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# Exploiting the Kinematic Redundancy of a (6+3)-DoF Parallel Manipulator to Produce Unlimited Rotation of the Platform 

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#### Abstract

Mechanical interference and singularities within the reachable workspace often restrict the orientational workspace of parallel robots. Introducing kinematic redundancy can alleviate this limitation. This paper discusses the possibility to produce unlimited rotation of the platform of a tripedal $(6+3)$-degree-of-freedom kinematically redundant parallel robot. The articulated platform of such a robot has three degrees of mobility. The platforms considered here are planar linkages that contain either revolute or prismatic joints. It is shown that at least two revolute joints are required to produce unlimited rotation with appropriate design and initial configuration, while the platforms with two prismatic joints cannot produce such rotations without crossing a singularity.


## 1 INTRODUCTION

Mechanical interference and type II (or parallel) singularities within the reachable workspace often limit the orientational capabilities of parallel robots. A typical example is the tilt angle of a Gough-Stewart platform which is limited to approximately $45^{\circ}$ [1]. While the size of the translational workspace of a robot can be increased

[^0]by scaling up the geometric parameters, the orientational workspace is scale invariant. In particular, designing a parallel robot capable of producing an unlimited rotation of the platform is a challenging problem. Existing solutions are lower mobility parallel manipulators, often providing Schöenflies motions (motions of the SCARA robot). Schöenflies motion parallel robots can achieve unlimited rotation by using kinematic redundancy [2]. Nonredundant designs also exist, but suffer from a limited workspace [3] or use motion transformation, which deteriorates the efficiency and the stiffness [4].

Typical applications of parallel robots usually do not require large ranges of orientations. However, significant orientational capabilities may be necessary for manipulation or advanced grasping techniques (e.g., scooping [5]). Introducing kinematic redundancy can alleviate this limitation (see, for instance, $[6,7,8,9,10]$ ). The reader may refer to recent reviews on kinematic redundancy for a better overview [11, 1]. As notable examples, kinematic redundancy was used (i) on a planar robot to obtain unlimited rotation and also to actuate a gripper at the endeffector using the redundant degree of freedom (DoF) [6] and (ii) on a spatial hybrid parallel $(6+3)$-DoF robot with very large - but not unlimited, because of mechanical
interference - orientational workspace and a remotely operated gripper [10]. Other 3-DoF planar platforms that can be used on the $(6+3)$-DoF robot mentioned previously were also studied [12].

The redundant DoFs must be controlled to avoid singular configurations [13]. Singularity avoidance on robots is a well-known topic that has been researched for several decades. Yoshikawa introduces and uses the manipulability measure in a multi-objective optimization problem for the trajectory control of a robot while avoiding singularities [14]. Klein and Blaho compare different dexterity measures for the design and control of manipulators with kinematic redundancy [15]: the manipulability measure, the condition number, the minimum singular value and the joint range availability [16]. Voglewede and Ebert-Uphoff give a general form for a singularity index that is equivalent to the solution of a generalized eigenvalue problem [17]. The generalized eigenvalue problem can be formulated such that it has a physical meaning. The physical meaning can be related to different quantities, such as the lowest natural frequency of the manipulator [17], the power developed by each leg [18, 19] or the potential energy associated to the stiffness of each actuator. Performance indices based on the static behaviour of a parallel manipulator have also been proposed [20, 21].

In the above-mentioned work by Wen et al. [10], the authors show that all type II singularities can be avoided by restricting the redundant joint angles to $\pm 60^{\circ}$ if the proximal links (with respect to the platform) are short enough to avoid any mutual intersection. While convenient, this method is nevertheless restrictive since (i) it prevents large ranges of motion and (ii) it limits the moment of the force generated by the leg that is proportional to the length of the proximal link.

In this work, we revisit the $(6+3)$-DoF tripedal parallel robots with planar configurable platforms proposed in $[10,12]$ in order to synthesize kinematic architectures that can perform unlimited rotation at the endeffector. Here, unlimited rotation means that the platform is able to rotate for any number of turns about some axis. To the best of our knowledge, this work presents the first 6-DoF parallel robot that can achieve unlimited torsion. Each leg of the parallel robot consists of a 3DoF robot controlling the Cartesian position of its endeffector. The legs are attached to the configurable platform through spherical joints. The platform has three degrees of mobility obtained by combining revolute and prismatic joints. First, the kinematic modelling of a platform with three revolute joints is presented (Section 2) and the singularities are analyzed (Section 3). Then, two different singularity indices are compared (Section 4). A path planning approach is presented in order to perform


Fig. 1. Hybrid $(6+3)$-DoF robot (adapted from Wen et al. [10]). The actuated joints are highlighted in one of the legs.
an unlimited rotation of the platform without crossing any singularity, assuming that one of the degrees of mobility of the platform is constrained (Section 5). The feasibility of the unlimited rotation is then discussed (Section 6). The influence of some geometric parameters is analyzed in order to guide the design of the platform (Section 7). Finally, two other platforms are explored, one with two revolute joints and one prismatic joint, the other one with one revolute joint and two prismatic joints (Section 8). The video accompanying the paper illustrates a remotely actuated gripper with unlimited rotation (https://youtu.be/D5txxesP51w).

## 2 MODELLING

### 2.1 Parameterization

Figure 1 (adapted from [10]) illustrates a $(6+3)$ DoF hybrid parallel robot. According to the nomenclature introduced by Gosselin and Schreiber in [7], the extra three DoFs correspond to the possible reorientation (or repositioning) of the links connecting the legs to the end-effector. Each of the three legs includes a fixed actuated revolute joint on which a planar five-bar linkage actuated by two base joints is mounted. The end of each leg consists of a passive spherical joint and the platform includes three passive revolute joints with parallel axes. Each leg includes three actuators located near the base. A fixed reference frame $\mathcal{R}_{f}=\left(O, \boldsymbol{x}_{f}, \boldsymbol{y}_{f}, \boldsymbol{z}_{f}\right)$ is attached to the base of the robot and a mobile reference frame $\mathcal{R}_{b}=\left(P, \boldsymbol{x}_{b}, \boldsymbol{y}_{b}, \boldsymbol{z}_{b}\right)$ is attached to the platform such that vector $\boldsymbol{z}_{b}$ is normal to the planar platform.

Figure 2 shows some geometric parameters of the planar platform. The platform is equipped with three revolute joints forming the vertices of a triangle. The posi-


Fig. 2. Parameterization of the platform.
tion vector of vertex $T_{i}$ in $\mathcal{R}_{b}$ is noted $\boldsymbol{t}_{i}$. The norm of $\boldsymbol{t}_{i}$ and its orientation angle with respect to $\boldsymbol{x}_{b}$ are noted respectively $L_{i}$ and $\phi_{i}$. With no loss of generality, it is assumed that $\phi_{1}=0$. A spherical joint of centre $S_{i}$ is attached to the revolute joint at $T_{i}$ through a link of length $l_{i}$. These spherical joints correspond to the end-effector of each leg of the tripedal parallel robot. The position of $S_{i}$ in $\mathcal{R}_{b}$ is noted $\boldsymbol{s}_{i}$, and we define $\boldsymbol{u}_{i}=\boldsymbol{s}_{i}-\boldsymbol{t}_{i}$. The angle from $\boldsymbol{t}_{i}$ to $\boldsymbol{u}_{i}$ is noted $\psi_{i}$. The position of the moving frame $\mathcal{R}_{b}$ with respect to a fixed base frame $\mathcal{R}_{f}$ is noted $p$, its orientation with respect to $\mathcal{R}_{f}$ is given by matrix $\boldsymbol{Q}$ and its angular velocity is noted $\omega$.

