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Multifractal analysis for spatial characterization of high resolution Sentinel-2/MAJA products in Southwestern France

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5 Highlights

- 6 Particular scale breaks were detected on Sentinel-2/MAJA products
- 7 Time variations in scaling regimes were noticed in series of surface reflectances
- 8 The large-scale break was related to the landscape spatial structure
- The fine-scale break was explained by the effect of the optical acquisition system

10 Abstract

11 Multifractal analysis was applied to evidence spatial properties, in particular scaling behaviour, of Sentinel-2 surface reflectances and optical indexes acquired over the Southwest region of France. 12 13 Considering the significant impact of clouds on the spatial properties of optical products, no study to our knowledge has focused on the multifractal analysis of a series of images. Images with minimal 14 15 cloud cover were identified for the whole year 2016 based on the MACCS-ATCOR Joint Algorithm (MAJA) cloud and shadow detection process. This allowed us to access the dynamics of the scaling 16 17 properties of surface reflectances, contributing to demonstrate the potential of multifractals for spatial 18 characterization. Different scaling regimes, separated by significant scale breaks, were evident in these 19 data. Through simulations from the Universal Multifractal model, we could explain the scaling 20 properties and especially the position of scale breaks observed in the images studied by two types of 21 phenomena. First, a transition scale located between hundreds of meters and one kilometre was 22 detected on power spectra, structure functions and statistical moments. This specific scale was related 23 to the average size of the agricultural plots of the studied region, revealing the direct link between the 24 scaling behaviour observed and the physical properties of the surface. Second, an instrumental 25 phenomenon was detected on the power spectra, resulting in a systematic transition scale observed 26 near the resolution of the sensor and revealing the impact of the acquisition system on the high 27 frequencies of optical images. In this research work, theoretical and/or empirical validation arguments

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- 28 were provided as to the origin of the scales detected, contributing to demonstrate the value of
- 29 multifractals to extract image texture from surface reflectance products.

30 Keywords

- 31 Surface reflectances; Optical indexes; Spatial characterization; Universal Multifractal; Sentinel-2;
- 32 MAJA
- 33

34 1 Introduction

35 The last generation of space-borne sensors enables the acquisition of Earth surface information with both high temporal and spatial resolutions, showing good capabilities for studying environmental 36 37 and anthropogenic evolutions at local to global spatial scales (Gutman et al., 2008; Huang et al., 2010). Among others, the most popular may be the U.S. Landsat-8 (Roy et al., 2014) and the recent ESA 38 39 Sentinel-2 mission (Martimort et al., 2007) satellites. They provide multispectral images at spatial 40 resolutions of tens of meters with a revisit time ranging between 5 and 16 days. Thanks to the benefits 41 of their thirteen spectral bands, Sentinel-2 satellites (S2) deliver surface reflectances acquired from 42 visible to short-wave infrared, allowing observation of a strong spatial variability in quite a wide range 43 of wavelengths (Sotian et al., 2019).

44 Spatial heterogeneity is often closely linked to the temporal dynamics which itself relies on both 45 natural phenomena associated with specific environmental conditions (meadows, forested areas, bare soils, etc.), and anthropogenic phenomena resulting from urban activities, agricultural practices, etc. 46 47 Because these factors are present at different space scales and may have different spectral signatures, 48 the information contained in acquired surface reflectances can be entangled in a very complex way 49 and difficult to analyse from an image processing stand point (De Cola, 1989; Lam, 1990). Depending 50 on the targeted applications, different techniques have been developed and used so far, to classify, 51 model or characterize the land surface spatial and temporal properties with the aim to account for its 52 complexity. In image classification, commonly used methods are based on the spectral signature of the 53 observed scene (Lu and Weng, 2007), some of them involve complex data processing algorithms such 54 as decision trees or neural networks (Hansen et al., 1996; Pal, 2005; Paola and Schowengerdt, 1995). 55 In object-oriented classifiers, segmentation techniques make use of the spatial context of each pixel 56 to improve the classification performance (Benz et al., 2004; Zucker, 1976). An embedding approach 57 insensitive to the initial conditions was also introduced recently (Mangiarotti et al., 2018). For the 58 modelling of the land surface, most of the developments have been based on optimization and data-59 assimilation techniques with the aim to improve the integration of remote sensed data into land 60 surface process models (Bach and Mauser, 2003; Mangiarotti et al., 2008; Rodriguez-Fernandez et al., 2019). Using the global modelling technique, the first chaotic attractor directly issued from satellite 61 data was obtained that illustrates the high complexity of the surface dynamics at large scale 62 63 (Mangiarotti et al. 2012; 2014) and which itself relies on the level of spatial synchronization 64 (Mangiarotti et al. 2016). To have a better description of the landscape changes, various linear and 65 nonlinear characterization methods were introduced to describe the spatial variability of land surface 66 (Gustafson, 1998; Puissant et al., 2005), some of them taking the spatial scale into account in their 67 analyses (Mangiarotti et al., 2010, 2012). However, a more systematic spatial characterization requires a statistical framework able to describe variability and image texture over a wide range of spatial
scales. Statistical approaches like the Fourier power spectrum and entropy (Weszka et al., 1976) or
variogram and fractal dimension (Ramstein and Raffly, 1989) can give relevant information on the
spatial texture of the image such as roughness, autocorrelation, or the noisy nature of the signal for
example.

73 Fractal analysis can be a relevant tool to quantitatively characterize the high degree of complexity 74 and the spatial structure of surface products (Verrier et al., 2020) such as new remote sensing 75 reflectances similar to S2 products (Alonso et al., 2017). Indeed, S2 products can be seen as images 76 that embed a hierarchy of spatial structures of different sizes and of different energy levels. Many 77 remote sensing images of natural variables show a fractal structure that can be demonstrated by 78 various geometrical and statistical tools (Lovejoy et al., 2001; Qiu et al., 1999; Schmitt et al, 1993; 79 Tessier et al., 1993). Studying such scaling features can contribute to identify specific set of scales on which specific properties apply. Kim and Barros (2002a) used such scaling analysis to characterize 80 space-time variability of remotely sensed soil moisture and relate it to landscape characteristics 81 82 (terrain, soil, vegetation). Two scale ranges were identified, a large one related to the hydrometeorological conditions (evapotranspiration, precipitation processes) and a fine one related 83 84 to the soil properties of the surface (texture, structure). Precisely knowing these sets of scales aims to 85 improve land surface modelling, considering both natural phenomena like topography, streams or 86 drainage basins (Aguado et al., 2014; Cheng et al., 2001; Lyu et al., 2017) and anthropogenic 87 phenomena like urban growth (Chen and Wang, 2013; Tan, 2021) or irrigation (Ko et al., 2016). 88 Specifically, in hydrology these scaling features led to the development of better calibrated 89 downscaling algorithms applied to surface soil moisture (Kim and Barros, 2002b; Ko et al., 2016; 90 Mascaro et al, 2010; Neuhauser et al., 2019) and precipitation (Gires et al., 2012; Rebora et al., 2006).

91 One of the major constraints when measuring surface reflectances from space is the atmosphere. 92 The latter is composed of different elements which have direct impact on reflectances. Indeed, clouds, 93 gas molecules and aerosols are responsible for scattering and absorbing effects on the sun light 94 reflected by the surface (Fraser et al., 1985; Vermote et al., 1997). These atmospheric components 95 must be taken into account when analysing optical images, especially since they can affect the scaling 96 properties of the signal measured from space. For example, several works have shown the fractal (and 97 multifractal) behaviour of clouds acquired from visible and infrared sensors (Lovejoy et al., 2001; Siebesma and Jonker, 2000; Tessier et al., 1993). 98

99 Thus, in order to avoid the impact of clouds, most of the scaling studies of surface reflectances
100 have been applied to images specifically chosen to be cloud-free (Alonso et al., 2017; Lovejoy et al.,

101 2008b). While these works prove the ability of multifractal analysis to investigate the scaling properties 102 of images of surface reflectances, no seasonal properties/differences were really explored yet. 103 Moreover, multifractal models were initially developed for modelling turbulence and atmospheric 104 processes (Parisi and Frisch, 1985; Schertzer and Lovejoy, 1987), leading to many studies on 105 atmospheric dynamics (Lovejoy et al., 2008a; Lovejoy and Schertzer, 2010c) but many fewer on 106 continental surface characteristics.

107 Recently, with the increasing use of spaceborne optical sensors, several algorithms have been 108 developed to take into account the effects of the atmosphere on the reflectances measured by satellite 109 (Hagolle et al., 2010; King et al., 1999; Liang et al, 2001; Rahman and Dedieu, 1994). One of the most 110 recent is the MACCS-ATCOR Joint Algorithm, called MAJA. It relies on the MACCS (Multi-sensor 111 Atmospheric Correction and Cloud Screening) method – dedicated to aerosol optical depth estimation, 112 clouds detection and atmospheric corrections (Hagolle et al, 2015; Lenoble et al., 2007; Rouquié et al., 113 2017) – and includes characteristics from the ATCOR algorithm (Atmospheric and Topographic Correction; Richter et al., 2006) such as the correction of slope effects. This processing chain is used 114 115 within the French THEIA land data center (http://www.theia-land.fr) on multispectral images from 116 Landsat-8, S2 and from the recent Venus satellite (Manivasagam et al., 2019).

117 Considering this context, key scientific questions still need to be addressed in the field of spatial 118 characterization of surface reflectances. The research work described here was realized in order to 119 deepen our understanding of multifractal analysis and, above all, to demonstrate its value for 120 improving spatial characterization of land surfaces. A particular reflection was made on the physical 121 origin of the specific scales identified, relating them to image texture and asking whether the latter 122 were due to:

- The seasonal specificities of the surface?
- The geographic specificities of the surface?
- The signal acquisition system of the satellite sensor?

126 To answer these questions, we investigated the occurrence of scaling properties in surface images 127 collected by S2 sensors and corrected by MAJA. The dataset consists of S2-MAJA maps located in 128 France at different dates within 2016. This research work was first focused on the southwestern region 129 of France, on which a selection of particular dates was made to minimize the impact of clouds. In this 130 region, particular consideration was given to the temporal dynamics of scaling properties due for 131 example to seasonal effects. In addition, other regions of France with a different land use were then selected to investigate the impact of geographic specificities on the scaling properties. For each 132 133 selected map, we analysed and compared the scaling properties of S2 images acquired in blue (B2),

- green (B3), red (B4), near infrared (B8) and short-wave infrared (B11, B12) bands. The latter were
 chosen considering their strong relation with Earth surface physical and biophysical processes and their
- 136 common use in previous studies involving multifractal analysis of surface reflectances (Alonso et al.,
- 137 2017; Lovejoy et al., 2008b). The several bands of S2 sensors enable to explore the possible impact of
- the acquisition wavelength on scaling behaviour. Moreover, vegetation and soil moisture indexes
- obtained from these bands were also analysed to investigate how the combination of such bands can
- 140 impact the scaling laws. This research study was done by applying multifractal analysis methods in the
- 141 framework of the Universal Multifractal model (Schertzer and Lovejoy, 1987).

