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Prediction Challenge: Simulating Rydberg Photoexcited Cyclobutanone with Surface Hopping Dynamics based on Different Electronic Structure Methods

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Abstract

This research examines the nonadiabatic dynamics of cyclobutanone after excitation into the $n \rightarrow 3s$ Rydberg S_2 state. It stems from our contribution to the Special Topic of the Journal of Chemical Physics to test the predictive capability of computational chemistry against unseen experimental data. Decoherence-corrected fewest-switches surface hopping (DC-FSSH) was used to simulate nonadiabatic dynamics with full and approximated nonadiabatic couplings. Several simulation sets were computed with different electronic structure methods, including a multiconfigurational wavefunction (MCSCF) specially built to describe dissociative channels, multireference semiempirical approach, time-dependent density functional theory, algebraic diagrammatic construction, and coupled cluster. MCSCF dynamics predicts a slow deactivation of the S_2 state (10 ps), followed by an ultrafast population transfer from S_1 to S_0 (<100 fs). CO elimination (C3 channel) dominates over C_2H_4 formation (C2 channel). These findings radically differ from the other methods, which predicted S_2 lifetimes 10 to 250 times shorter and C2 channel predominance. These results suggest that routine electronic structure methods may hold low predictive power for the outcome of nonadiabatic dynamics.

1 Introduction

Thanks to commendable scientific efforts globally over the last few decades, a handful of theoretical methods^{1–5} have emerged to simulate excited-state photoinduced dynamics. These developments span from excited-states electronic structure theory to nonadiabatic dynamics methodologies.⁶ However, a lack of comprehensive and uniform studies comparing the performance of these methods with recent experimental findings leaves the question of which method excels still unanswered. To address this very issue, The Journal of Chemical Physics proposed a double-blind prediction challenge to forecast the outcomes of an upcoming experiment performed using ultrafast electron diffraction,⁷ where a gas sample of cyclobutanone (CB) molecule will be irradiated with a 200 nm light source targeting the Rydberg excited state ($n \rightarrow 3s$) and electron diffraction images will be collected during the process in variable step sizes.

This article stems from our participation in this prediction challenge, offering theoretical insights and simulated electron diffraction spectra of the photochemical processes of CB in the gas phase. Our simulations employed decoherence-corrected fewest-switch surface hopping (DC-FSSH)¹ combined with different routine electronic structure methods, including MCSCF, MRCI/ODM3, TDDFT, ADC(2), and CC2 (acronyms are defined in Computational Methods).

We worked with an array of methods to gauge their quality performance. All approaches are well established, have low computational costs, and, except for MCSCF, are of straightforward usage. Each approach was expected to offer specific advantages and technical difficulties. MCSCF and MRCI/ODM3 should suffer from instabilities caused by changes in the active space orbital composition during dynamics. MCSCF mostly neglects dynamic electron correlation. MRCI/ODM3 should poorly describe diffuse states⁸ like S_2 in cyclobutanone. The single-reference character of TDDFT, CC2, and ADC(2) may cause trouble during dissociation. They are also not expected to work well near S_1/S_0 crossings.¹ CC2, as a non-Hermitian method, should not converge when S_2 and S_1 are degenerated.⁹ ADC(2) is well known for incorrectly describing the excited state of carbonyl compounds.¹⁰ Thus, given all these limitations in the excited state calculations, it is clear that the actual provocation of this prediction challenge is not so much about the choice of the nonadiabatic dynamics approach but about finding an electronic structure method inexpensive enough to be used for dynamics and still deliver robust results for states with strongly different diabatic characters.^{11,12}

Cyclobutanone mesmerizes chemists. It offers unique and exciting photochemical behavior compared to other noncyclic and cyclic ketones of larger ring sizes, mainly due to its ring strain. Upon S_1

photoexcitation, CB undergoes the Norrish type-I reaction with the cleavage of the C-C bond in the alpha to the carbonyl position (α -cleavage) in the excited state and subsequent multiple bond breaking and formation via diradical intermediates.^{13–15}

From the experimental absorption UV spectra,^{14,16,17} we learn that the S_1 state is a symmetry-forbidden $n \rightarrow \pi^*$ transition ranging from 330 to 240 nm, with its maximum around 280 nm. The S_2 state is a Rydberg excitation $n \rightarrow 3s$ spanning from 206 to 182 nm. Previous works have shown that photoexcited CB in the gas phase can follow two main photodissociation channels: the C2 and the C3 pathways. The C2 pathway forms ethylene (C_2H_4) and ketene (CH_2CO), as depicted in Figure 1. In contrast, the C3 pathway forms cyclopropane ($c\text{-}C_3H_6$) and carbon monoxide (CO). The branching ratio of C3 and C2 products depends sensitively on the excitation wavelength. The C3:C2 product ratio was recorded as 0.8 at 248 nm, 0.4 at 313 nm, 2 at 326 nm, and 7 at 344 nm in the gas phase.^{18–20}

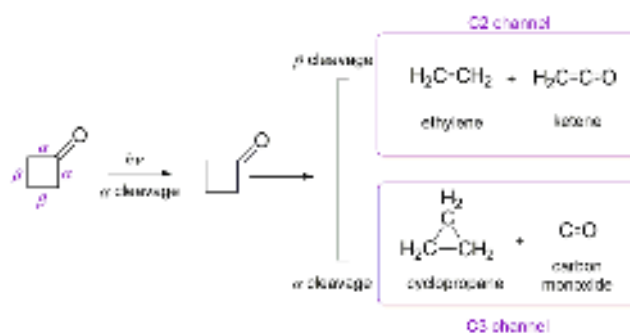


Figure 1: The schematics of two main photo-pathways of cyclobutanone (CB). Upon photoexcitation, CB forms a diradical intermediate through a C-C bond breaking in the α position. Subsequent cleavage of a β bond or the second α bond yields the C2 and C3 photoproducts, respectively.

Diau, Kötting, and Zewail¹⁷ employed femtosecond time-resolved mass spectrometry to explore the CB photoinduced dynamics. They reported that one-photon excitation at 307 nm to the S_1 state induces two different dynamical processes associated with the α -cleavage in the S_1 state. The fast decay component with a time constant of less than 50 fs is mainly responsible for the ultrafast relaxation of the initial wavepacket from the Franck-Condon to the S_1 minimum region. On the other hand, the wavepacket moves along the bond-breaking coordinates, generating the photoproducts in the slower 5 ps decay process. They further confirmed using DFT calculations that the CO stretching, CO out-of-plane wagging (pyramidalization), and the ring-puckering modes are the primary nuclear vibrations in the S_1 surface,

which should transfer the excess energy to the ring-opening modes through IVR to generate the photo-products.

A static theoretical study²¹ using CASSCF and MS-CASPT2 methods revealed that the ring-opening mechanism has a barrier of 0.29 eV in the S_1 state, which can be easily surmounted by the excess energy of the S_1 state after 4.1 eV photoexcitation. After that, the internal conversion to S_0 is downhill through an S_1/S_0 conical intersection of 3.5 eV energy. This offers a fast pathway that should dominate over an intersystem crossing followed by the ring opening in the T_1 state. Ab initio multiple spawning dynamics using SA(2)-CASSCF(12,11)/6-31G* was employed following photoexcitation to the S_1 state of CB. The S_1 lifetime was estimated to be 484 fs.²² An average time of 177 fs was estimated for the α -cleavage. After reaching the ground state, two comparable bond-breaking mechanisms, involving either a second α - or a β -cleavage, led to the C3 and C2 product formation with a 1.03 C3:C2 ratio.

Using a combination of time-resolved mass spectrometry (TR-MS) and time-resolved photoelectron spectroscopy (TR-PES), the decay of the $n \rightarrow 3s$ Rydberg state (S_2) to the $n\pi^*$ state was characterized²³ by two time constants of 0.08 ps and 0.74 ps. A low-frequency ring-puckering mode of 35 cm^{-1} was suspected to promote the transition from the S_2 to the S_1 state. A five-dimensional (ring-puckering, CO out-of-plane, CO stretching vibrations, and symmetric and asymmetric α C-C stretchings) vibronic coupling Hamiltonian model was constructed²⁴ using EOM-CCSD to model cyclobutanone's four lowest singlet electronic states. MCTDH quantum dynamics simulations using this model Hamiltonian predicted a biexponential S_2/S_1 internal conversion with time constants of 0.95 ps and 6.32 ps.