### 2.2 Jacobian Matrices

The Jacobian matrices $\boldsymbol{J}$ and $\boldsymbol{K}$ describe the relationship between the end-effector velocities $\boldsymbol{v}=$ $\left[\begin{array}{ll}\dot{\boldsymbol{p}}^{T} & \boldsymbol{\omega}^{T}\end{array}\right]^{T}$ and the actuated joint velocities $\dot{\boldsymbol{s}}=$ $\left[\begin{array}{lll}\dot{\boldsymbol{s}}_{1}^{T} & \dot{\boldsymbol{s}}_{2}^{T} & \dot{\boldsymbol{s}}_{3}^{T}\end{array}\right]^{T}:$

$$
\begin{equation*}
\boldsymbol{J} \boldsymbol{v}=\boldsymbol{K} \dot{\boldsymbol{s}} \tag{1}
\end{equation*}
$$

where $\dot{s}_{i}$ is the three-dimensional vector containing the actuated joint velocities of the $i$-th leg of the robot.

In the following, we use screw theory to find the expressions of $\boldsymbol{J}$ and $\boldsymbol{K}$. Let $\xi_{E E}=\left[\boldsymbol{\omega}^{T}(\dot{\boldsymbol{p}}+\boldsymbol{p} \times \boldsymbol{\omega})^{T}\right]_{O}^{T}$ be the twist of the platform expressed with respect to the base frame origin $O, \xi_{P_{i}}$ the twist of the $i$-th leg, $\xi_{S_{i}}$ the twist of the spherical joint at $S_{i}$ and $\xi_{T_{i}}$ the twist of the revolute joint at $T_{i}$. Then,

$$
\begin{equation*}
\xi_{E E}=\xi_{P_{i}}+\xi_{S_{i}}+\xi_{T_{i}} \tag{2}
\end{equation*}
$$

Let $\xi_{S_{i}, \boldsymbol{u}_{i}}$ be a zero-pitch twist of line $\left(S_{i}, \boldsymbol{u}_{i}\right)$ and amplitude $\left\|\boldsymbol{u}_{i}\right\|=l_{i}$. Then, $\xi_{S_{i}, \boldsymbol{u}_{i}}$ is reciprocal to $\xi_{S_{i}}$ and $\xi_{T_{i}}$, and

$$
\begin{equation*}
\xi_{S_{i}, \boldsymbol{u}_{i}} \cdot \xi_{E E}=\xi_{S_{i}, \boldsymbol{u}_{i}} \cdot \xi_{P_{i}} \tag{3}
\end{equation*}
$$

which gives

$$
\begin{equation*}
\boldsymbol{u}_{i}{ }^{T} \dot{\boldsymbol{p}}+\left(\boldsymbol{t}_{i} \times \boldsymbol{u}_{i}\right)^{T} \boldsymbol{\omega}=\boldsymbol{u}_{i}{ }^{T} \dot{\boldsymbol{s}}_{i} . \tag{4}
\end{equation*}
$$

Let $\xi_{S_{i}, \boldsymbol{z}_{b}}$ be a zero-pitch twist of line $\left(S_{i}, \boldsymbol{z}_{b}\right)$ and amplitude $\left\|\boldsymbol{z}_{b}\right\|=1$. Then, $\xi_{S_{i}, \boldsymbol{z}_{b}}$ is reciprocal to $\xi_{S_{i}}$ and $\xi_{T_{i}}$ (since the axis of the revolute joint at $T_{i}$ is also in the direction of $\boldsymbol{z}_{b}$ ) and

$$
\begin{equation*}
\xi_{S_{i}, \boldsymbol{z}_{b}} \cdot \xi_{E E}=\xi_{S_{i}, \boldsymbol{z}_{b}} \cdot \xi_{P_{i}} \tag{5}
\end{equation*}
$$

which gives

$$
\begin{equation*}
\boldsymbol{z}_{b}{ }^{T} \dot{\boldsymbol{p}}+\left(\boldsymbol{s}_{i} \times \boldsymbol{z}_{b}\right)^{T} \boldsymbol{\omega}=\boldsymbol{z}_{b}^{T} \dot{\boldsymbol{s}}_{i} . \tag{6}
\end{equation*}
$$

Combining (4) and (6) yields the expressions of $\boldsymbol{J}$ and $\boldsymbol{K}$

$$
\left[\begin{array}{l}
\boldsymbol{u}_{1}{ }^{T}\left(\boldsymbol{t}_{1} \times \boldsymbol{u}_{1}\right)^{T}  \tag{7}\\
\boldsymbol{u}_{2}{ }^{T}\left(\boldsymbol{t}_{2} \times \boldsymbol{u}_{2}\right)^{T} \\
\boldsymbol{u}_{3}^{T}\left(\boldsymbol{t}_{3} \times \boldsymbol{u}_{3}\right)^{T} \\
\boldsymbol{z}_{b}^{T}\left(\boldsymbol{s}_{1} \times \boldsymbol{z}_{b}\right)^{T} \\
\boldsymbol{z}_{b}^{T}\left(\begin{array}{l}
\boldsymbol{s}_{2} \times
\end{array} \boldsymbol{z}_{b}\right)^{T} \\
\boldsymbol{z}_{b}^{T} \\
\left(\boldsymbol{s}_{3} \times \boldsymbol{z}_{b}\right)^{T}
\end{array}\right]\left[\begin{array}{c}
\dot{\boldsymbol{p}} \\
\boldsymbol{\omega}
\end{array}\right]=\left[\begin{array}{ccc}
\boldsymbol{u}_{1}^{T} & \mathbf{0}^{T} & \mathbf{0}^{T} \\
\mathbf{0}^{T} & \boldsymbol{u}_{2}^{T} & \mathbf{0}^{T} \\
\mathbf{0}^{T} & \mathbf{0}^{T} & \boldsymbol{u}_{3}^{T} \\
\boldsymbol{z}_{b}^{T} & \mathbf{0}^{T} & \mathbf{0}^{T} \\
\mathbf{0}^{T} & \boldsymbol{z}_{b}^{T} & \mathbf{0}^{T} \\
\mathbf{0}^{T} & \mathbf{0}^{T} & \boldsymbol{z}_{b}^{T}
\end{array}\right]\left[\begin{array}{l}
\dot{\boldsymbol{s}}_{1} \\
\dot{\boldsymbol{s}}_{2} \\
\dot{\boldsymbol{s}}_{3}
\end{array}\right],
$$

where $\mathbf{0}^{T}$ stands for the zero three-dimensional line vector.

## 3 SINGULARITY ANALYSIS

Singularities restrict the orientational workspace of parallel robots. Their analysis is necessary in order to enable unlimited rotation of the platform.

### 3.1 Type I Singularities

The robot is in a type I singularity if the Jacobian matrix $\boldsymbol{K}$ is rank-deficient: $\operatorname{rank}(\boldsymbol{K})<6$. Since $\boldsymbol{u}_{i}$ and $\boldsymbol{z}_{b}$ are always linearly independent ( $\boldsymbol{u}_{i}$ has no component along the $\boldsymbol{z}_{b}$ axis), matrix $\boldsymbol{K}$ is trivially full row rank. Therefore, the platform has no type I singularity.

