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Using the Quartz Crystal Microbalance to Monitor the Curing of Drying Oils

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Abstract

Drying oils like linseed oil form a polymer network through a complex free-radical polymerization process. We have studied polymerization in this challenging class of polymers using the quartz crystal microbalance (QCM). The QCM is able to measure the evolution of polymer mass and mechanical properties as the oil transitions from a liquid-like to a solid-like state. Measurements using bulk materials and thin films provide information about the initial polymerization phase as well as the evolution of the mass and mechanical properties over the first two years of cure. The temperaturedependent response of cured linseed oil films was also measured. These results were combined with previously published results obtained from traditional dynamic mechanical analysis to give a unified picture of the properties of these materials across a very broad temperature range.

Introduction

Since the $15th$ century, linseed oil has been commonly used by painters as a binding medium for paints **?** . One of the material properties that drew painters to using linseed oil was its ability to mix with pigments and additives, coat surfaces well, and then undergo a curing process to form a solid film. To speed up this process, it was common practice to treat the oil before using it, in particular by heating in the presence of siccatives (usually lead compounds)^{1,2}. In drying oils such as linseed oil, a polymer film forms as a result of an auto-oxidation process where peroxide, ether, and alkyl crosslinks are formed as a result of oxygen reacting with unsaturated fatty acid chains in triglyceride molecules of the oil^{3–5}. While there are many chemical reactions that contribute to the curing process of drying oils, on a more macroscopic level the oil will transition from being liquid-like to solid-like as the polymer film forms. The stiffness of a polymer system increases as crosslinks form between polymer chains. Crosslink formation in drying oils is typically considered a slow curing process, while most synthetic polymer systems use ultraviolet light, temperature, or chemical initiators to induce and accelerate this process^{6–9}. It is important to learn how drying oils form networks in order to ultimately understand long-term degradation and the factors that influence network stability^{10?-12}. The network

formation process of linseed oil has been studied extensively through chemical investigations including infrared spectroscopy and chromatographic analysis $1,2,11,13,14$, but the mechanical properties during the initial curing step have never been properly characterized. The sol-gel transition associated with polymer network formation makes the most common approaches to measuring the mechanical properties of this film impractical to use. Rheology is most ideally used on liquid-like samples and has been used to understand the flow and application properties of paints and drying oils with an emphasis on the effects of historic preparation recipes for these oils $15-19$. Tensile testing, dynamic mechanical analysis (DMA) and micro-/nano-indentation can be used to investigate the mechanical properties of dried films, but the samples must be solid in order to be measured, making these techniques not feasible to use on oil films that have not reached an initial state of cure $^{10,20-29}.$

The quartz crystal microbalance (QCM) is a uniquely suitable technique to measure the mechanical properties of a polymer film experiencing a liquid to solid transition. This technique uses the piezoelectric properties of quartz to propagate a mechanical shear wave at high frequencies (MHz) through a film of interest, as illustrated in Figure $1^{30,31}.$ The resulting shifts in the resonant frequencies of the quartz crystal can be used to determine the mass and mechanical properties of the film³². A QCM sample typically has a surface of the film open to the surrounding environment, which can make it an ideal technique for looking at changes in the mechanical properties as a result of temperature, relative humidity, or exposure to solvent $8,33-36$. The films of interest are also fairly thin (on the order of $0.7 - 6 \,\mu\text{m}$), effectively removing any diffusion barrier that would be present for a bulk film. The QCM has been previously used to characterize the evolution of mechanical properties in many polymer systems, including an alkyd binding medium which undergoes a similar curing mechanism to that of linseed oil $8,32-34,36-43$. In the present work we use linseed oil as a case study to illustrate both the strengths and limitations of the use of the QCM to investigate time-dependent physical property changes for times ranging

from minutes to years.

