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Stability results for Sobolev, logarithmic Sobolev, and related inequalities

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Abstract. Obtaining explicit stability estimates in classical functional inequalities like the Sobolev inequality has been an essentially open question for 30 years, after the celebrated but non-constructive result [4] of G. Bianchi and H. Egnell in 1991. Recently, new methods have emerged which provide some clues on these fascinating questions. The goal of the course is to give an introduction to the topic for some fundamental functional inequalities and present several methods that can be used to obtain explicit estimates.

Keywords: Stability, functional inequalities, Sobolev inequality, logarithmic Sobolev inequality, Gagliardo-Nirenberg inequalities, constructive estimates, nonlinear diffusions, entropy methods, carré du champ

The most classical Sobolev inequality on \mathbb{R}^d with $d \geq 3$ is

$$\|\nabla f\|_{L^2(\mathbb{R}^d)}^2 \geq S_d \|f\|_{L^{2^*}(\mathbb{R}^d)}^2 \quad \forall f \in \mathcal{D}^{1,2}(\mathbb{R}^d) \quad (1)$$

where $\mathcal{D}^{1,2}(\mathbb{R}^d)$ is the Beppo-Levi space, S_d denotes the sharp Sobolev constant and $2^* = 2d/(d-2)$ is the critical Sobolev exponent. Equality in (1) holds if and only if f is in the $(d+2)$ -dimensional manifold \mathcal{M} of the *Aubin-Talenti functions*

$$g_{a,b,c}(x) = c(a + |x - b|^2)^{-\frac{d-2}{2}}, \quad a \in (0, \infty), \quad b \in \mathbb{R}^d, \quad c \in \mathbb{R}.$$

We refer to [39] for an up-to-date introduction to Ineq. (1). In 1991, G. Bianchi and H. Egnell proved in [4] that the *deficit* in the Sobolev inequality, that is, the difference of the two sides of (1), satisfies the *stability* inequality

$$\|\nabla f\|_{L^2(\mathbb{R}^d)}^2 - S_d \|f\|_{L^{2^*}(\mathbb{R}^d)}^2 \geq \kappa_d \min_{g \in \mathcal{M}} \|\nabla(f - g)\|_{L^2(\mathbb{R}^d)}^2 \quad (2)$$

for some constant $\kappa_d > 0$. The method of [4] uses compactness and provides no estimate on κ_d . Recently, several *constructive* results have been obtained, *i.e.*, results with explicit estimates of κ_d . The goal of these lecture notes is to make a short review based on two main strategies:

1. The strategy of G. Bianchi and H. Egnell made explicit in both of its steps: the global-to-local reduction and the local analysis, both relying on methods of the Calculus of Variations,
2. *Entropy methods* and *fast diffusion equations*, in which we can recognize also the same two steps, that are based on nonlinear parabolic flows and *carré du champ* techniques.

Entropy methods usually result in better estimates of the stability constant, but to the price of some limitations. We will also state some results for related inequalities, like some families of Gagliardo-Nirenberg interpolation inequalities and logarithmic Sobolev inequalities, on the Euclidean space \mathbb{R}^d or on the unit sphere $\mathbb{S}^d \subset \mathbb{R}^{d+1}$, and consider the use of the Legendre duality applied, for instance, to Hardy-Littlewood-Sobolev inequalities.

1 Sobolev and Hardy-Littlewood-Sobolev inequalities

Let us start with a first example of fast diffusion flows applied to inequalities in the context of the Legendre duality. On \mathbb{R}^d with $d \geq 3$, the *Hardy-Littlewood-Sobolev inequality*

$$\|g\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 \geq S_d \int_{\mathbb{R}^d} g(-\Delta)^{-1} g \, dx \quad \forall g \in L^{\frac{2d}{d+2}}(\mathbb{R}^d) \quad (3)$$

is obtained by Legendre's duality from (1), as observed in [49]. Indeed, by the Legendre transform, we obtain

$$\begin{aligned} \sup_{f \in \mathcal{D}^{1,2}(\mathbb{R}^d)} \left(\int_{\mathbb{R}^d} f g \, dx - \frac{1}{2} \|f\|_{L^{2^*}(\mathbb{R}^d)}^2 \right) &= \frac{1}{2} \|g\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 \\ \sup_{f \in \mathcal{D}^{1,2}(\mathbb{R}^d)} \left(\int_{\mathbb{R}^d} f g \, dx - \frac{1}{2} \|\nabla f\|_{L^2(\mathbb{R}^d)}^2 \right) &= \frac{1}{2} \int_{\mathbb{R}^d} g(-\Delta)^{-1} g \, dx \end{aligned}$$

where $q = \frac{d+2}{d-2}$ is the Hölder conjugate of 2^* , so that $\|f^q\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)} = \|f\|_{L^{2^*}(\mathbb{R}^d)}^q$. It is not difficult to check that S_d is also the optimal constant in (3).

