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Long time behaviour of the solution of Maxwell's equations in dissipative generalized Lorentz materials (II) A modal approach

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Abstract

This work concerns the analysis of electromagnetic dispersive media modelled by generalized Lorentz models. More precisely, this paper is the second of two articles dedicated to the long time behaviour of solutions of Maxwell's equations in dissipative Lorentz media, via the long time decay rate of the electromagnetic energy for the corresponding Cauchy problem. In opposition to the frequency dependent Lyapunov functions approach used in [5], we develop a method based on the spectral analysis of the underlying non selfadjoint operator of the model. Although more involved, this approach is closer to physics, as it uses the dispersion relation of the model, and has the advantage to provide more precise and more optimal results, leading to distinguish the notion of weak and strong dissipation.

Keywords: Maxwell's equations, passive electromagnetic media, weakly dissipative generalized Lorentz models, long time electromagnetic energy decay rate, spectral theory.

1 Introduction

1.1 Motivation

This paper is devoted to the study of the long time behaviour, of solutions of the Cauchy problem for Maxwell's equations in a large class dissipative and dispersive materials: the generalized Lorentz media. We refer for instance to [28, 24, 16, 2, 3, 4] for the interest of considering such media with respect to applications, in particular in the context of metamaterials. We are more precisely interested in the decay of the electromagnetic energy. The evolution problem we consider is the following

$$\text{Find } \begin{cases} \mathbf{E}(\mathbf{x}, t) : \mathbb{R}^3 \times \mathbb{R}^+ \longrightarrow \mathbb{R}^3 & \mathbf{H}(\mathbf{x}, t) : \mathbb{R}^3 \times \mathbb{R}^+ \longrightarrow \mathbb{R}^3 \\ \mathbf{P}_j(\mathbf{x}, t) : \mathbb{R}^3 \times \mathbb{R}^+ \longrightarrow \mathbb{R}^3, 1 \leq j \leq N_e, & \mathbf{M}_\ell(\mathbf{x}, t) : \mathbb{R}^3 \times \mathbb{R}^+ \longrightarrow \mathbb{R}^3, 1 \leq \ell \leq N_m, \end{cases}$$

¹POEMS (Propagation d'Ondes: Etude Mathématique et Simulation) is a mixed research team (UMR 7231) between CNRS (Centre National de la Recherche Scientifique), ENSTA Paris (Ecole Nationale Supérieure de Techniques Avancées) and INRIA (Institut National de Recherche en Informatique et en Automatique).

such that (for all $1 \leq j \leq N_e, 1 \leq \ell \leq N_m$)

$$\left\{ \begin{array}{ll} \varepsilon_0 \partial_t \mathbf{E} - \nabla \times \mathbf{H} + \varepsilon_0 \sum_{j=1}^{N_e} \Omega_{e,j}^2 \partial_t \mathbf{P}_j = 0, & (\mathbf{x}, t) \in \mathbb{R}^3 \times \mathbb{R}^{+,*}, & (1.1a) \\ \mu_0 \partial_t \mathbf{H} + \nabla \times \mathbf{E} + \mu_0 \sum_{\ell=1}^{N_m} \Omega_{m,\ell}^2 \partial_t \mathbf{M}_\ell = 0, & (\mathbf{x}, t) \in \mathbb{R}^3 \times \mathbb{R}^{+,*}, & (1.1b) \\ \partial_t^2 \mathbf{P}_j + \alpha_{e,j} \partial_t \mathbf{P}_j + \omega_{e,j}^2 \mathbf{P}_j = \mathbf{E}, & (\mathbf{x}, t) \in \mathbb{R}^3 \times \mathbb{R}^{+,*}, & (1.1c) \\ \partial_t^2 \mathbf{M}_\ell + \alpha_{m,\ell} \partial_t \mathbf{M}_\ell + \omega_{m,\ell}^2 \mathbf{M}_\ell = \mathbf{H}, & (\mathbf{x}, t) \in \mathbb{R}^3 \times \mathbb{R}^{+,*}, & (1.1d) \end{array} \right.$$

completed by the divergence free initial conditions

$$\left\{ \begin{array}{l} \mathbf{E}(\cdot, 0) = \mathbf{E}_0, \quad \mathbf{H}(\cdot, 0) = \mathbf{H}_0, \quad \mathbf{P}(\cdot, 0) = \mathbf{P}_0, \quad \partial_t \mathbf{P}(\cdot, 0) = \dot{\mathbf{P}}_0, \quad \mathbf{M}(\cdot, 0) = \mathbf{M}_0, \quad \partial_t \mathbf{M}(\cdot, 0) = \dot{\mathbf{M}}_0 \\ \text{with } \mathbf{E}_0, \mathbf{H}_0 \in \mathbf{L}^2(\mathbb{R}^3), \quad \mathbf{P}_0, \dot{\mathbf{P}}_0 \in \mathbf{L}^2(\mathbb{R}^3)^{N_e} \text{ and } \mathbf{M}_0, \dot{\mathbf{M}}_0 \in \mathbf{L}^2(\mathbb{R}^3)^{N_m}, \end{array} \right. \quad (1.2)$$

where $\mathbf{L}^2(\mathbb{R}^3) = L^2(\mathbb{R}^3)^3$ and, for simplicity, we have used the notations

$$\mathbf{P} = (\mathbf{P}_j) \quad \text{and} \quad \mathbf{M} = (\mathbf{M}_\ell) \quad \text{with} \quad (\mathbf{P}_j) := (\mathbf{P}_j)_{j=1}^{N_e} \quad \text{and} \quad (\mathbf{M}_\ell) := (\mathbf{M}_\ell)_{\ell=1}^{N_m}.$$

In (1.1), the fields \mathbf{E} and \mathbf{H} are respectively the electric and magnetic fields while the \mathbf{P}_j and the \mathbf{M}_ℓ are auxiliary unknowns (polarization and magnetization fields respectively). The coefficients $(\omega_{e,j}, \omega_{m,\ell})$ and $(\Omega_{e,j}, \Omega_{m,\ell})$ are supposed to satisfy

$$\begin{aligned} 0 < \Omega_{e,1} \leq \dots \leq \Omega_{e,N_e}, \quad 0 < \Omega_{m,1} \leq \dots \leq \Omega_{m,N_m}, \\ \omega_{e,j} > 0, \quad 1 \leq j \leq N_e, \quad \omega_{m,\ell} > 0, \quad 1 \leq \ell \leq N_m. \end{aligned} \quad (1.3)$$

These coefficients are responsible for dispersion effects. Finally, the coefficients $(\alpha_{e,\ell}, \alpha_{m,\ell})$ must be non negative

$$\alpha_{e,j} \geq 0, \quad 1 \leq j \leq N_e, \quad \alpha_{m,\ell} \geq 0 \quad 1 \leq \ell \leq N_m. \quad (1.4)$$

The presence of positive coefficients $(\alpha_{e,\ell}, \alpha_{m,\ell})$ is responsible for dissipation effects. The reader will notice that one can assume without any loss of generality that the couples $(\alpha_{e,j}, \omega_{e,j})$ (resp. $(\alpha_{m,\ell}, \omega_{m,\ell})$) are all distinct the $\omega_{m,\ell}$ from the others.

Remark 1.1. *Usually in physical applications, one assumes also that the initial conditions (1.2) of the evolution system (1.1) satisfy also*

$$\nabla \cdot \mathbf{E}_0 = \nabla \cdot \mathbf{H}_0 = 0, \quad \mathbf{P}_0 = \dot{\mathbf{P}}_0 = 0 \quad \text{and} \quad \mathbf{M}_0 = \dot{\mathbf{M}}_0 = 0. \quad (1.5)$$

In other words, one considers that the initial fields are divergence free and that the medium has no polarization, no magnetization and induced currents at the initial time.

Let us introduce the two rational functions

$$\varepsilon(\omega) = \varepsilon_0 \left(1 - \sum_{j=1}^{N_e} \frac{\Omega_{e,j}^2}{q_{e,j}(\omega)} \right) \quad \text{and} \quad \mu(\omega) = \mu_0 \left(1 - \sum_{\ell=1}^{N_m} \frac{\Omega_{m,\ell}^2}{q_{m,\ell}(\omega)} \right), \quad (1.6)$$

where the second order polynomials $q_{e,j}$ and $q_{m,\ell}$ are defined by

$$q_{e,j}(\omega) = \omega^2 + i\alpha_{e,j}\omega - \omega_{e,j}^2 \quad \text{and} \quad q_{m,\ell}(\omega) = \omega^2 + i\alpha_{m,\ell}\omega - \omega_{m,\ell}^2. \quad (1.7)$$

As functions of the frequency ω , ε and μ are respectively the permittivity and permeability of the generalized Lorentz media associated to the Maxwell's system (1.1) (see Section 1.1 of [5] for more details about their physical interpretation). We denote by \mathcal{P}_e (resp. \mathcal{P}_m) the set of poles of the function ε (resp. μ). We also introduce \mathcal{Z}_e (resp. \mathcal{Z}_m) the set of zeros of ε (resp. μ). Generalized Lorentz media are causal and passive electromagnetic materials, as defined in [5] for instance. The counterpart of causality and passivity in the frequency domain is that $\omega \mapsto \omega \varepsilon(\omega)$ and $\omega \mapsto \omega \mu(\omega)$ are Herglotz functions of the frequency. Herglotz functions are analytic functions f in the upper half-plane $\mathbb{C}^+ := \{\omega \in \mathbb{C} \mid \text{Im } \omega > 0\}$ whose imaginary part $\text{Im } f$ is non-negative on \mathbb{C}^+ . For more details about passive electromagnetic media and Herglotz functions, we refer to [2], [3] and [5].

We make the two following assumptions which are linked to the irreducibility of the dispersion relation associated to the system (1.1) that involves the product $\varepsilon(\omega)\mu(\omega)$ (see Section 3.2.1):

- (H₁): the electric polynomials $q_{e,j}$ (see (1.7)) with distinct indices j do not have common roots. The same holds for the magnetic polynomials $q_{m,\ell}$ with distinct indices ℓ .
- (H₂): The zeros of the permittivity ε are not poles of the permeability μ and vice versa. Namely, one assumes that $\mathcal{P}_e \cap \mathcal{Z}_m = \emptyset$ and $\mathcal{P}_m \cap \mathcal{Z}_e = \emptyset$.

Remark 1.2. *It is easy to see that, as $\alpha_{e,j} \geq 0$, the roots of $q_{e,j}$ have a non positive imaginary part. More precisely, as the discriminant of $q_{e,j}$ is $\delta_{e,j} = 4\omega_{e,j}^2 - \alpha_{e,j}^2$*

- *If $\alpha_{e,j} < 2\omega_{e,j}$, the two roots of $q_{e,j}$ are $-i\alpha_{e,j}/2 \pm \delta_{e,j}^{1/2}/2 \notin i\mathbb{R}^-$,*
- *If $\alpha_{e,j} \geq 2\omega_{e,j}$, the two roots of $q_{e,j}$ are $-i(\alpha_{e,j} \pm |\delta_{e,j}|^{1/2})/2 \in i\mathbb{R}^-$.*

Thus, when $\alpha_{e,j} < 2\omega_{e,j}$, the two roots of $q_{e,j}$, ω_ and $-\overline{\omega_*} \notin i\mathbb{R}^-$, are distinct. Moreover, as $\alpha_{e,j} = -2\text{Im}(\omega_*)$ and $\omega_{e,j} = |\omega_*|$, $q_{e,j}$ can not share a common root with an other electric polynomial $q_{e,j'}$ since by assumption $(\alpha_{e,j}, \omega_{e,j}) \neq (\alpha_{e,j'}, \omega_{e,j'})$ for $j \neq j'$. Therefore, one only needs to assume (H₁) for electric polynomials $q_{e,j'}$ for which $\alpha_{e,j'} \geq 2\omega_{e,j'}$. The same properties hold for the magnetic polynomials $q_{m,\ell}$ with obvious changes.*

The electromagnetic energy is defined as

$$\mathcal{E}(t) \equiv \mathcal{E}(\mathbf{E}, \mathbf{H}, t) := \frac{1}{2} \left(\varepsilon_0 \int_{\mathbb{R}^3} |\mathbf{E}(\mathbf{x}, t)|^2 d\mathbf{x} + \mu_0 \int_{\mathbb{R}^3} |\mathbf{H}(\mathbf{x}, t)|^2 d\mathbf{x} \right). \quad (1.8)$$

The present paper is the sequel of [5] in which we addressed the question of long time decay of $\mathcal{E}(t)$ with a (frequency dependent) Lyapunov function approach. We were able to prove some *polynomial stability* results under the so-called *strong dissipation assumption*:

Definition 1.3 (Strong Dissipation for Lorentz models).

$$\forall 1 \leq j \leq N_e, \quad \alpha_{e,j} > 0, \quad \forall 1 \leq \ell \leq N_m, \quad \alpha_{m,\ell} > 0. \quad (1.9)$$

By polynomial stability, we mean the energy decays more rapidly than a negative power of t :

$$\mathcal{E}(t) \leq C t^{-s} \text{ for some } s > 0. \quad (1.10)$$

In this paper, we address the same question as in [5] via a *modal approach*. Admittedly, this approach is technically more involved than the one in [5], but it presents many advantages:

- it is by many aspects more *physical* (in particular it refers explicitly to the dispersion relation of the medium) thus more easily understandable by physicists,
- it leads to *optimal* results in the sense that upper bounds of the type (1.10) can be completed by corresponding lower bounds,
- it allows us to obtain polynomial stability results under a condition on the dissipation coefficients which is strictly weaker than (1.9), namely the *weak dissipation* assumption:

Definition 1.4 (Weak Dissipation for Lorentz models).

$$\sum_{j=1}^{N_e} \alpha_{e,j} + \sum_{\ell=1}^{N_m} \alpha_{m,\ell} > 0. \quad (1.11)$$

The reader will notice that the *strong dissipation* assumption imposes that all the coefficients $(\alpha_{e,j}, \alpha_{m,\ell})$ are positive while the *weak dissipation* assumption means only that at least one of them is positive.

In [5], we made in Section 1.2 a rather extensive analysis of the literature addressing problems similar to the one we consider here, with application in viscoelasticity [6, 10, 8] or electromagnetism [11, 18]. We discussed in particular various existing techniques for obtaining *polynomial* or *exponential stability* results (Sections 1.2 and 1.3). The closest works to ours are the paper by Figotin-Schechter [12] and the papers [21, 22, 23] by S. Nicaise and her collaborator C. Pignotti. In [12], the authors prove that the electromagnetic energy of the solutions of a large class of linear dissipative models tends to zero but without addressing the question of the decay rate. In [21, 22, 23], the question of the decay rate of the solution is studied in detail in bounded domains for perfectly conducting materials. However, when specified for Maxwell's equations in generalized Lorentz media, the results of [12] and [21, 22, 23] only applied when the strong dissipation assumption (1.9) is satisfied. Furthermore, under the weak dissipation assumption (1.11), we enlighten for some cases (when (1.9) does not hold) new decay rates which are still polynomial but with a smaller exponent in $1/t$ than the ones observed in the literature [21, 22, 23, 5].

The outline of the rest of this paper is as follows. First in Section 1.2, we rewrite the Maxwell's time-dependent system as an evolution problem which involves a maximal dissipative operator \mathbb{A} . Then in Section 1.3, we state the main results of the article. In Section 2, we show how to reduce the analysis of the evolution problem to the study of an infinity of Ordinary Differential Equations (ODE) systems in time parametrized by one real parameter. This is done by exploiting the fact that we work with constant coefficients which allows us to use the Fourier transform in space and the isotropy of the model. The equations are then parametrized by

$|\mathbf{k}|$, the modulus of the wave vector \mathbf{k} , namely the dual variable of the space variable. Each of these ODE systems involves a finite dimensional dissipative operator $\mathbb{A}_{|\mathbf{k}|,\perp}$ and the rest of the analysis is devoted to estimates, for large t of the exponentials

$$e^{-i\mathbb{A}_{|\mathbf{k}|,\perp} t} \text{ for } |\mathbf{k}| \in \mathbb{R}^+, \quad (1.12)$$

which obviously depends on the eigenvalues of $\mathbb{A}_{|\mathbf{k}|,\perp}$ and more precisely on their imaginary parts. Section 3 is devoted to general spectral properties of $\mathbb{A}_{|\mathbf{k}|,\perp}$ at fixed $|\mathbf{k}|$: we characterise its eigenvalues as the solutions of a rational equation parametrised by $|\mathbf{k}|$, namely the dispersion relation of the medium, whose analysis allows us to provide a diagonalizability criterion for $\mathbb{A}_{|\mathbf{k}|,\perp}$. Then, we need to distinguish in Sections 4, 5 and 6 respectively the large, small and intermediate values of $|\mathbf{k}|$. In Section 4, devoted to large $|\mathbf{k}|$, we provide asymptotic expansions (when $|\mathbf{k}| \rightarrow +\infty$) of the solutions of dispersion relation, which shows the diagonalisability of $\mathbb{A}_{|\mathbf{k}|,\perp}$ for $|\mathbf{k}|$ large enough. After diagonalization of the exponential (1.12), we obtain optimal decay estimates in time of the high spatial frequency components (i. e. large values of $|\mathbf{k}|$) of the solution of the evolution problem, based on the behaviour of the imaginary part of the eigenvalues of $\mathbb{A}_{|\mathbf{k}|,\perp}$ and uniform bounds (in $|\mathbf{k}|$) of the associated (non-orthogonal) spectral projectors. A similar strategy is adopted in Section 5 to derive optimal estimates for the low frequency components (i.e. small values of $|\mathbf{k}|$) of the solution. Section 6 deals with estimates for intermediate values of $|\mathbf{k}|$. The method is quite different and somewhat less technical and precise than the method of Sections 4 and 5 but it is sufficient to show (in Section 7) that, the corresponding part of the solution decays exponentially fast in time. Sections 7.1 and 7.2 are devoted to the proofs of our main results, after regrouping the estimates of the three previous sections and deriving estimates in space variable norms via Plancherel's theorem: this is where the polynomial decay due to large and small values of $|\mathbf{k}|$ is put in evidence. Moreover, we show that the corresponding polynomial decay estimates are optimal. In Section 7.3, we detail some extensions of these results. Finally, the appendix Section A contains technical results used for the spectral analysis of $\mathbb{A}_{|\mathbf{k}|,\perp}$. In particular, in Section A.2, Lemma A.1 provides a useful implicit function-like result for some (scalar) nonlinear equations in the complex plane. This lemma is used for the analysis of the dispersion relation.

1.2 Maxwell's equations in dissipative generalized Lorentz media

The (Cauchy) problem (1.1,1.2) can be rewritten as an evolution problem

$$\frac{d\mathbf{U}}{dt} + i\mathbb{A}\mathbf{U} = 0 \text{ with } \mathbf{U}(0) = \mathbf{U}_0 \in \mathcal{H}, \quad (1.13)$$

where \mathbb{A} is an unbounded operator on the Hilbert-space

$$\mathcal{H} := \mathbf{L}^2(\mathbb{R}^3)^N = \mathbf{L}^2(\mathbb{R}^3) \times \mathbf{L}^2(\mathbb{R}^3) \times \mathbf{L}^2(\mathbb{R}^3)^{N_e} \times \mathbf{L}^2(\mathbb{R}^3)^{N_e} \times \mathbf{L}^2(\mathbb{R}^3)^{N_m} \times \mathbf{L}^2(\mathbb{R}^3)^{N_m} \quad (1.14)$$

where $N = 2 + 2(N_e + N_m)$. \mathcal{H} is endowed by the following inner product:

for any $\mathbf{U} = (\mathbf{E}, \mathbf{H}, \mathbf{P}, \dot{\mathbf{P}}, \mathbf{M}, \dot{\mathbf{M}}) \in \mathcal{H}$ and $\mathbf{U}' = (\mathbf{E}', \mathbf{H}', \mathbf{P}', \mathbf{M}', \dot{\mathbf{P}}', \dot{\mathbf{M}}') \in \mathcal{H}$, where (\mathbf{P}, \mathbf{M}) is

defined as in (1.1), (and the same for $(\mathbf{P}', \mathbf{M}')$, $(\dot{\mathbf{P}}, \dot{\mathbf{M}})$ and $(\dot{\mathbf{P}}', \dot{\mathbf{M}}')$)

$$\left| \begin{aligned} (\mathbf{U}, \mathbf{U}')_{\mathcal{H}} &= \frac{\varepsilon_0}{2} (\mathbf{E}, \mathbf{E}')_{\mathbf{L}^2} + \frac{\mu_0}{2} (\mathbf{H}, \mathbf{H}')_{\mathbf{L}^2} \\ &+ \frac{\varepsilon_0}{2} \sum_{j=1}^{N_e} \omega_{e,j}^2 \Omega_{e,j}^2 (\mathbf{P}_j, \mathbf{P}'_j)_{\mathbf{L}^2} + \frac{\varepsilon_0}{2} \sum_{j=1}^{N_e} \Omega_{e,j}^2 (\dot{\mathbf{P}}_j, \dot{\mathbf{P}}'_j)_{\mathbf{L}^2} \\ &+ \frac{\mu_0}{2} \sum_{\ell=1}^{N_m} \omega_{m,\ell}^2 \Omega_{m,\ell}^2 (\mathbf{M}_\ell, \mathbf{M}'_\ell)_{\mathbf{L}^2} + \frac{\mu_0}{2} \sum_{\ell=1}^{N_m} \Omega_{m,\ell}^2 (\dot{\mathbf{M}}_\ell, \dot{\mathbf{M}}'_\ell)_{\mathbf{L}^2} , \end{aligned} \right. \quad (1.15)$$

where $(\mathbf{u}, \mathbf{v})_{\mathbf{L}^2} = \int_{\mathbb{R}^3} \mathbf{u} \cdot \bar{\mathbf{v}} \, d\mathbf{x}$ denotes the standard inner product of $\mathbf{L}^2(\mathbb{R}^3)$ with \cdot the operation defined on $\mathbb{C}^3 \times \mathbb{C}^3$ by $\mathbf{a} \cdot \mathbf{b} := \mathbf{a}_1 \mathbf{b}_1 + \mathbf{a}_2 \mathbf{b}_2 + \mathbf{a}_3 \mathbf{b}_3$.

The operator $\mathbb{A} : D(\mathbb{A}) \subset \mathcal{H} \rightarrow \mathcal{H}$ is defined by

$$\forall \mathbf{U} = (\mathbf{E}, \mathbf{H}, \mathbf{P}, \dot{\mathbf{P}}, \mathbf{M}, \dot{\mathbf{M}}) \in D(\mathbb{A}), \quad \mathbb{A}\mathbf{U} := -i \begin{pmatrix} -\varepsilon_0^{-1} \nabla \times \mathbf{H} + \sum \Omega_{e,j}^2 \dot{\mathbf{P}}_j \\ \mu_0^{-1} \nabla \times \mathbf{E} + \sum \Omega_{m,\ell}^2 \dot{\mathbf{M}}_\ell \\ -\dot{\mathbf{P}} \\ \alpha_{e,j} \dot{\mathbf{P}}_j + \omega_{e,j}^2 \mathbf{P}_j - \mathbf{E} \\ -\dot{\mathbf{M}} \\ \alpha_{m,\ell} \dot{\mathbf{M}}_\ell + \omega_{m,\ell}^2 \mathbf{M}_\ell - \mathbf{H} \end{pmatrix} \quad (1.16)$$

with $D(\mathbb{A})$ (dense in \mathcal{H}) being given by

$$D(\mathbb{A}) := H(\text{rot}; \mathbb{R}^3) \times H(\text{rot}; \mathbb{R}^3) \times \mathbf{L}^2(\mathbb{R}^3)^{N_e} \times \mathbf{L}^2(\mathbb{R}^3)^{N_m} \times \mathbf{L}^2(\mathbb{R}^3)^{N_m}, \quad (1.17)$$

where $H(\text{rot}; \mathbb{R}^3) := \{\mathbf{u} \in \mathbf{L}^2(\mathbb{R}^3) \mid \nabla \times \mathbf{u} \in \mathbf{L}^2(\mathbb{R}^3)\}$.

Remark 1.5. *Let us point out that, in the sequel, as in the formula (1.16), we shall often omit the summation bounds 1 and N_e (resp. 1 and N_m) for sum on j (resp. ℓ) indices.*

Remark 1.6. *We notice that one can rewrite the evolution problem (1.1) with initial conditions satisfying (1.2, 1.5) as (1.13) with*

$$\mathbf{U}_0 = (\mathbf{E}_0, \mathbf{H}_0, 0, 0, 0, 0) \in \mathcal{H} \quad \text{satisfying} \quad \nabla \cdot \mathbf{E}_0 = \nabla \cdot \mathbf{H}_0 = 0. \quad (1.18)$$

We prove in Lemma A.2 of [5] that the operator $-i\mathbb{A}$ is maximal dissipative. Thus, it generates a contraction semi-group $\{\mathcal{S}(t)\}_{t \geq 0}$ of class \mathcal{C}^0 (see e.g. Theorem 4.3 page 14 of [25]). Hence, the Cauchy problem (1.13) is well-posed and stable. More precisely (see propositions A.3 of [5]), it admits a unique mild solution

$$\mathbf{U} = \mathcal{S}(t)\mathbf{U}_0 \in C^0(\mathbb{R}^+, \mathcal{H}),$$

which is a strong solution,

$$\mathbf{U} \in C^1(\mathbb{R}^+, \mathcal{H}) \cap C^0(\mathbb{R}^+, D(\mathbb{A})),$$

as soon as $\mathbf{U}_0 \in D(\mathbb{A})$, with $D(\mathbb{A})$ is endowed with the usual graph norm $\|\cdot\|_{D(\mathbb{A})}$ defined by

$$\|\mathbf{U}\|_{D(\mathbb{A})}^2 = \|\mathbf{U}\|_{\mathcal{H}}^2 + \|\mathbb{A}\mathbf{U}\|_{\mathcal{H}}^2, \quad \forall \mathbf{U} \in D(\mathbb{A}). \quad (1.19)$$

We introduce now the well-known Hodge decomposition in electromagnetism which consists in writing any $\mathbf{L}^2(\mathbb{R}^3)$ -fields as the orthogonal sum of a gradient (i.e. a curl free field) and a divergence free field (see e.g. [7]). Namely, one has

$$\mathbf{L}^2(\mathbb{R}^3) = \nabla W^1(\mathbb{R}^3) \oplus \mathbf{L}_{\text{div},0}^2, \quad (1.20)$$

where $\nabla W^1(\mathbb{R}^3) = \{\nabla u \mid u \in W^1(\mathbb{R}^3)\}$ with

$$W^1(\mathbb{R}^3) = \{u \in L_{loc}^2(\mathbb{R}^3) \mid (1 + |\mathbf{x}|^2)^{-1/2}u \in L^2(\mathbb{R}^3) \text{ and } \nabla u \in \mathbf{L}^2(\mathbb{R}^3)\}$$

the standard Beppo-Levy space on \mathbb{R}^3 , and $\mathbf{L}_{\text{div},0}^2 = \{\mathbf{u} \in \mathbf{L}^2(\mathbb{R}^3) \mid \nabla \cdot \mathbf{u} = 0\}$ the space of 3D free divergence fields in \mathbb{R}^3 .

Following the decomposition (1.20) for each copy $\mathbf{L}^2(\mathbb{R}^3)$ of $\mathcal{H} = \mathbf{L}^2(\mathbb{R}^3)^N$, one deduces immediately that the Hilbert \mathcal{H} admits the following orthogonal decomposition:

$$\mathcal{H} = \mathcal{H}_{\parallel} \oplus \mathcal{H}_{\perp} \text{ where } \mathcal{H}_{\parallel} = \nabla W^1(\mathbb{R}^3)^N \text{ and } \mathcal{H}_{\perp} = (\mathbf{L}_{\text{div},0}^2)^N. \quad (1.21)$$

Let us now introduce $D(\mathbb{A}_{\parallel}) = D(\mathbb{A}) \cap \mathcal{H}_{\parallel}$ and $D(\mathbb{A}_{\perp}) = D(\mathbb{A}) \cap \mathcal{H}_{\perp}$. As an immediate consequence of (1.16), (1.17), (1.19) and (1.21), one has

$$\mathbb{A}(D(\mathbb{A}_{\parallel})) \subset \mathcal{H}_{\parallel}, \quad \mathbb{A}(D(\mathbb{A}_{\perp})) \subset \mathcal{H}_{\perp} \text{ and } D(\mathbb{A}) = D(\mathbb{A}_{\parallel}) \oplus D(\mathbb{A}_{\perp}). \quad (1.22)$$

Thus one can reduce the operator \mathbb{A} to a sum of two closed and densely defined operators $\mathbb{A}_{\parallel} : D(\mathbb{A}_{\parallel}) \subset \mathcal{H}_{\parallel} \mapsto \mathcal{H}_{\parallel}$ and $\mathbb{A}_{\perp} : D(\mathbb{A}_{\perp}) \subset \mathcal{H}_{\perp} \mapsto \mathcal{H}_{\perp}$ in the sense that

$$\mathbb{A} = \mathbb{A}_{\parallel} \oplus \mathbb{A}_{\perp} \text{ with } \mathbb{A}_{\parallel}\mathbf{U} = \mathbb{A}\mathbf{U}, \quad \forall \mathbf{U} \in D(\mathbb{A}_{\parallel}) \text{ and } \mathbb{A}_{\perp}\mathbf{U} = \mathbb{A}\mathbf{U}, \quad \forall \mathbf{U} \in D(\mathbb{A}_{\perp}). \quad (1.23)$$

Moreover, since any element of \mathcal{H}_{\parallel} is made of curl free vector fields, it is readily seen on (1.16) and (1.17) that $D(\mathbb{A}_{\parallel}) = \mathcal{H}_{\parallel}$ and that \mathbb{A}_{\parallel} is bounded.

From (1.23), we deduce using the Hille-Yosida approximation of \mathbb{A} (see Corollary 3.5 page 11 of [25]) that a similar reduction holds for the semi-group $\{\mathcal{S}(t)\}_{t \geq 0}$. Namely, \mathbb{A}_{\parallel} and \mathbb{A}_{\perp} generate two contraction semi-groups $\{\mathcal{S}_{\parallel}(t)\}_{t \geq 0}$ and $\{\mathcal{S}_{\perp}(t)\}_{t \geq 0}$ on the Hilbert spaces \mathcal{H}_{\parallel} and \mathcal{H}_{\perp} such that for all $t \geq 0$:

$$S(t) = S_{\parallel}(t) \oplus S_{\perp}(t), \quad S_{\parallel}(t)\mathbf{U}_0 = S(t)\mathbf{U}_0, \quad \forall \mathbf{U}_0 \in \mathcal{H}_{\parallel}, \quad S_{\perp}(t)\mathbf{U}_0 = S(t)\mathbf{U}_0, \quad \forall \mathbf{U}_0 \in \mathcal{H}_{\perp}. \quad (1.24)$$

In more physical terms, it means that if the components of the initial conditions $\mathbf{U}_0 \in \mathcal{H}$ are curl free (i.e. if $\mathbf{U}_0 \in \mathcal{H}_{\parallel}$), then the solution $\mathbf{U}(t)$ of (1.13) remains curl free (i.e. $\mathbf{U}(t) \in \mathcal{H}_{\parallel}$ for all $t \geq 0$). In this case, the corresponding dynamics is trivial since it corresponds to a damped

harmonic oscillator in finite dimension (these are non propagative solutions of the Maxwell's equations (1.1)).

Oppositely, the interesting case is when $\mathbf{U}_0 \in \mathcal{H}_\perp$, in other words when the initial conditions are divergence free. Then, the solution $\mathbf{U}(t)$ of (1.13) is propagative and by (1.24), it satisfies

$$\mathbf{U}(t) = S(t)\mathbf{U}_0 = S_\perp(t)\mathbf{U}_0 \in \mathcal{H}_\perp, \quad \forall t \geq 0.$$

In other words, the solution $\mathbf{U}(t)$ remains divergence free for all positive time.

1.3 Statement of the main results

1.3.1 The long-time energy decay rate

To expose our main result (Theorem 1.10), we first need to define an adapted functional framework:

1. For each $p \geq 0$, we introduce the Banach space

$$\mathcal{L}_p(\mathbb{R}^3) = \{v \in \mathcal{S}'(\mathbb{R}^3) / |\mathbf{k}|^{-p}(1 + |\mathbf{k}|^p) \widehat{v} \in L^\infty(\mathbb{R}^3)\} \quad (1.25)$$

(where $\mathbf{k} \mapsto \widehat{v}(\mathbf{k})$ denotes here the Fourier transform of a scalar tempered distributions $v \in \mathcal{S}'(\mathbb{R}^3)$) equipped with the norm

$$\|v\|_{\mathcal{L}_p} := \left\| |\mathbf{k}|^{-p}(1 + |\mathbf{k}|^p) \widehat{v} \right\|_{L^\infty(\mathbb{R}^3)}. \quad (1.26)$$

We point out that if $v \in \mathcal{L}_p(\mathbb{R}^3)$, then one has $|\widehat{v}(\mathbf{k})| \leq \|v\|_{\mathcal{L}_p} |\mathbf{k}|^p$ for a.e. $\mathbf{k} \in \mathbb{R}^3$. The reader will also note that for $q > p$, $\mathcal{L}_q(\mathbb{R}^3) \subset \mathcal{L}_p(\mathbb{R}^3)$.

Identifying a function in $\mathcal{L}_p(\mathbb{R}^3)$ is not trivial since it requires to look at the same time at the function and its Fourier transform. See however Remark 1.7 for some more concrete examples of spaces of functions included in $\mathcal{L}_p(\mathbb{R}^3)$.

We shall introduce $\mathcal{L}_p(\mathbb{R}^3) = \mathcal{L}_p(\mathbb{R}^3)^3$ the vector valued version of $\mathcal{L}_p(\mathbb{R}^3)$ endowed with the norm $\|\cdot\|_{\mathcal{L}_p}$ defined by

$$\forall \mathbf{u} = (u_1, u_2, u_3) \in \mathcal{L}_p(\mathbb{R}^3), \quad \|\mathbf{u}\|_{\mathcal{L}_p}^2 = \|u_1\|_{\mathcal{L}_p}^2 + \|u_2\|_{\mathcal{L}_p}^2 + \|u_3\|_{\mathcal{L}_p}^2.$$

Then, we define on $\mathcal{L}_p(\mathbb{R}^3)^N$ the norm $\|\cdot\|_{\mathcal{L}_p^N}$ by: $\forall \mathbf{U} = (\mathbf{E}, \mathbf{H}, \mathbf{P}, \dot{\mathbf{P}}, \mathbf{M}, \dot{\mathbf{M}}) \in \mathcal{L}_p(\mathbb{R}^3)^N$,

$$\|\mathbf{U}\|_{\mathcal{L}_p^N}^2 = \|\mathbf{E}\|_{\mathcal{L}_p}^2 + \|\mathbf{H}\|_{\mathcal{L}_p}^2 + \sum_j (\|\mathbf{P}_j\|_{\mathcal{L}_p}^2 + \|\dot{\mathbf{P}}_j\|_{\mathcal{L}_p}^2) + \sum_\ell (\|\mathbf{M}_\ell\|_{\mathcal{L}_p}^2 + \|\dot{\mathbf{M}}_\ell\|_{\mathcal{L}_p}^2).$$

2. For $m \geq 0$, we introduce the space $\mathbf{H}^m(\mathbb{R}^3) = H^m(\mathbb{R}^3)^3$ where $H^m(\mathbb{R}^3)$ is the standard Sobolev space defined on scalar functions, endowed with the norm $\|\cdot\|_{\mathbf{H}^m}$ given by

$$\forall \mathbf{u} = (u_1, u_2, u_3) \in \mathbf{H}^m(\mathbb{R}^3), \quad \|\mathbf{u}\|_{\mathbf{H}^m(\mathbb{R}^3)}^2 = \|u_1\|_{H^m(\mathbb{R}^3)}^2 + \|u_2\|_{H^m(\mathbb{R}^3)}^2 + \|u_3\|_{H^m(\mathbb{R}^3)}^2,$$

with $\|\cdot\|_{H^m(\mathbb{R}^3)}$ the usual Sobolev norm on $H^m(\mathbb{R}^3)$ (see e.g. [7] page 500). Then, one defines the norm $\|\cdot\|_{\mathbf{H}^m(\mathbb{R}^3)^N}$ for any $\mathbf{U} = (\mathbf{E}, \mathbf{H}, \mathbf{P}, \dot{\mathbf{P}}, \mathbf{M}, \dot{\mathbf{M}}) \in \mathbf{H}^m(\mathbb{R}^3)^N$ by

$$\left| \begin{aligned} \|\mathbf{U}\|_{\mathbf{H}^m(\mathbb{R}^3)^N}^2 &= \|\mathbf{E}\|_{\mathbf{H}^m(\mathbb{R}^3)}^2 + \|\mathbf{H}\|_{\mathbf{H}^m(\mathbb{R}^3)}^2 \\ &+ \sum_j (\|\mathbf{P}_j\|_{\mathbf{H}^m(\mathbb{R}^3)}^2 + \|\dot{\mathbf{P}}_j\|_{\mathbf{H}^m(\mathbb{R}^3)}^2) \\ &+ \sum_\ell (\|\mathbf{M}_\ell\|_{\mathbf{H}^m(\mathbb{R}^3)}^2 + \|\dot{\mathbf{M}}_\ell\|_{\mathbf{H}^m(\mathbb{R}^3)}^2). \end{aligned} \right. \quad (1.27)$$

Remark 1.7. [On the space $\mathcal{L}_p(\mathbb{R}^3)$] A first trivial observation is that for any $p \geq 0$, $\mathcal{L}_p(\mathbb{R}^3)$ contains all integrable functions whose Fourier transform is supported outside a ball centered at the origin.

Moreover, for $p \in \mathbb{N}$, let us introduce the weighted L^1 space (already used in [5]):

$$L_p^1(\mathbb{R}^3) := \{v \in L^1(\mathbb{R}^3) / (1 + |\mathbf{x}|)^p v \in L^1(\mathbb{R}^3)\} \quad (L_0^1(\mathbb{R}^3) = L^1(\mathbb{R}^3)), \quad (1.28)$$

endowed with the norm

$$\|v\|_{L_p^1(\mathbb{R}^3)} := \|(1 + |\mathbf{x}|)^p v\|_{L^1(\mathbb{R}^3)}. \quad (1.29)$$

In particular, functions of $L_p^1(\mathbb{R}^3)$ have existing moments up to order p . Let us introduce, for $p \geq 1$, the closed subspace of $L_p^1(\mathbb{R}^3)$ of functions whose moments up to order $p - 1$ vanish:

$$L_{p,0}^1(\mathbb{R}^3) := \{u \in L_p^1(\mathbb{R}^3) / \forall |\alpha| \leq p - 1, \int \mathbf{x}^\alpha u \, d\mathbf{x} = 0\}, \quad (1.30)$$

where $\alpha = (\alpha_1, \alpha_2, \alpha_3) \in \mathbb{N}^3$ denotes a multi-index with “length” $|\alpha| = \alpha_1 + \alpha_2 + \alpha_3$ and where $\mathbf{x}^\alpha := \mathbf{x}_1^{\alpha_1} \mathbf{x}_2^{\alpha_2} \mathbf{x}_3^{\alpha_3}$ for all $\mathbf{x} = (x_1, x_2, x_3) \in \mathbb{R}^3$. Moreover, by convention, $L_{0,0}^1(\mathbb{R}^3) := L^1(\mathbb{R}^3)$.

Note that since any function v in $L_p^1(\mathbb{R}^3)$ belongs in particular to $L^1(\mathbb{R}^3)$, its Fourier transform $\widehat{v}(\mathbf{k})$ is well defined and belongs to $C_0^0(\mathbb{R}^3)$, namely the closed subspace of $L^\infty(\mathbb{R}^3)$ made of continuous functions that tend to 0 at infinity. We claim that

$$\forall p \in \mathbb{N}, \quad L_{p,0}^1(\mathbb{R}^3) \subset \mathcal{L}_p. \quad (1.31)$$

Indeed, from well-known properties of the Fourier transform, for all $v \in L_{p,0}^1(\mathbb{R}^3)$, one has

$$\partial^\alpha \widehat{v}(\mathbf{k}) \in C_0^0(\mathbb{R}^3) \text{ for } |\alpha| \leq p \quad \text{and, if } p > 0, \quad \partial^\alpha \widehat{v}(0) = 0, \text{ for } |\alpha| \leq p - 1. \quad (1.32)$$

Moreover, from the Taylor expansion of $\widehat{v}(\mathbf{k})$ at 0, one gets (by the Taylor-Lagrange formula):

$$\exists C(p) > 0 \quad | \widehat{v}(\mathbf{k}) | \leq C(p) \|v\|_{L_p^1(\mathbb{R}^3)} |\mathbf{k}|^p, \quad \forall \mathbf{k} \in \mathbb{R}^3 \quad \text{and} \quad \forall v \in L_{p,0}^1(\mathbb{R}^3). \quad (1.33)$$

Thus, using the fact that $\|\widehat{v}\|_{L^\infty(\mathbb{R}^3)} \leq \|v\|_{L^1(\mathbb{R}^3)} \leq \|v\|_{L_p^1(\mathbb{R}^3)}$, one deduces with (1.33) that

$$|\widehat{v}(\mathbf{k})| \lesssim \min(1, |\mathbf{k}|^p) \|v\|_{L_p^1(\mathbb{R}^3)}, \quad \forall \mathbf{k} \in \mathbb{R}^3 \quad \text{and} \quad \forall v \in L_{p,0}^1(\mathbb{R}^3). \quad (1.34)$$

Then, it follows from (1.26) and (1.34) that the embedding (1.31) is continuous, more precisely:

$$\exists C(p) > 0, \forall v \in L^1_{p,0}(\mathbb{R}^3), \quad \|v\|_{\mathcal{L}_p} \leq C(p) \|v\|_{L^1_p(\mathbb{R}^3)}. \quad (1.35)$$

It also follows from (1.31) that, in particular, \mathcal{L}_p contains all compactly supported integrable functions (in the \mathbf{x} variable) whose moments up to order $p - 1$ vanish.

