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Predictive 3D modelling of free oblique cutting introducing an ANN-based material flow law with experimental validation over a wide range of conditions

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Abstract

Modelling of the cutting process needs to move from 2D to 3D configurations to get closer to industrial applications. This study introduces a predictive 3D finite element model of free orthogonal and oblique cutting with an Artificial Neural Network (ANN)-based material flow law and experimental validation in strictly the same conditions (cutting and geometrical). The flow law based on a neural network allows simulating the cutting process based on data coming from the material characterization tests without requiring any postulate concerning the expression of the flow law. The developments are applied to the formation of continuous chips for the titanium alloy Ti6Al4V and an unseen broad range of 36 cutting conditions is considered: 2 cutting edge inclinations, 3 uncut chip thicknesses and 6 cutting speeds. The predictive performance of the model (i.e., the evaluation of the trends of fundamental variables with the absence of tuning of both numerical parameters and model features when cutting conditions are significantly modified) is high for the forces, mainly cutting and passive, and the chip thickness ratio on all 36 cutting conditions. The accuracy of the main cutting force is excellent: the average difference with the experiments is 4 %, within the experimental dispersion. No significant degradation of the results is brought by the apparition of the third, out-of-plane, force, which shows the ability of the model to handle orthogonal and oblique cutting configurations.

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Keywords:

Oblique cutting, Finite element method (FEM), Predictive model, Artificial Neural Network (ANN), Material flow law

1. Introduction

Selection of the tools and the cutting conditions in machining are still difficult to achieve because of the high level of complexity and the related nonlinear phenomena. Comprehension of the influence of the process parameters on the quality of a component and its optimization are also a challenge for the same reasons. In the frame of digital manufacturing and Industry 4.0, modelling the cutting process supports them, while remaining a challenging task. As highlighted by Arrazola et al. [1], most finite element (FE) models are developed in 2D (orthogonal cutting configuration usually) although industrial applications require 3D modelling.

The behaviour of the machined material is one of the key aspects of a FE model [1, 2]. Research is very intense in this area, leading to a growing number of constitutive material models ranging from empirical models to physical models, some including microstructure effects [2]. The empirical thermo-elasto-viscoplastic Johnson-Cook (JC) model [3] is still the most widely used to this day:

$$\sigma^y = \left(A + B \varepsilon^{p^n} \right) \left(1 + C \ln \frac{\dot{\varepsilon}^p}{\dot{\varepsilon}_0^p} \right) \left(1 - \left[\frac{T - T_{\text{room}}}{T_{\text{melt}} - T_{\text{room}}} \right]^m \right) \quad (1)$$

In this model, the flow stress, σ^y , is a function of the plastic strain, ε^p , the plastic strain rate, $\dot{\varepsilon}^p$, and the temperature, T . It is composed of 3 terms describing independently the plastic, viscous and thermal aspects. One of the points in favour of its adoption is the rather limited number of parameters to be identified, 5: A , B , C , m and n . Here, $\dot{\varepsilon}_0^p$ is the reference plastic strain rate, while T_{room} and T_{melt} are respectively the ambient (room) and melting temperatures. More recent models developed on this basis, such as that of Calamaz et al. [4], increase this number of parameters (for the particular Calamaz model to 9). Other authors have also used Zerilli-Armstrong model to simulate cutting processes [5]. The best description (in theory) of the behaviour is obtained at the cost of a greater complexity of the identification process and a reduction of the link with the physical meaning of the model.

One of the problems of modelling material behaviour for cutting simulation is the identification of parameters, especially as the experimental equipment does not allow the high levels of strain, strain rate and temperature of machining to be

31 achieved [2]. Inverse identification is an alternative, but the uniqueness of the so-
32 lution is not always guaranteed [1, 2]. Early work by Özel and Altan [6] used the
33 least squares method to identify the input parameters of a FE model in an inverse
34 manner. Shrot and Bäker [7] then used the Levenberg-Marquardt algorithm for
35 their identification of the material parameters. They showed that similar results
36 (cutting forces and chip morphology) could be obtained by different sets of pa-
37 rameters and thus highlighted the non-uniqueness of the solution of the inverse
38 problem. In addition to the flow stress parameters, Klocke et al. [8] also identi-
39 fied the damage parameters. In more recent work, such as Bosetti et al. [9] and
40 Denkena et al. [10], the approach to the inverse identification problem is shifting
41 from optimization to Artificial Intelligence (AI) based methods. The Downhill
42 Simplex Algorithm (DSA) is adopted by Bergs et al. [11] and by Hardt et al. [12]
43 for AISI 1045. Stampfer et al. [13] also chose DSA when treating AISI 4140
44 quenched at 3 different temperatures. In [14], Hardt et al. showed that Parti-
45 cle Swarm Optimization (PSO) was more efficient in solving the inverse problem
46 than DSA, even though the computational time is still significant. In order to re-
47 duce the computational time, an Efficient Global Optimization algorithm (EGO)
48 was recently introduced by Kugalur Palanisamy et al. [15]. They identified simul-
49 taneously the parameters of the material constitutive model and the friction model
50 for Ti6Al4V. The identified parameters showed good performance when applied
51 to a different FE model [16]. Most of these works highlight the non-uniqueness
52 of the identification and they all require the definition of the analytical expression
53 of the constitutive model.

54 ANN (Artificial Neural Network)-based material models have been introduced
55 to avoid postulating or knowing the analytical expression of the material be-
56 haviour. Gorji et al [17] recently reviewed the use of recurrent neural networks
57 for material models, while Jamli and Farid [18] reviewed their application in FE
58 simulation of material forming. When compared to classical analytical and em-
59 pirical models, such as JC model, they proved to be more powerful to represent
60 the experimental behaviour [19]. Use of these ANN-based models in FE simula-
61 tion of forming processes also turned out to provide better results than the classical
62 JC model [20] and to handle complex phenomena such as dynamic recrystallisa-
63 tion [21]. No application of these ANN-based models in FE simulation of cutting
64 currently exists.

65 Lagrangian and Eulerian formulations are the most used for FE modelling of
66 the cutting process. Combinations of formulations, such as Arbitrary Lagrangian-
67 Eulerian (ALE) and Coupled Eulerian-Lagrangian (CEL), are increasingly being
68 used to avoid (or reduce) mesh distortions [22]. The CEL formulation has recently

69 been successfully applied to the modelling of cutting (in 2D orthogonal configura-
70 tion): it provides accurate results with realistic chip shape and no mesh distortion.
71 The first 3D applications are found in recent works [23–27]. They cover orthog-
72 onal (free) cutting or a simple 3D operation, while free oblique cutting has yet to
73 be studied.

74 Experimental validation of a model is a crucial step in modelling the cutting
75 process. The experimental configuration should be as close as possible to the sim-
76 ulation. For the validation of orthogonal cutting, a rotational motion usually gen-
77 erates the cutting speed. This is often done in turning [28] or milling [23] and the
78 diameter of the rotating workpiece must be large enough to reduce the influence
79 of curvature on the results. Experimental configurations under strictly orthogonal
80 cutting conditions are less often adopted, for example on broaching machines [29]
81 or milling machines [30, 31]. If they remove the assumptions related to the rotary
82 cutting motion, they generally allow lower cutting speeds (except on a dedicated
83 machine, as in Afrasiabi et al. [32]). Free oblique cutting with a straight cutting
84 edge has not yet been studied: all efforts have been concentrated on orthogonal
85 cutting (mainly for validation of 2D FE models).

86 This paper fills the gap in the oblique cutting literature by investigating both
87 orthogonal and free oblique 3D cutting configurations, both experimentally and
88 numerically. An ANN, introduced in Pantalé et al. [33], is implemented in a FE
89 cutting model for the first time in place of the JC analytical law. A wide range of
90 cutting speeds (6), uncut chip thicknesses (3) and cutting edge inclination angles
91 (2) resulting in 36 different conditions are considered to demonstrate the predictive
92 capability of the FE model for the fundamental variables. The developments are
93 applied to the formation of continuous chips of the titanium alloy Ti6Al4V.

94 **2. Experimental setup**

95 A 3-axis GF Mikron VCE 600 Pro milling machine is used to perform dry or-
96 thogonal and oblique cutting tests on Ti6Al4V (grade 5 annealed at 750 °C for 1 h
97 followed by air cooling) with the same kinematics as a shaper. As shown in Fig-
98 ure 1, the tungsten carbide tool (modified LCGN160602-0600-FG, CP500 from
99 SECO) is fixed on a dedicated holder (modified CFHN-06 from SECO) and the
100 sample to be cut is clamped in the spindle (no rotation is allowed during the test).
101 The top of the sample has 3 ribs of 1 mm width (the width of the tool is 6 mm) and
102 10 mm length. The test consists of removing the top layer (its height is the uncut
103 chip thickness, h) of a rib at the prescribed cutting speed, v_c . The cutting speed is
104 provided by the feed rate, v_f , of the machine (maximum value of 40 m/min). The

105 tool cutting edge inclination, λ_s , results from the relative angular orientation of
 106 the tool and the sample. Table 1 shows the cutting conditions: 6 cutting speeds, 3
 107 uncut chip thicknesses and 2 inclination angles, each repeated 3 times. An incli-
 108 nation angle of 6° is the typical value when turning Ti6Al4V, while cutting speeds
 109 and uncut chip thicknesses values in accordance with recommended ranges by
 110 SECO for the standard tool [34] are adopted.