Regardless of previous investigations, a complete description following the S_2 photoexcitation of cyclobutanone to the end-product formation is yet to be seen. To deal with all this uncertainty, we simulated the photoinduced dynamics using surface hopping methodology combined with different electronic structure methods. These results allow us to discuss the S_2 lifetime, the internal conversion mechanisms, the photoproducts, and the dependence of the outputs on the method choice. Gas phase ultrafast electron diffraction (GUED) patterns are provided as a side product to compare the upcoming experimental results.

2 Computational Methods

2.1 DC-FSSH Dynamics

DC-FSSH is a mixed quantum-classical nonadiabatic dynamics methodology, where a swarm of classical and independent trajectories approximates the nuclear wavepacket evolution through multiple electronic states. Each trajectory is propagated on a single potential energy surface (PES) calculated by solving the time-independent Schrödinger equations for the electrons within the Born-Oppenheimer approximation. During the propagation, a stochastic algorithm,²⁵ considering the instantaneous probability of nonadiabatic transitions between electronic states, dictates whether the trajectory will hop to another electronic state or stay on the current one. The DC-FSSH algorithm generally needs the electronic energies, force (negative gradient of potential energy), and nonadiabatic couplings from the electronic structure calculations. It is structured so that global energy surfaces are never required, and an on-the-fly approach, computing electronic properties as needed, can propagate the trajectories.

The hop probabilities are computed with a local approximation of the time-dependent Schrödinger equation for the electron and depend on the time-derivative couplings $\sigma_{JL} = \left\langle J \left| \frac{d}{dt} L \right\rangle$, between adiabatic electronic states J and L . These couplings can be obtained from the first-order nonadiabatic coupling vector (NACV), $\mathbf{h}_{JL} = \left\langle J \left| \frac{\partial}{\partial \mathbf{R}} L \right\rangle$, or other approaches such as the Hammes-Schiffer-Tully²⁶ and local diabaticization methods,²⁷ both relying on wavefunction overlaps.²⁸ Alternatively, the coupling can be estimated from the PES topography through the time-dependent Baek-An (TDBA) approach,²⁹ as demonstrated in Refs.^{30–32}

It is well established that the original FSSH algorithm suffers from over-coherence, as state superpositions, eventually born during the time-dependent Schrödinger equation propagation, are never eliminated, yielding an inconsistency between trajectories lying on single electronic states and entangled electronic populations. Different approaches have been proposed to deal with this problem.^{33,34} In this work, we apply the simplified decay of mixing (SDM), which exponentially scales down the electronic population of unoccupied states with a decoherence lifetime estimated from the energy gaps and nuclear kinetic energies.³⁵

We carried out different sets of DC-FSSH nonadiabatic dynamics using NACVs as obtained by solving the time-independent Schrödinger equation and with approximated TDBA couplings to accelerate the dynamics. The TDBA model used analytical second derivatives from quadratic regression ($\Delta T =$

0.4 fs) and $\delta\eta = 0.1$ a.u. In both cases, the Newtonian equations of motion were integrated using the Velocity Verlet algorithm with a timestep of 0.5 fs. The time-dependent Schrödinger equation was integrated using interpolated electronic quantities between classical steps with a time step of 0.025 fs. The decoherence parameter in SDM was $\alpha = 0.1$ a.u.³⁵ After a successful hop, the total energy was balanced by rescaling the nuclear velocity in the direction of the NACV when available. In DC-FSSH with TDBA couplings, the nuclear velocity was adjusted in the momentum direction. In this case, to decide whether hopping from a lower into a higher electronic state was allowed or frustrated, we used the reduced kinetic energy reservoir discussed in Ref.³¹ to ensure size-extensivity. In the case of a frustrated hopping, the momentum direction was not changed.

The initial geometries for the dynamics were randomly sampled from a marginal Wigner distribution for the quantum harmonic oscillator in the electronic ground state equilibrium geometry. Initial momentum was selected to have the total energy of each trajectory equating to the zero point energy.³⁶ The initial energy excitation and number of trajectories in each set are shown later in Section 3.2. We have sampled one thousand initial conditions for each one of the electronic structure methods. These geometries were used to simulate the absorption spectra of CB in the gas phase with the nuclear ensemble approach.³⁷

We performed several sets of DC-FSSH nonadiabatic dynamics simulations in the gas phase, considering cyclobutanone's lowest three singlet electronic states (S_0 , S_1 , and S_2). For each of these sets, we used different electronic structure methods and nonadiabatic coupling approaches, as follows:

- Set-1: DC-FSSH with NACV and multiconfigurational self-consistent field (MCSCF) [MCSCF-NACV];
- Set-2: DC-FSSH with TDBA and MCSCF [MCSCF-TDBA];
- Set-3: DC-FSSH with TDBA and linear-response time-dependent density functional theory (TDDFT) for excited states and DFT for the ground state [TDDFT-TDBA];
- Set-4: DC-FSSH with NACV and multireference configuration interaction (MRCI) based on the semiempirical orthogonalization and dispersion corrections method ODM3 [MNDO-NACV].
- We also report our experience attempting to use DC-FSSH with TDBA in combination with either coupled cluster with approximated second-order (CC2) or algebraic diagrammatic construction to the second-order [ADC(2)].

The trajectories were initiated on the S_2 electronic state in all the simulations. Since the molecule has enough excess energy to surmount the fast ring-opening and deactivation channels in the S_1 state, we

can safely neglect²¹ any possible role of triplet states having intersystem crossing on nanosecond timescales¹⁷ in this photoinduced dynamics. Moreover, the other works^{38–42} submitted to this prediction challenge that considered the triplet states in the nonadiabatic dynamics showed that the intersystem crossing process is negligible in the relaxation process of photoexcited cyclobutanone from the S_2 state.

Dynamics based on active-space methods, such as CASPT2, MRCI, MCSCF, or CASSCF, are usually unstable due to orbital exchange between subspaces when the molecule explores the configurational space.⁴³ This instability commonly implies total energy conservation failure. Although our MCSCF reduces this problem by providing enough flexibility to describe even dissociated structures, we still observed active space instabilities in Set-1 and Set-2 trajectories with MCSCF near S_2/S_1 conical intersection featuring α -cleavage. As a result, trajectories entering the crossing region via α cleavage underwent a sudden change of potential energy, leading to total energy change. To overcome this problem, we leveraged these trajectories by imposing moderately loose energy conservation thresholds in our dynamics. This loosening protocol can be reasoned by safely assuming that the trajectories passed the CI region at the crashing point and effectively funneled down to the S_1 state. Consequently, the trajectory continued steady propagation with the altered total energy.

We also followed the same loose energy conservation protocol in Set-4 simulations with MRCI/ODM3, which also suffers from similar active-space instabilities. On the other hand, because TDDFT is not subjected to such a problem, Set-3 trajectories were propagated with strict energy conservation criteria.

Due to the limitation of TDDFT in dealing with the multireference ground states, trajectories were stopped whenever the energy difference between S_1 and S_0 states dropped below 0.1 eV. In several cases, however, the trajectories stopped with a small S_1/S_0 energy gap due to convergence errors in DFT calculations. In these cases, the trajectories were considered as they reached the energy difference criterion. The time at which this threshold was attained was considered as the S_1/S_0 internal conversion time.^{44,45} A few trajectories succeeded in hopping to the ground state and were continued for a few fs in this state. All trajectories were restarted in the S_0 at the DFT level using the velocities and positions from the last point they reached with the TDDFT dynamics.