Figure 1: The (a) bulk and (b) thin film sample geometries used for experiments with the quartz crystal microbalance (QCM) in this study. The quartz crystal sample substrate consists of a quartz disk (thickness of $330 \,\mu$ m and a diameter of $2.5 \,\text{cm}$) with two gold (Au) electrodes on either surface.

The detailed mathematical formulation that relates the measured QCM response to the properties of the thin film has been discussed previously^{32,44,45}, and is included in the supporting information for completeness. Briefly, we measure the resonant frequency and bandwidth (equivalent to the dissipation that is often reported in in experiments referred to as QCM-D experiments) at multiple harmonics of the crystal resonance, both before and after a thin film of material is deposited on the electrode surface of the QCM. In the most general 'thin film' analysis illustrated in Figure 1b, three quantities related to the behavior of the crystal resonances are obtained and the following three film properties are extracted:

- $d\rho$, the product of the film thickness and density, or equivalently, the mass per unit area of the film.
- $|G_3^*|$ ρ , the product of the density and magnitude of the complex shear modulus at the 3rd resonant harmonic (15 MHz in our case).
- ϕ , the viscolastic phase angle, assumed to be the same for the different harmonics that are used. This quantity is more commonly referred to as δ in the rheological

community, but we use ϕ here to remain consistent with our previous work, and to distinguish it from the decay length, defined below in Eq. 1.

In the simpler bulk analysis shown in Figure 1a we measure two quantities (typically the resonant frequency shift and bandwidth at the third harmonic) to obtain $|G_3^*| \rho$ and ϕ , obtaining no information about the total mass of the sample.

In the current paper we use the QCM to address the following questions as they pertain to linseed oil.

- How do the viscoelastic properties of linseed oil evolve during the very early stages of the curing process?
- How does an oil pre-treatment (heating in the presence of lead oxide, commonly reported in historic recipes) affect the curing kinetics and the mechanical properties of the fully cured systems?
- What is the temperature dependence of the mechanical properties of oil films, dried in different conditions? We are particularly interested in identifying the glass transition temperature, T_g , where the material behavior transitions from rigid, glassy behavior (below T_q) to soft, rubbery behavior (above T_q).

We also address a series of more general, practical issues pertaining to the use of the QCM to determine the mechanical properties of materials. These issues include the limitations of the relatively simple bulk analysis, the optimized sample thicknesses and harmonics to be used for a QCM measurement, and a quantitative **uncertaintly** analysis of the mechanical properties that are obtained.

Materials and Methods

Three different types of linseed oil were used in this study, which we refer to as SA, KP and WN, based on the supplier:

- SA: boiled linseed oil from Sigma Aldrich (without siccative)
- KP: cold-pressed linseed oil from Kremer Pigmente, also used by Baij *et al.* in their previously reported DMA experiments 46
- WN: refined linseed oil from Winsor & Newton

To better understand the influence of oil pre-treatment, we prepared oil according to historical recipes² by heating it with PbO. Litharge (PbO) was purchased from Sigma Aldrich and mixed with the KP cold-pressed linseed oil to give samples with PbO weight fractions of 0.05 and 0.2. These samples were prepared by placing the PbO into a mortar and grinding with a few drops of oil. The rest of the oil was added to this PbO paste in a beaker and heated at 150 °C for 2 hours under magnetic stirring. The linseed oils from Sigma Aldrich and Winsor & Newton were used as is. As preparation layers for the substrate, either polystyrene (192k MW, Sigma Aldrich) or 1-propanethiol were used. Toluene, turpentine, mineral spirits, and tetrahydrofuran were used as solvents during the spin casting for the pretreatment and oil films. The oils were cast on quartz crystal electrodes (AT-cut, 2.54 cm diameter, 330 μ m thickness) purchased from Renlux (Nanshan, China) and Advanced Wave Sensors (AWS, Valencia, Spain). Additional details of the sample preparation, including substrate pretreatments used to minimize complications arising from potential dewetting of the thin linseed oil film, are included in the supporting information.

