



HAL
open science

Long time behavior of a degenerate stochastic system modeling the response of a population face to environmental impacts

Pierre Collet, Claire Ecotièrè, Sylvie Méléard

► To cite this version:

Pierre Collet, Claire Ecotièrè, Sylvie Méléard. Long time behavior of a degenerate stochastic system modeling the response of a population face to environmental impacts. 2024. hal-04566978

HAL Id: hal-04566978

<https://hal.science/hal-04566978>

Preprint submitted on 2 May 2024

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Long time behavior of a degenerate stochastic system modeling the response of a population face to environmental impacts

Pierre Collet*, Claire Ecotière†, Sylvie Méléard‡

September 2023

Abstract

We study the asymptotics of a two-dimensional stochastic differential system with a degenerate diffusion matrix. This system describes the dynamics of a population where individuals contribute to the degradation of their environment through two different behaviors. We exploit the almost one-dimensional form of the dynamical system to compute explicitly the Freidlin-Wentzell action functional. That allows to give conditions under which the small noise regime of the invariant measure is concentrated around the equilibrium of the dynamical system having the smallest diffusion coefficient.

Keywords : Stochastic differential system - Large deviations - Small noise regime - Invariant measure.

1 Introduction

Reducing the effects of global change requires the adoption by most of the human population of consistent proenvironmental behaviors. But we, as individuals, tend to act in response to alarming events, and relax when things seem to get better. In [4] Ecotiere et al. have explored and numerically quantified how the tendency to behave inconsistently can be countered by social interactions and social pressure. In that aim, they developed a simple two-dimensional mathematical model describing the coupled dynamics of the perceived environmental state over time $(e_t, t \geq 0)$, and the repartition of a fixed size population between two behaviors A and B more or less active from an environmental perception. Expressing behavior A makes the agent reduce its environmental impact (compared to B). The macroscopic frequency dynamics of behavior A in the population is modeled by the deterministic function $(x_t, t \geq 0)$ taking values in $[0, 1]$. The model is as follows:

$$\begin{aligned} \frac{dx_t}{dt} &= p(x_t, e_t) = \kappa x_t(1 - x_t)(\lambda_A(x_t) - \lambda_B(x_t)) + \tau_A(e_t)(1 - x_t) - \tau_B(e_t)x_t \\ \frac{de_t}{dt} &= h(x_t, e_t) = \ell e_t(l_A x_t + l_B(1 - x_t) - e_t). \end{aligned} \tag{1}$$

The interpretation of the parameters has been carefully commented in [4] and we refer to that paper for more details. The deterministic system (1) is the large population approximation of a stochastic system $((X_t^N, E_t^N), t \geq 0)$, where the scaling parameter is the population size N which is assumed to tend to infinity. In this stochastic setting, the population dynamics is described by a pure jump Markov process whose jump rates model the individual behavior changes and depend on the perceived

*CPHT, Ecole Polytechnique, CNRS, Institut polytechnique de Paris, route de Saclay, 91128 Palaiseau Cedex-France; E-mail: pierre.collet@cphpt.polytechnique.fr

†CREST, ENSAE, CNRS, Institut polytechnique de Paris, 5 Avenue Henry Le Chatelier 91120 Palaiseau, France; E-mail: claire.ecotiere@ensae.fr

‡CMAP, Ecole Polytechnique, CNRS, Institut polytechnique de Paris, Inria, route de Saclay, 91128 Palaiseau Cedex-France and Institut Universitaire de France; E-mail: sylvie.meleard@polytechnique.edu

environmental state. The coupled stochastic system is a piecewise deterministic Markov process where the deterministic environmental component satisfies

$$E_t^N = E_0^N + \int_0^t h(X_s^N, E_s^N) ds.$$

We are interested in the long time behavior of the stochastic system $((X_t^N, E_t^N), t \geq 0)$, and also how it is related, for large N to the long time behavior of the deterministic system (1). To that goal, we have considered the approximation-diffusion of the process (see Ethier-Kurz p.354 for a precise definition) which consists in a two-dimensional process, solution of the stochastic differential system

$$\begin{cases} dX_t^N = P(X_t^N, E_t^N)dt + \frac{\sigma(X_t^N, E_t^N)}{\sqrt{2N}}dB_t \\ dE_t^N = h(X_t^N, E_t^N)dt, \end{cases} \quad (2)$$

where B is a Brownian motion and

$$\sigma^2(x, e) = \kappa x(1-x)(\lambda_A(x) + \lambda_B(x)) + \tau_A(e)(1-x) + \tau_B(e)x.$$

System (2) depends on so many parameters that we needed to simplify it by successive parametrizations. We have proven in Appendix 1 that, under some assumptions on the parameters and using different parametrizations and changes of variables, the system can be simplified in the following form:

$$\begin{cases} dX_t^\varepsilon = \lambda_1 X_t^\varepsilon (1 - (X_t^\varepsilon)^2) dt + \varepsilon \cos(\theta) (\sigma_0 + \sigma_1 X_t^\varepsilon) dB_t \\ dY_t^\varepsilon = \left(-\lambda_2 Y_t^\varepsilon + \lambda_3 X_t^\varepsilon (1 - (X_t^\varepsilon)^2) \right) dt + \varepsilon \sin(\theta) (\sigma_0 + \sigma_1 X_t^\varepsilon) dB_t, \end{cases} \quad (3)$$

where B is a standard Brownian motion and ε is a small parameter (ε of the order of $1/\sqrt{N}$). Note that this system is degenerate for two reasons: first, the two coordinates are driven by the same Brownian motion, and second, the diffusion coefficient in front of the Brownian motion cancels on the line $x = -\sigma_0/\sigma_1$; all this leading to a non-invertible diffusion matrix.

We are interested in the long time behaviour of the process $(X^\varepsilon, Y^\varepsilon)$ when ε is small. The limits ε tending to 0 and t tending to infinity a priori don't commute and one can ask about i) the long time behavior of the process $(X^\varepsilon, Y^\varepsilon)$ for fixed ε , ii) its limit when ε tends to 0, iii) the relation with the long time behavior of the deterministic system

$$\begin{cases} dx_t = b_1(x_t) = \lambda_1 x_t (1 - (x_t)^2) dt \\ dy_t = b_2(x_t, y_t) = \left(-\lambda_2 y_t + \lambda_3 x_t (1 - (x_t)^2) \right) dt, \end{cases} \quad (4)$$

obtained from (3) when $\varepsilon = 0$.

In Section 2, we state the different results answering to these questions. It's standard to prove the convergence of the stochastic differential system (3) when ε tends to 0, towards the deterministic system (4). The latter admits three equilibrium points: two stable points framing a saddle point. For a fixed ε , we show the existence, uniqueness and exponential convergence, of an invariant probability measure π^ε , absolutely continuous with respect to the Lebesgue measure. We also show that the sequence $(\pi^\varepsilon)_\varepsilon$ converges when ε tends to 0 to a linear combination of the Dirac measures on the stable equilibrium points of the deterministic system. In particular, we show that the sign of σ_1 determines on which of the two stable equilibria, the limit invariant measure π is concentrated.

