



HAL
open science

Weak solutions for a singular beam equation

Olena Atlasiuk, Arnaud Heibig, Adrien Petrov

► **To cite this version:**

Olena Atlasiuk, Arnaud Heibig, Adrien Petrov. Weak solutions for a singular beam equation. 2024.
hal-04578854

HAL Id: hal-04578854

<https://hal.science/hal-04578854v1>

Preprint submitted on 17 May 2024

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Weak solutions for a singular beam equation

Olena Atlasiuk*

Arnaud Heibig[†]

Adrien Petrov[†]

May 16, 2024

Abstract

This paper deals with a dynamic Gao beam of infinite length subjected to a moving concentrated Dirac mass. Under appropriate regularity assumptions on the initial data, the problem possesses a weak solution which is obtained as the limit of a sequence of solutions of regularized problems.

Key words. Gao beam equation, Dirac mass, moving load, existence result, Sobolev spaces.

AMS Subject Classification: 35B45, 35Q74, 74H20, 74K10, 74M15

1 Description of the problem

The behavior of a beam plays a crucial role in various applications as in railway track design. The railway companies aim to enhance the train rolling speed to meet the increased demands in both passenger and freight transportation worldwide. Identifying the factors contributing to the occurrence of railway track defects is rather important to maintain the required track quality. The oscillations amplitudes in railway tracks due to the trains movement is intensively studied in scientific literature (see [12, 10]). These oscillations lead to undesirable consequences such as a premature wear and deformation of railway tracks. Understanding the impact of these oscillations on system reliability is essential to maintain the track quality necessary for traffic safety and passenger comfort.

Most of mathematical models consider rails as beams under the influence of moving loads and their properties are intensively studied. A considerable engineering and mathematical literature is devoted to study to numerous mathematical models for beams, such as the Euler-Bernoulli (see [8, 7, 1, 9]), Timoshenko (see [17, 6]) and Gao beams. The Euler-Bernoulli and Timoshenko beam theories are the oldest ones and they are widely employed nowadays in various structural analysis methodologies in engineering. The Euler-Bernoulli beam is commonly used to model the bending beams behavior, the axial and shear deformations are neglected compared to bending. Thus, the beam cross-section remains planar during loading. The Timoshenko beam is suitable to study the shear beams behavior, thereby providing an accurate representation of deformation within the cross-section of the beam.

The current study involves analysis of the horizontal motion of a vertical point-load on a metallic rail, considering scenarios where the load has mass or is massless, and examining the resulting oscillations of the system. The mathematical model of the Gao beam was originally introduced in [13]. However the Gao beam is also studied in different contexts (see [3, 2, 5, 15, 14, 16]). Notice also that semilinear beam is studied in [19, 18]. Comparisons between simulations

*Institute of Mathematics Czech Academy of Sciences & Institute of Mathematics of the National Academy of Sciences of Ukraine, Žitná 25, CZ-11567 Praha 1, Czech Republic, (atlasiuk@math.cas.cz)

[†]Univ Lyon, INSA Lyon, UJM, UCBL, ECL, CNRS UMR 5208, ICJ, 43 boulevard du 11 novembre 1918, F-69622 Villeurbanne Cedex, France (arnaud.heibig@insa-lyon.fr, adrien.petrov@insa-lyon.fr)

of the Gao beam and the Euler-Bernoulli linear beam reveal significant differences, indicating that the linear beam is suitable only for small loads, whereas the Gao beam accommodates moderate loads (see [11]).

We consider in this paper a straight infinite Gao beam of thickness $2h$. A detailed derivation of the Gao beam model can be found in [13, 3]. An horizontal traction p , which is a time-dependent function, is applied at one end. We are focused in this work on examining the sideways movement of a concentrated load $P(t)$, where the load may change over time. This load is applied to a mobile mass m positioned at $\zeta(t)$ with a horizontal velocity $\dot{\zeta}(t)$, induced by a horizontal applied force. The transverse displacement of the Gao beam $u(t, x)$ for $(t, x) \in (0, T) \times \mathbb{R}$ is governed by a following partial differential equation:

$$\begin{aligned} \varrho u_{tt} + m\delta(x - \zeta(t))u_{tt} - m\delta'(x - \zeta(t))\dot{\zeta}(t)u_t + ku_{xxxx} - (eu_x^2 - \nu p)u_{xx} \\ = \varrho f + \delta(x - \zeta(t))P(t), \end{aligned} \quad (1.1)$$

where $\delta(\cdot)$ is the Dirac function, ϱ is the material density, f is the applied mechanical loading, $k \stackrel{\text{def}}{=} \frac{2h^3 E}{3(1-\bar{\nu}^2)}$, $\nu \stackrel{\text{def}}{=} (1 + \bar{\nu})$ and $e \stackrel{\text{def}}{=} 3hE$ where E and $\bar{\nu}$ denote the Young modulus and the Poisson ratio, respectively. Here and below $(\cdot)_t \stackrel{\text{def}}{=} \partial_t(\cdot)$, $(\cdot)_x \stackrel{\text{def}}{=} \partial_x(\cdot)$ denote the partial derivatives with respect to t and x , respectively, while (\cdot) and $(\cdot)'$ denote the derivatives with respect to t and x , respectively. We prescribe also initial data

$$u(0, \cdot) = u_0 \quad \text{and} \quad u_t(0, \cdot) = u_1. \quad (1.2)$$

For further explanations on this model, the reader is referred to [11] and to the references therein. Equation (1.1) can be considered the limit of the equation treated in [11] when the gamma viscosity tends to zero, and is related to the usual viscous beam equation without Dirac measure in the time derivative. Nevertheless, this leads to an entirely different mathematical problem since three-order estimates cannot be obtained using the viscous term, and what is more, since the usual energy estimates fail. In particular, the non-linear term should be addressed somewhat indirectly since no L^∞ estimate is available immediately.

The positive constants ϱ , m , k and e play any role in the mathematical analysis carried out below. Consequently, without loss of generality, we set them equal to 1. Notice, however, that the case of non constant coefficients (for example $k = k(x)$) could be processed using approximate square roots for elliptic operators.

Throughout this paper, we assume that $\zeta \in C^2([0, T])$, $P \in C^2([0, T])$, $p \in C^0([0, T]; H^2(\mathbb{R}))$ and $f \in C^0([0, T]; H^2(\mathbb{R}))$, but these assumptions could be weakened. The main result of this paper can be stated as follows.

Theorem 1.1 *Let $T > 0$. Assume that $u_0 \in H^2(\mathbb{R})$, $u_1 \in H^1(\mathbb{R})$, $\zeta \in C^2([0, T])$, $P \in C^2([0, T])$ and $f, p \in C^0([0, T]; H^2(\mathbb{R}))$. Then there exists a function $u \in C^0([0, T]; H_{\text{loc}}^{2-\alpha}(\mathbb{R})) \cap L^\infty([0, T]; H^2(\mathbb{R}))$, for any $\alpha \in]0, 2[$, and $u_t \in L^\infty([0, T]; L^2(\mathbb{R}))$ such that for any $v \in C^2([0, T]; H^2(\mathbb{R}))$ with a compact support in $[0, T] \times \mathbb{R}$, we have*

$$\left\{ \begin{aligned} & \int_0^T \int_{\mathbb{R}} u_t v_t dx dt + \int_0^T \int_{\mathbb{R}} u_{xx} v_{xx} dx dt + \frac{1}{3} \int_0^T \int_{\mathbb{R}} u_x^3 v_x dx dt \\ & - \int_0^T \int_{\mathbb{R}} \nu p u_{xx} v dx dt - u_0(\zeta(0))v_t(0, \zeta(0)) - \int_0^T u(t, \zeta(t))v_{tt}(t, \zeta(t)) dt \\ & - \int_0^T \dot{\zeta}(t)u_x(t, \zeta(t))v_t(t, \zeta(t)) dt - \int_0^T \dot{\zeta}(t)u(t, \zeta(t))v_{xt}(t, \zeta(t)) dt \\ & + \int_0^T \int_{\mathbb{R}} f v dx dt + \int_0^T P(t)v(t, \zeta(t)) dt + \int_{\mathbb{R}} u_1 v(0, \cdot) dx \\ & + u_1(\zeta(0))v(0, \zeta(0)) = 0. \end{aligned} \right. \quad (1.3)$$

Moreover, $u(0, \cdot) = u_0$ holds.

