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# The twin blow-up method for Hamilton-Jacobi equations in higher dimension 

Nicolas Forcadel, Cyril Imbert and Régis Monneau

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#### Abstract

In this paper, we show how to extend the twin blow-up method recently developped by the authors (Comptes Rendus. Math., 2024), in order to obtain a new comparison principle for an evolution coercive Hamilton-Jacobi equation posed in a domain of an Euclidian space of any dimension and supplemented with a boundary condition. The method allows dealing with the case where tangential variables and the variable corresponding to the normal gradient of the solution are strongly coupled at the boundary. We elaborate on a method introduced by P.-L. Lions and P. Souganidis (Atti Accad. Naz. Lincei, 2017). Their argument relies on a single blow-up procedure after rescaling the semi-solutions to be compared while two simultaneous blow-ups are performed in this work, one for each variable of the classical doubling variable technique. A one-sided Lipschitz estimate satisfied by a combination of the two blow-up limits plays a key role.


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## 1 Introduction

This work is concerned with strong uniqueness (comparison principle) of viscosity solutions to a HamiltonJacobi equation of evolution type of the form,

$$
\begin{equation*}
u_{t}+H(X, D u)=0 \quad \text { on } \quad(0, T) \times \Omega \tag{1.1}
\end{equation*}
$$

where $X:=(t, x)$, supplemented with the (desired) boundary condition

$$
u_{t}+F(X, D u)=0 \quad \text { on } \quad(0, T) \times \partial \Omega
$$

and the initial condition

$$
\begin{equation*}
u(0, \cdot)=u_{0} \quad \text { on } \quad\{0\} \times \bar{\Omega} \tag{1.2}
\end{equation*}
$$

The spatial domain $\Omega$ is a subset of the Euclidian space of dimension $d \geq 1$. We will first see how to deal with a half-space and we will then consider the case of a $C^{1}$ bounded domain.

Since the desired boundary condition can be lost when characteristics reach $\partial \Omega$, it has to be imposed in a weak sense. In the viscosity solution framework, the weak sense means that either the desired boundary condition is satisfied or the PDE is satisfied on the boundary. More precisely, subsolutions and supersolutions of (1.1) are assumed to satisfy at the boundary the following inequalities,

$$
\left\{\begin{array}{lll}
u_{t}+\min (F, H)(X, D u) \leq 0 & \text { on } & (0, T) \times \partial \Omega  \tag{1.3}\\
u_{t}+\max (F, H)(X, D u) \geq 0 & \text { on } & (0, T) \times \partial \Omega
\end{array} \quad\right. \text { (subsolutions) }
$$

We present in the introduction the comparison principle for (1.1), (1.2), (1.3) with $\Omega=\mathbb{R}^{d-1} \times(0,+\infty)$. In order to present the structure conditions imposed to the Hamiltonian $H$ and the nonlinearity $F$ associated with the boundary condition, we set $x=\left(x^{\prime}, x_{d}\right) \in \mathbb{R}^{d-1} \times[0,+\infty)$ and $p=\left(p^{\prime}, p_{d}\right) \in \mathbb{R}^{d-1} \times \mathbb{R}$ for the variable for the gradient. In particular, $p^{\prime}$ corresponds to the space tangential gradient of the solutions and $p_{d}$ to the normal gradient. In the following assumption, $\omega, \omega_{L}$ denote moduli of continuity.

$$
\left\{\begin{array}{l}
\text { i) }(\text { Continuity and bound) } \\
H:[0, T] \times \bar{\Omega} \times \mathbb{R}^{d} \rightarrow \mathbb{R} \text { is continuous } \\
\text { the map } X \mapsto H(X, 0) \text { is bounded. } \\
\text { ii) (Uniform continuity in the gradient) } \\
\text { For any } L>0 \text {, we have for all } X \in[0, T] \times \bar{\Omega} \text { and } p, q \in[-L, L]^{d} \\
|H(X, p)-H(X, q)| \leq \omega_{L}(|p-q|) \text {. }  \tag{1.4}\\
\text { iii) }(\text { Continuity in the tangential variables) } \\
\text { For } X=\left(t, x^{\prime}, x_{d}\right) \text { and } Y=\left(s, y^{\prime}, x_{d}\right) \text { with } t, s \in[0, T] \text { and } x^{\prime}, y^{\prime} \in \mathbb{R}^{d-1} \text { and } x_{d} \geq 0 \\
H(Y, p)-H(X, p) \leq \omega\left(|Y-X|\left(1+\left|p^{\prime}\right|+\max \{0, H(X, p)\}\right)\right) \text {. } \\
\text { iv) }(\mathbf{U n i f o r m} \text { normal coercivity) } \\
\text { For any } L>0, \text { we have } \\
\lim \inf \left\{H\left(X, p^{\prime}, p_{d}\right): X \in[0, T] \times \bar{\Omega}, p^{\prime} \in[-L, L]^{d-1}\right\}=+\infty . \\
\left|p_{d}\right| \rightarrow+\infty
\end{array}\right.
$$

and making artificially appear the dependence on $x \in \bar{\Omega}$ for $F$ (in order to unify the presentation of $H$ and $F)$, we consider
$\left\{\begin{array}{l}\text { i) (Continuity, bound and monotonicity) } \\ F:[0, T] \times \partial \Omega \times \mathbb{R}^{d} \rightarrow \mathbb{R} \text { is continuous, } \\ \text { the map } X \mapsto F(X, 0) \text { is bounded, } \\ \text { the map } p_{d} \mapsto F\left(X, p^{\prime}, p_{d}\right) \text { is nonincreasing. } \\ \text { ii) (Uniform continuity in the gradient) } \\ \text { For any } L>0, \text { we have for all } X \in[0, T] \times \partial \Omega \text { and } p, q \in[-L, L]^{d} \\ |F(X, p)-F(X, q)| \leq \omega_{L}(|p-q|) \text {. } \\ \text { iii) (Continuity in the tangential variables) } \\ \text { for all } X, Y \in[0, T] \times \partial \Omega \text { and } p \in \mathbb{R}^{d}, \\ F(Y, p)-F(X, p) \leq \omega\left(|Y-X|\left(1+\left|p^{\prime}\right|+\max \{0, \max (F, H)(X, p)\}\right)\right) . \\ \text { iv) (Uniform normal semi-coercivity) } \\ \text { For any } L>0, \text { we have } \\ \lim _{d} \rightarrow-\infty \\ \inf \left\{F\left(X, p^{\prime}, p_{d}\right): X \in[0, T] \times \partial \Omega, p^{\prime} \in[-L, L]^{d-1}\right\}=+\infty .\end{array}\right.$

## iii) (Continuity in the tangential variables)

for all $X, Y \in[0, T] \times \partial \Omega$ and $p \in \mathbb{R}^{d}$,
$F(Y, p)-F(X, p) \leq \omega\left(|Y-X|\left(1+\left|p^{\prime}\right|+\max \{0, \max (F, H)(X, p)\}\right)\right)$.
iv) (Uniform normal semi-coercivity)

For any $L>0$, we have
$\lim _{p_{d} \rightarrow-\infty} \inf \left\{F\left(X, p^{\prime}, p_{d}\right): X \in[0, T] \times \partial \Omega, p^{\prime} \in[-L, L]^{d-1}\right\}=+\infty$.
Under the previous structural conditions, sub and super-solutions of the Hamilton-Jacobi equation under study can be compared.

Theorem 1.1 (A comparison principle with strong tangential coupling). Let $T>0$ and assume that $H, F$ satisfy (1.4)-(1.5). Assume that the initial data $u_{0}$ is uniformly continuous. Let $u, v:[0, T) \times \bar{\Omega} \rightarrow \mathbb{R}$ be two functions with $u$ upper semi-continuous and $v$ lower semi-continuous. Assume that $u$ (resp. v) is a viscosity subsolution (resp. supersolution) of (1.1)-(1.2). Assume moreover that there exists a constant $C_{T}>0$ such that

$$
\begin{equation*}
u \leq u_{0}+C_{T} \quad \text { and } \quad v \geq u_{0}-C_{T} \quad \text { on } \quad[0, T) \times \bar{\Omega} \tag{1.6}
\end{equation*}
$$

If we have

$$
u(0, \cdot) \leq u_{0} \leq v(0, \cdot) \quad \text { on } \quad\{0\} \times \bar{\Omega}
$$

then we have

$$
u \leq v \quad \text { on } \quad[0, T) \times \bar{\Omega}
$$

Remark 1.2. A simplified version of Theorem 1.1 is presented in [5]. It was assumed in this note that dimension $d=1$ and that initial data are Lipschitz continuous. Some details were skept and they are presented in this new work.
Remark 1.3. In Section 5, we also extend this result to the case where $\Omega$ is a $C^{1}$ bounded open set.
Remark 1.4. Notice that, given (1.4), we can always define the state constraint boundary function

$$
H^{-}\left(X, p^{\prime}, p_{d}\right):=\inf _{q_{d} \leq p_{d}} H\left(X, p^{\prime}, q_{d}\right) \quad \text { for } \quad X \in[0, T] \times \partial \Omega \quad \text { and } \quad p=\left(p^{\prime}, p_{d}\right) \in \mathbb{R}^{d-1} \times \mathbb{R}
$$

and it satisfies (1.5). Up to our knowledge, the comparison principle was also an open problem for $F=H^{-}$ in this generality.

Remark 1.5. Notice that in Theorem 1.1, semi-coercivity of $F$ in condition (1.5) iv) can be replaced by the weak continuity of the subsolution $u$ on the boundary $(0, T) \times \partial \Omega$, using [4, Proposition 3.12]) and replacing $F$ by $F_{1}:=\max \left(F, H^{-}\right)$.

Main contribution. When comparing non-Lipschitz sub/supersolutions (for instance after constructing solutions by Perron's method), a strong coupling between tangential coordinates $\left(t, x^{\prime}\right) \in[0, T] \times \mathbb{R}^{d-1}$ and the normal gradient $\partial_{x_{d}} u$ is well identified in the literature as a technical difficulty, especially when this coupling arises in the boundary condition, see for instance $[1,2,3,10,11]$.

It is standard to make the (strong) assumption of uniform continuity in time $t$, uniformly in the gradient $D u$. Such an assumption is not satisfied by the following simple example,

$$
\begin{cases}u_{t}+a(X)|D u|=0 & \text { in } \quad(0, T) \times \Omega  \tag{1.7}\\ u_{t}+\max \left\{0,-b(X) \partial_{x_{d}} u\right\}=0 & \text { in } \quad(0, T) \times \partial \Omega\end{cases}
$$

when $a, b \geq 1$ are bounded Lipschitz continuous functions (here with $b(t, x)=b\left(t, x^{\prime}, 0\right)$ ).
Comparison with known results. J. Guerand [6] proved a comparison principle in our geometric setting in dimension $d=1$ in the case where $H$ and $F$ are independent of $(t, x)$. She also proved a comparison principle for non-coercive Hamiltonians.
P.-L. Lions and P. Souganidis [11] introduced a new method for proving comparison principles for bounded uniformly continuous sub/supersolutions for equations posed on junctions with several branches (or halfspaces). They use a blow-up argument that reduces the study to a 1D problem. They show the comparison principle in the case of Kirchoff-type boundary conditions and non-convex Hamiltonians. As far as $(t, x)$ dependence is concerned, their method allows them to handle Hamiltonians that are Lipschitz continuous in $t$, see [11, Assumption (4)].

This result is generalized by G. Barles and E. Chasseigne [3, Theorem 15.3.7, page 295] to the case of bounded semi-continuous sub/supersolutions under three different junction conditions. Even if they are presented for $N=2$ branches, we present their results in our geometric setting: a junction reduced to a single branch $N=1$ in dimension $d \geq 1$. The three cases are the following: (1) $F$ is independent on $p_{d}$,
(2) the Neumann problem and (3) general nonincreasing continuous $p_{d} \mapsto F\left(X, p^{\prime}, p_{d}\right)$. In the third case, the normal derivative is not coupled with the tangential coordinates $\left(t, x^{\prime}\right)$ in $F$ (see also the very end of [3, Subsection 13.2.2 and condition (GA-G-FLT) p. 247]).

As explained above, we improve these results, using the twin blow-up method introduced in [5]. A close look at the proof reveals that new ideas appear at the beginning of Step 4, when the reasoning focuses on the case where the point of maximum is on the boundary of the domain. Compared to the note [5], we also extend the result by considering uniformly continuous initial data (and not only Lipschitz continuous ones) and working in dimension greater than one.

Organization of the paper. In Section 2, we present two key boundary results stated for stationary problems in space dimension $d=1$. We also extend these results to the case of junctions (that will be used in future works). In Section 3, we recall two classical results which are suitable for our purpose. We first construct barriers. We next present some a priori estimates for the sup-convolution of subsolutions to coercive HJ equations. The proof of the comparison principle in the case of the half space (Theorem 1.1) is done in Section 4. Finally in Section 5, we show how to adapt our twin blow-up method to the case of a $C^{1}$ bounded open domain.

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## 2 Boundary lemmas

In this section, we work in dimension $d=1$ and set $\Omega:=(0,+\infty)$. We present some fundamental boundary results that will allow us to prove our comparison principle. At the end of this section, we also extend them naturally to the case of junctions (that will be useful for future works).

Before to state our result, we need to introduce the following notion of (limiting) semi-differentials.
Definition 2.1 ((Limiting) semi-differentials). Let $A \subset \bar{\Omega}$ and $x_{0} \in A$. For $(+/-)$, we define the (first order) super/subdifferential at $x_{0}$ of a function $u$ on $A$ as

$$
\begin{equation*}
D_{A}^{ \pm} u\left(x_{0}\right)=\left\{p \in \mathbb{R}, \quad \text { such that } \quad 0 \leq \pm\left\{u\left(x_{0}\right)+p \cdot\left(x-x_{0}\right)+o\left(x-x_{0}\right)-u(x)\right\} \quad \text { on } \quad A\right\} \tag{2.1}
\end{equation*}
$$

and the limit (first order) super/subdiffential at the boundary point $x_{0} \in \partial \Omega$ of $u$ as

$$
\begin{equation*}
\bar{D}_{\Omega}^{ \pm} u\left(x_{0}\right)=\left\{p \in \mathbb{R}, \text { there exists a sequence } p^{k} \in D_{\Omega}^{ \pm} u\left(x^{k}\right) \text { with } x^{k} \in \Omega \text { and }\left(x^{k}, p^{k}\right) \rightarrow\left(x_{0}, p\right)\right\} . \tag{2.2}
\end{equation*}
$$

Remark 2.2. Note that if $p \in \bar{D}_{\Omega}^{+} u\left(x_{0}\right)$ with $x_{0} \in \partial \Omega$, and if $u$ is a subsolution of $H(D u) \leq 0$ in $\Omega$, then $H(p) \leq 0$.