Results

Bulk analysis of early stage cure

A bulk sample of the SA linseed oil was measured to determine the initial viscosity, the change in mechanical properties over time, and the extent to which those changes could be observed using this experimental geometry. Values of $|G_3^*| \rho$ (or equivalently, $|\eta_3^*| \rho$

where $|\eta_3^*|$ is the magnitude of the complex shear viscosity at the third harmonic frequency) and ϕ are plotted in Figure 2. Assuming a density of 0.93 g/cm³ for linseed oil, we obtain the magnitude of the complex viscosity as it increases from 40 mPa·s at the beginning of the experiment to 125 mPa·s after about 3 days. The initial value is consistent with the viscosity of 44 mPa·s determined from rheological steady state flow experiments at shear rates (0.01-100 s⁻¹, see supporting information) which are much lower than the angular frequency of $\approx 10^8 \text{ s}^{-1}$ for the QCM experiments. The properties remained constant for the first 1.5 days. This initial period of constant viscosity is attributed to the presence of antioxidants in the linseed oil, which react with oxygen before the auto-oxidation process begins. Auto-oxidation induces crosslink formation in the oil and increases the complex shear modulus/viscosity as well as the decay length.

The QCM senses the properties of the material in the immediate vicinity of the electrode surface, within a distance determined by the decay length, δ , of the shear wave in the oil, which is determined by the frequency and the viscoelastic properties of the material:

$$
\delta = \frac{(|G^*|/\rho)^{1/2}}{2\pi f \sin(\phi/2)}
$$
(1)

When the viscosity in this portion of the sample becomes sufficiently large, the sample moves from the 'bulk' regime (where the sample thickness, d , is assumed to be much larger than δ) to the 'overdamped' regime and quantitative measurements are no longer possible⁴². Table 1 provides the decay length for water (a Newtonian liquid) as well as the SA bulk sample at early and later times of cure, along with the mechanical properties obtained from the QCM experiment. The most important outcome from these bulk experiments is the conclusion that they are restricted to materials with complex moduli below $\approx 2 \times 10^7$ Pa. This result is consistent with the range of moduli for which this same bulk analysis has been used by Szántó *et al.* in their QCM measurements of the rheological properties of polyethylene and copolymers of ethylene/vinyl acetate⁴⁷. For larger values of $|G^*| \Gamma$ is too large to give a well-defined resonance, as described in more detail

in the supporting info. For liquid samples at our measurement frequency of $1.5x10⁷$ Hz this limitation restricts us to materials with viscosities less than ≈ 0.2 Pa·s, which is 200 times the room-temperature viscosity of water. This limitation is rather severe in the context of studying polymerization processes, and motivates the development of the more complicated thin-film analysis utilized in the following section.

Figure 2: Measured time-dependent properties (a and b) and solution checks (c and d) for the bulk SA linseed oil sample. The solution checks compare the experimentally measured values of Δf_n and $\Delta \Gamma_n$ to the values of these quantities that are calculated from the measured properties. Because this is a 3:3 calculation, properties are determined from the requirement that experimental and calculated values are in agreement for Δf_3 and $\Delta\Gamma_3$. Uncertainties in Δf_n and $\Delta\Gamma_n$ are slightly larger than 5% of $\Delta\Gamma_n$, and are used to determine uncertainties in the measured properties as described in the supporting information.

	$ G_3^* \rho$			\mathcal{O}_3	$ \eta_3^* $	$\Delta\Gamma_3$
	$(Pa \cdot g/cm^3)$	(deg.)	(g/cm^3)	(μm)	$(mPa \cdot s)$	(Hz)
Water	9.4×10^{4}	90		0.145		1,230
Linseed oil, $t = 1$ hr	3.5×10^6	80	0.93	1.05	40	8,160
Linseed oil, $t = 72$ hr	1.1×10^{7}	80	0.93	1.86	125	14,500

Table 1: Extracted physical properties of the SA bulk sample at early and late stages of drying. Measurements are no longer possible when the bandwidth, Γ, is larger than \approx 20 kHz.

