



HAL
open science

Renormalised energy between boundary vortices in thin-film micromagnetics with Dzyaloshinskii-Moriya interaction

Radu Ignat, François L'official

► **To cite this version:**

Radu Ignat, François L'official. Renormalised energy between boundary vortices in thin-film micromagnetics with Dzyaloshinskii-Moriya interaction. *Nonlinear Analysis: Theory, Methods and Applications*, 2025, 250, pp.113622. 10.1016/j.na.2024.113622 . hal-04722730

HAL Id: hal-04722730

<https://hal.science/hal-04722730v1>

Submitted on 7 Oct 2024

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Renormalised energy between boundary vortices in thin-film micromagnetics with Dzyaloshinskii-Moriya interaction

Radu Ignat* François L'Official†

July 17, 2024

Abstract

We consider a three-dimensional micromagnetic model with Dzyaloshinskii-Moriya interaction in a thin-film regime for boundary vortices. In this regime, we prove a dimension reduction result: the nonlocal three-dimensional model reduces to a local two-dimensional Ginzburg-Landau type model in terms of the averaged magnetization in the thickness of the film. This reduced model captures the interaction between boundary vortices (so-called renormalised energy), that we determine by a Γ -convergence result at the second order and then we analyse its minimisers. They nucleate two boundary vortices whose position depends on the Dzyaloshinskii-Moriya interaction.

Contents

1	Introduction	2
1.1	Nondimensionalisation	2
1.2	Thin-film regime for boundary vortices	3
1.3	Dimension reduction	4
1.4	Gamma-convergence of the two-dimensional reduced functional $E_{\varepsilon,\eta}^\delta$	6
1.5	Minimisation of the renormalised energy	9
1.6	Gamma-convergence of the three-dimensional energy E_h	11
1.7	Related models	13
1.8	Outline of the paper	14
2	Two-dimensional reduced model for maps $v: \Omega \subset \mathbb{R}^2 \rightarrow \mathbb{R}^2$	14
2.1	Approximation by \mathbb{S}^1 -valued maps	14
2.2	Lifting	17
2.3	Gamma-convergence in terms of liftings	17
2.4	Gamma-convergence for vector-valued maps	21
2.5	Minimisers of the renormalised energy	23
3	Three-dimensional model for maps $m_h: \Omega_h \subset \mathbb{R}^3 \rightarrow \mathbb{S}^2$	31
3.1	Reduction from the 3D model to a reduced 2D model	31
3.2	Gamma-convergence of the three-dimensional energy	34

*Institut de Mathématiques de Toulouse & Institut Universitaire de France, UMR 5219, Université de Toulouse, CNRS, UPS IMT, F-31062 Toulouse Cedex 9, France. Email: Radu.Ignat@math.univ-toulouse.fr

†Institut de Mathématiques de Toulouse, UMR 5219, Université de Toulouse, CNRS, UPS, IMT, F-31062 Toulouse Cedex 9, France. Email: Francois.Lofficial@math.univ-toulouse.fr; francois.lof@gmail.com

1 Introduction

We consider a ferromagnetic sample of cylindrical shape of thickness t :

$$\Omega_t^\ell = \Omega^\ell \times (0, t) \subset \mathbb{R}^3,$$

where the horizontal section $\Omega^\ell \subset \mathbb{R}^2$ is a bounded, simply connected $C^{1,1}$ domain of typical length ℓ (for example, Ω^ℓ can be assumed to be a disk of diameter ℓ). The magnetization m is a three-dimensional unit-length vector field

$$m: \Omega_t^\ell \rightarrow \mathbb{S}^2,$$

where \mathbb{S}^2 is the unit sphere in \mathbb{R}^3 , the constraint $|m| = 1$ yielding the non-convexity of the problem. We consider the micromagnetic energy with Dzyaloshinskii-Moriya interaction (in the absence of anisotropy and applied magnetic field):

$$E(m) = A^2 \int_{\Omega_t^\ell} |\nabla m|^2 dx + \int_{\Omega_t^\ell} D : \nabla m \wedge m dx + \int_{\mathbb{R}^3} |\nabla u|^2 dx. \quad (1)$$

The first term in $E(m)$ is the exchange energy, generated by short-range interactions between magnetic spins, where the exchange length $A > 0$ is an intrinsic parameter of the ferromagnetic material, typically of the order of nanometers. The second term is the Dzyaloshinskii-Moriya interaction (DMI), see [16], taking into account the antisymmetric properties of the material. The DMI density is given here by

$$D : \nabla m \wedge m = \sum_{j=1}^3 D_j \cdot \partial_j m \wedge m, \quad (2)$$

where $D = (D_1, D_2, D_3) \in \mathbb{R}^{3 \times 3}$ is the DMI tensor, \cdot denotes the inner product in \mathbb{R}^3 , and \wedge denotes the cross product in \mathbb{R}^3 . The third term in $E(m)$ is called magnetostatic or stray-field energy, it is a nonlocal term generated by long-range spins interactions carried by the stray-field potential $u \in H^1(\mathbb{R}^3, \mathbb{R})$ satisfying

$$\Delta u = \nabla \cdot (m \mathbb{1}_{\Omega_t^\ell}) \quad \text{in the distributional sense in } \mathbb{R}^3,$$

where m is extended by 0 outside the sample Ω_t^ℓ . For more details, we refer to [2], [16] or [21].

1.1 Nondimensionalisation

The multiscale aspect of the micromagnetic model is carried by the DMI tensor D , the exchange length A , the horizontal length ℓ and the thickness t of the sample. We introduce the dimensionless parameters

$$h = \frac{t}{\ell} \quad \text{and} \quad \eta = \frac{A}{\ell};$$

the thin-film regimes correspond to the limit $h \rightarrow 0$ (sometimes denoted by $h \ll 1$).

In order to study the micromagnetic energy in a thin-film regime appropriate to boundary vortices, we nondimensionalize in length and set $\widehat{D} = \frac{1}{\ell} D$,

$$\Omega_h = \frac{\Omega_t^\ell}{\ell} = \Omega \times (0, h) \subset \mathbb{R}^3,$$

where $\Omega = \frac{\Omega^\ell}{\ell} \subset \mathbb{R}^2$ is a bounded, simply connected $C^{1,1}$ domain of typical length 1. To each $x = (x_1, x_2, x_3) \in \Omega^\ell$, we associate $\hat{x} = \frac{x}{\ell} \in \Omega_h$, and then $m_h: \Omega_h \rightarrow \mathbb{S}^2$ and $u_h: \mathbb{R}^3 \rightarrow \mathbb{R}$ given by

$$m_h(\hat{x}) = m(x), \quad u_h(\hat{x}) = \frac{1}{\ell}u(x),$$

that satisfy

$$\Delta u_h = \nabla \cdot (m_h \mathbb{1}_{\Omega_h}) \quad \text{in the distributional sense in } \mathbb{R}^3. \quad (3)$$

The micromagnetic energy (1) can then be written in terms of m_h :

$$E(m) = \widehat{E}(m_h) = \ell^3 \left[\eta^2 \int_{\Omega_h} |\nabla m_h|^2 d\hat{x} + \int_{\Omega_h} \widehat{D} : \nabla m_h \wedge m_h d\hat{x} + \int_{\mathbb{R}^3} |\nabla u_h|^2 d\hat{x} \right].$$

For simplicity of the notations, we write x instead of \hat{x} in the following.

1.2 Thin-film regime for boundary vortices

We consider the thin-film regime studied by Ignat-Kurzke [26] with appropriate scaling of the DMI tensor $\widehat{D} = \frac{D}{\ell}$. More precisely, we consider here the regime

$$\begin{aligned} h \ll 1, \quad \eta \ll 1, \quad \frac{1}{|\log h|} \ll \varepsilon \ll 1, \quad \frac{\widehat{D}_{13}}{\eta^2} \rightarrow 2\delta_1, \quad \frac{\widehat{D}_{23}}{\eta^2} \rightarrow 2\delta_2, \\ \frac{1}{\eta^2} \left(\sum_{j,k=1}^2 |\widehat{D}_{jk}| + \sum_{k=1}^3 |\widehat{D}_{3k}| \right) \ll \sqrt{|\log \varepsilon|}, \end{aligned} \quad (4)$$

where $\delta_1, \delta_2 \in \mathbb{R}$, $\widehat{D} = (\widehat{D}_{jk})_{1 \leq j, k \leq 3} \in \mathbb{R}^{3 \times 3}$ and

$$\varepsilon = \frac{\eta^2}{h |\log h|}$$

that corresponds to the core size of the boundary vortices. The parameters $\eta = \eta(h)$, $\varepsilon = \varepsilon(h)$ and $\widehat{D} = \widehat{D}(h)$ are assumed to be functions in h ; our first aim is to prove a Γ -convergence result at the first order in the limit $h \rightarrow 0$ that quantifies the number of boundary vortices in the domain.

The regime (4) implies¹:

$$\frac{\log |\log h|}{|\log h|} \ll \varepsilon |\log \varepsilon| \ll 1.$$

In fact, our second aim is to prove a Γ -convergence result at the second order as $h \rightarrow 0$, capturing the interaction energy between boundary vortices; for that, we restrict to a narrower regime than (4):

$$\begin{aligned} h \ll 1, \quad \eta \ll 1, \quad \frac{\log |\log h|}{|\log h|} \ll \varepsilon \ll 1, \quad \left| \frac{\widehat{D}_{13}}{\eta^2} - 2\delta_1 \right| + \left| \frac{\widehat{D}_{23}}{\eta^2} - 2\delta_2 \right| \ll \frac{1}{\sqrt{|\log \varepsilon|}}, \\ \frac{1}{\eta^2} \left(\sum_{j,k=1}^2 |\widehat{D}_{jk}| + \sum_{k=1}^3 |\widehat{D}_{3k}| \right) \ll \frac{1}{\sqrt{|\log \varepsilon|}}. \end{aligned} \quad (5)$$

¹We use that inequalities $a \ll b \ll 1$ imply $a |\log a| \ll b |\log b| \ll 1$, and then choose $a = \frac{1}{|\log h|}$ and $b = \varepsilon$.

We assume that $\Omega \subset \mathbb{R}^2$ is a bounded, simply connected $C^{1,1}$ domain. We consider the three-dimensional rescaled energy

$$E_h(m_h) = \frac{\widehat{E}(m_h)}{\ell^3 h \eta^2 |\log \varepsilon|},$$

for maps $m_h: \Omega_h = \Omega \times (0, h) \rightarrow \mathbb{S}^2$ and $u_h: \mathbb{R}^3 \rightarrow \mathbb{R}$ satisfying (3). More precisely, the energy $E_h(m_h)$ is given by

$$E_h(m_h) = \frac{1}{|\log \varepsilon|} \left(\frac{1}{h} \int_{\Omega_h} |\nabla m_h|^2 dx + \frac{1}{h \eta^2} \int_{\Omega_h} \widehat{D} : \nabla m_h \wedge m_h dx + \frac{1}{h \eta^2} \int_{\mathbb{R}^3} |\nabla u_h|^2 dx \right). \quad (6)$$

1.3 Dimension reduction

The thin-film regime is characterized by the assumption $h = t/\ell \rightarrow 0$, i.e., the variations of the magnetization m_h in the thickness direction x_3 are strongly penalised. Therefore m_h is expected to behave as its x_3 -average $\overline{m}_h: \Omega \rightarrow \mathbb{R}^3$, that is,

$$\overline{m}_h(x') = \frac{1}{h} \int_0^h m_h(x', x_3) dx_3 \quad \text{for every } x' \in \Omega, \quad (7)$$

where $|\overline{m}_h| \leq 1$ in Ω . With this in mind, we assume for a bit that m_h does not depend on x_3 , and that

$$m_h \text{ varies on length scales } \gg h.$$

We are interested in the scaling of the stray-field energy in this regime; within our assumption (i.e., $m_h \equiv \overline{m}_h$), the Maxwell equation (3) implies

$$\Delta u_h = \nabla \cdot (m_h \mathbf{1}_{\Omega_h}) = (\nabla' \cdot m'_h) \mathbf{1}_{\Omega_h} - (m_h \cdot \nu) \mathbf{1}_{\partial \Omega_h}$$

in the distributional sense in \mathbb{R}^3 , where ν is the outer unit normal vector on $\partial \Omega_h$, and **prime** corresponds to **2D quantities** such as the in-plane magnetization $m'_h = (m_{h,1}, m_{h,2})$ and the in-plane divergence $\nabla' \cdot m'_h = \partial_1 m_{h,1} + \partial_2 m_{h,2}$. In other words, $u_h \in H^1(\mathbb{R}^3)$ is the solution of the transmission problem

$$\begin{cases} \Delta u_h = \nabla' \cdot m'_h & \text{in } \Omega_h, \\ \Delta u_h = 0 & \text{in } \mathbb{R}^3 \setminus \Omega_h, \\ \left[\frac{\partial u_h}{\partial \nu} \right] = m_h \cdot \nu & \text{on } \partial \Omega_h, \end{cases}$$

where $[a] = a^+ - a^-$ stands for the jump of a with respect to the outer unit normal vector ν on $\partial \Omega_h$. From [18] (see also [1], [14], [22, Section 2.1.2]), we can express the stray-field energy by considering the Fourier transform in the horizontal variables:

$$\begin{aligned} \int_{\mathbb{R}^3} |\nabla u_h|^2 dx &= h \int_{\mathbb{R}^2} \frac{|\xi' \cdot \mathcal{F}(m'_h \mathbf{1}_{\Omega})(\xi')|^2}{|\xi'|^2} (1 - g_h(|\xi'|)) d\xi' \\ &\quad + h \int_{\mathbb{R}^2} |\mathcal{F}(m_{h,3} \mathbf{1}_{\Omega})(\xi')|^2 g_h(|\xi'|) d\xi', \end{aligned}$$

where \mathcal{F} stands for the Fourier transform in \mathbb{R}^2 , i.e., for every $f: \mathbb{R}^2 \rightarrow \mathbb{C}$ and for every $\xi' \in \mathbb{R}^2$,

$$\mathcal{F}(f)(\xi') = \int_{\mathbb{R}^2} f(x') e^{-2i\pi x' \cdot \xi'} dx', \quad \text{and} \quad g_h(|\xi'|) = \frac{1 - e^{-2\pi h |\xi'|}}{2\pi h |\xi'|}.$$

In the asymptotics $h \rightarrow 0$, we have $g_h(|\xi'|) \rightarrow 1$ and $1 - g_h(|\xi'|) \approx \pi h |\xi'|$, hence (see [10], [14], [32])

$$\begin{aligned} \int_{\mathbb{R}^3} |\nabla u_h|^2 dx &\approx h^2 \|(\nabla' \cdot m'_h) \mathbf{1}_\Omega\|_{\dot{H}^{-1/2}(\mathbb{R}^2)}^2 \\ &+ \frac{1}{2\pi} h^2 |\log h| \int_{\partial\Omega} (m'_h \cdot \nu')^2 d\mathcal{H}^1 + h \int_{\Omega} m_{h,3}^2 dx', \end{aligned}$$

where ν' is the outer unit normal vector on $\partial\Omega$. Hence, the stray-field energy is asymptotically decomposed in three terms in the thin-film regime. The first term is nonlocal and penalizes the volume charges, as an homogeneous $\dot{H}^{-1/2}$ seminorm, and favors Néel walls. The second term takes into account the lateral charges on the cylindrical sample and favors boundary vortices. The third term penalizes the surface charges on the top and bottom of the cylinder, and leads to interior vortices. For more details on the different types of singularities that may occur in thin-film regimes, we refer to [14] or [22].

Our dimension reduction result in Theorem 1.1 shows that the stray-field energy of a general magnetization m_h depending on x_3 reduces in our regime to the last two local terms in the x_3 -average \overline{m}_h (while the nonlocal term $h^2 \|(\nabla' \cdot \overline{m}'_h) \mathbf{1}_\Omega\|_{\dot{H}^{-1/2}(\mathbb{R}^2)}^2$ becomes negligible). More precisely, for a given vector $\delta = (\delta_1, \delta_2) \in \mathbb{R}^2$, we consider the two-dimensional reduced energy functional for a two-dimensional map $v: \Omega \subset \mathbb{R}^2 \rightarrow \mathbb{R}^2$ (standing for the in-plane average \overline{m}'_h):

$$\begin{aligned} E_{\varepsilon, \eta}^\delta(v) &= \int_{\Omega} |\nabla' v|^2 dx + 2 \int_{\Omega} \delta \cdot \nabla' v \wedge v dx \\ &+ \frac{1}{\eta^2} \int_{\Omega} (1 - |v|^2)^2 dx + \frac{1}{2\pi\varepsilon} \int_{\partial\Omega} (v \cdot \nu')^2 d\mathcal{H}^1, \end{aligned} \quad (8)$$

where $\varepsilon, \eta > 0$, $\nabla' = (\partial_1, \partial_2)$ and ν' is the outer unit normal vector on $\partial\Omega$. Note that by identifying the complex plane $\mathbb{C} \simeq \mathbb{R}^2$, we have $|\nabla' v|^2 + 2\delta \cdot \nabla' v \wedge v = |(\nabla' - i\delta)v|^2 - |\delta|^2 |v|^2$. Therefore, the functional $E_{\varepsilon, \eta}^\delta$ is similar to the Ginzburg-Landau model with magnetic potential, combining a boundary penalisation (favoring boundary vortices) with an interior penalisation (favoring interior vortices). The case $\delta = 0$ was analysed in Ignat-Kurzke [25] (see also Moser [39]).

Theorem 1.1. *Let $\Omega_h = \Omega \times (0, h)$ with $\Omega \subset \mathbb{R}^2$ a bounded, simply connected $C^{1,1}$ domain. In the regime (4), consider a family of magnetizations $\{m_h: \Omega_h \rightarrow \mathbb{S}^2\}_{h \downarrow 0}$ that satisfies*

$$\limsup_{h \rightarrow 0} E_h(m_h) < +\infty$$

and let $\overline{m}_h = (\overline{m}'_h, \overline{m}_{h,3}): \Omega \rightarrow \mathbb{R}^3$ be the average of m_h in (7). Then²

$$E_h(m_h) \geq \frac{1}{|\log \varepsilon|} E_{\varepsilon, \eta}^\delta(\overline{m}'_h) - o_h(1) \quad \text{as } h \rightarrow 0.$$

Moreover, in the more restrictive regime (5), we have

$$E_h(m_h) \geq \frac{1}{|\log \varepsilon|} E_{\varepsilon, \eta}^\delta(\overline{m}'_h) - o_h\left(\frac{1}{|\log \varepsilon|}\right) \quad \text{as } h \rightarrow 0.$$

Furthermore, if m_h is independent of x_3 (i.e., $m_h = m_h(x')$, $x' = (x_1, x_2)$) and $m_{h,3} = 0$ (i.e., $m_h = (m'_h, 0)$, $|m'_h| = 1$), then in the regime (4), we have

$$E_h(m_h) = \frac{1}{|\log \varepsilon|} E_{\varepsilon, \eta}^\delta(\overline{m}'_h) - o_h(1) \quad \text{as } h \rightarrow 0,$$

²We write $a = o_h(b)$ if $a/b \rightarrow 0$ as $h \rightarrow 0$.

while in the regime (5), we have

$$E_h(m_h) = \frac{1}{|\log \varepsilon|} E_{\varepsilon, \eta}^\delta(\overline{m}'_h) - o_h \left(\frac{1}{|\log \varepsilon|} \right) \quad \text{as } h \rightarrow 0.$$

Therefore, in order to determine the asymptotic behaviour of E_h as $h \rightarrow 0$, we need to analyse the behaviour of the two-dimensional reduced functional $E_{\varepsilon, \eta}^\delta$ in the regime (4), respectively in the regime (5).

1.4 Gamma-convergence of the two-dimensional reduced functional $E_{\varepsilon, \eta}^\delta$

In this section, we prove Γ -convergence results (at the first and second order) for the reduced functional $E_{\varepsilon, \eta}^\delta$ defined in (8) in the regimes (4) and (5), respectively. These results are reminiscent from Ignat-Kurzke [25], who studied the energy functional $E_{\varepsilon, \eta}^0$ (i.e. the case $\delta = 0$). At the first order, the Γ -limit energy is expected to quantify the number of boundary vortices detected by the global Jacobian introduced in [25], [26]. More precisely, for a two-dimensional map $v \in H^1(\Omega, \mathbb{R}^2)$ defined in a bounded $C^{1,1}$ domain $\Omega \subset \mathbb{R}^2$, the **global Jacobian** of v is given by the linear operator $\mathcal{J}(v): W^{1, \infty}(\Omega) \rightarrow \mathbb{R}$ defined as

$$\langle \mathcal{J}(v), \zeta \rangle = - \int_{\Omega} v \wedge \nabla' v \cdot \nabla'^{\perp} \zeta \, dx',$$

for every Lipschitz function $\zeta: \Omega \rightarrow \mathbb{R}$, where $v \wedge \nabla' v = (v \wedge \partial_1 v, v \wedge \partial_2 v)$, $\nabla'^{\perp} = (-\partial_2, \partial_1)$ and $\langle \cdot, \cdot \rangle$ stands for the algebraic dual pairing between $(W^{1, \infty}(\Omega))^*$ and $W^{1, \infty}(\Omega)$. In particular, the global Jacobian has zero average, i.e. $\langle \mathcal{J}(v), 1 \rangle = 0$. Moreover,

$$\mathcal{J}(v) = 2 \, \text{jac}(v) + \mathcal{J}_{\text{bd}}(v)$$

where $\text{jac}(v) = \det(\nabla' v) = \partial_1 v \wedge \partial_2 v \in L^1(\Omega)$ is the interior Jacobian of v , and $\mathcal{J}_{\text{bd}}(v)$ is the boundary Jacobian of v . For example, if $v \in C^1(\overline{\Omega}, \mathbb{R}^2)$, then $\langle \mathcal{J}_{\text{bd}}(v), \zeta \rangle = - \int_{\partial\Omega} v \wedge \partial_{\tau} v \zeta \, d\mathcal{H}^1$ where $\tau = (\nu')^{\perp}$ is the tangent unit vector at $\partial\Omega$ (for more details, see [25, Section 2]).

