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Article

Anthropogenic carbon pathways towards the North Atlantic interior revealed by Argo-O₂, neural networks and back-calculations

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The subpolar North Atlantic (SPNA) is a region of high anthropogenic CO₂ (Cant) storage per unit area. Although the average Cant distribution is well documented in this region, the Cant pathways towards the ocean interior remain largely unresolved. We used observations from three Argo-O2 floats spanning 2013-2018 within the SPNA, combined with existing neural networks and back-calculations, to determine the Cant evolution along the float pathways from a quasi-lagrangian perspective. Our results show that Cant follows a stepwise deepening along its way through the SPNA. The upper subtropical waters have a stratified Cant distribution that homogenizes within the winter mixed layer by Subpolar Mode Water formation in the Iceland Basin. In the Irminger and Labrador Basins, the high-C_{ant} footprint (> 55 µmol kg⁻¹) is mixed down to 1400 and 1800 dbar, respectively, by deep winter convection. As a result, the maximum C_{ant} concentration is diluted (<45 µmol kg⁻¹). Our study highlights the role of water mass transformation as a first-order mechanism for C_{ant} penetration into the ocean. It also demonstrates the potential of Argo-O₂ observations, combined with existing methods, to obtain reliable Cant estimates, opening ways to study the oceanic Cant content at high spatio-temporal resolution.

Since the beginning of the industrial revolution, human activities such as fossil fuel burning and changes in land-use have led to the emission of large amounts of carbon dioxide (CO₂) into the atmosphere. This excess of carbon, referred to as anthropogenic carbon (C_{ant}), represented an addition of 9.5 ± 0.5 GtC yr⁻¹ to the atmosphere between 2011 and 2020¹. The ocean has absorbed $25 \pm 2\%$ of C_{ant} emissions (170 ± 35 GtC out of 660 ± 65 GtC of total C_{ant} emissions) since the beginning of the industrial era¹. The ocean acts, therefore, as a net C_{ant} sink and a moderator of climate change. This net C_{ant} uptake occurs via air-sea exchange, driven by air-sea CO₂ disequilibria resulting from the difference between the steady increase in atmospheric *p*CO₂ over the years since the pre-industrial era and oceanic *p*CO₂². The latter is primarily regulated by seasonal variations of temperature in the subtropics. In contrast, at higher latitudes, its oscillation is typically

dominated by biological processes³ affecting dissolved inorganic carbon (DIC) and total alkalinity (A_T). A low DIC/ A_T ratio will lead to a large CO₂ uptake capability by the ocean in response to an increase in atmospheric pCO_2^4 . As a result of the accumulation of C_{ant}, the Revelle factor, measuring the buffer capacity of the ocean, decreases⁵. Ultimately, changes in oceanic circulation also have the capacity to modify the saturation of C_{ant} at the ocean surface, enhancing or dampening the ocean carbon uptake capability⁶.

The global C_{ant} distribution in the ocean is not homogeneous, neither vertically nor horizontally⁷. As expected from direct air-sea exchange, C_{ant} concentration is maximum at the surface and decreases with depth. Of the global ocean, the North Atlantic Ocean (Fig. 1) is the region with the largest C_{ant} inventory per surface area^{8,9}, storing up to 23–38% of the total oceanic $C_{ant}^{2,10}$. Within this region, C_{ant} uptake

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Fig. 1 | Main circulation patterns and C_{ant} distribution (µmol kg⁻¹) for the year 2015 in the subpolar North Atlantic gyre. Red arrows represent the upper limb of the Atlantic Meridional Oceanic Circulation (AMOC) while blue arrows indicate the lower limb of the AMOC. The purple arrow represents ocean circulation between the upper and lower limb of the AMOC. The black contours represent the isobath -500 m, -1000 m, and -2000 m. The section on this figure depicts the whole water

column, while our study focuses on the top 2000 m of the water column only. NAC North Atlantic Current, WBC Western Boundary Current, ISOW Iceland-Scotland Overflow Water, LSW Labrador Sea Water. C_{ant} concentrations come from the GLODAPv2 dataset⁵⁶ and were normalized to the year 2015 using the exponential equation of Carter et al.²⁵.

mostly occurs in the subtropics^{6,7} due to the low Revelle factor of subtropical surface waters². These subtropical C_{ant}-loaded waters are then transported northward toward subpolar latitudes^{6,11} by the upper branch of the Atlantic Meridional Overturning Circulation (AMOC), where additional Cant uptake takes place¹². The subpolar North Atlantic (SPNA) gyre has been identified as one of the most important regions for the injection of C_{ant} towards the deep ocean⁷. This deep C_{ant} penetration is mainly due to the formation of Labrador Sea Water (LSW) through deep winter convection in the Labrador and Irminger Seas^{7,13-16}. Considering the role of the ocean in moderating the ongoing climate change through its capacity to uptake and store Cant, any change in its uptake and storage rates might have drastic consequences on the climate system. Changes in the Cant storage rates in the North Atlantic Ocean are fundamentally related to the increased atmospheric pCO_2 , the changes in the AMOC strength, the changes in the North Atlantic Oscillation (NAO) index, and deep water formation^{6,16,17}. The current Cant increasing trends and related Cant storage rates might thus be modified in the future due to the increasing atmospheric Cant, and the projected weakening of both AMOC¹⁸ and deep-water formation^{19,20}.

