

Electron Delocalization and Electrochemical Potential Distribution Phenomena in Faradaic Electrode Materials for Understanding Electrochemical Behavior

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Yachao Zhu, Siraprapha Deebansok, Jie Deng, Xuanze Wang, Thierry Brousse, et al.. Electron Delocalization and Electrochemical Potential Distribution Phenomena in Faradaic Electrode Materials for Understanding Electrochemical Behavior. Advanced Energy Materials, In press, 10.1002/aenm.202304317. hal-04560118

HAL Id: hal-04560118 https://hal.science/hal-04560118

Submitted on 26 Apr 2024 $\,$

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| 2 | Faradaic Electrode Materials for understanding electrochemical behavior |
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17 Abstract

Electrochemical energy storage devices are built upon the foundations of batteries and 18 19 supercapacitors. In the past decade, new types of electrode materials, such as pseudocapacitor-20 like species, have been intensively developed to obtain superior performance in terms of energy 21 density and powder. Pseudocapacitive materials store charge through Faradaic redox processes, 22 while their electrochemical signature remains mystically similar to that of electrochemical 23 double-layer capacitors, even though it is caused by redox processes. To address this 24 controversy, we introduce an all-inclusive analytical model for evaluating the voltammograms 25 of electrode materials. In view of this, the present article focuses on understanding the origin 26 of the pseudocapacitive phenomena in the electrochemical signal (such as cyclic voltammetry, 27 CV). Based on the self-exchange electron transfer mechanism and tunneling effect within the 28 atomic structure of solid electrode materials, pseudocapacitive materials possess a high degree 29 of electron delocalization possibility related to the number of redox centers. This process 30 initiates the electrochemical potential distribution among the neighboring redox sites, which is 31 consequently observed in electrochemical signal as the plateau feature in CV like in EDLC. 32 We proposed for the first time the determination of capacitive tendency of various faradaic 33 electrode materials, which turned out to be relative to the number of redox center (n) and total 34 charge (z) as presented in this work as a functional type of electrode material classification 35 plot. Here, we proposed the theoretical link of electron hopping, likely occurring when the 36 number of redox centers are large. The developed all-inclusive model is versatile to predict 37 rate/potential regimes for the maximum faradaic storage, and then well differentiates the 38 pseudocapacitor- materials from battery- materials.

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43 Introduction

44 The ever-growing global demand for energy underwent a sharp and sudden increase, leading 45 to the expansion of energy storage technologies. Many of these storage technologies, most 46 notably electrochemical energy storage, rely on redox processes¹. The chemical composition 47 of electrodes has been the topic of many studies to improve the performance of electrode materials in terms of energy and power densities ^{2, 3}. Within this multiplicity of compositions, 48 49 nanostructures and properties, electrode materials are classified into two categories: non-50 Faradaic and Faradaic electrode materials. A non-Faradaic reaction points to one in which there 51 is an absence of electron exchange across the electrode-electrolyte interface. In this case, 52 charge storage occurs as a consequence of charging the electrochemical double layer at the 53 electrode-solution interface. When Faradaic processes are engaged, however, electrons are 54 transferred across the current collector to or from redox centers contained inside the electrode. Redox centers are buried in a solid phase, and the charging process thus involves the motion 55 of counterions. The difference between these electrode materials is electrochemically 56 57 detectable and is analyzed according to their physicochemical traits. Numerous characteristics 58 are determined using a variety of methods, including electrochemical impedance spectroscopy and cyclic voltammetry (CV)^{4,5}. The CV technique has the advantage of varying the cell 59 60 voltage (or potential, in a 3-electrode system) and observing the impact of these fluctuations on the cell dynamics (i.e., electron and ion transfer), while modulating the scan rate, *i.e.*, the 61 time scale of the material studied.⁶ While non-Faradaic electrode materials produce a box-62 63 shaped signal during CV measurements, Faradaic electrode materials present current peaks, indicating the potential value of their redox center.⁷ However, Faradaic electrode materials are 64 65 only seldom amenable to analysis through a perfect Gaussian peak signal, as idealized by the 66 notion of immobilized redox centers (Fig. S2a). Similarly, a thorough examination of the 67 i=f(E) curves reveals that Gaussian peaks are not perfectly symmetrical (see **Fig. S2b** for the 68 ideal electrochemical response).

