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# AN ESTIMATE FOR AN EQUATION WITH HOLDERIAN CONDITION AND WEIGHT AND BOUNDARY SINGULARITY.

SAMY SKANDER BAHOURA

ABSTRACT. We consider the following problem on open set  $\Omega$  of  $\mathbb{R}^2$ :

$$\begin{cases} -\Delta u_i = |x - x_0|^{2\beta} V_i e^{u_i} & \text{in } \Omega \subset \mathbb{R}^2, \\ u_i = 0 & \text{in } \partial\Omega. \end{cases}$$

Here,  $x_0 \in \partial\Omega$  and,  $\beta \in (0, 1)$ .

We assume, for example that:

$$\int_{\Omega} |x - x_0|^{+2\beta} V_i e^{u_i} dy \leq 16\pi - \epsilon, \quad \epsilon > 0$$

1) We give, a quantization analysis of the previous problem under the conditions:

$$\int_{\Omega} |x - x_0|^{+2\beta} e^{u_i} dy \leq C,$$

and,

$$0 \leq V_i \leq b < +\infty$$

2) In addition to the previous hypothesis we assume that  $V_i$   $s$ -holderian with  $1/2 + \beta/2 < s \leq 1$ , then we have a compactness result, namely:

$$\sup_{\Omega} u_i \leq c = c(b, C, A, s, \epsilon, \beta, x_0, \Omega).$$

where  $A$  is the holderian constant of  $V_i$ .

## 1. INTRODUCTION AND MAIN RESULTS

We set  $\Delta = \partial_{11} + \partial_{22}$  on open set  $\Omega$  of  $\mathbb{R}^2$  with a smooth boundary. Here we assume that we are on the unit disk.

We consider the following problem on  $\Omega \subset \mathbb{R}^2$ :

$$(P) \begin{cases} -\Delta u_i = |x - x_0|^{2\beta} V_i e^{u_i} & \text{in } \Omega \subset \mathbb{R}^2, \\ u_i = 0 & \text{in } \partial\Omega. \end{cases}$$

Here,  $x_0 \in \partial\Omega$  and,  $\beta \in (0, 1)$ .

We assume that,

$$0 \leq V_i \leq b < +\infty, \quad \int_{\Omega} |x - x_0|^{2\beta} e^{u_i} dy \leq C, \quad u_i \in W_0^{1,1}(\Omega)$$

The above equation is called, the Prescribed Scalar Curvature equation in relation with conformal change of metrics. The function  $V_i$  is the prescribed curvature.

Here, we try to find some a priori estimates for sequences of the previous problem.

Equations of this type (in dimension 2 and higher dimensions) were studied by many authors, see [1-24]. We can see in [8], different results for the solutions of those type of equations with or without boundaries conditions and, with minimal conditions on  $V$ , for example we suppose  $V_i \geq 0$  and  $V_i \in L^p(\Omega)$  or  $V_i e^{u_i} \in L^p(\Omega)$  with  $p \in [1, +\infty]$ .

Among other results, we can see in [8], the following important Theorem,

**Theorem A (Brezis-Merle [8]).** *If  $(u_i)_i$  and  $(V_i)_i$  are two sequences of functions relatively to the previous problem (P) with,  $0 < a \leq V_i \leq b < +\infty$ , then, for all compact set  $K$  of  $\Omega$ ,*

$$\sup_K u_i \leq c = c(a, b, K, \Omega).$$

A simple consequence of this theorem is that, if we assume  $u_i = 0$  on  $\partial\Omega$  then, the sequence  $(u_i)_i$  is locally uniformly bounded. We can find in [8] an interior estimate if we assume  $a = 0$ , but we need an assumption on the integral of  $e^{u_i}$ . We have in [8]:

**Theorem B (Brezis-Merle [8]).** *If  $(u_i)_i$  and  $(V_i)_i$  are two sequences of functions relatively to the previous problem (P) with,  $0 \leq V_i \leq b < +\infty$ , and,*

$$\int_{\Omega} e^{u_i} dy \leq C,$$

*then, for all compact set  $K$  of  $\Omega$ ,*

$$\sup_K u_i \leq c = c(b, C, K, \Omega).$$

If, we assume  $V$  with more regularity, we can have another type of estimates,  $\sup + \inf$ . It was proved, by Shafrir, see [23], that, if  $(u_i)_i, (V_i)_i$  are two sequences of functions solutions of the previous equation without assumption on the boundary and,  $0 < a \leq V_i \leq b < +\infty$ , then we have the following interior estimate:

$$C \left( \frac{a}{b} \right) \sup_K u_i + \inf_{\Omega} u_i \leq c = c(a, b, K, \Omega).$$

We can see in [12], an explicit value of  $C \left( \frac{a}{b} \right) = \sqrt{\frac{a}{b}}$ . In his proof, Shafrir has used the Stokes formula and an isoperimetric inequality, see [6]. For Chen-Lin, they have used the blow-up analysis combined with some geometric type inequality for the integral curvature.

