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CONSTRUCTIVE KREIN-RUTMAN RESULT FOR KINETIC FOKKER-PLANCK EQUATIONS IN A DOMAIN

K. CARRAPATOSO, P. GABRIEL, R. MEDINA, AND S. MISCHLER

ABSTRACT. We consider a general Kinetic Fokker-Planck (KFP) equation in a domain with Maxwell reflection condition on the boundary, not necessarily with conservation of mass. We establish the wellposedness in many spaces including Radon measures spaces, and in particular the existence and uniqueness of fundamental solutions. We also establish a Krein-Rutman theorem with constructive rate of convergence in an abstract setting that we use for proving that the solutions to the KFP equation converge toward the conveniently normalized first eigenfunction. Both results use the ultracontractivity of the associated semigroup in a fundamental way.

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1. INTRODUCTION

1.1. The KFP equation in a domain. In this paper, we consider the Kinetic Fokker-Planck (KFP) equation (also denominated sometimes as Kolmogorov equation or ultraparabolic equation)

$$(1.1) \quad \partial_t f + v \cdot \nabla_x f = \Delta_v f + b \cdot \nabla_v f + cf \quad \text{on } \mathcal{U}$$

on the function $f := f(t, x, v)$ depending on the time variable $t \geq 0$, the position variable $x \in \Omega$, where $\Omega \subset \mathbb{R}^d$ is a suitably smooth bounded domain, $d \geq 1$, and the velocity variable $v \in \mathbb{R}^d$. For $T \in (0, +\infty]$, we use the shorthands $\mathcal{U} := (0, T) \times \mathcal{O}$, $\mathcal{O} := \Omega \times \mathbb{R}^d$. We assume that

$$(1.2) \quad b = b(x, v) \in \mathbb{R}^d, \quad c = c(x, v) \in \mathbb{R},$$

each of these functions being at least in $L_{\text{loc}}^\infty(\mathcal{O})$. We complement the above KFP evolution equation with the Maxwell type reflection condition on the boundary

$$(1.3) \quad \gamma_- f = \mathcal{R} \gamma_+ f = \iota_S \mathcal{S} \gamma_+ f + \iota_D \mathcal{D} \gamma_+ f \quad \text{on } \Gamma_-,$$

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and with an initial condition

$$(1.4) \quad f(0, x, v) = f_0(x, v) \quad \text{on } \mathcal{O}.$$

Here Γ_- denotes the incoming part of the boundary, \mathcal{S} denotes the specular reflection operator, \mathcal{D} denotes the diffusive reflection operator, ι_S and ι_D are nonnegative coefficients. More precisely, we assume that $\Omega := \{x \in \mathbb{R}^d; \delta(x) > 0\}$ for a $W^{2,\infty}(\mathbb{R}^d)$ function δ such that $\delta(x) := \text{dist}(x, \partial\Omega)$ on a neighborhood of the boundary set $\partial\Omega$ and thus $n_x = n(x) := -\nabla\delta(x)$ coincides with the unit normal outward vector field on $\partial\Omega$. We next define $\Sigma_\pm^x := \{v \in \mathbb{R}^d; \pm v \cdot n_x > 0\}$ the sets of outgoing (Σ_+^x) and incoming (Σ_-^x) velocities at point $x \in \partial\Omega$, then the sets

$$\Sigma_\pm := \{(x, v); x \in \partial\Omega, v \in \Sigma_\pm^x\}, \quad \Gamma_\pm := (0, T) \times \Sigma_\pm,$$

and finally the outgoing and incoming trace functions $\gamma_\pm f := \mathbf{1}_{\Gamma_\pm} \gamma f$. The specular reflection operator \mathcal{S} is defined by

$$(1.5) \quad (\mathcal{S}g)(x, v) := g(x, \mathcal{V}_x v), \quad \mathcal{V}_x v := v - 2n(x)(n(x) \cdot v),$$

and the diffusive operator \mathcal{D} is defined by

$$(1.6) \quad (\mathcal{D}g)(x, v) := \mathcal{M}_x(v) \tilde{g}(x), \quad \tilde{g}(x) := \int_{\Sigma_+^x} g(x, w) n(x) \cdot w \, dw,$$

where \mathcal{M}_x stands for the Maxwellian function

$$(1.7) \quad \mathcal{M}_x(v) := (2\pi\Theta_x)^{-(d-1)/2} \exp(-|v|^2/(2\Theta_x)) > 0,$$

associated to the wall temperature Θ_x which is assumed to satisfy

$$(1.8) \quad \Theta_x \in W^{1,\infty}(\Omega), \quad 0 < \Theta_* \leq \Theta_x \leq \Theta^* < \infty.$$

It is worth observing that \mathcal{M}_x is conveniently normalized in such a way that $\widetilde{\mathcal{M}}_x = 1$. Denoting the accommodation coefficient $\iota := \iota_S + \iota_D$, we assume

$$\iota_S, \iota_D, \iota : \partial\Omega \rightarrow [0, 1].$$

Let us introduce some notations and then discuss some particular cases. In view of (1.1), we define the interior collisional operator

$$(1.9) \quad \mathcal{C}f := \Delta_v f + b \cdot \nabla_v f + cf$$

and next the (full) interior operator

$$(1.10) \quad \mathcal{L} := \mathcal{T} + \mathcal{C}, \quad \mathcal{T} := -v \cdot \nabla_x.$$

We name *microscopic or interior mass conservative* case, the case when

$$\mathcal{L}^* 1 = \mathcal{C}^* 1 = 0, \quad \text{or equivalently } c = \text{div } b,$$

and we name *macroscopic or boundary mass conservative* case, the case when

$$\mathcal{R}^* 1 = 1, \quad \text{or equivalently } \iota \equiv 1.$$

Here and below, the operators \mathcal{C}^* , \mathcal{L}^* and \mathcal{R}^* denote the (formal) dual operators. It is worth emphasizing that we always have $\mathcal{R}^* 1 \leq 1$ from the very definition (1.3), (1.5), (1.6) and the assumption $\iota \leq 1$, so that mass is never added from the boundary, it is only (possibly partially) returned. The boundary condition (1.3) corresponds to the *pure specular reflection* boundary condition when $\iota = \iota_S \equiv 1$ and it corresponds to the *pure diffusive* boundary condition when $\iota = \iota_D \equiv 1$. When both mass conservation conditions are fulfilled then equation (1.1)–(1.3) is mass conservative, meaning that any solution (at least formally) satisfies

$$\int_{\mathcal{O}} f(t, x, v) \, dx dv = \int_{\mathcal{O}} f_0(x, v) \, dx dv, \quad \forall t \geq 0.$$

We name *equilibrium or detailed balance condition* case when the maxwellian \mathcal{M} with constant temperature is a stationary state for each operator separately, namely

$$\mathcal{L}\mathcal{M} = \mathcal{C}\mathcal{M} = 0, \quad \mathcal{R}\mathcal{M} = \mathcal{M}.$$

When $\Theta \equiv 1$, that corresponds to the situation when $\iota \equiv 1$ (ι_S and ι_D can be space dependent) and \mathcal{C} is the usual harmonic Fokker-Planck operator with $b(x, v) = v$, $c(x, v) = d$, that is

$$\mathcal{C}f := \Delta f + \operatorname{div}(vf).$$

This very specific but physically motivated situation has been studied in the recent paper [19] where, in particular, a constructive exponential stability result is established.

On the other hand, in the general situation when at least one of the two above conservations fails, we rather look for an eigentriplet (λ_1, f_1, ϕ_1) satisfying

$$(1.11) \quad \lambda_1 \in \mathbb{R}, \quad \mathcal{L}f_1 = \lambda_1 f_1, \quad \gamma_- f_1 = \mathcal{R}\gamma_+ f_1, \quad \mathcal{L}^* \phi_1 = \lambda_1 \phi_1, \quad \gamma_+ \phi_1 = \mathcal{R}^* \gamma_- \phi_1.$$

This issue has been tackled recently in [28] with the restriction $\Theta_x = \Theta$ is a constant, where the existence and uniqueness of such eigenelements have been established as well a non-constructive exponential asymptotical stability of the associated eigenfunction $F = e^{\lambda_1 t} f_1$.

We refer to [40, 39, 55, 24, 13, 20, 12, 50, 60, 61] for a general discussion and mathematical analysis of the kinetic Fokker-Planck equation set in the whole space or in a domain and to related problems.

In the present paper, we carry on the analysis made in [28, 19] by establishing the following results.

(1) We prove the existence and uniqueness of solutions in many weighted Lebesgue spaces by establishing dissipativity estimates on the associated operator and next growth bound on the corresponding semigroup. We also establish the existence and uniqueness of a fundamental solution.

(2) We establish the ultracontractivity of the above semigroup associated to the evolution problem (1.1)–(1.3)–(1.4), that is some immediate gain of stronger Lebesgue integrability and even immediate gain of uniform bound.

(3) We prove a constructive version of a Krein-Rutman-Doblin-Harris theorem providing existence, uniqueness and exponential asymptotic stability with constructive rate of the first eigentriplet for a general class of positive semigroup in an abstract framework.

(4) We show that the KFP model addressed here satisfies the requirement of the above Krein-Rutman-Doblin-Harris theorem and thus give a clear and constructive understanding of the large time behavior of the solutions.

These results generalize or make more accurate some previous similar known results.

1.2. Confinement in the velocity variable and admissible weight functions. We introduce additional assumptions on b and c in order that the interior collisional operator \mathcal{C} provides a convenient velocity confinement mechanism. We first assume

$$(1.12) \quad \liminf_{|v| \rightarrow \infty} \inf_{\Omega} b \cdot \frac{v}{|v|} = +\infty, \quad \frac{|b|}{\langle v \rangle}, \operatorname{div}_v b, c = \mathcal{O}\left(b \cdot \frac{v}{|v|^2}\right).$$

These two conditions are fundamental in order that first the interior collisional operator \mathcal{C} and next the full operator are dissipative and even have a discrete spectrum in convenient functional spaces. In order to identify these spaces and to make the discussion simpler, we make more precise the confinement conditions by assuming that there exist $R_0, b_0, b_1 > 0$, $\gamma > 1$ and for any $p \in [1, \infty]$ there exists $k_p^* \geq 0$ such that

$$(1.13) \quad \forall x \in \Omega, v \in B_{R_0}^c, \quad b_0 |v|^\gamma \leq b \cdot v \leq b_1 |v|^\gamma, \quad c - \frac{1}{p} \operatorname{div}_v b \leq k_p^* b \cdot \frac{v}{|v|^2}.$$

We now introduce the class of the so-called *admissible* weight functions $\omega : \mathbb{R}^d \rightarrow (0, \infty)$ we will work with, which will be either a polynomial weight function

$$(1.14) \quad \omega = \langle v \rangle^k := (1 + |v|^2)^{k/2}, \quad k > k_* := \max(k'_1, k_\infty^*),$$

$k'_1 := k_1 + d/2 + \max(1, \gamma/2 - 1)$, $k_1 := \max(k_1^*, d + 2)$, and we set $s := 0$ in that case, or either an exponential weight function

$$(1.15) \quad \omega = \exp(\zeta \langle v \rangle^s),$$

with the restrictions

$$(1.16) \quad \begin{aligned} s < \min(\gamma, 2), \quad \zeta > 0; \quad s = \gamma < 2, \quad \zeta \in (0, b_0/2); \\ s = \gamma = 2, \quad \zeta \in (0, \min(1/\Theta^*, b_0)/2); \quad s = 2 < \gamma, \quad \zeta \in (0, 1/(2\Theta^*)). \end{aligned}$$

In order to explain that choice of weight functions, we introduce the function

$$(1.17) \quad \varpi = \varpi_{\varphi, p}^{\mathcal{C}}(x, v) := 2 \left(1 - \frac{1}{p}\right) \frac{|\nabla_v \varphi|^2}{\varphi^2} + \left(\frac{2}{p} - 1\right) \frac{\Delta_v \varphi}{\varphi} - b \cdot \frac{\nabla_v \varphi}{\varphi} + c - \frac{1}{p} \operatorname{div}_v b,$$

which is the key quantity in order to reveal the velocity confinement mechanism. We may notice that for $\omega := \langle v \rangle^\ell e^{\zeta |v|^s}$, with $\ell \in \mathbb{R}$ and $s, \zeta \geq 0$, and because of the second condition in (1.12), we have

$$(1.18) \quad \varpi_{\omega, p}^{\mathcal{C}} \underset{|v| \rightarrow \infty}{\sim} (s\zeta)^2 |v|^{2s-2} - s\zeta b \cdot v |v|^{s-2} \quad \text{if } s > 0,$$

$$(1.19) \quad \varpi_{\omega, p}^{\mathcal{C}} \underset{|v| \rightarrow \infty}{\sim} c - \frac{1}{p} \operatorname{div}_v b - \ell b \cdot v |v|^{-2} \quad \text{if } s = 0.$$

As a consequence, whatever is the value $\gamma > 1$, we have $(\varpi_{\omega, p}^{\mathcal{C}})_+ \in L^\infty(\mathcal{O})$ for any admissible weight function, what is the key information in order to establish a growth estimate in the corresponding weighted L^p space. Moreover, we have

$$(1.20) \quad \limsup_{|v| \rightarrow \infty} \sup_{\Omega} \varpi_{\omega, p}^{\mathcal{C}} = -\infty$$

for any admissible weight function when $\gamma > 2$ and for any exponential weight function with exponent $s \in (2 - \gamma, \gamma]$ when $\gamma \in (1, 2]$, what gives a key information on the spectrum of \mathcal{C} in the corresponding functional space. We will call *strongly confining* any admissible weight function satisfying (1.20). For further references, we also notice that

$$(1.21) \quad \sup_{\Omega} \mathcal{M}\omega \langle v \rangle \in (L^1 \cap L^\infty)(\mathbb{R}^d), \quad \omega^{-1} \langle v \rangle \in (L^1 \cap L^2)(\mathbb{R}^d),$$

for any admissible weight function because of the restrictions $k_* \geq d + 1$, $s \leq 2$ and $\zeta < 1/(2\Theta^*)$ when $s = 2$. The bound (1.21) provides the compatibility of the weight function with the boundary condition.

1.3. The main results. In order to state our main results, we need to introduce some functional spaces. For a given measure space (E, \mathcal{E}, μ) , a weight function $\rho : E \rightarrow (0, \infty)$ and an exponent $p \in [1, \infty]$, we define the weighted Lebesgue space L_ρ^p associated to the norm

$$\|g\|_{L_\rho^p} = \|\rho g\|_{L^p}.$$

We also define $M_{\omega, 0}^1$ as the space of Radon measures g on \mathcal{O} with vanishing mass at the boundary, that is such that $|g\omega|(\mathcal{O} \setminus \mathcal{O}_\varepsilon) \rightarrow 0$ as $\varepsilon \rightarrow 0$, where, for any $\varepsilon > 0$,

$$(1.22) \quad \Omega_\varepsilon := \{x \in \Omega \cap B_{\varepsilon^{-1}}, \delta(x) > \varepsilon\}, \quad \mathcal{O}_\varepsilon := \Omega_\varepsilon \times B_{\varepsilon^{-1}}.$$

We first state a general existence and uniqueness result for the kinetic Fokker-Planck equation (1.1), (1.3), (1.4).

Theorem 1.1 (Existence and uniqueness). *We make the above assumptions on Ω , Θ , b and c , in particular (1.2), (1.8), (1.12) and (1.13) hold. For any admissible weight function ω and any initial datum $f_0 \in L_\omega^p$, $p \in [1, \infty]$, or $f_0 \in M_{\omega,0}^1$, there exists a unique global weak solution f to the kinetic Fokker-Planck equation (1.1)–(1.3)–(1.4). In particular, for any $(x_0, v_0) \in \mathcal{O}$, there exists a unique fundamental solution associated to the initial datum $f_0 := \delta_{(x_0, v_0)}$.*

The precise sense of solution will be given in Proposition 3.3 (see also Theorem 3.5) in a L^2 framework, in Theorem 5.2 in a general L^p framework, and in Theorem 5.3 in a Radon measures framework. This result extends the existence and unique result of [28, Theorem 11.5] stated in a more restrictive L^2 framework (see also [24, 1] for further previous results). The L^2 framework is mainly based on Lions' variant of the Lax-Milgram theorem [43, Chap III, §1], as used in [24, 1], a trace theory developed in [49, 50, 28, 18] and boundary estimates in the spirit of [3, 50, 8]. The growth estimate is obtained by cooking up a modified but equivalent weight function for which the dissipativity of the full operator can be established. On the other hand, the general Lebesgue framework and the Radon measures framework are more involved and are also based on the ultracontractivity theorem below as well as some arguments adapted from the parabolic equation as developed in [11, 9, 10]. It is worth mentioning that the well-posedness and some regularity issues for the KFP equation set in the torus have been obtained in [1]. For the whole space setting, we refer to the recent works [5, 4] and the references therein. Finally, the KFP equation in a bounded domain has been considered in [61, 59, 56].

We next consider the first eigenvalue problem and the longtime behavior providing a quantitative answer to the first eigenelements issue.

Theorem 1.2 (Long time asymptotic). *Under the assumptions of Theorem 1.1, there exist two weight functions ω_1, m_1 and an exponent $r > 2$ with $L_{\omega_1}^r \subset (L_{m_1}^2)'$ such that there exists a unique eigentriplet $(\lambda_1, f_1, \phi_1) \in \mathbb{R} \times L_{\omega_1}^r \times L_{m_1}^2$ satisfying the first eigenproblem (1.11) together with the normalization condition $\|\phi_1\|_{L_{m_1}^2} = 1$, $\langle \phi_1, f_1 \rangle_{L_{m_1}^2, (L_{m_1}^2)'} = 1$. These eigenfunctions are continuous functions and they also satisfy*

$$(1.23) \quad 0 < f_1 \lesssim \omega^{-1}, \quad 0 < \phi_1 \lesssim \omega \quad \text{on } \mathcal{O},$$

for any admissible weight function ω . Furthermore, there exist some constructive constants $C \geq 1$ and $\lambda_2 < \lambda_1$ such that for any strongly confining admissible weight function ω , any exponent $p \in [1, \infty]$ and any initial datum $f_0 \in L_\omega^p$, the associated solution f to the kinetic Fokker-Planck equation (1.1), (1.3), (1.4) satisfies

$$(1.24) \quad \|f(t) - \langle f_0, \phi_1 \rangle f_1 e^{\lambda_1 t}\|_{L_\omega^p} \leq C e^{\lambda_2 t} \|f_0 - \langle f_0, \phi_1 \rangle f_1\|_{L_\omega^p},$$

for any $t \geq 0$.

This result improves the recent work [28, Section 11] (see in particular [28, Theorem 11.6], [28, Theorem 11.8] and [28, Theorem 11.11]) by slightly generalizing the framework to a position dependent wall temperature and by providing a fully constructive approach for the exponential stability of the first eigenfunction. We refer to the previous works [42, 33, 7] (partially based on [58, 38, 41]) where similar results are established for the same kind of equation in a bounded domain with no-flow boundary condition. We also emphasize that in the conservative case, many works have been done related to hypocoercivity and constructive rate of convergence to the steady state in [25, 34, 27, 37, 36, 60] or more recently in [26, 46, 17, 14, 1, 19]. From a technical point of view, this result is a consequence of the abstract version of the Krein-Rutman-Doblin-Harris theorem that will be presented in section 7.1 (and which is really in the spirit of the recent work [28]) together with the ultracontractivity property stated below and the Harnack estimates established in [30].

Both the above well-posedness and the longtime behavior results are based on the following ultracontractivity property.

Theorem 1.3 (Ultracontractivity). *There exist $\Theta, C > 0$ and $\kappa \geq 0$ such that any solution f to the KFP equation (1.1)–(1.3)–(1.4) satisfies*

$$(1.25) \quad \|f(T, \cdot)\|_{L^\infty} \leq C \frac{e^{\kappa T}}{T^\Theta} \|f_0\|_{L^1_\omega}, \quad \forall T > 0,$$

for any strongly confining admissible weight function ω .

This result slightly improves and generalizes [19, Theorem 1.1] which establishes a similar result in the conservative case. The proof is very alike the one of [19, Theorem 1.1] although some steps are slightly simplified. The strategy is based on Nash’s gain of integrability argument [54] which is performed however on a time integral inequality as in Moser’s work [53], and is then more convenient in order to use the interior gain of integrability deduced from Bouchut’s regularity result [12] following the way paved by [30, Theorem 6] for proving a somehow similar local version. Contrary to the last reference, the gain of integrability is not formulated locally in x, v and integrated in time but globally in x, v and pointwisely in time as in the the ultracontractivity theory of Davies and Simon [21, 22]. Exactly as in [19], the key argument consists in exhibiting a suitable twisted weight function which is somehow slightly more elaborated than the one used during the proof of the growth estimates in Theorem 1.1. In Theorem 1.1 and Theorem 1.3, the boundedness assumption on Ω is really only needed in the proof of the uniqueness of the solution in the L^p_ω framework and it is likely that it can be removed. Here, we do not try to generalize these results to the case of an unbounded domain, see however [15] for partial results in that direction.

1.4. Organization of the paper. Section 2 is dedicated to the proof of some weighted L^p a priori growth bounds for the primal and the dual problems. These estimates and the well-posedness in a L^2 framework for the same problems (and thus part of Theorem 1.1) are then established rigorously in Section 3. Section 4 is devoted to the proof of the ultracontractivity property as stated in Theorem 1.3. In section 5 we come back to the well-posedness in a general framework and we end the proof of Theorem 1.1. Section 6 is dedicated to the proof of the Harnack inequality associated to our equations. In Section 7 we state and prove a constructive version of the Krein-Rutman theorem and deduce Theorem 1.2.

2. WEIGHTED L^p A PRIORI GROWTH ESTIMATES

This section is devoted to the proof of some a priori growth estimates in weighted L^p spaces for solutions to the KFP equation (1.1)–(1.3)–(1.4) and its formal adjoint.

2.1. A priori estimates for the primal problem. We recall that for two functions $f, \varphi : \mathbb{R}^d \rightarrow \mathbb{R}_+$ and $p \in [1, \infty)$, we have

$$(2.1) \quad \int_{\mathbb{R}^d} (\mathcal{C}f) f^{p-1} \varphi^p dv = -\frac{4(p-1)}{p^2} \int_{\mathbb{R}^d} |\nabla_v (f\varphi)^{p/2}|^2 + \int f^p \varphi^p \varpi_{\omega,p}^{\mathcal{C}},$$

with $\varpi_{\omega,p}^{\mathcal{C}}$ defined in (1.17), what can be established by mere repeated integrations by part, see for instance [28, Lemma 7.7] and the references therein. From the definition of the admissible weight functions in Section 1.2 and for further references, we may observe that the large velocity asymptotic of $\varpi_{\omega,p}^{\mathcal{C}}$ is controlled by

$$(2.2) \quad \limsup_{|v| \rightarrow \infty} \left(\sup_{\Omega} \varpi_{\omega,p}^{\mathcal{C}} - \varpi_{\omega,p}^{\sharp} \right) \leq 0, \quad \varpi_{\omega,p}^{\sharp} := -b_0^{\sharp} \langle v \rangle^{\varsigma}, \quad \varsigma := \gamma + s - 2,$$

with $b_0^\sharp > 0$ given by

$$(2.3) \quad \begin{aligned} b_0^\sharp &:= (k - k_p)b_0 \quad \text{if } s = 0, \\ b_0^\sharp &:= b_0 s \zeta \quad \text{if } s \in (0, \gamma), \\ b_0^\sharp &:= b_0 s \zeta - (s \zeta)^2 \quad \text{if } s = \gamma. \end{aligned}$$

In a more quantitative way, for any $\vartheta \in (0, 1)$, there exists $\kappa', R' > 0$ such that

$$(2.4) \quad \sup_{\Omega} \varpi_{\omega, p}^{\mathcal{C}} \leq \kappa' \chi_{R'} + \chi_{R'}^c \vartheta \varpi_{\omega, p}^\sharp,$$

where $\chi_R(v) := \chi(|v|/R)$, $\chi \in C^2(\mathbb{R}_+)$, $\mathbf{1}_{[0,1]} \leq \chi \leq \mathbf{1}_{[0,2]}$, and $\chi_R^c := 1 - \chi_R$.

Lemma 2.1. *For any admissible weight function ω , there exist $\kappa \geq 0$ and $C \geq 1$ such that for both exponents $p = 1, 2$, any solution f to the KFP equation (1.1)–(1.3)–(1.4) satisfies, at least formally,*

$$(2.5) \quad \|f_t\|_{L_v^p} \leq C e^{\kappa t} \|f_0\|_{L_v^p}, \quad \forall t \geq 0.$$

The proof is based on moment estimates introduced in [18, Proposition 3.3] for the case $p = 2$ and in [19, Lemma 2.3] for the case $p = 1$, which are reminiscent of L^1 hypodissipativity techniques, see e.g. [48, 31, 8], and which are based on the usual multiplier used in order to control the diffusive reflection operator in previous works on the Boltzmann equation, see e.g. [3, 49, 50, 8]. For further references, we define the formal adjoints

$$(2.6) \quad \mathcal{L}^* := v \cdot \nabla_x + \mathcal{C}^*, \quad \mathcal{C}^* g := \Delta_v g - b \cdot \nabla_v g + (c - \operatorname{div} b)g.$$

Proof of Lemma 2.1. Consider $0 \leq f_0 \in L^p(\omega)$ and $f = f(t, x, v) \geq 0$ a solution to the Cauchy problem (1.1)–(1.3)–(1.4). We introduce the modified weight functions ω_A and $\tilde{\omega}$ defined by

$$(2.7) \quad \omega_A^p := \mathcal{M}_x^{1-p} \chi_A + \omega^p (1 - \chi_A), \quad \tilde{\omega}^p := \left(1 + \frac{1}{2} \frac{n_x \cdot v}{\langle v \rangle^4}\right) \omega_A^p,$$

with $A \geq 1$ to be chosen later, $\hat{v} := v/\langle v \rangle$ and $\tilde{v} := \hat{v}/\langle v \rangle$. It is worth emphasizing that

$$(2.8) \quad c_A^{-1} \omega \leq \frac{1}{2} \omega_A \leq \tilde{\omega} \leq \frac{3}{2} \omega_A \leq c_A \omega,$$

with $c_A \in (0, \infty)$. We then write

$$(2.9) \quad \frac{1}{p} \frac{d}{dt} \int_{\mathcal{O}} f^p \tilde{\omega}^p = \int_{\mathcal{O}} (\mathcal{C} f) f^{p-1} \tilde{\omega}^p + \frac{1}{p} \int_{\mathcal{O}} f^p \mathcal{T}^* \tilde{\omega}^p - \frac{1}{p} \int_{\Sigma} (\gamma f)^p \tilde{\omega}^p (n_x \cdot v),$$

and we estimate each term separately below.

