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► **To cite this version:**

Ye Mou, Xingyu Yang, Marlo Vega, Romeo Zapata, Bruno Gallas, et al.. Femtosecond Drift Photocurrents Generated by an Inversely Designed Plasmonic Antenna. Nano Letters, In press, 10.1021/acs.nanolett.4c00558 . hal-04609032

HAL Id: hal-04609032

<https://hal.science/hal-04609032>

Submitted on 12 Jun 2024

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Femtosecond drift photocurrents generated by an inversely designed plasmonic antenna

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Keywords: Ultrafast direct currents, Drift currents, Nanophotonics, Plasmonics, Optical Rectification, Inverse Faraday Effect, Inverse Design, Light-Matter Interactions.

Abstract

Photocurrents play a crucial role in various applications, including light detection, photovoltaics, and THz radiation generation. Despite the abundance of methods and materials for converting light into electrical signals, the use of metals in this context has been relatively limited. Nanostructures supporting surface plasmons in metals offer precise light manipulation and induce light-driven electron motion. Through inverse design optimization of a gold nanostructure, we demonstrate enhanced volumetric, unidirectional, intense, and ultrafast photocurrents via a magneto-optical process derived from the inverse Faraday effect. This is achieved through fine-tuning the amplitude, polarization, and their gradients in the local light

field. The virtually instantaneous process allows dynamic photocurrent modulation by varying optical pulse duration, potentially yielding nanosources of intense, ultrafast, planar magnetic fields, and frequency-tunable THz emission. These findings open avenues for ultrafast magnetic material manipulation and hold promise for nanoscale THz spectroscopy.

Introduction

A photocurrent is the flow of electric charges generated when photons are absorbed by a material. This phenomenon plays a critical role in various groundbreaking fields like light detection, solar photovoltaics, lightwave electronics, and THz spectroscopy and imaging.¹ For physicists, a direct photocurrent holds significant importance as it serves as a unique tool to investigate various processes across a wide range of spatial and temporal scales, ranging from femtoseconds to milliseconds and from the mesoscale to the microscale. Hence, understanding photocurrent is crucial to unveil physical and chemical effects due to the intimate connection between photons and electrons.

Photocurrents have been instrumental in advancing the development of low-dimensional structures,² topological materials,³ and the exploration of systems exhibiting symmetry breaking.^{4, 5} Their often ultrafast dynamics contribute to understanding out-of-equilibrium transient states.^{6, 7} Broadly speaking, the conversion of light into an electrical signal can result from various phenomena. Some, like the p-n junction, are commonly found in everyday technologies, while others, such as current generated by photothermal gradients, are less explored in the literature. There are fascinating prospects involving photo-induced transport of charge carriers in metals. Examples include optical rectification of visible photons through plasmonic nanoscale rectennas⁸⁻¹⁰ or the generation of photocurrents by structured light with variations in intensity, phase, and polarization.¹¹ What makes metals and plasmonics particularly appealing is their ability to fine-tune local optical fields while simultaneously generating electric currents on the same material platform, thereby eliminating the need for complex hetero-integration of materials.

Despite recent efforts to understand the interplay of phenomena in these systems,^{12, 13} a comprehensive framework for manipulating direct photocurrents using plasmonics at the nanoscale is still largely missing. A strategy to engineer their spatial distribution, control their strength, and dictate their direction remains to be rigorously formulated theoretically, optimized through numerical simulations, and verified experimentally.

There are physical processes that have been generally overlooked for photocurrent generation: the effect of nonlinear forces. A ponderomotive force is an example of such process enabling drift currents in the skin depth of a metal parallel to its surface. These drift currents have many advantages: they are relatively simple to model, their generation is almost instantaneous, and they last only as long as the light exposure, so they can be very short when generated by an optical pulse. Moreover, they can be manipulated at will by tuning the intensity of the optical field, its polarization, or its gradients.

Drift currents can be regarded as a type of optical rectification at the origin of the inverse Faraday effect in metals.¹⁴⁻¹⁶ This is a nonlinear process promoting the magnetization of matter by an optical excitation only. Just like other all-optical approaches,^{17, 18} its use recently spurred the theoretical¹⁹⁻²¹ and experimental^{22, 23} demonstration of very intense and very short magnetic fields at the nanoscale, with potential applications in data writing and processing at ultrafast time scales. In these studies, circular polarizations state of light are used to move electrons in structures with a rotational symmetry. This is the principle of the inverse Faraday effect (IFE) described in the 1960s. However, it was recently demonstrated that new physical effects can emerge from IFE by manipulating light at the nanoscale with plasmonic nanostructures. In particular, it was shown that IFE can be generated by linear polarization state of light,²⁴ that it can be chiral,^{25, 26} or even have its symmetry reversed.²⁷ Our study builds upon this body of literature. We demonstrate that a plasmonic nanostructure specifically designed by an inverse optimization manipulates light field at the nanoscale in such a way that a volumetric, unidirectional, intense, and ultra-short photocurrent emerges. We also demonstrate that this phenomenon is chiral, *i.e.*, occurring for only one polarization state of light and that it can

generate a magnetic field perpendicular to the direction of light propagation, which is impossible with a conventional inverse Faraday effect. Finally, by generating ultra-short, planar and unidirectional photocurrents, we show that this physical process opens the way to generate a linearly polarized and frequency-tunable THz nanosource.

Results and discussions

In this work, we use the optical rectification phenomenon at the origin of IFE to generate ultrafast direct photocurrents at the nanoscale by a light-induced magnetization of matter.¹⁴⁻

¹⁶ In a metal, this magnetization is made possible by the nonlinear forces that an electromagnetic wave applies to the free electrons. The average value of these forces being non-zero, a direct current called drift current is generated, and in turn, forms a stationary magnetic field.

The theory underlying the generation of drift currents was first developed by the plasma community^{28, 29} before being extrapolated to optics^{30, 31} and nanophotonics.^{19, 21-26, 32-34} The following equation gives the expression of the drift currents in a metal:^{21, 24}

$$\mathbf{J}_d = \frac{1}{2en} \text{Re} \left(\left(-\frac{\nabla \cdot (\sigma_\omega \mathbf{E})}{i\omega} \right) \cdot (\sigma_\omega \mathbf{E})^* \right) \quad (1)$$

e is the charge of the electron ($e < 0$), n is the charge density at rest, σ_ω is the dynamic conductivity of the metal and \mathbf{E} is the optical electric field. The terms in bold represent vectorial values, and the star * denotes the conjugate complex.

As one can see in equation 1, drift currents are proportional to the optical electric field and its divergence. Therefore, by its ability to enhance electric fields and to create strong field gradients, plasmonics is a particularly well-suited technology to generate and manipulate in a controlled manner strong drift currents with the aim to create ultrafast, intense magnetic field pulses at the nanoscale.²¹

In IFE, it is conventionally understood that drift currents have an azimuthal symmetry generating a magnetization perpendicular to the plane in which the electrons spin, as recently

demonstrated in a series of theoretical^{19, 21, 24, 26, 33, 34} and experimental^{22, 23} papers. However, as recognized by the plasma community, the nonlinear forces that light applies to the electron sea possess all kinds of symmetries. For instance, ponderomotive forces accelerating charges in a plasma hold radial symmetries.^{35, 36}

Based on this concept, we designed a plasmonic nanostructure capable of generating a direct drift photocurrent that is intense, unidirectional in the plane of the object, independent of the excitation polarization, and ultrafast. The origin of these exceptional properties can be found in equation 1. As shown in this expression, the drift currents can only be created when the product $\mathbf{E} \cdot \mathbf{E}^*$ is not zero, *i.e.*, for an elliptical or circular polarization state of the light. Therefore manipulating the local polarization state of the light with an adequate design enables to create different ellipticities around the object and a control of the drift currents.

Specifically, we have inversely designed a plasmonic nanostructure in a 30 nm thick gold layer deposited on a glass substrate using a genetic algorithm (GA) based on topological optimization (Figure 1).

