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A multi-scale multi-lane model for traffic regulation via autonomous vehicles

Paola Goatin¹ Benedetto Piccoli²

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Abstract

We propose a new model for multi-lane traffic with moving bottlenecks, e.g., autonomous vehicles (AV). It consists of a system of balance laws for traffic in each lane, coupled in the source terms for lane changing, and fully coupled to ODEs for the AVs' trajectories. More precisely, each AV solves a controlled equation depending on the traffic density, while the PDE on the corresponding lane has a flux constraint at the AV's location. We prove existence of entropy weak solutions, and we characterize the limiting behavior for the source term converging to zero (without AVs), corresponding to a scalar conservation law for the total density. The convergence in the presence of AVs is more delicate and we show that the limit does not satisfy an entropic equation for the total density as in the original coupled ODE-PDE model. Finally, we illustrate our results via numerical simulations.

Key words: Multi-scale traffic flow models; PDE-ODE system; moving bottlenecks, autonomous vehicles; wave-front tracking; finite volume schemes; control problems.

1 Introduction

Multi-lane traffic has attracted the attention of many researchers in transportation science, in particular applied mathematicians and engineers. Works on macroscopic models include papers on modeling [21] and analysis [6, 20]. Recently, extensions to road networks were proposed, see [16, 17]. Starting from these results, we aim at including the presence of autonomous vehicles, which act as moving bottlenecks to regulate traffic, e.g., dissipate traffic waves [1, 35]. Such regulation problems were addressed on a theoretical basis [7, 8, 11, 12, 29, 30, 31, 36, 40, 41], with machine learning approaches [22, 38] and also via real world experiments [18, 34, 39].

A number of approaches were developed to deal with moving bottlenecks [9, 13, 23, 25, 26, 32, 33]. Moreover, extensions to multiple bottlenecks [14] and to second order models [37] were proposed. In particular we follow the idea of coupling an ODE for the moving bottleneck to a PDE for the bulk traffic as in [9, 10, 13]. The coupling is realized in the ODE right-hand side, which depends on the PDE solution, and by imposing a flux limiter to the PDE at the location of the bottleneck. Convergence and continuous dependence for such models were achieved [27, 28].

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In this work, we consider a first order macroscopic multi-lane model [6, 20] of the form

$$\begin{cases} \partial_t \rho_j + \partial_x F_j(\rho_j) = \frac{1}{\tau} (S_{j-1}(\rho_{j-1}, \rho_j) - S_j(\rho_j, \rho_{j+1})), & x \in \mathbb{R}, t \geq 0, \\ \rho_j(0, x) = \rho_j^0(x), & j = 1, \dots, M. \end{cases} \quad (1.1)$$

Above, M is the number of lanes on the road, $\rho_j : [0, +\infty[\times \mathbb{R} \rightarrow [0, R_j]$ is the vehicle density on lane j , $F_j(\rho_j) = \rho_j v_j(\rho_j)$ is the flux function and the average speed $v_j = v_j(\rho_j) : [0, R_j] \rightarrow [0, V_j]$ is a strictly decreasing function such that $v_j(0) = V_j$ and $v_j(R_j) = 0$, so that F_j is concave. The source terms S_j account for mass exchanges from lane j to lane $j+1$ (obviously, $S_0 = S_M = 0$), scaled by a relaxation factor $1/\tau \in \mathbb{R}^+$. Letting $\tau \rightarrow 0$, this corresponds to letting to zero the lane-change relaxation parameter.

Throughout the paper, we will assume that

(S0) $S_j : \mathbb{R}_+^2 \rightarrow \mathbb{R}$, $j = 1, \dots, M-1$, are Lipschitz continuous in both variables, with Lipschitz constant \mathcal{S} , and $S_j(0, 0) = S_j(R_j, R_{j+1}) = 0$, $\partial_1 S_j(u, w) \geq 0$ and $\partial_2 S_j(u, w) \leq 0$, for all $u \in [0, R_j]$ and $w \in [0, R_{j+1}]$.

Some results will be subject to further hypotheses:

(S1) If $S_j(u, w) = 0$, then $u = w$ (in particular, by **(S0)** this implies $R_j = R$ for all $j = 1, \dots, M$);

(S2) There exists $c > 0$ such that $\partial_1 S_j(u, w) \geq c$ and $\partial_2 S_j(u, w) \leq -c$, for all $u \in [0, R_j]$ and $w \in [0, R_{j+1}]$.

To account for the presence of an autonomous vehicle (AV) in lane j , we model it as a moving bottleneck [9, 13] reducing to zero the capacity of the lane at the AV position, see also [24]. Let $y_j^\ell : [0, +\infty[\rightarrow \mathbb{R}$, $\ell = 1, \dots, N_j$, be the trajectory of the ℓ -th AV travelling on the lane j . The coupling with (1.1) is realized though the following microscopic ODE and constraint:

$$\dot{y}_j^\ell(t) = \min \left\{ u_j^\ell(t), v_j(\rho_j(t, y_j^\ell(t+))) \right\}, \quad t > 0, \quad (1.2a)$$

$$y_j^\ell(0) = y_{0,j}^\ell, \quad (1.2b)$$

$$\rho_j(t, y_j^\ell(t)) \left(v_j(\rho_j(t, y_j^\ell(t))) - \dot{y}_j^\ell(t) \right) \leq 0, \quad t > 0, \quad (1.2c)$$

where $u_j^\ell : [0, +\infty[\rightarrow [0, V_j]$ is the control law (desired speed) applied to the AV. Remark that, when the constraint (1.2c) is active, i.e. $\min \left\{ u_j^\ell(t), v_j(\rho_j(t, y_j^\ell(t+))) \right\} = u_j^\ell(t) < v_j(\rho_j(t, y_j^\ell(t+)))$, a non-classical shock arises and the downstream value of the density must be zero: $\rho_j(t, y_j^\ell(t+)) = 0$, while $v_j(\rho_j(t, y_j^\ell(t-))) = \dot{y}_j^\ell(t)$. In particular, $\rho_j(t, y_j^\ell(t-)) = \hat{\rho}_{u_j^\ell}$, where $\hat{\rho}_{u_j^\ell} > 0$ is defined by the unique root of the equation $v(\rho) = u_j^\ell$. We refer the reader to Figure 1 for a graphical representation.

We first provide existence results for the coupled system (1.1)-(1.2), under assumption **(S0)**, using the wave-front tracking approximation for a fixed $\tau > 0$ and assuming bounded variation for the initial datum and the AV control law. Bounds on the total variation are achieved via a careful analysis of the contribution of the source terms and non-classical shocks which may appear at AV locations.

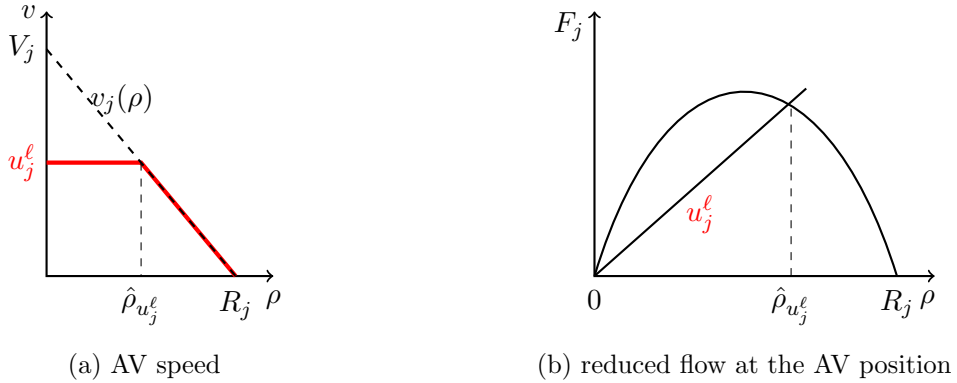


Figure 1: Graphical representation of the AV speed (1.2a) and the corresponding flow constraint (1.2c) introducing the corresponding notation.

Unfortunately, Lipschitz-type estimates cannot be achieved uniformly in $\tau > 0$. Therefore, we first study the limit as $\tau \rightarrow 0$ without AVs, which is given by the solution of a scalar conservation law for the sum of the densities over all the lanes. Such convergence is achieved under assumptions **(S0)** and **(S1)**, adapting a technique originally designed for a chromatography problem, see [4].

We then focus on the case of two lanes and a single AV, under assumptions **(S0)**, **(S1)**, and **(S2)**. Even in this simplified setting, we show that in the limit $\tau \rightarrow 0$ we do not recover the entropy weak solutions in the sense of [9], attesting the richness of the model.

To further illustrate our results, we provide numerical simulations via finite volume schemes. We start illustrating the dynamics on a 3-lane road with AVs and different speeds. An oscillatory pattern emerges due to the uneven distribution of traffic among lanes. We then pass to investigate numerically the convergence as $\tau \rightarrow 0$ with and without AVs. In particular, we show the difference between the limit of solutions with AVs and the scalar model proposed in [9] for the sum of densities.

This paper provide further comprehension of the complexities arising when coupling multi-lane traffic with moving bottlenecks, e.g. AVs. Natural next steps include characterizing further the limit as $\tau \rightarrow 0$, understanding continuous dependence, and extending the theory to networks.

2 Existence of solutions: Wave-front tracking approximations and compactness estimates

The multi-lane model including moving bottlenecks reads:

$$\begin{cases} \partial_t \rho_j + \partial_x F_j(\rho_j) = \frac{1}{\tau} (S_{j-1}(\rho_{j-1}, \rho_j) - S_j(\rho_j, \rho_{j+1})), \\ \rho_j(0, x) = \rho_j^0(x), \\ \dot{y}_j^\ell(t) = \min\{u_j^\ell(t), v_j(\rho_j(t, y_j^\ell(t)))\}, \\ y_j^\ell(0) = y_{0,j}^\ell, \\ \rho_j(t, y_j^\ell(t)) (v_j(\rho_j(t, y_j^\ell(t))) - \dot{y}_j^\ell(t)) \leq 0, \end{cases} \quad \begin{array}{l} x \in \mathbb{R}, t \geq 0, \\ \ell = 1, \dots, N_j, \\ j = 1, \dots, M, \end{array} \quad (2.3)$$

Solutions to (2.3) are intended in the following weak sense.

Definition 1. The $(M + \sum_j N_j)$ -tuple $(\boldsymbol{\rho}^\tau, \mathbf{y})$, with

$$\boldsymbol{\rho}^\tau = (\rho_1^\tau, \dots, \rho_M^\tau) \quad \text{and} \quad \mathbf{y} = (y_1^1, \dots, y_1^{N_1}, y_2^1, \dots, y_M^{N_M}),$$

provides a solution to (2.3) with $\boldsymbol{\rho}^0 = (\rho_1^0, \dots, \rho_M^0) \in \mathbf{L}^1(\mathbb{R}; [0, R_1] \times \dots \times [0, R_M])$ if the following conditions hold for $\ell = 1, \dots, N_j$, $j = 1, \dots, M$:

1. $\boldsymbol{\rho}^\tau \in \mathbf{C}^0([0, +\infty[; \mathbf{L}^1(\mathbb{R}; [0, R_1] \times \dots \times [0, R_M]))$ and $TV(\boldsymbol{\rho}^\tau(t)) < +\infty$ for all $t > 0$;
2. $y_j^\ell \in \mathbf{W}_{\text{loc}}^{1,1}(\mathbb{R}^+; \mathbb{R})$;
3. For every $\kappa \in \mathbb{R}$ and for all $\varphi \in \mathbf{C}_c^1(\mathbb{R}^2; \mathbb{R}^+)$ it holds

$$\begin{aligned} & \int_{\mathbb{R}^+} \int_{\mathbb{R}} (|\rho_j^\tau - \kappa| \partial_t \varphi + \text{sgn}(\rho_j^\tau - \kappa) (F_j(\rho_j^\tau) - F_j(\kappa)) \partial_x \varphi) dx dt + \int_{\mathbb{R}} |\rho_j^0 - \kappa| \varphi(0, x) dx \\ & + \frac{1}{\tau} \int_{\mathbb{R}^+} \int_{\mathbb{R}} \text{sgn}(\rho_j^\tau - \kappa) (S_{j-1}(\rho_{j-1}^\tau, \rho_j^\tau) - S_j(\rho_j^\tau, \rho_{j+1}^\tau)) \varphi dx dt \\ & + 2 \sum_{\ell=1}^{N_j} \int_{\mathbb{R}^+} (F_j(\kappa) - \dot{y}_j^\ell(t) \kappa)^+ \varphi(t, y_j^\ell(t)) dt \geq 0; \end{aligned} \quad (2.4)$$

4. For a.e. $t > 0$, $F_j(\rho_j^\tau(t, y_j^\ell(t) \pm)) - \dot{y}_j^\ell(t) \rho_j^\tau(t, y_j^\ell(t) \pm) \leq 0$;
5. For a.e. $t > 0$, $\dot{y}_j^\ell(t) = \min \left\{ u_j^\ell(t), v_j(\rho_j^\tau(t, y_j^\ell(t) +)) \right\}$.

Existence and uniqueness without moving bottlenecks have been provided in [16, 20]. To include moving bottlenecks, we can use the techniques developed in [13, 24].

From now on, we consider for simplicity only one AV, say on lane $i \in \{1, \dots, M\}$. We denote by $y = y(t)$ its trajectory, and by $u = u(t)$ its desired velocity. We therefore focus on the well-posedness of the strongly coupled problem

$$\begin{cases} \partial_t \rho_j + \partial_x F_j(\rho_j) = \frac{1}{\tau} (S_{j-1}(\rho_{j-1}, \rho_j) - S_j(\rho_j, \rho_{j+1})), \\ \rho_j(0, x) = \rho_j^0(x), \\ \dot{y}(t) = \min\{u(t), v_i(\rho_i(t, y(t) +))\}, \\ y(0) = y_0, \\ \rho_i(t, y(t)) (v_i(\rho_i(t, y(t))) - \dot{y}(t)) \leq 0, \end{cases} \quad \begin{array}{l} x \in \mathbb{R}, t \geq 0, \\ j = 1, \dots, M. \end{array} \quad (2.5)$$

We will prove the following existence result.

Theorem 1. Under hypotheses **(S0)**, let the initial conditions $\rho_j^0 \in \mathbf{L}^1(\mathbb{R}; [0, R_j])$ for $j = 1, \dots, M$, have finite total variation, and $u \in \mathbf{L}^1(\mathbb{R}^+; [0, V_i])$ also have finite total variation. Then, for any $\tau > 0$, there exists a solution $(\boldsymbol{\rho}^\tau, y^\tau)$ of (2.5) in the sense of Definition 1.