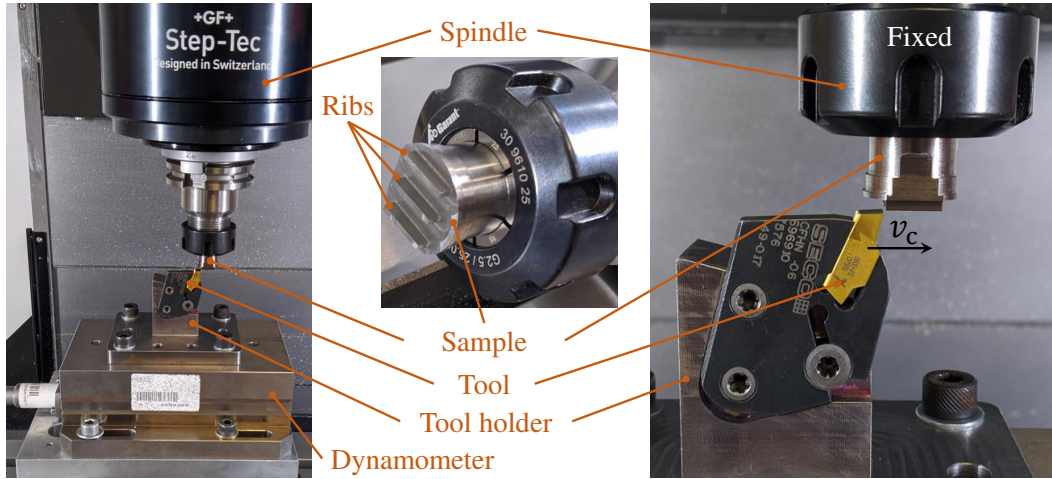


Figure 1: Experimental setup

Table 1: Cutting conditions of the study

Parameter	Values
Cutting speed, v_c (m/min)	5, 7.5, 10, 20, 30, 40
Uncut chip thickness, h (μm)	40, 60, 80
Cutting edge inclination, λ_s ($^\circ$)	0, 6
Width of the workpiece (mm)	1
Length of the workpiece (mm)	10
Width of the cutting edge (mm)	6 (1.1 in the model)
Cutting edge radius, r_β (μm)	20
Rake angle, γ_0 ($^\circ$)	15
Clearance angle, α_0 ($^\circ$)	2

111 Forces are measured with a 3-component Kistler 9257B dynamometer and

112 are amplified by a Kistler 5070A charge amplifier. Acquisition is performed at
113 3 kHz using a Kistler 5697A2 data acquisition system and DynoWare software.
114 The recorded forces are then filtered with a second-order low-pass Bessel filter at
115 750 Hz before calculating the average value of the steady state signal.

116 All chips are collected and observed with a Dino Lite digital microscope
117 AM7013MZT (5 MP, magnification 20× – 250×). Each chip is measured 3 times
118 along its length in order to obtain an average value representative of the whole
119 chip.

120 **3. Finite element model**

121 *3.1. Modelling choices*

122 The main objectives of a predictive model are the accurate modelling of trends
123 in results as conditions change and the good agreement of predicted values with
124 experimental values (exact values are not expected due to experimental disper-
125 sions of at least 10 % around the mean values). This type of model is intended to
126 support future choices and developments without the need for experimental data.
127 No assumptions are made about the geometry of the workpiece in the model (i.e.,
128 its width is the same as in experiments), while keeping the calculation time rel-
129 evant for industrial applications. The CEL formulation is adopted to model the
130 dry orthogonal and free oblique cutting tests with Abaqus/Explicit 2020. The
131 3D model is composed of a fixed Lagrangian tool and a Eulerian part (Figure 2).
132 Chip formation occurs by plastic flow through the Eulerian domain without mesh
133 distortion. The Eulerian formulation allows for chip formation without damage
134 properties, by removing modelling assumptions. These two features contribute to
135 the cutting models providing accurate results and realistic chips [22].

136 As shown in Figure 3, the full width of the workpiece (1 mm), i.e., one rib
137 in the experiments, is modelled. To allow for chip formation and lateral flow,
138 the Eulerian domain is wider (it includes the volume in which the material can
139 move). The volume above the initial part is also meshed with Eulerian elements
140 for the same reasons. As in the experiments and to satisfy the assumption of
141 an orthogonal and oblique free cut, the tool is wider than the workpiece (it is
142 1.1 mm in the model and 6 mm in the experiments). It is very important to note
143 that the models are the same for both inclination angles: they differ only in the
144 rotation of the tool by 6° around the Y axis as in the experiments (Figure 3). This,
145 together with the absence of assumptions when developing the models, contributes
146 to make the models predictive: no input is changed when the cutting conditions
147 are changed.

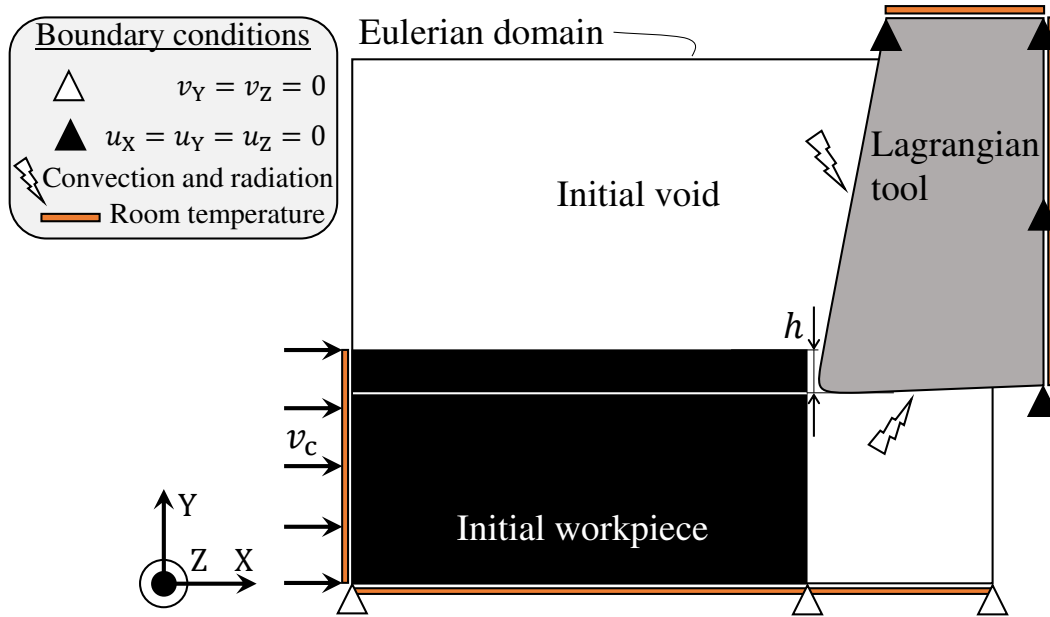


Figure 2: Boundary conditions and schematic initial geometry of the model

148 According to a previous sensitivity study of the mesh in orthogonal cutting
 149 with the CEL formulation [24], the edge size of the elements is $5\mu\text{m}$ in the
 150 plane parallel to the cutting speed. In the direction perpendicular to this plane,
 151 it is $5\mu\text{m}$ in the areas close to the lateral boundaries of the Eulerian domain
 152 and $50\mu\text{m}$ in the middle of the part. To reduce the computation time, the size
 153 of the model depends on the value of the uncut chip thickness. This results
 154 in a Eulerian domain (EC3D8RT 8-node 3D linear Eulerian elements, coupled
 155 mechanical-thermal behaviour and reduced integration) composed of 216 550
 156 273 350 nodes and a Lagrangian domain (C3D8T 8-node 3D linear Lagrangian
 157 elements, coupled mechanical-thermal behaviour) of 4650 nodes.

158 The Ti6Al4V part is assumed to be thermo-elasto-viscoplastic (isotropic) and
 159 the inelastic thermal fraction is 0.9. The JC parameters set of Seo et al. [35]
 160 is adopted because the value of A corresponds to the value of the typical yield
 161 strength of Ti6Al4V and this set was found to provide the best results among the
 162 20 sets available in the literature [36]. The TiN coated tungsten carbide (WC) tool
 163 is assumed to have linear elasticity. The material properties are given in Table 2.

164 According to the experimental results of Rech et al. [39], it is assumed that
 165 Coulomb friction occurs at the tool-piece interface and that the coefficients of

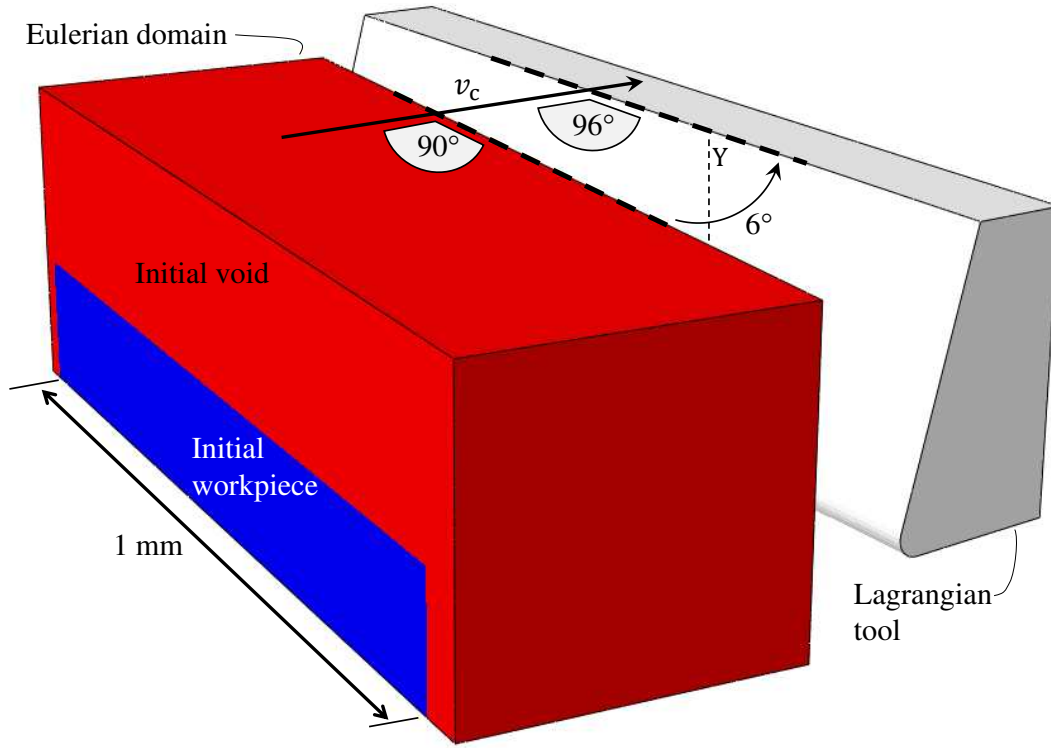


Figure 3: Configuration of the FE model for $\lambda_s = 6^\circ$

166 friction, μ , and heat partition, β , depend on the cutting speed. The limiting shear
 167 stress, τ_{\max} , is included and is given by:

$$\tau_{\max} = \frac{\text{yield stress}}{\sqrt{3}} = \frac{A}{\sqrt{3}} \quad (2)$$

168 All the friction energy is converted into heat. Table 3 shows the friction coeffi-
 169 cients adopted in this study. Gap heat conductance based on the distance between
 170 the two surfaces is not available in Abaqus/Explicit 2020. It is therefore not in-
 171 cluded in the modelling.