Theorem 1 ([33]). *Assume that $d \geq 3$. There exists a constant $\mathcal{C} \in [\frac{d}{d+4}, 1)$ such that*

$$\begin{aligned} \|f^q\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 - S_d \int_{\mathbb{R}^d} f^q (-\Delta)^{-1} f^q \, dx \\ \leq \mathcal{C} S_d^{-1} \|f\|_{L^{2^*}(\mathbb{R}^d)}^{\frac{8}{d-2}} \left(\|\nabla f\|_{L^2(\mathbb{R}^d)}^2 - S_d \|f\|_{L^{2^*}(\mathbb{R}^d)}^2 \right) \quad \forall f \in \mathcal{D}^{1,2}(\mathbb{R}^d). \end{aligned} \quad (4)$$

Theorem 1 can be considered as a stability result for (1), because the left-hand side in (4), which is the deficit associated with (3), measures a distance to \mathcal{M} . However, it is still a nonlinear functional and corresponds to a weaker norm than

the usual norm in $\mathcal{D}^{1,2}(\mathbb{R}^d)$. With $\mathcal{C} = 1$, there is an easy proof that goes as follows. Two integrations by parts show that

$$\int_{\mathbb{R}^d} |\nabla(-\Delta)^{-1} g|^2 dx = \int_{\mathbb{R}^d} g (-\Delta)^{-1} g dx$$

and, if $g = f^q$ with $q = \frac{d+2}{d-2}$,

$$\int_{\mathbb{R}^d} \nabla f \cdot \nabla(-\Delta)^{-1} g dx = \int_{\mathbb{R}^d} f g dx = \int_{\mathbb{R}^d} g^q dx = \int_{\mathbb{R}^d} f^{2^*} dx.$$

Hence the expansion of the square

$$0 \leq \int_{\mathbb{R}^d} \left| \|f\|_{L^{2^*}(\mathbb{R}^d)}^{\frac{4}{d-2}} \nabla f - S_d \nabla(-\Delta)^{-1} g \right|^2 dx$$

shows the result with $\mathcal{C} = 1$. To prove that the inequality holds with $\mathcal{C} < 1$ requires finer tools.

Let us consider the *fast diffusion* equation

$$\frac{\partial v}{\partial t} = \Delta v^m, \quad (t, x) \in \mathbb{R}^+ \times \mathbb{R}^d \quad (5)$$

and compute the evolution of the deficit in (3) defined as

$$-H := \|v\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 - S_d \int_{\mathbb{R}^d} v (-\Delta)^{-1} v dx \geq 0.$$

As in [27], we observe that

$$\frac{1}{2} H' = \left(\int_{\mathbb{R}^d} v^{\frac{2d}{d+2}} dx \right)^{\frac{2}{d}} \int_{\mathbb{R}^d} \nabla v^m \cdot \nabla v^{\frac{d-2}{d+2}} dx - S_d \int_{\mathbb{R}^d} v^{m+1} dx$$

where $v = v(t, \cdot)$ is a solution of (5). With the choice $m = (d-2)/(d+2)$, that is, $m+1 = 2d/(d+2) = q$, which corresponds to the exponent in the *Yamabe flow*, and $u = v^m$, we have

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \left(S_d \int_{\mathbb{R}^d} v (-\Delta)^{-1} v dx - \|v\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 \right) \\ = \left(\int_{\mathbb{R}^d} v^{m+1} dx \right)^{\frac{2}{d}} \left(\|\nabla u\|_{L^2(\mathbb{R}^d)}^2 - S_d \|u\|_{L^{2^*}(\mathbb{R}^d)}^2 \right) \geq 0. \end{aligned}$$

Eq. (5) admits a solution with *separation of variables* vanishing at $t = T > 0$,

$$\bar{v}_T(t, x) = c(T-t)^\alpha (g_\star(x))^{\frac{d+2}{d-2}}, \quad \alpha = \frac{1}{4}(d+2) \quad \text{and} \quad c = \left(4d \frac{d-2}{d+2}\right)^\alpha,$$

where $g_\star(x) = (1+|x|^2)^{-(d-2)/2}$ is the solution of $\Delta g_\star + d(d-2)g_\star^{(d+2)/(d-2)} = 0$. Up to conformal transformations, this solutions describes the asymptotic extinction profile of all solutions of (5) near the extinction time T .

Theorem 2 ([23]). *For any solution v with initial datum $v_0 \in L^{2d/(d+2)}(\mathbb{R}^d)$, $v_0 > 0$, there exists $T > 0$, $\lambda > 0$ and $x_0 \in \mathbb{R}^d$ such that*

$$\lim_{t \rightarrow T_-} (T-t)^{-\frac{1}{1-m}} \sup_{x \in \mathbb{R}^d} (1+|x|^2)^{d+2} \left| \frac{v(t,x)}{\bar{v}(t,x)} - 1 \right| = 0$$

with $\bar{v}(t,x) = \lambda^{-(d+2)/2} \bar{v}_T(t, (x-x_0)/\lambda)$.

The Hardy-Littlewood-Sobolev inequality (3) amounts to $H \leq 0$ and appears as a consequence of the Sobolev inequality (1) because $H' \geq 0$ for any $t \in (0, T)$ and $\lim_{t \rightarrow T_-} H(t) = 0$ by Theorem 2. By computing H'' , one can obtain refined estimates. Assume that $t \in [0, T)$ and let Y be such that

$$H(t) = -Y(J(t)) \quad \text{with} \quad J(t) := \int_{\mathbb{R}^d} v(t,x)^{\frac{2d}{d+2}} dx.$$

With $\kappa_0 = H'_0/J_0$, we obtain the differential inequality

$$Y' \left(S_d s^{1+\frac{2}{d}} + Y \right) \leq \frac{d+2}{2d} \kappa_0 S_d^2 s^{1+\frac{4}{d}}, \quad Y(0) = 0, \quad Y(J_0) = -H_0.$$

The relation $Y'(J)' + H' = 0$ and elementary integrations by parts provide the following refinement of (4) written with $\mathcal{C} = 1$.