Now, if we suppose  $(V_i)_i$  uniformly Lipschitzian with  $A$  the Lipschitz constant, then,  $C(a/b) = 1$  and  $c = c(a, b, A, K, \Omega)$ , see Brezis-Li-Shafrir [7]. This result was extended for Hölderian sequences  $(V_i)_i$  by Chen-Lin, see [12]. Also, we can see in [18], an extension of the Brezis-Li-Shafrir to compact Riemann surface without boundary. We can see in [19] explicit form,  $(8\pi m, m \in \mathbb{N}^*$  exactly), for the numbers in front of the Dirac masses, when the solutions blow-up. Here, the notion of isolated blow-up point is used.

In [8], Brezis and Merle proposed the following Problem:

**Problem (Brezis-Merle [8]).** *If  $(u_i)_i$  and  $(V_i)_i$  are two sequences of functions relatively to the previous problem (P) with,*

$$0 \leq V_i \rightarrow V \text{ in } C^0(\bar{\Omega}).$$

$$\int_{\Omega} e^{u_i} dy \leq C,$$

*Is it possible to prove that:*

$$\sup_{\Omega} u_i \leq c = c(C, V, \Omega) ?$$

Here, we assume more regularity on  $V_i$ , we suppose that  $V_i \geq 0$  is  $C^s$  ( $s$ -holderian)  $1/2 < s \leq 1$ ) and when we have a boundary singularity. We give the answer where  $bC < 16\pi$  for an equation with boundary singularity.

The following first theorem can be proved with the assumption  $\beta \in (0, +\infty)$  but we prove it in the case  $\beta \in (0, 1)$ .

Our main results are:

**Theorem 1.1.** *Assume  $\Omega = B_1(0)$ ,  $x_0 \in \partial\Omega$ ,  $\beta \in (0, 1)$ , and,*

$$\int_{B_1(0)} |x - x_0|^{2\beta} V_i e^{u_i} dy \leq 16\pi - \epsilon, \quad \epsilon > 0,$$

$$u_i(x_i) = \sup_{B_1(0)} u_i \rightarrow +\infty.$$

There is a sequences  $(x_i^0)_i, (\delta_i^0)$ , such that:

$$(x_i^0)_i \equiv (x_i)_i, \delta_i^0 = \delta_i = d(x_i, \partial B_1(0)) \rightarrow 0,$$

and,

$$u_i(x_i) = \sup_{B_1(0)} u_i \rightarrow +\infty,$$

$$u_i(x_i) + 2 \log \delta_i + 2\beta \log d(x_i, x_0) \rightarrow +\infty,$$

$$\forall \epsilon > 0, \limsup_{i \rightarrow +\infty} \int_{B(x_i, \delta_i \epsilon)} |x - x_0|^{+2\beta} V_i e^{u_i} dy \geq 4\pi > 0.$$

If we assume:

$$V_i \rightarrow V \text{ in } C^0(\bar{B}_1(0)),$$

then,

$$\forall \epsilon > 0, \sup_{B_1(0) - B(x_i, \delta_i \epsilon)} u_i \leq C_\epsilon$$

$$\forall \epsilon > 0, \limsup_{i \rightarrow +\infty} \int_{B(x_i, \delta_i \epsilon)} |x - x_0|^{+2\beta} V_i e^{u_i} dy = 8\pi.$$

And, thus, we have the following convergence in the sense of distributions:

$$\int_{B_1(0)} |x - x_0|^{2\beta} V_i e^{u_i} dy \rightarrow \int_{B_1(0)} |x - x_0|^{2\beta} V e^u dy + 8\pi \delta_{x_0^1}.$$

**Theorem 1.2.** Assume that,  $V_i$  is uniformly  $s$ -holderian with  $1/2 + \beta/2 < s \leq 1$ ,  $x_0 \in \partial\Omega$ ,  $\beta \in (0, 1)$ , and,

$$\int_{B_1(0)} |x - x_0|^{2\beta} V_i e^{u_i} dy \leq 16\pi - \epsilon, \quad \epsilon > 0,$$

then we have:

$$\sup_{\Omega} u_i \leq c = c(b, C, A, s, \beta, \epsilon, x_0, \Omega).$$

where  $A$  is the h\"olderian constant of  $V_i$ .

## 2. PROOFS OF THE RESULTS

### Proofs of the theorems:

Without loss of generality, we can assume that  $\Omega = B_1(0)$  the unit ball centered on the origin.