Next, we introduce the notion of critical configuration associated to the Maxwell's system (1.1), which will influence the long-time decay rate of the electromagnetic energy. It will be enlightened in our proof via the analysis of the dispersion curves associated to this system in the high frequency regime (see Section 4.1).

Definition 1.8. We say that the Maxwell's system (1.1) is in a critical configuration if the weak dissipation condition (1.11) holds and if one of the following conditions is satisfied:

1. $\forall \ell \in \{1, \dots, N_m\}, \alpha_{m,\ell} = 0$ and $\exists j \in \{1, \dots, N_e\} \mid \alpha_{e,j} = 0$ and $\omega_{e,j} \notin \{\omega_{m,\ell}\}$.
2. $\forall j \in \{1, \dots, N_e\}, \alpha_{e,j} = 0$ and $\exists \ell \in \{1, \dots, N_m\} \mid \alpha_{m,\ell} = 0$ and $\omega_{m,\ell} \notin \{\omega_{e,j}\}$.

When none of these condition holds, the system (1.1) is in a non-critical configuration.

We point out that, as the reader will easily check, under the weak dissipation assumption (1.11), the critical configuration can occur only if $N_e \geq 2$ or $N_m \geq 2$.

We are now able to state our main result under the weak dissipation assumption concerning the decay of the total $\|\mathbf{U}(t)\|_{\mathcal{H}}^2$ associated to the evolution problem (1.13) with divergence-free initial conditions $\mathbf{U}_0 \in \mathcal{H}_\perp$ which contain in particular the initial conditions given by (1.2,1.5). Thus, this result applies to the Maxwell's system (1.1) with initial conditions (1.2) in \mathcal{H}_\perp .

We shall use below, as well as in the rest of this paper, the following

Notation 1.9. To compare two non-negative functions $f(y)$ and $g(y)$, where $y \in Y$ and $y = \mathbf{x}, \mathbf{k}, t$ or any combination of the variables, we introduce the notation:

$$f \lesssim g \iff \exists C > 0 \quad | \quad f(y) \leq C g(y), \quad \forall y \in Y,$$

where the constant C depends only on N_e, N_m and the coefficients of the system (1.1).

Theorem 1.10. Let assume that the Maxwell's system (1.1) satisfies the weak dissipation assumption (1.11) and the irreducibility assumptions (H₁) and (H₂). Then, for any initial condition $\mathbf{U}_0 \in \mathcal{H}_\perp$, one has

$$\lim_{t \rightarrow +\infty} \|\mathbf{U}(t)\|_{\mathcal{H}}^2 = 0. \quad (1.36)$$

Moreover if for some real numbers $m > 0$ and $p \geq 0$, $\mathbf{U}_0 \in \mathbf{H}^m(\mathbb{R}^3)^N \cap \mathcal{L}_p^N \cap \mathcal{H}_\perp$ then the decay rate of $\|\mathbf{U}(t)\|_{\mathcal{H}}^2$ is polynomial. However, the associated exponents in $1/t$ depend on the configuration (see definition 1.8) of the corresponding Maxwell's system (1.1). More precisely,

1. If the Maxwell system is in a non-critical configuration, then one has

$$\|\mathbf{U}(t)\|_{\mathcal{H}}^2 \leq \frac{C_{\text{HF}}^m(\mathbf{U}_0)}{t^m} + \frac{C_{\text{LF}}^p(\mathbf{U}_0)}{t^{p+\frac{3}{2}}}, \quad \forall t > 0, \quad (1.37)$$

2. If the Maxwell system is in a critical configuration, then one has

$$\|\mathbf{U}(t)\|_{\mathcal{H}}^2 \leq \frac{C_{\text{HF}}^m(\mathbf{U}_0)}{t^{\frac{m}{2}}} + \frac{C_{\text{LF}}^p(\mathbf{U}_0)}{t^{p+\frac{3}{2}}}, \quad \forall t > 0, \quad (1.38)$$

where in (1.37) and (1.38) the constants satisfy

$$C_{\text{HF}}^m(\mathbf{U}_0) \lesssim \|\mathbf{U}_0\|_{\mathbf{H}^m(\mathbb{R}^3)^N}^2 \quad \text{and} \quad C_{\text{LF}}^p(\mathbf{U}_0) \lesssim \|\mathbf{U}_0\|_{\mathcal{L}_p^N}^2. \quad (1.39)$$

Remark 1.11. When the initial data \mathbf{U}_0 satisfy also (1.5), the estimates (1.37) and (1.38) give estimates involving the electromagnetic energy only since $\mathcal{E}(t) \leq \|\mathbf{U}(t)\|_{\mathcal{H}}^2$ while $\mathcal{E}(0) = \|\mathbf{U}_0\|_{\mathcal{H}}^2$.

Remark 1.12. The energy $\|\mathbf{U}(t)\|_{\mathcal{H}}^2$ is dominated in (1.37) and (1.38) by a sum of two terms:

- The first one involves the Sobolev regularity of the initial condition and is linked, as it will appear in the proof, to the decay of the high spatial frequency components of the energy.
- The second is related to the condition $\mathbf{U}_0 \in \mathcal{L}_p^N$ and thus to the decay of the low spatial frequency components of the energy.

The presence of the second term is directly related to the fact that the domain of propagation is unbounded and does not appear in the control of the energy in bounded domains for perfectly conducting materials (see, [21, 22, 23]).

Remark 1.13. For generalized Lorentz media, this Theorem generalizes under the weak dissipation assumption (1.11) the results obtained under the more restrictive strong dissipation assumption (1.9) in [21, 22, 23] for bounded domains and in [5] for the whole space \mathbb{R}^3 .

As (1.9) excludes the critical configuration, the weaker decay (1.38) could not be observed in [5, 21, 22, 23].

1.3.2 Optimality of the bounds

In Theorem 1.10, we prove a long-time polynomial decay rate for $\|\mathbf{U}(t)\|$ (cf. (1.37) and (1.38)) for solutions $\mathbf{U}(t)$ of the Cauchy problem (1.13) which satisfies divergence free initial conditions $\mathbf{U}_0 \in \mathcal{H}_\perp$. In this section, we analyse the question of the optimality of this decay rate.

The estimates (1.37) and (1.38) involved a sum of two terms which are related to the decay of the low and high spatial frequency components of the solution (see Remark 1.12). Thus, to speak about optimal polynomial exponents, we need to decouple these two terms by separating the high and low frequency behaviors of the solution \mathbf{U} . Therefore, we introduce the 3D spatial Fourier transform defined by

$$\mathbb{G}(\mathbf{k}) = \mathcal{F}(\mathbf{G})(\mathbf{k}) = \frac{1}{(2\pi)^{\frac{3}{2}}} \int_{\mathbb{R}^3} \mathbf{G}(\mathbf{x}) e^{-i\mathbf{k}\cdot\mathbf{x}} d\mathbf{x} \quad \forall \mathbf{G} \in \mathbf{L}^1(\mathbb{R}^3) \cap \mathbf{L}^2(\mathbb{R}^3),$$

which extends to a unitary transformation from $\mathbf{L}^2(\mathbb{R}_x^3)$ to $\mathbf{L}^2(\mathbb{R}_k^3)$. Applying \mathcal{F} to each copy $\mathbf{L}^2(\mathbb{R}^3)$ of $\mathcal{H} = \mathbf{L}^2(\mathbb{R}^3)^N$ defines a unitary transform, still denoted \mathcal{F} , in \mathcal{H} .

We now introduce for $(m, p) \in \mathbb{R}^{+,*} \times \mathbb{R}^+$ the spaces

$$\mathcal{H}_{\perp, \text{HF}}^m = \{\mathbf{V} \in \mathcal{H}_{\perp} \cap \mathbf{H}^m(\mathbb{R}^3)^N \mid \exists k_+ > 0 \mid \text{supp}(\mathcal{F}(\mathbf{V})) \subset \mathbb{R}^3 \setminus B(0, k_+)\}, \quad (1.40)$$

$$\mathcal{H}_{\perp, \text{LF}}^p = \{\mathbf{V} \in \mathcal{H}_{\perp} \cap \mathcal{L}_p(\mathbb{R}^3)^N \mid \exists k_- > 0 \mid \text{supp}(\mathcal{F}(\mathbf{V})) \subset \overline{B(0, k_-)}\}. \quad (1.41)$$

In the spirit of [21], we define for any $m > 0$ the high frequency optimal polynomial exponent of solutions $\mathbf{U}(t)$ of the Cauchy problem (1.13) with initial conditions $\mathbf{U}_0 \in \mathcal{H}_{\perp, \text{HF}}^m$ as follows:

$$\gamma_m^{\text{HF}} = \sup\{\gamma \in \mathbb{R}^+ \mid \forall \mathbf{U}_0 \in \mathcal{H}_{\perp, \text{HF}}^m, \exists C(\mathbf{U}_0) > 0 \mid \|\mathbf{U}(t)\|_{\mathcal{H}}^2 \lesssim \frac{C(\mathbf{U}_0)}{t^\gamma}, \quad \forall t \geq 1\}. \quad (1.42)$$

Similarly, we define for any $p \geq 0$ the low frequency optimal polynomial exponent:

$$\gamma_p^{\text{LF}} = \sup\{\gamma \in \mathbb{R}^+ \mid \forall \mathbf{U}_0 \in \mathcal{H}_{\perp, \text{LF}}^p, \exists C(\mathbf{U}_0) > 0 \mid \|\mathbf{U}(t)\|_{\mathcal{H}}^2 \lesssim \frac{C(\mathbf{U}_0)}{t^\gamma}, \quad \forall t \geq 1\}. \quad (1.43)$$

We are now able to state our result concerning the low frequency and high frequency optimal polynomial decay rates.

Theorem 1.14. *Let $(m, p) \in \mathbb{R}^{+,*} \times \mathbb{R}^+$. The exponents γ_m^{HF} and γ_p^{LF} are given as follows:*

1. *If the Maxwell system is in a non-critical configuration, $\gamma_m^{\text{HF}} = m$ and $\gamma_p^{\text{LF}} = p + 3/2$.*
2. *If the Maxwell system is in a critical configuration, $\gamma_m^{\text{HF}} = m/2$ and $\gamma_p^{\text{LF}} = p + 3/2$.*

We point out that in Section 7.1, we show that the upperbounds (1.37) and (1.38) give the right lower bounds on the exponents γ_m^{HF} and γ_p^{LF} . In Section 7.2, we show that these lower-bounds are also upper-bounds and thus deduce the value of these exponents.

2 Fourier reduction

2.1 The reduced operator $\mathbb{A}_{\mathbf{k}}$

The homogeneity of the propagation medium allows us to reduce the spectral analysis of the operator \mathbb{A} defined in (1.16) to the spectral analysis of a family (non-self-adjoint) operators $(\mathbb{A}_{\mathbf{k}})_{\mathbf{k} \in \mathbb{R}^3}$ on a finite dimensional space.

We introduce the space \mathbf{C}^N with $\mathbf{C} = \mathbb{C}^3$ endowed with the inner product $(\cdot, \cdot)_{\mathbf{C}^N}$ defined via the expression (1.15) by replacing $\mathbf{U} \in \mathcal{H}$ by $\mathbb{U} = (\mathbb{E}, \mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}}) \in \mathbf{C}^N$ with

$$\mathbb{P} = (\mathbb{P}_j), \dot{\mathbb{P}} = (\dot{\mathbb{P}}_j), \mathbb{M} = (\mathbb{M}_\ell), \dot{\mathbb{M}} = (\dot{\mathbb{M}}_\ell),$$

(the same for \mathbf{U}') and the \mathbf{L}^2 inner product $(\cdot, \cdot)_{\mathbf{L}^2}$ by the usual one in \mathbf{C} : $(\mathbb{E}, \mathbb{E}') = \mathbb{E} \cdot \overline{\mathbb{E}'}$.

The corresponding norms in \mathbf{C} or \mathbf{C}^N are both denoted $|\cdot|$ for simplicity.

In this functional framework, \mathbb{A} is unitarily equivalent via \mathcal{F} to a direct integral of operators $\mathbb{A}_{\mathbf{k}}$ defined on the finite dimension space \mathbf{C}^N . Namely, one has:

$$\mathbb{A} = \mathcal{F}^* \left(\int_{\mathbb{R}^3}^{\oplus} \mathbb{A}_{\mathbf{k}} d\mathbf{k} \right) \mathcal{F} \quad \text{i.e.} \quad \mathcal{F}(\mathbb{A} \mathbf{U})(\mathbf{k}) = \mathbb{A}_{\mathbf{k}} \mathcal{F} \mathbf{U}(\mathbf{k}), \quad \forall \mathbf{U} \in D(\mathbb{A}) \text{ and a.e. } \mathbf{k} \in \mathbb{R}^3, \quad (2.1)$$

where the (bounded) linear operator $\mathbb{A}_{\mathbf{k}} : \mathbf{C}^N \mapsto \mathbf{C}^N$ is given by

$$\forall \mathbf{U} = (\mathbb{E}, \mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}}) \in \mathbf{C}^N, \quad \mathbb{A}_{\mathbf{k}} \mathbf{U} = \begin{pmatrix} -\frac{\mathbf{k} \times \mathbb{H}}{\varepsilon_0} - i \sum \Omega_{e,j}^2 \dot{\mathbb{P}}_j \\ \frac{\mathbf{k} \times \mathbb{E}}{\mu_0} - i \sum \Omega_{m,\ell}^2 \dot{\mathbb{M}}_\ell \\ i \dot{\mathbb{P}} \\ -i \alpha_{e,j} \dot{\mathbb{P}}_j - i \omega_{e,j}^2 \mathbb{P}_j + i \mathbb{E} \\ i \dot{\mathbb{M}} \\ -i \alpha_{m,\ell} \dot{\mathbb{M}}_\ell - i \omega_{m,\ell}^2 \mathbb{M}_\ell + i \mathbb{H} \end{pmatrix}. \quad (2.2)$$

We point out that $\mathbb{A}_{\mathbf{k}}$ is deduced from the definition (1.16) of \mathbb{A} by replacing the curl operator $\nabla \times$ operator by its spatial Fourier counterpart $i \mathbf{k} \times$. The usual cross product on $\mathbb{R}^3 \times \mathbb{R}^3$ has been here extended to $\mathbf{C} \times \mathbf{C}$ via the formula $\mathbf{a} \times \mathbf{b} = (\mathbf{a}_2 \mathbf{b}_3 - \mathbf{a}_3 \mathbf{b}_2, \mathbf{a}_3 \mathbf{b}_1 - \mathbf{a}_1 \mathbf{b}_3, \mathbf{a}_1 \mathbf{b}_2 - \mathbf{a}_2 \mathbf{b}_1)$.

Remark 2.1. A simple computation shows that for any $\mathbf{U} = (\mathbb{E}, \mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}}) \in \mathbf{C}^N$:

$$\text{Im}(\mathbb{A}_{\mathbf{k}} \mathbf{U}, \mathbf{U})_{\mathbf{C}^N} = - \sum_{j=1}^{N_e} \frac{\varepsilon_0}{2} \alpha_{e,j} \Omega_{e,j}^2 |\dot{\mathbb{P}}_j|^2 - \sum_{\ell=1}^{N_m} \frac{\mu_0}{2} \alpha_{m,\ell} \Omega_{m,\ell}^2 |\dot{\mathbb{M}}_\ell|^2 \leq 0.$$

Thus, for all $\mathbf{k} \in \mathbb{R}^3$, the spectrum of $\mathbb{A}_{\mathbf{k}}$, $\sigma(\mathbb{A}_{\mathbf{k}})$ is included in $\overline{\mathbb{C}^-} = \{\omega \in \mathbb{C} \mid \text{Im}(\omega) \leq 0\}$.

Hence, $-i \mathbb{A}(\mathbf{k})$ is maximal dissipative and the resolvent of $\mathbb{A}(\mathbf{k})$, $R_{\mathbb{A}_{\mathbf{k}}}(\omega) = (\mathbb{A}_{\mathbf{k}} - \omega \text{Id})^{-1}$ is well-defined in \mathbb{C}^+ and satisfies:

$$\|R_{\mathbb{A}_{\mathbf{k}}}(\omega)\| \leq \text{Im}(\omega)^{-1}, \quad \forall \omega \in \mathbb{C}^+ \text{ and } \forall \mathbf{k} \in \mathbb{R}^3.$$

This estimate justifies the existence of the direct integral $\int_{\mathbf{k} \in \mathbb{R}^3}^{\oplus} \mathbb{A}_{\mathbf{k}} d\mathbf{k}$ (see [20] page 5).

From an operator point of view, one deduces from (2.1) (see e.g. Theorem 4.2 of [20]) that the contractions $S(t)$ of the semigroup $\{S(t)\}_{t \geq 0}$ are also unitarily equivalent via \mathcal{F} to a direct integral of contractions $e^{-i \mathbb{A}_{\mathbf{k}} t}$ on \mathbf{C}^N . More precisely, one has for any $t \geq 0$:

$$S(t) = \mathcal{F}^* \left(\int_{\mathbb{R}^3}^{\oplus} e^{-i \mathbb{A}_{\mathbf{k}} t} d\mathbf{k} \right) \mathcal{F}, \text{ i.e. } \mathcal{F}(S(t) \mathbf{U})(\mathbf{k}) = e^{-i \mathbb{A}_{\mathbf{k}} t} \mathcal{F} \mathbf{U}(\mathbf{k}), \quad \forall \mathbf{U} \in \mathcal{H} \text{ a.e. } \mathbf{k} \in \mathbb{R}^3. \quad (2.3)$$

From a more practical point of view, it means that the Fourier transform $\mathbb{U} = \mathcal{F} \mathbf{U}$ of the solution \mathbf{U} of the evolution problem (1.13) with initial condition $\mathbf{U}_0 \in \mathcal{H}$ satisfies a family of Ordinary Differential Equations (ODE) in \mathbf{C}^N parametrized by the wave vector \mathbf{k} . Namely, one has for a.e. $\mathbf{k} \in \mathbb{R}^3$:

$$\frac{d\mathbb{U}(\mathbf{k}, t)}{dt} + i \mathbb{A}_{\mathbf{k}} \mathbb{U}(\mathbf{k}, t) = 0 \quad \text{for } t \geq 0 \text{ with } \mathbb{U}(\mathbf{k}, 0) = \mathbb{U}_0(\mathbf{k}) = (\mathcal{F} \mathbf{U}_0)(\mathbf{k}). \quad (2.4)$$

The solution of each these ODE's is given by $\mathbb{U}(\mathbf{k}, t) = e^{-i \mathbb{A}_{\mathbf{k}} t} \mathbb{U}_0(\mathbf{k})$, $\forall t \geq 0$. We can now introduce the counterpart of the Hodge decomposition (1.21) in the (spatial) frequency domain. In this perspective, one decomposes orthogonally the space \mathbf{C}^N , for all wave vector $\mathbf{k} \neq 0$, as

$$\mathbf{C}^N = \mathbf{C}_{\mathbf{k}, \parallel}^N \oplus^{\perp} \mathbf{C}_{\mathbf{k}, \perp}^N \text{ where } \mathbf{C}_{\mathbf{k}, \parallel} := \text{span } \mathbf{k} \text{ and } \mathbf{C}_{\mathbf{k}, \perp} := \{\mathbf{k}\}^{\perp}. \quad (2.5)$$

One shows easily that $\mathbf{C}_{\mathbf{k},\parallel}^N$ and $\mathbf{C}_{\perp,\mathbf{k}}^N$ are stable by $\mathbb{A}_{\mathbf{k}}$. Thus, one can reduce $\mathbb{A}_{\mathbf{k}}$ as a sum of two operators

$$\mathbb{A}_{\mathbf{k},\parallel} : \mathbf{C}_{\mathbf{k},\parallel}^N \rightarrow \mathbf{C}_{\mathbf{k},\parallel}^N \quad \text{and} \quad \mathbb{A}_{\mathbf{k},\perp} : \mathbf{C}_{\mathbf{k},\perp}^N \rightarrow \mathbf{C}_{\mathbf{k},\perp}^N$$

in such way that

$$\mathbb{A}_{\mathbf{k}} = \mathbb{A}_{\mathbf{k},\parallel} \oplus \mathbb{A}_{\mathbf{k},\perp} \quad \text{with} \quad \mathbb{A}_{\mathbf{k},\parallel} \mathbf{U} = \mathbb{A}_{\mathbf{k}} \mathbf{U}, \forall \mathbf{U} \in \mathbf{C}_{\mathbf{k},\parallel}^N \quad \text{and} \quad \mathbb{A}_{\mathbf{k},\perp} \mathbf{U} = \mathbb{A}_{\mathbf{k}} \mathbf{U}, \forall \mathbf{U} \in \mathbf{C}_{\mathbf{k},\perp}^N. \quad (2.6)$$

The above relation is nothing but the counterpart of (1.23) for Fourier components. Indeed, following the decompositions (1.23) and (2.6), one proves easily that the operators \mathbb{A}_{\perp} and the contractions $S_{\perp}(t)$ of its associated semigroup $\{\mathcal{S}_{\perp}(t)\}_{t \geq 0}$ are also unitarily equivalent via \mathcal{F} to direct integral of operators $\mathbb{A}_{\mathbf{k},\perp}$ and $e^{-i\mathbb{A}_{\mathbf{k},\perp}t}$:

$$\begin{aligned} \mathbb{A}_{\perp} &= \mathcal{F}^* \left(\int_{\mathbb{R}^3}^{\oplus} \mathbb{A}_{\mathbf{k},\perp} d\mathbf{k} \right) \mathcal{F} \quad \text{i.e.} \quad \mathcal{F}(\mathbb{A}_{\perp} \mathbf{U})(\mathbf{k}) = \mathbb{A}_{\mathbf{k},\perp} \mathcal{F} \mathbf{U}(\mathbf{k}), \forall \mathbf{U} \in D(\mathbb{A}_{\perp}), \text{ a.e. } \mathbf{k} \in \mathbb{R}^3, \\ S_{\perp}(t) &= \mathcal{F}^* \left(\int_{\mathbb{R}^3}^{\oplus} e^{-i\mathbb{A}_{\mathbf{k},\perp}t} d\mathbf{k} \right) \mathcal{F} \quad \text{i.e.} \quad \mathcal{F}(S_{\perp}(t) \mathbf{U})(\mathbf{k}) = e^{-i\mathbb{A}_{\mathbf{k},\perp}t} \mathcal{F} \mathbf{U}(\mathbf{k}), \forall \mathbf{U} \in \mathcal{H}_{\perp}, \text{ a.e. } \mathbf{k} \in \mathbb{R}^3. \end{aligned}$$

Similarly, the operators \mathbb{A}_{\parallel} and $\mathcal{S}_{\parallel}(t)$ for $t \geq 0$ are also unitarily equivalent via \mathcal{F} to direct integral of operators and the latter relations hold also if one replaces all \perp -symbols by \parallel -symbols.

The direct integral decomposition of the operator $S_{\perp}(t)$ implies that the Fourier transform $\mathbb{U} = \mathcal{F} \mathbf{U}$ of the solution \mathbf{U} of the evolution problem (1.13) with initial divergence free conditions $\mathbf{U}_0 \in \mathcal{H}_{\perp}$ (as e.g. in (1.2)) satisfies for a.e. $\mathbf{k} \in \mathbb{R}^3$:

$$\mathbb{U}(\mathbf{k}, t) = e^{-i\mathbb{A}_{\mathbf{k},\perp}t} \mathbb{U}_0(\mathbf{k}), \quad \forall t \geq 0 \quad \text{where} \quad \mathbb{U}_0(\mathbf{k}) = (\mathcal{F} \mathbf{U}_0)(\mathbf{k}) \in \mathbf{C}_{\mathbf{k},\perp}^N. \quad (2.7)$$

2.2 From $\mathbb{A}_{\mathbf{k}}$ to $\mathbb{A}_{|\mathbf{k}|}$

In the expression of the solution (2.7), the space $\mathbf{C}_{\perp,\mathbf{k}}$ depends on \mathbf{k} which complicates slightly the analysis of the decay of the Fourier components $\mathbb{U}(\mathbf{k}, t)$. To remediate this point, one can use the isotropic character of the medium.

We introduce $(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$ the canonical orthonormal basis of \mathbf{C} and for $\mathbf{k} \neq 0$ the unit vector $\widehat{\mathbf{k}} = \mathbf{k}/|\mathbf{k}|$. We now construct a rotation $R_{\mathbf{k}} : \mathbf{C} \rightarrow \mathbf{C}$ that maps the vector $\widehat{\mathbf{k}}$ into \mathbf{e}_3 as follows

- If $\widehat{\mathbf{k}} \neq \pm \mathbf{e}_3$, we set $\mathbf{w}_{\mathbf{k}} := (\widehat{\mathbf{k}} \times \mathbf{e}_3)/|\widehat{\mathbf{k}} \times \mathbf{e}_3| \in \mathbf{C}_{\mathbf{k},\perp}$ so that $(\widehat{\mathbf{k}}, \mathbf{w}_{\mathbf{k}}, \widehat{\mathbf{k}} \times \mathbf{w}_{\mathbf{k}})$ is an orthonormal basis of \mathbf{C} . We then define $R_{\mathbf{k}}$ in this basis by

$$R_{\mathbf{k}}(\widehat{\mathbf{k}}) = \mathbf{e}_3, \quad R_{\mathbf{k}}(\mathbf{w}_{\mathbf{k}}) = \mathbf{e}_1, \quad R_{\mathbf{k}}(\widehat{\mathbf{k}} \times \mathbf{w}_{\mathbf{k}}) = \mathbf{e}_2.$$

- If $\widehat{\mathbf{k}} = \mathbf{e}_3$, we simply set $R_{\mathbf{k}} = \text{Id}_{\mathbf{C}}$.
- If $\widehat{\mathbf{k}} = -\mathbf{e}_3$, we define $R_{\mathbf{k}}$ by $R_{\mathbf{k}}(\mathbf{e}_1) = -\mathbf{e}_2$, $R_{\mathbf{k}}(\mathbf{e}_2) = -\mathbf{e}_1$ and $R_{\mathbf{k}}(\mathbf{e}_3) = -\mathbf{e}_3$.

Finally, we define the operator $\mathcal{R}_{\mathbf{k}} : \mathbf{C}^N \rightarrow \mathbf{C}^N$ defined for all $\mathbf{U} = (\mathbb{E}, \mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}}) \in \mathbf{C}^N$ by

$$\mathcal{R}_{\mathbf{k}} \mathbf{U} = (R_{\mathbf{k}} \mathbb{E}, R_{\mathbf{k}} \mathbb{H}, (R_{\mathbf{k}} \mathbb{P}_j), (R_{\mathbf{k}} \dot{\mathbb{P}}_j), (R_{\mathbf{k}} \mathbb{M}_{\ell}), (R_{\mathbf{k}} \dot{\mathbb{M}}_{\ell})). \quad (2.8)$$

Using the identity (we leave its proof to the reader):

$$\mathbf{R}_{\mathbf{k}}^*(\mathbf{R}_{\mathbf{k}}(\mathbf{u}) \times \mathbf{R}_{\mathbf{k}}(\mathbf{v})) = \mathbf{u} \times \mathbf{v}, \quad \forall \mathbf{u}, \mathbf{v} \in \mathbf{C}$$

applied to $\mathbf{u} = \mathbf{k}$ (and thus with $\mathbf{R}_{\mathbf{k}}(\mathbf{u}) = |\mathbf{k}| e_3$), one observes with (2.2) that

$$\mathbb{A}_{\mathbf{k}} = \mathcal{R}_{\mathbf{k}}^* \mathbb{A}_{|\mathbf{k}|} \mathcal{R}_{\mathbf{k}} \quad \text{where} \quad \mathbb{A}_{|\mathbf{k}|} := \mathbb{A}_{|\mathbf{k}|e_3}. \quad (2.9)$$

Thus, $\mathbb{A}_{\mathbf{k}}$ and $\mathbb{A}_{|\mathbf{k}|}$ are unitarily equivalent. Furthermore, as

$$\mathcal{R}_{\mathbf{k}}(\mathbf{C}_{\mathbf{k},\parallel}^N) = \mathbf{C}_{\parallel}^N \quad \text{and} \quad \mathcal{R}_{\mathbf{k}}(\mathbf{C}_{\perp,\mathbf{k}}^N) = \mathbf{C}_{\perp}^N \quad \text{where} \quad \mathbf{C}_{\parallel} := \mathbf{C}_{\parallel,e_3} \quad \text{and} \quad \mathbf{C}_{\perp} := \mathbf{C}_{\perp,e_3}, \quad (2.10)$$

one has by (2.6):

$$\mathbb{A}_{\mathbf{k},\parallel} = \mathcal{R}_{\mathbf{k}}^* \mathbb{A}_{|\mathbf{k}|,\parallel} \mathcal{R}_{\mathbf{k}} \quad \text{and} \quad \mathbb{A}_{\mathbf{k},\perp} = \mathcal{R}_{\mathbf{k}}^* \mathbb{A}_{|\mathbf{k}|,\perp} \mathcal{R}_{\mathbf{k}}. \quad (2.11)$$

In particular, it shows with (2.7) that the Fourier transform $\mathbb{U} = \mathcal{F}\mathbf{U}$ of the solution \mathbf{U} of (1.13) with initial divergence free conditions $\mathbf{U}_0 \in \mathcal{H}_{\perp}$ satisfies for a.e. $\mathbf{k} \in \mathbb{R}^3$:

$$\mathbb{U}(\mathbf{k}, t) = \mathcal{R}_{\mathbf{k}}^* e^{-i\mathbb{A}_{|\mathbf{k}|,\perp} t} \mathcal{R}_{\mathbf{k}} \mathbb{U}_0(\mathbf{k}), \quad \forall t \geq 0 \quad \text{where} \quad \mathbb{U}_0(\mathbf{k}) = (\mathcal{F}\mathbf{U}_0)(\mathbf{k}) \in \mathbf{C}_{\mathbf{k},\perp}^N. \quad (2.12)$$

Thus, to estimate the norm of $|\mathbb{U}(\mathbf{k}, t)|$, one only needs to analyse the spectral properties of $\mathbb{A}_{|\mathbf{k}|,\perp}$ on the space \mathbf{C}_{\perp}^N which is now independent of \mathbf{k} .

3 Modal analysis

The proofs of Theorems 1.10 and Theorem 1.14 are based on the spectral properties of $\mathbb{A}_{|\mathbf{k}|}$. These theorems are proved in Section 7 based on results established in Sections 3 to 6.

3.1 Spectrum and resolvent of the finite dimensional operators $\mathbb{A}_{|\mathbf{k}|}$

Using the expression (1.6), the rational functions ε and μ can be rewritten as

$$\varepsilon(\omega) = \varepsilon_0 \frac{P_e(\omega)}{Q_e(\omega)} \quad \text{and} \quad \mu(\omega) = \mu_0 \frac{P_m(\omega)}{Q_m(\omega)}, \quad (3.1)$$

where the unitary polynomials P_e and Q_e (resp. P_m and Q_m) are of degree $2N_e$ (resp. $2N_m$). More precisely, Q_e and P_e are explicitly given in term of the $q_{e,j}$ by

$$Q_e(\omega) = \prod_{j=1}^{N_e} q_{e,j}(\omega) \quad \text{and} \quad P_e(\omega) = Q_e(\omega) - \sum_{j=1}^{N_e} \Omega_{e,j}^2 \prod_{k=1, k \neq j}^{N_e} q_{e,k}(\omega). \quad (3.2)$$

The reader will verify that, with (3.2), the assumption (H₁) implies that P_e and Q_e do not share common zeroes. Therefore the representation (3.1) of the rational function ε is irreducible.

Thus, \mathcal{P}_e , the set of poles of ε , is exactly the set of roots of Q_e which is the union over j of the roots of $q_{e,j}$ and is therefore included in $\overline{\mathcal{C}}^-$ (see remark 1.2).

Moreover, by (H₁) and remark 1.2, the multiplicity of a root ω_* of Q_e is either 1 or 2.

Obviously, similar observations hold true for the polynomials P_m and Q_m .

We introduce the sets $\mathcal{P} := \mathcal{P}_e \cup \mathcal{P}_m$ and $\mathcal{Z} := \mathcal{Z}_e \cup \mathcal{Z}_m$, where we recall that \mathcal{P}_e (resp. \mathcal{P}_m) is the set of poles of the function ε (resp. μ) and \mathcal{Z}_e (resp. \mathcal{Z}_m) the set of zeros of ε (resp. μ) and the rational function $\mathcal{D} : \mathbb{C} \setminus \mathcal{P} \rightarrow \mathbb{C}$ defined by

$$\mathcal{D}(\omega) = \omega^2 \varepsilon(\omega) \mu(\omega), \quad \forall \omega \in \mathbb{C} \setminus \mathcal{P}. \quad (3.3)$$

Using (H₂), it is easy to see that $\mathcal{D}(\omega)$ is irreducible. From (3.1) and (3.2), one deduces immediately that its numerator is of degree $N = 2 + 2N_e + 2N_m$ and its denominator is of degree $2N_e + 2N_m$. Moreover, its poles p and zeros z are respectively the elements \mathcal{P} and $\mathcal{Z} \cup \{0\}$.

We introduce for any $\mathbf{k} \in \mathbb{R}^3$, the set

$$S(|\mathbf{k}|) := \{\omega \in \mathbb{C} \setminus \mathcal{P} \mid \mathcal{D}(\omega) = |\mathbf{k}|^2\}.$$

The following proposition establishes the link between $S(|\mathbf{k}|)$ and the spectrum of the operator $\mathbb{A}_{|\mathbf{k}|,\perp}$ for $\mathbf{k} \neq 0$.

Proposition 3.1. *Let assume (H₁) and (H₂). Then for any $\mathbf{k} \in \mathbb{R}^3 \setminus \{0\}$, one has*

$$\sigma(\mathbb{A}_{|\mathbf{k}|,\perp}) = S(|\mathbf{k}|). \quad (3.4)$$

Moreover, each eigenvalue $\omega \in \sigma(\mathbb{A}_{|\mathbf{k}|,\perp})$ has a geometric multiplicity of 2.

Finally, we conclude this section with the expression of the resolvent

$$R_{|\mathbf{k}|}(\omega) := (\mathbb{A}_{|\mathbf{k}|,\perp} - \omega \mathbf{I})^{-1}. \quad (3.5)$$

To obtain a readable expression, it appears useful to introduce some intermediate operators.

We first define four linear operators in $\mathcal{L}(\mathbf{C}_\perp^N, \mathbf{C}_\perp)$: given $\mathbb{F} := (\mathbf{e}, \mathbf{h}, \mathbf{p}_j, \dot{\mathbf{p}}_j, \mathbf{m}_\ell, \dot{\mathbf{m}}_\ell) \in \mathbf{C}_\perp^N$

$$\begin{aligned} \mathbb{A}_{e,j}(\omega) \mathbb{F} &:= \frac{(-i\alpha_{e,j} - \omega) \mathbf{p}_j - i \dot{\mathbf{p}}_j}{q_{e,j}(\omega)}, & \mathbb{A}_{m,\ell}(\omega) \mathbb{F} &:= \frac{(-i\alpha_{m,\ell} - \omega) \mathbf{m}_\ell - i \dot{\mathbf{m}}_\ell}{q_{m,\ell}(\omega)}, \\ \dot{\mathbb{A}}_{e,j}(\omega) \mathbb{F} &:= \frac{i\omega_{e,j}^2 \mathbf{p}_j - \omega \dot{\mathbf{p}}_j}{q_{e,j}(\omega)}, & \dot{\mathbb{A}}_{m,\ell}(\omega) \mathbb{F} &:= \frac{i\omega_{m,\ell}^2 \mathbf{m}_\ell - \omega \dot{\mathbf{m}}_\ell}{q_{m,\ell}(\omega)}, \end{aligned} \quad (3.6)$$

from which we define two more operators in $\mathcal{L}(\mathbf{C}_\perp^N, \mathbf{C}_\perp)$:

$$\mathbb{A}_e(\omega) \mathbb{F} = -\varepsilon_0 \left(\mathbf{e} + i \sum \Omega_{e,j}^2 \dot{\mathbb{A}}_{e,j}(\omega) \mathbb{F} \right), \quad \mathbb{A}_m(\omega) \mathbb{F} = -\mu_0 \left(\mathbf{h} + i \sum \Omega_{m,\ell}^2 \dot{\mathbb{A}}_{m,\ell}(\omega) \mathbb{F} \right). \quad (3.7)$$

Finally, we shall also define the finite subset of \mathbb{C} :

$$\mathcal{S}_\mathcal{T} := \mathcal{P} \cup \mathcal{Z}_m \cup \{0\}. \quad (3.8)$$

Proposition 3.2. *Let $\mathbf{k} \in \mathbb{R}^3$. For any $\omega \in \mathbb{C} \setminus (S(|\mathbf{k}|) \cup \mathcal{S}_{\mathcal{T}})$, the resolvent is given by*

$$R_{|\mathbf{k}|}(\omega) = \mathcal{V}_{|\mathbf{k}|}(\omega) \mathcal{S}_{|\mathbf{k}|}(\omega) + \mathcal{T}(\omega) \quad (3.9)$$

with $\mathcal{S}_{|\mathbf{k}|}(\omega) \in \mathcal{L}(\mathbf{C}_{\perp}^N, \mathbf{C}_{\perp})$ defined by: given $\mathbb{F} := (\mathbf{e}, \mathbf{h}, \mathbf{p}_j, \dot{\mathbf{p}}_j, \mathbf{m}_{\ell}, \dot{\mathbf{m}}_{\ell}) \in \mathbf{C}_{\perp}^N$

$$\mathcal{S}_{|\mathbf{k}|}(\omega) \mathbb{F} := \frac{\omega \mu(\omega) \mathbb{A}_e(\omega) \mathbb{F} - |\mathbf{k}| \mathbf{e}_3 \times \mathbb{A}_m(\omega) \mathbb{F}}{\mathcal{D}(\omega) - |\mathbf{k}|^2}, \quad (3.10)$$

$\mathcal{V}_{|\mathbf{k}|}(\omega) \in \mathcal{L}(\mathbf{C}_{\perp}, \mathbf{C}_{\perp}^N)$ defined by: given $\mathbb{X} \in \mathbf{C}^{\perp}$

$$\left| \begin{aligned} \mathcal{V}_{|\mathbf{k}|}(\omega) \mathbb{X} &:= \left(\mathbb{X}, 0, -\left(\frac{\mathbb{X}}{q_{e,j}(\omega)}\right), \left(\frac{i \omega \mathbb{X}}{q_{e,j}(\omega)}\right), 0, 0 \right) \\ &+ \frac{|\mathbf{k}|}{\omega \mu(\omega)} \left(0, \mathbf{e}_3 \times \mathbb{X}, 0, 0, -\left(\frac{\mathbf{e}_3 \times \mathbb{X}}{q_{m,\ell}(\omega)}\right), \left(\frac{i \omega \mathbf{e}_3 \times \mathbb{X}}{q_{m,\ell}(\omega)}\right) \right) \end{aligned} \right. \quad (3.11)$$

and finally $\mathcal{T}(\omega) \in \mathcal{L}(\mathbf{C}_{\perp}^N)$ defined by : given $\mathbb{F} := (\mathbf{e}, \mathbf{h}, \mathbf{p}_j, \dot{\mathbf{p}}_j, \mathbf{m}_{\ell}, \dot{\mathbf{m}}_{\ell}) \in \mathbf{C}_{\perp}^N$

$$\left| \begin{aligned} \mathcal{T}(\omega) \mathbb{F} &:= \left(0, \frac{\mathbb{A}_m(\omega) \mathbb{F}}{\omega \mu(\omega)}, 0, 0, -\left(\frac{\mathbb{A}_m(\omega) \mathbb{F}}{\omega \mu(\omega) q_{m,\ell}(\omega)}\right), \left(\frac{i \mathbb{A}_m(\omega) \mathbb{F}}{\mu(\omega) q_{m,\ell}(\omega)}\right) \right) \\ &+ \left(0, 0, (\mathbb{A}_{e,j}(\omega) \mathbb{F}), (\dot{\mathbb{A}}_{e,j}(\omega) \mathbb{F}), (\mathbb{A}_{m,\ell}(\omega) \mathbb{F}), (\dot{\mathbb{A}}_{m,\ell}(\omega) \mathbb{F}) \right). \end{aligned} \right. \quad (3.12)$$

The (purely computational) proofs of Propositions 3.1 and 3.2 are delayed to the Appendix.

Remark 3.3. *The function $\omega \mapsto R_{|\mathbf{k}|}(\omega)$ is well-defined and analytic for $\omega \in \mathbb{C} \setminus \sigma(\mathbb{A}_{|\mathbf{k}|,\perp})$. The set $\mathcal{S}_{\mathcal{T}}$ is the set of singularities for $\omega \mapsto \mathcal{T}(\omega)$ but they are removable singularities in the expression (3.9).*

3.2 The dispersion relation

3.2.1 General properties of the dispersion relation

By Proposition 3.1, for a fixed wave number $\mathbf{k} \neq 0$, the eigenvalues of $\mathbb{A}_{|\mathbf{k}|,\perp}$ are the solutions $\omega \in \mathbb{C} \setminus \mathcal{P}$ of the equation

$$\mathcal{D}(\omega) = |\mathbf{k}|^2, \quad (3.13)$$

(with $\mathcal{D}(\omega)$ defined by (3.3)) referred in physics as the dispersion relation. In this section, we provide several results and useful remarks about this equation.