The study of invariant measures appears in various fields, ranging from potential wells in physics [7, 8] to neural networks ([10, 9]). For non-degenerate stochastic differential systems, we refer to the book by [6] which details the steps to show that the invariant measure of the stochastic differential system is concentrated outside the unstable equilibrium points of the perturbed deterministic system such that col nodes and sources. This method is based on control theory (see [13]) and the calculation of the energy necessary to move from a compact containing one equilibrium point to another one. The energy translated by the action functional involves the inverse of the diffusion matrix of the stochastic

differential system. In our degenerate case, it is not possible to apply these theorems. We will therefore seek to extend these results to our case study. For this, we will use the particular form of our system which has a main component on its two dimensions.

In Section 3 and 4, we respectively prove the existence and uniqueness of the system's limit when t tends to infinity, then we study the limit of the invariant measures π^ε when ε tends to 0.

2 Main results

Let us state the main results of the paper. Let ε_0 positive and fixed through the paper, and we consider $\varepsilon \in]0, \varepsilon_0[$. Recall that the process $(X^\varepsilon, Y^\varepsilon)$ is solution of (3) with B a standard (unidimensional) Brownian motion, issued from $(x_0, y_0) \in \mathbb{R}^2$. The diffusion coefficient for both coordinates is proportional to

$$\sigma_0 + \sigma_1 x, \quad (5)$$

where $\sigma_0 \neq 0$ and $\sigma_1 \in \mathbb{R}$. Let us assume without restriction that $\sigma_0 > 0$. We also assume that the parameters appearing in (3) satisfy the following hypotheses.

Hypothesis 2.1. 1. $\sigma_1 \in]-\sigma_0, \sigma_0[$,

2. $\lambda_1, \lambda_2, \lambda_3$ are positive constants,

3. $\theta \in]-\pi, \pi[\setminus \left\{ \pm \frac{\pi}{2}, \arctan \left(\frac{2\lambda_3}{2\lambda_1 - \lambda_2} \right) \right\}$.

Note that under Assumption 2.1.1, the diffusion coefficient doesn't cancel between the lines $x = -1$ and $x = 1$.

Let us introduce the infinitesimal generator associated with the solution of (3). For $f \in \mathcal{C}^2(\mathbb{R}^2, \mathbb{R})$ and $(x, y) \in \mathbb{R}^2$, we define

$$\begin{aligned} \mathcal{L}^\varepsilon f(x, y) &= \lambda_1 x(1-x^2) \frac{\partial f}{\partial x} - (\lambda_2 y - \lambda_3 x(1-x^2)) \frac{\partial f}{\partial y} + \frac{\varepsilon^2 (\sigma_0 + \sigma_1 x)^2}{2} \left(\cos^2(\theta) \frac{\partial^2 f}{\partial x^2} + \sin^2(\theta) \frac{\partial^2 f}{\partial y^2} \right) \\ &\quad + \varepsilon^2 (\sigma_0 + \sigma_1 x)^2 \sin(\theta) \cos(\theta) \frac{\partial^2 f}{\partial x \partial y}. \end{aligned} \quad (6)$$

Lemma 2.2 (Lyapunov condition). *Under Hypothesis 2.1, there exist ε_0 and $\alpha > 0$ such that the infinitesimal generator \mathcal{L}^ε satisfies a Lyapunov condition for $W(x, y) = 1 + x^4 + \alpha y^2$, uniformly for $\varepsilon \leq \varepsilon_0$: there exist positive constants α_1 and α_2 , independent of $\varepsilon \leq \varepsilon_0$ such that*

$$\mathcal{L}^\varepsilon W \leq \alpha_1 - \alpha_2 W. \quad (7)$$

Proof. For all $0 < \varepsilon \leq \varepsilon_0$, we have

$$\mathcal{L}^\varepsilon W(x, y) = 4\lambda_1 x^4(1-x^2) - 2\alpha y(\lambda_2 y - \lambda_3 x(1-x^2)) + \varepsilon^2 (\sigma_0 + \sigma_1 x)^2 (6x^2 \cos^2(\theta) + \alpha \sin^2(\theta)).$$

The dominant term is

$$\begin{aligned} &-4\lambda_1 x^6 - 2\alpha \lambda_2 y^2 - 2\alpha \lambda_3 y x^3 \\ &= -2x^6 \left(2\lambda_1 + \alpha \lambda_2 \left(\frac{y}{x^3} \right)^2 + \alpha \lambda_3 \left(\frac{y}{x^3} \right) \right). \end{aligned}$$

We can observe that under Hypothesis 2.1.2, the polynomial $t \rightarrow 2\lambda_1 + \alpha \lambda_2 t^2 + \alpha \lambda_3 t$ will be always positive if and only if its discriminant is negative, i.e. for any α satisfying

$$\alpha \leq 8 \frac{\lambda_2 \lambda_1}{\lambda_3^2}.$$

Therefore, under Hypothesis 2.1.2 and for $\varepsilon \leq \varepsilon_0$, one can conclude that for $\alpha \in]0, 8\lambda_1 \lambda_2 / \lambda_3^2[$,

$$\mathcal{L}^\varepsilon W \leq \alpha_1 - \alpha_2 W,$$

for α_1 and α_2 (independent of ε), conveniently chosen. \square

It is then standard, using (7), to prove that for any $T > 0$ and any fixed $\varepsilon > 0$ and if $\mathbb{E}(W(X_0^\varepsilon, Y_0^\varepsilon)) < +\infty$, there exists a unique solution $((X_t^\varepsilon, Y_t^\varepsilon), t \in [0, T])$ of (3) in the space of continuous processes with $\mathbb{E}(\sup_{t \leq T} W(X_t^\varepsilon, Y_t^\varepsilon)) < +\infty$. Moreover (7) also implies that if $\sup_{\varepsilon \leq \varepsilon_0} \mathbb{E}(W(X_0^\varepsilon, Y_0^\varepsilon)) < +\infty$, then $\sup_{\varepsilon \leq \varepsilon_0} \mathbb{E}(\sup_{t \leq T} W(X_t^\varepsilon, Y_t^\varepsilon)) < +\infty$. In all what follows, we will assume that

$$\sup_{\varepsilon \leq \varepsilon_0} \mathbb{E}(W(X_0^\varepsilon, Y_0^\varepsilon)) < +\infty$$

and we will denote by $(P_t^\varepsilon)_{t \geq 0}$ the semigroup associated with $((X_t^\varepsilon, Y_t^\varepsilon), t \in \mathbb{R}_+)$.

