

On an equation of the Drinfeld-Sokolov hierarchy R. Conte

▶ To cite this version:

R. Conte. On an equation of the Drinfeld-Sokolov hierarchy. Open Communications in Nonlinear Mathematical Physics, 2024. hal-04549364

HAL Id: hal-04549364

https://hal.science/hal-04549364

Submitted on 17 Apr 2024

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



On an equation of the Drinfeld-Sokolov hierarchy

Robert Conte^{1,2}

¹ Centre Borelli, École normale supérieure de Paris-Saclay, France.

² Department of Mathematics, The University of Hong Kong. Robert.Conte@CEA.FR

ORCID https://orcid.org/0000-0002-1840-5095

17 April 2024

Abstract

A seventh order ordinary differential equation (ODE) arising by reduction of the Drinfeld-Sokolov hierarchy is shown to be identical to a similarity reduction of an equation in the hierarchy of Sawada-Kotera. We also exhibit its link with a particular F-VI, a fourth order ODE isolated by Cosgrove which is likely to define a higher order Painlevé function.

1 Introduction

In a recent article [6], the authors consider the tau cover of the Drinfeld-Sokolov hierarchy and, in order to obtain explicit solutions, perform a similarity reduction [6, Eq. (5.1)] which defines a system of nonlinear ODEs in the independent variable x. By construction, this system possesses a Lax pair (L, M) [6, Eq. (5.5)] whose zero-curvature condition is

$$[z\partial_z - M, \partial_x - L] \equiv zL_z - M_x + [L, M] = 0, \tag{1}$$

in which z is the spectral parameter.

For all their choices but one of the underlying affine Kac-Moody algebra \mathfrak{g} , the authors did succeed to explicitly integrate the nonlinear ODE system

in terms of various elliptic or Painlevé or higher Painlevé functions. The only system which could not be integrated results from the choice $\mathfrak{g}=A_2^{(2)}$, this is the seventh order nonautonomous system for $u(x), \omega(x)$ [6, Example 5.5 page 1487]

$$\begin{cases} \omega' - \frac{u}{3} = 0, \\ u^{(6)} + 14uu^{(4)} + 14u'u^{(3)} + 14u''^2 + 56u^2u'' + 28uu'^2 + \frac{56}{3}u^4 + 36xu + 108\omega = 0. \end{cases}$$
(2)

which can be viewed as a birational transformation between u(x) and $\omega(x)$, each variable obeying a seventh order ODE.

The purpose of this work is to explicitly integrate this system, i.e. to map it either to Painlevé equations (second order), or to one of the five "higher Painlevé equations" (fourth and fifth order) isolated by Cosgrove [2, 3], or to higher order (six and above) equations in the hierarchy of the previous ones.

The method, developed in next sections, is classical and it relies on three pieces of information: (i) the Lax pair, (ii) the singularity structure, (iii) exhaustive lists of ODEs possessing the Painlevé property.

In Section 2, by considering the invariants of the matrix Lax pair, we obtain a unique first integral, thus lowering the differential order only to six. This is an indication (not a proof) that the equations of Painlevé and Cosgrove should be insufficient to perform the integration.

In Section 3, we therefore investigate the singularity structure of the system (2). The three families of movable simple poles are then compared with sixth or serventh order members of various, already classified, hierarchies. This allows one to integrate (2) in terms of a higher member of the Sawada-Kotera hierarchy.

2 First integral

The system (2) admits a three-dimensional zero curvature representation, this is [courtesy of Wu Chao-Zhong],

$$L = \sqrt{2}(-L_0 + uL_1 - zL_2), M = 2\sum_{j=0}^{5} z^j M_j, zL_z - M_x + [L, M] = 0,$$

$$M_5 = 2\sqrt{2}L_2, M_4 = 2\sqrt{2}L_0 - \frac{4}{3}\sqrt{2}uL_1, M_3 = \frac{2}{3}\sqrt{2}uM_{3a} + \frac{2}{3}u'M_{3b},$$

$$M_{2} = \frac{\sqrt{2}}{9} \left(u'' + 2u^{2} \right) m_{2},$$

$$M_{1} = \frac{u^{(3)} + 4uu'}{9} M_{1a} + \frac{\sqrt{2}}{81} (3u^{(4)} + 9u'^{2} - 8u^{3} - 27x) L_{2}$$

$$M_{0} = A_{1} M_{0a} + A_{2} L_{1} + A_{3} L_{0},$$

$$A_{1} = -\frac{u^{(5)} + 12u'u'' + 6uu^{(3)} + 8u^{2} + 9}{54},$$

$$A_{2} = \frac{\sqrt{2}}{486} (20u^{4} + 60u^{2}u'' + 84uu'^{2} + 24uu^{(4)} + 33u''^{2} + 60u'u^{(3)} + 3u^{(6)} + 108xu - 162\omega),$$

$$A_{3} = -\frac{\sqrt{2}}{81} (3u^{(4)} + 18uu'' + 9u'^{2} - 8u^{3}).$$
(3)

in which the eight constant operators can be represented by third order matrices,