69 Surprisingly, there are no contemporary theories or formalist frameworks that make clear 70 the story of faradaic materials well. To compound the issue, many more questions arise when 71 addressing pseudocapacitive materials, which are characterized as Faradaic electrode materials that exhibit a box-shaped signature during a current-potential measurement (CV)⁸. Two 72 73 unsolved questions today are: firstly, how might it be possible to sign a Faradic phenomenon 74 with a curve comparable to that of an electric double-layer capacitor (EDLC) material when 75 the established signature must be Gaussian; secondly, which similarities and differences are 76 present when comparing a Gaussian ideal material to a material with pseudocapacitive 77 properties? The pseudocapacitor's box-shaped CV has been observed in a wide variety of transition metal oxides ⁹, conducting polymers ¹⁰, and other semiconducting redox materials. 78 79 The mechanisms by which Faradaic reactions may exhibit quasi-rectangular capacitive 80 electrochemical behavior remain unknown, and no cohesive model explaining the peculiarities 81 of Faradaic electrode materials exists in the literature. Numerous studies, however, serve to demonstrate this pseudocapacitive behavior (Fig. S1)^{2,4,11-44}. This thus leads one to consider 82 83 the possible reasons for a redox reaction, which is essentially a Faradaic occurrence, deviating 84 so far from the ideal peak response (as illustrated in Figure 1).





Fig. 1 (a) Conventional electrochemistry signals in ideal CVs with the conventional equation are limited to describe only simple CVs. (b) Potential distribution model within the electrode material structure based on electron transfer between redox centers (M_0 , M_1 , ..., M_n) with the redox potentials respectively to each redox site. And our new extended equation for the total current leads to the first-time complex CV fitting.

91

92 Fig. S3a illustrates the many and varied Faradaic electrode materials (i.e., those 93 concerned with electron transfer), including both those used in pseudocapacitors as well as 94 those used in batteries. Regardless of which technical device is concerned, the electrode component, including its active material, stays constant. The active material stores charges 95 through redox reactions involving the crystalline structure's redox center (as illustrated in 96 97 Figure 1). As a result, this active material may be regarded as a matrix, with redox centers 98 being connected by covalent bridges (Fig. S3b, e.g., oxygens) that lead these redox centers to 99 interact. Naturally, the true nature of materials remains multifaceted, including hierarchical 100 porosity, multicrystalline structures, the presence of binders, conductive additives, and the 101 percolation of highly conductive carbons near oxide particles.

102 In order to comprehend and rationalize the wide range of electrochemical reactions observed in Faradaic electrode materials (including pseudocapacitive materials), a three-step 103 104 strategy is provided herein. The first step involves categorizing the relevant materials according to their *n* (number of redox centers per nm³) and *Z* (electron included in crystalline structure) 105 106 values. These parameters bring the effect of the *n* and *Z* values on the electrochemical signature 107 to light. In the second step, the influence of n and Z is then rationalized using a certain 108 microscopic approach centered on electron transfer (*i.e.*, the transition-state model). The third 109 and final step, enabled by the first two in terms of creating a collection of parameters, makes 110 use of these to simulate and compare theoretical and experimental electrochemical signal 111 responses.

113 **Results and discussion**

114 As stated before, there is no available model that combines the behavior of a battery with that 115 of a pseudocapacitor currently. To complicate matters, these concepts are often separated into 116 two distinct categories. This means that researchers find themselves without a coherent, 117 universal method for analyzing and decoding electrochemical behavior (*i.e.*, the shape of the 118 cyclic voltammogram). A unified technique could be considered a significant step towards the 119 future development of innovative, high-power batteries. We thus propose an integrated model, 120 demonstrating that pseudocapacitive contributions indeed also exist in battery materials (Li-121 ion, Na-ion, Mg-ion). CVs of batteries demonstrating their cell dynamics (electron and ion 122 transport) are also shown to deviate significantly from the ideal signals published in the 1980s.

123 The number of redox centers per nm³ and the Z value as simple markers for 124 distinguishing between pseudocapacitive and battery materials