Assuming in addition that $(X_0^\varepsilon, Y_0^\varepsilon)$ converges in probability to the deterministic point (x_0, y_0) , classical arguments allow to show that for any $T > 0$, the processes $((X_t^\varepsilon, Y_t^\varepsilon), t \in [0, T])$ converge in probability to the solution $((x_t, y_t), t \in [0, T])$ of the deterministic system (4) issued from (x_0, y_0) , when ε tends to 0. System (4) admits three equilibria defined by

$$z_1 = (-1, 0), \quad z_2 = (0, 0) \text{ and } z_3 = (1, 0). \quad (8)$$

Let us note that z_2 is a saddle-node and z_1 and z_3 are attractive points. The basin of attraction of z_1 (resp. z_3) is the open left (resp. right) half plane. The basin of attraction of z_2 is the y -axis.

Making ε tend to 0 before time t tend to infinity yields

$$\lim_{t \rightarrow +\infty} \lim_{\varepsilon \rightarrow 0} (X_t^\varepsilon, Y_t^\varepsilon) = z_1 \mathbb{1}_{\{x_0 < 0\}} + z_2 \mathbb{1}_{\{x_0 = 0\}} + z_3 \mathbb{1}_{\{x_0 > 0\}}. \quad (9)$$

We are now interested in the limit $t \rightarrow +\infty$ first, when ε is fixed. In that aim, let us associate with the Lyapounov function W the norm $\|\cdot\|_W$ defined on the set of finite signed measures by

$$\|\mu\|_W = \int_{\mathbb{R}^2} W(x, y) |\mu|(dx, dy), \quad (10)$$

μ being a finite signed measure on \mathbb{R}^2 .

We can state the two main results of the paper.

Theorem 2.3. *Let $\varepsilon \in]0, \varepsilon_0[$ fixed. Under Hypothesis 2.1, the process $(X_t^\varepsilon, Y_t^\varepsilon)_{t \geq 0}$ solution of (3) admits a unique invariant probability measure π^ε with Lesbeque density. There exist a constant $C > 0$ and $\gamma > 0$ such that for any $(x, y) \in \mathbb{R}^2$,*

$$\|P_t^\varepsilon((x, y), \cdot) - \pi^\varepsilon(\cdot)\|_W \leq CW(x, y)e^{-\gamma t}. \quad (11)$$

The proof of Theorem 2.3 will be postponed to Section 3.

Theorem 2.4. *When ε tends to 0, the sequence $(\pi^\varepsilon)_{\varepsilon < \varepsilon_0}$ converges to the measure π defined by*

$$\pi = \begin{cases} \delta_{(1,0)} & \text{if } \sigma_1 < 0, \\ \frac{1}{2}\delta_{(1,0)} + \frac{1}{2}\delta_{(-1,0)} & \text{if } \sigma_1 = 0, \\ \delta_{(-1,0)} & \text{if } \sigma_1 > 0. \end{cases} \quad (12)$$

This result shows that the most stable equilibrium point is the one with the associated diffusion coefficient closest to zero.

The proof of Theorem 2.4 will be postponed to Section 4.

3 Proof of Theorem 2.3

The proof of Theorem 2.3 is based on Proposition 6.32, Theorems 6.34 and Theorem 8.15 in [1] (see also Corollary 7.2 in [12]). It will consist in proving first (Lemma 3.1) that the stable equilibria z_1 and z_3 of (4) satisfy the weak Hörmander condition and second (Lemma 3.2) that these two points are accessible. As explained in the introduction, the difficulty in this proof comes from the degeneracy of the stochasticity.

Lemma 3.1. *Under Hypothesis 2.1, the stable equilibria points z_1 and z_3 defined by (8) satisfy the weak Hörmander condition.*

Proof. Let us introduce the vector fields $F_0 = \lambda_1 x(1 - x^2)\partial_x + (\lambda_3 x(1 - x^2) - \lambda_2 y)\partial_y$ and $F_1 = \varepsilon(\sigma_0 + \sigma_1 x)(\cos(\theta)\partial_x + \sin(\theta)\partial_y)$. The infinitesimal generator (6) can be written as :

$$\mathcal{L}^\varepsilon = F_0 + \frac{1}{2}F_1^2, \quad (13)$$

where $F_1^2 = \varepsilon^2(\sigma_0 + \sigma_1 x)^2(\cos^2(\theta)\partial_{xx} + \sin^2(\theta)\partial_{yy} + 2\cos(\theta)\sin(\theta)\partial_{xy})$.

We introduce the sequence of vector fields (\mathcal{A}_k) , $k \geq 0$, defined by $\mathcal{A}_0 = \{F_0, F_1\}$ and $\mathcal{A}_{k+1} = \mathcal{A}_k \cup \{[F_i, G], G \in \mathcal{A}_k, i = 0, 1\}$ where $[U, V]$ is the Lie's bracket of U and V . We denote by $\mathcal{A}_k(x, y)$ the vector space spanned by the vector fields of \mathcal{A}_k for $(x, y) \in \mathbb{R}^2$.

As can be seen in [1, Section 6.5.2] (see also [11]), a point $(x, y) \in \mathbb{R}^2$ satisfies the weak Hörmander's condition if there exists $k > 0$ such that the vectoriel space $\mathcal{A}_k(x, y)$ spans \mathbb{R}^2 .

For $k = 0$, $\mathcal{A}_0(z_1)$ and $\mathcal{A}_0(z_3)$ only span \mathbb{R} , since F_0 is nul in this case. For $k = 1$, we compute

$$\begin{aligned} [F_0, F_1](x, y) &= F_0(F_1) - F_1(F_0) \\ &= \varepsilon \cos(\theta) \left(\lambda_1(\sigma_0 + \sigma_1 x)(3x^2 - 1) + \lambda_1 \sigma_1 x(1 - x^2) \right) \partial_x \\ &\quad + \varepsilon \left[\cos(\theta)\lambda_3(\sigma_0 + \sigma_1 x)(3x^2 - 1) + \sin(\theta) \left((\sigma_0 + \sigma_1 x)\lambda_2 + \lambda_1 \sigma_1 x(1 - x^2) \right) \right] \partial_y. \end{aligned}$$

Then

$$[F_0, F_1](z_1) = \varepsilon(\sigma_0 - \sigma_1) \left(2\lambda_1 \cos(\theta)\partial_x + (2\lambda_3 \cos(\theta) + \lambda_2 \sin(\theta))\partial_y \right),$$

and $[F_0, F_1](z_3)$ has a similar value with $(\sigma_0 - \sigma_1)$ replaced by $(\sigma_0 + \sigma_1)$. Let us study the colinearity of F_1 and $[F_0, F_1]$ at z_1 . The matrix

$$\begin{pmatrix} 2\lambda_1 \cos(\theta) & \cos(\theta) \\ 2\lambda_3 \cos(\theta) + \lambda_2 \sin(\theta) & \sin(\theta) \end{pmatrix} \quad (14)$$

has the determinant $\cos(\theta)((2\lambda_1 - \lambda_2)\sin(\theta) - 2\lambda_3 \cos(\theta))$ which is non zero under Hypothesis 2.1.3. Therefore, $\mathcal{A}_1(z_1)$ spans \mathbb{R}^2 and the same property holds for z_3 . Thus the two stable equilibria z_1 and z_3 satisfy the weak Hörmander condition. \square

Let us now prove their accessibility.

