidence hyperbola, we were able to estimate the associated LIBS uncertainties at around 3.7 % for a 5 % Tm molar concentration, 10 % for 1 mol.%Tm, and 100 % for a 0.1 mol.%Tm.

Finally, to investigate the distribution of both  $\text{Tm}^{3+}$  and  $\text{Tm}^{2+}$  cations along  $\text{CaF}_2$  crystal lengths and to calculate the effective segregation coefficient, the optical absorption method was used as previously described [45,46]. The optical absorption spectrum of each slice was measured, at room temperature, in UV-VIS spectral range using a Shimadzu 1650PC spectrophotometer.

### 3. Results and discussion

#### 3.1. Numerical modelling

The main parameters that influence the solid-liquid (S-L) interface shape in a crystal grown by the Bridgman technique are the growth rate and the temperature gradient. The radiative heat transfer will also affect the shape of the S-L interface, therefore we have set both the melt and crystal  $\text{CaF}_2$  as semi-transparent materials (transparent between  $0.15 - 9 \mu\text{m}$ ). We investigated the convection patterns of the melt and interface shape under different solidified fractions (fs) and temperature gradients during  $\text{CaF}_2$  growth using pseudostationary numerical simulations. In Fig. 2, we can see, respectively, axial and radial velocity distribution.

In general, the Bridgman configuration is hydrodynamically stable due to the temperature gradient (hot melt is in the upper part and cold melt is in the lower part of the crucible). When the growth starts, the deflection of the S-L interface generates radial temperature gradients, which in turn generates a stronger melt convection. From Fig. 2a, we can observe that there is a difference of two orders of magnitude between the maximum velocity in the cone and the maximum velocity in the cylindrical part of the crucible. The small velocities in the tip of the crucible ( $\sim 10^{-7} \text{ m}\cdot\text{s}^{-1}$ ) lead to a poor mixing of the melt in the early stages of the crystallization process. This may be caused by the vortex that forms near

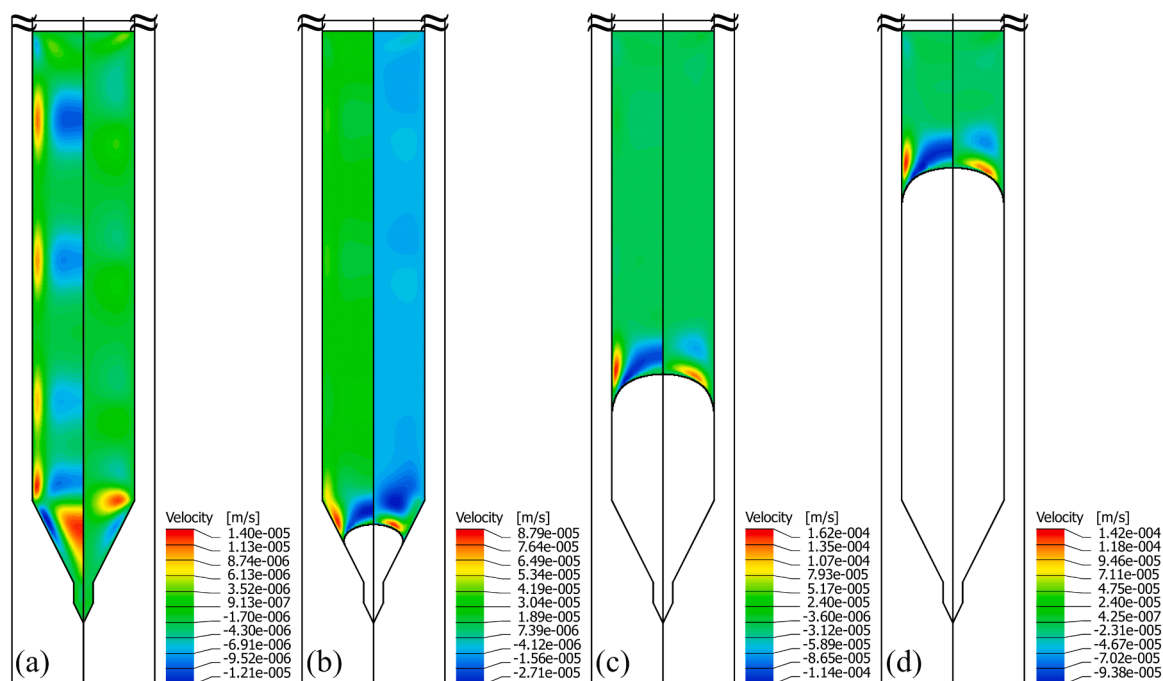


Fig. 2. Velocity field distribution along the axial (left side) and radial (right side) directions for different solidified fractions: (a) C0 – fs = 0, (b) C1 – fs = 0.25, (c) C2 – fs = 0.52 and (d) C3 – fs = 0.86.

the crucible shoulder, at the end of the conical part that does not reach the tip of the crucible and is unable to mix the melt in the tip with the rest of the melt. As for the crystallization front near the end of the conical part and further out (for  $f_s > 0.25$ , Fig. 2b, c and d), the velocities in the vortex generated near the S-L interface becomes one order of magnitude higher than in the start phase (Fig. 2a.). This vortex ensures a much better mixing in the melt in comparison with the start phase.

Interface deflection is defined as the perpendicular distance between the centre point of the S-L interface and the point where the melt, the crystal and the crucible come into contact. The evolution of the S-L interface can be seen in Fig. 3. The numerical results show that the interface presents a convex shape throughout the whole growth process. In the cylindrical region of the crucible, the convexity of the S-L interface (defined as the ratio between the maximum interface deflection and the crystal width) decreases from 0.73 (C2) to 0.68 (C3) due to the increase in crystal size that leads to a higher latent heat released during the solidification process.

### 3.2. Crystal growth

Centimetre-sized as-grown boules were obtained (Fig. 4) with lengths up to 90 mm and 10 mm in diameter. A thin graphite layer is observed at the outer surface of the as-grown boules. As shown by the numerical modelling (Fig. 3), a convex solid-liquid interface is observed during the whole growth process. This result is consistent with what was observed with  $\text{BaF}_2$ , as an isotype of  $\text{CaF}_2$ , boules and growth simulations in the same conditions [47,48]. Since carbon cannot enter the  $\text{CaF}_2$  matrix, radial growth occurred from the centre of the liquid towards the graphite wall of the crucible enabling the graphite impurities contained in the liquid to be rejected towards the periphery of the growing boule, leaving the core of the boule free of inclusions. However, few macroscopic black carbon inclusions at the late end of the boules are detected. We assume that because of the poor intrinsic late-end-of-growth conditions of Bridgman technique, the carbon impurities contained in the final part of the solidifying liquid are inevitably trapped. After graphite removal from the outer surface, the boules display homogeneously clear regions in their entire volume (Fig. 5) without macroscopic inclusions visible with a binocular under magnification X64.