The details of the electronic structure methods are given in the following sections. The initial conditions, absorption spectra, and dynamics simulations were done using the Newton-X CS suite of programs (version 2.7).³⁶

2.2 Electronic Structure Methods

Set-1 and Set-2 dynamics simulations were carried out with the state-averaged MCSCF method based on the graphical unitary group approach (GUGA) using Columbus software (version 7.2.2, Sept. 2022)^{46,47} utilizing analytic energy gradients and nonadiabatic coupling procedures as described in Refs.^{48–50} The MCSCF wave function was constructed within an extended general valence bond (GVB) perfect pairing (PP) multiconfigurational approach,^{51–53} combining a complete active space (CAS) with PP subspaces. Such an MCSCF scheme yields a compact representation of the wavefunction light enough to allow extended dynamics simulations but still robust enough to describe the excitation and occasional dissociation. This type of wavefunction has been successfully used previously for the photodynamics of ethylene,⁵⁴ azomethane,^{50,55} and in extended form for the accurate calculation of molecular bond lengths.⁵⁶

The active space in the MCSCF wavefunction comprises one complete active space (CAS) and five PP subspaces. The CAS subspace contains four electrons in four orbitals [CAS(4,4)], including a nonbonding lone pair orbital n , a 3s Rydberg orbital originating at the oxygen atom, and the C=O π and π^* orbitals. (The molecular orbitals are presented in Section S1 of the Supplementary Material (SM). This part of the wavefunction was designed to describe the electronic ground state and the two excited states ($n\pi^*$ and $n \rightarrow 3s$). Each of the five PP subspaces contains two electrons in two orbitals, allowing only double occupation in each orbital. These orbital pairs are the four C-C σ and σ^* orbitals and the C-O σ and σ^* orbitals. The four PP C-C subspaces allow independent bond breaking, enabling the description of the opening of the CB ring. The remaining electrons were kept in doubly occupied orbitals. This MCSCF wavefunction provides a smooth transition to the ground state and a continuation with the hot ground state dynamics. Besides that, it has only 640 configuration state functions (CSFs). A schematic representation of the MCSCF wavefunction is

$$\text{MCSCF}(14,14) = 4[\text{PP}_{\text{CC}}] \times [\text{PP}_{\text{CO}}] \times [\text{CAS}(4,4)_{n,\pi,\pi^*,3s}]$$

Three electronic states were computed with a state averaging procedure (SA-3). It was used with the cc-pVDZ basis set for hydrogen, the aug-cc-pVDZ for carbon, and the d-aug-cc-pVDZ for oxygen.⁵⁷ This hybrid basis set describes all orbitals well and delivers good excitation energies at reasonable computational demand. Such a balance is essential given the vast number of MCSCF calculations in the on-the-fly dynamics propagation.

Minima on the crossing seam were calculated with MCSCF using COLUMBUS's analytic energy gradient and nonadiabatic coupling vector features. These optimizations were performed with a modified

version of GDIIS as described using natural coordinates.^{49,58,59} For simplicity, these minima on the crossing seam are denoted conical intersections (CIs) in the following text.

The Set-3 simulations were performed with TDDFT followed by DFT dynamics in the ground state using the CAM-B3LYP functional.⁶⁰ The basis set consists of cc-pVDZ with an augmented set of *1s1p1d* diffuse functions for the oxygen atom. All the DFT (ground state optimization, frequency calculations) and TDDFT calculations were done using ORCA v5.0.4.⁶¹ Due to the limitation of the TDDFT method in dealing with the multireference ground states, trajectories were stopped whenever the energy difference between S_1 and S_0 states dropped below 0.1 eV. In several cases, however, the trajectories stopped with a small S_1/S_0 energy gap due to convergence errors in DFT calculations. In these cases, the trajectories were considered as they reached the energy difference criterion. The time at which this threshold was attained was considered as the S_1/S_0 internal conversion time.^{44,45} A few trajectories succeeded in hopping to the ground state and were continued for a few fs in this state. All trajectories were restarted in the S_0 at the DFT level using the velocities and positions from the last point they reached with the TDDFT dynamics.

In Set-4, we employed the multireference configuration interaction with single and double excitations (MRCI) based on ODM3 as implemented in the MNDO program.^{62,63} The MRCI calculation employed an active space of 3 occupied (6 electrons) and three unoccupied orbitals. The number of references was automatically selected so that these configurations contributed at least 85% of the weight in all three requested CI roots. Energy gradients and NACV were computed at the same level.

Additional calculations were performed with the resolution-of-identity (RI) CC2 and ADC(2) using the Turbomole v7.6 program.⁶⁴ For these calculations, the cc-pVDZ basis set was employed for hydrogen and carbon, and the d-aug-cc-pVDZ was used for the oxygen.

3 Results and Discussion

3.1 Potential Energy Topography

The ground state (S_0) equilibrium geometry of the gas phase CB computed by all methods is nearly planar. At the S_0 minimum, the first excited singlet state (S_1) has a dark $n\pi^*$ character, and the second excited singlet state (S_2) is a bright Rydberg state originating from a $n \rightarrow 3s$ transition. MCSCF shows an S_1 minimum nonplanar with the ring puckering dihedral angle of 14° , and the out-of-plane carbonyl dihedral angle is 146° (Figure 2). The carbonyl bond length increases from 1.2 Å in the ground state to 1.4 Å in the S_1 $n\pi^*$ state. Also using MCSCF, the S_2 minimum is again planar and similar to the ground

Table 1 collects the excitation energies and oscillator strengths of the first and second singlet excited states at the ground state minimum by different methods and compares them with previously computed values. Taking EOM-CCSD/cc-pVTZ²⁴ as the reference result, our MCSCF underestimates the vertical excitation into S₁ by about 0.2 eV and overestimates the S₂ energy by approximately the same amount. This can be rationalized by the unbalanced account of dynamic electron correlation in MCSCF.

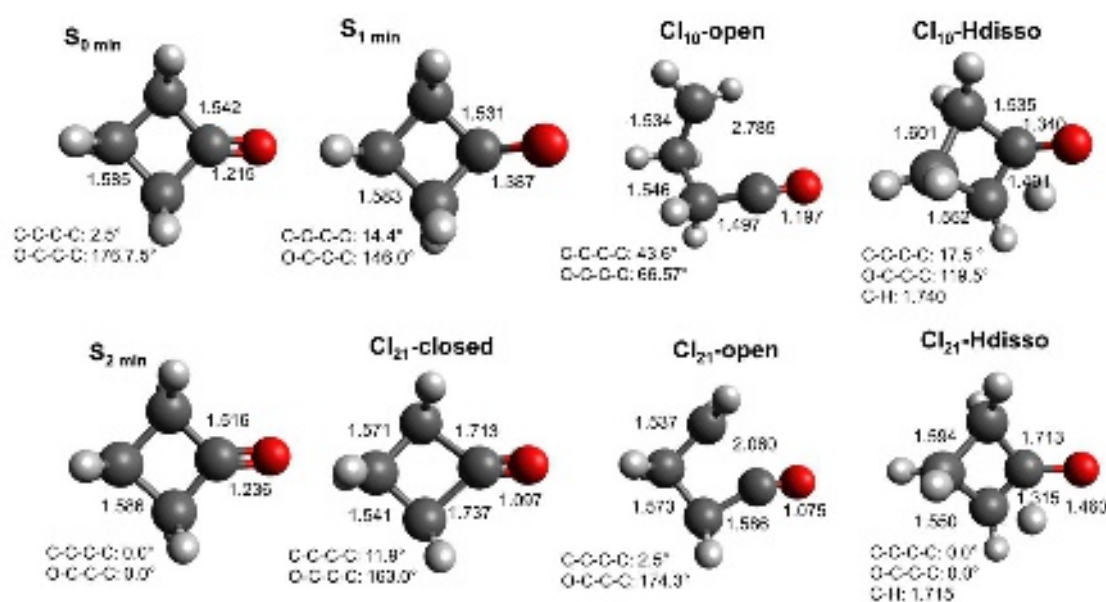


Figure 2: Geometries of minima and minima on the crossing seam (conical intersections) computed with MCSCF.

Table 1: Excitation energies (eV) and oscillator strengths (given in parentheses) at the S_0 minimum with MCSCF, TDDFT, and MRCI/ODM3. Previously reported computational results are shown as well.

Methods	$S_1(n\pi^*)$	$S_2(n \rightarrow 3s)$
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Present	MCSCF	4.24 (0.000)	6.77 (0.012)
	TDDFT	4.28 (0.000)	6.44 (0.043)
	MRCI/ODM3	4.09 (0.000)	6.35 (0.306)
Other works	SA(2)-CAS(12,11)/6-31G* [Ref. ²²]	4.39	
	MS-CASPT2//CASSCF(8,7)/6-31+G* [Ref. ²¹]	4.10	
	EOM-CCSD/cc-pVTZ+1s1p1d [Ref. ²⁴]	4.45	6.60
	XMS(3)-CASPT2(8/8) [Ref. ⁶⁵]	4.37 (0.000)	6.25 (0.017)

The excitation energies calculated at different critical points along the reaction coordinate, including S_2 and S_1 minima and several conical intersections, are given in Table 2. At the S_1 minimum, the MCSCF S_1 excitation energy matches perfectly with the EOM-CCSD result. Since the S_2 minimum structure is very similar to the ground state minimum, the electronic energies are also near their corresponding Franck-Condon values.