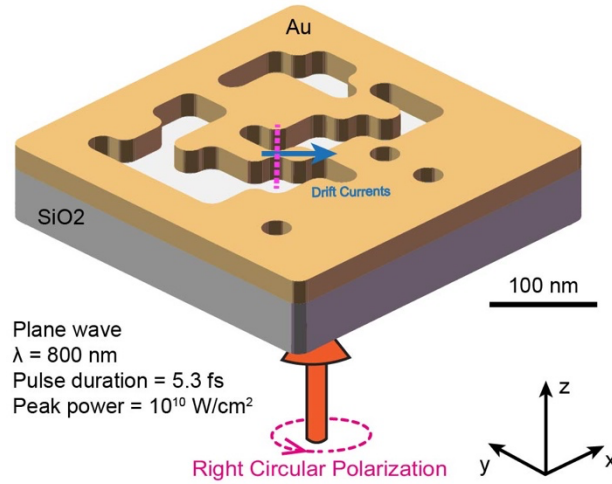


Figure 1. Optimized structure and excitation conditions. Example of a GA-optimized structure, made in a 30 nm thick gold layer. The 280 x 280 nm² nanostructure is based on a two-dimensional array of 10x10 elements, each element consisting of metal or air with a size of 28 nm. These dimensions are chosen to facilitate the fabrication of this structure by e.g. nanoscale lithography techniques. For the same reason and to avoid non-physical effects that a numerical approach can generate locally, the corners of the nanostructure are rounded (see figure S1 of the Supporting Information). The optical excitation of the nanostructure is modeled by a pulsed plane wave emitted at a wavelength of 800 nm. The pulse is launched from the substrate side and is right circularly polarized. We set the pulse duration at 5.3 fs and the peak power of 10¹⁰ W/cm², a value well below the damage threshold of the material.^{37, 38} The blue arrow represents the ultrafast direct currents flowing in a single direction, and the pink dashed line indicates the position of a cross-section.

In order to generate a strong unidirectional drift current in the plane of the metal layer, the GA optimization function consisted in maximizing the local stationary magnetic field created by the drift currents at the center of the structure and having a vectorial symmetry parallel to the gold surface. Figure S2 represents the evolutionary history of the nanostructures in connection with the generation of a strong magnetic field by IFE. After 87 generations the strength of the magnetic field reaches a plateau, and the structure is considered optimized. Figure 2a

represents schematically the result of the optimization procedure. The interesting part of the structure is certainly the bridge-like geometry at the center. Figure 2b shows the distribution and the increase of the optical electric field at the origin of the drift currents and **B** field in the center of this design (represented by a black square in figure 2a) in an XY plane taken at the middle of the metal thickness. From this electric field distribution and using Equation 1, we can then plot the distribution of drift currents in the structure as shown in Figure 2c. First, drift currents are maximal in a gold strip in the center of the structure but not necessarily where the electric field and field gradients are the strongest. Importantly, the arrows representing the direction of drift are all oriented in the same direction in the XY plane, demonstrating the possibility of generating, in a controlled manner, straight and unidirectional DC currents by optical excitation. Note that other physical processes may be responsible for generating DC currents, one of which is related to thermal effects, for instance. Although the message of this article is not to compare all these processes, a study of temperature generation in our optimized nano-system for our optical excitation characteristics is described in Supporting Information (Figure S3 and S4) based on two-temperature model.³⁹ As described in this supporting figure, the lattice temperature remains extremely modest and rather homogenous in the gold strip, implying a minimal contribution of this physical process to DC current generation. Also, the electronic temperature for the peak powers and pulse durations considered here are relatively low (Figure S4). Note that the dynamics of these temperatures do not change with increasing optical pulse duration, as shown in Figure S5. Conversely, at constant peak power, both lattice and electronic temperatures increase linearly with pulse duration, while they remain constant for constant pulse energy, as illustrated in Figure S6.

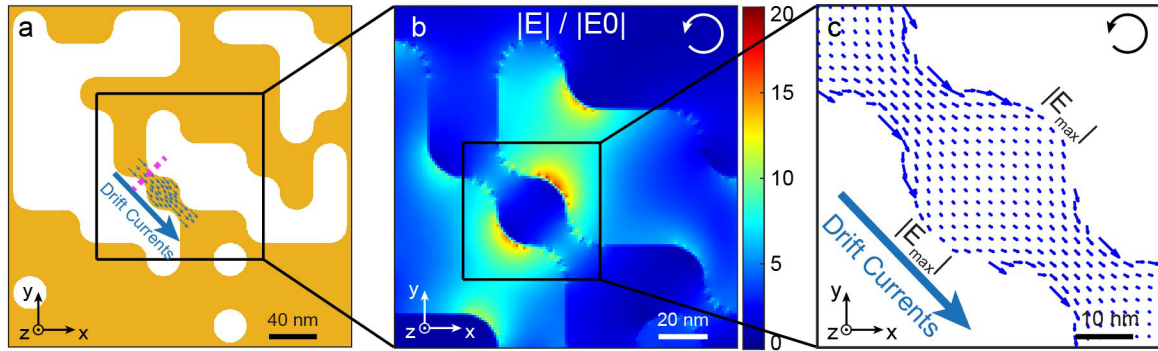


Figure 2. Optical response of the optimized plasmonic nanostructure and associated drift currents. a) Schematic, in an XY plane, of the GA-optimized structure. b) Spatial distribution of the electric field enhancement at the Z-center of the structure shown in a) for a right circular polarization of excitation. c) Spatial distribution of the drift currents generated at the surface of the structure by the electric field distribution shown in b). The length of the arrows refers to the amplitude of the drift currents.

One way to visualize the vectorial orientation of the DC currents in the metal is to image the resulting vectorial distribution of the magnetic field (**B**). For this purpose, Figure 3 represents the distribution of the magnetic field created by IFE in the same plane as Figure 2b,c. In particular, Figure 3a and b displays respectively the amplitude of the total **B**-field and the **B**-field component oriented along Z. Figure 3c shows the vectorial distribution in the y/z cross-sectional plane of the metal junction, indicated by a pink line in Figure 3a. This distribution of the magnetic field components around the thin gold strip is therefore characteristic of a current flowing unidirectionally in the direction shown by the arrow in Figure 2c (see also figure S7). Moreover, we note that in past studies, the orientations of the magnetic fields generated by optical rectification through IFE have all been collinear with the direction of light propagation.^{19, 21, 33} Here, by manipulating the light at the nanoscale, we demonstrate the creation of an azimuthally polarized magnetic field around the gold nanostructure (Figures 3b,c, and S7). This allows for generating a magnetic field above the gold layer parallel to the nanostructure's surface, and thus perpendicular to the direction of light propagation. This result has significant

implications for applying IFE to manipulate magnetic processes at the nano⁴⁰ and potentially ultrafast^{41, 42} scales since many magnetic layers have domains orientation parallel to the sample plane (i.e., XY here).

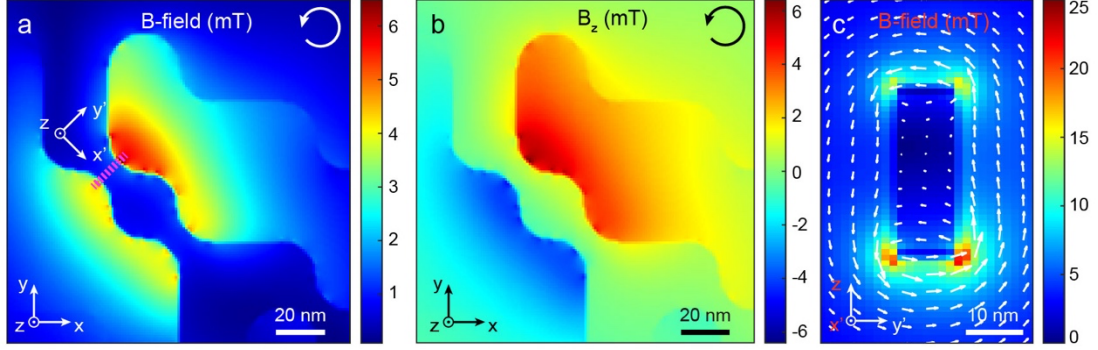


Figure 3. Magnetic response of the optimized plasmonic nanostructure under a right circular polarization of excitation. a) Spatial distribution of the amplitude of the total magnetic field generated at the Z-center of the structure represented in Figure 2a. b) Z-component of the spatial distribution of the magnetic field in a). c) Vectorial distribution in a cross section plane (y'z) of the metal junction symbolized by the pink line in a). The arrows represent the orientation of the \mathbf{B} field and the colors its amplitude.