### 3.2 Type II Singularities

The robot is in a type II singularity if the Jacobian matrix $\boldsymbol{J}$ is singular, i.e., $\operatorname{det}(\boldsymbol{J})=0$. By noticing that (i) $\boldsymbol{u}_{i}$ and $\boldsymbol{s}_{i} \times \boldsymbol{z}_{b}$ have no component along $\boldsymbol{z}_{b}$ and (ii) $\boldsymbol{t}_{i} \times \boldsymbol{u}_{i}$ is collinear to $\boldsymbol{z}_{b}$, the singularity condition can be written as the following: $\boldsymbol{J}$ is singular if and only if either $\boldsymbol{J}_{z}$ or $\boldsymbol{J}_{u}$ is singular, with

$$
\boldsymbol{J}_{z}=\left[\begin{array}{l}
1\left(\boldsymbol{s}_{1} \times \boldsymbol{z}_{b}\right)^{T} \boldsymbol{x}_{b}\left(\boldsymbol{s}_{1} \times \boldsymbol{z}_{b}\right)^{T} \boldsymbol{y}_{b}  \tag{8}\\
1\left(\boldsymbol{s}_{2} \times \boldsymbol{z}_{b}\right)^{T} \boldsymbol{x}_{b}\left(\boldsymbol{s}_{2} \times \boldsymbol{z}_{b}\right)^{T} \boldsymbol{y}_{b} \\
1\left(\boldsymbol{s}_{3} \times \boldsymbol{z}_{b}\right)^{T} \boldsymbol{x}_{b}\left(\boldsymbol{s}_{3} \times \boldsymbol{z}_{b}\right)^{T} \boldsymbol{y}_{b}
\end{array}\right]
$$

and

$$
\boldsymbol{J}_{u}=\left[\begin{array}{l}
\boldsymbol{u}_{1}^{T} \boldsymbol{x}_{b} \boldsymbol{u}_{1}^{T} \boldsymbol{y}_{b}\left(\boldsymbol{t}_{1} \times \boldsymbol{u}_{1}\right)^{T} \boldsymbol{z}_{b}  \tag{9}\\
\boldsymbol{u}_{2}^{T} \boldsymbol{x}_{b} \boldsymbol{u}_{2}^{T} \boldsymbol{y}_{b}\left(\boldsymbol{t}_{2} \times \boldsymbol{u}_{2}\right)^{T} \boldsymbol{z}_{b} \\
\boldsymbol{u}_{3}^{T} \boldsymbol{x}_{b} \boldsymbol{u}_{3}^{T} \boldsymbol{y}_{b}\left(\boldsymbol{t}_{3} \times \boldsymbol{u}_{3}\right)^{T} \boldsymbol{z}_{b}
\end{array}\right]
$$

### 3.2.1 Determinant 1

Matrix $\boldsymbol{J}_{z}$ can be rewritten as

$$
\boldsymbol{J}_{z}=\left[\begin{array}{c}
1 \boldsymbol{s}_{1}{ }^{T} \boldsymbol{y}_{b}-\boldsymbol{s}_{1}^{T} \boldsymbol{x}_{b}  \tag{10}\\
1 \boldsymbol{s}_{2}{ }^{T} \boldsymbol{y}_{b}-\boldsymbol{s}_{2}^{T} \boldsymbol{x}_{b} \\
1 \\
1
\end{array} \boldsymbol{s}_{3}^{T} \boldsymbol{y}_{b}-\boldsymbol{s}_{3}^{T} \boldsymbol{x}_{b} .\right] .
$$

The Laplace expansion with respect to the first column yields

$$
\begin{align*}
\operatorname{det}\left(\boldsymbol{J}_{z}\right)=\left|\begin{array}{ll}
\boldsymbol{s}_{2}^{T} \boldsymbol{x}_{b} \boldsymbol{s}_{2}^{T} \boldsymbol{y}_{b} \\
\boldsymbol{s}_{3}^{T} \boldsymbol{x}_{b} & \boldsymbol{s}_{3}^{T} \boldsymbol{y}_{b}
\end{array}\right| & -\left|\begin{array}{ll}
\boldsymbol{s}_{1}{ }^{T} \boldsymbol{x}_{b} & \boldsymbol{s}_{1}^{T} \boldsymbol{y}_{b} \\
\boldsymbol{s}_{3}^{T} \boldsymbol{x}_{b} & \boldsymbol{s}_{3}^{T} \boldsymbol{y}_{b}
\end{array}\right| \\
& +\left|\begin{array}{ll}
\boldsymbol{s}_{1}^{T} \boldsymbol{x}_{b} & \boldsymbol{s}_{1}^{T} \boldsymbol{y}_{b} \\
\boldsymbol{s}_{2}^{T} \boldsymbol{x}_{b} & \boldsymbol{s}_{2}^{T} \boldsymbol{y}_{b}
\end{array}\right| \tag{11}
\end{align*}
$$

which is equivalent to

$$
\begin{equation*}
\operatorname{det}\left(\boldsymbol{J}_{z}\right)=\left[\boldsymbol{s}_{1} \times \boldsymbol{s}_{2}-\boldsymbol{s}_{1} \times \boldsymbol{s}_{3}+\boldsymbol{s}_{2} \times \boldsymbol{s}_{3}\right]^{T} \boldsymbol{z}_{b} \tag{12}
\end{equation*}
$$

Grassmann line geometry gives a geometric interpretation: the singularities of matrix $J_{z}$ correspond to the alignment of $S_{i}$ [10]. As a consequence, matrix $J_{z}$ cannot be singular if $l_{i}$ is chosen small enough to prevent the alignment of joints $S_{i}$.

### 3.2.2 Determinant 2

After noticing that $\left(\boldsymbol{t}_{i} \times \boldsymbol{u}_{i}\right)^{T} \boldsymbol{z}_{b}=L_{i} l_{i} \sin \psi_{i}$, the Laplace expansion of $\boldsymbol{J}_{u}$ with respect to the third column
yields

$$
\begin{align*}
& \operatorname{det}\left(\boldsymbol{J}_{u}\right)=[ L_{1} l_{1} \\
& \sin \psi_{1}\left(\boldsymbol{u}_{2} \times \boldsymbol{u}_{3}\right) \\
&-L_{2} l_{2} \sin \psi_{2}\left(\boldsymbol{u}_{1} \times \boldsymbol{u}_{3}\right)  \tag{13}\\
&\left.+L_{3} l_{3} \sin \psi_{3}\left(\boldsymbol{u}_{1} \times \boldsymbol{u}_{2}\right)\right]^{T} \boldsymbol{z}_{b}
\end{align*}
$$

According to Grassmann line geometry, $\boldsymbol{J}_{u}$ is singular if and only if lines $\left(S_{i}, \boldsymbol{u}_{i}\right)$ have a common intersection or are parallel [10].

## 4 SINGULARITY INDEX

One possible approach to produce an unlimited rotation of the end-effector is to couple this rotation to one of the redundant links. For instance, if $\psi_{1}$ is performing an unlimited rotation and $\psi_{2}$ is constant, a singularity avoidance strategy is essential. Indeed, if $\psi_{3}$ is also constant, then at some point of the rotation of $\psi_{1}, \boldsymbol{J}_{u}$ will be singular. A possible solution is to choose a singularity index and control the orientation of the redundant links by maximizing the index. A singularity index is a positive semidefinite function $\mathcal{I}$ that vanishes only at singularities. Intuitively, it represents the "distance" to the singular configurations. However, there is no universal definition of the distance, since the Euclidean distance has no physical interpretation in most cases. In the following, we recall two singularity indices that are well known in the literature and compare their impact on joint forces for the manipulator considered here. These singularity indices are based on the Jacobian matrices from (1). These Jacobian matrices are independent from the pose (position and orientation) of the end-effector. This is clearly seen from (7). This is explained by the fact that the kinematics (1) are agnostic of the kinematics of the leg mechanism. Indeed, this result is generic (any serial or parallel manipulator that controls the Cartesian position of its endeffector can be used as a leg), but maintains the leg within a workspace far from singularities.

The analysis from Section 3 shows that type I singularities depend on the leg mechanisms and can even be completely avoided by using singularity-free 3-DoF robots. In order to stay as generic as possible, type I singularities are not considered here.