Proposition 1 ([33]). *Assume that $d \geq 3$. Then for any $f \in \mathcal{D}^{1,2}(\mathbb{R}^d)$ we have*

$$\begin{aligned} & \|f^q\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 - S_d \int_{\mathbb{R}^d} f^q (-\Delta)^{-1} f^q dx \\ & \leq \|f\|_{L^{2^*}(\mathbb{R}^d)}^{\frac{2d+2}{d-2}} \varphi \left(S_d^{-1} \|f\|_{L^{2^*}(\mathbb{R}^d)}^{-2} \left(\|\nabla f\|_{L^2(\mathbb{R}^d)}^2 - S_d \|f\|_{L^{2^*}(\mathbb{R}^d)}^2 \right) \right) \end{aligned}$$

with $q = (d+2)/(d-2)$ and $\varphi(s) := \sqrt{1+2s} - 1$ for any $s \geq 0$.

An asymptotic analysis of $v(t, \cdot)$ as $t \rightarrow T_-$ and some spectral estimates discard the case $\mathcal{C} = 1$ in Theorem 1, while the lower bound $\mathcal{C} = d/(d+4)$ is obtained by considering a sequence of test functions. See [33] for details. In dimension $d = 2$ similar results holds in which (1) and (3) are respectively replaced by Onofri's inequality and the logarithmic Hardy-Littlewood-Sobolev inequality while (5) has to be replaced by the logarithmic fast diffusion equation using the results of [55,50,54] instead of Theorem 2. Notice that stability results based on duality are known for (logarithmic) Hardy-Littlewood-Sobolev inequalities from [16,17]. The method based on (5) also applies to the fractional Sobolev inequality and the dual Hardy-Littlewood-Sobolev inequality as shown in [45]. Even if the optimal value of the stability constant \mathcal{C} is not known, a nice feature of Theorem 1 is that we have a rather precise estimate of \mathcal{C} . However, a significant drawback of the method is that the deficit of the Hardy-Littlewood-Sobolev inequality corresponds to a weaker notion of distance to \mathcal{M} than what can be expected.

2 Stability, fast diffusion equation and entropy methods

On \mathbb{R}^d , let us consider the *fast diffusion equation* written in self-similar variables,

$$\frac{\partial v}{\partial t} + \nabla \cdot \left(v (\nabla v^{m-1} - 2x) \right) = 0. \quad (6)$$

Up to a t -dependent rescaling, it is equivalent to (5). By *fast diffusion*, we mean $m < 1$ and reserve the expression *porous medium* to the regime $m > 1$ while the heat equation case corresponds to $m = 1$ (in this case, Eq (6) is known as the *Fokker-Planck equation*, and v^{m-1} should be replaced by $-\log v$). An introduction can be found in the two classical books [52,53]. Section 1 was devoted to $m = m_* := (d-2)/(d+2)$ with $d \geq 3$. Here we are interested in the range $m_c := (d-2)/d < m < 1$ in which nonnegative solution with initial data in $L^1(\mathbb{R}^d)$ at $t = 0$ exist for any $t \geq 0$ according to [42], and more specifically to the range $m_1 \leq m < 1$ where $m_1 := (d-1)/d$. Note that $m_* < m_c < m_1 < 1$ in dimension $d \geq 2$. In any case, (6) admits a stationary solution

$$\mathcal{B}(x) := (1 + |x|^2)^{\frac{1}{m-1}}$$

which is usually called a *Barenblatt solution* and has finite mass $M > 0$ for any $m \in (m_c, 1)$. If $m \in (m_c, 1)$, the mass of any nonnegative solution v is conserved and for simplicity, we shall assume that $\int_{\mathbb{R}^d} v(t, x) dx = M$ for any $t \geq 0$. See [54] for an interesting survey on mass conservation and the role of Barenblatt solutions in the range $(m_c, 1)$. With this notation, we can introduce the *generalized entropy* (or *free energy*) and *Fisher information* functionals respectively given by

$$\begin{aligned} \mathcal{F}[v] &:= \frac{1}{m-1} \int_{\mathbb{R}^d} (v^m - \mathcal{B}^m - m \mathcal{B}^{m-1} (v - \mathcal{B})) dx, \\ \mathcal{I}[v] &:= \frac{m}{1-m} \int_{\mathbb{R}^d} v |\nabla v^{m-1} - \nabla \mathcal{B}^{m-1}|^2 dx. \end{aligned}$$

A straightforward computation shows that, for a solution of (6), we have

$$\frac{d}{dt} \mathcal{F}[v(t, \cdot)] = -\mathcal{I}[v(t, \cdot)]. \quad (7)$$

Theorem 3 ([22]). *Assume that $m \in (m_1, 1)$ if $d = 1$ or $d = 2$, and $m \in [m_1, 1)$ if $d \geq 3$. If v is a nonnegative solution of (6) with initial datum $v_0 \in L^1(\mathbb{R}^d)$ of mass M such that $\mathcal{F}[v_0]$ is finite, then*

$$\mathcal{F}[v(t, \cdot)] \leq \mathcal{F}[v_0] e^{-4t} \quad \forall t \geq 0. \quad (8)$$