To further understand why this plasmonic structure creates a unidirectional DC current, we plot in Figure 4a the distribution of the Z-oriented component of the optical spin density in the same plane as Figure 2c, which is a parameter related to light polarization state. The spin density \mathbf{s} ,⁴³⁻⁴⁵ described by the following equation:

$$\mathbf{s} = \frac{1}{|\mathbf{E}_0|^2} \text{Im}(\mathbf{E}^* \times \mathbf{E}) \quad (2)$$

is a vectorial quantity describing the polarization state of the light. In our coordinate system, negative and positive spin densities correspond respectively to left elliptical and right elliptical polarization state of light; and $s = 0$ is equivalent to a linear polarization state of light. Now, as highlighted above, drift currents can emerge in the presence of an elliptical polarization. From equation 1, a left elliptical polarization state of light creates a drift current in one direction, while a right elliptical polarization state of light generates drift currents in the opposite direction. With

this understanding, the spin density distribution shown in Figure 4a sheds light on the vectorial distribution of DC currents shown in Figure 2c. On each side of the gold strip, where the currents are the strongest, we concomitantly observe the strongest spin densities of opposite signs. By symmetry, the generated drift currents are flowing in the same direction and are oriented in the same way as the red and blue arrows representing the local helicity of the light. These hot spots of so-called super circular light,^{24, 26} a concept similar to super chiral light related to elliptical polarization of enhanced optical fields⁴⁶, are thus at the origin of the strong currents and their straight and unidirectional orientation. The spin density distribution inducing this physical effect is the result of constructive and destructive interferences of the light in the plasmonic structure, interferences that the evolutionary approach of the GA has selected.

Because the nanostructure is excited by an ultrashort optical pulse, we anticipate that the resulting DC current to carry also an ultrafast dynamics. Figure 4b confirms this hypothesis. This figure represents the instantaneous intensity of the electric field as a function of time at a point on the surface of the structure, symbolized by a pink star in Figure 4a and for an optical pulse of 5.3 fs. In the same figure is shown the total drift current generated by this pulse and passing through a section of the gold strip represented by the pink dotted line in Figure 4a (Y'Z section in the X'Y'Z coordinate system rotated from the XYZ coordinate system by an angle 45° around Z, as shown in figure 4a). Clearly, an ultrashort photocurrent transient emerges across the section of the gold strip and reaches up to 400 μ A with the optical power used in the calculation. The duration of this transient is directly related to the duration of the optical pulse, since when there is no more energy to set the electrons in motion, they slow down immediately. Interestingly, this decay time can be observed in the tiny oscillations of the drift current, which occur when the instantaneous energy in the plasmonic nanostructure is 0 (taking into account a phase shift).

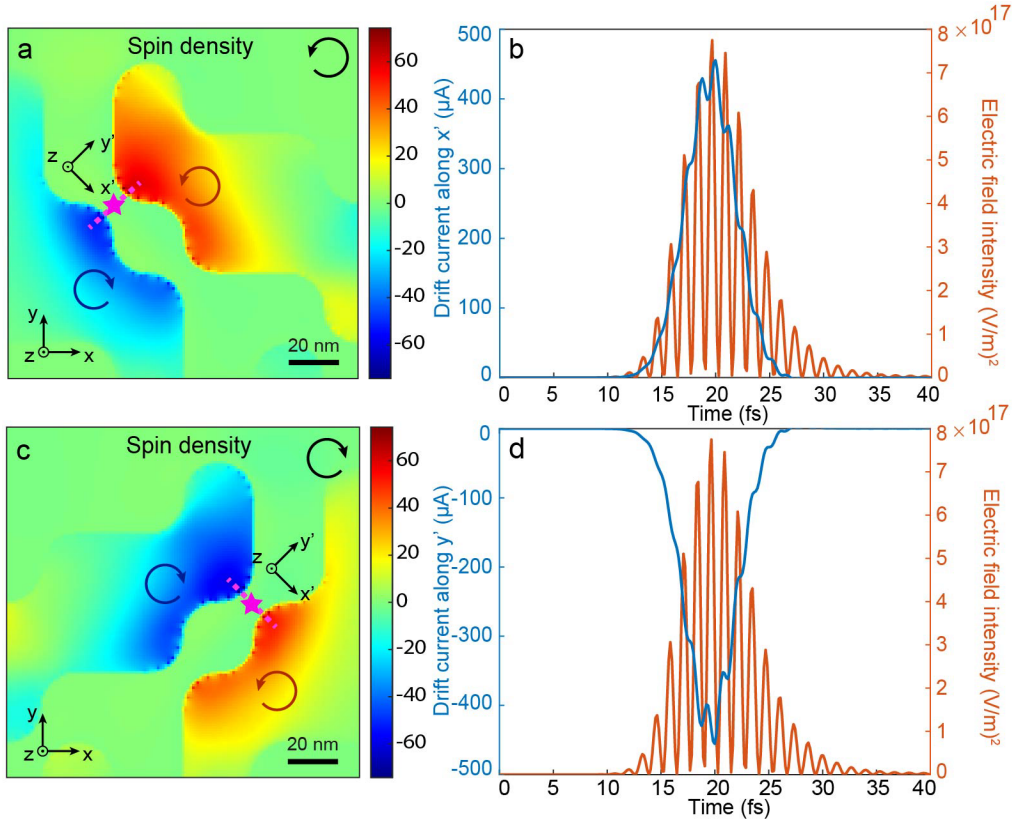


Figure 4. Spin density distribution and temporal responses. a) Local spin density at the Z-center of the optimized plasmonic nanostructure shown in Figure 2a for a right circular polarized excitation. b) Temporal response of the drift currents flowing along X' through the section shown in a) as a dashed line and the electric field intensity at the pink star point shown in a). c) Spatial distribution of the spin density in the same plane and for the mirror structure shown in a) excited by a left circular polarization. d) Temporal response of the drift currents and the electric field intensity, respectively, through the section symbolized by a dotted line and the pink star shown in c).

To further demonstrate the potentiality of controlling IFE-induced drift current at the nanoscale, an additional optimization condition was included in the definition of this nanostructure, consisting in making IFE chiral. During the selection process by the GA, the structures were selected to generate a magnetic field only for one helicity of the light, *i.e.* the right circular polarization state of light. Therefore, when a left circular polarization state of light illuminates the nanostructure, the drift currents and the associated stationary magnetic field are zero

(Figures S8 and S9). Reciprocally, the mirror structure of Figure 2a can generate a symmetric spin density only when excited by a left circular polarization state of light (Figure 4c) which results in an ultrafast direct current of perpendicular direction in the X'Y'Z coordinate system (Figures 4d and S10). These results are remarkable since this nanostructured landscape enables generating an ultrashort unidirectional current pulse even for an unpolarized excitation. Due to its chirality property, the GA-optimized structure naturally filters the polarizations from the unpolarized incoming field producing the direct current (principle of chirality and see supporting information for elliptical and linear polarizations of excitation).⁴⁷ Therefore, our algorithmic optimization approach isolates plasmonic nanostructures generating ultrafast photoinduced currents in one direction for a specific excitation polarization state of light and in another direction for an alternative polarization condition.

Based on these findings and drawing inspiration from the potential applications of ultrafast photocurrents for THz generation, we assess the suitability of the transient drift currents to act as a tunable THz nanosource. To achieve this, we excite the structure depicted in Figure 2a with femtosecond (fs) pulses of varying durations and observe temporal responses of drift photocurrents. The results are displayed in Figure 5a. Because of the ultrafast current dynamics already highlighted above (see Figure 4), an increase of the pulse duration leads to a longer current transient. Importantly, since the peak power of the optical pulses is kept identical, differences in photocurrent intensity are primarily attributed to the more effective power build-up of the structure at longer pulses.²¹

From the temporal responses calculated in Figure 5a, we can extract radiation generated at various frequencies by performing a Fourier transform of the time derivative of the currents. Figure 5b displays these Fourier transforms. The spectra reveal two frequency components associated with each photocurrent time response: one at low frequencies and another centered at 749.5 THz. The first component, displayed at higher resolution in Figure 5c, is the THz radiation produced by the current impulse. The second component corresponds to the

second harmonic generation resulting from the second-order nonlinear process described by equation 1.⁴⁸

Subsequently, we plot in Figure 5c the peak THz frequencies generated by the optimized nanostructure as a function of optical pulse durations based on the Fourier transform. As expected, these frequencies exhibit an inverse relationship with pulse lengths, ranging from 70 THz to approximately 4 THz for the selected pulse durations. An appealing feature of current transients produced by IFE is that they not only create a nanoscale, tunable source of THz radiation but also impart linear polarization to that radiation, despite the circular excitation of the incident wave. This is made possible by the unidirectionality of the photocurrents drifting in the gold plane. Considering the light energy absorbed and scattered by the optimized structure, it is estimated that approximately 63% of the incident energy interacts with it. Additionally, by comparing the amplitudes of the conduction currents to those of the drift currents, the conversion rate between these two types of currents is determined to be approximately 1.8%. Consequently, the THz radiation conversion efficiency is estimated to be slightly over 1%, aligning with the efficiencies reported for other THz sources in the literature.⁴⁹