The primary goal is to be able to compare Faradaic electrode materials using three simple 125 markers: (i) n, the number of redox centers per nm³ calculated as described in the Supporting 126 Information (Section 5), (ii) Z, the number of electrons included in the crystalline structure, 127 128 and (iii) capacitive tendency, the ratio of peak current to plateau current is named 129 pseudocapacitive percentage. The number of redox centers per cubic nanometer and the 130 number of electrons per crystalline structure were determined for a series of lithium-ion 131 batteries and pseudocapacitive materials. We compiled a list of fifty materials that are the most 132 often used in electrochemical energy storage devices. Furthermore, we established a new parameter, the capacitive tendency. Capacitive tendency is the parameter to describe the 133 134 electrochemical behavior of materials. This predictive model was published by some of us.¹¹⁵, 135 using deep learning to classify electrochemical signals. The specific research involves using 136 image classification to distinguish different curve shapes (rectangular shape and peak shape) 137 that act on electrochemical signals, such as cyclic voltammetry, to connect electrochemical 138 information (pseudocapacitor and batterie). When the capacitive tendency is close to 100%, 139 the curve is more rectangular, and the material is more like a supercapacitor. When the 140 capacitive tendency is close to 0%, the curve is more like a peak shape and the material is more 141 like a battery. The capacitive tendency analyzes the shape of the voltammogram.

The capacitive tendency and capacitive contribution are not the same concept. *Bruce Dunn et al.*¹¹⁶ proposed a model in 2007 that is proportional to current and scan rate, and used this model to describe the capacitive contribution and diffusion contribution in materials. The capacitive contribution is different from capacitive tendency proposed by some of us, because the capacitive tendency is a metric to describe the shape of the material's signal curve and classify the material, rather than to calculate the proportion of capacitive contribution inside the material or provide the percentage of EDLC, pseudocapacitor or battery.

Alternatively, there is an another simple way to calculate and express rectangular shape, which is called capacitive percentage, or simply *%capa* (see ESI section 5 and Fig. S5). Equation 1 defines this pseudocapacitive percentage, but this result is just an approximate result. This one can only explain the general trend, and capacitive tendency is the result more accurate.

154
$$\mathscr{H}_{capa} = 100 - \frac{i_{max} - i_{min}}{i_{max}} \cdot 100$$
 Eq. 1

155 Where i_{max} is the maximum current and i_{min} is the relative minimum current taken after the 156 faradaic contribution. Figure 2 shows a heatmap of capacitive tendency of all materials 157 researched in the work.



Figure 2. State-of-the-Art in Faradaic electrode materials for batteries and pseudocapacitors. (a) represents a list of materials classified according to capacitive tendency (ascribed to the color bar). The Y-axis represents the number of redox centers per nm³ and the X-axis represents the total number of electrons in the crystal structure. (b) the qualitative color bar representing capacitive tendency. (c) represent typical voltammograms of the corresponding electrode material.

165 This computation makes it possible to account for the nuances inherent in the 166 voltammogram shapes as measured in the CV experiments. The %*capa* is a useful indicator 167 when it comes to anticipating the degree of charge delocalization. On the one hand, if the local 168 minimum current is close to the maximum current, then the current is identical and the behavior

169 displayed is rectangular (noting that the global resistance of the device is not taken into account 170 in the calculation). This rectangular voltammogram represents the signature of so-called 171 supercapacitive substances, be they EDLCs or pseudocapacitors. On the other hand, if the 172 difference is significant, *i.e.*, the redox peak is highly prominent, then the behavior displayed 173 is battery-like. Indeed, in order to accurately depict the electrochemical behavior of the real-174 world voltammograms, calculating the %pseudo thus allows for a simple categorization of the 175 voltammogram shape, regardless of the material composition or lack thereof. The *capacitive* 176 *tendency* is a graphical descriptor of a curve i=f(E) that is independent of the usual material 177 classification (battery, EDLC, pseudocapacitor).

The result of combining the redox center number per nm^3 with the electron number Z of the crystalline structure is shown in Fig. 2. The cell count is based on the volume of each cell (see the CIF file, in ESI), and the cell unit's Z' value indicates the number of material crystalline units repeated in one cell volume. In addition, capacitive tendency is used for over forty materials involved in the manuscript and compared it with the %pseudo results calculated from Equation 1, the complete result table is Table S4 in ESI.