The estimate (8) is equivalent to the *entropy - entropy production* inequality

$$\mathcal{I}[v] \geq 4 \mathcal{F}[v], \quad (9)$$

which can also be rewritten as a family of *Gagliardo-Nirenberg-Sobolev inequalities*

$$\|\nabla f\|_{L^2(\mathbb{R}^d)}^\theta \|f\|_{L^{p+1}(\mathbb{R}^d)}^{1-\theta} \geq \mathcal{C}_{\text{GNS}}(p) \|f\|_{L^{2p}(\mathbb{R}^d)} \quad (10)$$

with optimal constant. The equivalence is obtained using

$$p = \frac{1}{2m-1} \iff m = \frac{p+1}{2p}$$

with $v = f^{2p}$ so that $v^m = f^{p+1}$ and $v|\nabla v^{m-1}|^2 = (p-1)^2 |\nabla f|^2$. The range $m \in (m_1, 1)$ means that $2p \in (2, 2^*)$ where $2^* = 2d/(d-2)$ if $d \geq 3$ and $2^* = +\infty$ otherwise. The case $m = m_1$, i.e., $2p = 2^*$, is also covered if $d \geq 3$.

Equality in (10) is achieved by $f_\star(x) = \mathcal{B}(x)^{\frac{1}{2p}} = (1+|x|^2)^{-\frac{1}{p-1}}$, which also means that 4 is the optimal constant in (9) and it is also the best possible decay rate in (8). The exponent $\theta = \frac{d(p-1)}{(d+2-p)(d-2)p}$ is determined by the invariance under scalings and it is easy to check that $\theta = 1$ for $p = 2^*/2$ and $d \geq 3$. If $d = 2$, we obtain the Euclidean Onofri inequality in the limit as $p \rightarrow +\infty$. On the other hand, $\lim_{p \rightarrow 1+} \mathcal{C}_{\text{GNS}}(p) = 1$ so that (10) becomes an equality for $p = 1$. By taking the derivative with respect to p at $p = 1$, we obtain the scale invariant form of the scale invariant *Euclidean logarithmic Sobolev inequality*

$$\frac{d}{2} \log \left(\frac{2}{\pi d e} \int_{\mathbb{R}^d} |\nabla f|^2 dx \right) \geq \int_{\mathbb{R}^d} |f|^2 \log |f|^2 dx \quad (11)$$

for any function $f \in H^1(\mathbb{R}^d)$ such that $\|f\|_{L^2(\mathbb{R}^d)} = 1$. Ineq. (11) can be found in [56, Theorem 2], [51, Inequality (2.3)] and [15, Inequality (26)]. See [35] for further comments and [40] for the classical Gaussian form and its Euclidean counterpart.

We obtain a *linearized free energy* and a *linearized Fisher information*

$$\mathbb{F}[w] := \frac{m}{2} \int_{\mathbb{R}^d} w^2 \mathcal{B}^{2-m} dx \quad \text{and} \quad \mathbb{I}[w] := m(1-m) \int_{\mathbb{R}^d} |\nabla w|^2 \mathcal{B} dx$$

by considering $f_\varepsilon := \mathcal{B}(1+\varepsilon \mathcal{B}^{1-m} w)$ and taking the $O(\varepsilon^2)$ in, respectively $\mathcal{F}[f_\varepsilon]$ and $\mathcal{I}[f_\varepsilon]$. The functionals \mathbb{F} and \mathbb{I} are related by a *Hardy-Poincaré inequality*.

Proposition 2 ([6, 34, 9]). *Let $m \in [m_1, 1)$ if $d \geq 3$, $m \in (1/2, 1)$ if $d = 2$, and $m \in (1/3, 1)$ if $d = 1$. If $w \in L^2(\mathbb{R}^d, \mathcal{B}^{2-m} dx)$ is such that $\nabla w \in L^2(\mathbb{R}^d, \mathcal{B} dx)$, $\int_{\mathbb{R}^d} w \mathcal{B}^{2-m} dx = 0$, then*

$$\mathbb{I}[w] \geq 4\alpha \mathbb{F}[w] \quad (12)$$

with $\alpha = 1$. Additionally, under the center of mass condition

$$\int_{\mathbb{R}^d} x w \mathcal{B}^{2-m} dx = 0, \quad (13)$$

then the inequality holds with $\alpha = 2 - d(1 - m)$.

The value of α in Proposition 2 is optimal and corresponds to the spectral gap of the linearized evolution operator \mathcal{L} associated with (6), and to an improved spectral gap under Condition (13) in the second case. In the first case, we recover that $\inf \mathcal{I}/\mathcal{F} = \inf I/F = 4$, which means that the worst decay rate is achieved in (8) in the *asymptotic regime* as $t \rightarrow +\infty$. The eigenspace of \mathcal{L} associated with $\alpha = 1$ corresponds to the equality case in (12) and determines the spectral gap. See [26] for a description of the spectrum of \mathcal{L} and [7] for an equivalence with the setting of Proposition 2. Under Condition (13), an improved entropy – entropy production inequality is obtained, at least for functions which are close enough to the manifold of the Barenblatt functions.