### 4.1 Maximum Actuated Joint Forces

In the neighbourhood of a type II singular configuration, the joint forces can become very large for a given external wrench applied to the platform. For a given external wrench $f$ acting on the platform, the actuated joint
forces (or torques) $\boldsymbol{f}_{j}$ can be obtained using the principle of virtual work as

$$
\begin{equation*}
\boldsymbol{f}_{j}=\left(\boldsymbol{J}^{-1} \boldsymbol{K}\right)^{T} \boldsymbol{f} \tag{14}
\end{equation*}
$$

The singularity index based on the maximum actuated joint forces, noted $\mathcal{I}_{n}$, is defined as the inverse of the maximum value of the actuated joint forces for any external wrench in some set $\mathbb{F}$ :

$$
\begin{equation*}
\frac{1}{\mathcal{I}_{n}}=\sup _{\boldsymbol{f} \in \mathbb{F}}\left\|\left(\boldsymbol{J}^{-1} \boldsymbol{K}\right)^{T} \boldsymbol{f}\right\|_{\infty} \tag{15}
\end{equation*}
$$

If $\mathbb{F}$ can be written as the Cartesian product of six intervals centred at 0 , that is, if there exists a diagonal matrix $\tilde{\boldsymbol{F}} \in \mathbb{R}^{6 \times 6}$ such that $\mathbb{F}=\left\{\tilde{\boldsymbol{F}} \boldsymbol{u},\|\boldsymbol{u}\|_{\infty}<1\right\}$, then the singularity index $\mathcal{I}_{n}$ can be written as

$$
\begin{equation*}
\frac{1}{\mathcal{I}_{n}}=\sup _{\|\boldsymbol{u}\|_{\infty} \leq 1}\left\|\left(\boldsymbol{J}^{-1} \boldsymbol{K}\right)^{T} \tilde{\boldsymbol{F}} \boldsymbol{u}\right\|_{\infty} \tag{16}
\end{equation*}
$$

Hence, the reciprocal of the singularity index $\mathcal{I}_{n}$ becomes

$$
\begin{equation*}
\frac{1}{\mathcal{I}_{n}}=\left\|\left(\boldsymbol{J}^{-1} \boldsymbol{K}\right)^{T} \tilde{\boldsymbol{F}}\right\|_{\infty} \tag{17}
\end{equation*}
$$

which corresponds to the maximum absolute row sum of matrix $\left(\boldsymbol{J}^{-1} \boldsymbol{K}\right)^{T} \tilde{\boldsymbol{F}}$. As a consequence, this singularity index is easy to compute numerically. This index is known as the infinity-norm-based kinematic sensitivity index proposed by Cardou et al. [22].

### 4.2 Determinant

Since type I singularities are not considered, the singularity index based on the determinant, noted $\mathcal{I}_{d}$, is

$$
\begin{equation*}
\mathcal{I}_{d}=|\operatorname{det}(\boldsymbol{J})| . \tag{18}
\end{equation*}
$$

This index is positive semidefinite and vanishes only at singular configurations. However, its value is not directly related to the actuated joint forces. An increasing value of the determinant does not guarantee a decrease of the maximum actuated joint forces, even though a crosscorrelation might be expected. Note that this index corresponds to Yoshikawa's manipulability measure [14].

An advantage of the index $\mathcal{I}_{d}$ compared to the index $\mathcal{I}_{n}$ is that $\mathcal{I}_{d}$ has a simple closed form, and so it may be possible to analytically determine if a feasible solution exists to the unlimited rotation problem.

### 4.3 Comparison

Figure 3 compares indices $\mathcal{I}_{d}$ and $\mathcal{I}_{n}$ for $\psi_{1} \in$ [ $0 ; 2 \pi\left[\right.$ and for some selected platform loads. Angle $\psi_{2}$ is fixed to $\frac{\pi}{2}$ and angle $\psi_{3}$ is then selected in order to maximize the singularity index. Matrix $\tilde{\boldsymbol{F}}$ is diagonal, the entries are 1 for forces and 0.1 for moments. Four scenarios are considered for the comparison: (i) a unit force along the $x_{b}$ direction, (ii) a unit force along the $\boldsymbol{z}_{b}$ direction, (iii) a unit torque around the $\boldsymbol{x}_{b}$ direction, (iv) a unit torque around the $\boldsymbol{z}_{b}$ direction. For each of these cases, the plots show the maximum actuator force obtained when maximizing $\mathcal{I}_{d}$ or $\mathcal{I}_{n}$. The platform considered here is similar to the one used in Wen et al. [10]. It is an equilateral triangle with $L_{i}=0.08 \mathrm{~m}$ and the proximal links have the same length $l_{i}=0.04 \mathrm{~m}$. Point $P$ is the barycentre of the platform.

Two main observations can be made. First, with each of the indices, the maximum actuated joint forces are discontinuous, especially with index $\mathcal{I}_{n}$. Indeed, arg max function that returns the argument of the function when it reaches its minimum is discontinuous. For example, if a function has two distinct local maxima, then an infinitesimal modification of this function may change the position of the global maximum. Therefore, the value of $\psi_{3}$ cannot be obtained by solely computing the maximum of the singularity index. A path planning approach is proposed in Section 5.

Second, it is not possible to conclude that one of the indices is better than the other, since the maximum actuator force is lower sometimes with one, sometimes with the other. This may seem counterintuitive since index $\mathcal{I}_{n}$ is designed to minimize the maximum actuated joint forces. However, this minimization is done in a worst-case scenario for a force belonging to $\mathbb{F}$. Hence, in cases other than the worst one, there is no reason for $\mathcal{I}_{n}$ to outdo $\mathcal{I}_{d}$. Yet, if $\mathbb{F}$ corresponds to a single wrench applied to the platform, then $\mathcal{I}_{n}$ will give the optimal solution if this wrench is indeed applied. This property may be useful in situations in which the external forces are already known, such as pick-and-place tasks with known objects.

The remaining of this paper considers $\mathcal{I}_{d}$ as the singularity index since (i) it can be written in a closed form and (ii) its maximum is less sensitive to small variations of the input $\psi_{1}$ (so there are fewer discontinuities).

## 5 PATH PLANNING

Consider the situation in which $\psi_{1} \in[0 ; 2 \pi[$ is performing a complete rotation and $\psi_{2}$ is constant. The feasibility of the unlimited rotation problem for $\psi_{2}$ varying in an interval $\left[\underline{\psi}_{2}, \bar{\psi}_{2}\right]$ is discussed in the Section 6. For now, consider the case where geometric parameters and


Fig. 3. Comparison of the maximum actuated joint force using each of the singularity indices.
$\psi_{2}$ are chosen such that the unlimited rotation problem is feasible and discuss how the path planning can be carried out. Note that it is not sufficient to find a value for $\psi_{3}$ for each value of $\psi_{1}$ to avoid singularities. This relation needs to be continuous, since $\psi_{3}$ cannot change instantaneously (it must be continuous).

The presented method can also be used if $\psi_{1}$ is constant and $\psi_{2}$ is varying. Therefore, it can be used for planning tasks requiring two DoFs, provided that $\psi_{1}$ and $\psi_{2}$ can be modified sequentially.

First, a simple case is presented for which a closedform solution exists for the singularity-free path planning problem. Then, a solution is proposed for the general case.

### 5.1 Simplified Case

Some simplifying assumptions are made. The platform is assumed to be an equilateral triangle with circumradius $L=L_{i}$, and the proximal links have the same length $l=l_{i}$. If the length of the proximal links is such that $l<\frac{3}{4} L$, then $\boldsymbol{J}_{z}$ can never be singular [10]. Therefore, it is possible to consider only $\boldsymbol{J}_{u}$ for singularity avoidance. Equation (13) can be rewritten as

$$
\begin{align*}
\operatorname{det}\left(\boldsymbol{J}_{u}\right) & =l L\left(\boldsymbol{u}_{1} \times \boldsymbol{u}_{2}\right)^{T} \boldsymbol{z}_{b} \sin \psi_{3} \\
& +\left[l L\left(\sin \psi_{1} \boldsymbol{u}_{2}-\sin \psi_{2} \boldsymbol{u}_{1}\right) \times \boldsymbol{u}_{3}\right]^{T} \boldsymbol{z}_{b} \tag{19}
\end{align*}
$$

For given values of $\psi_{1}$ and $\psi_{2}$, the first term of
the expression is a harmonic function of $\psi_{3}$ of frequency $\frac{1}{2 \pi}$. The second term is also a harmonic function of $\psi_{3}$ with the same frequency since $\left\|\boldsymbol{v}_{1} \times \boldsymbol{v}_{2}\right\|=$ $\left\|\boldsymbol{v}_{1}\right\|\left\|\boldsymbol{v}_{2}\right\|\left|\sin \left(\boldsymbol{v}_{1}, \boldsymbol{v}_{2}\right)\right|$ where $\sin \left(\boldsymbol{v}_{1}, \boldsymbol{v}_{2}\right)$ is the sine of the angle between $\boldsymbol{v}_{1}$ and $\boldsymbol{v}_{2}$. The sum of two terms with same frequency is also a harmonic function (they are solution to the same linear ordinary differential equation). Hence, one can write

$$
\begin{equation*}
\operatorname{det}\left(\boldsymbol{J}_{u}\right)=c_{1}\left(\psi_{1}, \psi_{2}\right) \sin \left(\psi_{3}+c_{2}\left(\psi_{1}, \psi_{2}\right)\right) \tag{20}
\end{equation*}
$$

where $c_{1}$ and $c_{2}$ are functions of angles $\psi_{1}$ and $\psi_{2}$ that can be readily obtained.