Here,  $G$  is the Green function of the Laplacian with Dirichlet condition on  $B_1(0)$ . We have (in complex notation):

$$G(x, y) = \frac{1}{2\pi} \log \frac{|1 - \bar{x}y|}{|x - y|},$$

Since  $u_i \in W_0^{1,1}(\Omega)$  and  $\beta \in (0, 1)$ , we have by the Brezis-Merle result and the elliptic estimates, (see [1]):

$$u_i \in C^2(\Omega) \cap W^{2,p}(\Omega) \cap C^{1,\epsilon}(\bar{\Omega})$$

for all  $2 < p < +\infty$ .

Set,

$$v_i(x) = \int_{B_1(0)} G(x, y) V_i(y) |x - x_0|^{2\beta} e^{u_i(y)} dy.$$

We decompose  $v_i$  in two terms (Newtonian potential):

$$v_i^1(x) = \int_{B_1(0)} -\frac{1}{2\pi} \log |x - y| V_i(y) |x - x_0|^{2\beta} e^{u_i(y)} dy,$$

and,

$$v_i^2(x) = \int_{B_1(0)} \frac{1}{2\pi} \log |1 - \bar{x}y| V_i(y) |x - x_0|^{2\beta} e^{u_i(y)} dy,$$

According to the proof in the book of Gilbarg-Trudinger see [15],  $v_i^1, v_i^2$  and thus  $v_i$  are  $C^1(\bar{\Omega})$ . Indeed, we use the same proof as in [15] (Chapter 4, Newtonian potential).

We use this fact and the maximum principle to have  $v_i = u_i$ .

Also, we can use integration by part (the Green representation formula, see its proof in the first chapter of [15]) to have in  $\Omega$  (and not  $\bar{\Omega}$ ):

$$u_i(x) = - \int_{B_1(0)} G(x, y) \Delta u_i(y) dy = \int_{B_1(0)} G(x, y) V_i(y) |x - x_0|^{2\beta} e^{u_i(y)} dy.$$

We write,

$$\begin{aligned} u_i(x_i) &= \int_{\Omega} G(x_i, y) |x - x_0|^{2\beta} V_i(y) e^{u_i(y)} dx = \int_{\Omega - B(x_i, \delta_i/2)} G(x_i, y) |x - x_0|^{2\beta} V_i e^{u_i(y)} dy + \\ &\quad + \int_{B(x_i, \delta_i/2)} G(x_i, y) |x - x_0|^{2\beta} V_i e^{u_i(y)} dy \end{aligned}$$

According to the maximum principle, the harmonic function  $G(x_i, \cdot)$  on  $\Omega - B(x_i, \delta_i/2)$  take its maximum on the boundary of  $B(x_i, \delta_i/2)$ , we can compute this maximum:

$$G(x_i, y_i) = \frac{1}{2\pi} \log \frac{|1 - \bar{x}_i y_i|}{|x_i - y_i|} = \frac{1}{2\pi} \log \frac{|1 - \bar{x}_i(x_i + \delta_i \theta_i)|}{|\delta_i/2|} = \frac{1}{2\pi} \log(2|(1 + |x_i|) + \theta_i|) < +\infty$$

with  $|\theta_i| = 1/2$ .

Thus,

$$u_i(x_i) \leq C + \int_{B(x_i, \delta_i/2)} G(x_i, y) |x - x_0|^{2\beta} V_i e^{u_i(y)} dy \leq C + e^{u_i(x_i) + 2\beta \log d(x_i, x_0)} \int_{B(x_i, \delta_i/2)} G(x_i, y) dy$$

Now, we compute  $\int_{B(x_i, \delta_i/2)} G(x_i, y) dy$   
we set in polar coordinates,

$$y = x_i + \delta_i t \theta$$

we find:

$$\begin{aligned} \int_{B(x_i, \delta_i/2)} G(x_i, y) dy &= \int_{B(x_i, \delta_i/2)} \frac{1}{2\pi} \log \frac{|1 - \bar{x}_i y|}{|x_i - y|} = \frac{1}{2\pi} \int_0^{2\pi} \int_0^{1/2} \delta_i^2 \log \frac{|1 - \bar{x}_i(x_i + \delta_i t \theta)|}{\delta_i t} t dt d\theta = \\ &= \frac{1}{2\pi} \int_0^{2\pi} \int_0^{1/2} \delta_i^2 (\log(|1 + |x_i| + t\theta|) - \log t) t dt d\theta \leq C \delta_i^2. \end{aligned}$$

Thus,

$$u_i(x_i) \leq C + C \delta_i^2 e^{u_i(x_i) + 2\beta \log d(x_i, x_0)},$$

which we can write, because  $u_i(x_i) \rightarrow +\infty$ ,

$$u_i(x_i) \leq C' \delta_i^2 e^{u_i(x_i) + 2\beta \log d(x_i, x_0)},$$

We can conclude that:

$$u_i(x_i) + 2 \log \delta_i + 2\beta \log d(x_i, x_0) \rightarrow +\infty.$$

Since in  $B(x_i, \delta_i \epsilon)$ ,  $d(x, x_0)$  is equivalent to  $d(x_i, x_0)$  we can consider the following functions:

$$v_i(y) = u_i(x_i + \delta_i y) + 2 \log \delta_i + 2\beta \log d(x_i, x_0), \quad y \in B(0, 1/2)$$