For the proof, we proceed similarly to [13, 15].

We approximate the initial data by a sequence of piece-wise constant functions $\rho_j^{0,\nu}$, $j = 1, \dots, M$, such that each $\rho_j^{0,\nu}$ has a finite number of discontinuities and

$$\lim_{\nu \rightarrow +\infty} \left\| \rho_j^{0,\nu} - \rho_j^0 \right\|_1 = 0 \quad \text{and} \quad \text{TV} \left(\rho_j^{0,\nu} \right) \leq \text{TV} \left(\rho_j^0 \right) \quad \text{for all } \nu \in \mathbb{N}.$$

Analogously, we approximate u with piece-wise constant functions u^ν with a finite number of discontinuities and such that

$$\lim_{\nu \rightarrow +\infty} \|u^\nu - u\|_1 = 0 \quad \text{and} \quad \text{TV} (u^\nu) \leq \text{TV} (u) \quad \text{for all } \nu \in \mathbb{N}.$$

Consider a sequence of time steps Δt^ν , $\nu \in \mathbb{N}$, such that $\lim_{\nu \rightarrow \infty} \Delta t^\nu = 0$. We construct a sequence $(\boldsymbol{\rho}^{\tau,\nu}, y_\nu^\tau)$ of approximate solutions to (1.1)-(1.2) using the Wave Front Tracking (WFT) and operator splitting algorithms as follows.

For any $\nu > 0$, we set $t^n := n\Delta t^\nu$, $n \in \mathbb{N}$. On the time interval $[t^n, t^{n+1}[$, $n \in \mathbb{N}$, we define $(\boldsymbol{\rho}^{\tau,\nu}, y_\nu^\tau)$ as the WFT approximation of

$$\begin{cases} \partial_t \rho_j + \partial_x F_j(\rho_j) = 0, \\ \rho_j(t^n, x) = \rho_j^{\tau,\nu}(t^n, x), \\ \dot{y}(t) = \min\{u^\nu(t), v_i(\rho_i(t, y(t)+))\}, \\ y(t^n) = y_\nu^\tau(t^n-), \\ \rho_i(t, y(t)) (v_i(\rho_i(t, y(t))) - \dot{y}(t)) \leq 0, \end{cases} \quad \begin{array}{l} x \in \mathbb{R}, t \in [t^n, t^{n+1}[\\ j = 1, \dots, M, \end{array} \quad (2.6)$$

constructed as described in [3, Chapter 7] or [1, Chapter A.7]. Note that, due to the presence of the source term, we cannot employ a fixed grid as in [3, Chapter 6], see also [13, Section 3.1]. At time $t = t^{n+1}$, we set

$$\rho_j^{\tau,\nu}(t^{n+1}, \cdot) = \rho_j^{\tau,\nu}(t^{n+1}-, \cdot) + \frac{\Delta t^\nu}{\tau} \left(S_{j-1}(\rho_{j-1}^{\tau,\nu}, \rho_j^{\tau,\nu}) - S_j(\rho_j^{\tau,\nu}, \rho_{j+1}^{\tau,\nu}) \right). \quad (2.7)$$

This sequence of approximate solutions satisfies the following compactness estimates.

Lemma 1. *Let (S0) hold and $\Delta t^\nu \leq \tau/2\mathcal{S}$. If $\rho_j^0(x) \in [0, R_j]$ for all $j = 1, \dots, M$, then the approximate solutions $\boldsymbol{\rho}^{\tau,\nu}$ satisfy*

$$\rho_j^{\tau,\nu}(t, x) \in [0, R_j] \quad \text{for all } x \in \mathbb{R}, t > 0, \tau > 0.$$

In particular, the domain $\Pi_{j=1}^M [0, R_j]$ is invariant for (2.5).

Proof. Since the homogeneous step (2.6) does not alter the invariant domain, we focus on (2.7). We proceed by induction and assume $0 \leq \rho_j^{\tau,\nu}(t^n-, x) \leq R_j$ for all $j = 1, \dots, M$. Denoting $\rho_j^\pm = \rho_j^{\tau,\nu}(t^n \pm, x)$, we get by hypotheses

$$\begin{aligned} \rho_j^+ &= \rho_j^- + \frac{\Delta t^\nu}{\tau} \left(S_{j-1}(\rho_{j-1}^-, \rho_j^-) - S_j(\rho_j^-, \rho_{j+1}^-) \right) \\ &= \rho_j^- + \frac{\Delta t^\nu}{\tau} \left(S_{j-1}(\rho_{j-1}^-, \rho_j^-) - S_{j-1}(0, 0) - S_j(\rho_j^-, \rho_{j+1}^-) + S_j(0, 0) \right) \\ &= \rho_j^- \left[1 - \frac{\Delta t^\nu}{\tau} (\partial_1 S_j - \partial_2 S_{j-1}) \right] + \frac{\Delta t^\nu}{\tau} \left(\partial_1 S_{j-1} \rho_{j-1}^- - \partial_2 S_j \rho_{j+1}^- \right) \geq 0, \end{aligned}$$

where we have used the Mean Value Theorem setting $\partial_1 S_\ell = \partial_1 S_\ell(\xi_\ell \rho_{\ell-1}^-, \xi_\ell \rho_\ell^-)$ and $\partial_2 S_\ell = \partial_2 S_\ell(\xi_\ell \rho_{\ell-1}^-, \xi_\ell \rho_\ell^-)$ for some $\xi_\ell \in [0, 1]$, $\ell = j-1, j$, and we have exploited **(S)** and the hypothesis $\Delta t^\nu \leq \tau/2\mathcal{S}$. Similarly,

$$\begin{aligned} \rho_j^+ - R_j &= \rho_j^- - R_j + \frac{\Delta t^\nu}{\tau} \left(S_{j-1}(\rho_{j-1}^-, \rho_j^-) - S_j(\rho_j^-, \rho_{j+1}^-) \right) \\ &= \rho_j^- - R_j + \frac{\Delta t^\nu}{\tau} \left(S_{j-1}(\rho_{j-1}^-, \rho_j^-) - S_{j-1}(R_{j-1}, R_j) - S_j(\rho_j^-, \rho_{j+1}^-) + S_j(R_j, R_{j+1}) \right) \\ &\leq (\rho_j^- - R_j) \left[1 - \frac{\Delta t^\nu}{\tau} (\partial_1 S_j - \partial_2 S_{j-1}) \right] \leq 0, \end{aligned}$$

with $\partial_1 S_j = \partial_1 S_j(\sigma_{j-1}, \sigma_j)$ and $\partial_2 S_{j-1} = \partial_2 S_{j-1}(\sigma_{j-2}, \sigma_{j-1})$ for some $\sigma_{\ell-1} = \rho_{\ell-1}^- + \xi_\ell (R_{\ell-1} - \rho_{\ell-1}^-)$, $\sigma_\ell = \rho_\ell^- + \xi_\ell (R_\ell - \rho_\ell^-)$ with $\xi_\ell \in [0, 1]$, $\ell = j-1, j$. \square

Lemma 2. *Let **(S0)** hold and $\Delta t^\nu \leq \tau/2\mathcal{S}$, then the approximate solutions $\rho^{\tau, \nu}$ satisfy*

$$TV\left(\rho_j^{\tau, \nu}(t, \cdot)\right) \leq \mathcal{T}, \quad \text{for all } t > 0, \tau > 0, \quad (2.8)$$

for $j \in \{1, \dots, M\}$, where $\mathcal{T} := \max_j TV\left(\rho_j^0\right) + 2R_i + CTV(u)$. In particular, the sequence total variation in space is uniformly bounded for all $\tau > 0$.

Proof. We consider the Glimm type functional

$$\Upsilon_i(t) = \Upsilon\left(\rho_i^{\tau, \nu}(t, \cdot), u_\nu\right) := TV\left(\rho_i^{\tau, \nu}(t, \cdot)\right) + 2R_i + \gamma(t) + CTV\left(u^\nu(\cdot); [t, +\infty[)\right), \quad (2.9)$$

where $C = -6/\max_{\rho \in [0, R_i]} F_i''(\rho)$ and γ is given by

$$\gamma(t) := \begin{cases} -2\hat{\rho}_{u^\nu(t)} & \text{if } \rho_i^{\tau, \nu}(t, y_\nu^\tau(t)-) = \hat{\rho}_{u^\nu(t)}, \rho_i^{\tau, \nu}(t, y_\nu^\tau(t)+) = 0, \\ 0 & \text{otherwise,} \end{cases}$$

we have that $\Upsilon_i(t^{n+1}-) \leq \Upsilon_i(t^n+)$. Of course, for $j \neq i$, where the AV is not present, we have $TV\left(\rho_j^{\tau, \nu}(t^{n+1}-, \cdot)\right) \leq TV\left(\rho_j^{\tau, \nu}(t^n+, \cdot)\right)$.

It is not restrictive to assume that jump locations in $\rho_j^{\tau, \nu}(t^{n+1}, \cdot)$ do not coincide for different $j \in \{1, \dots, M\}$. Denoting by $\rho_{j,L}^\pm$ and $\rho_{j,R}^\pm$ respectively the left and right traces of $\rho_j^{\tau, \nu}(t^{n+1} \pm, \cdot)$ at a jump discontinuity, we get from (2.7)

$$\begin{aligned} \rho_{j,L}^+ - \rho_{j,R}^+ &= \rho_{j,L}^- - \rho_{j,R}^- + \frac{\Delta t^\nu}{\tau} \left(S_{j-1}(\rho_{j-1}^{\tau, \nu}, \rho_{j,L}^-) - S_j(\rho_{j,L}^-, \rho_{j+1}^{\tau, \nu}) \right) \\ &\quad - \frac{\Delta t^\nu}{\tau} \left(S_{j-1}(\rho_{j-1}^{\tau, \nu}, \rho_{j,R}^-) - S_j(\rho_{j,R}^-, \rho_{j+1}^{\tau, \nu}) \right) \\ &= \left(\rho_{j,L}^- - \rho_{j,R}^- \right) \left[1 + \frac{\Delta t^\nu}{\tau} \left(\partial_2 S_{j-1}(\rho_{j-1}^{\tau, \nu}, \xi_j) - \partial_1 S_j(\tilde{\xi}_j, \rho_{j+1}^{\tau, \nu}) \right) \right], \end{aligned}$$

for some $\xi_j, \tilde{\xi}_j > 0$. By **(S)**, we have that

$$0 \leq 1 + \frac{\Delta t^\nu}{\tau} \left(\partial_2 S_{j-1}(\rho_{j-1}^{\tau, \nu}, \xi_j) - \partial_1 S_j(\tilde{\xi}_j, \rho_{j+1}^{\tau, \nu}) \right) \leq 1,$$

provided that $\Delta t^\nu \leq \tau/2\mathcal{S}$. Therefore we conclude that

$$\mathrm{TV}\left(\rho_j^{\tau,\nu}(t^{n+1+}, \cdot)\right) \leq \mathrm{TV}\left(\rho_j^{\tau,\nu}(t^{n+1-}, \cdot)\right) \quad \text{for } j = 1, \dots, M.$$

Finally, we get the uniform TV bounds (independent of τ)

$$\mathrm{TV}\left(\rho_j^{\tau,\nu}(t, \cdot)\right) \leq \mathrm{TV}\left(\rho_j^0\right), \quad \text{for } j \in \{1, \dots, M\} \setminus i, \quad (2.10)$$

$$\mathrm{TV}\left(\rho_i^{\tau,\nu}(t, \cdot)\right) \leq \mathrm{TV}\left(\rho_i^0\right) + 2R_i + C \mathrm{TV}(u). \quad (2.11)$$

Setting $\mathcal{T} := \max_j \mathrm{TV}\left(\rho_j^0\right) + 2R_i + C \mathrm{TV}(u)$ we get the desired estimate. \square

Lemma 3. *Let (S0) hold and $\Delta t^\nu \leq \tau/2\mathcal{S}$, then for any $t_1, t_2 \in [0, T]$, $T > 0$, the approximate solutions $\rho^{\tau,\nu}$ satisfy*

$$\sum_{j=1}^M \left\| \rho_j^{\tau,\nu}(t_1, \cdot) - \rho_j^{\tau,\nu}(t_2, \cdot) \right\|_1 \leq (C + C_\tau) |t_1 - t_2|, \quad \text{for all } \nu \in \mathbb{N},$$

where $C = ML\mathcal{T}$ is independent of τ (with $L := \max\left\{\max_j \|F'_j\|_\infty, \|u\|_\infty\right\}$) and $C_\tau = \frac{4\mathcal{S}}{\tau} \left(\sum_{j=1}^M \|\rho_j^0\|_1 + 2ML\mathcal{T}\right)$. In particular, the Lipschitz constant does not depend on ν but blows-up as $\tau \searrow 0$.