For a fixed vector $\delta \in \mathbb{R}^2$, we start by analysing the energy functional $E_{\varepsilon, \eta}^\delta$ in the asymptotic regime:

$$\eta \ll 1, \quad \varepsilon \ll 1, \quad |\log \varepsilon| \ll |\log \eta|, \quad (9)$$

which is implied³ by the regime (4). For a family $E_{\varepsilon, \eta}^\delta(v_\varepsilon) \leq C |\log \varepsilon|$, we prove that the global Jacobian $\mathcal{J}(v_\varepsilon)$ concentrates on Dirac masses at the boundary vortices (up to a diffuse measure carried by the curvature of $\partial\Omega$). Moreover, the lower bound of $E_{\varepsilon, \eta}^\delta(v_\varepsilon)$ quantifies the number of these boundary vortices. We also show $L^p(\partial\Omega)$ compactness of the traces $v_\varepsilon|_{\partial\Omega}$ converging to a map in $BV(\partial\Omega, \mathbb{S}^1)$.

Theorem 1.2 (Compactness at the boundary and lower bound at the first order). *Let $\delta \in \mathbb{R}^2$, $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain and κ be the curvature of $\partial\Omega$. Assume $\varepsilon \rightarrow 0$ and $\eta = \eta(\varepsilon) \rightarrow 0$ in the regime (9). Let $(v_\varepsilon)_\varepsilon$ be a family in $H^1(\Omega, \mathbb{R}^2)$ such that*

$$\limsup_{\varepsilon \rightarrow 0} \frac{1}{|\log \varepsilon|} E_{\varepsilon, \eta}^\delta(v_\varepsilon) < +\infty. \quad (10)$$

Then the following statements hold.

³ Indeed, the regime (4) implies $h \ll \eta^2 \ll h |\log h| \ll 1$ by using the definition of ε , hence $|\log h| \sim |\log \eta|$. It follows from (4) that $\frac{1}{\varepsilon} \ll |\log \eta|$, hence (9) is satisfied because $|\log \varepsilon| \ll \frac{1}{\varepsilon}$.

- (i) **Compactness of global Jacobians and $L^p(\partial\Omega)$ -compactness of $v_\varepsilon|_{\partial\Omega}$.**
 For a subsequence $\varepsilon \rightarrow 0$, $(\mathcal{J}(v_\varepsilon))_\varepsilon$ converges to a measure J on the closure $\overline{\Omega}$, in the sense that

$$\lim_{\varepsilon \rightarrow 0} \left(\sup_{|\nabla' \zeta| \leq 1 \text{ in } \Omega} |\langle \mathcal{J}(v_\varepsilon) - J, \zeta \rangle| \right) = 0. \quad (11)$$

Moreover, J is supported on $\partial\Omega$ and has the form⁴

$$J = -\kappa \mathcal{H}^1 \llcorner \partial\Omega + \pi \sum_{j=1}^N d_j \delta_{a_j} \quad (12)$$

for $N \geq 1$ distinct boundary vortices $a_j \in \partial\Omega$ carrying the multiplicities $d_j \in \mathbb{Z} \setminus \{0\}$, for $j \in \{1, \dots, N\}$, such that $\sum_{j=1}^N d_j = 2$.

Furthermore, for a subsequence, the family of traces $(v_\varepsilon|_{\partial\Omega})_\varepsilon$ converges to $e^{i\varphi_0} \in BV(\partial\Omega, \mathbb{S}^1)$ in $L^p(\partial\Omega, \mathbb{R}^2)$, for every $p \in [1, +\infty)$, where φ_0 is a BV lifting of the tangent field $\pm\tau$ on $\partial\Omega$ determined (up to a constant in $\pi\mathbb{Z}$) by

$$\partial_\tau \varphi_0 = \kappa \mathcal{H}^1 \llcorner \partial\Omega - \pi \sum_{j=1}^N d_j \delta_{a_j} \quad \text{as measure on } \partial\Omega.$$

- (ii) **Energy lower bound at the first order.**

If $(\mathcal{J}(v_\varepsilon))_\varepsilon$ satisfies the convergence assumption in (i) as $\varepsilon \rightarrow 0$, then the energy lower bound at the first order is the total mass of the measure $J + \kappa \mathcal{H}^1 \llcorner \partial\Omega$ on $\partial\Omega$:

$$\liminf_{\varepsilon \rightarrow 0} \frac{1}{|\log \varepsilon|} E_{\varepsilon, \eta}^\delta(v_\varepsilon) \geq \pi \sum_{j=1}^N |d_j| = |J + \kappa \mathcal{H}^1 \llcorner \partial\Omega|(\partial\Omega).$$

In order to get a lower bound at the second order for $E_{\varepsilon, \eta}^\delta$, we introduce the following renormalised energy, in the spirit of Brezis-Bethuel-Hélein [6], that captures the interaction energy between boundary vortices. In comparison with Ignat-Kurzke [25], [26], our renormalised energy includes the contribution of the DMI energy.

Definition 1.3. Let $\delta \in \mathbb{R}^2$, $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain and κ be the curvature of $\partial\Omega$. Consider $\varphi_0: \partial\Omega \rightarrow \mathbb{R}$ to be a BV function such that $e^{i\varphi_0} \cdot \nu' = 0$ in $\partial\Omega \setminus \{a_1, \dots, a_N\}$, and

$$\partial_\tau \varphi_0 = \kappa \mathcal{H}^1 \llcorner \partial\Omega - \pi \sum_{j=1}^N d_j \delta_{a_j} \quad \text{as measure on } \partial\Omega,$$

for $N \geq 1$ distinct points $a_j \in \partial\Omega$ carrying the multiplicities $d_j \in \{\pm 1\}$ for $j \in \{1, \dots, N\}$ with $\sum_{j=1}^N d_j = 2$. If φ_* is the harmonic extension of φ_0 to Ω , then the renormalised energy of $\{(a_j, d_j)\}_{j \in \{1, \dots, N\}}$ is defined as

$$W_\Omega^\delta(\{(a_j, d_j)\}) = \liminf_{r \rightarrow 0} \left(\int_{\Omega \setminus \bigcup_{j=1}^N B_r(a_j)} (|\nabla' \varphi_*|^2 - 2\delta \cdot \nabla' \varphi_*) \, dx - N\pi \log \frac{1}{r} \right), \quad (13)$$

where $B_r(a_j)$ is the disk of center a_j and radius $r > 0$.

⁴Note that the condition $\langle \mathcal{J}, 1 \rangle = 0$ and the Gauss-Bonnet formula $\int_{\partial\Omega} \kappa \, d\mathcal{H}^1 = 2\pi$ yield the constraint: $\sum_{j=1}^N d_j = 2$.

We will see that the \liminf in (13) is in fact a limit (see Proposition 2.10). In the following two theorems, we prove the asymptotic expansion at the second order for the energy $E_{\varepsilon,\eta}^\delta$ by the Γ -convergence method. In particular, under a prescribed level of energy (see (14)), we prove that the boundary vortices have multiplicities ± 1 and that the second order energy is given by the renormalised energy (13) up to a constant depending on the number of boundary vortices.

Theorem 1.4 (Compactness in the interior and lower bound at the second order). *Let $\delta \in \mathbb{R}^2$, $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain and κ be the curvature of $\partial\Omega$. Assume $\varepsilon \rightarrow 0$ and $\eta = \eta(\varepsilon) \rightarrow 0$ in the regime (9). Let $(v_\varepsilon)_\varepsilon$ be a family in $H^1(\Omega, \mathbb{R}^2)$ satisfying (10) and the convergence at (i) in Theorem 1.2 to the measure J given at (12) as $\varepsilon \rightarrow 0$. In addition, we assume the following more precise bound than (10):*

$$\limsup_{\varepsilon \rightarrow 0} \left(E_{\varepsilon,\eta}^\delta(v_\varepsilon) - \pi |\log \varepsilon| \sum_{j=1}^N |d_j| \right) < +\infty. \quad (14)$$

Then the following statements hold.

(i) **Single multiplicity and second order lower bound.**

The multiplicities satisfy $d_j \in \{\pm 1\}$ for every $j \in \{1, \dots, N\}$, so we have⁵ $\sum_{j=1}^N |d_j| = N$; moreover, the following second order energy lower bound holds:

$$\liminf_{\varepsilon \rightarrow 0} (E_{\varepsilon,\eta}^\delta(v_\varepsilon) - N\pi |\log \varepsilon|) \geq W_\Omega^\delta(\{(a_j, d_j)\}) + N\gamma_0$$

where $\gamma_0 = \pi \log \frac{e}{4\pi}$ and $W_\Omega^\delta(\{(a_j, d_j)\})$ is the renormalised energy defined in (13).

(ii) **$W^{1,q}(\Omega)$ -weak compactness of maps v_ε .**

For any $q \in [1, 2)$, $(v_\varepsilon)_\varepsilon$ is uniformly bounded in $W^{1,q}(\Omega, \mathbb{R}^2)$. Moreover, for a subsequence, $(v_\varepsilon)_\varepsilon$ converges weakly in $W^{1,q}(\Omega, \mathbb{R}^2)$, for every $q \in [1, 2)$, and strongly in $L^p(\Omega, \mathbb{R}^2)$, for every $p \in [1, +\infty)$, to $e^{i\widehat{\varphi}_0}$, where $\widehat{\varphi}_0 \in W^{1,q}(\Omega)$ is an extension (not necessarily harmonic) to Ω of the lifting $\varphi_0 \in BV(\partial\Omega, \pi\mathbb{Z})$ determined in Theorem 1.2(i).

Now we state the upper bound part in the Γ -expansion where the recovery sequence v_ε can be chosen with values into \mathbb{S}^1 :

Theorem 1.5 (Upper bound). *Let $\delta \in \mathbb{R}^2$, $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain and κ be the curvature of $\partial\Omega$. Let $\{a_j \in \partial\Omega\}_{1 \leq j \leq N}$ be $N \geq 1$ distinct points and $d_j \in \mathbb{Z} \setminus \{0\}$ be the corresponding multiplicities, for $j \in \{1, \dots, N\}$, that satisfy $\sum_{j=1}^N d_j = 2$. Assume $\varepsilon \rightarrow 0$ and $\eta = \eta(\varepsilon) \rightarrow 0$ in the regime (9). Then we can construct a family $(v_\varepsilon)_\varepsilon$ in $H^1(\Omega, \mathbb{S}^1)$ such that $(\mathcal{J}(v_\varepsilon))_\varepsilon$ converges as in (11) to the measure*

$$J = -\kappa \mathcal{H}^1 \llcorner \partial\Omega + \pi \sum_{j=1}^N d_j \delta_{a_j}.$$

Furthermore, $(v_\varepsilon)_\varepsilon$ converges strongly to $e^{i\varphi_*}$ in $L^p(\Omega, \mathbb{R}^2)$ and in $L^p(\partial\Omega, \mathbb{R}^2)$, for every $p \in [1, +\infty)$, where φ_* is the harmonic extension to Ω of a boundary lifting φ_0 of the tangent field $\pm\tau$ on $\partial\Omega$ that satisfies $\partial_\tau \varphi_0 = -J$ as measure on $\partial\Omega$, and the energy of v_ε satisfies

$$\lim_{\varepsilon \rightarrow 0} \frac{1}{|\log \varepsilon|} E_{\varepsilon,\eta}^\delta(v_\varepsilon) = \pi \sum_{j=1}^N |d_j|.$$

⁵Recall that $\sum_{j=1}^N d_j = 2$.

Furthermore, if $d_j \in \{\pm 1\}$ for every $j \in \{1, \dots, N\}$, then v_ε can be chosen such that

$$\lim_{\varepsilon \rightarrow 0} (E_{\varepsilon, \eta}^\delta(v_\varepsilon) - N\pi |\log \varepsilon|) = W_\Omega^\delta(\{(a_j, d_j)\}) + N\gamma_0$$

where $\gamma_0 = \pi \log \frac{e}{4\pi}$ and $W_\Omega^\delta(\{(a_j, d_j)\})$ is the renormalised energy defined in (13).

1.5 Minimisation of the renormalised energy

In this section, we compute explicitly the renormalised energy W_Ω^δ in terms of the location of boundary vortices (a_j) , their multiplicities (d_j) and the reduced DMI vector $\delta \in \mathbb{R}^2$. Then we will analyse the minimisers of W_Ω^δ showing the role of δ in the optimal location of boundary vortices. These results are reminiscent from [26] in the absence of the DMI vector δ .

Theorem 1.6. *Let $\delta \in \mathbb{R}^2$.*

- (i) *We denote by B_1 the unit disk in \mathbb{R}^2 . Let $\{a_j \in \partial B_1\}_{1 \leq j \leq N}$ be $N \geq 2$ distinct points and $d_j \in \{\pm 1\}$ be the corresponding multiplicities, for $j \in \{1, \dots, N\}$, that satisfy $\sum_{j=1}^N d_j = 2$. Then the renormalised energy of $\{(a_j, d_j)\}$ in B_1 satisfies*

$$W_{B_1}^\delta(\{(a_j, d_j)\}) = -2\pi \sum_{1 \leq j < k \leq N} d_j d_k \log |a_j - a_k| + 2\pi \sum_{j=1}^N d_j \delta \cdot a_j^\perp. \quad (15)$$

- (ii) *Let $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain with κ the curvature of $\partial\Omega$ and ν the outer unit normal vector on $\partial\Omega$. Let $\Phi: \overline{B_1} \rightarrow \overline{\Omega}$ be a C^1 conformal diffeomorphism with inverse $\Psi = \Phi^{-1}$. Let $\{a_j \in \partial\Omega\}_{1 \leq j \leq N}$ be $N \geq 2$ distinct points and $d_j \in \{\pm 1\}$ be the corresponding multiplicities, for $j \in \{1, \dots, N\}$, that satisfy $\sum_{j=1}^N d_j = 2$. Then the renormalised energy of $\{(a_j, d_j)\}$ in Ω satisfies*

$$\begin{aligned} & W_\Omega^\delta(\{(a_j, d_j)\}) \\ &= -2\pi \sum_{1 \leq j < k \leq N} d_j d_k \log |\Psi(a_j) - \Psi(a_k)| + \pi \sum_{j=1}^N (d_j - 1) \log |\partial_z \Psi(a_j)| \\ & \quad + \int_{\partial\Omega} (\kappa + 2\delta^\perp \cdot \nu')(w) \left(\sum_{j=1}^N d_j \log |\Psi(w) - \Psi(a_j)| - \log |\partial_z \Psi(w)| \right) d\mathcal{H}^1(w) \end{aligned} \quad (16)$$

where $\partial_z \Psi$ stands for the complex derivative of $\Psi: \overline{\Omega} \subset \mathbb{C} \rightarrow \overline{B_1} \subset \mathbb{C}$.

By Theorems 1.2, 1.4 and 1.5 together with the constraint $\sum_{j=1}^N d_j = 2$, any minimiser of $E_{\varepsilon, \eta}^\delta$ nucleates two boundary vortices of multiplicity 1 in the limit $\varepsilon \rightarrow 0$ in the regime (9). For such configurations, we prove the existence of optimal minimisers of the renormalised energy:

Corollary 1.7. *Let $\delta \in \mathbb{R}^2$ and $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain. There exists a pair $(a_1^*, a_2^*) \in \partial\Omega \times \partial\Omega$ of distinct points such that*

$$W_\Omega^\delta(\{(a_1^*, 1), (a_2^*, 1)\}) = \min \{W_\Omega^\delta(\{(a_1, 1), (a_2, 1)\}) : a_1 \neq a_2 \in \partial\Omega\}.$$

We will determine now the location of the optimal pair (a_1^*, a_2^*) in the unit disk $\Omega = B_1$. Recall that for the model with $\delta = 0$ studied by Ignat and Kurzke [25], [26], the pair of boundary vortices that minimises the renormalised energy corresponds to two diametrically opposed points on ∂B_1 and is unique (up to a rotation). For our model with $\delta \in \mathbb{R}^2$, the location of the optimal pair (a_1^*, a_2^*)

is influenced by the Dzyaloshinskii-Moriya interaction. Indeed, the renormalised energy for two vortices of multiplicity 1 at $a_1, a_2 \in \partial B_1$ is

$$W_{B_1}^\delta(\{(a_1, 1), (a_2, 1)\}) = -2\pi \log |a_1 - a_2| + 2\pi\delta \cdot (a_1^\perp + a_2^\perp).$$

For $\delta \neq 0$, the minimisation of the renormalised energy is different due to the competition between the two terms in $W_{B_1}^\delta$. In the next theorem, we show that the minimal configuration for the points $a_1, a_2 \in \partial B_1$ is unique (up to switching a_1 and a_2) and the vortices are not diametrically opposed if $\delta \neq 0$, but symmetric with respect to δ^\perp : the distance between a_1 and a_2 becomes smaller as $|\delta|$ gets larger.

Theorem 1.8. *Let B_1 be the unit disk in \mathbb{R}^2 and $\delta = |\delta|e^{i\theta} \in \mathbb{R}^2 \setminus \{(0,0)\}$ for some $\theta \in \mathbb{R}$. Then the pair of distinct points $(a_1^*, a_2^*) \in \partial B_1 \times \partial B_1$ that minimises the renormalised energy $W_{B_1}^\delta(\{(\cdot, 1), (\cdot, 1)\})$ in Corollary 1.7 is unique (up to switching a_1^* and a_2^*) and given by*

$$a_1^* = e^{i(\theta+\theta_\delta)} \quad \text{and} \quad a_2^* = e^{i(\theta+\pi-\theta_\delta)}$$

where $\theta_\delta = \arcsin\left(\sqrt{1 + \frac{1}{16|\delta|^2}} - \frac{1}{4|\delta|}\right)$ if $\delta \neq 0$. In particular, a_1^* and a_2^* are symmetric with respect to δ^\perp .

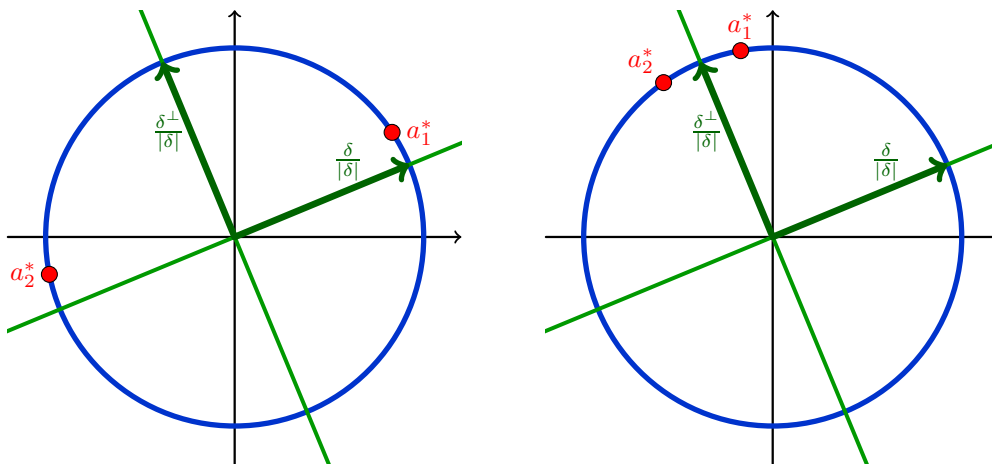


Figure 1: Minimising pair (a_1^*, a_2^*) for $\delta = \frac{1}{10}e^{i\pi/8}$ (left) and $\delta = 10e^{i\pi/8}$ (right). The distance between the boundary vortices a_1^* and a_2^* becomes smaller as $|\delta|$ gets larger.

The following result proves the asymptotic behaviour of minimisers⁶ of $E_{\varepsilon,\eta}^\delta$ as $\varepsilon \rightarrow 0$.

Corollary 1.9. *Let $\delta \in \mathbb{R}^2$, $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain and κ be the curvature of $\partial\Omega$. Assume $\varepsilon \rightarrow 0$ and $\eta = \eta(\varepsilon) \rightarrow 0$ in the regime (9). For every family $(v_\varepsilon)_\varepsilon$ of minimisers of $E_{\varepsilon,\eta}^\delta$ on $H^1(\Omega, \mathbb{R}^2)$, there exists a subsequence $\varepsilon \rightarrow 0$ such that $(\mathcal{J}(v_\varepsilon))_\varepsilon$ converges as in (11) to the measure*

$$J = -\kappa\mathcal{H}^1 \llcorner \partial\Omega + \pi(\delta_{a_1} + \delta_{a_2}),$$

where a_1 and a_2 are two distinct points in $\partial\Omega$ that minimise the renormalised energy (for the multiplicities $d_1 = d_2 = 1$), i.e.

$$W_\Omega^\delta(\{(a_1, 1), (a_2, 1)\}) = \min \{W_\Omega^\delta(\{(\tilde{a}_1, 1), (\tilde{a}_2, 1)\}) : \tilde{a}_1 \neq \tilde{a}_2 \in \partial\Omega\}.$$

⁶The existence of minimisers of $E_{\varepsilon,\eta}^\delta$ is proved in Lemma 2.9.

Moreover, for a subsequence $\varepsilon \rightarrow 0$, $(v_\varepsilon)_\varepsilon$ converges weakly to $e^{i\varphi_*}$ in $W^{1,q}(\Omega, \mathbb{R}^2)$, for every $q \in [1, 2)$, where φ_* is the harmonic extension to Ω of a boundary lifting $\varphi_0 \in BV(\partial\Omega, \pi\mathbb{Z})$ that satisfies $\partial_\tau \varphi_0 = -J$ on $\partial\Omega$ and $e^{i\varphi_0} \cdot \nu' = 0$ in $\partial\Omega \setminus \{a_1, a_2\}$.

Furthermore, there holds the following second order expansion of the minimal energy:

$$\min E_{\varepsilon, \eta}^\delta = 2\pi |\log \varepsilon| + W_\Omega^\delta(\{(a_1, 1), (a_2, 1)\}) + 2\gamma_0 + o_\varepsilon(1) \quad \text{as } \varepsilon \rightarrow 0,$$

where $\gamma_0 = \pi \log \frac{e}{4\pi}$.

In the case of the unit disk $\Omega = B_1$, combining Theorem 1.6 and Corollary 1.9, we deduce in the regime (9) for a minimiser v_ε of $E_{\varepsilon, \eta}^\delta$ that

$$\lim_{\varepsilon \rightarrow 0} (E_{\varepsilon, \eta}^\delta(v_\varepsilon) - 2\pi |\log \varepsilon|) = -2\pi \log |a_1^* - a_2^*| + 2\pi\delta \cdot ((a_1^*)^\perp + (a_2^*)^\perp) + 2\pi \log \frac{e}{4\pi}$$

where (a_1^*, a_2^*) is given in Theorem 1.8.