Step 1. We first compute separately each contribution of the boundary term in (2.9), namely we write

$$- \int_{\Sigma} (\gamma f)^p \tilde{\omega}^p (n_x \cdot v) = B_1 + B_2$$

with

$$B_1 := - \int_{\Sigma} (\gamma f)^p \omega_A^p n_x \cdot v, \quad B_2 := - \frac{1}{2} \int_{\Sigma} (\gamma f)^p (n_x \cdot \hat{v})^2 \omega_A^p.$$

On the one hand, we have

$$\begin{aligned} B_1 &= - \int_{\Sigma_+} (\gamma_+ f)^p \omega_A^p |n_x \cdot v| + \int_{\Sigma_-} \{\iota_S \mathcal{S} \gamma_+ f + \iota_D \mathcal{D} \gamma_+ f\}^p \omega_A^p |n_x \cdot v| \\ &\leq - \int_{\Sigma_+} (\gamma_+ f)^p \omega_A^p (n_x \cdot v)_+ + \int_{\Sigma_-} \iota_S (\mathcal{S} \gamma_+ f)^p \omega_A^p (n_x \cdot v)_- + \int_{\Sigma_-} \iota_D (\widetilde{\gamma_+ f})^p \mathcal{M}_x^p \omega_A^p (n_x \cdot v)_- \\ &\leq - \int_{\Sigma_+} \iota_D (\gamma_+ f)^p \omega_A^p (n_x \cdot v)_+ + \int_{\partial\Omega} \iota_D (\widetilde{\gamma_+ f})^p K_1(\omega_A), \end{aligned}$$

where we have used the convexity of the mapping $s \mapsto s^p$ in the first line, we have made the change of variables $v \mapsto \mathcal{V}_x v$ in the second integral in the second line and we have set

$$(2.10) \quad K_1(\omega_A) := \int_{\mathbb{R}^d} \mathcal{M}_x^p \omega_A^p (n_x \cdot v)_- dv < \infty.$$

For $p = 1$, we observe that $\omega_A \geq 1$ and we set $K_2(\omega_A) := 1$. For $p = 2$, using the Cauchy-Schwarz inequality, we have

$$(\widetilde{\gamma+f})^2 \leq K_2(\omega_A)^{-1} \int_{\mathbb{R}^d} (\gamma+f)^2 \omega_A^2 (n_x \cdot v)_+,$$

with

$$K_2(\omega_A)^{-1} := \int_{\mathbb{R}^d} \omega_A^{-2} (n_x \cdot v)_+ dv < \infty.$$

In both case we deduce

$$(2.11) \quad B_1 \leq \int_{\partial\Omega} \iota_D (K_1(\omega_A) - K_2(\omega_A)) (\widetilde{\gamma+f})^p.$$

On the other hand, using the boundary condition (1.3) and making change of variables $v \mapsto \mathcal{V}_x v$, there holds

$$\begin{aligned} B_2 &= -\frac{1}{2} \int_{\Sigma_+} (\gamma+f)^p \omega_A^p (n_x \cdot \hat{v})^2 - \frac{1}{2} \int_{\Sigma_-} (\iota_S \mathcal{S} \gamma+f + \iota_D \mathcal{D} \gamma+f)^p \omega_A^p (n_x \cdot \hat{v})^2 \\ &= -\frac{1}{2} \int_{\Sigma_+} (\gamma+f)^p \omega_A^p (n_x \cdot \hat{v})^2 - \frac{1}{2} \int_{\Sigma_+} (\iota_S \gamma+f + \iota_D \mathcal{D} \gamma+f)^p \omega_A^p (n_x \cdot \hat{v})^2. \end{aligned}$$

When $p = 1$, denoting

$$(2.12) \quad K_0(\omega_A) := \frac{1}{2} \int_{\mathbb{R}^d} \mathcal{M}_x (n_x \cdot \hat{v})_+^2 \omega_A dv < \infty,$$

we therefore have

$$B_2 \leq -\frac{1}{2} \int_{\Sigma_+} \iota_D (\mathcal{D} \gamma+f) (n_x \cdot \hat{v})^2 \omega_A = - \int_{\partial\Omega} \iota_D K_0(\omega_A) (\widetilde{\gamma+f}).$$

On the other hand, when $p = 2$, denoting

$$(2.13) \quad K_0(\omega_A)^{-1} := 2 \int_{\Sigma_+^x} \langle v \rangle^2 \omega_A^{-2} dv < \infty,$$

and thanks to the Cauchy-Schwarz inequality, we have

$$B_2 \leq -\frac{1}{2} \int_{\Sigma_+} (\gamma+f)^2 (n_x \cdot \hat{v})^2 \omega_A^2 \leq - \int_{\partial\Omega} K_0(\omega_A) (\widetilde{\gamma+f})^2.$$

In both cases $p = 1$ and $p = 2$, we have established

$$(2.14) \quad B \leq \int_{\partial\Omega} \iota_D [K_1(\omega_A) - K_2(\omega_A) - K_0(\omega_A)] (\widetilde{\gamma+f})^p,$$

and we observe that

$$\lim_{A \rightarrow \infty} K_1(\omega_A) = \lim_{A \rightarrow \infty} K_2(\omega_A) = 1,$$

thanks to the dominated convergence theorem, the normalization condition on \mathcal{M}_x and the condition (1.21). We similarly have

$$\lim_{A \rightarrow \infty} K_0(\omega_A) = \frac{1}{2} \int_{\mathbb{R}^d} \mathcal{M}_x (n_x \cdot \hat{v})_+^2 dv \geq C_1(\Theta_*, \Theta^*) > 0,$$

when $p = 1$, and

$$\lim_{A \rightarrow \infty} K_0(\omega_A)^{-1} = 2 \int_{\Sigma_x^-} \mathcal{M}_x \langle v \rangle^2 dv \leq C_2(\Theta_*, \Theta^*) < \infty,$$

when $p = 2$. All these convergences being uniform in $x \in \partial\Omega$, we can choose $A > 0$ large enough in such a way that

$$K_1(\omega_A) - K_2(\omega_A) - K_0(\omega_A) \leq 0,$$

and thus

$$- \int_{\Sigma} (\gamma f)^p \tilde{\omega}^p (n_x \cdot v) \leq 0.$$

Step 2. We now deal with the first term at the right-hand side of (2.9). On the one hand from (2.1), we have

$$\int_{\mathbb{R}^d} (\mathcal{C}f) f^{p-1} \tilde{\omega}^p = - \frac{4(p-1)}{p} \int_{\mathbb{R}^d} |\nabla_v (f^{p/2} \tilde{\omega}^{p/2})|^2 + \int_{\mathbb{R}^d} f^p \tilde{\omega}^p \varpi_{\tilde{\omega}, p}^{\mathcal{C}},$$

with

$$\varpi_{\tilde{\omega}, p}^{\mathcal{C}} := 2 \left(1 - \frac{1}{p}\right) \frac{|\nabla_v \tilde{\omega}|^2}{\tilde{\omega}^2} + \left(\frac{2}{p} - 1\right) \frac{\Delta_v \tilde{\omega}}{\tilde{\omega}} - b \cdot \frac{\nabla_v \tilde{\omega}}{\tilde{\omega}} + c - \frac{1}{p} \operatorname{div}_v b.$$

Defining, $\wp^p := 1 + \frac{1}{2} \frac{n_x \cdot v}{\langle v \rangle^4}$ and $\wp_A^p := 1 + \chi_A (\mathcal{M}^{1-p} \omega^{-p} - 1)$ so that $\tilde{\omega} = \wp \omega_A$ and $\omega_A = \wp_A \omega$, we compute

$$(2.15) \quad \varpi_{\tilde{\omega}, p}^{\mathcal{C}} = \varpi_{A, p} + 2 \frac{\nabla_v \omega_A}{\omega_A} \cdot \frac{\nabla_v \wp}{\wp} + 2 \left(1 - \frac{1}{p}\right) \frac{|\nabla_v \wp|^2}{\wp^2} + \left(\frac{2}{p} - 1\right) \frac{\Delta_v \wp}{\wp} - b \cdot \frac{\nabla_v \wp}{\wp},$$

with

$$\varpi_{A, p} := 2 \left(1 - \frac{1}{p}\right) \frac{|\nabla_v \omega_A|^2}{\omega_A^2} + \left(\frac{2}{p} - 1\right) \frac{\Delta_v \omega_A}{\omega_A} - b \cdot \frac{\nabla_v \omega_A}{\omega_A} + c - \frac{1}{p} \operatorname{div}_v b,$$

and next

$$(2.16) \quad \varpi_{A, p} = \varpi_{\omega, p}^{\mathcal{C}} + 2 \frac{\nabla_v \omega}{\omega} \cdot \frac{\nabla_v \wp_A}{\wp_A} + 2 \left(1 - \frac{1}{p}\right) \frac{|\nabla_v \wp_A|^2}{\wp_A^2} + \left(\frac{2}{p} - 1\right) \frac{\Delta_v \wp_A}{\wp_A} - b \cdot \frac{\nabla_v \wp_A}{\wp_A}.$$

Because χ_A has compact support in the velocity variable, the same holds for all the terms except the first one at the right-hand side of (2.16), and thus

$$|\varpi_{A, p} - \varpi_{\omega, p}^{\mathcal{C}}| \lesssim \frac{1}{\langle v \rangle^4} \lesssim \frac{|\varpi_{\omega, p}^{\#}|}{\langle v \rangle^{2+\gamma+s}}.$$

Similarly, observing that

$$\frac{\nabla_v \omega_A}{\omega_A} \cdot \frac{\nabla_v \wp}{\wp} = \frac{\nabla_v \omega}{\omega} \cdot \frac{\nabla_v \wp}{\wp} + \frac{\nabla_v \wp_A}{\wp_A} \cdot \frac{\nabla_v \wp}{\wp},$$

where $|\frac{\nabla_v \wp}{\wp}| \lesssim \langle v \rangle^{-4}$ and the second term is compactly supported, we have

$$|\varpi_{\tilde{\omega}, p}^{\mathcal{C}} - \varpi_{A, p}| \lesssim \left(1 + |b| + \frac{|\nabla \omega|}{\omega}\right) \frac{1}{\langle v \rangle^4} \lesssim \frac{|\varpi_{\omega, p}^{\#}|}{\langle v \rangle^2}.$$

Combining the last two estimates together with (2.4), we deduce that for any $\vartheta \in (0, 1)$, there exists $\tilde{\kappa}, \tilde{R} > 0$ such that

$$(2.17) \quad \sup_{\Omega} \varpi_{\tilde{\omega}, p}^{\mathcal{C}} \leq \tilde{\kappa} \chi_{\tilde{R}} + \chi_{\tilde{R}}^c \vartheta \varpi_{\omega, p}^{\#}.$$

Step 3. We finally deal with the second term at the right-hand side of (2.9). When $p = 1$ we have

$$v \cdot \nabla_x \tilde{\omega} = \frac{1}{2} (\hat{v} \cdot D_x n_x \hat{v}) \frac{\omega_A}{\langle v \rangle^2} \lesssim \frac{1}{\langle v \rangle^2} \tilde{\omega} \lesssim \frac{|\varpi_{\omega,p}^\#|}{\langle v \rangle^{s+\gamma}} \tilde{\omega},$$

thanks to the regularity assumption on Ω . On the other hand, when $p = 2$, we first compute

$$v \cdot \nabla_x (\tilde{\omega}^2) = \frac{1}{2} (\hat{v} \cdot D_x n_x \hat{v}) \frac{\omega_A^2}{\langle v \rangle^2} + \left(1 + \frac{1}{2} \frac{n_x \cdot v}{\langle v \rangle^4}\right) v \cdot \nabla_x (\omega_A^2).$$

Since

$$\nabla_x (\omega_A^2) = \chi_A \mathcal{M}_x^{-1} \left[\frac{(d-1)}{2} \frac{\nabla_x \Theta_x}{\Theta_x} - \frac{|v|^2}{2} \frac{\nabla_x \Theta_x}{\Theta_x^2} \right],$$

assumption (1.8) together with the fact that χ_A is compactly supported and the regularity assumption on Ω as above imply

$$v \cdot \nabla_x (\tilde{\omega}^2) \lesssim \frac{1}{\langle v \rangle^2} \tilde{\omega}^2 \lesssim \frac{|\varpi_{\omega,p}^\#|}{\langle v \rangle^{s+\gamma}} \tilde{\omega}.$$

Step 4. Coming back to (2.9) and using Step 1, we deduce that

$$(2.18) \quad \frac{1}{p} \frac{d}{dt} \int_{\mathcal{O}} f^p \tilde{\omega}^p \leq -\frac{4(p-1)}{p} \int_{\mathcal{O}} |\nabla_v (f \tilde{\omega})^{p/2}|^2 + \int_{\mathcal{O}} f^p \tilde{\omega}^p \varpi_{\tilde{\omega},p}^{\mathcal{L}}$$

for both $p = 1, 2$ and with

$$(2.19) \quad \varpi_{\tilde{\omega},p}^{\mathcal{L}} := \varpi_{\tilde{\omega},p}^{\mathcal{C}} + \frac{1}{\tilde{\omega}^p} v \cdot \nabla_x \tilde{\omega}^p.$$

Gathering the estimates (2.4) and those established in Step 2 and Step 3, we deduce that for any $\vartheta \in (0, 1)$, there exists $\kappa, R > 0$ such that

$$(2.20) \quad \varpi_{\tilde{\omega},p}^{\mathcal{L}} \leq \kappa \chi_R + \chi_R^c \vartheta \varpi_{\tilde{\omega},p}^\#.$$

In particular, $\varpi_{\tilde{\omega},p}^{\mathcal{L}} \leq \kappa$ and we immediately conclude thanks to Grönwall's lemma and the comparison (2.8) between ω and $\tilde{\omega}$. \square

Remark 2.2. For later references, we emphasize that the same proof works for establishing (2.5) with $p := 1$ and $\omega := \langle v \rangle^k$, $k \geq k_1 := \max(k_1^*, d + 2)$.

2.2. A priori estimates for the dual problem. We establish now a similar exponential growth a priori estimate in a general weighted L^q framework, $q = 1, 2$, for the dual backward problem associated to (1.1)–(1.3)–(1.4). More precisely we consider the equation

$$(2.21) \quad \begin{cases} -\partial_t g = v \cdot \nabla_x g + \mathcal{C}^* g & \text{in } (0, T) \times \mathcal{O}, \\ \gamma_+ g = \mathcal{R}^* \gamma_- g & \text{on } (0, T) \times \Sigma_+, \\ g(T) = g_T & \text{in } \mathcal{O}, \end{cases}$$

for any $T \in (0, \infty)$ and any final datum g_T . The adjoint Fokker-Planck operator \mathcal{C}^* is defined in (2.6), and the adjoint reflection operator \mathcal{R}^* is defined by

$$\mathcal{R}^* g(x, v) = \iota_S \mathcal{S} g(x, v) + \iota_D \mathcal{D}^* g(x),$$

with

$$\mathcal{D}^* g(x) = \widetilde{\mathcal{M}}_x g(x) := \int_{\mathbb{R}^d} g(x, w) \mathcal{M}_x(w) (n_x \cdot w)_- dw,$$

for any function g with support on Σ_- .

For two solutions f to the forward Cauchy problem (1.1)–(1.3)–(1.4) and g to the dual problem (2.21), the usual identity

$$(2.22) \quad \int_{\mathcal{O}} f(T) g_T = \int_{\mathcal{O}} f_0 g(0)$$

then holds at least formally, see also Theorem 3.5 below. We may indeed formally compute

$$\begin{aligned}
\int_{\mathcal{O}} f(T)g_T &= \int_{\mathcal{O}} f_0g(0) + \int_0^T \int_{\mathcal{O}} (\partial_t f g + f \partial_t g) ds \\
&= \int_{\mathcal{O}} f_0g(0) - \int_0^T \int_{\mathcal{O}} (v \cdot \nabla_x f g + f v \cdot \nabla_x g) ds \\
&= \int_{\mathcal{O}} f_0g(0) - \int_0^T \int_{\Sigma} (v \cdot n) \gamma f \gamma g ds \\
&= \int_{\mathcal{O}} f_0g(0) - \int_0^T \int_{\Sigma_+} (v \cdot n) (\gamma_+ f) (\mathcal{R}^* \gamma_- g) ds \\
&\quad + \int_0^T \int_{\Sigma_-} |v \cdot n| (\mathcal{R} \gamma_+ f) (\gamma_- g) ds,
\end{aligned}$$

by using the Green-Ostrogradski formula and the reflection conditions at the boundary. From the very definition of \mathcal{R} and \mathcal{R}^* , we then deduce (2.22).

Lemma 2.3. *For any admissible weight function ω and any exponent $q = 1$ or $q = 2$, there exist $\kappa \in \mathbb{R}$ and $C \geq 1$ such that for any $T > 0$ and any $g_T \in L_m^q$ with $m := \omega^{-1}$, the associated solution g to the backwards dual problem (2.21) satisfies*

$$(2.23) \quad \|g(0)\|_{L_m^q} \leq C e^{\kappa T} \|g_T\|_{L_m^q}.$$

Proof of Lemma 2.3. Without loss of generality we may suppose that $m \geq \mathcal{M}_x$, otherwise we replace m by cm where $c > 0$ is such that $m \geq c^{-1} \mathcal{M}_x$.

Consider a final time $T \in (0, \infty)$, a final datum $0 \leq g_T \in L_m^q$ and $g = g(t, x, v) \geq 0$ a solution to the backward dual Cauchy problem (2.21). For $A \geq 1$, we introduce the weight functions

$$(2.24) \quad m_A^q := \chi_A \mathcal{M}_x + (1 - \chi_A) m^q, \quad \tilde{m}^q := \left(1 - \frac{1}{2} \frac{n_x \cdot v}{\langle v \rangle^4}\right) m_A^q,$$

with the notations introduced in (2.4). It is worth emphasizing that

$$(2.25) \quad \mathcal{M}_x \leq m_A \leq m \quad \text{and} \quad c_A^{-1} m \leq \frac{1}{2} m_A \leq \tilde{m} \leq \frac{3}{2} m_A \leq \frac{3}{2} m,$$

with $c_A \in (0, \infty)$. Similarly as in the proof of Lemma 2.1, we compute

$$(2.26) \quad -\frac{1}{q} \frac{d}{dt} \int_{\mathcal{O}} g^q \tilde{m} = \int_{\mathcal{O}} g^{q-1} (\mathcal{C}^* g) \tilde{m}^q + \frac{1}{q} \int_{\Sigma} (\gamma g)^q \tilde{m}^q (n_x \cdot v) - \frac{1}{q} \int_{\mathcal{O}} g^q (v \cdot \nabla_x \tilde{m}^q),$$

and we estimate each term separately.

Step 1. In order to estimate the boundary term in (2.26), we split it into

$$\int_{\Sigma} (\gamma g)^q \tilde{m}^q (n_x \cdot v) = B_1 - B_2$$

with

$$B_1 = \int_{\Sigma} (\gamma g)^q m_A^q (n_x \cdot v) \quad \text{and} \quad B_2 = \frac{1}{2} \int_{\Sigma} (\gamma g)^q m_A^q (n_x \cdot \hat{v})^2.$$

For the first term, we have

$$\begin{aligned}
B_1 &= \int_{\Sigma_+} (\iota_S \mathcal{S} \gamma_- g + \iota_D \mathcal{D}^* \gamma_- g)^q m_A^q (n_x \cdot v)_+ - \int_{\Sigma_-} (\gamma_- g)^q m_A^q (n_x \cdot v)_- \\
&\leq \int_{\Sigma_+} \iota_S (\mathcal{S} \gamma_- g)^q m_A^q (n_x \cdot v)_+ + \int_{\Sigma_+} \iota_D (\mathcal{D}^* \gamma_- g)^q m_A^q (n_x \cdot v)_+ - \int_{\Sigma_-} (\gamma_- g)^q m_A^q (n_x \cdot v)_- \\
&\leq - \int_{\Sigma_-} \iota_D (\gamma_- g)^q m_A^q (n_x \cdot v)_- + \int_{\Sigma_+} \iota_D \left(\widetilde{\mathcal{M}_x \gamma_- g} \right)^q m_A^q (n_x \cdot v)_+,
\end{aligned}$$

where we have used the boundary condition in (2.21) on the first line, the convexity of the mapping $s \mapsto s^q$ on the second line, and the change of variables $v \rightarrow \mathcal{V}_x v$ on the third one. Defining

$$K_1(m_A) := \int_{\mathbb{R}^d} m_A^q (n_x \cdot v)_+ dv < \infty,$$

we equivalently have

$$B_1 \leq - \int_{\Sigma_-} \iota_D (\gamma - g)^q m_A^q (n_x \cdot v)_- + \int_{\partial\Omega} \iota_D K_1(m_A) (\widetilde{\mathcal{M}_x \gamma - g})^q.$$

When $q = 1$, we set $K_2(m_A) := 1$ and use the fact $m_A \geq \mathcal{M}_x$ in order to obtain

$$B_1 \leq \int_{\partial\Omega} \iota_D \{K_1(m_A) - 1\} \widetilde{\mathcal{M}_x \gamma - g}.$$

On the other hand, when $q = 2$, the Cauchy-Schwarz inequality yields

$$(\widetilde{\mathcal{M}_x \gamma - g})^2 \leq K_2(m_A)^{-1} \int_{\mathbb{R}^d} (\gamma - g)^2 m_A^2 (n_x \cdot v)_- dv,$$

where we have set

$$K_2(m_A)^{-1} := \int_{\mathbb{R}^d} \mathcal{M}_x^2 m_A^{-2} (n_x \cdot v)_- dv < \infty,$$

and we thus obtain

$$B_1 \leq \int_{\partial\Omega} \iota_D \{K_1(m_A) - K_2(m_A)\} (\widetilde{\mathcal{M}_x \gamma - g})^2.$$

We now deal with the second term B_2 , observing first that

$$B_2 = \frac{1}{2} \int_{\Sigma_-} (\gamma - g)^q m_A^q (n_x \cdot \hat{v})^2 + \frac{1}{2} \int_{\Sigma_-} \{\iota_S \gamma - g + \iota_D \mathcal{D}^* \gamma - g\}^q m_A^q (n_x \cdot \hat{v})^2,$$

where we have used the boundary condition in (2.21) and the change of variables $v \rightarrow \mathcal{V}_x v$. If $q = 1$, we through away the two first terms and we have

$$B_2 \geq \frac{1}{2} \int_{\Sigma_-} \iota_D (\widetilde{\mathcal{M}_x \gamma - g}) m_A (n_x \cdot \hat{v})^2 \geq \frac{1}{2} \int_{\partial\Omega} \iota_D K_0(m_A) (\widetilde{\mathcal{M}_x \gamma - g}),$$

where we have set

$$K_0(m_A) := \int_{\mathbb{R}^d} m_A (n_x \cdot \hat{v})_+^2 dv < \infty.$$

Otherwise if $q = 2$, we through away the last integral and using the Cauchy-Schwarz inequality and $\iota_D \leq 1$, we obtain

$$B_2 \geq \frac{1}{2} \int_{\Sigma_-} (\gamma - g)^q m_A^q (n_x \cdot \hat{v})^2 \geq \frac{1}{2} \int_{\partial\Omega} \iota_D K_0(m_A) (\widetilde{\mathcal{M}_x \gamma - g})^2,$$

where we have set

$$K_0(m_A)^{-1} := \int_{\Sigma_x^+} \mathcal{M}_x^2 m_A^{-2} \langle v \rangle^2 dv < \infty.$$

In both cases, gathering previous estimates yields

$$(2.27) \quad \int_{\Sigma} (\gamma g)^q \widetilde{m}^q (n_x \cdot v) \leq \int_{\partial\Omega} \iota_D \left\{ K_1(m_A) - K_2(m_A) - \frac{1}{2} K_0(m_A) \right\} (\widetilde{\mathcal{M}_x \gamma - g})^q.$$

We observe that in both cases $q = 1, 2$, we have

$$\lim_{A \rightarrow \infty} K_1(m_A) = \lim_{A \rightarrow \infty} K_2(m_A) = 1,$$

thanks to the dominated convergence theorem, the normalization condition on \mathcal{M}_x and the condition (1.21). We similarly have

$$\lim_{A \rightarrow \infty} K_0(m_A) = \int_{\mathbb{R}^d} \mathcal{M}_x(n_x \cdot \hat{v})_+^2 dv \geq C_1(\Theta_*, \Theta^*) > 0,$$

when $q = 1$, and

$$\lim_{A \rightarrow \infty} K_0(m_A)^{-1} = \int_{\Sigma_x^-} \mathcal{M}_x \langle v \rangle^2 dv \leq C_2(\Theta_*, \Theta^*) < \infty,$$

when $q = 2$. All these convergences being uniform in $x \in \partial\Omega$, we can choose $A > 0$ large enough in such a way that

$$K_1(m_A) - K_2(m_A) - \frac{1}{2}K_0(m_A) \leq 0,$$

which implies

$$(2.28) \quad \int_{\Sigma} (\gamma g)^q \tilde{m}^q (n_x \cdot v) \leq 0.$$

Step 2. We now estimate the first term at the right-hand side of (2.26). First of all, from (2.1)-(1.17), we have

$$\int_{\mathbb{R}^d} (\mathcal{E}^* g) g^{q-1} \tilde{m}^q = -\frac{4(q-1)}{q} \int_{\mathbb{R}^d} |\nabla_v (g^{q/2} \tilde{m}^{q/2})|^2 + \int g^q \tilde{m}^q \varpi_{\tilde{m},q}^{\mathcal{E}^*},$$

with

$$\varpi_{\varphi,q}^{\mathcal{E}^*} = 2 \left(1 - \frac{1}{q}\right) \frac{|\nabla_v \varphi|^2}{\varphi^2} + \left(\frac{2}{q} - 1\right) \frac{\Delta_v \varphi}{\varphi} + b \cdot \frac{\nabla_v \varphi}{\varphi} + c + \left(\frac{1}{q} - 1\right) \operatorname{div}_v b.$$

Arguing exactly as in Step 2 of the proof of Lemma 2.1, we can write

$$\varpi_{\tilde{m},q}^{\mathcal{E}^*} = \varpi_{\tilde{m},q}^{\mathcal{E}^*} + \mathfrak{W},$$

with

$$\varpi_{\tilde{m},q}^{\mathcal{E}^*} = \varpi_{\omega,p}^{\mathcal{E}^*}, \quad \frac{1}{p} + \frac{1}{q} = 1, \quad \mathfrak{W} = o(|\varpi_{\omega,p}^{\mathcal{E}^*}|).$$

Step 3. We finally deal with the second term at the right-hand side of (2.26). We compute

$$v \cdot \nabla_x (\tilde{m}^q) = \left(1 - \frac{1}{2} \frac{n_x \cdot v}{\langle v \rangle^4}\right) v \cdot \nabla_x (m_A^q) - \frac{1}{2} (\hat{v} \cdot D_x n_x \hat{v}) \frac{m_A^q}{\langle v \rangle^2}$$

and observe that

$$\nabla_x (m_A^q) = \chi_A \mathcal{M}_x \left[\frac{|v|^2}{2} \frac{\nabla_x \Theta_x}{\Theta_x^2} - \frac{(d-1)}{2} \frac{\nabla_x \Theta_x}{\Theta_x} \right].$$

Hence assumption (1.8) together with the fact that χ_A is compactly supported and the regularity assumption on Ω as above imply

$$\frac{1}{q} v \cdot \nabla_x (\tilde{m}^q) \lesssim \frac{1}{\langle v \rangle^2} \tilde{m}^q \lesssim \frac{|\varpi^{\#}|_{\omega,p}}{\langle v \rangle^{s+\gamma}} \tilde{m}^q.$$

Step 4. Coming back to (2.26) and gathering previous estimates, we deduce that

$$(2.29) \quad -\frac{1}{q} \frac{d}{dt} \int_{\mathcal{O}} g^q \tilde{m}^q \leq -\frac{4(q-1)}{q} \int_{\mathcal{O}} |\nabla_v (g \tilde{m})^{q/2}|^2 + \int_{\mathcal{O}} g^q \tilde{m}^q \varpi_{\tilde{m},q}^{\mathcal{L}^*}$$

for both $q = 1, 2$ and with

$$(2.30) \quad \varpi_{\tilde{m},q}^{\mathcal{L}^*} := \varpi_{\tilde{m},q}^{\mathcal{E}^*} + \frac{1}{\tilde{m}^q} v \cdot \nabla_x \tilde{m}^q.$$

Arguing as in the end of the proof of Lemma 2.1, we obtain that for any $\vartheta \in (0, 1)$, there exists $\kappa, R > 0$ such that

$$(2.31) \quad \varpi_{\tilde{m}, q}^{\mathcal{L}^*} \leq \kappa \chi_R + \chi_R^c \vartheta \varpi_{\omega, p}^{\sharp},$$

in particular, $\varpi_{\tilde{m}, q}^{\mathcal{L}^*} \leq \kappa$, and we immediately conclude thanks to Grönwall's lemma and the fact that $m \lesssim \tilde{m} \lesssim m$. \square

3. WELL-POSEDNESS IN A WEIGHTED L^2 FRAMEWORK

We briefly discuss the well-posedness in a weighted L^2 framework for both the primal and the dual Cauchy problems, using some material developed in [28, 18].

3.1. Trace results in a L^2 framework. In this section, we consider the kinetic Fokker-Planck equation

$$(3.1) \quad \partial_t g + v \cdot \nabla_x g = \mathbf{L}g + G, \quad \mathbf{L}g := \Delta_v g + b_i \partial_{v_i} g + \eta g$$

for a vector field $b = b(x, v)$, a function $\eta = \eta(x, v)$ and a source term $G = G(t, x, v)$. We formulate some trace results for solutions to the Vlasov-Fokker-Planck equation developed in [28, Sec. 11] and [18, Sec. 2] (see also [50, Section 4.1]) and which are mainly a consequence of the two following facts:

- (1) If $g \in L_{tx}^2 H_v^1$ is a weak solution to (3.1), then it is a renormalized solution;
- (2) If $g \in L_{txv}^\infty$ and $\nabla_v g \in L_{txv}^2$ is a weak solution to (3.1), then it admits a trace $\gamma g \in L^\infty$ in a renormalized sense.

We introduce some notations. We denote

$$(3.2) \quad d\xi_1 := |n_x \cdot v| dv d\sigma_x \quad \text{and} \quad d\xi_2 := (n_x \cdot \hat{v})^2 dv d\sigma_x$$

the measures on the boundary set Σ . We denote by \mathcal{B}_1 the class of renormalized functions $\beta \in W_{\text{loc}}^{2, \infty}(\mathbb{R})$ such that β'' has a compact support, by \mathcal{B}_2 the class of functions $\beta \in W_{\text{loc}}^{2, \infty}(\mathbb{R})$ such that $\beta'' \in L^\infty(\mathbb{R})$ and by $\mathcal{D}_0(\bar{\mathcal{U}})$ the space of test functions $\varphi \in \mathcal{D}(\bar{\mathcal{U}})$ such that $\varphi = 0$ on Γ_0 . We finally define the operators

$$\mathbf{M}_0^* \varphi := -\partial_t \varphi - v \cdot \nabla_x \varphi, \quad \mathbf{M}_i^* \varphi := \partial_{v_i} \varphi - b_i \varphi,$$

and we assume

$$(3.3) \quad b_i, \eta \in L_{tx}^\infty L_{\text{loc}, v}^\infty.$$

Theorem 3.1. *We consider $g \in L^2((0, T) \times \Omega; H_{\text{loc}}^1(\mathbb{R}^d))$, $G \in L_{tx}^2(H_{\text{loc}, v}^{-1})$, b_i, η satisfying (3.3) and we assume that g is a solution to the kinetic Fokker-Planck equation (3.1) in the sense of distribution $\mathcal{D}'(\mathcal{U})$.*

(1) *There exists $\gamma g \in L_{\text{loc}}^2(\Gamma; d\xi_2 dt)$, $g \in C([0, T]; L_{\text{loc}}^2(\mathcal{O}))$ and the following Green renormalized formula*

$$(3.4) \quad \int_{\mathcal{U}} (\beta(g) \mathbf{M}_0^* \varphi + \partial_{v_i} \beta(g) \mathbf{M}_i^* \varphi + \beta''(g) |\nabla_v g|^2 \varphi) dv dx dt \\ + \int_{\Gamma} \beta(\gamma g) \varphi n_x \cdot v dv d\sigma_x dt + \left[\int_{\mathcal{O}} (\beta(g) \varphi)(t, \cdot) dx dv \right]_0^T = \langle G + \eta g, \beta'(g) \varphi \rangle$$

holds for any renormalized function $\beta \in \mathcal{B}_1$ and any test functions $\varphi \in \mathcal{D}(\bar{\mathcal{U}})$, as well as for any renormalized function $\beta \in \mathcal{B}_2$ and any test functions $\varphi \in \mathcal{D}_0(\bar{\mathcal{U}})$. It is worth emphasizing that $\beta'(g) \varphi \in L_{tx}^2 H_v^1$ so that the duality product $\langle G, \beta'(g) \varphi \rangle$ is well defined.