172 An ambient temperature of 293 K is imposed on the top and right surfaces
 173 of the tool and on the left and bottom surfaces of the workpiece (Figure 2). It
 174 is assumed that radiation and convection occur on the rake and clearance faces
 175 of the tool. The initial temperature of the tool and workpiece is set to the room
 176 temperature (293 K). The heat transfer coefficients are provided in Table 3.

Table 2: Materials properties [35, 37, 38]

Young’s modulus, E (GPa)	Ti6Al4V	113.8 [†]
	WC	650
Poisson’s ratio, ν	Ti6Al4V	0.34
	WC	0.2
Density, ρ (kg/m ³)	Ti6Al4V	4430
	WC	14 850
Conductivity, k (W/m K)	Ti6Al4V	6.3 [†]
	WC	100
Expansion, α (1/K)	Ti6Al4V	8.6E−6 [†]
	WC	5E−6
Specific heat, c_p (J/kg K)	Ti6Al4V	531 [†]
	WC	202
JC flow stress	A (MPa)	997.9
	B (MPa)	653.1
	C	0.0198
	m	0.7
	n	0.45
	$\dot{\epsilon}_0$ (1/s)	1
	T_{room} (K)	293
	T_{melt} (K)	1873

[†]: Dependence on the temperature, value provided at 293 K

177 3.2. Material model of Ti6Al4V

178 In the numerical simulations presented in Section 4, a thermo-elasto-viscoplastic
179 material model for Ti6Al4V is employed, which utilizes a flow criterion based on
180 an Artificial Neural Network (ANN) identified for the material. This ANN is im-
181 plemented in the Abaqus/Explicit code through a Fortran VUHARD subroutine,
182 as proposed by Pantalé et al. [20, 33], to compute the flow stress σ^y as a func-
183 tion of the plastic strain ϵ^p , the plastic strain rate, $\dot{\epsilon}^p$, and the temperature T . The
184 approach replaces the analytical formulation of the flow law, typically based on
185 Johnson-Cook or Zerilli-Armstrong type models, with a multi-layer ANN serving
186 as a universal approximator. This enables the direct identification of the neu-
187 ral network parameters from experimental data without postulating a behavioral
188 model, simplifying the procedure and providing greater flexibility in model defi-
189 nition.

Table 3: Friction and heat transfer coefficients [13, 37, 39]

Cutting speed, v_c (m/min)	μ	β
5	0.24	1
7.5	0.22	0.89
10	0.21	0.80
20	0.19	0.63
30	0.18	0.55
40	0.17	0.50
Limiting shear stress, τ_{\max} (MPa)	576	
Convection, U (W/m ² K)	50	
Radiation, ϵ	0.3	

190 In contrast to the classic approach, which involves conducting experiments on
 191 a material, postulating an analytical form for the flow law, and identifying the
 192 parameters that best fit the experimental data, the use of ANN allows for direct
 193 identification of the law from experimental data without the need to postulate the
 194 analytical form of the flow law. This method also enables the computation of the
 195 three derivatives of the flow stress σ^y with respect to the three input variables of
 196 the model, which is necessary for implementing the model as a flow law in the
 197 form of a VUHARD subroutine in the FEM code Abaqus/Explicit. The same
 198 network architecture and identified trained parameters are used to compute the
 199 flow stress σ^y and the derivatives in a one-step procedure [20, 33].

200 In order to verify the influence of the neural network complexity on the nu-
 201 merical results of the simulation and on the computation time, several ANN ar-
 202 chitectures (i.e. hyperparameters of the ANN) are tested afterwards (in 3.4). The
 203 chosen global architecture has 2 hidden layers with a variable number of neurons
 204 for the first hidden layer ($\zeta = 9$ to 17) and 7 neurons for the second hidden layer, 3
 205 inputs (the plastic strain, ε^p , the plastic strain rate, $\dot{\varepsilon}^p$, and the temperature, T) and
 206 one output (the yield strength, σ^y). The global architecture of this type of ANN is
 207 given in Figure 4 for 9 neurons in the first hidden layer. According to Pantalé et
 208 al. [33], this ANN is referred to as ANN 3-9-7-1-sig, as it has 3 inputs, 9 neurons
 209 in the first hidden layer, 7 neurons in the second hidden layer, 1 output and a sig-
 210 moid activation function. The selection of an architecture with two hidden layers
 211 was made based on the conclusions drawn in Pantalé et al. [33]. Additionally,

212 the decision to use the sigmoid activation function was guided by the findings in
 213 Pantalé [40], who identified the most efficient and accurate activation functions
 214 for finite element simulations in thermomechanical forming.

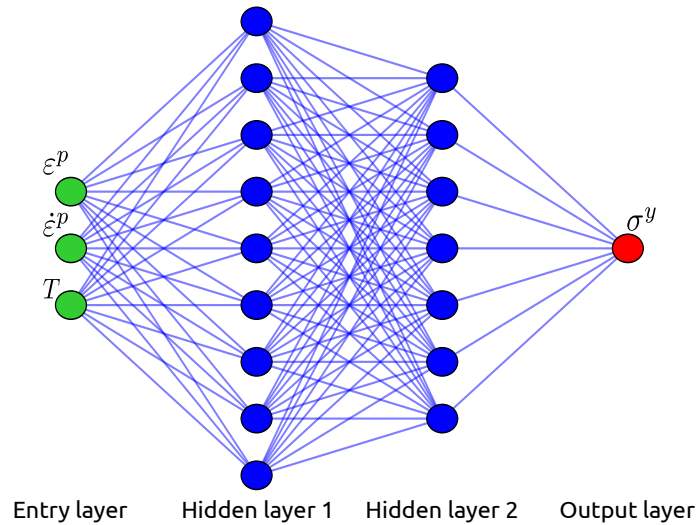


Figure 4: Architecture of the ANN 3-9-7-1-sig used for the flow law

215 In a preliminary phase, after having selected the global architecture of the neu-
 216 ral network, it is necessary to proceed to its training from some inputs. The inputs
 217 for this application were generated from the Johnson-Cook flow law expression
 218 reported in Equation (1) and the identified parameters reported in Table 2. This
 219 approach was chosen to demonstrate the ability of the neural network flow law to
 220 replace a classically formulated flow law such as Johnson-Cook’s for the simu-
 221 lation of metal cutting. In future developments, experimental tests on a Gleeble
 222 thermomechanical simulator associated with Taylor impact tests or Hopkinson
 223 bar tests will be used to generate this network training data. The training data,
 224 presented in the form of a data table containing the plastic strain ε^p , the plastic
 225 strain rate $\dot{\varepsilon}^p$, the temperature T and the flow stress σ^y , is processed by a learn-
 226 ing algorithm, developed at LGP, in Python, using the Tensorflow library [41].
 227 One hour of training on a Dell XPS13 7390 laptop running Ubuntu 20.04 64 bits
 228 with 16 GiB of Ram and an Intel 4-core i7-10510U processor allow obtaining the
 229 converged parameters of the ANN model.

230 Once this learning phase is completed, the neural network parameters result-
 231 ing from the learning process are used directly by a Python program, in charge
 232 of automatically generating the Fortran source code of the VUHARD subroutine

233 in order to compute the flow stress σ^y and its three derivatives, required for the
 234 explicit Abaqus FEM code.

235 The main advantage of this approach (the use of an ANN), after the learning
 236 phase, is that, for example, the output σ^y of the network is now linked to the inputs
 237 ε^p , $\dot{\varepsilon}^p$, and T by the equations (3) to (7) for a two hidden layers neural network
 238 with a sigmoid activation function as proposed previously.

239 Thus, in the VUHARD subroutine, the computation of the flow stress σ^y from
 240 the 3 input variables ε^p , $\dot{\varepsilon}^p$, and T is performed using the following procedure.
 241 The first step is to scale the input data to the interval $[0, 1]$ using the following
 242 equation:

$$\vec{x} = \begin{cases} x_1 = \frac{\varepsilon^p - [\varepsilon^p]_{min}}{[\varepsilon^p]_{max} - [\varepsilon^p]_{min}} \\ x_2 = \frac{\ln(\dot{\varepsilon}^p) - [\ln(\dot{\varepsilon}^p)]_{min}}{[\ln(\dot{\varepsilon}^p)]_{max} - [\ln(\dot{\varepsilon}^p)]_{min}} \\ x_3 = \frac{T - [T]_{min}}{[T]_{max} - [T]_{min}} \end{cases} \quad (3)$$

243 where quantities $[]_{min}$ and $[]_{max}$ are the boundaries of the range of the corre-
 244 sponding field during the training phase. Corresponding values, for the proposed
 245 case, are given in Appendix A. According to the architecture of the network,
 246 the outputs of the neurons of the first hidden layer \vec{y}_1 are given by the following
 247 equation:

$$\vec{y}_1 = \text{sig}(\mathbf{W}_1 \cdot \vec{x} + \vec{b}_1) \quad (4)$$

248 where, \mathbf{W}_1 and \vec{b}_1 are the weights and biases associated with the first hidden layer
 249 and $\text{sig}()$ is the sigmoid activation function defined by the equation (5) :

$$\text{sig}(x) = \frac{1}{1 + e^{-x}} \quad (5)$$

250 Then, the output of the neurons of the second hidden layer is given by the
 251 equation (6) :

$$\vec{y}_2 = \text{sig}(\mathbf{W}_2 \cdot \vec{y}_1 + \vec{b}_2) \quad (6)$$

252 where, \mathbf{W}_2 and \vec{b}_2 are the weights and biases associated with the second hidden
 253 layer. Finally, the σ^y output of the ANN is thus given by the equation (7) :

$$\sigma^y = ([\sigma^y]_{max} - [\sigma^y]_{min}) (\vec{w}^T \cdot \vec{y}_2 + b) + [\sigma^y]_{min} \quad (7)$$

254 where, \vec{w} and b are the weights and the bias associated with the output layer.