184 In Fig. 2, S7-S8, a scatter plot with a gradient of different colors denoting the *capacitive* tendency is used to position the 50 different materials ^{5, 45-91}. Each of these materials is 185 accompanied by voltammograms which were extracted from scientific publications. They were 186 187 selected due to the fact that they are the most frequently applied materials in batteries and 188 pseudocapacitors. For decades, these materials have been in constant evolution with the 189 synthesis of new materials, the design of new structures, adjustments to their inner nature, and 190 modifications to their morphology, all of which were geared towards maximizing the 191 capacitance/capacity to build advanced energy storage devices. As far as pseudocapacitive 192 materials are concerned, RuO₂ stands out as one that underwent years of trials before finally 193 shedding light on the mechanism of redox reactions in aqueous acidic electrolytes. This had to

do with the valence change from Ru^{4+} to Ru^{3+}/Ru^{2+} , associated with proton insertion ⁹²⁻⁹³. 194 195 Another prominent one is MnO₂, which possesses a smooth ionic diffusion but displays limited electronic conductivity ⁹. These two laid the groundwork for further research involving metal 196 197 oxide-based Faradaic electrode materials, such as TiO₂, Nb₂O₅, Co₃O₄, Mn₃O₄, Fe₃O₄, etc. In 198 the case of Co₃O₄, it was found that morphological changes and size could lead to distinctive electrochemical behavior, since when nanosized it triggers a high-rate redox reaction.¹⁰⁵ Aside 199 200 from metal oxides, carbides and nitrides (e.g., vanadium nitride) also become the subject of 201 research relating to redox reactions in pseudocapacitive materials. Among them, 2D MXenes, 202 such as Ti₃C₂, Ti₂C, V₂C, prevail due to their intriguing electronic conductivity and high-rate capability ^{106, 107}. Numerous battery materials have been studied for use as anodes and cathodes 203 204 in various battery cells. Current cathode materials that have already been commercialized include the well-known LiCoO₂ and LiFePO₄ ^(ref 94). On the anode side, the most notable are 205 LiTi₄O₅ and TiO₂ ^(ref 95). It has been ascertained that battery materials exert redox reactions 206 following phase changes after intercalating ¹⁰⁸⁻¹¹¹. The materials selected for this study all serve 207 208 to contribute to an overall picture of Faradaic electrode materials.

We observe an unmistakable trend whereby pseudocapacitive materials display a redox center density of between 20 and 60 per nm³ and a Z value of between 0 and 120. In other words, a pseudocapacitive material has a large number of redox centers with a low number of electrons inside its crystalline structure. For battery materials, the number of redox centers are never greater than 30 per nm³, while Z values can reach up to 140.

More specifically, MnO_2 and RuO_2 both show a *capacitive tendency* of close to 100% with a large number of redox centers per nm³ and low Z values. Mn-based materials can be battery-type ones such as Li_2MnO_3 (ref ⁸⁵), Li_2MnSiO_4 (ref ⁹¹) and $LiMnPO_4$ (ref ⁵⁹) or pseudocapacitor-type ones such as MnO_2 (ref ¹³). The same goes for Ti, with battery-type ones such as $LiTiS_2$ (ref ⁸⁰), Na_2TiSiO_5 (ref ⁵⁴), $NaTi_2(PO_4)_3$ (ref ⁷⁸), $Li_4Ti_5O_{12}$ (ref ⁶⁶) and $Na_2Ti_6O_{13}$ (ref ⁸¹),

and pseudocapacitor-type ones such as $Ti_2C^{(ref 29)}$ and $Ti_3C_2^{(ref 11)}$. Some materials are not a 219 good fit for this simple classification, such as $MoS_2^{(ref 51)}$ and MoO_3 that show around 17 redox 220 centers per nm^3 and a Z value of around 70, all the while displaying pseudocapacitive behavior. 221 222 Co₃O₄, Mn₃O₄ and Fe₃O₄ exhibit large densities of redox centers together with high Z values, with *capacitive tendency* values of around 50%. 2D MXene materials, including Ti₃C₂ and 223 V_2C , are pseudocapacitive materials with high redox center densities and low Z values. It 224 225 should be noted that the choice of electrolyte probably has an impact on the *capacitive tendency* value. For example, for $Ti_3C_2T_x$ -MXene ^(ref 29), the *capacitive tendency* value decreases from 226 80 to 50% (in KOH and H₂SO₄, respectively). V₂C-MXene ^(ref 72) shows similar results. 227 228 Concerning MnO₂ in an aqueous medium, the electrochemical reversibility can be highly 229 impacted depending on the pH. This can be explained by a confinement phenomenon, where a transition occurs from double-layer to Faradaic charge storage.¹¹⁴ The electrochemical 230 231 interface actually involves porous or layered cases instead of an ideal planar model. With 232 increasing interface complexity, charge transfer improves between an electrolyte ion and the 233 host regarding the extent of ion desolvation and confinement. When K⁺ intercalates in 234 birnessite metal oxide, the K⁺ is still hydrated without desolvation between interlayers, which 235 offers a weak ion-host interaction and reduced charge transfer.