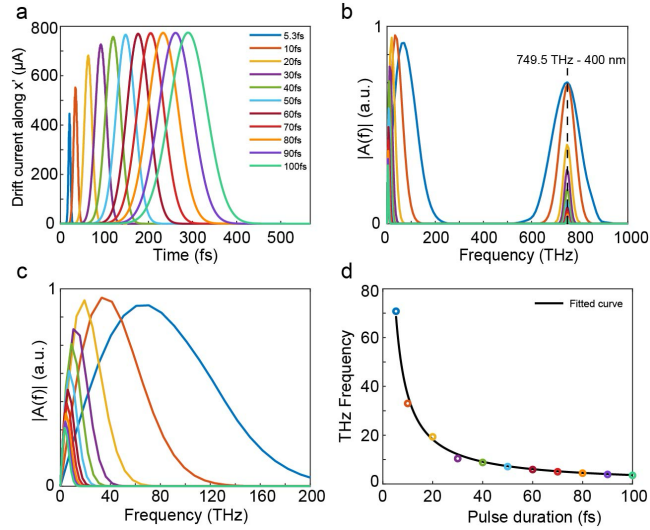


Figure 5. Generation of tunable THz radiation. a) Time response of the drift photocurrents as a function of the duration of the light pulse exciting the optimized nanostructure described in Figure 2a. b) Fourier transform of the time responses shown in a) over the frequency range from 0 to 1000 THz. c) Fourier transform of the time responses shown in a) over the frequency range from 0 to 200 THz. d) Distribution of the THz frequencies generated as a function of the pulse duration used.

In conclusion, we demonstrate that by manipulating the optical fields and field gradients in a plasmonic nanostructure optimized by an inverse design algorithm, unidirectional and direct photocurrents in the metal plane are generated. This is made possible by influencing the local polarization state of light around the photonic nanostructure. Also, thanks to the ability to assign several optimization objectives to the genetic algorithm, the optimized nanostructure creates a direct photocurrent only for one helicity, here the right circular polarization state of light. This chirality imparted to physical response enables generation of these currents even for unpolarized light. We then confirm that a direct current of the opposite direction is generated in the mirror structure for a left circular polarization state of light, implying that a control of the current flow is possible only by changing the polarization state of light incident on the structure. Also and importantly, we demonstrate that through the inverse Faraday effect, unidirectional current pulses in the plane of the metal are generated at ultrafast timescales. These currents

can be harnessed to produce ultrafast magnetic fields with unexpected symmetry, defying conventional expectations associated with the inverse Faraday effect.

In addition, the ultrafast drift photocurrents generated in this way offer the potential for creating a versatile and tunable nanosource for linearly polarized THz radiation. Hence, these findings pave the way to ultrafast in-plane manipulation of magnetic domains and the possibility of conducting ultra-localized THz spectroscopy at the nanoscale.

Supporting information

Simulation parameters, Figure S1-S13.

Acknowledgements

The authors declare no conflict of interest.

This work is supported by the ERC grant FemtoMagnet (grant no. 101087709). We acknowledge the financial support from the Agence national de la Recherche (ANR-20-CE09-0031-01, ANR-23-ERCC-0005), from the Institut de Physique du CNRS (Tremplin@INP 2020) and the China Scholarship Council. This work has been partially funded by the French Agence Nationale de la Recherche (ISITE-BFC ANR-15-IDEX-0003), the EIPHI Graduate School (ANR-17-EURE-0002) and the European Union through the PO FEDER-FSE Bourgogne Franche-Comté 2021/2027 programs. LN2 is a joint International Research Laboratory (IRL 3463) funded and co-operated in Canada by Université de Sherbrooke (UdeS) and in France by CNRS as well as ECL, INSA Lyon, and Université Grenoble Alpes (UGA). It is also supported by the Fonds de Recherche du Québec Nature et Technologie (FRQNT). This work was supported by the Natural Sciences and Engineering Research Council of Canada (NSERC) with Discovery Grants for JF Bryche. A doctoral scholarship from IDEX Paris-Saclay supports M. Vega. This work is supported by the “ADI” project funded by the IDEX Paris-Saclay, ANR-11-IDEX-0003-02.

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Supporting information

Femtosecond drift photocurrents generated by an inversely designed plasmonic antenna

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Keywords: Ultrafast direct currents, Drift currents, Nanophotonics, Plasmonics, Optical Rectification, Inverse Faraday Effect, Inverse Design, Light-Matter Interactions.

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Simulation parameters:

The simulations carried out in this study were done by the finite difference time domain (FDTD) method performed on the commercial software Lumerical from Ansys. This method solves Maxwell's equations in space and time using a finite difference technique. Indeed, the FDTD method solves these equations on a discrete spatial and temporal grid. The dimensions of the 3D computational window for the simulations were $750 \times 750 \times 900 \text{ nm}^3$. The boundary conditions of this window are made of a perfectly matched layer (PML), avoiding any parasitic reflection inside the calculation window. Several meshes are used for the discretization of the computational space, a coarser non-uniform mesh of 4 to 16 nm for the external unstructured parts, of the simulation, a finer mesh of 4 nm for a central nanostructured part of $288 \times 288 \times 36 \text{ nm}^3$ in X, Y, and Z, respectively containing the nanostructures, and an even finer mesh of 1 nm for the part where the drift currents and magnetic field are calculated of $140 \times 140 \times 32 \text{ nm}^3$ in X, Y, and Z (Figure S1c). The choice of this mesh size for the central part is chosen because convergence in the amplitude of the magnetic field is observed starting from this mesh size. The excitation of the nanostructures is performed by a pulsed plane wave of duration 5.3 fs spectrally centered at a wavelength of 800 nm. The peak power of the pulse is 10^{10} W/cm^2 , which corresponds to an energy slightly lower than 0.048 pJ applied to the plasmonic nanostructures (energy density of $53 \mu\text{J/cm}^2$). The convergence of the simulation was obtained when the energy inside the calculation window was lower than 10^{-5} of the initial injected energy. The textbook values of Johnson and Christy were used for the gold properties in these simulations.

The Genetic Algorithm (GA) we're employing utilizes a 10×10 binary matrix to represent each structure, where each cell in the matrix is a binary parameter indicating the presence (1) or absence (0) of material, with the material specified as nanocylinders of 28 nm in diameter and 30 nm in thickness. The entire topological process, including this matrix configuration, is illustrated in Figure S1.

Here are the details of the GA process:

1. **Initial Population Generation:** In the first generation of the GA, we randomly generate 200 structures. Each structure is represented by a 100-binary parameter matrix.
2. **Simulation and Evaluation:** We perform FDTD simulations for each structure to calculate the stationary magnetic field at the center point of each structure.

3. **Selection:** Using the objective function values outlined in Figure S2, you select the top 200 structures for reproduction. These consist of the top 100 structures from the most recent generation (N-1) and another 100 selected cumulatively from previous generations (1 to N-1).
4. **Reproduction:** The generation of new structures for the subsequent generation (N+1) is split into two processes:
 - **Mutation:** 100 new structures are generated by mutation at a rate of 10%, meaning ten of the 100 binary parameters in each structure are randomly changed.
 - **Breeding:** The remaining 100 structures are created through two-point scattering breeding. This involves randomly selecting and interchanging two binary parameters between two selected structures.

This GA approach is effective for exploring a vast design space by iterating over generations of structures, optimizing them based on the magnetic field calculations from the FDTD simulations. By employing both mutation and breeding, the algorithm promotes both the exploration of new areas in the design space and the exploitation of existing successful designs, aiming to find an optimal or near-optimal configuration that meets the objectives set forth in the algorithm's design.