255 On the other hand, the three derivatives of the yield stress σ^y with respect to
 256 the three input variables ε^p , $\dot{\varepsilon}^p$, and T are given by the equation (8):

$$\begin{cases} \partial\sigma^y/\partial\varepsilon^p = s'_1 \frac{[\sigma^y]_{max}-[\sigma^y]_{min}}{[\varepsilon^p]_{max}-[\varepsilon^p]_{min}} \\ \partial\sigma^y/\partial\dot{\varepsilon}^p = s'_2 \frac{[\sigma^y]_{max}-[\sigma^y]_{min}}{[\dot{\varepsilon}^p]_{max}-[\dot{\varepsilon}^p]_{min}} \dot{\varepsilon}^p \\ \partial\sigma^y/\partial T = s'_3 \frac{[\sigma^y]_{max}-[\sigma^y]_{min}}{[T]_{max}-[T]_{min}} \end{cases} \quad (8)$$

257 where s'_i is the i^{th} component of the vector \vec{s}' defined by the equation (9):

$$\vec{s}' = \mathbf{W}_1^T \cdot \left[\mathbf{W}_2^T \cdot \left(\frac{\vec{w} \circ e^{-\vec{y}_2}}{[1 + e^{-\vec{y}_2}]^2} \right) \circ \left(\frac{e^{-\vec{y}_1}}{[1 + e^{-\vec{y}_1}]^2} \right) \right] \quad (9)$$

258 and \circ is the elements-wise product, known as the Hadamard product. In equa-
 259 tions (3) to (9), quantities \mathbf{W}_1 , \mathbf{W}_2 , \vec{w} , \vec{b}_1 , \vec{b}_2 and b are evaluated by the training
 260 procedure of the ANN. Corresponding values for an ANN containing 9 neurons
 261 in the first hidden layer and 7 neurons in the second hidden layer are reported in
 262 Appendix A. The set of equations (3) to (9), together with the network param-
 263 eters identified in the learning phase, is automatically translated into a VUHARD
 264 Fortran subroutine used by the FEM code Abaqus to simulate the cutting model.

265 Because of the large number of identified parameters for all the ANN models
 266 (from 114 to 202 for 9 and 17 neurons for the first hidden layer, respectively), the
 267 other 4 sets of ANN parameters used in this publication can be found in [42].

268 3.3. Sensitivity study of the results to mass scaling

269 FE modelling of the cutting process is very expensive in terms of CPU time
 270 due to the coupling of many nonlinear phenomena and the large amount of tiny
 271 finite elements. Mass scaling (MS) is introduced into the model to reduce the CPU
 272 computation time while checking that it does not influence the results (forces and
 273 energies) via a mass scaling sensitivity study. MS factors, MS_f , ranging from
 274 1E6 (theoretical CPU time scale of $\sqrt{MS_f} = 1000$) to 1 (no scale) were used for
 275 a cutting condition ($\lambda_s = 0^\circ$, $v_c = 30$ m/min and $h = 60$ μ m). The same signal
 276 processing procedure is applied to the numerical forces as to the experimental
 277 forces (cf. 2): they are filtered with a second-order low-pass Bessel filter at 750 Hz
 278 before calculating the steady state average value. Table 4 gives the results of the
 279 model with MS normalized (\hat{F}_i) by those of the model without MS:

$$\hat{F}_i = \frac{F_i \text{ with MS}}{F_i \text{ without MS}} \quad (10)$$

280 with $i = c$ for the cutting force and $i = f$ for the feed force. As expected, the
 281 real speed-up does not increase linearly with the MS_f , but it remains significant.
 282 A MS_f of 1E6 leads to an unstable computation and a MS_f of 1E5 leads to erratic
 283 force evolutions. These results are confirmed by high values of the ratio of the
 284 kinetic (KE) to the internal (IE) energies (it should not exceed a few % [43, 44]).
 285 A value of MS_f of 1E3 is chosen as it offers a good balance between reducing the
 286 computation time and the impact on the forces, while keeping the $\frac{KE}{IE}$ below 1 %.
 287 To provide an order of magnitude of CPU computation time, between 10 h and
 288 50 h (depending on the value of h) are required on 4 cores of an Intel i7-5700HQ
 CPU at 2.7–3.5 GHz.

Table 4: MS sensitivity study (selected MS factor, MS_f , in bold, \hat{F}_c : normalized cutting force, \hat{F}_f : normalized feed force, KE : kinetic energy, IE : internal energy)

MS_f	CPU scaling	Speed-up	\hat{F}_c	\hat{F}_f	$\frac{KE}{IE}$ (%)
1	1	1	1	1	2.3E−4
1E2	10	9	1.006	0.982	2.2E−2
1E3	32	21	1.008	0.940	2.2E−1
1E4	100	61	1.012	0.921	2.4
1E5	316	173	Erratic	Erratic	22
1E6	1000	207	Unstable	Unstable	58

289

290 3.4. Sensitivity study of the results to the number of neurons

291 The number of neurons in the hidden layers may influence the results. A
 292 sensitivity study on the number of neurons of the first hidden layer, ζ , is performed
 293 in order to select the ANN offering the best balance between CPU computation
 294 time and quality of the results. The results of the study are provided in Table 5.
 295 \check{F}_i corresponds to the results of the model with ANN normalized by those of the
 296 model with the built-in JC model:

$$\check{F}_i = \frac{F_i \text{ with ANN}}{F_i \text{ with JC}} \quad (11)$$

297 They show no influence on the numerical results for the forces compared to the
 298 built-in Johnson-Cook model, only the computation time is influenced by the num-
 299 ber of neurons in the first hidden layer and increases with it. This increase in com-
 300 putation time is not only due to the increasing complexity of the neural network

301 with the number of neurons, but also to the need to go through a VUHARD user
 302 subroutine. A first hidden layer of 9 neurons is therefore selected as it leads to the
 303 smallest increase in CPU computation time, without influence on the final result.

Table 5: Sensitivity of the forces to the number of neurons of the first layer, ζ (selection in bold, \check{F}_c : normalized cutting force, \check{F}_f : normalized feed force)

ζ	Time increase (%)	\check{F}_c	\check{F}_f
Built-in	0	1.000	1.000
9	6	1.000	0.999
11	6	1.001	1.000
13	7	1.000	0.998
15	8	1.001	1.001
17	10	1.000	1.000

304 4. Experimental and numerical results

305 An example of the temporal evolution of the numerical and experimental
 306 forces is plotted for the 3 directions in Figure 5 at $\lambda_s = 6^\circ$, $v_c = 10$ m/min and
 307 $h = 40$ m/min. The FE models are calculated up to a few microseconds after the
 308 stationary state is reached. Then, a linear extrapolation (dashed line between the
 309 last two markers in Figure 5) is used to provide numerical values for the same
 310 time range as the experimental values. The average and standard deviation (2σ)
 311 are calculated from the 3 experimental values. The resulting dispersion is shown
 312 in Figure 5 around the average values of each force. Steady state takes longer to
 313 be reached for the experiments than for the numerical model, in particular for the
 314 cutting force. The dispersion around the evolution of the average force is greater
 315 for the feed force than for the cutting force, while the average value of the feed
 316 force is 46 % of the average value of the cutting force. The numerical cutting force
 317 is very close to the experimental average cutting force; it is only 4 % higher. This
 318 difference, Δj , is calculated by :

$$\Delta j = \frac{|j^{(\text{sim})} - j^{(\text{exp})}|}{j^{(\text{exp})}} \times 100 \quad (12)$$

319 where j is the cutting force, the feed force, the passive force or the chip thick-
 320 ness. $j^{(\text{sim})}$ is the average value from the simulation, while $j^{(\text{exp})}$ is the average

321 experimental value.

322 The numerical feed force is underestimated by the model, but is within the
 323 95 % experimental confidence interval. The numerical passive force difference is
 324 also underestimated and is not within the narrower experimental dispersion. The
 325 difference between the average values of the experimental and numerical feed and
 326 passive forces is 25 %. A less well modelled feed force than the cutting force is
 327 typical of FE models of the cutting process and the difference with the experi-
 328 mental value is similar to other studies for a narrower range of cutting conditions
 329 [32, 45–48]. Hardt and Bergs [27] also obtained larger differences for feed and
 330 passive force than for cutting force. The difference for passive force was higher
 331 than for feed force, which is the opposite observation of this work.

