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## Collaboration



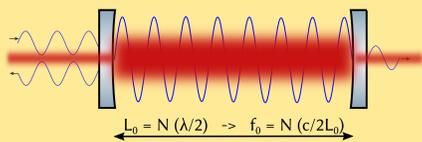
# Detection Techniques and Fundamental Limits in Frequency Stabilization via Spectral Hole Burning

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## Ultra-Stable Lasers The State of the Art

In order for optical clocks to reach their ultimate performance (quantum projection limit,  $\sim 10^{-17}$  fractional frequency stability at 1 s), they require an ultra-stabilized probe laser, setting a **goal of  $10^{-18}$  at 1 s**.



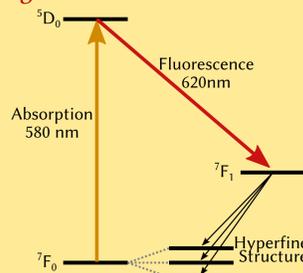
Traditionally, these lasers use length-stable optical cavities as a frequency reference.

Conventional cavities have a fundamental limit, due to thermal noise, at  $\sim 10^{-16}$  at 1 s. Efforts to mitigate thermal noise include alternate cavity geometries, use of high-Q materials, and cryogenic cavities.

- PTB, JILA<sup>[1]</sup>: Cryogenic Cavities:  $4 \times 10^{-17}$  at 1 s
- PTB<sup>[2]</sup>: Long Cavities:  $8 \times 10^{-17}$  at 1 s

While successful, this route faces considerable technological challenges.

- [1] D. G. Matei, et al., *Phys. Rev. Lett.* **118**, 263202 (2017).  
[2] S. Häfner, et al., *Opt. Lett.* **40**, 2112-2115 (2015).



We are researching a potential **paradigm shift in ultra-stable lasers: stabilization to an optical transition of a rare earth ion, Europium ( $\text{Eu}^{3+}$ ), doped into an Yttrium Orthosilicate ( $\text{Y}_2\text{SiO}_5$ )**.

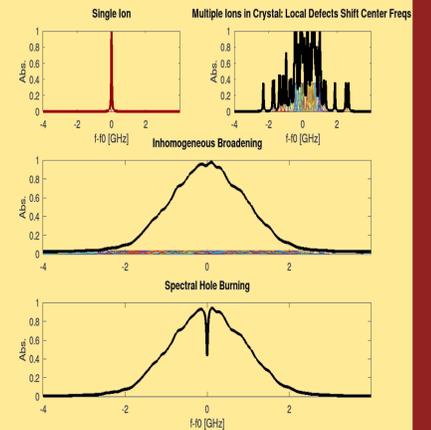
## Spectral Hole Burning (SHB)

The  ${}^7\text{F}_0 \rightarrow {}^5\text{D}_0$  transition in  $\text{Eu}^{3+}$  has a narrow homogeneous absorption line of order  $\sim 100$  Hz.

Within the crystal, inhomogeneous constraints lead to a broadening of the absorption spectrum.

A narrow linewidth pump beam at 580 nm excites the ions, which drop into a hyperfine state.

Pumping-out all the ions at one frequency produces a narrow transmissive 'spectral hole' in the absorption spectrum.

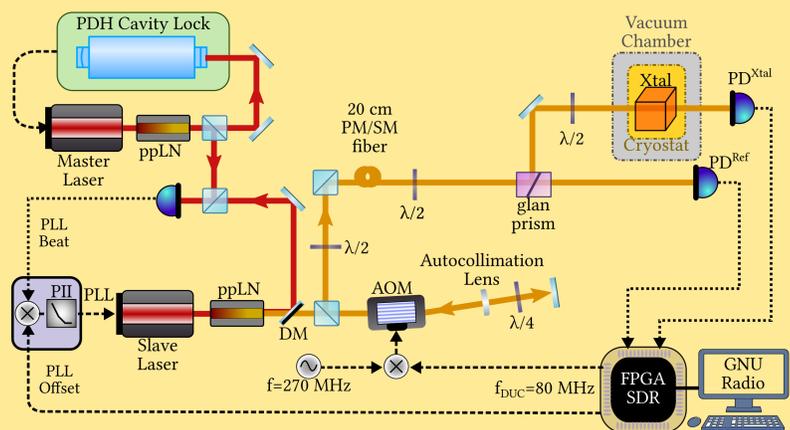


## Experimental Setup

A slave laser is phase-locked to a cavity-referenced pre-stabilized master laser at 1160 nm. The slave is frequency doubled and tuned to the absorption band of the  $\text{Eu}^{3+}:\text{Y}_2\text{SiO}_5$  crystal. The desired spectrum of beat notes is generated on the slave by an AOM driven by a GNU Radio controlled Ettus X310 FPGA-based SDR.

