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A THREE-DIMENSIONAL VECTOR POTENTIAL FORMULATION  
FOR SOLVING TRANSONIC FLOWS WITH MIXED FINITE ELEMENTS

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We present in this paper a three dimensional extension of the work done by AMARA. We firstly recall the physical model. The iterative numerical algorithm of resolution introduces an elliptic problem which a variational formulation involving a vector potential is given in section 2. Furthermore we describe its approximation using the NEDELEC's element conforming in  $H(\text{curl})$ . Finally preliminary numerical experiment is presented in a subsonic case.

### 1. Potential transonic flow.

We study an inviscid fluid in stationary irrotational and isentropic evolution. With those assumptions it is known [1] that only the two physical fields density  $\rho$  and velocity  $\vec{u}$  are necessary to its description. The corresponding partial differential equations are

$$\text{mass conservation} \quad (1.1) \quad \text{div}(\rho \vec{u}) = 0$$

$$\text{no vorticity} \quad (1.2) \quad \text{curl} \vec{u} = 0$$

furthermore the Bernoulli theorem can be written :

$$(1.3) \quad \rho = \left\{ 1 + \frac{\gamma - 1}{2} M_\infty^2 (1 - |\vec{u}|^2) \right\}^{1/\gamma - 1}$$

where  $M_\infty$  is the upstream Mach number and  $\gamma$  the specific heat ratio. This problem is non-linear, and the type of the equations is mixed : elliptic in the subsonic zone and hyperbolic in the supersonic one. Moreover an entropy condition has to be added to select the physical shock waves [2,3] .

The domain  $\Omega$  is discretized with a triangulation  $T_h$  of finite elements (we used tetrahedrons and prisms) and the density  $\rho$  is taken to be constant in each element. The set of equations (1.1), (1.2), (1.3) is solved numerically with a fixed point technique and an artificial density [4,5] .

### 2. Formulation of the linear problem.

In all this section we investigate the linear problem for the velocity ; so the density  $\rho$  is a fixed positive function in  $L^\infty(\Omega)$ .

$$(P) \quad \left\{ \begin{array}{ll} (2.1) & \text{div}(\rho \vec{u}) = 0 \quad \Omega \\ (2.2) & \text{curl} \vec{u} = 0 \quad \Omega \\ (2.3) & \rho \vec{u} \cdot \vec{n} = g \quad \Gamma \end{array} \right.$$

where  $\vec{n}$  is the exterior normal of the boundary  $\Gamma$ . Since we restrict our study to the case where both  $\Omega$  and  $\Gamma$  are connected domains, the equation (2.1) gives the existence of a vector potential  $\psi$  [6] such that

$$(2.4) \quad \rho \vec{u} - \text{curl} \psi = 0 \quad \Omega$$

The condition (2.3) is then equivalent to

$$(2.5) \quad \text{div}_\Gamma (\psi \times \vec{n}) = g$$

where  $\text{div}_\Gamma$  is the Laplace-Beltrami divergence operator [7] . Note that the condi-

tions (2.4), (2.5) do not insure the uniqueness of a potential  $\psi$  ; it is therefore necessary to add gauge conditions [8,9] like

$$(2.6) \quad \operatorname{div} \psi = 0 \quad \Omega$$

$$(2.7) \quad \operatorname{curl}_{\Gamma} (\psi \times \vec{n}) = 0 \quad \Gamma$$

• Let us introduce now some functional spaces :

$$H(\operatorname{curl}) = \{ \chi \in L^2(\Omega), \operatorname{curl} \chi \in L^2(\Omega) \}$$

$$H^0(\operatorname{curl}) = \{ \chi \in H(\operatorname{curl}), \chi \times \vec{n} = 0 \text{ on } \Gamma \}$$

$$V = \{ \chi \in H^0(\operatorname{curl}), \operatorname{div} \chi = 0 \}$$

$$U = \{ u \in H(\operatorname{div}), \operatorname{div} u = 0, u \cdot \vec{n} = 0 \text{ on } \Gamma \}$$

we know that there is some vector field  $\alpha \in H^1(\Omega)$  satisfying (2.8)  $\operatorname{div}_{\Gamma}(\alpha \times \vec{n}) = g$  and the problem (P) is equivalent to the variational problem

$$(Q) \quad \left\{ \begin{array}{l} \text{find } \phi \in V \\ (2.10) \quad \forall \chi \in V, \int_{\Omega} \frac{1}{\rho} \operatorname{curl} \phi \cdot \operatorname{curl} \chi \, dx = - \int_{\Omega} \frac{1}{\rho} \operatorname{curl} \alpha \cdot \operatorname{curl} \chi \, dx . \end{array} \right.$$

which admits a unique solution since Poincaré inequality holds in  $V$  [6] .