1.6 Gamma-convergence of the three-dimensional energy E_h

We come back to the nonlocal three-dimensional model for maps $m_h: \Omega_h \subset \mathbb{R}^3 \rightarrow \mathbb{S}^2$ and study the three-dimensional energy E_h given at (6). Thanks to the dimension reduction in Theorem 1.1 and the results obtained for the two-dimensional reduced model studied in Section 1.4, we prove the Γ -expansion at the second order of E_h as $h \rightarrow 0$ summarized in the following three theorems:

Theorem 1.10. *Let $\Omega_h = \Omega \times (0, h)$ with $\Omega \subset \mathbb{R}^2$ a bounded, simply connected $C^{1,1}$ domain, and κ be the curvature of $\partial\Omega$. In the regime (4), consider a family of magnetizations $\{m_h: \Omega_h \rightarrow \mathbb{S}^2\}_{h \downarrow 0}$ that satisfies*

$$\limsup_{h \rightarrow 0} E_h(m_h) < +\infty \quad (17)$$

and let $\overline{m}_h = (\overline{m}'_h, \overline{m}_{h,3}): \Omega \rightarrow \mathbb{R}^3$ be the average of m_h in (7).

- (i) **Compactness of the global Jacobians and $L^p(\partial\Omega)$ -compactness of the traces $\overline{m}_h|_{\partial\Omega}$.**
For a subsequence $h \rightarrow 0$, $(\mathcal{J}(\overline{m}'_h))_h$ converges to a measure J on the closure $\overline{\Omega}$, in the sense that

$$\lim_{h \rightarrow 0} \left(\sup_{|\nabla' \zeta| \leq 1 \text{ in } \Omega} |\langle \mathcal{J}(\overline{m}'_h) - J, \zeta \rangle| \right) = 0. \quad (18)$$

Moreover, J is supported on $\partial\Omega$ and has the form (12) for $N \geq 1$ distinct boundary vortices $a_j \in \partial\Omega$ carrying the multiplicities $d_j \in \mathbb{Z} \setminus \{0\}$, for $j \in \{1, \dots, N\}$, with $\sum_{j=1}^N d_j = 2$. Furthermore, for a subsequence, $(\overline{m}_h|_{\partial\Omega})_h$ converges to $(e^{i\varphi_0}, 0) \in BV(\partial\Omega, \mathbb{S}^1 \times \{0\})$ in $L^p(\partial\Omega, \mathbb{R}^3)$, for every $p \in [1, +\infty)$, where $\varphi_0 \in BV(\partial\Omega, \pi\mathbb{Z})$ is a lifting of the tangent field $\pm\tau$ on $\partial\Omega$ determined (up to a constant in $\pi\mathbb{Z}$) by $\partial_\tau \varphi_0 = -\mathcal{J}$ as measure on $\partial\Omega$.

- (ii) **Energy lower bound at the first order.**

If $(\mathcal{J}(\overline{m}'_h))_h$ satisfies the convergence assumption in (i) as $h \rightarrow 0$, then the energy lower bound at the first order is the total mass of the measure $J + \kappa \mathcal{H}^1 \llcorner \partial\Omega$ on $\partial\Omega$:

$$\liminf_{h \rightarrow 0} E_h(m_h) \geq \pi \sum_{j=1}^N |d_j| = |J + \kappa \mathcal{H}^1 \llcorner \partial\Omega|(\partial\Omega).$$

As in Theorem 1.4, within a more precise bound similar to (14), we prove the following lower bound at the second order for E_h in the narrower regime (5).

Theorem 1.11. Let $\Omega_h = \Omega \times (0, h)$ with $\Omega \subset \mathbb{R}^2$ a bounded, simply connected $C^{1,1}$ domain, and κ be the curvature of $\partial\Omega$. In the regime (5), consider a family of magnetizations $\{m_h: \Omega_h \rightarrow \mathbb{S}^2\}_{h \downarrow 0}$ that satisfies (17) and the convergence at (i) in Theorem 1.10 to the measure J given at (12) as $h \rightarrow 0$. In addition, we assume the following more precise bound than (17):

$$\limsup_{h \rightarrow 0} |\log \varepsilon| \left(E_h(m_h) - \pi \sum_{j=1}^N |d_j| \right) < +\infty. \quad (19)$$

(i) **Single multiplicity and second order lower bound.**

The multiplicities satisfy⁷ $d_j \in \{\pm 1\}$ for every $j \in \{1, \dots, N\}$, so we have $\sum_{j=1}^N |d_j| = N$, and there holds the following second order energy lower bound:

$$\liminf_{h \rightarrow 0} |\log \varepsilon| (E_h(m_h) - N\pi) \geq W_\Omega^\delta(\{(a_j, d_j)\}) + N\gamma_0$$

where $\gamma_0 = \pi \log \frac{e}{4\pi}$ and $W_\Omega^\delta(\{(a_j, d_j)\})$ is the renormalised energy defined in (13).

(ii) **$L^p(\Omega)$ -compactness of the rescaled magnetizations.**

For a subsequence, the family of rescaled magnetizations $\{\tilde{m}_h: \Omega_1 \rightarrow \mathbb{S}^2\}_{h \downarrow 0}$, defined as $\tilde{m}_h(x', x_3) = m_h(x', hx_3)$, converges strongly in $L^p(\Omega_1, \mathbb{R}^3)$, for every $p \in [1, +\infty)$, to a map $\tilde{m} = (\tilde{m}', 0) \in W^{1,q}(\Omega_1, \mathbb{R}^3)$, for every $q \in [1, 2)$, such that $|\tilde{m}| = |\tilde{m}'| = 1$ and $\partial_3 \tilde{m} = 0$, i.e. $\tilde{m} = \tilde{m}(x') \in W^{1,q}(\Omega, \mathbb{S}^1 \times \{0\})$. Furthermore, the global Jacobian $\mathcal{J}(\tilde{m}')$ coincides with the measure J on $\bar{\Omega}$ given at (12).

For the upper bound part in the Γ -expansion, the recovery sequence can be chosen to be invariant in the thickness direction x_3 and to be in-plane, i.e., they take values into $\mathbb{S}^1 \times \{0\}$.

Theorem 1.12. Let $\Omega_h = \Omega \times (0, h)$ with $\Omega \subset \mathbb{R}^2$ a bounded, simply connected $C^{1,1}$ domain, and κ be the curvature of $\partial\Omega$. Let $\{a_j \in \partial\Omega\}_{1 \leq j \leq N}$ be $N \geq 1$ distinct points and $d_j \in \mathbb{Z} \setminus \{0\}$ be the corresponding multiplicities, for $j \in \{1, \dots, N\}$, that satisfy $\sum_{j=1}^N d_j = 2$. Then in the regime (4), we can construct a family $\{m_h = (m'_h, 0)\}_{h > 0}$ of $H^1(\Omega_h, \mathbb{S}^1 \times \{0\})$ with the following properties.

(i) For every $h > 0$, m_h is independent of x_3 .

(ii) $(\mathcal{J}(m'_h))_h$ converges to $J = -\kappa \mathcal{H}^1 \llcorner \partial\Omega + \pi \sum_{j=1}^N d_j \delta_{a_j}$ as in (18).

(iii) We have $\lim_{h \rightarrow 0} E_h(m_h) = \pi \sum_{j=1}^N |d_j|$.

Furthermore in the regime (5), if $d_j \in \{\pm 1\}$ for every $j \in \{1, \dots, N\}$, then $\{m_h\}_{h \downarrow 0}$ can be chosen such that

$$\lim_{h \rightarrow 0} |\log \varepsilon| (E_h(m_h) - N\pi) = W_\Omega^\delta(\{(a_j, d_j)\}) + N\gamma_0$$

where $\gamma_0 = \pi \log \frac{e}{4\pi}$ and $W_\Omega^\delta(\{(a_j, d_j)\})$ is the renormalised energy defined in (13).

We conclude with the asymptotic behaviour of minimisers of E_h :⁸

Corollary 1.13. Let $\Omega_h = \Omega \times (0, h)$ with $\Omega \subset \mathbb{R}^2$ a bounded, simply connected $C^{1,1}$ domain, and κ be the curvature of $\partial\Omega$. In the regime (4), for every family $\{m_h\}_h$ of minimisers of E_h

⁷Recall that $\sum_{j=1}^N d_j = 2$.

⁸The existence of minimisers of E_h is proved in Corollary 3.2.

on $H^1(\Omega_h, \mathbb{S}^2)$, there exists a subsequence $h \rightarrow 0$ such that the global Jacobians $\mathcal{J}(\overline{m}'_h)$ of the in-plane averages \overline{m}'_h converge as in (18) to the measure

$$J = -\kappa \mathcal{H}^1 \llcorner \partial\Omega + \pi(\delta_{a_1} + \delta_{a_2})$$

where a_1 and a_2 are two points on $\partial\Omega$. Moreover, we have

$$\lim_{h \rightarrow 0} E_h(m_h) = 2\pi.$$

Furthermore, in the regime (5), we have that $a_1 \neq a_2$, the pair (a_1, a_2) minimises the renormalised energy (for the multiplicities $d_1 = d_2 = 1$) over the set $\{(\tilde{a}_1, \tilde{a}_2) \in \partial\Omega \times \partial\Omega : \tilde{a}_1 \neq \tilde{a}_2\}$, i.e.

$$W_\Omega^\delta(\{(a_1, 1), (a_2, 1)\}) = \min \{W_\Omega^\delta(\{(\tilde{a}_1, 1), (\tilde{a}_2, 1)\}) : \tilde{a}_1 \neq \tilde{a}_2 \in \partial\Omega\},$$

and

$$\lim_{h \rightarrow 0} |\log \varepsilon| (E_h(m_h) - 2\pi) = W_\Omega^\delta(\{(a_1, 1), (a_2, 1)\}) + 2\gamma_0$$

where $\gamma_0 = \pi \log \frac{e}{4\pi}$.

In the case of the unit disk $\Omega = B_1$, combining Theorem 1.6 and Corollary 1.13, we deduce that the minimal energy in the regime (5) is given by

$$\min_{H^1(\Omega_h, \mathbb{S}^2)} E_h = 2\pi |\log \varepsilon| - 2\pi \log |a_1^* - a_2^*| + 2\pi\delta \cdot ((a_1^*)^\perp + (a_2^*)^\perp) + 2\pi \log \frac{e}{4\pi} + o_h(1)$$

where (a_1^*, a_2^*) is the optimal pair given in Theorem 1.8.

1.7 Related models

Our thin-film regime (4), that is the same as in [26] (where $D = 0$), favors boundary vortices while taking into account the Dzyaloshinskii-Moriya interaction $D \neq 0$. However, assuming that $h \ll 1$, there are other convergence rates of $\eta(h) \rightarrow 0$ as $h \rightarrow 0$ (different than (4)), that lead to different thin-film regimes and thus to different effects on the magnetizations in the limit $h \rightarrow 0$. More precisely, three types of singular pattern of the magnetization can occur: Néel walls, interior vortices and boundary vortices. The choice of the asymptotic regime corresponds in general to the energy ordering of these three patterns (for more details, see [14]).

The micromagnetic energy has been studied in many thin-film regimes, often without taking into account the Dzyaloshinskii-Moriya interaction. In the regime of small films, where $\eta > 0$ is fixed, Gioia-James [20] proved that the Γ -limit of the micromagnetic energy is minimised by any constant and in-plane magnetization. Recently, Davoli et. al. [13] studied the Γ -convergence of the micromagnetic energy with DMI, and showed that in the limiting energy, the DMI energy contributes to increase the shape anisotropy of the thin film.

For relatively small films, i.e., $\eta^2 \gg h |\log h|$, Kohn-Slastikov [32] (see also Carbou [10]) derived a Γ -limit that reduces to a boundary penalisation term (coming from the stray-field energy, see Section 1.3) for magnetizations that are constant and in-plane. For slightly larger films, where $\eta^2 = \alpha h |\log h|$ with $0 < \alpha < +\infty$, also studied by Kohn-Slastikov [32], there is a competition between exchange and stray-field energies. The limiting energy, for maps $m \in H^1(\Omega, \mathbb{S}^1)$ is composed of the exchange term and a boundary penalisation term (as in the previous regime). The limiting magnetizations are in-plane but no longer constant. Recently, L'Official [36] considered the model of Kohn-Slastikov with DMI, assuming (after a rescaling) that the DMI density is of the same order as the exchange and boundary energy, and derived a Γ -limit for limiting in-plane

magnetizations. We also mention the recent work of Di Fratta et. al. [15] for the study of this regime with DMI.

In the regime $h \ll \eta^2 \ll h |\log h|$, Kurzke [33], [34], [35] and Ignat-Kurzke [26] showed that the limiting magnetization develops boundary vortices due to topological arguments. In the narrower regime $h \log |\log h| \ll \eta^2 \ll h |\log h|$, Ignat-Kurzke [26] derive a Γ -expansion at the second order of the limiting energy. Our aim in this paper consists in taking into account the DMI in these thin-film regimes, and more precisely to study the influence of the DMI on the location of the boundary vortices.

Several other thin-film regimes have been investigated before. In the regime $\eta^2 = O(h)$, Moser [39], [40], [41] showed the emergence of boundary vortices, again for topological reasons. In large thin-films corresponding to the regime $\eta^2 \ll h$, studied by Ignat-Otto [31] and Ignat-Knüpfer [24], other types of topological singularities occur as interior vortices or Néel walls.

1.8 Outline of the paper

The rest of the paper is organised as follows. In Section 2, we prove Theorems 1.2, 1.4 and 1.5, that give the Γ -convergence of the two-dimensional reduced energy $E_{\varepsilon,\eta}^\delta$. Moreover, we study minimisers of the renormalised energy W_Ω^δ , and prove Theorems 1.6 and 1.8, as well as Corollaries 1.7 and 1.9. In Section 3, we reduce the 3D energy E_h to the 2D energy $E_{\varepsilon,\eta}^\delta$, and prove Theorem 1.1. We also prove the Γ -convergence of the 3D energy E_h , i.e. Theorems 1.10, 1.11 and 1.12, and deduce Corollary 1.13.

Acknowledgment. R.I. is partially supported by the ANR projects ANR-21-CE40-0004 and ANR-22-CE40-0006-01.

2 Two-dimensional reduced model for maps $v: \Omega \subset \mathbb{R}^2 \rightarrow \mathbb{R}^2$

Since all quantities in this section are two-dimensional quantities, **we drop the primes ' in the notations**, e.g. $x = (x_1, x_2)$ instead of x' , $\nabla = (\partial_1, \partial_2)$ instead of ∇' , etc.

2.1 Approximation by \mathbb{S}^1 -valued maps

Given a map $v: \Omega \rightarrow \mathbb{R}^2$ such that $E_{\varepsilon,\eta}^\delta(v) = O(|\log \varepsilon|)$, we show in this section that v can be approximated by a \mathbb{S}^1 -valued map $V: \Omega \rightarrow \mathbb{S}^1$ in the regime (9). The idea is to prove that in this context, we have $E_{\varepsilon,\eta}^0(v) = O(|\log \varepsilon|)$, i.e., we reduce to the case $\delta = 0$ and then use the approximation result of Ignat-Kurzke [25, Theorem 3.1].

Lemma 2.1. *Let $\delta \in \mathbb{R}^2$ and $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain. Assume $\varepsilon \rightarrow 0$ and $\eta = \eta(\varepsilon) \rightarrow 0$ in the regime (9). For every $v = v_\varepsilon: \Omega \rightarrow \mathbb{R}^2$ satisfying $E_{\varepsilon,\eta}^\delta(v) = O(|\log \varepsilon|)$, we have $E_{\varepsilon,\eta}^0(v) = O(|\log \varepsilon|)$.*

Proof. We have by Young's inequality

$$|E_{\varepsilon,\eta}^\delta(v) - E_{\varepsilon,\eta}^0(v)| = 2 \left| \int_\Omega \delta \cdot \nabla v \wedge v \, dx \right| \leq \frac{1}{2} \int_\Omega (|\nabla v|^2 + 4|\delta|^2 |v|^2) \, dx.$$

Setting $S = \{x \in \Omega : |v(x)|^2 \geq 2\}$, it follows that for every $x \in S$, $|v(x)|^2 - 1 \geq \frac{|v(x)|^2}{2} \geq \frac{1}{\sqrt{2}} |v(x)|$,

thus

$$\begin{aligned}
|E_{\varepsilon,\eta}^\delta(v) - E_{\varepsilon,\eta}^0(v)| &\leq \frac{1}{2} \int_{\Omega} |\nabla v|^2 dx + 2|\delta|^2 \int_S |v|^2 dx + 2|\delta|^2 \int_{\Omega \setminus S} |v|^2 dx \\
&\leq \frac{1}{2} \int_{\Omega} |\nabla v|^2 dx + 4|\delta|^2 \int_S (1 - |v|^2)^2 dx + 4|\delta|^2 |\Omega \setminus S| \\
&\leq \frac{1}{2} \int_{\Omega} |\nabla v|^2 dx + \frac{1}{2\eta^2} \int_{\Omega} (1 - |v|^2)^2 dx + 4|\delta|^2 |\Omega| \\
&\leq \frac{1}{2} E_{\varepsilon,\eta}^0(v) + 4|\delta|^2 |\Omega|
\end{aligned}$$

because $|\delta|^2 \ll \frac{1}{\eta^2}$ if $\varepsilon > 0$ is small in the regime (9). It follows that

$$E_{\varepsilon,\eta}^0(v) \leq 2E_{\varepsilon,\eta}^\delta(v) + 8|\delta|^2 |\Omega| = O(|\log \varepsilon|).$$

□

Remark 2.2. We assumed in the above lemma that $\delta \in \mathbb{R}^2$ is constant. However, the above proof shows that Lemma 2.1 holds under the weaker assumption $|\delta| = O(|\log \varepsilon|^{1/2})$.

Notation 2.3. Let $\delta \in \mathbb{R}^2$, $\varepsilon > 0$ and $\eta = \eta(\varepsilon) > 0$. For any open set $G \subset \Omega$ and $v: \Omega \rightarrow \mathbb{R}^2$, we define the localised functional

$$\begin{aligned}
E_{\varepsilon,\eta}^\delta(v; G) &= \int_G |\nabla v|^2 dx + 2 \int_G \delta \cdot \nabla v \wedge v dx \\
&\quad + \frac{1}{\eta^2} \int_G (1 - |v|^2)^2 dx + \frac{1}{2\pi\varepsilon} \int_{\overline{G} \cap \partial\Omega} (v \cdot \nu)^2 d\mathcal{H}^1.
\end{aligned}$$

The next theorem gives the approximation of $v: \Omega \rightarrow \mathbb{R}^2$ in $L^2(\Omega)$ by a \mathbb{S}^1 -valued map $V: \Omega \rightarrow \mathbb{S}^1$ that keeps the energy and global Jacobian close to the ones of v . It is reminiscent from [25, Theorem 3.1] where the case $\delta = 0$ was treated.

Theorem 2.4. *Let $\delta \in \mathbb{R}^2$, $\beta \in (\frac{1}{2}, 1)$, $C > 0$ and $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain. Assume $\varepsilon \rightarrow 0$ and $\eta = \eta(\varepsilon) \rightarrow 0$ in the regime (9). There exist $\varepsilon_0 > 0$ and $\tilde{C} > 0$, depending only on β, C, δ, Ω , the function $\varepsilon \mapsto \eta(\varepsilon)$ and $\tilde{\beta} \in (0, \frac{1-\beta}{6})$ such that for every $\varepsilon \in (0, \varepsilon_0)$ and every $v = v_\varepsilon: \Omega \rightarrow \mathbb{R}^2$ satisfying $E_{\varepsilon,\eta}^\delta(v) \leq C|\log \varepsilon|$, we can construct a unit-length map $V = V_\varepsilon: \Omega \rightarrow \mathbb{S}^1$ that satisfies the following relations:*

$$\int_{\Omega} |V - v|^2 dx \lesssim \eta^{2\beta} E_{\varepsilon,\eta}^0(v), \quad \int_{\partial\Omega} |V - v|^2 d\mathcal{H}^1 \lesssim \eta^\beta E_{\varepsilon,\eta}^0(v), \quad (20)$$

and

$$E_{\varepsilon,\eta}^\delta(V) \leq E_{\varepsilon,\eta}^\delta(v) + \tilde{C}\eta^{\tilde{\beta}} \left(E_{\varepsilon,\eta}^0(v) + \sqrt{E_{\varepsilon,\eta}^0(v)} \right). \quad (21)$$

As a consequence, for every $p \in [1, +\infty)$,⁹

$$\begin{aligned}
\lim_{\varepsilon \rightarrow 0} \|V - v\|_{L^p(\Omega)} &= 0, \quad \lim_{\varepsilon \rightarrow 0} \|V - v\|_{L^p(\partial\Omega)} = 0, \\
\|\text{jac}(v)\|_{(W_0^{1,\infty})^*(\Omega)} &\lesssim \eta^\beta E_{\varepsilon,\eta}^0(v), \quad \text{and} \quad \|\mathcal{J}(V) - \mathcal{J}(v)\|_{(\text{Lip}(\Omega))^*} \lesssim \sqrt{\eta^\beta E_{\varepsilon,\eta}^0(v)}.
\end{aligned} \quad (22)$$

⁹Note that $\text{jac}(V) = 0$ as $V \in H^1(\Omega, \mathbb{S}^1)$. For $A \in (W^{1,\infty}(\Omega))^*$, we define $\|A\|_{(\text{Lip}(\Omega))^*} = \sup \{ \langle A, \zeta \rangle : \zeta \in W^{1,\infty}(\Omega), |\nabla \zeta| \leq 1 \}$ and $\|A\|_{(W_0^{1,\infty}(\Omega))^*} = \sup \{ \langle A, \zeta \rangle : \zeta \in W^{1,\infty}(\Omega), |\nabla \zeta| \leq 1, \zeta = 0 \text{ on } \partial\Omega \}$.

The map V also satisfies the following local estimate: for any open set $G \subset \Omega$ independent of ε , there exists a constant $\tilde{C}_G > 0$ such that

$$E_{\varepsilon,\eta}^0(V; G_\eta) \leq E_{\varepsilon,\eta}^0(v; G) + \tilde{C}_G \eta^{\tilde{\beta}} \left(E_{\varepsilon,\eta}^0(v; G) + \sqrt{E_{\varepsilon,\eta}^0(v; G)} \right) \quad (23)$$

where

$$G_\eta = \{x \in G : \text{dist}(x, \Omega \cap \partial G) > c_0 \eta^\beta\},$$

with $c_0 > 0$ depending only on $\partial\Omega$.