Proof. Let us fix $t_1, t_2 \in \mathbb{R}$ such that $0 \leq t_1 < t_2$ and suppose that there are $N + 1$ time steps between t_1 and t_2 :

$$t_1 \leq t^k < t^{k+1} < \dots < t^{k+N} \leq t_2 \quad \text{for some } k \geq 0.$$

Thus,

$$\begin{aligned} \left\| \rho_j^{\tau,\nu}(t_2, \cdot) - \rho_j^{\tau,\nu}(t_1, \cdot) \right\|_1 &\leq \left\| \rho_j^{\tau,\nu}(t_2, \cdot) - \rho_j^{\tau,\nu}(t^{k+N+}, \cdot) \right\|_1 + \left\| \rho_j^{\tau,\nu}(t^{k-}, \cdot) - \rho_j^{\tau,\nu}(t_1, \cdot) \right\|_1 \\ &+ \sum_{i=k}^{k+N} \left\| \rho_j^{\tau,\nu}(t^i+, \cdot) - \rho_j^{\tau,\nu}(t^i-, \cdot) \right\|_1 + \sum_{i=k}^{k+N-1} \left\| \rho_j^{\tau,\nu}(t^{i+1-}, \cdot) - \rho_j^{\tau,\nu}(t^i+, \cdot) \right\|_1. \end{aligned}$$

We observe that, due to finite wave-propagation speed, we have

$$\left\| \rho_j^{\tau,\nu}(t^{i+1-}, \cdot) - \rho_j^{\tau,\nu}(t^i+, \cdot) \right\|_1 \leq L\mathcal{T} \Delta t^\nu,$$

where $L := \max\left\{\max_j \|F'_j\|_\infty, \|u\|_\infty\right\}$ is the maximum of the Lipschitz constant of F_j and the maximal AV speed. Moreover, by (2.7),

$$\left\| \rho_j^{\tau,\nu}(t^i+, \cdot) - \rho_j^{\tau,\nu}(t^i-, \cdot) \right\|_1 = \frac{\Delta t^\nu}{\tau} \left\| \left(S_{j-1}(\rho_{j-1}^{\tau,\nu}, \rho_j^{\tau,\nu}) - S_j(\rho_j^{\tau,\nu}, \rho_{j+1}^{\tau,\nu}) \right) (t^i-, \cdot) \right\|_1.$$

Let us define

$$g_{j,i}^\pm := \left\| \left(S_{j-1}(\rho_{j-1}^{\tau,\nu}, \rho_j^{\tau,\nu}) - S_j(\rho_j^{\tau,\nu}, \rho_{j+1}^{\tau,\nu}) \right) (t^i\pm, \cdot) \right\|_1$$

and set

$$\begin{aligned}\rho_j^\pm &= \rho_{j,i}^\pm := \rho_{j,\nu}^\tau(t^i \pm, \cdot), \\ S_j^\pm &= S_{j,i}^\pm := S_j(\rho_j^\pm, \rho_{j+1}^\pm), \\ \partial_1 S_j &:= \partial_1 S_j(\sigma_j, \sigma_{j+1}), \quad \text{where } \sigma_j = \rho_j^- + \xi_j(\rho_j^+ - \rho_j^-), \\ \partial_2 S_j &:= \partial_2 S_j(\sigma_j, \sigma_{j+1}), \quad \text{where } \sigma_{j+1} = \rho_{j+1}^- + \xi_j(\rho_{j+1}^+ - \rho_{j+1}^-),\end{aligned}$$

for some $\xi_j \in [0, 1]$.

Using the above notations, we develop

$$\begin{aligned}\sum_{j=1}^M g_{j,i}^+ &= \sum_{j=1}^M \left\| S_{j-1}^+ - S_j^+ \right\|_1 \\ &= \sum_{j=1}^M \left\| S_{j-1}^- - S_j^- + \partial_1 S_{j-1}(\rho_{j-1}^+ - \rho_{j-1}^-) + (\partial_2 S_{j-1} - \partial_1 S_j)(\rho_j^+ - \rho_j^-) - \partial_2 S_j(\rho_{j+1}^+ - \rho_{j+1}^-) \right\|_1 \\ &= \sum_{j=1}^M \left\| S_{j-1}^- - S_j^- + \frac{\Delta t^\nu}{\tau} \left(\partial_1 S_{j-1}(S_{j-2}^- - S_{j-1}^-) + (\partial_2 S_{j-1} - \partial_1 S_j)(S_{j-1}^- - S_j^-) - \partial_2 S_j(S_j^- - S_{j+1}^-) \right) \right\|_1 \\ &\leq \sum_{j=1}^M \left[\left(1 - \frac{\Delta t^\nu}{\tau} (\partial_1 S_j - \partial_2 S_{j-1}) \right) \left\| S_{j-1}^- - S_j^- \right\|_1 + \frac{\Delta t^\nu}{\tau} \partial_1 S_{j-1} \left\| S_{j-2}^- - S_{j-1}^- \right\|_1 - \frac{\Delta t^\nu}{\tau} \partial_2 S_j \left\| S_j^- - S_{j+1}^- \right\|_1 \right] \\ &= \sum_{j=1}^M \left(1 - \frac{\Delta t^\nu}{\tau} (\partial_1 S_j - \partial_2 S_{j-1}) \right) g_{j,i}^- + \sum_{j=1}^M \frac{\Delta t^\nu}{\tau} \partial_1 S_{j-1} g_{j-1,i}^- - \sum_{j=1}^M \frac{\Delta t^\nu}{\tau} \partial_2 S_j g_{j+1,i}^- \\ &= \sum_{j=1}^M \left(1 - \frac{\Delta t^\nu}{\tau} (\partial_1 S_j - \partial_2 S_{j-1}) \right) g_{j,i}^- + \sum_{j=0}^{M-1} \frac{\Delta t^\nu}{\tau} \partial_1 S_j g_{j,i}^- - \sum_{j=2}^{M+1} \frac{\Delta t^\nu}{\tau} \partial_2 S_{j-1} g_{j,i}^- \\ &= \sum_{j=1}^M g_{j,i}^-, \end{aligned}$$

where, by abuse of notation, we set $g_{0,i}^- = g_{M+1,i}^- = 0$. Moreover, it holds

$$\begin{aligned}g_{j,i}^- - g_{j,i-1}^+ &= \left\| S_{j-1,i}^- - S_{j,i}^- \right\|_1 - \left\| S_{j-1,i-1}^+ - S_{j,i-1}^+ \right\|_1 \\ &\leq \left\| S_{j-1,i}^- - S_{j,i}^- - S_{j-1,i-1}^+ + S_{j,i-1}^+ \right\|_1 \\ &= \left\| \partial_1 S_{j-1}^i(\rho_{j-2,i}^- - \rho_{j-2,i-1}^+) + (\partial_2 S_{j-1}^i - \partial_1 S_j^i)(\rho_{j-1,i}^- - \rho_{j-1,i-1}^+) - \partial_2 S_j^i(\rho_{j,i}^- - \rho_{j,i-1}^+) \right\|_1 \\ &\leq 4SLT\Delta t^\nu,\end{aligned}$$

where we have set $\partial_l S_j^i := \partial_l S_j(s_{j-1}^i, s_j^i)$, $l = 1, 2$, with $s_{j-1}^i = \rho_{j-1,i-1}^+ + \xi_j(\rho_{j-1,i}^- - \rho_{j-1,i-1}^+)$ and $s_j^i = \rho_{j,i-1}^+ + \xi_j(\rho_{j,i}^- - \rho_{j,i-1}^+)$ for some $\xi_j \in [0, 1]$.

Therefore, we get

$$\sum_{j=1}^M \left\| \rho_j^{\tau,\nu}(t_2, \cdot) - \rho_j^{\tau,\nu}(t_1, \cdot) \right\|_1 \leq MLT(t_2 - t_1) + \frac{\Delta t^\nu}{\tau} \sum_{i=k}^{k+N} \sum_{j=1}^M g_{j,i}^-$$

$$\begin{aligned}
&\leq MLT(t_2 - t_1) + \frac{\Delta t^\nu}{\tau} \sum_{i=k}^{k+N} \left(\sum_{j=1}^M g_{j,0}^+ + i4SMLT\Delta t^\nu \right) \\
&\leq MLT(t_2 - t_1) + \frac{t_2 - t_1}{\tau} 4S \sum_{j=1}^M \|\rho_j^0\|_1 + 4SMLT\Delta t^\nu \frac{\Delta t^\nu}{\tau} \sum_{i=k}^{k+N} i \\
&\leq MLT(t_2 - t_1) + \frac{t_2 - t_1}{\tau} 4S \sum_{j=1}^M \|\rho_j^0\|_1 + 8SMLTT \frac{t_2 - t_1}{\tau},
\end{aligned}$$

which gives the desired \mathbf{L}^1 -Lipschitz in time estimate uniform in ν for any $\tau > 0$ fixed. We observe that the estimate blows-up as $\tau \searrow 0$, therefore not allowing to pass to the limit directly in (2.5). \square

Remark 1. *From the above estimates, we note that uniform Lipschitz continuity does not hold even if $\sum_j g_{j,0}^+ = 0$ (i.e. the initial data are at equilibrium).*

Proof of Theorem 1. For any $T > 0$, the previous lemmas provide the compactness estimates on $\{\rho^{\tau,\nu}\}_\nu$ that ensure the existence of a subsequence, still denoted by $\{\rho^{\tau,\nu}\}_\nu$, converging in \mathbf{L}^1 to some function $\rho^\tau \in \mathbf{C}^0([0, T]; \mathbf{L}^1(\mathbb{R}; [0, R_1] \times \dots \times [0, R_M]))$ such that $\text{TV}(\rho^\tau(t)) < +\infty$ for a.e. $t \in [0, T]$ (Helly's Theorem), so that point 1 in Definition 1 is satisfied.

Concerning $\{y_\nu^\tau\}_\nu$, by construction we have that $\dot{y}_\nu^\tau(t) \in [0, V_i]$ for a.e. $t \in [0, T]$ and $\nu \in \mathbb{N}$. Hence, by the Ascoli-Arzelà Theorem, there exists a subsequence, still denoted by $\{y_\nu^\tau\}_\nu$, converging uniformly to a function $y^\tau \in \mathbf{C}^0([0, T]; \mathbb{R})$, which is also Lipschitz of constant V_i . Following [13, Proof of Lemma 3.4], one can prove also that \dot{y}_ν^τ has uniformly bounded total variation on $[0, T]$. Here, in addition to jumps due to interactions with waves coming from the right and jumps in the control function u^ν , we have jumps induced by the relaxation step (2.7) at any $t^n = n\Delta t^\nu$. In this case, \dot{y}_ν^τ can be increasing only if $\rho_i^{\tau,\nu}(\cdot, y_\nu^\tau(\cdot+))$ decreases at $t = t^n$. This increment is bounded by $\frac{\Delta t^\nu}{\tau} \|v_i'\|_\infty (\|S_{i-1}\|_\infty + \|S_i\|_\infty)$, thus $\frac{T}{\tau} \|v_i'\|_\infty (\|S_{i-1}\|_\infty + \|S_i\|_\infty)$ in total (note that this bound is not uniform in τ). Therefore, \dot{y}_ν^τ converges to \dot{y}^τ in $\mathbf{L}_{\text{loc}}^1([0, T]; \mathbb{R})$ and $y^\tau \in \mathbf{W}_{\text{loc}}^{1,1}(\mathbb{R}^+; \mathbb{R})$, see Definition 1, point 2.

Proceeding as in [13, Proof of Theorem 3.1], we get that the (ρ^τ, y^τ) satisfies points 3 and 4 in Definition 1.

We are now left with Definition 1, point 5. To this end, we follow [28, Section 3.3] and [13, Section 3], but note that, in our case, $\hat{\rho}_u = \rho_u^*$, where ρ_u^* is implicitly defined by $v_i(\rho_u^*) = u$, and we have the additional source term in (2.7).

Fix $\bar{t} \in [0, T]$ and, possibly discarding a set of zero measure, without loss of generality assume:

1. $\lim_{\nu \rightarrow \infty} u^\nu(\bar{t}) = u(\bar{t})$, $u(\bar{t}-) = u(\bar{t}+) = \bar{u}$;
2. y^τ continuously differentiable at \bar{t} and $\lim_{\nu \rightarrow \infty} \dot{y}_\nu^\tau(\bar{t}) = \dot{y}^\tau(\bar{t})$;
3. $\dot{y}^{\tau,\nu}(\bar{t}) = \min \{u^\nu(\bar{t}), v_i(\rho_i^{\tau,\nu}(\bar{t}, y_\nu^\tau(\bar{t}+))\}$;
4. $\lim_{\nu \rightarrow \infty} \rho_i^{\tau,\nu}(\bar{t}, x) = \rho_i^\tau(\bar{t}, x)$ for a.e. $x \in \mathbb{R}$, and $t \neq n\Delta t^\nu$ for every $\nu, n \in \mathbb{N}$.

From 4. we have that step (2.7) does not occur at \bar{t} for any ν , thus $\rho_i^{\tau,\nu}$ is constructed only by wave-front tracking on a sufficiently small open interval containing \bar{t} . Therefore, we can

follow the same strategy of [28, Section 3.3] and [13, Section 3].

By 2. and 3., point 5 of Definition 1 is satisfied if

$$\lim_{\nu \rightarrow \infty} \min \{u^\nu(t), v_i(\rho_i^{\tau, \nu}(t, y_\nu^\tau(t)+))\} = \min \{u(t), v_i(\rho_i^\tau(t, y^\tau(t)+))\}. \quad (2.12)$$

Since by 1. $\lim_{\nu \rightarrow \infty} u^\nu(\bar{t}) = u(\bar{t}) = \bar{u}$, (2.12) holds true if $v_i(\rho_i^\tau(\bar{t}, y^\tau(\bar{t})+)) \geq \bar{u}$ and $v_i(\rho_i^{\tau, \nu}(\bar{t}, y_\nu^\tau(\bar{t})+)) \geq \bar{u}$ for infinitely many ν . From 4., if $v_i(\rho_i^{\tau, \nu}(\bar{t}, y_\nu^\tau(\bar{t})+)) < \bar{u}$ for ν sufficiently large, then we also have $v_i(\rho_i^\tau(\bar{t}, y^\tau(\bar{t})+)) < \bar{u}$. Define $\rho_\pm = \lim_{x \rightarrow y^\tau(\bar{t}) \pm} \rho_i^\tau(\bar{t}, x)$, then we can restrict to the case $v_i(\rho_i^\tau(\bar{t}, y^\tau(\bar{t})+)) = v_i(\rho_+) < \bar{u}$, which implies $\rho_+ > \rho_{\bar{u}}^*$.

We distinguish two cases:

1. $\rho_- > \rho_{\bar{u}}^*$: First, notice that Lemma 1 and Lemma 4 of [28] hold true since they are based only on TV bounds that still hold. For some $\epsilon > 0$ sufficiently small, we have that $\rho_- - 2\epsilon > \rho_{\bar{u}}^* = \hat{\rho}_{\bar{u}}$. Then the proof of Lemma 6 of [28] still holds, since it is based on Lemma 4 and the fact that $\rho_- - 2\epsilon > \hat{\rho}_{\bar{u}}$. We conclude in the same way as in [28, Section 3.3.1].
2. $\rho_- \leq \rho_{\bar{u}}^*$: In this case, we have $\rho_- \leq \rho_{\bar{u}}^* < \rho_+$. Following [28, Section 3.3.1], if $y^\tau(\bar{t}) \leq y_\nu^\tau(\bar{t})$ for an infinite number of indices ν , then we conclude by Lemma 4 that (2.12) holds true.