### 3. Conformial element in $H(\operatorname{curl})$ .

As the continuous problem (Q) is well posed in a subspace of  $H(\operatorname{curl})$ , a natural idea is to discretize this vector space. NEDELEC [10] proposed the following finite element  $(K, P_K, \Sigma_K)$  [fig.1] :  $K$  is a tetrahedron,

$$P_K = \{ \varphi, \exists \alpha, \beta \in \mathbb{R}^3, \forall x \in K, \varphi(x) = \alpha + \beta \times x \}$$

$$\Sigma_K = \{ \varphi \cdot \tau, \tau \text{ is an edge of } K \} .$$

Let us introduce also the finite dimensional spaces :

$$X_h = \{ \varphi_h \in H(\operatorname{curl}), \forall K \in \mathcal{T}_h, \varphi_h|_K \in P_K \}$$

$$X_h^0 = \{ \varphi_h \in X_h, \varphi_h \times \vec{n} = 0 \text{ on } \Gamma \}$$

$$U_h = \{ \operatorname{curl} \varphi_h, \varphi_h \in X_h^0 \}$$

$$H_{0,h}^1 = \{ v_h \in H_0^1, v_h \text{ polynomial of degree 1 in each } K \in \mathcal{T}_h \}$$

This was extended to prisms.

• We now focus on the boundary condition (2.8). In the continuous case the existence of  $\alpha$  satisfying (2.8) is given by a Laplace-Beltrami problem on the manifold  $\Gamma$ . We discretize the boundary condition (2.8) in  $X_h$  by the relations

$$(3.1) \quad \sum_{\tau \in \mathcal{T}} \alpha_h \cdot \tau = g_h(T) \quad \forall T \text{ triangle of } \Gamma$$

where  $g_h(T)$  is constant. We just need a procedure to compute explicitly such an  $\alpha_h$  : the edges of the triangulation lying on  $\Gamma$  allow to define a graph  $g$  [11] in

the set of all the vertices of  $\Gamma$ . Let  $A$  be a tree in this graph  $g$  [fig.2] . We choose the space  $W_h(\Gamma)$  containing  $\alpha_h$  as follow :

$$(3.2) \quad W_h(\Gamma) = \{ \varphi \in X_h, \varphi \cdot \tau = 0 \text{ if } \tau \in \Gamma \cap A \text{ or } \tau \notin \Gamma \}$$

and if  $\int_{\Gamma} g_h \, d\gamma = 0$  then (3.1) admits a unique solution  $\alpha_h \in W_h(\Gamma)$ .

• Assume now that  $\alpha_h$  is computed, we solve numerically the problem

$$(Q_h) \quad \left\{ \begin{array}{l} \text{find } \phi_h \in V_h \\ \forall \chi_h \in V_h, \int_{\Omega} \frac{1}{\rho} \text{curl} \phi_h \cdot \text{curl} \chi_h \, dx = - \int_{\Omega} \frac{1}{\rho} \text{curl} \alpha_h \cdot \text{curl} \chi_h \, dx \end{array} \right.$$

we have to choose the space  $V_h$ ,  $V_h \subset X_h^0$ , in order to adapt the gauge condition (2.6) in a finite dimensional space. The constraint on  $V_h$  is to have :

$$(3.3) \quad \text{the application } \text{curl} : V_h \rightarrow U_h \text{ is one to one}$$

(like in the continuous case for  $\text{curl} : V \rightarrow U$ ). Since the kernel of  $\text{curl} : X_h^0 \rightarrow U_h$  is not null,  $V_h$  must be a strict subspace of  $X_h^0$ . Then if  $V_h$  is any subspace of  $X_h^0$  satisfying (3.3), the solution  $\phi$  (resp  $\phi_h$ ) of (Q) (resp  $(Q_h)$ ) are such that the error  $\| \text{curl} \phi - \text{curl} \phi_h \|_{L^2}$  vanishes like  $h$ . Note that the fields  $\phi$  and  $\phi_h$  have no physical interpretation, furthermore it does not exist a priori any simple mathematical relation between them. NEDELEC in [12] proposed to use

$$\tilde{V}_h = \{ \varphi_h \in X_h^0, \int_{\Omega} \varphi_h \cdot \nabla \theta_h \, dx = 0 \quad \forall \theta_h \in H_{0,h}^1 \}$$

A basis of  $\tilde{V}_h$  is non local in space, and difficult to compute. For this reason we choose  $V_h$  in a purely algebraical way, following [13] and including references in an other context.

