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# On an equation of the Drinfeld-Sokolov hierarchy 

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#### Abstract

A seventh order ordinary differential equation (ODE) arising by reduction of the Drinfeld-Sokolov hierarchy is shown to be identical to a similarity reduction of an equation in the hierarchy of SawadaKotera. We also exhibit its link with a particular F-VI, a fourth order ODE isolated by Cosgrove which is likely to define a higher order Painlevé function.


## 1 Introduction

In a recent article [6], the authors consider the tau cover of the DrinfeldSokolov hierarchy and, in order to obtain explicit solutions, perform a similarity reduction [6, Eq. (5.1)] which defines a system of nonlinear ODEs in the independent variable $x$. By construction, this system possesses a Lax pair $(L, M)$ [6, Eq. (5.5)] whose zero-curvature condition is

$$
\begin{equation*}
\left[z \partial_{z}-M, \partial_{x}-L\right] \equiv z L_{z}-M_{x}+[L, M]=0 \tag{1}
\end{equation*}
$$

in which $z$ is the spectral parameter.
For all their choices but one of the underlying affine Kac-Moody algebra $\mathfrak{g}$, the authors did succeed to explicitly integrate the nonlinear ODE system
in terms of various elliptic or Painlevé or higher Painlevé functions. The only system which could not be integrated results from the choice $\mathfrak{g}=A_{2}^{(2)}$, this is the seventh order nonautonomous system for $u(x), \omega(x)$ [6, Example 5.5 page 1487]
$\left\{\begin{array}{l}\omega^{\prime}-\frac{u}{3}=0, \\ u^{(6)}+14 u u^{(4)}+14 u^{\prime} u^{(3)}+14 u^{\prime \prime 2}+56 u^{2} u^{\prime \prime}+28 u u^{\prime 2}+\frac{56}{3} u^{4}+36 x u+108 \omega=0 .\end{array}\right.$
which can be viewed as a birational transformation between $u(x)$ and $\omega(x)$, each variable obeying a seventh order ODE.

The purpose of this work is to explicitly integrate this system, i.e. to map it either to Painlevé equations (second order), or to one of the five "higher Painlevé equations" (fourth and fifth order) isolated by Cosgrove [2, 3], or to higher order (six and above) equations in the hierarchy of the previous ones.

The method, developed in next sections, is classical and it relies on three pieces of information: (i) the Lax pair, (ii) the singularity structure, (iii) exhaustive lists of ODEs possessing the Painlevé property.

In Section 2, by considering the invariants of the matrix Lax pair, we obtain a unique first integral, thus lowering the differential order only to six. This is an indication (not a proof) that the equations of Painlevé and Cosgrove should be insufficient to perform the integration.

In Section 3, we therefore investigate the singularity structure of the system (2). The three families of movable simple poles are then compared with sixth or senventh order members of various, already classified, hierarchies. This allows one to integrate (2) in terms of a higher member of the SawadaKotera hierarchy.

## 2 First integral

The system (2) admits a three-dimensional zero curvature representation, this is [courtesy of Wu Chao-Zhong],

$$
\begin{aligned}
& L=\sqrt{2}\left(-L_{0}+u L_{1}-z L_{2}\right), M=2 \sum_{j=0}^{5} z^{j} M_{j}, z L_{z}-M_{x}+[L, M]=0 \\
& M_{5}=2 \sqrt{2} L_{2}, M_{4}=2 \sqrt{2} L_{0}-\frac{4}{3} \sqrt{2} u L_{1}, M_{3}=\frac{2}{3} \sqrt{2} u M_{3 a}+\frac{2}{3} u^{\prime} M_{3 b}
\end{aligned}
$$

$$
\begin{align*}
& M_{2}=\frac{\sqrt{2}}{9}\left(u^{\prime \prime}+2 u^{2}\right) m_{2}, \\
& M_{1}=\frac{u^{(3)}+4 u u^{\prime}}{9} M_{1 a}+\frac{\sqrt{2}}{81}\left(3 u^{(4)}+9 u^{\prime 2}-8 u^{3}-27 x\right) L_{2} \\
& M_{0}=A_{1} M_{0 a}+A_{2} L_{1}+A_{3} L_{0}, \\
& A_{1}=-\frac{u^{(5)}+12 u^{\prime} u^{\prime \prime}+6 u u^{(3)}+8 u^{2}+9}{54}, \\
& A_{2}=\frac{\sqrt{2}}{486}\left(20 u^{4}+60 u^{2} u^{\prime \prime}+84 u u^{\prime 2}+24 u u^{(4)}+33 u^{\prime \prime 2}+60 u^{\prime} u^{(3)}+3 u^{(6)}+108 x u-162 \omega\right), \\
& A_{3}=-\frac{\sqrt{2}}{81}\left(3 u^{(4)}+18 u u^{\prime \prime}+9 u^{\prime 2}-8 u^{3}\right) . \tag{3}
\end{align*}
$$

in which the eight constant operators can be represented by third order matrices,