Conical intersections were computed only using MCSCF. We identified three S_2/S_1 conical intersections (CI) types with distinctly different structural and electronic features. (see Table 2). An S_2/S_1 CI with an α ring-opening and planar structure (CI_{21-open}) lies 0.1 eV above the S_2 minimum. The diabatic character of the crossing electronic states is $S_2(\sigma_{CC} \rightarrow 3s)$ and $S_1(\sigma_{CC} \rightarrow \pi^*)$. Another S_2/S_1 CI with a closed and slightly puckered ring (CI_{21-closed}) exists 0.5 eV above the S_2 minimum. In this CI, the diabatic character of the crossing electronic states is the same as at the S_2 minimum, $S_2(n \rightarrow 3s)/S_1(n \rightarrow \pi^*)$. A third S_2/S_1 CI, featuring a dissociated hydrogen initially bonded to an α carbon atom, lies 0.6 eV above the S_2 minimum (CI_{21-Hdisso}). The electronic transitions associated with this CI are $S_2(n \rightarrow \sigma_{CH}^*)/S_1(\pi_{CO} \rightarrow \sigma_{CH}^*)$. The structures, electronic states, and molecular orbitals of the CIs can be visualized in SM S3.

Two S₁/S₀ CIs were found: one showing a α -cleavage (CI₁₀-open) and the other with a dissociated-H (CI₁₀-Hdisso). CI₁₀-open lies 3.98 eV above the ground state minimum (see Table 2), 0.46 eV higher than the previously reported MS-CASPT2 estimate.²¹ It corresponds to an S₁($\sigma_{CC} \rightarrow \pi_{CO}^*$)/S₀(closed shell) crossing. CI₁₀-Hdisso lies at 6.1 eV, with an S₁($n \rightarrow \sigma_{CH}^*$)/S₀(closed shell) crossing. Both CIs are energetically accessible after the initial photoexcitation into S₂.

Table 2: Vertical excitation energies (eV) at the S_1 and S_2 minima and optimized conical intersections calculated with MCSCF. The energies are compared with the previously reported values with different levels of electronic structure theories.

	S ₁ min	S ₂ min	CI ₂₁ -open	CI ₂₁ -closed	CI ₂₁ -Hdisso	CI ₁₀ -open	CI ₁₀ -Hdisso
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S ₀	1.22 1.04 [⊥]	0.07 0.65 [⊥]	3.77 3.89 [⊥]	2.39 2.25 [⊥]	6.80	3.98 3.67 [⊥]	6.118
S ₁	3.20 3.71 [⊥] 3.84 [#] 3.50 ^{\$} 3.22 ^{##}	4.13 4.71 [⊥]	6.783 4.085 [⊥]	7.196 6.48 [⊥]	7.31	3.98 3.67 [⊥] 2.90 [#] 3.52 ^{\$}	6.123
S ₂	8.14 7.43 [⊥]	6.70 5.90 [⊥]	6.785 4.085 [⊥]	7.197 6.48 [⊥]	7.31	7.29 5.09 [⊥]	11.19

[#]SA(2)-CAS(12,11)/6-31G*;²² ^{##}CASSCF(10,8)/6-31G(d);¹⁷ ^{\$}MS-CASPT2//CASSCF(8,7)/6-31+G*;²¹ ^{\$\$}EOM-CCSD/cc-pVTZ+1s1p1d;²⁴ [⊥]XMS(3)-CASPT2(8,8)/aug-cc-pVDZ.⁶⁵

Considering the topography of the PESs computed by MCSCF (see Figure 3), we may hypothesize that after photoexcitation into the S₂ state (6.77 eV), CB will relax to the S₂ minimum (6.70 eV). From there, it may preferentially convert to S₁ through the CI₂₁-open, which is energetically nearby (Pathway 1). Then, CB would relax on S₁ via the ring-opening mechanism. Nevertheless, we cannot discard competing mechanisms where CB would convert to S₁ at the other two CI types by surmounting barriers of about 0.5 eV energy, which is well within the available energy after photoexcitation (Pathways 2 and 3). Since the S₁ minimum is a puckered closed-ring structure, the molecules relaxing from the S₂ state via the α -cleavage and H-dissociated CIs would not likely relax to the S₁ minimum. Instead, they may undergo rapid internal conversions to the ground state via S₁/S₀ CIs.

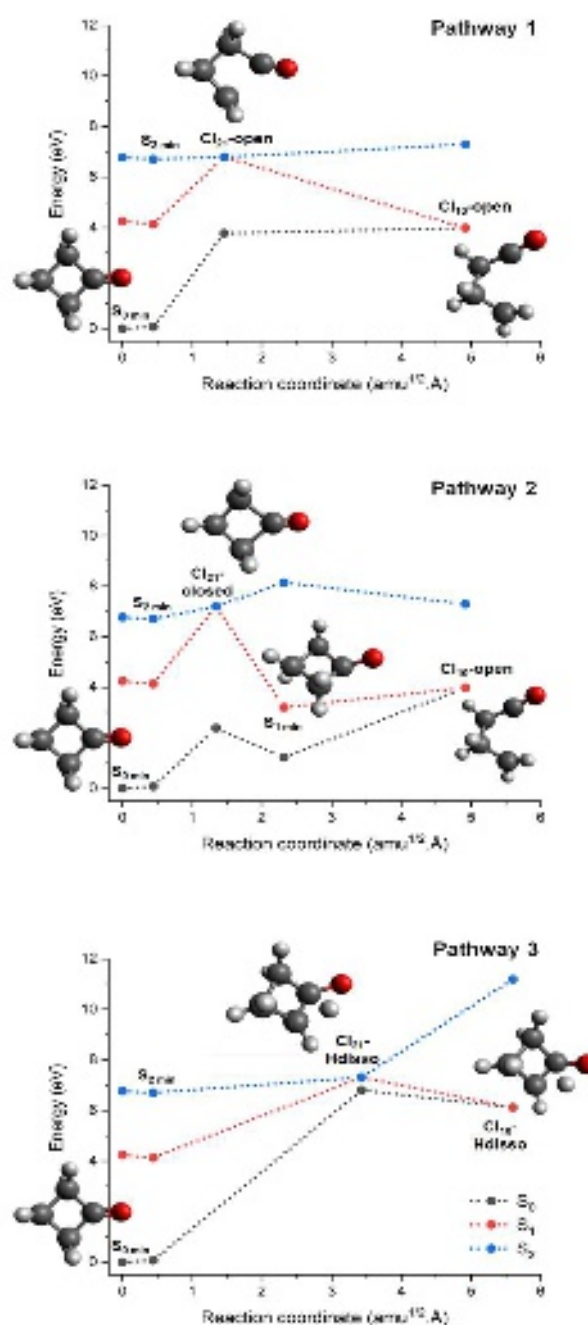


Figure 3: Topography of the S_0 , S_1 , and S_2 PESs involved in the dynamics calculated with MCSCF.

3.2 Absorption Spectra

Figure 4 shows the CB absorption spectrum in the gas phase up to the S_2 ($n \rightarrow 3s$) band computed with the nuclear ensemble approach. Compared to the experiments, MCSCF predicts an S_2 ($n \rightarrow 3s$) band

blue-shifted by 0.32 eV and underestimates the oscillator strength by a factor of 4.⁶⁶ MRCI/ODM3 strongly overestimates the oscillator strengths, likely due to the compactness of its wavefunction representation, but delivers reasonable S_2 band energies. TDDFT/CAM-B3LYP delivers an excellent S_2 ($n \rightarrow 3s$) band prediction thanks to the diffuse basis set and the range-separated functional. The S_1 band intensities of the experimental,⁶⁶ MCSCF, and TDDFT spectra are too small to appear on this scale. They are shown in the SM S4.