Lemma 3.2. *Under Hypothesis 2.1, the stable equilibria z_1 and z_3 are accessible for (P_t^ε) from any $(x_0, y_0) \neq (0, 0)$.*

The proof will be based on Proposition 6.32 in [1]. In that aim, we introduce the control system associated to the solution $(X_t^\varepsilon, Y_t^\varepsilon)$ of (3), defined by:

$$\begin{cases} dx_t^\varphi = \lambda_1 x_t^\varphi(1 - (x_t^\varphi)^2)dt + \varepsilon(\sigma_0 + \sigma_1 x) \cos(\theta)\varphi(t)dt \\ dy_t^\varphi = -\lambda_2 y_t^\varphi + \lambda_3 x_t^\varphi(1 - (x_t^\varphi)^2)dt + \varepsilon(\sigma_0 + \sigma_1 x) \sin(\theta)\varphi(t)dt, \end{cases} \quad (15)$$

where the control function φ is a piecewise continuous function defined on \mathbb{R}_+ .

Proof. Using Proposition 6.32 from [1], a point $z \in \mathbb{R}^2$ is accessible if and only if for every neighborhood U of z , there exist a control function φ and a time $t \geq 0$ such that $(x_t^\varphi, y_t^\varphi) \in U$. We will detail the proof for one of the two equilibria. Let's show that z_1 is accessible. If the initial condition is taken in z_1 's basin of attraction, i.e. $x_0 < 0$, then taking $\varphi \equiv 0$ makes (x^φ, y^φ) follow the flow and converge towards z_1 .

If the initial condition is taken in z_3 's basin of attraction, i.e. $x_0 > 0$, then similarly the flow (4) converges towards z_3 . Thus we can assume without loss of generality that $x_0 = 1 + \delta$, where $\delta \neq 0$ will be conveniently chosen. We are looking for a control function φ allowing the flow (15) to attain the attraction basin of z_1 . Since by Hypothesis 2.1, $|\sigma_1| < \sigma_0$, there exists $\delta \neq 0$ such that for any $x \in [-1/2; 1 + \delta]$, $|\sigma_0 + \sigma_1 x| > 0$. We fix the parameter δ and choose a real value k such that for any $x \in [-1/2; 1 + \delta]$,

$$\lambda_1 x(1 - x^2) - k\varepsilon(\sigma_0 + \sigma_1 x) \cos(\theta) < -1.$$

For example, any

$$k > \frac{1 + \lambda_1}{(\sigma_0 - |\sigma_1|(1 + \delta)) \cos \theta \varepsilon}$$

is suitable. Taking the constant control $\varphi(t) = -k$, one obtains that starting from $(1 + \delta, y)$ (for $y \in \mathbb{R}$), $dx_t^\varphi < 0$ as soon as $x(t) \in [-1/2; 1 + \delta]$. Therefore after a certain amount of time, $x(t)$ will become less than $-1/2$. In particular, there exists $t_1 > 0$ such that $(x_{t_1}^\varphi, y_{t_1}^\varphi) \in \mathbb{R}_*^- \times \mathbb{R}$. Once the dynamics is in z_1 's basin of attraction, the flow will make it converge to z_1 .

If $x_0 = 0$ and $y_0 \neq 0$, then the dynamics of (1) (i.e. $\varphi = 0$) will lead the system to be in the open half plan $\mathbb{R}_*^+ \times \mathbb{R}$ or $\mathbb{R}_*^- \times \mathbb{R}$. We are back to one of the previous situations. □

Proof of the Theorem 2.3. Using the lemmas 3.2, 3.1 and 2.2, we can use the Theorems 6.34 and 8.15 from [1], which concludes the proof. □

4 Limit and support of the invariant probability measure when $\varepsilon \rightarrow 0$

In this section, we explain the different steps of the proof of Theorem 2.4. First, we show the tightness of $(\pi^\varepsilon)_{\varepsilon \leq \varepsilon_0}$ and we identify the limiting values of $(\pi^\varepsilon)_{\varepsilon \leq \varepsilon_0}$. In a second time, we adapt the work of Freidlin and Wentzel in [6] in our degeneracy case to prove the uniqueness of the limit and to characterize it.

4.1 Existence and characterization of the limiting values of $(\pi^\varepsilon)_{\varepsilon \leq \varepsilon_0}$

Lemma 4.1. *The sequence $(\pi^\varepsilon)_{\varepsilon \leq \varepsilon_0}$ is tight.*

Proof. Let us denote by $\mathfrak{B}(R)$ the ball of \mathbb{R}^2 centered at 0 with radius R . In order to show that the sequence $(\pi^\varepsilon)_{\varepsilon \leq \varepsilon_0}$ is tight, we introduce the occupation measure r^ε defined for any $(x_0, y_0) \in \mathbb{R}^2$ and $t \in \mathbb{R}_+$ and A a Borel subset of \mathbb{R}^2 by

$$r^\varepsilon(t, (x_0, y_0), A) = \frac{1}{t} \int_0^t P_s^\varepsilon(\mathbb{1}_A)(x_0, y_0) ds = \frac{1}{t} \int_0^t \mathbb{E}_{(x_0, y_0)}(\mathbb{1}_A(X_s^\varepsilon, Y_s^\varepsilon)) ds.$$

Then

$$r^\varepsilon(t, (x_0, y_0), \mathbb{R}^2 \setminus \mathfrak{B}(R)) = \frac{1}{t} \int_0^t \mathbb{E}_{(x_0, y_0)}(\mathbb{1}_{\{\mathbb{R}^2 \setminus \mathfrak{B}(R)\}}(X_s^\varepsilon, Y_s^\varepsilon)) ds. \quad (16)$$

Let us prove that $\forall \delta > 0, \exists R > 0$ such that $\forall t > 0, \forall \varepsilon \in]0, \varepsilon_0]$,

$$r^\varepsilon(t, (x_0, y_0), \mathbb{R}^2 \setminus \mathfrak{B}(R)) \leq \delta.$$

We introduce for any $\varepsilon \leq \varepsilon_0$ the sequence of stopping times $(\tau_n^\varepsilon)_n$ defined by $\tau_n^\varepsilon = \inf\{t \geq 0 : \|(X_t^\varepsilon, Y_t^\varepsilon)\| \geq n\}$. The initial condition of the process being the fixed vector (x_0, y_0) , we have seen

that Section 2 that $\sup_{\varepsilon \leq \varepsilon_0} \mathbb{E}(\sup_{t \leq T} W(X_t^\varepsilon, Y_t^\varepsilon)) < +\infty$. It is then easy to prove that for any ε , the sequence $(\tau_n^\varepsilon)_n$ tends almost surely to infinity. We have that

$$\mathbb{E}_{(x_0, y_0)} \left(W(X_{t \wedge \tau_n}^\varepsilon, Y_{t \wedge \tau_n}^\varepsilon) \right) = W(x_0, y_0) + \mathbb{E}_{(x_0, y_0)} \left(\int_0^{t \wedge \tau_n} \mathcal{L}^\varepsilon W(X_s^\varepsilon, Y_s^\varepsilon) ds \right). \quad (17)$$