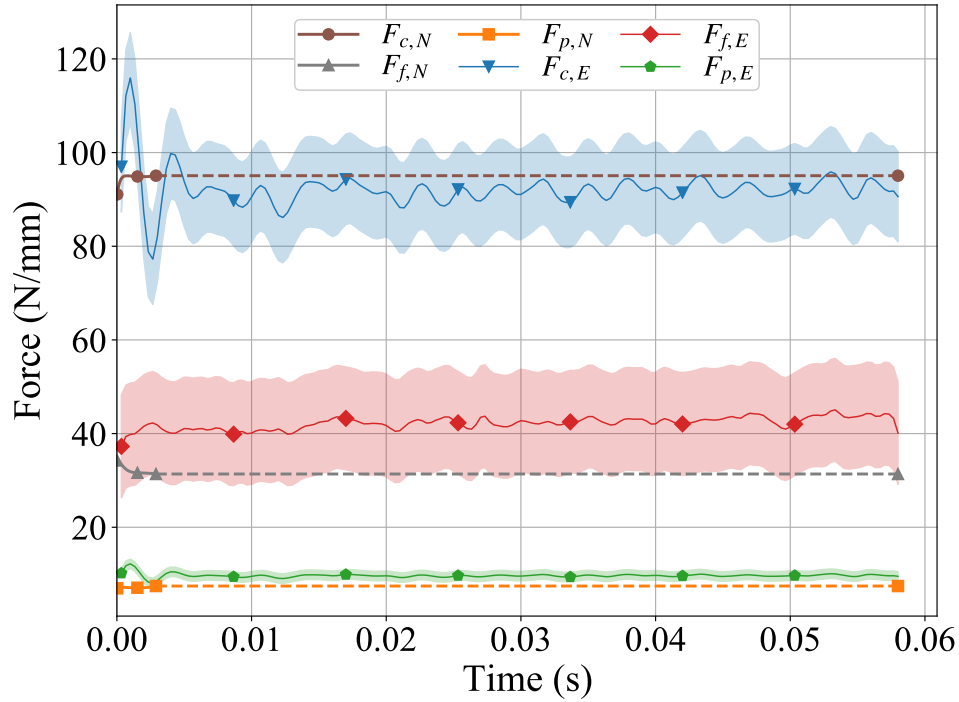


Figure 5: Temporal evolutions of experimental (E) and numerical (N) forces at $\lambda_s = 6^\circ$, $v_c = 10$ m/min and $h = 40 \mu\text{m}$ with dispersion around average experimental values (linear extrapolation of numerical values in dashed)

332 Numerical chips at $v_c = 10$ m/min and $h = 40 \mu\text{m}$ for $\lambda_s = 0^\circ$ and $\lambda_s = 6^\circ$

333 are provided in Figures 6 and 7. Due to the absence of heat gap generation in
334 the model, temperatures in the tool increase mainly by the heat generated by fric-
335 tion. They are therefore underestimated: maximal temperature in the tool is under
336 400 K (and all temperatures in the tools are in the blue colours with the scale of
337 Figure 6). When the inclination of the cutting edge is 0° , both sides of the chip
338 are identical and a symmetry plane can be drawn in the middle of the workpiece
339 (Figure 7 (a)). On the other hand, for an inclination of the cutting edge of 6° , the
340 chip is no longer aligned with the workpiece. The chip bends to one side due to
341 the orientation of the tool and symmetry is lost in both the geometry and the ther-
342 mal and mechanical fields, as shown in figure 7 (b). This produces helical chips
343 for the inclination angle of 6° as in the experiments. Figure 8 shows the variation
344 of the chip thickness across its width: it is thicker in the middle (i.e., the body of
345 the chip) than on its sides. This underlines the importance of 3D modelling, even
346 for the orthogonal cutting configuration as highlighted earlier [24]. The 3D mod-
347 elling also allows reproducing the lateral flow that occurs in the experiments for
348 both values of cutting edge inclination (Figure 6), unlike a 2D model [23–25]. Al-
349 though this leads to higher computation times, future cutting models should be in
350 3D, even when orthogonal cutting is considered. In this case, it is recommended
351 to take advantage of the symmetry of the configuration to reduce the computa-
352 tion time. This simplification has not been included in this study to avoid any
353 difference in the FE models between the 2 inclinations of the cutting edge.

354 Average values of the experimental forces and their dispersion are shown in
355 Figures 9 to 13 together with the average numerical values. Passive force values
356 are of course only plotted for $\lambda_s = 6^\circ$ as they are equal to zero when $\lambda_s = 0^\circ$.

357 The increase in cutting force with uncut chip thickness is clearly observed in
358 Figures 9 and 10 for both experimental and numerical results at the 2 inclination
359 angles, as well as the decrease in force with increasing cutting speed. This shows
360 that temperature softening dominates strain rate hardening for Ti6Al4V and is
361 accurately modelled. Increasing the inclination angle from 0° to 6° slightly re-
362 duces the cutting force; this is well captured by the model. For cutting speeds of
363 20–40 m/min and an inclination angle of 0° , F_c is almost constant with cutting
364 speed for uncut chip thicknesses of $40\ \mu\text{m}$ and $60\ \mu\text{m}$, while it decreases slightly
365 for $80\ \mu\text{m}$; this small stabilization is less marked for the model.

366 An increase in the deviation around the average value with the cutting speed
367 is noted for values above 10 m/min. All numerical values are within 95 % con-
368 fidence of the experiments (35 of the 36 conditions are within 68 % confidence).
369 The average difference with the experiments is 4 %, which is remarkable, also
370 considering the wide range of cutting conditions considered and the absence of

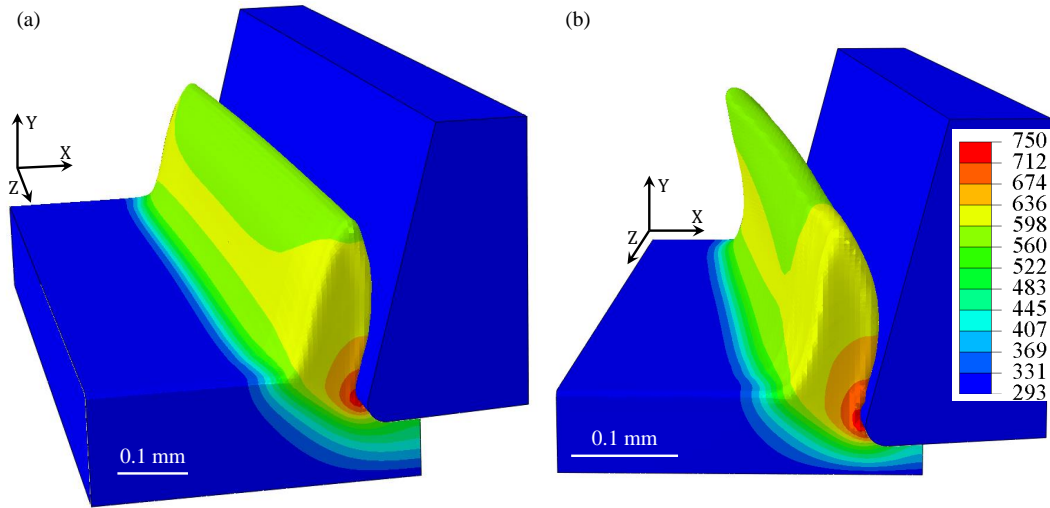


Figure 6: Temperature contours (in K) of the numerical chip after 1.5 ms at $v_c = 10$ m/min, $h = 40 \mu\text{m}$ and (a) $\lambda_s = 0^\circ$, (b) $\lambda_s = 6^\circ$

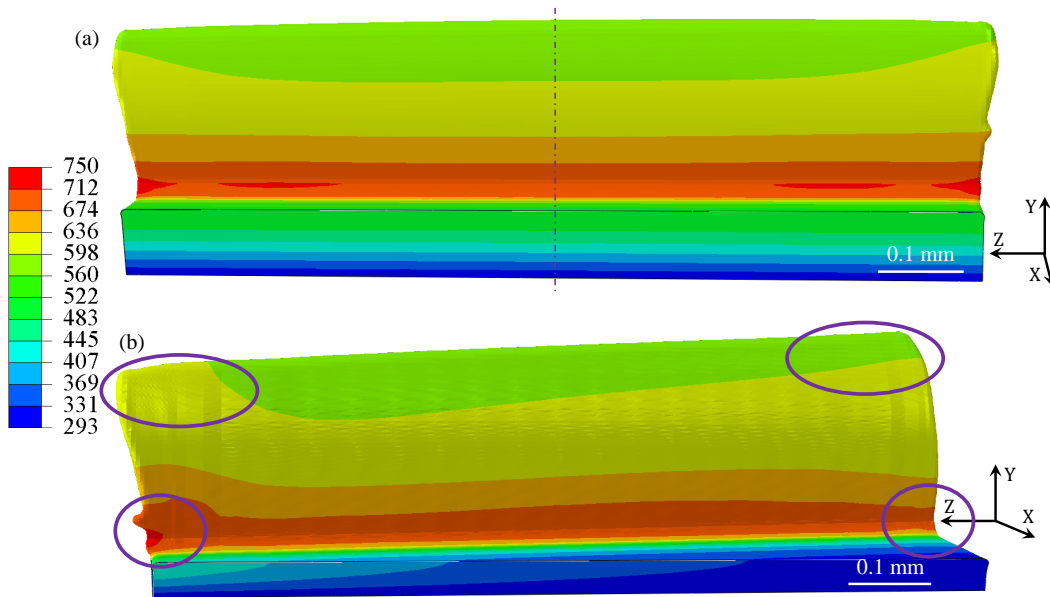


Figure 7: Temperature contours (in K) of the back of the numerical chip (tool is removed) after 1.5 ms at $v_c = 10$ m/min, $h = 40 \mu\text{m}$ and (a) $\lambda_s = 0^\circ$, (b) $\lambda_s = 6^\circ$

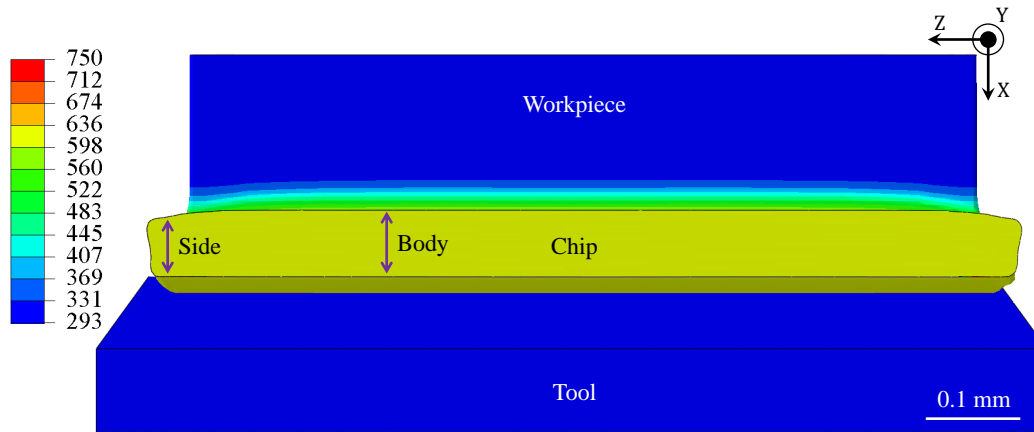


Figure 8: Temperature contours (in K) of the top of the numerical chip after 1.5 ms at $v_c = 10$ m/min, $h = 40$ μ m and $\lambda_s = 0^\circ$