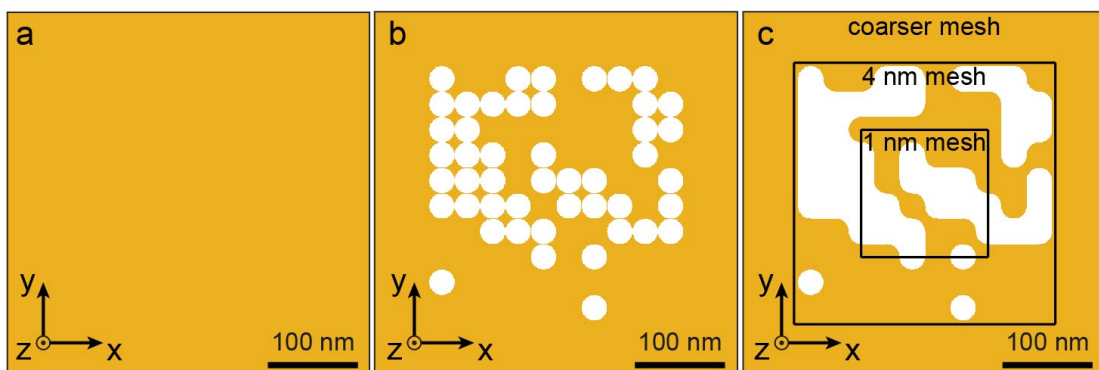


Figure S1. Construction of the elements constituting each generation of the genetic algorithm. Inside a) a 30 nm thick uniform layer of gold, b) holes are made according

to a binary matrix acting as the DNA in the evolutionary process. c) The obtained structure is then smoothed to avoid all the roughnesses not experimentally realistic and generating non-physical effects. The different mesh areas are shown in c.

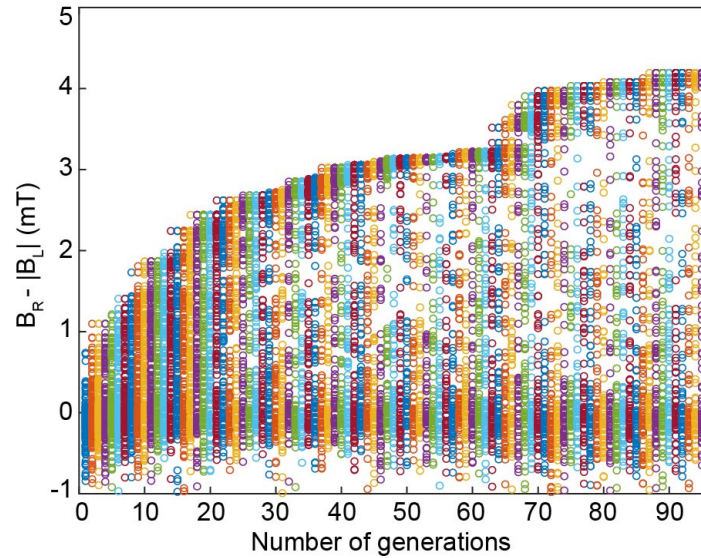


Figure S2. Evolutionary process. Evolution during the different generations of the optimization function that maximizes the difference $B_R - \text{abs}(B_L)$, with B_R and B_L being the magnetic fields B created by a right or left circular polarization, respectively. Each generation consists of 200 elements. The optimized structure appears after 87 generations.

Electronic and lattice temperature spatiotemporal evolution was carried out with a MATLAB finite elements model based on a two-temperature model, considering thermal conduction and out-of-equilibrium electron displacement¹.

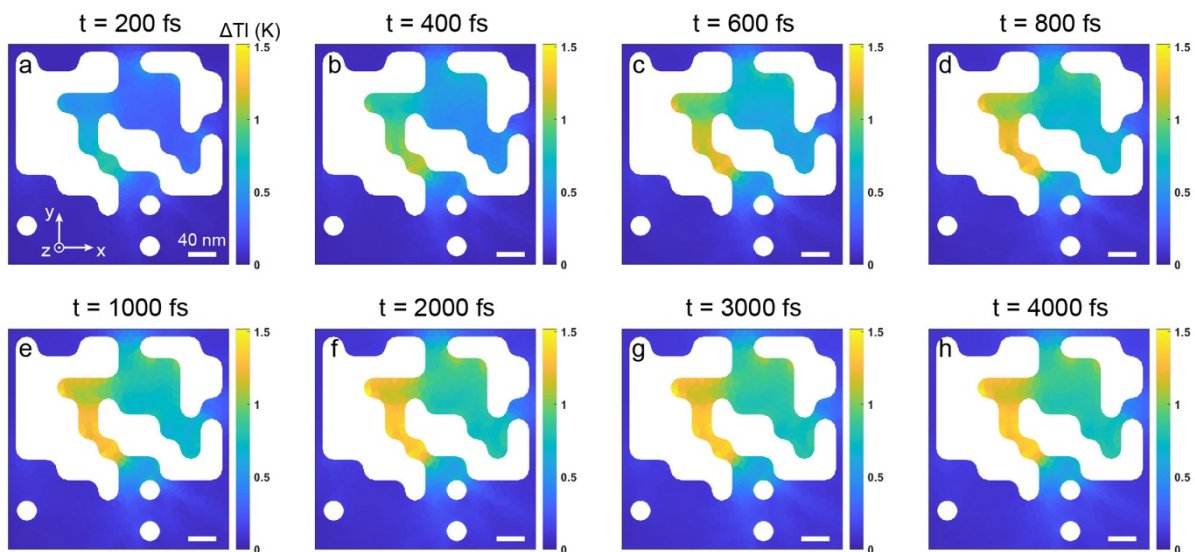


Figure S3. Spatial distribution of lattice temperature generation in an XY plane at the Z center of the optimized nanostructure for different characteristic times after the optical pulse, a) 200 fs, b) 400 fs, c) 600 fs, d) 800 fs, e) 1000 fs, f) 2000 fs, g) 3000 fs, h) 4000 fs.

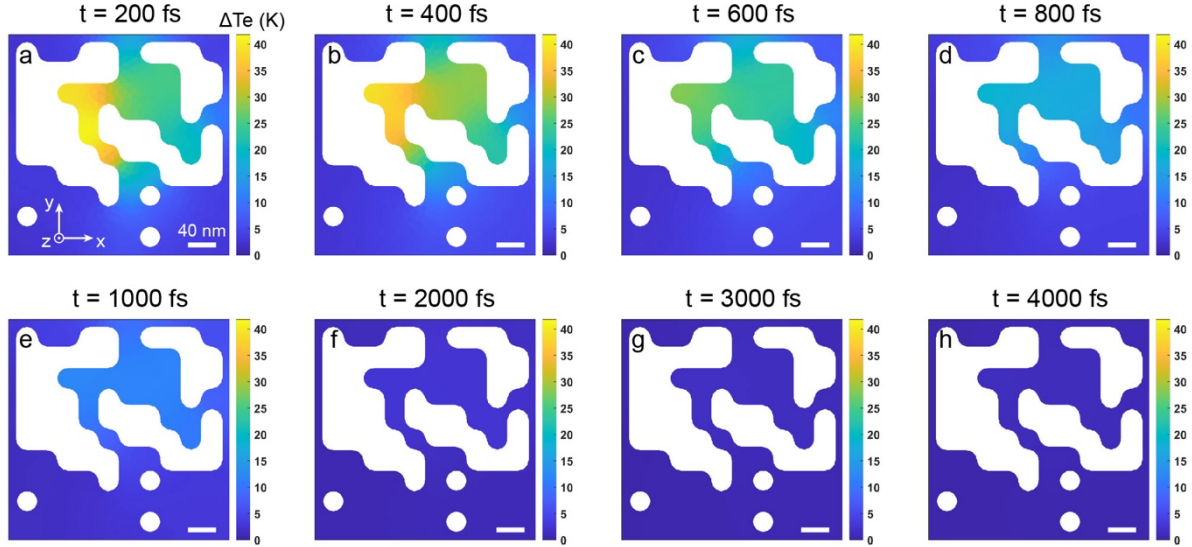


Figure S4. Spatial distribution of electronic temperature generation in an XY plane at the gold surface of the optimized nanostructure for different characteristic times after the optical pulse, a) 200 fs, b) 400 fs, c) 600 fs, d) 800 fs, e) 1000 fs, f) 2000 fs, g) 3000 fs, h) 4000 fs.