We then have the following result.
Lemma 2.3 (Critical slopes and semi-differentials). Let $\Omega:=(0,+\infty)$. We consider two functions $u, v$ : $\bar{\Omega} \rightarrow \mathbb{R} \cup\{-\infty,+\infty\}$ with $u$ upper semi-continuous and $v$ lower semicontinous satisfying $u(0)=0=v(0)$ with $u \leq v$ on $\bar{\Omega}$. We define the critical slopes by

$$
\begin{equation*}
\bar{p}:=\limsup _{\Omega \ni x \rightarrow 0} \frac{u(x)}{x}, \quad \underline{p}:=\liminf _{\Omega \ni x \rightarrow 0} \frac{v(x)}{x} . \tag{2.3}
\end{equation*}
$$

Then we have the following (limiting) semi-differential inclusions

$$
\begin{equation*}
\mathbb{R} \cap[\underline{p}, \bar{p}] \subset \bar{D}_{\Omega}^{+} u(0) \cap \bar{D}_{\Omega}^{-} v(0) \quad \text { if } \quad \bar{p} \geq \underline{p} \tag{2.4}
\end{equation*}
$$

$$
\begin{gather*}
\mathbb{R} \cap[\bar{p}, \underline{p}] \subset D_{\bar{\Omega}}^{+} u(0) \cap D_{\bar{\Omega}}^{-} v(0) \quad \text { if } \quad \bar{p} \leq \underline{p}  \tag{2.5}\\
\left\{\begin{array}{l}
\bar{p} \in \bar{D}_{\Omega}^{+} u(0) \\
\underline{p} \in \bar{D}_{\Omega}^{-} u(0)
\end{array} \quad \text { if } \quad \text { if } \quad \underline{p} \neq+\infty . \infty\right. \tag{2.6}
\end{gather*}
$$

Proof. The proof of this lemma is already contained in [5] but for sake of completeness, we give it here. We first notice that (2.5) is a straightforward consequence of the definition of sub and superdifferentials. In order to prove (2.4), we first focus on the proof of

$$
\begin{equation*}
\mathbb{R} \cap[\underline{p}, \bar{p}] \subset \bar{D}_{\Omega}^{+} u(0) \quad \text { in case } \quad \bar{p}>\underline{p} \tag{2.7}
\end{equation*}
$$

and we will even show the follower better result

$$
\begin{equation*}
\mathbb{R} \cap[\underline{q}, \bar{p}] \subset \bar{D}_{\Omega}^{+} u(0) \quad \text { in case } \quad \bar{p}>\underline{q}:=\liminf _{\Omega \ni x \rightarrow 0} \frac{u(x)}{x} \tag{2.8}
\end{equation*}
$$

Note that $u \leq v$ implies $\underline{q} \leq \underline{p}$ and so (2.7) is a consequence of (2.8). The claim is a variant of (18) in [11] and the proof is a variant of the one done in Barles, Chasseigne [3, Lemma 15.3.1]. We give the details for sake of completeness. We first assume that $p \in(\underline{q}, \bar{p})$. This implies that

$$
\limsup _{\Omega \ni x \rightarrow 0} \frac{u(x)}{x}=\bar{p}>p>\underline{q}=\liminf _{\Omega \ni x \rightarrow 0} \frac{u(x)}{x}
$$

and so for any $\varepsilon>0$, there exists $y_{\varepsilon} \in(0, \varepsilon)$ and $z_{\varepsilon} \in\left(0, y_{\varepsilon}\right)$ such that

$$
\frac{u\left(z_{\varepsilon}\right)}{z_{\varepsilon}}>p>\frac{u\left(y_{\varepsilon}\right)}{y_{\varepsilon}}
$$

Hence the function $\zeta(x):=u(x)-p x$ satisfies

$$
\zeta(0)=0>\zeta\left(y_{\varepsilon}\right) \quad \text { with } \quad M:=\sup _{\left[0, y_{\varepsilon}\right]} \zeta \geq \zeta\left(z_{\varepsilon}\right)>0
$$

Let $x_{\varepsilon} \in\left(0, y_{\varepsilon}\right)$ be a point of maximum of $\zeta$ in $\left[0, y_{\epsilon}\right]$. We see that the function $x \mapsto p x+M$ is a test function touching $u$ from above at $x_{\varepsilon}$, which implies that $p \in D_{\Omega}^{+} u\left(x_{\varepsilon}\right)$. In the limit $\varepsilon \rightarrow 0$, we recover $p \in \bar{D}_{\Omega}^{+} u(0)$ which proves the claim. In the case where $p \in[\underline{q}, \bar{p}]$, we get the result by the closedness of $\bar{D}_{\Omega}^{+} u(0)$. This proves (2.8). A similar inclusion for $v$ implies ( $\overline{2} .4$ ) in the special case where $\bar{p}>\underline{p}$. On the other hand, notice that (2.6) implies (2.4) in the case $\bar{p}=\underline{p}$.

Hence it remains to show (2.6). We claim that

$$
\begin{equation*}
\underline{p} \in \bar{D}_{\Omega}^{-} v(0) \quad \text { if } \quad \underline{p} \in \mathbb{R} . \tag{2.9}
\end{equation*}
$$

This result is a property of the critical slope for lower semi-continuous functions. Its proof follows exactly the lines of [7, Proof of Lemma 2.9] (where the proof does not use any Hamiltonian). A similar result holds for $u$ and proves (2.6). This ends the proof of the lemma.

Before to state the fundamental lemma for the comparison principle, we recall the definition of (semi-) coercive functions.

Definition 2.4 (Coercive and semi-coercive functions). Consider a function $G: \mathbb{R} \rightarrow \mathbb{R}$. Then $G$ is coercive if $\lim _{|p| \rightarrow+\infty} G(p)=+\infty$, and semi-coercive if $\lim _{p \rightarrow-\infty} G(p)=+\infty$.

As a consequence of Lemma 2.3, we have the following result which will be used to prove the comparison principle.

Corollary 2.5 (Boundary viscosity inequalities). Let $\Omega$ and $u, v$ be as in statement of Lemma 2.3. For $\gamma=\alpha, \beta$, consider continuous functions $H_{\gamma}, F_{\gamma}: \mathbb{R} \rightarrow \mathbb{R}$ with $H_{\alpha}$ coercive and $F_{\alpha}$ semi-coercive. Assume that we have the following viscosity inequalities for some $\eta>0$

$$
\left\{\begin{align*}
H_{\alpha}\left(u_{x}\right) \leq 0 & \text { on } & \Omega & \cap\{|u|<+\infty\}  \tag{2.10}\\
\min \left\{F_{\alpha}, H_{\alpha}\right\}\left(u_{x}\right) \leq 0 & \text { on } & \{0\} & \cap\{|u|<+\infty\} \\
H_{\beta}\left(v_{x}\right) \geq \eta & \text { on } & \Omega & \cap\{|v|<+\infty\} \\
\max \left\{F_{\beta}, H_{\beta}\right\}\left(v_{x}\right) \geq \eta & \text { on } & \{0\} & \cap\{|v|<+\infty\}
\end{align*}\right.
$$

For $\underline{p}, \bar{p}$ defined in (2.3), we set $a:=\min \{\underline{p}, \bar{p}\}$ and $b:=\max \{\underline{p}, \bar{p}\}$. Then $\bar{p} \in[a, b] \cap \mathbb{R}$ and there exists a real $\bar{n}$ umber $p \in[a, b]$ such that

$$
\begin{equation*}
\text { either } \quad H_{\alpha}(p) \leq 0<\eta \leq\left(H_{\beta}-H_{\alpha}\right)(p) \quad \text { or } \quad \max \left\{F_{\alpha}, H_{\alpha}\right\}(p) \leq 0<\eta \leq\left(F_{\beta}-F_{\alpha}\right)(p) \tag{2.11}
\end{equation*}
$$

Proof. The main steps of the proof is given in [5], but for sake of completeness, we give all the details here. We begin to explain why $\bar{p} \in \mathbb{R}$. Because $H_{\alpha}$ is coercive and $F_{\alpha}$ is semi-coercive, we know from [4, Lemma 3.8] that $u$ is weakly continuous at $x=0$, i.e.

$$
\begin{equation*}
0=u(0)=\limsup _{\Omega \ni x \rightarrow 0^{+}} u(x) \tag{2.12}
\end{equation*}
$$

Then [7, Proof of Lemma 2.10] shows additionally that $\bar{p}>-\infty$. Now we claim that we also have $\bar{p}<+\infty$. Indeed, assume by contradiction that $\bar{p}=+\infty$. Then, there exists $y_{n} \rightarrow 0$ such that $p_{n}:=u\left(y_{n}\right) / y_{n} \rightarrow+\infty$. For $b \in \mathbb{R}$, let us define $\phi_{b}(x):=p_{n} x+b$ and

$$
\bar{b}=\inf \left\{b, u \leq \phi_{b} \text { in }\left[0, y_{n}\right]\right\}
$$

In particular, there exists $x_{n} \in\left[0, y_{n}\right]$ such that $\phi_{\bar{b}}$ touches $u$ from above at $x_{n}$. If $x_{n}=0$, then $0=$ $u(0)-\phi_{\bar{b}}(0)=-\bar{b}$. In the same way, if $x_{n}=y_{n}$, then $u\left(y_{n}\right)=\phi_{\bar{b}}\left(y_{n}\right)=u\left(y_{n}\right)+\bar{b}$ and we recover again that $\bar{b}=0$. This implies that $u(x) \leq p_{n} x$ and so

$$
+\infty=\limsup _{x \rightarrow 0} \frac{u(x)}{x} \leq p_{n}<+\infty
$$

We then deduce that $x_{n} \in\left(0, y_{n}\right)$ and since $u$ is a sub-solution, we get

$$
H_{\alpha}\left(p_{n}\right) \leq 0
$$

which is absurd for $n$ large enough by coercivity of $H_{\alpha}$. This implies that $\bar{p}<+\infty$. We conclude that $\bar{p} \in \mathbb{R} \cap[a, b]$.

We now turn to the proof of (2.11). If $\underline{p} \leq \bar{p}$, then (2.4) shows, for all $p \in[\underline{p}, \bar{p}] \cap \mathbb{R}$, that

$$
H_{\alpha}(p) \leq 0<\eta \leq H_{\beta}(p)
$$

which implies in particular the desired conclusion.
We now assume that $\underline{p}>\bar{p}$. We have in particular $[a, b] \subset(-\infty,+\infty]$ with $a<b$ and

$$
\left\{\begin{array}{rlll}
H_{\alpha}(a) \leq 0 & \text { because } & a \in \mathbb{R}  \tag{2.13}\\
& 0<\eta \leq H_{\beta}(b) & \text { if } & b \in \mathbb{R} \\
\min \left\{H_{\alpha}, F_{\alpha}\right\} \leq 0<\eta \leq \max \left\{H_{\beta}, F_{\beta}\right\} & \text { on } & {[a, b] \cap \mathbb{R}}
\end{array}\right.
$$

where the last line follows from (2.5), and the first two lines follow from (2.6).
We now claim that for all $\varepsilon>0$ small enough, there exists some $p_{\varepsilon} \in[a, b] \cap \mathbb{R}$ such that we have at $p_{\varepsilon}$

$$
\begin{equation*}
\text { i) } \quad H_{\alpha} \leq \varepsilon<\eta-\varepsilon \leq H_{\beta}-H_{\alpha} \quad \text { or } \quad \text { ii) } \quad \max \left\{F_{\alpha}, H_{\alpha}\right\} \leq \varepsilon<\eta \leq F_{\beta}-F_{\alpha} \tag{2.14}
\end{equation*}
$$

By contradiction, we assume that there exists $\varepsilon>0$ (small enough) such that

$$
\left\{\left.\begin{array}{lll}
\text { i) } & H_{\beta}-H_{\alpha}<\eta-\varepsilon & \text { or } \varepsilon<H_{\alpha}  \tag{2.15}\\
\text { and } & & \text { for all } p \in[a, b] \cap \mathbb{R} . \\
\text { ii) } & F_{\beta}-F_{\alpha}<\eta \quad \text { or } & \varepsilon<\max \left\{F_{\alpha}, H_{\alpha}\right\}
\end{array} \right\rvert\, \quad\right. \text {. }
$$

Recall that the coercivity of $H_{\alpha}$ means $H_{\alpha}( \pm \infty):=\liminf _{p \rightarrow \pm \infty} H_{\alpha}(p)=+\infty$. We distinguish two cases.
Case 1: $H_{\alpha}(b)>\varepsilon$
Here $b$ can be finite or equal to $+\infty$. We get

$$
H_{\alpha}(b)>\varepsilon>0 \geq H_{\alpha}(a) .
$$

Therefore by continuity, there exists $p \in(a, b)$ such that $H_{\alpha}(p)=\varepsilon$. Hence in the last line of (2.13), the first inequality implies that $F_{\alpha}(p) \leq 0$. Because (2.15) i) and ii) hold true for $p$, we get

$$
H_{\beta}(p)<\eta \quad \text { and } \quad F_{\beta}(p)<\eta
$$

which leads to a contradiction with the second inequality in the last line of (2.13).
Case 2: $H_{\alpha}(b) \leq \varepsilon$
Then $b \in \mathbb{R}$ and (2.15) i) implies for $p=b$ that $H_{\beta}(b)<\eta$, which is in contradiction with the second line of (2.13).

In all the cases, we get a contradiction, which proves (2.14). Since $H_{\alpha}$ is coercive, we see in both cases i) or ii) of (2.14), that we can always extract a subsequence as $\varepsilon \rightarrow 0$ such that $p_{\varepsilon} \rightarrow p \in[a, b] \cap \mathbb{R}$. Passing to the limit in (2.14), we get the desired conclusion. This ends the proof of the corollary.