Note that the function $(x, y) \in \mathbb{R}^2 \rightarrow W(x, y)$ is continuous and lower bounded by $m_R > 0$ on the complementary of the open ball of radius R . Note also that m_R tends to infinity with R . Therefore, using Lemma 2.2, we obtain

$$\mathbb{E}_{(x_0, y_0)} \left(W(X_{t \wedge \tau_n}^\varepsilon, Y_{t \wedge \tau_n}^\varepsilon) \right) \leq W(x_0, y_0) + \alpha_1 \mathbb{E}_{(x_0, y_0)}(t \wedge \tau_n) - \alpha_2 m_R \mathbb{E}_{(x_0, y_0)} \left(\int_0^{t \wedge \tau_n} \mathbb{1}_{\{\|(X_s^\varepsilon, Y_s^\varepsilon)\| > R\}} ds \right),$$

from which we deduce that

$$\alpha_2 m_R \mathbb{E}_{(x_0, y_0)} \left(\int_0^{t \wedge \tau_n} \mathbb{1}_{\{\|(X_s^\varepsilon, Y_s^\varepsilon)\| > R\}} ds \right) \leq W(x_0, y_0) + \alpha_1 \mathbb{E}_{(x_0, y_0)}(t \wedge \tau_n),$$

since W is positive. Finally making n tend to infinity, we obtain that

$$r^\varepsilon \left(t, (x_0, y_0), \mathbb{R}^2 \setminus \mathfrak{B}(R) \right) \leq \frac{t\alpha_1 + W(x_0, y_0)}{t\alpha_2 m_R}. \quad (18)$$

Using Theorem 2.3, we have

$$\pi^\varepsilon(\mathbb{R}^2 \setminus \mathfrak{B}(R)) \leq r^\varepsilon \left(t, (x_0, y_0), \mathbb{R}^2 \setminus \mathfrak{B}(R) \right) + CW(x_0, y_0) \frac{1}{\gamma t} (1 - e^{-\gamma t}).$$

The rhs term can be made as small as required for t and R large enough, uniformly in ε . Therefore the sequence $(\pi^\varepsilon)_{\varepsilon \leq \varepsilon_0}$ is tight. □

Lemma 4.2. *The limiting values of $(\pi^\varepsilon)_{\varepsilon \leq \varepsilon_0}$ have the form*

$$\sum_{i=1}^3 a_i \delta_{z_i}, \quad (19)$$

where $z_i, i \in \{1, 2, 3\}$ are defined in (8) and the coefficients $a_i \geq 0$ satisfy $\sum_{i=1}^3 a_i = 1$.

Proof. We introduce the operator \mathcal{L}^0 defined for all $f \in \mathcal{C}_c^\infty(\mathbb{R}^2)$ by

$$\mathcal{L}^0 f(x, y) = \lambda_1 x(1 - x^2) \frac{\partial f}{\partial x}(x, y) - \left(\lambda_2 y - \lambda_3 x(1 - x^2) \right) \frac{\partial f}{\partial y}(x, y). \quad (20)$$

Let us consider an accumulation point π of the sequence $(\pi^\varepsilon)_{\varepsilon \leq \varepsilon_0}$. We will show that

$$\int_{\mathbb{R}^2} \mathcal{L}^0 f(x, y) \pi(dx, dy) = 0, \quad \forall f \in \mathcal{C}_c^\infty(\mathbb{R}^2). \quad (21)$$

Let be $f \in \mathcal{C}_c^\infty(\mathbb{R}^2)$. We decompose $\int_{\mathbb{R}^2} \mathcal{L}^0 f(x, y) \pi(dx, dy)$ as follows :

$$\begin{aligned} \int_{\mathbb{R}^2} \mathcal{L}^0 f(x, y) \pi(dx, dy) &= \int_{\mathbb{R}^2} \mathcal{L}^0 f(x, y) (\pi(dx, dy) - \pi^{\varepsilon_k}(dx, dy)) \\ &\quad + \int_{\mathbb{R}^2} \mathcal{L}^{\varepsilon_k} f(x, y) \pi^{\varepsilon_k}(dx, dy) \\ &\quad + \int_{\mathbb{R}^2} (\mathcal{L}^0 - \mathcal{L}^{\varepsilon_k}) f(x, y) \pi^{\varepsilon_k}(dx, dy). \end{aligned}$$

The sequence $(\pi^{\varepsilon_k})_{\varepsilon_k < \varepsilon_0}$ converges weakly to π and the first term in the rhs of previous equation tends to 0 when $\varepsilon_k \xrightarrow[k \rightarrow +\infty]{} 0$. The second term is null as π^{ε_k} is an invariant measure for the diffusion operator $\mathcal{L}^{\varepsilon_k}$.

Using a majoration of the third term, we obtain

$$\begin{aligned} \left| \int_{\mathbb{R}^2} \left(\mathcal{L}^{\varepsilon_k} f(x, y) - \mathcal{L}^0 f(x, y) \right) \pi^{\varepsilon_k}(dx, dy) \right| &\leq \frac{\varepsilon_k^2}{2} \int_{\mathbb{R}^2} (\sigma_0 + \sigma_1 x)^2 \left| \frac{\partial^2 f(x, y)}{\partial x^2} + \frac{\partial^2 f(x, y)}{\partial y^2} \right| \pi^{\varepsilon_k}(dx, dy) \\ &\leq C_f \varepsilon_k^2, \end{aligned}$$

where $C_f = \frac{1}{2} \sup_{(x, y) \in \mathbb{R}^2} \left((\sigma_0 + \sigma_1 x)^2 \left| \frac{\partial^2 f(x, y)}{\partial x^2} + \frac{\partial^2 f(x, y)}{\partial y^2} \right| \right) < +\infty$ since f has a compact support.

That concludes the proof of (21).

The operator \mathcal{L}^0 defined by (20) is associated with the deterministic system (1). We have for all $f \in \mathcal{C}_c^\infty(\mathbb{R}^2)$

$$f(x_t, y_t) = f(x_0, y_0) + \int_0^t \mathcal{L}^0 f(x_s, y_s) ds,$$

where $(x_0, y_0) \in \mathbb{R}^2$ is the initial condition.

Integrating the initial condition with respect to the measure π , we obtain

$$\int_{\mathbb{R}^2} f(x_t, y_t) \pi(dx_0, dy_0) = \int_{\mathbb{R}^2} f(x_0, y_0) \pi(dx_0, dy_0) + \int_{\mathbb{R}^2} \int_0^t \mathcal{L}^0 f(x_s, y_s) ds \pi(dx_0, dy_0),$$

where the last term vanishes because π is an invariant measure.