In this statement, we refrain from expressing the term $\int_0^T \int_{\mathbb{R}} \frac{d}{dt}(\delta(x-\xi(t))\partial_t u(t, \xi(t)))v(t, x) dx dt$ as $-\int_0^T \partial_t u(t, \xi(t))\partial_t v(t, \xi(t)) dt$. Indeed, the trace is not defined within our functional frame. To establish such an expression, higher-order regularity would be required. However, in the linear homogenous case, the formula $\partial_t^{(2)}u(t, x) + \partial_x^{(4)}u(t, x) = -\frac{d}{dt}(\delta(x-\zeta(t))\partial_t u(t, x))$ entails that $\partial_t^{(2)}u$ and $\partial_x^{(4)}u$ cannot both possess $L^2(0, T; L^2(\mathbb{R}))$ regularity. The same formula suggests that $\partial_t^{(2)}u \notin C^0([0, T]; L^2(\mathbb{R}))$ as the formal initial data $\partial_t^{(2)}u(0, x) = -(\partial_x^{(4)}u(0, x) - \frac{d}{dt}(\delta(x-\zeta(0))\partial_t u(0, x)))$ does not belong to $L^2(\mathbb{R})$ for a smooth data $u_0(x) = u(0, x)$. Lastly, the usual energy estimate, obtained by multiplying (1.1) by $\partial_t u$ and by integrating the resulting result over $[0, T] \times \mathbb{R}$, can not be performed due to the term $\delta(x-\zeta(t))$.

The paper is organized as follows. An existence result for a linear mollified problem is presented in Section 2. Section 3 addresses a mollified nonlinear equation; some uniform a priori estimates are derived. The proof of Theorem 1.1 is provided in Section 4.

2 A preliminary existence result

Under some regularity assumptions on the data, we present an existence and uniqueness result to the system (2.2) in a suitable function space. To this aim, for any $t > 0$, we introduce the following sets:

$$E_2(t) \stackrel{\text{def}}{=} C^0([0, t]; H^2(\mathbb{R})), \quad (2.1a)$$

$$E_1(t) \stackrel{\text{def}}{=} C^0([0, t]; H^2(\mathbb{R})) \cap C^1([0, t]; L^2(\mathbb{R})), \quad (2.1b)$$

$$E_0(t) \stackrel{\text{def}}{=} C^0([0, t]; H^4(\mathbb{R})) \cap C^1([0, t]; H^2(\mathbb{R})) \cap C^2([0, t]; L^2(\mathbb{R})). \quad (2.1c)$$

Hence, we have $E_0(t) \hookrightarrow E_1(t) \hookrightarrow E_2(t)$. Next, let $\theta \in \mathcal{D}(\mathbb{R})$ be an even density of probability. For any $(t, x) \in \mathbb{R}_+ \times \mathbb{R}$, we set

$$F(t, x) \stackrel{\text{def}}{=} \frac{1}{1 + \theta(x - \zeta(t))} \quad \text{and} \quad G(t, x) \stackrel{\text{def}}{=} -\frac{\dot{\zeta}(t)\theta'(x - \zeta(t))}{1 + \theta(x - \zeta(t))}.$$

We denote below by $C_\eta > 0$ a generic constant depending on η .

Proposition 2.1 *Let $T > 0$, $u_0 \in H^4(\mathbb{R})$, $u_1 \in H^2(\mathbb{R})$ and $g \in E_2(T)$. Then, there exists a unique $u \in E_0(T)$ solution to the following system:*

$$\begin{cases} u_{tt} + Fu_{xxxx} + Gu_t = g, & (2.2a) \\ u(0, \cdot) = u_0 \quad \text{and} \quad u_t(0, \cdot) = u_1. & (2.2b) \end{cases}$$

Moreover, u satisfies the following inequality:

$$\|u(t, \cdot)\|_{H^2(\mathbb{R})} + \|u_t(t, \cdot)\|_{L^2(\mathbb{R})} \leq C_T (\|u_0\|_{H^2(\Omega)} + \|u_1\|_{L^2(\Omega)} + \|g\|_{L^2(0, t; L^2(\mathbb{R}))}). \quad (2.3)$$

The proof is omitted. It follows from Proposition 2.1 that the application

$$\begin{aligned} \mathcal{A}_t : H^4(\mathbb{R}) \times H^2(\mathbb{R}) \times E_2(t) &\rightarrow E_0(t) \\ (u_0, u_1, g) &\mapsto u \end{aligned} \quad (2.4)$$

where u denotes the solution introduced in Proposition 2.1, can be continuously extended to a function, still denoted by \mathcal{A}_t :

$$\mathcal{A}_t : H^2(\mathbb{R}) \times L^2(\mathbb{R}) \times L^2([0, t]; L^2(\mathbb{R})) \rightarrow E_1(t).$$

For $(u_0, u_1) \in \mathbf{H}^2(\mathbb{R}) \times \mathbf{L}^2(\mathbb{R})$ fixed, we define

$$\begin{aligned} \mathcal{B}_t &: \mathbf{L}^2([0, t]; \mathbf{L}^2(\mathbb{R})) \rightarrow \mathbf{E}_1(t) \\ g &\mapsto \mathcal{A}_t(u_0, u_1, g). \end{aligned}$$

If we denote by $\mathcal{B}_t(g)$ by u , then u still satisfies (2.3). In the sequel, we always assume that $(u_0, u_1) \in \mathbf{H}^2(\mathbb{R}) \times \mathbf{L}^2(\mathbb{R})$.

3 The approximated problem

We mollify in system (2.2) the δ function, mollify and truncate the nonlinear term. Namely, we consider the following approximated problem:

$$\begin{cases} u_{tt} + Fu_{xxxx} + Gu_t - ((\varphi^R(u \star \theta') - \nu p)(u \star \theta'')) \star \theta + h)F = 0, & (3.1a) \\ u(0, \cdot) = u_0 \quad \text{and} \quad u_t(0, \cdot) = u_1, & (3.1b) \end{cases}$$

where \star denotes the convolution with respect to x , and $h \in \mathbf{C}^0([0, T]; \mathbf{H}^2(\mathbb{R}))$ is a given function. Here, for $R > 0$, $\varphi^R \in \mathbf{C}^\infty(\mathbb{R})$ is an even, increasing function on \mathbb{R}_+ such that $\varphi^R(x) \stackrel{\text{def}}{=} x^2$ for $|x| \leq R$ and $\varphi^R(x) \stackrel{\text{def}}{=} (R+1)^2$ for $|x| \geq R+2$. Clearly, the equation (3.1a) can formally be rewritten as follows:

$$u_{tt} + Fu_{xxxx} + Gu_t = \mathcal{C}_{t,R}(u),$$

where $\mathcal{C}_{t,R}$ is defined by

$$\begin{aligned} \mathcal{C}_{t,R} &: \mathbf{E}_1(t) \rightarrow \mathbf{E}_2(t), \\ v &\mapsto \mathcal{C}_{t,R}(v) \stackrel{\text{def}}{=} ([(\varphi^R(v \star \theta') - \nu p)(v \star \theta'')] \star \theta + h)F \end{aligned}$$

The fact that $\mathcal{C}_{t,R}$ is well defined with values in $\mathbf{E}_2(t)$ comes from usual convolution inequalities. Similarly, we have

Proposition 3.1 *Let $t \in [0, T]$. Then, for any $(v, \tilde{v}) \in \mathbf{E}_1(t) \times \mathbf{E}_1(t)$, the following inequality*

$$\begin{aligned} \|\mathcal{C}_{t,R}(v) - \mathcal{C}_{t,R}(\tilde{v})\|_{\mathbf{L}^2(0,t;\mathbf{L}^2(\mathbb{R}))} &\leq C_{T,R}(\|v - \tilde{v}\|_{\mathbf{L}^2([0,t];\mathbf{L}^2(\mathbb{R}))}) \\ &+ \|v_t - \tilde{v}_t\|_{\mathbf{L}^2([0,t];\mathbf{L}^2(\mathbb{R}))} + \|\tilde{v}\|_{\mathbf{C}^0([0,t];\mathbf{L}^2(\mathbb{R}))} \|v - \tilde{v}\|_{\mathbf{L}^2([0,t];\mathbf{L}^2(\mathbb{R}))}) \end{aligned} \quad (3.2)$$

holds true.