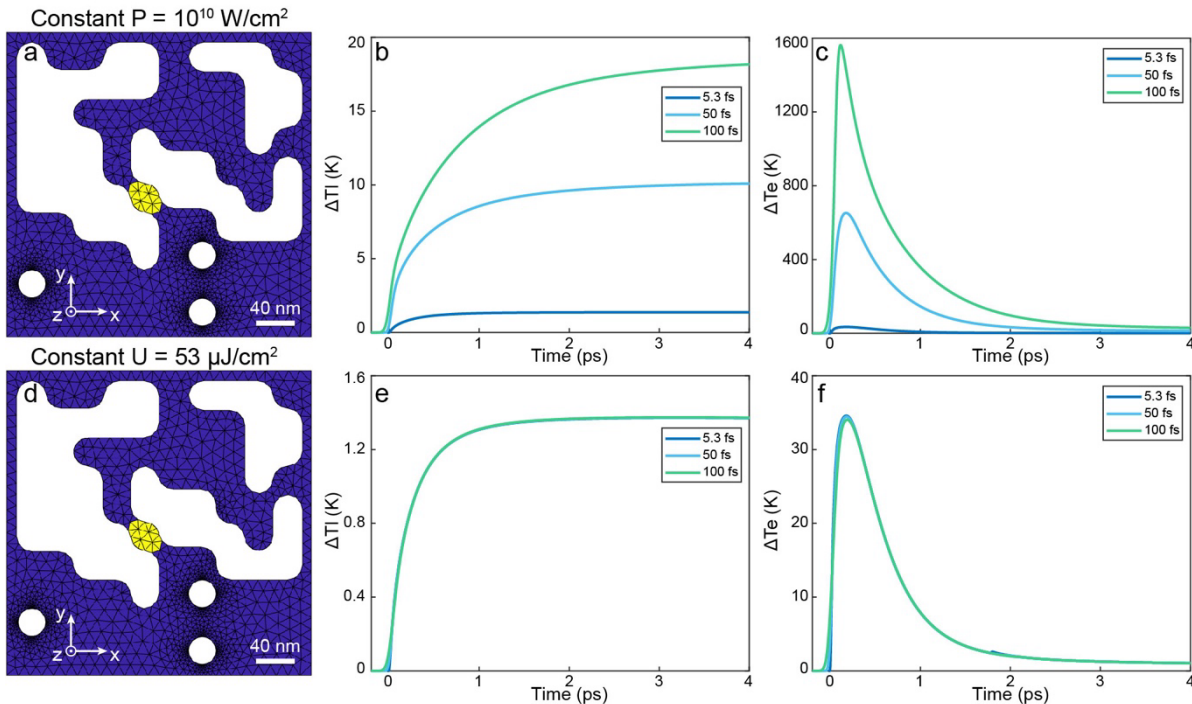


Figure S5. Time evolution of lattice and electronic temperatures for different optical pulse durations, for a) to c) the same peak power and d) to f) the same energy. a)

Schematic of the optimized structure with the yellow section indicating the area of temperature analysis. Time evolution of b) lattice temperatures and c) electronic temperatures for pulses of 5.3, 50, and 100 fs at a peak power of 10^{10} W/cm², showing similar dynamics. d) Schematic of the optimized structure, with the yellow area denoting the temperature study region. Time evolution of e) lattice and f) electronic temperatures for pulses of durations 5.3, 50, and 100 fs, each with an energy of 53 μ J/cm², displaying identical dynamics.

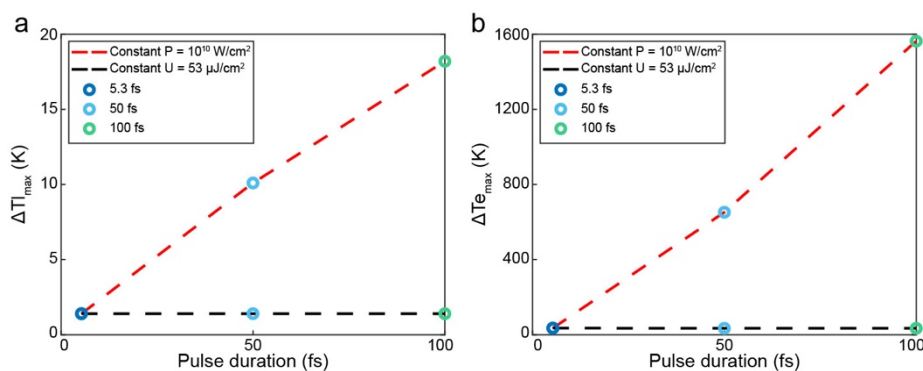


Figure S6. Evolution of maximum lattice and electronic temperatures based on the characteristics of the optical pulses used. a) Evolution of lattice temperatures for pulse durations of 5.3, 50, and 100 fs (represented by colored circles), analyzed under conditions of constant peak power (red dashed line) or constant pulse energy (black dashed line). It is observed that the temperature increases linearly in the studied range with constant peak powers, while remaining steady for constant pulse energies. b) Evolution of electronic temperatures for the same pulse durations (using the same color codes). The results are consistent: electronic temperatures rise linearly when the peak power is constant and remain unchanged with constant energy levels.

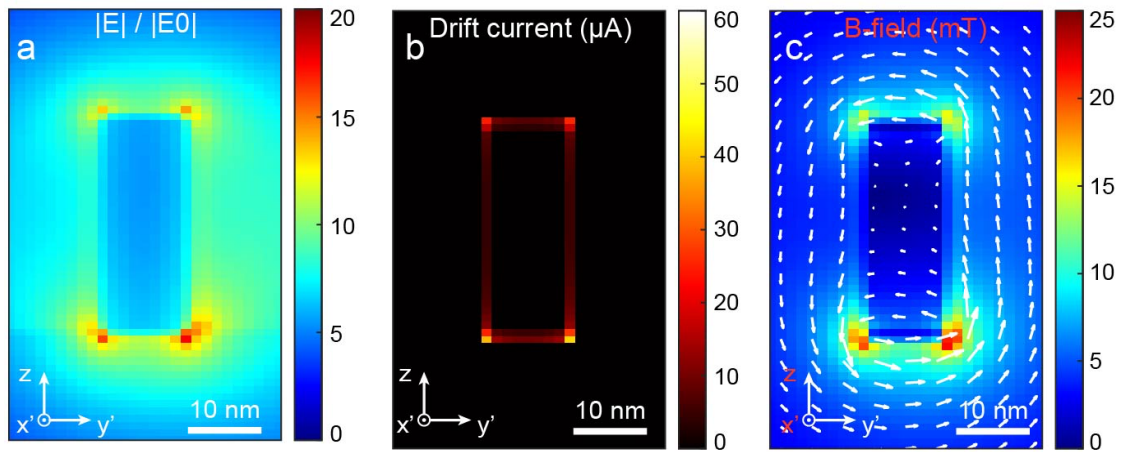


Figure S7. Spatial distribution of the electric field enhancement in a YZ plane of the X'Y'Z coordinate, symbolized by the pink dotted line and shown in Figure 4a of the manuscript. Distribution b) of the drift currents described by equation 1) and c) of the resulting magnetic field in the same plane as a).

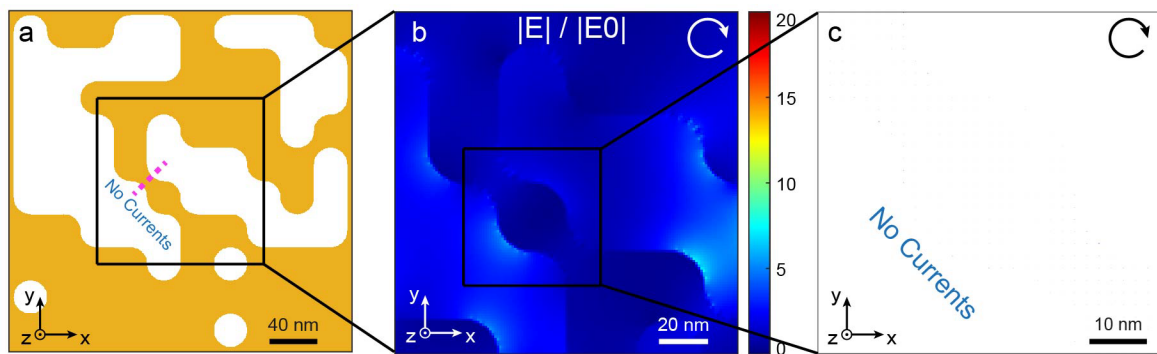


Figure S8. Optical response of the optimized plasmonic nanostructure and associated drift currents for excitation by a left circular polarization. a) Schematic, in an XY plane, of the GA-optimized structure. b) Spatial distribution of the electric field enhancement at the Z-center of the structure shown in a) for an excitation by a left circular polarization. c) Spatial distribution of the drift currents generated (or rather not generated) by the electric field distribution shown in b).