143 2 Case study and data

144 2.1 Study area

145 Southwestern France is a highly anthropized site, mainly composed of agricultural land. More than 146 half of the ground area is exclusively dedicated to the development and management of crops throughout the year. Among the most common are sunflower, corn, wheat, barley, soybean and 147 148 rapeseed crops. Considering also the urban areas which constitute ~ 2% of the landscape, the 149 Southwest region has a spatial structure which is largely characterized by anthropic factors. These have 150 an impact on the spatial heterogeneity of physical and biophysical variables such as surface 151 temperature, soil moisture or vegetation. Natural factors also play a significant role in the spatial 152 structuring of this landscape and its related processes, such as, for example, natural grasslands and forests (~ 10% of the surface) (Baup et al., 2015; <u>http://osr-cesbio.ups-tlse.fr/~oso</u>). 153

154 The climate in this region is temperate, inducing seasonal effects that impact significantly the 155 temporal dynamics of surface variables. Springs are characterized by heavy rain, and summers are hot 156 and dry with temperatures often reaching 35 °C (Battude, 2017). These climatic conditions have a 157 direct effect on the landscape structure, more particularly on the evolution of this structure over the 158 seasons. Indeed, the crop cycles – described by harvest and sowing periods – contribute to modify the 159 heterogeneity of the surface as the seasons go by. In addition, the differentiation between winter and 160 summer crops implies that not all crops have the same seasonal evolution and do not affect the 161 landscape structure in the same way over time.

162 The study area is part of a French observatory, named Regional Spatial Observatory (RSO) and 163 labelled by the French National Institute of Sciences of the Universe (INSU) since 2007. The purpose is to answer scientific questions related to the functioning and evolution of continental surfaces, but also 164 165 to deal with more applied issues related to land use planning and sustainable management of 166 resources (Dejoux et al., 2012). In order to ensure this, both in-situ (Béziat et al., 2009; Burel, 2018; 167 Tallec et al., 2013) and remote sensing observations (Baup et al., 2012) are combined. Thanks to these 168 data, several works related to crop monitoring by remote sensing have been carried out in this region 169 such as the study of wheat and rapeseed (Fieuzal et al., 2013) or even the estimation by agro-170 meteorological model of crop yields such as sunflower (Fieuzal and Baup, 2015), soybean (Baup et al., 171 2015; Betbeder et al., 2016) and corn (Fieuzal et al., 2017). These studies also allowed the preparation 172 of recent space missions such as Venus or S2, subsequently leading to high spatial resolution modelling 173 of biomass and crop yields by optical data (Ameline et al., 2018; Battude et al., 2016; Claverie et al., 174 2012). The optical products from these new satellites now benefit from particular attention in this region, since they are made available on the THEIA platform (<u>http://www.theia-land.fr</u>), corrected for
atmospheric effects by the MAJA processing chain.

177 Due to the heterogeneity of its surfaces and the growing number of data acquired at various 178 spatial scales, southwestern France is particularly well suited to the application of multifractal analysis 179 tools such as those we used in this study. In this context, the present work consists of the analysis of 180 MAJA-corrected surface reflectances acquired in the region of the 30TYP S2 tile (according to the tiles 181 format and nomenclature provided by ESA; Gatti and Naud, 2017) whose spatial extent is shown in 182 Figure 1. More specifically, our study was focused on a smaller study area of 40 x 40 km² covering part of the Gers French department (red square in Fig.1) and located within the 30TYP tile. This area in the 183 184 Gers was chosen due to the predominant place of agricultural surfaces, limiting the impact of urban 185 areas (cities, roads) on the spatial variability of physical and biophysical processes. In addition, due to 186 certain constraints related to analysis tools (sample size, computation time) and cloud cover, we 187 deliberately chose our study area by selecting a smaller zone than the area corresponding to the 30TYP 188 tile which covers an area of 100 x 100 km².



Figure 1 : Study area in the Southwest region of France.

This study investigates the period extending from January to December 2016, using the Sentinel-2A satellite. With a ten-day revisit time, S2A provides a sufficient temporal frequency to analyze the global temporal dynamics of the properties of continental surfaces over an entire year. Finally, we can also point out that the Southwest region has a relatively low average cloud cover throughout the year compared to other regions of France located further north and south. Considering the impact of cloudson optical products, this criterion is not negligible in the choice of the study area.

195 2.2 Sentinel-2/MAJA products

196 We analyzed surface reflectances acquired in six spectral bands by the S2-A Multi-Spectral 197 Instrument (MSI; Martimort et al., 2007) and then corrected by the MAJA chain within the CNES 198 processing center (Centre National d'Etudes Spatiales; <u>https://theia.cnes.fr</u>). The selected bands are: 199 three bands in the visible domain – blue B2 (0.50 μm), green B3 (0.56 μm), red B4 (0.66 μm), one band 200 in near infrared B8 (0.84 μ m) and two bands in the short-wave infrared B11 (1.6 μ m) and B12 (2.2 μ m). 201 Note that all these bands are defined at 10 m resolution (in terms of pixel size) except B11 and B12 202 which are defined at 20 m resolution. The six bands were chosen because of their important link with 203 the physical and biophysical processes of continental surfaces. The visible/near infrared has already 204 been the subject of studies aimed at transcribing the fractal (De Cola, 1989; Lam, 1990; Qiu et al., 1999) 205 and multifractal (Alonso et al., 2017; Lovejoy et al., 2008b) properties of continental surfaces.

In addition to reflectances, two optical indexes were also studied: the Normalized Difference
 Vegetation Index (NDVI; Tucker and Sellers, 1986) and the Normalized Soil Moisture Index (NSMI;
 Lampkin and Yool, 2004), which respectively account for the condition of the vegetation and the level
 of soil moisture. Since these variables are frequently used to investigate the functioning of continental
 surfaces, it is important to study their scaling properties. NDVI is constructed from surface reflectances
 (ρ) in the red (band B4) and near infrared (B8) wavelengths. NSMI is obtained from the combination
 of the bands in short-wave infrared B11 and B12:

214

 $NDVI = \frac{\rho_{B8} - \rho_{B4}}{\rho_{B8} + \rho_{B4}} \qquad (10 \ m)$

 $NSMI = \frac{\rho_{B11} - \rho_{B12}}{\rho_{B11} + \rho_{B12}} \qquad (20 m)$

213

216

215

217 These indexes are built in a non-linear way from reflectances in individual spectral bands, and it 218 is interesting to check if (and how) potential scaling properties of ρ could result in scaling properties 219 of the indexes.

1

220 Cloud masks from MAJA were used to identify images with less than 15% cloud cover. Fifteen 221 dates were then selected over 2016 in order to limit the impact of clouds on surface reflectances and 222 their scaling properties. These dates are shown in the chronogram in Figure 2, where we can see a 223 distribution of dates that covers all seasons, with a larger number of images in summer. For the fifteen 224 images, an average cloud cover of 2% was estimated over the period, with only 4 dates having more than 2% of clouds - the largest cloud cover being obtained on June 23 with 12.5%. Each of the 225 226 remaining cloud pixels in the selected images were considered as missing data and were interpolated 227 by bilinear interpolation. Note that no missing data other than that related to clouds was observed in 228 our data, and that the corresponding size of the selected images are 4096 x 4096 pixels and 2084 x 229 2048 pixels respectively for 10 m and 20 m resolution products.



Figure 2 : Cloud cover of Sentinel-2A products selected over the January-December 2016 period, and corrected for atmospheric effects by MAJA. The corresponding dates were selected with less than 15% cloud cover.

231 3 Multifractal analysis methodology

232 During the last century, researchers realized that many geophysical processes could present scale invariance or scaling properties (Kolmogorov, 1941; Mandelbrot, 1967; Richardson, 1922). Multifractal 233 234 analysis makes it possible to highlight these scaling behaviours, enabling spatial characterization of 235 surface processes. To do this, several scale sensitive statistical estimators are applied to the surface 236 reflectances acquired by satellite sensors. These estimators can be the power spectrum, the structure 237 functions or the statistical moments. Different multifractal models, like the Universal Multifractal 238 model (Schertzer and Lovejoy, 1987), are used to parametrize the scaling properties obtained from the 239 statistical estimators. In this paper, a focus was made on one estimator, the power spectrum, which 240 has the advantage to detect, in an easy and rapid way, some spatial properties of geophysical variables 241 over different space scales (Lovejoy et al., 2008b). Structure functions and statistical moments (in the 242 framework of Universal Model parametrization) were also investigated in this work. Since they confirm the properties already observed by power spectra, and for sake of simplicity and clarity, the theory 243 244 and the methodology relative to these estimators are not presented in this part but are detailed 245 respectively in Appendix A and B. More generally, we invite the reader to refer to Appendix A for a 246 deeper understanding of the theoretical framework of fractals and Universal Multifractal model.

Because our study treats time series of satellite images, the statistical estimators will be computed for each image available during the studied period. Therefore, from this analysis two types of spatial characterization will be studied: (1) a mean characterization resulting from the average of individual estimators computed within the period, and (2) a "date by date" characterization of each image, allowing us to have access to the temporal evolution of the scaling properties of the products studied. For the calculation of estimators, note that, since only satellite images will be treated in this study, we will focus on the two-dimensional versions of the following techniques.

254 3.1 Vocabulary : scale and resolution

For the sake of clarity, we primarily want to define the terms "resolution" and "scale" within the meaning of fractal mathematical formalism. Indeed, although their meaning is close to that usually used by geophysicists, in particular in remote sensing, it can be appreciably different and could lead to certain misinterpretations in the remainder of this paper.

If we consider a space of (geometric) dimension D = 2 discretized over a set of intervals or pixels, the resolution λ of any fractal entity can be defined as the number of pixels present along a dimension of space. In the case of a satellite image of size 4 x 4 pixels, the resolution of the fractal entity is therefore equal to $\lambda = 4$ (Fig.3). Inversely proportional to the resolution, the scale *l* is associated with the size of a pixel and is defined according to λ :

 $l = \frac{L}{\lambda}$

3

with *L* the extent or domain size along one dimension and corresponding to the greatest value of scale *l*. Unlike λ , the scale *l* is a dimensional variable (meters for an image). In the case of our 4 x 4 satellite image, the scale is l = L/4, which is the size of an elementary pixel. If we now aggregate our image on grids made up of new pixels twice as wide (or with an area four times as large), the resolution of the image will be taken equal to $\lambda = 2$, and the scale will be l = L/2.