Taking the limit when $t \rightarrow +\infty$, using the different basins of attraction and Lebesgue's theorem, we obtain that

$$\begin{aligned} &\int_{\mathbb{R}^2} f(x_0, y_0) \pi(dx_0, dy_0) \\ &= \int_{\mathbb{R}^2} f \left(z_1 \mathbb{1}_{(x_0 < 0)} + z_2 \mathbb{1}_{(x_0 = 0)} + z_3 \mathbb{1}_{(x_0 > 0)} \right) \pi(dx_0, dy_0) \\ &= f(z_1) \pi(\mathbb{R}_- \times \mathbb{R}) + f(z_2) \pi(\{0\} \times \mathbb{R}) + f(z_3) \pi(\mathbb{R}_+ \times \mathbb{R}), \end{aligned}$$

which concludes the proof. \square

4.2 Uniqueness and characterization of the limiting values of $(\pi^\varepsilon)_{\varepsilon \leq \varepsilon_0}$

We will adapt in our degeneracy framework the method introduced in [6]. Here the diffusion matrix of System (3) is not invertible. Our strategy will be uniquely based on the first component of the system. Indeed the latter is the principal component of the stochastic differential system (3), the drift of the second component pushing it toward 0. The equilibria of the associated dynamical system (4) are aligned on the horizontal axis and the cost of trajectories will be reduced to the movement on the horizontal axis, allowing a reduction of the two-dimensional problem to a one-dimensional study.

The diffusion matrix is given by

$$\sigma(x, y) = (\sigma_0 + \sigma_1 x) \begin{pmatrix} \cos \theta & 0 \\ \sin \theta & 0 \end{pmatrix}.$$

Using Theorem 5.6.7 in [3] and the specific form of the drift (b_1, b_2) , we obtain that the action functional is given by

$$\begin{aligned} I_{(x_0, y_0)}(f_1, f_2) = &\inf \left\{ \int_0^T |\dot{g}(s)|^2 ds ; g_1 \in H_1, f_1(t) = x_0 + \int_0^t b_1(f_1(s)) ds + \int_0^t \sigma_{11}(f_1(s)) g_1(s) ds \text{ and} \right. \\ &\left. f_2(t) = y_0 + \int_0^t b_2(f_1(s), f_2(s)) ds + \int_0^t \sigma_{21}(f_1(s)) g_1(s) ds \right\}, \end{aligned}$$

and $I_{(x_0, y_0)}(f_1, f_2) = +\infty$ otherwise. This minimization problem is solved by considering first the constraint on the first coordinate, and then by seeing if the second constraint is satisfied or not. Minimization under the first constraint is known to be equivalent to minimize the Freidlin-Wentzell functional (see [3], remark p. 214). If the second constraint is not satisfied for the minimizer, the action functional will be taken to be $+\infty$.

Therefore, as seen in [6] (see also [3] Theorem 5.6.12), the behavior of probabilities of large deviations from the most "probable" trajectory in x , i.e. the trajectory of the dynamical system (4), can be described thanks to the action functional S_{0T}^σ acting on smooth functions w from $[0, T]$ to \mathbb{R} , and defined for $\varepsilon \in]0, \varepsilon_0]$ and $\theta \in]-\pi, \pi[\setminus \{\pm \frac{\pi}{2}\}$ by

$$S_{0T}^\sigma((w_t)_{t \geq 0}) := \frac{1}{(\varepsilon \cos(\theta))^2} \int_0^T L^\sigma(w_s, \dot{w}_s) ds, \quad (22)$$

where the lagrangien $L^\sigma(t, w_t, \dot{w}_t)$ is defined as

$$L^\sigma(w_t, \dot{w}_t) = \frac{1}{2} \left(\frac{\dot{w}_t - \lambda_1 w_t (1 - w_t^2)}{\sigma_0 + \sigma_1 w_t} \right)^2.$$

Let introduce a parameter $\delta \in]0, 1/2[$. We consider three disjoint subsets of \mathbb{R}^2 , K_1^δ , K_2^δ and K_3^δ defined by

$$K_i^\delta = \{(x, y) \in \mathbb{R}^2, |x - x_i|^2 \leq \delta^2\}, \quad (23)$$

where $x_i, i \in \{1, 2, 3\}$ is the first component of z_i defined in (8). Note that these sets are bands and not compact as in [6], however their basis in x are compact intervals. Due to the particular form of the system (3), entrances and exits in these bands are described by the x coordinate only which obeys a stochastic differential equation independent of y . Let us fix $T > 0$.

An essential role will be played by the function characterizing the difficulty of passage from K_i^δ to a K_j^δ , defined by

$$V_{ij}^\sigma := \inf \left\{ S_{0T}^\sigma((w_t)_{t \geq 0}); (w_0, 0) \in K_i^\delta, (w_T, 0) \in K_j^\delta \right\}.$$

The global passage cost to K_j^δ is the quantity

$$W^\sigma(K_j^\delta) = \min_{g \in G(i)} \sum_{(m \rightarrow n) \in g} V_{mn}^\sigma,$$

where for $i \in \{1, 2, 3\}$, $G(i)$ is the set of graphes from $\{1, 2, 3\} \setminus \{i\}$ to $\{1, 2, 3\}$ without cycles and such that any point of $\{1, 2, 3\} \setminus \{i\}$ is the initial point of exactly one arrow. This quantity will play a main role in the following.

Let us also introduce the cost matrix $P^\sigma = ((V_{ij}^\sigma))_{i, j \in \{1, 2, 3\}}$. Our aim is to compute this matrix or at least to know which coefficients are positive. Indeed, if for any $j \neq i$, $V_{ij}^\sigma > 0$, the set K_i^δ is said stable. In the contrary, this set is unstable. These properties will be related to the long time behavior of the flow.

We recall the well known following property concerning the extrema of the action functional.

Proposition 4.3. *Let us assume that $\sigma_1 = 0$. The extrema of the functional $S_{0T}^{\sigma_0}$ are attained by the two flows solution of*

$$\dot{w}_t = \pm \lambda_1 w_t (1 - w_t^2).$$

When the sign is $+$, the extremum is attained by the solution of the dynamical system (4) and when the sign is $-$, the dynamics realizing the extremum has to go against the flow.