The space  $X_h^0$  is generated by the edges of  $T_h$  which are not lying in  $\Gamma$ . Consider the graph  $g'$  of the edges coupling two vertices of  $T_h$  which are not on  $\Gamma$  ( $g'$  is a strict subset of all the edges of  $X_h^0$ ), and  $A'$  a tree in this graph. Furthermore let  $\tau'$  be an edge joining up one vertex of  $\Gamma$  and one in the interior of  $\Omega$ . Then  $V_h$ , defined by

$$V_h = \{ \varphi_h \in X_h^0, \varphi_h \cdot \tau = 0 \text{ if } \tau \in A' \text{ or } \tau = \tau' \}$$

satisfies the condition (3.3) [fig.3] .

#### 4. Implementation.

The method described in paragraph 3 is beeing developped on the CRAY 1-S of CCVR. First results have been obtained in a subsonic case and compared with a 2D computation to validate the code [14] : from the channel of GAMM 79 workshop

[15] a 3D mesh using prisms has been generated [fig.4] ( $71 \times 20 \times 4$  elements, 6879 unknowns). The linear problem is solved by a preconditioning conjugate gradient algorithm. Several preconditioning matrices have been tested. A good and cheap choice seems to be the SSOR method. The figure [fig.5] shows isomach curves on the top and the bottom of the channel for both 2D and 3D computations.

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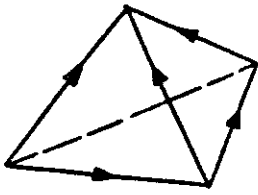


Figure 1

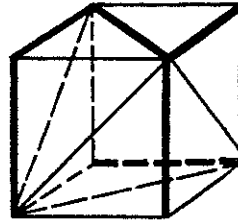


Figure 2 Example of a boundary tree if the domain is a cube divided in five tetrahedrons

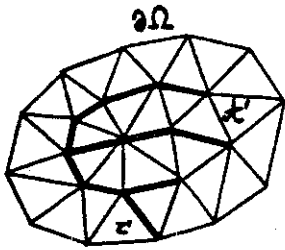


Figure 3

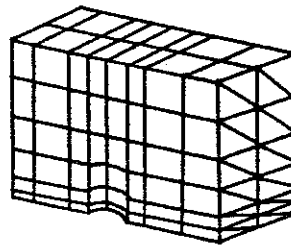


Figure 4

MACH EN ENTREE 0.500000 DEBIT TOTAL 2.073000  
NOMBRE DE MACH COUCHE 2

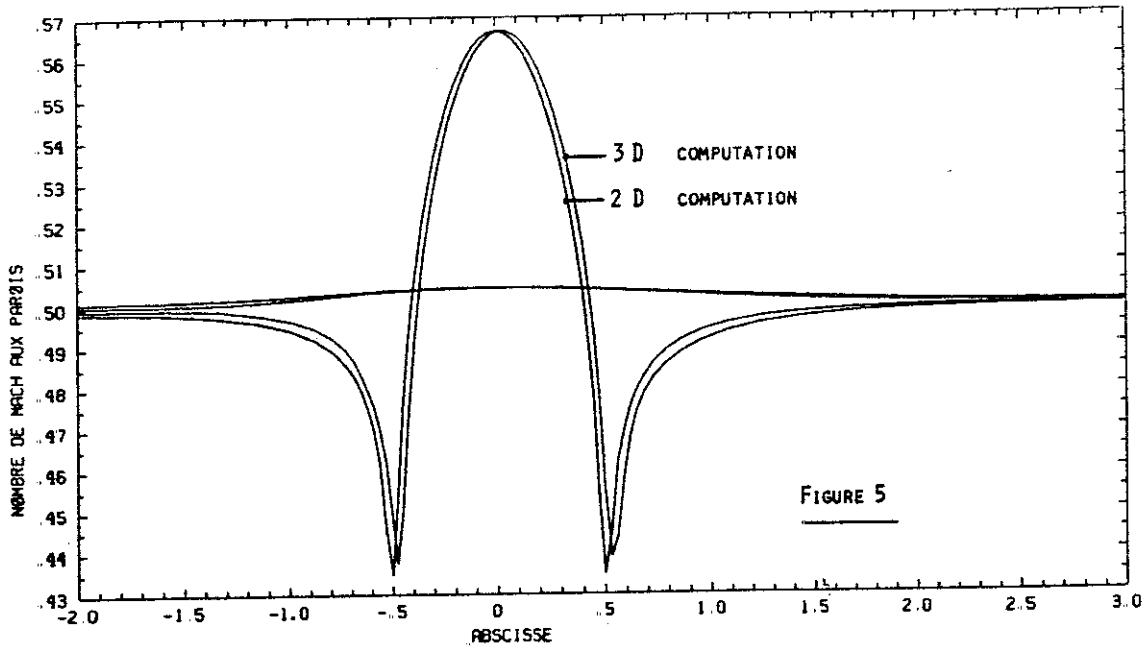


FIGURE 5