$$L_{0} = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, L_{1} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}, L_{2} = \begin{pmatrix} 0 & 0 & 0 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix},$$

$$M_{3a} = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{pmatrix}, M_{3b} = \begin{pmatrix} 0 & 0 & 0 \\ 1 & 0 & 0 \\ 0 & -1 & 0 \end{pmatrix}, m_{2} = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & -1 \end{pmatrix},$$

$$M_{1a} = \begin{pmatrix} 0 & -1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{pmatrix}, M_{0a} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}.$$

$$(4)$$

The traces of odd powers of M are zero, but the traces of even powers generate a single first integral K of the system (2),

$$K = v'^{2} + 4v \left[2uu^{(4)} - 2u'u^{(3)} + u''^{2} + 12u^{2}u'' + 4u^{4} + 54\omega \right], \quad (5)$$

$$v = u^{(4)} + 6uu'' + 3u'^{2} + \frac{8}{3}u^{3} + 9x.$$

The initial system (2) is then equivalent to a birational transformation between u(x) and v(x), each variable obeying a sixth order ODE,

$$\begin{cases} u^{(4)} + 6uu'' + 3u'^2 + \frac{8}{3}u^3 + 9x - v = 0, \\ 2vv'' - v'^2 + 8uv^2 + K = 0. \end{cases}$$
 (6)

3 Singularity structure

For the technical vocabulory, we refer for instance to Ref. [1]. Near a movable singularity x_0 , the system (2) has three families of movable poles, the first one being already mentioned in [6, page 1487], and the seven Fuchs indices of the linearized equation near this singularity are all relative integers,

$$\begin{cases}
\omega \sim a\chi^{-1}, u \sim -3a\chi^{-2}, \chi = x - x_0, \\
a = 1, \text{ Fuchs indices} = -1, 2, 3, 4, 7, 8, 12, \\
a = 2, \text{ Fuchs indices} = -2, -1, 3, 4, 8, 9, 14, \\
a = 5, \text{ Fuchs indices} = -5, -4, -1, 7, 8, 12, 18.
\end{cases}$$
(7)

This structure matches that of a seventh order equation in the Sawada-Kotera hierarchy mentioned by Gordoa and Pickering [5, Eq. (5.94)] (see also [4]),

$$\begin{cases} \frac{\mathrm{d}}{\mathrm{d}X} \left[U^{(6)} + (7/2)(UU^{(4)} + U'U^{(3)} + U''^2 + U^2U'') + 7/4UU'^2 + 7/24U^4 \right] \\ -q_1(2U + XU') = 0, q_1 = \text{ arbitrary nonzero constant,} \end{cases} (8)$$

with the correspondence

$$u = \frac{b^2}{4}U, X = bx, b^7 = -\frac{36}{q_1}.$$
 (9)

As to the scalar Lax pair [5, Eqs. (5.6), (5.96), (5.84)], as the interested reader can check, it is identical to that obeyed by the second component of the wave vector of the matrices in (3).

4 A link to F-VI

As a final remark, let us mention a link between this seventh order equation in the Sawada-Kotera hierarchy and the F-VI ODE isolated by Cosgrove [2],

(F-VI)
$$U'''' = 18UU'' + 9U'^2 - 24U^3 + \alpha U^2 + \frac{\alpha^2}{9}U + \kappa X + \beta.$$
 (10)

Indeed, when K = 0, v = 0, the system (6) admits a particular solution u(x) equal to a particular affine transform of F-VI with the correspondence

$$u = -3b^2U, X = b\left(x - \frac{\beta}{3}b^6\right), \kappa = 3b^{-7}, \alpha = 0.$$
 (11)

References

- [1] R. Conte and M. Musette, *The Painlevé handbook*, 2nd ed., 391 pages (Springer Nature, Switzerland, 2020). https://doi.org/10.1007/978-3-030-53340-3
- [2] C.M. Cosgrove, Higher order Painlevé equations in the polynomial class, I. Bureau symbol P2, Stud. Appl. Math. **104** (2000) 1–65. https://doi.org/10.1111/1467-9590.00130
- [3] C.M. Cosgrove, Higher order Painlevé equations in the polynomial class, II. Bureau symbol P1, Stud. Appl. Math. **116** (2006) 321–413. https://doi.org/10.1111/j.1467-9590.2006.00346.x
- [4] P. Gordoa and A. Pickering, Bäcklund transformations for two new integrable partial differential equations, Europhys. Lett. **47:1** (1999) 21–24. https://dx.doi.org/10.1209/epl/i1999-00344-9
- P. Gordoa and A. Pickering, Nonisospectral scattering problems: a key to integrable hierarchies, J. Math. Phys. 40 (1999) 5749–5786. https://doi.org/10.1063/1.533055
- [6] Si-Qi Liu, Chao-Zhong Wu and Youjin Zhang, Virasoro constraints for Drinfeld-Sokolov hierarchies and equations of Painlevé type, Journal of the London mathematical society **106:2** (2022) 1443–1500. https://doi.org/10.1112/jlms.12603 http://arXiv.org/abs/1908.06707