The proposed upcoming experiment will be conducted by exciting the S_2 state with a 200 nm (6.2 eV) laser. Comparing the shape and the position of the simulated S_2 band with the experimental one, we randomly sampled initial conditions for our dynamics simulations constrained to the 6.65 ± 0.03 eV excitation energy window (shaded area marked by A in Figure 4) for MCSCF-NACV (Set-1) and MCSCF-TDBA (Set-2) simulations. On the other hand, an energy window of 6.20 ± 0.05 eV (shaded area marked by B) is used to sample initial conditions for TDDFT-TDBA (Set-3) and MRCI/ODM3-NACV (Set-4) simulations. Hence, we have 93 independent trajectories for Set-1, 88 for Set-2, 78 for Set-3, and 190 for Set-4.

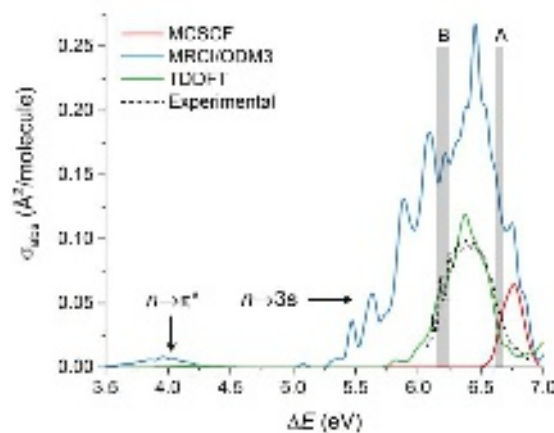


Figure 4: Simulated and experimental⁶⁶ absorption spectra of cyclobutanone in the gas phase, emphasizing the S_2 ($n \rightarrow 3s$) band. The shaded areas A and B indicate energy windows from where initial conditions for dynamics were selected (see text for details).

3.3 Population Time Evolution

Figure 5 shows the average trajectory occupation (fraction of trajectories in a particular state in a given timestep) in the three states as a function of time for different sets of simulations. We fitted the S_2

occupation with a mono-exponential decay function, $f(t) = e^{-(t-t_{lag})/t_{exp}}$, where t_{exp} is the exponential decay time constant of the S_2 state and t_{lag} refers to the time after the S_2 occupation starts to decay. The lifetimes is $\tau = t_{lag} + t_{exp}$. These parameters are tabulated in Table 3.

Table 3: S_2 lifetimes (τ), lag times (t_{lag}), and exponential decay times (t_{exp}) as obtained from different sets of simulations by a mono-exponential decay fit.

Simulations	t_{lag} (ps)	t_{exp} (ps)	τ (ps)
MCSCF-NACV (Set-1)	0.3	8.7	9.0
MCSCF-TDBA (Set-2)	0.3	10.1	10.4
TDDFT-TDBA (Set-3)	0.00	0.40	0.4
MRCI/ODM3-NACV (Set-4)	0.005	0.035	0.04
ADC(2)-TDBA	0.0	0.9	0.9

With MCSCF (Set-1 and Set-2), the S_2 state starts to decay after 300 fs, and its occupation reduces to ~30% after 10 ps. The predicted lifetimes of the S_2 state for Set-1 and Set-2 simulations are 9 and 10 ps, respectively. It is also interesting to note that the S_1 state is practically never populated. Once the trajectories hop down from the S_2 to the S_1 state, they immediately undergo another hopping event to the ground state. The mean time difference between the two successive hops (i.e., the difference between S_2/S_1 and S_1/S_0 hopping times) is 16 fs. Hence, the trajectories never relax to the S_1 minimum, as predicted by considering the topography of the PESs. The occupation profiles obtained by MCSCF-NACV (Set-1) and MCSCF-TDBA (Set-2) match closely considered the statistical uncertainty. The bootstrapped 95% confidence interval (with 10^5 resamplings) is 1 and 2 ps for Set-1 and Set-2, respectively. Therefore, these results validate using approximated TDBA couplings.

TDDFT-TDBA (Set-3) simulations show a starkly different result than the MCSCF dynamics. It predicts a much slower population transfer compared to Set-4. After photoexcitation to the S_2 state, CB immediately starts to decay (having no lag time) to the S_1 state. The population is entirely transferred to the ground state, yielding an S_2 lifetime of 0.4 ps. The population starts accumulating in the S_1 state from the beginning and attains a maximum of 35% around 200 fs. Afterward, it exponentially decays to the ground state in the S_1 state population.

MRCI/ODM3-NACV dynamics have features similar to TDDFT-TDBA but a ten times faster timescale. In this case, the S_2 population completely decays to the ground state within 200 fs, predicting its lifetime as 40 fs. The growth of the S_1 population reaches a maximum value of 60% around 50 fs and then exponentially decays to the ground state.

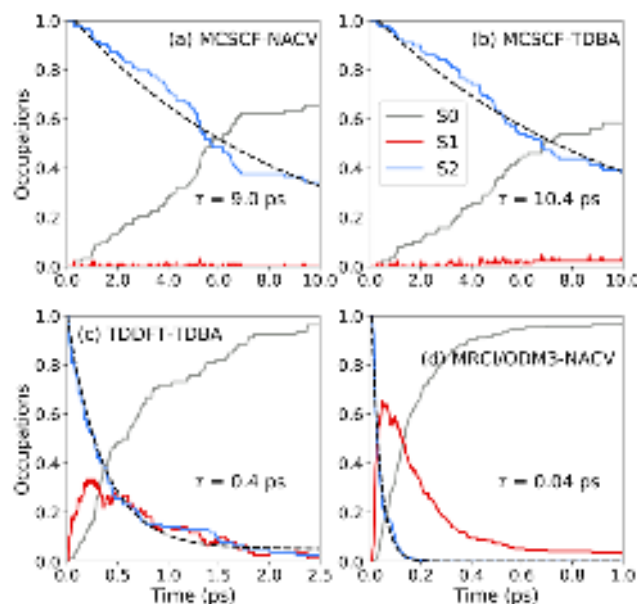


Figure 5: The time evolution of average occupations of the three states for different sets of simulations.

Our attempts to run the dynamics using CC2 and ADC(2) methods did not work well, as expected. Using CC2, most of the trajectories (77 out of 91) were finished too early due to convergence problems in the electronic structure calculations caused when S_2 and S_1 were energetically close. On the other hand, the trajectories using ADC(2) were less problematic in terms of convergence: only 8 out of the 80 trajectories stopped due to eigenvalue convergence failures, already in the S_1 state. However, although the hopping geometries from S_2 to S_1 were reasonably described, the S_1/S_0 hoppings were not correctly predicted, as most of the geometries show an abnormal stretching of the C=O bond above 1.5 Å (see SM S5). This problem with ADC(2) has been recently discussed in Ref.¹⁰ The fitting of the S_2 fraction of population decay provided this state a lifetime of 0.9 ps.

3.4 Internal Conversion Mechanisms

As speculated from the PES topography (see Section 3.1), indeed, after photoexcitation to the S_2 state, CB may convert to S_1 via one of the three different pathways related to the three S_2/S_1 conical

intersections (CI_{21} -open, CI_{21} -closed, and CI_{21} -Hdisso). Figure 6 depicts the distribution of the S_2 - S_1 energy gaps at the S_2/S_1 hopping geometries near each of these conical intersections using different methods.

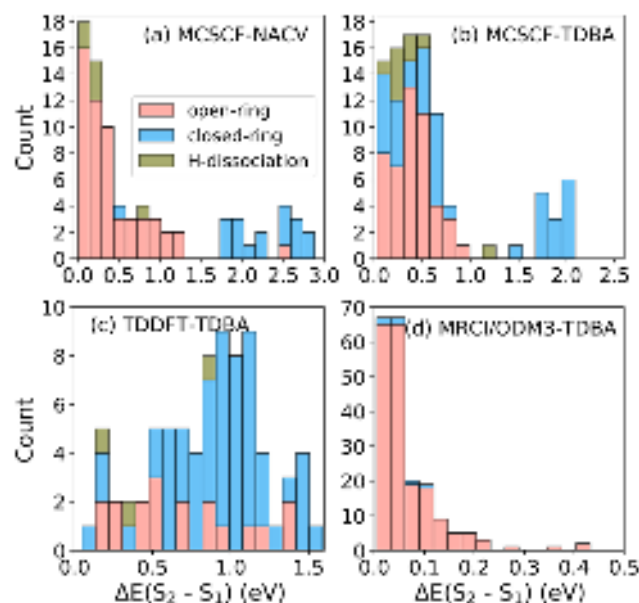


Figure 6: Distribution of energy gaps between S_2 and S_1 states at the S_2/S_1 hopping geometries for the four sets of simulated trajectories.