$$
\begin{align*}
& L_{0}=\left(\begin{array}{lll}
0 & 0 & 1 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right), L_{1}=\left(\begin{array}{lll}
0 & 0 & 0 \\
0 & 0 & 0 \\
1 & 0 & 0
\end{array}\right), L_{2}=\left(\begin{array}{lll}
0 & 0 & 0 \\
1 & 0 & 0 \\
0 & 1 & 0
\end{array}\right), \\
& M_{3 a}=\left(\begin{array}{lll}
0 & 1 & 0 \\
0 & 0 & 1 \\
0 & 0 & 0
\end{array}\right), M_{3 b}=\left(\begin{array}{ccc}
0 & 0 & 0 \\
1 & 0 & 0 \\
0 & -1 & 0
\end{array}\right), m_{2}=\left(\begin{array}{ccc}
-1 & 0 & 0 \\
0 & 2 & 0 \\
0 & 0 & -1
\end{array}\right), \\
& M_{1 a}=\left(\begin{array}{lll}
0 & -1 & 0 \\
0 & 0 & 1 \\
0 & 0 & 0
\end{array}\right), M_{0 a}=\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & 0 & 0 \\
0 & 0 & -1
\end{array}\right) . \tag{4}
\end{align*}
$$

The traces of odd powers of $M$ are zero, but the traces of even powers generate a single first integral $K$ of the system (2),

$$
\begin{align*}
& K=v^{\prime 2}+4 v\left[2 u u^{(4)}-2 u^{\prime} u^{(3)}+u^{\prime \prime 2}+12 u^{2} u^{\prime \prime}+4 u^{4}+54 \omega\right]  \tag{5}\\
& v=u^{(4)}+6 u u^{\prime \prime}+3 u^{\prime 2}+\frac{8}{3} u^{3}+9 x
\end{align*}
$$

The initial system (2) is then equivalent to a birational transformation between $u(x)$ and $v(x)$, each variable obeying a sixth order ODE,

$$
\left\{\begin{array}{l}
u^{(4)}+6 u u^{\prime \prime}+3 u^{\prime 2}+\frac{8}{3} u^{3}+9 x-v=0  \tag{6}\\
2 v v^{\prime \prime}-v^{\prime 2}+8 u v^{2}+K=0
\end{array}\right.
$$

## 3 Singularity structure

For the technical vocabulory, we refer for instance to Ref. [1]. Near a movable singularity $x_{0}$, the system (2) has three families of movable poles, the first one being already mentioned in [6, page 1487], and the seven Fuchs indices of the linearized equation near this singularity are all relative integers,

$$
\left\{\begin{array}{l}
\omega \sim a \chi^{-1}, u \sim-3 a \chi^{-2}, \chi=x-x_{0}  \tag{7}\\
a=1, \text { Fuchs indices }=-1,2,3,4,7,8,12 \\
a=2, \text { Fuchs indices }=-2,-1,3,4,8,9,14 \\
a=5, \text { Fuchs indices }=-5,-4,-1,7,8,12,18
\end{array}\right.
$$

This structure matches that of a seventh order equation in the SawadaKotera hierarchy mentioned by Gordoa and Pickering [5, Eq. (5.94)] (see also [4]),
$\left\{\begin{array}{c}\frac{\mathrm{d}}{\mathrm{d} X}\left[U^{(6)}+(7 / 2)\left(U U^{(4)}+U^{\prime} U^{(3)}+U^{\prime \prime 2}+U^{2} U^{\prime \prime}\right)+7 / 4 U U^{\prime 2}+7 / 24 U^{4}\right](8) \\ -q_{1}\left(2 U+X U^{\prime}\right)=0, q_{1}=\text { arbitrary nonzero constant },\end{array}\right.$ with the correspondence

$$
\begin{equation*}
u=\frac{b^{2}}{4} U, X=b x, b^{7}=-\frac{36}{q_{1}} . \tag{9}
\end{equation*}
$$

As to the scalar Lax pair [5, Eqs. (5.6), (5.96), (5.84)], as the interested reader can check, it is identical to that obeyed by the second component of the wave vector of the matrices in (3).

## 4 A link to F-VI

As a final remark, let us mention a link between this seventh order equation in the Sawada-Kotera hierarchy and the F-VI ODE isolated by Cosgrove [2],

$$
\begin{equation*}
(\mathrm{F}-\mathrm{VI}) \quad U^{\prime \prime \prime \prime}=18 U U^{\prime \prime}+9 U^{\prime 2}-24 U^{3}+\alpha U^{2}+\frac{\alpha^{2}}{9} U+\kappa X+\beta \tag{10}
\end{equation*}
$$

Indeed, when $K=0, v=0$, the system (6) admits a particular solution $u(x)$ equal to a particular affine transform of F-VI with the correspondence

$$
\begin{equation*}
u=-3 b^{2} U, X=b\left(x-\frac{\beta}{3} b^{6}\right), \kappa=3 b^{-7}, \alpha=0 . \tag{11}
\end{equation*}
$$

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