As a consequence of Remark 2.1, one has $\sigma(\mathbb{A}_{|\mathbf{k}|}) \subset \overline{\mathbb{C}^-}$, thus $\sigma(\mathbb{A}_{|\mathbf{k}|,\perp}) \subset \sigma(\mathbb{A}_{|\mathbf{k}|}) \subset \overline{\mathbb{C}^-}$. We give here an elementary proof that for $\mathbf{k} \neq 0$, the spectrum $\sigma(\mathbb{A}_{|\mathbf{k}|,\perp})$ is included in the lower open complex half plane

$$\sigma(\mathbb{A}_{|\mathbf{k}|,\perp}) \subset \mathbb{C}^-, \quad (3.14)$$

based on the dispersion relation (3.13). Assume by contradiction that there exists $\omega \in \overline{\mathbb{C}^+} \cap \sigma(\mathbb{A}_{|\mathbf{k}|,\perp})$. From Proposition 3.1, ω satisfies (3.13) and therefore $\omega \neq 0$ and $\omega \notin \mathcal{P}$. Taking the real part and the imaginary part of (3.13) leads to

$$\operatorname{Re}(\omega \varepsilon(\omega)) \operatorname{Re}(\omega \mu(\omega)) = |\mathbf{k}|^2 + \operatorname{Im}(\omega \varepsilon(\omega)) \operatorname{Im}(\omega \mu(\omega)) > 0 \quad (3.15)$$

$$\operatorname{Re}(\omega \varepsilon(\omega)) \operatorname{Im}(\omega \mu(\omega)) + \operatorname{Re}(\omega \mu(\omega)) \operatorname{Im}(\omega \varepsilon(\omega)) = 0 \quad (3.16)$$

where the rational Herglotz functions $\omega \mapsto \omega\varepsilon(\omega)$ and $\omega \mapsto \omega\mu(\omega)$ are analytic and have a positive imaginary part on \mathbb{C}^+ . Furthermore, one has for $\omega \in \mathbb{R} \setminus \mathcal{P}$:

$$\operatorname{Im}(\omega\varepsilon(\omega)) = \varepsilon_0|\omega|^2 \sum_{j=1}^{N_e} \frac{\Omega_{e,j}^2 \alpha_{e,j}}{|q_{e,j}(\omega)|^2} \geq 0 \quad \text{and} \quad \operatorname{Im}(\omega\mu(\omega)) = \mu_0|\omega|^2 \sum_{\ell=1}^{N_m} \frac{\Omega_{m,\ell}^2 \alpha_{m,\ell}}{|q_{m,\ell}(\omega)|^2} \geq 0. \quad (3.17)$$

Thus, by the weak dissipation assumption (1.11), at least one of the coefficients $\alpha_{e,j}$ or $\alpha_{m,\ell}$ is positive and it follows from (3.17) that

$$\operatorname{Im}(\omega\varepsilon(\omega)) > 0 \quad \text{or} \quad \operatorname{Im}(\omega\mu(\omega)) > 0 \quad \text{on} \quad \mathbb{R}^* \setminus \mathcal{P}. \quad (3.18)$$

Hence, (3.15) implies that $\operatorname{Re}(\omega\varepsilon(\omega))$ and $\operatorname{Re}(\omega\mu(\omega))$ have the same sign (and do not vanish), whereas (3.16) implies that they have opposite sign, which leads to a contradiction. Thus, for $\mathbf{k} \neq 0$, the solutions of the dispersion relation $\mathcal{D}(\omega) = |\mathbf{k}|^2$ all lie in \mathbb{C}^- .

Using the fact that $\varepsilon(-\bar{\omega}) = \overline{\varepsilon(\omega)}$ and $\mu(-\bar{\omega}) = \overline{\mu(\omega)}$, one observes that if ω is a solution of the dispersion relation then $-\bar{\omega}$ is also a solution of this equation. Thus for a fixed $\mathbf{k} \neq 0$, the set $\sigma(\mathbb{A}_{|\mathbf{k}|,\perp})$ is invariant by the transformation $\omega \rightarrow -\bar{\omega}$.

From the irreducible form (3.3) of \mathcal{D} , one deduces that the dispersion relation is equivalent to a polynomial equation of degree N . Namely, for a fixed $\mathbf{k} \in \mathbb{R}^3 \setminus \{0\}$, one has:

$$\mathcal{D}(\omega) = |\mathbf{k}|^2 \iff D_{|\mathbf{k}|}(\omega) = 0, \quad D_{|\mathbf{k}|}(\omega) := \varepsilon_0 \mu_0 \omega^2 P_e(\omega) P_m(\omega) - |\mathbf{k}|^2 Q_e(\omega) Q_m(\omega) \quad (3.19)$$

is a polynomial of degree $N = 2 + 2N_e + 2N_m$ with dominant coefficient $\varepsilon_0 \mu_0$. Thus, one has

$$D_{|\mathbf{k}|}(\omega) \sim \varepsilon_0 \mu_0 \omega^{2+2N_e+2N_m}, \quad (|\omega| \rightarrow +\infty). \quad (3.20)$$

This leads to the following corollary on the diagonalizability of $\mathbb{A}_{|\mathbf{k}|,\perp}$.

Corollary 3.4. *For $\mathbf{k} \in \mathbb{R}^3 \setminus \{0\}$, $\mathbb{A}_{|\mathbf{k}|,\perp}$ is diagonalizable on \mathbf{C}_\perp^N if and only if the roots of the polynomials $D_{|\mathbf{k}|}$ defined in (3.19) are simple.*

Proof. By Proposition 3.1 and relation (3.19), the eigenvalues of $\mathbb{A}_{|\mathbf{k}|,\perp}$ are the solutions the polynomial equation $D_{|\mathbf{k}|}(\omega) = 0$ of degree N . Furthermore, each distinct solution of this equation is an eigenvalue of geometric multiplicity 2. As the space \mathbf{C}_\perp^N is of dimension $2N$, see (2.10) the result follows immediately from a simple argument of dimension. \square

We end this section with two paragraphs: one on the poles \mathcal{P} and one on the zeros $\mathcal{Z} \cup \{0\}$ of the rational function \mathcal{D} associated to the dispersion relation. Indeed (as we will see in the Section 3.3 and in the parts 4 and 5) poles and zeros play a key roles for the asymptotics of eigenvalues of $\mathbb{A}_{|\mathbf{k}|}$ respectively for $|\mathbf{k}| \gg 1$ and $|\mathbf{k}| \ll 1$.

3.2.2 Properties of the poles of the rational function \mathcal{D}

As we saw in Section 3.1, the elements of \mathcal{P} are exactly the zeros of the denominator of the rational function \mathcal{D} , that is the polynomial of $Q_e Q_m$ of degree $2(N_e + N_m)$. Thus, \mathcal{P} contains $2(N_e + N_m)$ elements (counted with multiplicity) and this set is invariant by the transformation $\omega \mapsto -\bar{\omega}$ (see Remark 1.2).

The multiplicity \mathbf{m}_p of a pole $p \in \mathcal{P}$ can not take any arbitrary value. Indeed, the elements of \mathcal{P} are precisely (see Section 3.1) the roots of the 2-nd order polynomials $q_{e,j}$ or $q_{m,\ell}$ for $j \in \{1, \dots, N_e\}$ and $\ell \in \{1, \dots, N_m\}$. Thus, $\mathcal{P} \subset \overline{\mathbb{C}^-}$ (see Remark 1.2) and one has:

- If $p \notin i\mathbb{R}^{-,*}$ then $\mathbf{m}_p = 1$ if $p \notin \mathcal{P}_e \cap \mathcal{P}_m$ or $\mathbf{m}_p = 2$ if $\omega \in \mathcal{P}_e \cap \mathcal{P}_m$,
- Else if $p \in i\mathbb{R}^{-,*}$, then by (H₁), $\mathbf{m}_p \in \{1, 2\}$ if $p \notin \mathcal{P}_e \cap \mathcal{P}_m$ or $\mathbf{m}_p \in \{2, 3, 4\}$ if $p \in \mathcal{P}_e \cap \mathcal{P}_m$.

We will see in Sections 3.3 and 4.3 that the important poles for our analysis for $|\mathbf{k}| \gg 1$ are the ones that lie on the real axis. These poles are associated to polynomials $q_{e,j}$ or $q_{m,\ell}$ for which $\alpha_{e,j} = 0$ or $\alpha_{m,\ell} = 0$. Thus, they are of the form $p = \pm \omega_{e,\ell} \in \mathcal{P}_e$ or $p = \pm \omega_{m,\ell} \in \mathcal{P}_m$ with $\omega_{e,j}, \omega_{m,\ell} > 0$ and have multiplicity $\mathbf{m}_p = 1$ if $p \notin \mathcal{P}_e \cap \mathcal{P}_m$ or $\mathbf{m}_p = 2$ if $p \in \mathcal{P}_e \cap \mathcal{P}_m$.

3.2.3 Properties of the zeros of the rational function \mathcal{D}

As we saw in Section 3.1, the elements of $\mathcal{Z} \cup \{0\}$ are the zeros of the rational function \mathcal{D} , that is the zeros of the polynomial $\omega^2 P_e P_m$ of degree $N = 2(N_e + N_m) + 2$. Hence $\mathcal{Z} \cup \{0\}$ contains N elements (counted with multiplicity). \mathcal{D} is defined as the product of the two (non-constant) rational Herglotz functions $\omega \mapsto \omega \varepsilon(\omega)$ and $\omega \mapsto \omega \mu(\omega)$ which satisfy (using (1.6)) $\text{Im}(\omega \varepsilon(\omega)) > 0$ and $\text{Im}(\omega \mu(\omega)) > 0$ on \mathbb{C}^+ . Thus, these functions could not vanish in the upper-half plane \mathbb{C}^+ (indeed this property holds more generally for any non-constant Herglotz function as a consequence of the open mapping theorem for analytic functions, see for e.g. [3]). Therefore, $\mathcal{Z} \cup \{0\}$ is included in $\overline{\mathbb{C}^-}$. Furthermore, as $\varepsilon(-\bar{\omega}) = \overline{\varepsilon(\omega)}$ and $\mu(-\bar{\omega}) = \overline{\mu(\omega)}$, this set is also invariant by the transformation $\omega \mapsto -\bar{\omega}$.

Concerning the multiplicity \mathbf{m}_z of a zero $z \in \mathcal{Z} \cup \{0\}$, one observes that:

- 0 has a multiplicity $\mathbf{m}_0 = 2$ since (by (1.3) and (1.6)) in the vicinity of 0:

$$\mathcal{D}(\omega) \sim \omega^2 \varepsilon(0) \mu(0), \quad \varepsilon(0) = \varepsilon_0 \left(1 + \sum_{j=1}^{N_e} \frac{\Omega_{e,j}^2}{\omega_{e,j}^2} \right) > 0, \quad \mu(0) = \mu_0 \left(1 + \sum_{\ell=1}^{N_m} \frac{\Omega_{m,\ell}^2}{\omega_{m,\ell}^2} \right) > 0. \quad (3.21)$$

- If $z \notin i\mathbb{R}^-$, then $-\bar{z}$ is a distinct zero with the same multiplicity \mathbf{m}_z .
- From the property (3.18) of $\varepsilon(\omega)$ and $\mu(\omega)$, one immediately sees that
 1. If there exist two indices j_0 and ℓ_0 such that $\alpha_{e,j_0} > 0$ and $\alpha_{m,\ell_0} > 0$ then \mathcal{D} has no non-zero real zeros, i.e. $\mathbb{R}^* \cap \mathcal{Z} = \emptyset$.
 2. If not, either all $\alpha_{e,j}$ vanish in which case $\mathbb{R}^* \cap \mathcal{Z} = \mathcal{Z}_e$, either all $\alpha_{m,\ell}$ vanish in which case $\mathbb{R}^* \cap \mathcal{Z} = \mathcal{Z}_m$. Moreover, these zeros $z \in \mathcal{Z}_e$ (resp. \mathcal{Z}_m) have multiplicity $\mathbf{m}_z = 1$ (this is easily deduced from the graph of the function $\omega \in \mathbb{R} \setminus \mathcal{P}_e \mapsto \varepsilon(\omega)$ or $\omega \in \mathbb{R} \setminus \mathcal{P}_m \mapsto \mu(\omega)$ given by (1.6)).

We point out that the second scenario occurs in particular in the critical configurations described by Definition 1.8.

We will see in Sections 3.3 and 5 that only the real zeros $\mathbb{R} \cap (\mathcal{Z} \cup \{0\})$ can contribute at the main order to the asymptotic of the Fourier-components of $|\mathbf{U}(k, t)|$ for $|\mathbf{k}| \ll 1$.

3.3 Main lines of the analysis

The strategy for proving Theorem 1.10 is quite clear and simple:

- (i) For each $\mathbf{k} \in \mathbb{R}^3 \setminus \{0\}$, one estimates individually $\mathbb{U}(\mathbf{k}, t)$ using formula (2.7), or more precisely (2.12), via the estimation of the exponential $e^{-i\mathbb{A}_{|\mathbf{k}|, \perp} t}$.
- (ii) One gathers the above estimates to estimate the L^2 -norms of the various fields $\mathbf{E}, \mathbf{H}, \dots$ via the norm $\|\mathbf{U}(\cdot, t)\|_{\mathcal{H}}$ thanks to Plancherel's theorem.

For the first step, the key property is the fact that, as emphasized in Section 3.2.1 (see (3.14)), the spectrum of $\mathbb{A}_{|\mathbf{k}|, \perp}$ is included in the complex half-plane \mathbb{C}^- . As a consequence, each $|\mathbb{U}(\mathbf{k}, t)|$ will decay exponentially to 0 for large t .

The reason why the exponential decay is lost for $\|\mathbf{U}(\cdot, t)\|_{\mathcal{H}}$ and degenerates into a polynomial decay, is linked to the fact that the rate of decay of $|\mathbb{U}(\mathbf{k}, t)|$ depends on \mathbf{k} and degenerates when $|\mathbf{k}|$ tends to 0 or $+\infty$. This decay rate is of course linked to the distance of $\sigma(\mathbb{A}_{|\mathbf{k}|, \perp})$ to the real axis that can be deduced from the analysis of the dispersion relation $\mathcal{D}(\omega) = |\mathbf{k}|^2$, where $\mathcal{D}(\omega) = \omega^2 \varepsilon(\omega) \mu(\omega)$, that characterizes the eigenvalues of $\mathbb{A}_{|\mathbf{k}|, \perp}$ (see Proposition 3.1).

- When $|\mathbf{k}|$ is bounded from below and above, this distance is uniformly bounded from below by a positive number which results into a uniform exponential decay of the corresponding $|\mathbb{U}(\mathbf{k}, t)|$'s.
- When $|\mathbf{k}|$ tends to 0 or $+\infty$, this distance tends to 0 and obtaining sharp estimates of $|\mathbb{U}(\mathbf{k}, t)|$ requires to analyse the asymptotic behaviour of the imaginary parts $\text{Im } \omega(|\mathbf{k}|)$ of those eigenvalues $\omega(|\mathbf{k}|) \in \sigma(\mathbb{A}_{|\mathbf{k}|, \perp})$ whose distance to the real axis tends to 0.

From this observation, it is natural to split the analysis into three steps depending on the values of the “spatial frequency” $|\mathbf{k}|$:

- (a) For “mid-range frequencies”, typically $k_- \leq |\mathbf{k}| \leq k_+$ (with $k_- > 0$), one will essentially abandon the spectral approach to the profit of standard techniques for ODE's combined with compactness arguments (the region $k_- \leq |\mathbf{k}| \leq k_+$ is compact). This will be detailed in Section 6.
- (b) For “high frequencies”, $|\mathbf{k}| \geq k_+$, some eigenvalues can be arbitrarily close to the real axis when $|\mathbf{k}| \rightarrow +\infty$: considering the limit of the equation $\mathcal{D}(\omega) = |\mathbf{k}|^2$ when $|\mathbf{k}| \rightarrow +\infty$ it is natural to look at where, for real ω , the function $\mathcal{D}(\omega)$ tends to $+\infty$. This only occurs in one of the following situations when

- $\omega \rightarrow \pm\infty$ in which case $\mathcal{D}(\omega) \sim \omega^2/c^2$, so that one expects the existence of two branches of eigenvalues $\omega_\infty(|\mathbf{k}|)$ and $-\bar{\omega}_\infty(|\mathbf{k}|)$ where

$$\omega_\infty(|\mathbf{k}|) = c|\mathbf{k}| + o(|\mathbf{k}|), \quad |\mathbf{k}| \rightarrow +\infty,$$

where $c := \sqrt{\varepsilon_0 \mu_0}^{-1}$ is the speed of light in the vacuum.

- $\omega \rightarrow \pm \omega_{\nu,q}$ with $\alpha_{\nu,q} = 0$, $\nu = e$ or m where $\omega_{\nu,q}$ is a real pole of \mathcal{D} of multiplicity 1 if $\omega_{\nu,q} \notin \mathcal{P}_e \cap \mathcal{P}_m$ or 2 if $\omega_{\nu,q} \in \mathcal{P}_e \cap \mathcal{P}_m$ (see Section 3.2.2).

If $\omega_{\nu,q} \notin \mathcal{P}_e \cap \mathcal{P}_m$, one expects the existence of two branches of eigenvalues $\omega_{\nu,q}(|\mathbf{k}|)$ and $-\bar{\omega}_{\nu,q}(|\mathbf{k}|)$ such that

$$\omega_{\nu,q}(|\mathbf{k}|) = \omega_{\nu,q} + o(1), \quad |\mathbf{k}| \rightarrow +\infty.$$

Oppositely, if $\omega_{e,q_1} = \omega_{m,q_2} \in \mathcal{P}_e \cap \mathcal{P}_m$, one expects the existence of four branches of eigenvalues $\omega_{e,q_1}(|\mathbf{k}|)$, $-\bar{\omega}_{e,q_1}(|\mathbf{k}|)$, $\omega_{m,q_2}(|\mathbf{k}|)$, $-\bar{\omega}_{m,q_2}(|\mathbf{k}|)$ such that $\omega_{e,q_1}(|\mathbf{k}|)$ and $\omega_{m,q_2}(|\mathbf{k}|)$ tends to ω_{e,q_1} as $|\mathbf{k}| \rightarrow +\infty$.

In the above cases, a more precise analysis of the asymptotic behaviour $\text{Im } \omega_\infty(|\mathbf{k}|)$ and $\text{Im } \omega_{\nu,q}(|\mathbf{k}|)$ when $|\mathbf{k}| \rightarrow +\infty$ plays a crucial role.

- (c) For “low frequencies”, $0 < |\mathbf{k}| \leq k_-$, again some eigenvalues can be arbitrarily close to the real axis when $|\mathbf{k}| \rightarrow 0$: considering the limit of the equation $\mathcal{D}(\omega) = |\mathbf{k}|^2$ when $|\mathbf{k}| \rightarrow 0$ it is natural to look at where, for real ω , the function $\mathcal{D}(\omega)$ vanishes. This only occurs in one of the following situations when

- $\omega \rightarrow 0$ in which case by (3.21): $\mathcal{D}(\omega) \sim \omega^2/c_0^2$ with $c_0 := (\varepsilon(0)\mu(0))^{-1/2}$. Thus, one expects the existence of two branches of eigenvalues $\omega_0(|\mathbf{k}|)$ and $-\bar{\omega}_0(|\mathbf{k}|)$ such that

$$\omega_0(|\mathbf{k}|) = c_0 |\mathbf{k}| + o(|\mathbf{k}|), \quad |\mathbf{k}| \rightarrow 0.$$

- If we are in the second scenario described in the Section 3.2.3, non-zero real zeros exist and are of multiplicity $\mathbf{m}_z = 1$. Thus, in this case, one has $\omega \rightarrow \pm z_\nu$, for $\nu = e$ or m , where z_ν is a zero of $\mathcal{D}(\omega)$. One then expects the existence of two branches of eigenvalues $\omega_{z_\nu}(|\mathbf{k}|)$ and $-\bar{\omega}_{z_\nu}(|\mathbf{k}|)$ such that

$$\omega_{z_\nu}(|\mathbf{k}|) = z_\nu + o(1), \quad |\mathbf{k}| \rightarrow 0.$$

In the above cases, a more precise analysis of the asymptotic behaviour of $\text{Im } \omega_0(|\mathbf{k}|)$ and $\text{Im } \omega_{z_\nu}(|\mathbf{k}|)$ when $|\mathbf{k}| \rightarrow 0$ plays a crucial role.

The analysis of point (b) will explain why the polynomial stability is limited by the Sobolev regularity of the initial data with the first term in the right hand sides of the estimates (1.37) and (1.38), while the analysis of point (c) will explain why this polynomial stability in the second term in the right hand sides of (1.37) and (1.38) is related to the low frequency behavior of the Fourier components of the solution (and thus involved naturally the spaces $\mathcal{L}_p(\mathbb{R}^3)^N$.)

Let us mention that, since the matrices $\mathbb{A}_{|\mathbf{k}|,\perp}$ are not normal, the estimate of $e^{-i\mathbb{A}_{|\mathbf{k}|,\perp} t}$ in the regions $|\mathbf{k}| \geq k_+$ and $|\mathbf{k}| \leq k_-$ cannot be reduced to the study of their eigenvalues. That is why we shall complete the analysis by proving that

- for k_- and k_+ well chosen, the matrices $\mathbb{A}_{|\mathbf{k}|,\perp}$ are diagonalizable in the above regions,
- the associated spectral projectors can be bounded uniformly in $|\mathbf{k}|$ in each region.

4 Asymptotic analysis for large spatial frequencies $|\mathbf{k}| \gg 1$

As this section is the longest of the article, it seems useful to describe its structure. It is made of three main subsections:

- Section 4.1: Asymptotics of dispersion curves for $|\mathbf{k}| \gg 1$.
- Section 4.2: Spectral decomposition of the solution for $|\mathbf{k}| \gg 1$.
- Section 4.3: Large time estimate of $\mathbb{U}(\mathbf{k}, t)$ for $|\mathbf{k}| \gg 1$.

Section 4.3 puts together the results of Sections 4.1 and 4.2. It is decomposed into subsections dedicated to the estimate each of the components of the solution issued from Section 4.2.

From the technical point of view, the proof of the asymptotic expansions will be based on a lemma proved in the appendix A.2, namely the Lemma A.1, that can be seen as a kind of implicit function theorem for functions in the complex plane. This will also be the case in section 5.3.

4.1 Asymptotics of dispersion curves for $|\mathbf{k}| \gg 1$

In this section, we focus on long time estimates of the high (spatial) frequency components $\mathbb{U}(\mathbf{k}, t)$ (see (2.7)) of the solution. As explained in the above paragraph, the decay of $\mathbb{U}(\mathbf{k}, t)$ is related to the analysis of the solutions of the dispersion relation (3.13) for $|\mathbf{k}| \gg 1$. Roughly speaking, as $|\mathbf{k}|^2 \rightarrow +\infty$ when $|\mathbf{k}| \rightarrow +\infty$, the solutions of (3.13) must satisfy $|\mathcal{D}(\omega)| \rightarrow +\infty$ as $|\mathbf{k}| \rightarrow +\infty$. Thus, we observe two scenarios: either they diverge to ∞ or they converge to a pole $p \in \mathcal{P}_e \cup \mathcal{P}_m$ around which, if $p \in \mathcal{P}$ is a pole of multiplicity \mathfrak{m}_p , \mathcal{D} can be rewritten as

$$\mathcal{D}(\omega) = (\omega - p)^{-\mathfrak{m}_p} f(\omega) \text{ with } f \text{ analytic in a vicinity of } p \text{ and } f(p) = A_p \neq 0. \quad (4.1)$$

This leads to the following proposition.

Proposition 4.1. *There exists $k_+ > 0$ such that for $|\mathbf{k}| \geq k_+$, the solutions of the rational dispersion relation (3.13) (or of its equivalent polynomial form (3.19)) are all simple. These solutions form N distinct branches which are C^∞ -smooth functions (with respect to $|\mathbf{k}|$) characterized by their asymptotic expansion for large $|\mathbf{k}|$. More precisely:*

- For any $p \in \mathcal{P}$ with multiplicity \mathfrak{m}_p , for large enough $|\mathbf{k}|$, there exists \mathfrak{m}_p distinct branches of solutions $\omega_{p,n}$, $n = 1, \dots, \mathfrak{m}_p$ of (3.13) satisfying

$$\omega_{p,n}(|\mathbf{k}|) = p + a_{p,n} |\mathbf{k}|^{-\frac{2}{\mathfrak{m}_p}} (1 + o(1)), \quad a_{p,n} = |A_p|^{1/\mathfrak{m}_p} e^{i\frac{\theta_p}{\mathfrak{m}_p}} e^{\frac{2in\pi}{\mathfrak{m}_p}} \quad (|\mathbf{k}| \rightarrow +\infty) \quad (4.2)$$

where A_p is defined in (4.1) and $\theta_p \in (-\pi, \pi]$ is the principal argument of A_p .

- There are 2 distinct branches of solutions $\omega_{\pm\infty}$ of (3.13) that diverge to ∞ as

$$\omega_{\pm\infty}(|\mathbf{k}|) = \pm c |\mathbf{k}| (1 + o(1)) \quad \text{with } c = (\varepsilon_0 \mu_0)^{-\frac{1}{2}} \quad (|\mathbf{k}| \rightarrow +\infty). \quad (4.3)$$

Moreover, there are no other solutions of (3.13) (for $|\mathbf{k}| \geq k_+$) than those described above.

Proof. Step 1: construction of the branches of solutions $\omega_{p,n}(|\mathbf{k}|)$.

Let $p \in \mathcal{P}$ a pole of multiplicity \mathbf{m}_p . Then, the rational function \mathcal{D} can be factorized as in (4.1). Then p is a zero of multiplicity \mathbf{m}_p of the function $\omega \mapsto \mathcal{D}(\omega)^{-1} = (\omega - p)^{\mathbf{m}_p} f(\omega)^{-1}$ which is analytic in an open neighbourhood of p and satisfies $f(p)^{-1} = A_p^{-1} \neq 0$. Moreover, for $|\mathbf{k}| \neq 0$, solving the dispersion relation:

$$\mathcal{D}(\omega) = |\mathbf{k}|^2 \quad \text{is equivalent to solving the equation} \quad \mathcal{D}(\omega)^{-1} = |\mathbf{k}|^{-2} = (|\mathbf{k}|^{-\frac{2}{\mathbf{m}_p}})^{\mathbf{m}_p}.$$

Then, by applying the Lemma A.1 of the appendix A.2 with

$$\mathcal{G}(\omega) = \mathcal{D}(\omega)^{-1}, \quad z = p, \quad g(\omega) = f(\omega)^{-1}, \quad \mathbf{m} = \mathbf{m}_p, \quad A = A_p^{-1} \text{ and } \zeta = |\mathbf{k}|^{-2/\mathbf{m}_p},$$

we deduce for $|\mathbf{k}|$ large enough the existence of \mathbf{m}_p distinct branches of solutions: $|\mathbf{k}| \mapsto \omega_{p,n}(|\mathbf{k}|)$ of the equation $\mathcal{D}(\omega) = |\mathbf{k}|^2$ which are C^∞ functions with respect to $|\mathbf{k}|$ and satisfy (4.2).

Step 2: construction of the 2 distinct branches of solutions $\omega_{\pm\infty}(|\mathbf{k}|)$.

To reduce ourselves to the application of Lemma A.1 as in step 1, the trick consists in saying that $\mathcal{D}(\omega) = \omega^2 \varepsilon(\omega) \mu(\omega) = |\mathbf{k}|^2$ is equivalent to $\mathcal{D}(\omega)^{-1} = \omega^{-2} \varepsilon(\omega)^{-1} \mu(\omega)^{-1} = |\mathbf{k}|^{-2}$. Then we introduce the new unknown $\xi = 1/\omega$, so that $|\omega| \rightarrow +\infty \Leftrightarrow \xi \rightarrow 0$ and

$$\mathcal{D}(\omega) = \omega^2 \varepsilon(\omega) \mu(\omega) = |\mathbf{k}|^2 \quad \Longleftrightarrow \quad \xi^2 \varepsilon(1/\xi)^{-1} \mu(1/\xi)^{-1} = |\mathbf{k}|^{-2}, \text{ modulo } \xi = 1/\omega. \quad (4.4)$$

Then we introduce the rational function $\mathcal{G}(\xi)$:

$$\mathcal{G}(\xi) := \xi^2 g(\xi), \quad g(\xi) := \varepsilon(1/\xi)^{-1} \mu(1/\xi)^{-1} \quad \text{well defined for } \xi \in \mathbb{C} \setminus \{1/z, z \in \mathcal{Z}\}$$

after having remarked that $g(\xi)$ could be extended analytically at $\xi = 0$ via

$$g(0) = \lim_{\xi \rightarrow 0} \varepsilon(\xi^{-1})^{-1} \mu(\xi^{-1})^{-1} = (\varepsilon_0 \mu_0)^{-1} = c^2 \neq 0.$$

Thus, we can now apply the Lemma A.1, replacing ω par ξ , with

$$z = 0, \quad g(\xi) := \varepsilon(1/\xi)^{-1} \mu(1/\xi)^{-1}, \quad \mathbf{m} = 2, \quad A = g(0) = c^2 \text{ and } \zeta = |\mathbf{k}|^{-1}.$$

From Lemma A.1, we deduce the existence of two distinct analytic functions $\zeta \mapsto \xi_{\pm}(\zeta)$, defined in the vicinity of 0, such that $\mathcal{G}(\xi_{\pm}(\zeta)) = \zeta^2$ and such that

$$\xi_{\pm}(\zeta) = \pm c^{-1} \zeta (1 + o(1)), \quad \text{as } \zeta \rightarrow 0. \quad (4.5)$$

Thus, setting $\omega_{\pm\infty}(|\mathbf{k}|) = \xi_{\pm}(|\mathbf{k}|^{-1})^{-1}$, we construct two branches of solutions of $\mathcal{D}(\omega) = |\mathbf{k}|^2$ which admit by (4.5) the asymptotic expansion (4.3).

Step 3: conclusion. In step 1, since the sum of the \mathbf{m}_p 's over $p \in \mathcal{P}$ is equal to $2N_e + 2N_m$, we have constructed for $|\mathbf{k}|$ large enough, $2N_e + 2N_m$ solutions, namely $\{\omega_{p,n}(|\mathbf{k}|), p \in \mathcal{P}, n \leq \mathbf{m}_p\}$ which are all distinct due to the asymptotics (4.2).

In step 2, we have constructed for $|\mathbf{k}|$ large enough, 2 additional solutions $\omega_{\pm\infty}(|\mathbf{k}|)$ that are distinct thanks to (4.3).

From both asymptotics (4.2) and (4.3), none of this two solutions can coincide with any of those from step 1. Therefore, with $\{\omega_{p,n}(|\mathbf{k}|), p \in \mathcal{P}, n \leq \mathbf{m}_p\} \cup \{\omega_{\pm\infty}(|\mathbf{k}|)\}$, we have constructed $2N_e + 2N_m + 2$ distinct solutions of (3.13). Since (3.13) is equivalent to a polynomial equation of degree $2N_e + 2N_m + 2$, cf. (3.19), there are no other solutions. \square

4.2 Spectral decomposition of the solution for $|\mathbf{k}| \gg 1$

In the physics literature, the solutions of the dispersion relation (3.13) $|\mathbf{k}| \rightarrow \omega_{p,n}(|\mathbf{k}|)$ and $|\mathbf{k}| \rightarrow \omega_{\pm\infty}(|\mathbf{k}|)$ (given here for large $|\mathbf{k}|$ by Proposition 4.1) are referred to as the dispersion curves. The asymptotics (4.2) and (4.3) of these curves show that they do not cross each other for $|\mathbf{k}| \gg 1$. Thus, combining Proposition 4.1 and Corollary 3.4 immediately yields the following property of the operator $\mathbb{A}_{|\mathbf{k}|,\perp}$.

Corollary 4.2. *There exists $k_+ > 0$ such that for $|\mathbf{k}| \geq k_+$, $\mathbb{A}_{|\mathbf{k}|,\perp}$ is diagonalizable on \mathbf{C}_{\perp}^N .*

To express the diagonal decomposition of the solution $\mathbb{U}(\mathbf{k}, t)$, we split (see Section 3.2.2) the sets of poles $\mathcal{P} \subset \overline{\mathbb{C}^-}$ in three disjoint subsets : $\mathcal{P} = \mathcal{P}_- \cup \mathcal{P}_s \cup \mathcal{P}_d$ with

$$\mathcal{P}_- := \mathcal{P} \cap \mathbb{C}^-, \quad \mathcal{P}_s := \{p \in \mathcal{P} \cap \mathbb{R} \mid \mathbf{m}_p = 1\}, \quad \text{and} \quad \mathcal{P}_d := \{p \in \mathcal{P} \cap \mathbb{R} \mid \mathbf{m}_p = 2\}.$$

We point out that the weak dissipation condition (1.11) implies that at least one pole of \mathcal{P} lies in \mathbb{C}^- , thus $\mathcal{P}_- \neq \emptyset$ whereas the strong dissipation condition (1.9) implies that all poles lie in \mathbb{C}^- , that is $\mathcal{P}_- = \mathcal{P}$. From Section 3.2.2, one also has that $\mathcal{P}_d = \mathcal{P}_e \cap \mathcal{P}_m \cap \mathbb{R}$.

Using Propositions 3.1 and 4.1, we introduce the following partition of the spectrum of $\mathbb{A}_{|\mathbf{k}|,\perp}$ for $|\mathbf{k}| \geq k_+$:

$$\left| \begin{aligned} \sigma(\mathbb{A}_{|\mathbf{k}|,\perp}) &= \{\omega_{\pm\infty}(|\mathbf{k}|)\} \cup \{\omega_p(|\mathbf{k}|), p \in \mathcal{P}_s\} \cup \{\omega_{p,r}(|\mathbf{k}|), p \in \mathcal{P}_d, r \in \{1, 2\}\} \\ &\cup \{\omega_{p,n}(|\mathbf{k}|), p \in \mathcal{P}_-, n = 1, \dots, \mathbf{m}_p\}, \end{aligned} \right.$$

(where we set $\omega_p(|\mathbf{k}|) := \omega_{p,1}(|\mathbf{k}|)$ for a real pole $p \in \mathcal{P}_s$ of multiplicity $\mathbf{m}_p = 1$.) We refer to the figure 1 for an illustration of the behavior of the dispersion curves for $|\mathbf{k}| \gg 1$. For $|\mathbf{k}| \geq k_+$,

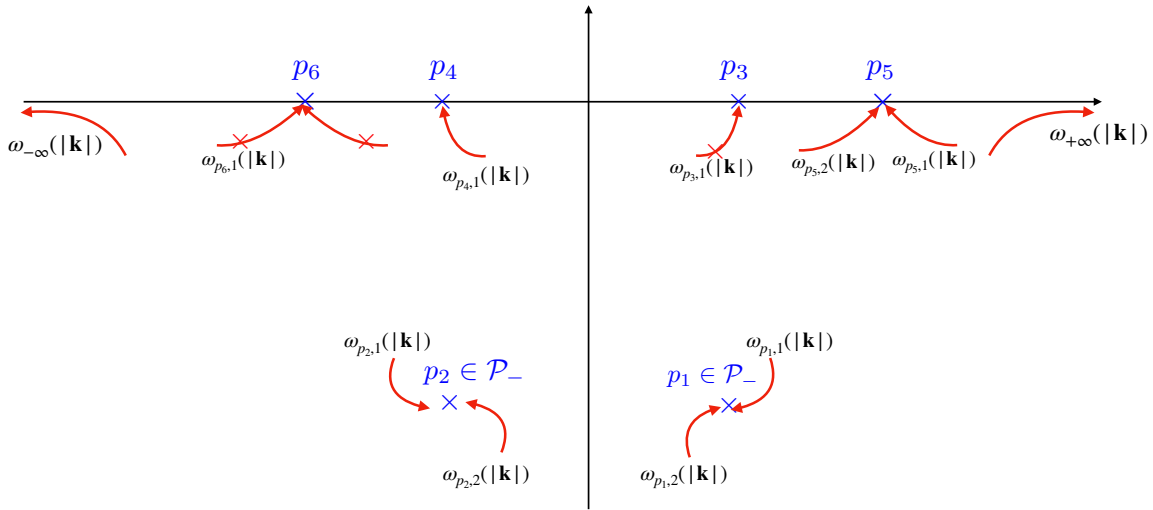


Figure 1: Sketch of a configuration of the dispersion curves for large $|\mathbf{k}|$ large in the case where $\mathcal{P} = \mathcal{P}_- \cup \mathcal{P}_s \cup \mathcal{P}_d$ with $\mathcal{P}_- = \{p_1, p_2\}$, $\mathcal{P}_s = \{p_3, p_4\}$ and $\mathcal{P}_d = \{p_5, p_6\}$.

as $\mathbb{A}_{|\mathbf{k}|,\perp}$ is diagonalizable by Corollary 3.4, \mathbf{C}_\perp^N can be decomposed as

$$\mathbf{C}_\perp^N = \bigoplus_{\pm} V_{|\mathbf{k}|,\pm\infty} \oplus \bigoplus_{p \in \mathcal{P}_s} V_{|\mathbf{k}|,p} \oplus \bigoplus_{p \in \mathcal{P}_d} \bigoplus_{r=1}^2 V_{|\mathbf{k}|,p,r} \oplus \bigoplus_{p \in \mathcal{P}_-} \bigoplus_{n=1}^{m_p} V_{|\mathbf{k}|,p,n} \quad (4.6)$$

$$\begin{cases} V_{|\mathbf{k}|,\pm\infty} = \ker(\mathbb{A}_{|\mathbf{k}|,\perp} - \omega_{\pm\infty}(|\mathbf{k}|) \text{Id}), & V_{|\mathbf{k}|,p,n} = \ker(\mathbb{A}_{|\mathbf{k}|,\perp} - \omega_{p,n}(|\mathbf{k}|) \text{Id}), \quad n \leq m_p, \\ V_{|\mathbf{k}|,p} = \ker(\mathbb{A}_{|\mathbf{k}|,\perp} - \omega_p(|\mathbf{k}|) \text{Id}), & V_{|\mathbf{k}|,p,r} = \ker(\mathbb{A}_{|\mathbf{k}|,\perp} - \omega_{p,r}(|\mathbf{k}|) \text{Id}), \quad r = 1, 2, \end{cases} \quad (4.7)$$

where the above direct sums are (in general) non-orthogonal. Following these direct sums, one decomposes uniquely any vector $x \in \mathbf{C}_\perp^N$ as

$$\left| \begin{array}{l} x = \sum_{\pm\infty} x_{|\mathbf{k}|,\pm\infty} + \sum_{p \in \mathcal{P}_s} x_{|\mathbf{k}|,p} + \sum_{p \in \mathcal{P}_d} \sum_{r=1}^2 x_{|\mathbf{k}|,p,r} + \sum_{p \in \mathcal{P}_-} \sum_{n=1}^{m_p} x_{|\mathbf{k}|,p,n} \\ x_{|\mathbf{k}|,\pm\infty} \in V_{|\mathbf{k}|,\pm\infty}, \quad x_{|\mathbf{k}|,p} \in V_{|\mathbf{k}|,p}, \quad x_{|\mathbf{k}|,p,r} \in V_{|\mathbf{k}|,p,r}, \quad \text{and} \quad x_{|\mathbf{k}|,p,n} \in V_{|\mathbf{k}|,p,n}. \end{array} \right.$$

Then, we define the spectral projectors $\Pi_{p,n}(|\mathbf{k}|)$, $p \in \mathcal{P}_-$ and $n \in \{1, \dots, m_p\}$, $\Pi_{\pm\infty}(|\mathbf{k}|)$, $\Pi_p(|\mathbf{k}|)$ for $p \in \mathcal{P}_s$ and $\Pi_{p,r}(|\mathbf{k}|)$ for $p \in \mathcal{P}_d$ and $r \in \{1, 2\}$ associated respectively to the eigenvalues $\omega_{p,n}(|\mathbf{k}|)$, $\omega_{\pm\infty}(|\mathbf{k}|)$, $\omega_p(|\mathbf{k}|)$ and $\omega_{p,r}(|\mathbf{k}|)$ by:

$$\begin{cases} \Pi_{\pm\infty}(|\mathbf{k}|)(x) = x_{|\mathbf{k}|,\pm\infty}, & \Pi_{p,n}(|\mathbf{k}|)(x) = x_{|\mathbf{k}|,p,n}, \\ \Pi_p(|\mathbf{k}|)(x) = x_{|\mathbf{k}|,p}, & \Pi_{p,r}(|\mathbf{k}|)(x) = x_{|\mathbf{k}|,p,r}. \end{cases} \quad (4.8)$$

From Proposition 3.1, the geometric multiplicity of each eigenvalues in $\sigma(\mathbb{A}_{|\mathbf{k}|,\perp})$ is two. Thus, all the $\Pi_{p,n}(|\mathbf{k}|)$, $\Pi_{\pm\infty}(|\mathbf{k}|)$, $\Pi_p(|\mathbf{k}|)$ and $\Pi_{p,r}(|\mathbf{k}|)$ are rank two projectors. We emphasize that the dissipative operator $\mathbb{A}_{|\mathbf{k}|,\perp}$ is not normal, thus its spectral projectors are not orthogonal. For $|\mathbf{k}| \geq k_+$, as $\mathbb{A}_{|\mathbf{k}|,\perp}$ is diagonalizable (by Corollary 3.4), one has

$$\begin{aligned} \mathbb{A}_{|\mathbf{k}|,\perp} &= \sum_{\pm} \omega_{\pm\infty}(|\mathbf{k}|) \Pi_{\pm\infty}(|\mathbf{k}|) + \sum_{p \in \mathcal{P}_s} \omega_p(|\mathbf{k}|) \Pi_p(|\mathbf{k}|) + \sum_{p \in \mathcal{P}_d} \sum_{r=1}^2 \omega_{p,r}(|\mathbf{k}|) \Pi_{p,r}(|\mathbf{k}|) \\ &+ \sum_{p \in \mathcal{P}_-} \sum_{n=1}^{m_p} \omega_{p,n}(|\mathbf{k}|) \Pi_{p,n}(|\mathbf{k}|). \end{aligned}$$

Thus, for $|\mathbf{k}| > k_+$, the solution $\mathbb{U}(\mathbf{k}, t)$ given by (2.12) can be expressed for all $t \geq 0$ as

$$\mathbb{U}(\mathbf{k}, t) = \mathbb{U}_\infty(\mathbf{k}, t) + \mathbb{U}_s(\mathbf{k}, t) + \mathbb{U}_d(\mathbf{k}, t) + \mathbb{U}_-(\mathbf{k}, t) \quad (4.9)$$

where

$$\left\{ \begin{array}{l} \mathbb{U}_\infty(\mathbf{k}, t) = \sum_{\pm} e^{-i\omega_{\pm\infty}(|\mathbf{k}|)t} \mathcal{R}_{\mathbf{k}}^* \Pi_{\pm\infty}(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}} \mathbb{U}_0(\mathbf{k}), \quad (i) \\ \mathbb{U}_s(\mathbf{k}, t) = \sum_{p \in \mathcal{P}_s} e^{-i\omega_p(|\mathbf{k}|)t} \mathcal{R}_{\mathbf{k}}^* \Pi_p(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}} \mathbb{U}_0(\mathbf{k}), \quad (ii) \\ \mathbb{U}_d(\mathbf{k}, t) = \sum_{p \in \mathcal{P}_d} \sum_{r=1}^2 e^{-i\omega_{p,r}(|\mathbf{k}|)t} \mathcal{R}_{\mathbf{k}}^* \Pi_{p,r}(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}} \mathbb{U}_0(\mathbf{k}), \quad (iii) \\ \mathbb{U}_-(\mathbf{k}, t) = \sum_{p \in \mathcal{P}_-} \sum_{n=1}^{m_p} e^{-i\omega_{p,n}(|\mathbf{k}|)t} \mathcal{R}_{\mathbf{k}}^* \Pi_{p,n}(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}} \mathbb{U}_0(\mathbf{k}). \quad (iv) \end{array} \right. \quad (4.10)$$

4.3 Estimates of $\mathbb{U}(\mathbf{k}, t)$ for $|\mathbf{k}| \gg 1$

In each subsection of this section, we shall provide a lower bound $k_+ > 0$ for giving a sense to $|\mathbf{k}| \gg 1$ via $|\mathbf{k}| \geq k_+$. A priori, the value of k_+ will change from one section to the other but we can always choose a value for k_+ that is larger than its previous values. This convention will be adopted systematically without being explicitly mentioned.