(2) *If furthermore $g_0 \in L_{\text{loc}}^2(\bar{\mathcal{O}})$ and $\gamma_- g \in L_{\text{loc}}^2(\Gamma; d\xi_1 dt)$, then $g \in C([0, T]; L_{\text{loc}}^2(\bar{\mathcal{O}}))$, $\gamma_+ g \in L_{\text{loc}}^2(\Gamma; d\xi_1 dt)$ and (3.4) holds for any renormalized function $\beta \in \mathcal{B}_2$ and any test functions $\varphi \in \mathcal{D}(\bar{\mathcal{O}})$.*

(3) Alternatively to point (2), if furthermore $g \in L_{\text{loc}}^\infty(\bar{\mathcal{U}})$, then $\gamma g \in L_{\text{loc}}^\infty(\Gamma)$ and (3.4) holds for any renormalized function $\beta \in \mathcal{B}_2$ and any test functions $\varphi \in \mathcal{D}(\bar{\mathcal{O}})$.

We will also use the following stability result in the spirit of [50, Theorem 5.2].

Theorem 3.2. *Let us consider four sequences (g^k) , (b^k) , (η^k) and (G^k) and four functions g , b , η , G which all satisfy the requirements of Theorem 3.1.*

(1) *If $g^k \rightharpoonup g$ weakly in $L_{\text{loc}}^2(\bar{\mathcal{U}})$, $b^k \rightarrow b$ strongly in $L_{\text{loc}}^2(\bar{\mathcal{O}})$, $\eta^k \rightarrow \eta$ strongly in $L_{\text{loc}}^2(\bar{\mathcal{O}})$ and $G^k \rightharpoonup G$ weakly in $L_{\text{loc},x}^2 H_{\text{loc},v}^{-1}$, then g satisfies (3.1) so that it admits a trace function $\gamma g \in L_{\text{loc}}^2(\Gamma; d\xi_2 dt)$ and $\gamma g^k \rightharpoonup \gamma g$ weakly in $L_{\text{loc}}^2(\Gamma; d\xi_2 dt)$.*

(2) *If $g^k \rightarrow g$ strongly in $L_{\text{loc}}^2(\bar{\mathcal{U}})$, $b^k \rightharpoonup b$ weakly in $L_{\text{loc}}^2(\bar{\mathcal{O}})$, $\eta^k \rightharpoonup \eta$ weakly in $L_{\text{loc}}^2(\bar{\mathcal{O}})$ and $G^k \rightarrow G$ weakly in $L_{\text{loc},x}^2 H_{\text{loc},v}^{-1}$, then g satisfies (3.1) so that it admits a trace function $\gamma g \in L_{\text{loc}}^2(\Gamma; d\xi_2 dt)$ and $\gamma g^k \rightharpoonup \gamma g$ weakly in $L_{\text{loc}}^2(\Gamma; d\xi_2 dt)$.*

3.2. Well-posedness for the primal equation. For further reference, for an admissible weight function ω , we define the Hilbert norm $\|\cdot\|_{\mathcal{H}} = \|\cdot\|_{\mathcal{H}_\omega}$ by

$$\|f\|_{\mathcal{H}}^2 := \|f\|_{L_\omega^2}^2 + \|f\|_{H_\omega^{1,+}}^2, \quad \|f\|_{H_\omega^{1,+}}^2 := \int_{\mathcal{U}} \{|\nabla_v g|^2 + \langle \varpi_{\omega,2}^\# \rangle g^2\} \omega^2 dv dx dt,$$

with $\varpi_{\omega,2}^\#$ defined in (2.2), and we denote by $\mathcal{H} = \mathcal{H}_\omega$ the associated Hilbert space. We now state the well-posedness result for the primal problem which is nothing but [18, Theorem 2.12].

Proposition 3.3. *We make the regularity assumptions on Ω , Θ , b and c as presented in Section 1, in particular (1.2), (1.8), (1.12) and (1.13) hold. For any admissible weight function ω and any $f_0 \in L_\omega^2(\mathcal{O})$, there exists a unique global solution $f \in X_T := L^\infty(0, T; L_\omega^2) \cap C([0, T]; L_{\text{loc}}^2) \cap \mathcal{H}$, $\forall T > 0$, to the kinetic Fokker-Planck equation (1.1) complemented with the Maxwell reflection boundary condition (1.3) and associated to the initial datum f_0 . More precisely, the function f satisfies equation (3.1) in the sense of distributions in $\mathcal{D}'(\mathcal{U})$ with trace functions, defined thanks to Theorem 3.1, satisfying $\gamma f \in L_\omega^2(\Gamma, d\xi_2)$ as well as the Maxwell reflection boundary condition (1.3) pointwisely and $f(t, \cdot) \in L_\omega^2$, $\forall t \in [0, T]$, and the initial condition $f(0, \cdot) = f_0$ pointwisely.*

Because we will need to adapt it in the next section, we allude the proof and we refer to [18] for details.

Proof of Proposition 3.3. We split the proof into four steps.

Step 1. Given $f \in L_\omega^2(\Gamma_-; d\xi_1)$, we solve the inflow problem

$$(3.5) \quad \begin{cases} \partial_t f + v \cdot \nabla_x f = \mathcal{C}f & \text{in } (0, \infty) \times \mathcal{O} \\ \gamma_- f = f & \text{on } (0, \infty) \times \Sigma_- \\ f|_{t=0} = f_0 & \text{in } \mathcal{O}, \end{cases}$$

thanks to Lions' variant of the Lax-Milgram theorem [43, Chap III, §1]. More precisely, we define $\tilde{\omega}$ as during the proof of Lemma 2.1 and the bilinear form $\mathcal{E} : \mathcal{H} \times C_c^1([0, T] \times \mathcal{O} \cup \Gamma_-) \rightarrow \mathbb{R}$ by

$$\mathcal{E}(f, \varphi) := \int_{\mathcal{U}} (\lambda f - \mathcal{C}f) \varphi \tilde{\omega}^2 - \int_{\mathcal{U}} f (\partial_t + v \cdot \nabla_x) (\varphi \tilde{\omega}^2).$$

Using the Green-Ostrogradski formula, we observe that

$$\begin{aligned} \mathcal{E}(\varphi, \varphi) &= \int_{\mathcal{U}} (\lambda \varphi - \mathcal{C}\varphi) \varphi \tilde{\omega}^2 + \frac{1}{2} \int_{\mathcal{O}} \varphi(0, \cdot)^2 \tilde{\omega}^2 \\ &\quad - \frac{1}{2} \int_{\mathcal{U}} \varphi^2 v \cdot \nabla_x \tilde{\omega}^2 + \frac{1}{2} \int_{\Gamma_-} (\gamma_- \varphi)^2 \tilde{\omega}^2 d\xi_1, \end{aligned}$$

for any $\varphi \in C_c^1([0, T] \times \mathcal{O} \cup \Gamma_-)$. The same computations as presented during the proof of Lemma 2.1 imply that

$$(3.6) \quad \mathcal{E}(\varphi, \varphi) \geq (\lambda - \lambda_0) \|\varphi\|_{L_\omega^2}^2 + \|\varphi\|_{H_\omega^{1,+}}^2 + \frac{1}{2} \|\varphi(0)\|_{L_\omega^2}^2 + \frac{1}{2} \|\gamma_- \varphi\|_{L_\omega^2(\Gamma_-; d\xi_1)}^2,$$

for some $\lambda_0 \in \mathbb{R}$. For $\lambda > \lambda_0$, the bilinear form \mathcal{E} is thus coercive and the above mentioned Lions' theorem implies the existence of a function $f_\lambda \in \mathcal{H}$ which satisfies the variational equation

$$(3.7) \quad \mathcal{E}(f_\lambda, \varphi) = \int_{\Gamma_-} \mathfrak{f} e^{-\lambda t} \varphi \tilde{\omega}^2 d\xi_1 + \int_{\mathcal{O}} f_0 \varphi(0, \cdot) \tilde{\omega}^2 dv dx, \quad \forall \varphi \in C_c^1([0, T] \times \mathcal{O} \cup \Gamma_-).$$

Defining $f := f_\lambda e^{\lambda t}$ and using Theorem 3.1, we deduce that $f \in \mathcal{H} \cap C([0, T]; L_\omega^2(\mathcal{O}))$ is a renormalized solution to the inflow problem (3.5) and that $\gamma f \in L_\omega^2(\Gamma; d\xi_1)$. From the renormalization formulation, we have the uniqueness of such a solution. Because of the trace Theorem 3.1-(2), we can take $\beta(s) = s^2$ in (3.4) and we get

$$\begin{aligned} & \int_{\mathcal{U}} (f^2 (-v \cdot \nabla_x \varphi^2 - 2\varphi^2 \varphi_{\varphi,2}^{\mathcal{C}}) + 2|\nabla_v(f\varphi)|^2) dv dx dt \\ & + \int_{\Gamma} (\gamma f)^2 \varphi^2 n_x \cdot v dv d\sigma_x dt + \left[\int_{\mathcal{O}} (f^2 \varphi^2)(t, \cdot) dx dv \right]_0^T = 0, \end{aligned}$$

for any $\varphi \in \mathcal{D}(\bar{\mathcal{O}})$. Taking $\varphi := \tilde{\omega}$ in that last identity thanks to an approximation procedure and next using (3.7), the same computations as presented during the proof of Lemma 2.1 and the Gronwall lemma, we also deduce the energy estimate

$$\begin{aligned} & \|f_t\|_{L_\omega^2}^2 + \int_0^t \left(\|\gamma_+ f_s\|_{L_\omega^2(\Gamma_+; d\xi_1)}^2 + 2\|f\|_{H_\omega^{1,+}}^2 \right) e^{\lambda_0(t-s)} ds \\ & \leq \|f_0\|_{L_\omega^2}^2 e^{\lambda_0 t} + \int_0^t \|\mathfrak{f}_s\|_{L_\omega^2(\Gamma_-; d\xi_1)}^2 e^{\lambda_0(t-s)} ds. \end{aligned}$$

Step 2. For any $\alpha \in (0, 1)$ and $h \in \mathcal{H} \cap C([0, T]; L_\omega^2(\mathcal{O}))$ solution to the problem (3.5) for some $\mathfrak{h} \in L_\omega^2(\Gamma_-; d\xi_1)$, and thus $\gamma h \in L_\omega^2(\Gamma; d\xi_1)$, we then consider the modified Maxwell reflection boundary condition problem

$$(3.8) \quad \begin{cases} \partial_t f + v \cdot \nabla_x f = \mathcal{C} f & \text{in } (0, T) \times \mathcal{O} \\ \gamma_- f = \alpha \mathcal{R} \gamma_+ h & \text{on } (0, T) \times \Sigma_- \\ f(t = 0, \cdot) = f_0 & \text{in } \mathcal{O}, \end{cases}$$

for which a solution $f \in \mathcal{H} \cap C([0, T]; L_\omega^2(\mathcal{O}))$ such that $\gamma f \in L_\omega^2(\Gamma; d\xi_1)$ is given by the first step. Repeating the arguments of Step 1 in the proof of Lemma 2.1, we have

$$\begin{aligned} \|\mathcal{R} \gamma_+ h\|_{L_\omega^2(\Sigma_-; d\xi_1)}^2 & \leq \int_{\Sigma_+} \iota_S (\gamma_+ h)^2 \omega_A^2 d\xi_1 + \int_{\partial\Omega} \iota_D (K_1(\omega_A) - K_0(\omega_A)) (\widetilde{\gamma_+ h})^2 \\ & \leq \int_{\Sigma_+} (1 - \iota_D) (\gamma_+ h)^2 \omega_A^2 d\xi_1 + \int_{\partial\Omega} \iota_D K_2(\omega_A) (\widetilde{\gamma_+ h})^2 \\ & \leq \|\gamma_+ h\|_{L_\omega^2(\Gamma_+; d\xi_1)}^2. \end{aligned}$$

Thanks to the energy estimate stated in the first step, we immediately deduce that the mapping $h \mapsto f$ is α -Lipschitz for the norm defined by

$$\sup_{t \in [0, T]} \left\{ \|f_t\|_{L_\omega^2}^2 e^{-\lambda_0 t} + \int_0^t \|\gamma f_s\|_{L_\omega^2(\Gamma_+; d\xi_1)}^2 e^{-\lambda_0 s} ds \right\}.$$

From the Banach fixed point theorem, we deduce the existence of a unique fixed point, and that provides a solution to (3.8).

Step 3. For a sequence $\alpha_k \in (0, 1)$, $\alpha_k \nearrow 1$, we next consider the sequence (f_k) obtained in Step 2 as the solution to the modified Maxwell reflection boundary condition problem

$$(3.9) \quad \begin{cases} \partial_t f_k + v \cdot \nabla_x f_k = \mathcal{C} f_k & \text{in } (0, T) \times \mathcal{O} \\ \gamma_- f_k = \alpha_k \mathcal{R} \gamma_+ f_k & \text{on } (0, T) \times \Sigma_- \\ f_k(t=0, \cdot) = f_0 & \text{in } \mathcal{O}. \end{cases}$$

From the fact that $\mathcal{R} : L_\omega^2(\Sigma_+; d\xi_1) \rightarrow L_\omega^2(\Sigma_+; d\xi_1)$ with norm less than 1 as established in Step 2 and the energy estimate stated at the end of Step 1, f_k satisfies

$$(3.10) \quad \|f_{kt}\|_{L_\omega^2}^2 + \int_0^t \left\{ (1 - \alpha_k^2) \|\gamma f_{ks}\|_{L_\omega^2(\Gamma_+; d\xi_1)}^2 + 2\|f_{ks}\|_{H_\omega^{1,\dagger}}^2 \right\} e^{\lambda_0(t-s)} ds \leq \|f_0\|_{L_\omega^2}^2 e^{\lambda_0 t},$$

for any $t \in (0, T)$ and any $k \geq 1$. Choosing $\beta(s) := s^2$ and $\varphi := (n_x \cdot v) \langle v \rangle^{-2} \omega^2(v)$ in the Green formula (3.4), we additionally have

$$(3.11) \quad \int_\Gamma (\gamma f_k)^2 \omega^2 d\xi_2 dt \lesssim \|f_0\|_{L_\omega^2}^2 e^{\lambda_0 T}.$$

From the above estimates, we deduce that, up to the extraction of a subsequence, there exist $f \in \mathcal{H} \cap L^\infty(0, T; L_\omega^2(\mathcal{O}))$ and $f_\pm \in L_\omega^2(\Gamma_\pm; d\xi_2 dt)$ such that

$$f_k \rightharpoonup f \text{ weakly in } \mathcal{H} \cap L^\infty(0, T; L_\omega^2(\mathcal{O})), \quad \gamma_\pm f_k \rightharpoonup f_\pm \text{ weakly in } L_\omega^2(\Gamma; d\xi_2 dt).$$

Because $\langle v \rangle \omega^{-1} \in L^2(\mathbb{R}^d)$, we have $L_\omega^2(\Gamma; d\xi_2) \subset L^1(\Gamma; d\xi_1)$. On the other hand, we recall that from the very definition (1.3), we have

$$(3.12) \quad \mathcal{R} : L^1(\Sigma_+; d\xi_1) \rightarrow L^1(\Sigma_-; d\xi_1), \quad \|\mathcal{R}\|_{L^1(\Sigma; d\xi_1)} \leq 1.$$

These three pieces of information together imply that $\mathcal{R}(\gamma f_{k+}) \rightharpoonup \mathcal{R}(f_+)$ weakly in $L^1(\Gamma_-; d\xi_1)$. On the other hand, from Theorem 3.2, we have $\gamma f_k \rightharpoonup \gamma f$ weakly in $L_{\text{loc}}^2(\Gamma; d\xi_2)$. Using both convergences in the boundary condition $\gamma_- f_k = \mathcal{R}(\gamma_+ f_k)$, we obtain $\gamma_- f = \mathcal{R}(\gamma_+ f)$. We may thus pass to the limit in equation (3.9) and we obtain that $f \in X_T$ is a renormalized solution to the KFP equation (1.1) complemented with the Maxwell reflection boundary condition (1.3) and associated to the initial datum f_0 . Passing to the limit in (3.10), we also have

$$(3.13) \quad \|f_t\|_{L_\omega^2}^2 + 2 \int_0^t \|f_{ks}\|_{H_\omega^{1,\dagger}}^2 e^{\lambda_0(t-s)} ds \leq \|f_0\|_{L_\omega^2}^2 e^{\lambda_0 t},$$

for any $t \in (0, T)$.

Step 4. We consider now two solutions f_1 and $f_2 \in X_T$ to the KFP equation (1.1)-(1.3) associated to the same initial datum f_0 , so that the function $f := f_2 - f_1 \in X_T$ is a solution to the KFP equation (1.1)-(1.3) associated to the initial datum $f(0) = 0$. We write (3.4) with the choice $\varphi := \tilde{\omega}_1 \chi_R$, with $\tilde{\omega}_1$ given by (2.7) associated to $p = 1$, $A > 0$ large enough, $\omega_1 := \langle v \rangle^{k_1}$ defined in Remark 2.2, with $\chi_R(v) := \chi(v/R)$, $\mathbf{1}_{B_1} \leq \chi \in \mathcal{D}(\mathbb{R}^d)$, and with the choice $\beta \in C^2(\mathbb{R})$, $\beta(0) = 0$, β'' with compact support. We get

$$\begin{aligned} & \int_{\mathcal{O}} \beta(f_T) \varphi dv dx + \int_\Gamma \beta(\gamma f) \varphi (n_x \cdot v) dv d\sigma_x dt + \int_{\mathcal{U}} \beta''(f) |\nabla_v f|^2 \varphi dv dx dt \\ &= \int_{\mathcal{U}} \{ \beta(f) \mathcal{T}^* \varphi + \beta(f) (\Delta_v \varphi - \partial_{v_i} (b_i \varphi)) + cf \beta'(f) \varphi \} dv dx dt. \end{aligned}$$

We assume $0 \leq \beta(s) \leq |s|$, $|\beta'(s)| \leq 1$, and $\beta'' \geq 0$ so that we may get rid of the last term at the left-hand side of the above identity. Recalling that Ω is bounded, we observe that

$$\begin{aligned} |\beta(f) (\Delta_v \varphi - \partial_{v_i} (b_i \varphi)) + cf \beta'(f) \varphi| &\leq |f| (|\Delta_v \varphi| + |\partial_{v_i} (b_i \varphi)| + |c\varphi|) \\ &\lesssim |f| \omega_1 (1 + |v|^{\gamma-2}) \in L^1(\mathcal{U}), \end{aligned}$$

because of the condition $k_* > k_1 + (d + \gamma - 2)/2$ and the bound $(1 + \langle v \rangle^{s/2}) \omega f \in L^2(\mathcal{U})$. The same uniform estimate holds for the term $\beta(f) \mathcal{F}^* \varphi$. We also observe that

$$|\beta(\gamma f) \varphi(n_x \cdot v)| \lesssim |\gamma f| \omega_1 |n_x \cdot v| \in L^1(\Gamma),$$

because of the condition $k_* > k_1 + d/2 + 1$ and the bound $\gamma f \omega \in L^2(\Gamma; d\xi_2 dt)$ provided by (3.11). We may thus pass to the limit $R \rightarrow \infty$ and $\beta(s) \nearrow |s|$ such that $0 \leq s\beta'(s) \nearrow |s|$, and we deduce

$$\int_{\mathcal{O}} |f_T| \tilde{\omega}_1 dv dx + \int_{\Gamma} |\gamma f| \tilde{\omega}_1(n_x \cdot v) dv d\sigma_x dt \leq \int_{\mathcal{U}} |f| \varpi_{\tilde{\omega}_1, 1}^{\mathcal{L}} \tilde{\omega}_1 dv dx dt.$$

Using finally the estimate of Step 1 in the proof of Lemma 2.1 in order to get rid of the boundary term as well as the estimates obtained in Step 2 and Step 3 in the proof of Lemma 2.1 in order to deal with the RHS term, we get

$$\int_{\mathcal{O}} |f_T| \tilde{\omega}_1 dv dx \leq \kappa \int_0^T \int_{\mathcal{O}} |f| \tilde{\omega}_1 dv dx dt,$$

and we conclude to $f = 0$ thanks to Grönwall's lemma. \square

3.3. Well-posedness for the dual equation and conclusions. We first establish the well-posedness of the dual KFP equation in a L^2 framework.

Proposition 3.4. *Under the assumptions of Proposition 3.3, for any admissible weight function ω , any final time $T > 0$ and any final datum $g_T \in L_m^2(\mathcal{O})$, $m := \omega^{-1}$, there exists a unique $g \in Y_T := L^\infty(0, T; L_m^2) \cap C([0, T]; L_{\text{loc}}^2) \cap \mathcal{H}_m$ solution to the backward dual kinetic Fokker-Planck equation (2.21) in a similar sense as stated in Proposition 3.3.*

Sketch of the proof of Proposition 3.4. We follow the same strategy as during the proof of Proposition 3.3.

Step 1. Given $\mathbf{g} \in L_m^2(\Gamma_-; d\xi_1)$, we consider the backward inflow problem

$$(3.14) \quad \begin{cases} -\partial_t g - v \cdot \nabla_x g = \mathcal{C}^* g & \text{in } (0, T) \times \mathcal{O} \\ \gamma_+ g = \mathbf{g} & \text{on } (0, T) \times \Sigma_+ \\ g|_{t=T} = g_T & \text{in } \mathcal{O}, \end{cases}$$

We define \tilde{m} as during the proof of Lemma 2.3 and the bilinear form $\mathcal{E} : \mathcal{H}_m \times C_c^1((0, T] \times \mathcal{O} \cup \Gamma_+) \rightarrow \mathbb{R}$, by

$$\mathcal{E}(g, \varphi) := \int_{\mathcal{U}} (\lambda g - \mathcal{C}^* g) \varphi \tilde{m}^2 + \int_{\mathcal{U}} g (\partial_t + v \cdot \nabla_x) (\varphi \tilde{m}^2),$$

which is coercive for λ large enough thanks to Lemma 2.3 (and more precisely (2.29)-(2.30)). Using Lions' variant of the Lax-Milgram theorem [43, Chap III, §1], we obtain a variational solution $g \in \mathcal{H}_m$ to (3.14), and more precisely

$$\mathcal{E}(g, \varphi) = \int_{\Gamma_+} \mathbf{g} \varphi \tilde{m}^2 d\xi_1 + \int_{\mathcal{O}} g_T \varphi(T, \cdot) \tilde{m}^2 dv dx, \quad \forall \varphi \in C_c^1((0, T] \times \mathcal{O} \cup \Gamma_+).$$

Thanks to the trace Theorem 3.1 and the dissipativity property (2.29) of \mathcal{L}^* , we deduce that $g \in C([0, T]; L_m^2) \cap \mathcal{H}_m$.

Step 2. For a sequence $\alpha_k \in (0, 1)$, $\alpha_k \nearrow 1$, we build a sequence (g_k) of solutions to the modified Maxwell reflection boundary condition problem

$$(3.15) \quad \begin{cases} -\partial_t g_k - v \cdot \nabla_x g_k = \mathcal{C}^* g_k & \text{in } (0, T) \times \mathcal{O} \\ \gamma_+ g_k = \alpha_k \mathcal{R}^* \gamma_- f_k & \text{on } (0, T) \times \Sigma_+ \\ g_k(t = T, \cdot) = g_T & \text{in } \mathcal{O}, \end{cases}$$

by using Step 1, the fact that $\mathcal{R}^* : L_m^2(\Sigma_-; d\xi_1) \rightarrow L_m^2(\Sigma_-; d\xi_1)$ from Step 1 in Lemma 2.3 and the Banach fixed point Theorem. This sequence satisfies

$$\sup_{[0, T]} \|g_{kt}\|_{L_m^2}^2 + \int_0^T \left\{ \|\gamma g_{ks}\|_{L_m^2(\Gamma_+; d\xi_2)}^2 + \|g_{ks}\|_{H_m^{1, \dagger}}^2 \right\} ds \leq C_T \|g_T\|_{L_m^2}^2$$

for some constant C_T and any $k \geq 1$. We may extract converging subsequences $(g_{k'})$ and $(\gamma g_{k'})$ with associated limits g and $\bar{\gamma}$, and passing to the limit in (3.15) with the help of Theorem 3.2, we deduce that $\bar{\gamma} = \gamma g$ and that g is a renormalized solution to (2.21).

Step 3. We may assume $\omega \lesssim e^{\zeta \langle v \rangle^{s_*}}$ with $s_* = \min(2, \gamma)$ and $\zeta \in (0, \zeta^*)$, defining $\zeta^* := b_0/2$ if $\gamma < 2$, $\zeta^* := \min(b_0, 1/\Theta^*)/2$ if $\gamma = 2$, $\zeta^* := 1/(2\Theta^*)$ if $\gamma > 2$. We set $m_1 := e^{-\zeta_1 \langle v \rangle^{s_*}}$, with $\zeta_1 \in (\zeta, \zeta_*)$. We consider two solutions g_1 and $g_2 \in C([0, T]; L_m^2) \cap \mathcal{H}_m$ to the backward dual KFP equation (2.21) associated to the same final datum g_T , so that the function $g := g_2 - g_1 \in C([0, T]; L_m^2) \cap \mathcal{H}_m$ is a solution to the KFP equation (1.1)-(1.3) associated to the initial datum $g(T) = 0$. Choosing $\varphi := \tilde{m}_1 \chi_R$ in (3.4), with the notations of Step 4 of the proof of Proposition 3.3, and proceeding similarly, we get

$$\begin{aligned} & \int_{\mathcal{O}} \beta(g(0)) \varphi \, dv \, dx + \int_{\Gamma} \beta(\gamma g) \varphi (n_x \cdot v) \, dv \, d\sigma_x \, dt + \int_{\mathcal{U}} \beta''(f) |\nabla_v f|^2 \varphi \, dv \, dx \, dt \\ & = \int_{\mathcal{U}} \beta(g) \mathcal{T} \varphi + \int_{\mathcal{U}} \{ \beta(g) (\Delta_v \varphi + \operatorname{div}_v(b\varphi)) + (c - \operatorname{div}_v b) f \beta'(f) \varphi \} \, dv \, dx \, dt. \end{aligned}$$

Taking advantage of the fact that all the terms in the interior and at the boundary are now well defined, we may argue as for the proof of the $L_{m_1}^1$ estimate performed in Lemma 2.3, and we conclude that $g \equiv 0$ as in Step 4 of the proof of Proposition 3.3. \square

We conclude the section by reformulating and slightly improving the two previous well-posedness results.

Theorem 3.5. *Consider an admissible weight function ω and set $m := \omega^{-1}$.*

(1) *There exists a semigroup $S_{\mathcal{L}}$ on $L_{\omega}^2(\mathcal{O})$ such that for any $f_0 \in L_{\omega}^2(\mathcal{O})$, the function $f_t := S_{\mathcal{L}}(t)f_0$ is the unique solution in $C([0, T]; L_{\omega}^2) \cap \mathcal{H}$, $\forall T > 0$, to the KFP equation (1.1)-(1.3)-(1.4). Furthermore (2.5) holds if additionally $f_0 \in L_{\omega}^p$ for some $p \in [1, \infty]$.*

(2) *Similarly, there exists a semigroup $S_{\mathcal{L}^*}$ on $L_m^2(\mathcal{O})$ such that for any $g_T \in L_m^2(\mathcal{O})$, the function $g_t := S_{\mathcal{L}^*}(T-t)g_T$ is the unique solution in $C([0, T]; L_m^2) \cap \mathcal{H}_m$, $\forall T > 0$, to the dual KFP problem (2.21). Furthermore (2.23) holds if additionally $g_T \in L_m^q$ for some $q \in [1, \infty]$.*

(3) *The semigroups $S_{\mathcal{L}}$ and $S_{\mathcal{L}^*}$ are dual one toward the other. In other words, the equation (2.22) holds for any $f_0 \in L_{\omega}^2(\mathcal{O})$ and $g_T \in L_m^2(\mathcal{O})$.*

Proof of Theorem 3.5. The proof is split into four steps.