Assume now that $y^\tau(\bar{t}) > y_\nu^\tau(\bar{t})$ for ν sufficiently big. Fix $\epsilon > 0$, then again by Lemma 4 of [28], there exists $\delta > 0$ such that the following holds true. The function $x \rightarrow \rho_i^{\tau, \nu}(\bar{t}, x)$ takes values in $[\rho_- - \epsilon, \rho_- + \epsilon]$ on the interval $[y_\nu^\tau(\bar{t}) - \delta, y_\nu^\tau(\bar{t})]$, and takes values in $[\rho_+ - \epsilon, \rho_+ + \epsilon]$ on the interval $[y^\tau(\bar{t}), y^\tau(\bar{t}) + \delta]$. Since $\rho_i^{\tau, \nu}$ is generated by wave-front tracking, on the interval $I = [y_\nu^\tau(\bar{t}), y^\tau(\bar{t})]$ it may contain only shocks, non-classical shocks, and rarefaction shocks. The only admissible non-classical shock is $(\rho_{\bar{u}}^*, 0)$. If $\rho_i^{\tau, \nu}$ contains a non-classical shock for infinitely many ν , then the non-classical shock must be located at $y^{\tau, \nu}(\bar{t})$ and by 3. $\rho_i^{\tau, \nu}(y^{\tau, \nu}(\bar{t})+) = 0$.

Now, by Lemma 1 of [28], for ν large enough, any rarefaction shock contained in I has strength less than ϵ , thus we deduce that $\rho_i^{\tau, \nu}$ takes values in $[\rho_- - 2\epsilon, \rho_+ + 2\epsilon]$ on the interval I , except possibly to the right of the non-classical shock located at $y^{\tau, \nu}(\bar{t})$. In this case, a classical big shock is present to the right of $y^{\tau, \nu}(\bar{t})$ connecting 0 to a value $\tilde{\rho} \in [\rho_- - 2\epsilon, \rho_+ + 2\epsilon]$. This big shock is followed by small waves, classical shocks or rarefaction shocks, with left and right-hand states in $[\rho_- - 2\epsilon, \rho_+ + 2\epsilon]$. Since these small waves have speed strictly lower than the non-classical shock and the big classical shock, they interact with both of them within time $\delta/(\bar{u} - \bar{\lambda})$, where \bar{u} is the speed of the non-classical shock and $\bar{\lambda}$ is the maximum speed of the small waves, thus $\bar{\lambda} \leq F_j'(\rho_- - 2\epsilon)$. In particular, we can follow the proof of Lemmas 13 and 14 of [28], thus also Lemma 10 of [28] holds true. Since ϵ is arbitrarily small, we can conclude as in [28, Section 3.3.3].

This concludes the proof of Theorem 1. □

3 Asymptotic behaviour without AVs

In this section we assume that no AV is present, thus we focus on the multi-lane model [20]:

$$\begin{cases} \partial_t \rho_j + \partial_x F_j(\rho_j) = \frac{1}{\tau} (S_{j-1}(\rho_{j-1}, \rho_j) - S_j(\rho_j, \rho_{j+1})), & x \in \mathbb{R}, t \geq 0, \\ \rho_j(0, x) = \rho_j^0(x), & j = 1, \dots, M. \end{cases} \quad (3.13)$$

In this case, the definition of entropy weak solution reduces to:

Definition 2. A function $\boldsymbol{\rho}^\tau = (\rho_1^\tau, \dots, \rho_M^\tau) \in \mathbf{C}^0([0, +\infty[; \mathbf{L}^1(\mathbb{R}; [0, R_1] \times \dots \times [0, R_M]))$ is a entropy weak solution of (3.13) with $\boldsymbol{\rho}^0 \in \mathbf{L}^1(\mathbb{R}; [0, R_1] \times \dots \times [0, R_M])$ if for every $\kappa \in \mathbb{R}$ and for all $\varphi \in \mathbf{C}_c^1(\mathbb{R}^2; \mathbb{R}^+)$ it holds

$$\begin{aligned} & \int_{\mathbb{R}^+} \int_{\mathbb{R}} (|\rho_j^\tau - \kappa| \partial_t \varphi + \text{sgn}(\rho_j^\tau - \kappa) (F_j(\rho_j^\tau) - F_j(\kappa)) \partial_x \varphi) dx dt + \int_{\mathbb{R}} |\rho_j^0 - \kappa| \varphi(0, x) dx \\ & + \frac{1}{\tau} \int_{\mathbb{R}^+} \int_{\mathbb{R}} \text{sgn}(\rho_j^\tau - \kappa) (S_{j-1}(\rho_{j-1}^\tau, \rho_j^\tau) - S_j(\rho_j^\tau, \rho_{j+1}^\tau)) \varphi dx dt \geq 0 \end{aligned} \quad (3.14)$$

for $j = 1, \dots, M$.

We extend the \mathbf{L}^1 stability result given in [20, Theorem 3.3] taking in account flux dependency. Besides (3.13), we consider

$$\begin{cases} \partial_t \sigma_j + \partial_x \tilde{F}_j(\sigma_j) = \frac{1}{\tau} (S_{j-1}(\sigma_{j-1}, \sigma_j) - S_j(\sigma_j, \sigma_{j+1})), & x \in \mathbb{R}, t \geq 0, \\ \sigma_j(0, x) = \sigma_j^0(x), & j = 1, \dots, M, \end{cases} \quad (3.15)$$

with \tilde{F} satisfying the same hypotheses as F ($\tilde{F}_j(0) = \tilde{F}_j(R_j) = 0$, Lipschitz and concave).

Lemma 4. Under hypothesis **(S0)**, let $\boldsymbol{\rho}^\tau$ and $\boldsymbol{\sigma}^\tau$ be respectively the entropy weak solutions of (3.13) and (3.15), and let $\mathfrak{T} := \min \{TV(\boldsymbol{\rho}^0), TV(\boldsymbol{\sigma}^0)\}$, $\mathfrak{L} := \max_{j=1, \dots, M} \|F'_j - \tilde{F}'_j\|_\infty$. Then, for a.e. $t > 0$ it holds

$$\sum_{j=1}^M \left\| \rho_j^\tau(t, \cdot) - \sigma_j^\tau(t, \cdot) \right\|_1 \leq \sum_{j=1}^M \left\| \rho_j^0 - \sigma_j^0 \right\|_1 + M \mathfrak{T} \mathfrak{L} t. \quad (3.16)$$

In particular, the above estimate is independent of τ .

Proof. Let us consider WFT approximate solutions constructed by fractional steps as in Section 2. Given a sequence of time steps Δt^ν , $\nu \in \mathbb{N}$, such that $\Delta t^\nu \leq \tau/2\mathcal{S}$, let $\boldsymbol{\rho}^{\tau, \nu}$ and $\boldsymbol{\sigma}^{\tau, \nu}$ the sequences of WFT approximations to (3.13) and (3.15) such that, setting $t^n := n\Delta t^\nu$, $n \in \mathbb{N}$ for any $\nu > 0$ fixed:

- In any time interval $[t^n, t^{n+1}[$, $n \in \mathbb{N}$, $\boldsymbol{\rho}^{\tau, \nu}$ is the WFT approximation of

$$\begin{cases} \partial_t \rho_j + \partial_x F_j(\rho_j) = 0, & x \in \mathbb{R}, t \in [t^n, t^{n+1}[\\ \rho_j(t^n, x) = \rho_j^{\tau, \nu}(t^n, x), & j = 1, \dots, M, \end{cases} \quad (3.17)$$

and $\boldsymbol{\sigma}^{\tau, \nu}$ is the WFT approximation of

$$\begin{cases} \partial_t \sigma_j + \partial_x \tilde{F}_j(\sigma_j) = 0, & x \in \mathbb{R}, t \in [t^n, t^{n+1}[\\ \sigma_j(t^n, x) = \sigma_j^{\tau, \nu}(t^n, x), & j = 1, \dots, M, \end{cases} \quad (3.18)$$

constructed as described e.g. in [19, Section 2.3].

- At time $t = t^{n+1}$, we define

$$\rho_j^{\tau,\nu}(t^{n+1}+, \cdot) = \rho_j^{\tau,\nu}(t^{n+1}-, \cdot) + \frac{\Delta t^\nu}{\tau} \left(S_{j-1}(\rho_{j-1}^{\tau,\nu}, \rho_j^{\tau,\nu}) - S_j(\rho_j^{\tau,\nu}, \rho_{j+1}^{\tau,\nu}) \right) (t^{n+1}-, \cdot), \quad (3.19)$$

$$\sigma_j^{\tau,\nu}(t^{n+1}+, \cdot) = \sigma_j^{\tau,\nu}(t^{n+1}-, \cdot) + \frac{\Delta t^\nu}{\tau} \left(S_{j-1}(\sigma_{j-1}^{\tau,\nu}, \sigma_j^{\tau,\nu}) - S_j(\sigma_j^{\tau,\nu}, \sigma_{j+1}^{\tau,\nu}) \right) (t^{n+1}-, \cdot). \quad (3.20)$$

Solutions of (3.17) and (3.18) satisfy

$$\left\| \rho_j^{\tau,\nu}(t^{n+1}-, \cdot) - \sigma_j^{\tau,\nu}(t^{n+1}-, \cdot) \right\|_1 \leq \left\| \rho_j^{\tau,\nu}(t^n+, \cdot) - \sigma_j^{\tau,\nu}(t^n+, \cdot) \right\|_1 + \mathfrak{T} \mathfrak{L} \Delta t^\nu, \quad (3.21)$$

see e.g. [2, Section 3]. Indeed, by [20, Corollary 3.4], we know that for all $t \geq 0$

$$\mathrm{TV}(\boldsymbol{\rho}^{\tau,\nu}(t, \cdot)) := \sum_{j=1}^M \mathrm{TV}(\rho_j^{\tau,\nu}(t, \cdot)) \leq \sum_{j=1}^M \mathrm{TV}(\rho_j^{0,\nu}) \leq \sum_{j=1}^M \mathrm{TV}(\rho_j^0) =: \mathrm{TV}(\boldsymbol{\rho}^0).$$

Subtracting (3.20) from (3.19) and setting $\rho_{j,\nu}^{\tau}(t^{n+1}-, \cdot) := \rho_j$ and $\sigma_{j,\nu}^{\tau}(t^{n+1}-, \cdot) := \sigma_j$ for $j = 1, \dots, M$, we get (using notations introduced earlier for the partial derivatives of the terms S_j)

$$\begin{aligned} & \rho_j^{\tau,\nu}(t^{n+1}+, \cdot) - \sigma_j^{\tau,\nu}(t^{n+1}+, \cdot) \\ &= \rho_j - \sigma_j + \frac{\Delta t^\nu}{\tau} \left(S_{j-1}(\rho_{j-1}, \rho_j) - S_j(\rho_j, \rho_{j+1}) - S_{j-1}(\sigma_{j-1}, \sigma_j) + S_j(\sigma_j, \sigma_{j+1}) \right) \\ &= \rho_j - \sigma_j + \frac{\Delta t^\nu}{\tau} \left(\partial_1 S_{j-1}(\rho_{j-1} - \sigma_{j-1}) + \partial_2 S_{j-1}(\rho_j - \sigma_j) - \partial_1 S_j(\rho_j - \sigma_j) - \partial_2 S_j(\rho_{j+1} - \sigma_{j+1}) \right) \\ &= (\rho_j - \sigma_j) \left[1 - \frac{\Delta t^\nu}{\tau} (\partial_1 S_j - \partial_2 S_{j-1}) \right] + \frac{\Delta t^\nu}{\tau} \left(\partial_1 S_{j-1}(\rho_{j-1} - \sigma_{j-1}) - \partial_2 S_j(\rho_{j+1} - \sigma_{j+1}) \right). \end{aligned}$$

Since $\partial_1 S_{j-1} \geq 0$ and $\partial_2 S_j \leq 0$ by hypothesis and $1 - \frac{\Delta t^\nu}{\tau} (\partial_1 S_j - \partial_2 S_{j-1}) \geq 0$ for Δt^ν sufficiently small, we conclude that if $\rho_j^{\tau,\nu}(t^{n+1}-, \cdot) \leq \sigma_j^{\tau,\nu}(t^{n+1}-, \cdot)$ for all $j = 1, \dots, M$, then $\rho_j^{\tau,\nu}(t^{n+1}+, \cdot) \leq \sigma_j^{\tau,\nu}(t^{n+1}+, \cdot)$ for all $j = 1, \dots, M$. Then, by the Crandall-Tartar lemma [19, Lemma 2.13], we get

$$\sum_{j=1}^M \left\| \rho_j^{\tau,\nu}(t^{n+1}+, \cdot) - \sigma_j^{\tau,\nu}(t^{n+1}+, \cdot) \right\|_1 \leq \sum_{j=1}^M \left\| \rho_j^{\tau,\nu}(t^{n+1}-, \cdot) - \sigma_j^{\tau,\nu}(t^{n+1}-, \cdot) \right\|_1. \quad (3.22)$$

Putting together (3.21) and (3.22), we obtain the desired estimate. \square

Let us now assume:

- (V) $v_j = v$ (in particular, the maximal speeds become $V_j = V$) and $R_j = R/M$ the maximal density on each lane (so that $v(R/M) = 0$), and thus $F_j(\rho) = F(\rho) = \rho v(\rho)$, for all $j = 1, \dots, M$.

Lemma 5. *Under hypotheses (S0), (S1) and (V), let the initial data $\rho_j^0 \in (\mathbf{L}^1 \cap BV)(\mathbb{R}; [0, R/M])$ for $j = 1, \dots, M$, be at equilibrium, i.e.*

$$S_{j-1}(\rho_{j-1}^0(x), \rho_j^0(x)) = S_j(\rho_j^0(x), \rho_{j+1}^0(x)) = 0, \quad x \in \mathbb{R}, \quad j = 1, \dots, M. \quad (3.23)$$

Then the solution stays at equilibrium, i.e.

$$S_{j-1}(\rho_{j-1}^\tau(t, x), \rho_j^\tau(t, x)) = S_j(\rho_j^\tau(t, x), \rho_{j+1}^\tau(t, x)) = 0, \quad j = 1, \dots, M,$$

for all $x \in \mathbb{R}$, $t > 0$. In particular, the solution is \mathbf{L}^1 -Lipschitz continuous in time, uniformly in τ :

$$\left\| \rho_j^\tau(t_1, \cdot) - \rho_j^\tau(t_2, \cdot) \right\|_1 \leq K |t_1 - t_2|, \quad j = 1, \dots, M, \quad (3.24)$$

for any $t_1, t_2 \in \mathbb{R}^+$, with $K = L\mathcal{T}$.

The proof is trivial observing that system (3.13) reduces to M identical equations with the same initial datum, and therefore $\rho_1^\tau(t, \cdot) = \dots = \rho_M^\tau(t, \cdot) := \bar{\rho}^\tau$ for all $t > 0$.