4.3.1 Orientation

In what follows we are going to estimate successively, in Lemmas 4.5, 4.9, 4.12 and 4.13, each of the terms appearing in the decomposition (4.9).

For $\mathbb{U}_\infty(\mathbf{k}, t)$, $\mathbb{U}_s(\mathbf{k}, t)$ and $\mathbb{U}_d(\mathbf{k}, t)$, which will be treated in Sections 4.3.2 to 4.3.4, we shall bound separately each of the terms of the sums in (4.10)(i) to (4.10)(iii). For each eigenvalue $\omega(|\mathbf{k}|) \in \{\omega_{\pm\infty}(|\mathbf{k}|), \omega_p(|\mathbf{k}|), \omega_{p,r}(|\mathbf{k}|)\}$, we shall first estimate in Lemmas 4.3, 4.6 and 4.10 the corresponding spectral projector $\Pi(|\mathbf{k}|)$. Many approaches are possible. It appeared useful to use here the expression of $\Pi(|\mathbf{k}|)$ provided by the Riesz-Dunford functional calculus (see e.g. [9], sections VII.1 and VII.3), in terms of a contour integral in the complex plane whose integrand involves the resolvent $R_{|\mathbf{k}|}(\omega)$ studied in Section 3.1. This contour will be taken as a (positively oriented) circle $\mathcal{C}_{|\mathbf{k}|}$ centered at $\omega(|\mathbf{k}|)$

$$\mathcal{C}_{|\mathbf{k}|} := \{\omega \in \mathbb{C} \mid |\omega - \omega(|\mathbf{k}|)| = \rho_{|\mathbf{k}|}\}, \quad (4.11)$$

whose radius $\rho_{|\mathbf{k}|}$ must be chosen in such a way that

$$\left\{ \begin{array}{l} (i) \mathcal{C}_{|\mathbf{k}|} \text{ does not enclose or intersect any other eigenvalue,} \\ (ii) \mathcal{C}_{|\mathbf{k}|} \text{ does not enclose or intersect any point of the set } \mathcal{S}_{\mathcal{T}}. \end{array} \right. \quad (4.12)$$

This will be automatically achieved if we take $\rho_{|\mathbf{k}|}$ equal to half of the distance of $\omega(|\mathbf{k}|)$ to all the points that one wants to avoid, namely

$$\rho_{|\mathbf{k}|} = 1/2 \text{ dist}(\omega(|\mathbf{k}|), \Omega(|\mathbf{k}|)), \quad \Omega(|\mathbf{k}|) = \left(\sigma(\mathbb{A}_{|\mathbf{k}|, \perp}) \cup \mathcal{S}_{\mathcal{T}} \right) \setminus \{\omega(|\mathbf{k}|)\}. \quad (4.13)$$

We then have the formula

$$\Pi(|\mathbf{k}|) = -\frac{1}{2i\pi} \int_{\mathcal{C}_{|\mathbf{k}|}} R_{|\mathbf{k}|}(\omega) d\omega = -\frac{1}{2i\pi} \int_{\mathcal{C}_{|\mathbf{k}|}} \mathcal{V}_{|\mathbf{k}|}(\omega) \mathcal{S}_{|\mathbf{k}|}(\omega) d\omega, \quad (4.14)$$

where the first equality is justified by (4.12)(i) and the second by the expression of the resolvent: $R_{|\mathbf{k}|}(\omega) = \mathcal{V}_{|\mathbf{k}|}(\omega) \mathcal{S}_{|\mathbf{k}|}(\omega) + \mathcal{T}(\omega)$ (given in Proposition 3.2) and by (4.12)(ii) since as $\mathcal{S}_{\mathcal{T}}$ is defined as the set of singularities of $\omega \mapsto \mathcal{T}(\omega)$, this function is analytic in $\mathbb{C} \setminus \mathcal{S}_{\mathcal{T}}$.

From (4.14), we deduce the inequality that we shall use systematically in the following, namely

$$\|\Pi(|\mathbf{k}|)\| \leq \rho_{|\mathbf{k}|} \sup_{\omega \in \mathcal{C}_{|\mathbf{k}|}} (\|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \|\mathcal{S}_{|\mathbf{k}|}(\omega)\|). \quad (4.15)$$

where $\|\cdot\|$ denotes the operator norm of $\mathcal{L}(\mathbb{C}_{\perp}^N)$.

In a second step, we shall concentrate on the exponentials appearing in each factor whose estimate for large $|\mathbf{k}|$ will rely on the asymptotic expansion of the eigenvalues (and more particularly their imaginary parts), see Lemmas 4.4, 4.11 and 4.7. These results will be transformed into sharp exponential decay estimates for $|\mathbb{U}_{\infty}(\mathbf{k}, t)|$, $|\mathbb{U}_s(\mathbf{k}, t)|$ and $|\mathbb{U}_d(\mathbf{k}, t)|$, in which the rate of decay degenerates for $|\mathbf{k}| \rightarrow +\infty$, see Lemmas 4.5, 4.9 and 4.12.

Finally, the last term of (4.10), $\mathbb{U}_{-}(\mathbf{k}, t)$, will not be treated by using (4.10)(iv) but an alternative expression directly issued from the Riesz-Dunford functional calculus (Section 4.3.5). In fact, the exponential decay of $|\mathbb{U}_{-}(\mathbf{k}, t)|$ will be, contrarily to the previous terms, uniform with respect to $|\mathbf{k}|$, so that it will not contribute at the end to the large time equivalent of $\mathbb{U}(\cdot, t)$. This is the reason why we can be satisfied with rough estimates.

4.3.2 Estimates of $\mathbb{U}_{\pm\infty}(\mathbf{k}, t)$ for $|\mathbf{k}| \gg 1$

In the following Lemma, we estimate the spectral projectors $\Pi_{\pm\infty}(|\mathbf{k}|)$.

Lemma 4.3. *The spectral projectors $\Pi_{\pm\infty}(|\mathbf{k}|)$ are uniformly bounded for large $|\mathbf{k}|$.*

Proof. We follow the approach described in Section 4.3.1 for $\omega(|\mathbf{k}|) = \omega_{\pm\infty}(|\mathbf{k}|)$ and denote $\mathcal{C}_{\pm\infty, |\mathbf{k}|}$ the corresponding contour, see (4.11) and figure 2.

Step 1: Estimate of $\rho_{|\mathbf{k}|}$. Defining $\rho_{|\mathbf{k}|}$ by (4.13) for $\omega(|\mathbf{k}|) = \omega_{\pm\infty}(|\mathbf{k}|)$ and using the asymptotic behaviour (4.3), it is clear that

$$\rho_{|\mathbf{k}|} \sim \frac{1}{2} c |\mathbf{k}|, \quad (|\mathbf{k}| \rightarrow +\infty), \quad \text{thus} \quad \rho_{|\mathbf{k}|} \lesssim |\mathbf{k}| \quad \text{for } |\mathbf{k}| \text{ large enough.} \quad (4.16)$$

Step 2: Estimate of $(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1}$. This term is of interest because it appears in the expression (3.10) of $\mathcal{S}_{|\mathbf{k}|}(\omega)$. Using (3.1), (3.3), we compute that

$$(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1} = \frac{Q_e(\omega) Q_m(\omega)}{D_{|\mathbf{k}|}(\omega)}, \quad \text{with } D_{|\mathbf{k}|}(\omega) \text{ given in (3.19).} \quad (4.17)$$

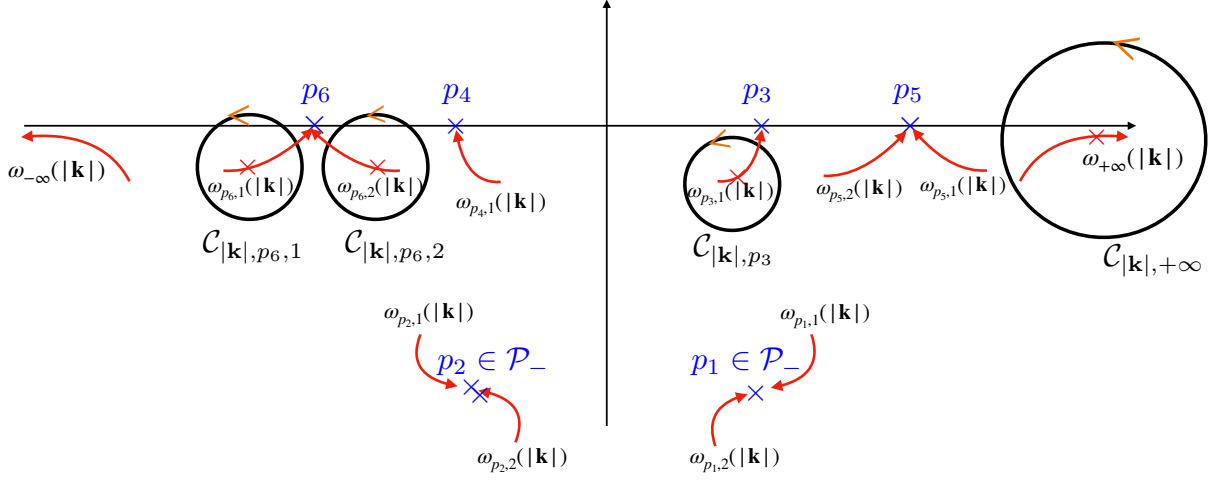


Figure 2: Contours of integration $\mathcal{C}_{|\mathbf{k}|,+\infty}$, $\mathcal{C}_{|\mathbf{k}|,p_3}$, $\mathcal{C}_{|\mathbf{k}|,p_6,1}$ and $\mathcal{C}_{|\mathbf{k}|,p_6,2}$ used for the estimate of $\Pi_{+\infty}(|\mathbf{k}|)$, $\Pi_{p_3}(|\mathbf{k}|)$ for $p_3 \in \mathcal{P}_s$, $\Pi_{p_6,1}(|\mathbf{k}|)$ and $\Pi_{p_6,2}(|\mathbf{k}|)$ for $p_6 \in \mathcal{P}_d$ (corresponding to the figure 1).

According to Proposition 4.1, we know that $D_{|\mathbf{k}|}(\omega)$ is a polynomial of degree N with simple roots so that, according to (3.20), it can be factorized as (note that we distinguish below $\omega_{\pm\infty}(|\mathbf{k}|)$ that go to ∞ with $|\mathbf{k}|$, from the other roots which remain bounded)

$$D_{|\mathbf{k}|}(\omega) = \varepsilon_0 \mu_0 D_{b,|\mathbf{k}|}(\omega) (\omega - \omega_{+\infty}(|\mathbf{k}|)) (\omega - \omega_{-\infty}(|\mathbf{k}|)) \quad (4.18)$$

with $D_{b,|\mathbf{k}|}(\omega)$ given by

$$D_{b,|\mathbf{k}|}(\omega) = \prod_{p \in \mathcal{P}_-} \prod_{n=1}^{m_p} (\omega - \omega_{p,n}(|\mathbf{k}|)) \prod_{p \in \mathcal{P}_s} (\omega - \omega_p(|\mathbf{k}|)) \prod_{p \in \mathcal{P}_d} \prod_{r=1}^2 (\omega - \omega_{p,r}(|\mathbf{k}|)). \quad (4.19)$$

In the same way, thanks to (4.2), for any $\omega(|\mathbf{k}|) \in \{\omega_{p,n}(|\mathbf{k}|), \omega_p(|\mathbf{k}|), \omega_{p,r}(|\mathbf{k}|)\}$,

$$|\omega - \omega(|\mathbf{k}|)| \geq |\omega_{\pm\infty}(|\mathbf{k}|) - \omega(|\mathbf{k}|)| - |\omega - \omega_{\pm\infty}(|\mathbf{k}|)| \sim c|\mathbf{k}|/2, \quad (|\mathbf{k}| \rightarrow +\infty). \quad (4.20)$$

Thus, by definition (4.19) of $D_{b,|\mathbf{k}|}(\omega)$, we deduce from (4.22) and (4.20), that

$$|D_{b,|\mathbf{k}|}(\omega)| \geq C |\mathbf{k}|^{2(N_e+N_m)}, \quad \forall \omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}, \quad \text{for some } C > 0. \quad (4.21)$$

Next, if $\omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}$, by the reverse triangular inequality

$$|\omega - \omega_{\mp\infty}(|\mathbf{k}|)| \geq |\omega_{\pm\infty}(|\mathbf{k}|) - \omega_{\mp\infty}(|\mathbf{k}|)| - |\omega - \omega_{\pm\infty}(|\mathbf{k}|)| \sim 3c|\mathbf{k}|/2, \quad (|\mathbf{k}| \rightarrow +\infty) \quad (4.22)$$

since, by (4.3), $|\omega_{\pm\infty}(|\mathbf{k}|) - \omega_{\mp\infty}(|\mathbf{k}|)| \sim 2c|\mathbf{k}|$ and $|\omega - \omega_{\pm\infty}(|\mathbf{k}|)| = \rho_{|\mathbf{k}|} \sim c|\mathbf{k}|/2$. Thus, using (4.22) and (4.21) in (4.18) yields, as $|\omega - \omega_{\pm\infty}(|\mathbf{k}|)| = \rho_{|\mathbf{k}|}$ on $\mathcal{C}_{\pm\infty,|\mathbf{k}|}$,

$$|D_{|\mathbf{k}|}(\omega)|^{-1} \lesssim \rho_{|\mathbf{k}|}^{-1} |\mathbf{k}|^{-(2(N_e+N_m)+1)}, \quad \forall \omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}. \quad (4.23)$$

For bounding $Q_e(\omega)Q_m(\omega)$, that appears in (4.17), we use an upper bound for $\omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}$

$$|\omega| \leq |\omega_{\pm\infty}(|\mathbf{k}|)| + \rho_{|\mathbf{k}|} \sim 3c|\mathbf{k}|/2, \quad |\mathbf{k}| \rightarrow +\infty, \quad \forall \omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}, \quad (4.24)$$

to deduce that, as $Q_e Q_m$ is a polynomial of degree $2(N_e + N_m)$,

$$|Q_e(\omega)Q_m(\omega)| \lesssim |\omega|^{2(N_e+N_m)} \lesssim |\mathbf{k}|^{2(N_e+N_m)}, \quad \forall \omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}. \quad (4.25)$$

Thus, using (4.25) and (4.23) in (4.17), we get, for $|\mathbf{k}|$ large enough,

$$|(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1}| \lesssim \rho_{|\mathbf{k}|}^{-1} |\mathbf{k}|^{-1}, \quad \forall \omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}. \quad (4.26)$$

Step 3: Estimates of $\mathcal{S}_{|\mathbf{k}|}(\omega)$ and $\mathcal{V}_{|\mathbf{k}|}(\omega)$. Owing to (3.10), to estimate $\mathcal{S}_{|\mathbf{k}|}(\omega)$, it remains to estimate $\omega\mu(\omega)\mathbb{A}_e(\omega) - |\mathbf{k}|\mathbf{e}_3 \times \mathbb{A}_m(\omega)$. We claim that,

$$\text{for } |\mathbf{k}| \text{ large enough, } \quad \|\omega\mu(\omega)\mathbb{A}_e(\omega) - |\mathbf{k}|\mathbf{e}_3 \times \mathbb{A}_m(\omega)\| \lesssim |\mathbf{k}|, \quad \forall \omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}, \quad (4.27)$$

which, combined with (4.26) provides, via (3.10),

$$\|\mathcal{S}_{|\mathbf{k}|}(\omega)\| \lesssim \rho_{|\mathbf{k}|}^{-1}, \quad \forall \omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}. \quad (4.28)$$

To prove (4.27), notice first that along $\mathcal{C}_{\pm\infty,|\mathbf{k}|}$,

$$|\omega| \geq |\omega_{\pm\infty}(|\mathbf{k}|)| - \rho_{|\mathbf{k}|} \sim c|\mathbf{k}|/2 \quad \text{when } |\mathbf{k}| \rightarrow +\infty. \quad (4.29)$$

Thus, as $\mu(\omega) = \mu_0 + o(1)$ as $|\omega| \rightarrow \infty$, with the upper bound (4.24), one gets

$$\|\omega\mu(\omega)\mathbb{A}_e(\omega) - |\mathbf{k}|\mathbf{e}_3 \times \mathbb{A}_m(\omega)\| \lesssim |\mathbf{k}| (\|\mathbb{A}_e(\omega)\| + \|\mathbb{A}_m(\omega)\|), \quad \forall \omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}.$$

Finally, we prove $\mathbb{A}_e(\omega)$ and $\mathbb{A}_m(\omega)$ are uniformly bounded (in \mathbf{k}) on $\mathcal{C}_{\pm\infty,|\mathbf{k}|}$. Towards this goal, we first estimate the operators $(\mathbb{A}_{e,j}(\omega), \dot{\mathbb{A}}_{e,j}(\omega), \mathbb{A}_{m,\ell}(\omega), \dot{\mathbb{A}}_{m,\ell}(\omega))$. From (3.6),

$$\|\mathbb{A}_{e,j}(\omega)\|, \|\dot{\mathbb{A}}_{e,j}(\omega)\| \lesssim \frac{1+|\omega|}{|q_{e,j}(\omega)|}, \quad \|\mathbb{A}_{m,\ell}(\omega)\|, \|\dot{\mathbb{A}}_{m,\ell}(\omega)\| \lesssim \frac{1+|\omega|}{|q_{m,\ell}(\omega)|}, \quad (4.30)$$

Thus, using (4.29), as $q_{e,j}(\omega), q_{m,\ell}(\omega) \sim \omega^2$ when $|\omega| \rightarrow +\infty$ (see (1.7)), one has

$$|q_{e,j}(\omega)|^{-1} \lesssim |\mathbf{k}|^{-2}, \quad |q_{m,\ell}(\omega)|^{-1} \lesssim |\mathbf{k}|^{-2} \quad \text{for } \omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}. \quad (4.31)$$

Thus, using the above inequalities and the upper bound (4.24) in (4.30), we get

$$\|\mathbb{A}_{e,j}(\omega)\|, \|\dot{\mathbb{A}}_{e,j}(\omega)\| \lesssim |\mathbf{k}|^{-1}, \quad \|\mathbb{A}_{m,\ell}(\omega)\|, \|\dot{\mathbb{A}}_{m,\ell}(\omega)\| \lesssim |\mathbf{k}|^{-1}.$$

The formulas (3.7) show that $\mathbb{A}_e(\omega)$ and $\mathbb{A}_m(\omega)$ are the sums of a fixed (independent of ω) operator with a (fixed) linear combination of $(\mathbf{e}, \mathbf{h}, \mathbb{A}_{e,j}(\omega), \dot{\mathbb{A}}_{e,j}(\omega), \mathbb{A}_{m,\ell}(\omega), \dot{\mathbb{A}}_{m,\ell}(\omega))$. They are thus uniformly bounded (in \mathbf{k}) on $\mathcal{C}_{\pm\infty,|\mathbf{k}|}$, proving (4.27).

Finally, using (4.24), (4.31) and that, on $\mathcal{C}_{\pm\infty,|\mathbf{k}|}$, $|\omega\mu(\omega)|^{-1} \lesssim \mu_0^{-1} |\mathbf{k}|^{-1}$, when $|\mathbf{k}| \rightarrow +\infty$, one sees on the definition (3.11) of the operator $\mathcal{V}_{|\mathbf{k}|}(\omega)$ that, for $|\mathbf{k}|$ large enough,

$$\|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \lesssim 1, \quad \forall \omega \in \mathcal{C}_{\pm\infty,|\mathbf{k}|}. \quad (4.32)$$

Conclusion. It suffices to use (4.28) and (4.32) in (4.15) for $\Pi(|\mathbf{k}|) = \Pi_{\pm\infty}(|\mathbf{k}|)$, to conclude that $\|\Pi_{\pm\infty}(|\mathbf{k}|)\| \lesssim 1$ for $|\mathbf{k}|$ large enough. \square

The next lemma is about the asymptotic expansion (in powers of $|\mathbf{k}|$) of $\omega_{\pm\infty}(|\mathbf{k}|)$ when $|\mathbf{k}| \rightarrow +\infty$. It is important to push the expansion up to the first apparition of a negative imaginary part, since it will govern the decay of $\mathbb{U}_\infty(|\mathbf{k}|, t)$ for large $|\mathbf{k}|$.

Lemma 4.4. *The eigenvalues $\omega_{\pm\infty}(|\mathbf{k}|)$ satisfy the following asymptotic expansion*

$$\omega_{\pm\infty}(|\mathbf{k}|) = \pm c |\mathbf{k}| \pm \frac{A_{1,\infty}}{2c} |\mathbf{k}|^{-1} - i \frac{A_{2,\infty}}{2c^2} |\mathbf{k}|^{-2} + o(|\mathbf{k}|^{-2}), \quad \text{as } |\mathbf{k}| \rightarrow +\infty, \quad (4.33)$$

where $A_{1,\infty}$ and $A_{2,\infty}$ are two real constants given by

$$A_{1,\infty} = \sum \Omega_{e,j}^2 + \sum \Omega_{m,\ell}^2 \quad \text{and} \quad A_{2,\infty} = \sum \alpha_{e,j} \Omega_{e,j}^2 + \sum \alpha_{m,\ell} \Omega_{m,\ell}^2, \quad (4.34)$$

with $A_{1,\infty} > 0$ and $A_{2,\infty} > 0$ thanks to the weak dissipation condition (1.11). In particular

$$\text{Im}(\omega_{\pm\infty}(|\mathbf{k}|)) = -\frac{A_{2,\infty}}{2c^2} |\mathbf{k}|^{-2} + o(|\mathbf{k}|^{-2}), \quad \text{as } |\mathbf{k}| \rightarrow +\infty. \quad (4.35)$$

Proof. Since $\omega_{-\infty}(|\mathbf{k}|) = -\overline{\omega_{+\infty}(|\mathbf{k}|)}$ (see Section 3.2.1), we only prove (4.33) for

$$\omega_{+\infty}(|\mathbf{k}|) = \xi_+ (|\mathbf{k}|^{-1})^{-1}, \quad (\text{cf. the proof of Proposition 4.1, step 2}).$$

Step 1: Proof of $\omega_{+\infty}(|\mathbf{k}|) = c|\mathbf{k}| + O(|\mathbf{k}|^{-1})$

The proof of Proposition 4.1 defines the functions $\zeta \mapsto \xi_\pm(\zeta)$ as the branches of solutions of the equation $G(\xi) = \zeta^2$ near $\xi = 0$, where $G(\xi) = \xi^2 g(\xi)$ and $g(\xi) = \varepsilon(1/\xi)^{-1} \mu(1/\xi)^{-1}$ is analytic near 0. In particular $g'(0) = 0$ and, by Lemma A.1 (with $z = 0, g(0) = c^2$ and $\mathbf{m} = 2$), and more precidely (A.22), we deduce that

$$\xi_+(\zeta) = c^{-1} \zeta (1 + O(\zeta^2)) \text{ near } 0, \quad \text{thus } \omega_{+\infty}(|\mathbf{k}|) = c|\mathbf{k}| + O(|\mathbf{k}|^{-1}) \text{ at } \infty. \quad (4.36)$$

Step 2: Computations of the terms of order $|\mathbf{k}|^{-1}$ and $|\mathbf{k}|^{-2}$

By analyticity of $\xi_+(\zeta)$, and using the step 1, we know that, for some coefficients (A_1, A_2) ,

$$\xi_+(\zeta) = c^{-1} \zeta (1 + A_1 \zeta^2 + A_2 \zeta^3 + O(\zeta^4)). \quad (4.37)$$

In order to compute A_1 and A_2 , we are going to proceed by identification in the equation

$$\xi_+(\zeta)^2 g(\xi_+(\zeta)) = \zeta^2. \quad (4.38)$$

Let us compute the Taylor expansion of the left hand side of (4.38). By (1.6,1.7), one has

$$\begin{aligned} \varepsilon(1/\xi) &= \varepsilon_0 \left[1 - \left(\sum \Omega_{e,j}^2 \right) \xi^2 + i \left(\sum \alpha_{e,j} \Omega_{e,j}^2 \right) \xi^3 + O(\xi^4) \right], \quad \text{as } \xi \rightarrow 0, \\ \mu(1/\xi) &= \mu_0 \left[1 - \left(\sum \Omega_{m,\ell}^2 \right) \xi^2 + i \left(\sum \alpha_{m,\ell} \Omega_{m,\ell}^2 \right) \xi^3 + O(\xi^4) \right], \quad \text{as } \xi \rightarrow 0. \end{aligned}$$

It follows that, with $(A_{1,\infty}, A_{2,\infty})$ the coefficients defined by (4.34),

$$g(\xi) = c^2 (1 + A_{1,\infty} \xi^2 - i A_{2,\infty} \xi^3 + O(\xi^4)), \quad \text{as } \xi \rightarrow 0, \quad (4.39)$$

that is to say, as $\xi(\zeta) = O(\zeta)$,

$$\xi_+(\zeta)^2 g(\xi_+(\zeta)) = c^2 (\xi_+(\zeta)^2 + A_{1,\infty} \xi_+(\zeta)^4 - i A_{2,\infty} \xi_+(\zeta)^5 + O(\zeta^6)), \quad \text{as } \zeta \rightarrow 0. \quad (4.40)$$

On the other hand, using (4.37), we have, truncating the expansions at the order 5 in ζ ,

$$\begin{aligned} \xi_+(\zeta)^2 &= c^{-2} \zeta^2 (1 + 2 A_1 \zeta^2 + 2 A_2 \zeta^3 + O(\zeta^4)), \\ \xi_+(\zeta)^4 &= c^{-4} \zeta^4 (1 + O(\zeta^2)), \quad \xi_+(\zeta)^5 = c^{-5} \zeta^5 (1 + O(\zeta^2)), \end{aligned}$$

which we can substitute into (4.40) to obtain

$$\xi_+(\zeta)^2 g(\xi_+(\zeta)) = \zeta^2 (1 + 2 A_1 \zeta^2 + 2 A_2 \zeta^3) + A_{1,\infty} c^{-2} \zeta^4 - i A_{2,\infty} c^{-3} \zeta^5 + O(\zeta^6). \quad (4.41)$$

Using the above in (4.38) leads to the equations

$$2 A_1 + A_{1,\infty} c^{-2} = 0, \quad 2 A_2 - i A_{2,\infty} c^{-3} = 0.$$

from which we conclude easily. The remaining details are left to the reader. \square

We conclude this section by estimating the term $\mathbb{U}_\infty(\mathbf{k}, t)$ in the decomposition (4.9).

Lemma 4.5. *For any constant $C \in (0, A_{2,\infty}/(2c^2))$ with $A_{2,\infty}$ given by (4.34), there exists $k_+ > 0$ such that for $|\mathbf{k}| \geq k_+$, the functions $\mathbb{U}_\infty(\mathbf{k}, t)$ defined by (4.9) and (4.10) satisfy*

$$|\mathbb{U}_\infty(\mathbf{k}, t)| \lesssim e^{-\frac{C}{|\mathbf{k}|^2} t} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0. \quad (4.42)$$

Proof. The functions $\mathbb{U}_\infty(\mathbf{k}, t)$ are defined by (4.9) and (4.10) for $|\mathbf{k}|$ large enough. As the operator $\mathcal{R}_\mathbf{k}$ and its inverse $\mathcal{R}_\mathbf{k}^*$ are unitary, one has

$$|\mathbb{U}_\infty(\mathbf{k}, t)| \leq \sum_{\pm} e^{\text{Im}(\omega_{\pm\infty}(|\mathbf{k}|)) t} \|\Pi_{\pm\infty}(|\mathbf{k}|)\| |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0.$$

Thus, the upper bound (4.42) follows immediately from (4.35) (Lemma 4.4) and the inequality $\|\Pi_{\pm\infty}(|\mathbf{k}|)\| \lesssim 1$ (Lemma 4.3). \square

4.3.3 Estimates of $\mathbb{U}_s(\mathbf{k}, t)$ for $|\mathbf{k}| \gg 1$

We estimate now the term $\mathbb{U}_s(\mathbf{k}, t)$ in (4.9) which involves in particular the projectors $\Pi_p(|\mathbf{k}|)$.

Lemma 4.6. *The spectral projectors $\Pi_p(|\mathbf{k}|)$, $p \in \mathcal{P}_s$ are uniformly bounded for large $|\mathbf{k}|$.*

Proof. We follow the approach described in Section 4.3.1 for $\omega(|\mathbf{k}|) = \omega_p(|\mathbf{k}|)$, $p \in \mathcal{P}_s$ and denote $\mathcal{C}_{p,|\mathbf{k}|}$ the corresponding contour, see (4.11) and figure 2.

Step 1: estimate of $\rho_{|\mathbf{k}|}$. Obviously, for $|\mathbf{k}|$ large enough, since $\omega_p(|\mathbf{k}|)$ tends to p when $|\mathbf{k}| \rightarrow +\infty$, the infimum defining $\rho_{|\mathbf{k}|}$ in (4.13) is attained at the point p , i.e.:

$$\rho_{|\mathbf{k}|} = \frac{1}{2} |\omega_p(|\mathbf{k}|) - p| \sim \frac{1}{2} |A_p| |\mathbf{k}|^{-2} \quad (|\mathbf{k}| \rightarrow +\infty), \quad \text{with } |A_p| > 0, \quad (4.43)$$

according to (4.2) for $\mathbf{m}_p = 1$ since $p \in \mathcal{P}_s$.

For the rest of the proof, we essentially follow the proof of Lemma 4.3.

Step 2: Estimate of $(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1}$.

First, due to (4.2) and (4.43), it is clear that there exists $C > 0$ such that, for $|\mathbf{k}|$ large enough,

$$\forall \tilde{p} \in \mathcal{P} \setminus \{p\}, \quad \forall \omega \in \mathcal{C}_{p,|\mathbf{k}|}, \quad |\omega - \omega_{\tilde{p},n}(|\mathbf{k}|)|, |\omega - \omega_{\tilde{p}}(|\mathbf{k}|)|, |\omega - \omega_{\tilde{p},r}(|\mathbf{k}|)| \geq C > 0.$$

Thus, for $|\mathbf{k}|$ large enough, the function $D_{b,|\mathbf{k}|}(\omega)$ given by the product (4.19) satisfies

$$\forall \omega \in \mathcal{C}_{p,|\mathbf{k}|}, \quad |D_{b,|\mathbf{k}|}(\omega)| \geq C \rho_{|\mathbf{k}|} \quad \text{thus} \quad |D_{b,|\mathbf{k}|}(\omega)|^{-1} \lesssim \rho_{|\mathbf{k}|}^{-1}. \quad (4.44)$$

On the other hand, due to (4.2) and (4.43),

$$|\omega - \omega_{\pm\infty}(|\mathbf{k}|)| = c|\mathbf{k}| + o(|\mathbf{k}|) \quad \text{uniformly for } \omega \in \mathcal{C}_{p,|\mathbf{k}|}.$$

It is then immediate to infer from the definition (4.18) of $D_{|\mathbf{k}|}(\omega)$ that, for $|\mathbf{k}|$ large enough,

$$|D_{|\mathbf{k}|}(\omega)|^{-1} \lesssim \rho_{|\mathbf{k}|}^{-1} |\mathbf{k}|^{-2}, \quad \forall \omega \in \mathcal{C}_{p,|\mathbf{k}|}. \quad (4.45)$$

Moreover, $p \in \mathcal{P}_s$ is a simple zero of the product $Q_e Q_m$. Hence, for any $\omega \in \mathcal{C}_{p,|\mathbf{k}|}$

$$|Q_e(\omega) Q_m(\omega)| \lesssim |\omega - p| \lesssim \rho_{|\mathbf{k}|} + |\omega_p(|\mathbf{k}|) - p| \sim \frac{3}{2} |A_p| |\mathbf{k}|^{-2} \quad (4.46)$$

As a consequence, from the expression (4.17) of $(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1}$, it is clear that

$$|(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1}| \lesssim \rho_{|\mathbf{k}|}^{-1} |\mathbf{k}|^{-4}, \quad \forall \omega \in \mathcal{C}_{p,|\mathbf{k}|}. \quad (4.47)$$

Step 3: Estimates of $\mathcal{S}_{|\mathbf{k}|}(\omega)$ and $\mathcal{V}_{|\mathbf{k}|}(\omega)$.

To estimate $\mathcal{S}_{|\mathbf{k}|}(\omega)$, it remains to estimate $\omega\mu(\omega)\mathbb{A}_e(\omega) - |\mathbf{k}|\mathbf{e}_3 \times \mathbb{A}_m(\omega)$. Due to this, we shall be led to distinguish two cases, $p \in \mathcal{P}_m$ and $p \in \mathcal{P}_e$, and show that, for $|\mathbf{k}|$ large enough,

$$\begin{cases} \text{If } p \in \mathcal{P}_m, \quad \forall \omega \in \mathcal{C}_{p,|\mathbf{k}|}, \quad \|\mathcal{S}_{|\mathbf{k}|}(\omega)\| \lesssim \rho_{|\mathbf{k}|}^{-1} |\mathbf{k}|^{-1}, \quad \|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \lesssim |\mathbf{k}|, & (i) \\ \text{If } p \in \mathcal{P}_e, \quad \forall \omega \in \mathcal{C}_{p,|\mathbf{k}|}, \quad \|\mathcal{S}_{|\mathbf{k}|}(\omega)\| \lesssim \rho_{|\mathbf{k}|}^{-1} |\mathbf{k}|^{-2}, \quad \|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \lesssim |\mathbf{k}|^2. & (ii) \end{cases} \quad (4.48)$$

Proof of (4.48)(i). As $p \in \mathcal{P}_m \cap \mathcal{P}_s$, $p \notin \mathcal{P}_e$ and by \mathbf{H}_1 , there is a unique index ℓ_0 such that $p = \pm\omega_{m,\ell_0}$. Thanks to (4.30), all the operators $(\mathbb{A}_{e,j}(\omega), \dot{\mathbb{A}}_{e,j}(\omega), \mathbb{A}_{m,\ell}(\omega), \dot{\mathbb{A}}_{m,\ell}(\omega))$ are uniformly bounded along $\mathcal{C}_{p,|\mathbf{k}|}$ except $(\mathbb{A}_{m,\ell_0}(\omega), \dot{\mathbb{A}}_{m,\ell_0}(\omega))$, as p is a simple pole of $\omega \mapsto \mathbb{A}_{m,\ell_0}(\omega)$ and $\omega \mapsto \dot{\mathbb{A}}_{m,\ell_0}(\omega)$. More precisely, as $|q_{m,\ell_0}(\omega)|^{-1} \sim |\omega - p|^{-1}$ when $\omega \rightarrow p$, for $|\mathbf{k}|$ large enough,

$$\forall \omega \in \mathcal{C}_{p,|\mathbf{k}|}, \quad |q_{m,\ell_0}(\omega)|^{-1} \lesssim |\omega - p|^{-1} = \rho_{|\mathbf{k}|}^{-1} \lesssim |\mathbf{k}|^2, \quad \text{according to (4.43)} \quad (4.49)$$

Thus, using the bounds (4.30) and the fact that $\mathcal{C}_{p,|\mathbf{k}|}$ lies in a bounded set,

$$\forall \omega \in \mathcal{C}_{p,|\mathbf{k}|}, \quad \|\mathbb{A}_{m,\ell_0}(\omega)\|, \|\dot{\mathbb{A}}_{m,\ell_0}(\omega)\| \lesssim |\mathbf{k}|^2. \quad (4.50)$$

Thanks to formulas (3.7) for $(\mathbb{A}_e(\omega), \mathbb{A}_m(\omega))$, we deduce that

$$\forall \omega \in \mathcal{C}_{p,|\mathbf{k}|}, \quad \|\mathbb{A}_e(\omega)\| \lesssim 1, \quad \|\mathbb{A}_m(\omega)\| \lesssim |\mathbf{k}|^2. \quad (4.51)$$

Moreover, as $p \in \mathcal{P}_m$ is a simple pole of $\mu(\omega)$ and $\mathcal{C}_{p,|\mathbf{k}|}$ is uniformly bounded:

$$|\mu(\omega)| \sim \mu_p |\omega - p|^{-1}, \quad \text{with } \mu_p > 0, \quad \text{thus } |\mu(\omega)| \lesssim |\omega - p|^{-1} \lesssim |\mathbf{k}|^2 \text{ along } \mathcal{C}_{p,|\mathbf{k}|}, \quad (4.52)$$

which, joined to (4.51), gives (the dominant term is the second one)

$$\forall \omega \in \mathcal{C}_{p,|\mathbf{k}|}, \quad \|\omega \mu(\omega) \mathbb{A}_e(\omega) - |\mathbf{k}| \mathbf{e}_3 \times \mathbb{A}_m(\omega)\| \lesssim |\mathbf{k}|^3. \quad (4.53)$$

The estimate for $\mathcal{S}_{|\mathbf{k}|}(\omega)$ in (4.48)(i) then follows from (4.47), (4.53) and (3.10).

For $\mathcal{V}_{|\mathbf{k}|}(\omega)$, we first observe that $|\mu(\omega)| \sim \mu_p |\omega - p|^{-1}$ when $\omega \rightarrow p$ also implies that

$$|\mu(\omega)|^{-1} \lesssim |\mathbf{k}|^{-2} \text{ along } \mathcal{C}_{p,|\mathbf{k}|}. \quad (4.54)$$

On the other hand one sees on formula (3.11) that the dominant term in $\mathcal{V}_{|\mathbf{k}|}(\omega)$ on $\mathcal{C}_{p,|\mathbf{k}|}$ is asymptotically proportional to the function

$$|\mathbf{k}| \mu(\omega)^{-1} q_{m,\ell_0}(\omega)^{-1}.$$

This explains the estimate for $\mathcal{V}_{|\mathbf{k}|}(\omega)$ in (4.48)(i) since, according to the inequalities (4.49) and (4.54), $|\mu(\omega)|^{-1} |q_{m,\ell_0}(\omega)^{-1}| \lesssim 1$ on $\mathcal{C}_{p,|\mathbf{k}|}$.

Proof of (4.48)(ii). The proof is very similar to that of (i) and we shall simply point out the difference with (i). As $p \in \mathcal{P}_e \cap \mathcal{P}_s$, $p \notin \mathcal{P}_m$ and there is a unique index j_0 such that $p = \pm \omega_{e,j_0}$.

This time, all the operators $(\mathbb{A}_{e,j}(\omega), \dot{\mathbb{A}}_{e,j}(\omega), \mathbb{A}_{m,\ell}(\omega), \dot{\mathbb{A}}_{m,\ell}(\omega))$ are uniformly bounded along $\mathcal{C}_{p,|\mathbf{k}|}$ except $(\mathbb{A}_{e,j_0}(\omega), \dot{\mathbb{A}}_{e,j_0}(\omega))$ and, similarly to (4.51), one shows that

$$\|\mathbb{A}_e(\omega)\| \lesssim |\mathbf{k}|^2, \quad \|\mathbb{A}_m(\omega)\| \lesssim 1. \quad (4.55)$$

The difference with (i) is that, this time, $\mu(\omega)$ remains bounded on $\mathcal{C}_{p,|\mathbf{k}|}$, which is the reason why we have

$$\|\omega \mu(\omega) \mathbb{A}_e(\omega) - |\mathbf{k}| \mathbf{e}_3 \times \mathbb{A}_m(\omega)\| \lesssim |\mathbf{k}|^2, \quad (4.56)$$

which leads to the estimate for $\mathcal{S}_{|\mathbf{k}|}(\omega)$ in (4.48)(ii).