Therefore, if the unlimited rotation problem is feasible, then $\psi_{3}=\frac{\pi}{2}-c_{2}\left(\psi_{1}, \psi_{2}\right)$ is a feasible solution. However, one may note that if $\psi_{2}$ is such that either $T_{1}$ or $T_{3}$ is aligned with $T_{2}$ and $S_{2}$, i.e., if $\psi_{2} \in\left\{-\frac{5 \pi}{6},-\frac{\pi}{6}, \frac{\pi}{6}, \frac{5 \pi}{6}\right\}$, then there is no feasible solution. Indeed, if the line $\left(T_{2}, \boldsymbol{u}_{2}\right)$ is coincident with either $\left(T_{1}, \boldsymbol{u}_{1}\right)$ or $\left(T_{3}, \boldsymbol{u}_{3}\right)$, the configuration is singular.

Figure 4 shows different solutions for $\left.\psi_{2} \in\right] \frac{\pi}{6}, \frac{5 \pi}{6}[$ (Fig. 5.1) and $\left.\psi_{2} \in\right]-\frac{\pi}{6}, \frac{\pi}{6}[$ (Fig. 5.1). Cases where $\left.\psi_{2} \in\right] \frac{5 \pi}{6}, \frac{7 \pi}{6}\left[\right.$ and $\left.\psi_{2} \in\right] \frac{7 \pi}{6}, \frac{11 \pi}{6}[$ are the same since the absolute value of the determinant of $J_{u}$ is conserved if $\psi_{2}$ is shifted by $\pi$. A notable consequence is that if $\psi_{1}$ performs a complete rotation, so does $\psi_{3}$, and the direction of the rotation depends on the value of $\psi_{2}$.

### 5.2 General Case

Knowing $\psi_{1}$ and $\psi_{2}$, one may be tempted to select the value of $\phi_{3}$ that maximizes the singularity index. However, as discussed previously in Section 4, the function arg max is not continuous, and, thus, cannot be used here.

The approach implemented in this paper is similar to obstacle avoidance with artificial potential fields. Singular configurations correspond to repulsive sources and the target position to an attractive source.

When $\psi_{2}$ is fixed, the set of configurations $\left\{\left(\psi_{1}, \psi_{3}\right) \in \mathbb{R}^{2}\right\}$ define a toric space: the configurations are conserved under a $2 \pi$ shift of the coordinates. Therefore, the distance between two configurations is not the Euclidean distance, but the shortest distance between two points on the surface of a torus.

Let $\mathbb{S}$ be the set of singular configurations. The repulsive potential $V_{\text {rep }}$ at a point $\boldsymbol{q}=\left(\psi_{1}, \psi_{3}\right)$ is defined as

$$
\begin{equation*}
V_{r e p}(\boldsymbol{q})=\frac{1}{d_{s}(\boldsymbol{q}, \mathbb{S})} \tag{21}
\end{equation*}
$$



Fig. 4. Solutions to simple path planning.
where the function $d_{s}(\cdot, \mathbb{S})$ is the distance of the configuration $\boldsymbol{q}$ to the set $\mathbb{S}$. It corresponds to a singularity index. The attractive potential $V_{a t t}$ is defined as

$$
\begin{equation*}
V_{a t t}(\boldsymbol{q})=\left\|\boldsymbol{q}-\boldsymbol{q}_{t}\right\| \tag{22}
\end{equation*}
$$

with $\boldsymbol{q}_{t}$ the target configuration. In order to perform a complete rotation from an initial configuration $\boldsymbol{q}_{0}$, the target configuration is either $\boldsymbol{q}_{t}=\boldsymbol{q}_{0}+(2 \pi,-2 \pi)$ or $\boldsymbol{q}_{t}=\boldsymbol{q}_{0}+(2 \pi, 2 \pi)$, depending on the value of $\psi_{2}$. The total potential $V$ is the sum of the repulsive and the attractive potentials: $V(\boldsymbol{q})=V_{\text {rep }}(\boldsymbol{q})+\lambda V_{\text {att }}(\boldsymbol{q})$, with $\lambda \geq 0$ a scaling coefficient.

Assuming a periodic sampling of $\psi_{1}$, the path from the initial configuration to the target configuration is obtained by gradient descent, namely

$$
\begin{equation*}
\Delta \psi_{3}=-\varepsilon \frac{\partial V}{\partial \psi_{3}}\left(\psi_{1}, \psi_{3}\right) \tag{23}
\end{equation*}
$$

Applying the gradient descent algorithm to variable $\psi_{3}$ only prevents the algorithm from getting trapped in a local
minimum. Furthermore, for a given value of $\psi_{1}$, there is only one corresponding value of $\psi_{3}$.

This approach is tested in simulation with two repulsive potentials: the first one uses the singularity index $\mathcal{I}_{d}=\operatorname{det}(\boldsymbol{J})$ (case A) and the other uses the shortest distance to the singularities considering the toric geometry (case B). The platform is an equilateral triangle with $L_{i}=L=8 \times 10^{-2} \mathrm{~m}$ and $l_{i}=l=7.3 \times 10^{-2} \mathrm{~m}$. For the first case, the input increment is $\Delta \psi_{1}=\frac{\pi}{36}$, the learning rate $\varepsilon=0.11$ and the scaling coefficient $\lambda=11.7 L^{3} l^{3}$. For the second case, two parameters are modified: $\lambda=5$ and $\varepsilon=0.05$. Figure 5 shows the result of the path planning from the configuration $\left(\psi_{1}, \psi_{3}\right)=(0,0)$ for $\psi_{2}=\pi / 3$ with both potentials. The colour map 'batlow" is used to show the values of the potential [23]. The value of the potential $V$ is saturated at 50 . Solid lines represent the singularity locus and the dotted line represents the generated trajectory.

In this example, case B (using the shortest distance to $\mathbb{S}$ ) results in a smoother trajectory. The derivative of $\psi_{3}$ with respect to $\psi_{1}$ is shown in Fig. 6. Case A induces abrupt variations to $\psi_{3}$ to avoid singular configurations, more than fifteen times higher with respect to $\psi_{1}$ in the worst case. Furthermore, the lowest value reached by $\mathcal{I}_{d}$ is $34 \%$ lower with case A. Therefore, higher actuation forces may be required. Hence, using the shortest distance to $\mathbb{S}$ is preferred.

Note that not all nonsingular configurations are reachable from the initial position without crossing a singularity. Indeed, to perform a complete rotation, there are two distinct homotopy classes. The class of the chosen solution is defined by the initial configuration. Furthermore, for some initial nonsingular configurations, unlimited rotation is not feasible without modifying $\psi_{2}$ (e.g., $\psi_{1}=4 \mathrm{rad}$ and $\psi_{3}=3 \mathrm{rad}$, corresponding to the plus sign in Fig. 5).

