

On extended model of Josephson junction, linear systems with polynomial solutions, determinantal surfaces and Painlevé 3 equations

Alexey Glutsyuk

To cite this version:

Alexey Glutsyuk. On extended model of Josephson junction, linear systems with polynomial solutions, determinantal surfaces and Painlevé 3 equations. Proceedings of the Steklov Institute of Mathematics, In press, $326.$ hal- 04786763

HAL Id: hal-04786763 <https://hal.science/hal-04786763v1>

Submitted on 16 Nov 2024

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

On extended model of Josephson junction, linear systems with polynomial solutions, determinantal surfaces and Painlevé 3 equations

Alexey Glutsyuk∗†‡

June 13, 2024

To dear Victor Matveevich Buchstaber on the occasion of his 80-th birthday

Abstract

We consider a three-parameter family of linear special double confluent Heun equations introduced and studied by V.M.Buchstaber and S.I.Tertychnyi, which is an equivalent presentation of a model of Josephson junction in superconductivity. Buchstaber and Tertychnyi have shown that the set of those complex parameters for which the Heun equation has a polynomial solution is a union of the so-called spectral curves: explicit algebraic curves in \mathbb{C}^2 indexed by $\ell \in \mathbb{N}$. In his paper with I.V.Netay, the author have shown that each spectral curve is irreducible in Heun equation parameters (consists of two irreducible components in parameters of Josephson junction model). Netay discovered numerically and conjectured a genus formula for spectral curves. He reduced it to the conjecture stating that each of them is regular in \mathbb{C}^2 with a coordinate axis deleted. Here we prove Netay's regularity and genus conjectures. For the proof we study a four-parameter family of linear systems on the Riemann sphere extending a family of linear systems equivalent to the Heun equations. They yield an equivalent presentation of the extension of model of Josephson junction introduced by the author in his paper with Yu.P.Bibilo. We describe the so-called determinantal surfaces, which consist of linear systems with polynomial solutions, as explicit affine algebraic hypersurfaces in \mathbb{C}^3 indexed by $\ell \in \mathbb{N}$. The spectral curves are their intersections with the hyperplane corresponding to the initial model. We prove that each determinantal surface is regular outside appropriate hyperplane and consists of two rational irreducible components. The proofs use Stokes phenomena theory, holomorphic vector bundle technique, the fact that each determinantal surface is foliated by isomonodromic families of linear systems (this foliation is governed by Painlevé 3 equation) and transversality of the latter foliation to the initial model.

[∗]CNRS, UMR 5669 (UMPA, ENS de Lyon), France. E-mail: aglutsyu@ens-lyon.fr †HSE University, Moscow, Russia

[‡]Kharkevich Institute for Information Transmission Problems (IITP, RAS), Moscow

Contents

1 Introduction

1.1 Introduction, brief description of main results and of their proofs

The three-parameter family of special double confluent Heun equations

$$
z^{2}E'' + ((-\ell + 1)z + \mu(1 - z^{2}))E' + (\lambda + \mu(\ell - 1)z)E = 0
$$
\n(1.1)

belongs to the well-known class of Heun equations, see [54]. It was studied by V. M. Buchstaber and S. I. Tertychnyi in [57, 18, 19, 20, 21, 22] and in joint papers [12, 13] by Buchstaber and the author. Buchstaber and Tertychnyi have shown in [18, 19] that its restriction to real parameters satisfying the inequality $\lambda + \mu^2 > 0$ is equivalent to a model of overdamped Josephson junction in superconductivity; see the corresponding material in Subsection 1.4. In [18] they have described those complex parameters (ℓ, λ, μ) with $\mu \neq 0$ for which equation (1.1) has a polynomial solution: this holds exactly, when $\ell \in \mathbb{N}$ and the point (λ, μ) lies on an algebraic curve $\Gamma_{\ell} \subset \mathbb{C}_{\lambda,\mu}^2$ called the *spectral curve*. Namely, Γ_{ℓ} is the zero locus of the determinant of a remarkable 3-diagonal $\ell \times \ell$ -matrix, see [18, formula (21)] and Theorem 1.2 below. It was shown by the author in [29, theorem 1.3] that each spectral curve Γ_{ℓ} is irreducible.

One of the main result (joint with I.V.Netay) of the present paper is the formula for genus of the spectral curve, see Subsection 1.2. Netay discovered it numerically, stated it as a conjecture, see [29, conjecture 1.5] and reduced it to the conjecture saying that the curve

$$
\Gamma_\ell^o := \Gamma_\ell \setminus \{\mu = 0\} \subset \mathbb{C}^2
$$

is regular, see [29, corollary 3.6]. Here we prove Netay's smoothness and hence, genus conjectures.

Family of Heun equations (1.1) is equivalent to the family of linear systems

$$
Y' = \left(\frac{1}{z^2} \operatorname{diag}(-\mu, 0) + \frac{1}{z} \begin{pmatrix} \ell & -\frac{a}{2} \\ \frac{a}{2} & 0 \end{pmatrix} + \operatorname{diag}(0, \mu)\right) Y, \tag{1.2}
$$

$$
\lambda = \frac{a^2}{4} - \mu^2,\tag{1.3}
$$

in the following sense: for every solution $Y(z) = (Y_1(z), Y_2(z))$ of system (1.2) the function

 $E(z) = Y_2(z)$

satisfies (1.1) and conversely: for each solution $E(z)$ of (1.1) the vector function

$$
Y = (Y_1, Y_2), \quad Y_2(z) = E(z), \quad Y_1(z) = \frac{2z}{a}(E'(z) - \mu E(z))
$$
\n(1.4)

is a solution of (1.2). See an implicit equivalent statement in [12, subsection 3.2, p. 3869].

Remark 1.1 A function $E(z)$ is a polynomial solution of double confluent Heun equation (1.1), if and only if the corresponding vector function $Y(z)$ given by (1.4) is a vector polynomial solution of system (1.2).

For the proof of the genus and smoothness conjectures we consider a four-parameter extension of the family of linear systems (1.2) depending on an additional parameter χ : the family

$$
Y' = \left(\frac{1}{z^2} \begin{pmatrix} -\frac{s}{2} & -s\chi \\ 0 & 0 \end{pmatrix} + \frac{1}{z} \begin{pmatrix} \ell - \chi a & -\frac{a}{2} \\ \frac{a}{2} & \chi a \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ s\chi & \frac{s}{2} \end{pmatrix} \right) Y. \tag{1.5}
$$

The initial family (1.2) is the subfamily of (1.5) with $\chi = 0$, written in the rescaled parameter $\mu := \frac{s}{2}$. The restriction of family (1.5) to real parameters is equivalent to the four-dimensional extension of model of Josephson junction introduced by the author in [8, subsection 6.2]. See the corresponding material in Subsection 1.5.

Our next series of main results, presented in Subsection 1.3, concern the set of those complex parameter values in (1.5) with $s \neq 0$ for which system (1.5) has a vector polynomial solution. We show that this is a union indexed by $\ell \in \mathbb{N}$ of affine algebraic surfaces $S_{\ell} \subset$ $\mathbb{C}^3_{\chi,a,s}$, called *determinantal surfaces*; each S_{ℓ} consists of two irreducible components $S_{\ell,\pm}$ permuted by the involution $(\chi, a, s) \mapsto (-\chi, -a, s)$. Each $S_{\ell, \pm}$ will be presented as the zero locus of an explicit polynomial, the determinant of an $(\ell + 1) \times (\ell + 1)$ -matrix. We show that the surface $S_{\ell, \pm}$ is rational and the complement

$$
S_{\ell,\pm}^o:=S_{\ell,\pm}\setminus\{s=0\}
$$

is a regular surface biholomorphically parametrized by the complement of $\mathbb{C}^* \times \overline{\mathbb{C}}$ to an analytic hypersurface (curve, which a priori may be neither irreducible, nor algebraic), and the surfaces $S_{\ell,\pm}^o$ are disjoint. The latter regularity and parametrization statements are proved in Subsections 2.8, 2.9 by using Stokes phenomena theory and holomorphic vector bundle technique analogous to the one used by A.A.Bolibruch in his famous works on the Riemann–Hilbert Problem and related topics: see [9]–[11] and references therein.

Using the above results on the determinantal surfaces S_{ℓ} we prove Netay's smoothness and genus conjectures on the spectral curves Γ_{ℓ} in Subsection 2.11 as follows. The curve Γ_ℓ^o written in the coordinates (a, s) instead of (λ, μ) , $s = 2\mu$, $\lambda = \frac{a^2 - s^2}{4}$ $\frac{-s^2}{4}$, see (1.3) , is known to consist of two irreducible components $\Gamma_{\ell,\pm}^o$ that are permuted by the involution $(a, s) \mapsto (-a, s)$ and are bijectively projected onto Γ_ℓ^o via the projection $(a, s) \mapsto (\lambda, \mu)$, as was shown in [29, section 3] using results of [29] and [18, section 3]. Smoothness of Γ_ℓ^o is equivalent to smoothness of any of the curves $\Gamma^o_{\ell, \pm}$. One has

$$
\Gamma_{\ell,\pm}^o=S_{\ell,\pm}^o\cap\{\chi=0\}.
$$

Thus, to prove smoothness of the curve $\Gamma^o_{\ell, \pm}$, and hence, the genus conjecture, it suffices to prove transversality of the (smooth) surface $S_{\ell,\pm}^o$ and the hyperplane $\{\chi=0\}$.

To prove transversality, we consider the line field on the parameter space introduced by the author in [8, subsection 6.2]:

$$
\begin{cases} \ell_s' = 0\\ \chi_s' = \frac{a - 2\chi(2\chi a - \ell)}{2s} \\ a_s' = -2s\chi + \frac{a}{s}(2\chi a - \ell) \end{cases} \tag{1.6}
$$

It was shown by the author [8, theorem 6.6] that each its integral curve (graph of solution of (1.6)) corresponds to an isomonodromic family of linear systems¹ (1.5) , and for every

¹In fact, this was proved in $[8]$ for equivalent linear systems, which are obtained from (1.5) by the space variable change $\widehat{Y} = e^{-\frac{s}{2}z}Y$.

 $\ell \in \mathbb{C}$ the function

$$
w(s) := \frac{a(s)}{2s\chi(s)} = \frac{a(s)}{\psi(s)}
$$
(1.7)

satisfies Painlevé 3 equation²

$$
w'' = \frac{(w')^2}{w} - \frac{w'}{s} + 2\ell \frac{w^2}{s} - (2\ell + 2)\frac{1}{s} + w^3 - \frac{1}{w}.
$$
\n(1.8)

along solutions of (1.6). As was shown in loc. cit. the above-mentioned isomonodromic families are induced from Jimbo isomonodromic deformations given in [35]. We show that isomonodromic deformation preserves the property of linear system (1.5) to have a polynomial solution, and hence, each determinantal surface $S_{\ell,\pm}^o$ is tangent to the line field on $\mathbb{C}^3_{\chi,a,s}$ induced by (1.6). The hyperplane $\{\chi=0\}$ is clearly transversal to line field (1.6) at its points where $s, a \neq 0$. We show that the surfaces $S^o_{\ell, \pm}$ do not intersect the s-axis $\{\chi = a = 0\}$. This together with their tangency to field (1.6) implies transversality, and hence, smoothness of the intersections $\Gamma^o_{\ell, \pm} = S^o_{\ell, \pm} \cap {\chi = 0}$. This will finish the proof of smoothness and genus conjectures.

The foliation of each surface $S^o_{\ell, \pm}$ by integral curves of the line field induced by (1.6) is called its *isomonodromic foliation*. The isomonodromic foliations of the surfaces $S^o_{\ell, \pm}$ with $\ell = 1, 2$ (and also an analogous foliation for $\ell = 0$) are discussed in Section 3. For $\ell = 0, 1$ we show that they are equivalently described by explicit Riccatti equations.

Plan of proof of main results is presented in Subsection 1.6.

Some open problems are presented in Subsection 1.7.

1.2 Genus and smoothness of the spectral curves

Let us recall the description of the parameters corresponding to equations (1.1) with polynomial solutions. To do this, consider the three-diagonal $\ell \times \ell$ -matrix

$$
H_{\ell} = \begin{pmatrix} 0 & \mu & 0 & 0 & 0 & 0 & \dots & 0 \\ \mu(\ell-1) & 1-\ell & 2\mu & 0 & 0 & \dots & 0 \\ 0 & \mu(\ell-2) & -2(\ell-2) & 3\mu & 0 & \dots & 0 \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ 0 & \dots & 0 & 0 & 2\mu & -2(\ell-2) & (\ell-1)\mu \\ 0 & \dots & 0 & 0 & 0 & \mu & 1-\ell \end{pmatrix} : (1.9)
$$

$$
H_{\ell;ij} = 0 \text{ if } |i - j| \geq 2; \quad H_{\ell;jj} = (1 - j)(\ell - j + 1);
$$

$$
H_{\ell; j, j+1} = \mu j; \quad H_{\ell; j, j-1} = \mu(\ell - j + 1).
$$

The matrix H_{ℓ} belongs to the class of the so called Jacobi matrices that arise in different questions of mathematics and mathematical physics [31].

²In system (1.6) and in Painlevé 3 equation (1.8) presented in [8] the parameter ℓ is changed to $-\ell$.

Theorem 1.2 [18, section 3] A special double confluent Heun equation (1.1) with $\mu \neq 0$ has a polynomial solution, if and only if $\ell \in \mathbb{N}$ and the three-diagonal matrix $H_{\ell} + \lambda \operatorname{Id}$ has zero determinant. For every $\ell \in \mathbb{N}$ the determinant $\det(H_{\ell} + \lambda \operatorname{Id})$ is a polynomial in $(u, v) = (\lambda, \mu^2)$ of degree ℓ (which also has degree ℓ in (λ, μ)):

$$
Q_{\ell}(\lambda, \mu^2) := \det(H_{\ell} + \lambda \operatorname{Id}).\tag{1.10}
$$

See also [13, remark 4.13] for non-existence of polynomial solutions for $\ell \notin \mathbb{N}$ and $s \neq 0$.

Definition 1.3 For every $\ell \in \mathbb{N}$ the ℓ -th spectral curve is

$$
\Gamma_{\ell} := \{ Q_{\ell}(\lambda, \mu^2) = 0 \} \subset \mathbb{C}^2_{\lambda, \mu}.
$$
\n(1.11)

The following results were proved by the author in [29].

Theorem 1.4 [29, theorems 1.2 and 1.3] For every $\ell \in \mathbb{N}$ the polynomial $Q_{\ell}(u, v)$ from (1.10) is irreducible, and the spectral curve Γ_{ℓ} is irreducible.

Recall that the geometric genus of an irreducible algebraic curve is the genus of the Riemann surface parametrizing it bijectively (except for possible self-intersections), i. e., the genus of its normalization.

One of the main results of the paper is the following theorem, which is a joint result of the author with I.V.Netay.

Theorem 1.5 The geometric genus of the curve Γ_{ℓ} equals

$$
\begin{cases} \left(\frac{\ell-2}{2}\right)^2, & \ell \text{ even;} \\ \\ \left(\frac{\ell-1}{2}\right)\left(\frac{\ell-3}{2}\right), & \ell \text{ odd.} \end{cases}
$$

I.V.Netay has discovered the above genus formula numerically and stated it as a conjecture, see [29, conjecture 1.5]. He reduced it to the conjecture stating that the curve

$$
\Gamma_\ell^o:=\Gamma_\ell\setminus\{\mu=0\}\subset\mathbb{C}^2
$$

is regular, see [29, corollary 3.6]. Here we prove this regularity conjecture, which together with [29, corollary 3.6] implies the genus formula:

Theorem 1.6 The curve Γ_{ℓ}^{o} is regular for every $\ell \in \mathbb{N}$.

Remark 1.7 Equation (1.10) defining the curve Γ_ℓ belongs to a remarkable class of determinantal representations of plane curves: equations

$$
\det(x_1L_1 + x_2L_2 + x_3L_3) = 0,
$$

where $(x_1 : x_2 : x_3) \in \mathbb{CP}^2$ and L_1, L_2, L_3 are $\ell \times \ell$ -matrices. Determinantal representations of curves arise in algebraic geometry and integrable systems, see [58, 59] and references therein. Complete description of determinantal representations of smooth complex irreducible projective curves was obtained in [58]. Self-adjoint determinantal representations of real smooth plane curves were described in [59].

1.3 The determinantal surfaces: formula, smoothness and parametrization

Definition 1.8 The *degree* of a polynomial vector function is the maximal degree of its component.

Proposition 1.9 Let $s \neq 0$. If a vector polynomial solution of system (1.5) exists, then $\ell \in \mathbb{Z}_{\geq 0}$, the polynomial solution has degree ℓ and is unique up to constant factor.

Proposition 1.10 Let $\ell = 0$ and $(a, s) \neq (0, 0)$. A system (1.5) has a polynomial solution, if and only if $\chi = \pm \frac{1}{2}$: then the solution is a constant vector $(1, \mp 1)$ up to constant factor.

Proof If a polynomial solution exists, then it is constant (Proposition 1.9). For $(a, s) \neq$ $(0,0), \ell = 0$, the intersection of kernels of the matrices in (1.5) is the subspace of those vectors (v_1, v_2) that satisfy the system of linear equations $\frac{1}{2}v_1 + \chi v_2 = 0$, $\chi v_1 + \frac{1}{2}$ $\frac{1}{2}v_2 = 0.$ The latter system has a nonzero solution, if and only if its determinant $\frac{1}{4} - \chi^2$ vanishes, i.e., $\chi = \pm \frac{1}{2}$, and then the solution is $(1, \pm 1)$ up to constant factor.

Remark 1.11 Let $s = 0$. Then system (1.5) is Fuchsian and has two singular points: zero and infinity. It has a polynomial solution, if and only if its residue matrix has at least one integer non-negative eigenvalue. In the case, when the residue matrix is diagonalizable and both its eigenvalues are integer non-negative, the system has two linearly independent polynomial solutions.

For every $\ell \in \mathbb{N}$ let us introduce the following $(\ell + 1) \times (\ell + 1)$ -matrix functions:

$$
G_{1,\ell} = \begin{pmatrix} \frac{s}{2} & \chi a - 1 & 0 & \dots & 0 & 0 \\ 0 & \frac{s}{2} & \chi a - 2 & \dots & 0 & 0 \\ \dots & \dots & \dots & \dots & \dots & \dots \\ 0 & 0 & 0 & \dots & \frac{s}{2} & \chi a - \ell \\ 0 & 0 & 0 & \dots & 0 & \frac{1}{2} \end{pmatrix}, \qquad (1.12)
$$

$$
G_{2,\ell} = \begin{pmatrix} 0 & 0 & \dots & 0 & \frac{a}{2} & \chi s \\ 0 & 0 & \dots & \frac{a}{2} & \chi s & 0 \\ \dots & \dots & \dots & \dots & \dots \\ \frac{a}{2} & \chi s & \dots & 0 & 0 & 0 \\ \chi & 0 & \dots & 0 & 0 & 0 \end{pmatrix}, \qquad (1.13)
$$

The next three theorems and corollary are main results of the paper.

Theorem 1.12 System (1.5) with $\ell \in \mathbb{N}$ and $s \neq 0$ has a vector polynomial solution, if and only if

 \mathcal{P}_{ℓ} + $(\chi, a, s) := \det(G_{1,\ell} \pm G_{2,\ell}) = 0$ for some choice of sign \pm .

Theorem 1.13 1) The above determinants $\mathcal{P}_{\ell, \pm}$ are irreducible polynomials in (χ, a, s) that are obtained one from the other by sign change of two variables: $(\chi, a) \mapsto (-\chi, -a)$.

2) Their zero loci

$$
S_{\ell, \pm} := \{ \mathcal{P}_{\ell, \pm}(\chi, a, s) = 0 \} \subset \mathbb{C}^3_{\chi, a, s},
$$
\n(1.14)

which are called the **determinantal surfaces**, are irreducible affine surfaces permuted by the involution

$$
\mathcal{I}: (\chi, a, s) \mapsto (-\chi, -a, s).
$$

Set

$$
S^o_{\ell, \pm} := S_{\ell, \pm} \setminus \{s = 0\}.
$$

Theorem 1.14 1) The surfaces $S_{\ell,\pm}^o$ are smooth and disjoint, and they are disjoint from the s-axis, i.e., the line $\{\chi = a = 0\}.$

2) Each surface $S^o_{\ell, \pm}$ is fibered over the s-axis (punctured at the origin) by smooth rational curves of degree $2\ell+1$ conformally parametrized by Riemann spheres punctured in at most $2\ell+1$ points; their projective closures in $\mathbb{CP}^3 \supset \mathbb{C}^3_{\chi,a,s}$ are parametrized by $\overline{\mathbb{C}}$.

3) Each surface $S^o_{\ell, \pm}$ is biholomorphically parametrized by the complement of the product $\mathbb{C}_{s}^{*} \times \overline{\mathbb{C}}$ to an analytic hypersurface (curve) so that the parametrization preserves the coordinate s: sends each s-fiber $\{s\} \times \overline{\mathbb{C}}$ to the s-fiber of the surface $S^o_{\ell, \pm}$.

Corollary 1.15 The projective closures in \mathbb{CP}^3 of the determinantal surfaces $S_{\ell, \pm} \subset \mathbb{C}^3$ \mathbb{CP}^3 are rational surfaces.

Proof The projection $\pi_s : S_{\ell, \pm} \to \mathbb{C}_s$ is rational, and the preimage of every $s \neq 0$ is a rational curve, by Statement 2) of Theorem 1.14. This together with [7, theorem III.4] implies that $S_{\ell, \pm}$ is a rational surface.

Let us now check that the intersection of each surface $S_{\ell,\pm} \subset \mathbb{C}^3_{\chi,a,s}$ with the hyperplane $\{\chi = 0\}$ is a component $\Gamma_{\ell, \pm}$ of the pullback of the spectral curve Γ_{ℓ} under the 2 : 1 projection $(a, s) \mapsto (\lambda = \frac{a^2 - s^2}{4})$ $\frac{-s^2}{4}, \mu = \frac{s}{2}$. On one hand, this follows "a posteriori" from Theorem 1.12 and Remark 1.1. On the other hand, we will present its next direct and self-contained proof. To do this, let us rescale the coordinates (a, s) to the new coordinates (μ, r) :

$$
\mu = \frac{s}{2}, r = \frac{a}{2}; \lambda = r^2 - \mu^2.
$$

Consider the following $\ell \times \ell$ -matrices:

$$
\mathcal{G}_{\ell} = \begin{pmatrix} 0 & \cdots & 0 & \mu \\ \vdots & & \ddots & -(\ell-1) \\ 0 & \cdots & & \ddots \\ \mu & -1 & & 0 \end{pmatrix}.
$$

The following relation (found in [18, section 3, eq. (30)]) holds:

$$
H_{\ell} + (r^2 - \mu^2) \operatorname{Id} = -(\mathcal{G}_{\ell} + r \operatorname{Id})(\mathcal{G}_{\ell} - r \operatorname{Id}),
$$

and it yields that

$$
Q_{\ell}(\lambda, \mu^2) = (-1)^{\ell} Q_{\ell,+}(\mu, r) Q_{\ell,-}(\mu, r), \quad Q_{\ell, \pm}(\mu, r) := \det(\mathcal{G}_{\ell} \pm r \operatorname{Id}). \tag{1.15}
$$

This implies that the preimage of the curve $\Gamma_{\ell} \subset \mathbb{C}^2_{\lambda,\mu}$ under the above projection consists of two curves

$$
\Gamma_{\ell,\pm} = \{Q_{\ell,\pm}(\mu,r) = 0\}.
$$

They are irreducible, as is Γ_{ℓ} , since the projection is 2 : 1, see [29, section 3].