To date, C_{ant} estimates are mainly based on methods that rely on scarce but valuable ship-based measurements of carbonate system parameters (carbon-based methods^{2,6,21-23}) or transient tracers such as CFCs (transient tracer-based methods^{9,16}). However, studying the spatio-temporal evolution of oceanic C_{ant} storage and understanding the processes involved is a crucial challenge that requires a more detailed view of the upper and deep C_{ant} distribution and of its main pathways into the ocean interior. Additionally, the distribution of C_{ant} on timescales shorter than GO-SHIP cruises is necessary to understand

the effect of CO₂ emissions reduction and carbon removal strategies on the ocean. Considering the unrivaled spatio-temporal sampling provided by the Argo-O₂ network²⁴, the purpose of this study is to demonstrate that Argo-O₂ data combined with existing neural networks (i.e., ESPER_NN²⁵, CANYON-B and its routine CONTENT²⁶) and a back-calculation method (φC_T^{0} method^{17,27}) can be used to obtain reliable C_{ant} estimate at the finest spatio-temporal scale to date. As a case study, we selected three Argo-O₂ floats to describe C_{ant} deepening as it enters and propagates along the SPNA region. We highlight the long journey of C_{ant} from the surface towards the interior of the SPNA gyre. Our results open up perspectives to investigate C_{ant} storage changes and transport at seasonal to decadal timescales with the growing global and seasonally unbiased Argo-O₂ data network compared to ship-based measurements.

Results

Distribution of C_{ant} associated with the North Atlantic Current Four main regions with the different vertical distributions of C_{ant} were identified along the floats' pathways: region 1, located East of the Grand Banks of Newfoundland, and regions 2, 3, and 4, comprising the Iceland, Irminger, and Labrador basins, respectively (Fig. 2a). Region 1 is located inside a dynamic transition zone separating the subpolar and subtropical domains of the North Atlantic Ocean²⁸. This transition zone is typically characterized by a sharp thermohaline and density front, which, in this case (Fig. 2c), indicates that float 5904988 entered the North Atlantic Current (NAC)²⁹. The thermohaline front is also accompanied by a sharp C_{ant} and oxygen gradient (Fig. 2b and Supplementary Fig. 5). In the first 400–600 dbar of the water column, region 1 is characterized by saline, warm, and poorly oxygenated North





(µmol kg⁻¹) along the float pathway. The black line represents the mixed layer depth (MLD). The white lines represent the limits of the Labrador Sea Water (LSW), defined by $O_2 \ge 290 \,\mu$ mol kg⁻¹ in the Labrador and Irminger Seas and by S < 34.94 outside these two basins. The brown lines represent the isopycnals (kg m⁻³). The black arrow indicates when the float is located within an eddy. LSW Labrador Sea Water, NACW North Atlantic Central Water. **c** Salinity section along the float pathway.

Atlantic Central Water (NACW) of subtropical origin. At this location, the NACW comprises the northern extension of the upper and thermocline subtropical waters, defined as water masses located between the surface and $\sigma_{\theta} = 27.0 \text{ kg m}^{-3.12}$. These waters are typically characterized by a low Revelle factor⁵, which favors atmospheric C_{ant} uptake along their journey from the subtropics. As a consequence, the NACW presents the highest values Cant (maximum $C_{ant} > 65.0 \,\mu\text{mol kg}^{-1}$) of all water masses identified along the pathway of float 5904988 (Fig. 2a). The values of Cant, def close to zero (Supplementary Fig. 3) indicate that the NACW is saturated in Cant due to the air-sea Cant uptake occurring during the journey of the water masses that form the NACW. Furthermore, just eastward of the outcrop of the isopycnal 27.2 kg m⁻³, the high C_{ant} concentrations $(C_{ant} > 55.0 \,\mu\text{mol kg}^{-1})$ reach the base of the winter mixed layer depth (MLD) at ~400 dbar (Fig. 2b). In this region, the high Cant signal is vertically transferred as the MLD deepens, highlighting the vertical Cant homogenization due to mixing during winter convection. Below the mixed layer, between August 2017 and March 2018, Cant concentration in region 1 decreased sharply to a relative minimum of $26.2 \pm 6.6 \,\mu\text{mol}\,\text{kg}^{-1}$ near 800 dbar (Fig. 2b). These concentrations increase again up to $36.1 \pm 7.5 \,\mu\text{mol kg}^{-1}$ at 1050 dbar, associated with the Labrador Sea Water (LSW). As inferred by the progression of the float, the Cant-loaded NACW is transported northwards by the NAC (Fig. 1), toward the Iceland Basin (region 2), following the general circulation pattern of the SPNA gyre³⁰.