Proof. Let $(v, \tilde{v}) \in \mathbf{E}_1(t) \times \mathbf{E}_1(t)$. Hence, we have

$$\begin{aligned} &\| [(\varphi^R(v \star \theta') - \nu p)(v \star \theta'')] \star \theta - [(\varphi^R(\tilde{v} \star \theta') - \nu p)(\tilde{v} \star \theta'')] \star \theta) F \|_{\mathbf{L}^2([0,t];\mathbf{L}^2(\mathbb{R}))} \\ &\leq C_{T,R}(\|(\varphi^R(v \star \theta') - \nu p)(v - \tilde{v}) \star \theta''\|_{\mathbf{L}^2([0,t];\mathbf{L}^2(\mathbb{R}))}) \\ &+ \|(\varphi^R(v \star \theta') - \varphi^R(\tilde{v} \star \theta'))(\tilde{v} \star \theta'')\|_{\mathbf{L}^2([0,t];\mathbf{L}^2(\mathbb{R}))}. \end{aligned}$$

Since φ^R is bounded and globally Lipschitz and by convolution inequalities, it comes that

$$\begin{aligned} &\| [(\varphi^R(v \star \theta') - \nu p)(v \star \theta'')] \star \theta - [(\varphi^R(\tilde{v} \star \theta') - \nu p)(\tilde{v} \star \theta'')] \star \theta) F \|_{\mathbf{L}^2([0,t];\mathbf{L}^2(\mathbb{R}))} \\ &\leq C_{T,R}(\|v - \tilde{v}\|_{\mathbf{L}^2([0,t];\mathbf{L}^2(\mathbb{R}))} + \|\tilde{v}\|_{\mathbf{C}^0([0,t];\mathbf{L}^2(\mathbb{R}))} \|v - \tilde{v}\|_{\mathbf{L}^2([0,t];\mathbf{L}^2(\mathbb{R}))}), \end{aligned}$$

which proves the result. \square

We are now looking for the fixed points of application:

$$\mathcal{B}_T \mathcal{C}_{T,R} : \mathbf{E}_1(T) \rightarrow \mathbf{E}_1(T).$$

Proposition 3.2 *The application $\mathcal{B}_T\mathcal{C}_{T,R}$ admits a fixed point.*

Proof. We use the Picard fixed point theorem. Due to Proposition 3.1 (i), we have mainly to bound $\|\tilde{v}\|_{C^0([0,t];L^2(\mathbb{R}))}$.

(a) An invariant set.

First, let us assume that $v \in E_1(T)$ with $(v(0, \cdot), v_t(0, \cdot)) = (u_0, u_1)$ and $\tilde{v} = 0$. It follows from (2.3) and (3.2) that

$$\begin{aligned} & \|(\mathcal{B}_T\mathcal{C}_{T,R}(v))(t)\|_{H^2(\mathbb{R})} + \|\partial_t(\mathcal{B}_T\mathcal{C}_{T,R}(v))(t)\|_{L^2(\mathbb{R})} \\ & \leq C_{T,R}(\|u_0\|_{H^2(\mathbb{R})} + \|u_1\|_{L^2(\mathbb{R})} + \|\mathcal{C}_{T,R}(v)\|_{L^2([0,T];L^2(\mathbb{R}))}) \\ & \leq C_{T,R}(\|u_0\|_{H^2(\mathbb{R})} + \|u_1\|_{L^2(\mathbb{R})} + \|\mathcal{C}_{T,R}(0)\|_{L^2([0,T];L^2(\mathbb{R}))} \\ & \quad + C_{T,R}(\|v\|_{L^2([0,t];L^2(\mathbb{R}))} + \|v_t\|_{L^2([0,t];L^2(\mathbb{R}))})). \end{aligned} \quad (3.3)$$

For $f \in C^0([0, t]; L^2(\mathbb{R}))$, we define

$$\|f\|_{\lambda,t} \stackrel{\text{def}}{=} \sup_{0 \leq s \leq t} (e^{-\lambda s} \|f(s)\|_{L^2(\mathbb{R})})$$

with $\lambda > 0$ and for $g \in E_1(t)$, we define

$$\|g\|_{1,\lambda,t} \stackrel{\text{def}}{=} \sup_{0 \leq s \leq t} (e^{-\lambda s} (\|g(s)\|_{H^2(\mathbb{R})} + \|g_t(s)\|_{L^2(\mathbb{R})})).$$

It comes from (3.3) that

$$\begin{aligned} & e^{-\lambda t} (\|(\mathcal{B}_T\mathcal{C}_{T,R}(v))(t)\|_{H^2(\mathbb{R})} + \|\partial_t(\mathcal{B}_T\mathcal{C}_{T,R}(v))(t)\|_{L^2(\mathbb{R})}) \\ & \leq C_{T,R}e^{-\lambda t} (\|u_0\|_{H^2(\mathbb{R})} + \|u_1\|_{L^2(\mathbb{R})} + \|\mathcal{C}_{T,R}(0)\|_{L^2([0,T];L^2(\mathbb{R}))}) \\ & \quad + C_{T,R} \left(\int_0^t e^{-2\lambda t} e^{2\lambda s} e^{-2\lambda s} (\|v(s)\|_{L^2(\mathbb{R})}^2 + \|v_t(s)\|_{L^2(\mathbb{R})}^2) ds \right)^{1/2} \\ & \leq C_{T,R}e^{-\lambda t} (\|u_0\|_{H^2(\mathbb{R})} + \|u_1\|_{L^2(\mathbb{R})} + \|\mathcal{C}_{T,R}(0)\|_{L^2([0,T];L^2(\mathbb{R}))}) \\ & \quad + \frac{C_{T,R}}{\sqrt{\lambda}} \|v\|_{1,\lambda,t}. \end{aligned} \quad (3.4)$$

Note that (3.4) can be rewritten for any $s \in [0, t]$ in place of t . Furthermore, in this new inequality, $\|\cdot\|_{1,\lambda,s}$ and $\|\cdot\|_{L^2(0,s;L^2(\mathbb{R}))}$ are respectively smaller than $\|\cdot\|_{1,\lambda,t}$ and $\|\cdot\|_{L^2(0,t;L^2(\mathbb{R}))}$. Hence, we get (3.4) with s in place of t in the left hand side. Taking $t = T$, we obtain

$$\|\mathcal{B}_T\mathcal{C}_{T,R}(v)\|_{1,\lambda,T} \leq M_R + \frac{C_{T,R}}{\sqrt{\lambda}} \|v\|_{1,\lambda,T}, \quad (3.5)$$

where $M_R \stackrel{\text{def}}{=} C_{T,R}(\|u_0\|_{H^2(\mathbb{R})} + \|u_1\|_{L^2(\mathbb{R})} + \|\mathcal{C}_{T,R}(0)\|_{L^2([0,T];L^2(\mathbb{R}))})$. Let $\lambda = 4C_{T,R}^2$ and define $\beta(0, 2M_R) \stackrel{\text{def}}{=} \{u \in E_1(T) : \|u\|_{1,\lambda,T} \leq 2M_R\}$. Then, we may deduce from (3.5) that $\mathcal{B}_T\mathcal{C}_{T,R}(\beta(0, 2M_R)) \subset \beta(0, 2M_R)$.

(b) Contraction of $\mathcal{B}_T\mathcal{C}_{T,R}$.

Let $(v, \tilde{v}) \in \beta(0, 2M_R) \times \beta(0, 2M_R)$ with $v(0, \cdot) = \tilde{v}(0, \cdot) = u_0$ and $v_t(0, \cdot) = \tilde{v}_t(0, \cdot) = u_1$. Since $v - \tilde{v}$ has 0 initial data, we have (see (2.3))

$$\begin{aligned} & \|(\mathcal{B}_T\mathcal{C}_{T,R}(v))(t) - (\mathcal{B}_T\mathcal{C}_{T,R}(\tilde{v}))(t)\|_{H^2(\mathbb{R})} + \|\partial_t((\mathcal{B}_T\mathcal{C}_{T,R}(v))(t) - (\mathcal{B}_T\mathcal{C}_{T,R}(\tilde{v}))(t))\|_{L^2(\mathbb{R})} \\ & \leq C_{T,R} \|\mathcal{C}_{T,R}(v) - \mathcal{C}_{T,R}(\tilde{v})\|_{L^2([0,t];L^2(\mathbb{R}))} \\ & \leq C_{T,R} (\|v - \tilde{v}\|_{L^2([0,T];L^2(\mathbb{R}))} + \|v_t - \tilde{v}_t\|_{L^2([0,T];L^2(\mathbb{R}))} \\ & \quad + \|\tilde{v}\|_{C^0([0,T];L^2(\mathbb{R}))} \|v - \tilde{v}\|_{L^2([0,T];L^2(\mathbb{R}))}), \end{aligned}$$

Since $\tilde{v} \in \beta(0, 2M_R)$, we have $\|\tilde{v}\|_{C^0([0,t];L^2(\mathbb{R}))} \leq 2M_R e^{\lambda T}$. It follows that

$$\begin{aligned} & \|(\mathcal{B}_T \mathcal{C}_{T,R}(v))(t) - (\mathcal{B}_T \mathcal{C}_{T,R}(\tilde{v}))(t)\|_{H^2(\mathbb{R})} + \|\partial_t((\mathcal{B}_T \mathcal{C}_{T,R}(v))(t) - (\mathcal{B}_T \mathcal{C}_{T,R}(\tilde{v}))(t))\|_{L^2(\mathbb{R})} \\ & \leq C_{T,R}(\|v - \tilde{v}\|_{L^2([0,T];L^2(\mathbb{R}))} + \|v_t - \tilde{v}_t\|_{L^2([0,T];L^2(\mathbb{R}))}). \end{aligned}$$