Notice that it is very easy to show the following extension of Corollary 2.5 to the case of junctions.
Proposition 2.6 (Junction viscosity inequalities). For $N \geq 1$, let $J^{i}:=(0,+\infty)$ for $i=1, \ldots, N$, and set

$$
J:=\{0\} \cup\left(\bigcup_{i=1, \ldots, N} J^{i}\right)
$$

with the topology of glued branches. For a piecewise $C^{1}$ function $u$ on $J$, and $u^{i}:=u_{\mid J^{i} \cup\{0\}}$, we set

$$
u_{x}(x)= \begin{cases}\left(u_{x}^{1}(0), \ldots, u_{x}^{N}(0)\right) & \text { if } x=0 \\ u_{x}^{i}(x) & \text { if } x \in J^{i}\end{cases}
$$

We consider two sets of functions $u, v: J \rightarrow \mathbb{R} \cup\{-\infty,+\infty\}$ with $u$ upper semi-continuous and $v$ lower semicontinous satisfying

$$
\begin{equation*}
u(0)=0=v(0) \quad \text { with } \quad u \leq v \quad \text { on } \quad J . \tag{2.16}
\end{equation*}
$$

For $i=1, \ldots, N$, we define

$$
\begin{equation*}
\bar{p}^{i}:=\limsup _{J^{i} \ni x \rightarrow 0} \frac{u(x)}{x}, \quad \underline{p}^{i}:=\liminf _{J^{i} \ni x \rightarrow 0} \frac{v(x)}{x} . \tag{2.17}
\end{equation*}
$$

We also set $a^{i}:=\min \left\{\underline{p}^{i}, \bar{p}^{i}\right\}, b^{i}:=\max \left\{\underline{p}^{i}, \bar{p}^{i}\right\}$ and

$$
[a, b] \cap \mathbb{R}^{N}:=\prod_{i=1, \ldots, N}\left[a^{i}, b^{i}\right] \cap \mathbb{R}
$$

For $\gamma=\alpha, \beta$, consider continuous functions $H_{\gamma}^{i}: \mathbb{R} \rightarrow \mathbb{R}$ and $F_{\gamma}: \mathbb{R}^{N} \rightarrow \mathbb{R}$ with $H_{\alpha}^{i}$ coercive and $F_{\alpha}$ semi-coercive. For $p=\left(p^{1}, \ldots, p^{N}\right) \in \mathbb{R}^{N}$, we set

$$
H_{\gamma ; \min }(p)=\min _{i=1, \ldots, N} H_{\gamma}^{i}\left(p^{i}\right), \quad H_{\gamma ; \max }(p)=\max _{i=1, \ldots, N} H_{\gamma}^{i}\left(p^{i}\right)
$$

Then assume that we have the following viscosity inequalities for some $\eta>0$

$$
\left\{\begin{align*}
H_{\alpha}^{i}\left(u_{x}\right) \leq 0 & \text { on } & J^{i} & \cap\{|u|<+\infty\}  \tag{2.18}\\
\min \left\{F_{\alpha}, H_{\alpha ; \min }\right\}\left(u_{x}\right) \leq 0 & \text { on } & \{0\} & \cap\{|u|<+\infty\} \\
H_{\beta}^{i}\left(v_{x}\right) \geq \eta & \text { on } & J^{i} & \cap\{|v|<+\infty\} \\
\max \left\{F_{\beta}, H_{\beta ; \max }\right\}\left(v_{x}\right) \geq \eta & \text { on } & \{0\} & \cap\{|v|<+\infty\} .
\end{align*} \quad \text { for } i=1, \ldots, N\right.
$$

Then there exists $p=\left(p^{1}, \ldots, p^{N}\right) \in[a, b] \cap \mathbb{R}^{N} \neq \emptyset$ such that

$$
\left\{\begin{array}{rll}
\text { either } & H_{\alpha}^{i}\left(p^{i}\right) & \leq 0<\eta \leq\left(H_{\beta}^{i}-H_{\alpha}^{i}\right)\left(p^{i}\right) \quad \text { for some } i \in\{1, \ldots, N\}  \tag{2.19}\\
\text { or } & \max \left(F_{\alpha}, H_{\alpha ; \max }\right)(p) & \leq 0<\eta \leq\left(F_{\beta}-F_{\alpha}\right)(p) .
\end{array}\right.
$$

## 3 Barriers and regularization

Lemma 3.1 (Barriers). Let $T>0$ and assume that $H, F$ satisfy (1.4)-(1.5), and that the initial data $u_{0}$ is uniformly continuous. Assume that $u$ (resp. v) is an upper semi-continuous subsolution (resp. a lower semi-continuous supersolution) of (1.1), (1.3), satisfying the a priori bounds (1.6) for some constant $C_{T}$.

Then there exists a continuous increasing function $f:[0, T] \rightarrow[0,+\infty)$ with $f(0)=0$ such that the functions

$$
u^{ \pm}(t, x):=u_{0}(x) \pm f(t)
$$

satisfy the following barrier properties:

- if $u \leq u_{0}$ in $\{0\} \times \bar{\Omega}$, then $u \leq u^{+}$in $[0, T) \times \bar{\Omega}$,
- if $v \geq u_{0}$ in $\{0\} \times \bar{\Omega}$, then $v \geq u^{-}$in $[0, T) \times \bar{\Omega}$.

Proof. The idea of the proof is somehow very standard. We first extend by continuity the initial data defined on $\mathbb{R}^{d-1} \times[0,+\infty)$ to a function defined on $\mathbb{R}^{d-1} \times \mathbb{R}$, setting $u_{0}\left(x^{\prime}, x\right):=u_{0}\left(x^{\prime}, 0\right)$ for all $x \leq 0$. Hence $u_{0}$ still satisfies

$$
\left|u_{0}(x)-u_{0}(y)\right| \leq \omega_{0}(|x-y|) \quad \text { for all } \quad x, y \in \mathbb{R}^{d}
$$

where $\omega_{0}$ is the modulus of continuity of $\left.\left(u_{0}\right)\right|_{\mathbb{R}^{d-1} \times[0,+\infty)}$. We do the proof to compare $u$ and $u^{+}$, the one to compare $v$ and $u^{-}$being similar.
Case 1: $u_{0} \in\left(C^{1} \cap \operatorname{Lip}\right)\left(\mathbb{R}^{d}\right)$
In this case, there exists some $L>0$ such that $\left|D u_{0}\right|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \leq L$. From assumptions i) and ii) of both (1.4) and (1.5), we see that there exists $\lambda=\lambda(T, L) \geq 0$ minimal such that, for $B_{L}:=B_{L}(0)$,

$$
\begin{equation*}
-\lambda \leq \inf _{[0, T] \times \bar{\Omega} \times \bar{B}_{2 L}} \min \{H, F\} \leq \sup _{[0, T] \times \bar{\Omega} \times \bar{B}_{2 L}} \max \{H, F\} \leq \lambda \tag{3.1}
\end{equation*}
$$

where we have extended the function $F$ as follows: $F\left(t, x^{\prime}, x_{d}, p\right):=F\left(t, x^{\prime}, 0, p\right)$ for all $x_{d} \geq 0$. Setting $f(t):=\lambda t$, we see that $u^{+}$is a supersolution of (1.1), (1.3). Assume now by contradiction that

$$
M:=\sup _{Q_{T}}\left(u-u^{+}\right)>0 \quad \text { with } \quad Q_{T}:=[0, T) \times \bar{\Omega}
$$

Now for $\eta, \alpha>0$, let us consider

$$
M_{\eta, \alpha}:=\sup _{Q_{T}} \Phi \quad \text { with } \quad \Phi(t, x):=u(t, x)-u^{+}(t, x)-\frac{\eta}{T-t}-\frac{\alpha}{2} x^{2}
$$

For $\eta, \alpha>0$ small enough, we have $M_{\eta, \alpha} \geq M / 2>0$. Moreover from the bound (1.6) on $u$, we see that the supremum in $M_{\eta, \alpha}$ is reached for some point $\bar{X}=(\bar{t}, \bar{x}) \in Q_{T}$. We also have

$$
\limsup _{(\eta, \alpha) \rightarrow(0,0)}\left\{\frac{\eta}{T-\bar{t}}+\frac{\alpha}{2} \bar{x}^{2}\right\}=0
$$

and then we can fix $\eta, \alpha>0$ small enough such that $|\alpha \bar{x}| \leq L$.
Assume that $\bar{t}=0$. Then $0<M / 2 \leq M_{\eta, \alpha}=\Phi(0, \bar{x}) \leq-\frac{\eta}{T}$ which leads to a contradiction. Hence $\bar{t}>0$ and $\bar{X}=(\bar{t}, \bar{x})=\left(\bar{t}, \bar{x}^{\prime}, \bar{x}_{d}\right) \in(0, T) \times \mathbb{R}^{d-1} \times[0,+\infty)$. Therefore we have the viscosity inequalities for $p:=D u_{0}(\bar{x}) \in \bar{B}_{L}$ and $\alpha \bar{x} \in \bar{B}_{L}$

$$
\left\{\begin{array}{lll}
\frac{\eta}{(T-t)^{2}}+\lambda+H(\bar{X}, p+\alpha \bar{x}) \leq 0 & \text { if } \quad \bar{x}_{d}>0 \\
\frac{\eta}{(T-t)^{2}}+\lambda+\max (F, H)(\bar{X}, p+\alpha \bar{x}) \leq 0 & \text { if } \quad \bar{x}_{d}=0
\end{array}\right.
$$

which leads to a contradiction from the choice of $\lambda$ in (3.1). This implies that $M \leq 0$ and then $u \leq u^{+}$.
Case 2: $u_{0}$ is only uniformly continuous
Let $\varphi$ be a smooth nonnegative function satisfying $\varphi=0$ on $\mathbb{R}^{n} \backslash B_{1}(0)$ and $\int_{\mathbb{R}^{n}} \varphi(x) d x=1$. For $\varepsilon>0$, we set the convolution $u_{0}^{\varepsilon}:=\varphi_{\varepsilon} \star u_{0}$ with $\varphi_{\varepsilon}(x)=\frac{1}{\varepsilon^{n}} \varphi\left(\frac{x}{\varepsilon}\right)$. Hence we have

$$
\left|u_{0}^{\varepsilon}-u_{0}\right|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \leq \omega_{0}(\varepsilon)
$$

and $D u_{0}^{\varepsilon}(x)=\frac{1}{\varepsilon} \int_{\mathbb{R}^{n}} d y \frac{1}{\varepsilon^{n}} D \varphi\left(\frac{y}{\varepsilon}\right)\left\{u_{0}(x-y)-u_{0}(x)\right\}$. Therefore, we get

$$
\left|D u_{0}^{\varepsilon}\right|_{L^{\infty}(\mathbb{R})} \leq|D \varphi|_{L^{1}(\mathbb{R})} \cdot \frac{\omega_{0}(\varepsilon)}{\varepsilon} \leq L_{\varepsilon}:=|D \varphi|_{L^{1}\left(\mathbb{R}^{n}\right)} \cdot \sup _{\delta \geq \varepsilon} \frac{\omega_{0}(\delta)}{\delta}
$$

We define $\lambda_{\varepsilon}=\lambda_{T, L_{\varepsilon}} \geq 0$ minimal such that

$$
-\lambda_{\varepsilon} \leq \inf _{[0, T] \times \bar{\Omega} \times \overline{B_{2 L_{\varepsilon}}}} \min (H, F) \leq \sup _{[0, T] \times \bar{\Omega} \times \overline{B_{2 L_{\varepsilon}}}} \max (H, F) \leq \lambda_{\varepsilon}
$$

where by construction the map $\varepsilon \mapsto \lambda_{\varepsilon}$ is nonincreasing, and we set $f_{\varepsilon}(t):=\lambda_{\varepsilon} t$. Using that $u_{0}(x) \leq$ $u_{0}^{\varepsilon}(x)+\omega_{0}(\varepsilon)$, we can show as in Case 1 that $u(t, x) \leq u_{0}^{\varepsilon}(x)+\omega_{0}(\varepsilon)+\lambda_{\varepsilon} t \leq u_{0}(x)+2 \omega_{0}(\varepsilon)+\lambda_{\varepsilon} t$. If we set

$$
f(t):=\inf _{\varepsilon>0}\left\{2 \omega_{0}(\varepsilon)+\lambda_{\varepsilon} t\right\}
$$

where $f$ is a (continuous) concave nondecreasing function satisfying $f(0)=0$, we get $u \leq u^{+}$. This ends the proof of the lemma.

We now consider a (classical) regularization of a subsolution $u$ by tangential sup-convolutions. Because we only assume a bound from above $u \leq u_{0}+C_{T}$, we have additionally to truncate $u$ from below by some function. We will use the function $\underline{u}_{0}:=u_{0}-C_{T}$ (which will also be later in the next section a bound from below for the supersolution $v$ ), where $u_{0}$ is the initial data, which is assumed to be Lipschitz continuous, in order to simplify the presentation. Then we have the following result.

Lemma 3.2 (Tangential regularization after truncation by Lipschitz initial data). Let $T>0$ and assume that $H$ satisfies (1.4) and that the initial data $u_{0}$ is Lipschitz continuous of Lipschitz constant $L_{0}$. Let $u$ be an upper semi-continuous subsolution of (1.1)-(1.2), satisfying moreover the a priori bound (1.6) for some constant $C_{T}$, namely

$$
\begin{equation*}
u(t, x) \leq u_{0}(x)+C_{T} \quad \text { for all } \quad(t, x) \in[0, T) \times \bar{\Omega} \tag{3.2}
\end{equation*}
$$

We define $u(T, x):=\limsup _{(s, y) \rightarrow(T, x), s<T} u(s, y)$ for all $x \in \bar{\Omega}=\mathbb{R}^{d-1} \times[0,+\infty)$, and extend $u$ to $\mathbb{R} \times \bar{\Omega}$, setting $u(t, x):=u(T, x)$ if $t \geq T$, and $u(t, x):=u(0, x)$ if $t \leq 0$. We set

$$
\tilde{u}:=\max \left(u, \underline{u}_{0}\right) \quad \text { with } \quad \underline{u}_{0}(t, x):=u_{0}(x)-C_{T} .
$$

We denote the tangential variable by $\xi=\left(\xi^{0}, \xi^{\prime}\right)=\left(s, x^{\prime}\right) \in \mathbb{R}^{d}$ and the normal variable by $x_{d} \in[0,+\infty)$ and we define for $\nu>0$ the tangential sup-convolution

$$
\tilde{u}^{\nu}\left(\xi, x_{d}\right):=\sup _{\zeta \in \mathbb{R}^{d}}\left\{\tilde{u}\left(\zeta, x_{d}\right)-\frac{|\xi-\zeta|^{2}}{2 \nu}\right\}=\tilde{u}\left(\bar{\zeta}, x_{d}\right)-\frac{|\xi-\bar{\zeta}|^{2}}{2 \nu}
$$

where each $\bar{\zeta}$ depends on $\left(\xi, x_{d}\right) \in \mathbb{R}^{d} \times[0,+\infty)$ with $|\bar{\zeta}-\xi| \leq \theta^{\nu}:=\sqrt{5\left(4 \nu C_{T}+\nu^{2} L_{0}^{2}\right)}<T / 2$, for $\nu$ small enough.