Distribution of C_{ant} in the Iceland Basin and over the Reykjanes Ridge

Region 2 encompasses the Iceland Basin, including the eastern flank of the Reykjanes Ridge, and is characterized by the presence of Subpolar Mode Water (SPMW) in the first 600 dbar of the water column. In region 2, the SPMW contains the largest Cant values over the water column, whose average is $57.2 \pm 5.8 \,\mu$ mol kg⁻¹. The SPMW is formed during winter convection^{31,32} by intense air-sea buoyancy loss, leading to the densification of the Cant-loaded NACW on its way toward subpolar latitudes³³. This process is well illustrated by the good agreement between the winter MLD and the depth of the SPMW base (Figs. 3b and 4b). In the Iceland Basin, the depth of the isoC_{ant} 50 μmol kg⁻¹, which mostly follows the winter MLD, reaches up to 600 dbar (Figs. 3b and 4b), that is, 200 dbar deeper than in region 1. Below the SPMW, Cant concentration decreases rapidly, highlighting trapping of C_{ant} within the SPMW and revealing that the formation of this water mass is a precursor in the deepening of Cant in the subpolar gyre. To further assess this deepening, Fig. 5 illustrates the mean Cant concentration along four density ranges describing the main water masses found in the four regions^{12,34}. In the second-density layer $(27.2 < \sigma_{\theta} < 27.52 \text{ kg m}^{-3})$ of region 2, the averaged C_{ant} concentration along the float 6901023 pathway increases from $40.1 \pm 6.2 \,\mu\text{mol}\,\text{kg}^{-1}$ in the south-easternmost part of the region to $55.5 \pm 6.5 \,\mu$ mol kg⁻¹ on the east side of the Reykjanes Ridge (Fig. 4b). This east-west Cant enrichment is likely due to the trajectory of the float, going from an area containing a thin SPMW layer, to an area with a thick SPMW layer (Fig. 3b). In this region, $C_{ant,def}$ fluctuates between –20 μ mol kg⁻¹ and $0 \,\mu\text{mol}\,\text{kg}^{-1}$ (Supplementary Fig. 3), suggesting that air-sea C_{ant} enrichment takes place, explaining part of the east-west Cant enrichment in the Iceland basin. In the western part of Region 2, the Reykjanes Ridge entails a natural topographic barrier separating the Iceland and Irminger basins, two basins with different C_{ant} distribution and water masses. The isoCant 35 µmol kg⁻¹, which fluctuated between 600 and 800 dbar east of the Reykjanes Ridge, deepens until ~1200 dbar on the western side of the ridge (Figs. 3b and 4b). The second density layer of region 2 (27.2 < σ_{θ} < 27.52 kg m⁻³) disappears west of the ridge during the winter months (Fig. 5, Supplementary Figs. 6c and 7c) due to outcropping and subsequent surface water densification in the vicinity of the ridge¹². As a consequence, the third oceanic layer

(27.52 < σ_{θ} < 27.71 kg m⁻³) shallows on the western side of the Reykjanes Ridge, illustrated by the shallowing of the isopycnal 27.71 kg m⁻³ from -1000 dbar east of the ridge to less than 500 dbar west of the ridge (Fig. 5b, c). Along with this densification of the surface water, C_{ant} increases by 34% and 31% in the third layer for float 6901023 and float 6901026, respectively, marking a rapid transfer of C_{ant} into the third oceanic layer. However, if it cannot be excluded that the C_{ant} increase is only due to vertical entrainment, this increase could partly be due to lateral entrainment of water masses from another origin. These results highlight the role of the Reykjanes Ridge, separating the C_{ant} distribution and dynamics between region 2 and region 3³⁴. In the fourth oceanic layer (27.71 < σ_{θ} < 27.80 kg m⁻³), C_{ant} concentrations are the lowest and slightly increase along the float trajectories (Fig. 5b, c). This increase stays within the averaged uncertainty range (±5.9 µmol kg⁻¹) and is therefore not significant.

Deep transport of Cant in the Irminger and Labrador basins

The Irminger Sea (region 3) is characterized by high concentrations of O₂ (Supplementary Figs. 5b, 6b, and 7b), delimiting the core of the LSW^{34} . These large O₂ values are indicative of the intense air-sea exchange that characterizes this region^{13,35,36}. In contrast to regions 1 and 2, here, the vertical Cant distribution has been subject to a large water-column homogenization. This leads to a weaker Cant gradient, with a gradual decrease of mean C_{ant} concentrations from 52.0 ± 5.8 $\mu mol~kg^{\mathchar`-1}$ at the surface to $41.2\pm7.9~\mu mol~kg^{\mathchar`-1}$ at the lower bound of the LSW (~1400 dbar) (Figs. 2b and 4b). The high Cant values are constrained within the LSW, meaning that the formation of this water mass traps and transports Cant toward the deeper ocean. Yet, the maximum Cant concentrations are still found at the surface, with values reaching $52.0 \pm 5.8 \,\mu\text{mol}\,\text{kg}^{-1}$ (Figs. 2b and 4b). As suggested by the negative Cant.def in the surface waters of the Irminger Sea (Supplementary Fig. 3b and 3c), these high surface concentrations are likely due to the intense air-sea CO₂ exchange, following the same behavior of O₂ exchanges³⁵. Following the Western Boundary Current (WBC) (Fig. 1), Cant is transported out of the Irminger Sea. In the WBC, we observe the occurrence of occasional high-Cant pulses throughout the water column (black arrows on Figs. 2b and 3b) that we relate to anticyclonic mesoscale eddies (Supplementary Fig. 8). These anticyclonic features lead to a punctual downward isopycnal displacement of surface waters containing high Cant concentration, explaining the Cant pulses identified in our sections. Such eddies could carry Cant away from their formation regions and when they collapse, they might contribute to the isopycnal mixing of this tracer. Once in the Labrador Sea, float 6901023 illustrates how the isoCant 35 µmol kg⁻¹ reaches a maximum pressure level of 1800 dbar in the inner part of the basin (Fig. 3b). That is, implying a ~ 200 dbar deepening of this isoCant compared to what is observed in the Irminger Sea. As in the Irminger Sea, in the Labrador Sea, the high Cant concentrations are also contained within the LSW, confirming that the formation of this water mass is a first-order deeptrapping mechanism for Cant.