Proof. Let us give an elementary proof (see also Theorem 3.1 of Chapter 4 in [6]). Denoting $F(w) = \lambda_1 w(1 - w^2)$, we have

$$\int_0^T (\dot{w}_t - F(w_t))^2 dt = \int_0^T (\dot{w}_t + F(w_t))^2 dt - 4 \int_0^T \dot{w}_t F(w_t) dt = \int_0^T (\dot{w}_t + F(w_t))^2 dt - 4(V(w_T) - V(w_0)),$$

where V is a primitive of F . The result follows. \square

Proposition 4.4. *Let us assume that $\sigma_1 = 0$. The solution of $\dot{w}_t = -\lambda_1 w_t(1 - (w_t)^2)$ is given by x^φ , with*

$$\varphi(t) = -\frac{2\lambda_1 w_0(1 - w_0^2)e^{-\lambda_1 t}}{\varepsilon\sigma_0 \cos(\theta) \left(1 - w_0^2 + w_0^2 e^{-2\lambda_1 t}\right)^{3/2}}.$$

We also have in this case

$$S_{0T}^{\sigma_0}(w) = \frac{\lambda_1^2(1 - w_0^2)^2}{2(\varepsilon\sigma_0 \cos(\theta))^2} \left(\frac{1}{(1 - w_0^2 + w_0^2 e^{-2\lambda_1 T})^2} - 1 \right).$$

Proof. A direct computation gives that

$$w_t = w_0 \frac{e^{-\lambda_1 t}}{\sqrt{1 - w_0^2 + w_0^2 e^{-2\lambda_1 t}}}$$

and then

$$\dot{w}_t = -\lambda_1 w_t(1 - w_t^2) = -\lambda_1 w_0(1 - w_0^2) \frac{e^{-\lambda_1 t}}{(1 - w_0^2 + w_0^2 e^{-2\lambda_1 t})^{3/2}},$$

from which we deduce

$$\varphi(t) = \frac{-2\lambda_1 w_t(1 - w_t^2)}{\varepsilon\sigma_0 \cos(\theta)} = -\frac{2\lambda_1 w_0(1 - w_0^2)e^{-\lambda_1 t}}{\varepsilon\sigma_0 \cos(\theta) (1 - w_0^2 + w_0^2 e^{-2\lambda_1 t})^{3/2}}.$$

We are now able to compute explicitly the value of the action functional on these extrema. Of course, it vanishes on the solution of (4). For w solution of $\dot{w}_t = -\lambda_1 w_t(1 - (w_t)^2)$,

$$\begin{aligned} S_{0T}^{\sigma_0}(w) &= \frac{1}{2(\varepsilon\sigma_0 \cos(\theta))^2} \int_0^T (-2\lambda_1 w_t(1 - w_t^2))^2 dt \\ &= \frac{4\lambda_1^2}{2(\varepsilon\sigma_0 \cos(\theta))^2} \int_0^T w_0^2(1 - w_0^2)^2 \frac{e^{-2\lambda_1 t}}{(1 - w_0^2 + w_0^2 e^{-2\lambda_1 t})^3} dt \\ &= \frac{\lambda_1^2(1 - w_0^2)^2}{2(\varepsilon\sigma_0 \cos(\theta))^2} \left(\frac{1}{(1 - w_0^2 + w_0^2 e^{-2\lambda_1 T})^2} - 1 \right). \end{aligned}$$

Note that another way to compute $S_{0T}^{\sigma_0}$ is as follows:

$$\begin{aligned} S_{0T}^{\sigma_0}(w) &= \frac{1}{2(\varepsilon\sigma_0 \cos(\theta))^2} \int_0^T (-2\lambda_1 w_t(1 - w_t^2))^2 dt \\ &= \frac{1}{2(\varepsilon\sigma_0 \cos(\theta))^2} \int_0^T 2\dot{w}_t(-2\lambda_1 w_t(1 - w_t^2)) dt \\ &= \frac{2\lambda_1}{2(\varepsilon\sigma_0 \cos(\theta))^2} \int_{w_0^2}^{w_T^2} du(u - 1)du \\ &= \frac{\lambda_1}{2(\varepsilon\sigma_0 \cos(\theta))^2} \left((w_T^2 - 1)^2 - (w_0^2 - 1)^2 \right). \end{aligned} \tag{24}$$

□

The explicit form of the extrema of $S_{0T}^{\sigma_0}$ allows us to compute or estimate the passage cost from one K_i^δ to another one. It is immediate to observe that for any $i \in \{1, 2, 3\}$,

$$V_{2i}^{\sigma_0} = V_{ii}^{\sigma_0} = 0,$$

since the flow (4) allows to go from K_2^δ to K_1^δ or K_3^δ , and of course to go from K_i^δ to K_i^δ .

The minimum on $w_T \in K_2^\delta$ and $w_0 \in K_3^\delta$ in the expression (24) is positive since it is attained as minimum of a positive function on a compact set. Then

$$V_{32}^{\sigma_0} > 0.$$

Furthermore following the notation in [6] p.150, we also have $\tilde{V}_{31}^{\sigma_0} = \infty$ and then from the expression in [6] p.152,

$$V_{31}^{\sigma_0} = V_{32}^{\sigma_0} > 0.$$

Similar arguments yield

$$V_{13}^{\sigma_0} = V_{12}^{\sigma_0} > 0$$

and since the diffusion coefficient σ_0 is constant, we also have by symmetry of the minimization problem that

$$V_{32}^{\sigma_0} = V_{12}^{\sigma_0}.$$

Therefore, we can state the following lemma.

Lemma 4.5. *Let us assume that $\sigma_1 = 0$. Under Hypothesis 2.1, the cost matrix of (3) is defined by*

$$P^\sigma = \begin{pmatrix} 0 & V_{12}^{\sigma_0} & V_{12}^{\sigma_0} \\ 0 & 0 & 0 \\ V_{12}^{\sigma_0} & V_{12}^{\sigma_0} & 0 \end{pmatrix}, \quad (25)$$

Let us now compute the global passage cost to the sets K_i^δ .

Proposition 4.6. *Let us assume that $\sigma_1 = 0$. The values of the global passage cost $W^{\sigma_0}(K_i^\delta)$ are given for $i \in \{1, 2, 3\}$ by*

$$W^{\sigma_0}(K_1^\delta) = W^{\sigma_0}(K_3^\delta) = V_{12}^{\sigma_0} ; W^{\sigma_0}(K_2^\delta) = 2V_{12}^{\sigma_0}.$$

Proof. Let us compute $W^{\sigma_0}(K_1^\delta)$. Consider the different graphs included in $G(1)$: $g_1 = \{2 \rightarrow 1, 3 \rightarrow 1\}$, $g_2 = \{2 \rightarrow 1, 3 \rightarrow 2\}$, $g_3 = \{2 \rightarrow 3, 3 \rightarrow 1\}$. We have

$$\sum_{(m \rightarrow n) \in g_1} V_{mn}^{\sigma_0} = 0 + V_{31}^{\sigma_0} ; \quad \sum_{(m \rightarrow n) \in g_2} V_{mn}^{\sigma_0} = V_{32}^{\sigma_0} ; \quad \sum_{(m \rightarrow n) \in g_3} V_{mn}^{\sigma_0} = V_{12}^{\sigma_0}.$$