In MCSCF-based dynamics (Set-1 and Set-2), most of those trajectories (47%) approached the CI_{21} -open via the ring-open mechanism. The energy gaps of hopping geometries peaking at lower energy signifies attaining the CI_{21} -open region in the MCSCF-NACV dynamics [see panel (a) of Figure 6]. A few trajectories (5%) went uphill, surmounted the higher-energy CI_{21} -Hdisso, and converted to the S_1 state. In MCSCF-TDBA dynamics, we also observe similar S_2/S_1 internal conversion mechanisms. The distribution of the ring-puckering angles at the hopping geometries is given in SM S6. The S_2/S_1 hoppings mainly occur with small ring-puckering geometries. However, some hoppings occur around a 60° C-C-C-C dihedral angle, revealing that the S_2/S_1 crossing seam extends into twisted configurations. A final observation about MCSCF simulations: 30% of the trajectories remained unreactive in the S_2 state until the end of 10 ps simulation time. This can also be seen in the remaining S_2 population in panels (a) and (b) of Figure 5.

In Set-3 of TDDFT-TDBA dynamics, most of the trajectories (73%) underwent closed-ring S_2/S_1 hops with minor and moderate energy gaps. In this case, the ring-opening mechanism (23%) at the S_2 state seems not dominant.

In Set-4 of MRCI/ODM3-TDBA dynamics, almost all trajectories (92%) were converted from S_2 to the S_1 state via the ring-opening pathway.

At the S_1 state, most of the trajectories (57% of the total and 86% of the S_1 state) were further converted to the ground state with ring-opened structures in the MCSCF-NACV (Set-1) simulation (see Figure 7). Only 7% of the total trajectories (11% of trajectories in the S_1 state) reached the ground state via the CI_{10} -Hdisso channel. It means that the closed-ring trajectories in the S_1 state also underwent a C-C bond cleavage to reach the S_0 state. The other simulations with MCSCF-TDBA, TDDFT-TDBA, and MRCI/ODM3-NACV also had the ring-open as the dominant path for the S_1/S_0 internal conversion.

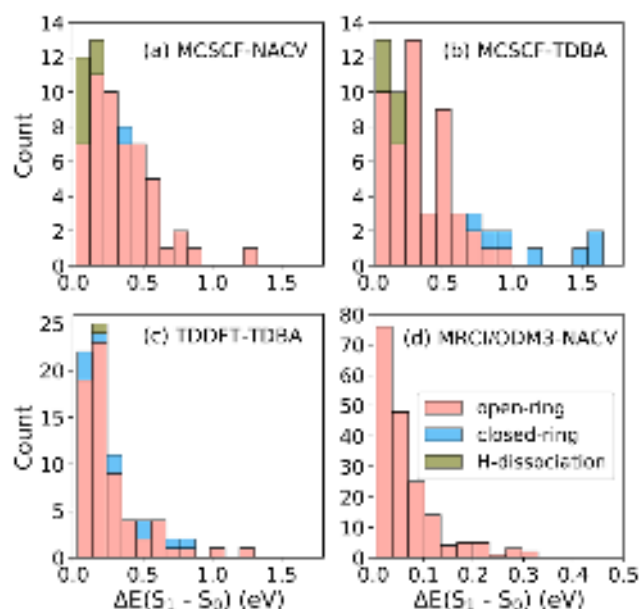


Figure 7: Distribution of energy gaps between S_1 and S_0 states at the S_1/S_0 hopping geometries for the four sets of simulated trajectories.

A final observation should not be missed here: 16% of MCSCF-NACV trajectories attaining closed-ring structures show S_2/S_1 hoppings with energy gaps between 2 and 3 eV (See the top panels in Figure 6). At first, it seems unusual but not impossible with the instantaneous probabilities in the framework of FSSH dynamics in the weak coupling regime.⁶⁷ Given the CI_{21} -closed is energetically uphill, the molecule never accessed the conical intersection, and the trajectories spent a long enough time

with moderate energy gaps, with low transition probability (small NACV) but with many hopping opportunities until a small enough random number triggered the stochastic hopping. This can be confirmed by the scatter plots in SM S9 between the random numbers drawn and the transition probabilities at the S_2/S_1 hopping events. They show that hoppings associated with high energy gaps always occur due to tiny random numbers against low hopping probabilities. Similar high-energy hoppings are also observed in Set-2 simulation with MCSCF-TDBA. However, the S_2/S_1 hoppings tend to occur at smaller energies with TDBA than with NACV. This is an artifact of the way TDBA was set. These approximated couplings were computed only when the energy gaps were smaller than 1 eV. Thus, the hops were constrained to be below this threshold. However, increasing the threshold to 2 eV, the S_2/S_1 CI-closed hoppings occur at higher energies, as with NACV. We show these high-energy hoppings for MCSCF-TDBA in Figure 6, too. These large-gap hoppings in the S_2/S_1 closed structures were also observed with TDDFT (Figure 6) and ADC(2) (SM S5).

3.5 Photoproducts

Working with different data sets and trajectories finishing at different times, it was crucial to fix a protocol to classify the photoproducts in terms of geometrical parameters and the time to observe their geometries. Based on the trajectory survival time in S_0 , we chose to classify the photoproducts 200 fs after the S_1/S_0 hopping. In a few cases, we took the last geometry, when the trajectories did not survive till that pre-defined time due to the failure of electronic structure convergence, usually due to dissociation events.

To check for single-bond cleavage, we defined parameters based on $r_{mean} + 3\sigma$ values for each bonded pair by fitting the particular bond distances with a Gaussian distribution, considering all the molecular geometries obtained in the dynamics. In this fitting, r_{mean} is the mean and σ is the standard deviation. If the atomic distances exceeded $r_{mean} + 3\sigma$, we considered it a single bond cleavage.

After that, the C2 products, which involve double bond cleavage, were assigned if a β and an α CC distance were greater than 2.5 Å. Similarly, the C3 products were assigned if both α CC distances were greater than 2.5 Å. Thus, C3 products contain CO + cyclopropane or propene molecules. Although we observed propene conversion in cyclopropane in a few trajectories propagated already in the ground state, we did not classify them as different channels. Analogously, C2 products containing ethylene + ketene molecules or ethylene + CO + CH₂ fragments were counted in the same channel. We also verified at times longer than 200 fs that the C3 and C2 products remained dissociated.

We observed an interesting mechanism in very few trajectories. A hydrogen atom was detached from an α - carbon position and followed a roaming pathway toward the oxygen atom. Since this happened only in a few trajectories, we are not in a position to predict the roaming mechanism confidently. Thus, they were generally classified as H-dissociation.

Following this protocol, the photoproducts were classified as a closed ring, open ring (single bond cleavage), C2 (α and β cleavages), C3 (two α cleavages), and H dissociation. The amounts of each obtained in the different sets of simulations are tabulated in Table 4.

Table 4: Amount (%) of different photoproducts obtained after the trajectory spends 200 fs in the ground state.

Simulations	closed-ring	open-ring	C3 products	C2 products	H-dissociation
MCSCF-NACV (Set-1)	0	20	32	8	6
MCSCF-TDBA (Set-2)	7	21	28	1	7
TDDFT-TDBA (Set-3)	49	22	5	19	5
MRCI/ODM3-NACV (Set-4)	3	23	28	46	0

Set-1 and Set 2 do not sum to 100% because this table does not include trajectories that remained excited at the end of the simulations.

In the MCSCF-NACV and MCSCF-TDBA dynamics (Set-1 and Set-2), about 32% of the CB molecules remained in the S_2 state after 10 ps of simulation, showing no internal conversion to lower states. On the contrary, in the TDDFT-TDBA and MRCI/ODM3-TDBA dynamics (Set-3 and Set-4), the S_2 state is entirely converted to S_0 within the simulation time.

The presented data encapsulates the outcomes of molecular simulations conducted across various computational methodologies, offering insights into the diverse chemical pathways and products emerging from distinct electronic structure methods. In the following analysis, given the number of trajectories in each set, we consider statistical uncertainty of about $\pm 10\%$ for MCSCF and TDDFT simulations and $\pm 5\%$ for MRCI/ODM3.