Discussion

This study describes the C_{ant} evolution along its pathways from the upper to the intermediate layers (up to 2000 dbar) of the SPNA gyre. It is based on a combination of existing neural networks and the back-calculation ϕC_T^{0} method, using Argo-O₂ observations as input data. Our results are based on three floats following the main branches of the subpolar gyre circulation. From this quasi-lagrangian perspective, we identified four main C_{ant} deepening stages and related them to their corresponding physical mechanisms. At the southernmost boundary of our study region, in the subtropical/subpolar transition zone near Flemish Cap, C_{ant} is mainly trapped within the NACW layer ($\sigma_0 < 27.2 \text{ kg m}^{-3}$). The NACW has the largest mean C_{ant} concentration (62.8 ± 5.9 µmol kg⁻¹) of all water masses intercepted by the studied floats. At this location, winter mixing favors vertical C_{ant}



Fig. 3 | Characteristics of the Argo float 6901023 in the North Atlantic

(47–65°N; 15–65°W). a Trajectory of the float with the four main regions. The color coding represents the date, with blue being the beginning of the trajectory and red being the end. The yellow arrow indicates when the float has passed the Reykjanes Ridge. The points with black outlines correspond to the profiles pointed out by the black arrows on the C_{ant} section. The points with red outlines represent the limits of the Reykjanes Ridge, defined by a depth shallower than 1700 m. **b** Section of

 C_{ant} (µmol kg⁻¹) along the float pathway. The black line represents the MLD. The white lines represent the limits of the Labrador Sea Water (LSW), defined by $O_2 \geq 290 \ \mu mol \ kg^{-1}$ in the Labrador and Irminger Seas and S < 34.94 outside these two basins. The blue lines represent the limits of the Subpolar Mode Water (SPMW) defined by a potential vorticity lower than $6 \times 10^{-11} \ m^{-1} \ s^{-1}$. The brown lines represent the isopycnals (kg m⁻³). Black arrows indicate when the float is located in an eddy. **c** Salinity section along the float pathway.



Fig. 4 \mid Characteristics of the Argo float 6901026 in the North Atlantic

(47–65°N; 15–65°W). a Trajectory of the float with the four main regions. The color coding represents the date, with blue being the beginning of the trajectory and red being the end. The green arrow indicates when the float has passed the Reykjanes Ridge, while the red arrow shows when the float passes Cape Farewell. The points with red outlines represent the limits of the Reykjanes Ridge, defined by a depth shallower than 1700 m. **b** Section of C_{ant} (µmol kg⁻¹) along the float pathway. The

black line represents the MLD. The white lines represent the limits of the Labrador Sea Water (LSW), defined by $O_2 \ge 290 \,\mu$ mol kg⁻¹ in the Labrador and Irminger Seas and *S* < 34.94 outside these two basins. The blue lines represent the limits of the Subpolar Mode Water (SPMW) defined by a potential vorticity lower than $6 \times 10^{-11} \,\text{m}^{-1} \,\text{s}^{-1}$. The brown lines represent the isopycnals (kg m⁻³). **c** Salinity section along the float pathway.



Fig. 5 | **Averaged C**_{ant} **concentrations (µmol kg**⁻¹) **over density layers. a** For float 5904988. **b** For float 6901023. **c** For float 6901026. Scatter points indicate raw C_{ant} data, while the solid lines represent the monthly moving mean of C_{ant}

concentrations. Dashed lines represent the pressure levels of the isopycnals. The gray rectangle illustrates the time period when the floats are above the Reykjanes Ridge.