Step 1. We may define the semigroup $S_{\mathcal{L}}$ by setting $S_{\mathcal{L}}(t)f_0 := f_t$ for any $f_0 \in L_{\omega}^2$ and $t \geq 0$, where f_t is the unique solution in X_T , $\forall T > 0$, to the KFP equation (1.1)-(1.3)-(1.4) provided by Proposition 3.3. In particular (2.5) holds for $p = 2$. Proceeding as in Step 4 during the proof of Proposition 3.3, we may justify the computations performed during the proof of Lemma 2.1 and we get that (2.5) holds for $p = 1$ when $f_0 \in L_{\omega}^2 \cap L_{\omega}^1$. By interpolation, we obtain that (2.5) holds for any $p \in [1, 2]$ when $f_0 \in L_{\omega}^2 \cap L_{\omega}^p$.

Step 2. From the well-posedness result of Proposition 3.4, we may define the semigroup $S_{\mathcal{L}^*}$ by setting $S_{\mathcal{L}^*}(t)g_T := g(T-t)$ for any $g_T \in L_m^2$ and any $0 \leq t \leq T$, where g is unique solution in Y_T to the KFP problem (2.21). We obtain as in Step 1 that furthermore (2.23) holds for any $q \in [1, 2]$ if $g_0 \in L_m^2 \cap L_m^q$.

Step 3. We change ι by $\iota_n := \iota_{S,n} + \iota_{D,n} \leq 1 - 1/n$, with $\iota_{S,n} := \iota_S(1 - 1/n)$, $\iota_{D,n} := \iota_D(1 - 1/n)$, and we denote by \mathcal{R}_n and \mathcal{R}_n^* the corresponding reflection operators. Denoting by f_n the solution associated to the KFP equation, the reflection operator \mathcal{R}_n and the initial datum f_0 given by

Step 1 (or Proposition 3.3), we have the additional property $\gamma f_n \in L_\omega^2(\Gamma; d\xi_1)$. The solution g_n associated to the dual problem (2.21) for the reflection operator \mathcal{R}_n^* and the final datum g_T given by Step 2 (or Proposition 3.4) also satisfies the additional property $\gamma g_n \in L_m^2(\Gamma; d\xi_1)$. Because of these additional estimates on the boundary, we may justify the computations leading to the identity (2.22) by starting from (3.4) applied with $g := f_n$, $\varphi := g_n$ and $\beta(s) := s$ or by applying (a variant of) Theorem 3.1 to the function $\beta(f_n g_n)$, noticing that

$$\partial_t(f_n g_n) + v \cdot \nabla_x(f_n g_n) = \Delta_v f_n g_n - f_n \Delta_v g_n + \operatorname{div}_v(b f_n g_n) \quad \text{in } \mathcal{D}'(\mathcal{U}).$$

In that way, we obtain

$$(3.16) \quad \int_{\mathcal{O}} f_n(T) g_T = \int_{\mathcal{O}} f_0 g_n(0), \quad \forall n \geq 1.$$

Because (f_n) is bounded in $L^\infty(0, T; L_\omega^2)$ and in $W^{1,\infty}(0, T; \mathcal{D}'(\mathcal{O}))$, we deduce that $f_n(T) \rightharpoonup f(T) := S_{\mathcal{L}}(T) f_0$ weakly in L_ω^2 . Similarly, we have $g_n(0) \rightharpoonup g(0) = S_{\mathcal{L}^*}(T) g_T$ weakly in L_m^2 . We may thus pass to the limit $n \rightarrow \infty$ in (3.16) and we deduce that (2.22) holds, which exactly means that $(S_{\mathcal{L}^*})^* = S_{\mathcal{L}}$.

Step 4. For any $p \in (2, \infty]$, we know from Step 2 that for any $g_T \in L_m^2 \cap L_m^{p'}$, $T > 0$, there holds

$$(3.17) \quad \|g(t, \cdot)\|_{L_m^{p'}} \leq C_1 e^{C_2(T-t)} \|g_T\|_{L_m^{p'}}, \quad \forall t \in [0, T].$$

Now, for $f_0 \in L_\omega^p$, we have

$$\begin{aligned} \|f(T)\|_{L_\omega^p} &= \sup_{g_T \in L_m^2, \|g_T\|_{L_m^{p'}} \leq 1} \int f(T) g_T \\ &= \sup_{g_T \in L_m^2, \|g_T\|_{L_m^{p'}} \leq 1} \int f_0 g(0) \\ &\leq \|f_0\|_{L_\omega^p} \sup_{g_T \in L_m^2, \|g_T\|_{L_m^{p'}} \leq 1} \|g(0)\|_{L_m^{p'}} \\ &\leq \|f_0\|_{L_\omega^p} \sup_{g_T \in L_m^2, \|g_T\|_{L_m^{p'}} \leq 1} C_1 e^{C_2 T} \|g_T\|_{L_m^{p'}} = C_1 e^{C_2 T} \|f_0\|_{L_\omega^p}, \end{aligned}$$

where we have successively used the Riesz representation theorem, the duality identity (2.22), the Holder inequality, the estimate (3.17) and the Riesz representation theorem again. We have thus established that (2.5) holds for any $f_0 \in L_\omega^p \cap L_\omega^2$, $p \in [1, \infty]$. We establish in the same way that (2.23) holds for any $g_T \in L_m^q \cap L_m^2$, $q \in [1, \infty]$.

Step 5. For $f_0 \in L_\omega^2$, let us introduce a sequence $f_{0,n} \in L_{\omega^\sharp}^2 \cap L_\omega^\infty$, such that ω^\sharp is an admissible weight function satisfying $\omega^\sharp/\omega \rightarrow \infty$ and such that $f_{0,n} \rightarrow f_0$ in L_ω^2 . From the previous analysis, the solution f_n to the KFP equation (1.1)–(1.3) associated to the initial datum $f_{0,n}$ satisfies $f_n \in L^\infty(0, T; L_{\omega^\sharp}^2 \cap L_\omega^\infty) \cap C([0, T]; L_{\text{loc}}^2) \subset C([0, T]; L_\omega^2)$. From (2.5), (f_n) is a Cauchy sequence in $C([0, T]; L_\omega^2)$ and thus converges to the solution f to the KFP equation (1.1)–(1.3) associated to the initial datum f_0 . We have established that $f \in C([0, T]; L_\omega^2)$ and the same argument holds for the dual problem. \square

4. ULTRA CONTRACTIVITY

In this section, we explain how to adapt to the KFP equation in a domain the De Giorgi-Nash-Moser theory developed for parabolic equations, in particular in [23, 54, 52, 53], and generalized recently to the KFP equation in the whole space, in particular in [57, 30]. The gain of integrability $L^1 \rightarrow L^2$ essentially follows and slightly simplifies the proofs presented in [19, 18] which are very in the spirit of Nash approach [54].

4.1. An improved weighted L^2 estimate at the boundary. Let us observe that for a solution f to the KFP equation (1.1), we may write

$$\mathcal{T} \frac{f^2}{2} = f \mathcal{T} f = f \mathcal{C} f,$$

where we define

$$(4.1) \quad \mathcal{T} := \partial_t + v \cdot \nabla_x,$$

and we recall that \mathcal{C} has been defined in (1.9). Multiplying that equation by $\Phi^2 := \varphi^2 \tilde{\omega}^2$ with a time truncation function $\varphi \in \mathcal{D}(0, T)$ and a weight function $\tilde{\omega} : \mathcal{O} \rightarrow (0, \infty)$, and next integrating in all the variables with the help of (2.1), we obtain

$$(4.2) \quad \frac{1}{2} \int_{\Gamma} (\gamma f)^2 \Phi^2 n_x \cdot v - \frac{1}{2} \int_{\mathcal{U}} f^2 \mathcal{T} \Phi^2 = - \int_{\mathcal{U}} |\nabla_v (f \Phi)|^2 + \int_{\mathcal{U}} f^2 \Phi^2 \tilde{\omega},$$

with $\mathcal{U} := (0, T) \times \mathcal{O}$, $\Gamma := (0, T) \times \Sigma$, $T \in (0, \infty)$ and $\tilde{\omega} := \varpi_{\tilde{\omega}, 2}^{\mathcal{C}}$ is defined in (1.17), or equivalently by

$$(4.3) \quad \tilde{\omega} = \frac{1}{4} \frac{|\nabla_v \tilde{\omega}^2|^2}{\tilde{\omega}^4} - \frac{1}{2} b \cdot \frac{\nabla_v \tilde{\omega}^2}{\tilde{\omega}^2} + c - \frac{1}{2} \operatorname{div}_v b.$$

We first establish a key estimate on the KFP equation (1.1)–(1.3)–(1.4) which makes possible to control a solution near the boundary. The proof is based on the introduction of an appropriate weight function which combines the twisting term used in the previous section and the twisting term used in [28, Section 11], that last one being in the spirit of moment arguments used in [45, 51].

Proposition 4.1. *For any admissible weight function ω there exists $C = C(\omega, \Omega) > 0$ such that for any solution f to the KFP equation (1.1)–(1.3), any $T > 0$ and any smooth function $0 \leq \varphi \in \mathcal{D}((0, T))$, there holds*

$$\int_{\mathcal{U}} f^2 \omega^2 \left\{ \frac{(n_x \cdot \hat{v})^2}{\delta^{1/2}} + \langle v \rangle^{\varsigma} \right\} \varphi^2 + \int_{\mathcal{U}} |\nabla_v (f \omega)|^2 \varphi^2 \leq C \int_{\mathcal{U}} f^2 \omega^2 [|\partial_t \varphi|^2 + \varphi^2],$$

with $\varsigma := \gamma + s - 2$.

It is worth emphasizing that an admissible weight function ω is strongly confining if and only if the associated parameter $\varsigma \in \mathbb{R}$ satisfies $\varsigma > 0$.

Proof of Proposition 4.1. We introduce the function

$$\Phi^2 := \varphi^2 \tilde{\omega}^2, \quad \tilde{\omega}^2 := \left(1 + \frac{1}{4} \frac{n_x \cdot v}{\langle v \rangle^4} + \frac{1}{4D^{1/2}} \delta(x)^{1/2} \frac{n_x \cdot v}{\langle v \rangle^2} \right) \omega_A^2,$$

where the weight function ω_A is defined in (2.7) for $A \geq 1$ large enough (to be fixed below) and where $D = \sup_{x \in \Omega} \delta(x)$ is half the diameter of Ω , so that in particular an estimate similar to (2.25) holds. From (4.2) we have

$$(4.4) \quad \begin{aligned} & 2 \int_{\mathcal{U}} |\nabla_v (f \Phi)|^2 + \int_{\Gamma} (\gamma f)^2 \Phi^2 (n_x \cdot v) - \int_{\mathcal{U}} f^2 v \cdot \nabla_x \Psi_2 \\ & = \int_{\mathcal{U}} f^2 v \cdot \nabla_x \Psi_1 + 2 \int_{\mathcal{U}} f^2 \Phi^2 \tilde{\omega} + \int_{\mathcal{U}} f^2 \partial_t (\Phi^2), \end{aligned}$$

for $\tilde{\omega}$ as defined on (4.3) and we denote

$$\Psi_1 := \varphi^2 \omega_A^2 \left(1 + \frac{1}{4} \frac{n_x \cdot v}{\langle v \rangle^4} \right), \quad \Psi_2 := \frac{\varphi^2 \omega_A^2}{4D^{1/2}} \delta(x)^{1/2} \frac{n_x \cdot v}{\langle v \rangle^2}.$$

We now compute each term separately.

Step 1. Observing that $\tilde{\omega}^2 = \left(1 + \frac{1}{4} n_x \cdot \frac{v}{\langle v \rangle^4}\right) \omega_A^2$ on the boundary $\Gamma = (0, T) \times \Sigma \times \mathbb{R}^d$, we can argue as in Step 1 in the proof of Lemma 2.1 to deduce that, choosing $A \geq 1$ large enough, the contribution of the boundary term in (4.4) is nonnegative, that is

$$\int_{\Gamma} (\gamma f)^2 \Phi^2(n_x \cdot v) \geq 0.$$

Step 2. In order to deal with the third term at the left-hand side of (4.4), we define $\psi := \delta(x)^{1/2} (n_x \cdot v) \langle v \rangle^{-2}$. Observing that $\langle v \rangle \psi \in L^\infty(\mathcal{O})$, $\nabla_v \psi \in L^\infty(\mathcal{O})$ and

$$-v \cdot \nabla_x \psi = \frac{1}{2} \frac{1}{\delta(x)^{1/2}} (n_x \cdot v)^2 \langle v \rangle^{-2} - \delta(x)^{1/2} (v \cdot D_x n_x v) \langle v \rangle^{-2},$$

we have

$$-\int_{\mathcal{U}} f^2 v \cdot \nabla_x \Psi_2 = \frac{1}{4D^{1/2}} \int_{\mathcal{U}} f^2 \varphi^2 \omega_A^2 \left\{ \frac{1}{2} \frac{1}{\delta(x)^{1/2}} (n_x \cdot \hat{v})^2 - \delta(x)^{1/2} (\hat{v} \cdot D_x n_x \hat{v}) \right\}.$$

For the first term at the right-hand side of (4.4), a direct computation gives

$$\int_{\mathcal{U}} f^2 v \cdot \nabla_x \Psi_1 = \frac{1}{4} \int_{\mathcal{U}} f^2 \varphi^2 \omega_A^2 \langle v \rangle^{-2} (\hat{v} \cdot D_x n_x \hat{v}).$$

Step 3. Writing $\tilde{\omega} \leq 2\langle \tilde{\omega}_+ \rangle - \langle \tilde{\omega}_- \rangle$ and gathering previous estimates yields

$$\begin{aligned} & 2 \int_{\mathcal{U}} |\nabla_v(f\tilde{\omega})|^2 \varphi^2 + \frac{1}{8D^{1/2}} \int_{\mathcal{U}} f^2 \varphi^2 \omega_A^2 \frac{(n_x \cdot \hat{v})^2}{\delta(x)^{1/2}} + 2 \int_{\mathcal{U}} f^2 \varphi^2 \tilde{\omega}^2 \langle \tilde{\omega}_- \rangle \\ & \leq \int_{\mathcal{U}} f^2 \varphi^2 \omega_A^2 (\hat{v} \cdot D_x n_x \hat{v}) \left\{ \frac{1}{4\langle v \rangle^2} + \frac{\delta^{1/2}}{4D^{1/2}} \right\} + 4 \int_{\mathcal{U}} f^2 \varphi^2 \omega_A^2 \langle \tilde{\omega}_+ \rangle + \int_{\mathcal{U}} f^2 \tilde{\omega}^2 \partial_t(\varphi^2) \\ & \leq C_{\Omega, A} \int_{\mathcal{U}} f^2 \omega^2 \langle \tilde{\omega}_+ \rangle \varphi^2 + C_A \int_{\mathcal{U}} f^2 \omega^2 |\partial_t \varphi^2|, \end{aligned}$$

where we recall that $\delta \in W^{2, \infty}(\Omega)$. Using that $\langle \tilde{\omega}_- \rangle \geq \kappa_0 \langle v \rangle^\varsigma$ with $\kappa_0 > 0$, because of (2.2)-(2.19), and also that $\tilde{\omega}_+$ is bounded, we deduce

$$\begin{aligned} & 2 \int_{\mathcal{U}} |\nabla_v(f\tilde{\omega})|^2 \varphi^2 + \frac{1}{8D^{1/2}} \int_{\mathcal{U}} f^2 \varphi^2 \omega_A^2 \frac{(n_x \cdot \hat{v})^2}{\delta(x)^{1/2}} + 2\kappa_0 \int_{\mathcal{U}} f^2 \varphi^2 \tilde{\omega}^2 \langle v \rangle^\varsigma \\ & \leq C_{\Omega, A} \int_{\mathcal{U}} f^2 \omega^2 \varphi^2 + C_A \int_{\mathcal{U}} f^2 \omega^2 |\partial_t \varphi^2|. \end{aligned}$$

We then conclude by observing that

$$\int_{\mathcal{O}} |\nabla_v(f\omega)|^2 + f^2 \omega^2 \lesssim \int_{\mathcal{O}} |\nabla_v(f\tilde{\omega})|^2 + f^2 \tilde{\omega}^2$$

and using that $\omega \lesssim \omega_A \lesssim \tilde{\omega} \lesssim \omega_A \lesssim \omega$. \square

We may write the above weighted L^2 estimate in a more convenient way, where the penalization of a neighborhood of the boundary is made clearer. For that purpose we state the following interpolation result which formalizes and improves some estimates used during the proof of [28, Lemma 11.9] (see also [19, Lemma 3.2]).

Lemma 4.2. *For $d \geq 1$, for any $\varsigma > 0$ and any function $g : \mathcal{O} \rightarrow \mathbb{R}$, there holds*

$$(4.5) \quad \int_{\mathcal{O}} \frac{g^2}{\delta^\beta} \lesssim \int_{\mathcal{O}} |\nabla_v g|^2 + \int_{\mathcal{O}} g^2 \left(\frac{(n_x \cdot \hat{v})^2}{\delta^{1/2}} + \langle v \rangle^\varsigma \right),$$

with $\beta = [2(d' + 1 + 2d'/\varsigma)]^{-1}$, $d' := \max(d, 3)$.

Proof of Lemma 4.2. For $\alpha, \beta, \eta > 0$, we start writing

$$\begin{aligned} \int_{\mathcal{O}} \frac{g^2}{\delta^\beta} &= \int_{\mathcal{O}} \frac{g^2}{\delta^\beta} \mathbf{1}_{\langle v \rangle^{\delta^\alpha} \geq 1} + \int_{\mathcal{O}} \frac{g^2}{\delta^\beta} \mathbf{1}_{(n_x \cdot v)^2 > \delta^{2\eta}} \mathbf{1}_{\langle v \rangle^{\delta^\alpha} < 1} \\ &\quad + \int_{\mathcal{O}} \frac{g^2}{\delta^\beta} \mathbf{1}_{|n_x \cdot v| \leq \delta^\eta} \mathbf{1}_{\langle v \rangle^{\delta^\alpha} < 1} =: T_1 + T_2 + T_3. \end{aligned}$$

For the first term, we have

$$T_1 \leq \int_{\mathcal{O}} \frac{g^2}{\delta^\beta} \langle v \rangle^\varsigma \delta^{\varsigma\alpha} = \int_{\mathcal{O}} g^2 \langle v \rangle^\varsigma,$$

by choosing $\alpha\varsigma = \beta$. For the second term, we have

$$T_2 \leq \int_{\mathcal{O}} \frac{g^2}{\delta^\beta} \frac{1}{\langle v \rangle^2 \delta^{2\beta/\varsigma}} \frac{(n_x \cdot v)^2}{\delta^{2\eta}} = \int_{\mathcal{O}} g^2 \frac{(n_x \cdot \hat{v})^2}{\delta^{1/2}},$$

by choosing $2\eta + \beta(1 + 2/\varsigma) = 1/2$. For the third term, we define $2^* := 2d'/(d' - 2)$ the Sobolev exponent in dimension $d' \geq 3$, and we compute

$$\begin{aligned} T_3 &\leq \int_{\Omega} \delta^{-\beta} \left(\int_{\mathbb{R}^d} |g|^{2^*} \right)^{2/2^*} \left(\int_{\mathbb{R}^d} \mathbf{1}_{|n_x \cdot v| \leq \delta^\eta} \mathbf{1}_{\langle v \rangle < \delta^{-\beta/\varsigma}} \right)^{2/d} \\ &\lesssim \int_{\Omega} \delta^{-\beta} (\delta^\eta \delta^{-(d-1)\beta/\varsigma})^{2/d} \int_{\mathbb{R}^d} (|\nabla_v g|^2 + g^2) \end{aligned}$$

where we have used the Hölder inequality in the first line and the Sobolev inequality in the second line together with the observation that $\langle v' \rangle \leq \langle v \rangle$ in the orthonormal representation $v = v_1 n_x + v'$. Choosing η such that

$$-\beta - 2(1 - 1/d)\beta/\varsigma + \eta 2/d = 0,$$

we deduce that $\eta = (d/2 + (d-1)/\varsigma)\beta$, then $\beta = [2(d+1+2d/\varsigma)]^{-1}$ and we conclude to (4.5). \square

As an immediate consequence of Proposition 4.1 and the interpolation inequality (4.5), we get the following estimate which holds for strongly confining weight functions.

Corollary 4.3. *Consider a strongly confining weight function ω and recall that $\varsigma := \gamma + s - 2 > 0$ and $\beta := [2(d' + 1 + 2d'/\varsigma)]^{-1}$. Under the conditions of Proposition 4.1, there exists C_Ω such that*

$$\int_{\mathcal{U}} \frac{f^2}{\delta^\beta} \omega^2 \varphi^2 + \int_{\mathcal{U}} f^2 \langle v \rangle^\varsigma \omega^2 \varphi^2 + \int_{\mathcal{U}} |\nabla_v(f\omega)|^2 \varphi^2 \lesssim \int_{\mathcal{U}} f^2 \omega^2 [|\partial_t \varphi^2| + C_\Omega \varphi^2].$$

For a weakly confining admissible weight function, we obtain the following weaker estimate which is similar to [19, Proposition 3.3] and [18, Proposition 5.3].

Corollary 4.4. *Let us consider a weakly confining admissible weight function ω corresponding to the case when $\varsigma := \gamma + s - 2 \leq 0$. We set $\beta := [2(d' + 1)]^{-1}$, recalling that $d' := \max(d, 3)$. Under the conditions of Proposition 4.1, there holds*

$$\int_{\mathcal{U}} \frac{f^2}{\delta^\beta} \frac{\omega^2}{\langle v \rangle^2} \varphi^2 + \int_{\mathcal{U}} |\nabla_v(f\omega)|^2 \varphi^2 \lesssim \int_{\mathcal{U}} f^2 \omega^2 [|\partial_t \varphi^2| + \varphi^2].$$

Proof of Corollary 4.4. We just use the inequality

$$\int_{\mathcal{O}} \frac{g^2}{\langle v \rangle^2 \delta^\beta} \lesssim \int_{\mathcal{O}} (|\nabla_v g|^2 + g^2) + \int_{\mathcal{O}} g^2 \frac{(n_x \cdot \hat{v})^2}{\delta^{1/2}},$$

that we may establish proceeding exactly as in [19, Lemma 3.2], with $g := f\omega$ and the conclusion of Proposition 4.1. \square

4.2. A downgraded weighted $L^2 - L^p$ estimate. Taking advantage of a known $L^2 - L^p$ estimate available for the KFP equation set in the whole space, and thus in the interior of the domain, we deduce a downgrade weighted $L^2 - L^p$ estimate.

Proposition 4.5. *We set $\nu := \max(2, \gamma - 1)$. There exists $p > 2$, $\alpha > p$ and $C \in (0, \infty)$ such that any solution f to the KFP equation (1.1)–(1.3) satisfies*

$$(4.6) \quad \left\| f \varphi \frac{\omega}{\langle v \rangle^\nu} \delta^{\alpha/p} \right\|_{L^p(\mathcal{U})} \leq C \|(\varphi + |\varphi'|)f\omega\|_{L^2(\mathcal{U})},$$

for any $0 \leq \varphi \in \mathcal{D}((0, T))$, any $T > 0$ and any admissible weight function ω .

We follow the proof of [30, Lemma 10] and of Step 2 in the proof of [19, Proposition 3.5].

Proof of Proposition 4.5. We split the proof into two steps.

Step 1. Consider a subset Ω' of Ω such that $\overline{\Omega'} \subset \Omega$. We introduce a truncation function $\chi \in \mathcal{D}(\Omega)$ such that $\mathbf{1}_{\Omega'} \leq \chi \leq 1$, and the function $\omega_0 := \langle v \rangle^{-\nu} \omega$. We define the function $\bar{f} := \varphi \chi \omega_0 f$ which satisfies

$$\partial_t \bar{f} + v \cdot \nabla_x \bar{f} = F \quad \text{in } \mathbb{R} \times \mathbb{R}^d \times \mathbb{R}^d,$$

where $F := F_0 + \text{div}_v F_1$ with

$$F_0 := f \omega_0 (\varphi' \chi + \varphi v \cdot \nabla_x \chi) + \varphi \chi (\omega_0 b - \nabla_v \omega_0) \cdot \nabla_v f + c \bar{f}$$

and

$$F_1 := \varphi \chi \omega_0 \nabla_v f.$$

From [12, Theorem 1.3] with $p = 2$, $r = 0$, $\beta = 1$, $m = 1$, $\kappa = 1$ and $\Omega = 1$, we have

$$\begin{aligned} \|D_t^{1/3} \bar{f}\|_{L^2(\mathbb{R}^{1+2d})}^2 + \|D_x^{1/3} \bar{f}\|_{L^2(\mathbb{R}^{1+2d})}^2 &\lesssim \|\bar{f}\|_{L^2(\mathbb{R}^{1+2d})}^2 + \|\nabla_v \bar{f}\|_{L^2(\mathbb{R}^{1+2d})}^2 \\ &\quad + \|\langle v \rangle F_0\|_{L^2(\mathbb{R}^{1+2d})}^2 + \|\langle v \rangle^2 F_1\|_{L^2(\mathbb{R}^{1+2d})}^2. \end{aligned}$$

On the one hand, a straightforward computation gives

$$\|\bar{f}\|_{L^2(\mathbb{R}^{1+2d})} \lesssim \|\varphi \omega \langle v \rangle^{-\nu} f\|_{L^2(\mathcal{U})}$$

and

$$\|\nabla_v \bar{f}\|_{L^2(\mathbb{R}^{1+2d})} \lesssim \|\varphi \omega \langle v \rangle^{-1-\nu} f\|_{L^2(\mathcal{U})} + \|\varphi \langle v \rangle^{-\nu} \nabla_v (f \omega)\|_{L^2(\mathcal{U})}.$$

On the other hand, we have

$$\begin{aligned} \|\langle v \rangle F_0\|_{L^2(\mathbb{R}^{1+2d})} &\lesssim \|\varphi' \omega \langle v \rangle^{1-\nu} f\|_{L^2(\mathcal{U})} + \|\nabla_x \chi\|_{L^\infty(\Omega)} \|\varphi \omega \langle v \rangle^{2-\nu} f\|_{L^2(\mathcal{U})} \\ &\quad + \|\varphi \omega \langle v \rangle^{\gamma-2-\nu} f\|_{L^2(\mathcal{U})} + \|\varphi \omega \langle v \rangle^{\max(\gamma-1, s)-\nu} \nabla_v f\|_{L^2(\mathcal{U})} \end{aligned}$$

using the growth conditions (1.12)–(1.13), and

$$\|\langle v \rangle^2 F_1\|_{L^2(\mathbb{R}^{1+2d})} \lesssim \|\varphi \omega \langle v \rangle^{2-\nu} \nabla_v f\|_{L^2(\mathcal{U})}.$$

Observing that $|\omega \nabla_v f| \lesssim |\nabla_v (f \omega)| + \langle v \rangle^{s-1} \omega |f|$, it follows

$$\begin{aligned} \|\langle v \rangle F_0\|_{L^2(\mathbb{R}^{1+2d})} + \|\langle v \rangle^2 F_1\|_{L^2(\mathbb{R}^{1+2d})} &\lesssim \|\varphi' \omega \langle v \rangle^{1-\nu} f\|_{L^2(\mathcal{U})} + \|\nabla_x \chi\|_{L^\infty(\Omega)} \|\varphi \omega \langle v \rangle^{2-\nu} f\|_{L^2(\mathcal{U})} \\ &\quad + \|\varphi \langle v \rangle^{\max(2, \gamma-1, s)-\nu} f\|_{L^2(\mathcal{U})} + \|\varphi \langle v \rangle^{\max(2, \gamma-1, s)-\nu} \nabla_v (f \omega)\|_{L^2(\mathcal{U})}. \end{aligned}$$

As $s \leq 2$, we have $\max(2, \gamma - 1, s) = \max(2, \gamma - 1) = \nu$. Therefore, Corollary 4.3 and the Sobolev embedding $H^{1/3}(\mathbb{R}^{1+2d}) \subset L^p(\mathbb{R}^{1+2d})$, with $p := 6d/(3d - 2) > 2$, yield

$$\begin{aligned} \|\bar{f}\|_{L^p(\mathbb{R}^{1+2d})} &\lesssim \|D_t^{1/3} \bar{f}\|_{L^2(\mathbb{R}^{1+2d})} + \|D_x^{1/3} \bar{f}\|_{L^2(\mathbb{R}^{1+2d})} + \|\nabla_v \bar{f}\|_{L^2(\mathbb{R}^{1+2d})} + \|\bar{f}\|_{L^2(\mathbb{R}^{1+2d})} \\ &\lesssim \|\varphi' \omega f\|_{L^2(\mathcal{U})} + \|\nabla_x \chi\|_{L^\infty(\Omega)} \|\varphi \omega f\|_{L^2(\mathcal{U})} + \|\varphi \omega f\|_{L^2(\mathcal{U})}. \end{aligned}$$

Step 2. Choosing $\chi_k \in \mathcal{D}(\Omega)$ such that $\mathbf{1}_{\Omega_{k+1}} \leq \chi_k \leq \mathbf{1}_{\Omega_k}$, with $\Omega_k := \{x \in \Omega \mid \delta > 2^{-k}\}$, and $2^{-k} \|\nabla_x \chi_k\|_{L^\infty} \lesssim 1$ uniformly in $k \geq 1$, we deduce from the last estimate that

$$\|f\varphi\omega_0\|_{L^p(\mathcal{U}_{k+1})} \lesssim \|\varphi'f\omega\|_{L^2(\mathcal{U})} + 2^k \|\varphi f\omega\|_{L^2(\mathcal{U})},$$

for any $k \geq 1$, where we denote $\mathcal{U}_k := (0, T) \times \Omega_k \times \mathbb{R}^d$. Summing up, we obtain

$$\begin{aligned} \int_{\mathcal{U}} \delta^\alpha (\varphi f\omega_0)^p &= \sum_k \int_{\mathcal{U}_{k+1} \setminus \mathcal{U}_k} \delta^\alpha (\varphi f\omega_0)^p \\ &\lesssim \sum_k 2^{-k\alpha} \int_{\mathcal{U}_{k+1}} (\varphi f\omega_0)^p \\ &\lesssim \sum_k 2^{k(p-\alpha)} (\|\varphi'f\omega\|_{L^2(\mathcal{U})} + \|\varphi f\omega\|_{L^2(\mathcal{U})})^p \\ &\lesssim (\|\varphi'f\omega\|_{L^2(\mathcal{U})} + \|\varphi f\omega\|_{L^2(\mathcal{U})})^p \end{aligned}$$

because $\alpha > p$, what is nothing but (4.6). \square

4.3. The $L^1 - L^r$ estimate up to the boundary. We are now in position for stating our weighted $L^1 - L^r$ estimate up to the boundary which is the well-known cornerstone step in the proof of De Giorgi-Nash-Moser gain of integrability estimate.