Then the function $\tilde{u}^{\nu}$ is Lipschitz continuous in $\mathbb{R} \times \bar{\Omega}$ with respect to the variable $\xi$. Moreover it is Lipschitz continuous in $I^{\nu} \times \Omega$ with respect to the variable $x_{d}$, with $I^{\nu}:=\left(\theta^{\nu}, T-\theta^{\nu}\right)$,

$$
\left|D_{\xi} \tilde{u}^{\nu}\right|_{L^{\infty}(\mathbb{R} \times \bar{\Omega})} \leq \frac{\theta^{\nu}}{\nu} \quad \text { and } \quad\left|\partial_{x_{d}} \tilde{u}^{\nu}\right|_{L^{\infty}\left(I^{\nu} \times \Omega\right)} \leq \max \left\{L^{\nu}, L_{0}\right\}
$$

where $L^{\nu}:=\sup \left\{p_{d} \in \mathbb{R}, \quad \inf _{\left(X, p^{\prime}\right) \in([0, T] \times \bar{\Omega}) \times \bar{B}_{\frac{\theta^{\nu}}{\nu}}} H\left(X, p^{\prime}, p_{d}\right) \leq \frac{\frac{\theta}{}^{\nu}}{\nu}\right\}$.
Assume furthermore that $u$ is a subsolution at the boundary $(0, T) \times \partial \Omega$, i.e. satisfies the first line of (1.3) for some $F$ satisfying (1.5). Then $\tilde{u}^{\nu}$ is Lipschitz continuous in space and time on $\overline{I^{\nu} \times \Omega}$ of Lipschitz constant $L_{\nu}:=\max \left\{\frac{\theta^{\nu}}{\nu}, L^{\nu}, L_{0}\right\}$.
Proof. The proof is splited into three steps.

## Step 1: first bounds using the 2-sided bound

We begin to show that

$$
\begin{equation*}
|\bar{\zeta}-\xi| \leq \theta^{\nu} \tag{3.3}
\end{equation*}
$$

From the 1 -sided bound (3.2) and the definition of $\underline{u}_{0}$ and $\tilde{u}$, we get the 2 -sided bound

$$
\begin{equation*}
\left|\tilde{u}(t, x)-u_{0}(x)\right| \leq C_{T} \quad \text { for all } \quad(t, x) \in \mathbb{R} \times \bar{\Omega} \tag{3.4}
\end{equation*}
$$

For $\xi=\left(t, \xi^{\prime}\right):=\left(t, x^{\prime}\right)$ and $z:=x_{d}$ (to simplify the notation), we have

$$
\tilde{u}(\xi, z) \leq \tilde{u}^{\nu}(\xi, z)=\tilde{u}(\bar{\zeta}, z)-\frac{|\xi-\bar{\zeta}|^{2}}{2 \nu}
$$

with $\bar{\zeta}=\left(\bar{t}, \bar{\zeta}^{\prime}\right)$. Using (3.4), we then get

$$
\frac{|\xi-\bar{\zeta}|^{2}}{2 \nu}-2 C_{T} \leq u_{0}\left(\bar{\zeta}^{\prime}, z\right)-u_{0}\left(\xi^{\prime}, z\right) \leq L_{0}\left|\bar{\zeta}^{\prime}-\xi^{\prime}\right|
$$

This implies

$$
\left\{\left|\xi^{\prime}-\bar{\zeta}^{\prime}\right|-\nu L_{0}\right\}^{2}+|t-\bar{t}|^{2} \leq 4 \nu C_{T}+\nu^{2} L_{0}^{2}=\frac{\left(\theta^{\nu}\right)^{2}}{5}
$$

We then deduce that $|t-\bar{t}| \leq \frac{\theta^{\nu}}{\sqrt{5}},\left|\xi^{\prime}-\bar{\zeta}^{\prime}\right| \leq 2 \frac{\theta^{\nu}}{\sqrt{5}}$ which implies (3.3).
We now prove that

$$
\begin{equation*}
\left|D_{\xi} \tilde{u}^{\nu}\right|_{L^{\infty}(\mathbb{R} \times \bar{\Omega})} \leq \frac{\theta^{\nu}}{\nu} \tag{3.5}
\end{equation*}
$$

For $\xi^{a} \in \mathbb{R}^{d}$, we set

$$
\tilde{u}^{\nu}\left(\xi^{a}, z\right):=\sup _{\zeta \in \mathbb{R}^{d}}\left\{\tilde{u}(\zeta, z)-\frac{\left|\xi^{a}-\zeta\right|^{2}}{2 \nu}\right\}=\tilde{u}\left(\bar{\zeta}^{a}, z\right)-\frac{\left|\xi^{a}-\bar{\zeta}^{a}\right|^{2}}{2 \nu}
$$

Hence, by definition, we have

$$
\tilde{u}^{\nu}\left(\xi^{a}, z\right) \geq \tilde{u}(\bar{\zeta}, z)-\frac{\left|\left(\xi^{a}-\xi\right)+\xi-\bar{\zeta}\right|^{2}}{2 \nu}=\tilde{u}^{\nu}(\xi, z)-\left(\xi^{a}-\xi\right) \cdot \frac{(\xi-\bar{\zeta})}{\nu}-\frac{\left|\xi^{a}-\xi\right|^{2}}{2 \nu}
$$

and also by symmetry

$$
\tilde{u}^{\nu}(\xi, z) \geq \tilde{u}^{\nu}\left(\xi^{a}, z\right)-\left(\xi-\xi^{a}\right) \cdot \frac{\left(\xi^{a}-\bar{\zeta}^{a}\right)}{\nu}-\frac{\left|\xi-\xi^{a}\right|^{2}}{2 \nu}
$$

i.e.

$$
\frac{\left|\tilde{u}^{\nu}\left(\xi^{a}, x\right)-\tilde{u}^{\nu}(\xi, x)\right|}{\left|\xi^{a}-\xi\right|} \leq \max \left\{\frac{|\xi-\bar{\zeta}|}{\nu}, \frac{\left|\xi^{a}-\bar{\zeta}^{a}\right|}{\nu}\right\}+\frac{\left|\xi^{a}-\xi\right|}{2 \nu} \leq \frac{\theta^{\nu}}{\nu}+\frac{\left|\xi^{a}-\xi\right|}{2 \nu}
$$

This implies that $\tilde{u}^{\nu}$ is Lipschitz continuous in the tangential coordinates and (3.5).

## Step 2: bounds on the normal gradient

It is easy to check that $\tilde{u}^{\nu}$ is upper semi-continuous (because this is the case for $u$ itself and the supremum in $\xi$ is locally taken in a compact set). Let $\varphi$ be a test function touching $\tilde{u}^{\nu}$ from above at $X_{0}:=\left(t_{0}, x_{0}\right) \in$ $\left(I^{\nu} \times \Omega\right) \cap\left\{\tilde{u}^{\nu}>\underline{u}_{0}^{\nu}\right\}$ and set $\xi_{0}:=\left(t_{0}, x_{0}^{\prime}\right)$ and $z_{0}:=\left(x_{0}\right)_{d}$. We have

$$
\underline{u}_{0}^{\nu}\left(X_{0}\right)<\tilde{u}^{\nu}\left(X_{0}\right):=\sup _{h \in \mathbb{R}^{d}}\left\{\tilde{u}\left(\xi_{0}+h, z_{0}\right)-\frac{|h|^{2}}{2 \nu}\right\}=\tilde{u}\left(\xi_{0}+\bar{h}_{0}, z_{0}\right)-\frac{\left|\bar{h}_{0}\right|^{2}}{2 \nu} \quad \text { with } \quad\left|\bar{h}_{0}\right| \leq \theta^{\nu}
$$

and

$$
\tilde{u}\left(\xi+\bar{h}_{0}, z\right)-\frac{\left|\bar{h}_{0}\right|^{2}}{2 \nu} \leq \tilde{u}^{\nu}(\xi, z) \leq \varphi(\xi, z) \quad \text { with equality at }(\xi, z)=\left(\xi_{0}, z_{0}\right)
$$

Setting

$$
\bar{\varphi}(\xi, z):=\varphi\left(\xi-\bar{h}_{0}, z\right)+\frac{\left|\bar{h}_{0}\right|^{2}}{2 \nu}, \quad \bar{X}_{0}:=\left(\xi_{0}+\bar{h}_{0}, z_{0}\right)
$$

we then get

$$
\left\{\begin{array}{l}
\tilde{u} \leq \bar{\varphi} \quad \text { with equality at } \bar{X}_{0} \\
\tilde{u}\left(\bar{X}_{0}\right)=\tilde{u}^{\nu}\left(X_{0}\right)+\frac{\left|\bar{h}_{0}\right|^{2}}{2 \nu}>\underline{u}_{0}^{\nu}\left(X_{0}\right)+\frac{\left|\bar{h}_{0}\right|^{2}}{2 \nu}=\sup _{h \in \mathbb{R}^{d}}\left\{\underline{u}_{0}\left(\xi_{0}+h, z_{0}\right)-\frac{|h|^{2}}{2 \nu}\right\}+\frac{\left|\bar{h}_{0}\right|^{2}}{2 \nu} \geq \underline{u}_{0}\left(\bar{X}_{0}\right)
\end{array}\right.
$$

Hence $\tilde{u}=u$ at $\bar{X}_{0} \in\left(\left[-\left|\bar{h}_{0}\right|,\left|\bar{h}_{0}\right|\right]+I^{\nu}\right) \times \Omega \subset(0, T) \times \Omega$. Because $u \leq \tilde{u}$, we have that $\bar{\varphi}$ touches $u$ from above at $\bar{X}_{0}$ and since $u$ satisfies the viscosity inequalities on $(0, T) \times \Omega$, we get

$$
\bar{\varphi}_{t}\left(\bar{X}_{0}\right)+H\left(\bar{X}_{0}, D \bar{\varphi}\left(\bar{X}_{0}\right)\right) \leq 0
$$

i.e.

$$
\varphi_{t}\left(X_{0}\right)+H\left(\bar{X}_{0}, D \varphi\left(X_{0}\right)\right) \leq 0
$$

Setting for $\theta>0$ and $X=\left(t, x^{\prime}, x_{d}\right)$

$$
H_{\theta}(X, p):=\min _{h \in \bar{B}_{\theta}} H\left(\left(t, x^{\prime}\right)+h, x_{d}, p\right)
$$

we see that we have

$$
\varphi_{t}\left(X_{0}\right)+H_{\theta^{\nu}}\left(X_{0}, D \varphi\left(X_{0}\right)\right) \leq 0
$$

which shows that $\tilde{u}^{\nu}$ satisfies this viscosity inequality on $\left(I^{\nu} \times \Omega\right) \cap\left\{\tilde{u}^{\nu}>\underline{u}_{0}^{\nu}\right\}$.
Recall that we have $\tilde{u} \geq \underline{u}_{0}$ and so $\tilde{u}^{\nu} \geq \underline{u}_{0}^{\nu}$ with $\left|D \underline{u}_{0}^{\nu}\right| \leq\left|D \underline{u}_{0}\right| \leq L_{0}$. Hence we get in the viscosity sense

$$
\min \left\{-\frac{\theta^{\nu}}{\nu}+H_{\theta^{\nu}}\left(X, D \tilde{u}^{\nu}\right),\left|D \tilde{u}^{\nu}\right|-L_{0}\right\} \leq 0 \quad \text { on } \quad I^{\nu} \times \Omega
$$

which implies that

$$
\left|\partial_{x_{d}} \tilde{u}^{\nu}\right| \leq \max \left\{L^{\nu}, L_{0}\right\} \quad \text { on } \quad I^{\nu} \times \Omega
$$

Step 4: Lispchitz bounds on $\left[\theta^{\nu}, T-\theta^{\nu}\right] \times \bar{\Omega}$
Because $H$ is coercive and $F$ is semi-coercive, we can use [4, Lemma 3.8] to get that $u$ is weakly continuous for all $(t, x) \in(0, T) \times \partial \Omega$, i.e.

$$
u^{*}(t, x)=\limsup _{(s, y) \rightarrow(t, x),} u(s, y)
$$

which is again the case for $\tilde{u}=\max \left\{u, \underline{u}_{0}\right\}$. By sup-convolution, it is then easy to check that this is also true for

$$
\tilde{u}^{\nu}\left(\xi, x_{d}\right)=\sup _{|\zeta-\xi| \leq \theta^{\nu}}\left\{\tilde{u}\left(\zeta, x_{d}\right)-\frac{|\xi-\zeta|^{2}}{2 \nu}\right\}
$$

at least for all $t \in I^{\nu}=\left(\theta^{\nu}, T-\theta^{\nu}\right)$. Because $\tilde{u}^{\nu}$ is uniformly Lipschitz continuous on $I^{\nu} \times \Omega$, we deduce that $\tilde{u}^{\nu}$ is also Lipschitz continuous on $I^{\nu} \times \bar{\Omega}$. Finally, because the bound on $\partial_{t} \tilde{u}^{\nu}$ is uniform in space and time, we deduce that $\tilde{u}^{\nu}$ is Lipschitz continuous on $\left[\theta^{\nu}, T-\theta^{\nu}\right] \times \bar{\Omega}$ with the same Lipschitz constants. This ends the proof of the lemma.