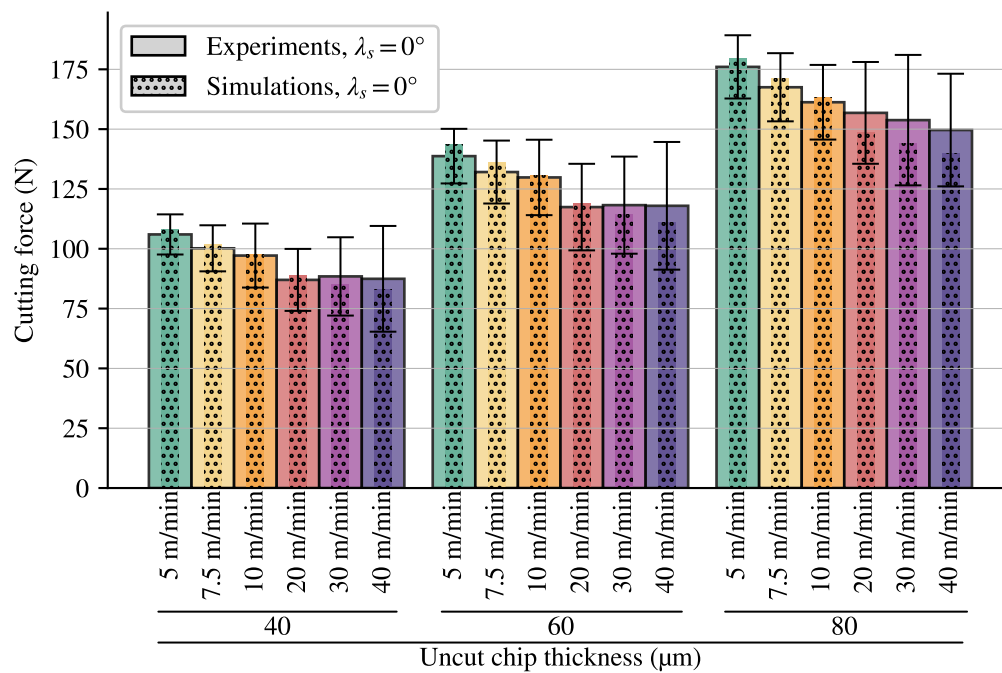


Figure 9: Comparison of experimental and numerical cutting forces at the cutting edge inclination of 0° for the 3 uncut chip thicknesses and the 6 cutting speeds

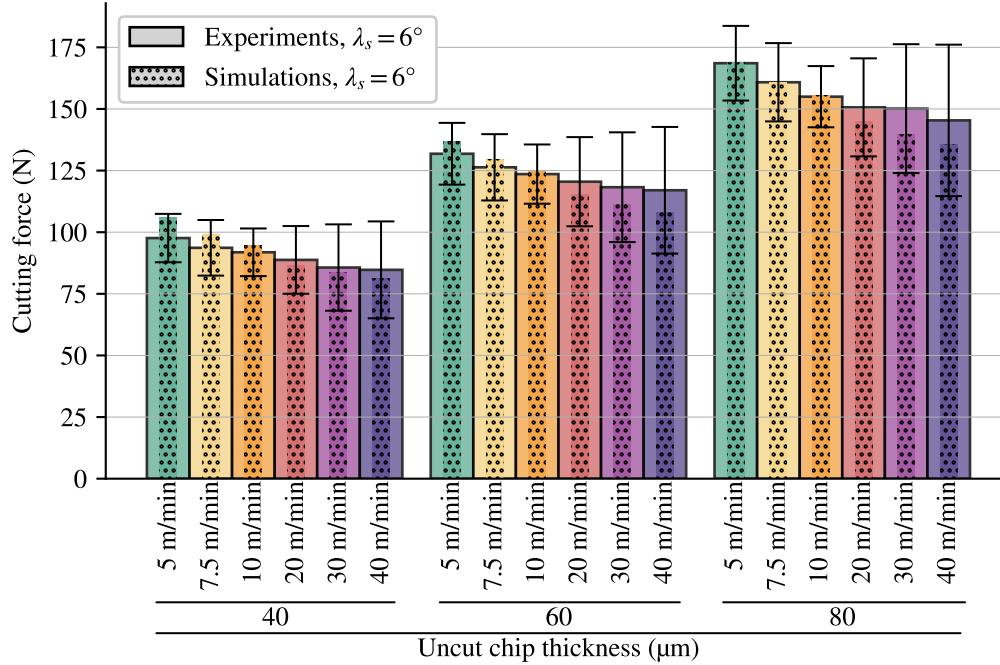


Figure 10: Comparison of experimental and numerical cutting forces at the cutting edge inclination of 6° for the 3 uncut chip thicknesses and the 6 cutting speeds

371 model tuning. This underlines the predictive ability and accuracy of the FE model
 372 for both inclination angles.

373 Figures 11 and 12 show the results for the feed force, where the two clearest
 374 trends for the experiments are its decrease with the inclination angle and its in-
 375 crease with the uncut chip thickness (even though it is lower than expected). For
 376 $80\ \mu\text{m}$, F_f decreases overall with v_c in the experiments. For $40\ \mu\text{m}$ and $60\ \mu\text{m}$, the
 377 force decreases at lower v_c , then increases for 0° , while a decrease is observed at
 378 all v_c for 6° (the experimental dispersion is high for both inclination angles, but the
 379 average trend with cutting speed is clear at 6° , not at 0°). For the numerical values,
 380 the overall trend is the same for the 3 uncut chip thicknesses and the two inclina-
 381 tion angles: a decrease for the lowest values of v_c and then an increase. It should
 382 be noted that the numerical model does not correctly handle the trends of the feed
 383 forces: as Figure 12 clearly shows, the numerical forces have an overall increas-
 384 ing trend with the cutting speed, while their average value mainly decreases when
 385 the uncut chip thickness increases. The differences between the average numeri-
 386 cal and experimental values increase with the uncut chip thickness: the forces are

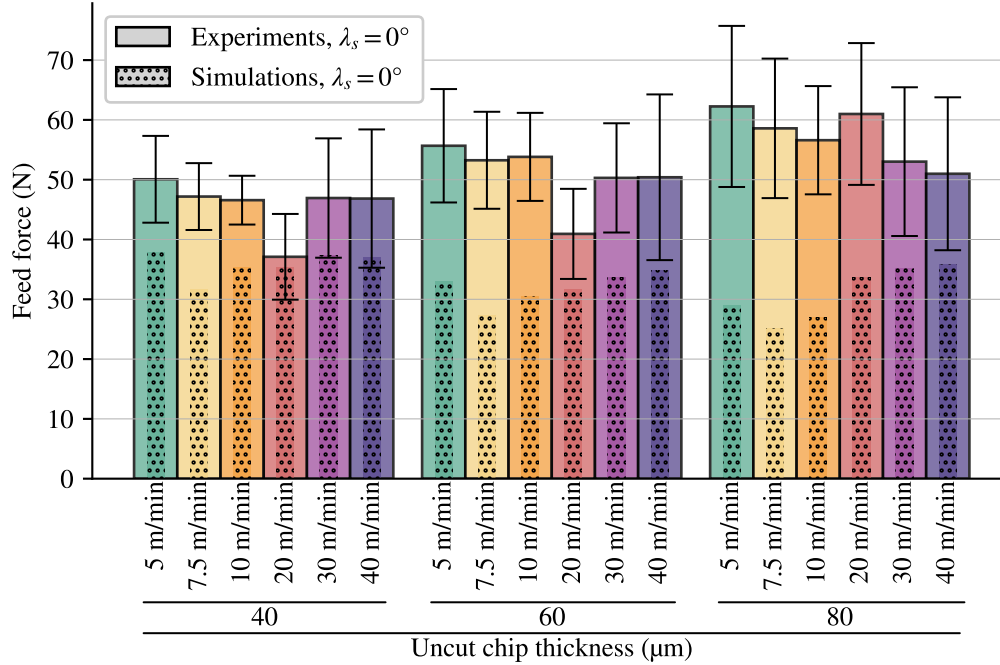


Figure 11: Comparison of experimental and numerical feed forces at the cutting edge inclination of 0° for the 3 uncut chip thicknesses and the 6 cutting speeds

387 closer at $40\ \mu\text{m}$ than at $80\ \mu\text{m}$. The numerical values are generally not within the
 388 95 % confidence interval (they do not clearly change with the cutting conditions).
 389 Coupled with the differences in trends, this shows that F_f is less well modelled
 390 (the average difference is 39 %) than F_c as usual in FE modelling of the cutting
 391 process and even more so in 3D [27]. The influence of the uncut chip thickness
 392 on the feed force should therefore be improved. The parameters of the material
 393 model are known to have an impact on the forces (and on the chip morphology)
 394 [15, 36]. The friction model should also be improved to strengthen the results
 395 [27].