The path generation could still fail because of local minima. More advanced techniques exist, such as modifying the potential to include virtual obstacles (i.e., generating a repulsive potential) [24] or implementing techniques that allow for escaping local minima [25]. Global path finding techniques could also be considered, since the dimension of the problem is low and the resolution could be carried out offline [26]. However, these are out of the scope of this paper.

## 6 FEASIBILITY

The feasibility of unlimited rotation depends on the geometric parameters of the platform $\left(L_{i}, l_{i}\right.$ and $\left.\phi_{i}\right)$ and the length and mean value of $\left[\underline{\psi}_{2}, \bar{\psi}_{2}\right]$, i.e., the interval of $\psi_{2}$. In the following, two conditions are presented. The

(a) Case A: using $\mathcal{I}_{d}$

(b) Case B: using the shortest distance to $\mathbb{S}$

Fig. 5. Path planning with artificial potential fields. Solid lines are the singularity locus. The dotted line is the generated trajectory. The point with a plus sign is an initial configuration from which unlimited rotation is not feasible.
first one is a geometric approach to avoid the singularity of matrix $\boldsymbol{J}_{z}$. The second one is a sufficient condition based on the path planning algorithm presented above in Section 5.


Fig. 6. Comparison of the rate of change of $\psi_{3}$ with respect to $\psi_{1}$ for each of the potentials.


Fig. 7. Feasibility condition.

### 6.1 Condition 1

Let $\mathcal{C}_{1}, \mathcal{C}_{2}$ and $\mathcal{C}_{3}$ be the sets of reachable positions respectively by $S_{1}, S_{2}$ and $S_{3}$ for a given position and orientation of the platform. Therefore, $\mathcal{C}_{1}$ and $\mathcal{C}_{3}$ describe circles and $\mathcal{C}_{2}$ a circular arc, whose respective centres are the vertices of the platform and the respective radii the length of the proximal links $l_{i}$.

From a design standpoint, a necessary and sufficient condition to avoid a singularity of matrix $J_{z}$ can be formulated as follows: no line intersects $\mathcal{C}_{1}, \mathcal{C}_{2}$ and $\mathcal{C}_{3}$. This condition is illustrated in Fig. 7. The shaded area is the set of all lines intersecting both $\mathcal{C}_{1}$ and $\mathcal{C}_{3}$. This area is delimited by the common tangent lines to $\mathcal{C}_{1}$ and $\mathcal{C}_{3}$. If $\mathcal{C}_{2}$ does not intersect this area, then $\boldsymbol{J}_{z}$ is never singular.

Note that this condition is only a consideration that can be taken into account during the design, and is neither necessary nor sufficient for the feasibility of the unlimited rotation problem. Indeed, even if a line can intersect $\mathcal{C}_{1}$,


Fig. 8. Feasibility test by path planning.
$\mathcal{C}_{2}$ and $\mathcal{C}_{3}$, there is no guarantee that a singular configuration will happen since the value of $\psi_{3}$ is not free and is rather obtained from $\psi_{1}$ and $\psi_{2}$. Moreover, this condition does not prevent the singularity of matrix $\boldsymbol{J}_{u}$.

### 6.2 Condition 2

If there is a periodic trajectory $\psi_{3}\left(\psi_{1}\right)$ with period $2 \pi$ performing a complete rotation for any value of $\psi_{2} \in$ $\left[\underline{\psi}_{2}, \bar{\psi}_{2}\right]$, then the unlimited rotation is feasible. This can be tested in simulation using the path planning method from Section 5 by discretizing $\left[\underline{\psi}_{2}, \bar{\psi}_{2}\right]$.

Simulation results using the same platform as in Section 5 with $\left[\underline{\psi}_{2}, \bar{\psi}_{2}\right]=\left[\frac{\pi}{3}-\frac{\pi}{8}, \frac{\pi}{3}+\frac{\pi}{8}\right], \Delta \psi_{1}=\frac{\pi}{36}$, $\lambda=5$ and $\varepsilon=0.04$ are shown in Fig. 8. In this case, the path generation succeeded. Therefore, the unlimited rotation problem is feasible with $\psi_{2} \in\left[\underline{\psi}_{2}, \bar{\psi}_{2}\right]$. However, situations in which path planning fails are not necessarily infeasible since (i) $\psi_{2}$ does not take different values at the same time and (ii) local minima may prevent the gradient descent to succeed close to singularities.

## 7 DESIGN CONSIDERATIONS

This section discusses the design of the platform in order to reduce the actuation forces.

A platform can be considered as a set of ten independent parameters

$$
\begin{equation*}
\rho=\left(L_{i}, l_{i}, \phi_{2}, \phi_{3}, \underline{\psi}_{2}, \bar{\psi}_{2}\right) \tag{24}
\end{equation*}
$$

with $i \in\{1,2,3\}$. Then, the singularity index $\mathcal{I}_{d}$ is a
function of $\psi_{1}, \psi_{2}, \psi_{3}$ and $\rho: \mathcal{I}_{d}=\mathcal{I}_{d}\left(\psi_{1}, \psi_{2}, \psi_{3}, \rho\right)$.
Let $\mathbb{P}$ be the set of all platforms. The performance index $\mathcal{M}$ of a platform $\rho \in \mathbb{P}$ is defined as the lowest value of $\mathcal{I}_{d}$ for any value of $\psi_{1}$ and $\psi_{2}$, where the value of $\psi_{3}$ maximizes $\mathcal{I}_{d}$ for a given $\psi_{1}$ and $\psi_{2}$. Hence,

$$
\begin{equation*}
\mathcal{M}(\rho)=\min _{\psi_{1}, \psi_{2}} \mathcal{I}_{d}\left(\psi_{1}, \psi_{2}, \psi_{3}, \rho\right) \tag{25}
\end{equation*}
$$

subject to the constraint

$$
\begin{equation*}
\psi_{3}=\arg \max _{\psi_{3}} \mathcal{I}_{d}\left(\psi_{1}, \psi_{2}, \psi_{3}, \rho\right) \tag{26}
\end{equation*}
$$

Practically speaking, the number of free parameters to optimize is much lower than ten. The design is illustrated with an example. Figure 9 shows a gripper whose opening and orientation can be controlled with the redundant links $R_{i} S_{i}$. Here, angle $\psi_{1}$ controls the orientation of the gripper and angle $\psi_{2}$ the opening through springs and a wire transmission. The timing belt with ratio $3: 1$ guarantees a $180^{\circ}$ rotation of the gripper in configurations where unlimited rotation is prevented due to mechanical interference between the configurable platform and the legs. The configurable platform is attached to the leg mechanisms using spherical joints proposed by Schreiber and Gosselin [8]. These joints use a kinematically redundant design and generate a very large range of motion, exceeding $\pm 150^{\circ}$. The CAD model suggests that $L_{i}$ is at least $8 \times 10^{-2} \mathrm{~m}$ and $l_{i}$ at most $6 \times 10^{-2} \mathrm{~m}$ to avoid interference between redundant links $R_{i} S_{i}$. The range of motion of angle $\psi_{2}$ is $\bar{\psi}_{2}-\underline{\psi}_{2}=60^{\circ}$ (modulo $360^{\circ}$ ). Therefore, the only parameters to optimize are $\psi_{2}$, $\phi_{3}$ and $\underline{\psi}_{2}$ (or, equivalently, $\bar{\psi}_{2}$ ).

In order to reduce the computational complexity, the approach is separated in two steps. First, the optimal value for $\underline{\psi}_{2}$ is obtained considering initially an equilateral triangle as the platform. Then, the platform is optimized with respect to $\phi_{2}$ and $\phi_{3}$. The impact of $\underline{\psi}_{2}$ is illustrated in Figure 10. The shaded area in Fig. 7 shows the optimal interval $\left[\underline{\psi}_{2}, \bar{\psi}_{2}\right]$, and the circle in Fig. 7 is the optimal value for $\underline{\underline{\psi}}_{2}$. Figure 11 shows the impact of $\phi_{2}$ and $\phi_{3}$ on the performance index $\mathcal{M}$. It can be seen that $\phi_{2}=\frac{2 \pi}{3}$ and $\phi_{3}=\frac{4 \pi}{3}$, corresponding to an equilateral triangle, are the optimal values. The video accompanying this paper shows an animation of the CAD model performing unlimited rotations. The robot used for the simulations is designed by our research group [27]. It can be seen that there is no mechanical interference.