### 3.3. XRD and Laue characterization

The X-ray diffraction patterns (Fig. 6) collected on crushed single crystals for 0.1 mol.%, 1 mol.% and 5 mol.%  $\text{TmF}_3$ -doped  $\text{CaF}_2$  samples at several stages of the growth feature a typical crystallized fluorite structure with Fm-3m space group. No secondary phases were detected. Calculated lattice parameters are mentioned in Table 3 along with those reported in the literature for undoped  $\text{CaF}_2$  samples [49]. The refined

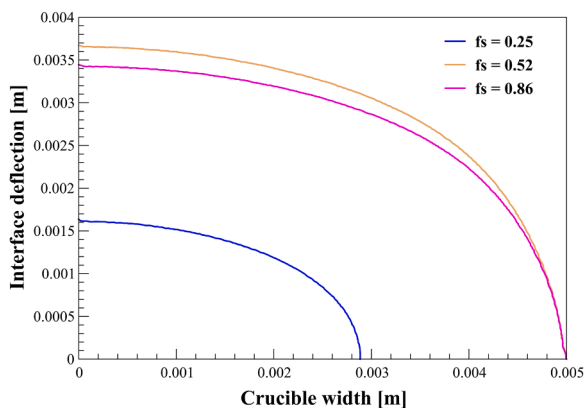


Fig. 3. Evolution of the S-L interface deflection during the growth process from numerical simulation.

parameters for our samples are in agreement with those previously reported for the fluorite structure [49], with slightly increased unit cell parameters/volume values due to the Tm doping into the  $\text{CaF}_2$  matrix. Moreover, the cell volume slightly increases from  $163.189(8) \text{ \AA}^3$  (0.1 mol.%  $\text{TmF}_3$ -doped  $\text{CaF}_2$  sample) to  $163.768(3) \text{ \AA}^3$  (5 mol.%  $\text{TmF}_3$ -doped  $\text{CaF}_2$  sample), pointing out the enrichment of Tm cations in Ca sites and fluorine anions in interstitial sites of the cubic  $\text{CaF}_2$  crystal structure. Lattice parameters calculated (Table 3) from Rietveld measurements were used for Laue experiments.

Natural crystallographic cleavage plane (111) was confirmed by Laue X-ray diffraction (Fig. 7a). Laue characterization of the boules has shown polycrystalline bulk in the 4-mm long seed location up to the beginning of the 5-mm long cone height. However, one single grain, overall the cross section of the crucible, is invariably detected for the 3 growth attempts from the end of the conical part. As no seed was used during the growth, numerous grains spontaneously nucleated in the seed location. These grains grew in the conical part until one remaining single grain was naturally selected. Growth attempts along this selected natural crystallographic direction enabled the growth of one single crystal from the end of the conical part and in the whole cylindrical volume of the crucible.

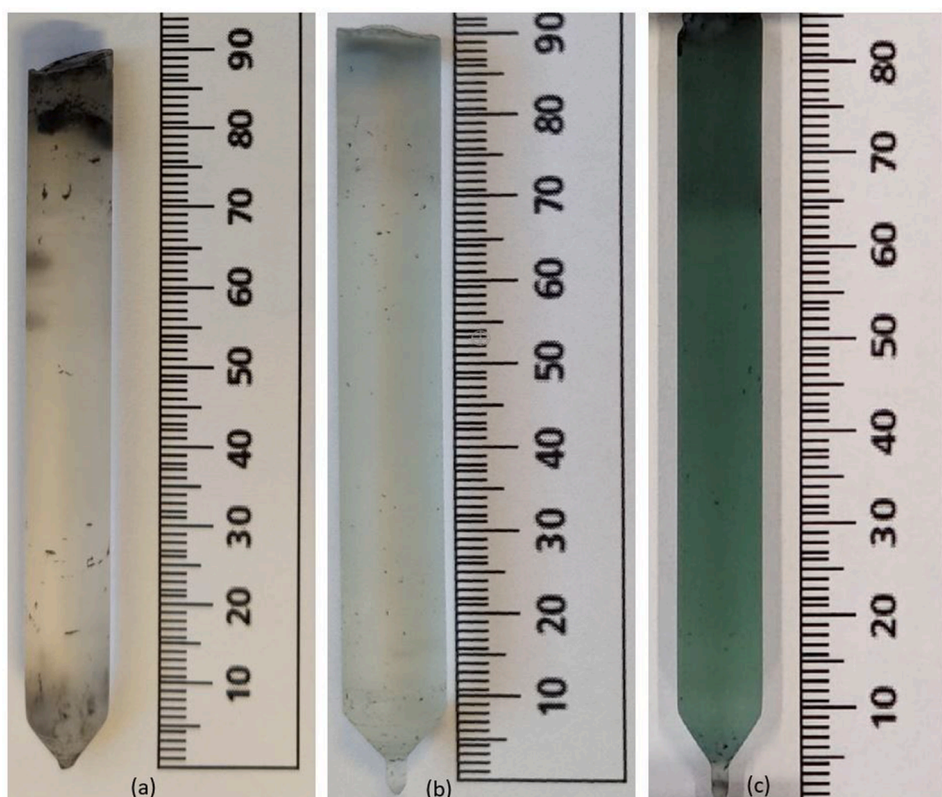
With respect to the thermokinetic parameters used in the Bridgman furnace of the present work (see paragraph 2), the natural growth directions of the as-grown crystals have been identified to be in the vicinity of the (110) crystallographic direction, where dense in-plane Ca-F chemical bonds are located, and (210) crystallographic direction. That is consistent with what was previously reported by Senguttuvan *et al.* [50] for  $\text{CaF}_2$  Bridgman growth by spontaneous nucleation with similar thermokinetic parameters, where a longitudinal gradient about  $G=10 \text{ K. cm}^{-1}$  and growth velocity about  $v=1 \text{ mm.h}^{-1}$  were used. (Fig. 7b and c). In addition, crystal habit, which refers to the crystal shape and natural crystallographic growth plans with respect to its structure and growth conditions, is strongly supported by the observed stability of the present growth attempts. Indeed, according to the Tiller criterion  $G/v$  [51] related to Mullins-Sekerka morphological instabilities [52,53], Fedorov *et al.* [54] calculated that Tiller criterion for the stable Bridgman growth of  $\text{CaF}_2$ , with a low amount of  $\text{SrF}_2$  impurities, must be  $G/v > 1.88 \times 10^4 \text{ K.s.cm}^{-2}$  in order to avoid morphological instabilities and constitutional supercooling. In the present work, as well as in the work of Senguttuvan *et al.* [50],  $G/v$  is calculated to be, respectively,  $G/v = 1.08 \times 10^5 \text{ K.s.cm}^{-2}$  and  $G/v = 3.60 \times 10^5 \text{ K.s.cm}^{-2}$ . They are greatly higher than the limit given by Fedorov *et al.* [54] enabling thus the habit of the crystal to develop naturally along these crystallographic directions with respect to the thermokinetic conditions employed. Finally, for each growth attempt, geometrical angles in between as-cleaved plan and growth direction were measured around  $40^\circ \pm 5^\circ$  and are consistent with theoretical calculated values of  $35.26^\circ$  and  $39.23^\circ$  in between (111) and, respectively, (110) and (210) directions for  $\text{CaF}_2$  cubic structure.