We have demonstrated that identifying the density of redox centers, n, and the corresponding Z values is enough to accurately predict the material type and the global shape of the cyclic voltammogram and vice versa. For the greater part of the researchers working in the field of electrochemical energy storage materials, determining the n and Z parameters will thus be quite straightforward. However, it is worth noting that the type of salt and nature of the solvent used in the electrolyte may sometimes affect the electrochemical behavior, and therefore the shape of the voltammogram.

To summarize, the *capacitive tendency*, and consequently the shape of the voltammogram, is chiefly governed by: (1) the number of redox centers per nm^3 , and (2) the *Z* value of the crystalline structure, bearing in mind that (3) the nature of the electrolyte may have a certain effect on the result. It is then necessary to correlate these pieces of information within a model that explains why these three parameters affect the voltammogram signal. To this end, it is necessary to understand the impact that the proximity of the density of redox centers, *n*, has on the electron moving through the Faradaic electrode material.

250 Electron movement as a cause of redox potential distribution

251 The initial Faradaic oxidative reactions that generate charges in the material are likely 252 connected to ionic mobility, independently of the pseudocapacitive or battery-like behavior of 253 the considered material. Therefore, the fundamental process of electron injection into a redox 254 center remains a redox reaction, which is a Faradaic phenomenon. However, the spatial 255 organization and density of the redox centers in battery and pseudocapacitive materials can 256 produce quite a different outcome. The density of redox centers has a direct impact on the 257 probability of an electron jumping from one redox center to another. Pseudocapacitive materials exhibit a large number of redox centers per nm^3 , with a total exceeding 30. Following 258 259 a charge injection (or extraction), the charge distribution then creates a distribution of redox 260 potentials (as demonstrated in the following section). To explore this effect, we will first 261 demonstrate that the probability of transferring an electron between two redox centers is greater 262 when *n* is large and *Z* is small. In a subsequent step, we will show that an electron exchanged 263 between two redox centers has an effect on the potential of neighboring centers.



Figure 3. Electron transfer between 2 redox centers. (a) represent the atomic changes associated with electron transfer for a solid structure. (b) represents the electron transfer probability as a function of the number of redox centers (Y-axis). The background color is respected to the transmission coefficient of electron transfer according to tunnelling effect.

270

To understand the possibility of an electron exchange between two redox centers, the probability of an electron hopping from one redox center to another may be expressed as the product of three separate probabilities. The first probability is equivalent to the chances of the electron encountering a neighboring redox center (noting that a defect in the crystalline structure will reduce this probability). The second probability is that a phonon (the vibration of a crystal lattice) causes the local distortion necessary for the temporary equalization of donor and acceptor levels (Fig 3a). The third probability is that tunneling occurs (Fig.3b), which is
exponentially dependent upon the distance between the two considered redox centers.

279 According to this formalism, the greater the density of redox centers, the greater the 280 probability of the electron being delocalized. This is consistent with the fact that a pseudocapacitive material has a large number of redox centers per nm³, so the probability of 281 charge delocalization is non-zero. It is worth noting that in the preliminary reports about the 282 redox behavior of RuO₂, Trasatti *et al.*⁽⁹⁶⁾ emphasized this fact from the very start, stating, with 283 284 respect to the metal cores, that the "Me-Me distance and the radius of the cation in these oxides 285 are such that overlap of the inner d orbitals is possible, and the d electrons in the d bands are 286 responsible for the metallic conduction", referring to Marcus studies on electron transfer to 287 substantiate this remark ⁽⁹⁷⁾.

One particular "chemical" description of the proximity of the redox center may be investigated as a reaction comprising electron hopping. There are two alternatives when an electron is injected into a redox center: (1) The injected electron is restricted to the vicinity of the redox center (*i.e.*, it is localized), or (2) the electron jumps to the nearest neighbor. After that, the charge is partially delocalized. This viewpoint is comparable to that of adiabatic vs. diabatic electron transfer ⁽⁹⁸⁾.

Now, the electron transfers or hops from one redox center to the neighboring one in the electrode material structure. The redox reaction can be transcribed, along with the associated ionic transfer (where *A* is the anion) with equilibrium constant with the equilibrium ($K_{hopping}$) as follows:

298
$$A \cdot M^{(a+1)+} + M^{(a-1)+} \stackrel{k_1}{\underset{k_{-1}}{\rightleftharpoons}} \underbrace{M^{(a)+} + M^{(a)+}}_{A}$$
 Eq. 2

where a is the reaction activity of a certain state.