Figure 3 : Illustration of the concepts of resolution and scale from a twodimensional fractal entity (satellite image).

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In this way, any variable with fractal properties is defined on a continuum of resolutions (or scales) which extends from the finest resolution (to which is associated the smallest scale *i.e.* the size of an elementary pixel), to the coarsest resolution ($\lambda = 1$) corresponding to the largest scale (l = L).

275 In the following, scale sensitive statistical estimators will be called as "scaling estimators", and will 276 be computed on a spatial support represented either by resolutions λ , scales l, varying window size 277 Δx or even spatial frequencies k, depending on the initial definitions of the estimators. Studying the 278 evolution of surface products according to these spatial variables (λ , l, Δx , k) provides access to 279 fractal/multifractal properties that will be called hereafter as "scaling properties". A scale range, 280 corresponding to a set of scales with similar scaling properties, will be called as "scaling regime".

281 3.2 Power spectrum: evidence of spatial properties

Power spectra are highly sensitive to the transition scales (also called "scale breaks") that are located between scaling regimes characterized by different scaling properties. In a first step, the twodimensional power spectral density $P(k_x, k_y)$ of the data under analysis, *X*, is estimated:

285
$$P(k_x, k_y) = |fft(X)|^2$$

286

with *P* the power spectral density (PSD) defined on both vertical and horizontal image axes, corresponding respectively to k_x and k_y wavenumbers (spatial frequencies). Here, the estimation of the PSD is done through a two-dimensional *fft* or Fast Fourier Transform. Then, the more convenient one-dimensional isotropic angle-integrated power spectrum E(k) is obtained (Lovejoy et al., 2008b; §8), that provides a 1D representation of the 2D *fft* power spectra:

292
$$E(k) = \int_{\|\vec{k}\|=k} P(\vec{k}) d\vec{k}$$

293

where k is the modulus of the wavenumber and $\|.\|$ is the Euclidean norm. Since it expresses space frequencies, k is directly related to the spatial resolution λ . If the process presents scaling properties, the spectrum should follow a power law, where β is the negative slope of E(k) on a log-log graph:

297

$$E(k) \approx k^{-\beta}$$

298

The symbol " \approx " indicates an equality within the limits of slowly varying functions. This method makes it possible to identify the range of scales (here the values of k) where there are scaling properties, and this through a first parameter called the "spectral exponent".

302 3.3 Scale break detection

303 When analysing the scaling estimators, a segmentation algorithm (D'Errico, 2017) was 304 systematically used to detect the scale breaks and thus the scale ranges in order to avoid subjective 305 detections. This algorithm is based on the non-linear regression of "spline" curves (piecewise 306 polynomial functions), making it possible to estimate the position of slope breaks within scaling 307 estimators. It requires three input parameters: (1) the position on the x-axis of each point of the cloud 308 to be segmented (support vector *i.e.* k, Δx or λ), (2) the value of each point on the y-axis (estimator 309 vector such as E) and (3) the number of breaks *nb_knots* to detect. Once applied, the function outputs 310 the position of the detected nb_knots breaks. Note that the number of breaks is a parameter which 311 is therefore defined by the user. In order to check if the position is reliable and independent of the

4

5

- number of breaks proposed, the algorithm was applied for several values of nb_knots ($nb_knots =$
- 313 1, 2, 3), ensuring the same break was well detected in each case.

315 4 Results

316 4.1 Mean characterization over the period

Figure 4 shows the mean power spectra estimated over the full period (January-December 2016) of the different S2/MAJA products presented in Section 2. This represents an average spectrum based on individual spectra obtained within the period. Each spectrum is plotted in log-log coordinates, with the horizontal axis converted into spatial scale $l (= k^{-1})$ expressed in kilometres. The scale break detection algorithm identified two scale breaks common to all products: a first break l_1 located between a hundred meters and one kilometre, and a second break l_2 located on smaller scales. l_1 and l_2 positions are summarized in Table 1 for both surface reflectances and optical indexes.



Figure 3 : Mean angle-integrated power spectra estimated from Sentinel-2 bands and indexes, over the period January-December 2016.

324

Considering all these products, the first scale break l_1 was estimated at about 400 m (computed from the geometric mean of individual breaks). The spectral slope β of each mean power spectrum (Table 2) was estimated for both ranges I ($l_1 > l$) and II ($l_1 > l > l_2$) (the last range (III) will be considered separately in Section 6). <u>On range I</u>, slopes are very close to zero and do not present significant differences within the different products ($0.2 < \beta < 0.3$). These results are in contradiction with those of Lovejoy et al. (2008b) who identified slopes greater than 1 ($\beta = 1.17 \pm 0.08$) on a similar 331 scale range (500 m – 25 km), from reflectances acquired in comparable spectral bands (MODIS sensor; 332 Justice et al., 1998). Similarly, our results are also not comparable to those of Alonso et al. (2017) who 333 observed only a single scaling regime over the range (120 m - 15 km), for reflectances from similar 334 bands (Ikonos and Landsat-7 satellites). However, because of the different multifractal formalism used 335 by Alonso et al. (2017), our spectral slopes and their results are difficult to compare. Nevertheless, 336 Alonso et al. (2017) have indeed shown the presence of scaling properties over this range of scales. 337 Thus, the differences between our results and those of Lovejoy et al. (2008b) and Alonso et al. (2017) can be linked to geographic specificities (central Spain in their cases) or to the fact that their analyzes 338 339 were carried out on one or two particular dates, not allowing the acquisition of the scaling behaviour 340 averaged over a period as we did here.

341

	Grid spacing	SPECTRA		STRUCTURES	MOMENTS	
	(m)	l ₁ (m)	l ₂ (m)	l ₁ (m)	l ₁ (m)	
B2	1	420	37	580	360	
B3		390	37	539	300	
B4	10	390	39	555	360	
B 8		290	37	495	320	
NDVI		380	48	575	340	
B11	C	320	75	503	370	
B12	20	370	75	565	390	
NSMI		420		695	430	

Table 1 : Positions of the scale breaks l_1 and l_2 estimated by segmentation (D'Errico, 2017) on mean Sentinel-2 power spectra, structure functions and statistical moments.

342

343 Regarding the behaviour of spectra on range II $(l_1 > l > l_2)$, β values are greater than one and 344 are relatively different between the products. The lowest slope value is estimated for the near infrared 345 band ($\beta \approx 1.4$) and the greatest ones are found for the two short-wave infrared bands ($\beta \approx 2$). 346 Considering visible bands (B2, B3, B4) and optical indexes (NDVI, NSMI), the β values obtained are 347 between 1.6 and 1.8. We can also note in the spectra an increase of slope with the wavelength for 348 reflectances acquired in visible domain.

As for the fine-scale break l_2 , the latter is located at around 40 m for products with a resolution of 10 m (visible bands, near infrared band and NDVI) and is shifted towards larger scales (75 m) for products with a 20 m grid resolution (short-wave infrared bands B11 and B12). It even disappears with regard to the NSMI soil moisture index (flat spectrum for the fine scales). If we analyze this last 353 spectrum in more detail, a rise in the curve is observed towards the points located at the finest scales 354 (~ 50 m). This rise reflects the presence of noise in the NSMI data which affects the highest frequencies 355 of the image, artificially creating an excess of energy at fine scales. Since bands B11 and B12, from 356 which the NSMI has been estimated, do not show such an increase in their spectra at fine scales, their 357 combination must be the source of the noise observed on the soil moisture index at high frequencies. 358 This phenomenon was expected since the ratio of two signals is known to contain more noise (due to a reduction in the signal to noise ratio). Similarly, we can note that the NDVI is characterized by a l_2 359 360 scale break (at 48 m scale) that is less pronounced than for the bands B3 and B8 from which it was 361 generated, with a very slight rise of the spectrum for highest frequencies (scales close to 20 m). 362 Although the phenomenon is much less pronounced than for the NSMI, the vegetation index also 363 appears to be somewhat impacted by the combination of bands.

364 Continuing the comparison between bands and indexes, these results tend to show that on fine 365 scales (range II) the vegetation index is characterized by an average spectral behaviour comparable to the band B4 (red) but different from band B8 (near infrared). However, the spectral slope of the soil 366 367 moisture index is found to be close to those obtained from both of the short-wave infrared bands B11 and B12. Another general remark on the spectra can be made on the quality of the scaling law which 368 is very good on range II, considering all products ($R^2 \approx 0.98$), while it is strongly reduced on range I 369 $(R^2 \approx 0.6)$. This can partly be explained by the sampling of power spectra which is sparser on low 370 frequencies (large scales) than on high frequencies (fine scales). 371

The results presented above were compared with the two other approaches (structure functions and statistical moments presented in Appendix B) and their results (figures in Appendix C) were found to be consistent with those of the power spectra. The positions of l_1 scale break obtained from these approaches are presented in Table 1, with the corresponding parameters detailed in Table 2 for both ranges.