### 3.4. Effective partition coefficients calculated from chemical and optical measurements

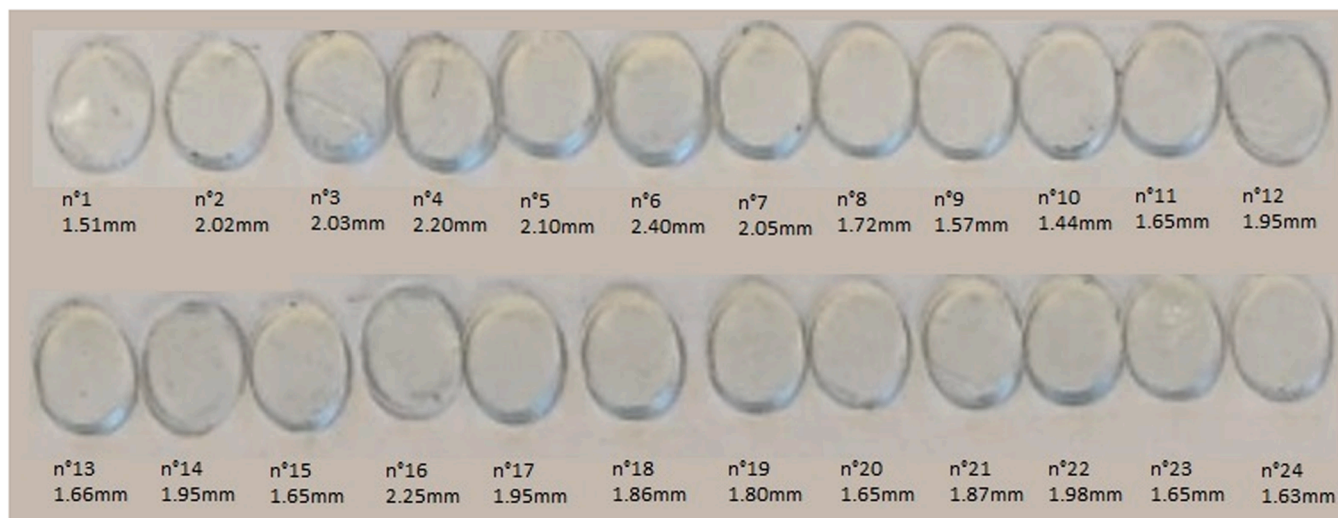
The effective partition coefficient  $k_{\text{eff}}$  is defined by  $k_{\text{eff}} = C_S / C_L^0$ , where  $C_L^0$  is the dopant concentration in the initial melt and  $C_S$  is the dopant concentration at the early beginning of the as-grown crystal. The value of  $C_S$  can be measured by various methods, from chemical analysis to optical absorption measurements [55–60].

#### 3.4.1. Chemical analysis

Several slices extracted from the seed location, the conical and the cylindrical parts were analysed by LIBS and ICP-AES. The average concentration of Tm molar content in each slice is calculated from every point of the LIBS mapping overall the whole sample surface. Note that the global content of [Tm] in the crystal measured by this method includes both  $\text{Tm}^{2+}$  and  $\text{Tm}^{3+}$  contents since these two cations cannot be distinguished with LIBS technique.



**Fig. 4.** (a) 0.1 mol.% Tm-doped  $\text{CaF}_2$  as-grown single crystal; (b) 1 mol.% Tm-doped  $\text{CaF}_2$  as-grown single crystal; and (c) 5 mol.% Tm-doped  $\text{CaF}_2$  as-grown single crystal.



**Fig. 5.** Example of cleaved defect-free slices, numbered from n°1 to n°24, obtained all along the growth axis for 1 mol.% Tm-doped  $\text{CaF}_2$  single crystal. The thickness of each slice is indicated below the number of the slice. No change in colour is observed in all the slices.

On the one hand, we consider (111)-orientated slices extracted from the cylindrical part. LIBS imaging revealed a highly homogeneous distribution of the dopant content on the whole surface of the 9 analysed slices (Fig. 8). In particular, no Tm radial segregation is detected. As shown in Fig. 3, a strong convex S-L interface occurs during the whole growth process, meaning that the crystallization occurs from the centre towards the periphery of the crucible and should lead to a radial chemical gradient of Tm content due to segregation phenomenon. However, as hereafter reported, partition coefficient of Tm is very close to unity. The observable segregation phenomenon of Tm in the radial

direction, with a maximum crystal radius of 0.5 cm, is substantially low compared to what can be slightly detected along the crystal axis, with a crystal length up to 9 cm, and has not been detected.

On the other hand, LIBS measurements performed in as-grown boules have shown a sudden drop of Tm content in the conical part compared to the cylindrical part where the distribution of [Tm] content evolves more steadily. The poor mixing within the liquid phase as pointed out by above-mentioned numerical simulation may be responsible for this sudden change.

It is worth mentioning that for the low initial [Tm] content growth



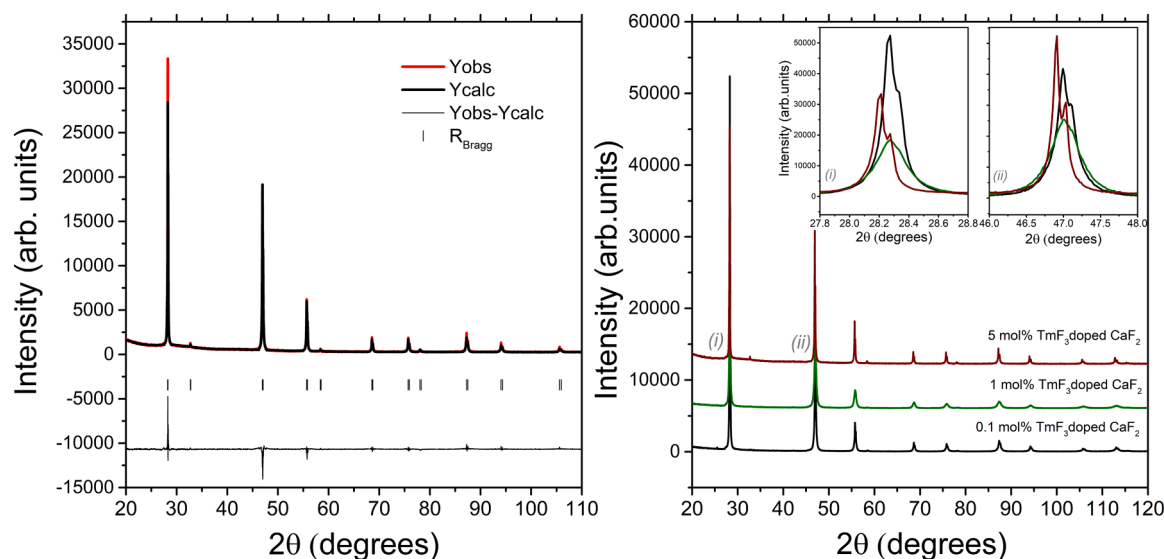


Fig. 6. (a) Calculated and measured X-Ray diffraction patterns for  $\text{CaF}_2$  cubic structure [49] (b) X-Ray diffraction patterns for the 0.1 mol.%, 1 mol.% and 5 mol.%  $\text{TmF}_3$ -doped  $\text{CaF}_2$  samples characteristic of the fluoride phase. (i) and (ii) Enlargements of the low-angle part of the X-ray diffraction patterns.