The function satisfies all conditions of the Brezis-Merle hypothesis, we can conclude that, on each compact set:

$$v_i \rightarrow -\infty$$

we can assume, without loss of generality that for  $1/2 > \epsilon > 0$ , we have:

$$v_i \rightarrow -\infty, \quad y \in B(0, 2\epsilon) - B(0, \epsilon),$$

**Lemma 2.1.** For all  $1/4 > \epsilon > 0$ , we have:

$$\sup_{B(x_i, (3/2)\delta_i \epsilon) - B(x_i, \delta_i \epsilon)} u_i \leq C_\epsilon.$$

Proof of the lemma

Let  $t'_i$  and  $t_i$  the points of  $B(x_i, 2\delta_i \epsilon) - B(x_i, (1/2)\delta_i \epsilon)$  and  $B(x_i, (3/2)\delta_i \epsilon) - B(x_i, \delta_i \epsilon)$  respectively where  $u_i$  takes its maximum.

According to the Brezis-Merle work, we have:

$$u_i(t'_i) + 2 \log \delta_i + 2\beta \log d(x_i, x_0) \rightarrow -\infty$$

We write,

$$\begin{aligned} u_i(t_i) &= \int_{\Omega} G(t_i, y) |x - x_0|^{2\beta} V_i(y) e^{u_i(y)} dx = \int_{\Omega - B(x_i, 2\delta_i \epsilon)} G(t_i, y) |x - x_0|^{2\beta} V_i e^{u_i(y)} dy + \\ &+ \int_{B(x_i, 2\delta_i \epsilon) - B(x_i, (1/2)\delta_i \epsilon)} G(t_i, y) |x - x_0|^{2\beta} V_i e^{u_i(y)} dy + \\ &+ \int_{B(x_i, (1/2)\delta_i \epsilon)} G(t_i, y) |x - x_0|^{2\beta} V_i e^{u_i(y)} dy \end{aligned}$$

But, in the first and the third integrale, the point  $t_i$  is far from the singularity  $x_i$  and we know that the Green function is bounded. For the second integrale, after a change of variable, we can see that this integrale is bounded by (we take the supremum in the annulus and use Brezis-Merle theorem)

$$\delta_i^2 e^{u_i(t'_i) + 2\beta \log d(x_i, x_0)} \times I_j$$

where  $I_j$  is a Jensen integrale (of the form  $\int_0^1 \int_0^{2\pi} (\log(|1 + |x_i| + t\theta|) - \log|\theta_i - t\theta|) t dt d\theta$  which is bounded).

we conclude the lemma.

From the lemma, we see that far from the singularity the sequence is bounded, thus if we take the supremum on the set  $B_1(0) - B(x_i, \delta_i \epsilon)$  we can see that this supremum is bounded and thus the sequence of functions is uniformly bounded or tends to infinity and we use the same arguments as for  $x_i$  to conclude that around this point and far from the singularity, the sequence is bounded.

The process will be finished, because, according to Brezis-Merle estimate, around each supremum constructed and tending to infinity, we have:

$$\forall \epsilon > 0, \quad \limsup_{i \rightarrow +\infty} \int_{B(x_i, \delta_i \epsilon)} |x - x_0|^{2\beta} V_i e^{u_i} dy \geq 4\pi > 0.$$

Finally, with this construction, we have a finite number of "exterior" blow-up points and outside the singularities the sequence is bounded uniformly, for example, in the case of one "exterior" blow-up point, we have:

$$u_i(x_i) \rightarrow +\infty$$

$$\forall \epsilon > 0, \quad \sup_{B_1(0) - B(x_i, \delta_i \epsilon)} u_i \leq C_\epsilon$$

$$\forall \epsilon > 0, \quad \limsup_{i \rightarrow +\infty} \int_{B(x_i, \delta_i \epsilon)} |x - x_0|^{2\beta} V_i e^{u_i} dy \geq 4\pi > 0.$$

$$x_i \rightarrow x_0 \in \partial B_1(0).$$

**Remark:** For the general case, the process of quantization can be extended to more than one blow-up points.

We have the following lemma:

**Lemma 2.2.** *Each  $\delta_i^k$  is of order  $d(x_i^k, \partial B_1(0))$ . Namely: there is a positive constant  $C > 0$  such that for  $\epsilon > 0$  small enough:*

$$\delta_i^k \leq d(x_i^k, \partial B_1(0)) \leq (2 + \frac{C}{\epsilon}) \delta_i^k.$$