Corollary 1 ([34,7,9]). *Let $m \in (m_1, 1)$ if $d \geq 2$, $m \in (1/3, 1)$ if $d = 1$, $\eta = 2(dm - d + 1)$ and $\chi = m/(266 + 56m)$. If $\int_{\mathbb{R}^d} v \, dx = M$, $\int_{\mathbb{R}^d} x v \, dx = 0$ and*

$$(1 - \varepsilon) \mathcal{B} \leq v \leq (1 + \varepsilon) \mathcal{B} \quad (14)$$

for some $\varepsilon \in (0, \chi \eta)$, then

$$\mathcal{I}[v] \geq (4 + \eta) \mathcal{F}[v].$$

If v solves (6) with $m \in (m_1, 1)$ and an initial datum v_0 with centred mass, i.e., such that $\int_{\mathbb{R}^d} x v_0(x) \, dx = 0$, then $\int_{\mathbb{R}^d} x v(t, x) \, dx = 0$ for any $t \geq 0$ so that $w = (v - \mathcal{B}) \mathcal{B}^{m-2}$ satisfies (13). If additionally v_0 satisfies (14), then

$$\mathcal{F}[v(t, \cdot)] \leq \mathcal{F}[v_0] e^{-(4+\eta)t} \quad \forall t \geq 0.$$

We have obtained an improved entropy – entropy production inequality and an improved decay rate of $\mathcal{F}[v(t, \cdot)]$ by comparing the nonlinear free energy \mathcal{F} and the Fisher information \mathcal{I} with their linearized counterparts F and I , and by getting rid of the eigenspace corresponding to $\alpha = 1$ in (12). The result of Corollary 1 can be extended to $m \notin [m_1, 1)$: see [6]. Under technical conditions, better rates can be achieved in some cases as in [25]. One of the issues is that Condition (14) is rather restrictive and determines only a small neighbourhood of the set of the Barenblatt functions. Here we are going to take advantage of a property of *relative uniform convergence* that goes back to [11]: Condition (14) is satisfied by any solution $v(t, \cdot)$ for $t \geq t_*$ if the *threshold time* t_* is taken large enough. Based on results of [10,9], we have the following result.

Theorem 4 ([9,8]). *Assume that $m \in (m_1, 1)$ if $d \geq 2$, $m \in (1/3, 1)$ if $d = 1$ and let $\varepsilon \in (0, 1/2)$, small enough, and $A > 0$ be given. There exists an explicit threshold time $t_* \geq 0$ such that, if v is a solution of (6) with nonnegative initial datum $v_0 \in L^1(\mathbb{R}^d)$ satisfying $\int_{\mathbb{R}^d} v_0 \, dx = \int_{\mathbb{R}^d} \mathcal{B} \, dx = M$ and*

$$A[v_0] = \sup_{r>0} r^{\frac{d(m-m_c)}{(1-m)}} \int_{|x|>r} v_0 \, dx \leq A < \infty, \quad (15)$$

and if $B(t, \cdot)$ solves (5) with initial datum \mathcal{B} , then

$$\sup_{x \in \mathbb{R}^d} \left| \frac{v(t, x)}{\mathcal{B}(x)} - 1 \right| \leq \varepsilon \quad \forall t \geq t_*.$$

Using the self-similar change of variables that transforms (5) into (6), we obtain that a solution of (6) satisfies Condition (14) for $t > t_*$ for some threshold time $t_* > 0$ which explicitly depends on T . See [9] for details. Under Condition (15), we have an improved decay rate for any t in the *asymptotic time range* $[t_*, +\infty)$.

The next question is to find an improvement in the *initial time layer* $[0, t_*)$ as well. A key ingredient is the *carré du champ* method for which we can refer to [9, Chapter 2]. We learn for instance that for some strictly convex function ψ with $\psi(0) = 0$, $\psi'(0) = 0$, there is an improved entropy – entropy production inequality

$$\mathcal{I} - 4\mathcal{F} \geq \psi(\mathcal{F}) \geq 0,$$

so that better decay rates can be expected if \mathcal{F} is large, which is precisely what can be expected in the initial time layer. In fact, a key estimate of the *carré du champ* method is the simpler observation that

$$\frac{d}{dt}\mathcal{I}[v(t, \cdot)] \leq -4\mathcal{I}[v(t, \cdot)]$$

if v solves (6), which allows us to prove a *backward in time* estimate. Let us define $\mathcal{Q}[v] := \frac{\mathcal{I}[v]}{\mathcal{F}[v]}$ and notice that, using (7), we have

$$\frac{d}{dt}\mathcal{Q}[v(t, \cdot)] \leq \mathcal{Q}[v(t, \cdot)] (\mathcal{Q}[v(t, \cdot)] - 4).$$

Lemma 1 ([9]). *Assume that $m > m_1$ and v is a solution to (6) with nonnegative initial datum v_0 . If for some $\eta > 0$ and $t_* \geq 0$, we have $\mathcal{Q}[v(t_*, \cdot)] \geq 4 + \eta$, then*

$$\mathcal{Q}[v(t, \cdot)] \geq 4 + \zeta \quad \forall t \in [0, t_*] \quad \text{with} \quad \zeta := \frac{4\eta e^{-4t_*}}{4 + \eta - \eta e^{-4t_*}}.$$

In the subcritical range with $m > m_1$, the first consequence of Lemma 1 is an improved decay rate of the free energy for the solution of (6).