Note that if the lengths $L_{i}$ and $l_{i}$ are not identical, the optimal values for $\phi_{2}$ and $\phi_{3}$ are different. Figure


Fig. 9. CAD model of the gripper.

12 shows the effect of $\phi_{2}$ and $\phi_{3}$ with a platform such that: $L_{1}=5 \times 10^{-2} \mathrm{~m}, L_{2}=6 \times 10^{-2} \mathrm{~m}, L_{3}=7 \times$ $10^{-2} \mathrm{~m}, l_{1}=3 \times 10^{-2} \mathrm{~m}, l_{2}=4 \times 10^{-2} \mathrm{~m}, l_{3}=5 \times$ $10^{-2} \mathrm{~m}$. In this case, choosing the optimal values rather than $\phi_{2}=\frac{2 \pi}{3}$ and $\phi_{3}=\frac{4 \pi}{3}$ yields a $0.9 \%$ improvement of the performance index, which may nevertheless not be that significant in practice.

## 8 OTHER PLATFORMS

There are three categories of planar platforms combining P (prismatic) and R (revolute) joints for unlimited rotations, namely: 3R (three revolute), 1P2R (one prismatic and two revolute) and 2P1R (two prismatic and one revolute). Platform 3R has been studied above. In the following, platforms 1P2R and 2P1R are discussed.

### 8.1 Platform 1P2R

Platform 1P2R uses a parameterization similar to the one used for platform 3R. Some geometric parameters specific to the platform 1 P 2 R are shown in Fig. 13. The axis of the prismatic joint is $\left(T_{3}, \boldsymbol{v}_{3}\right)$, with $\boldsymbol{v}_{3}=\cos \left(\phi_{3}+\psi_{3}\right) \boldsymbol{x}_{b}+\sin \left(\phi_{3}+\psi_{3}\right) \boldsymbol{y}_{b}$. Parameter $r$ is defined such that $s_{3}=\boldsymbol{t}_{3}+r \boldsymbol{v}_{3}$. It should be pointed


Fig. 10. Impact of $\underline{\psi}_{2}$ on $\mathcal{I}_{d}$. The shaded area corresponds to the optimal $\left[\underline{\psi}_{2}, \bar{\psi}_{2}\right]$.
out that for this platform, $\psi_{3}$ is a constant design parameter.

Let $\xi_{P_{r}}$ be the twist of the prismatic joint. Then,

$$
\begin{equation*}
\xi_{E E}=\xi_{P_{3}}+\xi_{S_{3}}+\xi_{P_{r}} \tag{27}
\end{equation*}
$$

Let $\boldsymbol{v}_{3 \perp}$ be a unit vector such that $\boldsymbol{v}_{3}^{T} \boldsymbol{v}_{3 \perp}=0$ and $\boldsymbol{v}_{3} \times \boldsymbol{v}_{3 \perp}=\boldsymbol{z}_{b}$. Then, the zero-pitch twist $\xi_{S_{3}, \boldsymbol{v}_{3 \perp}}$ of line $\left(S_{3}, \boldsymbol{v}_{3 \perp}\right)$ is reciprocal to $\xi_{S_{3}}$ and $\xi_{P_{r}}$. Hence,

$$
\begin{equation*}
\xi_{S_{3}, \boldsymbol{v}_{3 \perp}} \cdot \xi_{E E}=\xi_{S_{3}, \boldsymbol{v}_{3 \perp}} \cdot \xi_{P_{3}} \tag{28}
\end{equation*}
$$

which gives

$$
\begin{equation*}
\boldsymbol{v}_{3 \perp}^{T} \dot{\boldsymbol{p}}+\left(\boldsymbol{s}_{3} \times \boldsymbol{v}_{3 \perp}\right)^{T} \boldsymbol{\omega}=\boldsymbol{v}_{3 \perp}{ }^{T} \dot{\boldsymbol{s}}_{3} \tag{29}
\end{equation*}
$$


(a) Performance index $\mathcal{M}$ as a function of $\phi_{2}$

(b) Performance index $\mathcal{M}$ as a function of $\phi_{3}$

Fig. 11. Performance index $\mathcal{M}$ as a function of $\phi_{2}$ and $\phi_{3}$ with $L_{i}=L$ and $l_{i}=l$. The circles show the initial guess.

Combining (4), (6) and (29) yields

$$
\left[\begin{array}{cc}
\boldsymbol{u}_{1}{ }^{T} & \left(\boldsymbol{t}_{1} \times \boldsymbol{u}_{1}\right)^{T}  \tag{30}\\
\boldsymbol{u}_{2}^{T} & \left(\boldsymbol{t}_{2} \times \boldsymbol{u}_{2}\right)^{T} \\
\boldsymbol{v}_{3 \perp}{ }^{T} & \left(\boldsymbol{s}_{3} \times \boldsymbol{v}_{3 \perp}\right)^{T} \\
\boldsymbol{z}_{b}^{T} & \left(\boldsymbol{s}_{1} \times \boldsymbol{z}_{b}\right)^{T} \\
\boldsymbol{z}_{b}^{T} & \left(\boldsymbol{s}_{2} \times \boldsymbol{z}_{b}\right)^{T} \\
\boldsymbol{z}_{b}^{T} & \left(\boldsymbol{s}_{3} \times \boldsymbol{z}_{b}\right)^{T}
\end{array}\right]\left[\begin{array}{c}
\dot{\boldsymbol{p}} \\
\boldsymbol{\omega}
\end{array}\right]=\left[\begin{array}{ccc}
\boldsymbol{u}_{1}{ }^{T} & \mathbf{0}^{T} & \mathbf{0}^{T} \\
\mathbf{0}^{T} & \boldsymbol{u}_{2}^{T} & \mathbf{0}^{T} \\
\mathbf{0}^{T} & \mathbf{0}^{T} & \boldsymbol{v}_{3 \perp}^{T} \\
\boldsymbol{z}_{b}^{T} & \mathbf{0}^{T} & \mathbf{0}^{T} \\
\mathbf{0}^{T} & \boldsymbol{z}_{b}^{T} & \mathbf{0}^{T} \\
\mathbf{0}^{T} & \mathbf{0}^{T} & \boldsymbol{z}_{b}^{T}
\end{array}\right]\left[\begin{array}{l}
\dot{\boldsymbol{s}}_{1} \\
\dot{\boldsymbol{s}}_{2} \\
\dot{\boldsymbol{s}}_{3}
\end{array}\right] .
$$

The singularity analysis is similar to that of the platform with three R joints studied above. Since $s_{3} \times \boldsymbol{v}_{3 \perp}$ has no component along $\boldsymbol{z}_{b}$, type II singularities happen if and only if either $\boldsymbol{J}_{z}(10)$ or $\boldsymbol{J}_{v}$ is singular, with

$$
\boldsymbol{J}_{v}=\left[\begin{array}{ccc}
\boldsymbol{u}_{1}^{T} \boldsymbol{x}_{b} & \boldsymbol{u}_{1}^{T} \boldsymbol{y}_{b} & \left(\boldsymbol{t}_{1} \times \boldsymbol{u}_{1}\right)^{T} \boldsymbol{z}_{b}  \tag{31}\\
\boldsymbol{u}_{2}{ }^{T} \boldsymbol{x}_{b} & \boldsymbol{u}_{2}^{T} \boldsymbol{y}_{b} & \left(\boldsymbol{t}_{2} \times \boldsymbol{u}_{2}\right)^{T} \boldsymbol{z}_{b} \\
\boldsymbol{v}_{3 \perp}{ }^{T} \boldsymbol{x}_{b} & \boldsymbol{v}_{3 \perp}{ }^{T} \boldsymbol{y}_{b} & \left(s_{3} \times \boldsymbol{v}_{3 \perp}\right)^{T} \boldsymbol{z}_{b}
\end{array}\right]
$$

From Grassmann line geometry, matrix $\boldsymbol{J}_{v}$ is singular if and only if lines $\left(S_{1}, \boldsymbol{u}_{1}\right),\left(S_{2}, \boldsymbol{u}_{2}\right)$ and $\left(S_{3}, \boldsymbol{v}_{3 \perp}\right)$ have a common intersection or are parallel to each other.