Introducing a suitable $\mu > 0$, we derive as before the following inequality

$$\|\mathcal{B}_T \mathcal{C}_{T,R}(v) - \mathcal{B}_T \mathcal{C}_{T,R}(\tilde{v})\|_{1,\mu,T} \leq \frac{1}{2} \|v - \tilde{v}\|_{1,\mu,T},$$

which proves the Proposition. \square

An immediate consequence of Proposition 3.1 is the following Corollary:

Corollary 3.3 *Assume that $u_0 \in H^4(\mathbb{R})$, $u_1 \in H^2(\mathbb{R})$. Let $u^R \in E_1(T)$ such that $\mathcal{B}_T \mathcal{C}_{T,R}(u^R) = u^R$. Then, $u^R \in E_0(T)$. Moreover,*

$$\begin{cases} u_{tt}^R + F u_{xxxx}^R + G u_t^R - F((\varphi^R(u^R \star \theta') - \nu p)(u^R \star \theta'')) \star \theta + h = 0, & (3.6a) \end{cases}$$

$$\begin{cases} u^R(0, \cdot) = u_0 \quad \text{and} \quad u_t^R(0, \cdot) = u_1. & (3.6b) \end{cases}$$

Proof. We have $\mathcal{C}_{T,R}(E_1(T)) \subset E_2(T)$. Hence, by (2.4), we get $\mathcal{B}_T \mathcal{C}_{T,R}(u^R) \in E_0(T)$, which is $u^R \in E_0(T)$. \square

Let $\epsilon > 0$ and $\theta^\epsilon(x) \stackrel{\text{def}}{=} \frac{1}{\epsilon} \theta(\frac{x}{\epsilon})$ with $x \in \mathbb{R}$. We assume now that $u_0 \in H^2(\mathbb{R})$ and $u_1 \in L^2(\mathbb{R}) \cap L^\infty(\mathbb{R})$ and set $u_0^\epsilon \stackrel{\text{def}}{=} u_0 \star \theta^\epsilon$, $v_0^\epsilon \stackrel{\text{def}}{=} u_1 \star \theta^\epsilon$ and $h^\epsilon \stackrel{\text{def}}{=} -(f + \theta^\epsilon(\cdot - \zeta(t))P(t))$. Clearly, we have $u_0^\epsilon \in H^4(\mathbb{R})$ and $u_1^\epsilon \in H^2(\mathbb{R})$. We apply the Corollary 3.3 with $(u_0^\epsilon, u_1^\epsilon, \theta^\epsilon, h^\epsilon)$ in place of (u_0, u_1, θ, h) . This provides a function $u^{R,\epsilon}$ solution of the following problem:

$$\begin{cases} u_{tt}^{R,\epsilon} + \frac{d}{dt}(\theta^\epsilon(\cdot - \zeta(t))u_t^{R,\epsilon}) + u_{xxxx}^{R,\epsilon} & (3.7a) \\ -((\varphi^R(u^{R,\epsilon} \star \theta^{\epsilon'}) - \nu p)(u^{R,\epsilon} \star \theta^{\epsilon''})) \star \theta^\epsilon + h^\epsilon = 0, & \\ u^{R,\epsilon}(0, \cdot) = u_0^\epsilon \quad \text{and} \quad u_t^{R,\epsilon}(0, \cdot) = u_1^\epsilon. & (3.7b) \end{cases}$$

In the sequel, we set $Q_t \stackrel{\text{def}}{=} [0, t] \times \mathbb{R}$ where $t \in [0, T]$. We have the following (uniform in R and ϵ) estimates:

Proposition 3.4 *Assume that $u_0 \in H^2(\mathbb{R})$, $u_1 \in L^2(\mathbb{R}) \cap L^\infty(\mathbb{R})$. Then, there exists a constant $C > 0$, independent of R and ϵ , such that*

$$\sup_{0 \leq t \leq T} (\|u^{R,\epsilon}(t)\|_{H^2(\mathbb{R})} + \|u_t^{R,\epsilon}(t)\|_{L^2(\mathbb{R})}) \leq C.$$

Proof. We define $\frac{\partial}{\partial \tau} \stackrel{\text{def}}{=} \frac{\partial}{\partial t} + \dot{\zeta}(t) \frac{\partial}{\partial x}$. Then, we multiply (3.7a) by $u_\tau^{R,\epsilon}$ and we integrate this result over $[0, t] \times \mathbb{R}$ with $t \in [0, T]$. Since $u^{R,\epsilon} \in E_0(T)$, the boundary terms in $x = \pm\infty$ vanish. We observe that the term $u_{tt}^{R,\epsilon}$ in (3.7) gives

$$\begin{aligned} A_1(t) & \stackrel{\text{def}}{=} \frac{1}{2} \|u_t^{R,\epsilon}(t)\|_{L^2(\mathbb{R})}^2 - \frac{1}{2} \|u_1\|_{L^2(\mathbb{R})}^2 + \int_{Q_t} \dot{\zeta} u_{tt}^{R,\epsilon} u_x^{R,\epsilon} ds dx \\ & \geq \frac{1}{2} \|u_t^{R,\epsilon}(t)\|_{L^2(\mathbb{R})}^2 - \frac{1}{2} \|u_1\|_{L^2(\mathbb{R})}^2 - \frac{1}{2} \int_{Q_t} \partial_x(\dot{\zeta}(u_t^{R,\epsilon})^2) ds dx \\ & \quad - \int_{Q_t} \ddot{\zeta} u_x^{R,\epsilon} u_t^{R,\epsilon} ds dx + \int_{\mathbb{R}} [\dot{\zeta} u_t^{R,\epsilon} u_x^{R,\epsilon}]_0^t dx \\ & \geq \left(\frac{1}{2} - \eta\right) \|u_t^{R,\epsilon}(t)\|_{L^2(\mathbb{R})}^2 - C \int_0^t (\|u_t^{R,\epsilon}(s)\|_{L^2(\mathbb{R})}^2 + \|u_x^{R,\epsilon}(s)\|_{L^2(\mathbb{R})}^2) ds \\ & \quad - C_\eta \|u_x^{R,\epsilon}(t)\|_{L^2(\mathbb{R})}^2 - C(\|u_0\|_{H^2(\mathbb{R})}^2 + \|u_1\|_{L^2(\mathbb{R})}^2). \end{aligned}$$

Concerning the term $\frac{d}{dt}(\theta^\epsilon(x - \zeta(t))u_t^{R,\epsilon})$, we have

$$\begin{aligned}
A_2(t) &\stackrel{\text{def}}{=} \int_{Q_t} \frac{d}{ds} (\theta^\epsilon(x - \zeta(s))u_t^{R,\epsilon})u_\tau^{R,\epsilon} ds dx \\
&= - \int_{Q_t} \theta^\epsilon(x - \zeta(s))u_t^{R,\epsilon} \partial_\tau \partial_t u^{R,\epsilon} + \ddot{\zeta} u_x^{R,\epsilon} ds dx + \int_{\mathbb{R}} [\theta^\epsilon(x - \zeta(s))u_t^{R,\epsilon} u_\tau^{R,\epsilon}]_0^t dx \\
&= - \frac{1}{2} \int_{Q_t} \partial_\tau (\theta^\epsilon(x - \zeta(s))(u_t^{R,\epsilon})^2) ds dx - \int_{Q_t} \theta^\epsilon(x - \zeta(t)) \ddot{\zeta} u_t^{R,\epsilon} u_x^{R,\epsilon} ds dx \\
&\quad + \int_{\mathbb{R}} \theta^\epsilon(x - \zeta(t)) ((u_t^{R,\epsilon})^2 + \dot{\zeta}(t) u_t^{R,\epsilon}(t, x) u_x^{R,\epsilon}(t, x)) dx \\
&\quad - \int_{\mathbb{R}} \theta^\epsilon(x - \zeta(0)) ((u_1^\epsilon)^2 + \dot{\zeta}(0) u_1^\epsilon u_0^{\epsilon'}) dx \\
&\geq - \frac{1}{2} \int_{\mathbb{R}} (\theta^\epsilon(x - \zeta(t))(u_t^{R,\epsilon}(t, x))^2 - \theta^\epsilon(x - \zeta(0))(u_1^\epsilon)^2) dx \\
&\quad - \int_{Q_t} \theta^\epsilon(x - \zeta(s))(u_t^{R,\epsilon})^2 dt dx \\
&\quad - C \int_{Q_t} \theta^\epsilon(x - \zeta(s)) (\|u^{R,\epsilon}(s)\|_{L^2(\mathbb{R})}^2 + \|u_{xx}^{R,\epsilon}(s)\|_{L^2(\mathbb{R})}^2) ds dx + \int_{\mathbb{R}} \theta^\epsilon(x - \zeta(t))(u_t^{R,\epsilon}(t, x))^2 dx \\
&\quad - C_\eta \int_{\mathbb{R}} \theta^\epsilon(x - \zeta(t))(u_x^{R,\epsilon}(t, x))^2 dx - \eta \int_{\mathbb{R}} \theta^\epsilon(x - \zeta(t))(u_t^{R,\epsilon}(t, x))^2 dx \\
&\quad - \int_{\mathbb{R}} \theta^\epsilon(x - \zeta(0))(u_1^\epsilon)^2 dx - \int_{\mathbb{R}} \theta_\epsilon(x - \zeta(0)) \dot{\zeta}(0) u_1^\epsilon u_0^{\epsilon'} dx \\
&\geq \left(\frac{1}{2} - \eta\right) \int_{\mathbb{R}} \theta^\epsilon(x - \zeta(t))(u_t^{R,\epsilon}(t, x))^2 dx - \int_{Q_t} \theta^\epsilon(x - \zeta(s))(u_t^{R,\epsilon})^2 ds dx \\
&\quad - C \int_0^t \|u^{R,\epsilon}(s)\|_{H^2(\mathbb{R})}^2 ds - C_\eta \|u^{R,\epsilon}(t)\|_{L^2(\mathbb{R})}^2 \\
&\quad - \eta \|u^{R,\epsilon}(t)\|_{H^2(\mathbb{R})}^2 - C (\|u_1\|_{L^\infty(\mathbb{R})}^2 + \|u_0\|_{H^2(\mathbb{R})}^2)
\end{aligned}$$