Proof. As $E_{\varepsilon,\eta}^0(v) = O(|\log \varepsilon|)$ by Lemma 2.1, we apply [25, Theorem 3.1] to prove the existence of $V: \Omega \rightarrow \mathbb{S}^1$ that satisfies (20), (22) and

$$\int_{\Omega} (|\nabla V|^2 + |\nabla v|^2) dx \lesssim E_{\varepsilon,\eta}^0(v). \quad (24)$$

Moreover, by [25, Theorem 3.1, Equation (31)], we know that

$$E_{\varepsilon,\eta}^0(V) \leq E_{\varepsilon,c_0\eta}^0(v) + \tilde{C} \eta^{\tilde{\beta}} \left(E_{\varepsilon,c_0\eta}^0(v) + \sqrt{E_{\varepsilon,c_0\eta}^0(v)} \right)$$

for some $c_0 > 0$. As $|V| = 1$, note that $E_{\varepsilon,\eta}^0(V) = E_{\varepsilon,c_0\eta}^0(V)$, so we can replace η by $\hat{\eta} = c_0\eta$. We still denote η instead of $\hat{\eta}$ in the following. To prove (21), we compute

$$E_{\varepsilon,\eta}^\delta(V) = E_{\varepsilon,\eta}^0(V) + 2 \int_{\Omega} \delta \cdot (\nabla V \wedge V - \nabla v \wedge v) dx + 2 \int_{\Omega} \delta \cdot \nabla v \wedge v dx$$

so that

$$E_{\varepsilon,\eta}^\delta(V) \leq E_{\varepsilon,\eta}^\delta(v) + \tilde{C} \eta^{\tilde{\beta}} \left(E_{\varepsilon,\eta}^0(v) + \sqrt{E_{\varepsilon,\eta}^0(v)} \right) + 2 \int_{\Omega} \delta \cdot (\nabla V \wedge V - \nabla v \wedge v) dx.$$

By integration by parts, we have

$$\begin{aligned} \int_{\Omega} \delta \cdot (\nabla V \wedge V - \nabla v \wedge v) dx &= \int_{\Omega} \delta \cdot ((\nabla V - \nabla v) \wedge V - \nabla v \wedge (v - V)) dx \\ &= \int_{\partial\Omega} (\delta \cdot \nu)(V - v) \wedge V d\mathcal{H}^1 - \int_{\Omega} \delta \cdot (V - v) \wedge (\nabla V + \nabla v) dx. \end{aligned}$$

We deduce that

$$\begin{aligned} \left| \int_{\Omega} \delta \cdot (\nabla V \wedge V - \nabla v \wedge v) dx \right| &\leq \int_{\partial\Omega} |\delta \cdot \nu| |V - v| d\mathcal{H}^1 + |\delta| \int_{\Omega} |V - v| (|\nabla V| + |\nabla v|) dx \\ &\lesssim \sqrt{\int_{\partial\Omega} |V - v|^2 d\mathcal{H}^1} + \sqrt{\int_{\Omega} |V - v|^2 dx} \sqrt{\int_{\Omega} (|\nabla V|^2 + |\nabla v|^2) dx}, \\ &\lesssim \eta^\beta E_{\varepsilon,\eta}^0(v) + \eta^{\beta/2} \sqrt{E_{\varepsilon,\eta}^0(v)} \\ &\lesssim \eta^{\tilde{\beta}} \left(E_{\varepsilon,\eta}^0(v) + \sqrt{E_{\varepsilon,\eta}^0(v)} \right) \end{aligned}$$

where we used (20), (24) and $\frac{\beta}{2} \geq \frac{1}{4} > \tilde{\beta}$. We conclude to (21). Finally, (23) is reminiscent from [25, Theorem 3.1, Equation (32)] and the fact that $E_{\varepsilon,\eta}^0(V) = E_{\varepsilon,c_0\eta}^0(V)$ as $|V| = 1$. \square

2.2 Lifting

By Theorem 2.4, our study simplifies to the analysis of the energy functional $E_{\varepsilon,\eta}^\delta$ for \mathbb{S}^1 -valued maps V . Such maps have a lifting (see Bethuel-Zheng [7]) on which we focus in the following:

Lemma 2.5. *Let $\delta \in \mathbb{R}^2$, $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain and $V \in H^1(\Omega, \mathbb{S}^1)$. There exists a lifting $\varphi \in H^1(\Omega, \mathbb{R})$ such that $V = e^{i\varphi}$ and φ is unique up to an additive constant in $2\pi\mathbb{Z}$. Furthermore, for every $\varepsilon > 0$ and $\eta > 0$,*

$$E_{\varepsilon,\eta}^\delta(V) = \int_{\Omega} \left(|\nabla\varphi|^2 - 2\delta \cdot \nabla\varphi \right) dx + \frac{1}{2\pi\varepsilon} \int_{\partial\Omega} \sin^2(\varphi - g) d\mathcal{H}^1 =: \mathcal{G}_\varepsilon^\delta(\varphi), \quad (25)$$

where g is a lifting of the unit tangent vector field τ on $\partial\Omega$, i.e. as complex numbers

$$e^{ig} = \tau = i\nu \quad \text{on } \partial\Omega, \quad (26)$$

and g is continuous except at one point of $\partial\Omega$.

Proof. Existence and uniqueness of a lifting φ of V in Ω come from a well-known theorem of Bethuel and Zheng [7]. For the existence of g , we note that τ has winding number 1 on $\partial\Omega$ as Ω is simply connected, hence τ cannot be lifted continuously on $\partial\Omega$. However, if $\partial\Omega$ is $C^{1,1}$, we can choose a lifting g to be locally Lipschitz except at one point of $\partial\Omega$ where it jumps by 2π . Clearly, the curvature κ of $\partial\Omega$ is given by the absolutely continuous part of the derivative of g (as a BV function), i.e. $\kappa = (\partial_\tau g)_{ac}$ and $\int_{\partial\Omega} \kappa d\mathcal{H}^1 = 2\pi$, which is the Gauss-Bonnet formula on $\partial\Omega$. Therefore, $g \in BV(\partial\Omega, \mathbb{R})$ with $\partial_\tau g = \kappa \mathcal{H}^1 \llcorner \partial\Omega - 2\pi\delta_p$ for some point $p \in \partial\Omega$. As $|\nabla V| = |\nabla\varphi|$, $\nabla V \wedge V = -\nabla\varphi$ in Ω , and $V \cdot \nu = -\sin(\varphi - g)$ on $\partial\Omega$, we deduce (25). \square

The functional $\mathcal{G}_\varepsilon^\delta$ in the above lemma has been studied by Kurzke [33], [34] for $\delta = 0$. In the following, we will prove Γ -convergence for $\mathcal{G}_\varepsilon^\delta$ and use these result for proving Γ -convergence for $E_{\varepsilon,\eta}^\delta$.

2.3 Gamma-convergence in terms of liftings

We now present the Γ -convergence for the functional $\mathcal{G}_\varepsilon^\delta$ defined in (25) in terms of the lifting $\varphi_\varepsilon: \Omega \rightarrow \mathbb{R}$. These results will be useful to deduce similar statements for $E_{\varepsilon,\eta}^\delta(v_\varepsilon)$ in the next section. The first statement establishes the $L^p(\partial\Omega)$ -compactness of φ_ε and a lower bound for $\mathcal{G}_\varepsilon^\delta$ at the first order.

Theorem 2.6. *Let $\delta \in \mathbb{R}^2$, $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected and $C^{1,1}$ domain and κ be the curvature of $\partial\Omega$. Let $(\varphi_\varepsilon)_{\varepsilon \downarrow 0}$ be a family in $H^1(\Omega)$ such that*

$$\limsup_{\varepsilon \rightarrow 0} \frac{1}{|\log \varepsilon|} \mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon) < +\infty. \quad (27)$$

There exists a family $(z_\varepsilon)_\varepsilon$ of integers such that $(\varphi_\varepsilon - \pi z_\varepsilon)_\varepsilon$ is bounded in $L^p(\partial\Omega)$, for every $p \in [1, +\infty)$. Moreover, for a subsequence, $(\varphi_\varepsilon - \pi z_\varepsilon)_\varepsilon$ converges strongly in $L^p(\partial\Omega)$ to a limit φ_0 such that $\varphi_0 - g \in BV(\partial\Omega, \pi\mathbb{Z})$, with g given in (26), and

$$\partial_\tau \varphi_0 = \kappa \mathcal{H}^1 \llcorner \partial\Omega - \pi \sum_{j=1}^N d_j \delta_{a_j} \quad \text{as measure on } \partial\Omega$$

where $N \geq 1$, $\{a_j \in \partial\Omega\}_{1 \leq j \leq N}$ are distinct points, $d_j \in \mathbb{Z} \setminus \{0\}$, $\sum_{j=1}^N d_j = 2$, and $(\partial_\tau \varphi_\varepsilon)_\varepsilon$ converges to $\partial_\tau \varphi_0$ in $W^{-1,p}(\partial\Omega)$ for every $p \in [1, +\infty)$. Furthermore, we have the following first order lower bound for the energy:

$$\liminf_{\varepsilon \rightarrow 0} \frac{1}{|\log \varepsilon|} \mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon) \geq |\partial_\tau \varphi_0 - \kappa \mathcal{H}^1 \llcorner \partial\Omega|(\partial\Omega) = \pi \sum_{j=1}^N |d_j|. \quad (28)$$

Proof. Note that, for any $\varepsilon > 0$ and any $\sigma > 0$, Young's inequality yields

$$\int_{\Omega} |\nabla \varphi_{\varepsilon}|^2 dx = \int_{\Omega} |(\nabla \varphi_{\varepsilon} - \delta) + \delta|^2 dx \leq (1 + \sigma) \int_{\Omega} |\nabla \varphi_{\varepsilon} - \delta|^2 dx + \left(1 + \frac{1}{\sigma}\right) |\delta|^2 |\Omega|,$$

since $(a + b)^2 = a^2 + b^2 + 2ab \leq a^2 + b^2 + (\sqrt{\sigma}a)^2 + \left(\frac{1}{\sqrt{\sigma}}b\right)^2$ for every $a, b \in \mathbb{R}$. Hence,

$$\begin{aligned} \mathcal{G}_{\varepsilon}^0(\varphi_{\varepsilon}) &\leq (1 + \sigma) \int_{\Omega} |\nabla \varphi_{\varepsilon} - \delta|^2 dx + \frac{1}{2\pi\varepsilon} \int_{\partial\Omega} \sin^2(\varphi_{\varepsilon} - g) d\mathcal{H}^1 + \left(1 + \frac{1}{\sigma}\right) |\delta|^2 |\Omega| \\ &\leq (1 + \sigma) \int_{\Omega} (|\nabla \varphi_{\varepsilon}|^2 - 2\delta \cdot \nabla \varphi_{\varepsilon}) dx + \frac{1 + \sigma}{2\pi\varepsilon} \int_{\partial\Omega} \sin^2(\varphi_{\varepsilon} - g) d\mathcal{H}^1 + \left(2 + \sigma + \frac{1}{\sigma}\right) |\delta|^2 |\Omega| \\ &\leq (1 + \sigma) \mathcal{G}_{\varepsilon}^{\delta}(\varphi_{\varepsilon}) + \left(2 + \sigma + \frac{1}{\sigma}\right) |\delta|^2 |\Omega|. \end{aligned} \quad (29)$$

By (27), we deduce that $\limsup_{\varepsilon \rightarrow 0} \frac{\mathcal{G}_{\varepsilon}^0(\varphi_{\varepsilon})}{|\log \varepsilon|} < +\infty$. Hence, we can apply [25, Theorem 4.2] and deduce the desired compactness results. Moreover, by (29) and [25, Theorem 4.2], we have

$$\liminf_{\varepsilon \rightarrow 0} (1 + \sigma) \frac{\mathcal{G}_{\varepsilon}^{\delta}(\varphi_{\varepsilon})}{|\log \varepsilon|} \geq \liminf_{\varepsilon \rightarrow 0} \frac{\mathcal{G}_{\varepsilon}^0(\varphi_{\varepsilon})}{|\log \varepsilon|} \geq \pi \sum_{j=1}^N |d_j|.$$

Letting $\sigma \rightarrow 0$, we get (28). \square

The following theorem gives the lower bound of $\mathcal{G}_{\varepsilon}^{\delta}$ of second order together with the multiplicities ± 1 of boundary vortices provided a more precise energy estimate than (28).

Theorem 2.7. *Let $\delta \in \mathbb{R}^2$, $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected and $C^{1,1}$ domain and κ be the curvature of $\partial\Omega$. Let $(\varphi_{\varepsilon})_{\varepsilon \downarrow 0}$ be a family in $H^1(\Omega)$ satisfying the convergence in Theorem 2.6 with the limit φ_0 on $\partial\Omega$ as $\varepsilon \rightarrow 0$. Assume additionally that*

$$\limsup_{\varepsilon \rightarrow 0} \left(\mathcal{G}_{\varepsilon}^{\delta}(\varphi_{\varepsilon}) - \pi |\log \varepsilon| \sum_{j=1}^N |d_j| \right) < +\infty. \quad (30)$$

Then $d_j \in \{\pm 1\}$ for every $j \in \{1, \dots, N\}$ and, for a subsequence, $(\nabla \varphi_{\varepsilon})_{\varepsilon}$ converges weakly in $L^q(\Omega, \mathbb{R}^2)$ for any $q \in [1, 2)$ to $\nabla \widehat{\varphi}_0$, where $\widehat{\varphi}_0 \in W^{1,q}(\Omega)$ is an extension (not necessarily harmonic) of φ_0 to Ω . Furthermore, we have the following second order lower bound of the energy:

$$\liminf_{\varepsilon \rightarrow 0} (\mathcal{G}_{\varepsilon}^{\delta}(\varphi_{\varepsilon}) - N\pi |\log \varepsilon|) \geq W_{\Omega}^{\delta}(\{(a_j, d_j)\}) + N\gamma_0, \quad (31)$$

where $\gamma_0 = \pi \log \frac{e}{4\pi}$ and $W_{\Omega}^{\delta}(\{(a_j, d_j)\})$ is the renormalised energy defined in (13).

Proof. We adapt the proof of [25, Theorem 4.2] to the case of $\delta \in \mathbb{R}^2$.

Step 1: We prove that $d_j \in \{\pm 1\}$ for every $j \in \{1, \dots, N\}$ and the weak convergence of $(\nabla \varphi_{\varepsilon})_{\varepsilon \downarrow 0}$. For that, note that by Theorem 2.6, there exists $z_{\varepsilon} \in \mathbb{Z}$ such that

$$\int_{\Omega} \delta \cdot \nabla \varphi_{\varepsilon} dx = \int_{\Omega} \delta \cdot \nabla (\varphi_{\varepsilon} - \pi z_{\varepsilon}) dx = \int_{\partial\Omega} (\varphi_{\varepsilon} - \pi z_{\varepsilon}) \delta \cdot \nu d\mathcal{H}^1 \leq |\delta| \|\varphi_{\varepsilon} - \pi z_{\varepsilon}\|_{L^1(\partial\Omega)} \leq C.$$

As

$$\mathcal{G}_{\varepsilon}^{\delta}(\varphi_{\varepsilon}) = \mathcal{G}_{\varepsilon}^0(\varphi_{\varepsilon}) - 2 \int_{\Omega} \delta \cdot \nabla \varphi_{\varepsilon} dx \geq \mathcal{G}_{\varepsilon}^0(\varphi_{\varepsilon}) - 2C,$$

we deduce that (30) holds for $\mathcal{G}_\varepsilon^0(\varphi_\varepsilon)$ and by [25, Theorem 4.2, part 2], it yields the claim in Step 1. Moreover, by [25, Proof of Theorem 4.2], we have:

$$\int_{\Omega \setminus \bigcup_j B_r(a_j)} |\nabla \varphi_\varepsilon|^2 dx \leq N\pi \log \frac{1}{r} + C,$$

for every small $r > 0$ and $\varepsilon > 0$.

Step 2: We prove the second order lower bound (31). To do so, we replace φ_ε by the harmonic extension of $\varphi_\varepsilon|_{\partial\Omega}$ into Ω , denoted by φ_ε^* . More precisely, φ_ε^* is the minimiser of the Dirichlet energy in Ω under the Dirichlet boundary condition $\varphi_\varepsilon|_{\partial\Omega}$. In particular, φ_ε^* is harmonic in Ω . Since

$$\int_{\Omega} \delta \cdot \nabla \varphi_\varepsilon dx = \int_{\partial\Omega} \varphi_\varepsilon \delta \cdot \nu d\mathcal{H}^1 = \int_{\partial\Omega} \varphi_\varepsilon^* \delta \cdot \nu d\mathcal{H}^1 = \int_{\Omega} \delta \cdot \nabla \varphi_\varepsilon^* dx,$$

we deduce that $\mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon) \geq \mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon^*)$, thus it suffices to prove (31) for φ_ε^* . As $\{a_j, d_j\}_{1 \leq j \leq N}$ are determined by $\varphi_\varepsilon|_{\partial\Omega}$, the right-hand side in (31) remains the same when replacing φ_ε by φ_ε^* . Using Step 1, we know that (φ_ε^*) converges weakly in $W^{1,q}(\Omega)$, for every $q \in [1, 2)$, and weakly in $H^1(\omega)$, for any open set ω such that $\bar{\omega} \subset \bar{\Omega} \setminus \{a_1, \dots, a_N\}$, to the harmonic extension φ_* of φ_0 to Ω . Let $r > 0$ be small and $\Omega^r := \Omega \setminus \bigcup_{j=1}^N B_r(a_j)$. By weak lower semicontinuity of the Dirichlet integral, we have

$$\int_{\Omega^r} |\nabla \varphi_*|^2 dx \leq \liminf_{\varepsilon \rightarrow 0} \int_{\Omega^r} |\nabla \varphi_\varepsilon^*|^2 dx.$$

Also, by weak convergence of $(\nabla \varphi_\varepsilon^*)_\varepsilon$ to $\nabla \varphi_*$ in $L^1(\Omega, \mathbb{R}^2)$, we have

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega^r} \delta \cdot \nabla \varphi_\varepsilon^* dx = \int_{\Omega^r} \delta \cdot \nabla \varphi_* dx.$$

By definition of W_Ω^δ in (13), we have

$$\int_{\Omega^r} (|\nabla \varphi_*|^2 - 2\delta \cdot \nabla \varphi_*) dx \geq \pi N \log \frac{1}{r} + W_\Omega^\delta(\{a_j, d_j\}) + o_r(1), \text{ as } r \rightarrow 0.$$

In $\Omega \setminus \Omega^r$, we have by Hölder's inequality:

$$\int_{\Omega \setminus \Omega^r} \delta \cdot \nabla \varphi_\varepsilon^* dx = \sum_{j=1}^N \int_{\Omega \cap B_r(a_j)} \delta \cdot \nabla \varphi_\varepsilon^* dx \leq C |\delta| r^{2/3} \|\nabla \varphi_\varepsilon^*\|_{L^{3/2}(\Omega)} \leq Cr^{2/3},$$

for some constant $C > 0$, because $(\nabla \varphi_\varepsilon^*)_{\varepsilon \downarrow 0}$ is bounded (independently of ε) in $L^{3/2}(\Omega, \mathbb{R}^2)$. By [25, Equation (85)], we have

$$\begin{aligned} \liminf_{\varepsilon \rightarrow 0} \left[\sum_{j=1}^N \left(\int_{\Omega \cap B_r(a_j)} |\nabla \varphi_\varepsilon^*|^2 dx + \frac{1}{2\pi\varepsilon} \int_{\partial\Omega \cap B_r(a_j)} \sin^2(\varphi_\varepsilon^* - g) d\mathcal{H}^1 \right) - N\pi \log \frac{r}{\varepsilon} \right] \\ \geq -CNr^{1/2} + N\gamma_0. \end{aligned}$$

Therefore, we conclude

$$\liminf_{\varepsilon \rightarrow 0} (\mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon^*) - N\pi |\log \varepsilon|) \geq W_\Omega^\delta(\{a_j, d_j\}) - CNr^{1/2} + N\gamma_0 - O(r^{2/3}).$$

Taking the limit as $r \rightarrow 0$, we get (31) for φ_ε^* . \square

Now we prove the upper bound for the Γ -convergence of $\mathcal{G}_\varepsilon^\delta$ at the first order, respectively at the second order.

Theorem 2.8. Let $\delta \in \mathbb{R}^2$, $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain and κ be the curvature of $\partial\Omega$. Let $\varphi_0: \partial\Omega \rightarrow \mathbb{R}$ be such that

$$\partial_\tau \varphi_0 = \kappa \mathcal{H}^1 \llcorner \partial\Omega - \pi \sum_{j=1}^N d_j \delta_{a_j} \quad \text{as measure on } \partial\Omega,$$

where $d_j \in \mathbb{Z} \setminus \{0\}$ for every $j \in \{1, \dots, N\}$, $\sum_{j=1}^N d_j = 2$ and $e^{i\varphi_0} \cdot \nu = 0$ in $\partial\Omega \setminus \{a_1, \dots, a_N\}$ for $N \geq 1$ distinct points $a_1, \dots, a_N \in \partial\Omega$. There exists a family $(\varphi_\varepsilon)_\varepsilon$ in $H^1(\Omega)$ such that $(\varphi_\varepsilon)_\varepsilon$ converges to φ_0 in $L^p(\partial\Omega)$ and to φ_* in $L^p(\Omega)$, for every $p \in [1, +\infty)$, where φ_* is the harmonic extension of φ_0 to Ω , and we have

$$\lim_{\varepsilon \rightarrow 0} \frac{1}{|\log \varepsilon|} \mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon) \leq \pi \sum_{j=1}^N |d_j|. \quad (32)$$

Furthermore, if $d_j \in \{\pm 1\}$ for every $j \in \{1, \dots, N\}$, then

$$\lim_{\varepsilon \rightarrow 0} (\mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon) - N\pi |\log \varepsilon|) \leq W_\Omega^\delta(\{(a_j, d_j)\}) + N\gamma_0, \quad (33)$$

where $\gamma_0 = \pi \log \frac{e}{4\pi}$ and $W_\Omega^\delta(\{(a_j, d_j)\})$ is the renormalised energy defined in (13).