For $\mathcal{V}_{|\mathbf{k}|}(\omega)$, one sees on formula (3.11) that, this time, the dominant terms in $\mathcal{V}_{|\mathbf{k}|}(\omega)$ on $\mathcal{C}_{p,|\mathbf{k}|}$ is now proportional to

$$|q_{e,j_0}(\omega)^{-1}| \lesssim |\omega - p|^{-1} \lesssim |\rho_{\mathbf{k}}|^{-1} \lesssim |\mathbf{k}|^2.$$

This explains the estimate for $\mathcal{V}_{|\mathbf{k}|}(\omega)$ in (4.48)(ii).

Conclusion. The important observation is that, from (4.48), in all cases

$$\|\mathcal{S}_{|\mathbf{k}|}(\omega)\| \|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \lesssim \rho_{|\mathbf{k}|}^{-1}.$$

One concludes by applying (4.15) to $\Pi(|\mathbf{k}|) = \Pi_p(|\mathbf{k}|)$. □

We now give the asymptotic expansion of the eigenvalues $\omega_p(|\mathbf{k}|)$ for large $|\mathbf{k}|$. As the study of their imaginary part requires to distinguish two cases, for the sake of readability, we delay it to the corollary 4.8.

Lemma 4.7. *Let $p \in \mathcal{P}_s$. The eigenvalue $\omega_p(|\mathbf{k}|)$ satisfies the following asymptotic expansion*

$$\omega_p(|\mathbf{k}|) = p + A_{p,2} |\mathbf{k}|^{-2} + A_{p,4} |\mathbf{k}|^{-4} + o(|\mathbf{k}|^{-4}), \quad \text{as } |\mathbf{k}| \rightarrow +\infty, \quad (4.57)$$

where the complex numbers $A_{p,2}$ and $A_{p,4}$ are given by

$$\text{if } p = \pm \omega_{e,j_p} \in \mathcal{P}_e, \quad A_{p,2} = -\frac{1}{2} \varepsilon_0 p \mu(p) \Omega_{e,j_p}^2 \neq 0 \quad \text{and} \quad A_{p,4} = \frac{1}{2} \left((\omega^2 \mu h_{e,p})^2 \right)'(p). \quad (4.58)$$

$$\text{if } p = \pm \omega_{m,\ell_p} \in \mathcal{P}_m, \quad A_{p,2} = -\frac{1}{2} \mu_0 p \varepsilon(p) \Omega_{m,\ell_p}^2 \neq 0 \quad \text{and} \quad A_{p,4} = \frac{1}{2} \left((\omega^2 \varepsilon h_{m,p})^2 \right)'(p).$$

where for any $\omega \in (\mathbb{C} \setminus \mathcal{P}) \cup \{p\}$,

$$\text{if } p = \pm \omega_{e,j_p}, \quad h_{e,p}(\omega) = (\omega - p) \varepsilon(\omega), \quad \text{if } p = \pm \omega_{e,j_p}, \quad h_{m,p}(\omega) = (\omega - p) \mu(\omega). \quad (4.59)$$

Proof. Let $p \in \mathcal{P}_s$. Thus, either $p \in \mathcal{P}_e$ or $p \in \mathcal{P}_m$. We assume without a loss of generality that $p \in \mathcal{P}_e$. The proof is done by symmetric arguments if $p \in \mathcal{P}_m$.

We are in the situation covered by the step 1 of the proof of Proposition 4.1, in the case $\mathbf{m}_p = 1$, see also (4.1). As p is a simple pole of $\omega \varepsilon(\omega)$, there exists (by H₁) a unique index $1 \leq j_p \leq N_e$ such that $p = \pm \omega_{e,j_p}$. In that case the set $\{\omega_{p,n}(|\mathbf{k}|), 1 \leq n \leq \mathbf{m}_p\}$ is reduced to one single function denoted $|\mathbf{k}| \mapsto \omega_p(|\mathbf{k}|)$. It is constructed as the unique branch of solutions of the equation $\mathcal{D}(\omega)^{-1} = \zeta$ with $\zeta = |\mathbf{k}|^{-2}$ that converges to p when $\zeta \rightarrow 0$.

Next, we make more precise the factorization of $\mathcal{D}(\omega)^{-1} \equiv (\omega^2 \varepsilon \mu)^{-1}$ corresponding to (4.1) for $\mathbf{m}_p = 1$. By definition (4.59) of $h_{e,p}$, we can write

$$\mathcal{D}(\omega)^{-1} = (\omega - p) g_p(\omega) \quad \text{with} \quad g_p(\omega) := (\omega^2 \mu h_{e,p})(\omega)^{-1}.$$

From the definition of $\varepsilon(\omega)$, we also have

$$h_{e,p}(\omega) = \varepsilon_0 (\omega - p) \left(1 - \sum_{j \neq j_p} \frac{\Omega_{e,j}^2}{q_{e,j}(\omega)} \right) - \varepsilon_0 \frac{\Omega_{e,j_p}^2}{\omega + p}, \quad \left(\text{thus } h_{e,p}(p) = -\varepsilon_0 \frac{\Omega_{e,j_p}^2}{2p} \right) \quad (4.60)$$

which shows that g_p is analytic in the neighbourhood of p (by (H₂), $\mu(p) \neq 0$) and that

$$g_p(p) = -2 \left(\varepsilon_0 p \mu(p) \Omega_{e,j_p}^2 \right)^{-1} \neq 0.$$

By formula (A.22), with $\zeta = |\mathbf{k}|^{-2}$, $z = p$, $g = g_p$, $\mathbf{m} = 1$, $a_n^{-1} = g(p)^{-1}$, we get

$$\omega_p(|\mathbf{k}|) = p - \varepsilon_0 p \mu(p) \frac{\Omega_{e,j_p}^2}{2} |\mathbf{k}|^{-2} - \frac{g_p'(p) |\mathbf{k}|^{-4}}{g_p(p)^3} + o(|\mathbf{k}|^{-4}), \quad \text{as } |\mathbf{k}| \rightarrow +\infty. \quad (4.61)$$

This is nothing but (4.57) since $-g_p'/g_p^2 = (1/g_p) = (\omega^2 \mu h_{e,p})'$ by definition of g_p . Thus, one gets

$$-\frac{g_p'(p)}{g_p(p)^3} = (\omega^2 \mu h_{e,p})'(p) (\omega^2 \mu h_{e,p})(p) = \frac{1}{2} \left((\omega^2 \mu h_{e,p})^2 \right)'(p).$$

Finally, note that, since $\mathcal{P}_e \cap \mathcal{Z}_m = \emptyset$, (see (H₂)), $p \notin \mathcal{Z}_m$ so that $\mu(p) \neq 0$ thus $A_{p,2} \neq 0$. \square

The behaviour of the imaginary part of $A_{2,p}$ leads us to make the distinction between the critical and non critical configurations. For the ease of the reader, we recall below that the critical case corresponds to one of the following two situations (see definition 1.8):

1. $\forall \ell \in \{1, \dots, N_m\}$, $\alpha_{m,\ell} = 0$ and $\exists j \in \{1, \dots, N_e\} \mid \alpha_{e,j} = 0$ and $\omega_{e,j} \notin \{\omega_{m,\ell}\}$.
2. $\forall j \in \{1, \dots, N_e\}$, $\alpha_{e,j} = 0$ and $\exists \ell \in \{1, \dots, N_m\} \mid \alpha_{m,\ell} = 0$ and $\omega_{m,\ell} \notin \{\omega_{e,j}\}$.

Corollary 4.8. *Following the notations introduced in Lemma 4.7. In the non critical case,*

$$\forall p \in \mathcal{P}_s, \quad \text{Im } A_{2,p} < 0. \quad (4.62)$$

As a consequence, there exists $C > 0$ such that, for $|\mathbf{k}|$ large enough,

$$\forall p \in \mathcal{P}_s, \quad \text{Im } \omega_p(|\mathbf{k}|) < -C |\mathbf{k}|^{-2}. \quad (4.63)$$

In the critical case, in situation 1, $\mathcal{P}_e \cap \mathcal{P}_s \neq \emptyset$ and

$$\forall p \in \mathcal{P}_e \cap \mathcal{P}_s, \quad \text{Im } A_{2,p} = 0 \text{ and } \text{Im } A_{4,p} < 0, \quad \forall p \in \mathcal{P}_m \cap \mathcal{P}_s, \quad \text{Im } A_{2,p} < 0, \quad (4.64)$$

while, in situation 2, $\mathcal{P}_m \cap \mathcal{P}_s \neq \emptyset$ and

$$\forall p \in \mathcal{P}_m \cap \mathcal{P}_s, \quad \text{Im } A_{2,p} = 0 \text{ and } \text{Im } A_{4,p} < 0, \quad \forall p \in \mathcal{P}_e \cap \mathcal{P}_s, \quad \text{Im } A_{2,p} < 0. \quad (4.65)$$

As a consequence, there exists $C > 0$ such that, for $|\mathbf{k}|$ large enough

$$\forall p \in \mathcal{P}_s, \quad \text{Im } \omega_p(|\mathbf{k}|) < -C |\mathbf{k}|^{-4}. \quad (4.66)$$

Proof. Non critical case. We prove (4.62) for $p \in \mathcal{P}_s \cap \mathcal{P}_e$, the case $p \in \mathcal{P}_s \cap \mathcal{P}_m$ is similar.

According to formula (4.58), we simply have to check that $\text{Im}(p\mu(p)) > 0$. Thanks to (3.17), this will be true as soon as, at least one index ℓ , $\alpha_{m,\ell} > 0$. However, $p \in \mathcal{P}_s \cap \mathcal{P}_e$ means that $p = \pm \omega_{e,j}$ for some j with $\alpha_{e,j} = 0$ and also that $\omega_{e,j} \notin \{\omega_{m,\ell}\}$ (since p is a simple pole). Thus, if all $\alpha_{m,\ell}$ vanished, we would be in the situation 1 of the critical case, which is excluded.

Of course, (4.63) is a direct consequence of (4.62).

Critical case. We prove (4.64) in situation 1, the proof of (4.65) in situation 2 being similar.

First $\mathcal{P}_e \cap \mathcal{P}_s$ is non empty since any $\omega_{e,j}$ such that $\alpha_{e,j} = 0$ and $\omega_{e,j} \notin \{\omega_{m,\ell}\}$ (a set which is itself non empty by definition of the situation 1) belongs to $\mathcal{P}_e \cap \mathcal{P}_s$.

Let $p \in \mathcal{P}_e \cap \mathcal{P}_s$. In situation 1, all $\alpha_{m,\ell}$ vanish thus, by (3.17), $\mu(\omega)$ and $\mu'(\omega)$ are real-valued for $\omega \in \mathbb{R} \setminus \mathcal{P}_m$. In particular, $\text{Im}(p\mu(p)) = 0$ which implies by (4.58) that $\text{Im } A_{2,p} = 0$. Next, as $h_{e,p}(p) \in \mathbb{R}$ by (4.60), using (4.58) again yields

$$A_{p,4} = (\omega^2 \mu h_{e,p})'(p) (\omega^2 \mu h_{e,p})(p) \Rightarrow \text{Im } A_{p,4} = p^4 \mu^2(p) h_{e,p}(p) \text{Im } h'_{e,p}(p). \quad (4.67)$$

Moreover, $p \in \mathcal{P}_e \cap \mathcal{P}_s$ implies $p = \pm \omega_{e,k}$ for some k such that $\alpha_{e,k} = 0$. However, from the formulas for $\mu(\omega)$ (1.6), the polynomials $q_{e,j}$ (1.7) and $h_{e,p}$ (4.60), one computes that

$$h_{e,p}(p) = -\varepsilon_0 \frac{\Omega_{e,k}^2}{2p} \quad \text{and} \quad \text{Im } h'_{e,p}(p) = \varepsilon_0 p \sum_{j \neq k} \frac{\alpha_{e,j} \Omega_{e,j}^2}{|q_{e,j}(p)|^2}. \quad (4.68)$$

The weak dissipation condition (1.11) implies that at least one $\alpha_{e,j}$ for $j \neq k$ is positive so that $h_{e,p}(p) \operatorname{Im} h'_{e,p}(p) < 0$, implying by (4.67) that $\operatorname{Im} A_{p,4} < 0$.

Finally, for $p \in \mathcal{P}_s \cap \mathcal{P}_m$, the proof of $\operatorname{Im} A_{p,2} < 0$ uses (4.58), (3.17) and the fact that at least one $\alpha_{e,j}$ for $j \neq k$ is positive. The proof of (4.64) is thus complete.

Finally (4.66) is easily deduced from (4.57), (4.64) and (4.65) after remarking that, for $|\mathbf{k}|$ large enough, $|\mathbf{k}|^{-4} \lesssim |\mathbf{k}|^{-2}$ which implies $-|\mathbf{k}|^{-2} \lesssim -|\mathbf{k}|^{-4}$ (details are again left to the reader). \square

Thus, proceeding as in the proof of Lemma 4.5, one shows using Lemma 4.6 and Corollary 4.8 the following result.

Lemma 4.9. *There exists $k_+ > 0$ and $C > 0$ such that the function $\mathbb{U}_s(\mathbf{k}, t)$ defined by (4.9) and (4.34) satisfies the following estimates:*

1. *If the system (1.1) is in a non-critical configuration, then*

$$|\mathbb{U}_s(\mathbf{k}, t)| \lesssim e^{-\frac{Ct}{|\mathbf{k}|^2}} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0 \quad \text{and} \quad \forall |\mathbf{k}| \geq k_+. \quad (4.69)$$

2. *If the system (1.1) is in a critical configuration, then*

$$|\mathbb{U}_s(\mathbf{k}, t)| \lesssim e^{-\frac{Ct}{|\mathbf{k}|^4}} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0 \quad \text{and} \quad \forall |\mathbf{k}| \geq k_+. \quad (4.70)$$

4.3.4 Estimates of $\mathbb{U}_d(\mathbf{k}, t)$ for $|\mathbf{k}| \gg 1$

This time we estimate $\mathbb{U}_d(\mathbf{k}, t)$ in (4.9) which involves in particular the projectors $\Pi_{p,r}(|\mathbf{k}|)$.

Lemma 4.10. *The projectors $\Pi_{p,r}(|\mathbf{k}|)$, $p \in \mathcal{P}_d, r = 1, 2$ are uniformly bounded for large $|\mathbf{k}|$.*

Proof. We follow again the approach of Section 4.3.1 for $\omega(|\mathbf{k}|) = \omega_{p,r}(|\mathbf{k}|)$, $p \in \mathcal{P}_d, r = 1, 2$ and denote $\mathcal{C}_{p,r,|\mathbf{k}|}$ the corresponding contour (see (4.11) and figure 2). Without any loss of generality, we can restrict ourselves to $r = 1$.

Step 1: estimate of $\rho_{|\mathbf{k}|}$. Thanks to the asymptotic (4.2), it is clear that the distance from $\omega_{p,1}(|\mathbf{k}|)$ to any other eigenvalue that is different from $\omega_{p,2}(|\mathbf{k}|)$ remains bounded from below, for large $|\mathbf{k}|$, by a positive constant. The same observation holds true for distance from $\omega_{p,1}(|\mathbf{k}|)$ to any other point of the set $\mathcal{S}_{\mathcal{T}}$ than p .

Oppositely, the distances $|\omega_{p,1}(|\mathbf{k}|) - p|$ and $|\omega_{p,1}(|\mathbf{k}|) - \omega_{p,2}(|\mathbf{k}|)|$ tend to 0 when $|\mathbf{k}| \rightarrow +\infty$. More precisely, from (4.2) applied with $\mathfrak{m}_p = 2$, since $p \in \mathcal{P}_d$, one deduces that

$$|\omega_{p,1}(|\mathbf{k}|) - p| \sim |A_p|^{\frac{1}{2}} |\mathbf{k}|^{-1}, \quad |\omega_{p,1}(|\mathbf{k}|) - \omega_{p,2}(|\mathbf{k}|)| \sim 2 |A_p|^{\frac{1}{2}} |\mathbf{k}|^{-1}, \quad (|\mathbf{k}| \rightarrow +\infty) \quad (4.71)$$

with $|A_p| > 0$. As a consequence, for $|\mathbf{k}|$ large enough, we have

$$\rho_{|\mathbf{k}|} = \frac{1}{2} |\omega_p(|\mathbf{k}|) - p| \sim \frac{1}{2} |A_p|^{\frac{1}{2}} |\mathbf{k}|^{-1} \quad (|\mathbf{k}| \rightarrow +\infty), \quad \text{with} \quad |A_p| > 0. \quad (4.72)$$

Step 2: Estimate of $(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1}$. It is very similar to that of in Lemma 4.6. One observes thanks to (4.2) that, when $|\mathbf{k}| \rightarrow +\infty$, in the product (4.18) defining $D_{b,|\mathbf{k}|}(\omega)$, only two terms,

namely $\omega - \omega_{p,1}(|\mathbf{k}|)$ and $\omega - \omega_{p,2}(|\mathbf{k}|)$, are not bounded from below when ω describes $\mathcal{C}_{p,1,|\mathbf{k}|}$ and $|\mathbf{k}| \rightarrow +\infty$. Thus for some $C > 0$, $\omega \in \mathcal{C}_{p,1,|\mathbf{k}|}$ and $|\mathbf{k}|$ large enough

$$|D_{b,|\mathbf{k}|}(\omega)| \geq C \left| \omega - \omega_{p,1}(|\mathbf{k}|) \right| \left| \omega - \omega_{p,2}(|\mathbf{k}|) \right| = C \rho_{|\mathbf{k}|} \left| \omega - \omega_{p,2}(|\mathbf{k}|) \right|. \quad (4.73)$$

By the reverse triangular inequality, along $\mathcal{C}_{p,1,|\mathbf{k}|}$,

$$\left| \omega - \omega_{p,2}(|\mathbf{k}|) \right| \geq \left| \omega_{p,1}(|\mathbf{k}|) - \omega_{p,2}(|\mathbf{k}|) \right| - \left| \omega - \omega_{p,1}(|\mathbf{k}|) \right| = \left| \omega_{p,1}(|\mathbf{k}|) - \omega_{p,2}(|\mathbf{k}|) \right| - \rho_{|\mathbf{k}|},$$

thus, using the equivalent (4.72) for $\rho_{|\mathbf{k}|}$ and $|\omega_{p,1}(|\mathbf{k}|) - \omega_{p,2}(|\mathbf{k}|)| \sim 2|A_p| |\mathbf{k}|^{-1}$,

$$\left| \omega - \omega_{p,2}(|\mathbf{k}|) \right| \geq \left| \omega_{p,1}(|\mathbf{k}|) - \omega_{p,2}(|\mathbf{k}|) \right| - \rho_{|\mathbf{k}|} \sim \frac{3}{2} |A_p| |\mathbf{k}|^{-1}.$$

Therefore, from (4.73), we deduce that

$$|D_{b,|\mathbf{k}|}(\omega)|^{-1} \lesssim \rho_{|\mathbf{k}|}^{-1} |\mathbf{k}|, \quad \forall \omega \in \mathcal{C}_{p,1,|\mathbf{k}|}.$$

Thus, proceeding as in Lemma 4.6 (observe that by passing from $|D_{b,|\mathbf{k}|}(\omega)|^{-1}$ to $|D_{|\mathbf{k}|}(\omega)|^{-1}$, see (4.44) and (4.45), one loses two powers of $|\mathbf{k}|$),

$$|D_{|\mathbf{k}|}(\omega)|^{-1} \lesssim \rho_{|\mathbf{k}|}^{-1} |\mathbf{k}|^{-1}, \quad \forall \omega \in \mathcal{C}_{p,1,|\mathbf{k}|}. \quad (4.74)$$

Moreover, $p \in \mathcal{P}_d$ is a double zero of the product $Q_e Q_m$. Hence, for any $\omega \in \mathcal{C}_{p,1,|\mathbf{k}|}$

$$|Q_e(\omega) Q_m(\omega)| \lesssim |\omega - p|^2 \leq (\rho_{|\mathbf{k}|} + |\omega_p(|\mathbf{k}|) - p|)^2 \sim \frac{9}{4} |A_p| |\mathbf{k}|^{-2} \quad (\text{by (4.71, 4.72)}). \quad (4.75)$$

As a consequence, from (4.17), we finally get by product of (4.74) and (4.75),

$$|(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1}| \lesssim \rho_{|\mathbf{k}|}^{-1} |\mathbf{k}|^{-3}, \quad \forall \omega \in \mathcal{C}_{p,1,|\mathbf{k}|}. \quad (4.76)$$

Step 3: Estimates of $\mathcal{S}_{|\mathbf{k}|}(\omega)$ and $\mathcal{V}_{|\mathbf{k}|}(\omega)$

As $p \in \mathcal{P}_m \cap \mathcal{P}_e$, there is a unique pair of indices $\{\ell_0, j_0\}$ such that $p = \pm\omega_{e,j_0} = \pm\omega_{m,\ell_0}$. The situation is a kind of mix between the two situations (i) and (ii) met in the step 3 of the proof of Lemma 4.6.

Using (4.30) one sees that all the operators $(\mathbb{A}_{e,j}(\omega), \dot{\mathbb{A}}_{e,j}(\omega), \mathbb{A}_{m,\ell}(\omega), \dot{\mathbb{A}}_{m,\ell}(\omega))$ are uniformly bounded along $\mathcal{C}_{p,1,|\mathbf{k}|}$ except $(\mathbb{A}_{e,j_0}(\omega), \dot{\mathbb{A}}_{e,j_0}(\omega), \mathbb{A}_{m,\ell_0}(\omega), \dot{\mathbb{A}}_{m,\ell_0}(\omega))$.

Proceeding as for (4.51) in the proof of Lemma 4.6, one easily gets (we omit the details)

$$\|\mathbb{A}_e(\omega)\| \lesssim |\mathbf{k}|, \quad \|\mathbb{A}_m(\omega)\| \lesssim |\mathbf{k}|. \quad (4.77)$$

As p is a simple pole of $\mu(\omega)$, we obtain, similarly to (4.52) in the proof of Lemma 4.6,

$$\forall \omega \in \mathcal{C}_{p,1,|\mathbf{k}|}, \quad |\mu(\omega)| \lesssim |\omega - p|^{-1} \lesssim |\mathbf{k}|, \quad (4.78)$$

which, joined to (4.77), gives

$$\|\omega\mu(\omega)\mathbb{A}_e(\omega) - |\mathbf{k}| \mathbf{e}_3 \times \mathbb{A}_m(\omega)\| \leq |\mathbf{k}|^2. \quad (4.79)$$

Finally using (4.76) and (4.79) in the definition (3.10) of $\mathcal{S}_{|\mathbf{k}|}(\omega)$, we get

$$\|\mathcal{S}_{|\mathbf{k}|}(\omega)\| \lesssim \rho_{|\mathbf{k}|}^{-1} |\mathbf{k}|^{-1}, \quad \forall \omega \in \mathcal{C}_{p,1,|\mathbf{k}|}. \quad (4.80)$$

For $\mathcal{V}_{|\mathbf{k}|}(\omega)$, we first observe that $|\mu(\omega)| \sim \mu_p |\omega - p|^{-1}$ ($\omega \rightarrow p$) also implies that along $\mathcal{C}_{p,1,|\mathbf{k}|}$,

$$|\mu(\omega)|^{-1} \lesssim \mu_p^{-1} \rho_{|\mathbf{k}|}^{-1} \lesssim |\mathbf{k}|^{-1} \quad \text{as } |\mathbf{k}| \rightarrow +\infty.$$

Thus, the function $|\mathbf{k}| \mu(\omega)^{-1}$ is bounded for $|\mathbf{k}|$ large enough. As a consequence, one sees on formula (3.11) that $\mathcal{V}_{|\mathbf{k}|}(\omega)$ blows with the same rate than $q_{e,j_0}(\omega)^{-1}$ and $q_{m,\ell_0}(\omega)^{-1}$, that is to say proportionally to $\rho_{|\mathbf{k}|}^{-1}$, in other words proportionally to $|\mathbf{k}|$. Therefore

$$\|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \lesssim |\mathbf{k}|, \quad \forall \omega \in \mathcal{C}_{p,1,|\mathbf{k}|}. \quad (4.81)$$

Conclusion. Again, by (4.80) and (4.81), $\|\mathcal{S}_{|\mathbf{k}|}(\omega)\| \|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \leq \rho_{|\mathbf{k}|}^{-1}$ and one concludes with (4.15) for $\Pi(|\mathbf{k}|) = \Pi_p(|\mathbf{k}|)$. \square

We now give the asymptotic expansion of the eigenvalues $\omega_{p,r}(|\mathbf{k}|)$ for large $|\mathbf{k}|$.

Lemma 4.11. *For $p \in \mathcal{P}_d$, then $p = \pm \omega_{e,j_p} = \pm \omega_{m,\ell_p}$ for some (j_p, ℓ_p) and*

$$\omega_{p,r}(|\mathbf{k}|) = p + (-1)^r \frac{\Omega_{e,j_p} \Omega_{m,\ell_p}}{2c} |\mathbf{k}|^{-1} + A_{p,2} |\mathbf{k}|^{-2} + o(|\mathbf{k}|^{-2}), \quad \text{as } |\mathbf{k}| \rightarrow +\infty, \quad (4.82)$$

where the complex number $A_{p,2}$ is given by

$$A_{p,2} = \frac{1}{2} (\omega^2 h_{e,j_p} h_{m,\ell_p})'(p), \quad (4.83)$$

where the functions $h_{e,p}$ and $h_{m,p}$ are defined in (4.59). Moreover, one has

$$\text{Im } A_{p,2} < 0. \quad (4.84)$$

Proof. Let $p \in \mathcal{P}_d$. Then (by H_1), there exist two unique indices $j_p \in \{1, \dots, N_e\}$ and $\ell_p \in \{1, \dots, N_m\}$ such that $p = \pm \omega_{e,j_p} = \pm \omega_{m,\ell_p}$. Then the two functions $|\mathbf{k}| \mapsto \omega_{p,r}(|\mathbf{k}|)$ are defined in the proof of Proposition 4.1 for $|\mathbf{k}|$ large enough via the Lemma A.1 as the two branches of solutions of the equation $\mathcal{D}(\omega)^{-1} = \zeta^2$ with $\zeta = |\mathbf{k}|^{-1}$ in a vicinity of p . Here, $p \in \mathcal{P}_e \cap \mathcal{P}_m$ is a common simple pole of the rational functions $\omega \varepsilon$ and $\omega \mu$. More precisely, from the definition (4.58) of $(h_{e,p}, h_{m,p})$ (see Lemma 4.7) we have

$$\mathcal{D}(\omega) = \omega^2 h_{e,p}(\omega) h_{m,p}(\omega) (\omega - p)^{-2}.$$

The latter expression emphasizes that p is a double pole of the rational function $\mathcal{D}(\omega)$, thus a double zero of \mathcal{D}^{-1} : $\mathcal{D}(\omega)^{-1} = (\omega - p)^2 g(\omega)$ where $g(\omega) := (\omega^2 h_{e,p}(\omega) h_{m,p}(\omega))^{-1}$ is analytic in the vicinity of p and satisfies (using the expression (4.60) for $h_{e,p}$ and its equivalent for $h_{m,p}$):

$$g(p) = 4c^2 \Omega_{e,j_p}^{-2} \Omega_{m,\ell_p}^{-2} > 0.$$

Thus, using the asymptotic formula (A.22) of Lemma A.1 (with $\zeta = |\mathbf{k}|^{-1}$, $z = p$, $\mathbf{m} = 2$ and $a_1 = -\sqrt{g(p)} = -2c/(\Omega_{e,j_p} \Omega_{m,\ell_p})$, $a_2 = \sqrt{g(p)}$ the two roots of $X^2 = g(p)$) yields

$$\omega_{p,r}(|\mathbf{k}|) = p + a_r^{-1} |\mathbf{k}|^{-1} - \frac{a_r^{-2} g'(p)}{2g(p)} |\mathbf{k}|^{-2} + o(|\mathbf{k}|^{-2}), \text{ as } |\mathbf{k}| \rightarrow +\infty. \quad (4.85)$$

To conclude, it remains to remark that $a_r^{-2} = g(p)^{-1}$ and $g^{-1} = \omega^2 h_{e,p} h_{m,p}$ which gives

$$-\frac{a_r^{-2} g'(p)}{g(p)} = -\frac{g'(p)}{g^2(p)} = (g^{-1})'(p) = (\omega^2 h_{e,p} h_{m,p})'(p). \quad (4.86)$$

Finally, it remains to show that $\text{Im } A_{p,2} < 0$. From (4.68) and the expressions (4.59) of $h_{e,p}$ and $h_{m,p}$, it follows (as p^2 , $h_{e,p}(p)$, $h_{m,p}(p)$ are real) that

$$\text{Im } A_{p,2} = \frac{p^2}{2} [h_{m,p}(p) \text{Im } h'_{e,p}(p) + h_{e,p}(p) \text{Im } h'_{m,p}(p)].$$

Finally using that $h_{e,p}(p) = -\varepsilon_0 \Omega_{e,j_p}^2 / (2p)$ and $h_{m,p}(p) = -\mu_0 \Omega_{m,\ell_p}^2 / (2p)$, the expression (4.68) for $\text{Im}(h'_{e,p}(p))$ and its equivalent form for $\text{Im}(h'_{m,p}(p))$ gives that

$$\text{Im } A_{p,2} = -c^{-2} p^2 \left(\frac{\Omega_{m,\ell_p}^2}{4} \sum_{j=1, j \neq j_p}^{N_e} \frac{\Omega_{e,j}^2 \alpha_{e,j}}{|q_{e,j}(p)|^2} + \frac{\Omega_{e,j_p}^2}{4} \sum_{\ell=1, \ell \neq \ell_p}^{N_m} \frac{\Omega_{m,\ell}^2 \alpha_{m,\ell}}{|q_{m,\ell}(p)|^2} \right)$$

(where $c^{-2} = \varepsilon_0 \mu_0$). Thus, this term is negative by the weak dissipation condition (1.11) since at least one coefficient $\alpha_{e,j}$ or $\alpha_{m,\ell}$ is positive. Finally, combining (4.85) and (4.86) gives (4.82). \square

We now estimate the term $\mathbb{U}_d(\mathbf{k}, t)$ for large $|\mathbf{k}|$. Proceeding as in the proof of Lemma 4.5, one shows using Lemmas 4.10 and 4.11 the following result.

Lemma 4.12. *If $\mathcal{P}_d \neq \emptyset$, then there exists $k_+ > 0$ and $C > 0$ such that the function $\mathbb{U}_d(\mathbf{k}, t)$ defined by (4.9) and (4.34) satisfies the following estimate*

$$|\mathbb{U}_d(\mathbf{k}, t)| \lesssim e^{-\frac{Ct}{|\mathbf{k}|^2}} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0 \text{ and } \forall |\mathbf{k}| \geq k_+. \quad (4.87)$$

4.3.5 Estimates of $\mathbb{U}_-(\mathbf{k}, t)$ for $|\mathbf{k}| \gg 1$

As announced in Section 4.3.1, since we simply want to obtain a ‘‘rough’’ exponential decay estimate for $\mathbb{U}_-(\mathbf{k}, t)$, we do not need to separate the analysis in three lemmas as in Sections 4.3.2, 4.3.3 and 4.3.4 but give a direct proof using Riesz-Dunford functional calculus.

Lemma 4.13. *There exists $\delta > 0$ and $k_+ > 0$ such that $\mathbb{U}_-(\mathbf{k}, t)$, defined by (4.9) and (4.10), satisfies*

$$|\mathbb{U}_-(\mathbf{k}, t)| \lesssim e^{-\delta t} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0, \quad \forall |\mathbf{k}| \geq k_+. \quad (4.88)$$

Proof. We introduce a (positively oriented) simple closed contour Γ , included in \mathbb{C}^- such that all the poles of \mathcal{P}_- lies inside Γ (see figure 3). We denote δ the distance from Γ to the real axis:

$$\delta = \min\{-\text{Im}(\omega), \omega \in \Gamma\} > 0.$$

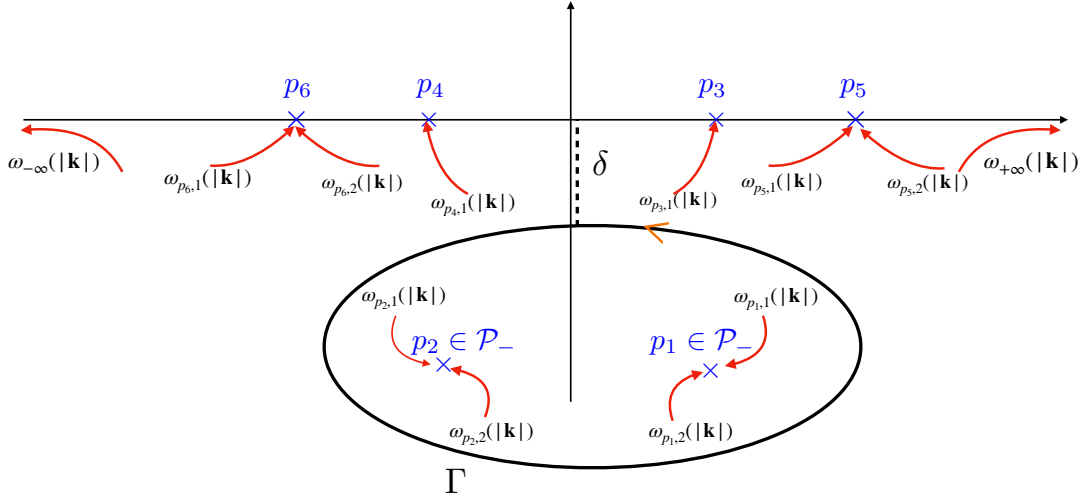


Figure 3: Contour integration for the estimate of $\mathbf{U}_-(\mathbf{k}, t)$ in the case where $\mathcal{P}_- = \{p_1, p_2\}$ (corresponding to the figure 1).

For $|\mathbf{k}|$ large enough, by (4.2) and (4.3), Γ encloses all eigenvalues $\omega_{p,n}(|\mathbf{k}|)$ for $p \in \mathcal{P}_-$ and $n \in \{1, \dots, \mathbf{m}_p\}$ but no other elements of the spectrum of $\mathbb{A}_{|\mathbf{k}|}$. Then, by the Riesz-Dunford functional calculus, there exists $k_+ > 0$ such that we have the formula:

$$\mathbf{U}_-(\mathbf{k}, t) = -\frac{\mathcal{R}_{\mathbf{k}}^*}{2i\pi} \int_{\Gamma} e^{-i\omega t} R_{|\mathbf{k}|}(\omega) \mathcal{R}_{\mathbf{k}} \mathbf{U}_0(\mathbf{k}) d\omega, \quad \text{for } |\mathbf{k}| \geq k_+. \quad (4.89)$$

As $\mathcal{R}_{\mathbf{k}}$ is unitary, it leads to

$$|\mathbf{U}_-(\mathbf{k}, t)| \lesssim e^{-\delta t} \left(\max_{\omega \in \Gamma} \|R_{|\mathbf{k}|}(\omega)\| \right) |\mathbf{U}_0(\mathbf{k})|. \quad (4.90)$$

It remains now to estimate $R_{|\mathbf{k}|}(\omega)$ from its expression $R_{|\mathbf{k}|}(\omega) = \mathcal{V}_{|\mathbf{k}|}(\omega) \mathcal{S}_{|\mathbf{k}|}(\omega) + \mathcal{T}(\omega)$ in Proposition 3.2 (choosing Γ such that $\Gamma \cap \mathcal{Z}_m = \emptyset$). The more involved step concerns the estimate of $\mathcal{S}_{|\mathbf{k}|}(\omega)$.

Step 1: Estimate of $\mathcal{S}_{|\mathbf{k}|}(\omega)$. The expression of $\mathcal{S}_{|\mathbf{k}|}(\omega)$, given in (3.10), involved the term $(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1}$. We first bound this term by using (4.17), (4.18) and (4.19). From the asymptotic behaviour (4.2), it follows that $D_{b,|\mathbf{k}|}(\omega)$ given by (4.19) satisfies:

$$|D_{b,|\mathbf{k}|}(\omega)^{-1}| \lesssim 1, \quad \forall \omega \in \Gamma. \quad (4.91)$$

From the asymptotic expansion (4.3), $|\omega - \omega_{\pm\infty}(|\mathbf{k}|)|^{-1} = c^{-1}|\mathbf{k}|^{-1} + o(|\mathbf{k}|^{-1})$. Hence, it yields with (4.18) and (4.91):

$$|D_{|\mathbf{k}|}(\omega)^{-1}| \lesssim |\mathbf{k}|^{-2}, \quad \text{and thus with (4.17)} \quad |(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1}| \lesssim |\mathbf{k}|^{-2} \quad \forall \omega \in \Gamma. \quad (4.92)$$

On the other hand, from the definitions (3.6) and (3.7), it is immediate that

$$\| -\omega\mu(\omega)\mathbb{A}_e(\omega) + |\mathbf{k}| \mathbf{e}_3 \times \mathbb{A}_h(\omega) \| \lesssim |\mathbf{k}|, \quad \forall \omega \in \Gamma. \quad (4.93)$$

Combining (4.92) and (4.93) with (3.10) yields

$$\|\mathcal{S}_{|\mathbf{k}|}(\omega)\| \lesssim |\mathbf{k}|^{-1}, \quad \forall \omega \in \Gamma \text{ and } |\mathbf{k}| \geq k_+. \quad (4.94)$$

Step 2: Final estimate. From the expression (3.11) and (3.12) of $\mathcal{V}_{|\mathbf{k}|}(\omega)$ and $\mathcal{T}(\omega)$ we deduce that

$$\|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \lesssim |\mathbf{k}|, \quad \|\mathcal{T}(\omega)\| \lesssim 1, \quad \forall \omega \in \Gamma, \quad (4.95)$$

thus by (3.9) and (4.94), $\|\mathcal{R}_{|\mathbf{k}|}(\omega)\| \lesssim 1, \forall \omega \in \Gamma$ which we substitute into (4.90) to get (4.88). \square

4.3.6 The global estimates

Theorem 4.14. *There exists $k_+ > 0$ such that for $|\mathbf{k}| \geq k_+$, the spatial Fourier components $|\mathbb{U}(\mathbf{k}, t)|$ of the solution of (1.13) with initial condition $\mathbf{U}_0 \in \mathcal{H}_\perp$ satisfy the following estimates:*

1. *If the Maxwell system is in a non-critical configuration, then there $\exists C, \tilde{C} > 0$ such that*

$$|\mathbb{U}(\mathbf{k}, t)| \leq \tilde{C} e^{-\frac{Ct}{|\mathbf{k}|^2}} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0. \quad (4.96)$$

2. *If the Maxwell system is in a critical configuration, then there $\exists C, \tilde{C} > 0$ such that*

$$|\mathbb{U}(\mathbf{k}, t)| \leq \tilde{C} e^{-\frac{Ct}{|\mathbf{k}|^4}} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0. \quad (4.97)$$

Proof. The inequalities (4.96) and (4.97) follow immediately from the expression (4.9) of $|\mathbb{U}(\mathbf{k}, t)|$ for $|\mathbf{k}|$ large enough and the estimates of the four terms $|\mathbb{U}_\infty(\mathbf{k}, t)|$, $|\mathbb{U}_s(\mathbf{k}, t)|$, $|\mathbb{U}_d(\mathbf{k}, t)|$ and $|\mathbb{U}_-(\mathbf{k}, t)|$ given respectively by Lemmas 4.5, 4.9, 4.12 and 4.13 (since for $|\mathbf{k}|$ large enough, $t \geq 0$ and $C_1, C_2 > 0$: $e^{-\delta t} \leq e^{-C_1 |\mathbf{k}|^{-2} t} \leq e^{-C_2 |\mathbf{k}|^{-4} t}$). \square

We prove in the following result that the estimates of Theorem 4.14 are optimal for an infinite family of well chosen initial conditions $\mathbf{U}_0 \in \mathcal{H}_{\perp, \text{HF}}^m$ for any fixed $m > 0$.

Theorem 4.15. *Let $m, k_+ > 0$ and $\phi : \mathbb{R}^+ \mapsto \mathbb{R}$ be any a measurable and bounded function satisfying:*

$$\text{supp } \phi \subset [k_+, +\infty[\text{ and } 0 < |\phi(|\mathbf{k}|)| \lesssim (1 + |\mathbf{k}|^2)^{-s/2} \text{ for } |\mathbf{k}| > k_+ + 1 \text{ and } s > \frac{3}{2} + m. \quad (4.98)$$

1. *If the Maxwell system is in a non-critical configuration, one defines (for k_+ sufficiently large) an initial condition \mathbf{U}_0 of (1.13) via its Fourier transform:*

$$\mathbb{U}_0(\mathbf{k}) = \mathcal{F}(\mathbf{U}_0)(\mathbf{k}) = \phi(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}}^* \frac{\mathcal{V}_{|\mathbf{k}|}(\omega_{+\infty}(|\mathbf{k}|)) \mathbf{e}_1}{|\mathcal{V}_{|\mathbf{k}|}(\omega_{+\infty}(|\mathbf{k}|)) \mathbf{e}_1}, \quad \forall \mathbf{k} \in \mathbb{R}^{3,*}. \quad (4.99)$$

Then, $\mathbf{U}_0 \in \mathcal{H}_{\perp, \text{HF}}^m$ and $\exists C, \tilde{C} > 0$ such that the associated solution \mathbf{U} of (1.13) satisfies

$$\tilde{C} e^{-\frac{Ct}{|\mathbf{k}|^2}} |\mathbb{U}_0(\mathbf{k})| \leq |\mathbb{U}(\mathbf{k}, t)|, \quad \forall t \geq 0 \text{ and } \forall \mathbf{k} \in \mathbb{R}^{3,*}. \quad (4.100)$$

2. If the Maxwell system is in a critical configuration, by Corollary 4.8, there exists (at least one) $p \in \mathcal{P}_s$ such that

$$\operatorname{Im} \omega_p(|\mathbf{k}|) = \operatorname{Im} A_{p,4} |\mathbf{k}|^{-4} + o(|\mathbf{k}|^{-4}), \text{ as } |\mathbf{k}| \mapsto +\infty, \text{ with } \operatorname{Im} A_{p,4} < 0. \quad (4.101)$$

Defining (for k_+ sufficiently large) as an initial condition \mathbf{U}_0 for (1.13) by

$$\mathbf{U}_0(\mathbf{k}) = \mathcal{F}(\mathbf{U}_0)(\mathbf{k}) = \phi(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}}^* \frac{\mathcal{V}_{|\mathbf{k}|}(\omega_p(|\mathbf{k}|)) \mathbf{e}_1}{|\mathcal{V}_{|\mathbf{k}|}(\omega_p(|\mathbf{k}|)) \mathbf{e}_1|}, \quad \forall \mathbf{k} \in \mathbb{R}^{3,*}, \quad (4.102)$$

then, $\mathbf{U}_0 \in \mathcal{H}_{\perp, \text{HF}}^m$ and $\exists C, \tilde{C} > 0$ such that the associated solution \mathbf{U} of (1.13) satisfy

$$\tilde{C} e^{-\frac{Ct}{|\mathbf{k}|^4}} |\mathbf{U}_0(\mathbf{k})| \leq |\mathbf{U}(\mathbf{k}, t)|, \quad \forall t \geq 0 \text{ and } \mathbf{k} \in \mathbb{R}^{3,*}. \quad (4.103)$$

In other words, the estimates (4.96) and (4.97) are optimal for an infinite family of solutions.