(a)	Grid spacing	Scale range	β	R^2_{β}	α	<i>C</i> ₁	$R^2_{K(q)}$	н	R^2_H
B2			0,30	0,62	1,96	0,007	0,89	-0,27	0,98
B3			0,31	0,64	1,86	0,006	0,91	-0,28	0,99
B4	10 m	range I	0,24	0,55	1,81	0,007	0,92	-0,30	0,99
B8			0,29	0,60	1,80	0,004	0,89	-0,27	0,99
NDVI			0,20	0,50	1,91	0,005	0,88	-0,31	0,99
B11			0,28	0,56	1,83	0,004	0,87	-0,29	0,99
B12	20 m	range I	0,25	0,53	1,81	0,005	0,88	-0,29	0,99
NSMI			0,22	0,55	2,23	0,005	0,83	-0,29	0,99
(b)	Grid spacing	Scale range	β	R^2_{β}	α	<i>C</i> ₁	$R^2_{K(q)}$	н	R^2_H
(b) B2	Grid spacing	Scale range	β 1,63	R² _B 0,99	α 1,97	C ₁ 0,11	R ² _{K(q)} 0,96	Н 0,32	R ² _H 0,98
(b) B2 B3	Grid spacing	Scale range	β 1,63 1,67	R ² β 0,99 0,99	α 1,97 1,78	C ₁ 0,11 0,11	R ² _{K(q)} 0,96 0,96	Н 0,32 0,32	R²_H 0,98 0,98
(b) B2 B3 B4	Grid spacing 10 m	Scale range range II	β 1,63 1,67 1,76	R²β 0,99 0,99 0,99	α 1,97 1,78 1,56	C ₁ 0,11 0,11 0,12	R ² _{K(q)} 0,96 0,96 0,96	Н 0,32 0,32 0,35	R²H 0,98 0,98 0,97
(b) B2 B3 B4 B8	Grid spacing 10 m	Scale range range II	β 1,63 1,67 1,76 1,37	R²β 0,99 0,99 0,99 0,98	α 1,97 1,78 1,56 1,53	C ₁ 0,11 0,11 0,12 0,09	R ² _{K(q)} 0,96 0,96 0,96 0,96	H 0,32 0,32 0,35 0,23	R ² _H 0,98 0,98 0,97 0,93
(b) B2 B3 B4 B8 NDVI	Grid spacing 10 m	Scale range range II	β 1,63 1,67 1,76 1,37 1,84	R ² β 0,99 0,99 0,99 0,98 0,98 0,99	α 1,97 1,78 1,56 1,53 1,61	C ₁ 0,11 0,11 0,12 0,09 0,13	R ² _{K(q)} 0,96 0,96 0,96 0,96 0,96	H 0,32 0,32 0,35 0,23 0,37	R²H 0,98 0,98 0,97 0,93 0,98
(b) B2 B3 B4 B8 NDVI B11	Grid spacing 10 m	Scale range	β 1,63 1,67 1,76 1,37 1,84 1,96	R ² _β 0,99 0,99 0,99 0,98 0,99 0,97	α 1,97 1,78 1,56 1,53 1,61 1,53	C ₁ 0,11 0,11 0,12 0,09 0,13 0,11	R ² _{K(q)} 0,96 0,96 0,96 0,96 0,96 0,97	H 0,32 0,32 0,35 0,23 0,23 0,37	R ² _H 0,98 0,98 0,97 0,93 0,98 0,95
(b) B2 B3 B4 B8 NDVI B11 B12	Grid spacing 10 m 20 m	Scale range range II range II	β 1,63 1,67 1,76 1,37 1,84 1,96 1,96	R ² _β 0,99 0,99 0,99 0,98 0,99 0,97 0,98	α 1,97 1,78 1,56 1,53 1,61 1,53 1,53	C ₁ 0,11 0,11 0,12 0,09 0,13 0,11 0,12	R ² _{K(q)} 0,96 0,96 0,96 0,96 0,96 0,97 0,96	H 0,32 0,32 0,35 0,23 0,37 0,35 0,37	R ² _H 0,98 0,98 0,97 0,93 0,98 0,95 0,97

Table 2: Scaling parameters obtained from multifractal analysis on range I (a) and range II (b) of Sentinel-2 products, over the period January-December 2016. The values of β , α , C_1 and H were obtained from linear regressions applied on mean statistical estimators. However, R^2 coefficients presented here corresponds to the average of all coefficients obtained on each date. Note that $R_{K(q)}^2$ is the average of the coefficients obtained on every moment curves. Similarly, H is obtained from the average of H_x and H_y , as R_H^2 is obtained from R_{Hx}^2 and R_{Hy}^2

379

380 4.2 Temporal and spatial variations of the scale breaks

Different scaling regimes were observed on the S2/MAJA products, separated by specific scale breaks: a first break l_1 detected on the scaling estimators between hundreds of meters and one kilometre, and a second break l_2 detected on the high frequencies of the mean power spectra. To explain the origins of these observations and their link with the spatial texture of the data, the evolution of these breaks was analysed according to the seasonal and spatial evolution of the surface.

To explore the <u>seasonal evolution of the scale breaks l_1 and l_2 </u>, the scale break detection algorithm was applied on each individual spectrum obtained between January and December 2016 (*i.e.* spectra computed for each available image). The time series of l_1 (Fig.5.a) is found to be largely variable for all bands, whereas l_2 (Fig.5.b) presents a quite constant evolution during the period. Therefore, these

two scale breaks detected on power spectra seem to be of a different nature: on the one hand the large-scale break l_1 seems to be related to the seasonal evolution of surface processes, on the other hand the fine-scale break l_2 position is found to be relatively independent of time. In addition, we note that the ratio between l_2 and the minimum scale (that is to say the grid resolution of the image for a given band) is comparable considering both 10 m and 20 m bands: $l_2 \approx 4 \times l_{min}$. This relationship appears to be constantly present throughout the period.



Figure 5: Time series of the position of the scale breaks l_1 (a) and l_2 (b) estimated from the power spectra of the Sentinel-2 bands, over the period January-December 2016.

396 To understand the origin of these two scales, it was investigated whether the geographic 397 specificities related to the study area could have an impact on the position of these scale breaks. Since 398 l_1 relies on the seasonal evolution of the surface, changes in land use should have an effect on the 399 position of l_1 , whereas they should not on l_2 . To so investigate this question, several regions of France 400 with a different land use were selected. These regions were identified using the land use products 401 supplied by the Scientific Expertise Center for Land Use CES OSO (http://osr-cesbio.ups-tlse.fr/~oso). 402 This analysis was done from the land use map for 2016, obtained from S2 products, and presented in 403 Figure 6. In addition to the original tile (30TYP), five new S2 tiles were added to our study, 404 corresponding to five types of surface or classes: crops, forest, urban, vineyards and grasslands. The areas were chosen so that the extracted S2 images contain for each class a majority of pixels belonging 405 406 to the class in question. For the six zones, we extracted the surface reflectances acquired in two 407 spectral bands defined at 10 m (B2, B8) at two dates in the year 2016 (summer, autumn). Table 3 408 specifies the position of the breaks l_1 and l_2 detected in the power spectra. In general, considering all 409 the zones, bands and dates provided, these results show a break l_2 which varies between 32 m and 42 410 m, and a break l_1 which is between 72 m and 1 km. Therefore, the position of the large-scale break l_1 411 seems much more impacted by the type of surface, the date of acquisition and the spectral band (when it is defined according to the same spatial resolution) than the fine-scale break l_2 . 412

- 413 These temporal and spatial behaviours confirm the different nature of the transition scales. In the
- following sections, demonstrations are made to relate these scale breaks with (1) the seasonal and
- 415 geographic specificities of the surface in the case of the scale l_1 and (2) with the acquisition system of
- 416 optical images for l_2 .



Figure 6: Six study areas corresponding to the five types of surface (crops, forest, urban, vineyards, grasslands) identified from the 2016 land use map of France CES OSO.

	Summer				Autumn					
	Band B2		Band B8		Band B2		Band B8		_	
	lı	l_2	<i>l</i> ₁	l_2	l_1	l_2	l ₁	l_2		
crops south	440	36	340	35	390	37	72	34		
forests	860	32	1000	39	910	34	210	38		
crops north	310	36	380	34	380	33	460	41		
urban	-	40	-	42	-	39	-	40		
vineyards	150	34	180	35	150	34	130	35		
grasslands	280	37	350	34	140	37	250	37		

Table 3: Positions (in meters) of the scale breaks l_1 and l_2 estimated from the power spectra corresponding to six types of surface, observed in bands B2 and B8, in summer and autumn 2016.

417 5 A large-scale break related to the seasonal and geographic specificities of the surface

418 5.1 Seasonal dynamics linked to crop cycle

To propose an explanation to the large-scale break l_1 between ranges I and II, first the <u>temporal</u> evolution of the scaling properties of the S2 products can be observed on the two ranges, for the studied period. The time evolution of the spectral exponent β estimated for the six bands (Fig.7) exhibits on the fine scale range (*range II*) a seasonal cycle of the spectral slopes, while no particular behaviour is detected on large scales (*range I*). In addition, the dynamics of the β values are much higher on fine scales ($0.8 < \beta < 2.4$) than on large scales ($0 < \beta < 0.9$).



Figure 7: Time series of spectral exponents of the Sentinel-2 bands, over the period January-December 2016 for range I (a) and range II (b).

425 If we analyze in more detail the cycle observed on fine scales (Fig.11.b), the dynamic detected on 426 the visible bands (B2, B3, B4) is characterized by minimum values in February, August and then December ($\beta \approx 1.5$) and peaks of values around May-June then October ($\beta \approx 1.9$). For the <u>short-</u> 427 428 wave infrared bands (B11, B12), the maxima are placed at the same dates as for the reflectances 429 acquired in the visible domain but they do not present the same intensity: the peak observed in May-430 June is very marked – with values reaching 2.3 (arrow on the graph) – while that of October is weaker 431 (even almost non-existent for B11). Considering the near infrared B8 band, a reversed cycle is found 432 on range II (compared to the other bands), characterized by minimum values for the months of May-433 June and then October. In addition, the spectral exponent of B8 shows a much more marked dynamic 434 in the second half of the period – values of β falling from 1.5 (July) to 0.9 (October, arrow on the graph) – than in the first half of the period $(1.4 < \beta < 1.7)$. Thus, for the month of October, the images 435 436 acquired in the near infrared seem to lose spatial autocorrelation, with a spectral slope approaching 437 zero. This temporal evolution of scaling properties is illustrated on Figure 8 where the individual power 438 spectra of four near infrared images are compared: two images acquired in summer (July 10 and July 439 30) and two images acquired in autumn (September 28, October 21). For these spectra, the fine-scale 440 break l_2 was estimated at around 35 m in all four cases, confirming its "instrumental" origin (see 441 Section 6). On the other hand, two different l_1 breaks were obtained: at around 350 m for the July 442 spectra (dark red and light red curves) and the September spectrum (light green curve), and at 72 m 443 for October (dark green curve). It appears that the l_1 break in July moves towards fine scales in 444 October, combining with l_2 . Thus, these results show that a switch takes place from two scaling 445 regimes in summer and early autumn to a single one in autumn. On this period, we observe the decrease of the spectral slope with time, confirming the evolution presented in Figure 7.b. For October 446 447 21, the regime seems to be found on a scale range close to the global range provided by our data (72 448 m - 41 km), and presenting a spectral slope less than 1.



Figure 8: Power spectra of surface reflectances acquired in summer and autumn 2016 by the near infrared band B8.