Thin film analysis

Film thickness considerations

Two linseed oils processed in different ways, SA and WN, were measured in ambient lab conditions and the evolution of mechanical properties was followed. Results obtained from the SA oil sample are particularly useful for illustrating the importance of the sample thickness. For this sample two complementary analysis approaches were used to obtain the properties shown in Figure $3(a,b,c)$. For the longer aging times ($t < 55 \text{ days}$) the dissipation for the 5th harmonic is low enough so that a suitably accurate value of Δf_5 is obtained. In these cases the properties were obtained from a 3,5:3 calculation, where values of Δf_3 , Δf_5 and $\Delta \Gamma_3$ are used. For shorter aging times the dissipation for the 5th harmonic is too large, meaning that Δf_5 cannot be accurately measured. In this case we used a 3:3 calculation involving only Δf_3 and $\Delta \Gamma_3$, assuming values of $d\rho$ determined from the longer aging times. We use two values of $d\rho$ that bracket the 55 day value obtained from the 3,5:3 calculation, and which differ from one another by 3%. In this way we assess the sensitivity of the extracted properties to this assumed thickness. As shown in Figure 3, the assumed value of $d\rho$ has a negligible effect on the measured values of $|G_3^*|$ ρ and ϕ at the shortest aging times. At the larger aging times, a 3% deviation in the assumed value of $d\rho$ results in a 25% deviation in the calculated values of $|G_3^*|$ ρ and ϕ .

Figure 3: Evolution of mechanical properties over time of the SA sample. The extracted properties ($|G_3^*| \rho$, ϕ , and $d\rho$) are shown in (a), (b) and (c). The 3:3 calculations use only the 3rd harmonic, assuming a fixed value of $d\rho$ listed in the legend. A 3.5:3 calculation was used for later times (t >55 days). Response maps for the SA LO sample are shown in (d) and (e). The color maps correspond to values of $\Delta f/n$ and $\Delta \Gamma/n$ calculated with the assumption that $d\rho = 3.35 \mu m \cdot g / cm^3$. This value of $d\rho$ is representative of the values obtained throughout the experiment. The specific choice of $d\rho$ here affects the detailed visualization of the response map, but does not affect the analysis itself.

The inability to measure film properties for intermediate aging times can be understood in terms of the response maps for these data, shown in Figure $3(d,e)$. An important parameter here is λ_n , the wavelength of the shear wave in the viscoelastic medium, which is related to the frequency and viscoelastic properties of the film as follows³²:

$$
\lambda_n = \frac{1}{f_n} \left(\frac{|\mathcal{G}_n^*|}{\rho} \right)^{1/2} \frac{1}{\cos(\phi/2)}.
$$
 (2)

Here we include the harmonic number n as a reminder that we are collecting data at multiple frequencies, accounting for a frequency dependence of G_n^* but assuming that ϕ is independent of n for the values of n that are used in the calculation of the properties. We further assume that $|G_n^*|$ increases with frequency (or *n*), according to the following power law:

$$
|G_n^*| \propto \phi^{n/90} \tag{3}
$$

with ϕ in degrees. Eq. 3 becomes exact, with a frequency-independent phase angle, if the power law is valid over a a very wide frequency range above and below the frequencies of interest^{48,49}, but is in general an approximation. This approximation is surprisingly good, standardizing and simplifying the analysis while introducing an error that remains small in comparison to errors originating from experimental uncertainties in the measured values of Δf and $\Delta \Gamma$. For example, it was previously shown using typical experimental data for a crosslinked polymer that in the regime where the thin-film analysis is necessary $(|G^*| \gtrapprox 2x10^7 \text{ Pa})$, this constant phase angle approximation introduces an error no larger than 10% in the measured values of both ϕ and- $|G^*| \rho$ ⁴⁵.