Proof. We separate it in two steps.

Step 1. Proof of the lower bound (4.100):

From Proposition 4.1, we know for $|\mathbf{k}| > k_+$, $|\mathbf{k}| \mapsto \omega_{+\infty}(|\mathbf{k}|)$ is well defined and satisfy (by Lemma 4.4) the asymptotic expansion (4.32).

Furthermore, the proof of Proposition 3.1 in appendix A.1.2, shows that any \mathbb{U} in the two-dimensional eigenspace $\ker(\mathbb{A}_{|\mathbf{k}|} - \omega_{+\infty}(|\mathbf{k}|) \operatorname{Id})$ is of the form $\mathbb{U} = (\mathbb{E}, \mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}})$ for some $\mathbb{E} \in \mathbf{C}_{\perp}$ and $(\mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}})$ deduced from \mathbb{E} by formula (A.16, A.17), with $\omega = \omega_{+\infty}(|\mathbf{k}|)$.

This can be expressed with the help of the operator $\mathcal{V}_{|\mathbf{k}|}(\omega) \in \mathcal{L}(\mathbf{C}_{\perp}, \mathbf{C}_{\perp}^N)$, see (3.11), as follows

$$\ker(\mathbb{A}_{|\mathbf{k}|} - \omega_{+\infty}(|\mathbf{k}|) \operatorname{Id}) = \mathcal{V}_{|\mathbf{k}|}(\omega_{+\infty}(|\mathbf{k}|))(\mathbf{C}_{\perp}). \quad (4.104)$$

Let $m \in \mathbb{N}^*$. Let us define \mathbf{U}_0 via its Fourier transform $\mathbb{U}_0(\mathbf{k})$ and formula (4.99) where $\mathcal{R}_{\mathbf{k}}$, see (2.8), is unitary. For $s > 3/2 + m$, $\mathbf{k} \mapsto (1 + |\mathbf{k}|^2)^{m/2} \mathbb{U}_0(\mathbf{k}) \in \mathbf{L}^2(\mathbb{R}^3)^N$, thus $\mathbf{U}_0 \in \mathbf{H}^m(\mathbb{R}^3)^N$ and therefore belongs by construction to $\mathcal{H}_{\perp, \text{HF}}^m$.

Next, thanks to formula (2.12) for $\mathbb{U}(\mathbf{k}, t)$, (4.99) and (4.104), one has

$$\mathbb{U}(\mathbf{k}, t) = e^{-i\mathbb{A}_{\mathbf{k}} t} \mathbb{U}_0(\mathbf{k}) = e^{-i\omega_{+\infty}(|\mathbf{k}|) t} \mathbb{U}_0(\mathbf{k}) \quad \text{and thus} \quad |\mathbb{U}(\mathbf{k}, t)| = e^{-\operatorname{Im} \omega_{+\infty}(|\mathbf{k}|) t} |\mathbb{U}_0(\mathbf{k})|.$$

Moreover, using Lemma 4.4 and more precisely the expansion (4.33) for $\omega_{+\infty}(|\mathbf{k}|)$, one knows that for large enough $|\mathbf{k}|$, with $A_{2,\infty} > 0$ defined by (4.34),

$$-\frac{A_{2,\infty}}{c^2 |\mathbf{k}|^2} \leq \operatorname{Im} \omega_{+\infty}(|\mathbf{k}|) < 0 \quad (\text{simply because } 1/3 < 1/2).$$

We thus get $e^{-\frac{A_{2,\infty} t}{c^2 |\mathbf{k}|^2}} |\mathbb{U}_0(\mathbf{k})| \lesssim |\mathbb{U}(\mathbf{k}, t)|$, which achieves the proof of (4.100).

Step 2 Proof of the lower bound (4.103):

If the Maxwell system is in the critical configuration, to show that the estimate (4.97), one only needs to define \mathbf{U}_0 by (4.102) (instead of (4.99)) replacing $\omega_{\infty}(|\mathbf{k}|)$ by $\omega_p(|\mathbf{k}|)$ satisfying (4.101) (instead of (4.35)). \square

5 Asymptotic analysis for small spatial frequencies $|\mathbf{k}| \ll 1$

As the reader can expect, the structure of this section is similar to that of Section 4 with three subsections:

- Section 5.1: Asymptotics of dispersion curves for $0 < |\mathbf{k}| \ll 1$.
- Section 5.2: Spectral decomposition of the solution for $0 < |\mathbf{k}| \ll 1$.
- Section 5.3: Large time estimate of $\mathbb{U}(\mathbf{k}, t)$ for $0 < |\mathbf{k}| \ll 1$.

This section is however shorter than Section 4 because much less particular cases appear.

5.1 Asymptotics of the dispersion curves for $|\mathbf{k}| \ll 1$

This section is the counterpart of Section 4.1 for low spatial frequencies. In other words, we focus here on long time estimates of the low (spatial) frequency components $\mathbb{U}(\mathbf{k}, t)$ (see (2.7)) of the solution. As explained in Section 3.3, the decay of $\mathbb{U}(\mathbf{k}, t)$ for $|\mathbf{k}| \ll 1$ is related to the analysis of the solutions of the dispersion relation (3.13) at low frequencies. Roughly speaking, as $|\mathbf{k}|^2 \rightarrow 0$ when $|\mathbf{k}| \rightarrow 0$, the solutions of (3.13) must satisfy $|\mathcal{D}(\omega)| \rightarrow 0$ as $|\mathbf{k}| \rightarrow 0$. Thus, they converge to a zero $z \in \mathcal{Z} \cup \{0\}$ of the rational function \mathcal{D} . Thus, as one can rewrite \mathcal{D} in the vicinity of a zero z of multiplicity \mathbf{m}_z as

$$\mathcal{D}(\omega) = (\omega - z)^{\mathbf{m}_z} g(\omega) \text{ with } g \text{ analytic in a vicinity of } z \text{ and } g(z) = A_z \neq 0. \quad (5.1)$$

This leads to the following proposition (the equivalent of Proposition 4.1 in Section 4.1).

Proposition 5.1. *There exists $k_- > 0$ such that for $0 < |\mathbf{k}| \leq k_-$, the solutions of the rational dispersion relation (3.13) (or of its equivalent polynomial form (3.19)) are all simple. These solutions form N distinct branches which are C^∞ -smooth functions of $|\mathbf{k}| \in (0, k_-]$. These branches are characterized by their asymptotic expansion for small $|\mathbf{k}|$. More precisely, if $z \in \mathcal{Z} \cup \{0\}$ is a zero of multiplicity \mathbf{m}_z of \mathcal{D} , there exists \mathbf{m}_z distinct branches of solutions $\omega_{z,n}$, $n = 1, \dots, \mathbf{m}_z$ of (3.13) satisfying*

$$\omega_{z,n}(|\mathbf{k}|) = z + a_{z,n}^{-1} |\mathbf{k}|^{\frac{2}{\mathbf{m}_z}} (1 + o(1)), \quad a_{z,n} = |A_z|^{1/\mathbf{m}_z} e^{i\frac{\theta_z}{\mathbf{m}_z}} e^{\frac{2in\pi}{\mathbf{m}_z}} \quad (|\mathbf{k}| \rightarrow 0), \quad (5.2)$$

where A_z is defined in (5.1) and $\theta_z \in (-\pi, \pi]$ is the principal argument of A_z .

Proof. Step 1: Construction of the branches of solutions $\omega_{z,n}(|\mathbf{k}|)$.

Let $z \in \mathcal{Z} \cup \{0\}$ a zero of multiplicity \mathbf{m}_z . Then, the rational function \mathcal{D} can be factorized as in (5.1). Moreover, the dispersion relation can be rewritten as

$$\mathcal{D}(\omega) = |\mathbf{k}|^2 = (|\mathbf{k}|^{\frac{2}{\mathbf{m}_z}})^{\mathbf{m}_z}.$$

Thus, by applying the Lemma A.1 of the appendix A.2 with

$$\mathcal{G}(\omega) = \mathcal{D}(\omega), \quad \mathbf{m} = \mathbf{m}_z, \quad A = A_z \text{ and } \zeta = |\mathbf{k}|^{2/\mathbf{m}_z},$$

we deduce the existence of $k_- > 0$ and \mathbf{m}_z distinct C^∞ branches of solutions $|\mathbf{k}| \in (0, k_-] \mapsto \omega_{z,n}(|\mathbf{k}|)$ of the dispersion $\mathcal{D}(\omega) = |\mathbf{k}|^2$. Furthermore, for the solutions $\omega_{z,n}(|\mathbf{k}|)$, we can use the

asymptotic expansion (A.22) of Lemma A.1 with $\zeta = |\mathbf{k}|^{2/m_z}$ which gives (5.2).

Step 2: Conclusion. From the asymptotic (5.2), $\omega_{z,n}(|\mathbf{k}|)$ for $z \in \mathcal{Z} \cup \{0\}$ and $n \in \{1, \dots, \mathbf{m}_z\}$ are all distinct for $|\mathbf{k}|$ small enough and positive. Since the sum of the \mathbf{m}_z 's, for $z \in \mathcal{Z} \cup \{0\}$ is equal to $2N_e + 2N_m + 2$, we have constructed $2N_e + 2N_m + 2$ distinct solutions of (3.13). As (3.13) is equivalent to a polynomial equation of degree $2N_e + 2N_m + 2$, cf. (3.19), these solutions are simple roots and there are no other solutions for small positive $|\mathbf{k}|$. \square

5.2 Spectral decomposition of the solution for $|\mathbf{k}| \ll 1$

Our goal here is to obtain for small $|\mathbf{k}|$ a decomposition of $\mathbb{U}(\mathbf{k}, t)$ similar to the one obtained for large $|\mathbf{k}|$ (see (4.9,4.10)).

For this, we first notice that, combining Proposition 5.1 and Corollary 3.4 immediately yields the following property of the operator $\mathbb{A}_{|\mathbf{k}|,\perp}$.

Corollary 5.2. *There exists $k_- > 0$ such that for $0 < |\mathbf{k}| \leq k_-$, $\mathbb{A}_{|\mathbf{k}|,\perp}$ is diagonalizable on \mathbf{C}_{\perp}^N .*

To obtain the decomposition of $\mathbb{U}(\mathbf{k}, t)$, we split (see Section 3.2.3) the sets of zeros $\{0\} \cup \mathcal{Z} \subset \overline{\mathbb{C}^-}$ of \mathcal{D} in three disjoint subsets : $\{0\} \cup \mathcal{Z} = \{0\} \cup \mathcal{Z}_s \cup \mathcal{Z}_-$ where the multiplicity of $z = 0$ is $\mathbf{m}_0 = 2$ and

$$\mathcal{Z}_- := \mathcal{Z} \cap \mathbb{C}^- \text{ and } \mathcal{Z}_s := \{z \in \mathcal{Z} \cap \mathbb{R} \mid \mathbf{m}_z = 1\}. \quad (5.3)$$

Remark 5.3. *[On the set \mathcal{Z}_-] We point out that the weak dissipation condition (1.11) implies (see (3.18) and (3.17)) that at least one element of \mathcal{Z} lies in \mathbb{C}^- , thus $\mathcal{Z}_- \neq \emptyset$.*

The strong dissipation condition (1.9) implies that all elements of \mathcal{Z} lie in \mathbb{C}^- , that is $\mathcal{Z}_- = \mathcal{Z}$.

Remark 5.4. *[On the set \mathcal{Z}_s] According to the analysis of the zeros of made in Section 3.2.3, the structure of the set \mathcal{Z}_s strongly relies on the condition*

$$\exists (j_0, \ell_0) \text{ such that } \alpha_{e,j_0} > 0 \text{ and } \alpha_{m,\ell_0} > 0. \quad (5.4)$$

More precisely, if (5.4) holds, $\mathcal{Z}_s = \emptyset$. On the contrary, under the weak dissipation condition (1.11), if (5.4) does not hold, either all $\alpha_{e,j}$ vanish in which case $\mathcal{Z}_s = \mathcal{Z}_e$ either all $\alpha_{\ell,m}$ vanish and $\mathcal{Z}_s = \mathcal{Z}_m$.

According to Propositions 3.1 and 5.1, we have the corresponding partition of the spectrum of $\mathbb{A}_{|\mathbf{k}|,\perp}$ for $0 < |\mathbf{k}| \leq k_-$:

$$\sigma(\mathbb{A}_{|\mathbf{k}|,\perp}) = \{\omega_{0,r}(|\mathbf{k}|), r = 1, 2\} \cup \{\omega_z(|\mathbf{k}|), z \in \mathcal{Z}_s\} \cup \{\omega_{z,n}(|\mathbf{k}|), z \in \mathcal{Z}_-, n = 1, \dots, \mathbf{m}_z\}, \quad (5.5)$$

where, for simplicity, we write $\omega_z(|\mathbf{k}|)$ instead of $\omega_{z,1}(|\mathbf{k}|)$ for the simple zeros $z \in \mathcal{Z}_s$. We point out that “by symmetry” of the dispersion relation $\omega_{0,2}(|\mathbf{k}|) = -\bar{\omega}_{0,1}(|\mathbf{k}|)$.

For a sketch of the small $|\mathbf{k}|$ behaviour of the dispersion curves, see figure 4.

For $0 < |\mathbf{k}| \leq k_-$, as $\mathbb{A}_{|\mathbf{k}|,\perp}$ is diagonalizable by Corollary 5.2, \mathbf{C}_{\perp}^N can be decomposed as

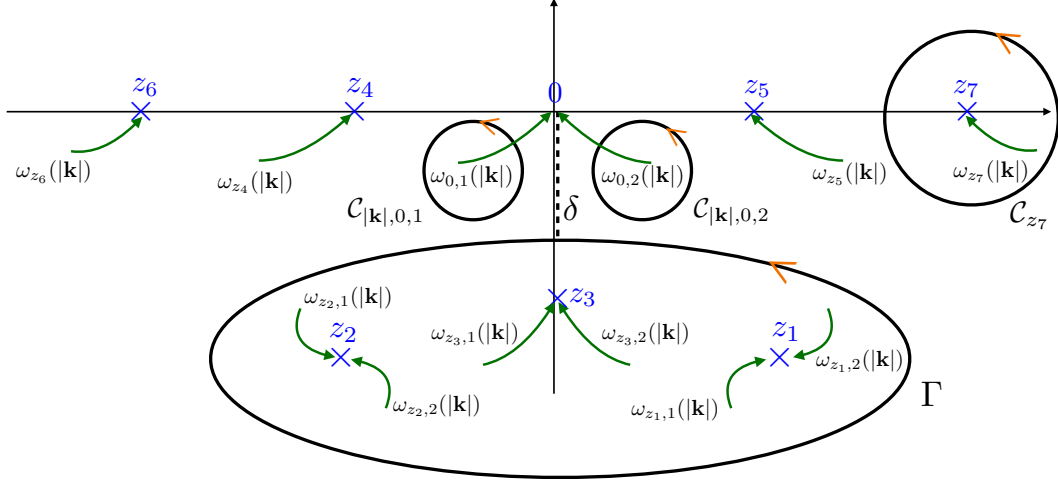


Figure 4: Sketch of a configuration of dispersion curves for small values of $|\mathbf{k}|$ in the case where $\{0\} \cup \mathcal{Z} = \{0\} \cup \mathcal{Z}_- \cup \mathcal{Z}_s$ with $\mathcal{Z}_- = \{z_1, z_2, z_3\}$ and $\mathcal{Z}_s = \{z_4, z_5, z_6, z_7\}$.

$$\mathbf{C}_{\perp}^N = \bigoplus_{r=1}^2 V_{|\mathbf{k}|,0,r} \bigoplus_{z \in \mathcal{Z}_s} V_{|\mathbf{k}|,z} \bigoplus \bigoplus_{z \in \mathcal{Z}_-} \bigoplus_{n=1}^{m_z} V_{|\mathbf{k}|,z,n} \quad (5.6)$$

where we have defined the 2D (see below) subspaces

$$\begin{cases} V_{|\mathbf{k}|,0,r} = \ker (\mathbb{A}_{|\mathbf{k}|,\perp} - \omega_{0,r}(|\mathbf{k}|) \text{Id}), \quad r = 1, 2 & V_{|\mathbf{k}|,z} = \ker (\mathbb{A}_{|\mathbf{k}|,\perp} - \omega_z(|\mathbf{k}|) \text{Id}), \\ V_{|\mathbf{k}|,z,n} = \ker (\mathbb{A}_{|\mathbf{k}|,\perp} - \omega_{z,n}(|\mathbf{k}|) \text{Id}), \quad 1 \leq n \leq m_z, \end{cases} \quad (5.7)$$

where the above direct sums are (in general) non-orthogonal. Thus, one decomposes uniquely any vector $x \in \mathbf{C}_{\perp}^N$ as

$$\begin{cases} x = \sum_{r=1}^2 x_{|\mathbf{k}|,0,r} + \sum_{z \in \mathcal{Z}_s} x_{|\mathbf{k}|,z} + \sum_{z \in \mathcal{Z}_-} \sum_{n=1}^{m_z} x_{|\mathbf{k}|,z,n} \\ x_{|\mathbf{k}|,0,r} \in V_{|\mathbf{k}|,0,r}, \quad x_{|\mathbf{k}|,z} \in V_{|\mathbf{k}|,z} \quad \text{and} \quad x_{|\mathbf{k}|,z,n} \in V_{|\mathbf{k}|,z,n}. \end{cases}$$

Then, we define the spectral projectors $\Pi_{0,r}(|\mathbf{k}|)$ for $r \in \{1, 2\}$, $\Pi_z(|\mathbf{k}|)$ for $z \in \mathcal{Z}_s$ and $\Pi_{p,n}(|\mathbf{k}|)$, $z \in \mathcal{Z}_-$ and $n \in \{1, \dots, m_z\}$ associated to the eigenvalues $\omega_{0,r}(|\mathbf{k}|)$, $\omega_z(|\mathbf{k}|)$ and $\omega_{z,n}(|\mathbf{k}|)$ by:

$$\Pi_{0,r}(|\mathbf{k}|)(x) = x_{|\mathbf{k}|,0,r}, \quad \Pi_z(|\mathbf{k}|)(x) = x_{|\mathbf{k}|,z} \quad \text{and} \quad \Pi_{z,n}(|\mathbf{k}|)(x) = x_{|\mathbf{k}|,z,n}. \quad (5.8)$$

From Proposition 3.1, the geometric multiplicity of each eigenvalues in $\sigma(\mathbb{A}_{|\mathbf{k}|,\perp})$ is two. Thus, all these projectors are rank two projectors. We again emphasize that the dissipative operator $\mathbb{A}_{|\mathbf{k}|,\perp}$ is not normal, thus its spectral projectors are not orthogonal.

For $0 < |\mathbf{k}| \leq k_-$, as $\mathbb{A}_{|\mathbf{k}|,\perp}$ is diagonalizable (by Corollary 3.4), one has

$$\mathbb{A}_{|\mathbf{k}|,\perp} = \sum_{r=1}^2 \omega_{0,r}(|\mathbf{k}|) \Pi_{0,r}(|\mathbf{k}|) + \sum_{z \in \mathcal{Z}_s} \omega_z(|\mathbf{k}|) \Pi_z(|\mathbf{k}|) + \sum_{z \in \mathcal{Z}_-} \sum_{n=1}^{m_z} \omega_{z,n}(|\mathbf{k}|) \Pi_{z,n}(|\mathbf{k}|). \quad (5.9)$$

Thus, for $0 < |\mathbf{k}| \leq k_-$, the solution $\mathbb{U}(\mathbf{k}, t)$ given by (2.12) can be expressed for all $t \geq 0$ as

$$\mathbb{U}(\mathbf{k}, t) = \mathbb{U}_{z,0}(\mathbf{k}, t) + \mathbb{U}_{z,s}(\mathbf{k}, t) + \mathbb{U}_{z,-}(\mathbf{k}, t) \quad (5.10)$$

where

$$\begin{cases} \mathbb{U}_{z,0}(\mathbf{k}, t) = \sum_{r=1}^2 e^{-i\omega_{0,r}(|\mathbf{k}|)t} \mathcal{R}_{\mathbf{k}}^* \Pi_{0,r}(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}} \mathbb{U}_0(\mathbf{k}), & (i) \\ \mathbb{U}_{z,s}(\mathbf{k}, t) = \sum_{z \in \mathcal{Z}_s} e^{-i\omega_z(|\mathbf{k}|)t} \mathcal{R}_{\mathbf{k}}^* \Pi_z(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}} \mathbb{U}_0(\mathbf{k}), & (ii) \\ \mathbb{U}_{z,-}(\mathbf{k}, t) = \sum_{z \in \mathcal{Z}_-} \sum_{n=1}^{m_z} e^{-i\omega_{z,n}(|\mathbf{k}|)t} \mathcal{R}_{\mathbf{k}}^* \Pi_{z,n}(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}} \mathbb{U}_0(\mathbf{k}). & (iii) \end{cases} \quad (5.11)$$

5.3 Estimates of the low frequency components of the solution

5.3.1 Orientation

Proceeding as for large $|\mathbf{k}|$ in Section 4.3, we estimate successively, in Lemmas 5.7, 5.10 and 5.11, the three terms appearing in the decomposition (5.10). An important difference with Section 4.3 is that we consider bounded values of $|\mathbf{k}|$, namely $0 < |\mathbf{k}| < k_-$, for some $k_- > 0$, which will simplify some estimates via compactness arguments.

Two cases will have to be distinguished.

Case 1: estimates of $\mathbb{U}_{z,0}(\mathbf{k}, t)$ and $\mathbb{U}_{z,s}(\mathbf{k}, t)$. This will be established in Sections 5.3.2 and 5.3.3, by bounding separately each of the terms of the sums in (5.11)(i) and (5.11)(ii).

For each eigenvalue $\omega(|\mathbf{k}|) \in \{\omega_{0,r}(|\mathbf{k}|), \omega_z(|\mathbf{k}|)\}$, we shall first estimate in Lemmas 5.5 and 5.8 the corresponding spectral projector $\Pi(|\mathbf{k}|)$ uniformly in $|\mathbf{k}|$ for $|\mathbf{k}|$ small enough and positive. To this aim as in Section 4.3, we use the Dunford-Riesz functional calculus that says that, if $\Gamma_{|\mathbf{k}|}$ is a (positively oriented) close contour in \mathbb{C} that encloses $\omega(|\mathbf{k}|)$ but no other eigenvalue of $\mathbb{A}_{|\mathbf{k}|,\perp}$, so that

$$\Pi(|\mathbf{k}|) = -\frac{1}{2i\pi} \int_{\Gamma_{|\mathbf{k}|}} R_{|\mathbf{k}|}(\omega) d\omega. \quad (5.12)$$

For $\Pi_z(|\mathbf{k}|)$ in Lemma 5.8, $\omega_z(|\mathbf{k}|)$ is a simple eigenvalues that converges to z while remaining well separated from the other eigenvalues. Hence, using Proposition 3.2, we can choose for $\Gamma_{|\mathbf{k}|}$ a fixed circle Γ (i.e. independent of $|\mathbf{k}|$), along which the operator $\mathcal{R}_{|\mathbf{k}|}(\omega)$, by continuity in $(|\mathbf{k}|, \omega)$ and compactness of $[0, k_-] \times \Gamma$, remain bounded. Thus, the uniform estimate of $\Pi_z(|\mathbf{k}|)$ becomes obvious.

For the projector $\Pi_{0,1}(|\mathbf{k}|)$ (this is the same for $\Pi_{0,2}(|\mathbf{k}|)$), this is more complicated because the two distinct eigenvalues $\omega_{0,r}(|\mathbf{k}|)$ ($r = 1, 2$) converge both to 0 while remaining far away from the other eigenvalues and any other point of the set $\mathcal{S}_{\mathcal{T}}$ than 0. For this reason, we shall take as $\Gamma_{|\mathbf{k}|}$ the circle $\mathcal{C}_{0,1,|\mathbf{k}|}$ centered at $\omega_{0,1}(|\mathbf{k}|)$ with radius $\rho_{|\mathbf{k}|} = |\omega_{0,2}(|\mathbf{k}|) - \omega_{0,1}(|\mathbf{k}|)|/4$ in order to prevent $\omega_{0,2}(|\mathbf{k}|)$ from being inside $\mathcal{C}_{0,1,|\mathbf{k}|}$. This circle will also not contain 0 or any other point of $\mathcal{S}_{\mathcal{T}}$ (by Proposition 5.1). As a consequence, in (5.12), $R_{|\mathbf{k}|}(\omega)$ can be replaced by the product $\mathcal{V}_{|\mathbf{k}|}(\omega)\mathcal{S}_{|\mathbf{k}|}(\omega)$, (see formula (3.9) for $R_{|\mathbf{k}|}(\omega)$ and the definition (3.8) of $\mathcal{S}_{\mathcal{T}}$) leading, similarly to (4.15), to

$$\|\Pi(|\mathbf{k}|)\| \leq \rho_{|\mathbf{k}|} \sup_{\omega \in \mathcal{C}_{|\mathbf{k}|}} (\|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \|\mathcal{S}_{|\mathbf{k}|}(\omega)\|). \quad (5.13)$$

Moreover the product $\|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \|\mathcal{S}_{|\mathbf{k}|}(\omega)\|$ will blow up when $|\mathbf{k}|$ tend to 0 but this will be compensated by the fact that the radius of $\rho_{|\mathbf{k}|}$ tends to 0.

In a second step, we concentrate on the exponentials appearing in each factor of the terms involved in $\mathbb{U}_{z,0}(\mathbf{k}, t)$ and $\mathbb{U}_{z,s}(\mathbf{k}, t)$ whose estimates for small $|\mathbf{k}|$ rely on the asymptotic expansion of the eigenvalues (and more particularly their imaginary parts), see Lemmas 5.6 and 5.9.

Case 2: estimate of $\mathbb{U}_{z,-}(\mathbf{k}, t)$. This term will be treated as the term $\mathbb{U}_{-}(\mathbf{k}, t)$ (in Section 4.3.5), because the set $\{\omega_{z,n}(|\mathbf{k}|)\}$ will remain far from the real axis for small enough $|\mathbf{k}|$. Thanks to this property, the exponential decay of $|\mathbb{U}_{z,-}(\mathbf{k}, t)|$ will be, contrarily to the previous terms $\mathbb{U}_{z,0}(\mathbf{k}, t)$ and $\mathbb{U}_{z,s}(\mathbf{k}, t)$, uniform with respect to $|\mathbf{k}|$, so that it will not contribute in the end to the large time equivalent of $\mathbf{U}(\cdot, t)$. For proving this, we shall not use (5.11)(ii) but an alternative formula directly issued from the Riesz-Dunford functional calculus. Compared to Section 4.3.5, its estimate in Lemma 5.11 will be simplified by a compactness argument.

5.3.2 Estimates of $\mathbb{U}_{z,0}(\mathbf{k}, t)$ for $0 < |\mathbf{k}| \leq k_-$

In the following Lemma, we estimate the spectral projectors $\Pi_{0,r}(|\mathbf{k}|)$ (for $r = 1, 2$) using (5.13) with a simple closed contour $\Gamma_{|\mathbf{k}|}$ satisfying the two following properties:

$$\begin{cases} (i) & \Gamma_{|\mathbf{k}|} \text{ encloses } \omega(|\mathbf{k}|) \text{ but no other eigenvalue of } \mathbb{A}_{|\mathbf{k}|,\perp}, \\ (ii) & \Gamma_{|\mathbf{k}|} \text{ does not enclose any point of the set } \mathcal{S}_{\mathcal{T}}. \end{cases} \quad (5.14)$$

Lemma 5.5. *There exists $k_- > 0$ such that the spectral projectors $\Pi_{0,r}(|\mathbf{k}|)$, $r \in \{1, 2\}$ are uniformly bounded for $0 < |\mathbf{k}| \leq k_-$.*

Proof. Let $r \in \{1, 2\}$. We follow the approach of the Section 4.3 and introduce as $\Gamma_{|\mathbf{k}|}$ the circle $\mathcal{C}_{0,r,|\mathbf{k}|}$ (positively oriented) of center $\omega_{0,r}(|\mathbf{k}|)$ and radius

$$\rho_{|\mathbf{k}|} = \frac{1}{4} |\omega_{0,1}(|\mathbf{k}|) - \omega_{0,2}(|\mathbf{k}|)|, \quad (5.15)$$

(see figure 4 for an example of contour $\mathcal{C}_{0,r,|\mathbf{k}|}$ with $r = 1, 2$.)

From the asymptotics (5.2) applied with $\mathbf{m}_0 = 2$ ($z = 0$ is a double zero), one deduces that

$$|\omega_{0,1}(|\mathbf{k}|)| \sim |A_0|^{\frac{1}{2}} |\mathbf{k}|, \quad |\omega_{0,1}(|\mathbf{k}|) - \omega_{0,2}(|\mathbf{k}|)| \sim 2 |A_0|^{\frac{1}{2}} |\mathbf{k}|, \quad (|\mathbf{k}| \rightarrow 0) \quad (5.16)$$

so that, for $|\mathbf{k}|$ small enough, $\Gamma_{|\mathbf{k}|}$ does not enclose $\omega_{0,2}(|\mathbf{k}|)$ nor 0 (and of course any other point of \mathcal{S}_T). As explained in Section 5.3.1, we can apply the inequality (5.13). As $\rho_{|\mathbf{k}|} \lesssim |\mathbf{k}|$, according to (5.15) and (5.16), we thus obtain

$$\|\Pi_{0,r}(|\mathbf{k}|)\| \lesssim |\mathbf{k}| \sup_{\omega \in \mathcal{C}_{0,r,|\mathbf{k}|}} (\|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \|\mathcal{S}_{|\mathbf{k}|}(\omega)\|). \quad (5.17)$$

It remains to estimate the operators $\mathcal{V}_{|\mathbf{k}|}(\omega)$ and $\mathcal{S}_{|\mathbf{k}|}(\omega)$ given by the formulas (3.10) and (3.11). Without any loss of generality, we can restrict ourselves to $r = 1$ for the rest of the proof.

Step 1: Estimate of $(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1}$.

This term is involved in the expression (3.10) of $\mathcal{S}_{|\mathbf{k}|}(\omega)$. Let us recall that (see (4.17))

$$(\mathcal{D}(\omega) - |\mathbf{k}|^2)^{-1} = Q_e(\omega) Q_m(\omega) D_{|\mathbf{k}|}(\omega)^{-1}.$$

where $D_{|\mathbf{k}|}(\omega)$ is the polynomial defined by (3.19). Using the proposition 5.1, we know that $D_{|\mathbf{k}|}(\omega)$ can be factorized as the following for $|\mathbf{k}|$ sufficiently small and positive:

$$D_{|\mathbf{k}|}(\omega) = \varepsilon_0 \mu_0 (\omega - \omega_{0,1}(|\mathbf{k}|)) (\omega - \omega_{0,2}(|\mathbf{k}|)) \prod_{z \in \mathcal{Z}_s} (\omega - \omega_z(|\mathbf{k}|)) \prod_{z \in \mathcal{Z}_-} \prod_{n=1}^{m_z} (\omega - \omega_{z,n}(|\mathbf{k}|)). \quad (5.18)$$

Since $\omega_z(|\mathbf{k}|), z \in \mathcal{Z}_s$ or $\omega_{z,n}(|\mathbf{k}|), z \in \mathcal{Z}_-$ converge to $z \neq 0$, cf. (5.2), the last two products in (5.18) remain bounded from below in modulus when ω describes $\mathcal{C}_{p,r,|\mathbf{k}|}$ and $|\mathbf{k}|$ is small enough.

On the other hand, it is easy to see that for $|\mathbf{k}|$ small enough ($\mathcal{C}_{0,1,|\mathbf{k}|}$ has been chosen for that purpose),

$$\forall \omega \in \mathcal{C}_{0,1,|\mathbf{k}|}, \quad |\omega - \omega_{0,1}(|\mathbf{k}|)| \geq C |\mathbf{k}|, \quad |\omega - \omega_{0,2}(|\mathbf{k}|)| \geq C |\mathbf{k}|,$$

Thus from (5.18), we infer that

$$\forall \omega \in \mathcal{C}_{0,1,|\mathbf{k}|}, \quad |D_{|\mathbf{k}|}(\omega)|^{-1} \lesssim |\mathbf{k}|^{-2}. \quad (5.19)$$

Thus, as the polynomial $Q_e(\omega)Q_m(\omega)$ remains obviously bounded for bounded ω , we deduce that there exists $k_- > 0$ such that

$$\forall \omega \in \mathcal{C}_{0,1,|\mathbf{k}|}, \quad |\mathcal{D}(\omega) - |\mathbf{k}|^2|^{-1} \lesssim |\mathbf{k}|^{-2}, \quad \text{for } 0 < |\mathbf{k}| \leq k_-. \quad (5.20)$$

Step 2: Estimates of $\mathcal{S}_{|\mathbf{k}|}(\omega)$ and $\mathcal{V}_{|\mathbf{k}|}(\omega)$

First, it follows from their definitions (3.6) that $\mathbb{A}_{e,j}(\omega), \dot{\mathbb{A}}_{e,j}(\omega), \mathbb{A}_{m,\ell}(\omega)$ and $\dot{\mathbb{A}}_{m,\ell}(\omega)$ are uniformly bounded in operator norm when ω does not approach 0 (indeed 0 is not a zero of the polynomials $q_{e,j}$ and $q_{m,\ell}$). Thus $\|\mathbb{A}_e(\omega)\| \lesssim 1$ and $\|\mathbb{A}_m(\omega)\| \lesssim 1$ according to (3.7).

Moreover, $\mu(\omega) = \mu(0) + o(1)$ as $\omega \rightarrow 0$ and $|\omega| \lesssim |\mathbf{k}|$ on $\mathcal{C}_{0,1,|\mathbf{k}|}$, thus, for $|\mathbf{k}|$ small enough,

$$\forall \omega \in \mathcal{C}_{0,1,|\mathbf{k}|}, \quad \|-\omega\mu(\omega)\mathbb{A}_e(\omega) + |\mathbf{k}| \mathbf{e}_3 \times \mathbb{A}_m(\omega)\| \lesssim |\omega| + |\mathbf{k}| \lesssim |\mathbf{k}|. \quad (5.21)$$

Combining (5.20) and (5.21) in (3.10) yields that there exists $k_- > 0$:

$$\forall \omega \in \mathcal{C}_{0,1,|\mathbf{k}|}, \quad \|\mathcal{S}_{|\mathbf{k}|}(\omega)\| \lesssim |\mathbf{k}|^{-1}, \quad \text{for } 0 < |\mathbf{k}| \leq k_-. \quad (5.22)$$

Finally, as, along $\mathcal{C}_{0,1,|\mathbf{k}|}$, $|\omega| \geq |\rho_{|\mathbf{k}|,1} - |\omega_{0,1}(|\mathbf{k}|)|| \sim \frac{1}{2} |A_0|^{\frac{1}{2}} |\mathbf{k}|$ as $|\mathbf{k}| \rightarrow 0$, we deduce that

$$\forall \omega \in \mathcal{C}_{0,1,|\mathbf{k}|}, \quad \frac{|\mathbf{k}|}{|\omega| |\mu(\omega)|} \lesssim \frac{|\mathbf{k}|}{|\omega|} \lesssim 1, \quad \text{for } 0 < |\mathbf{k}| \leq k_-.$$

Since $|q_{e,j}(\omega)|^{-1}$ and $|q_{m,\ell}(\omega)|^{-1}$ remain bounded along $\mathcal{C}_{0,1,|\mathbf{k}|}$, by the formula (3.11) for $\mathcal{V}_{|\mathbf{k}|}(\omega)$,

$$\|\mathcal{V}_{|\mathbf{k}|}(\omega)\| \lesssim 1, \quad \forall \omega \in \mathcal{C}_{0,1,|\mathbf{k}|} \quad \text{and } 0 < |\mathbf{k}| \leq k_-. \quad (5.23)$$

Conclusion: Combining (5.22) and (5.23) in (5.17) finally implies that there exists $k_- > 0$ such that $\Pi_{0,1}(|\mathbf{k}|)$ is uniformly bounded for $0 < |\mathbf{k}| \leq k_-$. \square

The next lemma is about the asymptotic expansion (in powers of $|\mathbf{k}|$) of $\omega_{0,r}(|\mathbf{k}|)$ when $|\mathbf{k}| \rightarrow 0$. It is important to push the expansion up to the first appearance of a negative imaginary part, since it will govern the decay of $\mathbb{U}_{z,0}(\mathbf{k}, t)$ for small values of $|\mathbf{k}|$.

Lemma 5.6. *For $r \in \{1, 2\}$, the eigenvalue $\omega_{0,r}(|\mathbf{k}|)$ satisfies the following asymptotic expansions, with $c_0 = (\varepsilon(0)\mu(0))^{-1/2} > 0$,*

$$\omega_{z,r}(|\mathbf{k}|) = (-1)^r c_0 |\mathbf{k}| - \frac{1}{2} (\varepsilon\mu)'(0) c_0^4 |\mathbf{k}|^2 + o(|\mathbf{k}|^2), \quad \text{as } |\mathbf{k}| \rightarrow 0, \quad (5.24)$$

with moreover

$$\text{Im} \left(-(\varepsilon\mu)'(0) \right) < 0. \quad (5.25)$$

Proof. Using (3.21), one has $\mathcal{D}(\omega) = \omega^2 g(\omega)$ where the rational function $g = \varepsilon(\cdot) \mu(\cdot)$ is analytic in the vicinity of 0 and satisfies $g(0) = \varepsilon(0)\mu(0) > 0$. Thus applying Lemma A.1 with $z = 0$, $\mathbf{m}_0 = 2$, $\zeta = |\mathbf{k}|$, $a_1 = -g(0)^{1/2} = -c_0^{-1}$ and $a_2 = c_0^{-1}$, the asymptotic formula (A.22) yields:

$$\omega_{z,r}(|\mathbf{k}|) = a_r^{-1} |\mathbf{k}|^1 - \frac{g'(0)}{2g(0)^2} |\mathbf{k}|^2 + o(|\mathbf{k}|^2), \quad \text{as } |\mathbf{k}| \rightarrow 0.$$

This is precisely the asymptotic (5.24). It remains to prove (5.25).

As $\mu(0) > 0$ and $\varepsilon(0) > 0$, it is equivalent to show that $\text{Im } g'(0) = \mu(0) \text{Im } \varepsilon'(0) + \varepsilon(0) \text{Im } \mu'(0)$ is positive. The Taylor expansion of the rational functions ε and μ (see (1.6) and (1.7)) gives

$$\text{Im } \varepsilon'(0) = \varepsilon_0 \sum_{j=1}^{N_e} \frac{\alpha_{e,j} \Omega_{e,j}^2}{\omega_{e,j}^4} \geq 0 \quad \text{and} \quad \text{Im } \mu'(0) = \mu_0 \sum_{\ell=1}^{N_m} \frac{\alpha_{m,\ell} \Omega_{m,\ell}^2}{\omega_{m,\ell}^4} \geq 0.$$

Furthermore, by the weak dissipation condition (1.11) at least one coefficient $\alpha_{e,j}$ or $\alpha_{e,m}$ is positive, thus one has $\text{Im } \varepsilon'(0) > 0$ or $\text{Im } \mu'(0) > 0$ and we can conclude. \square

Combining the Lemma 5.5 and Lemma 5.6 (formula (5.24) and (5.25)) and proceeding as in the proof of Lemma 4.5 gives the following estimate for $\mathbb{U}_{z,0}(\mathbf{k}, t)$.

Lemma 5.7. *There exists $k_- > 0$ and $C > 0$ such that the function $\mathbb{U}_{z,0}(\mathbf{k}, t)$ defined by (5.10) and (5.11), satisfies the following estimate*

$$|\mathbb{U}_{z,0}(\mathbf{k}, t)| \lesssim e^{-C|\mathbf{k}|^2 t} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0 \quad \text{and } 0 < |\mathbf{k}| \leq k_-. \quad (5.26)$$

5.3.3 Estimates of $\mathbb{U}_{z,s}(\mathbf{k}, t)$ for $0 < |\mathbf{k}| \leq k_-$

This time we estimate $\mathbb{U}_{z,s}(\mathbf{k}, t)$ in (5.10) which involves in particular the projectors $\Pi_z(|\mathbf{k}|)$.