Proposition 4.6. *There exists an exponent $r > 2$ such that any solution f to the KFP equation (1.1)–(1.3) satisfies*

$$(4.7) \quad \|f\varphi\omega\|_{L^2(\mathcal{U})} \lesssim \|(\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} f\omega\|_{L^1(\mathcal{U})},$$

for any $\varphi \in C_c^1((0, T))$, any $T > 0$ and any strongly confining admissible weight function ω .

The proof is a variant of the proof of [19, Proposition 3.7].

Proof of Proposition 4.6. We set $\Delta := \delta^{\alpha/(p\nu)} \langle v \rangle^{-1}$ and we observe that Proposition 4.5 writes

$$(4.8) \quad \|f\varphi\omega\Delta^\nu\|_{L^p(\mathcal{U})} \lesssim \|(\varphi + |\varphi'|)f\omega\|_{L^2(\mathcal{U})}.$$

From Corollary 4.3 and Hölder's inequality, we have next

$$\begin{aligned} \left\| f\varphi\omega \langle v \rangle^{\theta\varsigma/2} \delta^{-\beta(1-\theta)/2} \right\|_{L^2(\mathcal{U})} &\leq \left\| f\varphi\omega \langle v \rangle^{\varsigma/2} \right\|_{L^2(\mathcal{U})}^\theta \left\| f\varphi\omega \delta^{-\beta/2} \right\|_{L^2(\mathcal{U})}^{1-\theta} \\ &\lesssim \left\| (\varphi + \sqrt{|\varphi'|})f\omega \right\|_{L^2(\mathcal{U})}, \end{aligned}$$

for any $\theta \in (0, 1)$. Choosing $\theta = \theta_0$ such that

$$\frac{\beta}{\varsigma} \left(\frac{1}{\theta_0} - 1 \right) = \frac{\alpha}{p\nu},$$

and setting $\mu := \theta_0\varsigma/2$, we thus deduce

$$(4.9) \quad \|f\varphi\omega\Delta^{-\mu}\|_{L^2(\mathcal{U})} \lesssim \|(\varphi + \sqrt{|\varphi'|})f\omega\|_{L^2(\mathcal{U})}.$$

The Hölder inequality

$$(4.10) \quad \|h\|_{L^r(\mathcal{U})} \leq \|h\|_{L_{\sigma_p}^{1-\theta}(\mathcal{U})} \|h\|_{L_{\sigma_2}^\theta(\mathcal{U})},$$

with $1/r := (1-\theta)/p + \theta/2$ and $1 = \sigma_p^{1-\theta} \sigma_2^\theta$ for any $\theta \in (0, 1)$, implies

$$\begin{aligned} \|f\varphi\omega\|_{L^r(\mathcal{U})} &\lesssim \|f\varphi\omega\Delta^{\mu(\theta/(1-\theta))}\|_{L^p(\mathcal{U})}^{1-\theta} \|f\varphi\omega\Delta^{-\mu}\|_{L^2(\mathcal{U})}^\theta \\ &\lesssim \|(\varphi + |\varphi'|)f\omega\|_{L^2(\mathcal{U})}^{1-\theta_1} \|(\varphi + \sqrt{|\varphi'|})f\omega\|_{L^2(\mathcal{U})}^{\theta_1} \\ &\lesssim \|(\varphi + |\varphi'|)f\omega\|_{L^2(\mathcal{U})}, \end{aligned}$$

where we have chosen $\theta = \theta_1$ such that $\mu(\theta_1/(1 - \theta_1)) = \nu$, we have used the two inequalities (4.8) and (4.9) in the second line, and finally the most classical form of the Young inequality in the last line. We now use the Hölder inequality

$$\|(\varphi + |\varphi'|)f\omega\|_{L^2(\mathcal{U})} \leq \|\varphi f\omega\|_{L^r(\mathcal{U})}^{\theta_2} \|(\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} f\omega\|_{L^1(\mathcal{U})}^{1-\theta_2}$$

with $\theta_2 := r/(2r - 2)$ at the RHS of the last estimate. After simplification of the terms involving the L^r norm and taking the power $(1 - \theta_2)^{-1}$ of the resulting inequality, we conclude to (4.7). \square

For a weakly confining weight function ω , we have the following slightly weaker estimate.

Proposition 4.7. *Let $T > 0$ and ω a weakly confining admissible weight function, then there is $r > 2$ such that any solution f to the KFP equation (1.1)–(1.3) satisfies*

$$(4.11) \quad \|(\varphi + |\varphi'|)f\omega\|_{L^2(\mathcal{U})} \lesssim \|(\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} f\omega\langle v \rangle^K\|_{L^1(\mathcal{U})},$$

for any $\varphi \in C_c^1((0, T))$ and $K := 2(3d + 1)(2d + 3)$.

Proof of Proposition 4.6. We observe that for a weakly confining weight function ω , there always holds $\gamma \leq 2$, $s \leq 2 - \gamma$, and thus $\nu = \max(2, \gamma - 1) = 2$. The estimates established in Corollary 4.4 and Proposition 4.5 then write for instance

$$\left\| f\varphi\omega\langle v \rangle^{-1}\delta^{-\beta/2} \right\|_{L^2(\mathcal{U})} \leq C\|(\varphi + |\varphi'|)f\omega\|_{L^2(\mathcal{U})},$$

with $\beta := [2(d + 1)]^{-1}$, and

$$\left\| f\varphi\omega\langle v \rangle^{-2}\delta^{1+1/p} \right\|_{L^p(\mathcal{U})} \leq C\|(\varphi + |\varphi'|)f\omega\|_{L^2(\mathcal{U})},$$

with $1/p = 1/2 - 1/(2(2d + 1))$. Using an interpolation argument, we get

$$(4.12) \quad \|f\varphi\omega\langle v \rangle^{-\mu}\|_{L^r(\mathcal{U})} \leq C\|(\varphi + |\varphi'|)f\omega\|_{L^2(\mathcal{U})},$$

by choosing $r \in (2, p)$ and $\theta_1 \in (0, 1)$ in such a way that

$$\frac{1}{r} = \frac{\theta_1}{p} + \frac{1 - \theta_1}{2}, \quad (1 - \theta_1)\beta/2 = \theta_1(1 + 1/p)$$

and thus $\mu := 2\theta_1 + (1 - \theta_1) = \theta_1 + 1$. From the Hölder inequality, we also have

$$(4.13) \quad \|(\varphi + |\varphi'|)f\omega\|_{L^2(\mathcal{U})} \leq \|\varphi f\omega\langle v \rangle^{-\mu}\|_{L^r(\mathcal{U})}^{\theta_2} \|(\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} f\omega\langle v \rangle^K\|_{L^1(\mathcal{U})}^{1-\theta_2}$$

with $\theta_2 := r/(2r - 2)$ and $K = \frac{\theta_2}{1-\theta_2}\mu$. We now compute

$$p = \frac{2d + 1}{d}, \quad \theta_1 = \frac{2d + 1}{12d^2 + 20d + 5}, \quad r = \frac{12d^2 + 20d + 5}{6d^2 + 10d + 2},$$

so that

$$\begin{aligned} K &= \frac{\theta_2}{1 - \theta_2}\mu = \frac{r}{r - 2}(1 + \theta_1) = (12d^2 + 20d + 5) \left(1 + \frac{2d + 1}{12d^2 + 20d + 5} \right) \\ &= 12d^2 + 22d + 6 = 2(3d + 1)(2d + 3). \end{aligned}$$

The both last estimates together imply (4.11). \square

4.4. The $L^1 - L^p$ estimate on the dual problem. We consider the dual backward problem (2.21) for which we establish the same kind of estimate as for the forward KFP problem (1.1)–(1.3). In order to make the discussion simpler, we separate the analysis for strongly and weakly confining admissible weight function in this section.

Proposition 4.8. *There exist some exponent $r > 2$ such that any solution g to the dual backward problem (2.21) satisfies*

$$(4.14) \quad \|g\varphi m\|_{L^r(\mathcal{U})} \lesssim \|(\varphi + |\varphi'|)^{\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} gm\|_{L^1(\mathcal{U})},$$

for any test function $0 \leq \varphi \in \mathcal{D}((0, T))$, any $T > 0$ and any function $m = \omega^{-1}$ which is the inverse of a strongly confining admissible weight function ω .

Proof of Proposition 4.8. The proof follows the same steps as for the proof of Proposition 4.6 and we thus repeat it without too much details.

Step 1. An improved weighted L^2 estimate at the boundary. Let ω be an admissible weight function and define $m = \omega^{-1}$.

We define the modified weight function \tilde{m} by

$$\tilde{m}^2 := m_A^2 \left(1 + \frac{1}{4} \frac{n_x \cdot v}{\langle v \rangle^4} - \frac{1}{4D^{1/2}} \delta(x)^{1/2} \frac{n_x \cdot v}{\langle v \rangle^2} \right),$$

where m_A has been defined in (2.24) and $D = \sup_{x \in \Omega} \delta(x)$. Considering a solution g to the dual backward problem (2.21), multiplying the equation by $\Phi^2 := \varphi^2 \tilde{m}^2$ with $\varphi \in \mathcal{D}(0, T)$, and integrating in all the variables, we obtain

$$-\frac{1}{2} \int_{\Gamma} (\gamma g)^2 \Phi^2 (n_x \cdot v) + \frac{1}{2} \int_{\mathcal{U}} g^2 \mathcal{T} \Phi^2 = \int_{\mathcal{U}} g(\mathcal{E}^* g) \Phi^2,$$

with \mathcal{T} defined in (4.1).

Since $\tilde{m}^2 := m_A^2 (1 + \frac{1}{4} \frac{n_x \cdot v}{\langle v \rangle^4})$ on the boundary Γ , Step 1 of the proof of Lemma 2.3 implies that

$$-\frac{1}{2} \int_{\Gamma} (\gamma g)^2 \Phi^2 (n_x \cdot v) \geq 0.$$

Arguing exactly as in the proof of Proposition 4.1 and using the estimates from Step 2 of the proof of Lemma 2.3, we obtain

$$\int g^2 m^2 \left[\frac{(n_x \cdot v)^2}{\langle v \rangle \delta^{1/2}} + \langle v \rangle^\varsigma \right] \varphi^2 + \int |\nabla_v (gm)|^2 \varphi^2 \lesssim \int g^2 m^2 [|\partial_t \varphi^2| + \varphi^2].$$

Proceeding next as for Corollary 4.3 with the help of the interpolation Lemma 4.2, we deduce

$$(4.15) \quad \int g^2 \left(\frac{1}{\delta^\beta} + \langle v \rangle^\varsigma \right) m^2 \varphi^2 + \int |\nabla_v (gm)|^2 \varphi^2 \lesssim \int g^2 m^2 [|\partial_t \varphi^2| + \varphi^2],$$

for some $\beta > 0$.

Step 2. Downgraded weighted $L^2 - L^p$ estimate, $p > 2$. For $0 \leq \varphi \in \mathcal{D}(0, T)$, $0 \leq \chi \in \mathcal{D}(\Omega)$ and the weight function $m_0 = m \langle v \rangle^{-\nu}$ with $\nu > 0$ to be chosen later, the function $\bar{g} := g\varphi\chi m_0$ satisfies

$$-\partial_t \bar{g} - v \cdot \nabla_x \bar{g} = G$$

where $G = G_0 + \operatorname{div}_v G_1$ with

$$G_0 := -gm_0(\varphi' \chi + \varphi \nabla_x \chi \cdot v) - \varphi \chi (m_0 b - \nabla_v m_0) \cdot \nabla_v g + (c - \operatorname{div}_v b) \bar{g}$$

and

$$G_1 := \varphi \chi m_0 \nabla_v g.$$

Proceeding as in the proof of Proposition 4.5, we get first

$$\|\bar{g}\|_{L^p(\mathbb{R}^{2d+1})}^2 \lesssim \|\varphi' mg\|_{L^2(\mathcal{U})}^2 + \|\nabla \chi\|_{L^\infty}^2 \|\varphi mg\|_{L^2(\mathcal{U})}^2 + \|\varphi mg\|_{L^2(\mathcal{U})}^2.$$

for some $p \in (2, \infty)$, and next, by interpolation, we conclude with

$$(4.16) \quad \left\| g\varphi \frac{m}{\langle v \rangle^\nu} \delta^{\alpha/p} \right\|_{L^p(\mathcal{U})} \lesssim \|(\varphi + |\varphi'|)gm\|_{L^2(\mathcal{U})}$$

for some $\alpha > p > 2$ and $\nu \geq 0$.

Step 3. Up to the boundary $L^2 - L^r$ estimate, $r > 2$. Proceeding as during the proof of Proposition 4.6, we may use the estimates (4.15) and (4.16) together with the Hölder inequality (4.10) in order to obtain that there exists an exponent $r > 2$ such that any solution g to the dual problem (2.21) satisfies

$$\|g\varphi m\|_{L^r(\mathcal{U})} \lesssim \|(\varphi + |\varphi'|)gm\|_{L^2(\mathcal{U})}.$$

for any $\varphi \in C_c^1((0, T))$, $T > 0$. We conclude to (4.14) by using the Hölder inequality once more, exactly as during the proof of Proposition 4.6. \square

We now prove an estimate similar to Proposition 4.8 for weakly confining admissible weights in the spirit of Proposition 4.7

Proposition 4.9. *Let ω be a weakly confining admissible weight function and define $m = \omega^{-1}$. Any solution g to the dual backward problem (2.21) satisfies*

$$(4.17) \quad \|(\varphi + |\varphi'|)g\varphi m\|_{L^2(\mathcal{U})} \lesssim \|(\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} gm \langle v \rangle^K\|_{L^1(\mathcal{U})},$$

for any test function $0 \leq \varphi \in \mathcal{D}((0, T))$, any $T > 0$ and $K = 2(3d+1)(2d+3)$.

Proof of Proposition 4.9. The proof follows the same steps as for the proof of Proposition 4.8.

Step 1. A weighted L^2 estimate at the boundary. Let ω be an admissible weight function and define $m = \omega^{-1}$.

We define the modified weight function \tilde{m} by

$$\tilde{m}^2 := m_A^2 \left(1 + \frac{1}{4} \frac{n_x \cdot v}{\langle v \rangle^4} - \frac{1}{4D^{1/2}} \delta(x)^{1/2} \frac{n_x \cdot v}{\langle v \rangle^2} \right),$$

as done in the Step 1 of Proposition 4.8. Considering a solution g to the dual backward problem (2.21), indeed multiplying the equation by $\Phi^2 := \varphi^2 \tilde{m}^2$ with $\varphi \in \mathcal{D}(0, T)$, and integrating in all the variables, we obtain

$$-\frac{1}{2} \int_{\Gamma} (\gamma g)^2 \Phi^2 (n_x \cdot v) + \frac{1}{2} \int_{\mathcal{U}} g^2 \mathcal{T} \Phi^2 = \int_{\mathcal{U}} g(\mathcal{E}^* g) \Phi^2,$$

with \mathcal{T} defined in (4.1).

Since $\tilde{m}^2 := m_A^2 (1 + \frac{1}{4} \frac{n_x \cdot v}{\langle v \rangle^4})$ on the boundary Γ , Step 1 of the proof of Proposition 4.8 implies that

$$-\frac{1}{2} \int_{\Gamma} (\gamma g)^2 \Phi^2 (n_x \cdot v) \geq 0.$$

Arguing similarly as in the proof of Proposition 4.1 and using the estimates from Step 2 of the proof of Lemma 2.3 with the difference that, since ω is a weakly confining weight function, we will have $|\varpi_{\tilde{m}, 2}^{\mathcal{E}^*}| \lesssim |\varpi_{\tilde{m}, 2}^{\mathcal{E}^*}| < \infty$. Then we obtain

$$\int g^2 m^2 \frac{(n_x \cdot v)^2}{\langle v \rangle \delta^{1/2}} \varphi^2 + \int |\nabla_v(gm)|^2 \varphi^2 \lesssim \int g^2 m^2 [|\partial_t \varphi^2| + \varphi^2].$$

Proceeding next as during Corollary 4.4 we deduce

$$(4.18) \quad \int \frac{g^2}{\delta^\beta} m^2 \varphi^2 + \int |\nabla_v(gm)|^2 \varphi^2 \lesssim \int g^2 m^2 [|\partial_t \varphi^2| + \varphi^2],$$

for some $\beta = [2(d+1)]^{-1} > 0$.

Step 2. Downgraded weighted $L^2 - L^p$ estimate, $p > 2$. We remark that the computations from Proposition 4.8 hold for any admissible weight so proceeding similarly we get that there is some $p \in (2, \infty)$ such that

$$(4.19) \quad \left\| g\varphi \frac{m}{\langle v \rangle^\nu} \delta^{\alpha/p} \right\|_{L^p(\mathcal{U})} \lesssim \|(\varphi + |\varphi'|)gm\|_{L^2(\mathcal{U})}$$

for any $0 \leq \varphi \in \mathcal{D}(0, T)$ and some $\alpha > p > 2$ and $\nu \geq 0$.

Step 3. Up to the boundary $L^2 - L^r$ estimate, $r > 2$. Proceeding as during the proof of Proposition 4.7, we may use the estimates (4.18) and (4.19) together with the Hölder inequality (4.12) in order to obtain that there exists an exponent $r > 2$ such that any solution g to the dual problem (2.21) satisfies

$$\|\varphi gm \langle v \rangle^{-\mu}\|_{L^r(\mathcal{U})} \lesssim \|(\varphi + |\varphi'|)gm\|_{L^2(\mathcal{U})}.$$

for any $\varphi \in C_c^1((0, T))$, $T > 0$ and some $\mu > 0$. We conclude to (4.17) by using the Hölder inequality once more, exactly as in (4.13) during the proof of Proposition 4.7. \square

4.5. Conclusion of the proof of the ultracontractivity property. We now conclude the proof of Theorem 1.3 in several elementary and classical (after Nash's work) steps.

Proof of Theorem 1.3. Consider a strongly confining admissible weight function ω and denote by $m = \omega^{-1}$ its inverse. We split the proof into four steps.

Step 1. We first establish that there exist a constant $\eta > 0$ such that any solution f to the KFP equation (1.1)–(1.3)–(1.4) satisfies

$$(4.20) \quad \|f(T, \cdot)\|_{L_\omega^2(\mathcal{O})} \leq C_{12} T^{-\eta} \|f_0\|_{L_\omega^1(\mathcal{O})}, \quad \forall T \in (0, 1).$$

First indeed, from Proposition 4.6, there exist an exponent $r > 2$ such that

$$\|f\varphi\omega\|_{L^r(\mathcal{U})} \lesssim \left\| (\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} f\omega \right\|_{L^1(\mathcal{U})}.$$

Thanks to the estimate (2.5) from Lemma 2.1 for $p = 1$ and $p = r$ provided by Theorem 3.5, we have

$$\|\varphi\|_{L^r(0, T)} \|f_T\omega\|_{L^r(\mathcal{O})} \lesssim e^{\kappa T} \|f\varphi\omega\|_{L^r(\mathcal{U})}$$

and

$$\left\| (\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} f\omega \right\|_{L^1(\mathcal{U})} \lesssim e^{\kappa T} \left\| (\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} \right\|_{L^1(0, T)} \|f_0\omega\|_{L^1(\mathcal{O})}.$$

We choose $\varphi(t) := \psi(t/T)$ with $\psi \in C_c^1((0, 1))$ such that $0 \leq \psi \leq \mathbf{1}_{[1/4, 3/4]}$, $\psi \not\equiv 0$ and $|\psi'|^{2\frac{r-1}{r-2}} \psi^{-\frac{r}{r-2}} \in L^1(0, 1)$, which is possible by taking $\psi(s) := (s - 1/4)_+^n (3/4 - s)_+^n$ for $n > 0$ large enough, and we easily compute

$$\begin{aligned} \|\varphi\|_{L^r(0, T)} &= T^{1/r} \|\psi\|_{L^r(0, 1)}, & \|\varphi\|_{L^1(0, T)} &= T \|\psi\|_{L^1(0, 1)}, \\ \left\| |\varphi'|^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} \right\|_{L^1(0, T)} &= T^{1-2\frac{r-1}{r-2}} \left\| |\psi'|^{2\frac{r-1}{r-2}} \psi^{-\frac{r}{r-2}} \right\|_{L^1(0, 1)}. \end{aligned}$$

Gathering the three last estimates and the three last identities, we finally obtain

$$T^{1/r} \|f_T\omega\|_{L^r(\mathcal{O})} \lesssim \left(T + T^{1-2\frac{r-1}{r-2}} \right) \|f_0\omega\|_{L^1(\mathcal{O})},$$

which implies

$$\begin{aligned} \|f_T \omega\|_{L^r(\mathcal{O})} &\lesssim T^{-\frac{1}{r}-1-\frac{2}{r-2}} \left(T^{1+\frac{2}{r-2}} + 1 \right) \|f_0 \omega\|_{L^1(\mathcal{O})} \\ &\lesssim T^{-\frac{1}{r}-1-\frac{2}{r-2}} \|f_0 \omega\|_{L^1(\mathcal{O})}. \end{aligned}$$

Then by using again Lemma 2.1 for $p = 1$ and an interpolation argument as before, choosing $\theta := \frac{r}{2(r-1)} > 0$ such that $1/2 = 1 - \theta + \theta/r$, we deduce

$$\begin{aligned} \|f_T \omega\|_{L^2(\mathcal{O})} &\leq \|f_T \omega\|_{L^r(\mathcal{O})}^\theta \|f_T \omega\|_{L^1(\mathcal{O})}^{1-\theta} \\ &\lesssim T^{\theta(-\frac{1}{r}-1-\frac{2}{r-2})} \|f_0 \omega\|_{L^1(\mathcal{O})} \end{aligned}$$

from what we immediately conclude to (4.20) with $\eta := \frac{r}{2(r-1)} \left(\frac{1}{r} + 1 + \frac{2}{r-2} \right)$.

Step 2. Arguing in a similar fashion as above but using the estimates from Lemma 2.3 and Proposition 4.8 (instead of Lemma 2.1 and Proposition 4.6), we deduce a similar result for the dual problem. More precisely, there exists $\eta' > 0$ such that any solution g to the dual backward problem (2.21) satisfies

$$(4.21) \quad \|g(0, \cdot)\|_{L_m^2} \leq C_{12}^* T^{-\eta'} \|g_T\|_{L_m^1}, \quad \forall T \in (0, 1).$$

Step 3. As the dual counterpart of (4.21), we have that any solution f to the KFP equation (1.1)–(1.3)–(1.4) satisfies

$$(4.22) \quad \|f(T, \cdot)\|_{L_\omega^\infty} \leq C_{12}^* T^{-\eta'} \|f_0\|_{L_\omega^2}, \quad \forall T \in (0, 1).$$

Indeed we may argue by duality in the following way: For $f_0 \in L_\omega^2$ and $g_T \in L_m^1$ we denote respectively by f and g the solution to the primal (1.1)–(1.3)–(1.4) and dual (2.21) KFP problem, then we have

$$\begin{aligned} \|f(T, \cdot)\|_{L_\omega^\infty} &= \sup_{g_T \in L_m^1, \|g_T\|_{L_m^1} \leq 1} \int f(T, \cdot) g_T \\ &= \sup_{g_T \in L_m^1, \|g_T\|_{L_m^1} \leq 1} \int f_0 g(0, \cdot) \\ &\leq \|f_0\|_{L_\omega^2} \sup_{g_T \in L_m^1, \|g_T\|_{L_m^1} \leq 1} \|g(0, \cdot)\|_{L_m^2} \\ &\leq \|f_0\|_{L_\omega^2} C_{12}^* T^{-\eta'}, \end{aligned}$$

where we have used an usual representation formula in the first line, the duality formula (2.22) in the second line, the Cauchy-Schwarz inequality in the third line and estimate (4.21) in the last line.

Step 4. For $T \in (0, 1]$, the estimates (4.20) and (4.22) also write

$$\|f(T, \cdot)\|_{L_\omega^\infty} \leq C_{12}(T/2)^{-\eta'} \|f(T/2, \cdot)\|_{L_\omega^2}, \quad \|f(T/2, \cdot)\|_{L_\omega^2} \leq C_{12}^*(T/2)^{-\eta} \|f_0\|_{L_\omega^1},$$

so that

$$(4.23) \quad \|f(T, \cdot)\|_{L_\omega^\infty} \leq C_{12} C_{12}^* 2^{\eta+\eta'} T^{-\eta-\eta'} \|f_0\|_{L_\omega^1}.$$

which is nothing but (1.25) for $T \in (0, 1]$ with $\Theta := \eta + \eta'$, $\kappa = (\alpha + \alpha')/2$ and $C \geq C_{12} C_{12}^* 2^{\eta+\eta'}$. For $T > 1$ we write $f(T) = S_{\mathcal{L}}(1)f(T-1)$, then we compute

$$\begin{aligned} \|f(T, \cdot)\|_{L_\omega^\infty} &= \|S_{\mathcal{L}}(1)f(T-1, \cdot)\|_{L_\omega^\infty} \\ &\lesssim \|f(T-1, \cdot)\|_{L_\omega^1} \\ &\lesssim e^{\kappa(T-1)} \|f_0\|_{L_\omega^1}, \end{aligned}$$

where we have used (4.23) at the second line and Lemma 2.1 at the third one, which concludes the proof of (1.25) for $T > 1$. \square

We end this section by formulating a variant of Theorem 1.3 for weakly confining weight functions.

Proposition 4.10. *Let ω be a weakly confining admissible weight function such that $s > 0$ or $s = 0$ and $k > K + k^*$ for $K := 4(3d + 1)(2d + 3)$. Define $\omega_\infty := \omega^{1/2}$ if $s > 0$, $\omega_\infty := \omega \langle v \rangle^{-K}$ if $s = 0$. For any $T > 0$, there exists $\kappa, \eta > 0$ such that any solution f to the KFP equation (1.1)–(1.3) satisfies*

$$(4.24) \quad \|f(T, \cdot)\|_{L^\infty_\omega(\mathcal{O})} \lesssim e^{\kappa T} T^{-\eta} \|f_0\|_{L^1_\omega(\mathcal{O})}$$

Proof of Proposition 4.10. We adapt the first step of the proof of Theorem 1.3 by using the estimate established in Proposition 4.7 instead of Proposition 4.6 as well as the estimate (2.5) from Lemma 2.1 for $p = 1$ and $p = 2$ provided by Theorem 3.5.

Step 1. We set $\omega_2 = \omega^{3/4}$ if $s > 0$, $\omega_2 = \omega \langle v \rangle^{-K/2}$ if $s = 0$ and we prove first that

$$(4.25) \quad \|f(T, \cdot)\|_{L^2_{\omega_2}(\mathcal{O})} \lesssim T^{-\eta} \|f_0\|_{L^1_\omega(\mathcal{O})} \quad \forall T \in (0, T].$$

Indeed from Proposition 4.7 and the definitions of ω_∞ and ω_2 we have that there is $r > 2$, from Proposition 4.7, such that

$$(4.26) \quad \begin{aligned} \|(\varphi + |\varphi'|)f\omega_2\|_{L^2(\mathcal{U})} &\lesssim \left\| (\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} f\omega_2 \langle v \rangle^{K/2} \right\|_{L^1(\mathcal{U})}, \\ &\lesssim \left\| (\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} f\omega \right\|_{L^1(\mathcal{U})} \end{aligned}$$

for any $0 < \varphi \in C_c^1((0, T))$ and any $T > 0$.