homogenization, by which high Cant concentrations (>45.0 µmol kg⁻¹) reach down to ~400 dbar (1st deepening stage). The air-sea heat fluxes and mixing processes along the northward progression of the Cantloaded NACW lead to a densification of NACW and its transformation into SPMW as it reaches the Iceland Basin. This transformation causes a deepening of Cant down to ~600 dbar in this region (2nd deepening stage). The average Cant concentration within this newly formed SPMW $(57.2 \pm 5.8 \,\mu\text{mol kg}^{-1} \text{ for } \sigma_{\theta} < 27.52 \,\text{kg m}^{-3})$ is lower than upstream due to the C_{ant} content being spread out in the vertical. Further west, the Reykjanes Ridge is identified as a key topographic feature that marks a steep transition between the Iceland Basin and the Irminger Basin with markedly different vertical Cant distributions. In the Irminger Sea, the Cant signal (>35.0 µmol kg⁻¹) deepens until 1400 dbar (3rd deepening stage) and is constrained within the LSW layer, highlighting the role of the LSW formation on the Cant deepening. The deepening is even more pronounced in the Labrador Sea, where the Cant signal, still confined within the LSW, reaches 1800 dbar (4th deepening stage). The newly formed LSW joins the Deep Western Boundary Current as part of the North Atlantic Deep Water (NADW), spreading the Cant footprint downstream towards the ocean interior. Further evidence of this Cant trapping within the LSW and its propagation downstream is provided by our data east of the Flemish Cap transition zone, where a relative C_{ant} maximum can be identified within the LSW layer underneath the NACW (Fig. 2b). As surface waters spread through the different basins of the SPNA, we also inferred water masses that have favorable conditions to uptake additional atmospheric Cant, by means of the Cant, def proxy. Our results indicate that air-sea Cant fluxes are likely to occur all over the region, and as far west as in the Labrador Basin.

Deep water formation is a broadly known and acknowledged mechanism for Cant penetration in the interior of the North Atlantic Ocean^{2,6,7}. Relying on Argo-O₂ data and existing methods, we presented an original approach that allowed us to describe this Cant deepening at the finest spatiotemporal scale to date. Our results revealed the role of the SPMW formation as an early stage of subpolar Cant deepening. This result reinforces the idea that Cant deepening occurs throughout the whole SPNA gyre and is not only linked to the main convective sites (i.e., Irminger and Labrador Basins), as traditionally attributed^{7,16}. Mainly uptaken from the atmosphere in the subtropics², Cant is transported northwards towards the SPNA region by the NAC⁶, which at these latitudes is the main component of the upper AMOC limb. The quasi-lagrangian Argo-float trajectories are a suitable and timely means for describing this horizontal Cant journey. Despite lateral advection by the upper AMOC limb, being the main source of Cant in the region, the contribution from ocean C_{ant} uptake is also relevant^{6,12}. By means of the $C_{ant,def}$ proxy, the Argo-O₂ data also support the hypothesis of surface air-sea Cant enrichment all along the subpolar gyre domain by an outcropping of the low-carbon nutrient stream isopycnal range¹². Given the dominant role of lateral advection in Cant distribution, it is not surprising that AMOC variability affects the northward Cant transport towards the SPNA and its storage rates^{6,11}. Under current climate change, the AMOC is projected to weaken³⁷, and we could thus expect a concomitant decrease in Cant storage rates and content in the SPNA. Nevertheless, the northward oceanic Cant transport might still be subject to continuous increase in response to the rise in anthropogenic CO₂ emissions^{11,38}. Estimates on a finer spatiotemporal scale, such as those based on Argo-O₂ observations, should

be sustained in the long term to determine which effect will be dominant in the future.

Our case study has shown that neural networks combined with high-quality in situ Argo-O₂ measurements and a back-calculation method can effectively be used to retrieve Cant concentration through the three-dimensional estimates of oceanic variables (nutrients, DIC, total alkalinity). These oceanic variables are needed to compute Cant at a higher spatio-temporal resolution than that available from the current biogeochemical (BGC) observation networks. Hence, the methodology presented here, enchaining Argo-O₂ observations and existing methods, opens up a promising step forward to study Cant distribution, its variability, and the processes driving these variations, which are crucial challenges regarding the ocean carbon sink³⁹. The global nature of the Argo-O₂ array gives the opportunity to significantly increase the spatiotemporal coverage of current Cant estimates. However, we can only guarantee the validity of our approach in the case-study region and for the period (1970-2020) over which the ship-based data used for training and validation of the neural networks were collected (GLODAPv2.2020⁴⁰). Neural networks can only reproduce what they have learned. Therefore they must be trained with new data as the ocean changes, which is particularly true for the SPNA gyre where interannual to decadal changes are considerable (e.g., cold blob⁴¹ or great salinity anomaly⁴²). The importance of high-accuracy in situ data measurements for training the neural networks cannot be disregarded, as such, the continuation of ship-based observations should be assured. Another important consideration of our results and their interpretation lies within the averaged total uncertainty of our method, which is ±5.9 µmol kg⁻¹. Any spatio-temporal change in C_{ant} concentration below two times this value should be interpreted with caution. This current methodological uncertainty could possibly be reduced by using data from the growing BGC-Argo database^{43,44} to train the neural networks. Moreover, BGC-Argo floats, equipped with pH and nitrate sensors, could be directly used to derive Cant, hence providing independent validation Cant estimates when both BGC-Argo and Argo-O2 floats are concomitant in space-time. However, BGC-Argo floats were not included in our study because available floats did not match our selection criteria. In view of the above, there is, therefore, a complementarity between the scarce high-quality ship-based BGC measurements and the high spatiotemporal BGC-Argo and Argo-O2 observations. This complementarity is even more marked because Argo data need to be adjusted (QCed) against ship-based data. The Argo-O2 float data presented in this case study reached a maximum pressure level of 2000 dbar, limiting our result to the top 2000 dbar of the water column. The OneArgo mission calls for the deployment of 1250 floats capable of reaching 4000 or 6000 dbar45. Using Deep-Argo floats equipped with dissolved oxygen sensors could thus strengthen our findings and improve our knowledge about the mechanisms controlling the deep/ abyssal Cant transport. For instance, the role of the Iceland-Scotland Overflow Water (ISOW), flowing below 2000 dbar, could not be identified here even though this water mass is known to contain C_{ant}¹¹.