Lemma 5.8. *There exists $k_- > 0$ such that the spectral projectors $\Pi_z(|\mathbf{k}|)$, $z \in \mathcal{Z}_s$ are uniformly bounded for $0 < |\mathbf{k}| \leq k_-$.*

Proof. The proof of this lemma has already been sketched in Section 5.3.1 (paragraph about $\Pi_z(|\mathbf{k}|)$). The details are left to the reader. \square

We now give the asymptotic expansion of the eigenvalues $\omega_z(|\mathbf{k}|)$ for small $|\mathbf{k}|$.

Lemma 5.9. *Let $z \in \mathcal{Z}_s$. The eigenvalue $\omega_z(|\mathbf{k}|)$ satisfies the following asymptotic expansion*

$$\omega_z(|\mathbf{k}|) = z + A_z |\mathbf{k}|^2 + o(|\mathbf{k}|^2), \quad \text{as } |\mathbf{k}| \rightarrow 0, \quad (5.27)$$

where the complex number A_z is defined by (two disjoint cases have to be distinguished):

$$\text{if } z \in \mathcal{Z}_e, \quad A_z = \frac{(\omega\varepsilon)'(z)^{-1}}{z\mu(z)}, \quad \text{if } z \in \mathcal{Z}_m, \quad A_z = \frac{(\omega\mu)'(z)^{-1}}{z\varepsilon(z)}, \quad (5.28)$$

and satisfies in all cases $\text{Im } A_z < 0$.

Proof. The elements of \mathcal{Z}_s are simple zeroes ($m_z = 1$) of the rational function \mathcal{D} , it means that either $z \in \mathcal{Z}_e$ or $z \in \mathcal{Z}_m$. We will prove the Lemma by assuming that $z \in \mathcal{Z}_e$, the proof for $z \in \mathcal{Z}_m$ can be done with obvious ‘‘symmetric arguments’’.

Using Proposition 5.1 and more precisely the asymptotic expansion (5.2) (with here $m_z = 1$) yields (5.27) with $A_z = g(z)^{-1}$ where g is defined by (5.1), that is to say

$$g(\omega) = \frac{\mathcal{D}(\omega)}{\omega - z} = \frac{\omega\varepsilon(\omega)}{\omega - z} \omega\mu(\omega) \quad \text{for } \omega \in \mathbb{C} \setminus (\mathcal{P} \cup \{z\}).$$

Since $z \in \mathcal{Z}_e$ is a real simple zero of $\omega\varepsilon(\cdot)$ thanks to assumption (H₂), i.e. $\mathcal{Z}_e \cap \mathcal{P}_m = \emptyset$, it is not a zero nor a pole $\omega\mu(\cdot)$. Thus, g can indeed be extended analytically for $\omega = z$ with

$$g(z) = (\omega\varepsilon)'(z) z\mu(z).$$

Thus, one has $A_z = \frac{1}{g(z)} = \frac{1}{z\mu(z)(\omega\varepsilon)'(z)}$ that is (5.28).

We show now that $\text{Im } A_z < 0$. As $z \in \mathcal{Z}_e$, one observes on one hand (see remark 5.4) that $\mathcal{Z}_s = \mathcal{Z}_e$ and that all the $\alpha_{e,j}$ vanish. Thus, $\omega\varepsilon(\omega)$ is real-valued on the real axis (outside the poles \mathcal{P}_e of $\varepsilon(\omega)$) and from expression (1.6), one easily deduces that $(\omega\varepsilon)'(\omega) > 0$ on $\mathbb{R} \setminus \mathcal{P}_e$. In particular, as $z \in \mathbb{R} \setminus \mathcal{P}_e$, one has

$$(\omega\varepsilon)'(z) > 0.$$

On the other hand, as all the $\alpha_{e,j}$ vanish, one knows by the weak dissipation condition (1.11) that at least one coefficient $\alpha_{m,\ell} > 0$. Therefore by the formula (3.17), $\text{Im}(z\mu(z)) > 0$. Thus, one deduces that

$$\text{Im}((z\mu(z))^{-1}) = -\text{Im}(z\mu(z))/|z\mu(z)|^2 < 0.$$

Together with $(\omega\varepsilon)'(z) > 0$, this implies that $\text{Im}(A_z)$ is negative. \square

Combining the Lemma 5.8 and Lemma 5.9 (formula (5.27) and (5.28)) and proceeding as in the proof of Lemma 4.5 gives the following estimate for $\mathbb{U}_{z,s}(\mathbf{k}, t)$.

Lemma 5.10. *If $\mathcal{Z}_s \neq \emptyset$, then there exists $k_- > 0$ and $C > 0$ such that the function $\mathbb{U}_{z,s}(\mathbf{k}, t)$ defined by (5.10) and (5.11), satisfies the following estimate*

$$|\mathbb{U}_{z,s}(\mathbf{k}, t)| \lesssim e^{-C|\mathbf{k}|^2 t} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0 \quad \text{and} \quad 0 < |\mathbf{k}| \leq k_-. \quad (5.29)$$

5.3.4 Estimates of $\mathbb{U}_{z,-}(\mathbf{k}, t)$ for $0 < |\mathbf{k}| \leq k_-$

As announced in Section 5.3.1, since we simply want to obtain a “rough” exponential decay estimate for $\mathbb{U}_{z,-}(\mathbf{k}, t)$. We give a direct proof of it using Riesz-Dunford functional calculus.

Lemma 5.11. *There exists $\delta > 0$ and $k_- > 0$ such that $\mathbb{U}_{z,-}(\mathbf{k}, t)$, defined by (5.10) and (5.11), satisfies*

$$|\mathbb{U}_{z,-}(\mathbf{k}, t)| \lesssim e^{-\delta t} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0, \quad \forall 0 < |\mathbf{k}| \leq k_-. \quad (5.30)$$

Proof. The proof is similar to that of Lemma 4.13 and even shorter since $|\mathbf{k}|$ is bounded here. We introduce a fixed (positively oriented) simple closed contour Γ , included in \mathbb{C}^- such that all the zeros of \mathcal{Z}_- lie inside Γ (see figure 4). We choose Γ such that $\Gamma \cap \mathcal{P} = \emptyset$. We denote by

$$\delta = \min\{-\text{Im}(\omega), \omega \in \Gamma\} > 0 \quad (5.31)$$

the distance from Γ to the real axis.

For $|\mathbf{k}|$ positive and small enough, $\mathbb{U}_{z,-}(\mathbf{k}, t)$ is on one hand well-defined by (5.10) and (5.11) and on the other hand, by Proposition 5.1, Γ encloses all eigenvalues $\omega_{p,n}(|\mathbf{k}|)$ for $z \in \mathcal{Z}_-$ and $n \in \{1, \dots, \mathfrak{m}_z\}$ but no other elements of the spectrum of $\mathbb{A}_{|\mathbf{k}|,\perp}$. Thus by the Riesz-Dunford functional calculus, one gets that there exists $k_- > 0$ such that

$$\mathbb{U}_{z,-}(\mathbf{k}, t) = -\frac{\mathcal{R}_{\mathbf{k}}^*}{2i\pi} \int_{\Gamma} e^{-i\omega t} R_{|\mathbf{k}|}(\omega) \mathcal{R}_{\mathbf{k}} \mathbb{U}_0(\mathbf{k}) d\omega, \quad \text{for } 0 < |\mathbf{k}| \leq k_-, \quad (5.32)$$

where the resolvent $R_{\mathbf{k}}(\omega)$ is well-defined by Proposition 3.2 for $(|\mathbf{k}|, \omega) \in K = [0, k_-] \times \Gamma$. Moreover, the function $(|\mathbf{k}|, \omega) \mapsto R_{\mathbf{k}}(\omega)$, valued in $\mathcal{L}(\mathbb{C}_{\perp})$, is continuous on the compact K . Then, using the fact that $\mathcal{R}_{\mathbf{k}}$ is unitary and the definition (5.31) of δ (which implies that $|e^{-i\omega t}| \leq e^{-\delta t}$ on Γ), it follows from (5.32) that

$$|\mathbb{U}_{z,-}(\mathbf{k}, t)| \leq C e^{-\delta t} |\mathbb{U}_0(\mathbf{k})|, \quad \forall |\mathbf{k}| \in (0, k_-] \quad \text{where } C = \max_{(|\mathbf{k}|, \omega) \in K} \|R_{|\mathbf{k}|}(\omega)\|.$$

□

5.3.5 The global estimate

Theorem 5.12. *There exists $k_- > 0$, $C, \tilde{C} > 0$ such that for $0 < |\mathbf{k}| \leq k_-$, the spatial Fourier components $|\mathbb{U}(\mathbf{k}, t)|$ of the solution of (1.13) with initial condition $\mathbb{U}_0 \in \mathcal{H}_{\perp}$ satisfy:*

$$|\mathbb{U}(\mathbf{k}, t)| \leq \tilde{C} e^{-C|\mathbf{k}|^2 t} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0. \quad (5.33)$$

Proof. The proof is similar to that of Theorem 4.14 using Lemmas 5.7, 5.10 and 5.11. \square

We prove in the following result that the estimates of Theorem 5.12 is optimal for an infinite family of well chosen initial conditions $\mathbf{U}_0 \in \mathcal{H}_{\perp, \text{LF}}^p$ for any fixed $p \geq 0$.

Theorem 5.13. *Let $p \geq 0$ and $k_- > 0$ and $\phi : \mathbb{R}^+ \mapsto \mathbb{R}$ be any measurable and bounded function satisfying*

$$\text{supp } \phi \subset [0, k_-] \text{ and } 0 < |\phi(|\mathbf{k}|)| \lesssim |\mathbf{k}|^p \text{ for } |\mathbf{k}| \leq k_-. \quad (5.34)$$

If the initial condition \mathbf{U}_0 of (1.13) is defined (for k_- small enough) via its Fourier transform:

$$\mathbb{U}_0(\mathbf{k}) = \mathcal{F}(\mathbf{U}_0)(\mathbf{k}) = \phi(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}}^* \frac{\mathcal{V}_{|\mathbf{k}|}(\omega_{0,1}(|\mathbf{k}|)) \mathbf{e}_1}{|\mathcal{V}_{|\mathbf{k}|}(\omega_{0,1}(|\mathbf{k}|)) \mathbf{e}_1|}, \quad \forall \mathbf{k} \in \mathbb{R}^{3,*}, \quad (5.35)$$

then $\mathbf{U}_0 \in \mathcal{H}_{\perp, \text{LF}}^p$ and $\exists C, \tilde{C} > 0$ such that the associated solution \mathbf{U} of (1.13) satisfy

$$\tilde{C} e^{Ct|\mathbf{k}|^2} |\mathbb{U}_0(\mathbf{k})| \leq |\mathbb{U}(\mathbf{k}, t)|, \quad \forall t \geq 0 \text{ and } \forall \mathbf{k} \in \mathbb{R}^{3,*}. \quad (5.36)$$

In other words, the estimate (5.33) is optimal for an infinite family of solutions.

Proof. Using (5.34) and (5.35), it is clear that $\mathbf{U}_0 \in \mathcal{H}_{\perp, \text{LF}}^p$. The rest of the proof uses the following asymptotic expansion of $\text{Im } \omega_{0,1}(|\mathbf{k}|)$ (given by Lemma 5.6):

$$\text{Im } \omega_{0,1}(|\mathbf{k}|) = \text{Im} \left(-(\varepsilon\mu)'(0) |\mathbf{k}|^2 + o(|\mathbf{k}|^2) \right), \text{ as } |\mathbf{k}| \rightarrow 0 \text{ with } \text{Im} \left(-(\varepsilon\mu)'(0) \right) < 0$$

and is similar to the proof of Theorem 4.15. Therefore, the details are left to the reader. \square

6 Estimates of “mid-range frequencies” components of the solution

The goal is to prove (see Theorem 6.2) that for intermediate frequency components, i.e. for $0 < k_- \leq |\mathbf{k}| \leq k_+$, the Fourier components of the solution $\mathbb{U}(\mathbf{k}, t)$ decay exponentially with a uniform exponent depending only on the compact set $K = [k_-, k_+]$. The proof is based on three key ingredients: the Dunford decomposition in linear algebra, the following perturbation Lemma and finally a compactness argument.

Lemma 6.1. *Let $\mathbb{M} \in \mathcal{L}(\mathbb{C}_{\perp}^N)$. We assume that there exists two constants $C, \alpha > 0$ such that*

$$\|e^{t\mathbb{M}}\| \leq C e^{-\alpha t}, \quad \forall t \geq 0.$$

Then, for any perturbation $\Delta \in \mathcal{L}(\mathbb{C}_{\perp}^N)$, one has the following estimate:

$$\|e^{t(\mathbb{M}+\Delta)}\| \leq C e^{(-\alpha+C\|\Delta\|)t}, \quad \forall t \geq 0.$$

Proof. The proof, based on the Duhamel formula and the Grönwall’s lemma, is done e.g. in lemma 1.6 page 97 of [1] or in the Proposition 4.2.18 page 406 of [13]. \square

Theorem 6.2. *Let $K = [k_-, k_+]$ be a compact interval of $\mathbb{R}^{+,*}$, then there exist two constants $C > 0$ and $\beta > 0$ (depending only on K) such that the spatial Fourier components $|\mathbb{U}(\mathbf{k}, t)|$ of the solution of (1.13) with initial condition $\mathbb{U}_0 \in \mathcal{H}_\perp$ satisfy*

$$|\mathbb{U}(\mathbf{k}, t)| \leq C e^{-\beta t} |\mathbb{U}_0(\mathbf{k})|, \quad \forall t \geq 0 \text{ and a.e. } \mathbf{k} \in \mathbb{R}^3 \mid |\mathbf{k}| \in K. \quad (6.1)$$

Proof. Let $k \in K$ be fixed. By virtue of (3.14), $\sigma(\mathbb{A}_{k,\perp})$ is included in the lower half-plane \mathbb{C}^- , thus it follows from the Dunford Decomposition of $\mathbb{A}_{k,\perp}$ (see e.g. Corollary 2.26 page 106 of [13]) that there exist $\alpha_k > 0$ and $C_k > 0$ such that

$$\|e^{-i\mathbb{A}_{k,\perp}t}\| \leq C_k e^{-\alpha_k t}, \quad \forall t \geq 0. \quad (6.2)$$

As the function $|\mathbf{k}| \mapsto -i\mathbb{A}_{|\mathbf{k}|,\perp}$ is clearly continuous from $\mathbb{R}^{+,*}$ to $\mathcal{L}(\mathbf{C}_\perp^N)$, There exists an open interval $(k - \eta_k, k + \eta_k)$ such that

$$\| -i(\mathbb{A}_{|\mathbf{k}|,\perp} - \mathbb{A}_{k,\perp}) \| < \frac{\alpha_k}{2C_k}, \quad \forall |\mathbf{k}| \in (k - \eta_k, k + \eta_k) \cap K. \quad (6.3)$$

Hence, applying Lemma 6.1 with

$$\mathbb{M} = -i\mathbb{A}_{k,\perp} \quad \text{and} \quad \Delta = -i(\mathbb{A}_{|\mathbf{k}|,\perp} - \mathbb{A}_{k,\perp})$$

yields the following inequality for all $t \geq 0$ and $\forall |\mathbf{k}| \in (k - \eta_k, k + \eta_k) \cap K$:

$$\|e^{-i\mathbb{A}_{|\mathbf{k}|,\perp}t}\| \leq C_k e^{-\beta_k t} \quad \text{with} \quad \beta_k = \alpha_k - C_k \|\Delta\|. \quad (6.4)$$

Thus, one deduces immediately from (6.3) and (6.4) that

$$\|e^{-i\mathbb{A}_{|\mathbf{k}|,\perp}t}\| \leq C_k e^{-\frac{\alpha_k}{2}t}, \quad \forall t \geq 0 \quad \text{and} \quad \forall |\mathbf{k}| \in (k - \eta_k, k + \eta_k) \cap K. \quad (6.5)$$

As the compact $K \subset \bigcup_{k \in K} (k - \eta_k, k + \eta_k)$, there exists $N > 0$ and $k_1, k_2, \dots, k_N \in K$ such that

$$K \subset \bigcup_{1 \leq n \leq N} (k_n - \eta_{k_n}, k_n + \eta_{k_n}). \quad (6.6)$$

With $\alpha_{k_i} > 0$ and $C_{k_i} > 0$ given by (6.2) for $k = k_i$, we define

$$C = \max_{i=1,\dots,N} C_{k_i} > 0 \quad \text{and} \quad \alpha = \min_{i=1,\dots,N} \alpha_{k_i} > 0.$$

Then, combining (6.5) and (6.6) yields

$$\|e^{-i\mathbb{A}_{|\mathbf{k}|,\perp}t}\| \leq C e^{-\frac{\alpha}{2}t}, \quad \forall t \geq 0 \quad \text{and} \quad \forall |\mathbf{k}| \in K.$$

As the operator $\mathcal{R}_{\mathbf{k}}$ is unitary, one finally deduces, with (2.12), the estimate (6.1). \square

7 Proof of the main Theorems of the paper

7.1 Proof of Theorem 1.10 (decay rate estimates)

This section constitutes the last step of the proof of Theorem 1.10. For this final stage, we follow the approach developed in the proofs of Theorems 2.4 and 4.4 of [5]. For the sake of readability, we recall here the arguments, which rely on Plancherel's identity and the three decay estimates proved for high, low and intermediate frequencies respectively in Theorems 4.14, 5.12 and 6.2.

Let k_- and k_+ be two fixed positive real numbers satisfying $k_- < k_+$ such that the estimates (4.96) and (5.33) of Theorems 4.14 and 5.12 hold.

Using the Plancherel identity, any solution of (1.13) with initial condition $\mathbf{U}_0 \in \mathcal{H}_\perp$ satisfies

$$\|\mathbf{U}(t)\|_{\mathcal{H}}^2 = \int_{|\mathbf{k}| < k_-} |\mathbb{U}(\mathbf{k}, t)|^2 d\mathbf{k} + \int_{k_- \leq |\mathbf{k}| \leq k_+} |\mathbb{U}(\mathbf{k}, t)|^2 d\mathbf{k} + \int_{k_+ < |\mathbf{k}|} |\mathbb{U}(\mathbf{k}, t)|^2 d\mathbf{k}. \quad (7.1)$$

Step 1: proof of the convergence result (1.36)

In a non critical configuration, using the low frequency estimate (5.33), the intermediate frequency estimate (6.2) and the high frequency estimate (4.96) of $\mathbb{U}(\mathbf{k}, t)$, one sees that there exist three constants $C_1 > 0$, $C_2 > 0$ and $C_3 > 0$ (independent of \mathbf{U}) such that

$$\|\mathbf{U}(t)\|_{\mathcal{H}}^2 \lesssim \int_{|\mathbf{k}| < k_-} e^{-C_1 |\mathbf{k}|^2 t} |\mathbb{U}_0(\mathbf{k})|^2 d\mathbf{k} + e^{-C_2 t} \|\mathbf{U}_0\|_{\mathcal{H}}^2 + \int_{k_+ < |\mathbf{k}|} e^{-C_3 |\mathbf{k}|^{-2} t} |\mathbb{U}_0(\mathbf{k})|^2 d\mathbf{k}. \quad (7.2)$$

When $t \rightarrow +\infty$, the second term in (7.2) converges exponentially to 0, whereas the first and third terms tend to 0 by the Lebesgue's dominated convergence Theorem.

In a critical configuration, one obtains (1.36) by the same reasoning only after replacing the high frequency estimate (4.96) by (4.97) (thus it consists to substitute the factor $e^{-C_3 |\mathbf{k}|^{-2} t}$ by $e^{-C_3 |\mathbf{k}|^{-4} t}$ in the second integral of (7.2)).

Step 2: Estimate of the low frequency term in (7.2)

We assume now that the initial condition $\mathbf{U}_0 \in \mathcal{H}_\perp \cap \mathcal{L}_p^N$ for some integer $p \geq 0$.

Thus, by definition on the \mathcal{L}_p^N norm (see (1.26)), $|\mathbb{U}_0(\mathbf{k})| \leq |\mathbf{k}|^p \|\mathbf{U}_0\|_{\mathcal{L}_p^N}$ for a.e. $\mathbf{k} \in \mathbb{R}^3$ and consequently

$$\int_{|\mathbf{k}| < k_-} e^{-C_1 |\mathbf{k}|^2 t} |\mathbb{U}_0(\mathbf{k})|^2 d\mathbf{k} \leq \left(\int_{\mathbb{R}^3} e^{-C_1 |\mathbf{k}|^2 t} |\mathbf{k}|^{2p} d\mathbf{k} \right) \|\mathbf{U}_0\|_{\mathcal{L}_p^N}^2. \quad (7.3)$$

Since, with the change of variable $\xi = \sqrt{C_1 t} \mathbf{k}$,

$$\int_{\mathbb{R}^3} |\mathbf{k}|^{2p} e^{-C_1 |\mathbf{k}|^2 t} d\mathbf{k} = \frac{1}{(C_1 t)^{p+\frac{3}{2}}} \int_{\mathbb{R}^3} |\xi|^{2p} e^{-|\xi|^2} d\xi \equiv \frac{C(p)}{t^{p+\frac{3}{2}}},$$

one concludes from (7.3) that

$$\int_{|\mathbf{k}| < k_-} e^{-C_1 |\mathbf{k}|^2 t} |\mathbb{U}_0(\mathbf{k})|^2 d\mathbf{k} \leq \frac{C(p)}{t^{p+\frac{3}{2}}} \|\mathbb{U}_0\|_{\mathcal{L}^p}^2. \quad (7.4)$$

Step 3: Estimate of the high frequency term in (7.2)

In a non critical configuration:

We assume here that the Maxwell's system is in a non critical configuration and that the initial condition $\mathbf{U}_0 \in \mathcal{H}_\perp \cap \mathbf{H}^m(\mathbb{R}^3)^N$ for some $m > 0$. We estimate here the third term of (7.2) corresponding to the high frequency contribution. To this aim, we denote by $\langle \mathbf{k} \rangle$ the quantity: $\langle \mathbf{k} \rangle := (1 + |\mathbf{k}|^2)^{1/2}$. Then, using the fact that $e^{-C_3 |\mathbf{k}|^{-2} t} \leq e^{-C_3 \langle \mathbf{k} \rangle^{-2} t}$, we rewrite this third term as follows (we simply introduce artificially the factor $\langle \mathbf{k} \rangle^m / t^m$)

$$\begin{aligned} \int_{k_+ < |\mathbf{k}|} e^{-C_3 |\mathbf{k}|^{-2} t} |\mathbb{U}_0(\mathbf{k})|^2 d\mathbf{k} &\leq t^{-m} \int_{k_+ < |\mathbf{k}|} \langle \mathbf{k} \rangle^{2m} |\mathbb{U}_0(\mathbf{k})|^2 \left(\frac{t}{\langle \mathbf{k} \rangle^2} \right)^m e^{-C_3 \langle \mathbf{k} \rangle^{-2} t} d\mathbf{k} \\ &\leq t^{-m} \int_{\mathbb{R}^3} \langle \mathbf{k} \rangle^{2m} |\mathbb{U}_0(\mathbf{k})|^2 F_m(t/\langle \mathbf{k} \rangle^2) d\mathbf{k}, \end{aligned}$$

where we have set $F_m(r) := r^m e^{-C_3 r}$, $r \geq 0$ which satisfies: $\sup_{r \geq 0} F_m(r) = \tilde{C}_m := (m/(C_3 e))^m$.

Thus, by the Fourier characterization of Sobolev norms:

$$\int_{k_+ < |\mathbf{k}|} e^{-C_3 |\mathbf{k}|^{-2} t} |\mathbb{U}_0(\mathbf{k})|^2 d\mathbf{k} \lesssim \tilde{C}_m t^{-m} \|\mathbb{U}_0\|_{\mathbf{H}^m(\mathbb{R}^3)^N}^2, \quad \forall t > 0. \quad (7.5)$$

In a critical configuration:

We assume here that the Maxwell's system is in a critical configuration with again an initial condition $\mathbf{U}_0 \in \mathcal{H}_\perp \cap \mathbf{H}^m(\mathbb{R}^3)^N$. The proof is similar to that of the estimate (7.5). We simply enlightens the differences. In a critical configuration, one has to replace the estimate (4.96) by (4.97). Thus, in (7.2), one has to substitute the factor $e^{-C_3 |\mathbf{k}|^{-2} t}$ by $e^{-C_3 |\mathbf{k}|^{-4} t}$ in the second integral of (7.2). Then, one has

$$\begin{aligned} \int_{k_+ < |\mathbf{k}|} e^{-C_3 |\mathbf{k}|^{-4} t} |\mathbb{U}_0(\mathbf{k})|^2 d\mathbf{k} &\leq t^{-m/2} \int_{k_+ < |\mathbf{k}|} \langle \mathbf{k} \rangle^{2m} |\mathbb{U}_0(\mathbf{k})|^2 \left(\frac{t}{\langle \mathbf{k} \rangle^4} \right)^{m/2} e^{-C_3 \langle \mathbf{k} \rangle^{-4} t} d\mathbf{k} \\ &\leq t^{-m/2} \int_{\mathbb{R}^3} \langle \mathbf{k} \rangle^{2m} |\mathbb{U}_0(\mathbf{k})|^2 F_{m/2}(t/\langle \mathbf{k} \rangle^4) d\mathbf{k}. \end{aligned}$$

Setting $\tilde{C}_{m/2} := \sup_{r \geq 0} F_{m/2}(r)$, this yields, by definition (1.27) of the $\mathbf{H}^m(\mathbb{R}^3)^N$ norm,

$$\int_{k_+ < |\mathbf{k}|} e^{-C_3 |\mathbf{k}|^{-2} t} |\mathbb{U}_0(\mathbf{k})|^2 d\mathbf{k} \leq \tilde{C}_{m/2} t^{-m/2} \|\mathbb{U}_0\|_{\mathbf{H}^m(\mathbb{R}^3)^N}^2, \quad \forall t > 0. \quad (7.6)$$

Step 4: Proof of the estimate (1.37) and (1.38).

In the non critical case (resp. the critical case), it suffices to substitute (7.4) and (7.5) (resp. (7.4) and (7.6)) into (7.2) to obtain (1.37) (resp. (1.38)).

7.2 Proof of Theorem 1.14 (optimality decay rate estimates)

7.2.1 Optimality of the high frequency polynomial decay rate

Determination of γ_m^{HF} in a non-critical configuration

We assume first that the Maxwell system is in a non-critical configuration.

First, we emphasize that for any initial conditions $\mathbf{U}_0 \in \mathcal{H}_{\perp, \text{HF}}^m$ (see (1.40) for the definition of $\mathcal{H}_{\perp, \text{HF}}^m$) which does not contain any low frequency Fourier components (since as $\mathbf{U}_0 \in \mathcal{H}_{\perp, \text{HF}}^m$, $\text{supp}(\mathbf{U}_0) \subset \mathbb{R}^3 \setminus B(0, k_+)$), we have the upper-bound given by (7.2) and (7.5), namely

$$\|\mathbf{U}(t)\|_{\mathcal{H}}^2 \lesssim \int_{k_+ < |\mathbf{k}|} e^{-C_3 |\mathbf{k}|^{-4} t} |\mathbb{U}_0(\mathbf{k})|^2 d\mathbf{k} \lesssim \frac{\|\mathbf{U}_0\|_{\mathbf{H}^m(\mathbb{R}^3)^N}^2}{t^m}, \quad \forall t > 0. \quad (7.7)$$

Clearly, (7.7) implies that γ_m^{HF} (defined by (1.42)) exists and satisfies $\gamma_m^{\text{HF}} \geq m$. To show that $\gamma_m^{\text{HF}} = m$, we construct for any $\varepsilon > 0$ an initial condition $\mathbf{U}_{0, \varepsilon} \in \mathcal{H}_{\perp, \text{HF}}^m$ such that the associated solution of (1.13) satisfies

$$\|\mathbf{U}_{\varepsilon}(t)\|_{\mathcal{H}}^2 \geq \frac{C}{t^{m+\varepsilon}}, \quad \text{for some } C > 0, \quad \forall t \geq 1.$$

To this aim, according to Theorem 4.15, we choose $\mathbf{U}_{0, \varepsilon}$ in the form (4.99), more precisely

$$\begin{cases} \mathbb{U}_{0, \varepsilon}(\mathbf{k}) = \phi_{\varepsilon}(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}}^* \frac{\mathcal{V}_{|\mathbf{k}|}(\omega_{+\infty}(|\mathbf{k}|)) \mathbf{e}_1}{|\mathcal{V}_{|\mathbf{k}|}(\omega_{+\infty}(|\mathbf{k}|)) \mathbf{e}_1|}, \\ \phi_{\varepsilon}(|\mathbf{k}|) = (1 + |\mathbf{k}|^2)^{-(\frac{3}{4} + \frac{m}{2} + \frac{\varepsilon}{2})} \text{ if } |\mathbf{k}| \geq k_+, \text{ and } \phi_{\varepsilon}(|\mathbf{k}|) = 0 \text{ otherwise,} \end{cases} \quad (7.8)$$

where the exponent $\frac{3}{4} + \frac{m}{2} + \frac{\varepsilon}{2}$ has been chosen just above in order to ensure the H^m regularity of $\mathbf{U}_{0, \varepsilon}$ (see (4.98)). As moreover $\mathbb{U}_{0, \varepsilon}(\mathbf{k})$ is supported in $\{|\mathbf{k}| \geq k_+\}$, $\mathbf{U}_{0, \varepsilon} \in \mathcal{H}_{\perp, \text{HF}}^m$.

Therefore, applying the inequality (4.100) of Theorem 4.15, we get, as $|\mathbb{U}_{0, \varepsilon}(\mathbf{k})| = \phi_{\varepsilon}(|\mathbf{k}|)$,

$$e^{-\frac{2Ct}{|\mathbf{k}|^2}} (1 + |\mathbf{k}|^2)^{-(\frac{3}{2} + m + \varepsilon)} \lesssim |\mathbb{U}_{\varepsilon}(\mathbf{k}, t)|^2, \quad \text{for a.e. } \mathbf{k} \in \mathbb{R}^3 \mid |\mathbf{k}| \geq k_+ \text{ and } \forall t \geq 0. \quad (7.9)$$

By Plancherel identity, $\|\mathbf{U}_{\varepsilon}(t)\|_{\mathcal{H}}^2 = \int_{k_+ < |\mathbf{k}|} |\mathbb{U}_{\varepsilon}(\mathbf{k}, t)|^2 d\mathbf{k}$. Therefore, using (7.9), one gets

$$I_{\varepsilon}(t) = \int_{k_+ < |\mathbf{k}|} e^{-\frac{2Ct}{|\mathbf{k}|^2}} (1 + |\mathbf{k}|^2)^{-(3/2 + m + \varepsilon)} d\mathbf{k} \lesssim \|\mathbf{U}_{\varepsilon}(t)\|_{\mathcal{H}}^2. \quad (7.10)$$

that is to say, the integrand in (7.10) only depends on $|\mathbf{k}|$,

$$I_{\varepsilon}(t) = 4\pi \int_{k_+}^{+\infty} e^{-\frac{2Ct}{|\mathbf{k}|^2}} (1 + |\mathbf{k}|^2)^{-(3/2 + m + \varepsilon)} |\mathbf{k}|^2 d|\mathbf{k}|.$$

Using the change of variable $\xi = \sqrt{t}/|\mathbf{k}|$ in $I_{\varepsilon}(t)$ with $|\mathbf{k}|^2 d|\mathbf{k}| = -t^{\frac{3}{2}} \xi^{-4} d\xi$ yields

$$I_{\varepsilon}(t) = 4\pi t^{\frac{3}{2}} \int_0^{\frac{\sqrt{t}}{k_+}} e^{-2C\xi^2} \xi^{3+2m+2\varepsilon} (\xi^2 + t)^{-(3/2 + m + \varepsilon)} \frac{d\xi}{\xi^4}.$$

Then, using $\xi^2 \leq t/k_+^2$ and thus that $\xi^2 + t \leq (1 + k_+^{-2})t$ on the domain of integration gives

$$I_\varepsilon(t) \geq \frac{C_m^\varepsilon(t)}{t^{m+\varepsilon}}, \quad C_m^\varepsilon(t) := 4\pi(1 + k_+^{-2})^{-(m+\frac{3}{2}+\varepsilon)} \int_0^{\frac{\sqrt{t}}{k_+}} \frac{e^{-2C\xi^2}}{\xi^{1-2m-2\varepsilon}} d\xi.$$

As $C_m^\varepsilon(t) > 0$ is a strictly increasing, restricting ourselves to $t \geq 1$, (7.10) leads to

$$\forall t \geq 1, \quad \frac{C_m^\varepsilon(1)}{t^{m+\varepsilon}} \leq I_\varepsilon(t) \lesssim \|\mathbf{U}_\varepsilon(t)\|_{\mathcal{H}}^2.$$

Thus, $\gamma_m^{HF} \leq m + \varepsilon$, for all $\varepsilon > 0$ and this yields $\gamma_m^{HF} \leq m$. Thus, one concludes that $\gamma_m^{HF} = m$.

Determination of γ_m^{HF} in a critical configuration

We assume now that the Maxwell system is in a critical configuration.

The proof is very similar to that of the non critical case and we shall only point out the differences. First the inequality (7.6) shows that $\gamma_m^{HF} \geq m/2$.

To prove the reverse inequality, we proceed as for the non critical case. The main difference lies in the choice of the initial data $\mathbf{U}_{0,\varepsilon}$ which is now chosen as in (4.102), more precisely

$$\mathbb{U}_{0,\varepsilon}(\mathbf{k}) = \mathcal{F}(\mathbf{U}_0)(\mathbf{k}) = \phi_\varepsilon(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}}^* \frac{\mathcal{V}_{|\mathbf{k}|}(\omega_p(|\mathbf{k}|))\mathbf{e}_1}{|\mathcal{V}_{|\mathbf{k}|}(\omega_p(|\mathbf{k}|))\mathbf{e}_1|}, \quad \forall \mathbf{k} \in \mathbb{R}^{3,*}, \quad (7.11)$$

for some $p \in \mathcal{P}_s$ such that $\omega_p(|\mathbf{k}|)$ satisfies (4.101) (see again Theorem 4.15) and the function ϕ_ε defined in (7.8). In that case, we have to apply the inequality (4.103), instead of (4.100), which leads to (instead of (7.9))

$$e^{-\frac{2Ct}{|\mathbf{k}|^4}} (1 + |\mathbf{k}|^2)^{-(\frac{3}{2}+m+\varepsilon)} \lesssim |\mathbb{U}_\varepsilon(\mathbf{k}, t)|^2, \quad \text{for a.e. } \mathbf{k} \in \mathbb{R}^3 \mid |\mathbf{k}| \geq k_+ \text{ and } \forall t \geq 0. \quad (7.12)$$

the main difference with being (7.9) being that $|\mathbf{k}|^2$ is replaced by $|\mathbf{k}|^4$ in the exponential.

The rest of the calculations follows the same lines as in the non critical case modulo trivial adaptations : for instance the change of variable $\xi = \sqrt{t}/|\mathbf{k}|$ is replaced by $\xi = \sqrt{t}/|\mathbf{k}|^2$. The remaining details are left to the reader.

7.2.2 Optimality of the low frequency polynomial decay rate

We compute here the exponent γ_p^{LF} (defined by (1.43)) for $p \geq 0$. The approach is similar to that used in Section 7.2.1 to compute γ_m^{HF} .

First, we underline that for any initial condition $\mathbf{U}_0 \in \mathcal{H}_{\perp,LF}^p$ (see (1.41) for the definition of $\mathcal{H}_{\perp,LF}^p$) which does not contain any high frequency Fourier components (since the support of \mathbb{U}_0 is included in $\overline{B(0, k_-)}$), we have the upper-bounds (7.2) and (7.4), namely,

$$\|\mathbf{U}(t)\|_{\mathcal{H}}^2 \lesssim \int_{|\mathbf{k}| < k_-} e^{-C_1|\mathbf{k}|^2 t} |\mathbb{U}_0(\mathbf{k})|^2 d\mathbf{k} \lesssim \frac{C(p)}{t^{p+\frac{3}{2}}} \|\mathbf{U}_0\|_{\mathcal{L}_p^N}^2. \quad (7.13)$$

Clearly, (7.13) implies that γ_p^{LF} exists and satisfies $\gamma_p^{LF} \geq p + 3/2$. To show that $\gamma_p^{LF} \leq p + 3/2$, we construct an initial condition $\mathbf{U}_0 \in \mathcal{H}_{\perp, LF}^p$ such that

$$\|\mathbf{U}(t)\|_{\mathcal{H}}^2 \geq \frac{C(p)}{t^{p+\frac{3}{2}}}, \quad \text{for some } C > 0, \quad \forall t \geq 1. \quad (7.14)$$

To this aim, according to Theorem 5.13, we choose \mathbf{U}_0 in the form (5.35), more precisely

$$\begin{cases} \mathbb{U}_0(\mathbf{k}) = \phi(|\mathbf{k}|) \mathcal{R}_{\mathbf{k}}^* \frac{\mathcal{V}_{|\mathbf{k}|}(\omega_{0,1}(|\mathbf{k}|)) \mathbf{e}_1}{|\mathcal{V}_{|\mathbf{k}|}(\omega_{0,1}(|\mathbf{k}|)) \mathbf{e}_1}, \\ \phi(|\mathbf{k}|) = |\mathbf{k}|^p \text{ if } |\mathbf{k}| \leq k_-, = 0 \text{ otherwise.} \end{cases} \quad (7.15)$$

Therefore, applying the inequality (5.33) of Theorem 5.13, we get as $|\mathbb{U}_0(\mathbf{k})| = \phi(|\mathbf{k}|)$:

$$e^{-2C|\mathbf{k}|^2 t} |\mathbf{k}|^{2p} \lesssim |\mathbb{U}(\mathbf{k}, t)|^2, \quad \text{for a.e. } \mathbf{k} \in \mathbb{R}^3 \mid |\mathbf{k}| \leq k_- \text{ and } \forall t \geq 0. \quad (7.16)$$

By virtue of Plancherel identity: $\|\mathbf{U}(t)\|_{\mathcal{H}}^2 = \int_{|\mathbf{k}| < k_-} |\mathbb{U}(\mathbf{k}, t)|^2 d\mathbf{k}$, it follows from (7.16) that

$$\int_{|\mathbf{k}| < k_-} e^{-2C|\mathbf{k}|^2 t} |\mathbf{k}|^{2p} d\mathbf{k} = 4\pi \int_0^{k_-} e^{-2C|\mathbf{k}|^2 t} |\mathbf{k}|^{2p} |\mathbf{k}|^2 d|\mathbf{k}| \lesssim \|\mathbf{U}(t)\|_{\mathcal{H}}^2. \quad (7.17)$$

Hence, using the change of variable $\xi = \sqrt{2Ct} |\mathbf{k}|$, one obtains

$$\frac{1}{(2Ct)^{p+\frac{3}{2}}} \int_0^{\sqrt{t}k_-} \xi^{2p+2} e^{-\xi^2} d\xi \lesssim \|\mathbf{U}(t)\|_{\mathcal{H}}^2.$$

It leads to (7.14) and more precisely to

$$\frac{C(k_-, p)}{t^{p+\frac{3}{2}}} \lesssim \|\mathbf{U}(t)\|_{\mathcal{H}}^2 \quad \text{with } C(k_-, p) = \frac{1}{(2C)^{p+\frac{3}{2}}} \int_0^{k_-} \xi^{2p+2} e^{-\xi^2} d\xi, \quad \forall t \geq 1.$$

Thus, $\gamma_p^{LF} \leq p + 3/2$ and one concludes that $\gamma_p^{LF} = p + 3/2$.

7.3 Extensions of Theorems 1.10 and 1.14

7.3.1 The case of a non dispersive electric or magnetic media

For some passive dispersive materials (e.g. in metals, see for instance [17]), the magnetization of the material is weak so that for the corresponding physical model, the magnetic permeability μ is assumed to be a constant function of the frequency: $\mu(\omega) = \mu_0$ whereas the permittivity $\varepsilon(\omega)$ is given by a generalized Lorentz models. This situation corresponds to the case $N_m = 0$ and $N_e \geq 1$ in formula (1.6). In this setting, we would like to mention that our results still hold and to point out the simplifications in their proof.

Assume that $N_m = 0$ and $N_e \geq 1$ in the evolution problem (1.13). It implies in particular that

$\mathbf{U} = (\mathbf{E}, \mathbf{H}, \mathbf{P}, \dot{\mathbf{P}})$. Adapting the definition 1.8 to this particular case, the Maxwell's system is in a critical configuration if and only if one coefficient $\alpha_{e,j}$ vanishes (which imposes that $N_e \geq 2$ since by (1.11) at least one coefficient $\alpha_{e,j'}$ has to be positive). Thus, it is in a non-critical configuration if only if all the coefficients $\alpha_{e,j}$ are positive.

Under the same assumptions, the Theorems 1.10 and 1.14 still hold. The proof of these results remain the same as that developed in Sections 3 to 7 with the following two simplifications:

1. Concerning the Section 4, $\mathcal{P}_s = \mathcal{P}_e \cap \mathbb{R} \neq \emptyset$ if and only if one coefficient $\alpha_{e,j}$ vanishes which is precisely the definition of a critical configuration. Thus, in the Section 4.3.3 for the estimate of the term $\mathbb{U}_s(\mathbf{k}, t)$, one just needs to consider the critical configuration in Corollary 4.8 and Lemma 4.9.

As $\mathcal{P}_d = \emptyset$, $\mathbb{U}_d(\mathbf{k}, t) = 0$ in (4.10) and one can skip the Section 4.3.4 devoted to the estimate of this term.

