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#### Non-convex interferometric inversion for wave scatters recovery in complex media

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#### Abstract

In this work, we propose a new approach to recover sparse obstacles q from measurements bof the scattered wave field in a slowly varying medium. In the Born approximation the problem is linear (Aq = b) but the unknown wave speed variations induces a phase error on the data b. Classical approaches fail at recovering the scatters in this case. Hence we propose an alternative inversion that consists in considering an interferometric quadratic problem

$$(Aq)_k \overline{(Aq)}_\ell = b_k \overline{b_\ell}, \text{ for some } (k, \ell).$$

The interest of such system comes from the better robustness property of some cross-correlations  $b_k \overline{b_\ell}$  with respect to the medium uncertainties.

**Keywords:** Inverse scattering, random media, interferometric inversion, non-convex optimization.

#### 1 The inverse scattering problem

In a domain of interest  $\Omega \subset \mathbb{R}^d$ ,  $d \in \{2,3\}$ , we aim at recovering scatters  $q \in L^{\infty}(\Omega)$  of small support, from measurements of the harmonic scattered field u, solution of the Helmholtz equation in the Born approximation

$$\Delta u + \frac{\omega^2}{c^2} u = -\omega^2 q u_I \quad \text{ in } \mathbb{R}^d,$$

and satisfying the Sommerfeld radiation condition. The incident field  $u_I$  is generated by a controlled source at  $x_0$  and satisfies

$$\Delta u_I + \frac{\omega^2}{c^2} u_I = -\delta_{x_0} \quad \text{in } \mathbb{R}^d$$

Measurements are defined as a discrete evaluation of the scattered field:

$$b_k := \frac{1}{\omega_k^2} u(r_k, \omega_k), \quad k = 1, \dots, n,$$

at some receivers locations  $r_k \in \mathbb{R}^d$  and for chosen frequencies  $\omega_k \geq 0$ . The Green function  $G(x, y, \omega)$  (unknown) can be used to give an integral expression of the propagation operator. Th incident field  $u_I$  takes the form  $u_I(x, \omega) = G(x, x_0, \omega)$  and the measurements read

$$b_k = \int_{\Omega} G(r_k, y, \omega_k) G(y, x_0, \omega_k) q(y) \mathrm{d}y.$$

#### 2 Wave propagation in slowly varying random media

In [2], the authors described the behavior of a wave in a slowly varying random media, where the speed c is given by

$$\frac{1}{c^2(x)} = \frac{1}{c_0^2} \left( 1 + \sigma \mu \left( \frac{x}{\ell} \right) \right)$$

where  $\mu$  is a spatial random variable of autocorrelation  $\mathbb{E}(\mu(x)\mu(x')) = \exp(-|x-x'|^2/2)$  for all  $x, x' \in \mathbb{R}^d$ . In such media, the travel-time function is given by

$$\tau(x,y) := \frac{|x-y|}{c_0} + \nu(x,y), \quad x,y \in \mathbb{R}^d,$$

where  $\nu$  is a spatially smooth random fluctuation built by integrating  $\mu$  between x and y.

In a regime where  $\ell$  is large compared to the wavelength, the authors show that the Green function of the medium is approximated by

$$G(x, y, \omega) := G_0(x - y, \omega)e^{i\omega\nu(x,y)}, \quad x, y \in \mathbb{R}^d$$

where  $G_0$  is the well-known Helmholtz Green function in homogeneous media.

Assuming that  $\operatorname{supp}(q) \subset B(y_0, \varepsilon)$  with  $\varepsilon$ small compared to the coherence length of the medium  $\ell$ , one can approximate the travel time fluctuation from any  $x \in \mathbb{R}^d$  to  $y \in \operatorname{supp}(q)$  by

$$\nu(x,y) \approx \nu(x,y_0).$$

Hence, the measurements reads as

$$b_k = b_k^{\text{ex}} e^{i\omega_k \varphi_k}, \quad \varphi_k := \nu(r_k, y_0) + \nu(y_0, x_0)$$