Finally, it is worth mentioning that Earth system models tend to misrepresent the vertical penetration of C_{ant} due to the shallow and weak formation of NADW⁴⁶, even if this issue is better addressed in the latest generation of CMIP6 models^{18,47,48}. The results of this study provide insights subject to improving the representation of the key pathways for the deep C_{ant} propagation in Earth system models. Implementing a better representation of these pathways would definitely ameliorate the simulated vertical distribution of C_{ant} and, ultimately, improve climate predictions.

Methods

Argo float observations

To highlight the C_{ant} penetration into the deep ocean, we use Argo-O₂ floats, giving the opportunity to obtain one C_{ant} profile every 10 days. We selected the Argo-O₂ floats that (1) followed the cyclonic pathway

of the SPNA gyre, (2) had a lifetime longer than 3 years, and (3) crossed the A25 OVIDE hydrographic section⁴⁹, which will serve as a reference dataset for validation. Three floats matched these criteria, comprising a total of 651 profiles of pressure (P), temperature (T), salinity (S), and oxygen $(O_2)^{50}$ within the region 47–65°N; 15–65°W (Supplementary Fig. 1). Data span from July 2012 to April 2018. Only data adjusted in delayed mode with a quality flag of 1 or 2 (good or probably good data)⁵¹ were used in our analysis (Supplementary Table 1). The accuracy of the data is assumed to be 0.002 °C, 0.01, and 2.4 dbar for temperature, salinity, and pressure, respectively⁵², and better than 3 μ mol kg⁻¹ for oxygen⁵³.

Estimating biogeochemical variables with neural networks

To derive the biogeochemical variables (nutrients, DIC, and A_T) needed to compute C_{ant}, we rely on predictive neural networks. Neural networks (NN) are machine learning algorithms based on a multi-layer perceptron^{54,55}, that are trained and validated against observations. The input data of the neural networks are, in this study, Argo P/T/S/O₂ measurements. We used two neural networks, namely, ESPER NN²⁵ and CANYON-B. associated with its routine CONTENT²⁶. ESPER NN was used to obtain the macronutrients (phosphate, nitrate, and silicate) using P, T, S, O₂, location, and time as predictors. ESPER_NN was adopted over other ESPER methods (i.e., ESPER_LIR and ESPER_Mixed) because it gives the lowest biases and root mean square errors over the global ocean for the predicted macronutrients²⁵. ESPER_NN reproduces the validation dataset with average biases and errors of ~2% for macronutrients (Supplementary Table 1)²⁵. However, in the North Atlantic Ocean, ESPER NN gives uncertainties of ~1.3% for the predicted carbonate variables (A_T and DIC), which is higher than the previous NN. Consequently, A_T and DIC were computed with CANYON-B²⁶, and the outputs of CANYON-B were passed through the CONTENT routine. This routine ensures consistency between carbonate variables and thus reduces the uncertainties of the carbonate system variables to ~0.5% for A_T and DIC (Supplementary Table 1). We used neural networks rather than climatological products (such as, e.g., GLODAPv2 climatology⁵⁶) because gridded climatologies are unable to capture changes in carbon variables in regions where water masses move laterally due to mesoscale processes or rapid circulation changes. This is because climatologies represent the mean field in which the variability has been smoothed out. In contrast, neural networks, such as ESPER_NN and CANYON-B, derive biogeochemical variables based on water mass characteristics, hence they can cope with such changes. In particular, the use of Argo-O2 float data as input to the neural networks reinforces this water-mass change tracking capability due to the quasilagrangian behavior of the Argo-O₂ floats and their temporal resolution, three times greater than monthly climatological products.