Proposition 1.16 The restriction to the hyperplane $\{\chi = 0\}$ of the polynomial $\mathcal{P}_{\ell, \pm}$ coincides with the polynomial $Q_{\ell,\pm}$ up to non-zero constant factor. In particular, one has $S_{\ell,\pm} \cap {\chi = 0} = \Gamma_{\ell,\pm}.$

Proof Substituting $\chi = 0$ to the matrices $G_{1,\ell}$ and $G_{2,\ell}$ crosses out the χa in $G_{1,\ell}$ and transforms the matrix $G_{2,\ell}$ to the matrix whose lower line is zero and whose only non-zero elements are those equal to $\frac{a}{2}$. Therefore, each determinant polynomial $\mathcal{P}_{\ell,\pm}$ becomes $\frac{1}{2}$ times the determinant of the $\ell \times \ell$ -submatrix in $G_{1,\ell} \pm G_{2,\ell}$ formed by the first ℓ lines and columns. The latter submatrix is obtained from the matrix $\mathcal{G}_{\ell} \pm r$ Id by reflection transposition of lines, permuting the first and the last lines, the 2nd and the $\ell-1$ -th line etc. This permutation multiplies the determinant by a power of -1 . Finally, the restriction to the hyperplane $\{\chi = 0\}$ of the determinants $\mathcal{P}_{\ell, \pm} = \det(G_{1,\ell} \pm G_{2,\pm})$ and $Q_{\ell, \pm} = \det(\mathcal{G}_{\ell} \pm r \operatorname{Id})$ differ by a non-zero constant factor. This proves the proposition. non-zero constant factor. This proves the proposition.

Example 1.17 For $\ell = 1$, up to constant factor,

$$
\pm 4\mathcal{P}_{1,\pm}(\chi,a,s) := (a \pm s)(1 - 4\chi^2) + 4\chi; \tag{1.16}
$$

$$
S_{1,\pm} = \{\mathcal{P}_{1,\pm} = 0\} = \{(a \pm s)(1 - 4\chi^2) + 4\chi = 0\}.
$$

The surfaces $S_{1,\pm}$ are rational, parametrized by the coordinates (χ, s) and permuted by the involution I . This parametrization also implies rationality of the s-fibers of the surfaces $S_{1,\pm}$ given by Statement 2) of Theorem 1.14.

Example 1.18 For $\ell = 2$ one has

$$
\pm 8\mathcal{P}_{2,\pm}(\chi,a,s) = (2\chi \pm 1)^2 (2\chi \mp 1)(a^2 - s^2) - 2a(2\chi \pm 1)(6\chi \mp 1) + 16\chi. \tag{1.17}
$$

Each surface $S_{2,\pm} = \{P_{2,\pm} = 0\}$ is irreducible and rational. Indeed, for every fixed $\chi \in \mathbb{C}$ the equation $\mathcal{P}_{2,\pm} = 0$ defines a conic $\mathcal{C}_{\chi,\pm} \subset \mathbb{C}^2_{a,s}$, whose projective closure $\overline{\mathcal{C}}_{\chi,\pm}$ in $\mathbb{CP}^2_{a:s:t} \supset \mathbb{C}^2_{a,s} = \{t = 1\}$ passes through the point D at infinity with homogeneous coordinat $[1:-1:0]$. The projection of the projective conic $\overline{\mathcal{C}}_{\chi,\pm}$ from the point D to the a-axis yields its rational parametrization by $\overline{C} = C_a \cup \{\infty\}$ that depends rationally on χ . This yields a birational parametrization of $S_{2,\pm}$ by two parameters: the parameter χ and the above projection.

Remark 1.19 The above argument does not prove rationality of the s-fibers of the surface $S_{2,+}$. Each its s-fiber

$$
\gamma_{s,\pm} = \{(\chi, a) \mid \mathcal{P}_{2,\pm}(\chi, a, s) = 0\} \subset \mathbb{C}^2
$$

is a quintic that is the zero locus of the polynomial (1.17) with the fixed s. Though we prove rationality of the s-fibers in full generality for all ℓ , below we present yet another proof in the special case $\ell = 2$, together with the description of their branching points over the χ -axis.

In what follows for every $s \neq 0$ by $\widetilde{\gamma}_{s,\pm}$ we denote the normalization: the compact Riemann surface holomorphically parametrizing the projective closure $\overline{\gamma}_{s+}$ of the s-fiber $\gamma_{s,\pm}$ (bijectively up to self-intersections).

Proposition 1.20 Let $\ell = 2$. The projection to the coordinate χ induces a double branched covering $\pi_{\chi}: \widetilde{\gamma}_{s,\pm} \to \overline{\mathbb{C}}_{\chi}$ with two branching points (i.e., critical values)

$$
\chi_{\pm,+} = \pm \frac{1}{2} (1 + i s^{-1}), \quad \chi_{\pm,-} = \pm \frac{1}{2} (1 - i s^{-1}).
$$
\n(1.18)

In particular, the normalization $\tilde{\gamma}_{s,\pm}$ is the Riemann sphere, and $\gamma_{s,\pm}$ is rational.

Proof The fact that π_{χ} is a double covering follows from quadraticity in a of the polynomial (1.17). Its branching points lie in the zero locus of its discriminant $\Delta = \Delta_s(\chi)$ of the polynomial (1.17) as a quadratic polynomial in a. One has

$$
\Delta_s(\chi) = 4(2\chi \pm 1)^2((6\chi \mp 1)^2 - 16\chi(2\chi \mp 1) + s^2(4\chi^2 - 1)^2)
$$

= 4(2\chi \pm 1)^2(36\chi^2 \mp 12\chi + 1 - 32\chi^2 \pm 16\chi + s^2(4\chi^2 - 1)^2)
= 4(2\chi \pm 1)^4(1 + s^2(2\chi \mp 1)^2).

Therefore, the zeros of the discriminant are $\mp \frac{1}{2}$ $\frac{1}{2}$ (of multiplicity 4) and the points (1.18) (of multiplicity one). Its zero is a branching point, if and only if its multiplicity is odd. The infinity is not a branching point, since the polynomial Δ_s has even degree. Thus, the only branching points are the points (1.18). Hence, $\tilde{\gamma}_{s,\pm}$ is the Riemann sphere, by
Riemann-Hurwitz Theorem. Proposition 1.20 is proved. Riemann-Hurwitz Theorem. Proposition 1.20 is proved.

1.4 Relation to model of Josephson junction. Historical remarks

In 1962 B.Josephson [36] predicted tunnelling effect for the Josephson junction: a system of two superconductors separated by a narrow dielectric (Nobel Prize, 1973). It was confirmed experimentally by P.W.Anderson and J.M.Rowell in 1963 [1]. The overdamped Josephson junction, see [55, 47, 40, 50], [6, p. 306], [41, pp. 337–340], [42, p.193], [43, p. 88] is described by the family of nonlinear differential equations

$$
\frac{d\phi}{dt} = -\sin\phi + B + A\cos\omega t, \ \omega > 0, \ B \ge 0. \tag{1.19}
$$

Here ϕ is the difference of phases (arguments) of the complex-valued wave functions describing the quantum mechanic states of the two superconductors. Its derivative is equal to the voltage up to known constant factor.

The variable and parameter changes

$$
\tau := \omega t, \ \theta := \phi + \frac{\pi}{2}, \ \ell := \frac{B}{\omega}, \ a = \frac{1}{\omega}, \ s := \frac{A}{\omega}, \tag{1.20}
$$

transform (1.19) to a non-autonomous ordinary differential equation on the two-torus \mathbb{T}^2 = $S^1 \times S^1$ with coordinates $(\theta, \tau) \in \mathbb{R}^2 / 2\pi \mathbb{Z}^2$:

$$
\frac{d\theta}{d\tau} = a\cos\theta + \ell + s\cos\tau.
$$
 (1.21)

The rotation number of dynamical system (1.21), see [2, p. 104], is a function of the parameters a, ℓ, s . For every fixed a , the rotation number function is considered as a function of the parameters (B, A) . The two-dimensional phase-lock areas of family (1.21) are those level subsets of the rotation number function in $\mathbb{R}^2_{B,A}$ that have non-empty interiors. V.M.Buchstaber, O.V.Karpov and S.I.Tertychnyi have shown that they exists only for integer values of the rotation number $[16]$. The phase-lock areas of family (1.21) were studied by V.M.Buchstaber, O.V.Karpov, S.I.Tertychnyi, Yu.S.Ilyashenko, D.A.Filimonov, D.A.Ryzhov, A.V.Klimenko, O.L.Romaskevich, V.A.Kleptsyn, I.V.Schurov, Yu.P.Bibilo and the author, see [12]–[19], [8], [34, 38, 39], [26]–[29], [56, 57] and references therein.

Equations of type (1.19), (1.21) arise in different domains of mathematics, e.g., in planimeters, see [24, 25]. A subfamily of family (1.21) was studied by Yu.S.Ilyashenko and J.Guckenheimer [32] from the slow-fast system point of view. They obtained results on its limit cycles, as $\omega \to 0$.

Model (1.21) has the following equivalent description by a family of two-dimensional linear systems of differential equations on the Riemann sphere, see [14, 16, 19, 24, 25, 34], [12, subsection 3.2]. The variable change

$$
z = e^{i\tau} = e^{i\omega t}, \quad \Phi = e^{i\theta} = ie^{i\phi}
$$

transforms equation (1.21) on the function $\theta(\tau)$ to the Riccati equation

$$
\frac{d\Phi}{dz} = z^{-2}((\ell z + \frac{s}{2}(z^2 + 1))\Phi + \frac{a}{2}z(\Phi^2 + 1)).
$$
\n(1.22)

Equation (1.22) is the projectivization of the two-dimensional linear system

$$
Y' = \left(\frac{\text{diag}\left(-\frac{s}{2}, 0\right)}{z^2} + \frac{\mathcal{B}}{z} + \text{diag}\left(0, \frac{s}{2}\right)\right)Y, \quad \mathcal{B} = \left(\begin{matrix} -\ell & -\frac{a}{2} \\ \frac{a}{2} & 0 \end{matrix}\right),\tag{1.23}
$$

in the following sense: a function $\Phi(z)$ is a solution of (1.22), if and only if $\Phi(z) = \frac{Y_2}{Y_1}(z)$, where the vector function $Y(z) = (Y_1(z), Y_2(z))$ is a solution of system (1.23). For $s \neq 0$ system (1.23) has two irregular nonresonant singular points at 0 and at ∞ . In the new parametes

$$
\mu := \frac{s}{2} = \frac{A}{2\omega}, \quad \lambda := \frac{a^2 - s^2}{4} = \frac{1}{4\omega^2} - \mu^2 \tag{1.24}
$$

family of systems (1.23) is equivalent to the following family of special double confluent Heun equations on the function $E = Y_2(z)$, see [57], [18]–[22]:

$$
z^{2}E'' + ((\ell + 1)z + \mu(1 - z^{2}))E' + (\lambda - \mu(\ell + 1)z)E = 0.
$$
 (1.25)

We will deal with the so-called *conjugate* special double confluent Heun equation obtained from (1.25) by sign change at ℓ :

$$
z^{2}E'' + ((-\ell + 1)z + \mu(1 - z^{2}))E' + (\lambda + \mu(\ell - 1)z)E = 0.
$$
 (1.26)

Let $\rho = \rho(B, A) = \rho(B, A; \omega)$ denote the rotation number function of family (1.21). For every $r \in \mathbb{Z}$ let L_r denote the phase-lock area of family with the rotation number r:

$$
L_r = \{(B, A) \mid \rho(B, A) = r\} \subset \mathbb{R}_{B, A}^2.
$$

It is known that each phase-lock area L_r is an infinite garland of domains going to infinity "asymptotically vertically" where every two neighbor domains are separated by one point. In more detail, the boundary ∂L_r consists of two graphs $\{B = g_{r,\pm}(A)\}\$ of functions $g_{r,\pm}$ analytic on \mathbb{R}_A and having Bessel type asymptotics, as $A \to \infty$ (observed and proved on physics level in [51], see also [41, p. 338], [6, section 11.1], [15]; proved mathematically in [39]). The graphs have infinitely many intersection points (which separate adjacent domains). One of them lies in the B-axis, is called the growth point (it exists only for $r \in \mathbb{Z} \setminus \{0\}$, and its abscissa is equal to sign $r\sqrt{r^2\omega^2+1}$, see [17, corollary 3]. The other intersection (separation) points are called constrictions. It was shown in [8, theorem 1.4] that all the constrictions in L_r lie in the same vertical line

$$
\Lambda_r := \{ B = r\omega \},\
$$

which is called the *axis* of the phase-lock area L_r . See Fig. 1. The constrictions correspond exactly to those parameters $\ell = r$, $(\lambda, \mu) \in \mathbb{R}^2$, $\lambda + \mu^2 > 0$, see (1.24), for which the special double confluent Heun equation (1.25) has an entire solution, i.e., holomorphic on \mathbb{C}_z , as was shown in [12, theorem 3.3]. The set of those parameters (λ, μ) for which an entire solution exists is given by an explicit analytic equation in terms of an infinite matrix product introduced in [19, pp. 332, 337]. This was stated as a conjecture in loc. cit. and proved in [12, theorem 3.5].

For a given $\omega > 0$ a point $(B, A) \in L_r$ is called a *generalized simple intersection*, if it is an intersection point of the boundary ∂L_r with the axis $\Lambda_\ell = \{B = \ell \omega\}, \ell \in \mathbb{Z}$, of some phase-lock area L_{ℓ} with $\ell \equiv r \pmod{2}$, and it is not a constriction. See Fig. 2. It is known that the generalized simple intersections correspond exactly to those parameters ℓ , λ , μ with $\mu \neq 0$, $\lambda + \mu^2 > 0$, for which the conjugate Heun equation (1.26) has a polynomial

Figure 1: Phase-lock areas, their constrictions and growth points for $\omega = 1$. The abscissa is B, the ordinate is A. Figure taken from paper [13, fig. 1b)] with authors' permission, with coordinate axes added.

solution, as was shown in [13, theorem 1.15]. The proof of this result is based on the previous joint result by Buchstaber and the author [12, theorem 3.10] stating that if the conjugate Heun equation (1.26) has a polynomial solution, then Heun equation (1.25) has no entire solution. This was a solution of a series of conjectures stated and partially studied by V.M.Buchstaber and S.I.Tertychnyi in [19, 20].

Figure 2: Phase-lock areas for $\omega = 0.3$: constrictions (correspond to Heun equations (1.25) with entire solutions) and generalized simple intersections (correspond to conjugate Heun equations (1.26) with polynomial solutions).

1.5 Four-parameter extension of the Josephson family of dynamical systems on 2-torus and isomonodromic foliation

The four-parameter extension of the three-parameter "Josephson" family of dynamical systems (1.21) given in [8, subsection 6.2] is

$$
\frac{d\theta}{d\tau} = \nu + a\cos\theta + s\cos\tau + \psi\cos(\theta - \tau); \quad \nu, a, \psi \in \mathbb{R}, \quad s > 0, \ (a, \psi) \neq (0, 0). \tag{1.27}
$$

The variable changes $\Phi = e^{i\theta}, z = e^{i\tau}$ transform (1.27) to the Riccati equation

$$
\frac{d\Phi}{dz} = \frac{1}{z^2} \left(\frac{s}{2} \Phi + \frac{\psi}{2} \Phi^2 \right) + \frac{1}{z} \left(\nu \Phi + \frac{a}{2} (\Phi^2 + 1) \right) + \left(\frac{s}{2} \Phi + \frac{\psi}{2} \right). \tag{1.28}
$$

A function $\Phi(z)$ is a solution of the latter Riccati equation, if and only if $\Phi(z) = \frac{Y_2(z)}{Y_1(z)}$, where $Y = (Y_1, Y_2)(z)$ is a solution of the linear system

$$
Y' = \left(-s\frac{\mathbf{K}}{z^2} + \frac{\mathbf{R}}{z} + s\mathbf{N}\right)Y,
$$

\n
$$
\mathbf{K} = \begin{pmatrix} \frac{1}{2} & \chi \\ 0 & 0 \end{pmatrix}, \ \mathbf{R} = \begin{pmatrix} -(\ell + \chi a) & -\frac{a}{2} \\ \frac{a}{2} & \chi a \end{pmatrix}, \ \mathbf{N} = \begin{pmatrix} 0 & 0 \\ \chi & \frac{1}{2} \end{pmatrix};
$$

\n
$$
\chi = \frac{\psi}{2s}, \ \ell = \nu - \frac{\psi a}{s} = \nu - 2\chi a.
$$

\n(1.29)

The residue matrix of the formal normal forms of system (1.29) at 0 and at ∞ is the same and equal to

$$
diag(-\ell, 0), \quad \ell = \nu - 2\chi a. \tag{1.30}
$$

See the background material on formal normal forms, Stokes matrices and isomonodromic deformations in Subsection 2.3.

Theorem 1.21 [8, subsection 6.2] The four-parameter family of linear systems (1.29) is analytically foliated by one-dimensional isomonodromic families defined by the following non-autonomous system of differential equations:

$$
\begin{cases}\n\chi_s' = \frac{a - 2\chi(\ell + 2\chi a)}{2s} \\
a_s' = -2s\chi + \frac{a}{s}(\ell + 2\chi a) \\
\ell_s' = 0\n\end{cases} \tag{1.31}
$$

For every $\ell \in \mathbb{R}$ the function (1.7), i.e.,

$$
w(s) := \frac{a(s)}{2s\chi(s)} = \frac{a(s)}{\psi(s)}
$$
(1.32)

satisfies Painlevé 3 equation (1.8) with ℓ replaced by $-\ell$ along solutions of (1.31):

$$
w'' = \frac{(w')^2}{w} - \frac{w'}{s} - 2\ell \frac{w^2}{s} + (2\ell - 2)\frac{1}{s} + w^3 - \frac{1}{w}.
$$
\n(1.33)

Remark 1.22 In fact, isomonodromicity was proved in loc. cit. for systems obtained from (1.29) by adding the scalar matrix $-\frac{1}{2}$ $\frac{1}{2}Id$ to the matrix N in (1.29). But addition of the latter term changes neither projectivization (Riccati equation), nor isomonodromicity. The differential equation (1.31) is Hamiltonian with the time s-depending Hamiltonian function

$$
H(\chi, a, s) := -\frac{\chi^2 a^2}{s} + \frac{a^2}{4s} + s\chi^2 - \frac{\ell \chi a}{s},\tag{1.34}
$$

see $[27,$ proposition 1.14]. Representations of Painlevé equations as Hamiltonian systems were found by J.Malmquist $[46]$ (for all, except Painlevé 3) and by K.Okamoto $[48, p. 265]$ (for all, including Painlevé 3). The above Hamiltonian representation of Painlevé 3 equation (1.33) with Hamiltonian (1.34) can be obtained from Okamoto Hamiltonian representation by a variable change: see [27, remark 1.15].

We will be dealing with the linear systems obtained from (1.29) by change of sign at ℓ :

$$
Y' = \left(-s\frac{\mathbf{K}}{z^2} + \frac{\hat{\mathbf{R}}}{z} + s\mathbf{N}\right)Y,
$$
(1.35)

$$
\mathbf{K} = \begin{pmatrix} \frac{1}{2} & \chi \\ 0 & 0 \end{pmatrix}, \ \hat{\mathbf{R}} = \begin{pmatrix} \ell - \chi a & -\frac{a}{2} \\ \frac{a}{2} & \chi a \end{pmatrix}, \ \mathbf{N} = \begin{pmatrix} 0 & 0 \\ \chi & \frac{1}{2} \end{pmatrix}.
$$

Note that for $\chi = 0$ system (1.35) coincides with system (1.23) where ℓ is replaced by $-\ell$. Therefore, it is equivalent to the conjugate Heun equation (1.26) on $E = Y_2(z)$, analogously to the similar statement on (1.23) and Heun equation (1.25), see the end of Subsection 1.4.

Analogously to the above discussion, the parameter space of family (1.35) is foliated by isomonodromic families given by solutions of system of differential equations (1.6), along which the function $w(s)$ satisfies Painlevé 3 equation (1.8).

1.6 Plan of proof of main results

The proof of main results is based on analytic theory of linear systems with irregular non-resonant singularities: formal normal forms, Stokes phenomena (canonical solutions in Stokes sectors and monodromy-Stokes data) and isomonodromic deformations. The corresponding background material is recalled in Subsection 2.1.

In Subsection 2.2 we prove Proposition 1.9. We will deduce it from a more general proposition dealing with general linear systems of type

$$
Y' = \left(\frac{K}{z^2} + \frac{R}{z} + N\right)Y, \quad K, R, N \text{ are } 2 \times 2 - \text{matrices},\tag{1.36}
$$

having irregular non-resonant Poincaré rank 1 singular points at 0 and at ∞ . It states that a system has a vector polynomial solution, if and only if two appropriate canonical solutions coming from 0 and from ∞ coincide and form a solution holomorphic in \mathbb{C} , and then it is a unique polynomial solution. Its degree is equal to the corresponding residue eigenvalue of the formal normal form at ∞ .

In Subsection 2.4 we prove Theorem 1.12 describing the set of parameters of those systems (1.5) that have polynomial solutions as the union of the zero loci $S_{\ell,+} = \{P_{\ell,+} = 0\}$. To do this, in Subsection 2.3 we introduce and study a more general class of systems (1.36), the so-called class S, which contains systems (1.5), and show that class S coincides with the class of those systems (1.36) that have a symmetry of the type $(Y_1(z), Y_2(z)) \mapsto$ $z^{\ell}(Y_2(z^{-1}), Y_1(z^{-1}))$; for systems (1.5), the latter ℓ coincides with the ℓ in (1.5). We show that if a polynomial solution of a system (1.5) exists, then it is either invariant, or antiinvariant under the above symmetry, and then deduce that the parameters (χ, a, s) of the system lie in $S_{\ell,+}$, respectively $S_{\ell,-}$. And then we prove the converse.

For the proof of Theorem 1.14 we consider analytic extensions of appropriate canonical sectorial solutions of a system (1.5) at 0 and at ∞ to the base point $z_0 = 1$ along paths α_0^{-1} and α_{∞}^{-1} in \mathbb{C}^* going from appropriate Stokes sectors to z_0 ; the sectors and the paths are permuted by the involution $z \mapsto z^{-1}$ (such sector and path collections will be called symmetric). Taking projectivizations of the vector values at z_0 of thus extended solutions yields four points $q_{10}, q_{20}, q_{1\infty}, q_{2\infty} \in \mathbb{CP}^1 = \overline{\mathbb{C}}$. The collection of the four points q_{jp} taken together with the monodromy operator (up to natural $PSL_2(\mathbb{C})$ -action) is called the monodromy-Stokes data, see Subsection 2.1. We consider yet another point q'_{10} that is the image of the point q_{10} under the projectivized monodromy of system (1.5) acting on the space of projectivized initial conditions (local solutions) at z_0 . In other terms, take the analytic extension of the solution $\Phi(z)$ of the Riccatti equation (the priojectivization of our system (1.5)) with the initial condition $\Phi(z_0) = q_{10}$ along a counterclokwise circuit around the origin. Set q'_{10} to be the value of thus extended solution at z_0 .

Our goal is to construct a parametrization of each surface $S^o_{\ell, \pm}$ by the complement of $\mathbb{C}_s^* \times \overline{\mathbb{C}}$ to an analytic hypersurface (curve). In Subsection 2.5 we show that (1.5) has a polynomial solution, if and only if $\ell \in \mathbb{N}$ and the monodromy-Stokes data is parabolic and non-degenerate. Parabolicity means that some of the above points q_{ip} coincide, namely, $q_{20} = q_{1\infty}$, the monodromy is unipotent and its projectivization fixes q_{20} . Non-degeneracy means that the collection $q_{10'}$, q_{10} , q_{20} , $q_{2\infty}$ consists of at least three distinct points. We describe in similar terms all the systems (1.36) that are gauge equivalent to systems (1.5) with polynomial solutions.