From Lemma 2.1 for $p = 1$ and $p = 2$ we have

$$(4.27) \quad \|\varphi + |\varphi'|\|_{L^2(0, T)} \|f_T \omega_2\|_{L^2(\mathcal{O})} \leq e^{\kappa T} \|(\varphi + |\varphi'|)f\omega_2\|_{L^2(\mathcal{U})}$$

and

$$(4.28) \quad \left\| (\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} f\omega \right\|_{L^1(\mathcal{U})} \lesssim e^{\kappa T} \|(\varphi + |\varphi'|)^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}}\|_{L^1(0, T)} \|f_0 \omega\|_{L^1(\mathcal{O})}$$

As done during the proof of Theorem 1.3 we choose $\varphi(t) := \psi(t/T)$ with $\psi \in C_c^1((0, 1))$ such that $0 \leq \psi \leq \mathbf{1}_{[1/4, 3/4]}$, $\psi \not\equiv 0$ and $|\psi'|^2 \psi^{-\frac{r-1}{r-2}} \psi^{-\frac{r}{r-2}} \in L^1(0, 1)$, which is possible by taking $\psi(s) := (s - 1/4)_+^n (3/4 - s)_+^n$ for $n > 0$ large enough, and we easily compute

$$\begin{aligned} \|\varphi\|_{L^2(0, T)} &\geq T^{1/2} \|\psi + |\psi'|\|_{L^2(0, 1)}, \quad \|\varphi\|_{L^1(0, T)} = T \|\psi\|_{L^1(0, 1)}, \\ \left\| |\varphi'|^{2\frac{r-1}{r-2}} \varphi^{-\frac{r}{r-2}} \right\|_{L^1(0, T)} &= T^{1-2\frac{r-1}{r-2}} \left\| |\psi'|^{2\frac{r-1}{r-2}} \psi^{-\frac{r}{r-2}} \right\|_{L^1(0, 1)}. \end{aligned}$$

Then we deduce

$$(4.29) \quad \|f_T \omega_2\|_{L^2(\mathcal{O})} \lesssim T^{-\frac{1}{2} - \frac{r}{r-2}} (1 + T^{1 + \frac{r}{r-2}}) \|f_0 \bar{\omega}\|_{L^1(\mathcal{O})}$$

which is nothing but (4.25) for $T \in (0, 1]$ with $\eta = \frac{1}{2} + \frac{r}{r-2}$.

Step 2. Now we prove a similar result for the dual problem, for this we define $m_\infty = \omega_\infty^{-1}$ and $m_2 = \omega_2^{-1}$. Indeed by using Lemma 2.3 and Proposition 4.9 we prove, by arguing similarly as in the Step 1, that there exists $\eta' > 0$ such that any solution g to the dual backward problem (2.21) satisfies

$$(4.30) \quad \|g(0, \cdot)\|_{L^2_{m_2}} \leq C_{12}^* T^{-\eta'} \|g_T\|_{L^1_{m_\infty}}, \quad \forall T \in (0, 1).$$

Step 3. As the dual counterpart of (4.30), we have that any solution f to the KFP equation (1.1)–(1.3)–(1.4) satisfies

$$(4.31) \quad \|f(T, \cdot)\|_{L_{\omega_\infty}^\infty} \leq C_{12}^* T^{-\eta'} \|f_0\|_{L_{\omega_2}^2}, \quad \forall T \in (0, 1).$$

Indeed we may argue by duality as in the Step 3 of the proof of Theorem 1.3: For $f_0 \in L_{\omega_2}^2$ and $g_T \in L_{m_\infty}^1$ we denote respectively by f and g the solution to the primal (1.1)–(1.3)–(1.4) and dual (2.21) KFP problem, then we have

$$\begin{aligned} \|f(T, \cdot)\|_{L_{\omega_\infty}^\infty} &= \sup_{g_T \in L_{m_\infty}^1, \|g_T\|_{L_{m_\infty}^1} \leq 1} \int f(T, \cdot) g_T \\ &= \sup_{g_T \in L_{m_\infty}^1, \|g_T\|_{L_{m_\infty}^1} \leq 1} \int f_0 g(0, \cdot) \\ &\leq \|f_0\|_{L_{\omega_2}^2} \sup_{g_T \in L_{m_\infty}^1, \|g_T\|_{L_{m_\infty}^1} \leq 1} \|g(0, \cdot)\|_{L_{m_2}^2} \\ &\leq \|f_0\|_{L_{\omega_2}^2} C_{12}^* T^{-\eta'}, \end{aligned}$$

where we have used an usual representation formula in the first line, the duality formula (2.22) in the second line, the Cauchy-Schwarz inequality in the third line and estimate (4.30) in the last line.

Step 4. For $T \in (0, 1]$, the estimates (4.20) and (4.22) also write

$$\|f(T, \cdot)\|_{L_{\omega_\infty}^\infty} \leq C_{12} (T/2)^{-\eta'} \|f(T/2, \cdot)\|_{L_{\omega_2}^2}, \quad \|f(T/2, \cdot)\|_{L_{\omega_2}^2} \leq C_{12}^* (T/2)^{-\eta} \|f_0\|_{L_\omega^1},$$

so that

$$(4.32) \quad \|f(T, \cdot)\|_{L_{\omega_\infty}^\infty} \leq C_{12} C_{12}^* 2^{\eta+\eta'} T^{-\eta-\eta'} \|f_0\|_{L_\omega^1}.$$

which is nothing but (1.25) for $T \in (0, 1]$ with $\Theta := \eta + \eta'$ and $C \geq C_{12} C_{12}^* 2^{\eta+\eta'}$. For $T > 1$ we write $f(T) = S_{\mathcal{L}}(1) f(T-1)$, then we compute

$$\begin{aligned} \|f(T, \cdot)\|_{L_\omega^\infty} &= \|S_{\mathcal{L}}(1) f(T-1, \cdot)\|_{L_\omega^\infty} \\ &\lesssim \|f(T-1, \cdot)\|_{L_\omega^1} \\ &\lesssim e^{\kappa(T-1)} \|f_0\|_{L_\omega^1}, \end{aligned}$$

where we have used (4.32) at the second line and Lemma 2.1 at the third one, which concludes the proof of (4.24) for $T > 1$. \square

5. WELL-POSEDNESS IN A GENERAL FRAMEWORK

In this section, we establish the well-posedness of the KFP equation in a general weighted Lebesgue space framework and in a weighted Radon measures space framework. We deduce the existence and uniqueness of a family of fundamental solutions.

5.1. Additional a priori estimates in the L^1 framework. We recall that any solution f to the KFP equation (1.1) satisfies

$$(5.1) \quad \|f_t\|_{L_\omega^p(\mathcal{O})} \lesssim \|f_0\|_{L_\omega^p(\mathcal{O})}, \quad \forall t \in [0, T],$$

$$(5.2) \quad \|f_t\|_{L_{\omega(v)^\varsigma}^1(\mathcal{O})} \lesssim \|f_0\|_{L_\omega^1(\mathcal{O})},$$

$$(5.3) \quad \|\nabla_v f\|_{L_\omega^2(\mathcal{O})} \lesssim \|f_0\|_{L_\omega^2(\mathcal{O})},$$

$$(5.4) \quad \|f_t\|_{L_{\omega_r}^p(\mathcal{O})} \leq C_t \|f_0\|_{L_\omega^p(\mathcal{O})}, \quad \forall t \in (0, T],$$

for any admissible weight function ω , any exponent $1 \leq p \leq r \leq \infty$ and some admissible weight function ω_r from (2.5), from (2.18) and (2.20), from (3.13) and from Theorem 1.3 with

$\omega_r = \omega$ in the case of a strongly confining weight function ω or from Proposition 4.10 and a standard interpolation argument in the case of a weakly confining weight function ω . The two last estimates together immediately give

$$(5.5) \quad \|\nabla_v f\|_{L^2((t_0, T) \times \mathcal{O})} \leq C_{t_0, T} \|f_0\|_{L^2_\omega(\mathcal{O})}, \quad \forall T > t_0 > 0.$$

We will use an additional a priori estimate that we establish now. For further references, for $k > 0$, we define the functions T_k by

$$T_k(s) := \max(\min(s, k), -k).$$

Lemma 5.1. *Consider an admissible weight function ω and a solution f to the KFP equation (1.1)–(1.3)–(1.4) associated to an initial datum $0 \leq f_0 \in L^1_\omega$. There (at least formally) hold*

$$(5.6) \quad \|\nabla_v T_K(f/\mathcal{M})\mathcal{M}^{1/2}\|_{L^2(\mathcal{U})} \leq C_{T, K} \|f_0\|_{L^1_\omega}$$

and

$$(5.7) \quad \sup_{t \in [0, T]} \|f \mathbf{1}_{f \geq (k+1)\mathcal{M}}\|_{L^1(\mathcal{O})} \leq \|f_0 \mathbf{1}_{f_0 \geq k\mathcal{M}}\|_{L^1(\mathcal{O})} + C_T \|f_0\|_{L^1_\omega},$$

for any $T, K, k > 0$ and some constants $C_{T, K}$ and C_T .

Proof of Lemma 5.1. For a renormalizing function β , a positive weight function m and a nonnegative test function φ , we (at least formally) compute

$$\begin{aligned} \frac{d}{dt} \int_{\mathcal{O}} \beta\left(\frac{f}{m}\right) \varphi &= \int \beta'\left(\frac{f}{m}\right) \frac{\varphi}{m} \{-v \cdot \nabla_x f + \Delta_v f + b \cdot \nabla_v f + cf\} \\ &= \int_{\Sigma} (-v \cdot n_x) \beta\left(\frac{f}{m}\right) \varphi - \int \beta''\left(\frac{f}{m}\right) |\nabla_v\left(\frac{f}{m}\right)|^2 \varphi \\ &\quad + \int \left\{ (v \cdot \nabla_x \varphi) \beta\left(\frac{f}{m}\right) - (v \cdot \nabla_x m) \frac{f}{m} \beta'\left(\frac{f}{m}\right) \frac{\varphi}{m} \right\} + \int \frac{f}{m} \beta'\left(\frac{f}{m}\right) c \varphi \\ &\quad + \int \left\{ \alpha\left(\frac{f}{m}\right) \operatorname{div}_v\left(\frac{\varphi}{m} \nabla_v m\right) + \beta\left(\frac{f}{m}\right) \operatorname{div}_v\left(m \nabla_v \frac{\varphi}{m}\right) - \beta'\left(\frac{f}{m}\right) \frac{f}{m} \nabla_v \frac{\varphi}{m} \cdot \nabla_v m \right\} \\ &\quad + \int \left\{ \beta\left(\frac{f}{m}\right) \varphi [-\operatorname{div}_v b] - (b \cdot \nabla_v \varphi) \beta\left(\frac{f}{m}\right) + (b \cdot \nabla_v m) \frac{f}{m} \beta'\left(\frac{f}{m}\right) \frac{\varphi}{m} \right\}, \end{aligned}$$

with $\alpha'(s) = s\beta''(s)$. Assuming $\beta : \mathbb{R} \rightarrow \mathbb{R}_+$ even and convex such that $\beta(0) = \beta'(0) = 0$, the boundary term can be handled in the following way

$$\begin{aligned} \int_{\Sigma} (-v \cdot n_x) \beta\left(\frac{\gamma f}{m}\right) \varphi &= \int_{\Sigma_-} (v \cdot n_x)_- \beta\left(\frac{\mathcal{R}\gamma+f}{m}\right) \varphi - \int_{\Sigma_+} (v \cdot n_x)_+ \beta\left(\frac{\gamma+f}{m}\right) \varphi \\ &\leq \int_{\Sigma_-} (v \cdot n_x)_- \left\{ \iota_D \beta\left(\frac{\mathcal{D}\gamma+f}{m}\right) + \iota_S \beta\left(\frac{\mathcal{S}\gamma+f}{m}\right) \right\} \varphi - \int_{\Sigma_+} (v \cdot n_x)_+ \beta\left(\frac{\gamma+f}{m}\right) \varphi \\ &\leq \int_{\Sigma_+} (v \cdot n_x)_+ \iota_D \left\{ \beta\left(\mathcal{M} \frac{\widetilde{\gamma+f}}{m \circ \mathcal{V}_x}\right) \varphi \circ \mathcal{V}_x - \beta\left(\frac{\gamma+f}{m}\right) \varphi \right\} \\ &\quad + \int_{\Sigma_+} (v \cdot n_x)_+ \iota_S \left\{ \beta\left(\frac{\gamma+f}{m \circ \mathcal{V}_x}\right) \varphi \circ \mathcal{V}_x - \beta\left(\frac{\gamma+f}{m}\right) \varphi \right\}, \end{aligned}$$

where we have used the convexity of β in the second line and the change of variable $v \mapsto \mathcal{V}_x v$ on the last equality. Taking $m = \varphi := \mathcal{M}$, we get

$$\int_{\Sigma} (-v \cdot n_x) \beta\left(\frac{\gamma f}{\mathcal{M}}\right) \mathcal{M} \leq \int_{\Sigma_+} (v \cdot n_x)_+ \iota_D \left\{ \beta(\widetilde{\gamma+f}) - \beta\left(\frac{\gamma+f}{\mathcal{M}}\right) \right\} \mathcal{M} \leq 0,$$

where we have classically used the very definition of $\widetilde{\gamma_+ f}$ and the Jensen inequality in order to get the last inequality. With these choices of functions β , m and φ , the first identity simplifies

$$\begin{aligned} & \frac{d}{dt} \int_{\mathcal{O}} \beta\left(\frac{f}{\mathcal{M}}\right) \mathcal{M} \\ &= \int_{\Sigma} (-v \cdot n_x) \beta\left(\frac{f}{\mathcal{M}}\right) \mathcal{M} + \int [v \cdot \nabla_x \mathcal{M}] \left(\beta\left(\frac{f}{\mathcal{M}}\right) - \frac{f}{\mathcal{M}} \beta'\left(\frac{f}{\mathcal{M}}\right) \right) \\ & \quad - \int \beta''\left(\frac{f}{\mathcal{M}}\right) |\nabla_v\left(\frac{f}{\mathcal{M}}\right)|^2 \mathcal{M} + \int \alpha\left(\frac{f}{\mathcal{M}}\right) \Delta_v \mathcal{M} \\ & \quad + \int \left\{ \beta\left(\frac{f}{\mathcal{M}}\right) \mathcal{M} [-\operatorname{div}_v b] - (b \cdot \nabla_v \mathcal{M}) \beta\left(\frac{f}{\mathcal{M}}\right) + (b \cdot \nabla_v \mathcal{M} + c \mathcal{M}) \frac{f}{\mathcal{M}} \beta'\left(\frac{f}{\mathcal{M}}\right) \right\}. \end{aligned}$$

We finally particularize $\beta'' = \mathbf{1}_{[k, k+1]}$, $k \geq 0$, so that

$$0 \leq \beta(s), s\beta'(s), \alpha(s) \leq s, |s\beta'(s) - \beta(s)| \leq s \quad \forall s \in \mathbb{R}, k \geq 0.$$

Observing that

$$\frac{|v \cdot \nabla_x \mathcal{M}|}{\mathcal{M}} + \frac{|b \cdot \nabla_v \mathcal{M}|}{\mathcal{M}} + \frac{|\Delta_v \mathcal{M}|}{\mathcal{M}} + |c| + |\operatorname{div}_v b| \leq C \langle v \rangle^3,$$

we deduce with this choice of β that

$$\frac{d}{dt} \int_{\mathcal{O}} \beta\left(\frac{f}{\mathcal{M}}\right) \mathcal{M} + \int \beta''\left(\frac{f}{\mathcal{M}}\right) |\nabla_v\left(\frac{f}{\mathcal{M}}\right)|^2 \mathcal{M} \leq C \int |f| \langle v \rangle^3.$$

Because $\omega \gtrsim \langle v \rangle^3$, we may use (2.1) with $p = 1$ in order to bound the above RHS term and then both (5.6) and (5.7) follow. \square

5.2. Well-posedness in a L^p_ω framework. For further references, we note

$$\mathfrak{L}f := \partial_t f + \mathcal{L}f, \quad \mathfrak{L}^* \varphi := -\partial_t \varphi + \mathcal{L}^* \varphi,$$

where \mathcal{L} is defined in (1.10) and \mathcal{L}^* is defined in (2.6). We also denotes \mathcal{B}_3 the set of functions $\beta \in C^2(\mathbb{R})$ such that β' has compact support.

Theorem 5.2. *We consider an admissible weight function ω and an exponent $p \in [1, \infty]$. For any $f_0 \in L^p_\omega(\mathcal{O})$, there exists a unique function $f \in C(\mathbb{R}_+; L^p_\omega(\mathcal{O}))$ satisfying the estimates (5.1), (5.4), (5.5), (5.6), (5.7) and which is a renormalized solution to (1.1), that is*

$$(5.8) \quad \int_{\mathcal{U}} \left\{ \beta(f) \mathfrak{L}^* \varphi + \beta''(f) |\nabla_v f|^2 \varphi \right\} + \int_{\Gamma} \beta(\gamma f) \varphi n_x \cdot v \, dv d\sigma_x dt = \int_{\mathcal{O}} \beta(f_0) \varphi(0, \cdot) \, dv dx,$$

for any $\varphi \in \mathcal{D}(\bar{\mathcal{U}})$ and $\beta \in \mathcal{B}_3$. Furthermore, the one parameter family of mappings $S_{\mathcal{L}}(t) : L^p_\omega \rightarrow L^p_\omega$, defined by $S_{\mathcal{L}}(t)f_0 := f(t, \cdot)$ for $t \geq 0$ and $f_0 \in L^p_\omega$, is a positive semigroup of linear and bounded operators.

It is worth emphasizing that because $\beta \in \mathcal{B}_3$, we have $\operatorname{supp} \beta'' \subset [-K, K]$ for some $K > 0$ and thus

$$(5.9) \quad \beta''(g) |\nabla_v g|^2 = \beta''(g) \mathbf{1}_{|g| \leq K} |\nabla_v g|^2 = \beta''(g) |\nabla_v T_K(g)|^2.$$

Together with (5.6), that implies that the second term in (5.8) makes sense. Also observe that $\gamma f \in L^2_\omega(t_0, T) \times \Sigma, d\xi_2$ for any $0 < t_0 < T$, thanks to the trace Theorem 3.1, so that $\beta(\gamma f) \in L^\infty(\Gamma)$ and the boundary term makes sense. A similar result holds for the dual KFP equation (2.21).

Proof of Theorem 5.2. We split the proof into two steps.

Step 1. Existence part. Because of the linearity of the equation and the weak maximum principle, we may only consider nonnegative initial data (and solutions). We first assume $1 \leq p < 2$. We take $0 \leq f_{0,n} \in L^1_\omega \cap L^2_\omega$ such that $f_{n,0} \rightarrow f_0$ in L^p_ω . Thanks to Theorem 3.5, we may associate

a sequence (f_n) of solutions in $C_t(\omega) \cap L_t^2 H_\omega^1$ such that (f_n) satisfies uniformly in $n \geq 1$ the estimates (5.1), (5.3), (5.4) and (5.5). Because the equation is linear the function $f_n - f_m$ satisfies (5.1), namely

$$\sup_{t \in [0, T]} \|(f_n - f_m)(t)\|_{L_\omega^p} \leq \|f_{0,n} - f_{0,m}\|_{L_\omega^p},$$

and therefore (f_n) is a Cauchy sequence in $C([0, T]; L_\omega^p)$. We define $f := \lim f_n$. Similarly from (5.5), the sequence (f_n) is a Cauchy sequence in $L^2(t_0, T; H_\omega^1)$ for any $t_0 \in (0, T)$. We thus have

$$\int_{\mathcal{U}_{t_0}} \beta(f) \mathfrak{L}\varphi + \int_{\mathcal{U}_{t_0}} \beta''(f) |\nabla_v f|^2 \varphi + \int_{\Gamma_{t_0}} \beta(\gamma f) \varphi n_x \cdot v = \int_{\mathcal{O}} \beta(f(t_0)) \varphi(t_0, \cdot),$$

for any $t_0 > 0$ and $\varphi \in \mathcal{D}([0, T] \times \bar{\mathcal{O}})$, where $\mathcal{U}_{t_0} := (t_0, T) \times \mathcal{O}$ and $\Gamma_{t_0} := (t_0, T) \times \Sigma$. Because of (5.1) and (5.6), we may pass to the limit $t_0 \rightarrow 0$ and we get

$$\int_{\mathcal{U}} \beta(f) \mathfrak{L}\varphi + \int_{\mathcal{U}} \beta''(f) |\nabla_v f|^2 \varphi + \int_{\Gamma} \beta(\gamma f) \varphi n_x \cdot v = \int_{\mathcal{O}} \beta(f_0) \varphi(0, \cdot)$$

for any $\varphi \in \mathcal{D}([0, T] \times \bar{\mathcal{O}})$. The same holds for the dual KFP problem for $g_0 \in L_m^p$, $1 \leq p \leq 2$. By duality, we obtain the existence of a solution for any $1 \leq p \leq \infty$ for both problems.

Step 2. Uniqueness part. We consider two solutions f_i to the KFP equation in the sense of Theorem 5.2 and we set $f := f_2 - f_1$.

- Take $g_T \in C_c(\mathcal{O})$ and let us consider $g \in L^\infty([0, T] \times \mathcal{O}) \cap C([0, T] \times \mathcal{O})$ the solution associated to the backward dual problem (2.21) which existence is given by Theorem 3.5 and regularity is given by Theorem 3.5 and [30, Theorem 3]. Because $f_t \in L^2(\mathcal{O})$ for $t > 0$, we may use the Theorem 3.5 and thus write

$$\int_{\mathcal{O}} f(T) g_T dx dv = \int_{\mathcal{O}} f(t) g(t) dx dv, \quad \forall t \in (0, T).$$

- By construction, we have $f \in C([0, T]; L^1)$. We indeed have $f(t_k) \rightarrow f(t)$ a.e. on \mathcal{O} as a consequence of the continuity result in Theorem 3.1 applied to the function $g := \beta(f)$ with $0 \leq \beta(s) \leq |s|^{1/2}$. On the other hand, $\{f(t); t \in [0, T]\}$ is weakly L^1 relatively compact as a consequence of the L_ω^1 bound (5.1) (with $p = 1$) and the equi-integrability estimate (5.7) (recall that Ω is bounded). The claimed strong continuity follows. We may thus easily pass to the limit $t \rightarrow 0$ in the above formula in order to get

$$\int_{\mathcal{O}} f(T) g_T dx dv = \lim_{t \rightarrow 0} \int_{\mathcal{O}} f(t) g(t) dx dv = 0,$$

by using $f(0) = 0$ and $g \in L^\infty(0, T; L^\infty(\mathcal{O}))$. Because $g_T \in C_c(\mathcal{O})$ is arbitrary, we deduce that $f(T) = 0$ for any $T > 0$ and the uniqueness is proved. \square

5.3. Additional a priori estimates in a Radon measures framework. We present an additional a priori estimate which holds for nonnegative solutions in a $M_{\omega,0}^1$ framework. More precisely, we claim that any nonnegative solution f_t associated to an initial datum $0 \leq f_0 \in M_{\omega,0}^1$ satisfies, at least formally,

$$(5.10) \quad \limsup_{t \rightarrow 0} \int_{\mathcal{O}} f_t \phi^c \leq \int_{\mathcal{O}} f_0 \phi^c,$$

for any $\phi \in \mathcal{D}(\Omega)$ such that $0 \leq \phi \leq 1$, and we denote $\phi^c := 1 - \phi$. We indeed first observe that

$$\frac{d}{dt} \int_{\mathcal{O}} f \phi^c = h_1 - h_2,$$

with

$$h_1 := \int_{\mathcal{O}} \{(v \cdot \nabla_x \phi^c) + (c - \operatorname{div}_v b) \phi^c\} f$$

and

$$h_2 := \int_{\Sigma} \gamma f n_x \cdot v = \int_{\Sigma_+} (1 - \iota_S - \iota_D) \gamma_+ f n_x \cdot v,$$

so that $h_1 \in L^2(0, T)$ and $h_2 \geq 0$. For the first bound, on the one hand, we observe that $\omega \gtrsim \langle v \rangle^{r_0}$, $r_0 := 3d/2 + 1 + \gamma/2$ and $\omega \langle v \rangle^s \gtrsim \langle v \rangle^{r_2}$, $r_2 := 3\gamma/2$, so that from (5.1) and (5.2), we have

$$f \in L^\infty(0, T; L^1_{\omega}) \cap L^1(0, T; L^1_{\omega \langle v \rangle^s}) \subset L^\infty(0, T; L^1_{\omega_0}) \cap L^1(0, T; L^1_{\omega_2}) \subset L^2(0, T; L^1_{\omega_1}),$$

with $\omega_1 := \langle v \rangle^\gamma$. On the other hand, we have $\langle v \rangle + |c - \operatorname{div}_v b| \lesssim \omega_1$. Integrating in time the above differential equation and using the Cauchy-Schwarz inequality, we obtain

$$\int_{\mathcal{O}} f_t \phi^c = \int_{\mathcal{O}} f_0 \phi^c + H_1 - H_2,$$

with $H_1 \in C^{0,1/2}([0, T])$, $H_1(0) = 0$, and $H_2 \geq 0$. More precisely

$$\int_{\mathcal{O}} f_t \phi^c \leq \int_{\mathcal{O}} f_0 \phi^c + t^{1/2} C \|\phi\|_{W^{1,\infty}} \|f\|_{L^2_{\omega_1}(\mathcal{U})},$$

for any $t \in (0, T)$, and thus (5.10) follows. Now, for any weight function $m : \mathbb{R}^d \rightarrow (0, \infty)$ such that m is decreasing and $m\omega$ is increasing and any function $\psi \in \mathcal{D}(\mathbb{R}^d)$ such that $\mathbf{1}_{B_R} \leq \psi \leq 1$, we set $\chi := \phi\psi$, $\chi^c := 1 - \chi$ and we observe that

$$\frac{\chi^c}{m} = \frac{\phi^c}{m} \mathbf{1}_{B_R} + \left(\frac{\phi^c}{m} + \frac{\phi}{m} \psi^c \right) \mathbf{1}_{B_R^c},$$

so that

$$(5.11) \quad \int_{\mathcal{O}} f_t \frac{\chi^c}{m} \leq \frac{1}{m(R)} \int_{\mathcal{O}} f_t \phi^c + \frac{2}{(m\omega)(R)} \int_{\mathcal{O}} f_t \omega, \quad \forall t \geq 0$$

As a consequence of (5.10) and (5.1), we thus deduce

$$(5.12) \quad \limsup_{t \rightarrow 0} \int_{\mathcal{O}} f_t \frac{\phi^c}{m} \leq \frac{1}{m(R)} \int_{\mathcal{O}} f_0 \phi^c + \frac{2}{(m\omega)(R)} \int_{\mathcal{O}} f_0 \omega.$$

5.4. Radon measures solutions and fundamental solutions.

Theorem 5.3. *For any admissible weight function ω and any $f_0 \in M^1_{\omega,0}$, there exists a unique solution $f \in C(\mathbb{R}_+; M^1_{\omega}) \cap C((0, \infty); L^\infty)$ associated to the KFP equation (1.1)-(1.3) in a sense that we discuss below.*

As a consequence, for any $z_0 := (x_0, v_0) \in \mathcal{O}$, there exists a unique fundamental solution $F \in C(\mathbb{R}_+; M^1_{\omega}) \cap C((0, \infty); L^\infty)$ associated to the KFP equation (1.1)-(1.3) and the initial datum δ_{z_0} .