Response maps are a graphical representation of the following QCM master equation for the thin film regime 32 :

$$
\Delta f_n^* \equiv \Delta f_n + i \Delta \Gamma_n = -\Delta f_{sn} \left(\frac{\tan \{ (2\pi d/\lambda_n)(1 - i \tan(\phi_n/2)) \}}{(2\pi d/\lambda_n)(1 - i \tan(\phi_n/2))} \right)
$$
(4)

Here Δf_{sn} is the Sauerbrey frequency shift:

$$
\Delta f_{sn} = \frac{2nf_1^2}{Z_q} \rho d \tag{5}
$$

where Z_q is the acoustic impedance of quartz (8.84 $\times 10^6$ Kg/m²s for the AT quartz used here).

For this sample, the early times are in the upper right corner with high d/λ_n (equivalent to a lower complex shear modulus) and high phase angle. As the sample polymerizes over a period of several hours, the data shift to the lower left corner of the response maps. The areas of white on the response maps are where the film resonance phenomenon occurs, giving values for Γ that are too large for mechanical property data to be obtained. In order for this region to be avoided so that data can be obtained throughout the entire

curing process of the same sample, thinner samples must be used, as is the case for the remainder of the samples discussed in this paper.

Effect of Oil Treatment with PbO

The next comparison we investigated involved the addition of PbO to linseed oil, a common technique used by painters throughout history to reduce the drying time of the oil paints and improve the properties for paint application to a surface^{$2,50-52$}. This treatment leads to the partial saponification of the oil (among other reactions due to heating). For these measurements, a cold-pressed linseed oil (KP) was chosen. The time-dependent properties of the KP films are summarized in Figure 4, and involve the following samples:

- One film without PbO (KP0) that was monitored over 4 days.
- Three films of oil + 5% PbO (KP5-1,2,3) that were monitored for periods ranging from one hour to two years.
- One film of oil + 20% PbO (KP20) that was monitored from three days to three weeks.

Property uncertainties become large for the thinnest films, where d/λ_n becomes too small for accurate properties to be determined. Property uncertainties for the KP5-2 film are particularly large, and data from this film are not included in Figure 4 for this reason. Solution checks and response maps for each of the samples are included in the supporting information.

Figure 4: Evolution of mechanical properties for the KP films with no added PbO (KP0), with 5 wt% PbO (KP5-1 and KP5-3) and 20 wt% PbO (KP20).

Two of the KP samples capture the evolution of the mechanical properties during the initial polymerization of the film: KP0 and KP5-1, where $|G_3^*| \rho$ increases and ϕ decreases. These changes in the mechanical properties are due to oxygen uptake by the film, which fuels the auto-oxidative polymerization of the films. The main difference between these two films is the time at which this polymerization occurred, at approximately 25 hours and 3-5 hours for the KP0 and KP5-1 films, respectively. This difference in the rate of initial polymerization is a result of using PbO as a siccative and as part of a heat treatment for the linseed oil. To fully capture this initial polymerization regime for the KP5-1 sample, two different calculations were used, 3,5:5 for times less than 4.5 hr and 5,9:9 for longer times. A detailed comparison of these two calculations is included in the supporting information.

One important result from these samples is the similarity of $|G_3^*|$ ρ at times greater than 4 days for all the samples. These moduli are comparable to the values obtained for oil paints measured using tensile testing as well as well as the values determined by the QCM for a commercial alkyd binding medium^{20,21,34}. Note that because of the high frequency of the QCM measurement (15 MHz), this technique accesses the glassy regime of the cured films, even at temperatures above the glass transition temperature. We expand on this point below in the section on the effects of temperature on the mechanical properties.