In Subsection 2.6 we show that a parabolic non-degenerate monodromy-Stokes data is uniquely determined by the cross-ratio

$$
\mathcal{R} := \frac{(q_{20} - q'_{10})(q_{10} - q_{2\infty})}{(q_{20} - q_{2\infty})(q_{10} - q'_{10})}.
$$

Then we show that two systems of type (1.5) with the same ℓ having a polynomial solution coincide up to parameter sign change $(\chi, a) \mapsto (-\chi, -a)$, if and only if they have the same parameter s and the same cross-ratio R (constructed for given paths $\alpha_{0,\infty}$, the same for both systems). This already yields a local parametrization of surfaces $S_{\ell,\pm}^o$ by coordinates $(s, \mathcal{R}) \in \mathbb{C}^* \times \overline{\mathbb{C}}$ (though we will not prove that at this place). But \mathcal{R} is not uniquely defined: it depends on homotopic classes of paths $\alpha_{0,\infty}$.

In order to obtain a global parametrization of surfaces $S^o_{\ell,\pm}$, in Subsection 2.7 we first

construct a holomorphic family of vector bundles $\mathcal{E}_{t,\mathcal{R}} = \mathcal{E}_{t,\mathcal{R},\ell}$ with meromorphic connections on $\overline{\mathbb{C}}$ depending on two parameters $(t, \mathcal{R}) \in \mathbb{C} \times \overline{\mathbb{C}}$, set $s := e^t$, and a continuous family of pairs of paths $\alpha_{0,t}$ and $\alpha_{\infty,t}$ going from $z_0 = 1$ to appropriate Stokes sectors at 0 and at ∞ respectively and permuted by the involution $z \mapsto z^{-1}$ that satisfy the following statements:

1) the connections have irregular non-resonant Poincaré rank 1 singular points at 0 and at ∞ with formal normal forms

$$
Y' = \left(\frac{\text{diag}(-\frac{s}{2}, 0)}{z^2} + \frac{\text{diag}(\ell, 0)}{z}\right) Y \text{ at } 0,
$$
 (1.37)

$$
Y' = \left(\text{diag}(0, \frac{s}{2}) + \frac{\text{diag}(\ell, 0)}{z}\right) Y \quad \text{at } \infty; \tag{1.38}
$$

2) the above-mentioned Stokes sectors are those for the above normal forms;

3) the monodromy-Stokes data defined by the above paths $\alpha_{0,t}$, $\alpha_{\infty,t}$ are parabolic nondegenerate, and the corresponding cross-ratio is equal to \mathcal{R} ;

4) two bundles $\mathcal{E}_{t,\mathcal{R}}$ and $\mathcal{E}_{t',\mathcal{R}'}$ with connections are analytically gauge equivalent, if and only if

$$
(t',\mathcal{R}') = (t,\mathcal{R}) + (2\pi in, 2n), \quad n \in \mathbb{Z}.
$$

Parabolicity implies that each bundle with connection has a meromorphic flat section with only one pole, at infinity, which is a vector polynomial of degree ℓ in a trivializing chart on a neighborhood of infinity.

In Subsection 2.8 we take the quotient of the parameter space by the translation by the vector $(2\pi i, 2)$, which leaves invariant the coordinate $\xi := \pi i \mathcal{R} - t$. We get a holomorphic family of bundles with connections parametrized by $(s,\xi) \in \mathbb{C}^* \times \overline{\mathbb{C}}$, where all the bundles are distinct: pairwise gauge non-equivalent. It follows from a well-known theorem, see [45, proposition 4.1], [11, appendix 3], [49, theorem 2.2, p.449], that the subset $\text{Triv}_{\ell} \subset \mathbb{C}^* \times \overline{\mathbb{C}}$ of those parameters that correspond to holomorphically trivial bundles is the complement to an analytic hypersurface (curve). They correspond to linear systems of type (1.36) on the Riemann sphere. Parabolicity of the monodromy-Stokes data implies that the systems in question are gauge equivalent to systems of type S having a vector polynomial solution and the above formal normal forms. As is shown in Subsection 2.3, a system of type S with the above formal normal forms is gauge equivalent to a system of type (1.5) , if and only if it satisfies the following *line non-coincidence condition*: the eigenlines of the main term matrices at 0 and at ∞ corresponding to non-zero eigenvalues are distinct. Erasing yet another curve from the subset Triv_ℓ consisting of those systems that do not satisfy the line non-coincidence condition, we get a new open and dense subset $\text{Triv}_{\ell}^o \subset \mathbb{C}^* \times \overline{\mathbb{C}}$ of those points that correspond to systems gauge equivalent to systems (1.5) with polynomial solutions; for every point $(s, \xi) \in \text{Triv}_{\ell}^o$ the corresponding system (1.5) is uniquely defined up to sign change $(\chi, a) \mapsto (-\chi, a)$. This yields biholomorphic parametrizations of the surfaces $S^o_{\ell, \pm}$ by Triv_{ℓ}. Afterwards we show that the complement $\mathbb{C}^* \times \overline{\mathbb{C}} \setminus \text{Triv}_{\ell}^o$ is an analytic curve, by proving that potential non-extendability points of the closure of the second erased curve are isolated and hence, removable by Shiffman Theorem, see [52] and [23, subsection 4.4].

In Subsection 2.9 we deduce regularity of the surfaces $S^o_{\ell, \pm}$ and rationality of their sfibers from their parametrization and finish the proof of Theorem 1.14. Theorem 1.13 (irreducibility of surfaces $S_{\ell, \pm}$ and of the polynomials $\mathcal{P}_{\ell, \pm}$) is proved in Subsection 2.10. Smoothness of the spectral curves $\Gamma_{\ell,\pm}^o$ and the genus formula are proved in Subsection 2.11.

1.7 Open problems

A list of open problems on model of overdamped Josephson junction, on its four-dimensional extension (1.27) , its isomonodromic foliation given by (1.31) and related questions is presented in [8, section 6].

For every $r \in \mathbb{Z}_{\neq 0}$ set

 $L_r^+ = L_r \cap \{A > 0\},\ A_r := \text{the ordinate of the highest simple intersection in } L_r,$

 $\Lambda_r := \{ B = r \}, \quad Sr = \Lambda_r \cap \{ A > A_r \}.$

The connectivity conjecture, see [26, conjecture 1.14], states that the intersection $L_r^+ \cap \Lambda_r$ coincides with the ray Sr, and thus, is connected. By [8, corollary 6.1 and remark 6.2], it is reduced to a pure real-algebro-geometric conjecture on the real spectral curves Γ_{ℓ} , see [8, conjecture 6.3], which we reformulate here in the following slightly stronger form.

Conjecture 1.23 1) For every $\ell \in \mathbb{N}$ the intersection of the real spectral curve Γ_{ℓ} defined in (1.11) with the domain $\{\lambda + \mu^2 > 0\}$ contains a unique connected component Γ_{ℓ}^* whose points correspond to differential equations (1.21) on \mathbb{T}^2 with the rotation number ℓ . Its other components correspond to positive rotation numbers that are less than ℓ .

2) The projection $(\lambda, \mu) \mapsto \lambda + \mu^2$ maps the component Γ_{ℓ}^* diffeomorphically onto \mathbb{R}_+ .

Problem 1.24 Study real (complex) determinantal surfaces $S_{\ell,\pm}$ and their real (complex) isomonodromic foliations given by real (complex) solutions of system of differential equations (1.6). For which real initial conditions (a_0, s_0) in the hyperplane $\{\chi = 0\}$ (corresponding to model (1.21) of Josephson junction) the corresponding real solutions $(\chi(s), a(s))$ of (1.6) return back to the hyperplane $\{\chi = 0\}$ for some $s > s_0$ and for infinitely many $s > s_0$? In other words, study the definition domain of the Poincar´e first return map (and its iterates) to the above hyperplane for the real isomonodromic foliation (1.6).

It is known that returns to the hyperplane $\{\chi = 0\}$ corresponds to simple poles with residue 1 of the solution of Painlevé equation (1.8) governing the solution of (1.6) in question, see [8, remark 6.14]. Few solutions of Painlévé 3 equation, e.g., the *tronquée solutions*, see [44], are bounded on a semi-interval $[C, +\infty)$, and hence, have no poles there.

Problem 1.25 Describe those solutions of (1.6) whose graphs lie in $S_{\ell,+}$ and that correspond to tronquée solutions of the corresponding Painlevé 3 equation (1.8) .

2 Proof of main results

2.1 Background material: Stokes phenomena, monodromy–Stokes data and isomonodromic deformations

The following material on irregular singularities of linear systems and Stokes phenomena is contained in [3, 4, 5, 33, 37, 53].

Consider a two-dimensional linear system

$$
Y' = \left(\frac{K}{z^2} + \frac{R}{z} + N\right)Y, \quad Y = (Y_1, Y_2) \in \mathbb{C}^2,
$$
\n(2.1)

over a neighborhood of zero in \mathbb{C}_z . Here K, R, N are complex 2×2 -matrices. We consider that K has distinct eigenvalues $\lambda_{10} \neq \lambda_{20}$, and N has distinct eigenvalues $\lambda_{1\infty} \neq \lambda_{2\infty}$. By definition, the two latter conditions on eigenvalues are equivalent to the condition that the singular point 0 (respectively, ∞) of system (2.1) is *irregular non-resonant of Poincaré rank* 1. The matrix K is conjugate to $\widetilde{K} = \text{diag}(\lambda_{10}, \lambda_{20}), \widetilde{K} = \mathbf{H_0}^{-1} K \mathbf{H_0}, \mathbf{H_0} \in GL_2(\mathbb{C}),$ and one can achieve that $K = K$ by applying the constant linear change (gauge transformation) $Y = H_0 \dot{Y}$.

Recall that two systems of type (2.1) are analytically equivalent near the origin, if one can be transformed to the other by linear space coordinate change $Y = H_0(z)Y$, where $H_0(z)$ is a holomorphic $GL_2(\mathbb{C})$ -valued function on a neighborhood of the origin. Two systems (2.1) are *formally equivalent* at the origin, if the above $H_0(z)$ exists in the class of invertible formal power series with matrix coefficients. Analytic (formal) equivalence at infinity is defined analogously, with "power series" being Laurent series containing only non-positive degrees of z.

System (2.1) is formally equivalent at the origin to a unique formal normal form

$$
\widetilde{Y}' = \left(\frac{\widetilde{K}}{z^2} + \frac{\widetilde{R}_0}{z}\right) \widetilde{Y}, \ \widetilde{K} = \text{diag}(\lambda_{10}, \lambda_{20}), \ \widetilde{R}_0 = \text{diag}(b_{10}, b_{20}), \tag{2.2}
$$

 \widetilde{R}_0 is the diagonal part of the matrix $H_0^{-1} R H_0$. (2.3)

The matrix coefficient K in system (2.1) and the corresponding matrix K in (2.2) are called the main term matrices, and R , R_0 the residue matrices at the origin. However the normalizing series $H_0(z)$ bringing (2.1) to (2.2) generically diverges. At the same time, there exists a covering of a punctured neighborhood of zero by two sectors S_0^0 and S_1^0 with vertex at 0 in which there exist holomorphic $GL_2(\mathbb{C})$ -valued matrix functions $H_{j;0}(z)$, $j = 0, 1$, that are C^{∞} smooth on $\overline{S}_{j}^{0} \cap D_{r}$ for some $r > 0$, and such that the variable changes $Y = H_{i;0}(z)\tilde{Y}$ transform (2.1) to (2.2). This Sectorial Normalization Theorem holds for the so-called Stokes sectors. Namely, consider the rays issued from 0 and forming the set

$$
\{\operatorname{Re}\frac{\lambda_{10} - \lambda_{20}}{z} = 0\}.
$$
\n(2.4)

They are called *imaginary dividing rays* or *Stokes rays*. A sector S_j^0 is called a *Stokes sector*, if it contains one imaginary dividing ray and its closure does not contain the other one.

Let $W(z) = \text{diag}(e^{-\frac{\lambda_{10}}{z}} z^{b_{10}}, e^{-\frac{\lambda_{20}}{z}} z^{b_{20}})$ denote the canonical diagonal fundamental matrix solution of the formal normal form (2.2). The matrices X_0^j $\iint_0^j(z) := H_{j,0}(z)W(z)$ are fundamental matrix solutions of the initial system (2.1) defining solution bases in S_j^0 called the canonical sectorial solution bases. Here we choose the branches $W_0(z) = W_0^j$ $C_0^{\jmath}(z)$ of the matrix function $W_0(z)$ in S_j^0 so that $W_0^1(z)$ is obtained from $W_0^0(z)$ by counterclockwise analytic extension from S_0^0 to S_1^0 . And we define the branch $W_0^2(z)$ of $W_0(z)$ in $S_2^0 := S_0^0$ obtained from $W_0^1(z)$ by counterclockwise analytic extension from S_1^0 to S_0^0 . This yields another canonical matrix solution $X_0^2(z) := H_{0,0}(z)W_0^2(z)$ of system (2.1) in S_0^0 , which is obtained from $X_0^0(z)$ by multiplication from the right by the monodromy matrix $\exp(2\pi i \tilde{R}_0)$ of the formal normal form (2.2). Let $S^0_{j,j+1}$ denote the connected component of the intersection $S_{j+1}^0 \cap S_j^0$, $j = 0, 1$, that is crossed when one moves from S_j^0 to S_{j+1}^0 counterclockwise, see Fig. 3. The transition matrices C_{00} , C_{10} between thus defined canonical solution bases X_0^j $\frac{\jmath}{0},$

$$
X_0^1(z) = X_0^0(z)C_{00} \text{ on } S_{0,1}^0, \quad X_0^2(z) = X_0^1(z)C_{10} \text{ on } S_{1,2}^0,
$$
\n(2.5)

are called the Stokes matrices at the origin.

The above formal normal form and sectorial normalization theorems also hold at infinity. Namely, the formal normal form at infinity is

$$
\widetilde{Y}' = \left(\widetilde{N} + \frac{\widetilde{R}_{\infty}}{z}\right)\widetilde{Y}, \ \widetilde{N} = \text{diag}(\lambda_{1\infty}, \lambda_{2\infty}), \ \widetilde{R}_{\infty} = \text{diag}(b_{1\infty}, b_{2\infty}).\tag{2.6}
$$

Here $\lambda_{1\infty}$ and $\lambda_{2\infty}$ are the eigenvalues of the matrix N.

In more detail, let H_{∞} be a matrix such that $H_{\infty}^{-1}NH_{\infty} = \text{diag}(\lambda_{1\infty}, \lambda_{2\infty})$. Then

 \widetilde{R}_{∞} is the diagonal part of the matrix $\mathbf{H}_{\infty}^{-1} R \mathbf{H}_{\infty}$.

The Stokes rays at infinity are defined by the equation

$$
\{\operatorname{Re}(\lambda_{1\infty} - \lambda_{2\infty})z = 0\}.
$$
\n(2.7)

Example 2.1 Let $K = \text{diag}(\lambda_{10}, \lambda_{20})$, and let $\lambda_{20} - \lambda_{10} > 0$. Then the Stokes rays are the positive and negative imaginary semiaxes. The Stokes sectors S_0^0 and S_1^0 covering \mathbb{C}^* satisfy the following conditions:

- the sector S_0^0 contains the positive imaginary semiaxis, and its closure does not contain the negative one;

- the sector S_1^0 satisfies the opposite condition. See Fig. 3.

The Stokes matrices C_{00} and C_{10} are unipotent upper and lower triangular respectively. Case of singular point at infinity with $\lambda_{2\infty} - \lambda_{1\infty} > 0$ is treated analogously. The Stokes rays are the same, and the Stokes sectors also can be choosen the same, as above. But at infinity, the Stokes matrices $C_{0\infty}$, $C_{1\infty}$ are oppositely-triangular: unipotent lower (respectively, upper) triangular.

Figure 3: Stokes sectors at the origin in the case, when $\lambda_{10} - \lambda_{20} \in \mathbb{R}$.

Recall that the monodromy operator of system (2.1) acts on the space of local solutions (initial conditions) at a base point $z_0 \in \mathbb{C}^*$ (or on the space of solutions in a given simply connected domain (sector) in C ∗) by analytic extension along a counterclockwise circuit around the origin. It is well-known that in the canonical basis in the sector S_0^p p_0^p at $p \in 0, \infty$ the monodromy matrix is given by the following formula, see [33, p.290]:

$$
M = M_{norm,p} C_{1p}^{-1} C_{0p}^{-1},
$$
\n
$$
M_{norm,p} = \text{diag}(e^{2\pi i b_{1p}}, e^{2\pi i b_{2p}})
$$
\n(2.8)

is the monodromy matrix of the formal normal form at p . Formula (2.8) together with unipotence of the Stokes matrices implies that

$$
\det M = \det M_{norm,p} = \exp(2\pi i \operatorname{Tr} R). \tag{2.9}
$$

Remark 2.2 (see [8, remark 2.3]). The canonical solutions $Y(z)$ of system (2.1) on a Stokes sector S_i^p $_j^p$ (which are the columns of the fundamental matrix solution X_p^j) are numerated by the eigenvalues λ_{sp} of its main term matrix at p (or equivalently, by the eigenvalues b_{sp} of the residue matrix of the corresponding formal normal form). The projectivized canonical sectorial solutions $\Phi(z) = \frac{Y_2(z)}{Y_1(z)}$ (which take values in $\mathbb{CP}^1 = \overline{\mathbb{C}}_{\Phi}$) are the unique solutions of the Riccati equation corresponding to (2.1) that are holomorphic on S_i^p $_j^p$ and extend continuously to p. The value of an extended projectivized canonical solution at p is the tautological projection to $\mathbb{CP}^1 = \overline{\mathbb{C}}$ of the eigenline L_{sp} of the main term matrix corresponding to the eigenvalue λ_{sp} : the eigenvalue numerating the solution of linear system.

Definition 2.3 [8, definition 2.7]. Consider a linear system \mathcal{L} of type (2.1). Fix a point $z_0 \in \mathbb{C}^*$ and two pairs of Stokes sectors $(S_0^0, S_1^0), (S_0^{\infty}, S_1^{\infty})$ for the main term matrices K at 0 and N at ∞ respectively. Fix two paths α_p in \mathbb{C}^* numerated by $p = 0, \infty$, going from the point z_0 to a point in S_0^p \int_0^p . Let f_{1p} , f_{2p} be a canonical sectorial solution basis for the system $\mathcal L$ at p in S_0^p $_{0}^{p}$. Consider the analytic extensions of the basic functions f_{kp} to the point z_0 along the paths α_p^{-1} . Let $\pi : \mathbb{C}^2 \setminus \{0\} \to \mathbb{CP}^1$ denote the tautological projection. Set $\Phi := \frac{Y_2}{Y_1},$

$$
q_{kp} := \pi(f_{kp}(z_0)) \in \mathbb{CP}^1 = \overline{\mathbb{C}}_{\Phi}.
$$
\n(2.10)

Let M denote the monodromy operator of the system $\mathcal L$ acting on the local solution space at z_0 (identified with the space \mathbb{C}^2 of initial conditions at z_0) by analytic extension along counterclockwise circuit around zero. The tuple

$$
(q, M) := (q_{10}, q_{20}, q_{1\infty}, q_{2\infty}; M)
$$
\n
$$
(2.11)
$$

taken up to the next equivalence is called the monodromy–Stokes data of the system \mathcal{L} . Namely, two tuples $(q, M), (q', M') \in (\mathbb{CP}^1)^4 \times GL_2(\mathbb{C})$ are called *equivalent*³, if there exists a linear operator $H \in GL_2(\mathbb{C})$ whose projectivization sends q_{kp} to q'_{kp} and such that H^{-1} \circ $M' \circ H = M$.

Remark 2.4 [8, remark 2.8] The monodromy–Stokes data of a system (2.1) depends only on the homotopy class of the pair of paths $(\alpha_0, \alpha_\infty)$ in the space of pairs of paths in \mathbb{C}^* with a common (variable) starting point z_0 and with endpoints lying in given sectors S_0^0 and S_0^{∞} respectively.

Here and in what follows two linear systems on \overline{C} are said to be *gauge equivalent*, if they are "constant gauge equivalent": obtained one from the other by constant linear variable change $Y \mapsto HY$, $H \in GL_2(\mathbb{C})$. Gauge equivalence of holomorphic vector bundles with meromorphic connections is defined analogously, with H being a holomorphic bundle isomorphism.

Theorem 2.5 [8, theorem 2.11]. Two linear systems of type (2.1) are gauge equivalent, if and only if they have the same formal normal forms at each singular point and the same monodromy-Stokes data (defined by a given collection of sectors and paths). In this case each linear automorphism of the fiber $\mathbb{C}^2 \times \{z_0\}$ sending the monodromy-Stokes data of one system to that of the other system extends to a gauge equivalence of systems. Here both monodromy-Stokes data correspond to the same sectors and path collections.

Definition 2.6 A family of linear systems (2.1) with irregular non-resonant singular points at zero and at infinity is isomonodromic, if the eigenvalues of the residue matrices of formal normal forms at their singular points and the monodromy–Stokes data remain constant: independent on the parameter of the family.

³Here is an equivalent group-action definition. The group $PSL_2(\mathbb{C})$ acts on $\overline{\mathbb{C}}^4 \times GL_2(\mathbb{C})$ by action $h: q_{kp} \mapsto hq_{kp}$ on points in $\overline{\mathbb{C}} = \mathbb{CP}^1$ and conjugation $M \mapsto hMh^{-1}$ on matrices. The monodromy-Stokes data is the $PSL_2(\mathbb{C})$ -orbit of a collection (q, M) under this action.

Remark 2.7 In the case under consideration (two singularities of Poincaré rank 1) the above definition given in [8, definition 3.1] is known to be equivalent to the classical definition saying that a family of systems is isomonodromic, if the monodromy matrix (in appropriate canonical sectorial basis), Stokes matrices, the transition matrix between two canonical bases at zero and at infinity, and the above residue eigenvalues in normal forms are constant: see [8, proposition 2.9]. It is well-known that if a family of systems (2.1) depends continuously on a parameter from a connected manifold, then constance of the monodromy-Stokes data automatically implies constance of the residue eigenvalues of the formal normal forms.

2.2 Degree of polynomial solution. Proof of Proposition 1.9

Proposition 2.8 A system (2.1) with Poincaré rank 1 irregular non-resonant singular points at 0 and at ∞ has a vector polynomial solution, if and only if each its main term matrix at $0, \infty$ has a zero eigenvalue, the corresponding residue eigenvalue of the formal normal forms is zero at 0 and a number $\ell \in \mathbb{Z}_{\geq 0}$ at ∞ , and the corresponding canonical sectorial solutions at zero and at infinity coincide (up to constant factor) and holomorphic in \mathbb{C}^* . In this case the polynomial solution is unique, coincides with the above canonical solution and has degree ℓ .

Proof At each singular point $p = 0$, ∞ the corresponding main term matrix has at least one non-zero eigenvalue λ (non-resonance). In each Stokes sector S the corresponding canonical solution has type of an exponential function times a monomial (of some not necessary integer power) times a function holomorphic in S that extends continuously to p . The Stokes sector S contains a subsector S' (lying on appropriate side from the Stokes ray contained in S) where the exponential function tends to infinity exponentially, as $z \to p$, $z \in S'$. Therefore, if a polynomial solution exists, then its restriction to S' has smaller growth. In other words it is asymptotically dominated by the above "exponential" canonical solution. But it is a well-known fact from the Stokes phenomena theory that if in some sector S' one solution of a two-dimensional linear system (2.1) dominates another solution, then the smaller, dominated solution is a canonical sectorial solution, and it is canonical in each Stokes sector containing S' . Namely this fact explains unipotence and triangularity of the Stokes matrices. Therefore, the polynomial solution is canonical at both singular points $p = 0$, ∞ . It corresponds to zero eigenvalue of the main term matrices, being meromorphic, and the corresponding residue eigenvalues of the formal normal forms are clearly zero at 0 and a number $\ell \in \mathbb{Z}_{\geq 0}$ at infinity (equal to the degree of solution). Conversely, if there exists a solution holomorphic on \mathbb{C}^* that is canonical at each singular point and corresponds to zero eigenvalue of both main term matrices and the above-mentioned residue eigenvalues of the formal normal form, then it is clearly polynomial of degree ℓ . Uniqueness of polynomial solution follows from the above exponential growth and domination argument. Proposition 2.8 is proved. \Box

Proposition 1.9 follows immediately from Proposition 2.8.