449

450 These results confirm the link between the fine-scale break l_1 and the temporal and spatial 451 variations of the surface. Therefore, considering the predominant place of agricultural surfaces in the 452 study area, our hypothesis is that this time variation is related to the crop cycle (with the more abrupt 453 changes associated with harvesting and sowing), taking place during the mid-season periods (autumn, 454 spring). The evolution of crops, including the transitions between winter and summer crops, would explain the cycle that we detect here. This hypothesis can also explain the reverse nature of the cycles 455 456 between near infrared and visible/short-wave infrared behaviours. Indeed, the reflectances acquired 457 in the near infrared, being very sensitive to vegetation, should exhibit more spatial variability when

458 the crops are well developed, therefore for periods between the harvest and sowing phases. If we take 459 the example of corn, which is one of the main crops grown in Southwestern France, the harvesting 460 phase usually occurs between August and November. This corresponds to the fall observed on the β 461 parameter estimated from the band B8. Thus, a harvest phase would correspond to a loss of spatial 462 heterogeneity in these wavelengths, revealing a poorer and less structured signal than during a 463 vegetation growth phase. In addition, the gradual drop in spectral slope (beginning of decrease 464 observed from August) can be explained on the one hand by the staggering of the sowing phases 465 between August and November (according to the calendar of each farmer), and on the other hand by 466 the senescence phase occurring before harvest and helping to reduce the ability of vegetation to 467 reflect solar radiation (less chlorophyll). Conversely, visible reflectances behave differently by being sensitive to other elements of the surface such as roads, buildings, or even the bare soil that appears 468 469 once harvesting is complete. Thus, at mid-season, when the vegetation has just been sown, the surface 470 shows more spatial variability in the visible than in the near infrared domain.

471 To develop the interpretation proposed above, our results were compared with some 472 characteristics of the surface, and more particularly with its spatial structure. The study area is mainly 473 composed of agricultural plots of which the majority reveals an area between 0.3 and 130 ha, 474 corresponding to a plot width (assuming square areas) of around 50 to 1000 m. This could be identified 475 thanks to the Graphic Parcel Register (GPR; https://geoservices.ign.fr) for the year 2016, allowing 476 access to the delimitation, the surface area and the type of crops of all the agricultural plots in the 477 department of Gers. An average plot width of around 170 m was deduced from this data. Figure 9 478 below shows a part of the GPR for our study area, for which the plots were classified according to their 479 extent. We observe a large number of plots whose characteristic widths are between 100 m and 1 km 480 (areas in light blue, green and yellow). Although the agricultural plots reveal quite different structures, of more or less elongated shapes, even with rounded "edges", the more or less systematic cutting that 481 482 they generate on the surface could be at the origin of the large-scale break l_1 observed in our results. 483 In this way, the fine-scale regime observed on range II would reflect the intra-plot variability 484 governed by multifractal physical and biophysical processes linked to crops (soil moisture, vegetation, 485 etc.), while the large-scale regime on range I would represent the extra-plot variability resulting from 486 the aggregation of several fine-scale processes which together would not have particular multifractal 487 properties. This interpretation would explain the uncorrelated and less complex nature of the 488 reflectances on large scales (very low values of β and C_1 parameters on range I), whereas more 489 complex fractal and multifractal properties were obtained on range II tending to justify the interest 490 of applying a multifractal model like that of Universal Multifractals on these fine scales.



Figure 9: Classification of agricultural plots (parcels) according to their surface area estimated from the Graphic Parcel Register for the year 2016. The data represented here correspond to an area of $\sim 3 \times 3 \text{ km}^2$ extracted in our main study area. The plot widths were calculated from the surface areas by approximating them to square areas.

491

492 5.2 Simulation of the impact of agricultural plots on the scaling properties of surface 493 reflectances

494 In order to check if the previous explanations are realistic, an example of agricultural area was 495 simulated with the same scaling properties as the S2 reflectances. To do this, a synthetic multifractal 496 field was simulated by applying the Universal Multifractal model (see Appendix D), to which we 497 combined an image containing the delimitations of agricultural plots obtained from the GPR-2016 on 498 our study area (Fig.10.a). The multifractal field was generated by defining the values of the universal parameters of the model at $\alpha = 2$, $C_1 = 0.1$ and H = 0.3, corresponding to the values obtained for 499 500 the band B2 on range II (mean behaviour, Table 2.b). The combination method simply consists in 501 adding the spatial texture generated by the agricultural plots to the image of the simulated multifractal field. To reproduce a more realistic texture, the entities ("plot" objects) have been previously grouped 502 503 into six classes: the same pixel value between the min and max values of the initial field, *i.e.* 0 and 1, 504 has been assigned to each of these classes. These have been defined according to the area of the 505 entities: the plots have been grouped into object classes whose area is within a predefined interval. 506 Here, these intervals have been defined from the area quantiles, so as to have the same number of 507 objects in each class. Although this classification does not correspond in any way to reality, it allows to 508 reproduce artificially the differences in intensity of reflectances according to the type of soil (different 509 types of crop, bare soil, forest, etc.) which can be observed on the real data (Fig.10.b). The final surface 510 generated (i_{agri}) is based on an integrated multiplicative cascade (i_{MF}) defined at all scales of the 511 image, but to which is added the spatial texture of agricultural plots (*offset*_{parcels}) such as:

514

$$i_{aari} = i_{MF} + offset_{parcels}$$

513



Figure 10: Simulated and real surface reflectances corresponding to the study area (Gers, 30TYP tile). The simulation was obtained by combining a multifractal field and the GPR of the year 2016.

515

516	The graph in Figure 11 represents the power spectrum obtained for the simulated agricultural
517	surface i_{agri} (green curve), compared to that obtained from the original simulated multifractal field
518	i_{MF} (blue curve). The power spectrum estimated from an image i_{B2} of real surface reflectances (gray
519	curve) is also presented. $\rm i_{B2}$ is the image acquired in B2 band on July 10 (image presented in Fig.10.b).
520	For the sake of clarity, the finest scale of the simulated fields was set to the size of a S2 pixel (band B2)
521	so that $l_{min} = 10 m$.

522 As expected, a single scaling regime is observed on the original multifractal i_{MF} with no scale break 523 detected, whereas two scaling regimes are indeed observed for i_{agri} , with a scale break l_1 detected at 524 around 200 m. This break noticed on i_{agri} is located a few dozen scales higher than the finest scale, 525 which is comparable to the results obtained from S2 data. The two regimes exhibit behaviours close to those observed on the ranges I and II of i_{B2}, at least in terms of the shape of the curves. Indeed, on 526 527 large scales (range I) a spectral slope smaller than 1 is obtained, while a spectral slope greater than 1 is estimated on fine scales (range II). Therefore, the twofold scaling behaviour of S2 products is clearly 528 529 transcribed by our simulation : it was evidenced here that a definite link exists between the division of

530 the area by agricultural plots and the presence of two scaling regimes – with a transition scale which 531 is effectively close to the average plot scale. We note a flattening at the finest scales of the i_{agri} power spectrum, which reveals the presence of noise localized on the last octave (20 m - 40 m). A similar 532 533 flattening was observed on the NSMI spectrum (Fig. 4) due to the creation of noise by the application of the NSMI formula. For the i_{agri} spectrum, this phenomenon is probably due to the "local" effect of 534 535 the boundaries of agricultural plots at resolutions close to the pixel size. Indeed, these boundaries 536 generate abrupt transitions from one pixel of the image to its neighbour, leading to an increase of the 537 energy at this specific scale. The fact that this flattening is observed on the simulated agriculture 538 spectrum (i_{agri}) but is not on the real spectrum (i_{B2}) may be due to (1) the impact of our combination 539 method creating artificial boundaries that are more abrupt on the simulation than on the real case, 540 and/or (2) the impact of the optical acquisition system (Section 6) on the real spectrum which inversely 541 affects the fine scales (slope increased) but that is not considered in this simulation.

542



Figure 11: Impact of agricultural plots on power spectrum of simulated surface reflectances (Fig.10.a).

543

545 6 A fine-scale break related to the optical acquisition system of the satellite sensor

546 This section is dedicated to the interpretation of the scale break l_2 observed on power spectra 547 between range II and III.

548 6.1 The effect of the Modulation Transfer Function

549 Considering the relatively small variations of the position of the fine-scale break l_2 with time and 550 land use, this break could be a consequence of the acquisition system which affects the high spatial 551 frequencies of images. In any optical system, this phenomenon is modelled by a function called MTF 552 (Modulation Transfer Function; Norton et al., 1977). It defines the system ability to transmit the spatial 553 frequencies of the observed scene. It is affected by several components of the system such as the 554 objective (the lens) and detectors, but also by external factors such as temperature or atmospheric 555 diffusion. During the satellite launch in space, these last factors can modify the MTF. However, the 556 latter varies very little over time once the satellite is finally placed in its orbit. In this way, it could be 557 the reason for the systematic break observed at the maximum frequencies acquired by the MSI-S2 558 sensor.

As part of the validation and calibration procedures for S2 products, Gascon et al. (2017) determined the modulation transfer function of the MSI imager after the launching of the S2-A satellite. This has been estimated from an image acquired in the B2 band on Arizona (USA), applying the "slanted-edge" method (Viallefont-Robinet and Léger, 2010; Reichenbach et al., 1991) which has often been used to determine the MTF of space sensors like Landsat (Carnahan and Zhou, 1986) or lkonos (Choi, 2002).

565 6.2 Simulation of the impact of optical acquisition system on scaling properties

Given what has been described above, we made the hypothesis that the scale break l_2 detected on the power spectra is directly related to the impact of the MTF on the scaling properties of the "original" signal acquired by the MSI sensor.

To demonstrate this, a multifractal field was simulated (according to the method presented in Appendix D which is based on the Universal Multifractal model) representing the scaling behaviour of a hypothetical original signal, *i.e.* with fractal/multifractal properties not impacted by the optical system. Here, we suppose that the original signal should have the same scaling properties as those obtained on *range II* (Figure 4), with parameters set as $\alpha = 2$, $C_1 = 0.1$, H = 0.3 (values obtained from band B2, Table 2.b). Note that we did not consider in this simulation the impact of agricultural plots (Section 5.2). 576 Then, to simulate the impact of optical acquisition system, we generated an MTF according to the 577 model of Viallefont-Robinet and Léger (2010). The latter proposed a parametric model taking into 578 account various characteristics of the system such as the diffraction effects of the lens and the 579 sampling effects of the detector, as well as the impact of other factors such as the displacement of 580 electronic charges or correction of the focus that can be achieved after launching the satellite. Here, 581 for the sake of simplicity, only the major effects on MTF have been taken into account, namely 582 diffraction and sampling. Finally, this function was applied to the simulated multifractal field. The 583 operation was carried out in the Fourier domain by applying the following relation:

584 585

$$I_{filtered} = MTF \times I_{simu}$$

8

with $I_{filtered}$ and I_{simu} the Fourier transform modules of the images corresponding respectively to the field filtered by MTF and to the simulated original field.