## 4 The comparison principle on a half space

This Section is devoted to the proof of the comparison principle Theorem 1.1.
Proof of Theorem 1.1. The strategy of the proof is similar to the one of the note [5] but need technical adaptations. We first follow the proof of the comparison principle in [7], but then modify the proof on the boundary, introducing the twin blow-ups method. Let $\eta, \theta>0$ and consider

$$
\begin{equation*}
M(\theta):=\sup \left\{\Psi\left(\xi, \zeta, x_{d}\right), \quad x_{d} \in[0,+\infty), \quad \xi, \zeta \in[0, T) \times \mathbb{R}^{n-1}, \quad|\xi-\zeta| \leq \theta\right\} \tag{4.1}
\end{equation*}
$$

with $\zeta=\left(s, x^{\prime}\right), x=\left(x^{\prime}, x_{d}\right)$ and

$$
\Psi\left(\xi, \zeta, x_{d}\right):=\tilde{u}\left(\xi, x_{d}\right)-v\left(\zeta, x_{d}\right)-\frac{\eta}{T-s}, \quad \tilde{u}=\max \left\{u, \underline{u}_{0}\right\}, \quad \underline{u}_{0}(t, x):=u_{0}(X)-C_{T}=: \underline{u}_{0}(x)
$$

where we choose carefully $\frac{\eta}{T-s}$ instead of $\frac{\eta}{T-t}$, because we want to do later a doubling of variables which looks like a sup-convolution (in particular in time) to the function $\tilde{u}$.
We want to prove that $M:=\lim _{\theta \rightarrow 0} M(\theta) \leq 0$. Assume by contradiction that

$$
\begin{equation*}
M>0 \tag{4.2}
\end{equation*}
$$

## Step 0. Reduction to $u_{0}$ Lipschitz continuous

By assumption, the initial data $u_{0}$ is uniformly continuous. We follow the line of Case 2 of the proof of Lemma 3.1. We first extend $u_{0}$ by the value $u_{0}\left(x^{\prime}, 0\right)$ for $z=x_{d} \leq 0$ and $x^{\prime} \in \mathbb{R}^{n-1}$. For the ball $B_{1}=B_{1}(0)$, we then consider a smooth nonnegative function $\varphi$ satisfying $\varphi=0$ on $\mathbb{R}^{n} \backslash B_{1}$ with $\int_{\mathbb{R}^{n}} \varphi(x) d x=1$, and for $\beta>0$, we set the convolution $u_{0, \beta}:=\varphi_{\beta} \star u_{0}$ with $\varphi_{\beta}(x)=\frac{1}{\beta^{n}} \varphi\left(\frac{x}{\beta}\right)$. Then we can insure that

$$
\left|u_{0, \beta}-u_{0}\right| \leq \omega_{0}(\beta), \quad\left|D u_{0, \beta}\right| \leq L_{\beta}
$$

where $\omega_{0}$ is the modulus of continuity of $u_{0}$ and $L_{\beta}$ is some constant.
We have in particular

$$
u-\omega_{0}(\beta) \leq u_{0, \beta} \leq v+\omega_{0}(\beta) \quad \text { on } \quad\{0\}_{t} \times \bar{\Omega}
$$

Hence the problem is reduced to replace the quantities $\left(u, u_{0}, v\right)$ by $\left(u-\omega_{0}(\beta), u_{0, \beta}, v+\omega_{0}(\beta)\right.$ and $M$ by $M-2 \omega_{0}(\beta)$, and keep $C_{T}$ unchanged. Therefore fixing some $\beta:=\beta_{1}>0$ small enough such that $2 \omega_{0}(\beta)<M$, we see that we can redefine $u, u_{0}, v$ and assume without loss of generality that $u_{0}$ is Lipschitz continuous, say (forgetting now $\beta$ ) for some Lipschitz constant $L_{0}$, with $M>0$.

## Step 1. Doubling naively the space variables

We first consider a space penalization and standard doubling of variables in space and time, but we distinguish the tangential variables from the normal variable. To this end, we introduce parameters $\alpha, \nu, \delta>0$ and set

$$
M_{\nu, \alpha, \delta}:=\sup _{(t, x),(s, y) \in[0, T) \times \bar{\Omega}} \Psi_{\nu, \alpha, \delta}(t, x, s, y)
$$

with for $\xi=\left(t, x^{\prime}\right), \zeta=\left(s, y^{\prime}\right), x=\left(x^{\prime}, x_{d}\right), y=\left(y^{\prime}, y_{d}\right)$

$$
\begin{equation*}
\Psi_{\nu, \alpha, \delta}(t, x, s, y):=\tilde{u}(t, x)-v(s, y)-\alpha g(y)-\frac{\eta}{T-s}-\frac{\left|\left(t, x^{\prime}\right)-\left(s, y^{\prime}\right)\right|^{2}}{2 \nu}-\frac{\left|x_{d}-y_{d}\right|^{2}}{2 \delta}, \quad g(y):=\frac{y^{2}}{2} \tag{4.3}
\end{equation*}
$$

which satisfies $\liminf _{\alpha \rightarrow 0}\left\{\liminf _{\delta \rightarrow 0} M_{\nu, \alpha, \delta}\right\} \geq M>0$. Hence we see that (independently on $\nu>0$ ) for $\alpha>0$ small enough, and for $\delta>0$ small enough (say $\delta \in\left(0, \delta_{\alpha}\right]$ ), we get

$$
\begin{equation*}
M_{\nu, \alpha, \delta} \geq M / 2>0 \tag{4.4}
\end{equation*}
$$

In particular, the maximum is reached at some point that $\left(\bar{X}_{\delta}, \bar{Y}_{\delta}\right)=\left(\bar{t}_{\delta}, \bar{x}_{\delta}, \bar{s}_{\delta}, \bar{y}_{\delta}\right)$ and we claim that we have the following estimate.

Lemma 4.1 (Bounds on any optimizing sequence). Given $T, C_{T}>0$, there exists $\eta>0$ small enough such that the following holds true. Let $X, Y \in[0, T) \times \bar{\Omega}$ with $X=(t, x), Y=(s, y)$ be such that

$$
0<\Psi_{\nu, \alpha, \delta}(X, Y)
$$

Then for $\nu, \delta>0$ small enough (depending on $\eta>0$ ), we have

$$
\left\{\begin{array}{l}
t, s \in\left[\tau_{\eta}, T-\tau_{\eta}\right], \quad \tau_{\eta}:=\frac{\eta}{4 C_{T}} \\
\alpha g(y)+\frac{\eta}{T-s}+\frac{|t-s|^{2}}{2 \nu} \leq 3 C_{T} \\
\left|x^{\prime}-y^{\prime}\right| \leq \sqrt{2 \nu\left\{C_{T}+\frac{\delta L_{0}^{2}}{4}\right\}}+2 \nu L_{0} \\
\left|x_{d}-y_{d}\right| \leq \sqrt{2 \delta\left\{C_{T}+\frac{\nu L_{0}^{2}}{4}\right\}}+2 \delta L_{0}
\end{array}\right.
$$

This result is standard but since we need precise constants in the estimation, we postponed the proof.

## Step 2. When the doubled normal variable converges to a single variable

Step 1 shows that up to extract a subsequence, we have for $B_{\rho_{\alpha}}=B_{\rho_{\alpha}}(0)$ and for $\nu$ small enough,

$$
\begin{equation*}
\left(\bar{X}_{\delta}, \bar{Y}_{\delta}\right) \rightarrow(\bar{X}, \bar{Y}) \in\left(\left[\tau_{\eta}, T-\tau_{\eta}\right] \times \bar{B}_{\rho_{\alpha}}\right)^{2} \quad \text { as } \quad \delta \rightarrow 0 \quad \text { where } \quad \rho_{\alpha}:=2 \sqrt{\frac{6 C_{T}}{\alpha}} \tag{4.5}
\end{equation*}
$$

with $\bar{X}=\left(\bar{t}, \bar{x}^{\prime}, \bar{x}_{d}\right)=\left(\bar{\xi}, \bar{x}_{d}\right), \bar{Y}=\left(\bar{s}, \bar{y}^{\prime}, \bar{y}_{d}\right)=\left(\bar{\zeta}, \bar{y}_{d}\right), \bar{y}_{d}=\bar{x}_{d}$, and

$$
\alpha g(\bar{y})+\frac{\eta}{T-\bar{s}}+\frac{|\bar{t}-\bar{s}|^{2}}{2 \nu} \leq 3 C_{T} \quad \text { and } \quad\left|\bar{x}^{\prime}-\bar{y}^{\prime}\right| \leq \sqrt{2 \nu C_{T}}+2 \nu L_{0}=o_{\nu}(1) \rightarrow 0 \quad \text { as } \quad \nu \rightarrow 0
$$

where the last bound follows from estimate of Lemma 4.1. Moreover, we have

$$
0<M / 2 \leq M_{\nu, \alpha, \delta} \rightarrow M_{\nu, \alpha}:=\sup _{X, Y \in[0, T) \times \bar{\Omega}} \Psi_{\nu, \alpha}(X, Y)=\Psi_{\nu, \alpha}(\bar{X}, \bar{Y}) \quad \text { as } \quad \delta \rightarrow 0
$$

with

$$
\Psi_{\nu, \alpha}(t, x, s, y):= \begin{cases}\tilde{u}\left(t, x^{\prime}, x_{d}\right)-v\left(s, y^{\prime}, x_{d}\right)-\alpha g(y)-\frac{\eta}{T-s}-\frac{\left|\left(t, x^{\prime}\right)-\left(s, y^{\prime}\right)\right|^{2}}{2 \nu} & \text { if } x_{d}=y_{d} \\ -\infty & \text { if } x_{d} \neq y_{d}\end{cases}
$$

From the fact that $M_{\nu, \alpha} \rightarrow M_{\nu, 0}$ as $\alpha \rightarrow 0$ (with obvious definitions), we deduce that all maximizer in the definition of $M_{\nu, \alpha}$ satisfies

$$
\begin{equation*}
\lim _{\alpha \rightarrow 0} \alpha g(\bar{y})=0 \tag{4.6}
\end{equation*}
$$

Moreover we have

$$
\lim _{\nu \rightarrow 0}\left(\lim _{\alpha \rightarrow 0} M_{\nu, \alpha}\right)=M=\lim _{\theta \rightarrow 0} M(\theta)
$$

where $M(\theta)$ is defined in (4.1). This also implies that

$$
\begin{equation*}
\lim _{\nu \rightarrow 0}\left(\lim _{\alpha \rightarrow 0} \frac{\left|\left(\bar{t}, \bar{x}^{\prime}\right)-\left(\bar{s}, \bar{y}^{\prime}\right)\right|^{2}}{\nu}\right)=0 \tag{4.7}
\end{equation*}
$$

We now prove that $\bar{X} \in\left\{\tilde{u}>\underline{u}_{0}\right\}$. Assume by contradiction that $\tilde{u}(\bar{X})=\underline{u}_{0}(\bar{X})$. Then

$$
\Psi_{\nu, \alpha}(\bar{X}, \bar{Y}) \leq \underline{u}_{0}(\bar{X})-v(\bar{Y})-\frac{\left|\bar{x}^{\prime}-\bar{y}^{\prime}\right|^{2}}{2 \nu}
$$

which implies (using the a priori bound $v \geq \underline{u}_{0}$ )

$$
\begin{align*}
0<M / 4 & \leq M_{\nu, \alpha}=\Psi_{\nu, \alpha}(\bar{X}, \bar{Y}) \\
& \leq \underline{u}_{0}(\bar{x})-\underline{u}_{0}(\bar{y})-\frac{\left|\bar{x}^{\prime}-\bar{y}^{\prime}\right|^{2}}{2 \nu^{\prime}}  \tag{4.8}\\
& \leq L_{0}\left|\bar{x}^{\prime}-\bar{y}^{\prime}\right|-\frac{\left|\bar{x}^{\prime}-\bar{y}^{\prime}\right|^{2}}{2 \nu}
\end{align*}
$$

This leads to a contradiction as $\nu \rightarrow 0$. Hence

$$
\begin{equation*}
\bar{X} \in\left\{\tilde{u}>\underline{u}_{0}\right\} \tag{4.9}
\end{equation*}
$$

This implies also that for $\delta$ small enough $\bar{X}_{\delta} \in\left\{\tilde{u}>\underline{u}_{0}\right\}$.
Step 3: proof that $\bar{x}_{d}=0$
By contradiction, we assume that we are in the standard case $\bar{x}_{d}>0$. Then we also have $\left(\bar{x}_{\delta}\right)_{d},\left(\bar{y}_{\delta}\right)_{d}>0$ and the viscosity inequalities with $\bar{p}_{\delta}:=\left(\frac{\bar{x}_{\delta}^{\prime}-\bar{y}_{\delta}^{\prime}}{\nu}, \frac{\left(\bar{x}_{\delta}\right)_{d}-\left(\bar{y}_{\delta}\right)_{d}}{\delta}\right)$

$$
\left\{\begin{align*}
\frac{\bar{\epsilon}_{\delta}-\bar{s}_{\delta}}{\nu}+H\left(\bar{X}_{\delta}, \bar{p}_{\delta}\right) & \leq 0 \quad \text { because } \quad \bar{X}_{\delta} \in\left\{\tilde{u}>\underline{u}_{0}\right\}  \tag{4.10}\\
-\frac{\eta}{\left(T-\bar{s}_{\delta}\right)^{2}}+\frac{\overline{\bar{t}}_{\delta}-\bar{s}_{\delta}}{\nu}+H\left(\bar{Y}_{\delta},-\alpha D g\left(\bar{y}_{\delta}\right)+\bar{p}_{\delta}\right) & \geq 0
\end{align*}\right.
$$

We know from Lemma 4.1, that

$$
\left|\frac{\bar{x}_{\delta}^{\prime}-\bar{y}_{\delta}^{\prime}}{\nu}\right| \leq \nu^{-1}\left\{\sqrt{2 \nu\left\{C_{T}+\frac{\delta L_{0}^{2}}{4}\right\}}+2 \nu L_{0}\right\}
$$

and

$$
H\left(\bar{X}_{\delta}, \bar{p}_{\delta}\right) \leq-\left\{\frac{\bar{t}_{\delta}-\bar{s}_{\delta}}{\nu}\right\} \leq \sqrt{\frac{6 C_{T}}{\nu}}
$$

Moreover, the uniform coercivity of $H$ (see (1.4) iv)) implies the existence of some $\tilde{L}_{\nu}>0$ (independent on $\delta$, for $\delta>0$ small enough, and independent on $\alpha$ ) such that

$$
\left|\left(\frac{\bar{t}_{\delta}-\bar{s}_{\delta}}{\nu}, \bar{p}_{\delta}\right)\right| \leq \tilde{L}_{\nu}
$$

We can then subtract the two viscosity inequalities in (4.10), and get

$$
\frac{\eta}{\left(T-\bar{s}_{\delta}\right)^{2}} \leq H\left(\bar{Y}_{\delta},-\alpha D g\left(\bar{y}_{\delta}\right)+\bar{p}_{\delta}\right)-H\left(\bar{X}_{\delta}, \bar{p}_{\delta}\right)
$$

Passing to the limit $\delta \rightarrow 0$, we get (up to extraction of a subsequence) that $\left(\frac{\bar{t}_{\delta}-\bar{s}_{\delta}}{\nu}, \bar{p}_{\delta}\right) \rightarrow\left(\frac{\bar{t}-\bar{s}}{\nu}, \bar{p}\right)$ with

$$
\bar{p}=\left(\frac{\bar{x}^{\prime}-\bar{y}^{\prime}}{\nu}, \bar{p}_{d}\right) \in \mathbb{R}^{d-1} \times \mathbb{R} \quad \text { and } \quad\left(\frac{\bar{t}-\bar{s}}{\nu}, \bar{p}\right) \in \bar{D}_{t, x}^{1,+} u(\bar{X}) \quad \text { with } \quad\left|\left(\frac{\bar{t}-\bar{s}}{\nu}, \bar{p}\right)\right| \leq \tilde{L}_{\nu}
$$

and

$$
\begin{equation*}
\frac{\eta}{(T-\bar{s})^{2}} \leq H(\bar{Y},-\alpha D g(\bar{y})+\bar{p})-H(\bar{X}, \bar{p}) \quad \text { with } \quad \bar{y}_{d}=\bar{x}_{d} \tag{4.11}
\end{equation*}
$$

Using assumptions (1.4) ii) and iii), this implies that (say with $L:=2 \tilde{L}_{\nu}$ and $\alpha$ small enough)
$\frac{\eta}{T^{2}} \leq \omega_{L}(\alpha D g(\bar{y}))+\omega\left(|\bar{X}-\bar{Y}|\left(1+\left|\bar{p}^{\prime}\right|+\max (0, H(\bar{X}, \bar{p})) \leq \omega_{L}(\alpha D g(\bar{y}))+\omega\left(|\bar{X}-\bar{Y}|\left(1+\frac{\left|\bar{x}^{\prime}-\bar{y}^{\prime}\right|}{\nu}+\frac{|\bar{t}-\bar{s}|}{\nu}\right)\right)\right.\right.$.
Using the fact that $\alpha D g(\bar{y}) \rightarrow 0$ as $\alpha \rightarrow 0$ and estimate (4.7), we get a contradiction for $\alpha$ and $\nu$ small enough.