Proof of the lemma

Now, if we suppose that there is another "exterior" blow-up  $(t_i)_i$ , we have, because  $(u_i)_i$  is uniformly bounded in a neighborhood of  $\partial B(x_i, \delta_i \epsilon)$ , we have :

$$d(t_i, \partial B(x_i, \delta_i \epsilon)) \geq \delta_i \epsilon$$

If we set,

$$\delta_i' = d(t_i, \partial(B_1(0) - B(x_i, \delta_i \epsilon))) = \inf\{d(t_i, \partial B(x_i, \delta_i \epsilon)), d(t_i, \partial(B_1(0)))\}$$

then,  $\delta_i'$  is of order  $d(t_i, \partial B_1(0))$ . To see this, we write:

$$d(t_i, \partial B_1(0)) \leq d(t_i, \partial B(x_i, \delta_i \epsilon)) + d(\partial B(x_i, \delta_i \epsilon), x_i) + d(x_i, \partial B_1(0)),$$

Thus,

$$\frac{d(t_i, \partial B_1(0))}{d(t_i, \partial B(x_i, \delta_i \epsilon))} \leq 2 + \frac{1}{\epsilon},$$

Thus,

$$\delta_i' \leq d(t_i, \partial B_1(0)) \leq \delta_i' (2 + \frac{1}{\epsilon}).$$

Now, the general case follow by induction. We use the same argument for three, four, ...,  $n$  blow-up points.

We have, by induction and, here we use the fact that  $u_i$  is uniformly bounded outside a small ball centered at  $x_i^j, j = 0, \dots, k-1$ :

$$\delta_i^j \leq d(x_i^j, \partial B_1(0)) \leq C_1 \delta_i^j, \quad j = 0, \dots, k-1,$$

.

$$d(x_i^k, \partial B(x_i^j, \delta_i^j \epsilon / 2)) \geq \epsilon \delta_i^j, \quad \epsilon > 0, \quad j = 0, \dots, k-1,$$

.

and let's consider  $x_i^k$  such that:

$$u_i(x_i^k) = \sup_{B_1(0) - \cup_{j=0}^{k-1} B(x_i^j, \delta_i^j \epsilon)} u_i \rightarrow +\infty,$$

take,

$$\delta_i^k = \inf\{d(x_i^k, \partial B_1(0)), d(x_i^k, \partial(B_1(0) - \cup_{j=0}^{k-1} B(x_i^j, \delta_i^j \epsilon / 2)))\},$$

if, we have,

$$\delta_i^k = d(x_i^k, \partial B(x_i^j, \delta_i^j \epsilon / 2)), \quad j \in \{0, \dots, k-1\}.$$

Then,

$$\begin{aligned} \delta_i^k &\leq d(x_i^k, \partial B_1(0)) \leq \\ &\leq d(x_i^k, \partial B(x_i^j, \delta_i^j \epsilon / 2)) + d(\partial B(x_i^j, \delta_i^j \epsilon / 2), x_i^j) + d(x_i^j, \partial B_1(0)) \\ &\leq (2 + \frac{C_1}{\epsilon}) \delta_i^k. \end{aligned}$$

To apply lemma 2.1 for  $m$  blow-up points, we use an induction:

We do directly the same approach for  $t_i$  as  $x_i$  by using directly the Green function of the unit ball.

If we look to the blow-up points, we can see, with this work that, after finite steps, the sequence will be bounded outside a finite number of balls, because of Brezis-Merle estimate:

$$\forall \epsilon > 0, \quad \limsup_{i \rightarrow +\infty} \int_{B(x_i^k, \delta_i^k \epsilon)} |x - x_0|^{2\beta} V_i e^{u_i} dy \geq 4\pi > 0.$$

Here, we can take the functions:

$$u_i^k(y) = u_i(x_i^k + \delta_i^k y) + 2 \log \delta_i^k + 2\beta \log d(x_i, x_0).$$

Indeed, by corollary 4 of the paper of Brezis-Merle, if we have:

$$\limsup_{i \rightarrow +\infty} \int_{B(x_i^k, \delta_i^k \epsilon)} |x - x_0|^{2\beta} V_i e^{u_i} dy \leq 4\pi - \epsilon_0 < 4\pi,$$

then,  $(u_i^k)^+$  would be locally bounded and this contradict the fact that  $u_i^k(0) \rightarrow +\infty$ .