2.3 Symmetric systems of class S and their monodromy–Stokes data

For a 2 \times 2-matrix K let K^{tt} denote the matrix obtained from K by conjugation by the transposition matrix $\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, i.e., obtained from K by transposition of diagonal terms and transposition of off-diagonal terms.

Definition 2.9 A linear system is of class S , if it has the type

$$
Y' = \left(\frac{K}{z^2} + \frac{R}{z} - K^{tt}\right)Y, \quad R_{12} = -R_{21}.\tag{2.12}
$$

Example 2.10 Every system (1.35) is of type S.

Proposition 2.11 A linear system (2.1) has type S, if and only if it admits a symmetry of the type

$$
\mathbb{I}: (Y_1(z), Y_2(z)) \mapsto z^{\ell}(Y_2(z^{-1}), Y_1(z^{-1})). \tag{2.13}
$$

In this case one has

$$
\ell = \text{Tr}\,R. \tag{2.14}
$$

Proof The transformation $(Y_1, Y_2, z) \mapsto (Y_2, Y_1, z^{-1})$ sends graphs of solutions of a system (2.1) with matrices K, N, R to graphs of solutions of the new system of similar kind with the new matrices

$$
\widehat{K} = -PNP^{-1}, \ \widehat{N} = -PKP^{-1}, \ \widehat{R} = -PRP^{-1}; \quad P := \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}.
$$
 (2.15)

Multiplication by z^{ℓ} adds ℓId to R. The statement saying that I is a symmetry of the initial equation (2.1) is equivalent to the equality that the new matrices thus obtained are equal respectively to K , N and R :

$$
K = -PNP^{-1} = -N^{tt}, \ R = -PRP^{-1} + \ell Id = -R^{tt} + \ell Id. \tag{2.16}
$$

The last equation is equivalent to the fact that the off-diagonal elements of the matrix R are opposite and $\ell = \text{Tr } R$. This proves the proposition. \Box

Recall that the canonical basic solutions of a system (2.1) at a point $p \in \{0,\infty\}$ in a Stokes sector S_i^p ^p are numerated by the eigenvalues λ_{1p} , λ_{2p} of the corresponding main term matrix K (respectively, N). Or equivalently, by the eigenvalues b_{1p} , b_{2p} of the residue matrix⁴ of the corresponding formal normal form (2.2) , see Remark 2.2.

 4 The main term matrix eigenvalues are always distinct (the singularity in question in non-resonant), and hence, numerate the canonical basic solutions. But the residue matrix eigenvalues may coincide. One can say that they numerate the canonical basic solutions, only if they are distinct.

Definition 2.12 Let a linear system (2.1) have irregular non-resonant singular points of Poincaré rank 1 at the origin and at infinity. Let $(q; M) = (q_{10}, q_{20}, q_{1\infty}, q_{2\infty}; M)$ be its monodromy-Stokes data defined by the base point $z_0 = 1$, sectors S_0^0 and S_0^{∞} that are permuted by the involution $z \mapsto z^{-1}$, and paths α_0 , α_∞ going from z_0 to S_0^0 and S_0^∞ respectively and permuted by the same involution. (The latter pairs of sectors and paths will be called *symmetric*.) We say that the monodromy-Stokes data (q, M) is $(1, 1)$ -symmetric, if there exists a linear involution $\sigma : \mathbb{C}^2 \to \mathbb{C}^2$ with eigenvalues 1, -1 (i.e., conjugated to the above permutation matrix P) that conjugates the monodromy transformation with a scalar multiple of its inverse, and such that the projectivization of the involution σ permutes q_{10} and $q_{2\infty}$ and permutes q_{20} and $q_{1\infty}$. In this case

$$
\sigma M \sigma^{-1} = \exp(2\pi i \operatorname{Tr} R) M^{-1},
$$

where R is the residue matrix of the system (2.1) , which follows from (2.9) .

Remark 2.13 Let in a system (2.1) the differences $\lambda_{20} - \lambda_{10}$, $\lambda_{2\infty} - \lambda_{1\infty}$ of the eigenvalues in the main term matrices K and N be real-proportional. (This is true, e.g., for systems of type S, where the eigenvalue collections differ by sign.) Then the Stokes ray collections at the origin and at infinity are permuted by the involution $z \mapsto z^{-1}$. This implies that the Stokes sector collections at 0 and at ∞ can be chosen so that S_j^0 and S_j^{∞} be also permuted by the same involution for every $j = 0, 1$. Therefore, one can define the monodromy-Stokes data by paths α_0 and α_∞ permuted by the same involution, and the notion of (1,1)-symmetric monodromy-Stokes data is well-defined.

Theorem 2.14 A linear system (2.1) is gauge equivalent to a linear system of type S, if and only if

$$
(\lambda_{1\infty}, \lambda_{2\infty}; b_{1\infty}, b_{2\infty}) = (-\lambda_{20}, -\lambda_{10}; b_{10}, b_{20}), \qquad (2.17)
$$

and its monodromy-Stokes data $(q_{10}, q_{20}, q_{1\infty}, q_{2\infty}; M)$ (numerated by the eigenvalues λ_{10} , λ_{20} , $\lambda_{1\infty}$, $\lambda_{2\infty}$ and constructed from some given symmetric pairs of sectors and paths) is $(1,1)$ -symmetric. (In this case it is $(1,1)$ -symmetric for every symmetric pairs of sectors and paths.) By definition, equality (2.17) means that its formal normal forms at the origin and at the infinity are respectively

$$
Y' = \left(\frac{\operatorname{diag}(\lambda_{10}, \lambda_{20})}{z^2} + \frac{\operatorname{diag}(b_{10}, b_{20})}{z}\right) Y \quad at \quad 0,
$$

$$
Y' = \left(-\operatorname{diag}(\lambda_{20}, \lambda_{10}) + \frac{\operatorname{diag}(b_{10}, b_{20})}{z}\right) Y \quad at \quad \infty.
$$

Proof Let us prove the statements of the theorem for a system of type S. First let us prove (2.17). Equality $\{\lambda_{1\infty}, \lambda_{2\infty}\} = \{-\lambda_{20}, -\lambda_{10}\}$ follows from definition. Let us numerate the latter eigenvalues so that $\lambda_{1\infty} = -\lambda_{20}$, $\lambda_{2\infty} = -\lambda_{10}$. Let H be the matrix diagonalizing K: $H^{-1}KH = \text{diag}(\lambda_{10}, \lambda_{20})$, and let P be the permutation matrix, see (2.15). One has $P^2 = Id$. The matrix *PHP* diagonalizes the matrix $N = -K^{tt} = -PKP$:

$$
P^{-1}H^{-1}P^{-1}NPHP = -P^{-1}H^{-1}KHP = -\operatorname{diag}(\lambda_{20}, \lambda_{10}) = \operatorname{diag}(\lambda_{1\infty}, \lambda_{2\infty}).
$$

One has

 $diag(b_{10}, b_{20}) =$ the diagonal part of $H^{-1}RH$,

 $diag(b_{1\infty}, b_{2\infty}) =$ the diagonal part of $P^{-1}H^{-1}P^{-1}RPHP$.

One has $P^{-1}RP = -R + (\text{Tr } R)Id$, by definition and since $R_{12} = -R_{21}$. Therefore,

diag
$$
(b_{1\infty}, b_{2\infty})
$$
 = –(the diagonal part of $P^{-1}H^{-1}RHP$) + (Tr R)Id

$$
= -\operatorname{diag}(b_{20}, b_{10}) + (\operatorname{Tr} R)Id = \operatorname{diag}(b_{10}, b_{20}).
$$

Let us prove symmetry of the monodromy-Stokes data. First let us show that

$$
PMP^{-1} = e^{2\pi i\ell}M^{-1}, \quad \ell = \text{Tr}\,R. \tag{2.18}
$$

Indeed, the monodromy operator M acts on the space of initial conditions in $\mathbb{C}^2 = \mathbb{C}^2 \times \{1\}$ at $z_0 = 1$ so that an initial condition Y^0 is sent to the analytic extension of the corresponding germ of solution to 1 along a counterclockwise circuit around the origin. Let us write its matrix in the standard basis $(1, 0)$, $(0, 1)$. The transformation (2.13) is a symmetry of system (2.1) (Proposition 2.11). It is the composition of the three following operations:

(i) the space variable change $Y \mapsto PY$ (constant gauge transformation), it conjugates the monodromy by the involution P ;

(ii) the time variable change $z \mapsto z^{-1}$, it changes the monodromy to its inverse;

(iii) multiplication by the function z^{ℓ} , it multiplies the monodromy by the scalar factor $e^{2\pi i \ell}$.

Therefore, the transformation $\mathbb I$ given by (2.13) on the space of vector functions sends solutions of the initial system (with monodromy matrix M) to solutions of the new system whose monodromy matrix is equal to $e^{2\pi i \ell} P M^{-1} P^{-1}$. The latter matrix coincides with M, since both systems coincide ($\mathbb I$ is a symmetry). This proves (2.18).

Let us now prove that the transformation I sends canonical solutions to canonical solutions. For every $p = 0, \infty$ the transformation (i) sends canonical solutions at p of the initial system to canonical solutions of its image at the same point p , and the numeration by eigenvalues λ_{jp} is preserved. The transformation (ii) sends a canonical solution at 0 corresponding to the eigenvalue λ_{j0} in a Stokes sector S^0 to the canonical solution of the transformed system at ∞ in the Stokes sector S^{∞} that is the image of the sector S^0 under the involution $z \mapsto z^{-1}$, and the corresponding eigenvalue of the main term matrix at infinity is equal to $-\lambda_{i0} = \lambda_{i} - i\infty$. This follows from Remark 2.2. This implies that the projectivization of the operator P is an involution $\overline{\mathbb{C}}_{\Phi} \to \overline{\mathbb{C}}_{\Phi}$ permuting q_{10} and $q_{2\infty}$ and permuting q_{20} and $q_{1\infty}$. This together with (2.18) implies symmetry of the monodromy-Stokes data.

Let us now prove the converse. Let for a given system (2.1) the eigenvalue equalities (2.17) hold, and let the monodromy-Stokes data defined by some symmetric pairs of sectors and paths be (1,1)-symmetric. Let us prove that (2.1) has type S. Let $\sigma : \mathbb{C}^2 \to \mathbb{C}^2$ be the corresponding linear involution making their symmetry. Applying a gauge transformation we can and will consider that $\sigma = P$. Set $\ell = \text{Tr } R$. Then transformation (2.13) sends the initial system (2.1) to a system of similar kind, let us denote it by (2.1) ', with the same monodromy-Stokes data. This follows from the above discussion on action of the operations (i)–(iii). At each singular point the eigenvalues of the main term matrices in both systems are the same, and the eigenvalues of the residue matrices of the formal normal form are the same. In both systems they numerate the monodromy-Stokes data in the same way. This follows by construction and (2.17) . Finally, systems (2.1) and $(2.1)'$ have the same formal normal forms and monodromy-Stokes data. Hence, they are gauge equivalent: obtained one from the other by constant linear change $Y \mapsto HY$, $H \in GL_2(\mathbb{C})$ (Theorem 2.5).

We choose coordinates (Y_1, Y_2) in \mathbb{C}^2 , let $\Phi = \frac{Y_2}{Y_1}$ denote the corresponding coordinate on $\mathbb{CP}^1 = \overline{\mathbb{C}}$, so that the representatives $(q_{10}, q_{20}, q_{1\infty}, q_{2\infty}; M)$ of the monodromy-Stokes data of systems (2.1) and $(2.1)'$ coincide. Let us show that H is scalar: a multiple of the identity. The action of the operator H on the space of initial conditions $\mathbb{C}^2 = \mathbb{C}^2 \times \{1\}$ preserves the monodromy-Stokes data. This implies that its projectivization preserves the points q_{ip} and their images under the projectivized monodromy: the monodromy of the Riccati equation corresponding to (2.1). Let the total number of distinct points among the points q_{ip} and their latter images be greater than two. Then the projectivized operator H preserves at least three distinct points in $\overline{\mathbb{C}}$, and hence, is identity. Thus H is scalar, and the above variable change multiplies Y by a constant. This scalar change does not change the underlying linear system. Hence, systems (2.1) and $(2.1)'$ coincide, and thus, the transformation I, see (2.13) , is a symmetry of system (2.1) . Therefore, (2.1) is a system of type S (Proposition 2.11).

Let us consider the opposite case: the above number of distinct points is at most two. For every given $p = 0$, ∞ the points q_{1p} , q_{2p} are distinct, being projectivized initial conditions of two linearly independent basic solutions of (2.1). This implies that $\{q_{10}, q_{20}\} = \{q_{1\infty}, q_{2\infty}\},$ and the projectivized monodromy either preserves each q_{jp} , or permutes q_{1p} and q_{2p} . The second case is impossible. Indeed, otherwise, if q_{10} and q_{20} were permuted by the projectivized monodromy, then in the canonical basis at 0 the monodromy matrix of system (2.1) would have zero diagonal terms. But at least one of them should be nonzero and equal to an eigenvalue $e^{2\pi i b_{j0}}$ of the monodromy of the formal normal form (follows from (2.8) , opposite triangularity and unipotence of the Stokes matrices), – a contradiction. In the first case, when the projectivized monodromy preserves each q_{ip} , the monodromy of linear system is diagonal in the canonical sectorial basis at p. Hence, the corresponding Stokes operators are trivial, by their opposite triangularity and (2.8) (as in [28, proof of lemma 3.3]). And this holds at both points $p = 0, \infty$. This together with Remark 2.2 implies that the projectivized canonical solutions $g_{jp}(z) = \pi \circ f_{jp}(z)$, i.e., the solutions of the Riccati equation with initial conditions q_{i0} at $z_0 = 1$, are meromorphic on the whole Riemann sphere, and their values at each point $p = 0$, ∞ are the two distinct eigenlines of the corresponding main term matrix (K or N). This implies that the graphs $\{\Phi = g_{i0}(z)\}\$, $j = 1, 2$, are disjoint rational curves in $\overline{\mathbb{C}}_{\Phi} \times \overline{\mathbb{C}}_z$ parametrized by $\overline{\mathbb{C}}_z$ via projection. But they may be disjoint only if $g_{i0} \equiv const.$ Indeed, otherwise, if some g_{i0} is non-constant, then the intersection index of graphs would be positive, by K¨unneth Formula for product homology and bi-degree argument, see the proof of the same statement at the end of $[8, \text{proof of proposition } 4.8]$, – a contradiction. Hence, $g_{10}(z)$ and $g_{20}(z)$ are constant. Applying a gauge transformation, we can and will consider that $g_{10} \equiv 0$, $g_{2,0} \equiv \infty$. Then the matrices in linear system (2.1) are diagonal, as in loc cit, and the system is obviously of type S. Theorem 2.14 is proved in loc. cit., and the system is obviously of type S. Theorem 2.14 is proved.

Lemma 2.15 A system of type S where the main term matrix K has distinct eigenvalues (non-resonance condition) is gauge equivalent to a system of type (1.35) , if and only if in (2.17) one has $\lambda_{20} = b_{20} = 0$ and the following line non-coincidence condition holds:

(*) the eigenlines L_{10} , $L_{2\infty}$ of the main term matrices at 0 and at ∞ respectively corresponding to non-zero eigenvalues λ_{10} and $\lambda_{2\infty} = -\lambda_{10}$ are distinct.

Proof Each system (1.35) obviously satisfies equality $\lambda_{20} = b_{20} = 0$ and condition (*). Let us prove the converse. The permutation operator P permutes the eigenlines L_{10} and $L_{2\infty}$, see (2.16) and the above proof of the first part of Theorem 2.14. Applying a gauge transformation $Y \rightarrow HY$ with H commuting with P preserves the class S of linear systems, since it does not change relations (2.16) . We can apply H commuting with P that sends the eigenline L_{10} to the first coordinate line, and then $L_{2\infty}$ is sent to the second coordinate line, since they are distinct, by $(*)$. This yields another system of class S with upper-triangular matrix K and lower-triangular matrix $N = -K^{tt}$. Therefore, the new system has type (1.35), by (2.16). This proves the lemma. (1.35) , by (2.16) . This proves the lemma.

Corollary 2.16 A system (2.1) is gauge equivalent to a system of type (1.35) , if and only if the following conditions hold:

a) equalities (2.17) hold, and $\lambda_{20} = \lambda_{1\infty} = 0$, $b_{20} = b_{2\infty} = 0$;

b) the monodromy-Stokes data constructed from some given symmetric pairs of sectors and paths is $(1,1)$ -symmetric;

c) the above line non-coincidence condition $(*)$ holds.

In this case monodromy-Stokes data defined by arbitrary symmetric pairs of sectors and paths are $(1,1)$ -symmetric.

Proof Each system (1.35) obviously satisfies statements a) and c), and statement b) follows by Theorem 2.14. Conversely, let a system (2.1) satisfy conditions a)–c). Then it is gauge equivalent to a system of class S, by symmetry of monodromy-Stokes data and Theorem 2.14. The latter system is gauge equivalent to a system of type (1.35), by Lemma 2.15. The corollary is proved.

2.4 Polynomial solutions and determinantal surfaces. Proof of Theorem 1.12.

Everywhere below we consider that $s \neq 0$.

Proposition 2.17 Let $\ell \in \mathbb{Z}_{\geq 0}$. The transformation (2.13) is an involution sending each vector polynomial solution $Y(z)$ of system (1.35) (if any) to $\pm Y(z)$.

Proof The fact that (2.13) is an involution is obvious. It is a symmetry of system (1.35), by Proposition 2.11 and since (1.35) is a system of type S, see Example 2.10. If a polynomial solution exists, then it is unique (up to constant factor) and has degree ℓ (Proposition 2.8). The transformation (2.13) preserves the space of vector polynomials of degree at most ℓ and acts there as a linear involution sending a solution of (1.35) to a solution. Therefore, the unique polynomial solution is its eigenvector with eigenvalue ± 1 .

Proposition 2.18 A vector polynomial $Y(z) = (Y_1(z), Y_2(z))$ is a solution of linear system (1.5) , if and only if it satisfies the system of two following equations with at least one choice of sign \pm :

$$
(Y_1(z), Y_2(z)) = \pm z^{\ell} (Y_2(z^{-1}), Y_1(z^{-1}))
$$
\n(2.19)

$$
Y_2'(z) = \pm \left(\frac{a}{2z} + s\chi\right)z^{\ell}Y_2(z^{-1}) + \left(\frac{\chi a}{z} + \frac{s}{2}\right)Y_2(z)
$$
\n(2.20)

Proof Let a vector polynomial $Y(z)$ be a solution of system (1.5). Then equation (2.19) holds for some choice of sign \pm , by Proposition 2.17. Substituting thus found $Y_1(z)$ = $\pm z^{\ell} Y_2(z^{-1})$ to the second line of (1.5) yields (2.20). Conversely, let a vector polynomial $Y(z)$ be a solution of system of equations (2.19) and (2.20). Then the second line of system (1.5) holds, as in the above argument. It remains to show that its first line holds as well. Differentiating $Y_1(z) = \pm z^{\ell} Y_2(z^{-1})$ and substituting (2.20) yields

$$
Y_1'(z) = \frac{\ell}{z} Y_1(z) \mp z^{\ell-2} Y_2'(z^{-1}),
$$
\n
$$
\mp z^{\ell-2} Y_2'(z^{-1}) = -z^{\ell-2} (\frac{az}{2} + s\chi) z^{-\ell} Y_2(z) \mp z^{\ell-2} (\chi az + \frac{s}{2}) Y_2(z^{-1}).
$$
\n(2.21)

Substituting the equality $Y_2(z^{-1}) = \pm z^{-\ell} Y_1(z)$ to the latter right-hand side and substituting everything to (2.21) yields the first line in (1.5):

$$
Y_1'(z) = \left(\frac{\ell - \chi a}{z} - \frac{s}{2z^2}\right) Y_1(z) - \left(\frac{a}{2z} + \frac{s\chi}{z^2}\right) Y_2(z).
$$

Proof of Theorem 1.12. It suffices to show that existence of polynomial solution of system of equations (2.19) and (2.20) for a given choice of sign \pm is equivalent to the inclusion $(\chi, a, s) \in S_{\ell, \pm}$. The space of polynomial solutions of equation (2.19) is isomorphic to the space of polynomials $Y_2(z)$ of degree at most ℓ via the correspondence $Y_1(z) = \pm z^{\ell} Y_2(z^{-1})$. A given polynomial $Y_2(z) = \sum_{j=0}^{\ell} c_j z^j$ of degree at most ℓ is a solution of (2.20), if and only if its coefficients satisfy the following relations, corresponding to equality of degree j terms in (2.20) for $j = 0, \ldots, \ell$:

$$
\pm(\frac{a}{2}c_{\ell-(j+1)} + s\chi c_{\ell-j}) + (\chi a - (j+1))c_{j+1} + \frac{s}{2}c_j = 0.
$$
 (2.22)

Here we set $c_k = 0$ for $k \notin \{0, \ldots, \ell\}$. System of relations (2.22) is equivalent to the statement that the vector of coefficients (c_0, \ldots, c_ℓ) lies in the kernel of the matrix $G_{1,\ell} \pm G_{2,\ell}$. Therefore, existence of polynomial solution of system of equations (2.19), (2.20) is equivalent to the statement that the latter matrix has zero determinant, i.e., $(\chi, a, s) \in S_{\ell, \pm}$. This together with Proposition 2.18 implies the statement of Theorem 1.12. together with Proposition 2.18 implies the statement of Theorem 1.12.

Remark 2.19 The above proof of Theorem 1.12 remains valid for $\ell = 0$ and yields a proof of Proposition 1.10 for $s \neq 0$. Indeed, for $\ell = 0$ equation (2.22) becomes $\pm s\chi c_0 + \frac{s}{2}$ $\frac{s}{2}c_0 = 0$ and has a non-zero solution, if and only if $\chi = \pm \frac{1}{2}$ $rac{1}{2}$.

2.5 Systems with polynomial solutions. Parabolicity of the monodromy-Stokes data

Definition 2.20 We deal with an abstract monodromy-Stokes data: a collection $(q_{10}, q_{20}, q_{1\infty}, q_{2\infty}; M)$, where $q_{jp} \in \overline{\mathbb{C}}_{\Phi} = \mathbb{CP}^1_{[Y_1:Y_2]}, \ \Phi = \frac{Y_2}{Y_1}, \ q_{10} \neq q_{20}, \ q_{1\infty} \neq q_{2\infty}, \text{ and } M \in GL_2(\mathbb{C}) \text{ is a non$ degenerate linear operator acting on the space $\mathbb{C}^2_{Y_1,Y_2}$, called the *monodromy operator*. The above collection is taken up to equivalence from the definition of the monodromy-Stokes data. We say that an abstract monodromy-Stokes data is *parabolic*, if $q_{20} = q_{1\infty}$, and the projectivized monodromy operator is either a parabolic transformation $\overline{\mathbb{C}} \to \overline{\mathbb{C}}$ that fixes only the point q_{20} , or the identity. Set q'_{10} to be the image of the point q_{10} under the projectivized monodromy transformation. A parabolic monodromy-Stokes data is called *non-degenerate*, if among the four points q'_{10} , q_{10} , q_{20} , $q_{2\infty}$ at least three points are distinct.