Table 3

Sample composition, unit cell parameters and corresponding volume for the 0.1 mol.%, 1 mol.% and 5 mol.%  $\text{TmF}_3$  doped  $\text{CaF}_2$  samples obtained from the Rietveld refinement for X-ray diffraction data collected on crushed single crystals.

Sample	Unit cell parameters (Å)	Volume (Å <sup>3</sup> )	Reference
Undoped $\text{CaF}_2$	5.463	163.065	[49]
0.1 mol.% $\text{TmF}_3$ -doped $\text{CaF}_2$	5.46467(15)	163.189(8)	This work
1 mol.% $\text{TmF}_3$ -doped $\text{CaF}_2$	5.46619(5)	163.325(3)	This work
5 mol.% $\text{TmF}_3$ -doped $\text{CaF}_2$	5.47112(5)	163.768(3)	This work

attempt, with  $C_L^0=0.1$  mol.% and  $C_L^0=1$  mol.%, no trend in the evolution of Tm segregation along the length of the boules can be deduced, so that only boule with  $C_L^0=5$  mol.% initial Tm content is considered. Hence, the molar error of [Tm] content in the crystal is taken as that estimated by ICP-AES (Table 2) with  $\pm 0.60$  mol.%. The error of Tm content in the liquid corresponds to the accuracy of weighing  $\text{TmF}_3$  during the crucible loading. It is about  $\pm 1 \times 10^{-3}$  g ( $\sim 4.4 \times 10^{-4}$  mol.%) and is neglected in the following calculations. LIBS averaged content measurements have shown a slight evolution of the [Tm] content all along the boule length with a varying [Tm] average content from  $5.13 \pm 0.60$  mol.% to  $4.57 \pm 0.60$  mol.% at the end of, respectively, the conical and cylindrical parts. Besides, in the seed part, [Tm] content is measured about  $C_S = 5.09 \pm 0.60$  mol.% so close to that at the end of the conical part. As shown by numerical simulation, the convection is poor in the seed location. At the liquid interface during the crystallization, Tm content supply is mainly caused by Tm diffusion and poorly by convection. The [Tm] content in the crystallized seed part may be slightly underestimated compared to what is measured in the conical part. Hence, the effective partition coefficient of [Tm] in  $\text{CaF}_2$  is calculated with the highest value of  $C_S = 5.13 \pm 0.60$  mol.% and gives  $k_{\text{eff}}([\text{Tm}]_{\text{LIBS}}) = 1.026 \pm 0.120$ , close to unity, according to what was predicted by Fedorov *et al.* [61].

### 3.4.2. Spectroscopic measurements

The absorption spectra of  $\text{TmF}_3$ -doped  $\text{CaF}_2$  crystals reveal the characteristic absorption bands of both  $\text{Tm}^{2+}$  and  $\text{Tm}^{3+}$  cations (Fig. 9). The peak at 260 nm corresponding to  ${}^3\text{H}_6 \rightarrow {}^3\text{P}_2$  transition, and the absorption band at 409 nm corresponding to the transition in between  $4f^{13}({}^2\text{F}_{7/2})$  ground state to the  $4f^{12}5d^1$  excited state, were chosen for the determination of the effective partition coefficient of  $\text{Tm}^{2+}$  and  $\text{Tm}^{3+}$

respectively, by the optical absorption method. In Fig. 9, the  ${}^3\text{H}_6 \rightarrow {}^3\text{P}_2$  transition at 260 nm for  $\text{Tm}^{3+}$  and the  $4f^{13}({}^2\text{F}_{7/2}) \rightarrow 4f^{12}5d^1$  transition at 409 nm for  $\text{Tm}^{2+}$  were selected based on their characteristic absorption features and suitability for quantitative analysis. As demonstrated recently [21], the 260 nm band of  $\text{Tm}^{3+}$  ions exhibits a linear trend of increasing the absorption coefficient with  $\text{TmF}_3$  concentration, supporting its effective use here. To mitigate potential nonlinearity effects, such as self-absorption at higher concentrations, the stronger  $\text{Tm}^{2+}$  absorption band at 409 nm was also employed for accurate detection.

This method is based on the following two laws:

(a) according to the Beer–Lambert law the absorption coefficient,  $\alpha$  is proportional to the sample concentration,  $C$ ,  $\alpha = aC$ , where  $a$  is the absorption coefficient for unit dopant concentration and unit light path length. The dopant concentration of a each slice can be estimated using the equation:

$$C = \frac{\alpha}{a} = \frac{23,0258}{ad} \quad (1)$$

where  $d$  is the sample thickness in mm.

(b) The dopant concentration along the growth axis at the distance  $z$  from the origin of the crystal can be obtained by using the equation [62]:

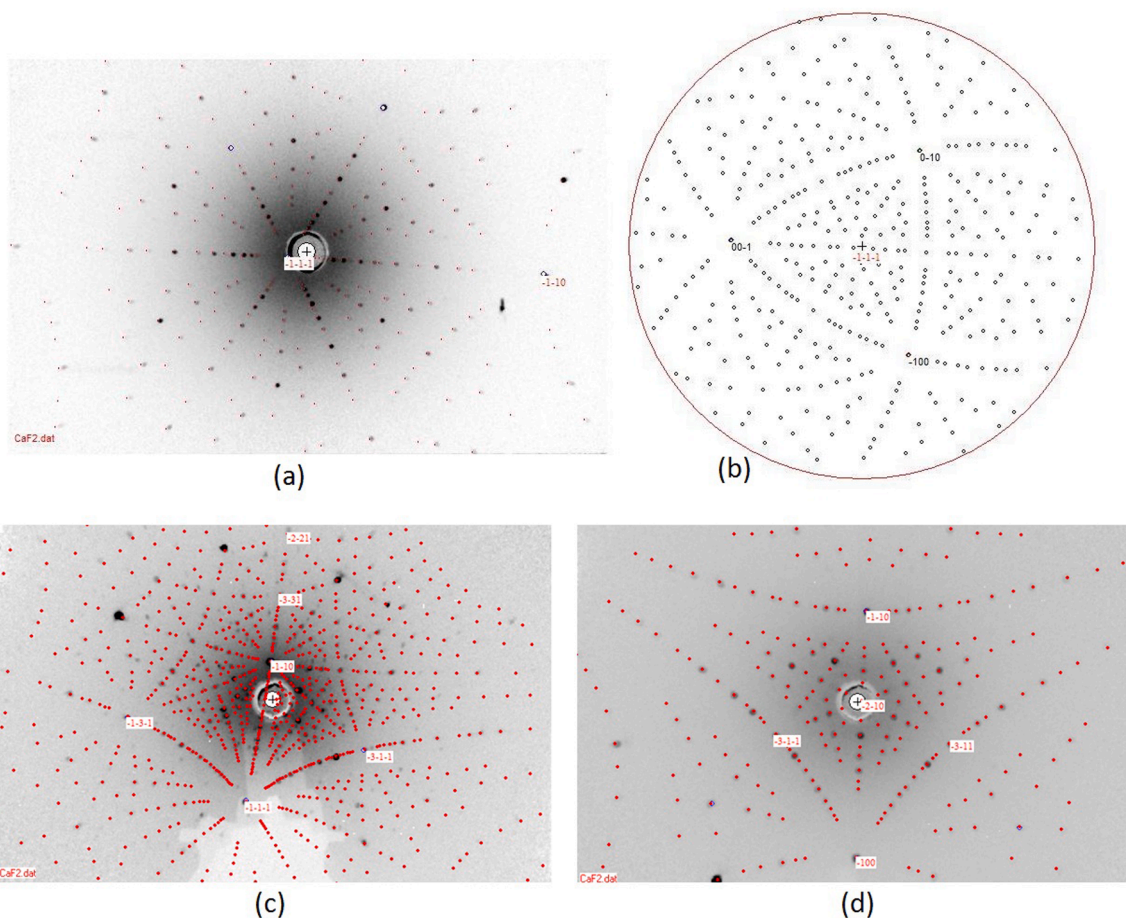
$$C_S(z) = C_L^0 k_{\text{eff}} [1 - g(z)]^{k_{\text{eff}} - 1} \quad (2)$$

where  $g(z)$  is the crystallized fraction of the melt.

Therefore, the dopant distribution along the crystal length can be estimated using the optical absorption method and the effective segregation coefficient can be calculated. The absorption coefficients  $\alpha(z)$  for  $\text{Tm}^{3+}$  and  $\text{Tm}^{2+}$  ions are measured from the optical absorption spectrum of every slice. Using the Eqs. (1) and (2), the following expression is obtained:

$$\log \alpha(z) = (k_{\text{eff}} - 1) \log \left( 1 - \frac{z}{L} \right) + \log (a k_{\text{eff}} C_L^0) \quad (3)$$

The effective partition coefficient was estimated from the slope  $m = k_{\text{eff}} - 1$ , of the fitting line of  $\log \alpha(z)$  vs.  $\log(1 - z/L)$ , where 'log' denotes the decimal logarithm. The calculated effective partition coefficients are shown in Fig. 10 and Table 4, together with the results provided by other authors [55–60]. As the concentration of the  $\text{TmF}_3$  in the initial melt increases, the effective partition coefficient,  $k_{\text{eff}}$  of both  $\text{Tm}^{3+}$  and  $\text{Tm}^{2+}$  ions decreases. This is in a good agreement with previous results reported for other rare earth-based trifluorides with close cationic radius



**Fig. 7.** Example of Laue patterns at the beginning of cylindrical part of as-grown boules. (a) Cleavage plan orientation in 5 mol.% Tm-doped  $\text{CaF}_2$  with (b) stereographic projection. Examples of (c) natural growth direction in 0.1 mol.% Tm-doped  $\text{CaF}_2$  and (d) natural growth direction 1 mol.% Tm-doped  $\text{CaF}_2$ . Red dots correspond to simulated Laue pattern superimposed to experimental one with black spots.

than those of  $\text{Tm}^{3+}$  or  $\text{Tm}^{2+}$  such as  $\text{Er}^{3+}$  cations [2] as well as  $\text{Yb}^{3+}$  and  $\text{Yb}^{2+}$  cations in fluorite crystals [63].

Considering the 1 mol.% and 5 mol.%  $\text{TmF}_3$ -doped  $\text{CaF}_2$  attempts, the global partition coefficients calculated from Optical Absorption (OA)  $k_{\text{eff}}([\text{Tm}]_{\text{OA}})$  (where  $[\text{Tm}]$  total content corresponds to  $\text{Tm}^{2+}$  and  $\text{Tm}^{3+}$  contents) evolve from a value higher to lower than unity. Whether both  $k_{\text{eff}}(\text{Tm}^{2+})_{\text{OA}}$  and  $k_{\text{eff}}(\text{Tm}^{3+})_{\text{OA}}$  are higher than one, then  $k_{\text{eff}}([\text{Tm}]_{\text{OA}})$  must be higher than one, and vice versa. These results obtained from optical measurements are in good agreement with LIBS results measured on 5 mol.%  $\text{TmF}_3$ -doped  $\text{CaF}_2$  crystal which allow a partition coefficient lower than unity with  $0.906 \leq k_{\text{eff}}([\text{Tm}]_{\text{LIBS}}) \leq 1.146$ . From both LIBS - ICP-AES and optical absorption measurements, for 5 mol.%  $\text{TmF}_3$ -doped  $\text{CaF}_2$ , partition coefficient of Tm is estimated to be  $0.906 \leq k_{\text{eff}}([\text{Tm}]_{\text{LIBS+OA}}) < 1$ .

### 3.4.3. Discussion

In analogy with Yb-doped fluorides, Tm ions may exhibit clustering behavior at higher concentrations, potentially impacting Rare Earth absorption intensities. Previous works on Yb-doped  $\text{CaF}_2$ , as reported by Nicoara *et al.* [46,64], showed that at low concentrations, the segregation coefficients for  $\text{Yb}^{3+}$  and  $\text{Yb}^{2+}$  ions were close to unity, resulting in a relatively homogeneous distribution. However, clustering increased with dopant concentration, affecting absorption characteristics. For the Tm-doped system studied here, we assume that the absorption coefficient  $\alpha$  remains proportional to the Tm concentration for the selected transition. This assumption is consistent with the linear trend observed for Tm absorption bands in similar low-concentration fluoride hosts [21], which suggests minimal influence of clustering effects under our

experimental conditions. Although clustering could introduce nonlinearity and measurement error at higher concentrations, our analysis is based on a concentration range where linearity is expected to hold, as observed in analogous systems [21,46].