588 For the calculation and generation of the MTF, the equations used are detailed in Appendix E. Two parameters are to be defined in the model: (1) the sampling frequency f_s of the detector which was 589 fixed here at the resolution of the simulated multifractal field, that is, at 1 pixel⁻¹ (here, the scale is 590 expressed in pixels, the corresponding unit is thus "pixel⁻¹" for frequency), (2) the cut-off frequency f_c 591 of the lens (also defined at 1 pixel⁻¹ for the sake of simplicity) which is responsible for the diffraction 592 593 effects. In reality, the value of f_c depends on the characteristics of the lens (diameter, focal length) and 594 on the wavelength of the incident radiation. However, since we did not have access to the value of the 595 MSI sensor focal length, we were unable to determine the actual cut-off frequency. According to 596 Gascon et al. (2017), the MTF they measured for band B2 has a value close to 0 at the sampling 597 frequency. This shows that, for this band, f_c should be positioned close to f_s , at slightly higher frequencies. According to this assumption, we will suppose thereafter that $f_s \approx f_c = 1$ pixel⁻¹. 598

599 Figure 12 presents the power spectra of the original field (blue curve), of the field affected by the 600 MTF by considering only the sampling effects of the detector (dark brown curve), only the diffraction 601 effects of the lens (light brown curve), and finally by considering both the sampling and diffraction 602 effects (red curve). For the sake of clarity, the scale axis of the figure was adapted to correspond to that of S2 spectra: the finest scale of the fields was set to the size of a S2 pixel (band B2) so that $l_s =$ 603 $\frac{1}{f_c} = 10 m$. The fine scales are found to be impacted by the MTF the more significantly: the energy of 604 605 the signal is attenuated for the high frequencies, resulting in an increase in the spectral slope for scales 606 lower than 100 m. This behaviour is obviously more pronounced when the two effects (diffraction and 607 sampling) are combined (spectrum in red). The scale break detection algorithm (Section 3.3) was 608 applied on this last spectrum, which revealed a scale break located at 48 m (= $4.8 l_s$). The relationship

- 609 found here between the position of the break and the resolution of the image is comparable to that
- 610 deduced from our S2 products, for which the coefficient linking the two scales was close to 4 over the
- 611 entire period (Fig.5.b).



Figure 12: Impact of the Modulation Transfer Function on the power spectra of simulated multifractal fields. The MTF was generated from the analytical model described in Viallefont-Robinet and Léger (2010).

612

613 The MTF model used in this study enables to reproduce the impact of the acquisition system on 614 the high frequencies of the image. Although the value of the cut-off frequency was chosen in a more 615 or less arbitrary way, the scale break generated on the original spectrum is similar to that of the finescale break l_2 observed on the spectra of our S2 bands. These results provide a physical explanation 616 617 as to the origin of this fine-scale break, which is the result of the acquisition system, *i.e.* MTF, that 618 reduces spatial variations (smooths the image) at high frequencies inducing a decrease of energy at 619 the finest scales. The strong similarity of the red and light brown curves in Figure 12 suggests that the rupture is mainly related to the diffraction effects of the lens. Further research may be needed to be 620 conclusive of the latter findings by integrating other factors influencing the MTF (electric charge 621 622 displacements, temperature, etc.). However, the interpretations given here show the potential of 623 power spectra estimators to easily detect the on-orbit effective resolution of the optical sensors, *i.e.* the actual resolution below which any signal will not correctly reflect the spatial properties of the 624 625 surface.

627 7 Conclusion and perspectives

628 In this study, the Universal Multifractal model was applied to analyze the scaling behaviour of Sentinel-2 surface reflectances and optical indexes acquired over the Southwest region of France, and 629 630 corrected for atmospheric effects by the MAJA processing chain. Images with minimal cloud cover were identified for the whole year 2016 based on the MAJA cloud and shadow detection process. This 631 632 allowed us to access the dynamics of the scaling properties of surface reflectances, which, to our 633 knowledge, had not yet been explored to date due to the significant impact of clouds on optical data. 634 Different scaling regimes, separated by significant scale breaks, were evidenced from these data. 635 Through the use of multifractal simulations, we could explain the scaling properties and especially the 636 position of scale breaks observed in the images studied by two types of phenomena.

637 First, a scale break between hundreds of meters and one kilometre was detected that was found 638 to be impacted mostly by the seasonal and geographic specificities of the surface. From this specific 639 break, two scaling regimes with distinct scaling properties were evidenced. Over large extents, 640 reflectances were found to be poorly autocorrelated, revealing spatial properties quite distant from 641 those commonly detected in the phenomena such as topography, rain and clouds. Over small extents, 642 higher autocorrelation were observed in our data, reflecting the presence of multifractal processes 643 comparable to those in turbulence. In addition, a seasonal cycle of scaling parameters was revealed 644 for the fine-scale regime, while no particular temporal behaviour was noticed for the large-scale 645 regime. We showed that this twofold scaling behaviour is directly linked to the physical properties of 646 the surface, and in particular to the agricultural plots that impact its spatial structure. In this way, the 647 fine-scale regime would reflect the intra-plot variability governed by multifractal physical and 648 biophysical processes related to crops (soil moisture, vegetation, etc.), while the large-scale regime 649 would represent the *extra-plot* variability resulting from the aggregation of several fine-scale processes 650 which together would not have particular multifractal properties. These interpretations were validated 651 by combining our knowledge of the surface and multifractal simulations, which confirmed that the 652 transition scale between the two observed regimes corresponds to the average size of the agricultural 653 plots for this region.

Second, a singular scaling behaviour was systematically detected at high frequencies on the various optical products: a systematic spectral drop was observed near the resolution of the sensor, reflecting an alteration of the signal measured on fine scales, independently of the observed surface type. This phenomenon was found to be an instrumental artefact generated by the Modulation Transfer Function of the optical system which affects the high frequencies of the acquired image. Therefore, this analysis highlighted the performance limits of the Sentinel-2 optical sensor MSI. The particular break detected at fine scales (*i.e.* about 40 m scale) would correspond to the *effective resolution* of the sensor, beyond which the fine-scale fractal properties of the observed surface cannot
 be detected.

663 To complete the interpretation of the spectral drop artifact observed near the resolution of the sensor, the multi-scale model proposed here could be deepened by integrating other factors 664 665 influencing the Modulation Transfer Function such as electric charge displacements or temperature. 666 In order to be in the real conditions of Sentinel-2 data acquisition, the real value of cut-off frequency 667 of the MSI sensor could be used in the MTF model. It would also be interesting to establish precisely 668 why the fine-scale break is located at $\sim 4 l_s$ (mathematical relation) and possibly propose a method to 669 correct this artefact statistically. Here, a fractal stochastic disaggregation could be applied to replace 670 the "impacted" field at fine scales ($l < 4 l_s$) by a simulated multiplicative cascade that would have the 671 scaling properties of the large-scale field $(l > 4 l_s)$. Based on scaling properties, this type of method 672 preserves the probability distribution from large to fine scales. Such fractal-based techniques have 673 been applied on geophysical variables such as rainfall (Rebora et al., 2006; Sharma et al., 2007; Deidda, 674 2000; Gires et al., 2012) or soil moisture (Bindlish and Barros, 2002a; Kim and Barros, 2002b; Mascaro 675 et al., 2010). Therefore, it would be interesting to apply this "stochastic" correction to the S2/MAJA 676 surface reflectances and to verify if the spectral drop is still observed near the resolution of the sensor. 677 However, since the disaggregation is based on a random generator, an ensemble of possible fields can 678 be proposed from just one set of scaling parameters. Therefore, this kind of methodology may not be 679 fully suitable in the case of operational hydro-agricultural applications, in particular for determining 680 the position of the extremes. Another possibility to correct the effects of the optical system would be 681 to estimate the original "non-impacted" image from the MTF model proposed in this study. It would 682 simply consist of reversing equation 8, and then estimating the inverse Fourier transform of the field. 683 This solution would respect the deterministic behaviour of the original signal, giving the correct 684 position of the extremes not only statistically valid but also punctually. The challenge would be to get 685 the exact MTF model corresponding to the MSI sensor. Although several methods have been proposed so far to estimate on-orbit MTF (Lei and Tiziani, 1988; Viallefont-Robinet and Léger, 2010; Xu et al., 686 687 2014), differences persist between empirical and modelled MTF (Gascon et al., 2017). In the case of an 688 improperly parametrized MTF, this methodology could further alter the signal and generate more 689 scaling artifacts.

690 This research work demonstrated the potential of the multifractal analysis for spatial 691 characterization of surface reflectances. The theoretical background on which this analysis is based 692 makes it possible to provide an interpretation that is both different and complementary to the more 693 usual texture analysis methods, thus improving our understanding of surface processes and their 694 spatial heterogeneity. Thanks to the Universal Multifractal model, we were able to identify specific 695 sets of scales ranging from tens of meters to tens of kilometres and on which specific statistical 696 properties were detected. The observed scaling regimes were found to be the results of the 697 combination at different scales of both natural and anthropogenic factors. This spatial characterization 698 could be useful to precisely define the set of scales to consider in information extraction from remote 699 sensing images and help calibrating current surface models which take little account of scaling 700 behaviour, operating on either local or global scales (Ko et al., 2016). It would be possible to verify 701 whether the products synthesized by these models comply with the expected scaling regimes, and if 702 not to constrain the statistics of these models on certain scale ranges. In addition, the present analysis 703 showed that the particular scale ranges detected are directly related to the seasonal (and geographic) 704 specificities of the surface and that these variations should be taken into account in this type of spatial 705 analysis.

706 Finally, this study showed the potential of multifractal analysis to assess the reliability of satellite 707 products to transcribe what was observed by the sensor. Due to its genericity, the method presented 708 here could be applied to other optical satellite data (Landsat-8, Venus), even defined at very high 709 resolution (Ikonos). It could also be used to analyze products from different technologies such as radar 710 (Sentinel-1) or passive microwaves (Neuhauser et al., 2019). From a more operational point of view, 711 this could be useful either for the preparation of future space missions in order to determine the 712 capacity of the sensors to detect the inherent scaling properties of the surface processes, or for post-713 launch assessment to detect in an easy and rapid way the effective resolution of satellite sensors. This 714 methodology complements more usual validation methods enabling a better estimation and 715 understanding of the effective resolution of different products derived from satellite acquisitions.

716

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724 **Appendix A: Theory of Multifractals**

725 During the last century researchers realized that many geophysical processes could present scale 726 invariance properties. This was first anticipated by Richardson (1922) in the case of turbulence: he 727 described turbulent flows as cascade processes that transfer kinetic energy from large to small scales. 728 Based on this approach, statistical models of turbulence were proposed such as the famous 729 Kolmogorov law (1941) to describe velocity increments. Later research generalized the study to take 730 into account the heterogeneity of the energy flux (Kolmogorov, 1962; Oboukhov, 1962; Yaglom, 1966). 731 Multi-scale models such as multiplicative cascades were proposed to reproduce scale invariance 732 properties through the use of fractal geometry. Later, scale invariance was noticed in other geophysical 733 fields. For example, in his study of the coast of Britain, Mandelbrot (1967) revealed the presence of 734 fractal properties in topography.