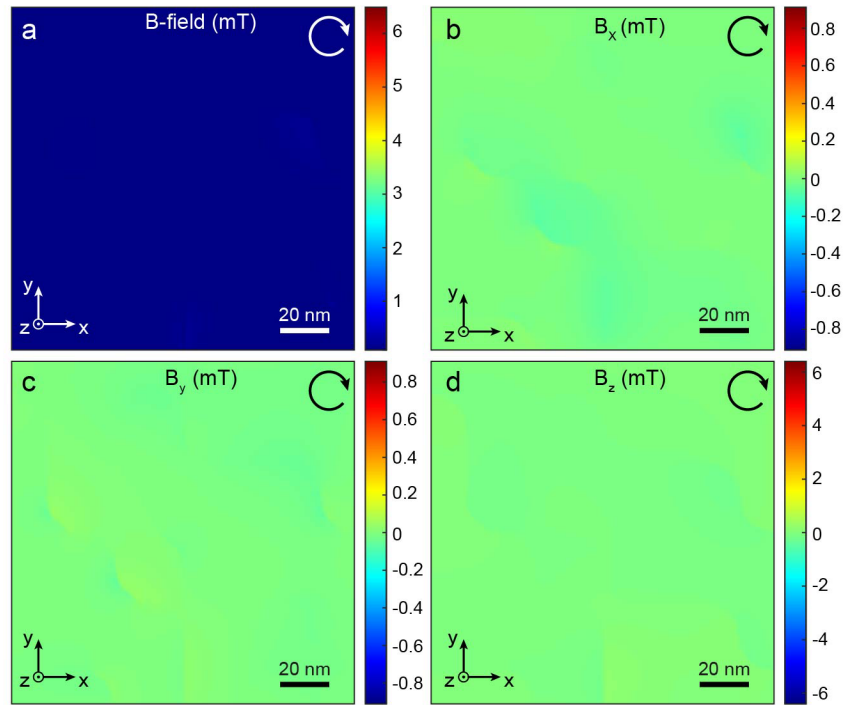


Figure S9. Magnetic response of the optimized plasmonic nanostructure under left circular polarized excitation. Spatial distribution of the a) total magnetic field and its different components along b) X, c) Y and d) Z, generated at the Z center of the structure shown in Figure 2a of the manuscript.

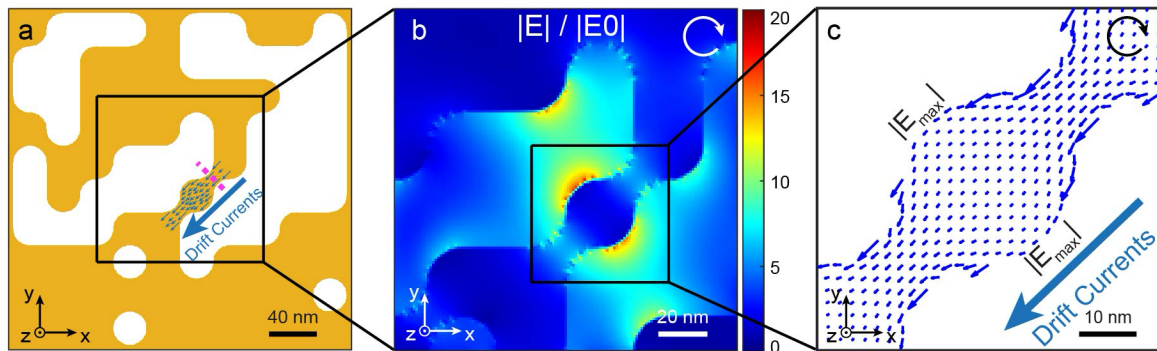


Figure S10. Optical response of the mirror plasmonic structure. a) Schematic representation of the GA-optimized mirror antenna. b) Electric field enhancement at the Z-center of the mirror structure in the area enclosed by a black square in a). c) Drift currents created by the electric field distribution described in b) via equation 1 at the surface of the structure.

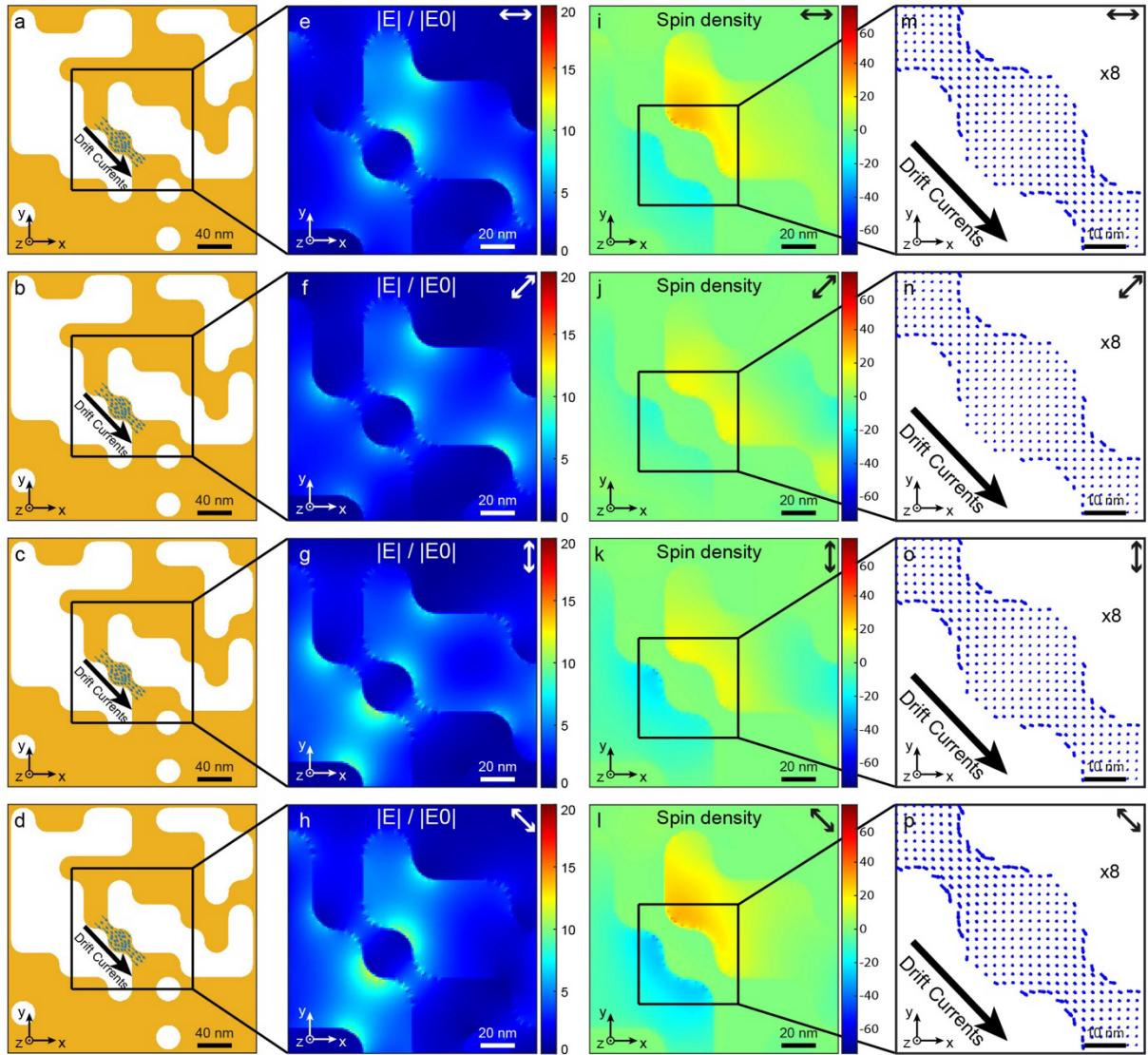


Figure S11: Optical response and associated drift currents in a GA-optimized plasmonic nanostructure for linear polarizations. a-d) Schematic representation of the structure in an XY plane. e-h) Spatial distribution of electric field enhancement at the Z -center of the structure for linear polarizations at 0° , 45° , 90° , and 135° orientations. i-l) Corresponding spatial distributions of spin density at the Z -center of the structure. m-p) Spatial distributions of drift currents generated on the surface of the structure by the preceding electric field and spin density distributions. The length of the blue arrows in m-p) indicates the relative amplitude of the drift currents (multiplied by 8 with respect to figure 2), and the double-ended arrows in the top right of e-p) show the polarization states of the incoming light.