Remark 2.21 Consider a monodromy-Stokes data of linear system (2.1) defined by symmetric pairs of sectors and paths. If the monodromy-Stokes data is parabolic, then it is automatically (1,1)-symmetric. This follows from the fact that every conformal involution $\sigma : \overline{\mathbb{C}} \to \overline{\mathbb{C}}$ fixing $q_{20} = q_{1\infty}$ (e.g., the one permuting q_{10} and $q_{2\infty}$) conjugates each parabolic transformation fixing q_{20} (e.g., the projectivized monodromy) with its inverse. Namely, if we choose the coordinate w on $\overline{\mathbb{C}}$ so that $q_{20} = \infty$ and the origin is the second fixed point of the involution σ , then $\sigma(w) = -w$, the projectivized monodromy is a translation $w \mapsto w+u$, and hence, σ conjugates it with its inverse $w \mapsto w - u$.

Here we prove the following theorem.

Theorem 2.22 A system (2.1) with irregular non-resonant singular points of Poincaré rank 1 at 0 and at ∞ is gauge equivalent to a system of type (1.5) having a non-constant polynomial solution, if and only if it satisfies the following conditions:

a) $(\lambda_{1\infty}, \lambda_{2\infty}; b_{1\infty}, b_{2\infty}) = (-\lambda_{20}, -\lambda_{10}; b_{10}, b_{20}), \lambda_{20} = \lambda_{1\infty} = 0, b_{20} = b_{2\infty} = 0.$

b) $\ell := b_{10} = b_{1\infty} \in \mathbb{N}$.

c) The monodromy-Stokes data defined by sectors and paths permuted by the involution $z \mapsto \frac{1}{z}$ is parabolic non-degenerate.

 $d\tilde{d}$) The above line non-coincidence condition $(*)$ holds.

In this case Statement c) holds for monodromy–Stokes data defined by arbitrary symmetric pairs of sectors and paths.

Proof First let us prove that each system (1.5) with a non-constant polynomial solution satisfies statements a $(-d)$. Statements a) and d) hold for each system (1.5) . Statement b) follows from Proposition 1.9. Let us now prove parabolicity of the monodromy-Stokes data. The monodromy operator has unit determinant, since it is the product of unipotent matrices and the monodromy of the formal normal form, i.e., the formal monodromy, see (2.8), and the formal monodromy is equal to diag($e^{2\pi i \ell}$, 1) = Id, since $\ell \in \mathbb{N}$. Hence, the monodromy is unipotent: it has unit determinant and at least one unit eigenvalue, which corresponds to the polynomial solution (fixed by the monodromy). The value at $z = 1$ of the projectivized polynomial solution is equal to $q_{20} = q_{1\infty}$, and it is a fixed point of the projectivized monodromy, by the above similar statement on the monodromy. This together with unipotence of the monodromy implies parabolicity of the monodromy-Stokes data. Let us prove its non-degeneracy.

Suppose the contrary: among the points q_{jp}, q'_{10} at most two are distinct. Then $q_{10} =$ $q_{2\infty} = q'_{10}$, since $q_{10} \neq q_{20} = q_{1\infty}, q_{1\infty} \neq q_{2\infty}, q_{20}$ is a fixed point of the projectivized monodromy transformation, and thus, the image q'_{10} of the point q_{10} under the projectivized monodromy is distinct from q_{20} . Therefore, $q_{10} = q_{2\infty}$, $q_{20} = q_{1\infty}$, and they are fixed by the projectivized monodromy. Thus, the corresponding linear system is gauge equivalent to a diagonal system, as at the end of the proof of Theorem 2.14. Hence, $\chi = a = 0$, and thus, system (1.5) is diagonal. But then it has no non-constant polynomial solution, since it is the direct sum of one-dimensional equations $Y'_1 = \left(-\frac{s}{2z}\right)$ $rac{s}{2z^2} + \frac{\ell - \chi a}{z}$ $(\frac{x}{z})Y_1, Y_2' = (\frac{x}{z} + \frac{s}{2})$ $\frac{s}{2}$) Y_2 , and none of them has a non-constant polynomial solutuon, – a contradiction.

Let us prove the converse: assyming conditions a)-d), let us show that the system (2.1) in question is gauge equivalent to a system (1.5) having a non-constant polynomial solution. We already know that the system is gauge equivalent to a system (1.5), by Corollary 2.16: the monodromy-Stokes data is (1,1)-symmetric, see Remark 2.21. Thus, we consider that our system is of type (1.5). Its projectivized canonical solutions corresponding to the eigenvalues $\lambda_{20} = \lambda_{1\infty} = 0$ at the origin and at infinity paste together to the projectivization of just one solution $Y(z)$ of linear system, since $q_{20} = q_{1\infty}$ (parabolicity). The projectivized solution $Y(z)$ is fixed by the projectivized monodromy, since the latter fixes q_{20} (parabolicity). Hence, $Y(z)$ is a canonical solution of linear system, and the monodromy multiplies $Y(z)$ by a constant factor. The formal residue eigenvalues at 0 and at ∞ corresponding to $Y(z)$ are respectively $b_{20} = 0$ and $b_{1\infty} = \ell$, by definition, and $\ell \in \mathbb{N}$ by assumption. This together with (2.8) implies that in the canonical sectorial solution basis at the origin the monodromy is triangular and $Y(z)$ is its eigenvector with unit eigenvalue. Thus, $Y(z)$ is holomorphic on all of \mathbb{C}^* . Recall that the canonical solution $Y(z)$ corresponds to zero main term eigenvalues at 0 and ∞ . This together with the above formal residue eigenvalue statement imply that $Y(z)$ extends holomorphically to the origin and has growth of order z^{ℓ} (times a non-zero constant vector), as $z \to \infty$. Hence, it is a nonconstant vector polynomial solution. Theorem 2.22 is proved. \Box

In what follows we use the following general property of bundles with connections having parabolic monodromy-Stokes data.

Proposition 2.23 Let a holomorphic vector bundle $\mathcal E$ on the Riemann sphere $\overline{\mathbb C}_z$ be equipped with a meromorphic connection with irregular non-resonant singular points at 0 and at ∞ with formal normal forms

$$
\dot{Y} = \left(\frac{\text{diag}\left(-\frac{s}{2}, 0\right)}{z^2} + \frac{\text{diag}(\ell, 0)}{z}\right) Y \quad at \ 0,
$$
\n
$$
\dot{Y} = \left(\text{diag}\left(0, \frac{s}{2}\right) + \frac{\text{diag}(\ell, 0)}{z}\right) Y \quad at \ \infty.
$$

Let S_0^0 and S_0^{∞} be Stokes sectors at 0 and at ∞ respectively that are permuted by the involution $z \mapsto z^{-1}$. Let α_0 and α_∞ be paths from 1 to S_0^0 and S_0^∞ respectively that are also permuted by the same involution.

1) The connection has a meromorphic solution (i.e., a flat section), if and only if $\ell \in \mathbb{Z}$ and the monodromy-Stokes data⁵ $(q_{10}, q_{20}, q_{1\infty}, q_{2\infty})$ defined by the paths α_p , $p = 0, \infty$ and indexed by the main term eigenvalues $-\frac{s}{2}$ $\frac{s}{2}$, 0, 0, $\frac{s}{2}$ $\frac{s}{2}$, is parabolic. In this case the meromorphic solution is unique up to constant factor and is holomorphic on C.

2) Let a meromorphic solution exist. Then the bundle with connection is diagonal, i.e., a direct sum of one-dimensional holomorphic vector bundles with connections, if and only if the monodromy-Stokes data is degenerate.

Proof A meromorphic solution f should coincide with a canonical solution at the origin and at infinity that corresponds to zero eigenvalue of the main term matrix: thus, these solutions are pasted together, and hence, $q_{20} = q_{1\infty}$. This together with the equality $b_{20} = 0$ implies holomorphicity on $\mathbb C$ and uniqueness. The corresponding eigenvalue of the residue matrix of the formal normal form should be integer, by meromorphicity. Hence, $\ell \in \mathbb{Z}$. Therefore, at each singular point $0, \infty$ the formal monodromy is trivial, and at least one of the Stokes matrices is also trivial, since the canonical solution f is meromorphic. Hence, the monodromy of the connection is unipotent, by (2.8). This together the equality $q_{20} = q_{1\infty}$ implies parabolicity of the monodromy-Stokes data. The converse is proved analogously to the above proof of Theorem 2.22. Statement 1) is proved.

Let in the above assumptions the bundle with connection be diagonal. Then the solutions of the connections in one-dimensional subbundles are canonical solutions at both singular points: the origin and the infinity. This implies that the connection is analytically equivalent to its formal normal form, and its monodromy is trivial, as is that of the normal form. The projectivizations of the above solutions are holomorphic sections of the projectivized ambient bundle on \mathbb{C}^* . Their values at $z = 1$ are the points q_{jp} from the monodromy-Stokes data. This together with the previous statement implies that $q_{10} = q_{2\infty}, q_{20} = q_{1\infty}$. Hence, the monodromy-Stokes data is degenerate. The converse is proved analogously to the proof of Theorem 2.22. \Box

⁵The notion of monodromy-Stokes data is well-defined for connections in not necessarily trivial bundles over the Riemann sphere that have irregular non-resonant singular points of Poincaré rank 1 at 0 and at ∞ . In our case, when the eigenvalue collections of the main term matrices differ by sign, the notion of parabolic monodromy-Stokes data is also well-defined.

2.6 Local parametrization by formal invariant s and cross-ratio $\mathcal R$

Definition 2.24 Let $(q_{10}, q_{20} = q_{1\infty}, q_{2\infty}; M)$ be an abstract non-degenerate parabolic monodromy-Stokes data. Its *characterizing cross-ratio* is the cross-ratio of the four points $q_{20}, q_{10}, q'_{10}, q_{2\infty}$, see Definition 2.20:

$$
\mathcal{R} := \frac{(q_{20} - q'_{10})(q_{10} - q_{2\infty})}{(q_{20} - q_{2\infty})(q_{10} - q'_{10})}.
$$
\n(2.23)

Proposition 2.25 Each abstract non-degenerate parabolic monodromy-Stokes data is uniquely determined by its characterizing cross-ratio, up to multiplication of the monodromy operator by scalar factor.

Proof The cross-ratio R defines the four points q_{10} , q'_{10} , q_{20} , $q_{2\infty}$ on the Riemann sphere $\overline{\mathbb{C}}$ up to Möbius transformation. Except may be for some of the cases, when some of them coincide. Let us check that this statement remains valid in case of coincidences as well. Indeed, $q_{10}, q_{2\infty} \neq q_{20} = q_{1\infty}$. One has $q'_{10} \neq q_{20}$: the point q'_{10} is the image of the point $q_{10} \neq q_{20}$ under the projectivized monodromy, which fixes the point q_{20} . Therefore, the only a priori possible coincidences are either $q_{10} = q_{2\infty}$, or $q'_{10} = q_{2\infty}$, or $q'_{10} = q_{10}$: only one coincidence may take place, since the total number of distinct points among q_{10} , q'_{10} , $q_{20}, q_{2\infty}$ is at least three (non-degeneracy). In the case, when $q_{10} = q_{2\infty}$, one has $\mathcal{R} = 0$. In the case, when $q'_{10} = q_{2\infty}$, one has $\mathcal{R} = 1$. In the case, when $q'_{10} = q_{10}$, one has $\mathcal{R} = \infty$. This yields three distinct values for R for the three possible coincidences.

Let us first assume for simplicity that $q_{10} \neq q_{2\infty}$. Then applying a Möbius transformation we can and will consider that $q_{20} = \infty$, $q_{10} = 0$, $q_{2\infty} = 1$. Then in the coordinate on $\mathbb{C} = \overline{\mathbb{C}} \setminus \{q_{20}\}\$, which will be denoted by w, the projectivized monodromy operator is a translation $w \mapsto w + u$, $u \in \mathbb{C}$. One should have $q'_{10} = q_{10} + u = u$, since q'_{10} is the image of the point $q_{10} = 0$. Therefore,

$$
\mathcal{R} = \frac{q_{2\infty} - q_{10}}{q'_{10} - q_{10}} = \frac{1}{u}.
$$

Thus, the correspondence $\mathcal{R} \mapsto u$ is bijective, and the projectivized monodromy can be restored from R, as does u. The case, when $q_{10} = q_{2\infty}$, i.e., $\mathcal{R} = 0$, is treated analogously. Namely, in this case $q'_{10} \neq q_{10}, q_{20}$ (non-degeneracy). Applying a Möbius transformation, we consider that $q_{20} = q_{1\infty} = \infty$, $q_{10} = q_{2\infty} = 0$ and $q'_{10} = 1$ in the standard affine coordinate w on $\mathbb{C} = \overline{\mathbb{C}} \setminus \{q_{20}\}.$ Then the projectivized monodromy is the translation $w \mapsto w + 1$. Hence, the value $\mathcal{R} = 0$ determines the monodromy-Stokes data uniquely up to multiplication of the monodromy by scalar factor. Proposition 2.25 is proved. the monodromy by scalar factor. Proposition 2.25 is proved.

Theorem 2.26 Two systems of type (1.5) with $s \neq 0$ having a polynomial solution of degree $\ell \geq 1$ coincide up to the parameter change $(\chi, a) \mapsto (-\chi, -a)$, if and only if the corresponding values of the paramer s and the cross-ratio R (constructed for given symmetric pair of homotopic classes of paths α_{ip}) are the same for both systems.

Proof If two systems of type (1.5) are gauge equivalent, then the gauge equivalence is given by a matrix of the type $diag(1, \pm 1)$. Indeed, the gauge equivalence matrix should be simultaneously upper- and lower-triangular, since the conjugation by it should preserve the (opposite) triangular types of the matrices K and S , and each one of the latter matrices has distinct eigenvalues. Therefore, the gauge transformation is given by a diagonal matrix $\Lambda = (\Lambda_1, \Lambda_2)$. It acts on the residue matrix by multiplying its off-diagonal terms $\pm a$ by ratios of eigenvalues: by the numbers $\left(\frac{\Lambda_1}{\Lambda_2}\right)$ Λ_2 $\int_{0}^{\pm 1}$. Thus transformed residue matrix has opposite off-diagonal elements, as does the initial matrix, if and only if the latter ratios are equal to ± 1 . That is, if and only if $\Lambda = \text{diag}(1, \pm 1)$ up to scalar factor. The gauge transformation in question does not change s, and it multiplies both χ , a by the same number ± 1 . It does not change the monodromy-Stokes data, and hence, the cross-ratio \mathcal{R} . Conversely, let two systems (1.5) with polynomial solutions of degree $\ell > 1$ have the same parameter value s, and let the corresponding monodromy-Stokes data have the same crossratio. Then the corresponding monodromy operators are the same. Indeed, one of the Stokes matrices at the origin is trivial: this holds for every irregular non-resonant Poincaré rank 1 singularity, whenever one of the canonical solutions is holomorphic on \mathbb{C}^* ; in our case the polynomial solution is such a solution. Therefore, in a canonical sectorial solution basis at 0 the monodromy is a triangular matrix with the same eigenvalues, as the formal monodromy, by (2.8). The eigenvalues of the formal monodromy are $e^{2\pi i \ell} = 1$ and $e^{0} = 1$. Thus, the monodromy of system (1.5) in question is unipotent. Then it is uniquely determined by the projectivized monodromy, which in its turn is uniquely determined by the cross-ratio $\mathcal R$ (Proposition 2.25). Finally, the complete monodromy-Stokes data is uniquely determined by the cross-ratio R . Thus, the systems in question have the same formal normal forms at each singular point $p = 0, \infty$ and the same monodromy-Stokes data. Hence, they are gauge equivalent, by Theorem 2.5. This together with the above characterization of gauge equivalent systems (1.5) implies that the systems in question are obtained one from the other by the simultaneous sign change $(\chi, a) \mapsto (-\chi, -a)$. This proves Theorem 2.26. \Box

2.7 Realization of parabolic monodromy-Stokes data by vector bundles with connections

Our goal is to prove Theorem 1.14, which states that each determinantal surface $S_{\ell,+}$ is biholomorphically parametrized (outside the locus $\{s = 0\}$) by the complement of $\mathbb{C}^* \times \overline{\mathbb{C}}$ to an analytic hypersurface (curve), and its fibers $\{s = const \neq 0\}$ are rational curves parametrized by finitely punctured fibers $\{s\} \times \overline{\mathbb{C}}$. The results of the previous subsection together with [11, appendix 3, lemma 1] imply that the parameters (s, \mathcal{R}) yield a local biholomorphic parametrization of each surface $S^o_{\ell, \pm}$. But the value of the parameter R depends on the homotopy class of its defining pair of paths α_0 , α_∞ and the corresponding Stokes sectors S_0^0 , S_0^{∞} ; the paths and the sectors should be permuted by the involution $z \mapsto z^{-1}$. As s makes one turn around the origin, the Stokes sectors turn, the endpoints of the paths (which lie in the Stokes sectors) also turn, and the homotopy class of each path changes. Therefore, as s makes a turn while R remains constant, the monodromy-Stokes data defined by the initial paths (and hence, the underlying system) changes in general.

To construct a global biholomorphic parametrization of the surfaces $S^o_{\ell, \pm}$, we pass to the universal covering \mathbb{C}_t over the punctured line \mathbb{C}_s^* , $s = e^t$, and construct continuous families of Stokes sectors $S_{0,t}^0$, $S_{0,t}^{\infty}$ and paths $\alpha_{0,t}$, $\alpha_{\infty,t}$. Then we construct a holomorphic family $\mathcal{E}_{t,\mathcal{R}}$ of holomorphic vector bundles with connections over $\overline{\mathbb{C}}_z$, parametrized by (t,\mathcal{R}) , with formal normal forms (1.37) and (1.38) at 0 and at ∞ respectively, whose monodromy-Stokes data defined by the above paths and sectors is parabolic and has cross-ratio \mathcal{R} . We show that two bundles $\mathcal{E}_{t,\mathcal{R}}$ and $\mathcal{E}_{t',\mathcal{R}'}$ with connections are gauge equivalent, if and only if

$$
(t', \mathcal{R}') = (t, \mathcal{R}) + (2\pi i n, 2n), \quad n \in \mathbb{Z}.
$$

Afterwards, in the next subsection we pass to the quotient of the space $\mathbb{C}_t \times \overline{\mathbb{C}}_{\mathcal{R}}$ by the cyclic group generated by the translation by the vector $(2\pi i, 2)$ and get a family of bundles with connections parametrized by $\mathbb{C}_{s}^{*} \times \overline{\mathbb{C}}$. We show that all the parameters in $\mathbb{C}^{*} \times \overline{\mathbb{C}}$ except for an analytic hypersurface (curve) correspond to trivial bundles equipped with linear systems of type (1.5) with polynomial solutions. This will yield a biholomorphic parametrization of each surface $S_{\ell,\pm}^o$ by the above parameters.

The above-mentioned families of sectors and paths are

$$
S_{0,t}^0 = \{ \text{Im } t - \arg z \in (-\frac{\pi}{2} + \delta, \frac{3\pi}{2} - \delta) \}, \quad S_{0,t}^{\infty} = \{ \text{Im } t + \arg z \in (-\frac{\pi}{2} + \delta, \frac{3\pi}{2} - \delta) \},
$$

$$
\alpha_{0,t}(\tau) = e^{i\tau(Im t + \pi)}, \quad \alpha_{\infty,t}(\tau) = \overline{\alpha_{0,t}(\tau)}, \quad \tau \in [0, 1]. \tag{2.25}
$$

They are permuted by the involution $z \mapsto z^{-1}$, and the sectors are Stokes for the formal normal forms (1.37) and (1.38) respectively. This follows from definition.

Theorem 2.27 For every $\ell \in \mathbb{N}$ there exists a two-dimensional holomorphic vector bundle $\mathcal{E} = \mathcal{E}_{\ell}$ over $\overline{\mathbb{C}}_z \times \mathbb{C}_t \times \overline{\mathbb{C}}_{\mathcal{R}}$ such that for every $t \in \mathbb{C}$, set $s := e^t$, and every $\mathcal{R} \in \overline{\mathbb{C}}$ the restriction $\mathcal{E}_{t,\chi} = \mathcal{E}_{t,\chi,\ell}$ of the bundle \mathcal{E} to $\overline{\mathbb{C}}_z = \overline{\mathbb{C}}_z \times \{(t,\chi)\}\$ is equipped with a meromorphic connection depending holomorphically on (t, χ) that has formal normal forms

$$
Y' = \left(\frac{\text{diag}\left(-\frac{s}{2}, 0\right)}{z^2} + \frac{\text{diag}(\ell, 0)}{z}\right) Y \quad at \ 0,
$$

$$
Y' = \left(\text{diag}(0, \frac{s}{2}) + \frac{\text{diag}(\ell, 0)}{z}\right) Y \quad at \ \infty
$$
 (2.26)

and satisfies the following statements:

a) The monodromy-Stokes data of the connection defined by the above Stokes sectors and paths is parabolic and has the given cross-ratio R.

b) Two bundles $\mathcal{E}_{t,\mathcal{R}}$, $\mathcal{E}_{t',\mathcal{R}'}$ with connections corresponding to two distinct (t,\mathcal{R}) and (t', \mathcal{R}') are analytically equivalent, if and only if relation (2.24) holds.

Proof Fix $a \ell \in \mathbb{N}$. Set

$$
d_{\ell} = d_{\ell}(s) := \frac{\ell!}{2\pi i} \left(\frac{2}{s}\right)^{\ell} e^{\frac{s}{2}}.
$$

We consider the following model linear systems depending on parameters $s \in \mathbb{C}^*$, $u \in \mathbb{C}$:

$$
Y' = \left(\frac{\text{diag}(-\frac{s}{2}, 0)}{z^2} + \frac{\begin{pmatrix} \ell & 0\\ud_{\ell} & 0\end{pmatrix}}{z}\right) Y \quad \text{over the disk } D_2 := \{|z| < 2\};\tag{2.27}
$$

$$
Y' = \left(\text{diag}(0, \frac{s}{2}) + \frac{\begin{pmatrix} \ell & u d_{\ell} \\ 0 & 0 \end{pmatrix}}{z}\right) Y \quad \text{over } \overline{\mathbb{C}} \setminus \overline{D_{\frac{1}{2}}}.
$$
 (2.28)

These systems have the prescribed formal normal forms (2.26) at the origin and at the infinity respectively. It is easy to check that their monodromy operators in appropriate local solution bases are given by lower (upper) triangular unipotent matrices with triangular elements being equal to $\pm u$. We paste them together over the annulus

$$
\mathcal{A} := \{ \frac{1}{2} < |z| < 2 \} = D_2 \cap (\overline{\mathbb{C}} \setminus \overline{D_{\frac{1}{2}}})
$$

in appropriate way depending on (t, \mathcal{R}) , preserving the differential equations as connections on trivial bundles, in order to obtain a holomorphic vector bundle $\mathcal{E}_{t,\mathcal{R}}$ with connection having the desired cross-ratio R of the monodromy-Stokes data. Recall that R is the crossratio between appropriate analytic extensions of canonical sectorial solutions. To define the gluing, we first find the canonical solutions and their analytic extensions to $z_0 = 1$ along the paths $\alpha_{0,t}^{-1}$, $\alpha_{\infty,t}^{-1}$ and find monodromy matrices of systems (2.27), (2.28). To this end, let us introduce the paths

$$
\beta_{0,t}(\tau) = (1 - \tau)\alpha_{0,t}(1) = (1 - \tau)e^{i(\text{Im }t + \pi)},
$$

$$
\beta_{\infty,t}(\tau) = \frac{1}{\beta_{0,t}(\tau)} = \frac{1}{1 - \tau}e^{-i(\text{Im }t + \pi)},
$$

$$
\gamma_{p,t} := \beta_{p,t}^{-1} \circ \alpha_{p,t}^{-1}
$$

Each $\gamma_{p,t}$ is a path from the point p to 1.