The term $u_x^{R,\epsilon}$ provides

$$\begin{aligned}
A_3(t) &\stackrel{\text{def}}{=} \int_{Q_t} u_{xxxx}^{R,\epsilon} u_\tau^{R,\epsilon} ds dx \\
&\geq \frac{1}{2} \|u_{xx}^{R,\epsilon}(t)\|_{L^2(\mathbb{R})}^2 - \frac{1}{2} \|u_0''\|_{L^2}^2 + \frac{1}{2} \int_0^t \dot{\zeta}(s) \int_{\mathbb{R}} \partial_x (u_{xx}^{R,\epsilon})^2 ds dx \\
&= \frac{1}{2} (\|u_{xx}^{R,\epsilon}(t)\|_{L^2(\mathbb{R})}^2 - \|u_0\|_{H^2(\mathbb{R})}^2).
\end{aligned}$$

In the sequel, we define ψ^R and μ^R by $\psi^{R'} \stackrel{\text{def}}{=} \varphi^R$ and $\mu^{R'} \stackrel{\text{def}}{=} \psi^R$ such that $\mu^R(0) = \psi^R(0) = 0$. Since φ^R is even, so is μ^R . Furthermore, notice that $\mu^{R''} = \varphi^R \geq 0$. Consequently, μ^R is convex, even with $\mu^R(0) = 0$. It follows that $\mu^R \geq 0$. Now, the term $-(\varphi^R(u_x^{R,\epsilon} \star \theta^\epsilon) - \nu p)(u_{xx}^{R,\epsilon} \star \theta^\epsilon)$ provides

$$\begin{aligned}
A_4(t) &\stackrel{\text{def}}{=} - \int_{Q_t} (\varphi^R(u_x^{R,\epsilon} \star \theta^\epsilon) - \nu p)(u_{xx}^{R,\epsilon} \star \theta^\epsilon) \partial_\tau (u^{R,\epsilon} \star \theta^\epsilon) ds dx \\
&= \int_{Q_t} (\psi_R(u_x^{R,\epsilon} \star \theta^\epsilon) \partial_\tau (u_x^{R,\epsilon} \star \theta^\epsilon) ds dx + \int_0^t [\psi^R(u_x^{R,\epsilon} \star \theta^\epsilon) \partial_\tau (u^{R,\epsilon} \star \theta^\epsilon)]_{-\infty}^{+\infty} ds \\
&\quad + \int_{Q_t} \nu p(u_{xx}^{R,\epsilon} \star \theta^\epsilon) \frac{\partial}{\partial \tau} (u^{R,\epsilon} \star \theta^\epsilon) ds dx.
\end{aligned}$$

Since $|\psi^R(u^{R,\epsilon} \star \theta^{\epsilon'})| \leq C|u^{R,\epsilon} \star \theta^{\epsilon'}| \leq C_{R,\epsilon}$ and $\lim_{|x| \rightarrow +\infty} |\frac{\partial}{\partial \tau}(u^{R,\epsilon} \star \theta^\epsilon)| = 0$, we may deduce that $\int_0^t [\psi^R(u_x^{R,\epsilon} \star \theta^\epsilon) \partial_\tau(u_x^{R,\epsilon} \star \theta^\epsilon)]_{-\infty}^{+\infty} dt$ vanishes. Hence, we have

$$\begin{aligned} A_4(t) &\geq \int_{Q_t} \frac{\partial}{\partial \tau} (\mu^R(u_x^{R,\epsilon} \star \theta^\epsilon)) ds dx - C \int_0^t (\|u^{R,\epsilon}(s)\|_{\mathbb{H}^2(\mathbb{R})}^2 + \|u_t^{R,\epsilon}(s)\|_{\mathbb{L}^2(\mathbb{R})}^2) ds \\ &= \int_{\mathbb{R}} (\mu^R(u_x^{R,\epsilon}(t) \star \theta^\epsilon(t)) - \mu^R(u_x^{R,\epsilon}(0, x) \star \theta^\epsilon(0))) dx \\ &\quad - C \int_0^t (\|u^{R,\epsilon}(s)\|_{\mathbb{H}^2(\mathbb{R})}^2 + \|u_t^{R,\epsilon}(s)\|_{\mathbb{L}^2(\mathbb{R})}^2) ds. \end{aligned}$$

Since $\mu^R(u_x^{R,\epsilon}(t) \star \theta_\epsilon(t)) \geq 0$, it comes that

$$\begin{aligned} A_4(t) &\geq -C \int_{\mathbb{R}} |u_0^{\epsilon'}|^4 dx - C \int_0^t (\|u^{R,\epsilon}(s)\|_{\mathbb{H}^2(\mathbb{R})}^2 + \|u_t^{R,\epsilon}(s)\|_{\mathbb{L}^2(\mathbb{R})}^2) ds \\ &\geq -C \|u_0\|_{\mathbb{H}^2(\mathbb{R})}^4 - C \int_0^t (\|u^{R,\epsilon}(s)\|_{\mathbb{H}^2(\mathbb{R})}^2 + \|u_t^{R,\epsilon}(s)\|_{\mathbb{L}^2(\mathbb{R})}^2) ds. \end{aligned}$$

The term h_ϵ provides

$$\begin{aligned} A_5(t) &\stackrel{\text{def}}{=} - \int_{Q_t} f u_\tau^{R,\epsilon} ds dx - \int_{Q_t} P(s) \theta_\epsilon(x - \zeta(t)) u_\tau^{R,\epsilon} ds dx \\ &\geq -C \int_{Q_t} f^2 ds dx - C \int_{Q_t} |u_t^{R,\epsilon}|^2 ds dx - C \int_0^t \|u^{R,\epsilon}(s)\|_{\mathbb{H}^2(\mathbb{R})}^2 ds \\ &\quad + \int_{Q_t} P(s) \partial_\tau(\theta^\epsilon(x - \zeta(s))) u^{R,\epsilon}(s, x) ds dx - \int_{\mathbb{R}} [P(s) \theta^\epsilon(x - \zeta(s)) u^{R,\epsilon}(s, x)]_0^t dx \\ &\geq -C \int_{Q_t} f^2 - C \int_{Q_t} |u_t^{R,\epsilon}|^2 ds dx - C \int_0^t \|u^{R,\epsilon}(s)\|_{\mathbb{H}^2(\mathbb{R})}^2 ds - C \|u_0\|_{\mathbb{H}^2(\mathbb{R})}^2 \\ &\quad - C_\eta \|u^{R,\epsilon}(t)\|_{\mathbb{L}^2(\mathbb{R})}^2 - \eta \|u^{R,\epsilon}(t)\|_{\mathbb{H}^2(\mathbb{R})}^2. \end{aligned}$$