## Step 4: The key one-sided Lipschitz estimate

In the remaining of the proof we then have $\bar{x}_{d}=0$. For $\xi=\left(t, x^{\prime}\right), \zeta=\left(s, y^{\prime}\right)$, and $x_{d} \in[0,+\infty)$, we set

$$
\Psi_{\nu}^{\alpha}\left(\xi, \zeta, x_{d}\right):=\Psi_{\nu, \alpha}\left(\xi, x_{d}, \zeta, x_{d}\right)=\tilde{u}\left(\xi, x_{d}\right)-v\left(\zeta, x_{d}\right)-\frac{\eta}{T-s}-\alpha g\left(y^{\prime}, x_{d}\right)-\frac{|\xi-\zeta|^{2}}{2 \nu}
$$

and consider

$$
\begin{equation*}
0<M / 2 \leq M_{\nu, \alpha}=\sup _{\xi, \zeta \in\left([0, T) \times \mathbb{R}^{d-1}\right)^{2}, x_{d} \in[0,+\infty)} \Psi_{\nu}^{\alpha}\left(\xi, \zeta, x_{d}\right)=\Psi_{\nu}^{\alpha}\left(\bar{\xi}, \bar{\zeta}, \bar{x}_{d}\right) \tag{4.12}
\end{equation*}
$$

We define

$$
V(s, y):=v(s, y)+\frac{\eta}{T-s}+\alpha g(y)
$$

so that we have

$$
\left\{\begin{aligned}
\partial_{t} \tilde{u}+H(X, D \tilde{u}) \leq 0 & \text { in } \quad((0, T) \times \Omega) \quad \cap\left\{\tilde{u}>\underline{u}_{0}\right\} \\
\partial_{t} \tilde{u}+\min \{F, H\}(X, D \tilde{u}) \leq 0 & \text { on } \quad((0, T) \times \partial \Omega) \cap\left\{\tilde{u}>\underline{u}_{0}\right\} \\
-\frac{\eta}{(T-s)^{2}}+\partial_{s} V+H(Y, D V-\alpha D g) \geq 0 & \text { in } \quad(0, T) \times \Omega \\
-\frac{\eta}{(T-s)^{2}}+\partial_{s} V+\max \{F, H\}(Y, D V-\alpha D g) \geq 0 & \text { on } \quad(0, T) \times \partial \Omega
\end{aligned}\right.
$$

We now claim the following one-sided "Lipschitz" estimate

$$
\begin{equation*}
\tilde{u}\left(\xi, x_{d}\right)-V\left(\zeta, y_{d}\right) \leq \tilde{u}\left(\bar{\xi}, \bar{x}_{d}\right)-V\left(\bar{\zeta}, \bar{x}_{d}\right)+\frac{|\xi-\zeta|^{2}}{2 \nu}-\frac{|\bar{\xi}-\bar{\zeta}|^{2}}{2 \nu}+L_{\nu}\left|x_{d}-y_{d}\right| \tag{4.13}
\end{equation*}
$$

where $L_{\nu}$ is given in Lemma 3.2, and where equality holds for $t=\bar{t}, s=\bar{s}, x^{\prime}=\bar{x}^{\prime}, y^{\prime}=\bar{y}^{\prime}$ and $x_{d}=y_{d}=\bar{x}_{d}$, with $\bar{X}=\left(\bar{\xi}, \bar{x}_{d}\right), \bar{Y}=\left(\bar{\zeta}, \bar{x}_{d}\right)$. For clarity, the proof of (4.13) is postponed at the end of the proof of the theorem.

## Step 5: the twin blow-ups.

We then consider the following twin blow-ups with small parameter $\varepsilon>0$ : one blow-up for $\tilde{u}$ at the point $\bar{X}=\left(\bar{\xi}, \bar{x}_{d}\right)$ and one blow-up for $V$ at the point $\bar{Y}=\left(\bar{\zeta}, \bar{x}_{d}\right)$,

$$
\begin{cases}U^{\varepsilon}(\hat{X}):=\varepsilon^{-1}\{\tilde{u}(\bar{X}+\varepsilon \hat{X})-\tilde{u}(\bar{X})\}, & U^{\varepsilon}(0)=0  \tag{4.14}\\ V^{\varepsilon}(\hat{Y}):=\varepsilon^{-1}\{V(\bar{Y}+\varepsilon \hat{Y})-V(\bar{Y})\}, & V^{\varepsilon}(0)=0\end{cases}
$$

Before passing to the limit $\varepsilon \rightarrow 0$, they satisfy for $\hat{X}=(\hat{t}, \hat{x})$ and $\hat{Y}=(\hat{s}, \hat{y})$

$$
\left\{\begin{array}{rlll}
\partial_{\hat{t}} U^{\varepsilon}+H\left(\bar{X}+\varepsilon \hat{X}, D_{\hat{x}} U^{\varepsilon}\right) \leq 0 & & \text { in } & I_{\bar{t}}^{\varepsilon} \times \Omega  \tag{4.15}\\
\partial_{\hat{t}} U^{\varepsilon}+\min (F, H)\left(\bar{X}+\varepsilon \hat{X}, D_{\hat{x}} U^{\varepsilon}\right) \leq 0 & & \text { on } & I_{\bar{t}}^{\varepsilon} \times \partial \Omega \\
& & \\
-\bar{\eta}^{\varepsilon}+\partial_{\hat{s}} V^{\varepsilon}+H\left(\bar{Y}+\varepsilon \hat{Y}, D_{\hat{y}} V^{\varepsilon}-\alpha D_{y} g(\bar{y}+\varepsilon \hat{y})\right) \geq 0 & & \text { in } & I_{\hat{s}}^{\varepsilon} \times \Omega \\
-\bar{\eta}^{\varepsilon}+\partial_{\hat{s}} V^{\varepsilon}+\max (F, H)\left(\bar{Y}+\varepsilon \hat{Y}, D_{\hat{y}} V^{\varepsilon}-\alpha D_{y} g(\bar{y}+\varepsilon \hat{y})\right) \geq 0 & & \text { on } & I_{\bar{s}}^{\varepsilon} \times \partial \Omega
\end{array}\right.
$$

with

$$
\bar{\eta}^{\varepsilon}(\hat{s}):=\frac{\eta}{(T-(\bar{s}+\varepsilon \hat{s}))^{2}} \quad \text { and } \quad I_{\bar{r}}^{\varepsilon}:=\left(-\frac{\bar{r}}{\varepsilon}, \frac{T-\bar{r}}{\varepsilon}\right) \quad \text { for } \quad \bar{r}=\bar{t}, \bar{s}
$$

From (4.13), they also satisfy

$$
\begin{equation*}
U^{\varepsilon}\left(\hat{\xi}, \hat{x}_{d}\right)-V^{\varepsilon}\left(\hat{\zeta}, \hat{y}_{d}\right) \leq L_{\nu}\left|\hat{x}_{d}-\hat{y}_{d}\right|+\bar{b} \cdot(\hat{\xi}-\hat{\zeta})+\varepsilon \frac{|\hat{\xi}-\hat{\zeta}|^{2}}{2 \nu} \quad \text { with } \quad \bar{b}:=\frac{\bar{\xi}-\bar{\zeta}}{\nu} \tag{4.16}
\end{equation*}
$$

We then define the following half-relaxed limits

$$
\begin{cases}U^{0}:=\limsup _{\varepsilon \rightarrow 0}^{*} U^{\varepsilon}, & U^{0}(0) \geq 0 \\ V^{0}:=\liminf _{\varepsilon \rightarrow 0} V^{\varepsilon}, & V^{0}(0) \leq 0\end{cases}
$$

Passing to the limit in (4.16), we get

$$
\begin{equation*}
U^{0}\left(\hat{\xi}, \hat{x}_{d}\right)-V^{0}\left(\hat{\zeta}, \hat{y}_{d}\right) \leq L_{\nu}\left|\hat{x}_{d}-\hat{y}_{d}\right|+\bar{b} \cdot(\hat{\xi}-\hat{\zeta}) \tag{4.17}
\end{equation*}
$$

which implies in particular that $U^{0}(0)=0=V^{0}(0)$. Passing to the limit in (4.15) and using the discontinuous stability of viscosity solutions, we also get

$$
\left\{\begin{array}{rll}
\partial_{\hat{t}} U^{0}+H\left(\bar{X}, D_{\hat{x}} U^{0}\right) \leq 0 & \text { in } & (\mathbb{R} \times \Omega) \cap\left\{\left|U^{0}\right|<+\infty\right\}  \tag{4.18}\\
\partial_{\hat{t}} U^{0}+\min (F, H)\left(\bar{X}, D_{\hat{x}} U^{0}\right) \leq 0 & \text { on } & (\mathbb{R} \times \partial \Omega) \cap\left\{\left|U^{0}\right|<+\infty\right\} \\
-\bar{\eta}+\partial_{\hat{s}} V^{0}+H\left(\bar{Y}, D_{\hat{y}} V^{0}-\alpha D_{y} g(\bar{y})\right) \geq 0 & \text { in } & (\mathbb{R} \times \Omega) \cap\left\{\left|V^{0}\right|<+\infty\right\} \\
-\bar{\eta}+\partial_{\hat{s}} V^{0}+\max (F, H)\left(\bar{Y}, D_{\hat{y}} V^{0}-\alpha D_{y} g(\bar{y})\right) \geq 0 & \text { on } & (\mathbb{R} \times \partial \Omega) \cap\left\{\left|V^{0}\right|<+\infty\right\}
\end{array}\right.
$$

with $\bar{\eta}:=\frac{\eta}{(T-\bar{s})^{2}}$.

## Step 6: the 1D problem.

We now define the following functions on $[0,+\infty)$ as the supremum/infimum in the tangential variables of the functions defined in $\mathbb{R} \times \bar{\Omega}$,

$$
\bar{u}\left(\hat{x}_{d}\right):=\sup _{\hat{\xi} \in \mathbb{R}^{d}}\left\{U^{0}\left(\hat{\xi}, \hat{x}_{d}\right)-\bar{b} \cdot \hat{\xi}\right\}, \quad \underline{v}\left(\hat{y}_{d}\right):=\inf _{\hat{\zeta} \in \mathbb{R}^{d}}\left\{V^{0}\left(\hat{\zeta}, \hat{y}_{d}\right)-\bar{b} \cdot \hat{\zeta}\right\} .
$$

From (4.17), these functions satisfy

$$
-\infty \leq-L_{\nu}\left|\hat{x}_{d}-\hat{y}_{d}\right|+\bar{u}\left(\hat{x}_{d}\right) \leq \underline{v}\left(\hat{y}_{d}\right) \leq+\infty, \quad 0 \leq \bar{u}(0) \leq \underline{v}(0) \leq 0
$$

In particular, this implies that $\bar{u}(0)=0=\underline{v}(0)$. Because of this one-sided Lipschitz inequality, this is also the case for their semi-continuous envelopes, i.e. we have (and this is important)

$$
\begin{equation*}
-\infty \leq-L_{\nu}\left|\hat{x}_{d}-\hat{y}_{d}\right|+\bar{u}^{*}\left(\hat{x}_{d}\right) \leq \underline{v}_{*}\left(\hat{y}_{d}\right) \leq+\infty, \quad \bar{u}^{*}(0)=0=\underline{v}_{*}(0) . \tag{4.19}
\end{equation*}
$$

We set $H_{\alpha}(Y, p)=H(Y, p-\alpha D g(\bar{y}))$ and $F_{\alpha}(Y, p)=F(Y, p-\alpha D g(\bar{y}))$. From (4.18), we get (again from stability) that these functions satisfy in particular for $\bar{X}:=(\bar{t}, \bar{x}), \bar{Y}:=(\bar{s}, \bar{x})$ and $\bar{b}=\left(\bar{b}_{0}, \bar{b}^{\prime}\right) \in \mathbb{R} \times \mathbb{R}^{d-1}$

$$
\left\{\begin{array}{rlrr}
\bar{b}_{0}+H\left(\bar{X}, \bar{b}^{\prime}, \partial_{\hat{x}_{d}} \bar{u}^{*}\right) \leq 0 & \text { in } & (0,+\infty) & \cap\left\{\left|\bar{u}^{*}\right|<+\infty\right\}  \tag{4.20}\\
\bar{b}_{0}+\min (F, H)\left(\bar{X}, \bar{b}^{\prime}, \partial_{\hat{x}_{d}} \bar{u}^{*}\right) \leq 0 & \text { in } & \{0\} & \cap\left\{\left|\bar{u}^{*}\right|<+\infty\right\} \\
-\bar{\eta}+\bar{b}_{0}+H_{\alpha}\left(\bar{Y}, \bar{b}^{\prime}, \partial_{\hat{y}_{d}} \underline{v}_{*}\right) \geq 0 & \text { in } & (0,+\infty) & \cap\left\{\left|\underline{v}_{*}\right|<+\infty\right\} \\
-\bar{\eta}+\bar{b}_{0}+\max \left(F_{\alpha}, H_{\alpha}\right)\left(\bar{Y}, \bar{b}^{\prime}, \partial_{\hat{y}_{d} \underline{v}_{*}}\right) \geq 0 & \text { in } & \{0\} & \cap\left\{\left|\underline{v}_{*}\right|<+\infty\right\} .
\end{array}\right.
$$