2. In the Section 5, since by (1.11) at least one coefficient $\alpha_{e,j}$ is positive, we deduce from (3.17) and (3.18), that the set \mathcal{Z}_s of real zeros (see (5.3)) is empty : $\mathcal{Z}_s = \mathcal{Z}_e \cap \mathbb{R} = \emptyset$. Hence, $\mathbb{U}_{z,s}(\mathbf{k}, t) = 0$ in (5.11) and one can skip the Section 5.3.3 devoted to the estimate of of this term.

We finally point out all our results still hold in the ‘‘symmetric case’’ (more rare in physical models) where $N_e = 0$ and $N_m \geq 1$ with obvious changes compared to the above discussion.

7.3.2 The case of bounded domains

In [5] (Section 4.1), we show how to generalize our Lyapunov estimates technique to the case of bounded propagation domains under the strong dissipation assumption (1.9). Similarly, we show here how to extend our modal approach to bounded electromagnetic cavities but under the weak dissipation assumption (1.11).

One considers here the Maxwell's evolution system (1.1,1.2) in a bounded Lipschitz-continuous domain $\Omega \subset \mathbb{R}^3$ completed with perfectly conducting boundary conditions

$$\mathbf{E} \times \mathbf{n} = 0, \quad \text{where } \mathbf{n} \text{ is the outward normal vector to } \partial\Omega. \quad (7.18)$$

We introduce the functional spaces (see e.g. [7])

$$\begin{aligned} H(\text{rot}; \Omega) &= \{\mathbf{u} \in \mathbf{L}^2(\Omega) \mid \nabla \times \mathbf{u} \in \mathbf{L}^2(\Omega)\}, & H_0(\text{rot}; \Omega) &= \{\mathbf{u} \in H(\text{rot}; \Omega) \mid \mathbf{u} \times \mathbf{n} = 0 \text{ on } \partial\Omega\}, \\ H(\text{div } 0; \Omega) &= \{\mathbf{u} \in \mathbf{L}^2(\Omega) \mid \nabla \cdot \mathbf{u} = 0\}, & H_0(\text{div } 0; \Omega) &= \{\mathbf{u} \in H(\text{div } 0; \Omega) \mid \mathbf{u} \cdot \mathbf{n} = 0 \text{ on } \partial\Omega\}. \end{aligned}$$

In this setting, one can rewrite the Maxwell's system as the evolution equation (1.13) with the Hilbert space $\mathcal{H} = \mathbf{L}^2(\Omega)^N$ (where $\mathbf{L}^2(\Omega) = L^2(\Omega)^3$) and a maximal dissipative unbounded operator $\mathbb{A} : D(\mathbb{A}) \subset \mathcal{H} \rightarrow \mathcal{H}$ defined by the relation (1.16) but with domain

$$D(\mathbb{A}) = H_0(\text{rot}; \Omega) \times H(\text{rot}; \Omega) \times \mathbf{L}^2(\Omega)^{N_e} \times \mathbf{L}^2(\Omega)^{N_e} \times \mathbf{L}^2(\Omega)^{N_m} \times \mathbf{L}^2(\Omega)^{N_m}.$$

Similarly, one introduces a reduced evolution equation on the Hilbert space

$$\mathcal{H}_\perp = H(\text{div } 0; \Omega) \times H_0(\text{div } 0; \Omega) \times H(\text{div } 0; \Omega)^{N_e} \times H(\text{div } 0; \Omega)^{N_e} \times H_0(\text{div } 0; \Omega)^{N_m} \times H_0(\text{div } 0; \Omega)^{N_m}$$

with a reduced maximal dissipative operator $\mathbb{A}_\perp : D(\mathbb{A}_\perp) \subset \mathcal{H}_\perp \rightarrow \mathcal{H}_\perp$ defined on its domain $D(\mathbb{A}_\perp) = \mathcal{H}_\perp \cap D(\mathbb{A})$ (see (1.22, 1.23)). We point out that $\ker(\mathbb{A}_\perp)$ has a finite dimension. Moreover, \mathbb{A}_\perp is injective (or equivalently $\ker(\mathbb{A}_\perp)^\perp = \mathcal{H}_\perp$) if Ω is simply connected.

As in [5], the main difference in the analysis, with respect to the case of \mathbb{R}^3 , is that the use of the Fourier transform in space has to be replaced by an adequate modal expansion. More precisely, we introduce the closed subspace of $\mathbf{L}^2(\Omega) \times \mathbf{L}^2(\Omega)$ given by

$$\mathcal{H}_0 := H(\operatorname{div} 0; \Omega) \times H_0(\operatorname{div} 0; \Omega).$$

Then (see e.g. [7] chapter IX), the operator $\mathcal{M} : D(\mathcal{M}) \subset \mathcal{H}_0 \mapsto \mathcal{H}_0$, defined by,

$$\mathcal{M} \begin{pmatrix} \mathbf{u} \\ \mathbf{v} \end{pmatrix} = \mathbf{i} \begin{pmatrix} \nabla \times \mathbf{v} \\ -\nabla \times \mathbf{u} \end{pmatrix}, \quad \forall (\mathbf{u}, \mathbf{v}) \in D(\mathcal{M}) := \mathcal{H}_0 \cap (H_0(\operatorname{rot}; \Omega) \times H(\operatorname{rot}; \Omega)),$$

is self-adjoint. One proves, see e.g. [7] chapter IX, that, as $D(\mathcal{M})$ is compactly embedded in $\mathbf{L}^2(\Omega) \times \mathbf{L}^2(\Omega)$, \mathcal{M} has a compact resolvent and a finite dimensional kernel (which reduces to $\{0\}$ if Ω is simply connected). From the theory of self-adjoint operators with compact resolvent, and the symmetries of Maxwell's equations, which imply that the spectrum of \mathcal{M} is symmetric with respect to the origin, one can construct a countable infinity of normalized modes indexed by $p \in \mathbb{N}^*$:

$$(\pm k_p, (\mathbf{u}_p^\pm, \mathbf{v}_p^\pm)) \in \mathbb{R} \times D(\mathcal{M}) \mid \mathcal{M} \begin{pmatrix} \mathbf{u}_p^\pm \\ \mathbf{v}_p^\pm \end{pmatrix} = \pm k_p \begin{pmatrix} \mathbf{u}_p^\pm \\ \mathbf{v}_p^\pm \end{pmatrix} \text{ with } k_p > 0$$

where (k_p) is an increasing sequences satisfying $k_p \rightarrow +\infty$ (as $p \rightarrow \infty$). These cavity modes form an orthonormal basis of $\ker(\mathcal{M})^\perp$. Using this basis, one decomposes the components of the solution $\mathbf{U} = (\mathbf{E}, \mathbf{H}, \mathbf{P}, \dot{\mathbf{P}}, \mathbf{M}, \dot{\mathbf{M}})$ of the equation (1.13) with an initial condition $\mathbf{U}_0 \in \mathcal{H}_\perp$ as follows:

$$\mathbf{E}(\cdot, t) = \sum_{\pm} \sum_{p=1}^{+\infty} \mathbb{E}(\pm k_p; t) \mathbf{u}_p^\pm, \quad \mathbf{H}(\cdot, t) = \sum_{\pm} \sum_{p=1}^{+\infty} \mathbb{H}(\pm k_p; t) \mathbf{v}_p^\pm,$$

and the auxiliary field \mathbf{P}_j and \mathbf{M}_ℓ accordingly

$$\mathbf{P}_j(\cdot, t) = \sum_{\pm} \sum_{p=1}^{+\infty} \mathbb{P}_j(\pm k_p; t) \mathbf{u}_p^\pm, \quad \mathbf{M}_\ell(\cdot, t) = \sum_{\pm} \sum_{p=1}^{+\infty} \mathbb{M}_\ell(\pm k_p; t) \mathbf{v}_p^\pm.$$

Injecting these modal expansions in the equation (1.13) yields that

$$\mathbf{U}(\pm k_p; t) = (\mathbb{E}(\pm k_p; t), \mathbb{H}(\pm k_p; t), \mathbb{P}(\pm k_p; t), \dot{\mathbb{P}}(\pm k_p; t), \mathbb{M}_\ell(\pm k_p; t), \dot{\mathbb{M}}_\ell(\pm k_p; t)) \in \mathbb{C}^N$$

satisfies the following system of ODE's indexed by $\pm k_p$ (which plays the same role here as the Fourier parameter \mathbf{k}):

$$\frac{d\mathbf{U}(\pm k_p, t)}{dt} + \mathbf{i} \mathbb{A}_{\perp, \pm k_p} \mathbf{U}(\pm k_p, t) = 0 \quad \text{for } t \geq 0 \text{ with } \mathbf{U}(\pm k_p, 0) = \mathbf{U}_0(\pm k_p).$$

where $\mathbb{A}_{\perp, \pm k_p} \in \mathcal{L}(\mathbb{C}^N)$ is defined for all $\mathbf{U} = (\mathbb{E}, \mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}}) \in \mathbb{C}^N$ as $\mathbb{A}_{\mathbf{k}}$ in (2.2) but by substituting in the two first lines $-(\mathbf{k} \times \mathbb{H})/\varepsilon_0$ and $(\mathbf{k} \times \mathbb{E})/\mu_0$ by $\pm k_p \mathbb{H}/\mu_0$ and $\pm k_p \mathbb{E}/\varepsilon_0$.

One also remarks on (2.2) that \mathbb{A}_{\perp, k_p} and $\mathbb{A}_{\perp, -k_p}$ are unitarily equivalent via the unitary map

$$\mathcal{R} : (\mathbb{E}, \mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}}) \in \mathbb{C}^N \longmapsto (\mathbb{E}, -\mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, -\mathbb{M}, -\dot{\mathbb{M}}) \in \mathbb{C}^N.$$

Thus, one only has to estimate the solution $e^{-i\mathbb{A}_{\perp, k_p} t} \mathbb{U}_0(k_p)$ of (7.3.2) for $k_p > 0$. A computation similar to that in the Appendix A.1.2 shows the spectrum $\sigma(\mathbb{A}_{k_p})$ is given by the set of solutions of the dispersion relation

$$\omega^2 \varepsilon(\omega) \mu(\omega) = k_p^2.$$

As the set $\{k_p, p \in \mathbb{N}^*\}$ is discrete and contains only non-zero elements, one does not need to derive low frequencies estimates (as in Section 5) and the exponential uniform decay of the mid-range frequencies is trivial since there are only finitely many of them. Thus, one only estimates the high frequencies components $\mathbb{U}(\pm k_p, t)$ for $k_p \gg 1$. To this aim, one uses the results of Sections 3 and 4 by replacing $|\mathbf{k}|$ by k_p .

Let m be a positive integer. One defines the Hilbert space $D(\mathbb{A}^m)$ by $\{\mathbf{U} \in \mathcal{H} \mid \mathbb{A}^k \mathbf{U} \in \mathcal{H} \text{ for } k = 1, \dots, m\}$ endowed with its usual norm $\|\cdot\|_{D(\mathbb{A}^m)}$ defined by

$$\|\mathbf{U}\|_{D(\mathbb{A}^m)}^2 = \sum_{k=0}^m \|\mathbb{A}^k \mathbf{U}\|_{\mathcal{H}}^2, \quad \forall \mathbf{U} \in D(\mathbb{A}^m).$$

Then, the main result is as follows.

Theorem 7.1. *Let assume that the Maxwell's system (1.1, 1.2, 7.18) satisfies the weak dissipation assumption (1.11) and the irreducibility assumptions (\mathbf{H}_1) and (\mathbf{H}_2) . Then, for any initial condition $\mathbf{U}_0 \in \mathcal{H}_{\perp}$ of this system, one has*

$$\lim_{t \rightarrow +\infty} \|\mathbf{U}(t)\|_{\mathcal{H}}^2 = 0. \quad (7.19)$$

In addition, if the initial data satisfies

$$\mathbf{U}_0 \in D(\mathbb{A}^m) \cap \ker(\mathbb{A}_{\perp})^{\perp}, \quad \text{for any integer } m > 0 \quad (\text{see remark 7.2}),$$

then, in the non-critical (resp. critical) case, see definition 1.8, one has, for any $t > 0$,

$$\|\mathbf{U}(t)\|_{\mathcal{H}}^2 \lesssim \frac{\|\mathbf{U}_0\|_{D(\mathbb{A}^m)}^2}{t^m}, \quad \left(\text{resp. } \|\mathbf{U}(t)\|_{\mathcal{H}}^2 \lesssim \frac{\|\mathbf{U}_0\|_{D(\mathbb{A}^m)}^2}{t^{\frac{m}{2}}} \right). \quad (7.20)$$

Moreover, the decay estimate (7.20) are optimal since in the non-critical case

$$m = \sup\{\gamma \in \mathbb{R}^+ \mid \forall \mathbf{U}_0 \in D(\mathbb{A}^m) \cap \ker(\mathbb{A}_{\perp})^{\perp}, \exists C(\mathbf{U}_0) > 0 \mid \|\mathbf{U}(t)\|_{\mathcal{H}}^2 \lesssim \frac{C(\mathbf{U}_0)}{t^{\gamma}}, \forall t \geq 1\}$$

and in the critical case, the above equality holds with m replaced by $m/2$ in the left hand side.

Remark 7.2. For C^{∞} domains Ω , the condition $\mathbf{U}_0 \in D(\mathbb{A}^m) \cap \ker(\mathbb{A}_{\perp})^{\perp}$ (see e.g. [7], chapter IX) is equivalent to $\mathbf{U}_0 \in \mathbf{H}^m(\Omega)^N \cap \ker(\mathbb{A}_{\perp})^{\perp}$ for any integer $m > 0$ and the norms $\|\cdot\|_{D(\mathbb{A}^m)}$ and $\|\cdot\|_{\mathbf{H}^m(\Omega)^N}$ on these spaces are equivalent.

Appendix A.

Appendix A.1. Spectrum and resolvent of $\mathbb{A}_{|\mathbf{k}|,\perp}$

Appendix A.1.1. Proof of Proposition 3.2: resolvent of $\mathbb{A}_{|\mathbf{k}|,\perp}$

Given $\mathbb{F} = (\mathbf{e}, \mathbf{h}, \mathbf{p}, \dot{\mathbf{p}}, \mathbf{m}, \dot{\mathbf{m}})^t \in \mathbf{C}_\perp^N$ and $\omega \in \mathbb{C} \setminus (S(|\mathbf{k}|) \cap \mathcal{S}_T)$, we look at the problem

$$\text{Find } \mathbb{U} = (\mathbb{E}, \mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}}) \in \mathbf{C}_\perp^N \text{ such that } (\mathbb{A}_{|\mathbf{k}|,\perp} - \omega \mathbb{I}) \mathbb{U} = \mathbb{F}, \quad (\text{A.1})$$

which is equivalent to finding $\mathbb{U} = (\mathbb{E}, \mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}}) \in \mathbf{C}_\perp^N$ solution of the system

$$\begin{cases} -\frac{|\mathbf{k}| \mathbf{e}_3 \times \mathbb{H}}{\varepsilon_0} - i \sum \Omega_{e,j}^2 \dot{\mathbb{P}}_j - \omega \mathbb{E} = \mathbf{e}, & (i) \\ \frac{|\mathbf{k}| \mathbf{e}_3 \times \mathbb{E}}{\mu_0} - i \sum \Omega_{m,\ell}^2 \dot{\mathbb{M}}_\ell - \omega \mathbb{H} = \mathbf{h}, & (ii) \end{cases} \quad (\text{A.2})$$

$$\begin{cases} i \dot{\mathbb{P}}_j - \omega \mathbb{P}_j = \mathbf{p}_j, & (i) \\ -i \alpha_{e,j} \dot{\mathbb{P}}_j - i \omega_{e,j}^2 \mathbb{P}_j + i \mathbb{E} - \omega \dot{\mathbb{P}}_j = \dot{\mathbf{p}}_j, & (ii) \end{cases} \quad (\text{A.3})$$

$$\begin{cases} i \dot{\mathbb{M}}_\ell - \omega \mathbb{M}_\ell = \mathbf{m}_\ell, & (i) \\ -i \alpha_{m,\ell} \dot{\mathbb{M}}_\ell - i \omega_{m,\ell}^2 \mathbb{M}_\ell + i \mathbb{H} - \omega \dot{\mathbb{M}}_\ell = \dot{\mathbf{m}}_\ell. & (ii) \end{cases} \quad (\text{A.4})$$

We want to show that this problem admits a unique solution \mathbb{U} defining $\mathbb{U} = R_{|\mathbf{k}|}(\omega) \mathbb{F}$. The proof of Proposition 3.2 is thus purely computational and reduces to solving explicitly the system (A.2, A.3, A.4). We provide some details for the ease of the reader.

Substituting the expression of $\dot{\mathbb{P}}_j$ and $\dot{\mathbb{M}}_\ell$ given by (A.3)(i) and (A.4)(i), i.e.

$$\dot{\mathbb{P}}_j = -i \omega \mathbb{P}_j - i \mathbf{p}_j, \quad \dot{\mathbb{M}}_\ell = -i \omega \mathbb{M}_\ell - i \mathbf{m}_\ell \quad (\text{A.5})$$

into (A.3)(ii) and (A.4)(ii), we get, by definition of $q_{e,j}(\omega)$ and $q_{m,\ell}(\omega)$ and by definition (3.6) of the operators $\mathbb{A}_{e,j}(\omega), \mathbb{A}_{m,\ell}(\omega)$ (the division by $q_{e,j}(\omega)$ and $q_{m,\ell}(\omega)$ is allowed because $\omega \notin \mathcal{P}$),

$$\mathbb{P}_j = -\frac{\mathbb{E}}{q_{e,j}(\omega)} + \mathbb{A}_{e,j}(\omega) \mathbb{F} \quad \text{and} \quad \mathbb{M}_\ell = -\frac{\mathbb{H}}{q_{m,\ell}(\omega)} + \mathbb{A}_{m,\ell}(\omega) \mathbb{F}. \quad (\text{A.6})$$

Thus, going back to $\dot{\mathbb{P}}_j$ and $\dot{\mathbb{M}}_\ell$, via (A.3)(i) and (A.4)(i), gives by definition (3.6) of $\dot{\mathbb{A}}_{e,j}(\omega), \dot{\mathbb{A}}_{m,\ell}(\omega)$:

$$\dot{\mathbb{P}}_j = \frac{i \omega \mathbb{E}}{q_{e,j}(\omega)} + \dot{\mathbb{A}}_{e,j}(\omega) \mathbb{F} \quad \text{and} \quad \dot{\mathbb{M}}_\ell = \frac{i \omega \mathbb{H}}{q_{m,\ell}(\omega)} + \dot{\mathbb{A}}_{m,\ell}(\omega) \mathbb{F}. \quad (\text{A.7})$$

Substituting $\dot{\mathbb{P}}_j$, from (A.7), into (A.2)(ii) we get, by definitions (1.6, 3.6) of $\mu(\omega)$ and $\mathbb{A}_m(\omega)$:

$$|\mathbf{k}| \mathbf{e}_3 \times \mathbb{E} - \omega \mu(\omega) \mathbb{H} = \mu_0 \left(\mathbf{h} + i \sum \Omega_{m,\ell}^2 \dot{\mathbb{A}}_{m,\ell}(\omega) \mathbb{F} \right) = -\mathbb{A}_m(\omega) \mathbb{F}. \quad (\text{A.8})$$

Similarly, from $\dot{\mathbb{M}}_\ell$ in (A.7), we obtain, by definitions (1.6, 3.6) of $\varepsilon(\omega)$ and $\mathbb{A}_e(\omega)$,

$$-|\mathbf{k}| \mathbf{e}_3 \times \mathbb{H} - \omega \varepsilon(\omega) \mathbb{E} = \varepsilon_0 \left(\mathbf{e} + i \sum \Omega_{e,j}^2 \dot{\mathbb{A}}_{e,j}(\omega) \mathbb{F} \right) = -\mathbb{A}_e(\omega) \mathbb{F}. \quad (\text{A.9})$$

In order to eliminate \mathbb{H} between (A.8) and (A.9), we perform the combination

$$|\mathbf{k}| \mathbf{e}_3 \times (\text{A.8}) - \omega \mu(\omega) (\text{A.9}).$$

This gives, since $-|\mathbf{k}| \mathbf{e}_3 \times (|\mathbf{k}| \mathbf{e}_3 \times \mathbb{E}) = |\mathbf{k}|^2 \mathbb{E}$, (as $\mathbb{U} \in \mathbf{C}_\perp^N$, \mathbf{e}_3 and \mathbb{E} are orthogonal)

$$(\mathcal{D}(\omega) - |\mathbf{k}|^2) \mathbb{E} = \omega \mu(\omega) \mathbb{A}_e(\omega) \mathbb{F} - |\mathbf{k}| \mathbf{e}_3 \times \mathbb{A}_m(\omega) \mathbb{F}. \quad (\text{A.10})$$

As $\omega \notin S(|\mathbf{k}|)$, one has $\mathcal{D}(\omega) - |\mathbf{k}|^2 \neq 0$ and therefore, by definition (3.10) of $\mathcal{S}_{|\mathbf{k}|}(\omega)$,

$$\mathbb{E} = \frac{-|\mathbf{k}| \mathbf{e}_3 \times \mathbb{A}_m(\omega) \mathbb{F} + \omega \mu(\omega) \mathbb{A}_e(\omega) \mathbb{F}}{\mathcal{D}(\omega) - |\mathbf{k}|^2} = \mathcal{S}_{|\mathbf{k}|}(\omega) \mathbb{F}. \quad (\text{A.11})$$

Thus, since $\omega \mu(\omega) \neq 0$ (as $\omega \notin \mathcal{Z}_m \cup \{0\}$), from (A.8) and by definition (3.7) of $\mathbb{A}_h(\omega)$:

$$\mathbb{H} = \frac{1}{\omega \mu(\omega)} \left(|\mathbf{k}| \mathbf{e}_3 \times \mathcal{S}_{|\mathbf{k}|}(\omega) \mathbb{F} + \mathbb{A}_m(\omega) \mathbb{F} \right). \quad (\text{A.12})$$

Finally, substituting (A.10, A.12) in (A.6, A.7) we get the expression of the other

$$\begin{cases} \mathbb{P}_j = -\frac{\mathcal{S}_{|\mathbf{k}|}(\omega) \mathbb{F}}{q_{e,j}(\omega)} + \mathbb{A}_{e,j}(\omega) \mathbb{F}, & \mathbb{M}_\ell = \mathbb{A}_{m,\ell}(\omega) \mathbb{F} - \frac{\mathbb{A}_m(\omega) \mathbb{F} + |\mathbf{k}| \mathbf{e}_3 \times \mathcal{S}_{|\mathbf{k}|}(\omega) \mathbb{F}}{\omega \mu(\omega) q_{m,\ell}(\omega)}, \\ \dot{\mathbb{P}}_j = \frac{i \omega \mathcal{S}_{|\mathbf{k}|}(\omega) \mathbb{F}}{q_{e,j}(\omega)} + \dot{\mathbb{A}}_{e,j}(\omega) \mathbb{F}, & \dot{\mathbb{M}}_\ell = \dot{\mathbb{A}}_{m,\ell}(\omega) \mathbb{F} + i \frac{\mathbb{A}_m(\omega) \mathbb{F} + |\mathbf{k}| \mathbf{e}_3 \times \mathcal{S}_{|\mathbf{k}|}(\omega) \mathbb{F}}{\mu(\omega) q_{m,\ell}(\omega)}. \end{cases} \quad (\text{A.13})$$

The reader will then easily verify that the formulas (A.10, A.12, A.13) lead to the expression (3.9, 3.10, 3.10, 3.12) given in Proposition 3.2 for the resolvent $R_{|\mathbf{k}|}(\omega)$.

Appendix A.1.2. The Spectrum of $\mathbb{A}_{|\mathbf{k}|,\perp}$: proof of Proposition 3.1

From Proposition 3.2, we already know that

$$\sigma(\mathbb{A}_{|\mathbf{k}|,\perp}) \subset S(|\mathbf{k}|) \cup \mathcal{S}_\mathcal{T}.$$

To prove the Proposition 3.1, it suffices to show that

- (i) $\sigma(\mathbb{A}_{|\mathbf{k}|,\perp}) \cap \mathcal{S}_\mathcal{T} = \emptyset$.
- (ii) $S(|\mathbf{k}|) \subset \sigma(\mathbb{A}_{|\mathbf{k}|,\perp})$ and for any $\omega \in S(|\mathbf{k}|)$, $\dim \ker (\mathbb{A}_{|\mathbf{k}|,\perp} - \omega \text{Id}) = 2$.

Preamble. For both steps, we shall use the fact that $\mathbb{U} = (\mathbb{E}, \mathbb{H}, (\mathbb{P}_j), (\dot{\mathbb{P}}_j), (\mathbb{M}_\ell), (\dot{\mathbb{P}}_\ell)) \in \ker (\mathbb{A}_{|\mathbf{k}|,\perp} - \omega \text{Id})$ means that \mathbb{U} is a solution of (A.2, A.3, A.4) with $\mathbb{F} = (\mathbf{e}, \mathbf{h}, \mathbf{p}, \dot{\mathbf{p}}, \mathbf{m}, \dot{\mathbf{m}}) = 0$. Therefore, proceeding as for obtaining (A.6) and (A.7), before division by $q_{e,j}(\omega)$ and $q_{m,\ell}(\omega)$, we deduce from (A.3, A.4) that

$$q_{e,j}(\omega) \mathbb{P}_j = -\mathbb{E}, \quad q_{m,\ell}(\omega) \mathbb{M}_\ell = -\mathbb{H}, \quad q_{e,j}(\omega) \dot{\mathbb{P}}_j = i \omega \mathbb{E}, \quad q_{m,\ell}(\omega) \dot{\mathbb{M}}_\ell = i \omega \mathbb{H}. \quad (\text{A.14})$$

Proof of Step (i). Assume that $\omega \in \mathcal{S}_\mathcal{T} = \mathcal{P}_e \cup \mathcal{P}_m \cup \mathcal{Z}_m \cup \{0\}$.

(a) Let $\omega \in \mathcal{P}_e$ and $\mathbb{U} \in \ker(\mathbb{A}_{|\mathbf{k}|,\perp} - \omega \text{Id})$, for some j_0 , $q_{e,j_0}(\omega) = 0$, thus, by (A.14), $\mathbb{E} = 0$. Thus, by (A.14) and assumption (H₁) that yields $q_{e,j}(\omega) \neq 0$ for $j \neq j_0$, we conclude that $\mathbb{P}_j = 0$ and $\dot{\mathbb{P}}_j = 0$ for any $j \neq j_0$.

Next, according to assumption (H₁) again,

- Either $q_{m,\ell}(\omega) \neq 0$ for any m , we conclude from (A.14) $\mathbb{M} = 0, \dot{\mathbb{M}} = 0$. Thus, going back to (A.2)(ii) with $\mathbf{h} = 0$, we have $\mathbb{H} = 0$ (since $\omega \neq 0$). Then, from (A.2)(i) with $\mathbf{e} = 0$, we have $\dot{\mathbb{P}}_{j_0} = 0$. Thus, one gets $\mathbb{P}_{j_0} = 0$ since, by (A.5) for $\mathbf{p} = 0$, $-i\omega \mathbb{P}_{j_0} = \dot{\mathbb{P}}_{j_0}$ and $\omega \neq 0$ (as $0 \notin \mathcal{P}_e$).
- Or there exists a unique ℓ_0 for which $q_{m,\ell}(\omega) = 0$. In that case (A.14) implies that $\mathbb{H} = 0$ and $\mathbb{M}_\ell = \dot{\mathbb{M}}_\ell = 0$ for any $\ell \neq \ell_0$. Finally, going back to (A.2) (i) and (ii) with $\mathbf{e} = \mathbf{h} = 0$, we have $\dot{\mathbb{P}}_{j_0} = \dot{\mathbb{M}}_{\ell_0} = 0$, thus $\mathbb{P}_{\ell_0} = \mathbb{M}_{\ell_0} = 0$ since (A.5) for $\mathbf{m} = 0$ and $\mathbf{p} = 0$ gives $-i\omega \mathbb{P}_{j_0} = \dot{\mathbb{P}}_{j_0}$ and $-i\omega \mathbb{M}_{\ell_0} = \dot{\mathbb{M}}_{\ell_0}$ with $\omega \neq 0$.

In all cases, $\mathbb{U} = 0$, thus $\omega \notin \sigma(\mathbb{A}_{|\mathbf{k}|,\perp})$.

(b) In a symmetric manner, we prove that if $\omega \in \mathcal{P}_m$, $\omega \notin \sigma(\mathbb{A}_{|\mathbf{k}|,\perp})$.

Now, for the rest of the proof, we point out that for $\omega \notin \mathcal{P}$, one shows by proceeding exactly as for obtaining (A.8) and (A.9) that

$$(1) \quad |\mathbf{k}| \mathbf{e}_3 \times \mathbb{E} - \omega \mu(\omega) \mathbb{H} = 0, \quad (2) \quad -|\mathbf{k}| \mathbf{e}_3 \times \mathbb{H} - \omega \varepsilon(\omega) \mathbb{E} = 0. \quad (\text{A.15})$$

More precisely, one checks that as soon as $\omega \notin \mathcal{P}$, (A.14, A.15) is equivalent to (A.2, A.3, A.4) for $\mathbb{F} = 0$.

(c) It remains to look at $\omega \in \mathcal{Z}_m \cup \{0\}$, i.e. $\omega \mu(\omega) = 0$. By (H₂), one has $\mathcal{P} \cap (\mathcal{Z}_m \cup \{0\}) = \emptyset$. Thus, from (A.15)(1) and since $\mathbb{E} \cdot \mathbf{e}_3 = 0$, we deduce that $\mathbb{E} = 0$ and it follows from (A.15)(2) that $\mathbb{H} = 0$ (as $\mathbb{H} \cdot \mathbf{e}_3 = 0$). Then, going back to (A.14), since $\omega \notin \mathcal{P}$, we deduce that $\mathbb{P} = \dot{\mathbb{P}} = 0$ and $\mathbb{M} = \dot{\mathbb{M}} = 0$, thus $\mathbb{U} = 0$. This proves that $\omega \notin \sigma(\mathbb{A}_{|\mathbf{k}|,\perp})$.

Proof of Step (ii). Let $\omega \in S(|\mathbf{k}|)$ and $\mathbb{U} \in \ker(\mathbb{A}_{|\mathbf{k}|,\perp} - \omega \text{Id})$. First note that, since by assumption $\mathbf{k} \neq 0$ and thus $\omega \mu(\omega) \neq 0$ (since ω satisfies the dispersion relation (3.4)), according to (A.15)(a), we have

$$\mathbb{H} = \frac{|\mathbf{k}| \mathbf{e}_3 \times \mathbb{E}}{\omega \mu(\omega)}, \quad (\text{A.16})$$

so that, it yields with (A.14):

$$\mathbb{P}_j = -\frac{\mathbb{E}}{q_{e,j}(\omega)}, \quad \dot{\mathbb{P}}_j = \frac{i\omega \mathbb{E}}{q_{e,j}(\omega)}, \quad \mathbb{M}_\ell = -\frac{|\mathbf{k}| \mathbf{e}_3 \times \mathbb{E}}{\omega \mu(\omega) q_{m,\ell}(\omega)}, \quad \dot{\mathbb{M}}_\ell = \frac{i\omega |\mathbf{k}| \mathbf{e}_3 \times \mathbb{E}}{\omega \mu(\omega) q_{m,\ell}(\omega)}. \quad (\text{A.17})$$

Of course (A.16) and (A.17) mean that $\dim \ker(\mathbb{A}_{|\mathbf{k}|,\perp} - \omega \text{Id}) \leq \dim \mathbf{C}_\perp = 2$.

To prove (ii), it suffices to check that $\forall \mathbb{E} \in \mathbf{C}_\perp$, $\mathbb{U} = (\mathbb{E}, \mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}}) \in \ker(\mathbb{A}_{|\mathbf{k}|,\perp} - \omega \text{Id})$,

when $\omega \in S(|\mathbf{k}|)$ and $(\mathbb{H}, \mathbb{P}, \dot{\mathbb{P}}, \mathbb{M}, \dot{\mathbb{M}})$ is given by (A.16, A.17). Thus, it consists to show that such vectors \mathbb{U} are solutions to (A.2, A.3, A.4) with $\mathbb{F} = 0$. This is equivalent (since $\omega \notin \mathcal{P}$) to checking (A.14, A.15). This is rather immediate since (A.17) is nothing but (A.14) with \mathbb{H} given by (A.16) while (A.16) is nothing but (A.15)(1). It remains to check (A.15)(2). However, substituting (A.16) into the left hand side of (A.15)(2) gives, as $-|\mathbf{k}| \mathbf{e}_3 \times (|\mathbf{k}| \mathbf{e}_3 \times \mathbb{E}) = |\mathbf{k}|^2 \mathbb{E}$ for \mathbb{E} in \mathbf{C}_\perp ,

$$(|\mathbf{k}|^2 - \mathcal{D}(\omega)) \mathbb{E} = 0, \quad (\text{A.18})$$

which is true for any $\mathbb{E} \in \mathbf{C}_\perp$ since $\mathcal{D}(\omega) = |\mathbf{k}|^2$ for $\omega \in S(|\mathbf{k}|)$.

Appendix A.2. Technical result for the analysis of the dispersion curves

Let $z \in \mathbb{C}$. We are interested in the solutions $\omega \in \mathbb{C}$ of the nonlinear equation

$$(\omega - z)^{\mathbf{m}} g(\omega) = \zeta^{\mathbf{m}}, \quad (\text{A.19})$$

parametrized by ζ , where g is a given analytic function near z satisfying $g(z) \neq 0$. More precisely, we want to prove that (A.19) implicitly defines ω in function of ζ , via \mathbf{m} analytic branches of solutions $\omega_n(\zeta)$, for small $|\zeta|$.

Lemma A.1. *Let \mathcal{G} be an analytic function on a domain $\Omega \subset \mathbb{C}$ and $z \in \Omega$ a zero of multiplicity $\mathbf{m} \in \mathbb{N}^*$ of \mathcal{G} , so that \mathcal{G} can be rewritten as*

$$\mathcal{G}(\omega) = (\omega - z)^{\mathbf{m}} g(\omega) \text{ with } g \text{ analytic in a vicinity of } z \text{ and } g(z) = A \neq 0. \quad (\text{A.20})$$

Then, there exist an open neighborhood U of $\zeta = 0$ and \mathbf{m} distinct analytic functions $\zeta \mapsto \omega_n(\zeta)$ for $n = 1 \dots, \mathbf{m}$ defined on U such that

$$\mathcal{G}(\omega_n(\zeta)) = \zeta^{\mathbf{m}} \quad \text{on } U. \quad (\text{A.21})$$

Moreover the functions $\omega_n(\zeta)$ have the following Taylor expansion

$$\omega_n(\zeta) = z + a_n^{-1} \zeta - \frac{a_n^{-2} g'(z)}{\mathbf{m} g(z)} \zeta^2 + O(\zeta^3), \text{ as } \zeta \rightarrow 0, \text{ where } a_n = |A|^{1/\mathbf{m}} e^{i \frac{\theta}{\mathbf{m}}} e^{2i \frac{n\pi}{\mathbf{m}}}, \quad (\text{A.22})$$

$\theta \in (-\pi, \pi]$ is the principal value of the argument of A and the complex coefficients a_n are the \mathbf{m} distinct roots of the polynomial equation $X^{\mathbf{m}} = A$.

Furthermore, one can find an open neighborhood D of z included in Ω such that the set of solutions of $\mathcal{G}(\omega) = \zeta^{\mathbf{m}}$ in D for $\zeta \in U$ is exactly given by $\{\omega_n(\zeta) \text{ for } n = 1, \dots, \mathbf{m}\}$.

Proof. We want to find branches of solutions $\zeta \mapsto \omega_n(\zeta)$ of the equation

$$\mathcal{G}(\omega) = \zeta^{\mathbf{m}}, \quad (\text{A.23})$$

where ζ is a parameter that lies in an open neighborhood of 0. To this aim, fix $1 \leq n \leq \mathbf{m}$ and make the following change of unknown $\omega \rightarrow \eta$, for $\zeta \neq 0$:

$$\eta = -1 + (\omega - z) a_n \zeta^{-1} \iff \omega = z + a_n^{-1} \zeta (1 + \eta) \quad (\text{A.24})$$

where the coefficient a_n is given by (A.22). Therefore by replacing (A.24) into (A.23) and using (A.20) and the fact that $g(z) = A = a_n^{\mathfrak{m}}$, we notice that on an open neighborhood of $\zeta = 0$, finding an analytic branches $\zeta \mapsto \omega_n(\zeta)$ of solutions of (A.23) is equivalent to find an analytic branch $\zeta \mapsto \eta_n(\zeta)$ satisfying

$$H_n(\zeta, \eta) = 0 \quad \text{with} \quad H_n(\zeta, \eta) := (1 + \eta)^{\mathfrak{m}} g(z + a_n^{-1} \zeta (1 + \eta)) - g(z), \quad (\text{A.25})$$

with ω_n given in term of η_n by (A.24). As g is analytic in a vicinity of z , by the Hartogs Theorem (see [19], Theorem 36.8 page 271), the function $H_n(\zeta, \eta)$ is well-defined and analytic in a neighbourhood of $(0, 0)$ in $\mathbb{C} \times \mathbb{C}$ and satisfies $H_n(0, 0) = 0$. We then want to use the analytic implicit function Theorem (see [14], Theorem 7.6 page 34) to solve $H_n(\zeta, \eta) = 0$. Thus, one first has to check that $\partial_\eta H_n(0, 0) \neq 0$. Indeed, from (A.25), we compute

$$\partial_\eta H_n(\zeta, \eta) = \mathfrak{m} (1 + \eta)^{\mathfrak{m}-1} g(z + a_n^{-1} \zeta (1 + \eta)) + (1 + \eta)^{\mathfrak{m}} a_n^{-1} \zeta g'(z + a_n^{-1} \zeta (1 + \eta))$$

so that in particular

$$\partial_\eta H_n(0, 0) = \mathfrak{m} g(z) \neq 0. \quad (\text{A.26})$$

Hence, the analytic implicit function Theorem proves the existence of two open neighbourhoods of the origin $U_n \subset \mathbb{C}$ (the ζ -complex plane) and $V_n \subset \mathbb{C}$ (the η -complex plane) and of an analytic function $\eta_n(\zeta) : U_n \rightarrow V_n$ satisfying $\eta(0) = 0$ such that

$$\{(\zeta, \eta) \in U_n \times V_n \mid H_n(\zeta, \eta_n) = 0\} = \{(\zeta, \eta_n(\zeta)), \zeta \in U_n\}.$$

Moreover, one has

$$\eta'(0) = -\frac{\partial_\zeta H_n(0, 0)}{\partial_\eta H_n(0, 0)} = -\frac{a_n^{-1} g'(z)}{\mathfrak{m} g(z)}. \quad (\text{A.27})$$

Thus $\zeta \mapsto \omega_n(\zeta) := z + a_n^{-1} \zeta (1 + \eta_n(\zeta))$ is an analytic branch of solution of (A.23) in U_n . Since $\eta_n(0) = 0$, as $\zeta \rightarrow 0$, $\eta_n(\zeta) = \zeta \eta'_n(0) + O(\zeta^2)$, which, put in the definition of $\omega_n(\zeta)$, gives (A.22).

We point out that, as the coefficients a_n^{-1} in (A.22) are all distinct, we have constructed \mathfrak{m} distinct analytic functions $\zeta \rightarrow \omega_n(\zeta)$ for $n = 1 \dots, \mathfrak{m}$ defined on the open neighborhood \tilde{U} , given by the intersection of the sets U_n for $n \in \{1, \dots, \mathfrak{m}\}$, which satisfy (A.21) and (A.22).

The last point is proven with Rouché Theorem (see for e.g. Theorem 10.43 page 225 of [26]). As g does not vanish in the vicinity of z , we can find $r > 0$ such that, if D is the open disk of radius r and center z , $\overline{D} \subset \Omega$, g is analytic in an open set containing \overline{D} and g does not vanish in \overline{D} . Hence, by (A.20), the only zero of $G(\omega)$ in D is z , with multiplicity \mathfrak{m} . Let us define the function $G_\zeta : \omega \mapsto G(\omega) - \zeta^{\mathfrak{m}}$ which is analytic in Ω and ∂D the circle of center z and radius r . Then, one has:

$$\forall \omega \in \partial D, \quad |G_\zeta(\omega) - G(\omega)| = |\zeta|^{\mathfrak{m}} \quad \text{and by (A.20),} \quad |G(\omega)| \geq r^{\mathfrak{m}} \min_{\omega \in \partial D} |g(\omega)| > 0.$$

Thus for $R = r \min_{\omega \in \partial D} |g(\omega)|^{1/\mathfrak{m}}$, one gets

$$\forall |\zeta| < R, \quad \forall \omega \in \partial D, \quad |G_\zeta(\omega) - G(\omega)| < |G(\omega)|.$$

Therefore by the Rouché Theorem, G and G_ζ have the same number of zeros (counted with multiplicity) in D . Hence G_ζ has \mathfrak{m} zeros on $D(z, r)$ for $|\zeta| < R$. Let $D(0, R)$ be the open disk of center 0 and radius R . Thus, if we define U as $\tilde{U} \cap D(0, R)$, then there are no other solutions of $G(\omega) = \zeta^m$ in D when $\zeta \in U$ than $\omega_n(\zeta)$ for $n \in \{1, \dots, \mathfrak{m}\}$. \square

Remark A.2. *The result of Lemma A.1 is related to Puiseux series. Puiseux series have been derived in [13, 15, 27] for the perturbation analysis of eigenvalues of non-selfadjoint matrices. However the results of [27] do not apply directly here since in some situations we are perturbing derogatory eigenvalues (i.e. non simple eigenvalues). Moreover, we need to compute explicitly the two first coefficients of the Puiseux series which is not done in [15], Section 1.2 nor in [13], see Corollary 4.2.9.*

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