Since $\sum_{k=1}^5 A_k = 0$, the previous estimates lead to

$$\begin{aligned} &\|u_t^{R,\epsilon}(t)\|_{\mathbb{L}^2(\mathbb{R})}^2 + \|u_{xx}^{R,\epsilon}(t)\|_{\mathbb{L}^2(\mathbb{R})}^2 + \int_{\mathbb{R}} \theta^\epsilon(x - \zeta(s)) (u_t^{R,\epsilon}(t, x))^2 dx \\ &\leq C_T \left(\|u_0\|_{\mathbb{H}^2(\mathbb{R})}^2 + \|u_0\|_{\mathbb{H}^2(\mathbb{R})}^4 + \|u_1\|_{\mathbb{L}^2(\mathbb{R}) \cap \mathbb{L}^\infty(\mathbb{R})}^2 + \int_0^t \|u^{R,\epsilon}(s)\|_{\mathbb{H}^2(\mathbb{R})}^2 ds \right) \\ &\quad + \int_0^t \|u_t^{R,\epsilon}(s)\|_{\mathbb{L}^2(\mathbb{R})}^2 ds + \int_{Q_t} \theta^\epsilon(x - \zeta(s)) (u_t^{R,\epsilon}(s, x))^2 ds dx + 1) \\ &\quad + 2\eta \|u^{R,\epsilon}(t)\|_{\mathbb{H}^2(\mathbb{R})}^2 + C_\eta \|u^{R,\epsilon}(t)\|_{\mathbb{L}^2(\mathbb{R})}^2. \end{aligned} \tag{3.8}$$

In order to handle the terms $\eta \|u^{R,\epsilon}(t)\|_{\mathbb{H}^2(\mathbb{R})}^2$ and $C_\eta \|u^{R,\epsilon}(t)\|_{\mathbb{L}^2(\mathbb{R})}^2$, we write

$$u^{R,\epsilon}(t, x) = u_0(x) + \int_0^t u_t^{R,\epsilon}(s, x) ds,$$

which implies that

$$\|u^{R,\epsilon}(t)\|_{\mathbb{L}^2(\mathbb{R})} \leq C \left(\|u_0\|_{\mathbb{L}^2(\mathbb{R})} + \int_0^t \|u_t^{R,\epsilon}(s)\|_{\mathbb{L}^2(\mathbb{R})} ds \right), \tag{3.9}$$

and

$$\|u^{R,\epsilon}(t)\|_{L^2(\mathbb{R})}^2 \leq C_T \left(\|u_0\|_{L^2(\mathbb{R})}^2 + \int_0^t \|u_t^{R,\epsilon}(s)\|_{L^2(\mathbb{R})}^2 ds \right). \quad (3.10)$$

It follows from (3.10) that

$$\|u(t)\|_{H^2(\mathbb{R})}^2 \leq C_T \left(\|u_0\|_{L^2(\mathbb{R})}^2 + \int_0^t \|u_t^{R,\epsilon}(s)\|_{L^2(\mathbb{R})}^2 ds + \|u_{xx}^{R,\epsilon}(t)\|_{L^2(\mathbb{R})}^2 \right). \quad (3.11)$$

As a consequence of (3.10) and (3.11), (3.8) remains valid with $\|u^{R,\epsilon}(t)\|_{H^2(\mathbb{R})}^2$ in place of $\|u_{xx}^{R,\epsilon}(t)\|_{L^2(\mathbb{R})}^2$ in the left hand side, and $\eta = 0$ and $C_\eta = 0$ in the right hand side. Finally, the Grönwall's lemma gives the desired estimate. \square

Now, from Proposition 3.4 and the Sobolev embeddings, there exists a constant $C > 0$, independent of R and ϵ , such that

$$\|u^{R,\epsilon}\|_{L^\infty([0,T] \times \mathbb{R})} \leq C.$$

Hence, $\|u^{R,\epsilon} \star \theta^{\epsilon'}\|_{L^\infty([0,T] \times \mathbb{R})} \leq \frac{C}{\epsilon}$. It follows that for $R = \frac{C}{\epsilon}$, we have $\varphi_R(u^{R,\epsilon} \star \theta^{\epsilon'}) = (u^{R,\epsilon} \star \theta^{\epsilon'})^2$. Consequently, we replace $\varphi^R(u^{R,\epsilon} \star \theta^{\epsilon'})$ by $(u^{R,\epsilon} \star \theta^{\epsilon'})^2$ in the identity (3.7a) and drop the index R in (3.7a). From (3.7), we induce the following weak formulation:

$$\left\{ \begin{aligned} & \int_0^T \int_{\mathbb{R}} u_t^\epsilon v_t dx dt + \int_0^T \int_{\Omega} u_{xx}^\epsilon v_{xx} dx dt \\ & + \frac{1}{3} \int_0^T \int_{\mathbb{R}} ((u_x^\epsilon \star \theta^\epsilon)^3 (v_x \star \theta^\epsilon) + \nu p(u_x^\epsilon \star \theta^\epsilon)(v_x \star \theta^\epsilon) + (\nu p)_x(u_x^\epsilon \star \theta^\epsilon)) dx dt \\ & + \int_0^T \int_{\mathbb{R}} \theta^\epsilon(x - \xi(t)) u_t^\epsilon v_t dx dt + \int_0^T \int_{\mathbb{R}} f v dx dt + \int_0^T P(t) v(t, \xi(t)) dt \\ & + \int_{\mathbb{R}} u_1^\epsilon(x) v(0, x) dx + \int_{\mathbb{R}} \theta^\epsilon(x - \xi(0)) u_1^\epsilon(x) v(0, x) dx = 0, \end{aligned} \right. \quad (3.12)$$

for any $v \in C^2([0, T]; H^2(\Omega))$ with compact support in $[0, T] \times \mathbb{R}$.

4 Proof of Theorem 1.1

We derive from Proposition 3.4 the following convergences.

Corollary 4.1 *Assume that $u_0 \in H^2(\mathbb{R})$, $u_1 \in L^2(\mathbb{R}) \cap L^\infty(\mathbb{R})$. Then, there exists a subsequence $\{u_{\epsilon_k}\}_{k \in \mathbb{N}}$ of u_ϵ such that*

$$(i) \quad u^{\epsilon_k}, u_t^{\epsilon_k}, u_{xx}^{\epsilon_k} \rightharpoonup u, u_t, u_{xx} \text{ weakly in } L^2(0, T; L^2(\mathbb{R})).$$

(ii) *For any $\alpha \in]0, 2]$ and any compact set $K \subset \mathbb{R}$, we have*

$$u^{\epsilon_k} \xrightarrow[k \rightarrow +\infty]{} u \text{ strongly in } C^0([0, T]; H^{2-\alpha}(K)), \quad (4.1a)$$

$$u_x^{\epsilon_k} \xrightarrow[k \rightarrow +\infty]{} u_x \text{ strongly in } L^\infty([0, T] \times K). \quad (4.1b)$$

(iii) *For any $(a, b) \in \mathbb{R}^2$ such that $a < b$, $T > 0$, we have*

$$u_x^{\epsilon_k} \star \theta^{\epsilon_k} \xrightarrow[k \rightarrow +\infty]{} u_x \text{ strongly in } L^2([0, T] \times [a, b]), \quad (4.2a)$$

$$(u_x^{\epsilon_k} \star \theta^{\epsilon_k})^3 \xrightarrow[k \rightarrow +\infty]{} u_x^3 \text{ strongly in } L^2([0, T] \times [a, b]). \quad (4.2b)$$

Proof. (i) follows from Proposition 3.4. Note that (4.1a) comes from the Sobolev embeddings and Aubin-Lions lemma while (4.1b) follows from (4.1a) and $H^{1-\alpha}(K) \hookrightarrow L^\infty(K)$ for $\alpha \in]0, \frac{1}{2}[$. It remains to prove (iii). Let $t \in [0, T]$. We have

$$\begin{aligned} \|u_x^{\epsilon_k}(t) \star \theta^{\epsilon_k} - u_x(t)\|_{L^2([a,b])} &\leq \|u_x^{\epsilon_k}(t) \star \theta^{\epsilon_k} - u_x(t) \star \theta^{\epsilon_k}\|_{L^2([a,b])} + \|u_x(t) \star \theta^{\epsilon_k} - u_x(t)\|_{L^2([a,b])} \\ &\leq \|u_x^{\epsilon_k}(t) - u_x(t)\|_{L^2([a,b])} + \|u_x(t) \star \theta^{\epsilon_k} - u_x(t)\|_{L^2([a,b])}. \end{aligned}$$

Since

$$u_x^{\epsilon_k} \xrightarrow[k \rightarrow +\infty]{} u_x \text{ strongly in } L^\infty([0, T] \times [a, b]), \quad (4.3)$$

we have $\|u_x^{\epsilon_k}(t) - u_x(t)\|_{L^2([a,b])} \xrightarrow[k \rightarrow +\infty]{} 0$. On the other hand, $\|u_x(t) \star \theta^{\epsilon_k} - u_x(t)\|_{L^2([a,b])} \xrightarrow[k \rightarrow +\infty]{} 0$.