Proof of Theorem 5.3. Step 1. Existence. Because of the a priori estimates (5.1), (5.4), (5.5), (5.6), (5.7), we may proceed exactly as during the proof of Theorem 5.2, and we obtain without difficulty the existence of

$$f \in C([0, T]; \mathcal{D}'(\mathcal{O})) \cap L^\infty(0, T; L^1_{\omega}(\mathcal{O})) \cap L^\infty(t_0, T; L^\infty_{\omega}(\mathcal{O})), \quad \nabla_v f \in L^2(t_0, T; L^2_{\omega}(\mathcal{O})),$$

for any $0 < t_0 < T$ and any admissible weight function ω , which is a renormalisation solution on $(0, T) \times \bar{\mathcal{O}}$ and a weak solution $[0, T) \times \mathcal{O}$ corresponding to the initial condition f_0 . More precisely, we have both

$$(5.13) \quad \int_{\mathcal{U}} \{ \beta(f) \mathfrak{L}^* \varphi + \beta''(f) |\nabla_v f|^2 \varphi \} + \int_{\Gamma} \beta(\gamma f) \varphi n_x \cdot v \, dv d\sigma_x dt = 0,$$

for any $\varphi \in \mathcal{D}((0, T) \times \bar{\mathcal{O}})$ and any $\beta \in \mathcal{B}_3$, and

$$(5.14) \quad \int_{\mathcal{U}} f \mathfrak{L}^* \varphi = \int_{\mathcal{O}} \varphi(0, \cdot) f_0(dx dv),$$

for any $\varphi \in \mathcal{D}([0, T] \times \mathcal{O})$. By construction, we may also assume that f satisfies (5.10). Because of Theorem 3.5, for any weight function $m : \mathbb{R}^d \rightarrow (0, \infty)$ such that m is decreasing and $m\omega$ is increasing and converges to ∞ , and any

$$(5.15) \quad \varphi \in C([0, T]; L_m^1(\mathcal{O})) \cap L^\infty(0, T; L_m^\infty(\mathcal{O})), \quad \nabla_v \varphi \in L^2(0, T; L_m^2(\mathcal{O})),$$

solution to the backward dual KFP equation associated to a final datum $\varphi_T \in C_c(\mathcal{O})$, we have

$$(5.16) \quad \int_{\mathcal{O}} f(T)\varphi_T = \int_{\mathcal{O}} f(t)\varphi(t), \quad \forall t \in (0, T).$$

Step 2. Improved identity. We claim that (5.16) also holds for $t = 0$. From the weak formulation, we have $t \mapsto f(t) \in C([0, T]; \mathcal{D}'(\mathcal{O}))$. Because of the $L^\infty(0, T; M_\omega^1)$ bound, we deduce that

$$(5.17) \quad f(t) \rightharpoonup f_0 \quad \text{in } (C_c(\mathcal{O}))' \quad \text{as } t \rightarrow 0.$$

Because the solution to the backward dual problem also satisfies $0 \leq \varphi \in C([0, T] \times \mathcal{O})$ as a consequence of [30, Theorem 3] when $0 \leq \varphi_T \in C_c(\mathcal{O})$ and with the notations of Section 5.3, we may next write

$$\left| \int_{\mathcal{O}} f(t)\varphi(t) - \int_{\mathcal{O}} f_0\varphi(0) \right| \leq \left| \int_{\mathcal{O}} f(t)\varphi(t)\chi - \int_{\mathcal{O}} f_0\varphi(0)\chi \right| + \int_{\mathcal{O}} f(t)\varphi(t)\chi^c + \int_{\mathcal{O}} f_0\varphi(0)\chi^c,$$

for any $\chi \in \mathcal{D}(\mathcal{O})$ such that $0 \leq \chi \leq 1$. For the first term, we have

$$\left| \int_{\mathcal{O}} f(t)\varphi(t)\chi - \int_{\mathcal{O}} f_0\varphi(0)\chi \right| \leq \left| \int_{\mathcal{O}} f(t)(\varphi(t)\chi - \varphi(0)\chi) \right| + \left| \int_{\mathcal{O}} (f(t) - f_0)\varphi(0)\chi \right| \rightarrow 0,$$

as $t \rightarrow \infty$, thanks to $\varphi\chi \in C_c([0, T] \times \mathcal{O})$ and (5.17). Particularizing $\chi := \phi\psi$, with ϕ and ψ defined in Section 5.3, and using (5.15) and (5.12), we have

$$\begin{aligned} \limsup_{t \rightarrow 0} \int_{\mathcal{O}} f(t)\varphi(t)\chi^c &\lesssim \limsup_{t \rightarrow 0} \int_{\mathcal{O}} f(t)m^{-1}\chi^c \\ &\lesssim \frac{1}{m(R)} \int_{\mathcal{O}} f_0\phi^c + \frac{1}{(m\omega)(R)} \|f_0\|_{L_\omega^1}. \end{aligned}$$

Using the above estimates as well as (5.11) at time $t = 0$, we conclude with

$$\limsup_{t \rightarrow 0} \left| \int_{\mathcal{O}} f(t)\varphi(t) - \int_{\mathcal{O}} f_0\varphi(0) \right| \leq \frac{1}{m(R)} \int_{\mathcal{O}} f_0\phi^c + \frac{1}{(m\omega)(R)} \|f_0\|_{L_\omega^1}.$$

Assuming now that $\phi \geq \mathbf{1}_{\Omega_\varepsilon}$ and using the very definition of $f_0 \in M_{\omega,0}^1$, the RHS term vanishes in the limit $\varepsilon \rightarrow 0$ and $R \rightarrow \infty$. We may thus pass to the limit in (5.16), and we obtain

$$(5.18) \quad \int_{\mathcal{O}} f(T)\varphi_T = \int_{\mathcal{O}} f_0\varphi(0).$$

Step 3. Uniqueness. We consider two solutions f_1, f_2 associated to the same initial datum $0 \leq f_0 \in M_{\omega,0}^1$. From Step 2, we have

$$\int_{\mathcal{O}} f_1(T)\varphi_T = \int_{\mathcal{O}} f_2(T)\varphi_T,$$

for any $T > 0$ and $\varphi_T \in C_c(\mathcal{O})$, and thus $f_1 = f_2$. \square

6. ABOUT THE HARNACK INEQUALITY

In this section we establish a strong maximum principle for the solutions of the kinetic Fokker-Planck equation in the form of a Harnack inequality, which is very similar to those in [41, Theorem 2.15] and [2].

Theorem 6.1. *Consider a weak solution $0 \leq f \in L^2((0, T) \times \mathcal{O}) \cap L^2((0, T) \times \Omega; H^1(\mathbb{R}^d))$ to the Kinetic Fokker-Planck equation (1.1). For any $0 < T_0 < T_1 < T$ and $\varepsilon > 0$, there holds*

$$(6.1) \quad \sup_{\mathcal{O}_\varepsilon} f_{T_0} \leq C \inf_{\mathcal{O}_\varepsilon} f_{T_1},$$

for some constant $C = C(T_0, T_1, \varepsilon) > 0$, where we recall that \mathcal{O}_ε is defined in (1.22).

The proof will be obtained in two steps. In a first step we shall obtain a local version of the Harnack inequality, and then in a second step we shall use a chain argument in order to get (6.1). The local Harnack inequality is a direct consequence of [32, Theorem 5 & Proposition 12], see also [30] for previous results in that direction, which apply to super- and sub-solutions to the Kinetic Fokker-Planck equation with vanishing damping term

$$(6.2) \quad \partial_t g = \mathfrak{M}g := -v \cdot \nabla_x g + \Delta_v g + b \cdot \nabla_v g.$$

For the reader's convenience, we state these results now. For that purpose, for $r > 0$ and $z_0 := (t_0, x_0, v_0) \in \mathbb{R}^{1+2d}$ we define the set

$$Q_r(z_0) := \{(t, x, v) \in \mathbb{R}^{1+2d} \mid -r^2 < t - t_0 \leq 0, |x - x_0 - (t - t_0)v_0| < r^3, |v - v_0| < r\}$$

as well as the map

$$\mathcal{T}_{r, z_0} : (t, x, v) \mapsto (t_0 + r^2 t, x_0 + r^3 x + r^2 t v_0, v_0 + r v),$$

and we observe that $\mathcal{T}_{R, z_0}(Q_r(0, 0, 0)) = Q_{rR}(z_0)$.

Theorem 6.2. *Let $T > 0$.*

(1) *There exist $\zeta \in (0, 1)$ and $C_1 \in (0, \infty)$ such that for any $z_0 \in \mathcal{U}$, any $0 < R \leq 1$ with $\overline{Q_R(z_0)} \subset \mathcal{U}$, and any nonnegative weak super-solution g to (6.2), there holds*

$$(6.3) \quad \|g\|_{L^\zeta(\tilde{Q}_{\eta R}^-(z_0))} \leq C_1 \inf_{Q_{\eta R}(z_0)} g,$$

where $\eta = 1/40$ and $\tilde{Q}_{\eta R}^-(z_0) := \mathcal{T}_{R, z_0}(Q_\eta(-\tau, 0, 0)) = Q_{\eta R}(z_0 - (R^2 \tau, R^2 \tau v_0, 0))$ with $\tau := 19\eta^2/2$.

(2) *For any $z_0 \in \mathcal{U}$, any $0 < r < r' \leq 1$ such that $\overline{Q_{r'}(z_0)} \subset \mathcal{U}$, and any $\zeta > 0$, there is $C_3 > 0$ such that any nonnegative weak sub-solution h to (6.2) satisfies*

$$(6.4) \quad \|h\|_{L^\infty(Q_r(z_0))} \leq C_3 \|h\|_{L^\zeta(Q_{r'}(z_0))}.$$

Proof of Theorem 6.1. We split the proof into two steps.

Step 1: Local Harnack inequality. We claim that for any $z_0 \in \mathcal{U}$ and $0 < R \leq 1$ such that $\overline{Q_R(z_0)} \subset \mathcal{U}$, there exists a constant $C > 0$ such that

$$(6.5) \quad \sup_{\tilde{Q}_{\eta R/2}^-(z_0)} f \leq C \inf_{Q_{\eta R/2}(z_0)} f,$$

where $\tilde{Q}_{\eta R/2}^-(z_0) := Q_{\eta R/2}(z_0 - (R^2 \tau, R^2 \tau v_0, 0))$.

On the one hand, we take $\lambda > \|c\|_{L^\infty(Q_R(z_0))}$ and we set $g := e^{\lambda t} f$. The function g satisfies

$$\partial_t g = \mathfrak{M}g + (\lambda + c)g \geq \mathfrak{M}g \quad \text{in } Q_R(z_0),$$

so that g is a nonnegative weak super-solution to (6.2). We deduce from Theorem 6.2–(1) that

$$e^{\lambda t_0} e^{-\lambda R^2(\eta^2 + \tau)} \|f\|_{L^\zeta(\tilde{Q}_{\eta R}^-(z_0))} \leq \|g\|_{L^\zeta(\tilde{Q}_{\eta R}^-(z_0))} \leq C_1 \inf_{Q_{\eta R}(z_0)} g \leq C_1 e^{\lambda t_0} \inf_{Q_{\eta R}(z_0)} f.$$

On the other hand, the function $h := e^{-\lambda t} f$ with $\lambda > \|c\|_{L^\infty(Q_{\eta R}(z_0 - (R^2\tau, R^2\tau v_0, 0)))}$ satisfies

$$\partial_t h = \mathfrak{M}h + (c - \lambda)h \leq \mathfrak{M}h \quad \text{in } \tilde{Q}_{\eta R}^-(z_0) = Q_{\eta R}(z_0 - (R^2\tau, R^2\tau v_0, 0)).$$

Therefore h is a nonnegative weak sub-solution to (6.2), and thus we deduce from Theorem 6.2–(3) that

$$\begin{aligned} e^{-\lambda(t_0 - R^2\tau)} \|f\|_{L^\infty(Q_{\eta R/2}(z_0 - (R^2\tau, R^2\tau v_0, 0)))} &\leq \|h\|_{L^\infty(Q_{\eta R/2}(z_0 - (R^2\tau, R^2\tau v_0, 0)))} \\ &\leq C_3 \|h\|_{L^c(Q_{\eta R}(z_0 - (R^2\tau, R^2\tau v_0, 0)))} \\ &\leq e^{-\lambda(t_0 - R^2\tau)} e^{\lambda\eta^2 R^2} C_3 \|f\|_{L^c(Q_{\eta R}(z_0 - (R^2\tau, R^2\tau v_0, 0)))}. \end{aligned}$$

We conclude the local Harnack inequality (6.5) by gathering the two previous estimates.

Step 2: Proof of (6.1). Once the local Harnack inequality (6.5) holds, one can deduce (6.1) by following the second step in the proof of [41, Theorem 2.15], which uses the Harnack chain from [2]. \square

7. CONSTRUCTIVE ASYMPTOTIC ESTIMATE

7.1. An abstract constructive Krein-Rutman-Doblin-Harris theorem. We formulate a general abstract constructive Krein-Rutman-Doblin-Harris theorem in the spirit of the ones presented in the recent work [28, Section 6].

We consider a positive semigroup $S = (S_t) = (S(t))$ on a Banach lattice X , which means that X is a Banach space endowed with a closed positive cone X_+ (we write $f \geq 0$ if $f \in X_+$) and that S_t is a bounded linear mapping such that $S_t : X_+ \rightarrow X_+$ for any $t \geq 0$. We also assume that S is in duality with a dual semigroup S^* defined on a dual Banach lattice Y , with closed positive cone Y_+ . More precisely, we assume that $X \subset Y'$ or $Y \subset X'$, so that the bracket $\langle \phi, f \rangle$ is well defined for any $f \in X$, $\phi \in Y$, that $f \in X_+$ (resp. $\phi \geq 0$) iff $\langle \psi, f \rangle \geq 0$ for any $\psi \in Y_+$ (resp. iff $\langle \phi, g \rangle \geq 0$ for any $g \in X_+$) and that $\langle \phi, S(t)f \rangle = \langle S^*(t)\phi, f \rangle$, for any $f \in X$, $\phi \in Y$ and $t \geq 0$. We denote by \mathcal{L} the generator of S with domain $D(\mathcal{L})$ and by \mathcal{L}^* the generator of S^* with domain $D(\mathcal{L}^*)$. We are interested in the existence of positive eigenvectors for both \mathcal{L} and \mathcal{L}^* , and in their quantified exponential stability.

When $\|\cdot\|_k$ is a norm on X (resp. Y), we denote $X_k := (X, \|\cdot\|_k)$ (resp. $Y_k := (Y, \|\cdot\|_k)$). For $\psi \in Y_+$ and $g \in X_+$, we define the seminorms

$$[f]_\psi := \langle \psi, |f| \rangle, \quad \forall f \in X, \quad [\phi]_g := \langle |\phi|, g \rangle, \quad \forall \phi \in Y.$$

In order to obtain a very accurate and constructive description of the longtime asymptotic behavior of the semigroup S , we introduce additional assumptions.

- We first make the strong dissipativity assumption

$$(7.1) \quad \|S(t)f\|_k \leq C_0 e^{\lambda t} \|f\|_k + C_1 \int_0^t e^{\lambda(t-s)} \|S(s)f\|_0 ds,$$

$$(7.2) \quad \|S^*(t)\phi\|_k \leq C_0 e^{\lambda t} \|\phi\|_k + C_1 \int_0^t e^{\lambda(t-s)} \|S^*(s)\phi\|_0 ds,$$

for any $f \in X$, $\phi \in Y$, $t > 0$ and $k = 0, 1$, where $\lambda \in \mathbb{R}$, $C_i \in (0, \infty)$ and $\|\cdot\|_k$, $k = 0, 1$ denote two families of dual norms on X and Y such that $X_1 \subset X_0$ and $Y_1 \subset Y_0$. More precisely, we assume

$$(7.3) \quad \|f\|_0 \leq \|f\|_1, \quad \|\phi\|_0 \leq \|\phi\|_1, \quad |\langle \phi, f \rangle| \leq \|\phi\|_0 \|f\|_1, \quad |\langle \phi, f \rangle| \leq \|\phi\|_1 \|f\|_0,$$

for any $f \in X$ and $\phi \in Y$.

- We also assume for instance one of the two following conditions

$$(7.4) \quad \exists \lambda_0 \in \mathbb{R}, \lambda_0 > \lambda, \exists t_0 > 0, \exists f_0 \in X_+ \setminus \{0\}, \quad S(t_0)f_0 \geq e^{\lambda_0 t_0} f_0$$

or

$$(7.5) \quad \exists \lambda_0 \in \mathbb{R}, \lambda_0 > \lambda, \exists t_0 > 0, \exists \phi_0 \in Y_+ \setminus \{0\}, \quad S^*(t_0)\phi_0 \geq e^{\lambda_0 t_0} \phi_0,$$

and we refer to [28, Lemma 2.4] for variants of these conditions regarding the existence of positive supereigenvectors.

- Next, we make a slightly relaxed Doblin-Harris positivity assumption

$$(7.6) \quad S_T f \geq \eta_{\varepsilon, T} g_{\varepsilon} [S_{T_0} f]_{\psi_{\varepsilon}}, \quad \forall f \in X_+,$$

$$(7.7) \quad S_T^* \phi \geq \eta_{\varepsilon, T} \psi_{\varepsilon} [S_{T_0}^* \phi]_{g_{\varepsilon}}, \quad \forall \phi \in Y_+,$$

for any $T \geq T_1 > T_0 \geq 0$ and $\varepsilon > 0$, where $\eta_{\varepsilon, T} > 0$ and where (g_{ε}) and (ψ_{ε}) are two bounded and decreasing families of X_+ and Y_+ . We say that the above condition is relaxed because we possibly have $T_0 > 0$ while in the usual Doblin-Harris condition (7.6) or (7.7) holds with $T_0 = 0$.

- We finally assume the following compatibility interpolation like conditions

$$(7.8) \quad \|f\|_0 \leq \xi_{\varepsilon} \|f\|_1 + \Xi_{\varepsilon} [f]_{\psi_{\varepsilon}}, \quad \forall f \in X, \quad \varepsilon \in (0, 1],$$

$$(7.9) \quad \|\phi\|_0 \leq \xi_{\varepsilon} \|\phi\|_1 + \Xi_{\varepsilon} [\phi]_{g_{\varepsilon}}, \quad \forall \phi \in Y, \quad \varepsilon \in (0, 1],$$

for two positive families (ξ_{ε}) and (Ξ_{ε}) such that $\xi_{\varepsilon} \searrow 0$ and $\Xi_{\varepsilon} \nearrow \infty$ as $\varepsilon \searrow 0$.

It is worth pointing out that the above assumptions are written in a very symmetric way between the primal and dual spaces and semigroups. They are yet too rough for addressing the issue of the existence of positive eigenvectors. This existence problem is not our main purpose since it has been widely treated for instance in [28] (see also the references therein). Nevertheless, for keeping the presentation as self-contained as possible, we consider some strengthened (and quite natural) assumptions that allow us to derive the existence part, keeping in mind that many variants are possible and referring the interested reader to Sections 2 and 3 in [28].

- On the one hand, we assume that X_1 is a Banach space and

$$(7.10) \quad \|S(t)f\|_1 \leq C'_0 e^{\lambda' t} \|f\|_1,$$

for some $\lambda' \in \mathbb{R}$, some $C'_0 \geq 1$, any $f \in X$ and any $t \geq 0$. Of course this is a consequence of (7.1) and the Gronwall inequality with $\lambda' := \lambda + C_1$ and $C'_0 := C_0$ under the mild assumption that $t \mapsto \|S(t)f\|_1$ is a (everywhere defined) measurable and locally bounded function on \mathbb{R}_+ .

- On the other hand, instead of (7.2), we rather assume that

$$(7.11) \quad S^* = V + W * S^*$$

with

$$(7.12) \quad \|V(t)\phi\|_0 \leq C_0 e^{\lambda t} \|\phi\|_0, \quad \|W(t)\phi\|_1 \leq C_1 e^{\lambda t} \|\phi\|_0, \quad W \geq 0,$$

which is a variant of (7.2) for $k = 1$ and which obviously implies (7.2) for $k = 0$. We also assume that bounded sequences of Y_1 are weakly compact sequences in the $\sigma(Y_0, X_1)$ sense.

Theorem 7.1. *Consider a semigroup S on a Banach lattice X which satisfies the above conditions. Then, there exists a unique eigentriplet $(\lambda_1, f_1, \phi_1) \in \mathbb{R} \times X \times Y$ such that*

$$\mathcal{L}f_1 = \lambda_1 f_1, \quad f_1 \geq 0, \quad \mathcal{L}^* \phi_1 = \lambda_1 \phi_1, \quad \phi_1 \geq 0,$$

together with the normalization conditions $\|\phi_1\|_0 = 1$, $\langle \phi_1, f_1 \rangle = 1$. Furthermore, there exist some constructive constants $C \geq 1$ and $\lambda_2 < \lambda_1$ such that

$$(7.13) \quad \|S(t)f - \langle f, \phi_1 \rangle f_1 e^{\lambda_1 t}\|_1 \leq C e^{\lambda_2 t} \|f - \langle f, \phi_1 \rangle f_1\|_1$$

for any $f \in X$ and $t \geq 0$.

Let us make some few comments.

- The above result is a variant, and in some sense a consequence, of [28, Theorem 6.3], see also [48, Theorem 5.3] and [6, Theorem 2.1]. However, the set of assumptions here only involves the semigroups S and S^* and not the eigenelements (λ_1, f_1, ϕ_1) as it was the case in [28, Theorem 6.3]. That makes clearer the properties on the semigroup S really necessary to get the conclusions. The framework is very general and in particular it is not restricted to the measures space in duality with the bounded measurable functions space as it is the case in [6]. Our result is truly constructive what was not the case in the approach developed in [48].

- In the conservative case, namely $\lambda_1 = 0$, $\phi_1 \equiv 1 \in Y_1 \subset L^\infty$, and solely assuming (7.1) with $\lambda < 0$, (7.6), (7.8) and X_1 is a Banach space, the same conclusion (7.13) holds true by just following the same proof. Such a result is a general Banach lattice variant of the classical Doblin-Harris theorem available in the measures space in duality with the bounded measurable functions space framework, see [35, 16, 6, 29] for more details and references.

- It is also worth emphasizing that (7.7) with $\phi := \psi_\varepsilon$ implies

$$S_T^* \psi_\varepsilon \geq \eta_{\varepsilon, T} [S_{T_0}^* \psi_\varepsilon]_{g_\varepsilon} \psi_\varepsilon =: e^{\lambda'_0 T} \psi_\varepsilon,$$

what is a condition similar to (7.5). We may thus alternatively first assume (7.6), (7.7) and next assume that (7.1), (7.2) hold for some $\lambda < \lambda'_0$. In other words, our constructive Krein-Rutman-Doblin-Harris theorem is really a consequence of a suitable strong dissipativity condition and of a suitable positivity condition on both primal and dual semigroups together with a compatibility conditions over the several involved norms and seminorms. This strong dissipativity condition is automatically satisfied when the semigroup has appropriate smoothing effects (measured in terms of gain of regularity, exponent of integrability or weight function) as it is the case here for the kinetic Fokker-Planck equation (see Section 7.2 below) but can be not true for less regularizing semigroup as for the linear Boltzmann model for instance.

- An alternative natural way to formulate the Doblin-Harris positivity conditions (7.6), (7.7) is to rather assume a family of weak Harnack conditions

$$(7.14) \quad S_T f \geq g_\varepsilon \int_{T_0}^T [S_t f]_{\psi_\varepsilon} dt \quad \text{if } f \geq 0,$$

$$(7.15) \quad S_T^* \phi \geq \psi_\varepsilon \int_{T_0}^T [S_t^* \phi]_{g_\varepsilon} dt \quad \text{if } \phi \geq 0,$$

for some constants $T > T_0 \geq 0$, together with a family of supereigenvalues (or barrier) conditions

$$(7.16) \quad S_t^* \psi_\varepsilon \geq e^{\nu_\varepsilon t} \psi_\varepsilon, \quad S_t g_\varepsilon \geq e^{\nu_\varepsilon t} g_\varepsilon, \quad \forall t \geq T_0,$$

for any $\varepsilon > 0$ and some $\nu_\varepsilon \in \mathbb{R}$. Using the first inequality in (7.16), we find for $f \in X_+$

$$\int_{T_0}^T [S_t f]_{\psi_\varepsilon} dt = \int_{T_0}^T \langle f, S_t^* \psi_\varepsilon \rangle dt \geq \int_{T_0}^T \langle f, e^{\nu_\varepsilon t} \psi_\varepsilon \rangle dt = \int_{T_0}^T e^{\nu_\varepsilon t} dt [f]_{\psi_\varepsilon},$$

and we thus immediately deduce (7.6) (with $T_0 = 0$) from (7.14) and (7.16). We may similarly deduce (7.7) (with $T_0 = 0$) from (7.15) and (7.16).

- We briefly discuss the link between our set of hypotheses and the strong maximum principle which is also classically used in the Krein-Rutman theory. For that purpose, we introduce the notion of strict positivity by writing $f \in X_{++}$ or $f > 0$ (resp. $\phi \in Y_{++}$ or $\phi > 0$) if $\langle \psi, f \rangle > 0$ for any $\psi \in Y_+ \setminus \{0\}$ (resp. $\langle \phi, g \rangle > 0$ for any $g \in X_+ \setminus \{0\}$). Under assumptions (7.8) and (7.9), we claim that (7.6) with $T_0 = 0$, or (7.6) and (7.7) together with (7.4), imply the classical strong maximum principle, and we recall that this last one classically writes

$$(7.17) \quad f \in D(\mathcal{L}) \cap X_+ \setminus \{0\}, \quad \mu \in \mathbb{R}, \quad (\mu - \mathcal{L})f =: g \geq 0 \quad \text{implies} \quad f > 0.$$

Before proving this claim we establish the following elementary facts

- (i) $f \in X_+ \setminus \{0\}$ iff $f \in X_+$ and $\langle \psi_\varepsilon, f \rangle > 0$ for any $\varepsilon \in (0, \varepsilon_f)$, $\varepsilon_f > 0$ small enough,
- (ii) $f \geq \alpha_\varepsilon g_\varepsilon$, $\alpha_\varepsilon > 0$, for any $\varepsilon > 0$, imply $f > 0$,

as consequences of (7.8) and (7.9).

One the one hand, for any fixed $f \in X \setminus \{0\}$ the family of interpolation estimates (7.8) implies

$$0 < \frac{1}{2} \|f\|_0 \leq \|f\|_0 - \xi_\varepsilon \|f\|_1 \leq \Xi_\varepsilon [f]_{\psi_\varepsilon},$$

$\varepsilon \in (0, \varepsilon_f)$, $\varepsilon_f > 0$ small enough, which gives (i). In particular, $\psi_\varepsilon \neq 0$ for $\varepsilon > 0$ small enough (what can also be added as an assumption in the definition of ψ_ε !). We similarly have $[\phi]_{g_\varepsilon} > 0$ for any $\phi \in Y_+ \setminus \{0\}$ and any $\varepsilon \in (0, \varepsilon_\phi)$, $\varepsilon_\phi > 0$ small enough, and thus $g_\varepsilon \neq 0$ for $\varepsilon > 0$ small enough. In particular, we have

$$\langle \psi_\varepsilon, g_\varepsilon \rangle \geq \langle \psi_\varepsilon, g_{\varepsilon_0} \rangle \geq \Xi_\varepsilon^{-1} \frac{1}{2} \|g_{\varepsilon_0}\|_0 > 0,$$

for any $\varepsilon \in (0, \varepsilon_{g_{\varepsilon_0}})$ and $\varepsilon_0 > 0$ such that $g_{\varepsilon_0} \neq 0$. Assume now $f \geq \alpha_\varepsilon g_\varepsilon$, $\alpha_\varepsilon > 0$, for any $\varepsilon > 0$. For any $\phi \in Y_+ \setminus \{0\}$, we then have

$$\langle \phi, f \rangle \geq \alpha_{\varepsilon_\phi} \langle \phi, g_{\varepsilon_\phi} \rangle > 0.$$

We have established that $f > 0$, and thus (ii) is proved.

We come now to the proof of the strong maximum principle and thus consider (μ, f, g) satisfying the requirements of (7.17). We fix ν strictly larger than μ and strictly larger than the growth bound of S so that we may write

$$f = (\nu - \mathcal{L})^{-1}((\nu - \mu)f + g) = \int_0^\infty e^{-\nu t} S_t((\nu - \mu)f + g) dt$$

to get by positivity of g

$$(7.18) \quad f \geq (\nu - \mu) \int_0^\infty e^{-\nu t} S_t f dt.$$

From (7.6) with $T_0 = 0$, we deduce

$$(7.19) \quad f \geq g_\varepsilon (\nu - \mu) \int_{T_1}^\infty e^{-\nu t} \eta_{\varepsilon, t} dt \langle \psi_\varepsilon, f \rangle, \quad \forall \varepsilon > 0.$$

From (i) above and because $f \in X_+ \setminus \{0\}$, there exists $\varepsilon_f > 0$ such that $\langle \psi_\varepsilon, f \rangle > 0$ for any $\varepsilon \in (0, \varepsilon_f)$. Together with (7.19), we deduce that $f \geq \alpha_\varepsilon g_\varepsilon$, with $\alpha_\varepsilon > 0$, for any $\varepsilon \in (0, \varepsilon_f)$, and that in turn implies $f > 0$ from (ii) above.

When (7.6) and (7.7) are satisfied with $T_0 > 0$, we may derive (7.17) by using the additional condition (7.4). We apply (7.6) with $T = t - T_1 + T_0$ to the vector $S_{T_1 - T_0} f$ to get that

$$S_t f \geq \eta_{\varepsilon, t - T_1 + T_0} g_\varepsilon \langle \psi_\varepsilon, S_{T_1} f \rangle$$

for any $t \geq 2T_1 - T_0$. Injecting this inequality in (7.18), we obtain

$$\begin{aligned} f &\geq g_\varepsilon (\nu - \mu) \int_{2T_1 - T_0}^\infty e^{-\nu t} \eta_{\varepsilon, t - T_1 + T_0} dt \langle \psi_\varepsilon, S_{T_1} f \rangle \\ &\geq g_\varepsilon (\nu - \mu) \int_{2T_1 - T_0}^\infty e^{-\nu t} \eta_{\varepsilon, t - T_1 + T_0} dt \langle S_{T_1}^* \psi_\varepsilon, f \rangle, \end{aligned}$$

for any $\varepsilon > 0$. Together with (7.7), that implies

$$(7.20) \quad f \geq g_\varepsilon (\nu - \mu) \int_{2T_1 - T_0}^\infty e^{-\nu t} \eta_{\varepsilon, t - T_1 + T_0} dt \eta_{\varepsilon', T_1} \langle S_{T_0}^* \psi_\varepsilon, g_{\varepsilon'} \rangle \langle \psi_{\varepsilon'}, f \rangle.$$

On the other hand, taking $n \in \mathbb{N}$ large enough so that $nt_0 \geq T_0$ and iterating (7.4), we get that $S_{nt_0} f_0 \geq e^{nt_0} f_0$ and as a consequence $S_{nt_0} f_0 \in X_+ \setminus \{0\}$. We infer that necessarily $S_{T_0} f_0 \in X_+ \setminus \{0\}$, and the existence of ε_0 such that $\langle S_{T_0}^* \psi_\varepsilon, f_0 \rangle = \langle \psi_\varepsilon, S_{T_0} f_0 \rangle > 0$ for all

$\varepsilon \in (0, \varepsilon_0)$. We thus deduce in particular $S_{T_0}^* \psi_\varepsilon \in X_+ \setminus \{0\}$, and since we also have $f \in X_+ \setminus \{0\}$, we deduce the existence of $\varepsilon' > 0$ such that $\langle S_{T_0}^* \psi_\varepsilon, g_{\varepsilon'} \rangle > 0$ and $\langle \psi_{\varepsilon'}, f \rangle > 0$. Coming back to (7.20), we have proved, for any $\varepsilon \in (0, \varepsilon_0)$, the existence of $\alpha_\varepsilon > 0$ such that $f \geq \alpha_\varepsilon g_\varepsilon$ and this guarantees that $f > 0$.