If $k_1 \gg k_{-1}$ ($K_{hopping} \gg 1$) then the probability of an electron exchange between 2 redox centers is high. And inversely, if $k_1 \ll k_{-1}$ ($K_{hopping} \ll 1$), then the probability of an electron exchange is low. According to this equation, the way in which the redox charge is shared is based on an equilibrium constant, $K_{hopping}$. This equilibrium constant can then be expressed as the difference in redox potential required in order to oxidize both redox sites. E₀ is the potential to oxidize the first one, the so-called standard redox center, and E₁ is the redox potential of the closest neighbor.

307 E_n corresponds to the redox potential of the umpteenth neighbor, then the umpteenth neighbour, 308 $K_{hopping,n}$ is:

$$K_{hopping,n} = exp\left[\frac{F \cdot (E_n - E_0)}{R \cdot T}\right]$$
Eq. 3

310 Where *F* is the Faraday constant (96485 in $C \cdot mol^{-1}$), *R* is the gas constant (8.3145 *J* · 311 $mol^{-1} \cdot K^{-1}$), and *T* is temperature (293.15 K).

The equation Eq. 3 highlights two critical points: (i) Firstly, when two redox centers easily 312 313 exchange electrons ($K_{hopping,1} \gg 1$), then the difference in potential between the two redox centers is quite large $(E_n - E_0 >> 0)$. In other words, injecting a charge into one of the redox 314 centers has an effect on the amount of work required to inject a charge into the redox center 315 closest to it. (ii) Secondly, the further away a neighbor is from the standard redox center, the 316 317 less likely it is for them to share an electron. The potential distributes from the first redox center (M_0) to the nearest (M_1) and other neighboring redox centers (M_n) with hopping potential (ε) . 318 319 n is the total number of potential hopping in the system, which corresponds to the number of neighbors nearest to the area where tunneling is likely to occur. ε_{max} and ε_{min} are the 320 321 maximum and minimum potential difference between two redox centers as demonstrated in Figure 4. 322

323 The neighboring redox centers (M_i , when i = 1, 2, 3, ..., n and i is the neighbor according to 324 its position with respect to the standard redox center), the potential distributed is in descending 325 order:

$$E_1 > E_2 \gg E_n > E_{n+1} \approx E_0$$

327 Then,
$$K_{hopping,1} > K_{hopping,2} > K_{hopping,3} \gg K_{hopping,n}$$

328 Firstly, among n number of redox center, started from M_0 .

329
$$M_0, M_1, \dots, M_{n-4}, M_{n-3}, M_{n-2}, M_{n-1}, M_n$$

At the farthest redox center (M_n), the lowest energy needed for the electron to jump, ε_{min} . By recurrence sequence, the hopping potential from the farthest to the nearest neighboring redox site is in ascending order:

333
$$n\varepsilon_{min}, \dots, 5\varepsilon_{min}, 4\varepsilon_{min}, 3\varepsilon_{min}, 2\varepsilon_{min}, \varepsilon_{min}$$

334 So,
$$E_{n-i} = E_0 + (i+1)\varepsilon_{min}$$
 Eq. 4

335 When i = 1, 2, 3, ..., n.

In this model of the redox potential distribution, there are three notable variables: ε_{max} , ε_{min} and *n*. So, ε_{max} is the work difference between the two nearest redox centers $(E_1 - E_0)$. When a charge is injected into the standard redox center, its nearest neighbour has the most impacted electron affinity, so its redox potential is increased by a maximum jump (ε_{max}). Whereas the minimum potential jump corresponds to the furthest redox center influenced by the charge, noted by ε_{min} . Then *n* represents the number of neighbors influenced by the charge injection. Note that **Eq. 3** and **Eq. 4** then become, for the general form:

343
$$K_{ET,i} = exp\left[\frac{F \cdot \varepsilon_{max}}{i \cdot R \cdot T}\right]$$
 Eq. 5,

344 and, for the nearest neighbor (i = 1):

$$K_{ET,1} = exp\left[\frac{F \cdot \varepsilon_{max}}{R \cdot T}\right]$$
 Eq. 6

and, for the most distant neighbor (i = n):

347

$$K_{ET,n} = exp\left[\frac{F \cdot \varepsilon_{min}}{R \cdot T}\right]$$
 Eq. 7.