Corollary 2 ([9]). *Let $m \in (m_1, 1)$ if $d \geq 2$, $m \in (1/2, 1)$ if $d = 1$, and $A > 0$. If v is a solution of (6) with nonnegative initial datum $v_0 \in L^1(\mathbb{R}^d)$ such that $\int_{\mathbb{R}^d} v_0 dx = M$, $\int_{\mathbb{R}^d} x v_0 dx = 0$ and v_0 satisfies (15), then there is an explicit $\zeta > 0$ such that*

$$\mathcal{F}[v(t, \cdot)] \leq \mathcal{F}[v_0] e^{-(4+\zeta)t} \quad \forall t \geq 0.$$

The constant ζ can be explicitly computed in terms of d , m and A using Corollary 1, Theorem 4, and Lemma 1. The second consequence of Lemma 1 is a *stability result for the entropy – entropy production inequality*, which reads as

$$\mathcal{I}[v] - 4\mathcal{F}[v] \geq \frac{\zeta}{4 + \zeta} \mathcal{I}[v]$$

where the left-hand side is the deficit in (9), while the distance in the right-hand side is measured by the Fisher information, which is a quantity that controls strong norms.

In the critical case $m = m_1$ corresponding to the Sobolev inequality for $d \geq 3$, we also obtain a *constructive stability result*. Let $\mathbf{g}(x) := (1 + |x|^2)^{-(d-2)/2}$ be the Aubin–Talenti function such that $\mathbf{g}^{2^*} = \mathcal{B}$.

Theorem 5 ([9]). *Let $d \geq 3$, $A > 0$, and $p = 2^*/2$. Then for any nonnegative $f \in L^{p+1}(\mathbb{R}^d)$ such that $\nabla f \in L^2(\mathbb{R}^d)$ and $|x| f^p \in L^2(\mathbb{R}^d)$,*

$$\int_{\mathbb{R}^d} (1, x, |x|^2) f^{2^*} dx = \int_{\mathbb{R}^d} (1, x, |x|^2) \mathbf{g} dx \quad \text{and} \quad \sup_{r>0} r^d \int_{|x|>r} f^{2^*} dx \leq A,$$

then we have

$$\|\nabla f\|_{L^2(\mathbb{R}^d)}^2 - S_d^2 \|f\|_{L^{2^*}(\mathbb{R}^d)}^2 \geq \mathcal{C} \int_{\mathbb{R}^d} |(p-1)\nabla f + f^p \nabla \mathbf{g}^{1-p}|^2 dx$$

for some explicit $\mathcal{C} > 0$ which depends only on d and A .

Extending the subcritical result of Corollary 2 to the critical case uses the introduction of a time delay, as shown in [9, Chapter 6]. One can also obtain a stability result in the subcritical range $m \in (m_1, 1)$, that is, $1 < p < 2^*/2$ for Ineq. (10), but the statement is more delicate as several scales are involved. A remarkable feature in Theorem 5 is that we measure the distance to a specific Aubin–Talenti function and not to the manifold \mathcal{M} .

Summing up, we have obtained a construction of an explicit stability result using the *fast diffusion equation* and *entropy methods* based on a parabolic regularity theory made explicit, spectral estimates and various features of the *carré du champ* method. However, estimates of the stability constants are way smaller than the expected values of the optimal stability constants and are available only under the restriction (15) on the tails of the functions, which is inherent to the method, as one can prove that t_* diverges if $A \rightarrow +\infty$ in (15).

3 Explicit stability results for Sobolev and log-Sobolev inequalities, with optimal dimensional dependence

In this section, we just state the main results and refer the reader to [28] for detailed results and to [29] for a shorter introduction to the main tools. The most important fact is that the method gives an explicit value to the constant in the *stability* inequality (2), but again one can suspect that the lower bound is not very good, even if it has an optimal dimensional dependence.

Theorem 6 ([28]). *There is a constant $\beta > 0$ with an explicit lower estimate which does not depend on d such that for all $d \geq 3$ and all $f \in H^1(\mathbb{R}^d)$ we have*

$$\|\nabla f\|_{L^2(\mathbb{R}^d)}^2 - S_d \|f\|_{L^{2^*}(\mathbb{R}^d)}^2 \geq \frac{\beta}{d} \inf_{g \in \mathcal{M}} \|\nabla f - \nabla g\|_{L^2(\mathbb{R}^d)}^2. \quad (16)$$

Using the inverse stereographic projection we can rewrite the result on \mathbb{S}^d as

$$\begin{aligned} & \|\nabla F\|_{L^2(\mathbb{S}^d)}^2 - \frac{1}{4} d(d-2) \left(\|F\|_{L^{2^*}(\mathbb{S}^d)}^2 - \|F\|_{L^2(\mathbb{S}^d)}^2 \right) \\ & \geq \frac{\beta}{d} \inf_{G \in \mathcal{M}(\mathbb{S}^d)} \left(\|\nabla F - \nabla G\|_{L^2(\mathbb{S}^d)}^2 + \frac{1}{4} d(d-2) \|F - G\|_{L^2(\mathbb{S}^d)}^2 \right). \end{aligned}$$