The uniform Lipschitz estimate (3.24) and the uniform TV bound (2.8) allow to apply Helly's Theorem to the sequence $\{\rho^\tau\}_{\tau > 0}$, proving the existence of a subsequence converging in $\mathbf{L}_{\text{loc}}^1$ (indeed, the sequence is constant for all $\tau > 0$, so there is nothing to prove). In particular, the limit $\bar{\rho} = (\bar{\rho}_1, \dots, \bar{\rho}_M)$ satisfies $\bar{\rho}_j = \bar{\rho}^\tau$ for any $\tau > 0$ and it holds

$$\partial_t \bar{\rho}_j + \partial_x F(\bar{\rho}_j) = 0, \quad j = 1, \dots, M,$$

so that, setting $r := \sum_{j=1}^M \bar{\rho}_j = M\bar{\rho}^\tau$, r is a entropy weak solution to the scalar Cauchy problem

$$\begin{cases} \partial_t r + \partial_x f(r) = 0, & t > 0, \quad x \in \mathbb{R}, \\ r(0, x) = \sum_{j=1}^M \rho_j^0(x), & x \in \mathbb{R}, \end{cases} \quad (3.25)$$

where we have set $f(r) := MF(r/M)$.

We can now give the following result about the convergence of (3.13) to (3.25) as $\tau \searrow 0$.

Theorem 2. *Under hypotheses (S0), (S1) and (V), let $\rho_j^0 \in (\mathbf{L}^1 \cap BV)(\mathbb{R}; [0, R/M])$ and $\rho^\tau = (\rho_1^\tau, \dots, \rho_M^\tau)$ be the corresponding solution of (3.13). Then, for each $t > 0$,*

$$\rho_j^\tau(t, \cdot) \longrightarrow \frac{r(t, \cdot)}{M} \quad \text{in } \mathbf{L}_{\text{loc}}^1(\mathbb{R}; \mathbb{R})$$

as $\tau \searrow 0$, where $r \in \mathbf{C}^0(\mathbb{R}^+; \mathbf{L}^1(\mathbb{R}; [0, R]))$ is the entropy weak solution of (3.25).

Proof. We adapt an argument by Bressan and Shen [4, Section 5]. Denote by $\tilde{\rho}^\tau = (\tilde{\rho}_1^\tau, \dots, \tilde{\rho}_M^\tau)$ the solution of

$$\partial_t \rho_j = \frac{1}{\tau} (S_{j-1}(\rho_{j-1}, \rho_j) - S_j(\rho_j, \rho_{j+1})), \quad j = 1, \dots, M, \quad (3.26)$$

with the same initial data ρ_j^0 . Linearizing (3.26) around the equilibrium initial condition $\rho(0, x) = (\rho_1(0, x), \dots, \rho_M(0, x))$ with $\rho_j(0, x) := r(0, x)/M$ for $j = 1, \dots, M$ we obtain

$$\partial_t \rho = \frac{1}{\tau} (\mathbb{S}(\rho(0, \cdot)) + D\mathbb{S}(\rho(0, \cdot))(\rho - \rho(0, \cdot))),$$

where we have set $\mathbb{S}_j(\boldsymbol{\rho}(0, \cdot)) := S_{j-1}(r(0, \cdot)/M, r(0, \cdot)/M) - S_j(r(0, \cdot)/M, r(0, \cdot)/M) = 0$ and

$$D\mathbb{S}(\boldsymbol{\rho}) = \begin{bmatrix} -\partial_1 S_1 & -\partial_2 S_1 & 0 & \cdots & \cdots & 0 \\ \partial_1 S_1 & \partial_2 S_1 - \partial_1 S_2 & -\partial_2 S_2 & 0 & \cdots & \vdots \\ 0 & \partial_1 S_2 & \partial_2 S_2 - \partial_1 S_3 & -\partial_2 S_3 & 0 & \vdots \\ \vdots & 0 & \partial_1 S_3 & \ddots & & \\ 0 & \cdots & \cdots & 0 & \partial_1 S_{M-1} & \partial_2 S_{M-1} \end{bmatrix}$$

is a tridiagonal matrix such that

$$D\mathbb{S}(\boldsymbol{\rho})_{i,j} = \begin{cases} \partial_1 S_j & i = j + 1, \\ \partial_2 S_{j-1} - \partial_1 S_j & i = j, \\ -\partial_2 S_{j-1} & i = j - 1. \end{cases}$$

The similarity transformation to a symmetric matrix $\mathbb{A}(\boldsymbol{\rho})$ gives

$$\mathbb{A}(\boldsymbol{\rho})_{i,j} = \begin{cases} \sqrt{-\partial_1 S_j \partial_2 S_j} & i = j + 1, \\ \partial_2 S_{j-1} - \partial_1 S_j & i = j, \\ \sqrt{-\partial_1 S_{j-1} \partial_2 S_{j-1}} & i = j - 1, \end{cases}$$

which, by the monotonicity hypotheses in (\mathbf{S}) , turns out to be negative definite (this can be easily seen applying Sylvester's criterion). Therefore,

$$\sum_{j=1}^M \left| \tilde{\rho}_j^\tau(t, x) - \frac{r(0, x)}{M} \right| = \mathcal{O}(1) \cdot e^{-t/\tau}, \quad x \in \mathbb{R}. \quad (3.27)$$

Moreover, for any interval $[a, b] \subset \mathbb{R}$, comparison between (3.13) and (3.26) gives, by Lemma 4,

$$\sum_{j=1}^M \left\| \tilde{\rho}_j^\tau(t, \cdot) - \rho_j^\tau(t, \cdot) \right\|_{\mathbf{L}^1([a,b])} = \mathcal{O}(1) \cdot (b - a) t. \quad (3.28)$$

Taking $t = \sqrt{\tau}$ in (3.27) and (3.28) we get

$$\sum_{j=1}^M \left\| \rho_j^\tau(\sqrt{\tau}, \cdot) - \frac{r(0, \cdot)}{M} \right\|_{\mathbf{L}^1([a,b])} = \mathcal{O}(1) \cdot (b - a) \left(\sqrt{\tau} + e^{-1/\sqrt{\tau}} \right).$$

Taking $\boldsymbol{\rho}^\tau(\sqrt{\tau}, \cdot)$ as initial datum in (3.13), the stability estimate (3.16) gives

$$\sum_{j=1}^M \left\| \rho_j^\tau(t, \cdot) - \frac{r(t - \sqrt{\tau}, \cdot)}{M} \right\|_{\mathbf{L}^1([a,b])} = \mathcal{O}(1) \cdot (b - a) \left(\sqrt{\tau} + e^{-1/\sqrt{\tau}} \right)$$

for all $t \geq \sqrt{\tau}$. By the \mathbf{L}^1 -Lipschitz continuity in time of r guaranteed by Lemma 5, we finally get

$$\sum_{j=1}^M \left\| \rho_j^\tau(t, \cdot) - \frac{r(t, \cdot)}{M} \right\|_{\mathbf{L}^1([a,b])} = \mathcal{O}(1) \cdot (b - a) \left(\sqrt{\tau} + e^{-1/\sqrt{\tau}} \right)$$

for all $t > 0$. □

4 Asymptotic behaviour for $M = 2$ with AV

We now focus on the case of a two-lane road with one AV, say on lane $i = 1$ (without loss of generality by symmetry), which reads:

$$\begin{cases} \partial_t \rho_1 + \partial_x F_1(\rho_1) = -\frac{1}{\tau} S(\rho_1, \rho_2), \\ \partial_t \rho_2 + \partial_x F_2(\rho_2) = \frac{1}{\tau} S(\rho_1, \rho_2), \\ \rho_j(0, x) = \rho_j^0(x), \quad j = 1, 2, \\ \dot{y}(t) = \min \{u(t), v_1(\rho_1(t, y(t)+))\}, \\ y(0) = y_0, \\ \rho_1(t, y(t)) (v_1(\rho_1(t, y(t))) - \dot{y}(t)) \leq 0, \end{cases} \quad x \in \mathbb{R}, \quad t \geq 0, \quad (4.29)$$

where we have set $S(\rho_1, \rho_2) := S_1(\rho_1, \rho_2)$.

In this case, Lemma 3 becomes:

Lemma 6. *Let (S0) and (S2) hold and $\Delta t^\nu \leq \tau/2S$, then for any $t_1 \leq t_2 \in [0, T]$, $T > 0$, the approximate solutions $\rho^{\tau, \nu}$ satisfy*

$$\sum_{j=1}^2 \left\| \rho_j^{\tau, \nu}(t_1, \cdot) - \rho_j^{\tau, \nu}(t_2, \cdot) \right\|_1 \leq (C + C_\tau(t_1)) |t_1 - t_2|, \quad \text{for all } \nu \in \mathbb{N}, \quad (4.30)$$

where $C = 2LT(1 + S/c)$ is independent of τ and

$$C_\tau(t_1) = \frac{4S}{\tau} \exp\left(-\frac{2c}{\tau} t_1\right) \sum_{j=1}^2 \left\| \rho_j^0 \right\|_1.$$

Proof. We follow the same procedure and notations as in the proof of Lemma 3. Setting

$$g_{1,i}^\pm = g_{2,i}^\pm := \left\| S(\rho_{1,\nu}^\tau, \rho_{2,\nu}^\tau)(t^i \pm, \cdot) \right\|_1,$$

we compute

$$\begin{aligned} g_{1,i}^+ + g_{2,i}^+ &= 2 \left\| S(\rho_1^+, \rho_2^+) \right\|_1 \\ &= 2 \left\| S(\rho_1^-, \rho_2^-) + \partial_1 S \cdot (\rho_1^+ - \rho_1^-) + \partial_2 S \cdot (\rho_2^+ - \rho_2^-) \right\|_1 \\ &= 2 \left\| S(\rho_1^-, \rho_2^-) - \frac{\Delta t^\nu}{\tau} \partial_1 S \cdot S(\rho_1^-, \rho_2^-) + \frac{\Delta t^\nu}{\tau} \partial_2 S \cdot S(\rho_1^-, \rho_2^-) \right\|_1 \\ &= 2 \left(1 - \frac{\Delta t^\nu}{\tau} (\partial_1 S - \partial_2 S) \right) \left\| S(\rho_1^-, \rho_2^-) \right\|_1 \\ &= \left(1 - \frac{\Delta t^\nu}{\tau} (\partial_1 S - \partial_2 S) \right) (g_{1,i}^- + g_{2,i}^-) \\ &\leq \left(1 - \frac{2c}{\tau} \Delta t^\nu \right) (g_{1,i}^- + g_{2,i}^-). \end{aligned}$$

Moreover, for $j = 1, 2$, it holds

$$\begin{aligned}
g_{j,i}^- - g_{j,i-1}^+ &= \left\| S(\rho_{1,i}^-, \rho_{2,i}^-) \right\|_1 - \left\| S(\rho_{1,i-1}^+, \rho_{2,i-1}^+) \right\|_1 \\
&\leq \left\| S(\rho_{1,i}^-, \rho_{2,i}^-) - S(\rho_{1,i-1}^+, \rho_{2,i-1}^+) \right\|_1 \\
&= \left\| \partial_1 S^i \cdot (\rho_{1,i}^- - \rho_{1,i-1}^+) + \partial_2 S^i \cdot (\rho_{2,i}^- - \rho_{2,i-1}^+) \right\|_1 \\
&\leq 2SLT\Delta t^\nu,
\end{aligned}$$

where we have set $\partial_l S^i := \partial_l S(s_1^i, s_2^i)$, $l = 1, 2$, with $s_1^i = \rho_{1,i-1}^+ + \xi(\rho_{1,i}^- - \rho_{1,i-1}^+)$ and $s_2^i = \rho_{2,i-1}^+ + \xi(\rho_{2,i}^- - \rho_{2,i-1}^+)$ for some $\xi \in [0, 1]$.

Therefore, we get

$$\begin{aligned}
\sum_{j=1}^M \left\| \rho_j^{\tau,\nu}(t_2, \cdot) - \rho_j^{\tau,\nu}(t_1, \cdot) \right\|_1 &\leq 2LT(t_2 - t_1) + \frac{\Delta t^\nu}{\tau} \sum_{i=k}^{k+N} \sum_{j=1}^2 g_{j,i}^- \\
&\leq 2LT(t_2 - t_1) + \frac{\Delta t^\nu}{\tau} \sum_{i=k}^{k+N} \left[\left(1 - \frac{2c}{\tau} \Delta t^\nu\right)^{i-1} \sum_{j=1}^2 g_{j,0}^+ + 4SLT\Delta t^\nu \sum_{\ell=0}^{i-1} \left(1 - \frac{2c}{\tau} \Delta t^\nu\right)^\ell \right] \\
&\leq 2LT(t_2 - t_1) + \frac{\Delta t^\nu}{\tau} \sum_{i=k}^{k+N} \left[\left(1 - \frac{2c}{\tau} \Delta t^\nu\right)^{i-1} \sum_{j=1}^2 g_{j,0}^+ + 2SLT\frac{\tau}{c} \right] \\
&\leq 2LT(t_2 - t_1) + 2SLT\frac{t_2 - t_1 + \Delta t^\nu}{c} + \frac{\Delta t^\nu}{\tau} \sum_{i=k}^{k+N} \left[\left(1 - \frac{2c}{\tau} \Delta t^\nu\right)^{i-1} \sum_{j=1}^2 g_{j,0}^+ \right] \\
&\leq 2LT(t_2 - t_1 + \Delta t^\nu) \left(1 + \frac{\mathcal{S}}{c}\right) + \frac{1}{2c} \left(1 - \frac{2c}{\tau} \Delta t^\nu\right)^{k-1} \left[1 - \left(1 - \frac{2c}{\tau} \Delta t^\nu\right)^{N+1}\right] \sum_{j=1}^2 g_{j,0}^+ \\
&\leq 2LT(t_2 - t_1 + \Delta t^\nu) \left(1 + \frac{\mathcal{S}}{c}\right) + \frac{1}{c} \left(1 - \frac{2c}{\tau} \Delta t^\nu\right)^{\frac{k\Delta t^\nu}{\Delta t^\nu}} (N+1) \frac{2c}{\tau} \Delta t^\nu \sum_{j=1}^2 g_{j,0}^+ \\
&\leq (t_2 - t_1 + \Delta t^\nu) \left[2LT \left(1 + \frac{\mathcal{S}}{c}\right) + \frac{4\mathcal{S}}{\tau} \exp\left(-\frac{2c}{\tau} t_1\right) \sum_{j=1}^2 \left\| \rho_j^0 \right\|_1 \right],
\end{aligned}$$

where we have used the inequality $\left(1 - \frac{2c}{\tau} \Delta t^\nu\right) \geq \frac{1}{2}$, which holds provided that $\Delta t^\nu \leq \tau/4c$. \square

Remark 2. By Lemma 6, the Lipschitz constant $C_\tau(t_1)$ is uniformly bounded as $\tau \searrow 0$ for any strictly positive time. Indeed, taking $\frac{1}{n} \leq t_1 \leq t_2$ with $n \in \mathbb{N}$ in (4.30), we have

$$C_\tau(t_1) \leq \frac{4\mathcal{S}}{\tau} \exp\left(-\frac{2c}{n\tau}\right) \sum_{j=1}^2 \left\| \rho_j^0 \right\|_1, \quad (4.31)$$

which goes to zero with τ .