Anthropogenic carbon estimates

To determine the anthropogenic carbon fraction (C_{ant}) from DIC, we used the carbon-based back-calculation φC_T^{0} method^{17,27}. This method has been widely applied to study the inventory of C_{ant} , its storage rates, its variability 14,17,57 , and the influence of C_{ant} on ocean acidification 58 . In any back-calculation method, Cant is computed as a two-step approach: (1) the changes in DIC due to biological activity occurring since a water parcel has left the ocean surface are removed, and (2) the preindustrial preformed DIC is also removed, with the residual being interpreted as C_{ant} . The ϕC_T^{o} method (Supplementary Equation 1) presents two main advantages compared to other back-calculation methods. First, this method accounts for the spatiotemporal variability of the preformed A_T. Second, the parameterization of the preformed A_T and disequilibrium terms are determined using the subsurface layer as a reference²⁷. These two modifications improve the C_{ant} estimates in cold and deep water formation regions subject to strong mixing processes, such as the SPNA gyre. A study comparing observational methods to estimate C_{ant} in the Atlantic Ocean, including the TTD⁹, the TrOCA⁵⁹, the C^o_{IPSL}⁶⁰, the Δ C^{*61}, and the ϕ C_T^O method¹⁷, demonstrated that the latter provided the closest value to the average of all methods for the whole latitudinal range²⁷. Based on that, we selected the ϕC_T^{O} method to perform our analysis. The input variables for this method are date, geographical location, T, S, O₂ (in this study, the Argo variables), the macronutrients (obtained from ESPER NN), plus A_T and DIC (computed from CANYON-B and CONTENT). The ϕC_T^{0} method requires date as input because it adjusts the oceanic Cant concentration against the corresponding atmospheric pCO_2 . It also requires the geographical location (latitude and longitude) to define the water masses involved in the computation of the preformed A_T and disequilibrium terms. All the input variables have associated uncertainties that propagate through Cant calculations. The uncertainties on Argo pressure, temperature, and salinity data are 2.4 dbar, 0.002 °C, and 0.01, respectively⁵². These uncertainties relate to the instrument's performance. The uncertainties on O₂ data are about 3.5 µmol kg⁻¹ (Supplementary Table 1). The uncertainties of the predicted variables A_T, DIC, and macronutrients are, on average, 10.85, 10.48, and 0.73 µmol kg⁻¹, respectively (Supplementary Table 1). These NNs uncertainties are provided as part of the NNs predictive outputs and they result from the combination of training measurement uncertainties, the Bayesian method uncertainties, and the NNs input sensitivities²⁶. Taking all of them into account, we quantified the uncertainty associated with our Cant estimate by randomly generating 100 C_{ant} fields (Supplementary Table 1) using a Monte Carlo method⁶². The standard deviation from these Cant fields fluctuates between $\pm 5.4\,\mu mol\,kg^{\text{--}1}$ and $\pm 10.2\,\mu mol\,kg^{\text{--}1}$ (Supplementary Fig. 2) and the overall average value (for the three floats) is $\pm 5.9 \,\mu$ mol kg⁻¹. It is worth highlighting that the uncertainties of the modeled tracer fields (Supplementary Table 1) are evidently greater than the corresponding in situ measured variables (i.e., by 1.5 µmol kg⁻¹ for oxygen). However, the uncertainties of the predicted variables are yet reasonably small to serve as inputs to the back-calculation method. Consequently, they allow an estimation of Cant that is statistically significant (and whose final uncertainties lie within 9% to 19% of the signal). To further validate our approach, we compared these ArgoTSO2-NN-based Cant estimates with ship-based C_{ant} estimates obtained during OVIDE cruises^{49,63}. We identified Argo-O₂ profiles within the 35 km radius and less than 7 month-timespan of an OVIDE CTD cast (Supplementary Fig. 2). For each CTD cruise cast, Cant is estimated from in situ bottle-sample measurements of nutrients, oxygen, A_T and pH, via the ϕC_T^0 method (standard procedure¹¹). Emulating the approach proposed in this study for the Argo-O₂ floats (TSO₂-NN procedure), we re-estimated C_{ant} by using the ship-based P/T/S/O₂ data as input to the neural networks to obtain nutrients, DIC, and A_T , and then applied the ϕC_T^{o} method. These C_{ant} profiles (ship-based standard and ship-based TSO₂-NN) were then compared to our Argo-based Cant estimates (Argo-based TSO₂-NN procedure), and all of them agreed well within the averaged method uncertainty of $\pm 5.9 \,\mu\text{mol}\,\text{kg}^{-1}$ (Supplementary Fig. 2). This comparison indicates that the Cant values obtained from ship-based measurements and Cant estimates calculated via Argo-O2 data and neural networks give consistent results.

We also computed a proxy that allows the indirect assessment of the air-sea C_{ant} uptake capacity, the C_{ant} deficit (C_{ant,def}, Supplementary Fig. 3). This proxy was defined by Ridge & McKinley¹² as the difference between the C_{ant} concentration of a water parcel and a reference ($\langle C_{ant} \rangle_{\sigma_{\theta} < 26.5}$) computed as the average C_{ant} concentration between the surface and the potential density (σ_{θ}) of 26.5 kg m⁻³. This isopycnal delimits the subtropical surface waters assumed to be in equilibrium with the current atmosphere. A water parcel with a negative C_{ant,def} value indicates that this parcel has a deficit of C_{ant}. If transported to the surface, this water parcel will have a lower pCO_2 compared to a water parcel without this deficit, hence favoring the air-sea C_{ant} uptake. Usually, C_{ant,def} becomes more and more negative with depth (Supplementary Fig. 3), meaning that the deep waters would uptake a higher amount of C_{ant} compared to the subsurface waters if they were transported to the ocean surface. The reference value is computed for 2002 with the GLODAPv2 dataset⁵⁶, considering the averaged C_{ant} concentration in surface waters ($\sigma_{\theta} < 26.5 \text{ kg m}^{-3}$) offshore the Venezuelan coasts (10–20°N; 70–64°W), and is $\langle C_{ant} \rangle^{2002}_{\sigma_{\theta} < 26.5} = 51.5 \,\mu\text{mol kg}^{-1}$. We assume that our reference water mass offshore the Venezuelan coasts takes 5 years to arrive in the Iceland and Irminger basins⁶⁴. For a given float, we use that reference value rescaled in time to 5 years before the time of its first profile, using the exponential equation of Carter et al.²⁵. Assuming that due to its lagrangian behavior, the float follows the transformation of the same water mass, we kept the same reference value along the float trajectory within the subpolar gyre. However, this assumption is not valid for float 5904988 when it drifted from Region 4 to Region 1 and entered the North Atlantic Current, where we re-initialized the reference value.