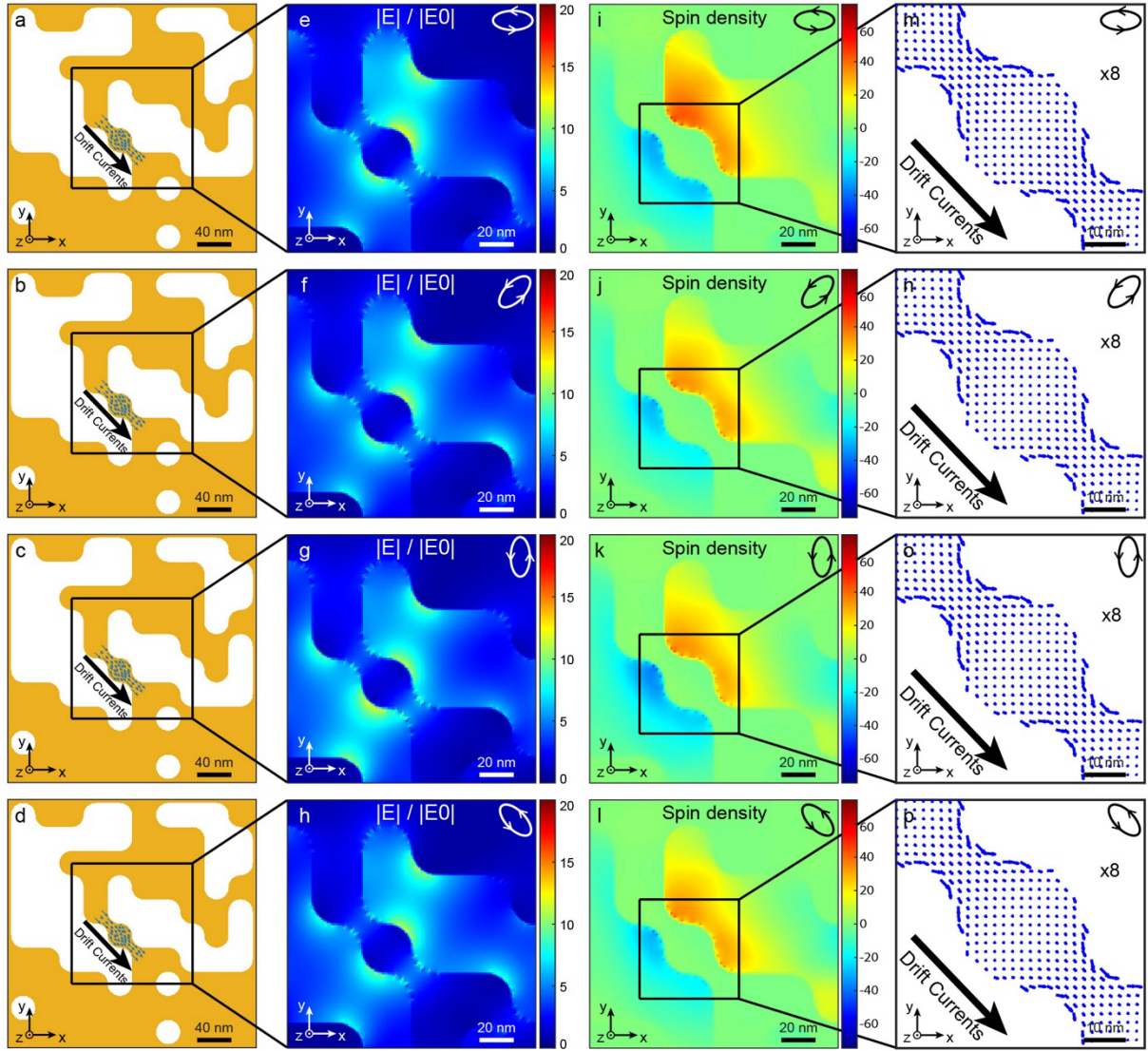


Figure S12: Optical response of the optimized plasmonic nanostructure and associated drift currents for right elliptical polarizations. a-d) Schematic representation in the XY plane of the GA-optimized structure. e-h) Spatial distribution of electric field enhancement at the Z-center of the structure for right elliptical polarizations at orientations of 0° , 45° , 90° , and 135° with $\pi/8$ ellipticity. i-l) Corresponding spatial distributions of spin density at the Z-center of the structure. m-p) Spatial distributions of drift currents generated on the structure's surface due to the electric field and spin density distributions depicted in e-h). The length of the blue arrows in m-p) indicates the relative amplitude of the drift currents (multiplied by 8 with respect to figure 2) and the ellipses in the top right of e-p) denote the polarization states of the incoming light.

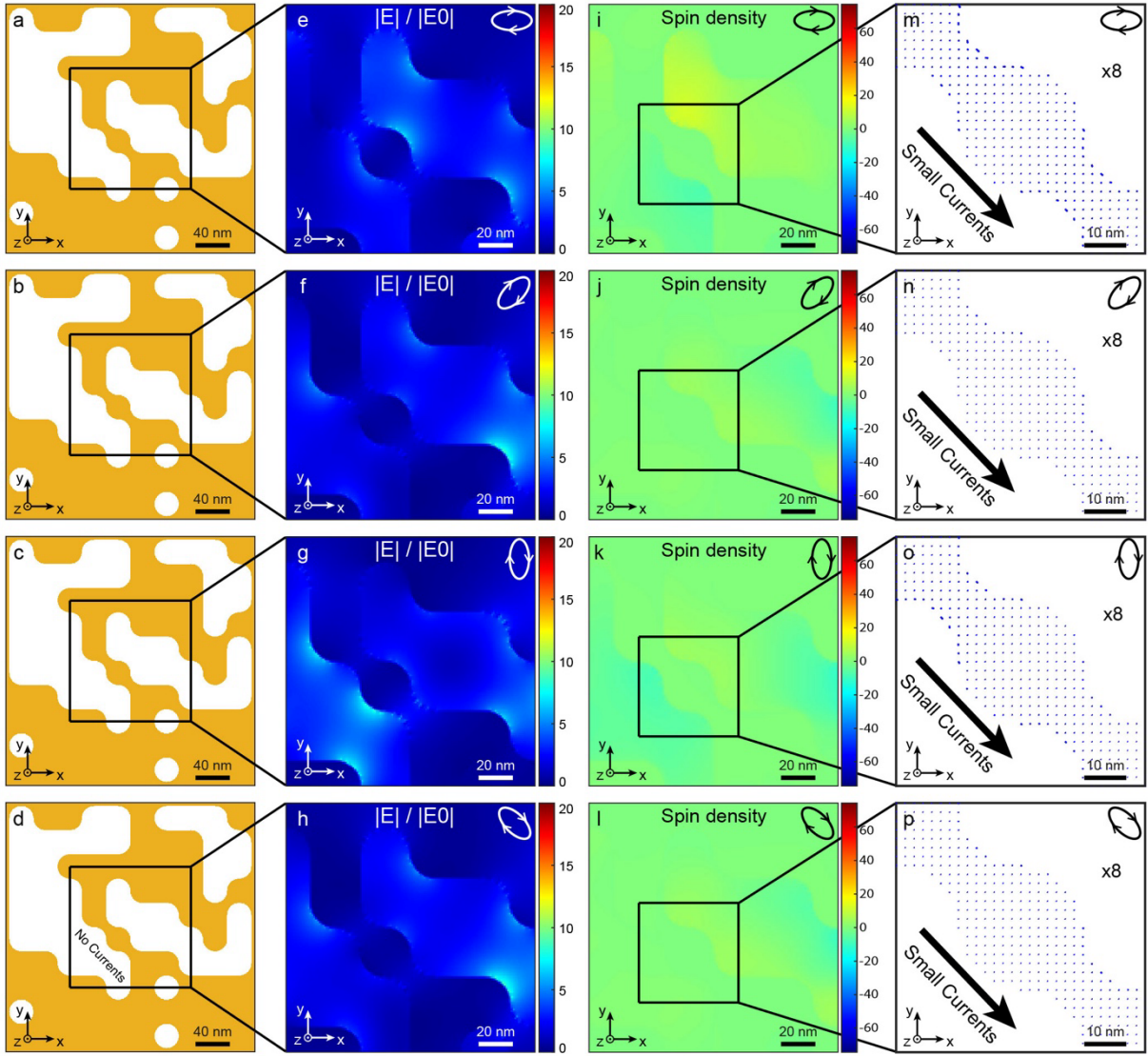


Figure S13: Optical response of the optimized plasmonic nanostructure and associated drift currents for left elliptical polarizations. a-d) Schematic representation in the XY plane of the GA-optimized structure. e-h) Spatial distribution of electric field enhancement at the Z-center of the structure for left elliptical polarizations at orientations of 0° , 45° , 90° , and 135° with $\pi/8$ ellipticity. i-l) Corresponding spatial distributions of spin density at the Z-center of the structure. m-p) Spatial distributions of drift currents generated on the structure's surface due to the electric field and spin density distributions depicted in e-h). The length of the blue arrows in m-p) indicates the relative amplitude of the drift currents (multiplied by 8 with respect to figure 2) and the ellipses in the top right of e-p) denote the polarization states of the incoming light.

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