735

1) From fractal sets to multifractal fields

736 The concept of fractal has been used in many works related to multi-scale analysis and geophysical 737 modelling. Indeed, the term "fractal" refers to any entity (time series or 2D/3D random field) in which 738 each part presents similar properties, geometrically or statistically, to the ensemble. In this manner, 739 the structure of a fractal entity is characterized by scale invariance.

740 Monofractal stochastic models aim to represent the simple scaling behaviour of geophysical 741 processes. In this context, the fractal dimension or scaling parameter is assumed to be unique, 742 restricting multi-scale modelling to a specific class of variability. However, most geophysical processes 743 are characterized by more complex statistics. In case of operational hydrology, rare and extreme 744 events, present in precipitation or soil moisture for example, correspond to high-order statistics and 745 need to be detected (Hubert et al., 1993). Therefore, multifractal models, characterized by an infinite 746 spectrum of fractal dimensions, have been proposed to account for a more exhaustive set of statistics. 747 Schertzer and Lovejoy (1987), based on the findings of Parisi and Frisch (1985), initially established the 748 multifractal formalism through the fundamental equation:

- 749
- 750

 $Pr(\bigoplus_{\lambda} > \lambda^{\gamma}) \approx \lambda^{-c(\gamma)}$

9

751 where \oplus_{λ} is a positive normalized random scalar process, time series or random field defined on R^2 or R^3 . The symbol " \approx " indicates an equality within the limits of slowly varying functions. The mean of 752 the process is assumed to be statistically conserved across scales ($\langle \bigoplus_{\lambda} \rangle = const$). λ is the observation 753 resolution, here defined as the number of intervals (or pixels) along a dimension of space. Note that 754 755 the definition given here is different from that usually encountered in remote sensing. In our case, it is a dimensionless factor that is inversely proportional to the scale *l*. The latter can be seen as the sampling time or pixel size for time and space domain processes respectively (see Section 3.1 for a more detailed definition of scale and resolution concepts).

Eq.9 expresses the fact that for a multifractal process, the probability of exceeding a threshold varies as a power law of the resolution with exponent $c(\gamma)$. This exponent is called as fractal codimension of the process, depending on the amplitude of thresholds. The thresholds are defined by the following power law:

 $T_{\lambda} = \lambda^{\gamma}$

763

764

10

765 with γ the notion of singularity, characterizing the amplitude of the process independently of the 766 resolution. Each singularity is associated with a fractal codimension $c(\gamma)$, corresponding to a measure 767 of the sparseness of the area where the process exceeds the thresholds defined by γ . From a more 768 physical point of view, high singularities (detected by high thresholds) are related to rare and extreme 769 events, with high fractal codimensions and inversely low (box-counting) fractal dimensions $D_f(\gamma)$ 770 (Mandelbrot, 1967). Indeed, the latter are related to the dimension of space D through the relation 771 $c(\gamma) = D - D_f(\gamma)$. Therefore, $c(\gamma)$ can be described as a codimension function, increasing with γ , 772 which completely characterizes the multi-scale statistical properties of the field \oplus_{λ} . In general, if the 773 field is multifractal, $c(\gamma)$ is found to be a convex and positive function (with a fixed point C_1 imposed 774 by the condition of canonical conservation $\langle \oplus_{\lambda} \rangle = const$), whereas monofractality is associated to 775 the trivial case $c(\gamma) = const.$

Since probability distributions and statistical moments are related by a Mellin transform,
 Schertzer and Lovejoy (1987) proposed an equivalent equation to Eq.9:

779 778 $\langle igoplus_{\lambda}^{q}
angle pprox \lambda^{K(q)}$

where $\langle \cdot \rangle$ is the statistical averaging operator, q is the order of the moment ($q \ge 0$), and K(q) is the moment scaling function. Eq.11 expresses that, for any fixed moment order, statistical moments and resolution are linked through a power law. Singularities and moment orders are directly linked, since the moment scaling function K(q) is the Legendre transform of the codimension function $c(\gamma)$. Similarly to $c(\gamma)$, K(q) is a convex function (with the special case K(1) = 0 related to the conservation of the mean across scales), which entirely characterizes the multifractal field.

786 2) Multiplicative cascades

787 Multiplicative cascades are stochastic models that can be used to build multifractal fields. 788 Cascades are multiplicative processes because they are defined by an iterative multiplicative 789 construction, from larger scales to smaller scales. Considering a two dimensional random signal (field), 790 the cascade consists of dividing each pixel at resolution λ_n (with *n* the construction level of the 791 cascade) into sub-pixels at finer resolution λ_{n+1} , and then to multiply the latter by a random variable 792 $\mu \varepsilon$. This is described by the following equation:

794

 $\oplus_{\lambda_{n+1}} = \mu \varepsilon \times \oplus_{\lambda_n}$

In this manner, the statistical properties of the field $\bigoplus_{\lambda_{n+1}}$ are directly related to the statistical properties of the coarser field \bigoplus_{λ_n} . If all the multiplicative random variables used for each step of the iterative construction are independent and identically distributed, and distributed independently of the scale, the final random field (obtained after multiple iterations) presents scale invariant properties.

799 Several models of cascades have been developed so far. First models were built within the 800 framework of turbulence, such as the α -model (Schertzer and Lovejoy, 1984) which corresponds to 801 discrete construction of cascades: the multiplicative random variables are limited to two possible fixed 802 values, respectively leading to increasing or decreasing pixel value when the resolution is refined. 803 Later, more elaborate models were constructed generalizing the discrete case to continuous cascades 804 (Dubrulle, 1994; Schertzer and Lovejoy, 1987, 1991, 1997; She and Levêque, 1994). The latter are based 805 on an infinite number of steps between any pair of resolutions, leading to continuity in scale. The 806 benefit of continuous cascades is twofold. First, they can represent possibly more realistic structures 807 by avoiding any arbitrary discretization of scales. Moreover, they often converge towards random 808 processes that are characterized by a small number of degrees of freedom (special cases of log-809 infinitely divisible distributions). This is interesting considering that multifractal fields built by 810 multiplicative cascade processes would otherwise need an infinite number of scaling parameters (one 811 for each fractal dimension). For example, She and Levêque (1994) proposed the Log-Poisson model 812 which constructs multifractal fields from the generation of the random variable $\mu\varepsilon$ whose distribution 813 law is of log-Poisson type. The (log-) Poisson laws are part of the (log-) infinitely-divisible laws, just as 814 the (log-) stable laws (including the Gaussian case). The Universal Multifractal (UM) model proposed 815 by Schertzer and Lovejoy (1987) uses random variables $\mu\epsilon$ of log-stable laws to generate the 816 multifractal field (see Appendix D for more details on the generation of such a field). In both models, 817 only two fundamental parameters are needed to fully define multifractality, i.e. to fully characterize 818 the $c(\gamma)$, K(q) functions.

819 3) Universal Multifractals

820 Physically, multifractal fields built by Log-Poisson or Log-stable cascades are found in many 821 geophysical applications. The Log-Poisson model has been successfully applied to different geophysical 822 variables such as rain (Deidda, 2000), or even soil moisture (Mascaro et al., 2010). However, the Log-823 Poisson model can have the disadvantages of representing a somewhat restricted range of variabilities, 824 since they are not suitable for modeling processes with unbounded singularities. On the other hand, 825 by assuming the stability of the random variables and suitable renormalization, the UM model is likely 826 adaptable for characterizing a wide range of processes: topography (Lavallée et al., 1993), rain and 827 clouds (Tessier et al., 1993) and more recently soil moisture and vegetation optical indexes (Lovejoy et 828 al., 2008b). Moreover, a possibly more immediate physical interpretation of the parameters is found 829 in this model. For mathematical and physical arguments supporting the universality of the UM model, 830 see Schertzer and Lovejoy (1997); see also Gupta and Waymire (1997) for discussion about its 831 generality. The UM model defines the moment scaling function using two "universal" parameters, 832 through the following equation (Schertzer and Lovejoy, 1987):

$$K(q) = \frac{C_1}{(\alpha - 1)}(q^\alpha - q)$$
13

834 833

where α is the degree of multifractality of the field. It varies between 0 (monofractality) and 2 (lognormality) and expresses how fast the codimension evolves as a function of the singularity. The second parameter C_1 is the codimension giving the dominant contribution to the mean value of the field (related to moment of order 1): $C_1 = K'(1)$. Physically, it indicates inhomogeneity (dispersion) of the field: it varies from 0 (homogeneous field) to the dimension D of the embedding space (very intermittent field). Because of the Legendre transform, $c(C_1) = K'(1)$ is also defined as the fixed point of the codimension function.

842 4) *FIF model*

843 Generally, most geophysical fields are "non-conservative", i.e. integrated processes defined by a 844 certain degree of fractional integration (classical integration with non-integer orders). This appellation 845 comes from multifractal cascade models: see Gagnon et al., 2006 for detailed explanations on the 846 notion of fractional integration and its use in 2D multifractal models. Thus, to account for a wider range 847 of processes (like filtered fields caracterized by steeper spectra), an extension of the UM model to non-848 conservative fields has been proposed (Schertzer and Lovejoy, 1991): the Fractionally Integrated Flux 849 (FIF) model. It expresses the degree of fractional integration of the UM field using a third parameter, 850 H. The latter is called the order of integration and defines the non-conservativity of the field: in plain 851 words, the larger is H, the smoother is the field. The integrated flux is noted R_{λ} and is characterized 852 by a power law variation of its stationary increments:

$$\Delta R_{\lambda} \approx \bigoplus_{\lambda} \Delta x^{H}$$
853 14

853

855 where ΔR_{λ} are the increments (fluctuations of the flux) estimated over a varying window Δx , which is equivalent to the spatial scale l. Note that when H = 0, the equation corresponds to the conservative 856 case \bigoplus_{λ} , and that when H < 0 the equation corresponds to the specific "fractional derivation" non-857 858 conservative case. Additionally, in the case of two dimensional fluxes, Eq.14 also applies for other 859 directions (*i.e.* Δy increments), with the same exponent *H* if the process is isotropic.