Consider a platform with the following parameters: $L_{1}=L_{2}=8 \times 10^{-2} \mathrm{~m}, L_{3}=0, l_{1}=l_{2}=7.3 \times 10^{-2} \mathrm{~m}$,


Fig. 12. Effect of $\phi_{2}$ and $\phi_{3}$ on the performance index $\mathcal{M}$ for an asymmetrical platform. The circles show the initial guess.
$\phi_{1}=0, \phi_{2}=\frac{2 \pi}{3}, \phi_{3}=\frac{4 \pi}{3}, \psi_{3}=0$ and $\left[\underline{r}_{3}, \bar{r}_{3}\right]=$ $[0.05, \quad 0.1]$. The potential $V\left(\psi_{1}, \psi_{2}\right)$ for $r_{3} \in\left[\underline{r}_{3}, \bar{r}_{3}\right]$ with the initial point $\psi_{1}=\psi_{2}=0$ is shown in Fig. 14. As it can be seen, unlimited rotation is feasible without crossing any singularity.

Design considerations presented in Section 7 can also be used for the 1P2R platform. Additionally, parameter $\psi_{3}$ needs to be chosen appropriately. Indeed, the effect of $\psi_{3}$ on the performance index $\mathcal{M}$ is illustrated in Fig. 15. The value of the performance index is multiplied by 1.6 between $\psi_{3}=\frac{\pi}{2}$ (the worst case) and $\psi_{3}=0$ (the best case).

### 8.2 Platform 2P1R

The parameterization of the 2P1R platform is similar to that of the 1P2R platform. The revolute joint associated with variable $\psi_{2}$ is replaced by a prismatic joint of axis $\boldsymbol{v}_{3}=\cos \left(\psi_{2}+\psi_{2}\right) \boldsymbol{x}_{b}+\sin \left(\psi_{2}+\psi_{2}\right) \boldsymbol{y}_{b}$ and vector $\boldsymbol{v}_{3 \perp}$ is such that $\left(\boldsymbol{v}_{3}, \boldsymbol{v}_{3 \perp}, \boldsymbol{z}_{b}\right)$ is a right-handed orthonormal basis.

The following parameters are defined from Fig. 16. The intersection of lines $\left(S_{1}, \boldsymbol{u}_{1}\right)$ and $\left(S_{2}, \boldsymbol{v}_{2 \perp}\right)$ (respectively $\left(S_{3}, \boldsymbol{v}_{3 \perp}\right)$ ) is noted $M_{12}$ (respectively $M_{13}$ ). Let


Fig. 13. Parameterization of the platform 1P2R.


Fig. 14. Potential $V$ for the 1 P 2 R platform.


Fig. 15. Effect of $\psi_{3}$ on the performance index $\mathcal{M}$ for the 1P2R platform.


Fig. 16. Parameterization of the platform 2P1R.
$\boldsymbol{m}_{13}$ be the position vector of point $M_{13}$ and $\boldsymbol{m}_{23}$ the position vector of point $M_{23}$. Parameters $\lambda_{13}$ and $\lambda_{23}$ are defined such that $\boldsymbol{m}_{13}=\lambda_{13} \boldsymbol{v}_{3 \perp}+\boldsymbol{s}_{3}$ and $\boldsymbol{m}_{23}=$ $\lambda_{23} v_{3 \perp}+s_{3}$.

According to Grassmann line geometry, the platform is in a singular configuration if lines $\left(S_{1}, \boldsymbol{u}_{1}\right),\left(S_{2}, \boldsymbol{v}_{2 \perp}\right)$ and $\left(S_{3}, \boldsymbol{v}_{3 \perp}\right)$ have a common intersection or are parallel. The case where $\boldsymbol{v}_{2 \perp}$ and $\boldsymbol{v}_{3 \perp}$ are collinear is not considered since unlimited rotation is trivially unfeasible (there exists a value of $\psi_{1}$ such that $\boldsymbol{u}_{1}, \boldsymbol{v}_{2 \perp}$ and $\boldsymbol{v}_{3 \perp}$ are collinear).

Consider the initial situation of Fig. 16, in which $\lambda_{13}<\lambda_{23}$ and $\dot{\psi}_{1}>0$. The case $\lambda_{13}>\lambda_{23}$ is equivalent to this one considering a rotation of $\psi_{1}$ in the opposite direction, i.e., $\dot{\psi}_{1}<0$. Let $\psi_{o}$ be the initial value of $\psi_{1}$ and $\psi_{f}>\psi_{o}$ be the value of $\psi_{1}$ such that $\left(S_{1}, \boldsymbol{u}_{1}\right)$ is parallel to $\left(S_{3}, \boldsymbol{v}_{3 \perp}\right)$ and $\boldsymbol{u}_{1}{ }^{T} \boldsymbol{v}_{3 \perp}>0$.

We now show that unlimited rotation of $\psi_{1}$ is not feasible for constant $r_{3}$ using a proof by contradiction. If unlimited rotation of $\psi_{1}$ is feasible, then $\psi_{1}$ can move from $\psi_{o}$ to $\psi_{f}$. Let $f:\left[\psi_{o}, \psi_{f}\right] \rightarrow \mathbb{R}$ be a continuous function associating to each value of $\psi_{1}$ a feasible value of $\lambda_{23}$. In other words, $f$ is a feasible path for $M_{23}$. The existence of such a function is guaranteed since it is necessary for unlimited rotation. Since $f$ is continuous, there exists a segment (a closed and bounded interval) $\left[\underline{\lambda}_{23}, \bar{\lambda}_{23}\right]$ such that $\left[\underline{\lambda}_{23}, \bar{\lambda}_{23}\right]=f\left(\left[\psi_{o}, \psi_{f}\right]\right)$. However, $\lambda_{13}$ goes to $+\infty$ when $\psi_{1}$ converges towards $\psi_{f}$. Since $\lambda_{23}>\lambda_{13}$, interval $\left[\underline{\lambda}_{23}, \bar{\lambda}_{23}\right]$ cannot be bounded, which is contradictory. Indeed, moving $\psi_{1}$ from $\psi_{o}$ to $\psi_{f}$ is assumed to be feasible. As a consequence, unlimited rotation is not feasible. Note that the same reasoning can be applied to the platform 1P2R if $\psi_{2}$ is constant and $r_{3}$ is used to avoid singularities.

## 9 CONCLUSION AND PERSPECTIVES

With appropriate design and initial configuration, the planar platform of a $(6+3)$-DoF tripedal parallel robot can carry out unlimited rotation if it is equipped with at least two revolute joints (and at most one prismatic joint). A method based on artificial potential fields is proposed in order to find a feasible path for unlimited rotation. The success of the path finding approach can also be used as a sufficient condition for the feasibility of the unlimited rotation problem. Two singularity indices are compared in order to resolve the redundancy while reducing the actuated joint torques. A performance index based on a singularity index is proposed and the impact of geometric parameters are discussed in order to identify rules for the design of the platform. While the presented methods are applied to determine the architecture parameters that can produce unlimited rotation, they can also be used for tasks requiring a smaller (finite) range of motion.

An immediate perspective to this work is the generalization to spatial platforms. Therefore, any 3-DoF parallel mechanism could be used as a configurable platform without requiring transmission components. As an example, it could be possible to have three revolute joints with concurrent axes in order to produce a spherical movement [28].

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