Mixed layer depth and water mass definition

To determine the mixed layer depth (MLD) along the Argo float trajectories, we use the threshold method, which is based on a density difference of 0.01 kg m⁻³⁶⁵ between the surface and the base of the mixed layer⁶⁶. Since this method tends to overestimate the MLD during deep convection events in winter⁶⁷, we conducted a visual inspection⁶⁸ to verify the MLD values during these events. When the MLD values disagreed between the threshold method and the visual inspection, we kept the value determined by visual inspection.

The North Atlantic is composed of several water masses (Fig. 1), most of which can be identified by their potential temperature (θ), *S*, σ_{θ} , O₂, and potential vorticity. Here, the potential vorticity is calculated via the Brunt-Väisälä frequency, the gravity, and the Coriolis parameter. In the upper layers, the North Atlantic Current (NAC) flowing from the subtropics to the subpolar regions (Fig. 1), carries North Atlantic Central Water (NACW)³⁰. The lower limit of this water mass is defined by $\sigma_{\theta} = 27.2 \text{ kg m}^{-3.69}$ (Supplementary Fig. 4). Continuous airsea interaction and winter convection along the NAC path cools and freshens the NACW. leading to its transformation into Subpolar Mode Waters (SPMWs)^{31,32}. SPMWs are characterized by nearly homogeneous properties, with S > 34.98, $\sigma_{\theta} < 27.71 \text{ kg/m}^3$, and a potential vorticity lower than $6 \times 10^{-11} \text{ m}^{-1} \text{ s}^{-1} \frac{31,70}{10}$ (Figs. 3 and 4). At intermediate depths (400-2000 dbar), in the subpolar gyre, the Labrador Sea Water (LSW) is located between $27.71 < \sigma_{\theta} < 27.80 \text{ kg/m}^3$. The LSW is formed by deep winter convection in the Labrador Sea⁷¹ and Irminger Sea^{36,65}, trapping large amounts of O₂ and C_{ant} during its formation⁷². In those basins, we defined the LSW by $O_2 \ge 290 \mu \text{mol kg}^{-1.34}$. However, along its path, the LSW loses O₂ due to mixing and biological activity. Therefore the O₂ criteria is not fully reliable to identify the LSW core away from its source regions. As a consequence, outside the Irminger and Labrador Seas we identify the LSW by S < 34.94⁵³. At deeper levels (up to 2000 dbar), the Iceland-Scotland Overflow Water (ISOW) flows along the Reykjanes Ridge and is identified by $\sigma_{\theta} > 27.80 \text{ kg m}^{-3}$ and $S > 34.94^{34}$ (Supplementary Fig. 4).

Eddy identification

To detect the eddies along the Argo-float pathways, we run the autonomous eddy identification scheme of Faghmous et al.⁷³, which monitors mesoscale ocean eddy activity in global sea level anomaly (SLA) dataset (http://www.aviso.altimetry.fr/duacs/). The scheme starts from the simple notion that every eddy has a single extremum, defined as a grid cell whose SLA is higher or lower than its 24 neighbors in a 5×5 grid. To determine the contour of the eddy, the algorithms search for the largest possible contour that would allow the feature not to violate the assumption that an eddy can have only a single extremum within its interior. For each eddy feature identified at time *t*, the eddies at time t + 1 are searched to find the closest feature within a predefined search space. When a feature at time t is associated with another feature at time t + 1, their sizes are compared to ensure that

they are reasonably similar from a physical point of view. There are two major uncertainties associated with this eddy identification method. First, the use of a geometric eddy definition, as opposed to a physical one, introduces uncertainties in the eddy boundaries because they are not necessarily associated with the physical properties of the eddy. Second, as the method constrains features to have a single extremum, small features that are in close proximity might cancel each other out and not be detected. Despite these limitations, the success discovery rate of the autonomous eddy identification algorithm is 96.4%⁷³.

Reporting summary

Further information on research design is available in the Nature Portfolio Reporting Summary linked to this article.

Data availability

The original Argo data can be freely downloaded on the Euro Argo Data Selection platform (https://dataselection.euro-argo.eu/). Our C_{ant} estimates, needed to evaluate the conclusion of the paper, can be downloaded on Zenodo (https://doi.org/10.5281/zenodo.7071614).

Code availability

The Matlab codes used for the simulation are available from the corresponding author upon request.

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Author contributions

R.A., L.I.C., and V.T. designed and developed the concept of the study. R.A. conducted the data analysis with inputs from L.I.C., V.T., H.M., R.B., and F.F.P. R.A. drafted the first version of the paper. All co-authors read and reviewed the paper, and all co-authors agreed on the final version of the paper.

Competing interests

The authors declare no competing interests.

Additional information

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