Hence, we get $\|u_x^{\epsilon_k}(t) \star \theta^{\epsilon_k} - u_x(t)\|_{L^2([a,b])} \xrightarrow[k \rightarrow +\infty]{} 0$. Besides, observe that

$$\|u_x^{\epsilon_k}(t) \star \theta^{\epsilon_k} - u_x(t)\|_{L^2([a,b])} \leq \|u_x^{\epsilon_k}(t)\|_{L^2([a,b])} + \|u_x(t)\|_{L^2([a,b])} \stackrel{(4.3)}{\leq} C.$$

Finally, we deduce from the dominated convergence theorem that (4.2a) holds. It remains to prove (4.2b). To this aim, we may assume that $\{u_x^{\epsilon_k}\}_{k \in \mathbb{N}}$ is bounded on $L^2([0, T] \times [a-1, b+1])$. Hence, we have

$$\begin{aligned} &\|(u_x^{\epsilon_k} \star \theta^{\epsilon_k})^3 - u_x^3\|_{L^2([a,b])} \\ &\leq \|(u_x^{\epsilon_k} \star \theta^{\epsilon_k} - u_x)((u_x^{\epsilon_k} \star \theta^{\epsilon_k})^2 + (u_x^{\epsilon_k} \star \theta^{\epsilon_k})u_x + u_x^2)\|_{L^2([a,b])} \\ &\stackrel{(4.3)}{\leq} C \|u_x^{\epsilon_k} \star \theta^{\epsilon_k} - u_x\|_{L^2([a,b])}. \end{aligned}$$

The result follows from (4.2a). \square

Proposition 4.2 *Assume that $u_0 \in H^2(\mathbb{R})$ and $u_1 \in H^1(\mathbb{R})$. Let $v \in C^2([0, T]; H^2(\mathbb{R}))$ with a compact support $[0, T] \times \mathbb{R}$. Then, we have*

$$\int_{Q_T} u_t^{\epsilon_k} v_t dt dx \xrightarrow[k \rightarrow +\infty]{} \int_{Q_T} u_t v_t dt dx, \quad (4.4a)$$

$$\int_{Q_T} u_{xx}^{\epsilon_k} v_{xx} dt dx \xrightarrow[k \rightarrow +\infty]{} \int_{Q_T} u_{xx} v_{xx} dt dx, \quad (4.4b)$$

$$\int_{Q_T} \left(-\frac{1}{3} (u_x^{\epsilon_k} \star \theta^{\epsilon_k})^3 (v_x \star \theta^{\epsilon_k}) + \nu p (u_x^{\epsilon_k} \star \theta^{\epsilon_k}) (v_x \star \theta^{\epsilon_k}) \right. \quad (4.4c)$$

$$\left. + (\nu p)_x (u_x^{\epsilon_k} \star \theta^{\epsilon_k}) (v \star \theta^{\epsilon_k}) \right) dt dx \xrightarrow[k \rightarrow +\infty]{} \int_{Q_T} \left(-\frac{1}{3} u_x^3 v_x - \nu p u_{xx} v \right) dt dx,$$

$$\int_{Q_T} (f + \theta^\epsilon(x - \zeta(t))P(t))v dt dx \xrightarrow[\epsilon \rightarrow 0]{} \int_{Q_T} f v dx dt + \int_0^T P(t)v(t, \zeta(t)) dt, \quad (4.4d)$$

$$\int_{\mathbb{R}} u_1^{\epsilon_k}(x)v(0, x) dx + \int_{\mathbb{R}} \theta^{\epsilon_k}(x - \zeta(0))u_1^{\epsilon_k}(x)v(0, x) dx \quad (4.4e)$$

$$\xrightarrow[k \rightarrow +\infty]{} \int_{\mathbb{R}} u_1(x)v(0, x) dx + u_1(\zeta(0))v(0, \zeta(0)),$$

$$\int_{Q_T} \theta^{\epsilon_k}(x - \zeta(t))u_t^{\epsilon_k} v_t dt dx \xrightarrow[k \rightarrow +\infty]{} -u_0(\zeta(0))v_t(0, \zeta(0)) \quad (4.4f)$$

$$- \int_0^T u(t, \zeta(t))v_{tt}(t, \zeta(t)) dt dx - \int_0^T \dot{\zeta}(t)u_x(t, \zeta(t))v_t(t, \zeta(t)) dt$$

$$- \int_0^T \dot{\zeta}(t)u(t, \zeta(t))v_{xt}(t, \zeta(t)) dt.$$

Proof. We observe that (4.4a) and (4.4b) follow from Corollary 4.1 (i). Notice that all the functions v_{xx} , $v \star \theta^{\epsilon_k}$ and $v_x \star \theta^{\epsilon_k}$ have their supports included in a fixed bounded product $[0, T] \times [a, b]$. Hence, we may deduce from (4.2) that

$$\begin{aligned} & \int_{Q_T} \left(-\frac{1}{3}(u_x^{\epsilon_k} \star \theta^{\epsilon_k})^3 (v_x \star \theta^{\epsilon_k}) + \nu p (u_x^{\epsilon_k} \star \theta^{\epsilon_k})(v_x \star \theta^{\epsilon_k}) + (\nu p)_x (u_x^{\epsilon_k} \star \theta^{\epsilon_k})(v \star \theta^{\epsilon_k}) \right) dt dx \\ & \xrightarrow{k \rightarrow +\infty} \int_{Q_T} \left(-\frac{1}{3}u_x^3 v_x + \nu p u_x v_x + (\nu p)_x u_x v \right) dt dx = \int_{Q_T} \left(-\frac{1}{3}u_x^3 v_x - \nu p u_{xx} v \right) dt dx, \end{aligned}$$

which proves (4.4c). Property (4.4d) is straightforward. Next, $u_1 \in H^1(\mathbb{R}) \subset (L^2(\mathbb{R}) \cap L^\infty(\mathbb{R}))$ and $u_1^\epsilon = u_1 \star \theta^\epsilon$. Hence, we have

$$u_1^\epsilon \xrightarrow{\epsilon \rightarrow 0} u_1 \text{ strongly in } L^2(\mathbb{R}) \cap L^\infty(\mathbb{R}),$$

and (4.4e) follows. In order to prove (4.4f), we first establish the following result:

Lemma 4.3 *Let $u_0 \in H^2(\mathbb{R})$ and $u_1 \in H^1(\mathbb{R})$. Then, for any $v \in C^2([0, T]; H^2(\mathbb{R}))$ with a compact support in $[0, T[\times \mathbb{R}$, we have*

$$\begin{aligned} & \int_{Q_T} \theta^{\epsilon_k}(x - \zeta(t)) u_t^{\epsilon_k} v_t dt dx = - \int_{\mathbb{R}} \theta^{\epsilon_k}(x - \zeta(0)) u_0^{\epsilon_k} v_t(0, x) dx \\ & - \int_{Q_T} \theta^{\epsilon_k}(x - \zeta(t)) u^{\epsilon_k} v_{tt} dt dx - \int_{Q_T} \dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) u_x^{\epsilon_k} v_t dx dt \\ & - \int_{Q_T} \dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) u^{\epsilon_k} v_{xt} dx dt. \end{aligned}$$

Proof. Observe that

$$\begin{aligned} & \int_{Q_T} \theta^{\epsilon_k}(x - \zeta(t)) u_t^{\epsilon_k} v_t dt dx \\ & = - \int_{Q_T} (-\dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) v_t + \theta^{\epsilon_k}(x - \zeta(t)) v_{tt}) u^{\epsilon_k} dt dx \\ & + \int_{\mathbb{R}} [\theta^{\epsilon_k}(x - \zeta(t)) u^{\epsilon_k} v_t]_0^T dx \\ & = - \int_{\mathbb{R}} \theta^{\epsilon_k}(x - \zeta(0)) u_0^{\epsilon_k} v_t(0, x) dx - \int_{Q_T} \theta^{\epsilon_k}(x - \zeta(t)) u^{\epsilon_k} v_{tt} dt dx \\ & - \int_{Q_T} \dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) \partial_x (u^{\epsilon_k} v_t) dt dx, \end{aligned}$$

and the desired result follows. \square

We now pass to the limit in the various terms appearing in the equality of Lemma 4.3.