396 The passive force is non-zero for the inclination angle of 6° (Figure 13). Like
 397 the cutting force, it increases with the uncut chip thickness and decreases with
 398 the cutting speed. The comparison with experiments is broadly the same as for
 399 F_c , except for a greater difference in the magnitude of F_p (the average difference
 400 is 26 %, but it is small in absolute terms – less than 5 N). Most of the numerical
 401 values do not fall within the experimental 95 % confidence interval. A lower mag-
 402 nitude of the passive force from the simulation than from the experiments with

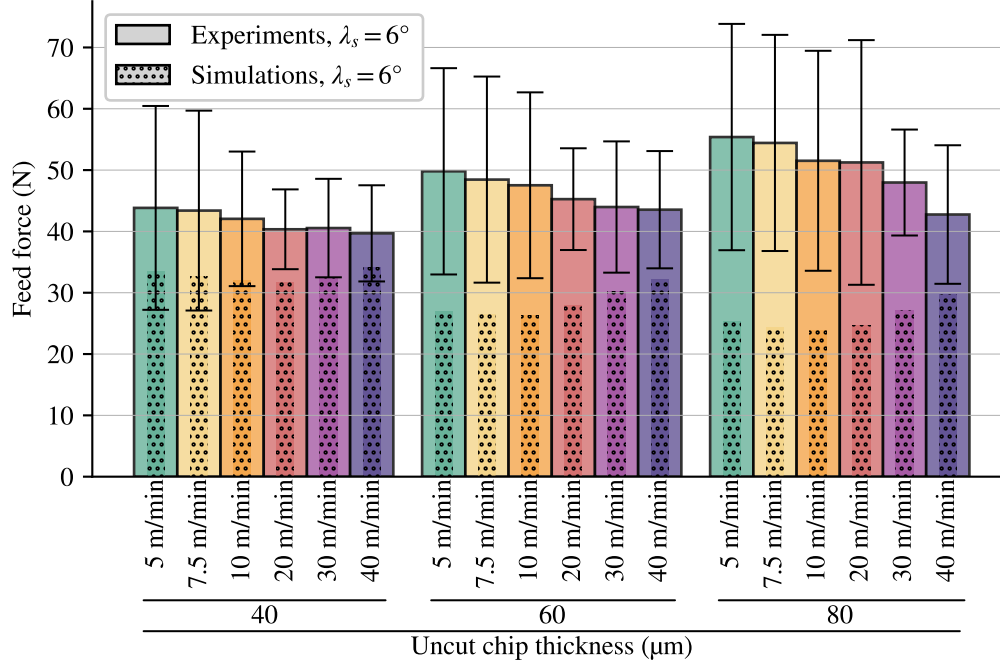


Figure 12: Comparison of experimental and numerical feed forces at the cutting edge inclination of 6° for the 3 uncut chip thicknesses and the 6 cutting speeds

403 the correct trends when the cutting conditions change was also observed by Hardt
 404 and Bergs [27]. The differences were mainly attributed to differences in cutting
 405 edge radius, friction modelling and material model. In this work, the impact of
 406 the cutting edge radius can be neglected as it is the same in the model as in the
 407 experiments.

408 As far as the chip morphology is concerned, all chips are continuous. For both
 409 the simulation and the experiments, the chip thickness ratio, λ_h :

$$\lambda_h = \frac{h'}{h} \quad (13)$$

410 with h the uncut chip thickness and h' the chip thickness, is almost independent of
 411 the uncut chip thickness (Figures 14 and 15). It is slightly reduced from $\lambda_s = 0^\circ$ to
 412 $\lambda_s = 6^\circ$, which means that the chip thickness decreases with the inclination angle.
 413 This influence is underestimated by the model: the reduction of λ_h is smaller than
 414 in the experiments. The average difference between the experimental and numeri-
 415 cal λ_h is 17 % over the whole range of cutting conditions. The chip thickness ratio

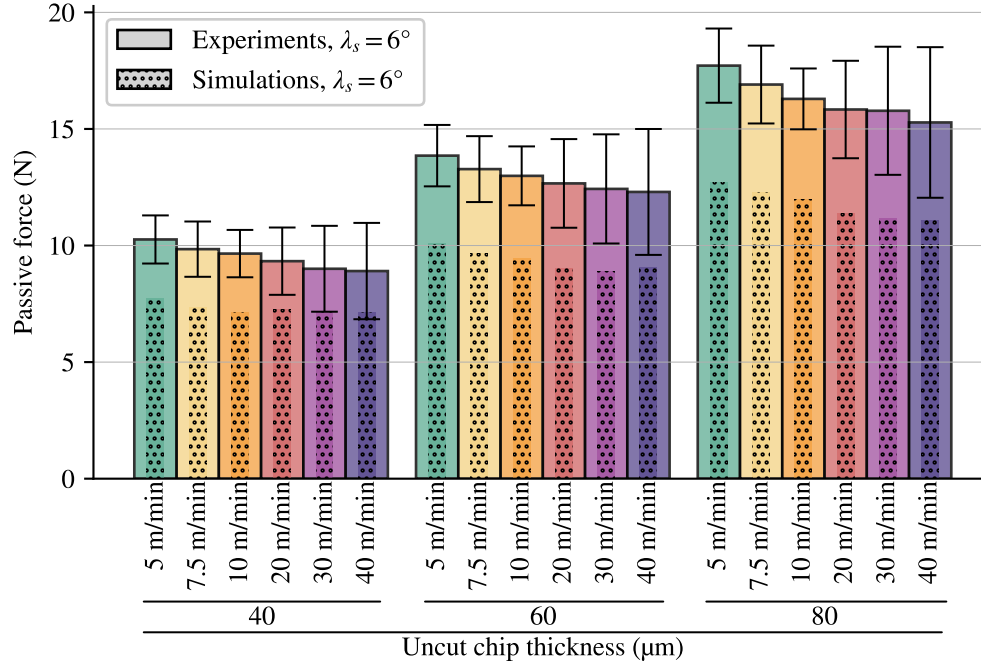


Figure 13: Comparison of experimental and numerical passive forces at the cutting edge inclination of 6° for the 3 uncut chip thicknesses and the 6 cutting speeds

416 decreases with cutting speed due to the reduction in friction, which is correctly ac-
 417 counted for by the model. As with the feed force, the results should be improved
 418 by more complex friction models and a set of material parameters for which the
 419 identification includes forces and chip thickness: [15].

420 The differences calculated according to the equation (12) are presented in Ta-
 421 ble 6 to provide a quantitative overview of the results. The cutting force is the best
 422 modelled quantity as observed in the literature. This result was to be expected
 423 as the parameter set of the material model was selected mainly due to its good
 424 approximation of the cutting force [36]. As this selection was made with a 2D
 425 model, the results show the ability of the model to correctly handle the third (pas-
 426 sive) force. Based on the average differences, the performance of the model is very
 427 close for the cutting and feed forces for both cutting edge inclinations, although
 428 a small degradation (1 % and 2 %, respectively) is noted for 6° . This degradation
 429 is more important (7 %) for the chip thickness ratio and must be linked to the dif-
 430 ference in passive force. Indeed, the chip thickness and out-of-plane force models
 431 are deeply linked. Improving the friction at the tool-workpiece interface should

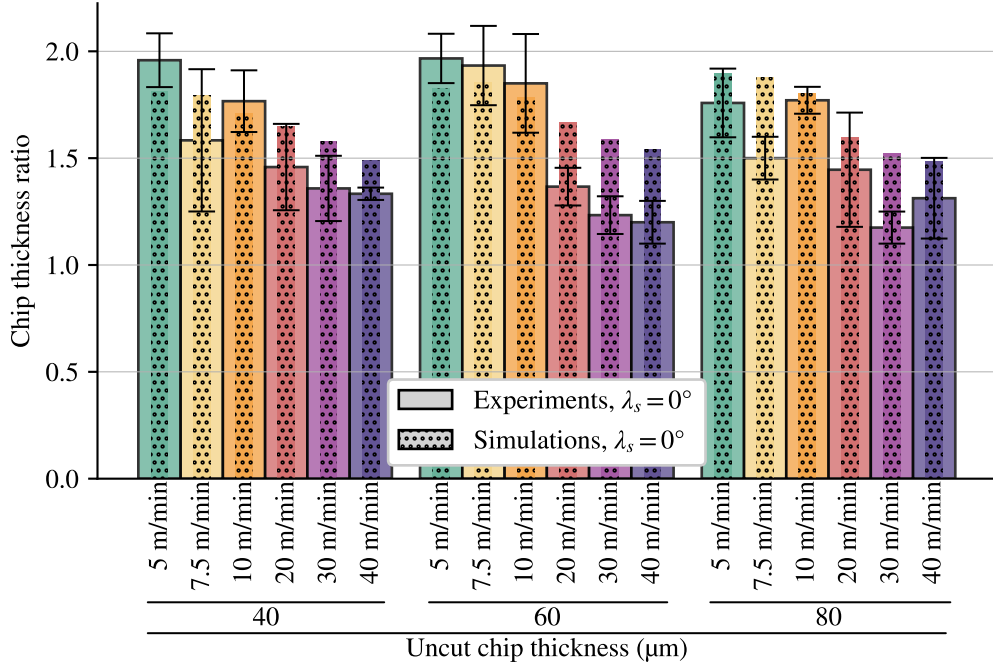


Figure 14: Comparison of experimental and numerical chip thickness ratios at the cutting edge inclination of 0° for the 3 uncut chip thicknesses and the 6 cutting speeds

432 be a key point. It should be noted that the chip thickness is very well modelled under
 433 certain cutting conditions with a minimum difference of 2 %. The difference
 434 is larger for the feed force than for the passive force, a trend opposite to that of
 435 Hardt and Bergs [27]. The average and range (min – max) of the differences are
 436 larger for the feed force. The smaller range of the passive force confirms a shift
 437 for all cutting conditions, similar to the results of Hardt and Bergs [27]. Again,
 438 the friction modelling should be the first aspect of the model to be improved in
 439 future developments.