Step 7: getting a contradiction from structural assumptions.
We now apply Corollary 2.5. In order to do so, we now set $z=\hat{x}_{d}=\hat{y}_{d}$ and consider

$$
\bar{p}_{d}:=\limsup _{[0,+\infty) \ni z \rightarrow 0} \frac{\bar{u}^{*}(z)}{z}, \quad \underline{p}_{d}:=\liminf _{[0,+\infty) \ni z \rightarrow 0} \frac{\underline{v}_{*}(z)}{z}, \quad a_{d}:=\min \left(\underline{p}_{d}, \bar{p}_{d}\right), \quad b_{d}:=\max \left(\underline{p}_{d}, \bar{p}_{d}\right)
$$

and we get that there exists $p_{d} \in\left[a_{d}, b_{d}\right] \cap \mathbb{R} \neq \emptyset$ such that either

$$
\bar{b}_{0}+H\left(\bar{X}, \bar{b}^{\prime}, p_{d}\right) \leq 0<\bar{\eta} \leq H_{\alpha}\left(\bar{Y}, \bar{b}^{\prime}, p_{d}\right)-H\left(\bar{X}, \bar{b}^{\prime}, p_{d}\right)
$$

or

$$
\bar{b}_{0}+\max (F, H)\left(\bar{X}, \bar{b}^{\prime}, p_{d}\right) \leq 0<\bar{\eta} \leq F_{\alpha}\left(\bar{Y}, \bar{b}^{\prime}, p_{d}\right)-F\left(\bar{X}, \bar{b}^{\prime}, p_{d}\right) .
$$

One of these facts are true along a subsequence $\nu \rightarrow 0$. In the first case, we get from the assumption on the Hamiltonian $H$, see (1.4) ii), that (using $\bar{p}^{\prime}=\bar{b}^{\prime}$ ) and again $L:=2 \tilde{L}_{\nu}$ for $\alpha$ small enough,

$$
\begin{aligned}
\bar{\eta} \leq H_{\alpha}\left(\bar{Y}, \bar{b}^{\prime}, p_{d}\right)-H\left(\bar{X}, \bar{b}^{\prime}, p_{d}\right) & \leq \omega\left(|\bar{X}-\bar{Y}| \cdot\left[1+\left|\bar{b}^{\prime}\right|+\max \left\{0, H\left(\bar{X}, \bar{b}^{\prime}, p_{d}\right)\right\}\right]\right)+\omega_{L}(\alpha D g(\bar{y})) \\
& \leq \omega\left(|\bar{\xi}-\bar{\zeta}| \cdot\left[1+\left|\bar{b}^{\prime}\right|+\max \left\{0,-\bar{b}_{0}\right\}\right]\right)+\omega_{L}(\alpha D g(\bar{y})) \\
& \leq \omega\left(2 \frac{|\bar{\xi}-\bar{\zeta}|^{2}}{\nu}+|\bar{\xi}-\bar{\zeta}|\right)+\omega_{L}(\alpha D g(\bar{y})) \rightarrow 0 \quad \text { as } \quad \alpha \rightarrow 0, \text { and then } \nu \rightarrow 0
\end{aligned}
$$

where we have used the expression of $\bar{b}=\frac{\bar{\xi}-\bar{\zeta}}{\nu}$ in the third line, and (4.7) in the last line. Contradiction because $\bar{\eta} \geq \eta / T^{2}>0$.
From the assumption on the function $F$, see (1.5) ii), we get a similar contradiction in the second case,

$$
\begin{aligned}
\bar{\eta} \leq F_{\alpha}\left(\bar{Y}, \bar{b}^{\prime}, p_{d}\right)-F\left(\bar{X}, \bar{b}^{\prime}, p_{d}\right) & \leq \omega\left(|\bar{\xi}-\bar{\zeta}| \cdot\left[1+\left|\bar{b}^{\prime}\right|+\max \left\{0, \max \{F, H\}\left(\bar{X}, \bar{b}^{\prime}, p_{d}\right)\right\}\right]\right)+\omega_{L}(\alpha D g(\bar{y})) \\
& \leq \omega\left(|\bar{\xi}-\bar{\zeta}| \cdot\left[1+\left|\bar{b}^{\prime}\right|+\max \left\{0,-\bar{b}_{0}\right\}\right]\right)+\omega_{L}(\alpha D g(\bar{y})) \\
& \leq \omega\left(2 \frac{|\bar{\xi}-\bar{\zeta}|^{2}}{\nu}+|\bar{\xi}-\bar{\zeta}|\right)+\omega_{L}(\alpha D g(\bar{y})) \rightarrow 0 \quad \text { as } \quad \alpha \rightarrow 0, \text { and then } \nu \rightarrow 0 .
\end{aligned}
$$

We conclude that $M \leq 0$. Recalling that

$$
M=\sup _{t \in[0, T), x \in \bar{\Omega}}\left\{\tilde{u}(t, x)-v(t, x)-\frac{\eta}{T-t}\right\} \leq 0,
$$

it is enough to let $\eta \rightarrow 0$ to get $u \leq \tilde{u} \leq v$ as desired.
Back to Step 3: proof of the key one-sided Lipschitz estimate (4.13)
We now justify (4.13). Following Lemma 3.2, we extend $\tilde{u}$ and consider

$$
\tilde{U}^{\nu}\left(\zeta, x_{d}\right):=\sup _{\xi \in \mathbb{R} \times \mathbb{R}^{d-1}}\left\{\tilde{u}\left(\xi, x_{d}\right)-\frac{|\zeta-\xi|^{2}}{2 \nu}\right\}
$$

and there exists some (possibly non unique) $\bar{\xi}_{\zeta} \in\left[s-\theta^{\nu}, s+\theta^{\nu}\right] \times \mathbb{R}^{d-1}$ such that $\tilde{U}^{\nu}\left(\zeta, x_{d}\right)=\tilde{u}\left(\bar{\xi}_{\zeta}, x_{d}\right)-\frac{\left|\bar{\xi}_{\zeta}-\zeta\right|^{2}}{2 \nu}$. If $s \in\left[\theta^{\nu}, T-\theta^{\nu}\right]$, then we see that $\bar{\xi}_{\zeta} \in(0, T) \times \mathbb{R}^{d-1}$ and we also have

$$
\tilde{U}^{\nu}\left(\zeta, x_{d}\right):=\sup _{\xi \in[0, T) \times \mathbb{R}^{d-1}}\left\{\tilde{u}\left(\xi, x_{d}\right)-\frac{|\xi-\zeta|^{2}}{2 \nu}\right\} .
$$

In particular for $\left(\zeta, x_{d}\right)=\left(\bar{\zeta}, \bar{x}_{d}\right)$, we can choose $\bar{\xi}_{\bar{\zeta}}=\bar{\xi}$ where $\bar{\xi}_{\bar{\zeta}}$ is given by Lemma 3.2 and $\bar{X}=\left(\bar{\xi}, \bar{x}_{d}\right), \bar{Y}=$ $\left(\bar{\zeta}, \bar{x}_{d}\right)$ appear in (4.5). Now we choose $\nu>0$ small enough such that $\theta^{\nu}<\tau_{\eta}$, and we set $I^{\nu}:=\left(\theta^{\nu}, T-\theta^{\nu}\right)$. Moreover we have for all $\zeta \in I^{\nu} \times \mathbb{R}^{d-1}, y_{d} \in[0,+\infty)$,

$$
\tilde{U}^{\nu}\left(\zeta, y_{d}\right)-V\left(\zeta, y_{d}\right) \leq \sup _{\xi \in[0, T) \times \mathbb{R}^{d-1}, x_{d} \in[0,+\infty)} \Psi_{\nu, \alpha}\left(\xi, \zeta, y_{d}\right)=\Psi_{\nu, \alpha}\left(\bar{\xi}, \bar{\zeta}, \bar{x}_{d}\right)=\tilde{U}^{\nu}\left(\bar{\zeta}, \bar{x}_{d}\right)-V\left(\bar{\zeta}, \bar{x}_{d}\right) .
$$

Now from Lemma 3.2, we also know that $\tilde{U}^{\nu}$ is $L_{\nu}$-Lipschitz, and then $\tilde{U}^{\nu}\left(\zeta, x_{d}\right)-\tilde{U}^{\nu}\left(\zeta, y_{d}\right) \leq L_{\nu}\left|x_{d}-y_{d}\right|$, which implies

$$
\tilde{U}^{\nu}\left(\zeta, x_{d}\right)-V\left(\zeta, y_{d}\right) \leq \tilde{U}^{\nu}\left(\bar{\zeta}, \bar{x}_{d}\right)-V\left(\bar{\zeta}, \bar{x}_{d}\right)+L_{\nu}\left|x_{d}-y_{d}\right|
$$

which gives exactly (4.13). This ends the proof of the theorem.

We now turn to the proof of Lemma 4.1.
Proof of Lemma 4.1. Recall that we have

$$
\tilde{u}(t, x) \leq u_{0}(x)+C_{T}, \quad v(t, x) \geq u_{0}(x)-C_{T} .
$$

Hence

$$
\Psi_{\nu, \alpha, \delta}(t, x, s, y) \leq 2 C_{T}+B_{\delta}(x, y)-\alpha g(y)-\frac{\eta}{T-s}-\frac{|t-s|^{2}}{2 \nu}
$$

with

$$
B_{\delta}(x, y):=\left\{u_{0}(x)-u_{0}(y)-\frac{\left|x_{d}-y_{d}\right|^{2}}{2 \delta}-\frac{\left|x^{\prime}-y^{\prime}\right|^{2}}{2 \nu}\right\} \leq \phi_{\nu}\left(\left|x^{\prime}-y^{\prime}\right|\right)+\phi_{\delta}\left(\left|x_{d}-y_{d}\right|\right)
$$

and

$$
\phi_{\delta}(r):=L_{0} r-\frac{r^{2}}{2 \delta} \leq \frac{\delta L_{0}^{2}}{2}
$$

Here $\phi_{\delta}$ is concave with $\phi_{\delta}\left(r_{\delta}\right)=0$ for $r_{\delta}:=2 \delta L_{0}$. Moreover

$$
\phi_{\delta}(r) \leq\left(r-r_{\delta}\right) \phi_{\delta}^{\prime}\left(r_{\delta}\right)=\left(r-r_{\delta}\right)\left(L_{0}-\frac{r}{\delta}\right)
$$

i.e.

$$
\begin{equation*}
\phi_{\delta}(r) \leq-\delta^{-1}\left(r-r_{\delta}\right)^{2} \quad \text { for } \quad r \geq r_{\delta}=2 \delta L_{0} \tag{4.21}
\end{equation*}
$$

We get in particular

$$
0<\Psi_{\nu, \alpha, \delta}(X, Y) \leq 2 C_{T}+\phi_{\delta}\left(\left|x_{d}-y_{d}\right|\right)+\phi_{\nu}\left(\left|x^{\prime}-y^{\prime}\right|\right)
$$

and then

$$
0<\Psi_{\nu, \alpha, \delta}(X, Y) \leq 2 C_{T}+\phi_{\delta}\left(\left|x_{d}-y_{d}\right|\right)+\frac{\nu L_{0}^{2}}{2}
$$

which implies from (4.21) that

$$
\begin{equation*}
\left|x_{d}-y_{d}\right| \leq \sqrt{2 \delta\left\{C_{T}+\frac{\nu L_{0}^{2}}{4}\right\}}+2 \delta L_{0} \tag{4.22}
\end{equation*}
$$

and symmetrically that

$$
\begin{equation*}
\left|x^{\prime}-y^{\prime}\right| \leq \sqrt{2 \nu\left\{C_{T}+\frac{\delta L_{0}^{2}}{4}\right\}}+2 \nu L_{0} \tag{4.23}
\end{equation*}
$$

We also deduce from $0<\Psi_{\nu, \alpha, \delta}(X, Y)$ that

$$
\begin{equation*}
\alpha g(y)+\frac{\eta}{T-s}+\frac{|t-s|^{2}}{2 \nu} \leq 2 C_{T}+\frac{(\nu+\delta) L_{0}^{2}}{2} \leq 2 C_{T}+\frac{\eta}{T} \leq 3 C_{T} \tag{4.24}
\end{equation*}
$$

for $\eta>0$ small enough (the size of $\eta$ depending on $C_{T}$ and $T$, but not on $\nu, \alpha, \delta$ ), and for $\delta, \nu>0$ small enough (for a size depending on $\eta$ ). Therefore we have

$$
\begin{equation*}
\left|t_{k}-s_{k}\right| \leq \bar{\theta}^{\nu}:=3 \sqrt{\nu C_{T}} \tag{4.25}
\end{equation*}
$$

and

$$
T-s_{k}>\frac{\eta}{3 C_{T}}, \quad T-t_{k}>\frac{\eta}{3 C_{T}}-\bar{\theta}^{\nu} \geq \frac{\eta}{4 C_{T}}
$$

for $\nu>0$ small enough (for a size depending on $\eta$ ).
Similarly, from Lemma 3.1 on the barriers (in particular using Case 1 of the proof, for Lipschitz initial data $u_{0}$ ), we know that there exists some $\lambda>0$ such that

$$
u(t, x) \leq u_{0}(x)+\lambda t, \quad v(s, y) \geq u_{0}(y)-\lambda s
$$

Hence

$$
\Psi_{\nu, \alpha, \delta}\left(P_{k}\right) \leq \lambda(t+s)+L_{0}|x-y|-\frac{\eta}{T} \leq \lambda\left(t_{k}+s_{k}\right)-\frac{2 \eta}{3 T}
$$

where we have used bound (4.22)-(4.23) for $\delta, \nu>0$ small enough (for a size depending on $\eta>0$ ). Therefore

$$
\max (t, s)>\frac{\eta}{3 \lambda T}, \quad \min (t, s)>\frac{\eta}{3 \lambda T}-\bar{\theta}^{\nu} \geq \frac{\eta}{4 \lambda T}
$$

for $\nu>0$ small enough (for a size depending on $\eta$ ). Up to increase $\lambda$ or $C_{T}$ (and decrease $\nu>0$ if necessary), we can assume that

$$
\lambda T \equiv C_{T}
$$

Setting

$$
\tau_{\eta}:=\frac{\eta}{4 \lambda T}=\frac{\eta}{4 C_{T}}
$$

and for $\nu>0$ small enough, we see that

$$
X, Y \in\left[\tau_{\eta}, T-\tau_{\eta}\right] \times \bar{\Omega}
$$

This gives the result with (4.22), (4.23) and (4.24). This ends the proof of the lemma.