Finally, we can say that, there is a finite number of sequences  $(x_i^k)_i, (\delta_i^k), 0 \leq k \leq m$ , such that:

$$(x_i^0)_i \equiv (x_i)_i, \quad \delta_i^0 = \delta_i = d(x_i, \partial B_1(0)),$$

$$(x_i^1)_i \equiv (t_i)_i, \quad \delta_i^1 = \delta_i' = d(t_i, \partial(B_1(0) - B(x_i, \delta_i \epsilon))),$$

and each  $\delta_i^k$  is of order  $d(x_i^k, \partial B_1(0))$ .

and,

$$u_i(x_i^k) = \sup_{B_1(0) - \cup_{j=0}^{k-1} B(x_i^j, \delta_i^j \epsilon)} u_i \rightarrow +\infty,$$

$$u_i(x_i^k) + 2 \log \delta_i^k + 2\beta \log d(x_i^k, x_0) \rightarrow +\infty,$$

$$\forall \epsilon > 0, \quad \sup_{B_1(0) - \cup_{j=0}^m B(x_i^j, \delta_i^j \epsilon)} u_i \leq C_\epsilon$$

$$\forall \epsilon > 0, \quad \limsup_{i \rightarrow +\infty} \int_{B(x_i^k, \delta_i^k \epsilon)} |x - x_0|^{2\beta} V_i e^{u_i} dy \geq 4\pi > 0.$$

#### The work of YY.Li-I.Shafrir

Since in  $B(x_i, \delta_i \epsilon)$ ,  $d(x, x_0)$  is equivalent to  $d(x_i, x_0)$  we can consider the following functions:

$$v_i(y) = u_i(x_i + \delta_i y) + 2 \log \delta_i + 2\beta \log d(x_i, x_0).$$

With the previous method, we have a finite number of "exterior" blow-up points (perhaps the same) and the sequences tend to the boundary. With the aid of proposition 1 of the paper of Li-Shafrir, we see that around each exterior blow-up, we have a finite number of "interior"



blow-ups. Around, each exterior blow-up, we have after rescaling with  $\delta_i^k$ , the same situation as around a fixed ball with positive radius. If we assume:

$$V_i \rightarrow V \text{ in } C^0(\bar{B}_1(0)),$$

then,

$$\forall \epsilon > 0, \limsup_{i \rightarrow +\infty} \int_{B(x_i^k, \delta_i^k \epsilon)} |x - x_0|^{2\beta} V_i e^{u_i} dy = 8\pi m_k, \quad m_k \in \mathbb{N}^*.$$

And, thus, we have the following convergence in the sense of distributions:

$$\int_{B_1(0)} |x - x_0|^{2\beta} V_i e^{u_i} dy \rightarrow \int_{B_1(0)} |x - x_0|^{2\beta} V e^u dy + \sum_{k=0}^m 8\pi m'_k \delta_{x_0^k}, \quad m'_k \in \mathbb{N}^*, \quad x_0^k \in \partial B_1(0).$$

Consequence: using a Pohozaev-type identity. proof of theorem 2

By a conformal transformation, we can assume that our domain  $\Omega = B^+$  is a half ball centered at the origin,  $B^+ = \{x, |x| \leq 1, x_1 \geq 0\}$ , and,  $x_0 = 0$ . In this case the normal at the boundary is  $\nu = (-1, 0)$  and  $u_i(0, x_2) \equiv 0$ . Also, we set  $x_i$  the blow-up point and  $x_i^2 = (0, x_i^2)$  and  $x_i^1 = (x_i^1, 0)$  respectively the second and the first part of  $x_i$ . Let  $\partial B^+$  the part of the boundary for which  $u_i$  and its derivatives are uniformly bounded and thus converge to the corresponding function.

The case of one blow-up point:

**Theorem 2.3.** *If  $V_i$  is  $s$ -Holderian with  $1/2 + \beta/2 < s \leq 1$  and,*

$$\int_{\Omega} |x|^{2\beta} V_i e^{u_i} dy \leq 16\pi - \epsilon, \quad \epsilon > 0,$$

we have :

$$2V_i(x_i) \int_{B(x_i, \delta_i \epsilon)} |x|^{2\beta} e^{u_i} dy = o(1),$$

which means that there are no blow-up points.

Proof of the theorem

In order to use the Pohozaev identity we need to have a good function  $u_i$ , since  $\beta \in (0, 1)$ , we have a function  $u_i$  such that:

$$u_i \in C^{2, \epsilon}(\bar{\Omega})$$

Thus, we can use integration by parts. The Pohozaev identity gives us the following formula:

$$\int_{\Omega} \langle (x - x_2^i) | \nabla u_i \rangle (-\Delta u_i) dy = \int_{\Omega} \langle (x - x_2^i) | \nabla u_i \rangle |x|^{2\beta} V_i e^{u_i} dy = A_i$$

$$A_i = \int_{\partial B^+} \langle (x - x_2^i) | \nabla u_i \rangle \langle \nu | \nabla u_i \rangle d\sigma + \int_{\partial B^+} \langle (x - x_2^i) | \nu \rangle |\nabla u_i|^2 d\sigma$$

We can write it as:

$$\begin{aligned} \int_{\Omega} \langle (x - x_2^i) | \nabla u_i \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy &= A_i + V_i(x_i) \int_{\Omega} \langle (x - x_2^i) | \nabla u_i \rangle |x|^{2\beta} e^{u_i} dy = \\ &= A_i + V_i(x_i) \int_{\Omega} \langle (x - x_2^i) | x \rangle |x|^{2\beta} |\nabla(e^{u_i})| dy \end{aligned}$$