In the MCSCF-NACV set, a notable prevalence of open-ring and C3 products is observed. MCSCF-TDBA simulations exhibit a similar trend within the statistical uncertainty, once more confirming the excellent performance of TDBA-approximated couplings. Thus, both MCSCF sets suggest a propensity toward CO elimination in the ground-state photoproducts.

The TDDFT-TDBA simulations portray a distinctive distribution profile characterized by a substantial predominance of closed-ring products alongside notable contributions from open-ring and C2 channels. Therefore, different from MCSCF, TDDFT predicts a predominance of ethylene formation.

The MRCI/ODM3-NACV set exhibits another pattern, characterized by a pronounced abundance of C2 products and a notable representation of C3 species while exhibiting minimal closed-ring occurrences. This distinctive distribution predicts equally large amounts of CO and ethylene eliminations.

3.6 Electron Diffraction Spectra

We simulated the GUED signal using the independent atom model (IAM) as prescribed by Centurion et al.⁷ This model assumes that the electron arrangement within a molecule mimics that of independent atoms with no interactions among them. As a result, we can safely work with the predetermined atomic force factors (AFFs) to describe the arrangements. However, the approximation overlooks the alterations in electron distribution resulting from the formation of chemical bonds. The working equations to calculate the signals can be seen in SM S7.

The elastic scattering signal is the incoherent sum over an ensemble of all the molecular structures obtained from the nonadiabatic dynamics simulation. To calculate the time-dependent real-space pair-distribution function (ΔPDF), we used a range of 0 to 12 Å⁻¹ as the integration limits of the momentum transfer (s) and a damping factor $\alpha = 0.03$ Å². We used a locally modified version of the “Diffraction_simulation” code,⁶⁸ originally developed by Wolf et al., to simulate electron diffraction signals. The AFFs were taken from the git repository (without verification) and were simulated using the ELSEPA program.⁶⁹ The GUED diffraction patterns obtained from different sets of simulations are displayed in Figure 8.

Figure 8a shows the PDF signal for the ground state geometry and the most critical contribution of specific atomic interactions to the signal. It is justified to use the same reference to evaluate all sets of simulations, given how well the signal of the initial conditions matches (see section SM S8). In all the methods, it is possible to identify a loss of signal at the distances of 1.5, 2.5, and 3.0 Å associated with the CC and CH distances, suggesting an increase in those distances. It is also possible to see a signal gain around 2 Å, particularly in the MCSCF sets, implying a relaxation of the molecule associated with longer CC bonds.

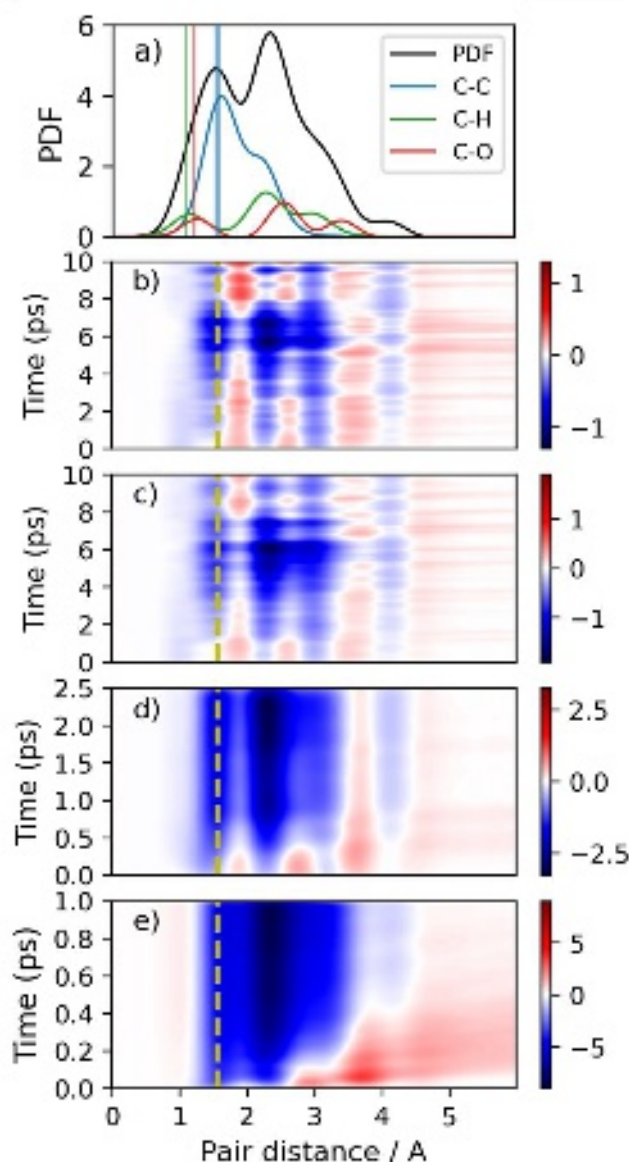


Figure 8: GUED patterns for different sets of simulation, a) shows the PDF for the ground state geometry and the contribution of different bonds. Vertical lines indicate the equilibrium distance; b-e) shows the Δ PDF for the dynamics with different methods, convoluted with 150 fs FWHM. b) MCSCF-NACV, c) MCSCF-TDBA, d) TDDFT-TDBA, e) MRCI/ODM3-NACV. The yellow dashed line marks the α and β CC bond distances of the equilibrium geometry of CB.

In the MCSCF sets, this initial profile is somewhat conserved during the whole simulation, caused by a combination of factors. Most importantly, the electronic structure used showed some instability in treating the system after hops, which led to the trajectories stopping after some time, varying from a few

hundred femtoseconds to a few picoseconds. This means that at any given time, most of the trajectories are still in the relaxed structure trapped in the S_2 , and only a few trajectories yield dissociation. This is complemented by the relatively long timescale of the simulated process, which causes the hops of different trajectories to happen at distant times from each other, leading to the visible horizontal lines at longer pair distances [panels a) and b) of Figure 8].

The calculations with TDDFT show an intermediate timescale associated with more stable trajectories that survived longer after hop. The first difference is in the fast disappearance of the positive signal around 2 Å and smoother gain in signal at longer pair distances.

Finally, MRCI/ODM3 also shows a smooth signal due to the short decay time and the higher number of trajectories allowed by this inexpensive method. The signal also starts with a relaxation phase with signal loss at 1.5, 2.5, and 3.0 Å and an almost unnoticeable signal gain around 2 Å. In this case, however, this relaxation is almost instantaneously followed by the dissociation process, which can be seen by the deep negative signal from 1 to 2.5 Å and a strong positive signal above 3 Å already at 100 fs. In this simulation, almost all dissociation has already occurred in the first half of the dynamics, after which the signal at longer distances appears in lighter red.

4 Conclusion

In this work, we explored the dynamics of cyclobutanone after photoexcitation to the S_2 ($n \rightarrow 3s$) Rydberg state using DC-FSSH nonadiabatic dynamics combined with different routine electronic structure methods, namely, MCSCF, TDDFT, MRCI/ODM3, CC2, and ADC(2). We also tested DC-FSSH using full and approximated nonadiabatic couplings.

We characterized the S_2 lifetime, internal conversion mechanisms, and photoproducts. The results radically depend on the electronic structure method used to propagate the dynamics, as summarized in Table 5.

MCSCF predicts a long-lived S_2 state, with internal conversions via the ring-open mechanism and predominantly following the C3 channel. Using approximated couplings with TDBA does not change this picture. These outcomes contrast with the TR-PES results from Ref.²³, which reported sub-ps S_2 lifetime with the dominance of the C2 channel.

All other methods predict much shorter S_2 lifetimes. TDDFT and ADC(2) also diverge from MCSCF, predicting S_2/S_1 internal conversion via the ring-closed mechanism. The S_1/S_0 internal conversion was consensually predicted to be the ring-open mechanism. TDDFT and MRCI/ODM3 also

diverge from MCSCF regarding the photoproduct. Both predict significant C2 product formation. Methodological limitations turned ADC(2) and CC2 inadequate for these simulations.