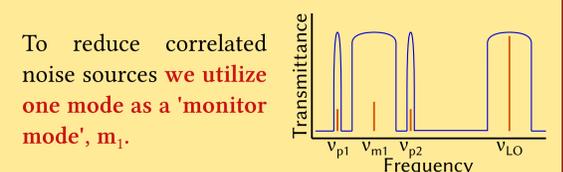
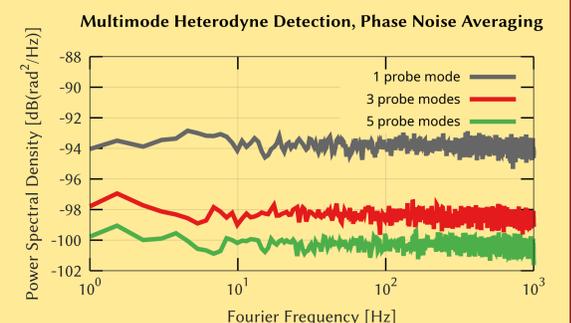
At high power (1 uW) the slave laser acts as the pump, burning the spectral holes. Switching to low power (10 nW), it probes the holes, experiencing a phase retardation proportional to the frequency difference between the laser modes and spectral holes.

This phase difference,  $\phi_{p_i}(t) = \phi_{p_i}^{\text{Xtal}}(t) - \phi_{p_i}^{\text{Ref}}(t)$ , is measured in the FPGA, comparing the transmitted beat note on  $\text{PD}^{\text{Xtal}}$  with a reference beat on  $\text{PD}^{\text{Ref}}$ . The FPGA generates the correction signal tuning the PLL offset.



## Multi-mode Heterodyne Detection<sup>[3]</sup>

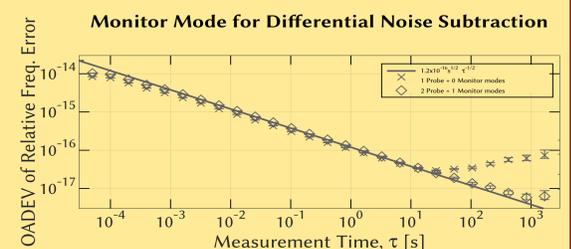
We implemented a novel multi-mode detection scheme to reduce sensing noise. In this scheme we probe multiple narrow spectral holes simultaneously using corresponding laser probe modes,  $p_i$ :  $\bar{\phi}_{p_i}(t) = \frac{1}{N} \sum_{i=1}^N \phi_{p_i}(t)$   
Below: We average over the modes, reducing uncorrelated noise.



To reduce correlated noise sources we utilize one mode as a 'monitor mode',  $m_1$ .

Above: Propagating on the same beam, the monitor mode passes through a broad 'square' spectral hole, accumulating common-mode phase noise from sources such as parasitic interferometers.

Below: Subtracting path-length noise in real-time has yielded high common-mode noise rejection into the thousands of seconds of integration time.

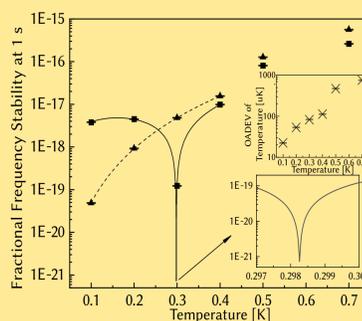


[3] X. Lin, *Opt. Express*, **31**(23), 38475-38493 (2023).

## Spectral Hole Line Shifts at $< 1\text{K}$ <sup>[5]</sup>

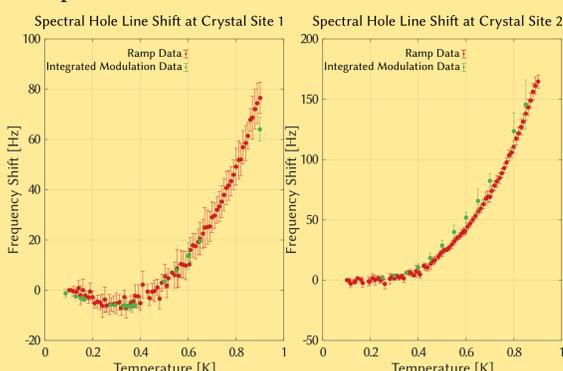
Early measurements in our experiment<sup>[4]</sup> found frequency noise that is compatible with spectral hole (SH) temperature-induced line shifts at  $T \sim 4$  K combined with the temperature stability of our cryostat. This motivated the addition of a custom dilution refrigerator (MyCryoFirm) where, at temperatures  $< 1$  K, we expected a **4 order-of-magnitude reduction in temperature sensitivity**. With the laser frequency locked to the SH, we measured the SH temperature sensitivity in two ways<sup>[5]</sup>: 1) we ramp the temperature and measure the frequency shift of the SH-locked laser against the SYRTE frequency comb, the temperature sensitivity is the derivative of this shift. 2) we modulate the temperature while recording the laser frequency against the comb. Demodulation of these signals reveals the temperature sensitivity directly, the integration of this signal indicates the temperature shift.

Currently, we are collaborating with LENA (U. Braunschweig) to measure the mechanical loss angle of  $\text{Y}_2\text{SiO}_5$ . This, combined with our temperature shift measurements, will allow us to develop a suite of thermal noise models to understand the fundamental stability limit.



Left: The measured frequency shift in a spectral hole at dilution temperatures. At spectroscopic site 2, the line shift roughly follows the expected  $T^4$  dependence, however at site 1 a large deviation from this trend reveals a local minimum around 300 mK **providing a region of temperature insensitivity to first order**. Above: Measured cryostat temperature stability and SH temperature sensitivities reveals a region with temperature induced fractional frequency instabilities  $< 10^{-20}$ .

- [4] N. Galland, et al., *Opt. Lett.*, **45**, 1930-1933 (2020).  
[5] Manuscript in Preparation.



## Thanks To