And, if we integrate by part the second term, we have (because  $x_1 = 0$  on the boundary and  $\nu_2 = 0$ ):

$$\int_{\Omega} \langle (x - x_2^i) | \nabla u_i \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy = -2(1 + \beta) V_i(x_i) \int_{\Omega} |x|^{2\beta} e^{u_i} dy +$$

$$+ 2\beta V_i(x_i) \int_{B(x_i, \delta_i \epsilon)} x_2 x_2^i |x|^{2\beta-2} e^{u_i} dy - 2\beta V_i(x_i) \int_{\Omega - B(x_i, \delta_i \epsilon)} x_2 x_2^i |x|^{2\beta-2} e^{u_i} dy + B_i$$

where  $B_i$  is,

$$B_i = V_i(x_i) \int_{\partial B^+} \langle (x - x_2^i) | \nu \rangle |x|^{2\beta} e^{u_i} dy$$

applying the same procedure to  $u$ , we can write:

$$\int_{\Omega} \langle (x - x_2^i) | \nabla u \rangle (V - V(0)) |x|^{2\beta} e^u dy = -2(1 + \beta) V(0) \int_{\Omega} |x|^{2\beta} e^u dy +$$

$$- 2\beta V(0) \int_{B(x_i, \delta_i \epsilon)} x_2 x_2^i |x|^{2\beta-2} e^u dy - 2\beta V(0) \int_{\Omega - B(x_i, \delta_i \epsilon)} x_2 x_2^i |x|^{2\beta-2} e^u dy + B,$$

with,

$$B = V(0) \int_{\partial B^+} \langle (x - x_2^i) | \nu \rangle |x|^{2\beta} e^u dy$$

we use the fact that,  $u_i$  is bounded outside  $B(x_i, \delta_i \epsilon)$  and the convergence of  $u_i$  to  $u$  on compact set of  $\bar{\Omega} - \{0\}$ , and the fact that  $\beta \in (0, +\infty)$ , to write the following:

$$2V_i(x_i) \int_{B(x_i, \delta_i \epsilon)} |x|^{2\beta} e^{u_i} dy + o(1) =$$

$$= \int_{\Omega} \langle (x - x_2^i) | \nabla u_i \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy - \int_{\Omega} \langle (x - x_2^i) | \nabla u \rangle (V - V(0)) |x|^{2\beta} e^u dy +$$

$$+ (A_i - A) + (B_i - B),$$

where  $A$  and  $B$  are,

$$A = \int_{\partial B^+} \langle (x - x_2^i) | \nabla u \rangle \langle \nu | \nabla u \rangle d\sigma + \int_{\partial B^+} \langle (x - x_2^i) | \nu \rangle |\nabla u|^2 d\sigma$$

$$B = V(0) \int_{\partial B^+} \langle (x - x_2^i) | \nu \rangle |x|^{2\beta} e^u dy$$

and, because of the uniform convergence of  $u_i$  and its derivatives on  $\partial B^+$ , we have:

$$A_i - A = o(1) \quad \text{and} \quad B_i - B = o(1)$$

which we can write as:

$$2V_i(x_i) \int_{B(x_i, \delta_i \epsilon)} |x|^{2\beta} e^{u_i} dy + o(1) =$$

$$= \int_{\Omega} \langle (x - x_2^i) | \nabla (u_i - u) \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy +$$

$$+ \int_{\Omega} \langle (x - x_2^i) | \nabla u \rangle (V_i - V_i(x_i)) |x|^{2\beta} (e^{u_i} - e^u) dy +$$

$$+ \int_{\Omega} \langle (x - x_2^i) | \nabla u \rangle (V_i - V_i(x_i) - (V - V(0))) |x|^{2\beta} e^u dy + o(1)$$

We can write the second term as:

$$\int_{\Omega} \langle (x - x_2^i) | \nabla u \rangle (V_i - V_i(x_i)) |x|^{2\beta} (e^{u_i} - e^u) dy = \int_{\Omega - B(0, \epsilon)} \langle (x - x_2^i) | \nabla u \rangle (V_i - V_i(x_i)) (e^{u_i} - e^u) |x|^{2\beta} dy +$$

$$+ \int_{B(0, \epsilon)} \langle (x - x_2^i) | \nabla u \rangle (V_i - V_i(x_i)) (e^{u_i} - e^u) |x|^{2\beta} dy = o(1),$$

because of the uniform convergence of  $u_i$  to  $u$  outside a region which contain the blow-up and the uniform convergence of  $V_i$ . For the third integral we have the same result:

$$\int_{\Omega} \langle (x - x_2^i) | \nabla u \rangle (V_i - V_i(x_i) - (V - V(0))) |x|^{2\beta} e^u dy = o(1),$$

because of the uniform convergence of  $V_i$  to  $V$ .