We will deal with functions defined on a disk $U \subset \mathbb{C}^*$ centered at $z_0 = 1$. For every $z \in U$ by $\gamma_{p,t,z}$ we will denote the composition of the path $\gamma_{p,t}$ from p to 1 with the straight path from 1 to z.

The canonical sectorial basic solutions of system (2.27) at 0 in $S_{0,t}^0$ numerated by the main term matrix eigenvalues $-\frac{s}{2}$ and 0 will be denoted by f_{10} , f_{20} . Similarly the canonical basic solutions of system (2.28) at ∞ in $S_{0,t}^{\infty}$ numerated by the main term matrix eigenvalues 0 and $\frac{s}{2}$ will be denoted by $f_{1\infty}$, $f_{2\infty}$.

Proposition 2.28 1) The analytic extensions to the disk U of sectorial solutions f_{10} , f_{20} of system (2.27) and $f_{1\infty}$, $f_{2\infty}$ for system (2.28) along paths $\alpha_{0,t}$ and $\alpha_{\infty,t}$ respectively (after their appropriate normalizations by constant factors) are given by the following formulas:

$$
f_{10}(z) = (e^{\frac{s}{2}(\frac{1}{z}-1)}z^{\ell}, d_{\ell}u \int_{\gamma_{0,t,z}} \zeta^{\ell-1} e^{\frac{s}{2}(\frac{1}{\zeta}-1)} d\zeta), \quad f_{20}(z) = (0,1), \tag{2.29}
$$

$$
f_{1\infty}(z) = (z^{\ell}, 0), \quad f_{2\infty}(z) = (d_{\ell}uz^{\ell} \int_{\gamma_{\infty, t, z}} \zeta^{-(\ell+1)} e^{\frac{s}{2}(\zeta - 1)} d\zeta, \ e^{\frac{s}{2}(z - 1)}).
$$
 (2.30)

2) In the above canonical bases the monodromy matrices M_0 , M_{∞} of systems (2.27) and (2.28) respectively are equal to

$$
M_0 = \begin{pmatrix} 1 & 0 \\ u & 1 \end{pmatrix}, \ M_\infty = \begin{pmatrix} 1 & u \\ 0 & 1 \end{pmatrix}.
$$
 (2.31)

Proof The vector functions f_{20} and $f_{1\infty}$ are clearly solutions of systems (2.27) and (2.28) respectively. They are obviously canonical, since they are respectively holomorphic at zero and meromorphic at infinity. Let us show that the other vector functions f_{10} , $f_{2\infty}$ are canonical. First let us prove their well-definedness: convergence of the corresponding integrals. The integral in (2.29) has subintegral expression $\zeta^{\ell-1}e^{\frac{s}{2}(\frac{1}{\zeta}-1)}$, which has singular point at the origin. It is taken along the path $\gamma_{0,t,z}$, which goes out from the origin in the straightline direction with azimuth $\text{Im } t + \pi$ with respect to the real axis (as does the path $\beta_{0,t}^{-1}$). Therefore, along the latter straightline part of the integration path one has

$$
\arg \zeta = \operatorname{Im} t + \pi, \quad \frac{s}{2\zeta} = \frac{e^t}{2|\zeta|e^{i(\operatorname{Im} t + \pi)}} = -\frac{e^{\operatorname{Re} t}}{2|\zeta|}.
$$

Therefore, the subintegral exponent is a flat function along the path $\gamma_{0,t,z}$ at the origin, and hence, the integral in (2.29) converges. Convergence of the integral in (2.30) is proved analogously. The fact that the vector functions f_{10} and $f_{2\infty}$ are solutions of systems (2.27) and (2.28) respectively is proved by straightforward differentiation. Let us prove that their analytic extensions along the paths $\alpha_{0,t}$ and $\alpha_{\infty,t}$ are the canonical sectorial solutions in $S_{0,t}^{0}$ and $S_{0,t}^{\infty}$ respectively that are numerated by the main term matrix eigenvalues $-\frac{s}{2}$ $\frac{s}{2}$ and s $\frac{s}{2}$. Let us prove this statement for f_{10} : for the vector function $f_{2\infty}$ the proof is analogous. It suffices to prove the latter statement for the backward analytic extension along $\gamma_{0,t,z}^{-1}$ of the vector function f_{10} towards a point $w \in \gamma_{0,t,z}$ close to the origin with $\arg w = \text{Im } t + \pi$. The backward extension is given by the integral (2.29) taken along the straightline segment $[0, w]$. The integral is flat in w, as is the subintegral function, see the above discussion, and the first component in f_{10} is also flat for the same reason. Therefore, $f_{10}(w) = o(1)$, as $w \to 0$ with $\arg w = \text{Im } t + \pi$, while $f_{20} \equiv (0, 1)$. The canonical solution in $S_{0,t}^0$ is uniquely determined by the condition that it should tend to zero, as $w \to 0$ along a ray along which $e^{\frac{s}{2w}}$ is flat, since the other canonical basic solution $(0,1)$ does not tend to zero. Hence, f_{10} is a canonical solution.

Now for the proof of the proposition it remains to calculate the monodromy. Clearly the monodromy acts trivially on the functions f_{20} and $f_{1\infty}$. Let us calculate the monodromy images of the other functions f_{10} and $f_{2\infty}$. As z makes one courterclockwise circuit around the origin, the integral in (2.29) changes by the additive constant

$$
2\pi i e^{-\frac{s}{2}} Res_0(z^{\ell-1} e^{\frac{s}{2z}}) = \frac{2\pi i}{\ell!} e^{-\frac{s}{2}} (\frac{s}{2})^{\ell} = d_{\ell}^{-1}.
$$

This together with (2.29) implies the formula for the monodromy matrix M_0 in (2.31) . Similarly, as z makes a counterclockwise circuit, the integral in (2.30) is changed by the additive constant equal to $2\pi i e^{-\frac{s}{2}}$ times the Laurent coefficient at ζ^{-1} of the function $\zeta^{-(\ell+1)}e^{\frac{s}{2}\zeta}$. The latter Laurent coefficient is equal to $\frac{(\frac{s}{2})^{\ell}}{\ell!}$ $\frac{2}{\ell!}$. Thus, the resulting additive constant is equal to d_{ℓ}^{-1} ℓ^{-1} . This implies the formula for the monodromy matrix M_{∞} in (2.31) . Proposition 2.28 is proved. \Box

Let

$$
\mathcal{Y}^{0}(z) = (f_{10}, f_{20})(z), \ \mathcal{Y}^{\infty}(z) = (f_{1\infty}, f_{2\infty})(z)
$$

denote the fundamental matrices formed by the canonical solution bases (2.29) and (2.30) of systems (2.27) and (2.28) respectively. Here the vector functions $f_{jp}(z)$ are treated as columns of the corresponding matrix \mathcal{Y}^p . Their monodromy matrices M_0 and M_∞ given by (2.31) are conjugated:

$$
M_0 = \Psi_{\zeta}^{-1} M_{\infty} \Psi_{\zeta}, \ \Psi_{\zeta} := \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -\zeta & 1 \end{pmatrix}.
$$
 (2.32)

Here ζ is an arbitrary complex number. Therefore, the rescaled fundamental matrix

$$
\widetilde{\mathcal{Y}}^{\infty}(z) = \mathcal{Y}^{\infty}(z)\Psi_{\zeta}
$$
\n(2.33)

of system (2.28) has the same monodromy matrix M_0 , as the fundamental matrix $\mathcal{Y}^0(z)$.

Now we paste the fundamental matrices $\mathcal{Y}^0(z)$ and $\widetilde{\mathcal{Y}}^{\infty}(z)$ together over the annulus $\mathcal{A} := D_2 \setminus D_{\frac{1}{2}}$ in order to paste the systems (2.27) and (2.28) to get a holomorphic vector bundle with connection induces by the pasted systems. Namely, the direct products

$$
\mathcal{F}_0 := \mathbb{C}^2 \times D_2, \quad \mathcal{F}_{\infty} := \mathbb{C}^2 \times (\overline{\mathbb{C}} \setminus \overline{D_{\frac{1}{2}}})
$$

are identified as follows: for every $z \in \mathcal{A}$,

 $(v, z) \in \mathcal{F}_0$ is equivalent to $(w, z) \in \mathcal{F}_{\infty}$, if and only if

$$
v = X_{t,u,\zeta}(z)w, \ X_{t,u,\zeta}(z) := \mathcal{Y}^0(z) \left(\widetilde{\mathcal{Y}}^{\infty}(z)\right)^{-1}.
$$
 (2.34)

Here the matrix function $X_{t,u,\zeta}(z)$ is defined by (2.34) for z close to 1, and for the other z it is defined by its analytic extension along a path going from 1 to z in the annulus \mathcal{A} .

Proposition 2.29 1) The matrix function $X_{t,u,\zeta}(z)$ is holomorphic on the annulus A, and it depends holomorphically on the parameters $(t, u, \zeta) \in \mathbb{C}^3$.

2) The corresponding gluing (2.34) pastes the direct products $\mathcal{F}_0 \times \mathbb{C}^3_{t,u,\zeta}$ and $\mathcal{F}_{\infty} \times \mathbb{C}^3_{t,u,\zeta}$ to a holomorphic vector bundle over $\overline{C}_z \times C^3_{t,u,\zeta}$, which will be denoted by $\widetilde{\mathcal{E}}$ and treated as a holomorphic family of holomorphic vector bundles $\mathcal{E}_{t,u,\zeta}$ on $\overline{\mathbb{C}}_z$

3) Gluing (2.34) sends the connection on \mathcal{F}_{∞} defined by system (2.28) to the connection on \mathcal{F}_0 defined by (2.27), and hence, induces a meromorphic connection $\nabla_{t,u,\zeta}$ on each bundle $\mathcal{E}_{t,u,\zeta}$; the family of connections $\nabla_{t,u,\zeta}$ depends holomorphically on the parameters (t, u, ζ) .

4) The monodromy-Stokes data of each connection $\nabla_{t,u,\zeta}$ defined by the paths $\alpha_{0,t}$ and $\alpha_{\infty,t}$ is parabolic. It is non-degenerate, if and only if $(u,\zeta) \neq (0,0)$, and in this case the corresponding cross-ratio is equal to

$$
\mathcal{R} = \frac{\zeta}{u}.\tag{2.35}
$$

Proof Holomorphicity of the matrix function $X_{t,u,\zeta}(z)$ on the annulus A, i.e., its singlevaluedness (independence on path of analytic extension) follows from coincidence of monodromy matrices of the fundamental matrices \mathcal{Y}^0 and $\widetilde{\mathcal{Y}}^{\infty}$: as z makes a counterclockwise circuit around the origin, both fundamental matrices are multiplied by M_0 , and in the ratio $X_{t,u,\zeta}(z) = \mathcal{Y}^0(\widetilde{\mathcal{Y}}^{\infty})^{-1}$, this M_0 cancels out. This proves Statement 1), which in its turn (together with holomorphicity of the fundamental matrices in the parameters t, u, ζ) implies Statement 2). Gluing (2.34) sends each column of the fundamental matrix solution $\widetilde{\mathcal{Y}}^{\infty}$ of system (2.28) to the corresponding column of the fundamental matrix solution \mathcal{Y}^0 of system (2.27), by definition. Hence, it pastes the systems to one meromorphic connection in the bundle $\mathcal{E}_{t,u,\zeta}$, which depends holomorphically on the parameters, as do the systems and the fundamental matrices. The second columns of both fundamental matrices are canonical solutions $f_{1\infty}$ and f_{20} , and they are pasted together. This implies parabolicity of the monodromy-Stokes data of the pasted connection. Its non-degeneracy criterium $(u, \zeta) \neq (0, 0)$ follows from construction.

Let us now calculate the cross-ratio R of the monodromy-Stokes data in the case, when $(u, \zeta) \neq (0, 0)$. To do this, let us find the values $f_{ip}(1)$ of the canonical solutions after identification (2.34) in the trivial bundle chart \mathcal{F}_0 and their projectivizations q_{jp} and the point q'_{10} . Note that

$$
\widetilde{\mathcal{Y}}^{\infty} = (f_{2\infty} - \zeta f_{1\infty}, \ f_{1\infty}) : \tag{2.36}
$$

the columns of the fundamental matrix $\widetilde{\mathcal{Y}}^{\infty}$ are the vector functions $f_{2\infty}-\zeta f_{1\infty}$ and $f_{1\infty}$, by (2.32) and (2.33). And they are identified with f_{10} and f_{20} . Therefore, $f_{1\infty}(1)$ is identified with $f_{20}(1)$, and $f_{2\infty}(1)$ is identified with $f_{10}(1) + \zeta f_{20}(1)$. Let $f_{10}(z)$ denote the result of counterclockwise analytic extension of the vector function $f_{10}(z)$ considered as a germ of vector function at $z = 1$. One has

$$
f_{10}(1) = f_{10}(1) + uf_{20}(1),
$$

by formula (2.31) for the monodromy matrix M_0 . Thus, in the chart \mathcal{F}_0 , the four points $q_{20}, q_{10}, q'_{10}, q_{2\infty}$ are the projections to $\mathbb{CP}^1 = \overline{\mathbb{C}}$ of the four vectors $f_{20}(1), f_{10}(1), \tilde{f}_{10}(1) =$ $f_{10}(1) + uf_{20}(1)$, $f_{10}(1) + \zeta f_{20}(1) \simeq f_{2\infty}(1)$. In the basis $f_{10}(1)$, $f_{20}(1)$ (let (w_1, w_2)) denote the corresponding coordinates on \mathbb{C}^2) the latter vectors are $(0, 1), (1, 0), (1, u), (1, \zeta)$. Their projections to $\mathbb{CP}^1 = \overline{\mathbb{C}}_{\Phi}$, $\Phi = \frac{w_2}{w_1}$, are the points ∞ , 0, u, ζ . Their cross-ratio (2.23) is equal to $\frac{\zeta}{u}$. This proves (2.35) and finishes the proof of Proposition 2.29.

Proposition 2.30 For every given $t \in \mathbb{C}$ two holomorphic vector bundles with connections $(\mathcal{E}_{t,u,\zeta},\nabla_{t,u,\zeta}),$ $(\mathcal{E}_{t,u',\zeta'},\nabla_{t,u',\zeta'})$ corresponding to the same t and some $(u,\zeta), (u',\zeta') \neq (0,0)$ are holomorphically gauge equivalent, if and only if the corresponding cross-ratio are the same, i.e., if and only if $[u : \zeta] = [u' : \zeta']$ as points of the projective line \mathbb{CP}^1 .

Proof The bundles in question have the same formal normal forms at each singular point. The paths α_{tp} , $p = 0$, ∞ , defining their monodromy-Stokes data are the same for both systems, since they depend only on t . Therefore, the bundles are gauge equivalent, if and only if their monodromy-Stokes data are the same. Thus, gauge equivalence implies the cross-ratio equality. Let us prove the converse. Recall that the cross-ratio defines the monodromy-Stokes data uniquely up to multiplication of the monodromy operator by constant factor (Proposition 2.25). In both connections the monodromy operators are unipotent. Therefore, equality of cross-ratios implies coincidence of the monodromy-Stokes data, and hence, gauge equivalence. Proposition 2.30 is proved. \Box

In what follows the gauge equivalent bundles $\mathcal{E}_{t,u,\zeta}$ with the given value $\mathcal{R} := \frac{\zeta}{u}$ will be denoted by one and the same symbol

$$
\mathcal{E}_{t,\mathcal{R}}:=\mathcal{E}_{t,u,\zeta},\ \mathcal{R}=\frac{\zeta}{u}.
$$

Proposition 2.31 The family of vector bundles $\mathcal{E}_{t,R}$ form a holomorphic vector bundle over $\overline{\mathbb{C}}_z \times \mathbb{C}_t \times \overline{\mathbb{C}}_{\mathcal{R}}$, which will be denoted by \mathcal{E} , with a holomorphic family of meromorphic connections over \overline{C}_z with irregular nonresonant singular points at zero and at infinity.

Proof The vector bundle family $\mathcal{E}_{t,u,\zeta}$ form a holomorphic vector bundle (denoted $\widetilde{\mathcal{E}}$) over $\overline{\mathbb{C}} \times \mathbb{C}_t \times (\mathbb{C}_{u,\zeta}^2 \setminus \{(0,0)\})$ with a holomorphic family of connections over $\overline{\mathbb{C}}$. Namely, the bundle $\widetilde{\mathcal{E}}$ is defined by gluing (2.34) of two trivial bundles $\mathbb{C}^2 \times D_2 \times \mathbb{C}_t \times (\mathbb{C}^2 \setminus \{(0,0)\})$ and $\mathbb{C}^2 \times (\overline{\mathbb{C}} \setminus \overline{D_{\frac{1}{2}}}) \times \mathbb{C}_t \times (\mathbb{C}^2 \setminus \{(0,0)\})$: the gluing matrix function $X_{t,u\zeta}(z)$ is holomorphic in (z, t, u, ζ) . Consider the tautological bundle $\pi : \mathbb{C}^2_{u,\zeta} \setminus \{(0,0)\} \to \mathbb{CP}^1 = \overline{\mathbb{C}}_{\mathcal{R}}, \mathcal{R} = \frac{\zeta}{u}$ $\frac{\varsigma}{u}$, and its two sections σ_0 , σ_∞ over the two affine charts $\mathbb{C} = {\mathcal{R} \neq \infty}$ and ${\mathcal{R} \neq 0}$:

$$
\sigma_0: \mathcal{R} \mapsto (1, \mathcal{R}), \quad \sigma_1: \mathcal{R} \mapsto (\mathcal{R}^{-1}, 1).
$$

Let \mathcal{E}^0 and \mathcal{E}^{∞} denote the pullbacks of the bundle $\widetilde{\mathcal{E}}$ under the maps $(z, t, \overline{\mathcal{R}}) \mapsto (z, t, \sigma_p(\mathcal{R})),$ $p = 0, \infty$. This yields two holomorphic vector bundle families $\mathcal{E}_{t,\mathcal{R}}^p$ over $\overline{\mathbb{C}}_z$, with $p = 0, \infty$, $\mathcal{R} \neq 0,\infty$. For every $(t,\mathcal{R}) \in \mathbb{C}^2$ with $\mathcal{R} \neq 0$ the bundles $\mathcal{E}_{t,\mathcal{R}}^0$ and $\mathcal{E}_{t,\mathcal{R}}^{\infty}$ are isomorphic as holomorphic bundles with connections. The isomorphism is uniquely determined by normalization so that in the trivializing chart \mathcal{F}^0 it sends the solution $f_{20}(z) = (0,1)$ of the connection in one bundle to the same solution of the connection in the other bundle. Uniqueness follows from non-degeneracy of the monodromy-Stokes data. Indeed, the restriction to the fiber $\overline{\mathbb{C}} \times \{z=1\}$ of the projectivized isomorphism should map the set $\{q_{20}, q_{10}, q'_{10}, q_{2\infty}\}$ corresponding to one connection to the similar set corresponding to the other connection. The collection of points in each set is at least three, by non-degeneracy. This defines the projectivized isomorphism uniquely. Thus, the bundle isomorphism is uniquely defined up to constant factor, and the latter factor is uniquely defined by the above normalization. The isomorphism thus constructed depends holomorphically on the parameters, as does the monodromy-Stokes data. Therefore, the isomorphism family thus constructed pastes the bundles \mathcal{E}^0 and \mathcal{E}^{∞} together to one holomorphic vector bundle \mathcal{E} . Proposition 2.31 is \Box

Lemma 2.32 Two bundles $\mathcal{E}_{t,\mathcal{R}}$ and $\mathcal{E}_{t',\mathcal{R}'}$ with connections are holomorphically gauge equivalent, if and only if relation (2.24) holds.

Proof Clearly, the bundles with connections are gauge equivalent, if and only if they have the same formal normal forms and the same monodromy-Stokes data (or equivalently, the same cross-ratio \mathcal{R}) defined by one and the same pairs of paths and Stokes sectors. Coincidence of formal normal forms is equivalent to the equality saying that $s = e^t$ is equal to $s' = e^{t'}$. This holds, if and only if $t' - t = 2\pi n$, $n \in \mathbb{Z}$. The above monodromy-Stokes data of bundles $\mathcal{E}_{t,\mathcal{R}}$ and $\mathcal{E}_{t',\mathcal{R}'}$ are defined by two different pairs of paths: $(\alpha_{0,t}, \alpha_{\infty,t})$ and $(\alpha_{0,t'}, \alpha_{\infty,t'})$ Let us show that the monodromy-Stokes data defined by the first pair of paths is the same for both connections, if and only if $\mathcal{R}' - \mathcal{R} = 2n$. To do this, we use the next proposition, which shows how does the cross-ratio change when we replace its defining paths $\alpha_{p,t}$ by $\alpha_{p,t'}$, $p=0,\infty$.

Proposition 2.33 Let $\mathcal E$ be a holomorphic vector bundle with connection on $\overline{\mathbb C}$ having formal normal forms (2.26) at the origin and at infinity that is not "diagonal", i.e., not a direct sum of one-dimensional holomorphic bundles with connections. Let the connection have a meromorphic solution. Let s be the corresponding parameter from the formal normal forms. For every value $t = \ln s$ of the logarithm let $\alpha_{p,t}$, $p = 0, \infty$, be the paths defined in (2.25) , and let \mathcal{R}_t denote the characteristic cross-ratio of the corresponding monodromy-Stokes data. One has

$$
\mathcal{R}_{t+2\pi i} = \mathcal{R}_t + 2. \tag{2.37}
$$

Proof When we add $2\pi i$ to t, the new path $\alpha_{0,t}$ is obtained from the old one by multiplication by counterclockwise circuit. This means that the analytic continuation along the new path $\alpha_{0,t}^{-1}$ of the canonical sectorial solution f_{10} (flat section with respect to the connection) is obtained from that for the old path by applying the inverse of the monodromy. Similar statement holds for $\alpha_{\infty,t}$ and $f_{2\infty}$: the new $f_{2\infty}$ is obtained from the old one by applying the monodromy. Recall that the monodromy-Stokes data defined by the paths $\alpha_{0,t}$ and $\alpha_{\infty,t}$ are parabolic and non-degenerate, by Proposition 2.23. We consider that $f_{2\infty}$ is not proportional to f_{10} ; the opposite case is treated analogously (and can be also deduced by passing to limit in a family of pasted bundles with connections). Thus, $q_{2\infty} \neq q_{10}$, and q_{10} , q_{20} , $q_{2\infty}$ are distinct. Choosing appropriate coordinate on $\overline{\mathbb{C}}$ we consider that the monodromy-Stokes data for the old t are given by

$$
q_{10} = 0
$$
, $q_{20} = q_{1\infty} = \infty$, $q'_{10} = u$, $q_{2\infty} = 1$.

Then for the new t one has

$$
q_{10} = -u
$$
, $q_{20} = q_{1\infty} = \infty$, $q'_{10} = 0$, $q_{2\infty} = 1 + u$,

by the above path and monodromy argument. Therefore, the old cross-ratio (2.23) is equal to $\frac{1}{u}$, and the new one is equal to $\frac{1}{u} + 2$. This proves the proposition.