Temperature dependent mechanical properties of cured oil films

The WN film was of sufficient quality so that temperature sweeps could be performed from -40 \degree C to 80 \degree C. Results from the QCM measurements after room temperature aging at 5 months are shown in Figure 5a. Also included in this plot are published data from Baij *et al.* for linseed oil films that were cured overnight at 150 °C.⁴⁶. Measured values of E ∗ (the complex dynamic modulus in extension) from Baij *et al.* were converted to values of $G^*\rho$ as described in the supporting information. In addition to the different processing conditions (long term cure at room temperature vs. shorter cure at 150 ◦C) and linseed oil source (KP vs. WN), an important difference between our measurements and those of Baij *et al.* concerns the frequency of our measurements. The measurements of Baij *et al.* were performed at frequencies of 0.1 Hz, 1 Hz and 10 Hz, whereas our QCM data correspond to a measurement frequency of $1.5x10⁷$ Hz. Comparison of the data obtained at these very different frequencies requires the use of time-temperature superposition, where the frequencies are multiplied by temperature-dependent shift factors, a_T . These shift factors are determined by the requirement that superposable master curves be obtained when the viscoelastic properties are plotted as a function of the reduced frequency, fa_T . These shift factors are important in that they describe the temperature dependence of the mechanical relaxation times within the material.

Figure 5: a) Temperature dependence of $|G^*| \, \rho$ at 1.5x10⁷ Hz from the QCM data, and at much lower frequencies from the DMA data from Baij *et al.* ⁴⁶; b) Temperature-dependent shift factors with an Arrhenius fit (Eq. 6, $E_a = 300 \text{ kJ/mol}$) for temperatures less than -7.4° C and a Vogel-Fulcher-Tammann fit (Eqs. 7) for temperatures above -7.4°; c), master curve for $|G^*|$ ρ using the temperature shift factors from part (b).

The shift factors generated from our analysis of the data are shown in Figure 5b. The corresponding rheological master curves, where we plot $|G^*| \rho$ as a function of the temperature-shifted frequency $(f a_T)$ for all of the data, are shown in Figure 5c. An important parameter here is the glass transition temperature, T_g . Several different metrics can be applied to specify T_q for viscoelastic materials. Use of the maximum in the phase angle at a specified frequency (typically 1 Hz) is one example, which in this case gives $T_g \approx 8\,^{\circ}\text{C}^{46}.$ In synthetic polymers the onset of material softening at 1 Hz as the sample is heated is more closely correlated with the calorimetrically determined glass transition temperature⁴¹. With this criterion (described more quantitatively in the supporting information), we obtain $T_g \approx -7.4\,^{\circ}\text{C}$ for these materials. This is the value of T_g that we will assume in the remainder of our discussion.

Below T_g the low-frequency DMA results from Baij *et al.* are best suited for determination of the shift factors, whereas the QCM results are needed to obtain shift factors at higher temperatures. For temperatures between -35 °C and T_g the shift factors obey an Arrhenius temperature dependence, with a_T given by the following:

$$
\ln a_T = \frac{E_a}{RT} - \frac{E_a}{RT_{ref}}\tag{6}
$$

Here R is the gas constant, T is the absolute temperature, E_a is the activation energy and T_{ref} is an arbitrarily defined reference temperature, for which we use T_g (-7.4 °C). The dashed curve in Figure 5b is a representation of Eq. 6, with $E_a = 300 \text{ J/mol}$. This Arrhenius temperature dependence for segmental relaxation processes below T_g , and the specific value of E_a that we obtain, is consistent with results obtained for other synthetic polymers⁵³.