Here we take on \mathbb{S}^d the uniform probability measure. Let $v \in H^1(\mathbb{R}^n, dx)$ with compact support, $d \geq n$ large enough and consider

$$u_d(\omega) = v\left(\frac{1}{r_d} \omega_1, \frac{1}{r_d} \omega_2, \dots, \frac{1}{r_d} \omega_n\right), \quad r_d = \sqrt{\frac{d}{2\pi}}$$

where $\omega = (\omega_1, \omega_2, \dots, \omega_{d+1}) \in \mathbb{S}^d \subset \mathbb{R}^{d+1}$. If $d\gamma = e^{-\pi|x|^2} dx$ denotes the Gaussian measure on \mathbb{R}^n , by taking the limit as $d \rightarrow +\infty$, we obtain

$$\begin{aligned} & \lim_{d \rightarrow +\infty} d \left(\|\nabla u_d\|_{L^2(\mathbb{S}^d)}^2 - \frac{1}{4} d(d-2) \left(\|u_d\|_{L^2(\mathbb{S}^d)}^2 - \|u_d\|_{L^p(\mathbb{S}^d)}^2 \right) \right) \\ & = \|\nabla u\|_{L^2(\mathbb{R}^n, d\gamma)}^2 - \pi \int_{\mathbb{R}^n} |u|^2 \log \left(\frac{|u|^2}{\|u\|_{L^2(\mathbb{R}^n, d\gamma)}^2} \right) d\gamma, \end{aligned}$$

where the right-hand side is the deficit in the *Gaussian logarithmic Sobolev inequality* of [40], which is equivalent to (11). We can also let $d \rightarrow +\infty$ in (16). The delicate part is to bound the coefficients (a, b, c) of the Aubin–Talenti functions which minimize the distance in (2) uniformly with respect to d . Because of the rescaling, the gradient term in the distance is lost and stability is measured only by a $L^2(\mathbb{R}^n, d\gamma)$ norm. See [28] for a detailed proof.

Corollary 3 ([28]). *With $\beta > 0$ as in Theorem 6, we have*

$$\begin{aligned} & \|\nabla u\|_{L^2(\mathbb{R}^n, d\gamma)}^2 - \pi \int_{\mathbb{R}^n} |u|^2 \log \left(\frac{|u|^2}{\|u\|_{L^2(\mathbb{R}^n, d\gamma)}^2} \right) d\gamma \\ & \geq \frac{\beta \pi}{2} \inf_{a \in \mathbb{R}^d, c \in \mathbb{R}} \int_{\mathbb{R}^n} |u - c e^{a \cdot x}|^2 d\gamma. \end{aligned}$$

The proof of Theorem 6 relies on three main steps.

1. A sequence of nonnegative functions built using the competing symmetries of E. Carlen and M. Loss in [18], which alternate a Schwarz symmetric decreasing rearrangement and a transformation which corresponds, up to the stereographic projection, to a rotation on the sphere. This step reduces the stability problem to a local problem, on a neighbourhood of \mathcal{M} . We also use a continuous Steiner symmetrization in the discussion of one of the cases and an argument of [21].
2. A Taylor expansion with explicitly controlled remainder terms. To obtain the d -dependence, one has to perform a delicate analysis which involves an argument of [36].

3. The result for sign-changing functions is deduced from the result for the positive and negative parts by a convexity estimate.

For full details on the proof of Theorem 6, one has to refer to [28]. How much is lost in the first two steps is unclear, and this is why a proof based, *e.g.*, on a fast diffusion equation would be a significant added value. There are several results which complement Theorem 6 and Corollary 3. We know from [46] that $\beta < 4d/(d+4)$. The equality case in (2) written for the optimal value of κ_d is achieved according to [47]. A stability result for the Hardy-Littlewood-Sobolev inequality follows from the duality method of [16] and has been generalized in [20] to cover various exponents. This result also proves an explicit stability estimate for the fractional Sobolev inequality, which uses duality but otherwise follows the very same lines of proof as in [28].

The case of the logarithmic Hardy-Littlewood-Sobolev inequality is also obtained in [17] using duality with the Onofri inequality, whose explicit stability has been obtained in [41] by other methods.

4 More explicit stability results for the logarithmic Sobolev and Gagliardo-Nirenberg inequalities on \mathbb{S}^d

Here the goal is to reconcile the flow methods of Section 2 with stability estimates without restrictions as in Section 3 and get estimate which may have the right order of magnitude compared to the optimal stability constants. Let us consider the family of *subcritical Gagliardo-Nirenberg inequalities*

$$\|\nabla F\|_{L^2(\mathbb{S}^d)}^2 \geq d \mathcal{E}_p[F] := \frac{d}{p-2} \left(\|F\|_{L^p(\mathbb{S}^d)}^2 - \|F\|_{L^2(\mathbb{S}^d)}^2 \right) \quad \forall F \in H^1(\mathbb{S}^d, d\mu) \quad (17)$$

for any $p \in [1, 2) \cup (2, 2^*)$, with $2^* := \frac{2d}{d-2}$ if $d \geq 3$ and $2^* = +\infty$ if $d = 1$ or 2 . If $d \geq 3$, the limit case $p = 2^*$ of (17) is the Sobolev inequality on the sphere. As another limit case when $p \rightarrow 2$, we obtain the *logarithmic Sobolev inequality on the sphere*

$$\int_{\mathbb{S}^d} |\nabla F|^2 d\mu \geq d \mathcal{E}_2[F] := \frac{d}{2} \int_{\mathbb{S}^d} F^2 \log \left(\frac{F^2}{\|F\|_{L^2(\mathbb{S}^d)}^2} \right) d\mu \quad \forall F \in H^1(\mathbb{S}^d, d\mu) \quad (18)$$

where $d\mu$ is the uniform probability on \mathbb{S}^d . These inequalities were obtained in [1, 2, 5]. While equality is always achieved by constant functions, the optimal constant d can be identified with the first positive eigenvalue of the Laplace-Beltrami operator on \mathbb{S}^d . Let X_1 be the corresponding eigenspace of spherical harmonic functions and denote by Π_1 the orthogonal projection onto X_1 . If $\varphi \in X_1$ and $F_\varepsilon = 1 + \varepsilon \varphi$, a simple Taylor expansion shows that $\|\nabla F_\varepsilon\|_{L^2(\mathbb{S}^d)}^2 - d \mathcal{E}_p[F_\varepsilon] = O(\varepsilon^4)$, so that there is no hope to control the deficit from below by the square of a distance based on usual norm, at least in the direction of X_1 .