Then, it is immediate that $W^{\sigma_0}(K_1^\delta) = V_{12}^{\sigma_0}$. □

Proposition 4.7. *Let us assume that $\sigma_1 = 0$. The sequence $(\pi^\varepsilon)_{\varepsilon \leq \varepsilon_0}$ converges when ε tends to 0 to the probability measure*

$$\frac{1}{2}(\delta_{(-1,0)} + \delta_{(1,0)}).$$

Proof. Using Theorem 4.1 in [6] for the action functional defined in (22), we know that there exists $\gamma > 0$ such that $\pi^\varepsilon(K_i^\delta)$ belongs to the interval

$$\left[\exp \left(-\varepsilon^2 \left(W^{\sigma_0}(K_i^\delta) - \min_{j \in \{1,2,3\}} W^{\sigma_0}(K_j^\delta) + \gamma \right) \right), \exp \left(-\varepsilon^2 \left(W^{\sigma_0}(K_i^\delta) - \min_{j \in \{1,2,3\}} W^{\sigma_0}(K_j^\delta) - \gamma \right) \right) \right].$$

When ε tends to 0, the accumulation points of (π^ε) will concentrate on the sets K_i^δ attaining the minimum of W^{σ_0} , i.e. K_1^δ and K_3^δ . Symmetry arguments on the stochastic system (3) (with $\sigma_1 = 0$) with respect to the axis $x = 0$, yield $\pi^\varepsilon(x < 0) = \pi^\varepsilon(x > 0)$, and then, $\pi^\varepsilon(K_1^\delta) = \pi^\varepsilon(K_3^\delta)$. Lemmas 4.1 and 4.2 allow then to conclude that there is a unique limiting value of (π^ε) given by $\frac{1}{2}(\delta_{(-1,0)} + \delta_{(1,0)})$. □

We can now consider the general case where $\sigma_1 \neq 0$ and satisfies Hypothesis 2.1.1 and prove Theorem 2.4.

Proof of Theorem 2.4. We will bound the passage costs from below and above. Let $\delta \in]0, \frac{\sigma_0 - \sigma_1}{\sigma_1}[$ and to fix ideas, assume that $\sigma_1 > 0$ (the other case is similar). We first consider the passage from K_1^δ to K_2^δ . Let $x \in]-(1 + \delta), 0[$, then

$$\sigma_0 - (1 + \delta)\sigma_1 < \sigma_0 + \sigma_1 x < \sigma_0.$$

In particular, if $w = (w_1, w_2)$ is a trajectory such that for any $t \in [0, T]$, $w_1(t) \in]-(1 + \delta), 0[$, then

$$S_{0T}^\sigma((w_t)_{t \geq 0}) > S_{0T}^{\sigma_0}((w_t)_{t \geq 0}),$$

and

$$V_{12}^\sigma > V_{12}^{\sigma_0}.$$

For the cost passage from K_3^δ to K_2^δ , we assume that $w_1(t) \in]0, (1 + \delta)[$, leading to

$$V_{32}^\sigma < V_{32}^{\sigma_0} = V_{12}^{\sigma_0}.$$

Then we deduce that $V_{12}^\sigma > V_{32}^\sigma > 0$. By similar arguments, $\sigma_1 < 0$ yields $V_{32}^\sigma > V_{12}^\sigma > 0$. That allows to compute $W^\sigma(K_i^\delta)$ as in the proof of Proposition 4.6. That gives

$$W^\sigma(K_1^\delta) = V_{32}^\sigma; W^\sigma(K_2^\delta) = V_{32}^\sigma + V_{12}^\sigma; W^\sigma(K_3^\delta) = V_{12}^\sigma.$$

The quantity $\min_{j \in \{1, 2, 3\}} W^{\sigma_0}(K_j^\delta)$ depends on the sign of σ_1 .

Using again [6, Theorem 4.2] for the action functional defined by (22), we obtain that when ε tends to 0, any converging subsequence of (π^ε) converges to a measure π which concentrates on K_1^δ if $\sigma_1 < 0$ and on K_3^δ if $\sigma_1 > 0$ (and on both sets if $\sigma_1 = 0$). Using Lemma 4.2, we deduce that there is a unique limiting measure, either $\delta_{(-1, 0)}$ if $\sigma_1 > 0$ or $\delta_{(1, 0)}$ if $\sigma_1 < 0$. \square

Acknowledgments: This work has been supported by the Chair Modélisation Mathématique et Biodiversité of Veolia - Ecole polytechnique - Museum national d'Histoire naturelle - Fondation X. It is also funded by the European Union (ERC AdG SINGER, 101054787). Views and opinions expressed are however those of the author(s) only and do not necessarily reflect those of the European Union or the European Research Council. Neither the European Union nor the granting authority can be held responsible for them.

References

- [1] M. Benaïm and T. Hurth, *Markov Chains on Metric Spaces: A Short Course*. (2022), Springer.
- [2] B. Dacorogna. Direct methods in the calculus of variations, 78, (2007), Springer Science & Business Media.
- [3] A. Dembo, O. Zeitouni. Large deviations techniques and applications, Applications of Mathematics, 38, 2011.
- [4] C. Ecotière, S. Billiard, J.B. André, P. Collet, R. Ferrière, S. Méléard. Human-environment feedback and the consistency of proenvironmental behavior, Plos Computational biology (2023), <https://doi.org/10.1371/journal.pcbi.1011429>.
- [5] S.N. Ethier and T.G. Kurtz. *Markov processes: characterization and convergence*, (1986), John Wiley & Sons
- [6] M.I. Freidlin, A.D. Wentzell. *Random perturbations of dynamical systems*, Second Edition, (1998), Springer-Verlag, New York.
- [7] M. Hairer. How hot can a heat bath get?, Communications in Mathematical Physics, 292, no 1, p.131–177, (2009), Springer.

- [8] M. Hairer, J.C. Mattingly. Slow energy dissipation in anharmonic oscillator chains, *Communications on Pure and Applied Mathematics: A Journal Issued by the Courant Institute of Mathematical Sciences*, 62, no 8, p. 999–1032, (2009), Wiley Online Library.
- [9] S. Holbach. Positive Harris recurrence for degenerate diffusions with internal variables and randomly perturbed time-periodic input, *Stochastic Processes and their Applications*, 130, no 11, p. 6965–7003, (2020), Elsevier.
- [10] R. Höpfner, E. Löcherbach, M. Thieullen. Strongly degenerate time inhomogeneous SDEs: Densities and support properties. Application to Hodgkin–Huxley type systems, *Bernoulli*, (2017).
- [11] L. Hörmander. Hypoelliptic second order differential equations, 119, *Acta Mathematica*, Institut Mittag-Leffler, p. 147 – 171, (1967), <https://doi.org/10.1007/BF02392081>.
- [12] L. Rey-Bellet. Ergodic properties of Markov processes, *Open quantum systems II*, p. 1–39, (2006), Springer.
- [13] D. Stroock, S.R.S. Varadhan. On the support of diffusion processes with applications to, *Proceedings of the Berkeley Symposium on Mathematical Statistics and Probability*, 1, p. 333, (1972), University of California Press.