Proof. Let φ_* be the harmonic extension of φ_0 to Ω , that satisfies (13). We consider the family $(\widehat{\varphi}_\varepsilon)$ constructed in [25, Theorem 4.2, part 3] that satisfies the required convergence to φ_* and estimates (32) and (33) for $\delta = 0$. We will show that this family $(\widehat{\varphi}_\varepsilon)$ satisfies (32) and (33) for general $\delta \in \mathbb{R}^2$.

Case 1: $d_j \in \mathbb{Z}$ are not necessarily equal to ± 1 . Then for $\delta \in \mathbb{R}^2$, recall that

$$\mathcal{G}_\varepsilon^\delta(\widehat{\varphi}_\varepsilon) = \mathcal{G}_\varepsilon^0(\widehat{\varphi}_\varepsilon) - 2 \int_\Omega \delta \cdot \nabla \widehat{\varphi}_\varepsilon \, dx = \mathcal{G}_\varepsilon^0(\widehat{\varphi}_\varepsilon) + 2 \int_{\partial\Omega} \widehat{\varphi}_\varepsilon \, \delta \cdot \nu \, d\mathcal{H}^1.$$

As $\widehat{\varphi}_\varepsilon \rightarrow \varphi_*$ in $L^1(\partial\Omega)$, we deduce that $|\int_{\partial\Omega} \widehat{\varphi}_\varepsilon \, \delta \cdot \nu \, d\mathcal{H}^1| \leq C$ as $\varepsilon \downarrow 0$, which yields (32) for general $\delta \in \mathbb{R}^2$.

Case 2: $|d_j| = 1$ for every $j \in \{1, \dots, N\}$. By the definition of $W_\Omega^0(\{(a_j, d_j)\})$ and the construction in [25, Theorem 4.2, part 3] (satisfying (33) above for $\delta = 0$), we have:

$$\mathcal{G}_\varepsilon^0(\widehat{\varphi}_\varepsilon) \leq N\pi |\log \varepsilon| + \left(\int_{\Omega \setminus \bigcup_{j=1}^N B_r(a_j)} |\nabla \varphi_*|^2 \, dx - N\pi \log \frac{1}{r} \right) + N\gamma_0 + o_\varepsilon(1) + o_r(1).$$

Moreover, we have

$$\int_\Omega \delta \cdot \nabla \widehat{\varphi}_\varepsilon \, dx = \int_{\Omega \setminus \bigcup_j B_r(a_j)} \delta \cdot \nabla \widehat{\varphi}_\varepsilon \, dx + \int_{\Omega \cap \bigcup_j B_r(a_j)} \delta \cdot \nabla \widehat{\varphi}_\varepsilon \, dx.$$

Let $r > 0$ be sufficiently small. Then by Hölder's inequality,

$$\int_{\Omega \cap \bigcup_j B_r(a_j)} \delta \cdot \nabla \widehat{\varphi}_\varepsilon \, dx = \sum_{j=1}^N \int_{\Omega \cap B_r(a_j)} \delta \cdot \nabla \widehat{\varphi}_\varepsilon \, dx \leq C |\delta| r^{2/3} \|\nabla \widehat{\varphi}_\varepsilon\|_{L^{3/2}(\Omega)} = O(r^{2/3})$$

because $(\nabla \widehat{\varphi}_\varepsilon)_\varepsilon$ is bounded in $L^{3/2}(\Omega)$ by construction (see [25, Theorem 4.2, part 2]). Moreover, by [25, Theorem 4.2, part 2] again, $(\nabla \widehat{\varphi}_\varepsilon)$ converges weakly to $\nabla \varphi_*$ in $L^1(\Omega, \mathbb{R}^2)$, thus

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega \setminus \bigcup_j B_r(a_j)} \delta \cdot \nabla \widehat{\varphi}_\varepsilon \, dx = \int_{\Omega \setminus \bigcup_j B_r(a_j)} \delta \cdot \nabla \varphi_* \, dx$$

Combining all above, and taking the lim sup as $\varepsilon \rightarrow 0$, we deduce

$$\limsup_{\varepsilon \rightarrow 0} \left(\mathcal{G}_\varepsilon^\delta(\widehat{\psi}_\varepsilon) - N\pi |\log \varepsilon| \right) \leq \left(\int_{\Omega \setminus \bigcup_{j=1}^N B_r(a_j)} \left(|\nabla \varphi_*|^2 - 2\delta \cdot \nabla \varphi_* \right) dx - N\pi \log \frac{1}{r} \right) + N\gamma_0 + o_r(1).$$

Taking the lim inf as $r \rightarrow 0$, we get the desired upper bound. \square

2.4 Gamma-convergence for vector-valued maps

We now prove Theorems 1.2, 1.4 and 1.5.

Proof of Theorem 1.2. Let (v_ε) be a family in $H^1(\Omega, \mathbb{R}^2)$ such that $E_{\varepsilon, \eta}^\delta(v_\varepsilon) \leq C |\log \varepsilon|$, for some constant $C > 0$. Using Theorem 2.4 and Lemma 2.1, we can construct a family (V_ε) in $H^1(\Omega, \mathbb{S}^1)$ such that in the regime (9):

$$E_{\varepsilon, \eta}^\delta(V_\varepsilon) \leq E_{\varepsilon, \eta}^\delta(v_\varepsilon) + o_\varepsilon(1), \quad (34)$$

$$\lim_{\varepsilon \rightarrow 0} \|V_\varepsilon - v_\varepsilon\|_{L^p(\partial\Omega)} = 0, \quad \text{for every } p \in [1, +\infty), \quad (35)$$

and

$$\lim_{\varepsilon \rightarrow 0} \|\mathcal{J}(V_\varepsilon) - \mathcal{J}(v_\varepsilon)\|_{(\text{Lip}(\Omega))^*} = 0. \quad (36)$$

Using Lemma 2.5, for every $\varepsilon > 0$, there exists a lifting $\varphi_\varepsilon \in H^1(\Omega)$ such that $V_\varepsilon = e^{i\varphi_\varepsilon}$ and $E_{\varepsilon, \eta}^\delta(V_\varepsilon) = \mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon)$. Moreover, for every $\varepsilon > 0$, the global Jacobian of V_ε is given by

$$\mathcal{J}(V_\varepsilon) = \mathcal{J}_{\text{bd}}(V_\varepsilon) = -\partial_\tau \varphi_\varepsilon \text{ as a distribution in } H^{-1/2}(\partial\Omega).$$

Using (34) and $\mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon) = E_{\varepsilon, \eta}^\delta(V_\varepsilon)$, we deduce from Theorem 2.6 that, for a subsequence, there exists a family of integers (z_ε) – that we can all assume to be either even or odd, up to take a further subsequence – such that $(\varphi_\varepsilon - \pi z_\varepsilon)$ converges strongly in $L^p(\partial\Omega)$, for every $p \in [1, +\infty)$, to a limit ϕ_0 that satisfies $\phi_0 - g \in BV(\partial\Omega, \pi\mathbb{Z})$ with g such that $e^{ig} = \tau = i\nu$ on $\partial\Omega$. Let φ_0 be such that $\varphi_0 = \phi_0$ if the integers z_ε are all even, and $\varphi_0 = \phi_0 - \pi$ if the integers z_ε are all odd. By definition of ϕ_0 and g , φ_0 is a BV lifting of $\pm\tau$. Moreover, as $|e^{is} - e^{it}| \leq |s - t|$ for every $s, t \in \mathbb{R}$, we have, for every $\varepsilon > 0$,

$$|V_\varepsilon - e^{i\varphi_0}| = \left| e^{i(\varphi_\varepsilon - \pi z_\varepsilon)} - e^{i\varphi_0} \right| \leq |(\varphi_\varepsilon - \pi z_\varepsilon) - \varphi_0|.$$

It follows that $(V_\varepsilon)_{\varepsilon \downarrow 0}$ converges strongly to $e^{i\varphi_0}$ in $L^p(\partial\Omega)$, for every $p \in [1, +\infty)$. Combining this with (35), we deduce that (v_ε) converges strongly to $e^{i\varphi_0}$ in $L^p(\partial\Omega)$, for every $p \in [1, +\infty)$. By Theorem 2.6, we also have

$$\partial_\tau \varphi_0 = \partial_\tau \phi_0 = \kappa \mathcal{H}^1 \llcorner \partial\Omega - \pi \sum_{j=1}^N d_j \delta_{a_j} = -J \quad \text{as measure on } \partial\Omega$$

where, for every $j \in \{1, \dots, N\}$, $a_j \in \partial\Omega$ are distinct points, $d_j \in \mathbb{Z} \setminus \{0\}$ and $\sum_{j=1}^N d_j = 2$. We also get, again by Theorem 2.6, the convergence of $(\partial_\tau \varphi_\varepsilon)$ to $\partial_\tau \phi_0$ in $W^{-1,p}(\partial\Omega)$ for every $p \in (1, +\infty)$. For every Lipschitz function $\zeta \in W^{1,\infty}(\Omega)$, using that $V_\varepsilon = e^{i\varphi_\varepsilon}$ and integrating by parts, we have

$$\begin{aligned} \langle \mathcal{J}(V_\varepsilon), \zeta \rangle &= - \int_{\Omega} V_\varepsilon \wedge \nabla V_\varepsilon \cdot \nabla^\perp \zeta \, dx = - \int_{\Omega} \nabla \varphi_\varepsilon \cdot \nabla^\perp \zeta \, dx \\ &= - \langle \partial_\tau \varphi_\varepsilon, \zeta \rangle_{H^{-1/2}(\partial\Omega), H^{1/2}(\partial\Omega)} = - \langle \partial_\tau \varphi_\varepsilon, \zeta \rangle_{W^{-1,2}(\partial\Omega), W^{1,2}(\partial\Omega)}. \end{aligned}$$

Thus,

$$\sup_{|\nabla\zeta|\leq 1} |\langle \mathcal{J}(v_\varepsilon) - J, \zeta \rangle| \leq C \|\partial_\tau \varphi_\varepsilon - \partial_\tau \varphi_0\|_{W^{-1,2}(\partial\Omega)} \rightarrow 0 \text{ as } \varepsilon \rightarrow 0.$$

Combining this with (36), we deduce that $(\mathcal{J}(v_\varepsilon))$ converges to J in $(\text{Lip}(\Omega))^*$. Finally, since

$$E_{\varepsilon,\eta}^\delta(v_\varepsilon) \geq E_{\varepsilon,\eta}^\delta(V_\varepsilon) - o_\varepsilon(1) = \mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon) - o_\varepsilon(1)$$

thanks to (34), the lower bound (28) for $\mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon)$ given by Theorem 2.6 gives the expected lower bound for $E_{\varepsilon,\eta}^\delta(v_\varepsilon)$ at the first order. \square

We will now prove Theorem 1.4. Within those assumptions, we additionally prove the following statements.

(a) **Penalty bound.**

The penalty terms in the energy are of order $O(1)$, i.e.

$$\limsup_{\varepsilon \rightarrow 0} \left(\frac{1}{\eta^2} \int_\Omega (1 - |v_\varepsilon|^2)^2 dx + \frac{1}{2\pi\varepsilon} \int_{\partial\Omega} (v_\varepsilon \cdot \nu)^2 d\mathcal{H}^1 \right) < +\infty. \quad (37)$$

(b) **Lower bound for the energy near boundary vortex cores.**

There exist $r_0 > 0$, $\varepsilon_0 > 0$ and $C > 0$ such that the Dirichlet energy of v_ε near the singularities $\{a_j\}_{j \in \{1, \dots, N\}}$ satisfies, for all $\varepsilon \in (0, \varepsilon_0)$ and $r \in (0, r_0)$,

$$\int_{\Omega \cap \bigcup_j B_r(a_j)} |\nabla v_\varepsilon|^2 dx - N\pi \log \frac{r}{\varepsilon} \geq -C. \quad (38)$$

(c) **DMI bound.**

The Dzyaloshinskii-Moriya interaction energy is of order $O(1)$, i.e.

$$\limsup_{\varepsilon \rightarrow 0} \int_\Omega |\delta \cdot \nabla v_\varepsilon \wedge v_\varepsilon| dx < +\infty. \quad (39)$$

Proof of Theorem 1.4. Continuing as in the proof of Theorem 1.2 (with the same notations), we now assume the stronger condition (14). By definition of V_ε , Theorem 2.4 and Lemma 2.1,

$$\mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon) = E_{\varepsilon,\eta}^\delta(V_\varepsilon) \leq E_{\varepsilon,\eta}^\delta(v_\varepsilon) + o_\varepsilon(1),$$

hence by (14),

$$\limsup_{\varepsilon \rightarrow 0} \left(\mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon) - \pi |\log \varepsilon| \sum_{j=1}^N |d_j| \right) < +\infty. \quad (40)$$

Step 1: Proof of (i). We deduce immediately from the above assumption and Theorem 2.7 that (30) holds true and $d_j \in \{\pm 1\}$ for every $j \in \{1, \dots, N\}$. Since $E_{\varepsilon,\eta}^\delta(v_\varepsilon) \geq \mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon) - o_\varepsilon(1)$, it follows that the desired lower bound for $E_{\varepsilon,\eta}^\delta(v_\varepsilon)$ holds true.

Step 2: Estimating the DMI term in (39). By Theorem 2.4 and Lemma 2.1, we have

$$\begin{aligned} \left| \int_\Omega \delta \cdot \nabla v_\varepsilon \wedge v_\varepsilon dx \right| &\leq \left| \int_\Omega \delta \cdot \nabla V_\varepsilon \wedge V_\varepsilon dx \right| + \left| \int_\Omega \delta \cdot (\nabla V_\varepsilon \wedge V_\varepsilon - \nabla v_\varepsilon \wedge v_\varepsilon) dx \right| \\ &\leq \left| \int_\Omega \delta \cdot \nabla V_\varepsilon \wedge V_\varepsilon dx \right| + |\delta| \|\mathcal{J}(V_\varepsilon) - \mathcal{J}(v_\varepsilon)\|_{(\text{Lip}(\Omega))^*} \\ &\leq \left| \int_\Omega \delta \cdot \nabla \varphi_\varepsilon dx \right| + o_\varepsilon(1). \end{aligned} \quad (41)$$

By Theorem 2.7, $(\nabla\varphi_\varepsilon)$ is bounded in $L^1(\Omega)$ and we conclude with (39).

Step 3: Compactness and proofs of (37) and (38). By (39), one checks that

$$E_{\varepsilon,\eta}^\delta(v_\varepsilon) = E_{\varepsilon,\eta}^0(v_\varepsilon) + O_\varepsilon(1).$$

In particular, $E_{\varepsilon,\eta}^0(v_\varepsilon)$ satisfies (14) for $\delta = 0$. Therefore, we apply Theorem 1.4 in [25] and conclude this step. \square

Proof of Theorem 1.5. Let $\varphi_0: \partial\Omega \rightarrow \mathbb{R}$ be such that $\partial_\tau\varphi_0 = \kappa\mathcal{H}^1\llcorner\partial\Omega - \pi\sum_{j=1}^N d_j\delta_{a_j}$ and $e^{i\varphi_0} \cdot \nu = 0$ in $\partial\Omega \setminus \{a_1, \dots, a_N\}$. Let φ_* be the harmonic extension of φ_0 to Ω . For any $\varepsilon > 0$, we consider $(\varphi_\varepsilon)_{\varepsilon\downarrow 0}$ as in Theorem 2.8 and we set $v_\varepsilon = e^{i\varphi_\varepsilon}$. Then, for every $\varepsilon > 0$, $v_\varepsilon \in H^1(\Omega, \mathbb{S}^1)$ and $\mathcal{J}(v_\varepsilon) = -\partial_\tau\varphi_\varepsilon$ as measure on $\partial\Omega$. Since, for every $\varepsilon > 0$,

$$|v_\varepsilon - e^{i\varphi_*}| = |e^{i\varphi_\varepsilon} - e^{i\varphi_*}| \leq |\varphi_\varepsilon - \varphi_*|,$$

it follows from Theorem 2.8 that $(v_\varepsilon)_{\varepsilon\downarrow 0}$ converges strongly to φ_* in $L^p(\Omega)$ and in $L^p(\partial\Omega)$ for every $p \in [1, +\infty)$, and that $(\mathcal{J}(v_\varepsilon))_{\varepsilon\downarrow 0}$ converges to

$$-\partial_\tau\varphi_0 = -\kappa\mathcal{H}^1\llcorner\partial\Omega + \pi\sum_{j=1}^N d_j\delta_{a_j}$$

in $(\text{Lip}(\Omega))^*$. Finally, the expected upper bounds at first and at second order for $E_{\varepsilon,\eta}^\delta(v_\varepsilon)$ follow from (32) for the first order, from (33) for the second order, combined with the equality $E_{\varepsilon,\eta}^\delta(v_\varepsilon) = \mathcal{G}_\varepsilon^\delta(\varphi_\varepsilon)$. \square

2.5 Minimisers of the renormalised energy

For proving Corollary 1.9, we first need to prove that $E_{\varepsilon,\eta}^\delta$ admits minimisers in $H^1(\Omega, \mathbb{R}^2)$, and secondly that the renormalised energy in (13) admits minimisers corresponding to two boundary vortices of multiplicities 1, i.e. $N = 2$ and $d_1 = d_2 = 1$ (see Corollary 1.7).

Lemma 2.9. *Let $\delta \in \mathbb{R}^2$ and $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain. Assume $\varepsilon \rightarrow 0$ and $\eta = \eta(\varepsilon) \rightarrow 0$ in the regime (9). There exists a minimiser of $E_{\varepsilon,\eta}^\delta$ over the set $H^1(\Omega, \mathbb{R}^2)$ for small $\varepsilon, \eta > 0$.*

Proof. We use the direct method in the calculus of variations. First, we show that $E_{\varepsilon,\eta}^\delta$ is coercive in H^1 . Indeed, if $v \in H^1(\Omega, \mathbb{R}^2)$, then choosing η small such that $4|\delta|^2 \leq \frac{1}{2\eta^2}$, we estimate as in the proof of Lemma 2.1:

$$\left| 2 \int_\Omega \delta \cdot \nabla v \wedge v \, dx \right| \leq \frac{1}{2} \int_\Omega |\nabla v|^2 \, dx + 4|\delta|^2 \int_\Omega (1 - |v|^2)^2 \, dx + 4|\delta|^2 |\Omega|,$$

and we conclude

$$E_{\varepsilon,\eta}^\delta(v) \geq \frac{1}{2} \int_\Omega \left(|\nabla v|^2 + \frac{1}{\eta^2} (1 - |v|^2)^2 \right) \, dx - C \gtrsim \|v\|_{H^1}^2 - 1.$$

Second, one easily checks that $E_{\varepsilon,\eta}^\delta$ is lower semicontinuous in the weak H^1 topology for two-dimensional domains Ω . Therefore, we conclude to the existence of minimisers of $E_{\varepsilon,\eta}^\delta$ for small ε and η . \square

The following statement extends [26, Proposition 20] by giving a formula for the renormalised energy $W_\Omega^\delta(\{(a_j, d_j)\})$ defined in (13). In particular, it shows that the limit in (13) exists. The formula for the renormalised energy is computed using the solution of a Neumann problem as in [6].

Proposition 2.10. *Let $\delta \in \mathbb{R}^2$, $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected $C^{1,1}$ domain, and κ be the curvature of $\partial\Omega$. Let $\{a_j \in \partial\Omega\}_{1 \leq j \leq N}$ be $N \geq 2$ distinct points and $d_j \in \{\pm 1\}$ be the corresponding multiplicities, for $j \in \{1, \dots, N\}$, that satisfy $\sum_{j=1}^N d_j = 2$. Then the limit in (13) exists and the renormalised energy of $\{(a_j, d_j)\}$ satisfies*

$$\begin{aligned} W_{\Omega}^{\delta}(\{(a_j, d_j)\}) &= -2\pi \sum_{1 \leq j < k \leq N} d_j d_k \log |a_j - a_k| \\ &\quad - \int_{\partial\Omega} \psi(\kappa + 2\delta^{\perp} \cdot \nu) d\mathcal{H}^1 + \pi \sum_{j=1}^N d_j R(a_j), \end{aligned} \quad (42)$$

where ν is the outer unit normal vector on $\partial\Omega$ and ψ denotes the unique solution (up to an additive constant) in $W^{1,q}(\Omega)$, for every $q \in [1, 2)$, of the inhomogeneous Neumann problem

$$\begin{cases} \Delta\psi = 0 & \text{in } \Omega, \\ \frac{\partial\psi}{\partial\nu} = -\kappa + \pi \sum_{j=1}^N d_j \delta_{a_j} & \text{on } \partial\Omega, \end{cases} \quad (43)$$

and R is the harmonic function given by

$$R(z) = \psi(z) + \sum_{j=1}^N d_j \log |z - a_j|, \quad (44)$$

for every $z \in \Omega$. Moreover, we have $R \in C^{0,\alpha}(\bar{\Omega}) \cap W^{s,p}(\Omega)$ for every $\alpha \in (0, 1)$, $p \in [1, +\infty)$ and $s \in [1, 1 + \frac{1}{p})$.

Proof. By the definition of $W_{\Omega}^{\delta}(\{(a_j, d_j)\})$ in (13), we have

$$W_{\Omega}^{\delta}(\{(a_j, d_j)\}) = W_{\Omega}^0(\{(a_j, d_j)\}) - \lim_{r \rightarrow 0} \int_{\Omega^r} 2\delta \cdot \nabla\varphi_* \, dx,$$

where $\Omega^r = \Omega \setminus \bigcup_{j=1}^N B_r(a_j)$ for $r > 0$, and φ_* is the harmonic extension of φ_0 (given in Definition 1.3) to Ω . First, by [26, Proposition 20], we deduce all the stated properties for ψ and R ; moreover

$$W_{\Omega}^0(\{(a_j, d_j)\}) = -2\pi \sum_{1 \leq j < k \leq N} d_j d_k \log |a_j - a_k| - \int_{\partial\Omega} \psi \kappa \, d\mathcal{H}^1 + \pi \sum_{j=1}^N d_j R(a_j). \quad (45)$$

Moreover, any solution ψ of (43) is clearly a harmonic conjugate of φ_* , in particular $\nabla\varphi_* = -\nabla^{\perp}\psi$, and $\nabla\varphi_* \in L^q(\Omega)$, for every $q \in [1, 2)$. It follows by dominated convergence theorem that

$$\lim_{r \rightarrow 0} \int_{\Omega^r} 2\delta \cdot \nabla\varphi_* \, dx = 2 \int_{\Omega} \delta \cdot \nabla\varphi_* \, dx = -2 \int_{\Omega} \delta \cdot \nabla^{\perp}\psi \, dx = 2 \int_{\partial\Omega} \psi(\delta^{\perp} \cdot \nu) d\mathcal{H}^1$$

and (42) follows. \square

We now prove Theorem 1.6:

Proof of Theorem 1.6. By Proposition 2.10, we have

$$W_{\Omega}^{\delta}(\{(a_j, d_j)\}) = W_{\Omega}^0(\{(a_j, d_j)\}) - 2 \int_{\partial\Omega} \psi(\delta^{\perp} \cdot \nu) d\mathcal{H}^1, \quad (46)$$

for any bounded, simply connected $C^{1,1}$ domain $\Omega \subset \mathbb{R}^2$, with boundary curvature κ .