Theorem 7 ([12]). *Let $d \geq 1$ and $p \in (1, 2^*)$. There is an explicit stability constant $\mathcal{S}_{d,p} > 0$ such that, for any $F \in H^1(\mathbb{S}^d, d\mu)$, we have*

$$\begin{aligned} & \int_{\mathbb{S}^d} |\nabla F|^2 d\mu - d\mathcal{E}_p[F] \\ & \geq \mathcal{S}_{d,p} \left(\frac{\|\nabla \Pi_1 F\|_{L^2(\mathbb{S}^d)}^4}{\|\nabla F\|_{L^2(\mathbb{S}^d)}^2 + \|F\|_{L^2(\mathbb{S}^d)}^2} + \|\nabla(\text{Id} - \Pi_1)F\|_{L^2(\mathbb{S}^d)}^2 \right). \end{aligned}$$

This result is a consequence of the stability analysis of [38], of explicit improvements of (17) and (18) obtained for instance in [30,32] based on [24,31], which involve a generalized *carré du champ* method for the flow

$$\frac{\partial u}{\partial t} = u^{-p(1-m)} \left(\Delta u + (mp - 1) \frac{|\nabla u|^2}{u} \right)$$

with exponents $m \in [m_-(d,p), m_+(d,p)]$ where

$$m_{\pm}(d,p) := \frac{1}{(d+2)p} \left(dp + 2 \pm \sqrt{d(p-1)(2d - (d-2)p)} \right),$$

and of spectral properties taken from the proof of [3]. As in Section 3 and with the same notation, we can consider the large dimensional limit.

Theorem 8 ([13]). *Let $n \geq 1$ and $u \in H^1(\mathbb{R}^n, dx)$ with compact support. For any $p \in (1, 2)$, we have*

$$\begin{aligned} & \lim_{d \rightarrow +\infty} d \left(\|\nabla u_d\|_{L^2(\mathbb{S}^d)}^2 - \frac{d}{2-p} \left(\|u_d\|_{L^2(\mathbb{S}^d)}^2 - \|u_d\|_{L^p(\mathbb{S}^d)}^2 \right) \right) \\ & = \|\nabla u\|_{L^2(\mathbb{R}^n, d\gamma)}^2 - \frac{1}{2-p} \left(\|u\|_{L^2(\mathbb{R}^n, d\gamma)}^2 - \|u\|_{L^p(\mathbb{R}^n, d\gamma)}^2 \right). \end{aligned}$$

In Theorem 7 we have the satisfaction to obtain a stability result which distinguishes the X_1 eigenspace, is based on a fast diffusion equation and does not involve any restriction on the function space, but the estimate is not uniform as $p \rightarrow 2^*$, as could be expected from Theorem 6.

There is still much more to understand, particularly concerning (18). In Corollary 3, the stability of the *Gaussian logarithmic Sobolev inequality* is measured only by a $L^2(\mathbb{R}^n, d\gamma)$ norm. It is known from [44,43] that an $H^1(\mathbb{R}^n, d\gamma)$ stability result cannot be expected without additional conditions. The quest for optimal conditions is open but as an illustration, we can conclude with a last result.

Theorem 9 ([14]). *Let $d \geq 1$. For any $\varepsilon > 0$, there is some explicit $\mathcal{C} > 1$ depending only on ε such that, for any $u \in H^1(\mathbb{R}^d, d\gamma)$ with*

$$\int_{\mathbb{R}^d} (1, x) |u|^2 d\gamma = (1, 0), \quad \int_{\mathbb{R}^d} |u|^2 e^{\varepsilon|x|^2} d\gamma < \infty$$

for some $\varepsilon > 0$, then we have

$$\|\nabla u\|_{L^2(\mathbb{R}^d, d\gamma)}^2 \geq \frac{\mathcal{C}}{2} \int_{\mathbb{R}^d} |u|^2 \log |u|^2 d\gamma.$$

Moreover, if u is compactly supported in a ball of radius $R > 0$, we know from [48] that with $\mathcal{C} = 1 + \frac{\mathcal{C}_*(K_*)-1}{1+R^2 \mathcal{C}_*(K_*)}$, $\mathcal{C}_*(K) = 1 + \frac{1}{432K}$ and $K_* := \max(d, \frac{(d+1)R^2}{1+R^2})$.

Here we use the heat flow case corresponding to the limit $m = 1$ of the fast diffusion flow. The proof is based on a stability result of [37] for a restricted class of functions, which contains the solution of the heat flow after an initial time layer recently determined in [19]. For a compactly supported initial datum, we obtain a very explicit expression using a result of [48]. One can certainly object that $1/432 \approx 0.00231481$ does not provide us with a glorious improvement of the constant, but this is definitely a step in the understanding of *explicit* stability results and, as a consequence, better estimates of decay rates in nonlinear evolution problems.

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