Fig. S3 illustrates the matrix model, including the redox potential distribution, using Eq.
4. When redox centers are tightly packed together, the electron affinity decreases (*i.e.*, it is
more difficult to inject the electron into the nearest neighbor). This model involves a large
redox potential distribution within the material.

A Faradaic electrode material is thus represented by the matrix, for which the continuous phase represents the overall charge in the crystalline structure and where the color of the redox centers is the distribution of redox potentials. Therefore, the proximity of the redox centers – high-density n - results in a broad redox potential distribution, as illustrated by the matrix model.

The rate constant, k_1 , which is the kinetic constant of the electron transfer, is expressed by 357 358 considering two rate-limiting steps. The first step is to consider that the electron transfer is sequential to the ion transfer. Therefore, k_1 will either be the electron transfer or the transfer of 359 360 the counter-ion to be split. However, an alternative to this could be that the electron transfer is 361 concomitant with the ion transfer, so that the rate constant k_1 will reflect a new mathematical 362 expression. Recently, Bazant et al. incorporated ion transfer effects into the kinetic theory of electron transfer⁹⁹. The nature of the counterion can be varied: Such as Li⁺, Na^{+,} Mg²⁺, Zn²⁺ 363 364 etc.

Three phenomena influence the charge storage process in electrochemical energy storage materials: (1) the tunneling effect, (2) the chemical environment of the redox center, and (3) the effect of the counterion from the electrolyte. By analogy with the electron transfer in solution, therefore, a link on charge processes exist. The crystalline structure, the counterions and the distance between the redox centers all influence the ease with which charges can be 370 delocalized. The effect of the counterion also deserves careful consideration, since the 371 kinetics of electron exchange will be influenced by its probability of being shared or not, resulting in charge delocalization. Significantly, the charge delocalization induced by these 372 373 three phenomena results in a redox potential distribution within the Faradaic electrode 374 materials. This framework thus allows one to pinpoint the fundamental differences between a 375 wide range of Faradaic electrode materials. Unexpectedly, however, battery materials display 376 a partial charge distribution, implying that their charging mechanism must include a 377 pseudocapacitive contribution. This last point is crucial, as these two types of materials are 378 frequently pitted against each other in the scientific literature. Nonetheless, we demonstrate 379 herein, for the first time, that by using our inclusive model one sees a strong connection 380 between these two types of materials and can observe the source of this distinction. This is 381 achieved by associating the redox potential distribution model with an experimentally 382 exploitable electrochemical signal.

383 Using the quantification of redox potential distribution to determine the electrochemical 384 behavior of Faradaic electrode materials

Based on the probabilistic description of electron exchange through many redox centers, it is then necessary to create a model that incorporates the redox potential distributions of the electrochemical signals.

In this case, the signal would be the sum of the individual signals from *n* populations of redox centers, *n* being the variable developed in the previous section, *i.e.*, the number of neighbors impacted by the redox reaction of the so-called standard redox center. In a similar way to the EDLC current and the battery current, the current resulting from a pseudocapacitive contribution is expressed as:

393
$$\varphi_{pseudo} = \sum_{i=0}^{i=n=\frac{E_{pseudo}^{J}-E_{pseudo}^{i}}{\varepsilon}} \varphi_{pseudo}^{i}$$
 Eq. 8

The contribution of the neighbor *i* would then be similar to the mathematical expression of the pure battery current as follows:

$$396 \qquad \varphi_{pseudo}^{i} = \frac{R \cdot T}{2 \cdot F \cdot G} \cdot e^{\frac{F}{R \cdot T} \left(E_{apply} - \left(E_{pseudo}^{i} + n \varepsilon \right) \right)} \left(\frac{e^{\frac{F \cdot G}{R \cdot T}}}{1 + e^{\frac{F \cdot G}{R \cdot T}} \cdot e^{\frac{F}{R \cdot T}} \left(E_{apply} - \left(E_{pseudo}^{i} + n \varepsilon \right) \right)} - \frac{e^{\frac{-F \cdot G}{R \cdot T}}}{1 + e^{\frac{-F \cdot G}{R \cdot T}} \cdot e^{\frac{F}{R \cdot T}} \left(E_{apply} - \left(E_{pseudo}^{i} + n \varepsilon \right) \right)} \right) Eq.9$$