At temperatures above T_g the temperature shift factors obey the Vogel-Fulcher-Tamaan (VFT) form:

$$
\ln a_T = \frac{B}{T - T_{\infty}} - \frac{B}{T_{ref} - T_{\infty}} \tag{7}
$$

We handle the crossover between the Arrhenius and VFT forms at $T = T_g$ by requiring that d ln a_T /dT be continuous at T_g , so that B is related to E_a in the following manner:

$$
B = \frac{E_a}{R} \left(\frac{T_g - T_\infty}{T_g} \right)^2 \tag{8}
$$

We use T_{∞} as the remaining adjustable parameter to force agreement between high frequency QCM data and the low-frequency DMA data. As shown in Figure 5b, excellent agreement between the DMA and QCM data are obtained if $T_{\infty} = -65$ °C, which from Eq. 8 gives B=1700K. This value for T_{∞} is consistent with results obtained for a wide variety of polymers, where T_{∞} is found to be about 50K below $T_g^{\,53}$. Overall, we find that by combining results obtained with traditional dynamic mechanical analysis with our QCM results, an accurate set of shift factors are obtained from -45 ◦C to 85 ◦C.

While the analysis presented above does not necessarily prove the similarity of properties for linseed oil aged under very different conditions, the internal consistency of the analysis is compelling evidence for the similarity of the properties in the glass transition regime ($G > 10^8$ Pa) where data are obtained from both the DMA and QCM experiments. However, it is important to keep in mind that this similarity is restricted to the glassy and near-glassy regimes. Because of limitations imposed by the very high frequency, the QCM is only able to access the material behavior in the behavior of cured linseed oil in the regime where the modulus is greater than 10^8 Pa. The detailed degree of crosslinking and overall network topology affect the relaxed, rubbery modulus at values of fa_T that are too low to be accessed by the QCM, even at elevated temperatures. While the glassy dynamics of the different linseed oil samples cured at very different temperatures appear to be quite similar, it is entirely possible that the network architectures of the two samples differ in more substantial ways.

Conclusions

The quartz crystal microbalance can be utilized to track curing and aging processes of thin films as they are exposed to the ambient environment. In this work we have demonstrated the capability with linseed oil films that are of significant interest to the conservation science community. Two types of analyses were performed. The 'bulk' analysis is the most easily employed, and requires that accurate measurements be obtained for a single crystal resonance, in our case the third harmonic at a frequency of 15 MHz. This analysis can be utilized when the magnitude of the complex shear modulus at 15 MHz is less than about $2x10^7$ Pa, corresponding to a complex viscosity of about 200 mPa·s (200 times the viscosity of water). This limitation is a substantial one, and more complicated 'thin film' analysis must be utilized in most cases. Data from two different harmonics are generally required, and the film thickness must also be in a relatively narrow range, typically near 1 micrometer for the linseed oil films investigated here. By working with a film of the appropriate thickness and choosing the most appropriate harmonics for the analysis it is possible to obtain data from the uncured to fully cured states from the same film. A useful general guideline for the analysis is that the thickness of the film should be from 5% to 15% of the wavelength of the shear wave at the harmonics utilized, although useful data outside this range can sometimes be utilized. Also, when the film mass can be assumed to be constant (as with a temperature-dependent study, for example) data from a single harmonic is sufficient, providing access to a broader range of material properties. With regard to the linseed oil samples investigated here, we showed that the addition of PbO accelerates the cure kinetics by a factor of 5. Also, we showed that the glassy polymer dynamics obtained from the QCM are consistent with properties obtained from similar linseed oil films with traditional dynamic mechanical analysis at frequencies from 0.1 Hz to 10 Hz. Addition of the QCM data obtained at 15 MHz to this previous dataset enables us to accurately measure the temperature dependence of the polymer relaxation times (quantified by the temperature shift factors a_T) to much higher temperatures than can be

accomplished with the DMA data alone.

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Supporting Information: QCM theory and sample regimes, Sample preparation and measurements, Comparison of QCM calculation methods, Ellipsometric measurements, Bulk rheological measurements, Solution checks and response maps for KP thin films, Analysis of temperature-dependent properties, Generalized power law model and T_q deterimation, Reflection Fourier transform infrared spectroscopy measurements.

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