Note also that, under the hypotheses of Lemma 6, if the initial datum is at equilibrium, i.e. $S(\rho_1^0(x), \rho_2^0(x)) = 0$ for a.e. $x \in \mathbb{R}$, it holds $C_\tau(t) = 0$ for all $t \geq 0$. Therefore, the

approximate solutions $\rho^{\tau, \nu}$ are uniformly Lipschitz continuous for all $t \geq 0$:

$$\sum_{j=1}^2 \left\| \rho_j^{\tau, \nu}(t_1, \cdot) - \rho_j^{\tau, \nu}(t_2, \cdot) \right\|_1 \leq C |t_1 - t_2|, \quad \text{for all } \nu \in \mathbb{N}, \tau > 0. \quad (4.32)$$

We can now prove the following convergence result.

Theorem 3. *Under hypotheses (S0), (S1) and (S2), let $v_1 = v_2 = v$ (thus $F_1(\rho) = F_2(\rho) = F(\rho) = \rho v(\rho)$) and $\rho_1^0, \rho_2^0 \in (\mathbf{L}^1 \cap BV)(\mathbb{R}; [0, R/2])$ and $\rho^\tau = (\rho_1^\tau, \rho_2^\tau)$ and y^τ be the corresponding entropy weak solution of (4.29). Then*

$$\begin{aligned} \rho_j^\tau &\longrightarrow \frac{r}{2} && \text{in } \mathbf{L}_{\text{loc}}^1(]0, +\infty[\times \mathbb{R}; \mathbb{R}), \quad i = 1, 2, \\ y^\tau &\longrightarrow y && \text{in } \mathbf{L}_{\text{loc}}^\infty(\mathbb{R}^+; \mathbb{R}), \end{aligned}$$

as $\tau \searrow 0$, where $y \in \mathbf{C}^0(\mathbb{R}^+; \mathbb{R})$ and $r \in \mathbf{C}^0(\mathbb{R}^+; \mathbf{L}^1(\mathbb{R}; [0, R]))$ is a weak solution of (3.25) with $M = 2$. If moreover $\dot{y}^\tau \longrightarrow \dot{y}$ in $\mathbf{L}_{\text{loc}}^1(\mathbb{R}^+; \mathbb{R})$, then (r, y) satisfies

$$\begin{aligned} &\int_{\mathbb{R}^+} \int_{\mathbb{R}} (|r - k| \partial_t \varphi + \text{sgn}(r - k)(f(r) - f(k)) \partial_x \varphi) dx dt \\ &+ \int_{\mathbb{R}^+} (f(k) - \dot{y}(t)k)^+ \varphi(t, y(t)) dt \geq 0, \end{aligned} \quad (4.33)$$

with $f(r) := 2F(r/2)$.

Remark 3. *We notice that, as pointed out in the proof of Theorem 1, we need to assume a-priori the $\mathbf{W}^{1,1}$ convergence of the AV trajectories y^ν to y .*

Theorem 3 shows in particular that (r, y) cannot be an entropy weak solution of the moving bottleneck model [9, 13] with $\alpha = 1/2$, which writes

$$\begin{cases} \partial_t r + \partial_x f(r) = 0, \\ r(0, x) = \rho_1^0(x) + \rho_2^0(x), \\ \dot{y}(t) = \min \{u(t), \tilde{v}(r(t, y(t)+))\}, \\ y(0) = y_0, \\ f(r(t, y(t))) - \dot{y}(t) r(t, y(t)) \leq \mathcal{F}_{1/2}(\dot{y}(t)) := \max_{r \in [0, R]} \left(\frac{1}{2} f(2r) - r \dot{y}(t) \right), \end{cases} \quad x \in \mathbb{R}, t \geq 0, \quad (4.34)$$

where $\tilde{v}(r) := f(r)/r = v(r/2)$, whose entropy weak solution satisfies

$$\begin{aligned} &\int_{\mathbb{R}^+} \int_{\mathbb{R}} (|r - k| \partial_t \varphi + \text{sgn}(r - k)(f(r) - f(k)) \partial_x \varphi) dx dt \\ &+ 2 \int_{\mathbb{R}^+} \left(f(k) - k \dot{y}(t) - \mathcal{F}_{1/2}(\dot{y}(t)) \right)^+ \varphi(t, y(t)) dt \geq 0. \end{aligned} \quad (4.35)$$

Observe that we can estimate

$$\begin{aligned} &2 \left(f(k) - k \dot{y}(t) - \mathcal{F}_{1/2}(\dot{y}(t)) \right)^+ \\ &= 2 \left(f(k) - k \dot{y}(t) - \max_{r \in [0, R]} \left(\frac{1}{2} f(2r) - r \dot{y}(t) \right) \right)^+ \end{aligned}$$

$$\begin{aligned}
&= \left(2f(k) - 2k \dot{y}(t) - \max_{r \in [0, R]} (f(2r) - 2r \dot{y}(t)) \right)^+ \\
&= \left(2f(k) - 2k \dot{y}(t) + \min_{r \in [0, R]} (-f(2r) + 2r \dot{y}(t)) \right)^+ \\
&\leq \left(2f(k) - 2k \dot{y}(t) + (-f(k) + k \dot{y}(t)) \right)^+ \\
&= (f(k) - k \dot{y}(t))^+
\end{aligned}$$

taking $r = k/2$. This shows that (4.35) implies (4.33), but the converse is not true. See also the numerical examples in Section 5.3.

Proof of Theorem 3. By Helly's Theorem, for any $n \in \mathbb{N}$ there exists a subsequence $\{\tau_n\} \subseteq \{\tau_{n-1}\}$ such that $\boldsymbol{\rho}^{\tau_n}$ converges to some function $\bar{\boldsymbol{\rho}}$ in $\mathbf{L}_{\text{loc}}^1([1/n, +\infty[\times \mathbb{R})$. We can then construct by diagonal process a sequence, still labeled $\{\boldsymbol{\rho}^{\tau_n}\}$, converging to $\bar{\boldsymbol{\rho}}$ in $\mathbf{L}_{\text{loc}}^1(]0, +\infty[\times \mathbb{R})$. From (4.30) and (4.31), the limit satisfies

$$\sum_{j=1}^2 \|\bar{\rho}_j(t_1, \cdot) - \bar{\rho}_j(t_2, \cdot)\|_1 \leq C|t_1 - t_2|, \quad \text{for any } t_1, t_2 > 0. \quad (4.36)$$

From (2.4) we deduce that

$$S(\bar{\rho}_1, \bar{\rho}_2) = 0 \quad \text{a.e. in }]0, +\infty[\times \mathbb{R}.$$

By **(S1)**, this implies $\bar{\rho}_1 = \bar{\rho}_2 = \bar{\rho}$. Moreover, since $\{\boldsymbol{\rho}^{\tau_n}\}$ is equi-bounded, it actually converges in $\mathbf{L}_{\text{loc}}^1(]0, +\infty[\times \mathbb{R})$. By (4.36) we can then pass to the limit in $\mathbf{L}_{\text{loc}}^1$ concluding that

$$\lim_{t \rightarrow 0^+} S(\bar{\rho}_1(t, \cdot), \bar{\rho}_2(t, \cdot)) = 0.$$

Summing (2.4) for $j = 1, 2$ and $\tau = \tau_n$ with $\varphi \in \mathbf{C}_c^1(]0, +\infty[\times \mathbb{R}; \mathbb{R}^+)$, we obtain

$$\begin{aligned}
&\int_{\mathbb{R}^+} \int_{\mathbb{R}} (|\rho_1^\tau - \kappa| + |\rho_2^\tau - \kappa|) \partial_t \varphi \, dx \, dt \\
&\quad + \int_{\mathbb{R}^+} \int_{\mathbb{R}} \left(\text{sgn}(\rho_1^\tau - \kappa) (F(\rho_1^\tau) - F(\kappa)) + \text{sgn}(\rho_2^\tau - \kappa) (F(\rho_2^\tau) - F(\kappa)) \right) \partial_x \varphi \, dx \, dt \\
&\quad + \frac{1}{\tau} \int_{\mathbb{R}^+} \int_{\mathbb{R}} (\text{sgn}(\rho_2^\tau - \kappa) - \text{sgn}(\rho_1^\tau - \kappa)) S_j(\rho_1^\tau, \rho_2^\tau) \varphi \, dx \, dt \\
&\quad + 2 \int_{\mathbb{R}^+} (F(\kappa) - \dot{y}^\tau(t) \kappa)^+ \varphi(t, y^\tau(t)) \, dt \geq 0
\end{aligned} \quad (4.37)$$

for all $\kappa \in \mathbb{R}$. We note that

$$(\text{sgn}(\rho_2^\tau - \kappa) - \text{sgn}(\rho_1^\tau - \kappa)) S_j(\rho_1^\tau, \rho_2^\tau) \leq 0 \quad \text{for every } \kappa \in \mathbb{R}.$$

Indeed, this is obviously true if $\text{sgn}(\rho_2^\tau - \kappa) = \text{sgn}(\rho_1^\tau - \kappa)$. If $\rho_1^\tau \leq \kappa \leq \rho_2^\tau$, we get

$$(\text{sgn}(\rho_2^\tau - \kappa) - \text{sgn}(\rho_1^\tau - \kappa)) S_j(\rho_1^\tau, \rho_2^\tau) = 2S_j(\rho_1^\tau, \rho_2^\tau) \leq S_j(\kappa, \kappa) = 0,$$

due to the monotonicity in **(S0)** and **(S1)**. The same holds for $\rho_2^\tau \leq \kappa \leq \rho_1^\tau$. Therefore, (4.37) implies

$$\begin{aligned} & \int_{\mathbb{R}^+} \int_{\mathbb{R}} (|\rho_1^\tau - \kappa| + |\rho_2^\tau - \kappa|) \partial_t \varphi \, dx \, dt \\ & + \int_{\mathbb{R}^+} \int_{\mathbb{R}} \left(\operatorname{sgn}(\rho_1^\tau - \kappa) (F(\rho_1^\tau) - F(\kappa)) + \operatorname{sgn}(\rho_2^\tau - \kappa) (F(\rho_2^\tau) - F(\kappa)) \right) \partial_x \varphi \, dx \, dt \\ & + 2 \int_{\mathbb{R}^+} (F(\kappa) - \dot{y}^\tau(t) \kappa)^+ \varphi(t, y^\tau(t)) \, dt \geq 0. \end{aligned} \quad (4.38)$$

Passing to the limit as $\tau_n \searrow 0$ in (4.38), thanks to the hypothesis of **L¹** convergence of \dot{y}^ν to \dot{y} , we get that $\bar{\rho}$ satisfies

$$\begin{aligned} & \int_{\mathbb{R}^+} \int_{\mathbb{R}} (|2\bar{\rho} - 2\kappa| \partial_t \varphi + \operatorname{sgn}(\bar{\rho} - \kappa) (2F(2\bar{\rho}/2) - 2F(2\kappa/2)) \partial_x \varphi) \, dx \, dt \\ & + \int_{\mathbb{R}^+} (2F(2\kappa/2) - \dot{y}(t) 2\kappa)^+ \varphi(t, y(t)) \, dt \geq 0, \end{aligned}$$

which, setting $r = 2\bar{\rho}$, $k = 2\kappa$ and $f(r) = 2F(r/2)$, is equivalent to (4.33).

Finally, taking $\kappa \notin [0, R/2]$ in (2.4) we get for $j = 1, 2$ and $\tau = \tau_n$,

$$\begin{aligned} & \int_{\mathbb{R}^+} \int_{\mathbb{R}} \left(\rho_j^\tau \partial_t \varphi + F(\rho_j^\tau) \partial_x \varphi \right) \, dx \, dt + \int_{\mathbb{R}} \rho_j^0(x) \varphi(0, x) \, dx \\ & + \frac{1}{\tau} \int_{\mathbb{R}^+} \int_{\mathbb{R}} \left(S_{j-1}(\rho_{j-1}^\tau, \rho_j^\tau) - S_j(\rho_j^\tau, \rho_{j+1}^\tau) \right) \varphi \, dx \, dt = 0 \end{aligned}$$

for all $\varphi \in \mathbf{C}_c^1(\mathbb{R}^2; \mathbb{R})$. Passing to the limit as $\tau_n \searrow 0$ and summing over $j = 1, 2$, we get

$$\begin{aligned} & \int_{\mathbb{R}^+} \int_{\mathbb{R}} (2\bar{\rho} \partial_t \varphi + 2F(2\bar{\rho}/2) \partial_x \varphi) \, dx \, dt + \int_{\mathbb{R}} (\rho_1^0(x) + \rho_2^0(x)) \varphi(0, x) \, dx \\ & = \int_{\mathbb{R}^+} \int_{\mathbb{R}} (r \partial_t \varphi + f(r) \partial_x \varphi) \, dx \, dt + \int_{\mathbb{R}} (\rho_1^0(x) + \rho_2^0(x)) \varphi(0, x) \, dx = 0, \end{aligned}$$

concluding the proof. \square

5 Numerical simulations

To illustrate the behaviour of solutions of model (2.3), we compute approximate solutions obtained via a finite volume scheme with time splitting developed merging the numerical schemes proposed in [16] for multi-lane models and in [5] for correctly capturing the dynamics around moving bottlenecks. See also [14, 17, 37] for further extensions.