440 5. Conclusions

441 An experimental and numerical study of the orthogonal and oblique free cut-
 442 ting of Ti6Al4V was carried out for a wide range of cutting conditions using an
 443 ANN-based flow law. The following main conclusions are drawn:

- 444 • The experimental study was carried out with the same set-up in free or-
 445 thogonal and free oblique cutting for the titanium alloy Ti6Al4V (the only

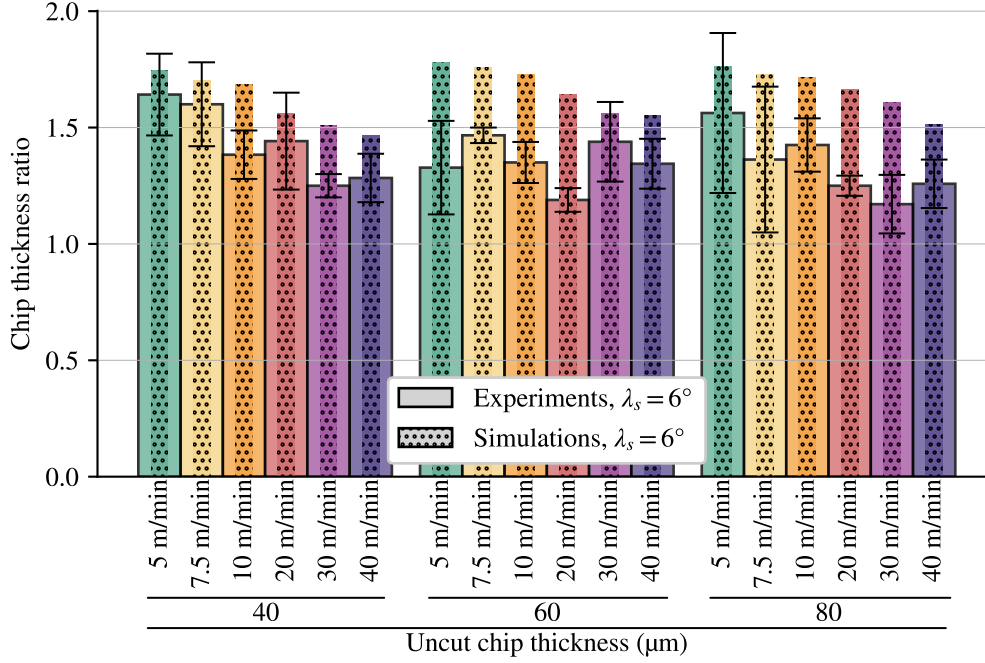


Figure 15: Comparison of experimental and numerical chip thickness ratios at the cutting edge inclination of 6° for the 3 uncut chip thicknesses and the 6 cutting speeds

Table 6: Synthetic quantitative overview of the results: differences between the experimental and the numerical results (average difference for each cutting edge inclination, and maximal, minimal and average differences for all the conditions) for the cutting force, ΔF_c , the feed force, ΔF_f , the passive force, ΔF_p , and the chip thickness ratio, $\Delta \lambda_h$

Difference	ΔF_c (%)	ΔF_f (%)	ΔF_p (%)	$\Delta \lambda_h$ (%)
Average $\lambda_s = 0^\circ$	3	38	–	14
Average $\lambda_s = 6^\circ$	4	40	26	21
Max. global	10	60	29	38
Min. global	1	10	19	2
Average global	4	39	26	17

446 change is the cutting edge inclination). This is a reference to evaluate the
 447 performance of the FE 3D model introducing an ANN-based flow law de-
 448 veloped under the same conditions. An unpreviously seen wide range of

449 cutting conditions, 36, is considered, including 2 cutting edge inclinations.

450 • A major novelty of this work is the accurate evaluation of the fundamental
451 variables and their trends in 3D, without the need to adjust the numerical
452 parameters and the model characteristics when the cutting conditions and
453 the inclination angle are changed significantly. The mere fact of changing
454 the inclination angle from free orthogonal cutting to oblique cutting while
455 maintaining the quality of the results has no equivalent in the current litera-
456 ture, especially since no studies (experimental or numerical) on free oblique
457 cutting are available.

458 • Taking into account the material's flow law by means of a neural network
459 makes it possible to overcome the limitations of conventional flow laws and
460 to reduce the approximations associated with the establishment of an analyt-
461 ical formulation of the flow law as conventionally adopted. The numerical
462 model is then able to better reproduce the real behaviour of the material and
463 to take into account thermomechanical transformations which are sources
464 of non-linearities, difficult to take into account with an analytical flow law
465 model. Current work, using a Gleeble thermomechanical simulator, on the
466 behaviour of a modified carbon alloy AISI P20 shows the advantages of
467 this approach compared to models in the literature such as Johnson-Cook,
468 Zerilli-Armstrong [5] or Hansel-Spittel [49], insofar as one is then able to
469 better reproduce more complex material behaviours.

470 • The cutting force is the best modelled quantity with an average difference
471 of 4 % with the experiments. Chip thickness ratio and passive force show
472 a larger deviation from the experiments (17 % and 26 %, respectively), but
473 their trends as the cutting conditions change are accurate. This is in line with
474 the expected results provided by a predictive model. The deviation for feed
475 force is higher (39 %), and opposite trends compared to the experimental
476 reference are observed. The lack of influence of uncut chip thickness on
477 friction in the model seems to be one of the aspects to be included as a
478 priority in future work. The model is found to handle the occurrence of the
479 third force, out of plane, well without significant degradation of the results.

480 • The predictive capabilities of the model make it suitable for the development
481 of straight-edged tools, for example. This work also demonstrates the ability
482 to model material behaviour with ANN and opens up possibilities in this
483 promising direction.

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487

488 **Competing Interests**

489 The authors have no relevant financial or non-financial interests to disclose.

490

491 **Author Contributions**

492 François Ducobu contributed to Data curation, Formal analysis, Investigation,
493 Methodology, Software, Supervision, Validation, Visualization, Writing – orig-
494 inal draft and review & editing (focussing on non ANN-related aspects). Olivier
495 Pantalé contributed to Data curation, Formal analysis, Investigation, Methodol-
496 ogy, Software, Validation, Visualization, Writing – original draft and review &
497 editing (focussing on ANN-related aspects). Bert Lauwers contributed to Super-
498 vision and Writing – review & editing. All authors read and approved the final
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660 **Appendix A. Coefficients of the ANN 3-9-7-1-sig**

661 In this appendix, we present the values obtained after the training phase of an
662 ANN containing 9 neurons in the first hidden layer and 7 neurons in the second
663 hidden layer. Conforming to [33], this one is referred ANN-3-9-7-1-sig. The
664 training of the neural network was performed using a dataset containing 3430
665 data points defined by:

- 666 • 70 equidistant values for $\varepsilon^p \in [0, 3]$, so that $[\varepsilon^p]_{min} = 0$ and $[\varepsilon^p]_{max} = 3$.
- 667 • 7 plastic strain rates $\dot{\varepsilon}^p \in [1/s, 10/s, 50/s, 500/s, 5000/s, 50\,000/s, 500\,000/s]$,
668 so that $[\ln(\dot{\varepsilon}^p)]_{min} = 0$ and $[\ln(\dot{\varepsilon}^p)]_{max} = 13.12236$.
- 669 • 7 temperatures $T \in [293\text{ K}, 400\text{ K}, 500\text{ K}, 700\text{ K}, 900\text{ K}, 1200\text{ K}, 1500\text{ K}]$, so
670 that $[T]_{min} = 293\text{ K}$ and $[T]_{max} = 1500\text{ K}$.

671 Stresses in the training dataset ranges from $[\sigma^y]_{min} = 171.4$ MPa to $[\sigma^y]_{max} =$
672 2606.1 MPa. The results of the training process are given here after for the ANN
673 quantities \mathbf{W}_1 , \mathbf{W}_2 , \vec{w} , \vec{b}_1 , \vec{b}_2 and b . The weight matrix for the first hidden layer
674 \mathbf{W}_1 is a 9×3 matrix:

$$\mathbf{W}_1 = \begin{bmatrix} -0.87229 & -0.47675 & -1.50771 \\ -0.95762 & -0.25619 & 1.65222 \\ -10.61660 & 0.22003 & -0.11539 \\ 3.67883 & 0.37146 & -1.51069 \\ -63.39468 & 0.15466 & -0.95431 \\ 0.54807 & 0.25959 & -5.44355 \\ -1.33883 & 0.36089 & -1.66735 \\ -0.68125 & 1.02121 & 0.34242 \\ 0.08740 & 0.18764 & -41.32542 \end{bmatrix}$$

675 The weight matrix for the second hidden layer \mathbf{W}_2 is a 7×9 matrix:

$$\mathbf{W}_2^T = \begin{bmatrix} 1.66285 & -0.59645 & -3.17333 & 0.20706 & 1.18760 & 2.01250 & -0.82147 \\ -0.26237 & -2.50330 & -1.45941 & -1.59833 & 4.05169 & -1.21146 & 1.05610 \\ -0.12958 & 0.67119 & -5.85989 & -2.55061 & 4.85245 & 4.31876 & 3.24070 \\ -2.12890 & 0.68296 & 0.71183 & 0.81706 & -0.09405 & 0.34919 & -1.41223 \\ 2.33631 & -0.08089 & 14.65789 & 0.12531 & 23.66363 & 2.55872 & 2.15338 \\ 0.11567 & 1.77629 & -1.80448 & 0.77825 & -1.58254 & 1.90442 & 1.23152 \\ 1.49265 & 0.41821 & -3.53803 & -0.48705 & -0.23671 & 0.75887 & -0.37441 \\ 0.95990 & 0.69041 & 0.43870 & 0.28393 & -1.40101 & -0.64569 & -0.38964 \\ 5.89937 & -0.13015 & 2.99264 & 1.78534 & -3.90189 & 1.17494 & -3.78854 \end{bmatrix}$$

676 The weight vector for the output layer \vec{w} is a 7 components vector:

$$\vec{w} = \begin{bmatrix} 0.34701 \\ 1.42079 \\ -0.96564 \\ 0.62467 \\ -0.56322 \\ 0.40960 \\ -0.42810 \end{bmatrix}$$

677 The biases of the first hidden layer \vec{b}_1 is a 9 components vector:

$$\vec{b}_1 = \begin{bmatrix} 2.57141 \\ 0.22673 \\ -1.16985 \\ -0.11246 \\ -0.82210 \\ -2.13264 \\ 0.78794 \\ 1.20434 \\ -3.48681 \end{bmatrix}$$

678 The biases of the second hidden layer \vec{b}_2 is a 7 components vector:

$$\vec{b}_2 = \begin{bmatrix} -0.36566 \\ -1.14445 \\ -0.79065 \\ -0.50670 \\ 1.30136 \\ 0.04521 \\ -0.29995 \end{bmatrix}$$

679 The bias of the output layer b is a scalar:

$$b = 0.04213$$

680 The corresponding coefficients for the other networks identified during this
681 work (ANN-3-11-7-1-sig, ANN-3-13-7-1-sig, ANN-3-15-7-1-sig and ANN-3-17-
682 7-1-sig) can be found in [42].