## 5 The comparison principle on a bounded domain

Let us consider an open set $\Omega$ satisfying for, $d \geq 1$,

$$
\begin{equation*}
\Omega \subset \mathbb{R}^{d} \text { is a bounded open set with } C^{1} \text { boundary and outward unit normal } n(x) \tag{5.1}
\end{equation*}
$$

Let $T>0$. We consider the following equation for $u(t, x)$ with $X:=(t, x) \in[0, T) \times \bar{\Omega}$

$$
\begin{equation*}
u_{t}+H(X, D u)=0 \quad \text { on } \quad(0, T) \times \Omega \tag{5.2}
\end{equation*}
$$

and the boundary condition

$$
u_{t}+F(X, D u)=0 \quad \text { on } \quad(0, T) \times \partial \Omega
$$

We also consider an initial boundary condition

$$
\begin{equation*}
u(0, \cdot)=u_{0} \quad \text { on } \quad\{0\} \times \bar{\Omega} \tag{5.3}
\end{equation*}
$$

The rigorous meaning of desired boundary conditions is the following,

$$
\left\{\begin{array}{lll}
u_{t}+\min (F, H)(X, D u) \leq 0 & \text { on } & (0, T) \times \partial \Omega  \tag{5.4}\\
\text { (for subsolutions) } \\
u_{t}+\max (F, H)(X, D u) \geq 0 & \text { on } & (0, T) \times \partial \Omega
\end{array} \quad\right. \text { (for supersolutions) }
$$

As far as Hamiltonians are concerned, we assume the following structure conditions, where $\omega, \omega_{L}$ are moduli of continuity.

```
i) (Continuity)
\(H:[0, T] \times \bar{\Omega} \times \mathbb{R}^{d} \rightarrow \mathbb{R}\) is continuous
ii) (Uniform continuity in the gradient)
For any \(L>0\), we have
\(|H(X, p)-H(X, q)| \leq \omega_{L}(|p-q|) \quad\) for all \(\quad X \in[0, T] \times \bar{\Omega}, \quad p, q \in[-L, L]^{d}\)
iii) (Quantified continuity in time-space variables)
\(H(Y, p)-H(X, p) \leq \omega(|Y-X|(1+\max \{0, H(X, p)\})) \quad\) for all \(\quad\left\{\begin{array}{l}X, Y \in[0, T] \times \bar{\Omega} \\ p \in \mathbb{R}^{d}\end{array}\right.\)
iv) (Uniform coercivity)
\(\lim _{|p| \rightarrow+\infty} \inf _{X \in[0, T] \times \bar{\Omega}} H(X, p)=+\infty\)
```

and as previously, making artificially appear the dependence on $x \in \bar{\Omega}$ for $F$ (in order to unify the presentation of $H$ and $F$ ), we consider for $X=(t, x)$

```
i) (Continuity)
\(F:[0, T] \times \partial \Omega \times \mathbb{R}^{d} \rightarrow \mathbb{R}\) is continuous
and the map \(q \mapsto F(X, p-q n(x))\) is nonincreasing
ii) (Uniform continuity in the gradient)
For any \(L>0\), we have
\(|F(X, p)-F(X, q)| \leq \omega_{L}(|p-q|) \quad\) for all \(\quad X \in[0, T] \times \partial \Omega, \quad p, q \in[-L, L]^{d}\)
iii) (Continuity in the tangential variables)
\(F(Y, p)-F(X, p) \leq \omega(|Y-X|(1+\max \{0, \max (F, H)(X, p)\})) \quad\) for all \(\quad\left\{\begin{array}{l}X, Y \in[0, T] \times \partial \Omega \\ p \in \mathbb{R}^{d}\end{array}\right.\)
iv) (Uniform normal semi-coercivity)
For any \(L>0\), we have
\(\left(\lim _{q \rightarrow-\infty} \inf _{X \in[0, T] \times \partial \Omega, p \in[-L, L]^{d}} F(X, p-q n(x))=+\infty\right.\).
```

We then have the following theorem.
Theorem 5.1 (Comparison principle on a bounded open set $\Omega$ ). Let $T>0$ and assume that $H, F$ satisfy respectively (5.5) and (5.6). Assume that the initial data $u_{0}$ is continuous. Let $u, v:[0, T) \times \bar{\Omega} \rightarrow \mathbb{R}$ be two functions with $u$ upper semi-continuous and $v$ lower semi-continuous. Assume that $u$ (resp. v) is a viscosity subsolution (resp. supersolution) of (5.2)-(5.3). Assume moreover that there exists a constant $C_{T}>0$ such that

$$
u \leq u_{0}+C_{T} \quad \text { and } \quad v \geq u_{0}-C_{T} \quad \text { on } \quad[0, T) \times \bar{\Omega}
$$

If

$$
u(0, \cdot) \leq u_{0} \leq v(0, \cdot) \quad \text { on } \quad\{0\} \times \bar{\Omega}
$$

then

$$
u \leq v \quad \text { in } \quad[0, T) \times \bar{\Omega}
$$

In order to give the proof of Theorem 5.1, we need the following lemma which proof is left to the reader.

Lemma 5.2 (Action of a diffeomorphism on the structural conditions satisfied by $H, F)$. Assume that $\Omega$ has the regularity given in assumption (5.1). For $T>0$, we set $Q_{T}:=(0, T) \times \Omega$ with $\Omega \subset \mathbb{R}^{d}$ and for $P=\left(p_{0}, p\right) \in \mathbb{R} \times \mathbb{R}^{d}$, we set

$$
H_{0}(Y, P):=p_{0}+H(Y, p), \quad F_{0}(Y, P):=p_{0}+F(Y, p)
$$

We say by extension that $H_{0}, F_{0}$ satisfy respectively (5.5) and (5.6), if $H, F$ do it.
Assume that $H_{0}, F_{0}$ satisfy respectively (5.5) and (5.6) and for $x_{0} \in \partial \Omega$, consider, locally around $x_{0}$, a $C^{1}$-diffeomorphism $\Phi$ from $\bar{\Omega}$ to $\bar{\Omega}$ with $\Phi\left(x_{0}\right)=y_{0} \in \partial \tilde{\Omega}$, that we extend by the identity on the time variable. Still denoting by $\Phi$ this diffeomorphism, we assume that $\Phi$ maps locally $\bar{Q}_{T}$ to $\overline{\tilde{Q}}_{T} \subset \mathbb{R}^{1+d}$ with locally $\Phi\left(\partial Q_{T}\right)=\partial \tilde{Q}_{T}$. For $Y \in \tilde{Q}_{T}$ and $P \in \mathbb{R}^{1+d}$, we set

$$
\begin{cases}\tilde{H}_{0}(Y, P):=H_{0}\left(\Phi^{-1}(Y), P \cdot B(Y)\right) & \text { locally around }[0, T] \times\left\{y_{0}\right\}, \text { on } \\ \tilde{F}_{T}(Y, P):=F_{0}\left(\Phi^{-1}(Y), P \cdot B(Y)\right) & \text { locally around }[0, T] \times\left\{y_{0}\right\}, \text { on } \\ \partial \tilde{Q}_{T}\end{cases}
$$

with

$$
(P \cdot B)_{j}=\sum_{i=0}^{d} P_{i}\left\{\left(D_{j} \Phi_{i}\right) \circ \Phi^{-1}\right\} \quad \text { for } \quad j=0, \ldots, d
$$

Then $\tilde{H}_{0}, \tilde{F}_{0}$ satisfy respectively (5.5) and (5.6) locally around $[0, T] \times\left\{y_{0}\right\}$, with some suitable moduli.
We now turn to the proof of Theorem 5.1.
Proof of Theorem 5.1. Up to proceed as in Step 0 of the proof of Theorem 1.1, we can assume that $u_{0}$ belongs to $C^{1}(\bar{\Omega})$.
We set

$$
\tilde{u}:=\max \left\{u, \underline{u}_{0}\right\}, \quad \underline{u}_{0}(t, X)=u_{0}(X)-C_{T}=\underline{u}_{0}(X)
$$

and

$$
M:=\sup _{(t, x) \in[0, T) \times \bar{\Omega}} \Psi(t, x) \quad \text { with } \quad \Psi(t, X)=\tilde{u}(t, x)-v(t, x)-\frac{\eta}{T-t}
$$

Assume by contradiction that

$$
0<M=\Psi\left(X_{0}\right) \quad \text { with } \quad X_{0}:=\left(t_{0}, x_{0}\right) \in[0, T) \times \bar{\Omega}
$$

By assumption, we have $t_{0}>0$. If $x_{0} \in \Omega$, then we can localize, and then get a contradiction by standard method of doubling of variables. Hence assume that $x_{0} \in \partial \Omega$. Up to modify slithly the functions, we can assume that the suppremum is strict at $X_{0}$. Up to change the coordinates, we can also assume that

$$
x_{0}=0, \quad \Omega \cap B_{r}\left(x_{0}\right)=\left\{x=\left(x^{\prime}, x_{d}\right) \in \mathbb{R}^{d-1} \times \mathbb{R}, \quad x_{d}>h\left(x^{\prime}\right)\right\} \cap B_{r}\left(x_{0}\right), \quad h(0)=0=\operatorname{Dh}(0)
$$

Setting
$x=\Phi(y) \quad$ with $\quad y=\left(y^{\prime}, y_{d}\right)=\Phi^{-1}(x):=\left(x^{\prime}, x_{d}-h\left(x^{\prime}\right)\right) \quad$ and $\quad \tilde{U}(t, y):=\tilde{u}(t, \Phi(y)), \quad V(t, y):=v(t, \Phi(y))$
we see that $\Phi$ is locally invertible and its inverse is a $C^{1}$ map, given, for some $\rho>0$ small enough, by

$$
\Phi: K_{\rho}^{+} \rightarrow \bar{\Omega} \cap B_{r}\left(x_{0}\right) \quad \text { with } \quad K_{\rho}^{+}:=[-\rho, \rho]^{d-1} \times[0, \rho]
$$

Hence we have

$$
x=\left(x^{\prime}, x_{d}\right)=\left(y^{\prime}, y_{d}+h\left(y^{\prime}\right)\right)=\Phi(y), \quad \tilde{u}(t, x)=\tilde{U}\left(t, \Phi^{-1}(x)\right), \quad v(t, x)=V\left(t, \Phi^{-1}(x)\right)
$$

and

$$
\begin{cases}D_{x_{d}} \tilde{u}(t, x) & =D_{y_{d}} \tilde{U}\left(t, \Phi^{-1}(x)\right) \\ D_{x^{\prime}} \tilde{u}(t, x) & =D_{y^{\prime}} \tilde{U}\left(t, \Phi^{-1}(x)\right)-\left\{D_{y_{d}} \tilde{U}\left(t, \Phi^{-1}(x)\right)\right\} \cdot D_{x^{\prime}} h\left(x^{\prime}\right) \\ \tilde{u}_{t}(t, x) & =\tilde{U}_{t}\left(t, \Phi^{-1}(x)\right)\end{cases}
$$

This gives the new Hamiltonian $\tilde{H}$ and boundary function $\tilde{F}$ for $Y=(t, y), X=(t, x)$ and $x=\Phi(y)$

$$
\begin{aligned}
& \tilde{H}(Y, D \tilde{U}(Y))=H(X, D \tilde{u}(X)), \\
& \tilde{F}(Y, D \tilde{U}(Y))=F(X, D \tilde{u}(X)), \quad \text { for } \quad y=\left(y^{\prime}, 0\right)
\end{aligned}
$$

which are defined by $\left(\right.$ for $\left.y:=\left(y^{\prime}, y_{d}\right)\right)$

$$
\begin{aligned}
& \tilde{H}\left(t, y, p^{\prime}, p_{d}\right):=H\left(t, \Phi(y), p^{\prime}-p_{d} D_{y^{\prime}} h\left(y^{\prime}\right), p_{d}\right), \\
& \tilde{F}\left(t, y, p^{\prime}, p_{d}\right):=F\left(t, \Phi(y), p^{\prime}-p_{d} D_{y^{\prime}} h\left(y^{\prime}\right), p_{d}\right), \quad \text { for } \quad y=\left(y^{\prime}, 0\right)
\end{aligned}
$$

Hence $\tilde{U}$ and $V$ are respectively sub/supersolutions of

$$
\begin{cases}W_{t}+\tilde{H}(Y, D W)=0 & \text { on } \quad(0, T) \times[-\rho, \rho]^{d-1} \times[0, \rho] \\ W_{t}+\tilde{F}(Y, D W)=0 & \text { on } \quad(0, T) \times[-\rho, \rho]^{d-1} \times\{0\}\end{cases}
$$

We now apply Lemma 5.2 to insure that $\tilde{H}$ and $\tilde{F}$ satisfy (locally) the same structural conditions than $H$ and $F$. Moreover, we have

$$
M=\sup _{(t, y) \in[0, T) \times K_{\rho}^{+}} \tilde{\Psi}(t, y)=\tilde{\Psi}\left(t_{0}, 0\right) \quad \text { with } \quad \tilde{\Psi}(t, y):=\tilde{U}(t, y)-V(t, y)-\frac{\eta}{T-t}
$$

Up to add some small and smooth tangential correction term $\left|t-t_{0}\right|^{2}+\left|y^{\prime}\right|^{2}$ to $V$ (here we neglect this correction which can be treated in a very classical way), we can assume that

$$
\tilde{\Psi}(t, y)<M \quad \text { for all } \quad(t, y) \in\left([0, T) \times K_{\rho}^{+}\right) \backslash\left\{\left(t_{0}, 0\right)\right\}
$$

This implies that for $\xi=\left(t, x^{\prime}\right), \zeta=\left(s, y^{\prime}\right)$

$$
M(\theta):=\sup \left\{\tilde{U}\left(\xi, y_{d}\right)-V\left(\zeta, y_{d}\right)-\frac{\eta}{T-s}, \quad \xi, \zeta \in[0, T) \times[-\rho, \rho]^{d-1}, y_{d} \in[0, \rho], \quad|\xi-\zeta| \leq \theta\right\}
$$

with

$$
\lim _{\theta \rightarrow 0^{+}} M(\theta)=M>0
$$

We are then back to the begining of the proof of Theorem 1.1, which leads to a contradiction. Again, we conclude that $M \leq 0$ for all $\eta \rightarrow 0^{+}$, and then deduce that $\tilde{U} \leq V$, and then $u \leq v$. This ends the proof of the theorem.

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