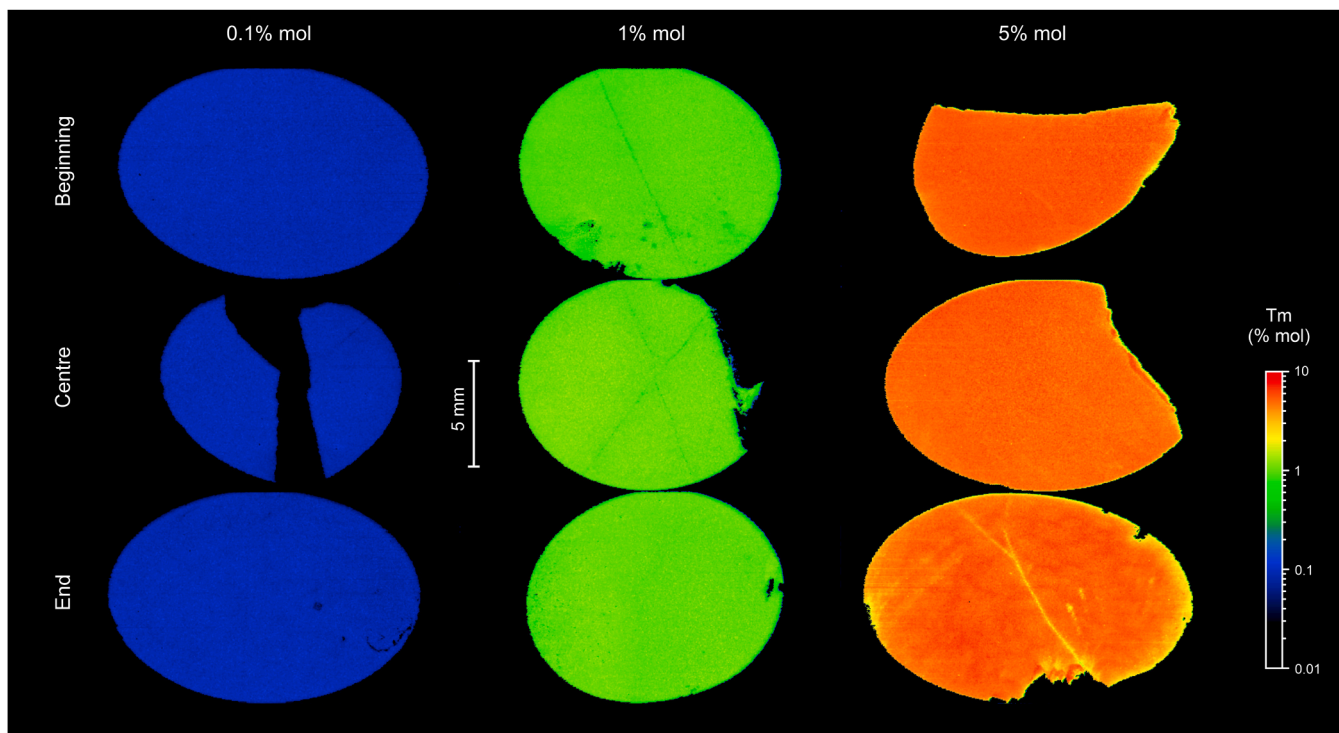
#### • Congruent melting point in $\text{CaF}_2$ - $\text{TmF}_3$ phase diagram

As reported by Sobolev *et al.* [65], the phase diagram  $\text{CaF}_2$ - $\text{TmF}_3$  does not exhibit a maximum of the liquidus curve corresponding to a congruent melting point in its  $\text{CaF}_2$ -rich region, although it should exist as previously mentioned [61]. LIBS as well as optical absorption spectroscopy measurements have shown a partition coefficient of Tm in  $\text{CaF}_2$  higher than unity for 1 mol.% Tm initial content and lower than unity for 5 mol.% Tm initial content. That highlights a change in the slope signs of both solidus and liquidus curves in the  $\text{CaF}_2$ - $\text{TmF}_3$  phase diagram. Indeed, when the slope signs of both liquidus and solidus curves are positive, during the crystallization, the Tm content is always higher in the crystal than that in the liquid leading, thus to  $k_{\text{eff}} > 1$ . On the contrary, when the slope signs of both liquidus and solidus curves are negative, the concentration of Tm in the liquid is higher than that in the crystal leading, thus, to  $k_{\text{eff}} < 1$ . We conclude that the expected congruent melting point previously predicted [61,65], ranges higher than 1 mol.% and below 5 mol.%  $\text{TmF}_3$  molar fraction.

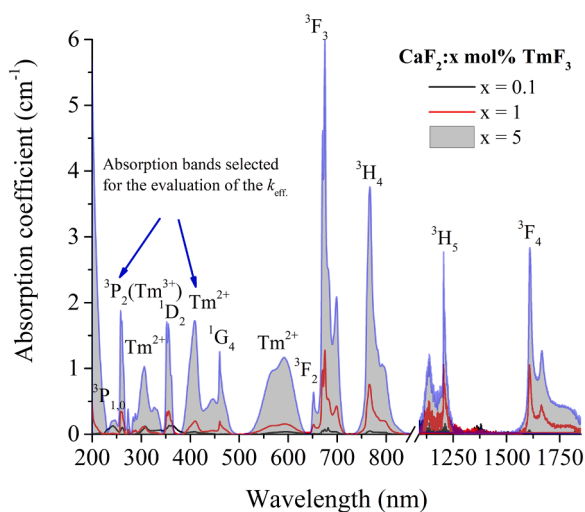
#### • Carbon environment influence on $\text{Tm}^{3+}$ reduction into $\text{Tm}^{2+}$

The presence of  $\text{Tm}^{2+}$  peaks underlines the partial reduction of  $\text{Tm}^{3+}$  into  $\text{Tm}^{2+}$  during the  $\text{CaF}_2$  growth process. The high purity of the





**Fig. 8.** LIBS imaging on three samples of each of the three attempts at different stages of the 3 boules: beginning, centre and end of the cylinder. Colours illustrate the relative concentration of Tm content with increasing Tm content from blue colour (lower content) to green colour (intermediate content) and to red colour (higher content).



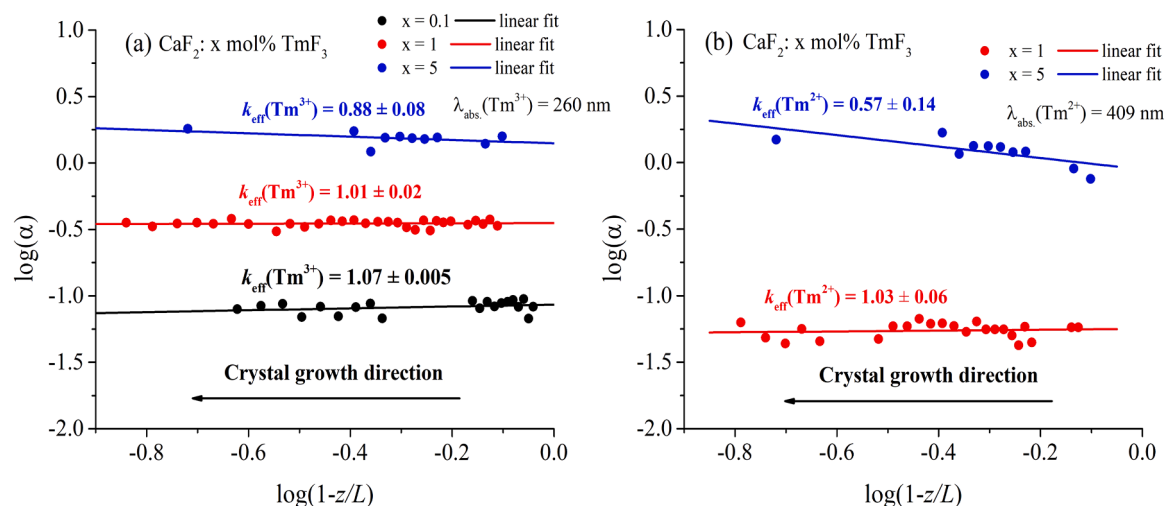
**Fig. 9.** Room temperature absorption spectra of TmF<sub>3</sub>-doped CaF<sub>2</sub> crystals where the transitions of the observed absorption bands are depicted. The absorption bands selected for evaluating the effective partition coefficients ( $k_{\text{eff}}$ ) of Tm<sup>2+</sup> and Tm<sup>3+</sup> are indicated with blue arrows.