We consider a uniform space mesh of width Δx and a time step Δt , subject to the stability condition

$$\Delta t \leq \min \left\{ \frac{\Delta x}{\|F_j'\|_\infty}, \frac{\tau}{2S} \right\},$$

see Lemma 1. We set $x_k = (k + 1/2) \Delta x$ and $x_{k-1/2} = k \Delta x$ for $k \in \mathbb{Z}$, where x_k is the center of cell $C_k = [x_{k-1/2}, x_{k+1/2}]$ and $x_{k\pm 1/2}$ are the interfaces. Set $\lambda = \Delta t / \Delta x$.

We approximate the initial data as

$$\rho_{j,k}^0 = \frac{1}{\Delta x} \int_{x_{k-1/2}}^{x_{k+1/2}} \rho_j^0(x) dx, \quad \text{for } j = 1, \dots, M.$$

The approximate solutions to (1.1) are obtained through a Godunov type scheme with fractional step accounting for the contribution of the source term: for $j = 1, \dots, M$, and $n \in \mathbb{N}$, we set

$$\rho_{j,k}^{n+1/2} = \rho_{j,k}^n - \lambda \left[\mathcal{F}_j(\rho_{j,k}^n, \rho_{j,k+1}^n) - \mathcal{F}_j(\rho_{j,k-1}^n, \rho_{j,k}^n) \right], \quad (5.39)$$

$$\rho_{j,k}^{n+1} = \rho_{j,k}^{n+1/2} + \frac{\Delta t}{\tau} \left[S_{j-1}(\rho_{j-1,k}^{n+1/2}, \rho_{j,k}^{n+1/2}) - S_j(\rho_{j,k}^{n+1/2}, \rho_{j+1,k}^{n+1/2}) \right],$$

where

$$\mathcal{F}_j(u, w) = \min \{ D_j(u), S_j(w) \}$$

with

$$D_j(u) = F_j(\min\{u, \theta_j\}), \quad S_j(u) = F_j(\max\{u, \theta_j\}),$$

$\theta_j \in [0, R_j]$ being the point of maximum of F_j .

To capture the presence of an AV, let's say at position $y_i^n \in C_m$ in lane $i \in \{1, \dots, M\}$, we modify the Godunov fluxes at the interfaces $x_{m\pm 1/2}$ as follows. We denote by \mathcal{R}_i the standard Riemann solver for $\partial_t \rho + \partial_x F_i(\rho) = 0$, i.e. $\mathcal{R}_i(\rho_L, \rho_R)(\xi)$ gives the value of the self-similar entropy weak solution corresponding to the Riemann initial data (ρ_L, ρ_R) at $x = \xi t$. We also set

$$u_i^n = \frac{1}{\Delta t} \int_{t^n}^{t^{n+1}} u_i(t) dt, \quad \text{for } n \in \mathbb{N},$$

the discretization of the AV's desired speed. If

$$F_i(\mathcal{R}(\rho_{i,m-1}^n, \rho_{i,m+1}^n)(u_i^n)) > u_i^n \mathcal{R}(\rho_{i,m-1}^n, \rho_{i,m+1}^n)(u_i^n), \quad (5.40)$$

we expect the flux constraint to be active at $\bar{x}_m = x_{m-1/2} + d_m^n \Delta x$ with $d_m^n = \rho_{i,m}^n / \hat{\rho}_{u_i^n}$ to ensure mass conservation. If $d_m^n \in [0, 1]$, then $\bar{x}_m \in C_m$ and, setting $\Delta t_m^n := \Delta x(1 - d_m^n) / u_i^n$, we replace the numerical fluxes in (5.39) by

$$\begin{aligned} \tilde{\mathcal{F}}_i(\rho_{i,m-1}^n, \rho_{i,m}^n) &:= \mathcal{F}_i(\rho_{i,m-1}^n, \hat{\rho}_{u_i^n}), \\ \tilde{\mathcal{F}}_i(\rho_{i,m}^n, \rho_{i,m+1}^n) &:= \max \{ 1 - \Delta t_m^n / \Delta t, 0 \} F_i(\hat{\rho}_{u_i^n}). \end{aligned}$$

Besides, to track the AV's trajectory, at each time step we implement an explicit Euler scheme:

$$y_i^{n+1} = y_i^n + \begin{cases} u_i^n \Delta t, & \text{if (5.40) holds,} \\ v_i(\rho_m^n) \Delta t, & \text{otherwise.} \end{cases}$$

In the following numerical tests we consider (2.3) with

$$v_j(\rho_j) = V_j(1 - \rho_j/R_j), \quad (5.41a)$$

$$F_j(\rho_j) = \rho_j v_j(\rho_j) \quad (5.41b)$$

$$S_j(U_j, U_{j+1}) = (v_{j+1}(\rho_{j+1}) - v_j(\rho_j))^+ \rho_j - (v_{j+1}(\rho_{j+1}) - v_j(\rho_j))^- \rho_{j+1}. \quad (5.41c)$$

5.1 A 3-lane road with heterogeneous speed limits and AVs

As an illustration of the behaviour of model (2.3), we consider $M = 3$ and $N_j = 1$ for $j = 1, 2, 3$. We set $R_j = 1$ for $j = 1, 2, 3$ and $V_1 = 50$ km/h, $V_2 = 80$ km/h, $V_3 = 100$ km/h, thus considering different speed limits on each lane. AVs' desired speeds are set constantly to $u_1^1 = u_2^1 = u_3^1 = 30$ km/h and the relaxation parameter $\tau = 0.05$ km⁻¹.

We consider a road stretch of length $L = 10$ km with $\Delta x = 0.02$ and initial conditions

$$\begin{aligned}\rho_1^0(x) &= 0.5 + 0.5 \sin(0.5\pi x), \\ \rho_2^0(x) &= 0.5 + 0.5 \cos(0.5\pi x), \\ \rho_3^0(x) &= 0.5 + 0.5 \sin(\pi x),\end{aligned}\tag{5.42}$$

for the (normalized) traffic densities on each lane, and

$$y_{0,1}^1 = 1, \quad y_{0,2}^1 = 2, \quad y_{0,3}^1 = 3,$$

for the AVs' positions.

In Figure 2, we can observe that, due to the different speed limits, vehicles tend to move and accumulate into the fastest lanes 2 and 3, while lane 1 remains almost empty. Also, due to low traffic, on lane 1 there is no interaction between the AV and the bulk traffic. On the contrary, in Figure 2c we can observe the appearance of non-classical shocks. Moreover, we may observe that AVs slow down to adapt to the downstream traffic speed both in Figures 2b and 2c. Figure 2d gives an overview of the total traffic density on the three lanes, showing that an oscillatory pattern in the density can still be observed due to the uneven distribution of the traffic density across lanes.

5.2 Relaxation limit $\tau \searrow 0$ without AVs

To illustrate the convergence result proved in Section 3 (see Theorem 2), we consider the same scenario as in the previous Section 5.1, i.e. $M = 3$, $R_j = 1$ for $j = 1, 2, 3$ and initial density data as in (5.42). Now, no AVs are present.

We consider two cases:

- (C1) $V_1 = V_2 = V_3 = 80$ km/h, so that the speed function is the same on all lanes and we fulfill the hypotheses of Theorem 2;
- (C2) $V_1 = 60$ km/h, $V_2 = 80$ km/h, $V_3 = 100$ km/h, to be compared with the solution of the conjectured limit LWR equation (3.25) with speed flux function $f(r) = rv(r)$, $v(r) = V(1 - r/3)$, where $V = 80$ km/h is the average maximal speed among lanes 1 to 3. Note that this is indeed the relaxation limit in case (C1).

Figure 3 shows the profiles of the total density for model (3.13) for $\tau = 1$ km⁻¹ and $\tau = 0.1$ km⁻¹ and the solution of the limit LWR problem (3.25) at time $t = 1$ min. In particular, in Figure 3a we can observe an illustration of relaxation limit stated in Theorem 2. Moreover, Figure 3b suggests that the limit also holds for non homogeneous speeds on the different lanes, as conjectured in the design of test (C2).

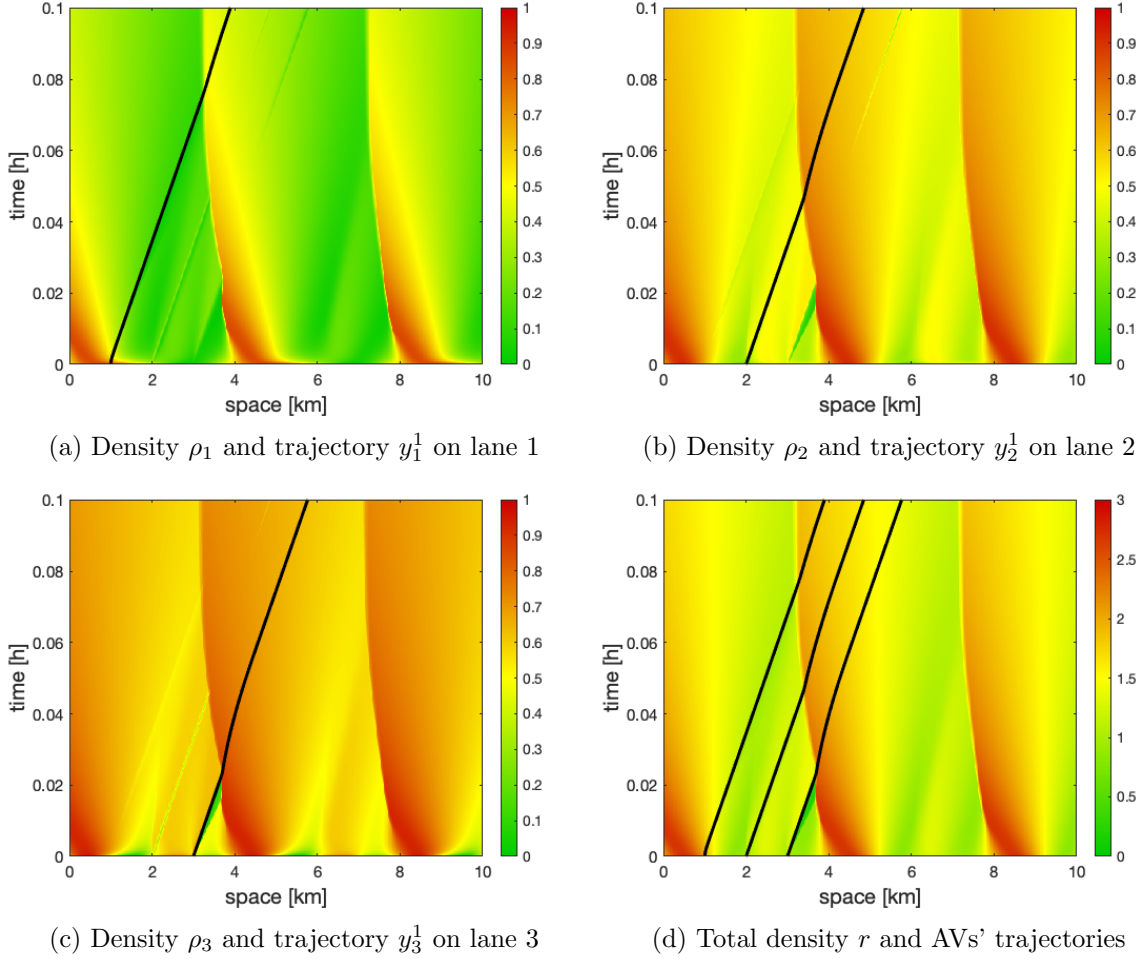


Figure 2: Heat-map of the traffic densities on each lane and the total density on the road segment $[0, 10]$ km and the time interval $[0, 0.1]$ h, with the corresponding AV's trajectories (black lines).

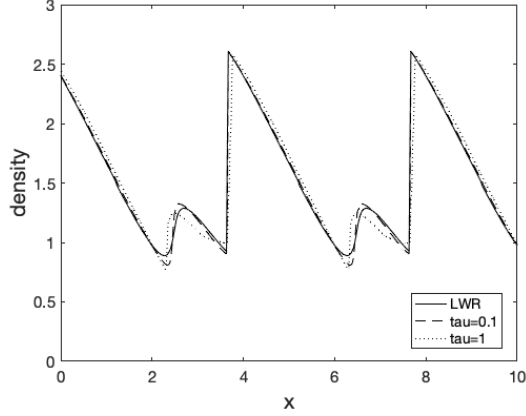
5.3 Relaxation limit $\tau \searrow 0$ with AV

Referring to the results in Section 4, we aim at comparing the relaxation limit of model (4.29), i.e. the multi-lane model with $M = 2$ lanes and a single AV, with the solution of the corresponding moving bottleneck model (4.34).

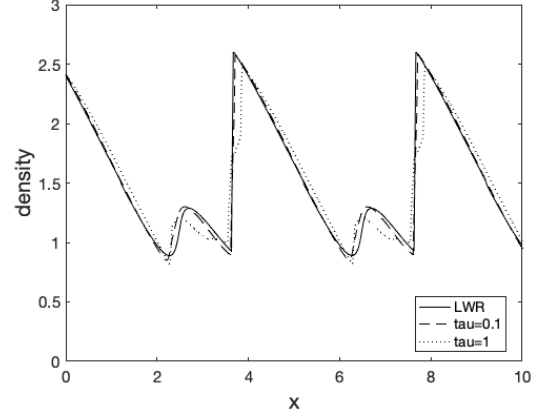
We take $V = V_1 = V_2 = 2$, $R_1 = R_2 = 1$ (thus $R = 2$), so that $F(\rho) = 2\rho(1 - \rho)$ and $f(r) = 2r - r^2$. Moreover, we set $\tau = 0.01$. We consider different density initial data and AV speeds, to see when non-classical shocks appear in the solution of (4.29) and to understand their nature. In the following, we denote by \hat{r} and \check{r} respectively the left and right traces of the non-classical shock in (4.34), ρ^* the density satisfying $F(\rho^*) = u$ and r^* the density satisfying $f(r^*) = u$. Setting $u(t) = 1$, we get $\check{r} \simeq 0.15$, $\hat{r} \simeq 0.85$, $\rho^* = 0.5$ and $r^* = 1$. Accordingly, we consider the following constant initial data:

$$\text{(IC1)} \quad \rho_1^0(x) = \rho_2^0(x) \equiv 0.05, \text{ i.e. } r^0(0) \equiv 0.1;$$

$$\text{(IC2)} \quad \rho_1^0(x) = \rho_2^0(x) \equiv 0.15, \text{ i.e. } r^0(0) \equiv 0.3;$$



(a) $V_1 = V_2 = V_3 = 80$ km/h



(b) $V_1 = 60, V_2 = 80, V_3 = 100$ km/h

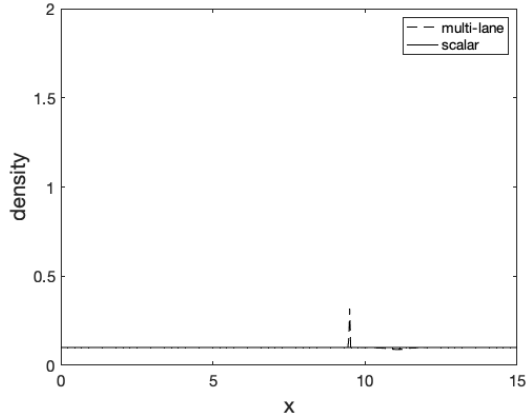
Figure 3: Density profiles at $t = 1$ min of the total densities $\rho_1 + \rho_2 + \rho_3$ of model (3.13) with $\tau = 0.1, 1$ km $^{-1}$ and the solution r of LWR model (3.25) for cases **(C1)** (left) and **(C2)** (right).