(i) We assume that $\Omega = B_1$. By [26, Theorem 6], we have

$$W_{B_1}^0(\{(a_j, d_j)\}) = -2\pi \sum_{1 \leq j < k \leq N} d_j d_k \log |a_j - a_k|, \quad (47)$$

and, for every $z \in B_1$, the solution ψ in (43) is given (up to an additive constant) by

$$\psi(z) = -\sum_{j=1}^N d_j \log |z - a_j|. \quad (48)$$

By Green's formula,

$$\int_{\partial B_1} \psi \delta^\perp \cdot \nu \, d\mathcal{H}^1 = \int_{B_1} \operatorname{div}(\psi(z) \delta^\perp) \, dz = -\sum_{j=1}^N d_j \int_{B_1} \operatorname{div}(\delta^\perp \log |z - a_j|) \, dz.$$

For any $j \in \{1, \dots, N\}$ and $z \in B_1$,

$$\operatorname{div}(\delta^\perp \log |z - a_j|) = \delta^\perp \cdot \frac{z - a_j}{|z - a_j|^2} = \Re \left(\frac{\overline{z - a_j}}{|z - a_j|^2} \delta^\perp \right) = \Re \left(\frac{1}{z - a_j} \delta^\perp \right)$$

with the identification $\delta^\perp = \begin{pmatrix} -\delta_2 \\ \delta_1 \end{pmatrix} = -\delta_2 + i\delta_1$ and $\Re(z)$ is the real part of $z \in \mathbb{C}$, so that

$$\int_{B_1} \operatorname{div}(\delta^\perp \log |z - a_j|) \, dz = \Re \left(\delta^\perp \int_{B_1} \frac{1}{z - a_j} \, dz \right) = -\pi \Re(\delta^\perp \overline{a_j}) = \pi \delta \cdot a_j^\perp;$$

above we used that for $a \in \partial B_1$, the function $f : B_1 \rightarrow \mathbb{C}$, $f(z) = \frac{1}{z-a}$ for every $z \in B_1$ is holomorphic and integrable in B_1 so that the mean value theorem yields $f(0) = \frac{1}{\pi} \int_{B_1} f(z) \, dz$. We conclude

$$\int_{\partial B_1} \psi \delta^\perp \cdot \nu \, d\mathcal{H}^1 = -\pi \sum_{j=1}^N d_j \delta \cdot a_j^\perp. \quad (49)$$

(ii) By [26, Theorem 6], we have

$$\begin{aligned} & W_\Omega^0(\{(a_j, d_j)\}) \\ &= -2\pi \sum_{1 \leq j < k \leq N} d_j d_k \log |\Psi(a_j) - \Psi(a_k)| + \pi \sum_{j=1}^N (d_j - 1) \log |\partial_z \Psi(a_j)| \\ & \quad + \int_{\partial \Omega} \kappa(w) \left(\sum_{j=1}^N d_j \log |\Psi(w) - \Psi(a_j)| - \log |\partial_z \Psi(w)| \right) \, d\mathcal{H}^1(w), \end{aligned}$$

and, for every $w \in \Omega$, the solution ψ in (43) is given (up to an additive constant) by

$$\psi(w) = -\sum_{j=1}^N d_j \log |\Psi(w) - \Psi(a_j)| + \log |\partial_z \Psi(w)|.$$

The expected identity (16) is a direct consequence of (46) and the above identities. \square

We now prove Corollary 1.7, that gives the existence of a minimising pair (a_1^*, a_2^*) for the renormalised energy when $N = 2$ and $d_1 = d_2 = 1$.

Proof of Corollary 1.7. Let $\Phi: \overline{B_1} \rightarrow \overline{\Omega}$ be a C^1 conformal diffeomorphism with inverse $\Psi = \Phi^{-1}$. Let $\mathfrak{D} = (\partial\Omega \times \partial\Omega) \setminus \{(a, a) : a \in \partial\Omega\}$. By Theorem 1.6, for every $(a_1, a_2) \in \mathfrak{D}$,

$$W_{\Omega}^{\delta}(\{(a_1, 1), (a_2, 1)\}) = -2\pi \log |\Psi(a_1) - \Psi(a_2)| + F(a_1, a_2)$$

where for every $(a_1, a_2) \in \partial\Omega \times \partial\Omega$,

$$F(a_1, a_2) = \int_{\partial\Omega} (\kappa(w) + 2\delta^{\perp} \cdot \nu(w)) (\log |\Psi(w) - \Psi(a_1)| + \log |\Psi(w) - \Psi(a_2)| - \log |\partial_z \Psi(w)|) d\mathcal{H}^1(w).$$

Setting $b_1 = \Psi(a_1) \in \partial B_1$ and $b_2 = \Psi(a_2) \in \partial B_1$, we get, after changing variables,

$$F(a_1, a_2) = \int_{\partial B_1} (\kappa(\Phi(z)) + 2\delta^{\perp} \cdot \nu(\Phi(z))) (\log |z - b_1| + \log |z - b_2| + \log |\partial_z \Phi(z)|) |\partial_z \Phi(z)| d\mathcal{H}^1(z).$$

Thus F is bounded on $\partial\Omega \times \partial\Omega$, since $\kappa + 2\delta^{\perp} \cdot \nu \in L^{\infty}(\partial\Omega)$ and the functions $z \mapsto \log |z - b_1|$ and $z \mapsto \log |z - b_2|$ are in $L^1(\partial B_1)$. Moreover, the function $(a_1, a_2) \in \mathfrak{D} \mapsto W_{\Omega}^{\delta}(\{(a_1, 1), (a_2, 1)\})$ is continuous on \mathfrak{D} . Let $(a_1^{(n)}, a_2^{(n)}) \subset \mathfrak{D}$ be a minimising sequence for $W_{\Omega}^{\delta}(\{(\cdot, 1), (\cdot, 1)\})$, i.e.

$$\lim_{n \rightarrow +\infty} W_{\Omega}^{\delta}(\{(a_1^{(n)}, 1), (a_2^{(n)}, 1)\}) = \inf_{\mathfrak{D}} W_{\Omega}^{\delta}(\{(\cdot, 1), (\cdot, 1)\}).$$

Note that such a sequence exists because F is bounded in \mathfrak{D} . As $\partial\Omega \times \partial\Omega$ is compact, we can assume (up to a subsequence) that $(a_1^{(n)}, a_2^{(n)})$ converges to some $(a_1^*, a_2^*) \in \partial\Omega \times \partial\Omega$. Since $(W_{\Omega}^{\delta}(\{(a_1^{(n)}, 1), (a_2^{(n)}, 1)\}))$ is bounded, then using the boundedness of F , we deduce that the sequence $(\log |\Psi(a_1^{(n)}) - \Psi(a_2^{(n)})|)$ is bounded. By taking the limits as $n \rightarrow +\infty$, we deduce that $\Psi(a_1^*) \neq \Psi(a_2^*)$, and since Ψ is injective, $a_1^* \neq a_2^*$, i.e. $(a_1^*, a_2^*) \in \mathfrak{D}$. Finally, by continuity of $W_{\Omega}^{\delta}(\{(\cdot, 1), (\cdot, 1)\})$ over \mathfrak{D} , we get

$$W_{\Omega}^{\delta}(\{(a_1^*, 1), (a_2^*, 1)\}) = \inf_{\mathfrak{D}} W_{\Omega}^{\delta}(\{(\cdot, 1), (\cdot, 1)\}).$$

□

We prove Theorem 1.8, that gives the configuration of the vortices that minimises the renormalised energy in the unit disk B_1 in \mathbb{R}^2 .

Proof of Theorem 1.8. By Theorem 1.6 and Corollary 1.7, there exists a pair of distinct points $(a_1^*, a_2^*) \in \partial B_1 \times \partial B_1$ that minimises the renormalised energy

$$W_{B_1}^{\delta}(\{(a_1, 1), (a_2, 1)\}) = -2\pi \log |a_1 - a_2| + 2\pi\delta \cdot (a_1^{\perp} + a_2^{\perp})$$

defined for $(a_1, a_2) \in \partial B_1 \times \partial B_1$ such that $a_1 \neq a_2$. We write $a_1^* = e^{i\varphi_1}$ and $a_2^* = e^{i\varphi_2}$ for some $\varphi_1, \varphi_2 \in \mathbb{R}$. Since $a_1^* \neq a_2^*$, we have $\varphi_1 - \varphi_2 \notin 2\pi\mathbb{Z}$. Computing the renormalised energy above in terms of φ_1, φ_2 and $\delta = (\delta_1, \delta_2)$, we deduce that

$$W_{B_1}^{\delta}(\{(a_1^*, 1), (a_2^*, 1)\}) = 2\pi f(\varphi_1, \varphi_2)$$

where

$$f(\varphi_1, \varphi_2) = -\frac{1}{2} \log 2 - \frac{1}{2} \log(1 - \cos(\varphi_1 - \varphi_2)) - \delta_1(\sin \varphi_1 + \sin \varphi_2) + \delta_2(\cos \varphi_1 + \cos \varphi_2)$$

is $C^\infty(\mathbb{R}^2 \setminus S)$ where $S = \{(\varphi_1, \varphi_2) \in \mathbb{R}^2 : \varphi_1 - \varphi_2 \in 2\pi\mathbb{Z}\}$. Hence, the problem of minimising the renormalised energy turns in minimising f over $\mathbb{R}^2 \setminus S$.

Step 1 : We compute the critical points of f .
For every $(\varphi_1, \varphi_2) \in \mathbb{R}^2 \setminus S$,

$$\nabla f(\varphi_1, \varphi_2) = \left(-\frac{\sin(\varphi_1 - \varphi_2)}{2(1 - \cos(\varphi_1 - \varphi_2))} - a_1^* \cdot \delta \right) = 0 \iff \begin{cases} \frac{\sin(\varphi_1 - \varphi_2)}{1 - \cos(\varphi_1 - \varphi_2)} = -2a_1^* \cdot \delta \\ (a_1^* + a_2^*) \cdot \delta = 0 \end{cases}. \quad (50)$$

Moreover,

$$\begin{aligned} (a_1^* + a_2^*) \cdot \delta = 0 &\iff \delta_1(\cos \varphi_1 + \cos \varphi_2) + \delta_2(\sin \varphi_1 + \sin \varphi_2) = 0 \\ &\iff 2 \cos\left(\frac{\varphi_1 - \varphi_2}{2}\right) \left(\delta_1 \cos \frac{\varphi_1 + \varphi_2}{2} + \delta_2 \sin \frac{\varphi_1 + \varphi_2}{2} \right) = 0 \\ &\iff \cos \frac{\varphi_1 - \varphi_2}{2} = 0 \quad \text{or} \quad \delta \cdot b = 0 \quad \text{where} \quad b = e^{i(\varphi_1 + \varphi_2)/2}. \end{aligned}$$

Case 1: $\cos \frac{\varphi_1 - \varphi_2}{2} = 0$. Then we have $\varphi_1 - \varphi_2 = \pi \pmod{2\pi}$, thus a_1^* and a_2^* are diametrically opposed (i.e. $a_2^* = -a_1^*$). By the first equation in (50) and since $\delta = |\delta| e^{i\theta}$, we deduce that

$$a_1^* \cdot \delta = 0 \iff \cos(\varphi_1 - \theta) = 0 \iff \varphi_1 = \theta + \frac{\pi}{2} \pmod{\pi}.$$

Hence, $a_1^* = e^{i\varphi_1} = \pm i e^{i\theta} = \pm \frac{1}{|\delta|} \delta^\perp$.

Case 2: $\delta \cdot b = 0$ where $b = e^{i(\varphi_1 + \varphi_2)/2}$ and $\varphi_1 - \theta \neq \frac{\pi}{2} \pmod{\pi}$. Since $\delta = |\delta| e^{i\theta}$, we have

$$\delta \cdot b = 0 \iff \delta \perp b \iff \frac{\varphi_1 + \varphi_2}{2} = \theta + \frac{\pi}{2} \pmod{\pi} \iff \varphi_1 + \varphi_2 = 2\theta + \pi \pmod{2\pi},$$

thus a_1^* and a_2^* are symmetric with respect to δ^\perp . By the first equation in (50), we have

$$\begin{aligned} -2|\delta| a_1^* \cdot e^{i\theta} = \frac{\sin(\varphi_1 - (2\theta + \pi - \varphi_1))}{1 - \cos(\varphi_1 - (2\theta + \pi - \varphi_1))} &\iff -2|\delta| \cos(\varphi_1 - \theta) = \frac{-\sin(\varphi_1 - \theta)}{\cos(\varphi_1 - \theta)} \\ &\iff 2|\delta| \sin^2(\varphi_1 - \theta) + \sin(\varphi_1 - \theta) - 2|\delta| = 0 \\ &\iff \begin{cases} 2|\delta| X^2 + X - 2|\delta| = 0 \\ X = \sin(\varphi_1 - \theta) \end{cases}. \end{aligned}$$

As $X = \sin(\varphi_1 - \theta) \geq -1$, we deduce that

$$\sin(\varphi_1 - \theta) = \sqrt{1 + \frac{1}{16|\delta|^2}} - \frac{1}{4|\delta|} \in [-1, 1],$$

thus $\varphi_1 = \theta + \theta_\delta \pmod{2\pi}$ or $\varphi_1 = \theta + \pi - \theta_\delta \pmod{2\pi}$ where $\theta_\delta = \arcsin\left(\sqrt{1 + \frac{1}{16|\delta|^2}} - \frac{1}{4|\delta|}\right)$. We obtain that $\varphi_2 = 2\theta + \pi - \varphi_1 = \theta + \pi - \theta_\delta \pmod{2\pi}$ or $\varphi_2 = \theta + \theta_\delta \pmod{2\pi}$. Up to interchanging φ_1 and φ_2 , we will assume in the following that $\varphi_1 = \theta + \theta_\delta \pmod{2\pi}$.

Step 2 : We study the nature of the critical points of f .
For every $(\varphi_1, \varphi_2) \in \mathbb{R}^2 \setminus S$, the Hessian matrix of f is

$$\text{Hess}(f)(\varphi_1, \varphi_2) = \begin{pmatrix} \frac{1}{2(1 - \cos(\varphi_1 - \varphi_2))} + a_1^* \cdot \delta^\perp & \frac{-1}{2(1 - \cos(\varphi_1 - \varphi_2))} \\ \frac{-1}{2(1 - \cos(\varphi_1 - \varphi_2))} & \frac{1}{2(1 - \cos(\varphi_1 - \varphi_2))} + a_2^* \cdot \delta^\perp \end{pmatrix}$$

and its determinant is

$$h(\varphi_1, \varphi_2) = \frac{1}{2(1 - \cos(\varphi_1 - \varphi_2))} (a_1^* + a_2^*) \cdot \delta^\perp + (a_1^* \cdot \delta^\perp)(a_2^* \cdot \delta^\perp).$$

Case 1: $\cos \frac{\varphi_1 - \varphi_2}{2} = 0$. Here, we have diametrically opposed points $a_2^* = -a_1^*$ with $a_1^* = \pm \frac{1}{|\delta|} \delta^\perp$, and

$$h(\varphi_1, \varphi_2) = (a_1^* \cdot \delta^\perp)(a_2^* \cdot \delta^\perp) = -(a_1^* \cdot \delta^\perp)^2 = -|\delta|^2 < 0.$$

Hence, (φ_1, φ_2) cannot be a minimiser for our renormalised energy.¹⁰

Case 2: $\delta \cdot b = 0$ and $\cos \frac{\varphi_1 - \varphi_2}{2} \neq 0$. Here, a_1^* and a_2^* are symmetric with respect to δ^\perp , with (by our convention) $\varphi_1 = \theta + \theta_\delta \pmod{2\pi}$ and $\varphi_2 = \theta + \pi - \theta_\delta \pmod{2\pi}$ where θ_δ is given above. Then $h(\varphi_1, \varphi_2) > 0$, because

$$\begin{aligned} a_1^* \cdot \delta^\perp &= e^{i\varphi_1} \cdot |\delta| e^{i(\theta + \pi/2)} = |\delta| \cos\left(\varphi_1 - \theta - \frac{\pi}{2}\right) \\ &= |\delta| \sin(\varphi_1 - \theta) = |\delta| \left(\sqrt{1 + \frac{1}{16|\delta|^2}} - \frac{1}{4|\delta|} \right) > 0, \end{aligned}$$

and $a_2^* \cdot \delta^\perp = a_1^* \cdot \delta^\perp > 0$, as a_2^* and a_1^* are symmetric with respect to δ^\perp . Moreover, the trace of $Hess(f)(\varphi_1, \varphi_2)$ is positive, hence $(\varphi_1, \varphi_2) = (\theta + \theta_\delta, \theta + \pi - \theta_\delta)$ minimizes f . \square

To conclude this section, we prove Corollary 1.9.

Proof of Corollary 1.9. By Lemma 2.9, there exists a minimiser v_ε of $E_{\varepsilon, \eta}^\delta$ on $H^1(\Omega, \mathbb{R}^2)$ for small $\varepsilon > 0$. By Corollary 1.7, there exist two points $a_1^* \neq a_2^* \in \partial\Omega$ such that

$$W_\Omega^\delta(\{(a_1^*, 1), (a_2^*, 1)\}) = \min \{W_\Omega^\delta(\{(\tilde{a}_1, 1), (\tilde{a}_2, 1)\}) : \tilde{a}_1 \neq \tilde{a}_2 \in \partial\Omega\}. \quad (51)$$

By Theorem 1.5 applied to $\{(a_1^*, 1), (a_2^*, 1)\}$, the minimisers v_ε must satisfy

$$E_{\varepsilon, \eta}^\delta(v_\varepsilon) \leq 2\pi |\log \varepsilon| + W_\Omega^\delta(\{(a_1^*, 1), (a_2^*, 1)\}) + 2\gamma_0 + o_\varepsilon(1). \quad (52)$$

By Theorem 1.2(i), for a subsequence, $(\mathcal{J}(v_\varepsilon))_{\varepsilon > 0}$ converges as in (11) to

$$J = -\kappa \mathcal{H}^1 \llcorner \partial\Omega + \pi \sum_{j=1}^N d_j \delta_{a_j}$$

for $N \geq 2$ distinct boundary points $a_1, \dots, a_N \in \partial\Omega$, with $d_1, \dots, d_N \in \mathbb{Z} \setminus \{0\}$ such that $\sum_{j=1}^N d_j = 2$. By Theorem 1.2(ii), we also have

$$\liminf_{\varepsilon \rightarrow 0} \frac{1}{|\log \varepsilon|} E_{\varepsilon, \eta}^\delta(v_\varepsilon) \geq \pi \sum_{j=1}^N |d_j|. \quad (53)$$

Combining (52) and (53), we get $\sum_{j=1}^N |d_j| \leq 2$, hence $\sum_{j=1}^N (|d_j| - d_j) \leq 0$ so that $d_j = |d_j| > 0$ for $1 \leq j \leq N$. In particular, (14) holds thanks to (52) and applying Theorem 1.4(i), we deduce that $N = 2$, $d_1 = d_2 = 1$ and

$$E_{\varepsilon, \eta}^\delta(v_\varepsilon) \geq 2\pi |\log \varepsilon| + W_\Omega^\delta(\{(a_1, 1), (a_2, 1)\}) + 2\gamma_0 + o_\varepsilon(1). \quad (54)$$

¹⁰Note that the diametrically opposed configuration $(a_1^*, -a_1^*)$ with $a_1^* = \pm \frac{1}{|\delta|} \delta^\perp \in \partial B_1$ is an unstable critical point of the renormalised energy when $\delta \neq 0$.

Combining (52) and (54), and letting $\varepsilon \rightarrow 0$, we get

$$W_{\Omega}^{\delta}(\{(a_1, 1), (a_2, 1)\}) \leq W_{\Omega}^{\delta}(\{(a_1^*, 1), (a_2^*, 1)\}),$$

so by definition of $W_{\Omega}^{\delta}(\{(a_1^*, 1), (a_2^*, 1)\})$, we deduce that (a_1, a_2) is also a minimiser in (51). It follows then from (52) and (54) that

$$E_{\varepsilon, \eta}^{\delta}(v_{\varepsilon}) = 2\pi |\log \varepsilon| + W_{\Omega}^{\delta}(\{(a_1, 1), (a_2, 1)\}) + 2\gamma_0 + o_{\varepsilon}(1). \quad (55)$$

By (39), for a subsequence, (v_{ε}) converges weakly in $W^{1,q}(\Omega, \mathbb{R}^2)$ for every $q \in [1, 2)$, and strongly in $L^p(\Omega, \mathbb{R}^2)$ for every $p \in [1, +\infty)$, to $e^{i\widehat{\varphi}_0}$, where $\widehat{\varphi}_0$ is an extension to Ω of a function $\varphi_0 \in BV(\partial\Omega, \pi\mathbb{Z})$ that satisfies

$$\partial_{\tau}\varphi_0 = \kappa\mathcal{H}^1 \llcorner \partial\Omega - \pi(\boldsymbol{\delta}_{a_1} + \boldsymbol{\delta}_{a_2}) \quad \text{as measure on } \partial\Omega$$

and $e^{i\varphi_0} \cdot \nu = 0$ in $\partial\Omega \setminus \{a_1, a_2\}$.