commercial TmF<sub>3</sub> raw material used in this work is not questioned. Ellingham diagram calculations related to F<sub>2</sub> gas [66] is considered up to the maximum growth working temperature 1785 K. The Gibbs free enthalpies of reaction  $\Delta_r G^\circ$  of fluorides are calculated from enthalpies of formation  $\Delta_f H^\circ$  and standard entropies  $S^\circ$  (298 K) (Table 5). Note that  $\Delta_r G^\circ(\text{TmF}_{3(l)})$  has been extrapolated beyond TmF<sub>3(s)</sub> melting point from  $\Delta_r G^\circ(\text{TmF}_{3(s)})$  by assuming similar  $\Delta_f H^\circ$  and  $S^\circ$  values for both solid and liquid phases such as observed with CaF<sub>2(s,l)</sub>. It is consistent with the values reported by Greis *et al.* [67] who mentioned the same trend for both TmF<sub>3(s)</sub> and TmF<sub>3(l)</sub>. Ellingham calculations pointed out that C<sub>(s)</sub>,

CaF<sub>2(s,l)</sub> and TmF<sub>3(s,l)</sub> are the stable compounds in the present system under vacuum. However, the partial pressures of F<sub>2(g)</sub> have been calculated in the range 1691K-1785K for TmF<sub>3</sub> and CaF<sub>2</sub> formation equilibrium. At these temperatures, respectively, partial vapour pressure P(F<sub>2(g)</sub>) has been estimated ranging from  $\sim 1 \times 10^{-29}$  to  $5 \times 10^{-26}$  atm. for CaF<sub>2</sub> and P(F<sub>2(g)</sub>)  $\sim 1 \times 10^{-27}$  to  $1 \times 10^{-24}$  atm. for TmF<sub>3</sub>. As the growth is performed under vacuum atmosphere in a C<sub>(s)</sub> graphite crucible, the reaction between the continuously released F<sub>2(g)</sub> vapour and C<sub>(s)</sub> can form continuously CF<sub>4(g)</sub>. This F<sub>2</sub> depletion in the liquid phase leads to the Tm<sup>3+</sup> to Tm<sup>2+</sup> conversion driven by the charge compensation phenomenon within the melt and the as-grown crystal.

#### • Tm<sup>3+</sup> reduction by synproportionation reaction

Although TmF<sub>3</sub> volatilization is considered as congruent [70,71], stable non-stoichiometric thulium-based fluorides such as Tm<sub>13</sub>F<sub>32.5</sub> [72], TmF<sub>2.37</sub> [73], TmF<sub>2.40</sub> [67,70,74], TmF<sub>2.38</sub> and TmF<sub>2.42</sub> [75] are reported when TmF<sub>3</sub> is in the presence of metallic Tm. These compounds are formed by synproportionation reaction [76] and exhibit a mixed valence of Tm. They are chemically similar to other well-known rare-earth trifluorides, with a higher divalent state stability [76], that reduce to stoichiometric or non-stoichiometric compounds such as EuF<sub>3</sub> into EuF<sub>2</sub> [70], SmF<sub>3</sub> into SmF<sub>2</sub>, SmF<sub>2.35</sub> and SmF<sub>2.40</sub> [70,77], and YbF<sub>3</sub> into YbF<sub>2</sub>, YbF<sub>2.375</sub> and YbF<sub>2.40</sub> [70,78]. Such redox reaction for Tm can be written theoretically as  $\text{Tm} + 2\text{TmF}_3 \rightleftharpoons 3\text{TmF}_2$  where a small Tm amount can be originated from the formation equilibrium of TmF<sub>3</sub>. However, TmF<sub>2</sub> having a hypothetical existence [67,70,74], in the case of CaF<sub>2</sub>:Tm<sup>2+</sup>, thermodynamic instability of divalent state of Tm has been reported [76,79] causing Tm<sup>2+</sup> to slowly oxidize into Tm<sup>3+</sup> via the reverse redox reaction known as disproportionation reaction:  $3\text{TmF}_2 \rightleftharpoons \text{Tm} + 2\text{TmF}_3$ . Hence, the qualitative but significant spectroscopic detection of both Tm<sup>2+</sup> and Tm<sup>3+</sup> in as-grown CaF<sub>2</sub> single crystal of the present work is assumed to arise partially from the synproportionation reaction of TmF<sub>3</sub>.



**Fig. 10.** partition coefficients of (a)  $\text{Tm}^{3+}$  cations and (b)  $\text{Tm}^{2+}$  cations in  $\text{TmF}_3$ : doped  $\text{CaF}_2$  crystals. Due to inaccuracy of measurements and the very low  $\text{Tm}^{2+}$  content for  $x=0.1$  %, no line could be fitted with experimental data.

**Table 4**

Calculated effective segregation coefficient of  $\text{CaF}_2$ :  $x$  mol%  $\text{TmF}_3$  crystals ( $x = 0.1$  %, 1 % and 5 %).

Ions	$x = 0.1$ This work	$x = 1$ This work	$x = 5$ This work	Tm: YAG Crystal [57]	Tm: Ca(Gd,Lu)AlO <sub>4</sub> crystal [55]	Tm: YAP [58]	Tm: Fe-garnet crystal [60]	Tm: Ca <sub>3</sub> (VO <sub>4</sub> ) <sub>2</sub> [59]	Tm: CGA [56]
$\text{Tm}^{3+}$	$1.070 \pm 0.005$	$1.01 \pm 0.02$	$0.88 \pm 0.08$	1.04	0.90	0.88	0.85	0.8	0.61
$\text{Tm}^{2+}$	-	$1.03 \pm 0.06$	$0.57 \pm 0.14$	-	-	-	-	-	-

**Table 5**

Thermodynamics values of C, Ca and Tm fluorides used in Ellingham diagram.