860 Appendix B: Estimation of Structure Functions and Statistical Moments

861

1) Structure functions: some evidence of non-conservativity

The structure functions were initially used in turbulence by Kolmogorov (1941) to highlight the 862 scaling behaviour of velocity increments. Then, these methods were widely applied in geophysics, in 863 864 particular thanks to their simplicity of construction and interpretation. Indeed, unlike the spectrum 865 which implies a transform in the Fourier domain, the structure functions simply represent the 866 fluctuations of the observations in real space, as a function of their distance. In addition, these 867 methods allow direct access to the degree of smoothness (or degree of integration) of the data, also 868 called in the multifractal formalism the "non-conservative" nature of the data studied. This is possible through the estimation of the parameter H that characterizes any integrated multifractal data R_{λ} . In 869 870 the multifractal formalism, R_{λ} is defined by the power law variation of its stationary increments described in Appendix A (Eq.14). 871

872 There are several ways to estimate these fluctuations. The simplest is to calculate the average of 873 the absolute differences (or gradient) of R_{λ} . In general, this method is very suitable for processes 874 characterized by a parameter H between 0 and 1 (Lovejoy and Schertzer, 2012a), thus corresponding 875 to increasing fluctuations with the distance Δx . However, one constraint with this definition is that it 876 does not allow decreasing fluctuations to be represented with distance: for multifractal data defined 877 by H < 0, the structure function converges to a constant "parasitic" value, which gives an inaccurate 878 estimate of H (which would be considered equal to 0). To be able to extend the range of H to which 879 the structure functions have access, one solution is to change the shape of the wave estimating the 880 fluctuations. Lovejoy and Schertzer (2012b) proposed a simple formalism of definitions of fluctuations 881 which is based on that of wavelets. In particular, Haar's wavelet is a definition that turns out to be 882 simple to calculate and interpret, while allowing access to negative H values. The fluctuations 883 estimated from this method are calculated according to the following equation:

884
$$\Delta R_{\lambda}(\Delta x) = \left\langle \frac{2}{\Delta x} \left| S(x) + S(x + \Delta x) - 2S\left(x + \frac{\Delta x}{2}\right) \right| \right\rangle$$

885

886 where x and Δx are vectors and $\langle \cdot \rangle$ corresponds to the empirical average. S is the cumulative sum of 887 the integrated data R_{λ} from which the global average has been subtracted. Thus, although this 888 definition seems complex at first glance, it simply consists in estimating the average fluctuations from 889 x to $x + \frac{\Delta x}{2}$ and from $x + \frac{\Delta x}{2}$ to $x + \Delta x$. We can note that, from an algorithmic point of view, equation 890 8 involves aggregating the field at half the resolution as to allow the calculation of $\frac{\Delta x}{2}$. This implies that, 891 in the case of Haar wavelets, the structure function will not be defined for the finest observation scale.

This technique was used in this study to estimate the structure functions and the corresponding Hparameters of our surface S2 products. In particular, in the case of 2D processes, fluctuations were estimated on vertical Δx and horizontal Δy distances separately, allowing us to obtain H_x and H_y parameters respectively. If the observations present indeed multifractal properties and are nonconservative in nature, the order of integration H_x (H_y) should be the slope of the increments ΔR_{λ} , plotted in a log-log graph as a function of space scale Δx (Δy).

898 2) Statistical moments: multifractal properties

Statistical moments analysis deepens and extends spectral or structure functions analysis by focusing on the extreme variability of the signal. By comparison with the spectral analysis which represents the evolution of a single order of statistics, this method analyzes the behaviour of a set of statistical moments on the range of scales studied.

903 To test the presence of multifractal properties in the data (Eq.11), statistical moments and 904 moment scaling function need to be estimated. To do this, different steps must be followed. First, the 905 underlying flux $\bigoplus_{\lambda_{max}}$ has to be reconstructed from the data, *i.e.* random field *R*, at the maximum observation resolution λ_{max} . Because of the possible existence of a fractional integration of order H 906 907 (Eq.14), a fractional derivative of the same order should be done. In this study, the modulus of the 908 gradient was applied to the data. Indeed, this operator provides a simple and good numerical 909 approximation of the fractional derivation without prior knowledge of H order (Lavallée et al., 1993; 910 see also Renosh, 2015 for discussion about this approximation):

911
$$\bigoplus_{\lambda_{max}} = \sqrt{\left(\frac{\partial R_{\lambda_{max}}}{\partial x}\right)^2 + \left(\frac{\partial R_{\lambda_{max}}}{\partial y}\right)^2}$$

912

16

913 Once the flux (conservative field) is retrieved, $\phi_{\lambda_{max}}$ is normalized by its mean.

The second step involves the degradation of the flux at lower resolutions $\lambda < \lambda_{max}$. It aims to approximate the inversion of the stochastic multiplicative cascade by iteratively averaging the flux at coarser scales: each coarse pixel (level *n* of the cascade) is obtained by a simple average of neighboring finer pixels (level *n* + 1). Note that each roughened pixel size is a power of two multiplied greater than the observation scale l_{min} (= $\frac{L}{\lambda_{max}}$).

Finally, empirical moments (*i.e.* computing *q*-th order moments in Eq.11 while replacing statistical
averages by empirical averages) are then computed for various orders and all accessible resolutions.
Because of sample size limitations and/or possible divergence of high-order moments (Hubert et al.,

922 2007), in this study moments were computed for orders set from 0 to 3, by steps of 0.1. On a log-log 923 graph, the different moments are plotted as a function of the resolution. If linearity is observed for 924 each moment curve, at least over a significant range of resolutions, Eq.11 is therefore verified, which 925 is the signature of multifractality. The empirical moment scaling function can be estimated, with K(q)926 corresponding to the slope of each linear fit of *q*th order moment. Afterwards, the universal 927 parameters α and C_1 may be obtained by optimization: here the Nelder-Mead method is used to find 928 the couple of parameters minimizing the (quadratic) distance between the empirical function K(q)929 and the UM model form (Eq.13).



Figure 13: Mean Haar structure functions estimated from Sentinel-2 bands and indexes, over the period January-December 2016. For each band, structure functions are computed in two orthogonal directions (X and Y axis of the image).



Figure 14: Mean moments estimated from Sentinel-2 bands and indexes, over the period January-December 2016. The straight lines in red and green correspond to the linear regressions carried out respectively on range I (large scales) and range II (small scales).



Figure 15: Scaling functions estimated from the mean statistical moments in figure 14, for range I (a) and range II (b). Note the different range of values on the vertical axes.

935 Appendix D: Simulation of multifractal random fields

The simulator that we present here is used to generate multifractal random fields knowing the three parameters α , C_1 and H of the UM model (FIF version) as well as the size of the simulation domain corresponding to the maximum resolution. This simulator includes several steps, which are described in detail in various works on multifractal processes such as Lovejoy and Schertzer (2010b, 2010c), Pecknold et al. (1993), and Verrier (2011). We therefore invite the reader to refer to this work for more information on these processing steps.

In general, this simulator aims to reproduce the mathematical construction of the multiplicative
cascades corresponding to the UM/FIF model, defined in Appendix A. This construction is based on five
main steps that we describe hereafter:

945 1) Generation of the log-stable random variable Υ_{α} :

The construction of the UM model is based on the generation of a log-stable random variable, 946 947 corresponding here to an α -stable white noise Υ_{α} . The parameter α designates the degree of 948 multifractality of the UM model, and also corresponds to the stability index relating to stable laws 949 (Schertzer et al., 2002). Thus, a first step consists in generating Υ_{α} thanks to the algorithm of Chambers 950 et al. (1976). This algorithm makes it possible to simulate stable variables quickly and efficiently, simply 951 from realizations of independent random variables and according to uniform and exponential laws. 952 Note here that some additional constraints must be considered to generate Υ_{α} , relatively to its 953 normalisation (unit asymmetrical α-stable white noise) and its asymmetry (asymmetry parameter set 954 to -1).

955

2) Multiplication by the amplitude factor σ :

956 The random variable Υ_{α} is then multiplied by an amplitude factor σ allowing to integrate the 957 parameter C_1 in the construction of the cascade:

958
$$\sigma = \left(\frac{c_1}{|\alpha - 1|}\right)^{\frac{1}{\alpha}}$$

959

17

960 3) Generator creation Γ_{λ} :

961 In order to obtain the multifractal field, we must first construct the random variable r_{λ} called 962 "generator". This is obtained by convolution of the log-stable variable σ . Υ_{α} and a filter g_{α} (in negative 963 power law of the distance) such that:

964
$$g_{\alpha}(\vec{x}) = \frac{1}{\|\vec{x}\|^{\frac{D}{\alpha}}}$$

965

where $\|\cdot\|$ designates the Euclidian norm, *i.e.* in the 2D case $\|\vec{x}\| = \|(x, y)\| = \sqrt{x^2 + y^2}$. For the sake of speed, the convolution operations present in this simulator are performed numerically in the Fourier domain by means of the two-dimensional *fft* or "Fast Fourier Transform". *D* is here taken equal to 2 for a two-dimensional simulation.

970

4) Generation of the conservative process
$$\oplus_{\lambda}$$
 :

971 The conservative process (or flux) is finally obtained by exponentiation of the generator, then 972 normalization:

974

975 5) Generation of the integrated process R_{λ} :

976 In the case of the simulation of a non-conservative process (integrated process named here as 977 "field"), a last step consists in applying a convolution to \bigoplus_{λ} by the filter in decreasing power law g_H , 978 defined by :

979
$$g_H(\vec{x}) = \frac{1}{\|\vec{x}\|^{D-H}}$$

980

This convolution amounts to carrying out a fractional integration of order *H* (less than *D*) (Gagnon et al., 2006). In the particular case where H = 0, step 5 is omitted since here the random process to be generated is conservative, that is to say $R_{\lambda} = \bigoplus_{\lambda}$.

984

18

19

985 <u>Appendix E</u>: Modelling of the Modulation Transfer Function of an optical sensor from the
 986 characteristics of the lens (diffraction) and detector (sampling)

987
988
$$MTF = MTF_{diffract} \times MTF_{sampling}$$

990 $MTF_{diffract} = \frac{2}{n} \left\{ \cos^{-1} \left(\frac{f}{f_c} \right) - \left(\frac{f}{f_c} \right) \sqrt{1 - \left(\frac{f}{f_c} \right)^2} \right\}$
91
92
93
94
94
95
96
97
98
996
1001 $MTF_{sampling} = sinc \left(\prod \frac{f_x}{f_s} \right) . sinc \left(\prod \frac{f_y}{f_s} \right)$
101 $MTF_{sampling} = sinc \left(\prod \frac{f_x}{f_s} \right) . sinc \left(\prod \frac{f_y}{f_s} \right)$
102
103
1004

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1006

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