End of the proof of (4.4f). We only treat the case of $\int_{Q_T} \dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) u_x^{\epsilon_k} v_t dt dx$, which we replace by $\int_0^T \int_a^b \dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) u_x^{\epsilon_k} v_t dt dx$. The other terms in Lemma 4.3 can be handled similarly. Notice that for $\alpha \in]0, \frac{1}{2}[$, $H^{1-\alpha}([a, b])$ is a multiplicative algebra. Since $u_{\epsilon_k, x} \in C^0([0, T]; H^{1-\alpha}([a, b]))$ and $v_t \in C^1([0, T]; H^1([a, b])) \hookrightarrow C^0([0, T]; H^{1-\alpha}([a, b]))$, it follows that $u_x^{\epsilon_k} v_t \in C^0([0, T]; H^{1-\alpha}([a, b]))$ for any $\alpha \in]0, \frac{1}{2}[$. Notice also, since $\{u_x^{\epsilon_k}\}_{n \in \mathbb{N}^*}$ is bounded in $L^\infty([0, T] \times [a, b])$ and $v_t \in C^1([0, T]; H^1([a, b])) \hookrightarrow L^\infty([0, T] \times [a, b])$, that $\{u_x^{\epsilon_k} v_t\}_{n \in \mathbb{N}^*}$ is

bounded in $L^\infty([0, T] \times [a, b])$. Now, we write

$$\begin{aligned} & \int_{Q_T} \dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) u_x^{\epsilon_k} v_t \, dt \, dx \\ &= \int_{Q_T} \dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) (u_x^{\epsilon_k} - u_x) v_t \, dt \, dx + \int_{Q_T} \dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) u_x v_t \, dt \, dx. \end{aligned} \quad (4.5)$$

Since $u_x v_t \in C^0([0, T]; H^{1-\alpha}([a, b]))$ and $H^{1-\alpha}([a, b]) \hookrightarrow C^0([a, b])$ with $\alpha \in]0, \frac{1}{2}[$, we have

$$\int_{\mathbb{R}} \theta^{\epsilon_k}(x - \zeta(t)) u_x(t, x) v_t(t, x) \, dx \xrightarrow{k \rightarrow +\infty} u_x(t, \zeta(t)) v_t(t, \zeta(t))$$

for any $t \in [0, T]$ and since

$$\left| \dot{\zeta}(t) \int_a^b \theta^{\epsilon_k}(x - \zeta(t)) u_x(t, x) v_t(t, x) \, dx \right| \leq \|\dot{\zeta}\|_{L^\infty(0, T)} \|u_x v_t\|_{L^\infty([0, T] \times [a, b])},$$

it follows from the dominated convergence theorem that

$$\int_{Q_T} \dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) u_x v_t \, dt \, dx \xrightarrow{k \rightarrow +\infty} \int_0^T \dot{\zeta}(t) u_x(t, \zeta(t)) v_t(t, \zeta(t)) \, dt.$$

For the other term in (4.5), we have

$$\begin{aligned} & \left| \int_{Q_T} \dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) (u_x^{\epsilon_k}(t, x) - u_x(t, x)) v_t(t, x) \, dt \, dx \right| \\ & \leq C_T \|\dot{\zeta}\|_{L^\infty(0, T)} \|u_x^{\epsilon_k} - u_x\|_{L^\infty([0, T] \times [a, b])} \xrightarrow{k \rightarrow +\infty} 0. \end{aligned}$$

Finally, we have

$$\int_{Q_T} \dot{\zeta}(t) \theta^{\epsilon_k}(x - \zeta(t)) u_x^{\epsilon_k}(t, x) v_t(t, x) \, dt \, dx \xrightarrow{k \rightarrow +\infty} \int_0^T \dot{\zeta}(t) u_x(t, \zeta(t)) v_t(t, \zeta(t)) \, dt.$$

□

Remark 4.4 *The right hand side of (4.4f) is formally equal to*

$$\begin{aligned} & -u_0(\zeta(0)) v_t(0, \zeta(0)) - \int_0^T u(t, \zeta(t)) v_{tt}(t, \zeta(t)) \, dt \, dx \\ & - \int_0^T \dot{\zeta}(t) u_x(t, \zeta(t)) v_t(t, \zeta(t)) \, dt - \int_0^T \dot{\zeta}(t) u(t, \zeta(t)) v_{xt}(t, \zeta(t)) \, dt \\ & = -u_0(\zeta(0)) v_t(0, \zeta(0)) - \int_0^T \frac{d}{dt} (u(t, \zeta(t)) v_t(t, \zeta(t))) \, dt \\ & + \int_0^T u_t(t, \zeta(t)) v_t(t, \zeta(t)) \, dt = \int_0^T u_t(t, \zeta(t)) v_t(t, \zeta(t)) \, dt. \end{aligned} \quad (4.6)$$

Nevertheless, with regularity of Corollary 4.1, the trace $\int_0^T u_t(t, \zeta(t)) v_t(t, \zeta(t)) \, dt$ is not well defined.

End of the proof of Theorem 1.1. The proof of (1.3) follows from (3.12) and Proposition 4.2. Details are omitted. It remains to show that $u(0, \cdot) = u_0$. To this aim, notice that by (4.1a), $u^{\epsilon_k}(0, \cdot) \xrightarrow{k \rightarrow +\infty} u(0, \cdot)$ in $L^2_{\text{loc}}(\mathbb{R})$, and that $u^{\epsilon_k}(0, \cdot) = u_0 \star \theta^{\epsilon_k} \xrightarrow{k \rightarrow +\infty} u_0$ in $L^2(\mathbb{R})$. This proves the result. □

Acknowledgment

O.A. was funded by the Czech Academy of Sciences within Grant RVO: 67985840 and by the "EU Strategy for the Danube Region" Grant: 8X23001. A.H. and A.P. were supported by the Ministry for Europe and Foreign Affairs (France) within Programme Danube 2023-2025 (Project PHC Danube n° 49916UM).

References

- [1] J. Ahn, D. Stewart. An Euler-Bernoulli beam with dynamic contact: discretization, convergence, and numerical results. *SIAM J. Numer. Anal.*, 43 (2005), no.4, 1455–1480.
- [2] J. Ahn, K. L. Kuttler, M. Shillor. Dynamic contact of two Gao beams. *Electron J. Differential Equations*, 194 (2012) 1–42.
- [3] K. T. Andrews, M. F. M'Bengue, M. Shillor. Vibrations of a nonlinear dynamic beam between two stops. *Discrete Contin. Dyn. Syst.*, 12 (1) (2009): 23–38.
- [4] K. T. Andrews, K. L. Kuttler, M. Shillor. Dynamic Gao Beam in Contact with a Reactive or Rigid Foundation. *Advances in variational and hemivariational inequalities*, 225–248, Adv. Mech. Math., 33 Springer, Cham, 2015
- [5] K. T. Andrews, J. Li, M. Shillor. Existence for dynamic contact of a stochastic viscoelastic Gao beam. *Nonlinear Anal. RWA*, 22 (4) (2015) 568–580.
- [6] L. Beirão da Veiga, C. Lovadina, A. Reali. Avoiding shear locking for the Timoshenko beam problem via isogeometric collocation methods. *Comput. Methods Appl. Mech. Engrg.*, 241/244 (2012), 38–51.
- [7] C. Castro, E. Zuazua. Exact boundary controllability of two Euler-Bernoulli beams connected by a point mass. *Math. Comput. Modelling*, 32 (2000), no.9, 955–969.
- [8] B. Dekoninck, S. Nicaise. Control of networks of Euler-Bernoulli beams. *ESAIM Control Optim. Calc. Var.*, 4 (1999), 57–81.
- [9] A. Heibig, P. Krejčí, A. Petrov. Solvability for a dynamical beam problem on a frictionally damped foundation under a moving load. *Discrete Contin. Dyn. Syst. Ser. B*, 28 (2023), n $^{\circ}$ 5, 3277–3293.
- [10] Z. Dimitrovová, A. Rodrigues. Critical velocity of a uniformly moving load. *Adv. Eng. Softw.*, 50 (2012) 44–56.
- [11] B. Dyniewicz, C. I. Bajer, K. L. Kuttler, M. Shillor. Vibrations of a Gao beam subjected to a moving mass. *Nonlinear Analysis: Real World Appl.*, 50 (2019): 342–364.
- [12] L. Frýba. *Vibration of Solids and Structures under Moving Loads*. Prague: Research Institute of Transport, 1972.
- [13] A. D. Y. Gao. Nonlinear elastic beam theory with application in contact problems and variational approaches. *Mech. Res. Commun.*, 23 (1) (1996): 11–17.
- [14] A. D. Y. Gao, J. Machalová, H. Netuka. Mixed finite element solutions to contact problems of nonlinear Gao beam on elastic foundation. *Nonlinear Anal. Real World Appl.*, 22 (2015), 537–550.

- [15] J. Machalová, H. Netuka. Solution of contact problems for nonlinear Gao beam and obstacle. *J. Appl. Math.*, (2015), Art. ID 420649, 12 pp.
- [16] J. Machalová, H. Netuka. Solution of contact problems for Gao beam and elastic foundation. *Math. Mech. Solids*, 23 (2018), no.3, 473–488.
- [17] J. U. Kim, Y. Renardy. Boundary control of the Timoshenko beam. *SIAM J. Control Optim.* 25 (1987), no.6, 1417–1429.
- [18] H. Takeda, S. Yoshikawa. On the initial value problem of the semilinear beam equation with weak damping I: Smoothing effect. *J. Math. Anal. Appl.*, 401(2013), no.1, 244–258.
- [19] H. Takeda, S. Yoshikawa. On the initial value problem of the semilinear beam equation with weak damping II: Asymptotic profiles. *J. Differential Equations*, 253 (2012), no.11, 3061–3080.