Element or compound (s=solid, l=liquid, g=gas)	$\Delta_f H^0$ (kJ. mol <sup>-1</sup> )	$S^0$ (J.K <sup>-1</sup> . mol <sup>-1</sup> )	Melting or boiling point (K)	Chemical equilibrium	$\Delta_r H^0$ (kJ. mol <sup>-1</sup> )	$\Delta_r S^0$ (J.K <sup>-1</sup> . mol <sup>-1</sup> )	Reference
$\text{F}_2$ (g)	0	202.80	53.33	X	X	X	[66–69]
C (s)	0	5.74	3823	X	X	X	
$\text{CF}_4$ (g)	-933.60	261.60	89.4	$\frac{1}{2} \text{C(s)} + \text{F}_2(\text{g}) \rightleftharpoons \frac{1}{2} \text{CF}_4(\text{g})$	-466.8	-74.87	
Ca (s)	0	41.59	1115	X	X	X	
Ca (l)	7.79	45.51	1757	X	X	X	
Ca (g)	177.80	154.89	X	X	X	X	
$\text{CaF}_2$ (s)	-1228	68.50	1691	$\text{Ca(s)} + \text{F}_2(\text{g}) \rightleftharpoons \text{CaF}_2(\text{s})$	-1228.00	-175.89	
$\text{CaF}_2$ (s)	"	"	"	$\text{Ca(l)} + \text{F}_2(\text{g}) \rightleftharpoons \text{CaF}_2(\text{s})$	-1235.79	-179.81	
$\text{CaF}_2$ (l)	-1184.07	92.59	1757	$\text{Ca(l)} + \text{F}_2(\text{g}) \rightleftharpoons \text{CaF}_2(\text{l})$	-1191.86	-155.72	
$\text{CaF}_2$ (l)	"	"	"	$\text{Ca(g)} + \text{F}_2(\text{g}) \rightleftharpoons \text{CaF}_2(\text{l})$	-1361.87	-265.097	
Tm (s)	0	74	1818	X	X	X	
$\text{TmF}_3$ (s,l)	-1695	118	1431	$\frac{2}{3} \text{Tm(s)} + \text{F}_2(\text{g}) \rightleftharpoons \frac{2}{3} \text{TmF}_3(\text{s,l})$	-1130.00	-173.47	

#### 4. Conclusion

Centimetre-sized Tm-doped  $\text{CaF}_2$  crystals were successfully grown by the Bridgman-Stockbarger technique with different  $\text{TmF}_3$  initial contents: 0.1 mol.%, 1 mol.% and 5 mol.%.  $\text{CaF}_2$  crystal habit was highlighted by Laue diffraction which have shown that natural growth orientation is preferentially along the (110) and (210) crystallographic directions with  $G=12 \text{ K.cm}^{-1}$  as thermal gradient and  $v=4 \text{ mm.h}^{-1}$  as pulling speed. A drop of global [Tm] content in the conical part of the as-grown boule was highlighted by LIBS measurements. As suggested by numerical modelling, it is caused by a poor mixing of elements in the liquid phase at the beginning of the growth due to an one order of magnitude lower convection in comparison with the melt convection in the later growth stages, where the bending of solid-liquid interface generates a higher convection. Chemical characterization has shown that single crystals exhibit a highly homogenous [Tm] content distribution throughout the whole boule volume but optical absorption

measurements have displayed the presence of both  $\text{Tm}^{3+}$  and  $\text{Tm}^{2+}$  cations. The simultaneous presence of these cations is assumed to be likely due to the carbon environment and synproportionation partial reaction of  $\text{TmF}_3$  which leads to Tm-mixed-valency-doped  $\text{CaF}_2$  single crystals since  $\text{Tm}^{2+}$  can enter easily in the  $\text{CaF}_2$  matrix because no charge compensation is needed.

Through the results from growth attempt with  $C_L^0=5$  mol.%  $\text{TmF}_3$  initial content, effective partition coefficient of [Tm] has been estimated. LIBS coupled with ICP-AES gave  $0.906 < k_{\text{eff}}([\text{Tm}])_{\text{LIBS}} < 1.126$  where [Tm] total content corresponds to both  $\text{Tm}^{2+}$  and  $\text{Tm}^{3+}$  contents. Furthermore, optical absorption measurements gave  $k_{\text{eff}}(\text{Tm}^{3+})_{\text{OA}} = 0.88 \pm 0.08$  and  $k_{\text{eff}}(\text{Tm}^{2+})_{\text{OA}} = 0.52 \pm 0.14$  leading to a global partition coefficient  $k_{\text{eff}}([\text{Tm}])_{\text{OA}} < 1$ . From these results, it is finally concluded that  $0.906 \leq k_{\text{eff}}([\text{Tm}])_{\text{LIBS+OA}} < 1$  suggesting, thus, that the expected congruent point in the  $\text{CaF}_2$ -rich region of  $\text{CaF}_2$ - $\text{TmF}_3$  phase diagram does exist and lie within a range higher than 1 mol.% and lower than 5 mol.%  $\text{TmF}_3$  molar fraction. Further physical characterization of  $\text{Tm}^{2+}$

and Tm<sup>3+</sup> content, such as with Electron Paramagnetic Resonance (EPR) technique and X-ray Photoelectron Spectroscopy (XPS), should have to be carried out in order to quantify their amount in both the raw material and as-grown crystals.

### CRedit authorship contribution statement

**Frédéric Pelascini:** Writing – original draft, Resources, Formal analysis, Data curation. **Daniel Vizman:** Writing – original draft, Software, Methodology, Investigation, Formal analysis, Data curation. **Jérôme Debray:** Writing – original draft, Resources, Investigation, Formal analysis, Data curation. **Vincent Motto-Ros:** Writing – original draft, Resources, Investigation, Formal analysis, Data curation. **Dragos Tatomirescu:** Writing – original draft, Software, Investigation, Formal analysis, Data curation. **Alexandra Popescu:** Writing – original draft, Software, Methodology, Investigation, Formal analysis, Data curation. **Maria Poienar:** Writing – original draft, Software, Investigation, Formal analysis, Data curation. **Carla Schornig:** Writing – original draft, Methodology, Investigation, Formal analysis, Data curation. **Marius Stef:** Writing – original draft, Methodology, Investigation, Formal analysis, Data curation. **Philippe Roger Veber:** Writing – original draft, Validation, Supervision, Resources, Project administration, Methodology, Investigation, Funding acquisition, Formal analysis, Data curation. **Gabriel Buse:** Writing – original draft, Resources, Investigation, Formal analysis, Data curation.

### Declaration of Competing Interest

The authors declare the following financial interests/personal relationships which may be considered as potential competing interests: Philippe Veber reports financial support was provided by European Commission. Philippe Veber reports a relationship with European Commission that includes: funding grants. If there are other authors, they declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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### Appendix A. Supporting information

Supplementary data associated with this article can be found in the online version at doi:10.1016/j.jallcom.2024.177784.

### Data availability

Data will be made available on request.

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