Symmetrically, the assumptions (7.6), (7.7) and (7.5) imply that the dual operator \mathcal{L}^* satisfies the strong maximum principle.

In particular, we deduce that the first eigenvectors exhibited in Theorem 7.1 satisfies $f_1 > 0$ and $\phi_1 > 0$.

7.2. Application to the KFP equation: proof of Theorem 1.2. In this section, we consider the kinetic Fokker-Planck equation (1.1), (1.3), (1.4) and we prove Theorem 1.2 by using Theorem 7.1. We define $X_0 := L_{\omega_2}^2$ for a strongly confining admissible exponential weight function ω_2 and $X = X_1 := L_{\omega_r}^r$ with $r \in (2, \infty)$ given by Proposition 4.6 and an admissible exponential weight function ω_r such that $L_{\omega_r}^r \subset L_{\omega_2}^2$. We next define $Y_0 := L_{m_r}^{r'}$ with r' the conjugate exponent associated to r , $m_r := \omega_r^{-1}$ and $Y_1 := L_{m_2}^2$, with $m_2 := \omega_2^{-1}$. Many other choices are possible. This choice however contrasts with the usual $L^1 - L^\infty$ framework considered when using Doblin-Harris type arguments.

We now check that the assumptions of Theorem 7.1 are met.

We recall that \mathcal{O}_ε has been defined in (1.22) and we denote $\varepsilon_0 > 0$ such that $\mathcal{O}_{\varepsilon_0} \neq \emptyset$.

• *A positive supereigenvector condition.* For a given function $0 \leq h_0 \in C_c^2(\mathcal{O})$ normalized by $\|h_0\|_{L_{\omega_2}^2} = 1$ and such that $\text{supp } h_0 \subset \mathcal{O}_{\varepsilon_0}$, and for $\lambda > \omega(S_{\mathcal{L}})$ the growth bound of $S_{\mathcal{L}}$, we define $f_0 \in D(\mathcal{L})$ as the solution to

$$(\lambda - \mathcal{L})f_0 = h_0 \text{ in } \mathcal{O}, \quad \gamma_- f_0 = \mathcal{R}\gamma_+ f_0 \text{ on } \Sigma_-.$$

The existence and uniqueness of such a solution is a classical consequence of the existence of the semigroup $S_{\mathcal{L}}$ given by Proposition 3.3. Repeating the proof of the condition (H2) in [28, Section 11.4], there exists a constructive constant $c > 0$ such that $f_0 \geq c h_0$. Coming back to the equation, we have

$$\mathcal{L}f_0 = \lambda f_0 - h_0 \geq (\lambda - c^{-1})f_0 \text{ in } \mathcal{O},$$

which is a variant of (7.4), and in particular from [28, Remark 2.5], it implies (7.4) with $\lambda_0 := (\lambda - c^{-1})$.

• *Strong dissipativity conditions.* We define

$$\mathcal{B} := \mathcal{L} - \mathcal{A}, \quad \mathcal{A}f := M\chi_R f, \quad \mathbf{1}_{B_R} \leq \chi_R \leq \mathbf{1}_{B_{2R}},$$

with $B_r := \{v \in \mathbb{R}^d; |v| \leq r\}$ and χ_R a smooth function. We then define the semigroup $S_{\mathcal{B}}$ associated to \mathcal{B} and the reflection condition (1.3) which existence is given by Proposition 3.3. We claim that for any $a^* \in \mathbb{R}$, we may choose $M, R > 0$ large enough in such a way that $S_{\mathcal{B}}$ satisfies

$$(7.21) \quad \|S_{\mathcal{B}}(t)f\|_{L_\omega^p} \leq e^{a^* t} \|f\|_{L_\omega^p}, \quad \forall t \geq 0, \forall f \in L_\omega^p,$$

for any Lebesgue space L_ω^p with admissible exponential weight function ω . Coming back indeed to the proof of Lemma 2.1, for $p = 1, 2$, and more precisely to (2.18), the function $f(t) := S_{\mathcal{B}}(t)f_0$ satisfies

$$\frac{d}{dt} \int_{\mathcal{O}} f^p \tilde{\omega}^p \leq \int_{\mathcal{O}} f^p \tilde{\omega}^p \varpi^{\mathcal{B}},$$

with $\varpi^{\mathcal{B}} := \varpi_{\tilde{\omega}, p}^{\mathcal{L}} - M\chi_R$. Because of (2.20), we may thus fix $M, R > 0$ large enough, in such a way that $\varpi^{\mathcal{B}} \leq a^*$. That implies (7.21) for $p = 1, 2$. We deduce that (7.21) holds for any $p \in [1, \infty]$ as we proved the similar growth estimate for $S_{\mathcal{L}}$, and we thus refer to Theorem 3.5 for more details.

On the other hand, from Theorem 1.3 applied to the semigroup $S_{\mathcal{B}}$, we know that

$$(7.22) \quad \|S_{\mathcal{B}}(t)f\|_{L_{\omega_r}^2} \lesssim \frac{e^{Ct}}{t^{\Theta}} \|f\|_{L_{\omega_2}^2}.$$

We finally recall that from Theorem 1.1, we have

$$(7.23) \quad \|S_{\mathcal{A}}(t)f\|_{L_{\omega_q}^q} \leq e^{Ct} \|f\|_{L_{\omega_q}^q},$$

for the exponents $q = 2, r$, for any admissible weight function ω_q , for any $f \in L_{\omega_q}^q$ and any $t \geq 0$.

Iterating the Duhamel formula

$$S_{\mathcal{A}} = S_{\mathcal{B}} + S_{\mathcal{B}}\mathcal{A} * S_{\mathcal{A}},$$

we get

$$S_{\mathcal{A}} = \mathcal{V} + \mathcal{W} * S_{\mathcal{A}},$$

with

$$\mathcal{V} := S_{\mathcal{B}} + \dots + (S_{\mathcal{B}}\mathcal{A})^{*(N-1)} * S_{\mathcal{B}}, \quad \mathcal{W} := (S_{\mathcal{B}}\mathcal{A})^{*N}.$$

Here $*$ denotes the usual convolution operation for functions defined on \mathbb{R}_+ and we define recursively $U^{*1} = U$, $U^{*k} = U^{*(k-1)} * U$. Combining (7.22) and (7.23), we may use [47, Proposition 2.5] (see also [31, 46]) and we deduce that

$$\|\mathcal{V}(t)f\|_{L_{\omega_r}^r} \lesssim e^{at} \|f\|_{L_{\omega_r}^r}, \quad \|\mathcal{W}(t)f\|_{L_{\omega_r}^r} \lesssim e^{at} \|f\|_{L_{\omega_2}^2},$$

for $N \geq 1$ large enough, any $a > a^*$, any $t \geq 0$ and any $f \in L_{\omega_r}^r$, and thus

$$(7.24) \quad \|S_{\mathcal{A}}(T)f_0\|_{L_{\omega_r}^r} \leq Ce^{aT} \|f_0\|_{L_{\omega_r}^r} + Ce^{CT} \|f_0\|_{L_{\omega_2}^2},$$

for any $t \geq 0$ and $f_0 \in L_{\omega_r}^r$. That is nothing but (7.1) for $k = 1$. The same estimate for $k = 0$ is clear, it is nothing but Lemma 2.1. The proof of (7.11)-(7.12) is similar and it is thus skipped.

• *Doblin-Harris condition.* Let us fix $0 \leq f_0 \in L_{\omega_2}^2$ and denote $f_t := S_{\mathcal{A}}(t)f_0$. For $T_0 > 0$ and $\varepsilon \in (0, \varepsilon_0)$, we know from Theorem 6.1 that for any $T_1 > T_0 > 0$ and for every $T \geq T_1$, we have

$$\sup_{\mathcal{O}_\varepsilon} f_{T_0} \leq C \inf_{\mathcal{O}_\varepsilon} f_T,$$

for a constant C independent of f . We deduce

$$\begin{aligned} f_T &\geq (\inf_{\mathcal{O}_\varepsilon} f_T) \mathbf{1}_{\mathcal{O}_\varepsilon} \\ &\geq \frac{1}{C} (\sup_{\mathcal{O}_\varepsilon} f_{T_0}) \mathbf{1}_{\mathcal{O}_\varepsilon} \\ &\geq \frac{1}{C} \frac{1}{|\mathcal{O}_\varepsilon|} \langle S_{T_0} f_0, \mathbf{1}_{\mathcal{O}_\varepsilon} \rangle \mathbf{1}_{\mathcal{O}_\varepsilon} \end{aligned}$$

what is (7.6) with $g_\varepsilon = \psi_\varepsilon := \mathbf{1}_{\mathcal{O}_\varepsilon}$. The proof of (7.7) is identical.

• *The interpolation condition.* Let us consider two exponents $p > q$ and two locally bounded weight functions ω_p, ω_q such that $\omega_p \geq \omega_q$ and $\omega_q/\omega_p \in L^{qr'}$ with $r := p/q \in (1, \infty)$, and in particular $L_{\omega_p}^p \subset L_{\omega_q}^q$. We have

$$\begin{aligned} \|f\|_{L_{\omega_q}^q} &\leq \|f\mathbf{1}_{\mathcal{O}_\varepsilon}\|_{L_{\omega_q}^q} + \|f\omega_q\mathbf{1}_{\mathcal{O}_\varepsilon^c}\|_{L^q} \\ &\leq \|f\mathbf{1}_{\mathcal{O}_\varepsilon}\|_{L_{\omega_q}^p}^\theta \|f\mathbf{1}_{\mathcal{O}_\varepsilon}\|_{L_{\omega_q}^1}^{1-\theta} + \|f\omega_p\|_{L^p} \|\omega_q/\omega_p\mathbf{1}_{\mathcal{O}_\varepsilon^c}\|_{L^{r'q}} \\ &\leq (\varepsilon^{\frac{1}{\theta}} + \|\omega_q/\omega_p\mathbf{1}_{\mathcal{O}_\varepsilon^c}\|_{L^{r'q}}) \|f\|_{L_{\omega_p}^p} + \varepsilon^{\frac{1}{\theta-1}} \|f\mathbf{1}_{\mathcal{O}_\varepsilon}\|_{L_{\omega_q}^1}, \end{aligned}$$

where we have used the classical interpolation inequality (with $1/q = \theta/p + 1 - \theta$, $\theta \in (0, 1)$) and the Holder inequality in the second line, and the Young inequality in the third line. That implies both (7.8) and (7.9).

The same conditions hold for the dual problem, so that we may apply Theorem 7.1 in order to conclude that Theorem 7.1 holds in the space $X_1 = L^1_\omega$. We deduce that Theorem 7.1 holds in any weighted Lebesgue spaces associated to admissible weight functions by using the extension trick as developed in [31, 46] to which we refer for details. It is also worth emphasizing that the uniform estimates in (1.23) directly follow from the ultracontractivity estimate established in Theorem 1.3 for the primal and the dual semigroups and that the strict positivity properties in (1.23) directly follow from the discussion about the strong maximum principle just after the statement of Theorem 7.1. Furthermore, $f_1, \phi_1 \in C(\mathcal{O})$ as a direct consequence of Theorem 3.5 and [30, Theorem 3].

7.3. Proof of Theorem 7.1. This section is devoted to the proof of Theorem 7.1 which is split into six steps. We closely follow the material presented in [28, Section 2,3] (see also [44]) in Steps 1, 2, 3, 4 and the material presented in [16] in Steps 5 and 6.

Step 1. Existence of ϕ_1 . From the fact that S^* is a positive semigroup, (7.5) and [28, Lemma 2.6], we know that

$$\lambda_1 := \inf\{\kappa \in \mathbb{R}; z - \mathcal{L} \text{ is invertible for any } z \geq \kappa\} \geq \lambda_0$$

and

$$(7.25) \quad \exists \lambda_n \searrow \lambda_1, \exists \hat{\phi}_n \in D(\mathcal{L}^*) \cap Y_+, \varphi_n := \lambda_n \hat{\phi}_n - \mathcal{L}^* \hat{\phi}_n \geq 0, \|\hat{\phi}_n\|_0 = 1, \|\varphi_n\|_0 \rightarrow 0.$$

Because $\lambda_n > \lambda_1$, the following representation formula

$$\hat{\phi}_n := (\lambda_n - \mathcal{L}^*)^{-1} \varphi_n = \int_0^\infty S^*(t) e^{-\lambda_n t} \varphi_n dt$$

holds true. Introducing the sequences

$$v_n := \mathcal{V}_n \varphi_n, \quad \mathcal{V}_n := \int_0^\infty V(t) e^{-\lambda_n t} dt,$$

and

$$\begin{aligned} w_n &:= \int_0^\infty (W * S^*)(t) e^{-\lambda_n t} \varphi_n dt \\ &= \int_0^\infty W(t) e^{-\lambda_n t} dt \int_0^\infty S^*(t) e^{-\lambda_n t} \varphi_n dt \\ &= \mathcal{W}_n \hat{\phi}_n, \quad \mathcal{W}_n := \int_0^\infty W(t) e^{-\lambda_n t} dt, \end{aligned}$$

and using (7.11), we deduce that

$$(7.26) \quad \hat{\phi}_n = v_n + w_n.$$

By construction and (7.12), we have $\|v_n\|_0 \rightarrow 0$, (w_n) is bounded in Y_1 and thus weakly compact in Y_0 . That implies that $(\hat{\phi}_n)$ is weakly compact in Y_0 . There thus exist a subsequence $(\hat{\phi}_{n_\ell})$ and $\phi_1 \in Y_0$ such that $\hat{\phi}_{n_\ell} \rightharpoonup \phi_1$ weakly in Y_0 . In particular, $\phi_1 \geq 0$ and $\mathcal{L}^* \phi_1 = \lambda_1 \phi_1$. On the other hand, from (7.26), we have

$$1 = \|\hat{\phi}_n\|_0 \leq \|v_n\|_0 + \|w_n\|_0,$$

with $\|v_n\|_0 \rightarrow 0$ again and

$$\begin{aligned} \|w_n\|_0 &\leq \xi_\varepsilon \|\mathcal{W}_n \hat{\phi}_n\|_1 + \Xi_\varepsilon [\mathcal{W}_n \hat{\phi}_n]_{g_\varepsilon} \\ &\leq \xi_\varepsilon C \|\hat{\phi}_n\|_0 + \Xi_\varepsilon [\mathcal{W}_n \hat{\phi}_n]_{g_\varepsilon}, \end{aligned}$$

where we have used (7.9) in the first line and (7.12) in the second line. Choosing $n \geq 1$ large enough so that $\|v_n\|_0 \leq 1/4$ and $\varepsilon > 0$ small enough so that $\xi_\varepsilon C \leq 1/4$, we deduce from the two above estimates and the fact that $\mathcal{W}_n \geq 0$ that

$$\frac{1}{2} \leq \Xi_\varepsilon[\mathcal{W}_n \hat{\phi}_n]_{g_\varepsilon} = \Xi_\varepsilon \langle w_n, g_\varepsilon \rangle, \quad \forall n \geq 0.$$

Using that $g_\varepsilon \in Y'_0$, we may pass to the limit in the above inequality and we deduce

$$(7.27) \quad \frac{1}{2\Xi_\varepsilon} \leq \langle \phi_1, g_\varepsilon \rangle,$$

in particular $\phi_1 \neq 0$ and that concludes the proof of the existence of a dual eigenelement. We have thus established the existence of a first dual eigenelement, that is $(\lambda_1, \phi_1) \in \mathbb{R} \times Y$ such that

$$(7.28) \quad \mathcal{L}^* \phi_1 = \lambda_1 \phi_1, \quad \phi_1 \geq 0, \quad \phi_1 \neq 0, \quad \lambda_1 \geq \lambda_0.$$

It is worth emphasizing that we only have $\|\phi_1\|_0 \leq 1$ from the lsc property of the norm $\|\cdot\|_0$.

Step 2. More about the dual eigenfunction. From (7.2) applied to ϕ_1 , we have

$$e^{\lambda_1 t} \|\phi_1\|_1 = \|S^*(t)\phi_1\|_1 \leq C_0 e^{\lambda t} \|\phi_1\|_1 + C_1 \int_0^t e^{\lambda(t-s) + \lambda_1 s} ds \|\phi_1\|_0,$$

so that

$$(1 - C_0 e^{(\lambda - \lambda_1)t}) \|\phi_1\|_1 \leq C_1 \int_0^t e^{(\lambda - \lambda_1)\tau} d\tau \|\phi_1\|_0.$$

We recall that $\lambda_1 \geq \lambda_0$ from (7.28), thus

$$\frac{\lambda_0 - \lambda}{C_1} (1 - C_0 e^{(\lambda - \lambda_0)t}) \|\phi_1\|_1 \leq \|\phi_1\|_0,$$

and finally, passing to the limit $t \rightarrow \infty$, we deduce that

$$(7.29) \quad \|\phi_1\|_1 \leq \frac{C_1}{\lambda_0 - \lambda} \|\phi_1\|_0.$$

We normalize the dual eigenfunction with the norm $\|\cdot\|_0$. Note that since for the eigenvector ϕ_1 built in the step 1 we had $\|\phi_1\|_0 \leq 1$, the lower bound (7.27) remains valid for the new normalization $\|\phi_1\|_0 = 1$. Using (7.7), we thus deduce

$$(7.30) \quad \phi_1 \geq e^{\lambda_1(T_0 - T_1)} [\phi_1]_{g_\varepsilon} \psi_\varepsilon \geq \frac{e^{\lambda_1(T_0 - T_1)}}{2\Xi_\varepsilon} \psi_\varepsilon.$$

Step 3. The Lyapunov condition. We define $\tilde{S}(t) := e^{-\lambda_1 t} S(t)$. From (7.1), we have

$$\begin{aligned} \|\tilde{S}(t)f\|_1 &\leq C_0 e^{(\lambda - \lambda_0)t} \|f\|_1 + C_1 \int_0^t e^{(\lambda - \lambda_0)(t-s)} \|\tilde{S}(s)f\|_0 ds \\ &\leq \gamma'_L \|f\|_1 + K' \|f\|_0, \end{aligned}$$

for any $t \geq 0$, and with $\gamma'_L := C_0 e^{(\lambda - \lambda_0)t}$, $K' := C_0 e^{(\lambda - \lambda_0)t} (e^{C_1 t} - 1)$. From (7.8) and (7.30), we have

$$\|f\|_0 \leq \xi_\varepsilon \|f\|_1 + \Xi_\varepsilon [f]_{\psi_\varepsilon} \leq \xi_\varepsilon \|f\|_1 + 2\Xi_\varepsilon^2 e^{\lambda_1(T_1 - T_0)} [f]_{\phi_1}.$$

We then fix $T \geq T_1 > 0$ such that $\gamma'_L < 1$ and next $\varepsilon > 0$ such that $\gamma_L := \gamma'_L + \xi_\varepsilon K' < 1$ and we deduce that S satisfies the Lyapunov condition

$$(7.31) \quad \|\tilde{S}_T f\|_1 \leq \gamma_L \|f\|_1 + K [f]_{\phi_1},$$

with $K := 2\Xi_\varepsilon^2 e^{\lambda_1(T_1 - T_0)} K'$.

Step 4. Take $f \geq 0$ such that $\|f\|_1 \leq A[f]_{\phi_1}$ with $A > K/(1 - \gamma_L)$. We have

$$\begin{aligned} \|\tilde{S}_{T_0} f\|_1 &\leq C'_0 e^{(\lambda' - \lambda_0)T_0} \|f\|_1 \\ &\leq C'_0 e^{(\lambda' - \lambda_0)T_0} A[f]_{\phi_1} \\ &= C'_0 e^{(\lambda' - \lambda_0)T_0} A[\tilde{S}_{T_0} f]_{\phi_1} \\ &\leq C'_0 e^{(\lambda' - \lambda_0)T_0} A\|\phi_1\|_1 \|\tilde{S}_{T_0} f\|_0 \\ &\leq C'_0 e^{(\lambda' - \lambda_0)T_0} A\|\phi_1\|_1 (\xi_\varepsilon \|\tilde{S}_{T_0} f\|_1 + \Xi_\varepsilon [\tilde{S}_{T_0} f]_{\psi_\varepsilon}) \end{aligned}$$

for any $\varepsilon > 0$, where we have used successively the growth estimate (7.10) in the first line, the condition on f in the second line, the eigenfunction property of ϕ_1 in the third line, the duality bracket estimate (7.3) in the fourth line and the interpolation inequality (7.8) in the last line. Choosing $\varepsilon > 0$ small enough, we immediately obtain

$$\|\tilde{S}_{T_0} f\|_1 \leq 2C'_0 e^{(\lambda' - \lambda_0)T_0} A\|\phi_1\|_1 \Xi_\varepsilon [\tilde{S}_{T_0} f]_{\psi_\varepsilon}.$$

Together with

$$[f]_{\phi_1} = [\tilde{S}_{T_0} f]_{\phi_1} \leq \|\tilde{S}_{T_0} f\|_1$$

and the relaxed Doblin-Harris positivity condition (7.6), we conclude to the conditional Doblin-Harris positivity estimate

$$(7.32) \quad \tilde{S}_T f \geq c g_\varepsilon [f]_{\phi_1}$$

for all $T \geq T_1$, with $c^{-1} = 2C'_0 e^{(\lambda' - \lambda_0)T_0} A\|\phi_1\|_1 \Xi_\varepsilon e^{\lambda'(T - T_0)} \eta_{\varepsilon, T}^{-1}$.

Step 5. We define $\mathcal{N} := \{f \in X_1; \langle f, \phi_1 \rangle = 0\}$. As a consequence of the last conditional Doblin-Harris positivity estimate, we show that there exists $\gamma_H \in (0, 1)$ such that holds the following local coupling condition

$$(7.33) \quad (f \in \mathcal{N}, \|f\|_1 \leq A[f]_{\phi_1}) \quad \text{implies} \quad [\tilde{S}_T f]_{\phi_1} \leq \gamma_H [f]_{\phi_1},$$

still under the same condition $A > K/(1 - \gamma_L)$. Take indeed $f \in \mathcal{N}$. Because $\langle f, \phi_1 \rangle = 0$, the Doblin-Harris condition (7.32) tells us that

$$(7.34) \quad \tilde{S}_T f_\pm \geq c g_\varepsilon [f_\pm]_{\phi_1} = r g_\varepsilon, \quad r := c [f]_{\phi_1} / 2,$$

and we may thus write

$$\begin{aligned} |\tilde{S}_T f| &= |\tilde{S}_T f_+ - r g_\varepsilon - \tilde{S}_T f_- + r g_\varepsilon| \\ &\leq |\tilde{S}_T f_+ - r g_\varepsilon| + |\tilde{S}_T f_- - r g_\varepsilon| \\ &= \tilde{S}_T f_+ - r g_\varepsilon + \tilde{S}_T f_- - r g_\varepsilon = \tilde{S}_T |f| - 2r g_\varepsilon, \end{aligned}$$

where we have used the inequality (7.34) in the third line. We deduce

$$[\tilde{S}_T f]_{\phi_1} \leq [\tilde{S}_T |f|]_{\phi_1} - 2r [g_\varepsilon]_{\phi_1} = (1 - c [g_\varepsilon]_{\phi_1}) [f]_{\phi_1},$$

where we have used $\tilde{S}_T^* \phi_1 = \phi_1$, and that ends the proof of (7.33).

We now introduce a new norm $\|\!\| \cdot \|\!\|$ on X_1 defined by

$$(7.35) \quad \|\!\| f \|\!\| := [f]_{\phi_1} + \beta \|f\|_1,$$

and we claim that there exist $\beta > 0$ small enough and $\gamma \in (0, 1)$ such that

$$(7.36) \quad \|\!\| \tilde{S}_T f \|\!\| \leq \gamma \|\!\| f \|\!\|, \quad \text{for any } f \in \mathcal{N}.$$

Note that $\|\!\| \cdot \|\!\|$ and $\|\cdot\|_1$ are equivalent norms, with

$$(1 + \beta)^{-1} \|\!\| f \|\!\| \leq \|f\|_1 \leq \beta^{-1} \|\!\| f \|\!\|.$$

In order to establish the contraction estimate (7.36), we fix $f \in \mathcal{N}$ and estimate the norm $\|\!\| \tilde{S}_T f \|\!\|$ in two alternative cases:

First case. *Contractivity for small X_1 norm.* When

$$(7.37) \quad \|f\|_1 < A[f]_{\phi_1},$$

the local coupling condition (7.33) implies

$$[\tilde{S}_T f]_{\phi_1} \leq \gamma_H [f]_{\phi_1}.$$

Together with the Lyapunov condition (7.31), we have

$$\|\tilde{S}_T f\| \leq (\gamma_H + \beta K)[f]_{\phi_1} + \beta \gamma_L \|f\|_1 \leq \gamma_1 \|f\|,$$

with

$$\gamma_1 := \max\{\gamma_H + \beta K, \gamma_L\}.$$

Choosing $\beta > 0$ small enough such that $\beta K < 1 - \gamma_H$, we get $\gamma_1 < 1$ and that gives the contractivity property (7.36) in this case.

Second case. *Contractivity for large X_1 norm.* Assume on the contrary that

$$(7.38) \quad \|f\|_1 \geq A[f]_{\phi_1}.$$

Directly from (7.31) we deduce that

$$\|\tilde{S}_T f\|_1 \leq \gamma_L \|f\|_1 + K[f]_{\phi_1} \leq (\gamma_L + K/A)\|f\|_1,$$

with $\gamma_L + K/A < 1$ by assumption. On the other hand, we have

$$[\tilde{S}_T f]_{\phi_1} \leq \langle \tilde{S}_T |f\rangle, \phi_1 \rangle = \langle |f\rangle, \phi_1 \rangle,$$

by using the positivity property of \tilde{S}_T and the eigenvector property of ϕ_1 . Using both last estimates together, we deduce

$$\begin{aligned} \|\tilde{S}_T f\| &= [\tilde{S}_T f]_{\phi_1} + \beta \|\tilde{S}_T f\|_1 \\ &\leq [f]_{\phi_1} + \beta(\gamma_L + K/A)\|f\|_1 \\ &\leq (1 - \beta\delta_0)[f]_{\phi_1} + \beta(\gamma_L + K/A + \delta_0)\|f\|_1, \end{aligned}$$

for any $\delta_0 \geq 0$. We thus get

$$\|\tilde{S}_T f\| \leq \gamma_2 \|\tilde{S}_T f\|,$$

with $\gamma_2 := \max(1 - \beta\delta_0, \gamma_L + K/A + \delta_0)$. We get the contractivity property (7.36) in this case by choosing $\delta_0 > 0$ small enough (and keeping the choice of $\beta > 0$ made in the previous case) so that $\gamma_2 \in (0, 1)$. The proof of (7.36) is completed by setting $\gamma := \max\{\gamma_1, \gamma_2\}$.

Step 6. In order to prove the existence and uniqueness of the eigenvector $f_1 \in X_1$, we fix $g_0 \in \mathcal{M} := \{g \in X_1, g \geq 0, \langle g, \phi_1 \rangle = 1\}$, and we define recursively $g_k := \tilde{S}_T g_{k-1}$ for any $k \geq 1$. Thanks to (7.36), we get


$$\sum_{k=1}^{\infty} \|g_k - g_{k-1}\| \leq \sum_{k=0}^{\infty} \gamma^k \|g_1 - g_0\| < \infty,$$

so that (g_k) is a Cauchy sequence in \mathcal{M} . We set $f_1 := \lim g_k \in \mathcal{M}$ which is a stationary state for the mapping \tilde{S}_T , as seen by passing to the limit in the recursive equations defining (g_k) . From (7.36) again, this is the unique stationary state for this mapping in \mathcal{M} . From the semigroup property, we have $\tilde{S}_t f_1 = \tilde{S}_t \tilde{S}_T f_1 = \tilde{S}_T(\tilde{S}_t f_1)$ for any $t > 0$, so that $\tilde{S}_t f_1$ is also a stationary state in \mathcal{M} , and thus $\tilde{S}_t f_1 = f_1$ for any $t > 0$, by uniqueness. That precisely means that f_1 is a positive eigenvector associated to λ_1 for the original problem.

For $f \in X$, we see that $h := f - \langle f, \phi_1 \rangle \phi_1 \in \mathcal{N}$, and using recursively (7.36), we deduce

$$\|\tilde{S}_{nT} h\| \leq \gamma^n \|h\|, \quad \forall n \geq 0,$$

from what (7.13) follows by standard arguments. \square

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