Now, we look to the first integral:

$$\int_{\Omega} \langle (x - x_2^i) | \nabla(u_i - u) \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy,$$

we can write it as:

$$\begin{aligned} \int_{\Omega} \langle (x - x_2^i) | \nabla(u_i - u) \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy &= \int_{\Omega} \langle (x - x_i) | \nabla(u_i - u) \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy + \\ &+ \int_{\Omega} \langle x_1^i | \nabla(u_i - u) \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy, \end{aligned}$$

Thus, we have proved by using the Pohozaev identity the following equality:

$$\begin{aligned} &\int_{\Omega} \langle (x - x_i) | \nabla(u_i - u) \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy + \\ &+ \int_{\Omega} \langle x_1^i | \nabla(u_i - u) \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy = \\ &= 2V_i(x_i) \int_{B(x_i, \delta_i \epsilon)} |x|^{2\beta} e^{u_i} dy + o(1) \end{aligned}$$

We can see, because of the uniform boundedness of  $u_i$  outside  $B(x_i, \delta_i \epsilon)$  and the fact that :

$$\|\nabla(u_i - u)\|_1 = o(1),$$

it is sufficient to look to the integral on  $B(x_i, \delta_i \epsilon)$ .

Assume that we are in the case of one blow-up, it must be  $(x_i)$  and isolated, we can write the following inequality as a consequence of YY.Li-I.Shafrir result:

$$u_i(x) + 2 \log |x - x_i| + 2\beta \log d(x, 0) \leq C,$$

We use this fact and the fact that  $V_i$  is s-holderian to have that, on  $B(x_i, \delta_i \epsilon)$ ,

$$|(x - x_i)(V_i - V_i(x_i))|x|^{2\beta} e^{u_i} \leq \frac{C}{|x - x_i|^{1-s}} \in L^{(2-\epsilon')/(1-s)}, \quad \forall \epsilon' > 0,$$

and, we use the fact that:

$$\|\nabla(u_i - u)\|_q = o(1), \quad \forall 1 \leq q < 2$$

to conclude by the Hölder inequality that:

$$\int_{B(x_i, \delta_i \epsilon)} \langle (x - x_i) | \nabla(u_i - u) \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy = o(1),$$

For the other integral, namely:

$$\int_{B(x_i, \delta_i \epsilon)} \langle x_1^i | \nabla(u_i - u) \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy,$$

We use the fact that, because our domain is a half ball, and the sup + inf inequality to have:

$$x_1^i = \delta_i,$$

$$u_i(x) + 4 \log \delta_i + 4\beta \log d(x, 0) \leq C$$

and,

$$|x|^{s\beta} e^{(s/2)u_i(x)} \leq |x - x_i|^{-s},$$

$$|V_i - V_i(x_i)| \leq |x - x_i|^s,$$

Finally, we have:

$$\begin{aligned} & \left| \int_{B(x_i, \delta_i \epsilon)} \langle x_1^i |\nabla(u_i - u) \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy \right| \leq \\ & \leq C \int_{B(x_i, \delta_i \epsilon)} |\nabla(u_i - u)| |x|^{2\beta(1/2-s/2)} (e^{u_i})^{(3/4-s/2)}, \end{aligned}$$

But in the second member, for  $1/2 < s \leq 1$ , we have  $q_s = 1/(3/4 - s/2) > 2$  and thus  $q'_s < 2$  and,

$$(|x|^{2\beta} e^{u_i})^{3/4-s/2} \in L^{q_s}$$

$$|x|^{2\beta(-1/4)} \in L^r, \quad r = \frac{4}{\beta + \epsilon}, \quad \epsilon > 0,$$

with,

$$\frac{1}{r} + \frac{1}{q_s} = \frac{1}{t} < 1/2,$$

and,

$$\|\nabla(u_i - u)\|_{t'} = o(1),$$

This is possible if,

$$\beta < 2(s - 1/2).$$

or,

$$s > \frac{\beta}{2} + \frac{1}{2},$$

This is possible for  $\beta \in (0, 1)$ .

one conclude that:

$$\int_{B(x_i, \delta_i \epsilon)} \langle x_1^i |\nabla(u_i - u) \rangle (V_i - V_i(x_i)) |x|^{2\beta} e^{u_i} dy = o(1)$$

Finally, with this method, we conclude that, in the case of one blow-up point and  $V_i$  is s-Holderian with  $1/2 < s \leq 1$  :

$$2V_i(x_i) \int_{B(x_i, \delta_i \epsilon)} |x|^{2\beta} e^{u_i} dy = o(1)$$

which means that there is no blow-up, which is a contradiction.

Finally, for one blow-up point and  $V_i$  is s-Hölderian with  $1/2 < s \leq 1$ , the sequence  $(u_i)$  is uniformly bounded on  $\Omega$ .

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