Proposition 2.33 together with the above argument implies Lemma 2.32. \Box

Propositions 2.29, 2.30, 2.31 and Lemma 2.32 together imply Theorem 2.27. \Box

2.8 Quotient bundle family. Global parametrization of surfaces $S^o_{\ell, \pm}$

Consider the quotient of the space $\mathbb{C}_t \times \overline{\mathbb{C}}_{\mathcal{R}}$ by the action of the cyclic group generated by the translation F by the vector $(2\pi i, 2)$. The quotient space $(\mathbb{C} \times \overline{\mathbb{C}})/F$ is identified with $\mathbb{C}_s^* \times \overline{\mathbb{C}}_{\xi}$, and the quotient projection is

$$
\pi_{tr} : (t, \mathcal{R}) \mapsto (s = e^t, \ \xi = \pi i \mathcal{R} - t).
$$

Lemma 2.34 There exists a holomorphic vector bundle $\widehat{\mathcal{E}}$ over $\overline{\mathbb{C}}_z \times \mathbb{C}^*_s \times \overline{\mathbb{C}}_\xi$ with a family of meromorphic connections over \overline{C}_z such that the bundle $\mathcal{E} = (\tilde{\mathcal{E}}_{t,\mathcal{R}})$ is its pullback under the propection $z \times (t, \mathcal{R}) \mapsto z \times \pi_{tr}(t, \mathcal{R}).$

Proof For every (t, \mathcal{R}) and $(t', \mathcal{R}') = F^{n}(t, \mathcal{R})$, $n \in \mathbb{Z}$, there is a unique isomorphism of the corresponding bundles $\mathcal{E}_{t,\mathcal{R}}$ and $\mathcal{E}_{t',\mathcal{R}'_n}$ with connections up to scalar factor. This isomorphism can be normalized as in the proof of Proposition 2.31, and then it depends locally holomorphically on the parameters (t, \mathcal{R}) , as in the same proof. This implies the lemma analogously to the same proof lemma analogously to the same proof.

Consider the bundle $\hat{\mathcal{E}}$ as a family of holomorphic vector bundles $\widehat{\mathcal{E}}_{s,\xi}$ over $\overline{\mathbb{C}}_z$ with connections. Let $\ell \in \mathbb{N}$ be the corresponding parameter in the formal normal forms of the connections. Set

Triv_{$\ell := \{ (s, \xi) \in \mathbb{C}^* \times \overline{\mathbb{C}} \mid \text{ the holomorphic bundle } \widehat{\mathcal{E}}_{s, \xi} \text{ is trivial} \},$}

$$
\Sigma_\ell := (\mathbb{C}^* \times \overline{\mathbb{C}}) \setminus \operatorname{Triv}_\ell.
$$

Let $\Sigma'_{\ell} \subset \text{Triv}_{\ell}$ denote the subset of those (s,ξ) for which in the corresponding connection in the trivial bundle on $\overline{\mathbb{C}}$ the line non-coincidence condition does not hold: that is, the eigenline of the main term matrix at 0 with the eigenvalue $-\frac{s}{2}$ coincides with the eigenline of the main term matrix at infinity with the eigenvalue $\frac{s}{2}$.

Theorem 2.35 1) The subset $\text{Triv}_\ell \subset \mathbb{C}^* \times \overline{\mathbb{C}}$ is the complement to a complex analytic hypersurface (curve) Σ_{ℓ} : a complex analytic subset of pure codimension 1.

2) The subset $\Sigma' \subset \text{Triv}_\ell$ is also a complex analytic hypersurface (curve), in particular, the subsets Triv_ℓ and

$$
\operatorname{Triv}_{\ell}^o := \operatorname{Triv}_{\ell} \setminus \Sigma' \subset \mathbb{C}_s^* \times \overline{\mathbb{C}}_{\xi}
$$

are both open connected and dense in $\mathbb{C}_{s}^{*} \times \overline{\mathbb{C}}_{\xi}$.

3) For every $(s,\xi) \in \text{Triv}_{\ell}^o$ the trivial bundle $\widehat{\mathcal{E}}_{s,\xi}$ with connection is analytically gauge equivalent to a system of type (1.5) having a polynomial solution. The parameters (χ, a) of the corresponding system (1.5) are uniquely determined up to sign change $(\chi, a) \mapsto (-\chi, -a)$ and $(\chi, a) \neq (0, 0)$.

4) Thus defined two-valued map $\mathcal{T}: \text{Triv}_{\ell}^o \to \mathbb{C}^3_{\chi,a,s}$ associating to each (s,ξ) the parameters of the corresponding system (1.5) from Statement 3), has two holomorphic branches \mathcal{T}_\pm bijectively parametrizing the two determinantal surfaces $S^o_{\ell,\pm} = \{ \det(G_{1,\ell}(\chi,a,s) \pm \ell_s) \}$ $G_{2,\ell}(\chi,a,s)) = 0\} \setminus \{s=0\}$ by the domain $\text{Triv}_{\ell}^o \subset \mathbb{C}^* \times \overline{\mathbb{C}}$.

Proof The subset $\text{Triv}_{\ell}^o \subset \mathbb{C}^* \times \overline{\mathbb{C}}$ is non-empty. Indeed, the set $S_{\ell,+} \cup S_{\ell,-}$ of systems (1.5) with polynomial solutions is non-empty, their monodromy-Stokes data corresponding to symmetric sectors and paths are parabolic, and the corresponding parameters (s, ξ) clearly lie in Triv_{ℓ}. This implies that Σ_{ℓ} is an analytic hypersurface, by [45, proposition 4.1] (see also its short proof in [11, appendix 3]); see also a more general [49, theorem 2.2, p.449]. The set Σ'_{ℓ} is the set of those parameters $(s,\xi) \in \text{Triv}_{\ell}$ that correspond to those connections in trivial bundle where appropriate main term matrix eigenlines coincides. This is a non-trivial codimension one condition, since the complement $\text{Triv}_{\ell}^o = \text{Triv}_{\ell} \setminus \Sigma'_{\ell}$ has a non-empty interior. Therefore, Σ'_{ℓ} is a complex hypersurface in Triv_{ℓ}. Statements 1) and 2) are proved.

Proposition 2.36 The surfaces $S^o_{\ell, \pm}$ are disjoint from the s-axis $\{\chi = a = 0\}$.

Proof Way 1. System (1.5) with $\chi = a = 0$ is diagonal with main term matrices $\operatorname{diag}(-\frac{s}{2})$ $(\frac{s}{2},0)$ and diag(0, $\frac{s}{2}$) $\frac{s}{2}$ at zero and at infinity respectively. Therefore, if $s \neq 0$, then it has no vector polynomial solution, being the direct sum of one-dimensional equations with solutions $z^{\ell}e^{\frac{s}{2z}}$ and $e^{\frac{s}{2}z}$.

Way 2. For $\chi = a = 0$ one has $G_{2,\ell} = 0$, and $G_{1,\ell} = G_{1,\ell} \pm G_{2,\ell}$ is upper-triangular with -zero diagonal elements $\frac{s}{2}$ and $\frac{1}{2}$. Hence, its determinant is non-zero. non-zero diagonal elements $\frac{s}{2}$ and $\frac{1}{2}$. Hence, its determinant is non-zero.

Statement 3) follows from Theorems 2.22 and 2.26; inequality $(\chi, a) \neq (0, 0)$ follows from Proposition 2.36.

Let us prove Statement 4). First, each point $(s_0, \xi_0) \in \text{Triv}_{\ell}^o$ has a neighborhood where the map $\mathcal{T} : (s, \xi) \mapsto (\chi, a, s)$ has two local holomorphic branches. This follows from [11, appendix 3, theorem 2], which implies holomorphicity of normalized trivialization of the bundle $\mathcal{E}_{s,\xi}$ in the parameter (s,ξ) and hence, holomorphicity in (s,ξ) of the family of linear systems thus obtained (images of the connections) in the trivialized bundle. The above two holomorphic branches clearly differ by sign (Statement 3)), and hence, their images lie in different surfaces $S_{\ell,\pm}$, since the latter surfaces are permuted by the sign change $(\chi, a) \mapsto (-\chi, -a)$. The parameter subset Triv_{ℓ} is open and connected, being the complement of the manifold $\mathbb{C}^* \times \overline{\mathbb{C}}$ to the union $\Sigma_{\ell} \cup \Sigma'_{\ell}$. Therefore, analytic extension of each initial holomorphic branch of the map $\mathcal T$ along every path in Triv_ℓ^o may lie only in one of the surfaces $S_{\ell, \pm}$. This together with the above statement implies that each branch yields a global holomorphic map

$$
\mathcal{T}_{\pm} : \text{Triv}_{\ell}^{o} \to S_{\ell, \pm}^{o} = S_{\ell, \pm} \setminus \{s = 0\}. \tag{2.38}
$$

Its bijectivity follows from Statement 3) and the fact that each system (1.5) with $s \neq 0$ having a polynomial solution has non-degenerate parabolic monodromy-Stokes data and hence, corresponds to a connection defined by some $(s, \xi) \in \text{Triv}_{\ell}^o$ (Theorem 2.22). Theorem 2.35 is proved. \Box

2.9 Regularity and rationality. Proof of Theorem 1.14

Lemma 2.37 The surfaces $S^o_{\ell, \pm}$ are disjoint regular submanifolds in $\mathbb{C}^3 \setminus \{s = 0\}$ diffeomorphically parametrized by the bijective maps \mathcal{T}_\pm given by (2.38). In particular, they are connected and hence, irreducible.

Proof The maps \mathcal{T}_{\pm} are holomorphic and bijective. The surfaces $S_{\ell,\pm}^o$ are analytic subsets in $\mathbb{C}^3 \setminus \{s = 0\}$. To prove that they are submanifolds, it suffices to show that the maps \mathcal{T}_{\pm} are immersions, and it suffices to prove this for one sign, say "+". (Their connectivity (hence, irreducibility) follows from connectivity of the parametrizing domain Triv_ℓ^o .) To this end, we fix some $(s_0, \xi_0) \in \text{Triv}_{\ell}^o$ and consider the new local coordinates (s, \mathcal{R}) on its neighborhood, where $\mathcal{R} = -i\pi^{-1}(\xi + t)$ is the characteristic cross-ratio of the monodromy-Stokes data corresponding to the paths $\alpha_{0,t}$, $\alpha_{\infty,t}$, $t = \ln s$. Set $(\chi_0, a_0, s_0) = \mathcal{T}_+(s_0, \xi_0) \in$ $S_{\ell,+}^o$. We claim that the value of the cross-ratio R is well-defined on all systems (1.5) with parameters (χ, a, s) close enough to (χ_0, a_0, s_0) , and it is a holomorphic function of (χ, a, s) . Indeed, recall that R is the cross-ratio of the points q_{20} , q_{10} , q'_{10} , $q_{2\infty}$, where q_{10} , $q_{20}, q_{2\infty}$ are projectivizations of the values at $z = 1$ of the canonical solutions $f_{10}, f_{20}, f_{2\infty}$, and q'_{10} is the projectivization of the value at 1 of the image of the solution f_{10} under the monodromy. All the above four points q depend holomorphically on the parameters of the system, as do the canonical solutions. Therefore, R is holomorphic. This yields an inverse map $(\chi, a, s) \mapsto (s, \mathcal{R})$ holomorphic on a neighborhood of the point (χ_0, a_0, s_0) in \mathbb{C}^3 . It is inverse to \mathcal{T}_+ : its pre-composition with \mathcal{T}_+ is the identity. This implies that \mathcal{T}_+ is an immersion.

Let us show that the surfaces $S_{\ell,\pm}^o$ are disjoint. Suppose the contrary: they are intersected at some point $(\chi_0, a_0, s_0), s_0 \neq 0$. One has $(\chi_0, a_0) \neq (0, 0)$, by Proposition 2.36, and

$$
(\chi_0, a_0, s_0) = \mathcal{T}_{\pm}(s_0, \xi_{\pm}) \text{ for some } \xi_{\pm} \in \overline{\mathbb{C}}.
$$

But $\xi_{+} = \xi_{-}$, since they are equal (up to known multiplicative and additive constants) to the characteristic cross-ratio of parabolic monodromy-Stokes data of one and the same linear system (1.5) with parameters (χ_0, a_0, s_0) , defined by the same pair of paths. Thus, two distinct maps \mathcal{T}_\pm take the same value (χ_0, a_0, s_0) at the point (s_0, ξ_+) , and they differ by signs of (χ, a) -components. Hence, $\chi_0 = a_0 = 0$. The contradiction thus obtained proves Lemma 2.37. \Box

Proposition 2.38 The surfaces $S^o_{\ell, \pm}$ are fibered by regular rational curves (finitely punctured Riemann spheres) of degree $2l+1$ over the s-axis punctured at the origin. There exists $a \ d \in \mathbb{N}, d \leq 2\ell + 1$, such that for all but a discrete subset $Coll_{\ell} \subset \mathbb{C}^*_s$ of values s the s-fiber of the surface $S_{\ell,\pm}^o$ has exactly d distinct punctures, and for $s \in Coll_{\ell}$ the corresponding number of punctures is less than d.

Proof For every $s \neq 0$ let $\mathcal{X}_{s,\pm}$ denote the s-fiber: the preimage under the projection $\pi_s : S^o_{\ell, \pm} \to \mathbb{C}_s$. It is a non-empty algebraic curve. Indeed, it is the zero locus of the polynomial $\mathcal{P}_{\ell,\pm}(\chi,a,s)$ with fixed s and variable (χ,a) . The latter polynomial in (χ,a) is non-constant for every s: substituting $\chi = 0$ we get the polynomial $Q_{\ell, \pm}(\mu, r) = \det(\mathcal{G}_{\ell} \pm$ r Id) defining the curve $\Gamma_{\ell, \pm}$ (up to constant factor), $\mu = \frac{s}{2}$ $\frac{s}{2}, r = \frac{a}{2}$ $\frac{a}{2}$. This is a polynomial of degree $\ell > 0$ in r, and hence, is non-constant. Thus, the set $\mathcal{X}_{s,\pm}$ is a non-empty zero locus of a polynomial, and hence, an algebraic curve. Therefore, is parametrized (bijectively except for possible singularities) by a punctured compact Riemann surface. The punctures correspond to the intersections of the projective closure $\overline{\mathcal{X}_{s,\pm}}$ with the infinity hyperplane, which will be denoted by H . Each puncture comes with multiplicity equal to the corresponding intersection index of the compact curve $\mathcal{X}_{s,\pm}$ with \mathcal{H} . Therefore, the total number of punctures with multiplicities is bounded by the degrees of the curves $\overline{\mathcal{X}_{s,\pm}}$. Their degrees are no greater than the degree of the polynomials $\mathcal{P}_{\ell,\pm}$, which is no greater than $2\ell+1$.

On the other hand, the map \mathcal{T}_{\pm} parametrizes the curve $\mathcal{X}_{s,\pm}$ biholomorphically by an open subset $X_s \subset \overline{\mathbb{C}}_{\xi}$, which is the s-fiber of the open subset $\text{Triv}_{\ell}^o \subset \mathbb{C}^* \times \overline{\mathbb{C}}$. Therefore, X_s should be Riemann sphere punctured in no more than 2 ℓ points. Hence, each $\mathcal{X}_{s,\pm}$ is a rational curve, and it is regular by biholomorphicity of the parametrization \mathcal{T}_{\pm} . Distinct punctures may collide only at a discrete subset $Coll_{\ell} \subset \mathbb{C}^*$ of values s. Indeed, resolution of singularities of the surface $S^o_{\ell, \pm}$, which exists by Hironaka Theorem [30], transforms the (singular) fibration by s-fibers into a singular fibration of a regular surface by rational curves. Their punctures correspond to their intersections with another algebraic curve: the image of the intersection of the surface $S_{\ell,\pm}^o$ with the infinity plane under the resolution of singularities. Collisions of the above-mentioned intersections obviously correspond to a discrete subset of values s. This implies the last statement of Proposition 2.38. Each curve $\mathcal{S}_{s,\pm}$ is irreducible, being rational. Let us show that it has degree $2\ell+1$. Indeed, it is the zero locus of the restriction to the given value s of the polynomial $\mathcal{P}_{\ell,\pm}(\chi,a,s) = \det(G_{1,\ell} \pm G_{2,\ell}).$ The highest degree monomial of the latter polynomial in the variables (χ, a) with fixed s is equal to $\chi^{\ell+1}a^{\ell}$ up to sign, which follows from formulas (1.12), (1.13) for the matrices $G_{1,\ell}$ and $G_{2,\ell}$. Therefore, the curve $\mathcal{X}_{s,\pm}$ has degree $2\ell+1$, unless the polynomial $\mathcal{P}_{\ell,\pm}(\chi,a,s)$ with fixed $s \neq 0$ has a multiple zero along its irreducible zero locus $\mathcal{X}_{s,\pm}$. The latter multiple zero case is impossible. Indeed, in the latter case there would exist another polynomial

 $\Psi_{s,\pm}(\chi, a)$ vanishing on $\mathcal{X}_{s,\pm}$ and an $m \in \mathbb{N}$, $m \geq 2$, such that $\Psi_{s,\pm}^m(\chi, a) = \mathcal{P}_{\ell,\pm}(\chi, a, s)$ up to constant factor. Let $\chi^{\alpha} a^{\beta}$ be the highest degree monomial of the polynomial $\Psi_{s,\pm}$. Then $\alpha m = \ell + 1$, $\beta m = \ell$, $m(\alpha - \beta) = 1$, thus $m = 1$, – a contradiction. The proposition is proved. is proved. \Box

Statement 1) of Theorem 1.14 follows from Lemma 2.37 and Proposition 2.36. Its Statement 2) follows from Proposition 2.38. It remains to prove Statement 3). The preimage $\mathcal{T}_{\pm}^{-1}(S_{\ell,\pm}^o)$ is the complement of the product $\mathbb{C}^* \times \overline{\mathbb{C}}$ to the union $\Sigma_{\ell} \cup \Sigma'_{\ell}$. Let us show that the latter union is an analytic hypersurface. We already know that Σ_{ℓ} is an analytic hypersurface, and Σ'_{ℓ} is an analytic hypersurface in $(\mathbb{C}^* \times \overline{\mathbb{C}}) \setminus \Sigma_{\ell}$. It remains to show that the closure $\overline{\Sigma'_{\ell}}$ is an analytic hypersurface (if non-empty), i.e., that its a priori potential singularities contained in Σ_{ℓ} are removable. Here by non-removable singularity we mean a point where the analytic subset in question does not extend as an analytic subset: a point where the analytic subset is defined but not regular is considered as removable by definition. Indeed, it follows from definition and analyticity that the only potential non-removable singularities are points of the intersection $\overline{\Sigma'_{\ell}} \cap \Sigma_{\ell}$. The latter intersection points have scoordinates lying in the discrete subset $Coll_{\ell}$ from Proposition 2.38, and each intersection point corresponds to a puncture: lies in the finite complement $\overline{\mathbb{C}}\setminus X_{s,\pm}$ of the corresponding fiber of Triv $_{\ell}^o = (\mathbb{C}^* \times \overline{\mathbb{C}}) \setminus (\Sigma_{\ell} \cup \Sigma_{\ell}')$. Hence, each intersection point is isolated. An isolated intersection point is a removable singularity of the hypersurface Σ'_{ℓ} , by Shiffman Theorem, see [52] and [23, subsection 4.4]. Therefore, $\overline{\Sigma'_{\ell}}$ is a hypersurface, and hence, so is $\Sigma_{\ell} \cup \Sigma'_{\ell}$. Theorem 1.14 is proved.

2.10 Irreducibility of polynomials and surfaces. Proof of Theorem 1.13

Irreducibility of the surface $S_{\ell,\pm}$ is equivalent to connectivity of its regular part. Its open subset $S_{\ell, \pm}^o = S_{\ell, \pm} \setminus \{s = 0\}$ is connected and lies in its regular part (Lemma 2.37). Thus, to prove irreducibility, it remains to show that $S_{\ell, \pm}$ does not contain the hyperplane $\{s = 0\}.$

Proposition 2.39 The polynomial $\mathcal{P}_{\ell, \pm}$ does not vanish identically on the hyperplane {s = 0}.

Proof The polynomial $\mathcal{P}_{\ell,\pm}(0, a, 0)$ is a polynomial in one variable a with highest degree monomial $\frac{a^{\ell}}{2\ell+1}$ $\frac{a^{\ell}}{2^{\ell+1}}$ up to sign, by formulas (1.12), (1.13) for the matrices $G_{1,\ell}$ and $G_{2,\ell}$. \Box

The surfaces $S_{\ell, \pm}$ do not contain the hyperplane $\{s = 0\}$, by Proposition 2.39. This together with the above argument implies their irreducibility.

Let us prove irreducibility of the polynomial $\mathcal{P}_{\ell,\pm}$. Its zero locus is an irreducible surface $S_{\ell,\pm}$. Thus, it remains to show that $S_{\ell,\pm}$ is not its multiple zero locus. This follows from the fact that for every fixed $s \neq 0$ the zero locus of the polynomial $\mathcal{P}_{\ell,\pm}(\chi, a, s)$ with fixed s is the s-fiber of the surface $S^o_{\ell, \pm}$, and it is not its multiple zero locus, as was shown in the proof of Proposition 2.38. Theorem 1.13 is proved.

2.11 Smoothness of the spectral curves and the genus formula. Proof of Theorems 1.6 and 1.5

As was shown by I.Netay, Theorem 1.6 (smoothness of the curve $\Gamma_\ell^o = \Gamma_\ell \setminus {\mu = 0} \subset \mathbb{C}^2_{\lambda,\mu}$) implies the genus formula given by Theorem 1.5. Thus, for the proof of the genus formula, it remains to prove Theorem 1.6. To do this, we consider the line field on $\mathbb{C}^3_{\chi,a,s}$ given by the system of the second and third differential equations in (1.6). It is generated by the vector field

$$
v := (v_{\chi}, v_a, v_s), \ v_{\chi} = \frac{1}{2}(a(1 - 4\chi^2) + 2\ell\chi), \ v_a = 2\chi(a^2 - s^2) - \ell a, \ v_s = s,
$$
 (2.39)

which will be called the P3 isomonodromic vector field.

Theorem 2.40 Each surface $S_{\ell, \pm}$ is tangent to the P3 isomonodromic vector field v.

Proof The surface $S_{\ell, \pm}$ is described in terms of monodromy-Stokes data: parabolicity and non-degeneracy conditions, see Theorem 2.22, which are invariant under isomonodromic deformations. Therefore, it is tangent to the P3 isomonodromic vector field v. \Box

Smoothness of the curve Γ_{ℓ}^{o} is equivalent to smoothness of one (or equivalently, both) components $\Gamma^o_{\ell, \pm} \subset \mathbb{C}^2_{a,s}$ of its pullback under the map $(a, s) \mapsto (\lambda = \frac{1}{4}(a^2 - s^2), \mu = \frac{s}{2}).$ The curve $\Gamma_{\ell,\pm}^o$ is the complement of the zero locus $\Gamma_{\ell,\pm} = \{Q_{\ell,\pm} = 0\}$ to the line $\{s = 0\}$. It is the intersection of the smooth surface $S^o_{\ell, \pm}$ with the hyperplane $\{\chi = 0\}$. The latter intersection is transversal, since the line field (1.6) is tangent to $S^o_{\ell, \pm}$ and transversal to the hyperplane $\{\chi = 0\}$ outside the line $\{a = \chi = 0\}$. But we already know that $S_{\ell, \pm}^o$ is disjoint from the line $\{\chi = a = 0\}$, by Theorem 1.14, Statement 1). Therefore, the line field on $\mathbb{C}^3_{\chi,a,s}$ given by (1.6) is transversal to the hyperplane $\{\chi=0\}$ at points of its intersection with $S_{\ell,\pm}^o$, which are exactly the points of the curve $\Gamma_{\ell,\pm}^o$, by Proposition 1.16. Thus, their intersection is transversal. Therefore, the curve $\Gamma^o_{\ell,\pm}$ is smooth, and hence, so is Γ_ℓ^o . Theorem 1.6, and hence, Theorem 1.5 are proved.