It remains to prove that $\widehat{\varphi}_0$ is harmonic in Ω . Let $r > 0$ be small. As for a subsequence, $(\nabla v_{\varepsilon})_{\varepsilon}$ converges weakly to $\nabla(e^{i\widehat{\varphi}_0})$ in $L^{3/2}(\Omega \setminus (B_r(a_1) \cup B_r(a_2)))$, we have for every test function $\zeta \in C_c^{\infty}(\Omega \setminus (B_r(a_1) \cup B_r(a_2)), \mathbb{R}^2)$,

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} \nabla v_{\varepsilon} \cdot \zeta \, dx = \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} \nabla(e^{i\widehat{\varphi}_0}) \cdot \zeta \, dx$$

yielding

$$\int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} |\nabla \widehat{\varphi}_0|^2 \, dx \leq \liminf_{\varepsilon \rightarrow 0} \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} |\nabla v_{\varepsilon}|^2 \, dx,$$

since $|\nabla(e^{i\widehat{\varphi}_0})| = |\nabla \widehat{\varphi}_0|$. Moreover, using that $v_{\varepsilon} \rightarrow e^{i\widehat{\varphi}_0}$ in $L^3(\Omega)$ for a subsequence $\varepsilon \rightarrow 0$, we get

$$\int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} -\delta \cdot \nabla \widehat{\varphi}_0 \, dx = \lim_{\varepsilon \rightarrow 0} \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} \delta \cdot \nabla v_{\varepsilon} \wedge v_{\varepsilon} \, dx,$$

since $\nabla(e^{i\widehat{\varphi}_0}) \wedge e^{i\widehat{\varphi}_0} = -\nabla \widehat{\varphi}_0$. From these observations, we deduce

$$\begin{aligned} & \liminf_{r \rightarrow 0} \left(\int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} (|\nabla \widehat{\varphi}_0|^2 - 2\delta \cdot \nabla \widehat{\varphi}_0) \, dx - 2\pi \log \frac{1}{r} \right) \\ & \leq \liminf_{r \rightarrow 0} \liminf_{\varepsilon \rightarrow 0} \left(\int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} (|\nabla v_{\varepsilon}|^2 + 2\delta \cdot \nabla v_{\varepsilon} \wedge v_{\varepsilon}) \, dx - 2\pi \log \frac{1}{r} \right). \end{aligned} \quad (56)$$

Inside the disks $\Omega \cap (B_r(a_1) \cup B_r(a_2))$, by the compactness of $(v_{\varepsilon})_{\varepsilon}$, we have as above

$$\begin{aligned} & \lim_{\varepsilon \rightarrow 0} \left| \int_{\Omega \cap (B_r(a_1) \cup B_r(a_2))} 2\delta \cdot \nabla v_{\varepsilon} \wedge v_{\varepsilon} \, dx \right| = \left| \int_{\Omega \cap (B_r(a_1) \cup B_r(a_2))} 2\delta \cdot \nabla \widehat{\varphi}_0 \, dx \right| \\ & \leq 2|\delta| |\Omega \cap (B_r(a_1) \cup B_r(a_2))|^{1/3} \|\nabla \widehat{\varphi}_0\|_{L^{3/2}(\Omega)} \leq Cr^{2/3} \end{aligned}$$

for some constant $C > 0$ independent of ε and r , and thus

$$\liminf_{r \rightarrow 0} \liminf_{\varepsilon \rightarrow 0} \int_{\Omega \cap (B_r(a_1) \cup B_r(a_2))} 2\delta \cdot \nabla v_{\varepsilon} \wedge v_{\varepsilon} \, dx = 0. \quad (57)$$

Moreover, by [25, Equation (86)], we have

$$\liminf_{r \rightarrow 0} \liminf_{\varepsilon \rightarrow 0} \left(E_{\varepsilon, \eta}^0(v_{\varepsilon}; \Omega \cap (B_r(a_1) \cup B_r(a_2))) - 2\pi \log \frac{r}{\varepsilon} - 2\gamma_0 \right) \geq 0. \quad (58)$$

Note that, for every small $r > 0$,

$$\begin{aligned}
E_{\varepsilon, \eta}^{\delta}(v_{\varepsilon}) - 2\pi |\log \varepsilon| - 2\gamma_0 &= E_{\varepsilon, \eta}^0(v_{\varepsilon}; \Omega \cap (B_r(a_1) \cup B_r(a_2))) - 2\pi \log \frac{r}{\varepsilon} - 2\gamma_0 \\
&+ \int_{\Omega \cap (B_r(a_1) \cup B_r(a_2))} 2\delta \cdot \nabla v_{\varepsilon} \wedge v_{\varepsilon} \, dx + E_{\varepsilon, \eta}^{\delta}(v_{\varepsilon}; \Omega \setminus (B_r(a_1) \cup B_r(a_2))) - 2\pi \log \frac{1}{r} \\
&\geq E_{\varepsilon, \eta}^0(v_{\varepsilon}; \Omega \cap (B_r(a_1) \cup B_r(a_2))) - 2\pi \log \frac{r}{\varepsilon} - 2\gamma_0 \\
&+ \int_{\Omega \cap (B_r(a_1) \cup B_r(a_2))} 2\delta \cdot \nabla v_{\varepsilon} \wedge v_{\varepsilon} \, dx + \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} (|\nabla v_{\varepsilon}|^2 + 2\delta \cdot \nabla v_{\varepsilon} \wedge v_{\varepsilon}) \, dx - 2\pi \log \frac{1}{r}.
\end{aligned}$$

By (55), (56), (57) and (58), it follows by passing to the limit as $\varepsilon \rightarrow 0$ and then $r \rightarrow 0$:

$$\begin{aligned}
W_{\Omega}^{\delta}(\{(a_1, 1), (a_2, 1)\}) &= \liminf_{\varepsilon \rightarrow 0} (E_{\varepsilon, \eta}^{\delta}(v_{\varepsilon}) - 2\pi |\log \varepsilon| - 2\gamma_0) \\
&\geq \liminf_{r \rightarrow 0} \left(\int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} (|\nabla \widehat{\varphi}_0|^2 - 2\delta \cdot \nabla \widehat{\varphi}_0) \, dx - 2\pi \log \frac{1}{r} \right). \tag{59}
\end{aligned}$$

Let φ_* be the harmonic extension of φ_0 to Ω . Then $\widehat{\varphi}_0 - \varphi_* \in W_0^{1,q}(\Omega)$, for every $q \in [1, 2)$. For every small $r > 0$,

$$\begin{aligned}
\int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} 2\delta \cdot (\nabla \widehat{\varphi}_0 - \nabla \varphi_*) \, dx &= \int_{\Omega} 2\delta \cdot \nabla (\widehat{\varphi}_0 - \varphi_*) \, dx \\
&- 2 \int_{\Omega \cap (B_r(a_1) \cup B_r(a_2))} 2\delta \cdot \nabla (\widehat{\varphi}_0 - \varphi_*) \, dx.
\end{aligned}$$

As

$$\left| -2 \int_{\Omega \cap (B_r(a_1) \cup B_r(a_2))} 2\delta \cdot \nabla (\widehat{\varphi}_0 - \varphi_*) \, dx \right| \leq Cr^{2/3} \|\nabla (\widehat{\varphi}_0 - \varphi_*)\|_{L^{3/2}(\Omega)} \leq Cr^{2/3},$$

for some $C > 0$ independent of r , by Green's formula and using that $\widehat{\varphi}_0 = \varphi_*$ on $\partial\Omega$, we get

$$\liminf_{r \rightarrow 0} \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} 2\delta \cdot (\nabla \widehat{\varphi}_0 - \nabla \varphi_*) \, dx = 0.$$

Combining this observation with (59) and the definition of $W_{\Omega}^{\delta}(\{(a_1, 1), (a_2, 1)\})$, we deduce that

$$\begin{aligned}
W_{\Omega}^{\delta}(\{(a_1, 1), (a_2, 1)\}) &\geq \liminf_{r \rightarrow 0} \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} (|\nabla \widehat{\varphi}_0|^2 - |\nabla \varphi_*|^2) \, dx \\
&+ \liminf_{r \rightarrow 0} \left(\int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} (|\nabla \varphi_*|^2 - 2\delta \cdot \nabla \varphi_*) \, dx - 2\pi \log \frac{1}{r} \right) \\
&= \liminf_{r \rightarrow 0} \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} (|\nabla \widehat{\varphi}_0|^2 - |\nabla \varphi_*|^2) \, dx + W_{\Omega}^{\delta}(\{(a_1, 1), (a_2, 1)\}),
\end{aligned}$$

i.e.

$$\liminf_{r \rightarrow 0} \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} (|\nabla \widehat{\varphi}_0|^2 - |\nabla \varphi_*|^2) \, dx \leq 0.$$

To deduce that $\widehat{\varphi}_0 = \varphi_*$, thus $\widehat{\varphi}_0$ is harmonic in Ω , we proceed as follows:

$$|\nabla \widehat{\varphi}_0|^2 - |\nabla \varphi_*|^2 = |\nabla (\widehat{\varphi}_0 - \varphi_*)|^2 + 2\nabla \varphi_* \cdot \nabla (\widehat{\varphi}_0 - \varphi_*).$$

Since φ_* behaves as the sum of an angular function around a_1 and a_2 and a harmonic function $h \in W^{1,p}(\Omega)$ for every $p \in [1, +\infty)$ (as in Step 3 of the proof of [25, Theorem 4.2]), then integration by parts yields for small $r > 0$:

$$\begin{aligned} \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} \nabla \varphi_* \cdot \nabla (\widehat{\varphi}_0 - \varphi_*) \, dx &= \int_{\partial(\Omega \setminus (B_r(a_1) \cup B_r(a_2)))} \partial_\nu \varphi_* (\widehat{\varphi}_0 - \varphi_*) \, d\mathcal{H}^1 \\ &= \int_{\Omega \cap \partial(B_r(a_1) \cup B_r(a_2))} \partial_\nu h (\widehat{\varphi}_0 - \varphi_*) \, d\mathcal{H}^1 = \int_{\Omega \cap (B_r(a_1) \cup B_r(a_2))} \nabla h \cdot \nabla (\widehat{\varphi}_0 - \varphi_*) \, dx. \end{aligned}$$

By Hölder's inequality, we deduce that

$$\begin{aligned} \left| \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} \nabla \varphi_* \cdot \nabla (\widehat{\varphi}_0 - \varphi_*) \, dx \right| \\ \leq \| \nabla h \|_{L^3(\Omega \cap (B_r(a_1) \cup B_r(a_2)))} \| \nabla (\widehat{\varphi}_0 - \varphi_*) \|_{L^{3/2}(\Omega)} \rightarrow 0 \quad \text{as } r \rightarrow 0. \end{aligned}$$

As a consequence,

$$0 \geq \liminf_{r \rightarrow 0} \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} (|\nabla \widehat{\varphi}_0|^2 - |\nabla \varphi_*|^2) \, dx = \liminf_{r \rightarrow 0} \int_{\Omega \setminus (B_r(a_1) \cup B_r(a_2))} |\nabla (\widehat{\varphi}_0 - \varphi_*)|^2 \, dx.$$

We deduce that $\widehat{\varphi}_0 = \varphi_* + c$ for some constant $c \in \mathbb{R}$, but $\widehat{\varphi}_0 = \varphi_*$ on $\partial\Omega$, hence $c = 0$ and $\widehat{\varphi}_0 = \varphi_*$ is harmonic in Ω . \square

3 Three-dimensional model for maps $m_h: \Omega_h \subset \mathbb{R}^3 \rightarrow \mathbb{S}^2$

3.1 Reduction from the three-dimensional model to a reduced two-dimensional model

We begin by relating the three-dimensional energy $E_h(m_h)$ given at (6) with the corresponding energy in the absence of the Dzyaloshinskii-Moriya interaction:

$$E_h^0(m_h) = \frac{1}{|\log \varepsilon|} \left(\frac{1}{h} \int_{\Omega_h} |\nabla m_h|^2 \, dx + \frac{1}{h\eta^2} \int_{\mathbb{R}^3} |\nabla u_h|^2 \, dx \right) \quad (60)$$

defined for maps $m_h: \Omega_h = \Omega \times (0, h) \rightarrow \mathbb{S}^2$ and $u_h: \mathbb{R}^3 \rightarrow \mathbb{R}$ satisfying (3). This latter energy has been studied by Ignat and Kurzke [26], hence the next lemma will allow us to use in this section the statements they obtained on $E_h^0(m_h)$.

Lemma 3.1. *Let $\Omega_h = \Omega \times (0, h)$ with $\Omega \subset \mathbb{R}^2$ a bounded, simply connected $C^{1,1}$ domain. Assume that $\frac{1}{\eta^2} |\widehat{D}| = O_h(\sqrt{|\log \varepsilon|})$ and consider a family of magnetizations $\{m_h: \Omega_h \rightarrow \mathbb{S}^2\}_{h \downarrow 0}$ that satisfies*

$$\limsup_{h \rightarrow 0} E_h(m_h) < +\infty.$$

Then $\limsup_{h \rightarrow 0} E_h^0(m_h) < +\infty$.

Proof. Using $|m_h| = 1$ and Young's inequality, we obtain

$$\begin{aligned} |\log \varepsilon| |E_h(m_h) - E_h^0(m_h)| &= \frac{1}{h\eta^2} \left| \int_{\Omega_h} \widehat{D} : \nabla m_h \wedge m_h \, dx \right| \leq \frac{1}{h\eta^2} \int_{\Omega_h} |\widehat{D}| |\nabla m_h| \, dx \\ &\leq \frac{1}{2h} \int_{\Omega_h} \left(|\nabla m_h|^2 + \left(\frac{|\widehat{D}|}{\eta^2} \right)^2 \right) \, dx \leq \frac{|\log \varepsilon|}{2} E_h^0(m_h) + \frac{|\Omega|}{2} \left(\frac{|\widehat{D}|}{\eta^2} \right)^2. \end{aligned}$$

Thus,

$$E_h^0(m_h) \leq 2E_h(m_h) + \frac{|\Omega|}{|\log \varepsilon|} \left(\frac{|\widehat{D}|}{\eta^2} \right)^2 \quad (61)$$

and the conclusion follows in the regime $\frac{|\widehat{D}|}{\eta^2} \lesssim \sqrt{|\log \varepsilon|}$. \square

As a consequence, we deduce the existence of minimisers for the energy E_h :

Corollary 3.2. *Let $\Omega_h = \Omega \times (0, h)$ with $\Omega \subset \mathbb{R}^2$ a bounded, simply connected $C^{1,1}$ domain. Then E_h admits a minimiser over the set $H^1(\Omega_h, \mathbb{S}^2)$ for every $h > 0$.*

Proof. We use the direct method in the calculus of variations. Indeed, for fixed $h > 0$, consider a minimising sequence $(m_n)_{n \in \mathbb{N}}$ of E_h in $H^1(\Omega_h, \mathbb{S}^2)$; by (61), $(m_n)_{n \in \mathbb{N}}$ is bounded in $H^1(\Omega_h, \mathbb{R}^3)$, in particular, for a subsequence, $(m_n)_n$ converges weakly in $H^1(\Omega_h, \mathbb{R}^3)$, strongly in $L^2(\Omega_h, \mathbb{R}^3)$ and a.e. in Ω_h . As the Helmholtz projection $P: m \in L^2(\Omega_h, \mathbb{R}^3) \rightarrow \nabla u \in L^2(\mathbb{R}^3, \mathbb{R}^3)$ is a linear bounded operator (via (3)), the relative compactness of $(m_n)_{n \in \mathbb{N}}$ and the lower semicontinuity of E_h in the weak $H^1(\Omega_h, \mathbb{R}^3)$ topology yield the conclusion. \square

In the next lemma, we estimate the DMI in $E_h(m_h)$ with respect to the DMI corresponding to the x_3 -average of m_h .

Lemma 3.3. *Let $\Omega_h = \Omega \times (0, h)$ with $\Omega \subset \mathbb{R}^2$ a bounded, simply connected $C^{1,1}$ domain. In the regime $\frac{1}{\eta^2} |\widehat{D}| = O_h(\sqrt{|\log \varepsilon|})$, consider a family of magnetizations $\{m_h: \Omega_h \rightarrow \mathbb{S}^2\}_{h \downarrow 0}$ that satisfies $\limsup_{h \rightarrow 0} E_h(m_h) < +\infty$. If \overline{m}_h is the x_3 -average of m_h defined in (7), then*

$$\frac{1}{|\log \varepsilon|} \frac{1}{h\eta^2} \int_{\Omega_h} \widehat{D} : \nabla m_h \wedge m_h \, dx = \frac{1}{|\log \varepsilon|} \frac{1}{\eta^2} \int_{\Omega} \widehat{D}' : \nabla' \overline{m}_h \wedge \overline{m}_h \, dx' + O(R(h)) \quad \text{as } h \rightarrow 0,$$

where $\widehat{D}' = (\widehat{D}_1, \widehat{D}_2) \in \mathbb{R}^{3 \times 2}$ and

$$R(h) = h \frac{|\widehat{D}_1| + |\widehat{D}_2|}{\eta^2} + \frac{|\widehat{D}_3|}{\eta^2} \cdot \frac{1}{\sqrt{|\log \varepsilon|}}. \quad (62)$$

Proof. Writing

$$\widehat{D} : \nabla m_h \wedge m_h = \widehat{D}' : \nabla' m_h \wedge \overline{m}_h + \widehat{D}' : \nabla' m_h \wedge (m_h - \overline{m}_h) + \widehat{D}_3 \cdot \partial_3 m_h \wedge m_h,$$

we obtain the decomposition

$$\int_{\Omega_h} \widehat{D} : \nabla m_h \wedge m_h \, dx = I_1 + I_2 + I_3$$

with

$$I_1 = \int_{\Omega_h} \widehat{D}' : \nabla' m_h \wedge \overline{m}_h \, dx = \sum_{j=1}^2 \int_{\Omega_h} \widehat{D}_j \cdot \partial_j m_h \wedge \overline{m}_h \, dx,$$

$$I_2 = \int_{\Omega_h} \widehat{D}' : \nabla' m_h \wedge (m_h - \overline{m}_h) \, dx = \sum_{j=1}^2 \int_{\Omega_h} \widehat{D}_j \cdot \partial_j m_h \wedge (m_h - \overline{m}_h) \, dx,$$

and

$$I_3 = \int_{\Omega_h} \widehat{D}_3 \cdot \partial_3 m_h \wedge m_h \, dx.$$

For calculating I_1 , by Fubini's theorem, we have

$$\begin{aligned} I_1 &= \sum_{j=1}^2 \int_{\Omega} \int_0^h \left(\widehat{D}_j \cdot \partial_j m_h(x', x_3) \wedge \overline{m}_h(x') \right) dx_3 dx' \\ &= \sum_{j=1}^2 \int_{\Omega} \widehat{D}_j \cdot \partial_j \left(\int_0^h m_h(x', x_3) dx_3 \right) \wedge \overline{m}_h(x') dx' = h \sum_{j=1}^2 \int_{\Omega} \widehat{D}_j \cdot \partial_j \overline{m}_h \wedge \overline{m}_h dx'. \end{aligned}$$

For I_2 , by the Cauchy-Schwarz and Poincaré-Wirtinger inequalities, we have

$$\begin{aligned} |I_2| &\leq \sum_{j=1}^2 \int_{\Omega_h} \left| \widehat{D}_j \cdot \partial_j m_h \wedge (m_h - \overline{m}_h) \right| dx \\ &\leq \sum_{j=1}^2 |\widehat{D}_j| \left(\int_{\Omega_h} |\partial_j m_h|^2 dx \right)^{1/2} \left(\int_{\Omega_h} |m_h - \overline{m}_h|^2 dx \right)^{1/2} \\ &\leq h \sum_{j=1}^2 |\widehat{D}_j| \left(\int_{\Omega_h} |\partial_j m_h|^2 dx \right)^{1/2} \left(\int_{\Omega_h} |\partial_3 m_h|^2 dx \right)^{1/2} \\ &\leq h \sum_{j=1}^2 |\widehat{D}_j| \int_{\Omega_h} |\nabla m_h|^2 dx \leq h^2 (|\widehat{D}_1| + |\widehat{D}_2|) |\log \varepsilon| E_h^0(m_h). \end{aligned}$$

Finally, for I_3 , since $|m_h| = 1$, the Cauchy-Schwarz inequality yields

$$|I_3| \leq \int_{\Omega_h} |\widehat{D}_3 \cdot \partial_3 m_h \wedge m_h| dx \leq \sqrt{|\Omega|} |\widehat{D}_3| \sqrt{h} \left(\int_{\Omega_h} |\partial_3 m_h|^2 dx \right)^{1/2} \leq \sqrt{|\Omega|} |\widehat{D}_3| h \sqrt{|\log \varepsilon| E_h^0(m_h)},$$

where E_h^0 is defined in (60). Combining the above estimates, we obtain

$$\begin{aligned} &\frac{1}{|\log \varepsilon|} \frac{1}{h\eta^2} \left| \int_{\Omega_h} \widehat{D} : \nabla m_h \wedge m_h dx - h \int_{\Omega} \widehat{D}' : \nabla' \overline{m}_h \wedge \overline{m}_h dx' \right| \\ &\leq h \frac{|\widehat{D}_1| + |\widehat{D}_2|}{\eta^2} E_h^0(m_h) + \sqrt{|\Omega|} \frac{|\widehat{D}_3|}{\eta^2} \sqrt{\frac{E_h^0(m_h)}{|\log \varepsilon|}}. \end{aligned}$$

Since $\limsup_{h \rightarrow 0} E_h^0(m_h) < +\infty$ by Lemma 3.1, the conclusion follows. \square

We have now all the ingredients to prove Theorem 1.1.

Proof of Theorem 1.1. As $\frac{1}{\eta^2} |\widehat{D}| = O_h(\sqrt{|\log \varepsilon|})$, by Lemma 3.3, we have

$$\begin{aligned} E_h(m_h) &= E_h^0(m_h) + \frac{1}{|\log \varepsilon|} \frac{1}{h\eta^2} \int_{\Omega_h} \widehat{D} : \nabla m_h \wedge m_h dx \\ &= E_h^0(m_h) + \frac{1}{|\log \varepsilon|} \frac{1}{\eta^2} \int_{\Omega} \widehat{D}' : \nabla' \overline{m}_h \wedge \overline{m}_h dx' - O(R(h)) \end{aligned}$$

where $E_h^0(m_h)$ is given in (60). By (62), we have that $O(R(h)) = o_h(1)$ in the regime (4) and $O(R(h)) = o_h(1/|\log \varepsilon|)$ in the regime (5) because $h \ll \frac{1}{|\log h|} \ll \varepsilon \ll \frac{1}{|\log \varepsilon|}$ in the regime (4).