The reason behind the discrepancy between MCSCF and other methods are two-fold: First, while our specially designed MCSCF wavefunction is more suitable to describe dissociations, it lacks excitations between the PP and CAS subspaces; moreover, MCSCF delivers an unbalanced treatment of dynamic electron correlation, as discussed in Ref.⁷⁰ in another context. The combination of both shortcomings results in overestimating the S₂-S₁ energy gaps. To gaze both these effects, we have compared the relative energies (see Table 1 and Table 2) of a few critical points calculated by Janoš and coworkers⁶⁵ (also submitted to this prediction challenge) at XMS(3)-CASPT2(8,8)/aug-cc-pVDZ level of theory. Please note that the geometries are different in the MCSCF and CASPT2 calculations. Still, for the sake of argument, they should be similar and can provide a sensible estimate of CASPT2 results. The CASPT2 calculation substantially stabilizes the S₂ state and predicts higher S₁ energies than the MCSCF method. Hence, the S₂ and S₁ states are close to each other in CASPT2 compared to the MCSCF method, yielding faster dynamics in the former case.

Compared to the experimental results from Ref.²³, TDDFT-TDBA dynamics performed the best. However, such comparison defeats the purpose of the Prediction Challenge of making predictions on unseen empirical data. Suppose we had not viewed the prior time-resolved experimental data. In that case, we might question how to assign credence to our various theoretical simulations. The least credence would likely be assigned to the MRCI/ODM3 simulations, which should be regarded, at best, as exploratory of the configuration space. While TDDFT's outstanding predictive power for absorption spectra is a plus, describing photo-dissociation events using a single reference method would not be accurate enough. On the other hand, our specially designed MCSCF should correctly account for the dissociation processes, but it overestimates the S₂-S₁ energy gaps. Therefore, we do not have enough information to attribute higher credence to either method.

Table 5 Summary of predictions obtained from different dynamics methodologies with different electronic structure methods. Relevant secondary channels are not mentioned to be concise.

Dynamics ¹	Electronic structure	S ₂ lifetime (ps)	Dominant mechanism		Dominant photoproducts	Ref.
			S ₂ /S ₁	S ₁ /S ₀		
DC-FSSH (NACV)	MCSCF(14,14)	9.0	open ring	open ring	C3	Present work

DC-FSSH (TDBA)	MCSCF(14,14)	10.4	open ring	open ring	C3	Present work
DC-FSSH (TDBA)	TDDFT	0.4	closed ring	open ring	C2	Present work
DC-FSSH (TDBA)	MRCI/ODM3	0.04	open ring	open ring	C2	Present work
DC-FSSH (TDBA)	ADC(2)	0.9	closed ring	-	-	Present work
MCTDH (5D)	CCSD	0.95 / 6.32				Ref. ²⁴
MCTDH (8D)	TDA-TDDFT / PBE0	~0.5	closed ring	-	-	Ref. ³⁸
MCTDH (up to 27D)	CASPT2 / SA(6)-CASSCF(8,10)	~0.05 (diabatic)	-	-	closed ring	Ref. ⁷¹
DD-vMCG	SA(6)-CASSCF(8,10)	~0.05	-	-	closed ring	Ref. ⁷¹
AIMS	TDA-TDDFT / LRC- ω PBE	> 2	C α -H stretch	-	-	Ref. ⁷²
AIMS	EOM-CCSD	> 2	α C-C stretch	-	-	Ref. ⁷²
AIMC	SA(3)-CASSCF(12,12)	~0.2	-	-	C2	Ref. ⁷³
Generalized DC-FSSH	XMS-CASPT2(8,8) / CASSCF(12,12)	0.67	open ring	open ring	C3 and C2	Ref. ³⁹
DC-SHARC	XMS-CASPT2 / SA(4/4)-CAS(8,8)	~0.3	-	-	C3 and C2	Ref. ⁴⁰
Modified FSSH	MRCIS(6,6) / CASSCF(6,6)	0.025	closed ring	open ring	C3	Ref. ⁴¹
DC-SHARC	TDDFT / B3LYP-D3	~0.2	-	-	closed ring	Ref. ⁴²
DC-FSSH	SA(6)-CASSCF(8,11)	~0.025	-	-	open ring	Ref. ⁴²
DC-FSSH (TDBA)	TDA-TDDFT / PBE0	0.36	closed ring	open ring	C2	Ref. ³⁸
DC-FSSH	XMS-CASPT2(8,8)	0.33	open ring	open ring	C3	Ref. ⁶⁵
DC-FSSH	XMS(3)-CASPT2(10,8)	~0.05	closed ring	open ring	C3 and C2	Ref. ⁷⁴
FSSH	TDDFT / PBE0-D3	7	CC α H bend	α C-C stretch	closed ring	Ref. ⁷⁵
Ehrenfest- TAB	FOMO-CASCI (8,11)	~0.05	-	-	C3	Ref. ⁷⁶
MS-MASH	SA(3)-CASSCF(12,12)	~0.1	-	-	open ring and C2	Ref. ⁷⁷
unSMASH	CASSCF(12,11)	0.08	-	-	C3	Ref. ⁷⁸

¹Main acronyms definition: DC-FSSH: decoherence corrected fewest switches surface hopping; MCTDH: multiconfigurational time-dependent Hartree; DD-vMCG: direct dynamics variational multi-configurational Gaussian; AIMS: ab initio multiple spawning; AIMC: ab initio multiple cloning; SHARC: surface hopping with arbitrary couplings; Ehrenfest-TAB: Ehrenfest dynamics collapse to a block; MASH: mapping approach to surface hopping.

We commend the Journal of Chemical Physics for implementing this Prediction Challenge initiative. Metanalyses in psychology, ecology, and evolutionary biology have demonstrated that when different research groups are invited to analyze the same statistical data set, they may arrive at radically different conclusions.^{79–81} We see no reason that would differ in our field, maybe even more in the Prediction Challenge, considering that each group used different methodologies, software, and analysis protocols. At the time of writing, we still did not know how different the predictions of the diverse groups answering this challenge would be. However, based only on the striking differences in our simulations following the same protocols, we expect they would fit on a wide range.

During the manuscript revision, we saw the other computational works submitted for this Prediction Challenge^{38,39,75–78,40–42,65,71–74} but not the experimental results. In Table 5, we compiled a comparative summary of ours, Ref. ²⁴, and those works' key findings. We refrain from analyzing those differences here. Nevertheless, our expectations were confirmed by the wide range of results from different computational chemistry groups, even using similar electronic structure methods. Another interesting fact: Among the 15 computational works on CB photodynamics, some surface hopping variation (including FSSH, GTSH, SHARC, and MASH) was employed in 11, confirming its expected role as the most popular class of nonadiabatic molecular dynamics methods.

Nonadiabatic dynamics simulations have proved invaluable for helping assign experimental features, rationalize empirical data, and explore the configurational space to reveal excited state reaction mechanisms far from our chemical intuition. Nevertheless, these conflicting dynamics results we report here (and others have reported before; for instance, Refs. ^{11,12}) draw a worrisome picture of a field where routine nonadiabatic dynamics methods may have low prediction power. We hope the debate on these divergent results will lead us to more accurate and precise methodologies.

The philosopher of science, Imre Lakatos, proposed we should judge a research program as progressive if it predicted novel empirical facts and at least some novel facts could be tested.⁸² If not, the program is degenerating. Nonadiabatic dynamics constitutes a research program in Lakatos' sense of a core set of theoretical assumptions (the Born-Oppenheimer separability, for example) surrounded by a protective belt of auxiliary hypotheses (the need for nonadiabatic corrections). Therefore, the utmost importance of this Prediction Challenge is to unveil the best practices and development directions we must follow to guarantee that our research field remains progressive.

Supplementary Material

The supplementary material contains molecular orbitals at the S_0 minimum, optimized geometric parameters, conical intersections characterization, S_1 -band absorption spectrum, ADC(2) dynamics results, additional results from MCSCF dynamics, working equations to simulate GUED patterns, and GUED signals for initial conditions.

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Author Declarations

Conflict of Interest

MB is an advisory board member of J. Chem. Phys. and guest editor of this issue. The other authors have no conflicts to disclose.

Author Contributions

Supervision, Project Administration, Funding Acquisition: MB; Methodology: MB, HL; Analysis, Investigation: SM, RSM, JMT; Visualization: SM, RSM; Writing – Original Draft: SM, MB; Writing – Review & Editing: SM, RSM, HL, JMT, MB.

Data availability

The data supporting this study's findings are openly available in Zenodo (DOI: 10.5281/zenodo.10662306).

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