(IC3) $\rho_1^0(x) = \rho_2^0(x) \equiv 0.3$, i.e. $r^0(0) \equiv 0.6$;

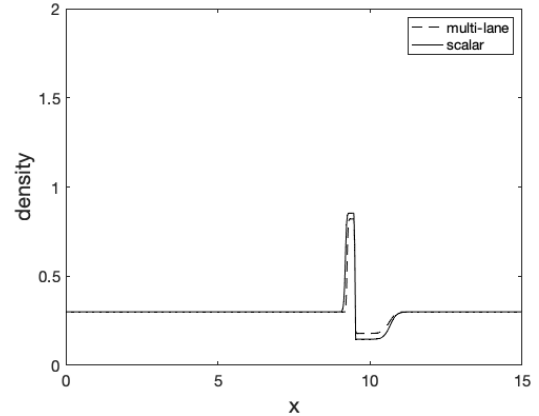
(IC4) $\rho_1^0(x) = \rho_2^0(x) \equiv 0.45$, i.e. $r^0(0) \equiv 0.9$;

(IC5) $\rho_1^0(x) = \rho_2^0(x) \equiv 0.75$, i.e. $r^0(0) \equiv 1.5$.

The corresponding density profiles at $t = 2$ are displayed in Figure 4. It appears clearly that, when the AV acts as a flux constraint, the density traces at the AV location differ between model (4.29) and (4.34). In particular, we observe that we always have $r(t, y(t)-) > (\rho_1 + \rho_2)(t, y(t)-) > (\rho_1 + \rho_2)(t, y(t)+) > r(t, y(t)+)$.



(a) **(IC1)** $r^0 \in]0, \check{r}[$



(b) **(IC2)** $r^0 \in]\check{r}, \rho^*[$

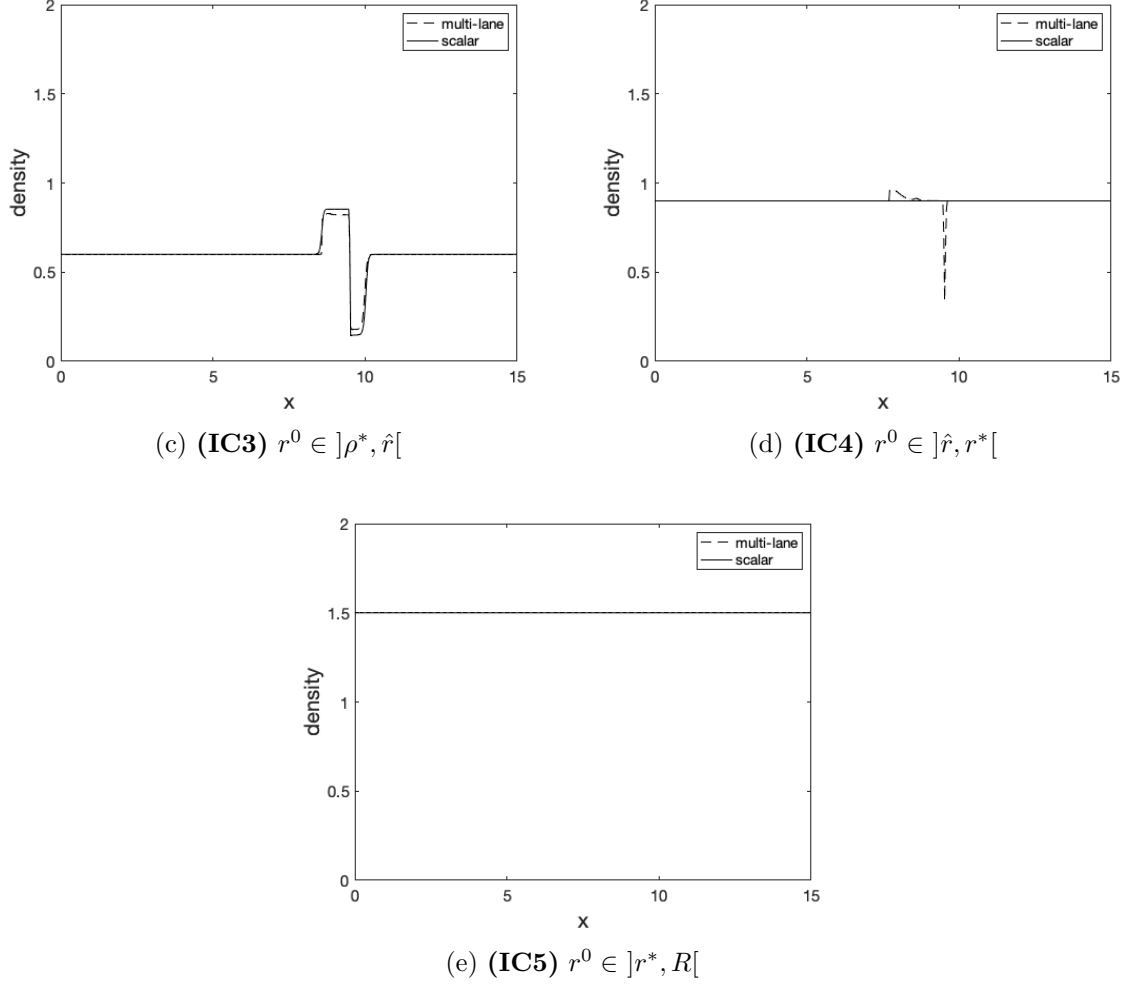


Figure 4: Comparison between $\rho_1(2, \cdot) + \rho_2(2, \cdot)$ solutions of (4.29) with $\tau = 0.01$ and r solution of (4.34), with $r_0 := \rho_1^0 + \rho_2^0$, for $u(t) = 1$.

A L^1 stability for systems of weakly coupled balance laws

Let us consider the $M \times M$ systems of balance laws coupled in the source term

$$\begin{cases} \partial_t \boldsymbol{\rho} + \partial_x F(\boldsymbol{\rho}) = G(\boldsymbol{\rho}), & x \in \mathbb{R}, t \geq 0, \\ \boldsymbol{\rho}(0, x) = \boldsymbol{\rho}^0(x), \end{cases} \quad (\text{A.43a})$$

and

$$\begin{cases} \partial_t \boldsymbol{\sigma} + \partial_x \tilde{F}(\boldsymbol{\sigma}) = G(\boldsymbol{\sigma}), & x \in \mathbb{R}, t \geq 0, \\ \boldsymbol{\sigma}(0, x) = \boldsymbol{\sigma}^0(x), \end{cases} \quad (\text{A.43b})$$

for some $\boldsymbol{\rho}^0, \boldsymbol{\sigma}^0 \in \text{BV}(\mathbb{R}; \mathbb{R}^M)$, where

$$\begin{aligned} F(\boldsymbol{\rho}) &= \left(F_1(\rho_1), \dots, F_M(\rho^M) \right)^T, \\ \tilde{F}(\boldsymbol{\rho}) &= \left(\tilde{F}_1(\rho_1), \dots, \tilde{F}_M(\rho^M) \right)^T, \end{aligned}$$

$$G(\boldsymbol{\rho}) = (G_1(\boldsymbol{\rho}), \dots, G_M(\boldsymbol{\rho}))^T.$$

Note that system (3.13) corresponds to setting

$$\begin{aligned} F_j(\rho_j) &= \rho_j v_j(\rho_j), \\ G_j(\boldsymbol{\rho}) &= \frac{1}{\tau} (S_{j-1}(\rho_{j-1}, \rho_j) - S_j(\rho_j, \rho_{j+1})), \end{aligned}$$

in (A.43a) for $j = 1, \dots, M$.

The following stability result holds.

Theorem 4. *Let (A.43a) and (A.43b) admit unique BV entropy weak solutions $\boldsymbol{\rho}, \boldsymbol{\sigma} \in \mathbf{C}^0([0, +\infty[; \mathbf{L}^1(\mathbb{R}; \mathbb{R}^M))$ and let $\mathfrak{T} := \min \left\{ \max_{s \in [0, t]} TV(\boldsymbol{\rho}(s, \cdot)), \max_{s \in [0, t]} TV(\boldsymbol{\sigma}(s, \cdot)) \right\}$ and $\mathfrak{L} := \max_{j=1, \dots, M} \|F'_j - \tilde{F}'_j\|_\infty$. Then for all $t \geq 0$ it holds*

$$\|\boldsymbol{\rho}(t, \cdot) - \boldsymbol{\sigma}(t, \cdot)\|_1 \leq e^{\|\nabla G\|_\infty t} \|\boldsymbol{\rho}^0 - \boldsymbol{\sigma}^0\|_1 + \frac{M \mathfrak{T} \mathfrak{L}}{\|\nabla G\|_\infty} (e^{\|\nabla G\|_\infty t} - 1). \quad (\text{A.44})$$

Proof. Let us consider WFT approximate solutions constructed by fractional steps as in Section 3:

- In any time interval $[t^n, t^{n+1}[$, $n \in \mathbb{N}$, $\boldsymbol{\rho}^{\tau, \nu}$ is the WFT approximation of

$$\begin{cases} \partial_t \rho_j + \partial_x F_j(\rho_j) = 0, & x \in \mathbb{R}, t \in [t^n, t^{n+1}[, \\ \rho_j(t^n, x) = \rho_j^{\tau, \nu}(t^n+, x), & j = 1, \dots, M, \end{cases} \quad (\text{A.45})$$

and $\boldsymbol{\sigma}^{\tau, \nu}$ is the WFT approximation of

$$\begin{cases} \partial_t \sigma_j + \partial_x \tilde{F}_j(\sigma_j) = 0, & x \in \mathbb{R}, t \in [t^n, t^{n+1}[, \\ \sigma_j(t^n, x) = \sigma_j^{\tau, \nu}(t^n+, x), & j = 1, \dots, M, \end{cases} \quad (\text{A.46})$$

constructed as described e.g. in [19, Section 2.3].

- At time $t = t^{n+1}$, we define

$$\rho_j^{\tau, \nu}(t^{n+1}+, \cdot) = \rho_j^{\tau, \nu}(t^{n+1}-, \cdot) + \Delta t^\nu G(\boldsymbol{\rho}^{\tau, \nu}(t^{n+1}-, \cdot)), \quad (\text{A.47})$$

$$\sigma_j^{\tau, \nu}(t^{n+1}+, \cdot) = \sigma_j^{\tau, \nu}(t^{n+1}-, \cdot) + \Delta t^\nu G(\boldsymbol{\sigma}^{\tau, \nu}(t^{n+1}-, \cdot)), \quad (\text{A.48})$$

Solutions of A.45 and A.46 satisfy

$$\left\| \rho_j^{\tau, \nu}(t^{n+1}-, \cdot) - \sigma_j^{\tau, \nu}(t^{n+1}-, \cdot) \right\|_1 \leq \left\| \rho_j^{\tau, \nu}(t^n+, \cdot) - \sigma_j^{\tau, \nu}(t^n+, \cdot) \right\|_1 + \mathfrak{T} \mathfrak{L} \Delta t^\nu. \quad (\text{A.49})$$

Subtracting (A.48) from (A.47) and taking \mathbf{L}^1 -norms, we get

$$\left\| \rho_j^{\tau, \nu}(t^{n+1}+, \cdot) - \sigma_j^{\tau, \nu}(t^{n+1}+, \cdot) \right\|_1 \leq \left\| \rho_j^{\tau, \nu}(t^{n+1}-, \cdot) - \sigma_j^{\tau, \nu}(t^{n+1}-, \cdot) \right\|_1 (1 + \Delta t^\nu \|\nabla G\|_\infty),$$

which, combined with (A.49), gives

$$\sum_{j=1}^M \left\| \rho_j^{\tau, \nu}(t^n+, \cdot) - \sigma_j^{\tau, \nu}(t^n+, \cdot) \right\|_1$$

$$\begin{aligned}
&\leq (1 + \Delta t^\nu \|\nabla G\|_\infty)^{t^n/\Delta t^\nu} \sum_{j=1}^M \left\| \rho_j^0 - \sigma_j^0 \right\|_1 + \Delta t^\nu M \mathfrak{L} \sum_{k=1}^{t^n/\Delta t^\nu} (1 + \Delta t^\nu \|\nabla G\|_\infty)^k \\
&= (1 + \Delta t^\nu \|\nabla G\|_\infty)^{t^n/\Delta t^\nu} \sum_{j=1}^M \left\| \rho_j^0 - \sigma_j^0 \right\|_1 + \frac{M \mathfrak{L}}{\|\nabla G\|_\infty} \left[(1 + \Delta t^\nu \|\nabla G\|_\infty)^{t^n/\Delta t^\nu + 1} - 1 \right],
\end{aligned}$$

which converges to (A.44) as $\Delta t^\nu \searrow 0$. \square

Remark 4. Note that, setting $G \equiv 0$ in (A.44), one recovers the classical stability estimate for scalar equations. Instead, (3.16) is stronger than (A.44), due to the specific monotonicity assumption **(S0)**. In particular, estimate (A.44) does not suffice to conclude in the proof of Theorem 2, since in this case $\|\nabla G\|_\infty = 2\mathcal{S}/\tau$ and we get

$$\sum_{j=1}^M \left\| \rho_j^\tau(t, \cdot) - \frac{r(t, \cdot)}{M} \right\|_{\mathbf{L}^1([a,b])} = \mathcal{O}(1) \cdot (b-a) \left(\tau e^{2\mathcal{S}/\sqrt{\tau}} + e^{-1/\sqrt{\tau}} \right) \xrightarrow{\tau \rightarrow 0} \infty.$$

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