Step 1: Estimating $E_h^0(m_h)$. We denote by $\overline{m}_h = (\overline{m}'_h, \overline{m}_{h,3})$ the x_3 -average of m_h defined in (7). By [26, Theorem 1 and Inequality (1.11)], we have in the absence of DMI:

$$E_h^0(m_h) \geq \frac{1}{|\log \varepsilon|} \left(E_{\varepsilon, \eta}^0(\overline{m}'_h) + \int_{\Omega} |\nabla' \overline{m}_{h,3}|^2 dx' \right) - o_h(R_1(h))$$

where $R_1(h) = 1$ in the regime (4) and $R_1(h) = 1/|\log \varepsilon|$ in the regime (5). Moreover, if m_h is independent of x_3 (i.e., $m_h = \overline{m}_h$) and $m_{h,3} = 0$ (i.e., $m_h = (m'_h, 0): \Omega \rightarrow \mathbb{S}^1 \times \{0\}$), then the above inequality becomes equality.

Step 2: Estimating the DMI term. Since $\limsup_{h \rightarrow 0} E_h^0(m_h) < +\infty$ by Lemma 3.1, Step 1 implies that $\int_{\Omega} |\nabla' \overline{m}_h|^2 dx' = O(|\log \varepsilon|)$. We claim that

$$\frac{1}{|\log \varepsilon|} \int_{\Omega} \left(\frac{1}{\eta^2} \widehat{D}' : \nabla' \overline{m}_h \wedge \overline{m}_h - 2\delta \cdot \nabla' \overline{m}'_h \wedge \overline{m}'_h \right) dx = o_h(R_1(h)). \quad (63)$$

As $|\overline{m}_h| \leq 1$, the term involving \widehat{D}_{11} is estimated using the Cauchy-Schwarz inequality:

$$\left| \int_{\Omega} \frac{\widehat{D}_{11}}{\eta^2} (\overline{m}_{h,3} \partial_1 \overline{m}_{h,2} - \overline{m}_{h,2} \partial_1 \overline{m}_{h,3}) dx' \right| \leq \frac{|\widehat{D}_{11}|}{\eta^2} \int_{\Omega} |\nabla' \overline{m}_h| dx' = \frac{|\widehat{D}_{11}|}{\eta^2} O(\sqrt{|\log \varepsilon|}).$$

Similar estimates hold for the DMI terms involving \widehat{D}_{12} , \widehat{D}_{21} and \widehat{D}_{22} . For the term involving \widehat{D}_{k3} with $k = 1, 2$, we have by the Cauchy-Schwarz inequality:

$$\left| \int_{\Omega} \left(\frac{\widehat{D}_{k3}}{\eta^2} - 2\delta_k \right) \partial_k \overline{m}'_h \wedge \overline{m}'_h dx' \right| \leq \left| \frac{\widehat{D}_{k3}}{\eta^2} - 2\delta_k \right| \int_{\Omega} |\nabla' \overline{m}_h| dx' = \left| \frac{\widehat{D}_{k3}}{\eta^2} - 2\delta_k \right| O(\sqrt{|\log \varepsilon|}).$$

The choice of our regimes (4) and (5) yields the claim (63). \square

3.2 Gamma-convergence of the three-dimensional energy

In this section, we prove the Γ -convergence for $E_h(m_h)$, i.e. Theorems 1.10, 1.11, 1.12 and Corollary 1.13. Recall that the regime (4) (and also (5)) implies regime (9), see Footnote 3.

Proof of Theorem 1.10. Let $\{m_h: \Omega_h \rightarrow \mathbb{S}^2\}_{h \downarrow 0}$ be a family of magnetizations such that $\limsup_{h \rightarrow 0} E_h(m_h) < \infty$. Denoting $\overline{m}_h = (\overline{m}'_h, \overline{m}_{h,3}): \Omega \rightarrow \mathbb{R}^3$ the x_3 -average of m_h given in (7), we have by Theorem 1.1 in the regime (4),

$$\limsup_{\varepsilon \rightarrow 0} \frac{1}{|\log \varepsilon|} E_{\varepsilon, \eta}^{\delta}(\overline{m}'_h) \leq \limsup_{h \rightarrow 0} E_h(m_h) < \infty.$$

As $|\log \varepsilon| \ll |\log \eta|$, we can apply Theorem 1.2 to $v_{\varepsilon} := \overline{m}'_h: \Omega \rightarrow \mathbb{R}^2$. More precisely, by Theorem 1.2(i), for a subsequence, $(\mathcal{J}(\overline{m}'_h))$ converges (in the sense of (11)) to the measure J given in (12). Moreover, for a subsequence, $(\overline{m}'_h|_{\partial\Omega})$ converges to $e^{i\varphi_0} \in BV(\partial\Omega, \mathbb{S}^1)$ in $L^p(\partial\Omega)$, for every $p \in [1, +\infty)$, where $\varphi_0 \in BV(\partial\Omega, \pi\mathbb{Z})$ is a lifting of the tangent field $\pm\tau$ on $\partial\Omega$ determined (up to a constant in $\pi\mathbb{Z}$) by $\partial_{\tau}\varphi_0 = -J$ as measure on $\partial\Omega$. Since $\overline{m}_{h,3}^2 \leq 1 - |\overline{m}'_h|^2$ and $(\overline{m}'_h)_h$ converges to $e^{i\varphi_0}$ in $L^2(\partial\Omega, \mathbb{R}^2)$ with $|e^{i\varphi_0}| = 1$, we deduce that $(\overline{m}_{h,3})_h$ converges to zero in $L^2(\partial\Omega)$ and almost everywhere on $\partial\Omega$ (up to a subsequence). As $|\overline{m}_{h,3}| \leq 1$, by dominated convergence theorem, we get that $(\overline{m}_{h,3})_h$ converges to zero in $L^p(\partial\Omega)$ for every $p \in [1, +\infty)$. For proving (ii), we apply Theorem 1.2(ii) and Theorem 1.1 to get

$$\pi \sum_{j=1}^N |d_j| \leq \liminf_{h \rightarrow 0} \frac{1}{|\log \varepsilon|} E_{\varepsilon, \eta}^{\delta}(\overline{m}'_h) \leq \liminf_{h \rightarrow 0} E_h(m_h).$$

\square

Proof of Theorem 1.11. By (19) and Theorem 1.1 in the regime (5), we have

$$\infty > \limsup_{h \rightarrow 0} |\log \varepsilon| \left(E_h(m_h) - \pi \sum_{j=1}^N |d_j| \right) \geq \limsup_{h \rightarrow 0} \left(E_{\varepsilon, \eta}^\delta(\overline{m}'_h) - \pi |\log \varepsilon| \sum_{j=1}^N |d_j| \right).$$

As $|\log \varepsilon| \ll |\log \eta|$, we can apply Theorem 1.4 to $v_\varepsilon := \overline{m}'_h$. More precisely, by Theorem 1.4(i), we have $d_j \in \{\pm 1\}$ for every $j \in \{1, \dots, N\}$, so that $\sum_{j=1}^N |d_j| = N$, and

$$W_\Omega^\delta(\{(a_j, d_j)\}) + N\gamma_0 \leq \liminf_{h \rightarrow 0} (E_{\varepsilon, \eta}^\delta(\overline{m}'_h) - N\pi |\log \varepsilon|) \leq \liminf_{h \rightarrow 0} |\log \varepsilon| (E_h(m_h) - N\pi).$$

Let us prove (ii). By the proof of Theorem 1.1, we have

$$\begin{aligned} E_h(m_h) &= E_h^0(m_h) + \frac{1}{|\log \varepsilon|} \frac{1}{\eta^2} \int_\Omega \widehat{D}' : \nabla' \overline{m}_h \wedge \overline{m}_h \, dx' - O(R(h)) \\ &= E_h^0(m_h) + \frac{1}{|\log \varepsilon|} \int_\Omega 2\delta \cdot \nabla' \overline{m}'_h \wedge \overline{m}'_h \, dx - o_h(R_1(h)) - O(R(h)) \end{aligned}$$

where $o_h(R_1(h)) + O(R(h)) = o_h(\frac{1}{|\log \varepsilon|})$ in the regime (5). By (39), $\int_\Omega 2\delta \cdot \nabla' \overline{m}'_h \wedge \overline{m}'_h \, dx = O(1)$. By (19), we deduce that

$$\limsup_{h \rightarrow 0} |\log \varepsilon| (E_h^0(m_h) - N\pi) \leq \limsup_{h \rightarrow 0} |\log \varepsilon| (E_h(m_h) - N\pi) + O(1) < \infty.$$

Hence, we can apply [26, Theorem 9(iv)] and we obtain the expected compactness of $(m_h)_h$ and properties of their limit points. \square

Proof of Theorem 1.12. In the regime (4), we have $|\log \varepsilon| \ll |\log \eta|$ and we construct the family $\{v_\varepsilon : \Omega \rightarrow \mathbb{S}^1\}$ as in Theorem 1.5. Set $m_h : (x', x_3) \in \Omega_h \mapsto (v_\varepsilon(x'), 0) \in \mathbb{S}^1 \times \{0\}$. For every $h > 0$, m_h is clearly independent of x_3 , thus $m_h = \overline{m}_h$ and by Theorem 1.5, it follows that the global Jacobian $\mathcal{J}(m'_h) = \mathcal{J}(v_\varepsilon)$ converges (in the sense (11)) to the measure J given in (12). Also, $(m'_h)_h$ converges strongly to $e^{i\varphi_*}$ in $L^p(\Omega, \mathbb{R}^2)$ and in $L^p(\partial\Omega, \mathbb{R}^2)$, for every $p \in [1, +\infty)$, where φ_* is the harmonic extension to Ω of a boundary lifting φ_0 of the tangent field $\pm\tau$ on $\partial\Omega$ that satisfies $\partial_\tau \varphi_0 = -J$ as measure on $\partial\Omega$. By Theorem 1.1 in the regime (4) combined with Theorem 1.5, we have

$$E_h(m_h) = \frac{1}{|\log \varepsilon|} E_{\varepsilon, \eta}^\delta(v_\varepsilon) - o_\varepsilon(1) = \pi \sum_{j=1}^N |d_j| - o_\varepsilon(1).$$

Furthermore, if $d_j \in \{\pm 1\}$ for every $j \in \{1, \dots, N\}$, then by Theorem 1.1 in the regime (5) and Theorem 1.5,

$$\begin{aligned} |\log \varepsilon| (E_h(m_h) - N\pi) &= |\log \varepsilon| \left(\frac{1}{|\log \varepsilon|} E_{\varepsilon, \eta}^\delta(v_\varepsilon) - N\pi - o\left(\frac{1}{|\log \varepsilon|}\right) \right) \\ &= W_\Omega^\delta(\{(a_j, d_j)\}) + N\gamma_0 - o_h(1). \end{aligned}$$

\square

Proof of Corollary 1.13. By Corollary 1.7, there exist two points $a_1^* \neq a_2^* \in \partial\Omega$ such that

$$W_\Omega^\delta(\{(a_1^*, 1), (a_2^*, 1)\}) = \min \{W_\Omega^\delta(\{(\tilde{a}_1, 1), (\tilde{a}_2, 1)\}) : \tilde{a}_1 \neq \tilde{a}_2 \in \partial\Omega\}.$$

Let $(m_h)_h$ be a family of minimisers of E_h on $H^1(\Omega_h, \mathbb{R}^3)$ (such a family exists by Corollary 3.2). In the regime (4), by Theorem 1.12 applied to $\{(a_1^*, 1), (a_2^*, 1)\}$, the minimisers m_h must satisfy

$$\limsup_{h \rightarrow 0} E_h(m_h) \leq 2\pi. \quad (64)$$

Hence, we can apply Theorem 1.10(i) and we deduce that, for a subsequence, the global Jacobian $(\mathcal{J}(\overline{m'_h}))_h$ converges as in (18) to

$$J = -\kappa \mathcal{H}^1 \llcorner \partial\Omega + \sum_{j=1}^N d_j \delta_{a_j}$$

for $N \geq 1$ distinct boundary points $a_1, \dots, a_N \in \partial\Omega$, with $d_1, \dots, d_N \in \mathbb{Z} \setminus \{0\}$ such that $\sum_{j=1}^N d_j = 2$. Moreover, by Theorem 1.10(ii), we also have

$$\liminf_{h \rightarrow 0} E_h(m_h) \geq \pi \sum_{j=1}^N |d_j|. \quad (65)$$

Combining (64) and (65), we get $\sum_{j=1}^N |d_j| \leq 2 = \sum_{j=1}^N d_j$, hence $\sum_{j=1}^N (|d_j| - d_j) \leq 0$ so that for every $j \in \{1, \dots, N\}$, $|d_j| = d_j$. It follows that

$$\lim_{h \rightarrow 0} E_h(m_h) = 2\pi,$$

and two cases can occur: either $N = 1$ and $d_1 = 2$, or $N = 2$ and $d_1 = d_2 = 1$. Hence, there are two boundary points $a_1, a_2 \in \partial\Omega$ (that might a-priori coincide) such that

$$J = -\kappa \mathcal{H}^1 \llcorner \partial\Omega + \pi(\delta_{a_1} + \delta_{a_2}).$$

We now assume that the regime (5) holds. By Theorem 1.12, we necessarily have

$$|\log \varepsilon| (E_h(m_h) - 2\pi) \leq W_\Omega^\delta(\{(a_1^*, 1), (a_2^*, 1)\}) + 2\gamma_0 + o_\varepsilon(1) \quad \text{as } h \rightarrow 0. \quad (66)$$

Hence, we can apply Theorem 1.11(i) and we deduce that $d_1 = d_2 = 1$, $N = 2$ (thus $a_1 \neq a_2$) and

$$|\log \varepsilon| (E_h(m_h) - 2\pi) \geq W_\Omega^\delta(\{(a_1, 1), (a_2, 1)\}) + 2\gamma_0 + o_\varepsilon(1) \quad \text{as } h \rightarrow 0. \quad (67)$$

Combining (66) and (67), we deduce that

$$W_\Omega^\delta(\{(a_1, 1), (a_2, 1)\}) \leq W_\Omega^\delta(\{(a_1^*, 1), (a_2^*, 1)\}).$$

By definition of $W_\Omega^\delta(\{(a_1^*, 1), (a_2^*, 1)\})$, we have

$$W_\Omega^\delta(\{(a_1, 1), (a_2, 1)\}) = W_\Omega^\delta(\{(a_1^*, 1), (a_2^*, 1)\}),$$

and thus

$$\lim_{\varepsilon \rightarrow 0} |\log \varepsilon| (E_h(m_h) - 2\pi) = W_\Omega^\delta(\{(a_1, 1), (a_2, 1)\}) + 2\gamma_0.$$

□

References

- [1] A. Aharoni, Energy of one dimensional domain walls in ferromagnetic films, *J. Appl. Phys.* **37** (1966), 3271-3279.
- [2] A. Aharoni, *Introduction to the Theory of Ferromagnetism, Second edition*, Oxford University Press, 2001.
- [3] S. Alama, L. Bronsard, and D. Golovaty, Thin film liquid crystals with oblique anchoring and boojums, *Ann. Inst. Henri Poincaré, Anal. Non Linéaire* **37** (2020), 817-853.
- [4] C. J. Amick and J. F. Toland, Uniqueness and related analytic properties for the Benjamin-Ono equation – a nonlinear Neumann problem in the plane, *Acta Math.* **167** (1991), 107-126.
- [5] T. B. Benjamin, Internal waves of finite amplitude and permanent form, *Journal of Fluid Mechanics* **25** (1966), 241-270.
- [6] F. Bethuel, H. Brezis, and F. Hélein, *Ginzburg-Landau vortices*, Progress in Nonlinear Differential Equations and Their Applications, vol.13, Birkhäuser Boston Inc, Boston, MA, 1994.
- [7] F. Bethuel and X. Zheng, Density of smooth functions between two manifolds in Sobolev spaces, *Journal of Functional Analysis* **80** (1988).
- [8] J. Bourgain, H. Brezis, and P. Mironescu, Lifting in Sobolev spaces, *Journal d'Analyse Mathématique* **80** (2000), 37-86.
- [9] X. Cabré and J. Solà-Morales, Layer solutions in a half-space for boundary reactions, *Communications on Pure and Applied Mathematics* **58** (2005), 1678-1732.
- [10] G. Carbou, Thin layers in micromagnetism, *Mathematical Models and Methods in Applied Sciences* **11** (2001), 1529-1546.
- [11] R. Côte and R. Ignat, Asymptotic stability of precessing domain walls for the Landau-Lifshitz-Gilbert equation in a nanowire with Dzyaloshinskii-Moriya interaction, *Communications in Mathematical Physics* **401** (2023), 2901-2957.
- [12] G. Dal Maso, *An Introduction to Γ -convergence*, Progress in Nonlinear Differential Equations and Their Applications, vol.8, Birkhäuser Boston Inc, Boston, MA, 1993.
- [13] E. Davoli, G. Di Fratta, D. Praetorius, and M. Ruggeri, Micromagnetics of thin films in the presence of Dzyaloshinskii-Moriya interaction, *Mathematical Models and Methods in Applied Sciences* **32** (2022), 911-939.
- [14] A. DeSimone, R. V. Kohn, S. Müller, and F. Otto, Recent analytical developments in micromagnetics *The science of hysteresis II*, Elsevier, 2006, 269-281.
- [15] G. Di Fratta, C. B. Muratov and V. V. Slastikov, Reduced energies for thin ferromagnetic films with perpendicular anisotropy, arXiv:2306.07634, 2023.
- [16] I. Dzyaloshinskii, A thermodynamic theory of "weak" ferromagnetism of antiferromagnetics, *Journal of Physics and Chemistry of Solids* **4** (1957), 241-255.
- [17] L. C. Evans, *Partial Differential Equations, Second edition*, Graduate Studies in Mathematics, 1998.

- [18] C. García-Cervera, *Magnetic Domains and Magnetic Domain Walls*, Ph.D thesis, Courant Institute of Mathematical Sciences, New York University, 1999.
- [19] D. Gilbarg and N. S. Trudinger, *Elliptic Partial Differential Equations of Second Order*, Springer, 2001.
- [20] G. Gioia and R.D. James, Micromagnetics of very thin films, *Proceedings of the Royal Society of London. Series A: Mathematical, Physical and Engineering Sciences* **453** (1997), 213-223.
- [21] A. Hubert and R. Schäfer, *Magnetic Domains*, Springer-Verlag Berlin Heidelberg, 1998.
- [22] R. Ignat, Singularities of divergence-free vector fields with values into S^1 or S^2 . Applications to micromagnetics, *Confluentes Mathematici* **4** (2012), 1-80.
- [23] R. Ignat, A Gamma-convergence result for Néel walls in micromagnetics, *Calculus of Variations and Partial Differential Equations* **36** (2009), 285-316.
- [24] R. Ignat and H. Knüpfer, Vortex energy and 360° Néel walls in thin-film micromagnetics, *Comm. Pure Appl. Math.* **63** (2010), 1677-1724.
- [25] R. Ignat and M. Kurzke, Global Jacobian and Gamma-convergence in a two-dimensional Ginzburg-Landau model for boundary vortices, *Journal of Functional Analysis* **280** (2021).
- [26] R. Ignat and M. Kurzke, An effective model for boundary vortices in thin-film micromagnetics, *Mathematical Models and Methods in Applied Sciences* **33** (2023), 1929-1973.
- [27] R. Ignat, M. Kurzke and X. Lamy, Global uniform estimate for the modulus of two-dimensional Ginzburg-Landau vortexless solutions with asymptotically infinite boundary energy, *SIAM J. Math. Anal.* **52** (2020), 524-542.
- [28] R. Ignat and R. Moser, Interaction energy of domain walls in a nonlocal Ginzburg-Landau type model from micromagnetics, *Archive for Rational Mechanics and Analysis* **221** (2016), 419-485.
- [29] R. Ignat and R. Moser, Energy minimisers of prescribed winding number in an S^1 -valued nonlocal Allen-Cahn type model, *Adv. Math.* **357** (2019).
- [30] R. Ignat and F. Otto, A compactness result in thin-film micromagnetics and the optimality of the Néel wall, *J. Eur. Math. Soc. (JEMS)* **10** (2008), 909-956.
- [31] R. Ignat and F. Otto, A compactness result for Landau state in thin-film micromagnetics, *Ann. Inst. H. Poincaré, Anal. Non linéaire* **28** (2011), 247-282.
- [32] R. V. Kohn and V. V. Slastikov, Another Thin-Film Limit of Micromagnetics, *Archive for Rational Mechanics and Analysis* **178** (2005), 227-245.
- [33] M. Kurzke, Boundary vortices in thin magnetic films, *Calculus of Variations and Partial Differential Equations* **26** (2006), 1-28.
- [34] M. Kurzke, A nonlocal singular perturbation problem with periodic potential well, *ESAIM Control Optim. Calc. Var.* **12** (2006), 52-63.
- [35] M. Kurzke, The gradient flow motion of boundary vortices, *Ann. Inst. H. Poincaré Anal. Non Linéaire* **24** (2007), 91-112.
- [36] F. L'Official, The micromagnetic energy with Dzyaloshinskii-Moriya interaction in a thin-film regime relevant for boundary vortices, *Communications on Pure and Applied Analysis* **22** (2023), 1271-1295.

- [37] C. Melcher, The logarithmic tail of Néel walls, *Archive for Rational Mechanics and Analysis* **168** (2003), 83-113.
- [38] C. Melcher, *Logarithmic lower bounds for Néel walls*, *Calculus of Variations and Partial Differential Equations* **21** (2004), 209-219.
- [39] R. Moser, Ginzburg-Landau vortices for thin ferromagnetic films, *Applied Mathematics Research eXpress*, 2003, 1-32.
- [40] R. Moser, Boundary vortices for thin ferromagnetic films, *Archive for Rational Mechanics and Analysis* **174** (2004), 267-300.
- [41] R. Moser, Moving boundary vortices for a thin-film limit in micromagnetics, *Comm. Pure Appl. Math.* **58** (2005), 701-721.
- [42] F. R. N. Nabarro, Dislocations in a simple cubic lattice, *Proceedings of the Physical Society* **59** (1947), 256-272.
- [43] J. F. Toland, The Peierls-Nabarro and Benjamin-Ono Equations, *Journal of Functional Analysis* **145** (1997), 136-150.
- [44] H. Weyl, The method of orthogonal projection in potential theory, *Duke Math. J.* **7** (1940), 414-444.