3 Isomonodromic foliations of the determinantal surfaces

Here we study the isomonodromic foliations of the complex surfaces $S_{\ell,\pm}$ for $\ell = 0,1$: foliations by phase curves of the P3 isomonodromic vector field v given by (2.39) .

Theorem 2.40 implies the following corollary.

Corollary 3.1 For every $\ell \in \mathbb{N}$ one has

$$
\frac{d\mathcal{P}_{\ell,\pm}}{dv} = h_{\ell,\pm}\mathcal{P}_{\ell,\pm}, \ h_{\ell,\pm} = h_{\ell,\pm}(\chi,a,s) \ \text{is a polynomial.} \tag{3.1}
$$

Proof The surface $S_{\ell, \pm}$ being zero locus of an irreducible polynomial $\mathcal{P}_{\ell, \pm}$, and the vector field *v* being polynomial and tangent to $S_{\ell, \pm}$ (Theorem 2.40), we get (3.1). field v being polynomial and tangent to $S_{\ell, \pm}$ (Theorem 2.40), we get (3.1).

Problem 3.2 For every ℓ find the polynomial multiplier $h_{\ell,\pm}$ from (3.1).

Below we present formulas for the polynomial multipliers $h_{\ell,\pm}$ for $\ell = 1,2$.

3.1 Case $\ell = 0$

For $\ell = 0$ the set of those parameters (χ, a, s) for which the corresponding system (1.5) has a vector polynomial solution is the union of two planes $\{\chi = \pm \frac{1}{2}\}\$ $\frac{1}{2}$, both being equipped with the coordinates (a, s) . The isomonodromic foliations of the latter planes are given by the Riccati equation derived from the second equation in (1.6):

$$
a'_s = \pm \frac{1}{s}(a^2 - s^2).
$$

3.2 Case $\ell = 1$

Let us calculate $h_{1,\pm}$. Afterwards we show that the corresponding isomonodromic foliation is given by a Riccati equation. We also present a result on monodromy and Stokes operators of the linear system whose projectivization is the latter Riccati equation.

Proposition 3.3 For $\ell = 1$ one has

$$
\frac{d\mathcal{P}_{1,\pm}}{dv} = h_{1,\pm}\mathcal{P}_{1,\pm}, \quad h_{1,\pm}(\chi, a, s) = \pm \frac{1}{4}(1 - 2\chi(a \pm s)).\tag{3.2}
$$

Proof Let us multiply $\mathcal{P}_{1,\pm}$ by ± 4 . We get

$$
\widetilde{\mathcal{P}}_{1,\pm} := \pm 4\mathcal{P}_{1,\pm} = (a \pm s)(1 - 4\chi^2) + 4\chi,
$$
\n
$$
v_a \frac{\partial \widetilde{\mathcal{P}}_{1,\pm}}{\partial a} + v_s \frac{\partial \widetilde{\mathcal{P}}_{1,\pm}}{\partial s} = (1 - 4\chi^2)(2\chi(a^2 - s^2) - a \pm s)
$$
\n
$$
= (\widetilde{\mathcal{P}}_{1,\pm} - 4\chi)(2\chi(a \mp s) + 1) - 2a(1 - 4\chi^2),
$$
\n(3.3)

$$
v_{\chi} \frac{\partial \tilde{P}_{1,\pm}}{\partial \chi} = (-8\chi(a \pm s) + 4) \frac{1}{2} (a(1 - 4\chi^2) + 2\chi) = (-4\chi(a \pm s) + 2)(a(1 - 4\chi^2) + 2\chi)
$$

= -8\chi^2(a \pm s) + 4\chi + 2a(1 - 4\chi^2) - 4a\chi(1 - 4\chi^2)(a \pm s)
= -8\chi^2(a \pm s) + 4\chi + 2a(1 - 4\chi^2) - 4a\chi(\tilde{P}_{1,\pm} - 4\chi). (3.4)

Summing expressions (3.4) and (3.3) yields

$$
\frac{d\mathcal{P}_{1,\pm}}{dv} = -8\chi^2 a \mp 8\chi^2 s + 4\chi + \tilde{\mathcal{P}}_{1,\pm}(1 - 2\chi(a \pm s)) + 16a\chi^2 - 4\chi
$$

$$
-8\chi^2 a \pm 8\chi^2 s = \tilde{\mathcal{P}}_{1,\pm}(1 - 2\chi(a \pm s)).
$$

This proves (3.2) .

In what follows we treat only the case " $-$ "; the case " $+$ " is then obtained by the sign change $(\chi, a) \mapsto (-\chi, -a)$.

Proposition 3.4 In the coordinates (χ, s) on the surface $\Gamma_{1,-} = {\mathcal{P}}_{1,-} = 0$ the vector field (2.39) generates the line field defined by the Riccati equation

$$
\chi_s' = \frac{1 - 4\chi^2}{2} - \frac{\chi}{s}.\tag{3.5}
$$

Proof The equation $\mathcal{P}_{1,-} = 0$ is equivalent to the formula

$$
a = s - \frac{4\chi}{1 - 4\chi^2}.
$$
\n(3.6)

Substituting (3.6) to the first differential equation in (1.6) yields

$$
\chi'_s = \frac{a - 2\chi(-1 + 2\chi a)}{2s} = \frac{a(1 - 4\chi^2) + 2\chi}{2s} = \frac{s(1 - 4\chi^2) - 2\chi}{2s} = \frac{1}{2}(1 - 4\chi^2) - \frac{\chi}{s}.
$$

This proves the statement of the proposition. \Box

The Riccati equation (3.5) is the projectivization of the linear system

$$
\Psi_s' = \left(\frac{1}{s} \begin{pmatrix} 0 & 0 \\ 0 & -1 \end{pmatrix} + \begin{pmatrix} 0 & 2 \\ \frac{1}{2} & 0 \end{pmatrix} \right) \Psi : \tag{3.7}
$$

a function $\chi(s)$ is a solution of (3.5), if and only if it can be presented as a ratio $\frac{\Psi_2}{\Psi_1}(s)$, where $\Psi(s) = (\Psi_1(s), \Psi_2(s))$ is a vector solution of (3.7).

System (3.7) has Fuchsian singularity at the origin with residue matrix diag $(0, -1)$ and irregular singularity at infinity of Poincaré rank 1.

Proposition 3.5 The formal normal form at infinity of system (3.7) is

$$
\widetilde{\Psi}'_s = \begin{pmatrix} \frac{1}{s} \begin{pmatrix} -\frac{1}{2} & 0\\ 0 & -\frac{1}{2} \end{pmatrix} + \begin{pmatrix} 1 & 0\\ 0 & -1 \end{pmatrix} \end{pmatrix} \widetilde{\Psi}.
$$
 (3.8)

2) The Stokes matrices of system (3.7) at infinity are both non-trivial.

3) Its monodromy along circuit around the origin is unipotent Jordan cell.

Proof The free term matrix in (3.7) has eigenvalues 1, -1 ; the corresponding eigenvectors are $(2, 1)$, $(2, -1)$. Therefore, it is diagonalized by conjugacy as follows:

$$
\begin{pmatrix} 2 & 2 \ 1 & -1 \end{pmatrix}^{-1} \begin{pmatrix} 0 & 2 \ \frac{1}{2} & 0 \end{pmatrix} \begin{pmatrix} 2 & 2 \ 1 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \ 0 & -1 \end{pmatrix}.
$$

The same conjugacy brings the residue matrix in (3.7) to

$$
\begin{pmatrix} 2 & 2 \ 1 & -1 \end{pmatrix}^{-1} \begin{pmatrix} 0 & 0 \ 0 & -1 \end{pmatrix} \begin{pmatrix} 2 & 2 \ 1 & -1 \end{pmatrix} = \begin{pmatrix} -\frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & -\frac{1}{2} \end{pmatrix}.
$$

This implies that the residue matrix of the formal normal form at infinity is diag($-\frac{1}{2}$ $\frac{1}{2}, -\frac{1}{2}$ $(\frac{1}{2})$. Statement 1) is proved.

The eigenvalues of the monodromy operator of system (3.7) along a positive circuit around the origin are equal to the exponents of $2\pi i$ times the residue matrix eigenvalues 0, -1 , since the singular point at the origin is Fuchsian with residue matrix diag(0, -1). Thus, both monodromy eigenvalues are equal to one, the monodromy is unipotent, and its trace is equal to 2. On the other hand, the monodromy matrix in an appropriate canonical sectorial solution basis at infinity is the product of three matrices: the monodromy matrix of the formal normal form $\exp(-2\pi i \operatorname{diag}(\frac{1}{2}, \frac{1}{2}))$ $(\frac{1}{2})$) = $-Id$ and the inverses of the Stokes matrices; the latter inverses are unipotent respectively lower and upper triangular of the type

$$
C_1^{-1} = \begin{pmatrix} 1 & 0 \\ -c_1 & 1 \end{pmatrix}, C_0^{-1} = \begin{pmatrix} 1 & -c_0 \\ 0 & 1 \end{pmatrix}; c_{0,1} \text{ are called the Stokes multipliers,
$$

see (2.8). The total product of the three matrices in question is equal to

$$
M = \begin{pmatrix} -1 & c_0 \\ c_1 & -(1 + c_0 c_1) \end{pmatrix} . \tag{3.9}
$$

Its trace is equal to $-(2 + c_0 c_1)$. But on the other hand, it is equal to 2, see the above discussion. Thus, $c_0c_1 = -4$. This proves Statement 2) of the proposition. Statement 3), which is equivalent to nontriviality of the monodromy (unipotence), follows from nontriviality of the Stokes multipliers and (3.9) , as in [28]. The proposition is proved. \Box

3.3 Case $\ell = 2$

Proposition 3.6 For $\ell = 2$ one has

$$
\frac{d\mathcal{P}_{2,\pm}}{dv} = h_{2,\pm}\mathcal{P}_{2,\pm}, \quad h_{2,\pm}(\chi, a, s) = \pm \frac{1}{8}(2 - a(2\chi \mp 1)).\tag{3.10}
$$

Proof Set $\widetilde{\mathcal{P}}_{2,\pm} := \pm 8\mathcal{P}_{2,\pm}$. One has

$$
\widetilde{\mathcal{P}}_{2,\pm} = (2\chi \pm 1)^2 (2\chi \mp 1)(a^2 - s^2) - 2a(2\chi \pm 1)(6\chi \mp 1) + 16\chi,
$$

see (1.17),

$$
v_s \frac{\partial \mathcal{P}_{2,\pm}}{\partial s} = -2s^2(2\chi \pm 1)^2(2\chi \mp 1)
$$

= $2(a^2 - s^2)(2\chi \pm 1)^2(2\chi \mp 1) - 2a^2(2\chi \pm 1)^2(2\chi \mp 1),$ (3.11)

$$
v_a \frac{\partial \widetilde{\mathcal{P}}_{2,\pm}}{\partial a} = (2\chi(a^2 - s^2) - 2a)(2\chi \pm 1)(2a(4\chi^2 - 1) - 12\chi \pm 2),
$$

$$
v_\chi \frac{\partial \widetilde{\mathcal{P}}_{2,\pm}}{\partial \chi} = \frac{1}{2}(a(1 - 4\chi^2) + 4\chi)((a^2 - s^2)(2\chi \pm 1)(12\chi \mp 2) - 8a(6\chi \pm 1) + 16).
$$

Therefore

$$
\frac{d\widetilde{P}_{2,\pm}}{dv} = v_s \frac{\partial \widetilde{P}_{2,\pm}}{\partial s} + v_a \frac{\partial \widetilde{P}_{2,\pm}}{\partial a} + v_\chi \frac{\partial \widetilde{P}_{2,\pm}}{\partial \chi} = A(a^2 - s^2) + B,\tag{3.12}
$$

$$
A = (2\chi \pm 1)(2(4\chi^2 - 1) + 4\chi a(4\chi^2 - 1) - 24\chi^2 \pm 4\chi + (a(1 - 4\chi^2) + 4\chi)) (6\chi \mp 1))
$$

= (2\chi \pm 1)(4\chi^2 - 1)(2 - a(2\chi \mp 1)). (3.13)

$$
B = -2a^2(2\chi \pm 1)^2(2\chi \mp 1) - 4a(2\chi \pm 1)(a(4\chi^2 - 1) - 6\chi \pm 1) + 4(a(1 - 4\chi^2) + 4\chi)(-a(6\chi \pm 1) + 2)
$$

= $2a^2(4\chi^2 - 1)(6\chi \mp 1) + 4a(-20\chi^2 + 1) + 32\chi.$ (3.14)

In order to prove (3.10), taking into account formula (3.13), it remains to show that

$$
B = (2 - a(2\chi \mp 1))(-2a(2\chi \pm 1)(6\chi \mp 1) + 16\chi). \tag{3.15}
$$

Opening brackets in (3.15) , one gets (3.14) . This proves (3.15) , and hence, (3.10) .

4 Acknowledgements

I am grateful to Yulia Bibilo, Victor Buchstaber, Igor Netay, Victor Novokshenov, Constantin Shramov, Vladimir Sokolov and Ilya Vyugin for helpful discussions.

References

- [1] Anderson P. W.; Rowell J. M. Probable observation of the Josephson tunnel effect. Phys. Rev. Lett. 10 (6) (1963): 230–232.
- [2] Arnold V. I. Geometrical Methods in the Theory of Ordinary Differential Equations. Second edition. – Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences], 250. Springer-Verlag, New York, 1988.
- [3] Arnold V. I.; Ilyashenko Yu. S. Ordinary differential equations. In: Dynamical Systems I. – Encyclopaedia Math. Sci. (1988), 1–148.
- [4] Balser W.; Jurkat W.B,; Lutz D.A. A general theory of invariants for meromorphic differential equations. I. Formal invariants. – Funkcialaj Ekvacioj, 22:2 (1979), 197–221.
- [5] Balser W.; Jurkat W.B.; Lutz D.A. Birkhoff invariants and Stokes' multipliers for meromorphic linear differential equations. $-$ J. Math. Anal. Appl. 71 (1979), No. 1, 48–94.
- [6] Barone A.; Paterno G. Physics and Applications of the Josephson Effect. John Wiley and Sons, New York–Chichester–Brisbane–Toronto–Singapore, 1982.
- [7] Beauville A. Complex Algebraic Surfaces. Second edition. Cambridge University Press, 1996.
- [8] Bibilo Yu.; Glutsyuk A. On families of constrictions in model of overdamped Josephson junction and Painlevé 3 equation. – Nonlinearity, 35 (2022), 5427–5480.
- [9] Bolibruch A.A. The Riemann–Hilbert problem on the complex projective line. [In Russian.] – Mat. Zametki, 46 (1989), No. 3, 118–120.
- [10] Bolibruch A. Inverse problems for linear differential equations with meromorphic co $efficients. - Isomonodromic Deformations and Applications in Physics (Montréal, 2000),$ CRM Proceeding and Lecture Notes, 31 (2002), 3–25.
- [11] Bolibruch A.A. Inverse monodromy problems in analytic theory of differential equations. [In Russian.] – Moscow, MCCME, 2018.
- [12] Buchstaber V.M.; Glutsyuk A.A. On determinants of modified Bessel functions and entire solutions of double confluent Heun equations. – Nonlinearity, 29 (2016), 3857– 3870.
- [13] Buchstaber V.M.; Glutsyuk A.A. On monodromy eigenfunctions of Heun equations and boundaries of phase-lock areas in a model of overdamped Josephson effect. – Proc. Steklov Inst. Math., 297 (2017), 50–89.
- [14] Buchstaber V.M.; Karpov O.V.; Tertychnyi S.I. On properties of the differential equation describing the dynamics of an overdamped Josephson junction. – Russ. Math. Surveys, 59:2 (2004), 377–378.
- [15] Buchstaber V.M.; Karpov O.V.; Tertychnyi S.I. Peculiarities of dynamics of a Josephson junction shifted by a sinusoidal SHF current. [In Russian.] – Radiotekhnika i Elektronika, 51:6 (2006), 757–762.
- [16] Buchstaber V.M.; Karpov O.V.; Tertychnyi S.I. The rotation number quantization *effect.* – Theoret and Math. Phys., **162** (2010), No. 2, 211–221.
- [17] Buchstaber V.M.; Karpov O.V.; Tertychnyi S.I. The system on torus modeling the dynamics of Josephson junction. – Russ. Math. Surveys, **67** (2012), No. 1, 178–180.
- [18] Buchstaber V.M.; Tertychnyi S.I. Explicit solution family for the equation of the resistively shunted Josephson junction model. – Theoret. and Math. Phys., 176 (2013), No. 2, 965–986.
- [19] Buchstaber V.M.; Tertychnyi S.I. Holomorphic solutions of the double confluent Heun equation associated with the RSJ model of the Josephson junction. – Theoret. and Math. Phys., 182:3 (2015), 329–355.
- [20] Buchstaber V.M.; Tertychnyi S.I. A remarkable sequence of Bessel matrices. Mathematical Notes, 98 (2015), No. 5, 714–724.
- [21] Buchstaber V.M.; Tertychnyi S.I. Automorphisms of solution space of special double confluent Heun equations. – Funct. Anal. Appl., $50:3$ (2016), 176–192.
- [22] Buchstaber V.M.; Tertychnyi S.I. Representations of the Klein group determined by quadruples of polynomials associated with the double confluent Heun equation. – Math. Notes, 103:3 (2018), 357–371.
- [23] Chirka E.M., Complex analytic sets. Math. Appl. (Soviet Ser.), 46, Kluwer Acad. Publ., Dordrecht, 1989, xx+372 pp.
- [24] Foote R.L., Geometry of the Prytz Planimeter. Reports on Math. Phys. 42:1/2 (1998), 249–271.
- [25] Foote R.L.; Levi M.; Tabachnikov S. Tractrices, bicycle tire tracks, hatchet planimeters, and a 100-year-old conjecture. – Amer. Math. Monthly, 120 (2013), 199–216.
- [26] Glutsyuk A. On constrictions of phase-lock areas in model of overdamped Josephson effect and transition matrix of the double-confluent Heun equation. – J. Dyn. Control Syst. 25 (2019), Issue 3, 323–349.
- [27] Glutsyuk A. On germs of constriction curves in model of overdamped Josephson junction, dynamical isomonodromic foliation and Painlevé 3 equation. Moscow Math. J. 23 (2023), No. 4, 479–513.
- [28] Glutsyuk, A.A.; Kleptsyn, V.A.; Filimonov, D.A.; Schurov, I.V. On the adjacency quantization in an equation modeling the Josephson effect. Funct. Analysis and Appl., 48 (2014), No. 4, 272–285.
- [29] Glutsyuk A.; Netay I. On spectral curves and complexified boundaries of phase-lock areas in a model of Josephson junction. – J. Dyn. Control Systems, 26 (2020), 785–820.
- [30] Hironaka H. On the resolution of singularities (characteristic zero). Proc. Int. Cong. Math. (1962), 507–521.
- [31] Il'in V.P.; Kuznetsov Yu.I. Tridiagonal Matrices and Their Applications [in Russian]. – Nauka, Moscow (1985).
- [32] Ilyashenko Yu.; Guckenheimer J. The duck and the devil: canards on the staircase. Moscow Math. J., 1 (2001), No. 1, 27–47.
- [33] Ilyashenko Yu.S.; Khovanskii A.G. Galois groups, Stokes operators, and a theorem of Ramis. – Functional Anal. Appl., 24:4 (1990), 286–296.
- [34] Ilyashenko Yu.S.; Filimonov D.A.; Ryzhov D.A. Phase-lock effect for equations modeling resistively shunted Josephson junctions and for their perturbations. – Funct. Analysis and its Appl. 45 (2011), No. 3, 192–203.
- [35] Jimbo M. Monodromy Problem and the Boundary Condition for Some Painlevé Equations. – Publ. RIMS, Kyoto Univ. 18 (1982), Issue 3, 1137–1161.
- [36] Josephson B.D. Possible new effects in superconductive tunnelling. Phys. Lett., 1 (1962), No. 7, 251–253.
- [37] Jurkat W.B.; Lutz D.A.; Peyerimhoff A. Birkhoff invariants and effective calculations for meromorphic linear differential equations. $-$ J. Math. Anal. Appl. 53 (1976), No. 2, 438–470.
- [38] Kleptsyn V.A.; Romaskevich O.L.; Schurov I.V. Josephson effect and slow-fast systems. [In Russian.] – Nanostuctures. Mathematical physics and Modelling, 8 (2013), 31–46.
- [39] Klimenko A.V; Romaskevich O.L. Asymptotic properties of Arnold tongues and Josephson effect. – Mosc. Math. J., 14:2 (2014), 367–384.
- [40] Levinson Y. Quantum noise in a current-biased Josephson junction. Phys. Rev. B 67 (2003), 184504.
- [41] Likharev K.K. Dynamics of Josephson junctions and circuits. Gordon and Breach Science Publishers, 1986.
- [42] Likharev K.K. Introduction to the dynamics of Josephson junctions. [In Russian.] Moscow, Nauka, 1985.
- [43] Likharev, K.K.; Ulrikh, B.T. Systems with Josephson junctions: Basic Theory. [In Russian.] – Izdat. MGU, Moscow, 1978.
- [44] Lin Y., Dai D., Tibboel P., *Existence and uniqueness of tronquée solutions of the third* and fourth Painlevé equations. – Nonlinearity 27 (2014), 171–186.
- $[45]$ Malgrange B. Sur les déformations isomonodromiques. I. Singularités réqulières. Cours. de l'Institut Fourier, 17 (1982), 1–26.
- [46] Malmquist J. Sur les équations différentielles du second ordre dont l'intégrale générale a ses points critiques fixes. – Ark. Mat. Astr. Fys., 17 (1922-23), 1–89.
- [47] McCumber D.E. Effect of ac impedance on dc voltage-current characteristics of superconductor weak-link junctions. $-$ J. Appl. Phys., **39** (1968), No. 7, 3113–3118.
- [48] Okamoto K. Polynomial Hamiltonians associated with Painlevé equations. I. Proc. Japan Acad., 56, Ser. A (1980), 264–268.
- [49] Röhrl H. On holomorphic families of fiber bundles over the Riemannian sphere. Mem. Coll. Sci. Univ. Kyoto, Ser. A, XXXIII (1961), No. 3, 435–477.
- [50] Schmidt V.V. Introduction to physics of superconductors. [In Russian.] MCCME, Moscow, 2000.
- [51] Shapiro S.; Janus A.; Holly S. Effect of microwaves on Josephson currents in superconducting tunneling. – Rev. Mod. Phys., 36 (1964), 223–225.
- [52] Shiffman B. On the removal of singularities of analytic sets. Michigan Math. J., 15 (1968), 111–120.
- [53] Sibuya Y. Stokes phenomena. Bull. Amer. Math. Soc., 83 (1977), 1075–1077.
- [54] Slavyanov S.Yu.; Lay W. Special functions: a unified theory based on singularities. Oxford Univ. Press, Oxford, 2000.
- [55] Stewart W.C., Current-voltage characteristics of Josephson junctions. Appl. Phys. Lett., 12 (1968), No. 8, 277–280.
- [56] Tertychnyi S.I. Long-term behavior of solutions of the equation $\dot{\phi} + \sin \phi = f$ with periodic f and the modeling of dynamics of overdamped Josephson junctions. – Preprint https://arxiv.org/abs/math-ph/0512058.
- [57] Tertychnyi S.I. The modeling of a Josephson junction and Heun polynomials. Preprint https://arxiv.org/abs/math-ph/0601064.
- [58] Vinnikov V. Complete description of determinantal representations of smooth irreducible curves. – Lin. Alg. Appl. 125 (1989), 103–140.
- [59] Vinnikov V. Self-adjoint determinantal representations of real plane curves. Math. Annalen, 296:1 (1993), 453–479.