397 where n is the number of potential hopping between redox centers (per nm³), ε is the hopping potential (in V), E_{pseudo}^{i} is the standard redox potential of the pseudocapacitive current (in V), 398 E_{pseudo}^{f} is the final potential of the pseudocapacitive current (in V), and φ_{pseudo}^{i} is the 399 400 neighboring pseudocapacitive current (dimensionless current). For the first time, herein we 401 thus propose a homogeneous and universal formalism for simulating a voltammogram, 402 irrespective of the origin of its redox dynamics, for all electrochemical energy storage 403 materials, be they electrochemical double layer, purely battery-type or pseudocapacitive in 404 nature. Moreover, via this formalism, we can show that all electrode materials possess a 405 combination of these three contributions. However, the contribution of each of the components will then be more or less intense, depending on the value of the maximum current (i_{EDLC}^{max}) 406 $i_{pseudo}^{max}, i_{batt}^{max}), i.e.:$ 407

408 $i_{tot} = i_{EDLC}^{max} \cdot \varphi_{EDLC} + i_{pseudo}^{max} \cdot \varphi_{pseudo} + i_{batt}^{max} \cdot \varphi_{batt}$ Eq.10.

409

410 Fig. S11 shows CV profiles of various electrochemical systems, including i_{peak} + 411 i_{EDLC} , pure i_{pseudo} , i_{pseudo} + i_{EDLC} , and the contour plots of % pseudo of these systems related 412 to the resistance and scan rate. These results were obtained by using the Wolfram Language with its Mathematica tool to plot the functions according to the equation Eq. 10, where $A_{batt} =$ 413 $A_{EDLC} = A_{pseudo} = 1$ was considered. Firstly, the different current profiles of the first two 414 415 models in Fig. S11a and b, with different values for R_s and ν , can be observed. These characteristics relate to different %pseudo as indicated in the contour plot (Fig. S11c). One can 416 417 clearly see that, even when increasing the scan rate, the %pseudo still barely reaches 90% as a result of the peak current contribution. This result agrees with experimental observations (see the ESI). Secondly, the rectangular CV profile (**Fig. S11d** and **e**) is distorted by the presence of i_{EDLC} current with a large $R_s = 50\Omega$. Similarly, **Fig. S11f** demonstrates %pseudo as a function of the R_s and the scan rate for a pseudocapacitive current with the contribution of i_{EDLC} . It was found that the %*pseudo* is high, even when increasing the scan rate. This result also agrees with experimental observations (Fig. S13-S60).





425 **Fig. 4** Formalism of the novel pseudocapacitor CV fitting method. (a) Pseudocapacitor CV 426 of various numbers of electron hopping (or the number of redox centers, n) and the potential 427 distribution (or hopping potential, ϵ). (b) The equation of pseudocapacitor signal. (c) 428 Experimental CV fitted by our formalism.

Fig. 4a shows the proposed pseudocapacitor model where the redox peaks start to overlap with
the number of electron hopping. The extended formalism from Eq. 8 (as demonstrated in Fig.
432 4b) was finally introduced into the total current with other current components (Eq. 10). Using
this model, an experimental CV was well fitted, as shown in Fig. 4c. It was also found that,

even in the battery materials, there is a certain pseudocapacitive contribution resulting from the
distribution of the potential to the nearby redox centers when the number of redox centers per
nm³ is high, as presented in the corresponding matrix models with numerous redox centers
(Fig. S13-S60).

The conventional method is to use the peak current to determine whether a device is a battery or a pseudocapacitor. While the peak current model is a good fit, the signal including all potential ranges is incomplete. Moreover, the model proposed herein has shown a strong correlation between state-of-the-art experience and reliable theory. Besides, this proposed theory is based on equilibrium exchange between two redox centers, so it cannot be used to analyse the reversibility of the charge-discharge process.

444 Conclusion

445 Batteries and supercapacitors are both widely used electrochemical energy storage technologies. When it comes to batteries, redox processes display unusual electrochemical 446 447 behavior, for which there is no clear explanation to date. Complicating matters further, no 448 universal analytical formalism for assessing voltammograms has thus far been proposed. This 449 article began by discussing the electrochemical properties of batteries and pseudocapacitive 450 materials, as well as what sets them apart while focusing on electron transport pathways. We 451 have shown that battery and pseudocapacitor materials, called Faradaic electrode materials, 452 both exhibit strong electron delocalization resulting in large redox potential distributions due 453 to tunneling and the polaron effect. Perhaps more significantly, battery materials display the 454 same potential dispersion, albeit weaker. Faradaic electrode materials are very much in the 455 scientific limelight of late, and our research serves to provide a link between them all, thereby 456 culminating in a ground-breaking, all-inclusive model for assessing such materials.

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