

Séverine Demeyer, V. Le Sant, A. Koenen, N. Fischer, Julien Waeytens, Rémi

Bouchié

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| 4 | Séverine Demeyer [*] , V. Le Sant, A. Koenen and N. Fischer |
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| 5 | Laboratoire National de Métrologie et d'Essais, 29 avenue Roger Hennequin, 78197 |
| 6 | Trappes Cedex |
| 7 | *E-mail: severine.demeyer@lne.fr |
| 8 | Julien Waeytens |
| 9 | Univ. Gustave Eiffel, IFSTTAR. |
| 10 | 14-20 boulevard Newton, 77447, Marne la Vallée Cedex 2, France |
| 11 | E-mail: julien.waeytens@univ-eiffel.fr |
| 12 | Rémi Bouchié |
| 13 | Centre Scientifique et Technique du Bâtiment. |
| 14 | 84 avenue Jean Jaurès, Champs-sur-Marne, 77447 Marne-la-Vallée, France |
| 15 | E-mail: remi.bouchie@cstb.fr |
| 16 | Abstract. |
| 17 | In this work, we propose a fully Bayesian uncertainty analysis of the indirect |
| 18 | measurement of thermal properties of walls from in-situ temperature and flux |
| 19 | measurements, obtained with an active method, using a one dimensional transient |
| 20 | thermal model. We show that this approach is able to take into account the uncertainty |
| 21 | of the inputs of the thermal model and the uncertainty of the output observations, for |
| 22 | a more reliable uncertainty estimation of the <i>calibration parameters</i> and any derived |
| 23 | quantity. For this problem, we improve the classical Bayesian inversion model by taking |
| 24 | into account underestimated uncertainty on reported output observations, which is a |
| 25 | frequently encountered issue in practice. We provide some recommendations for a wider |
| 26 | applicability of the method. We illustrate the principles of uncertainty evaluation of the |
| 27 | Guide to the Expression of Uncertainty in Measurement [BIPM et al., 2008a] in terms |
| 28 | of a real case study to evaluate the thermal resistance of a multilayer wall placed in a |
| 29 | climatic chamber. For this application, we compare results of the Bayesian inversion |
| 30 | with classical steady-state results in comparable experimental conditions. We perform |
| 31 | a sensitivity analysis to study the effect of duration and input uncertainties and we |
| 32 | make recommendations. R code is made available that enables a Bayesian uncertainty |
| 33 | evaluation of inversion models for related applications. |

34 1. Introduction

In the context of building energy management or retrofit interventions, thermal resistance of walls is used in standard heat transfer models as a parameter for building performance simulations [Iglesias et al., 2018]. Due to the cost of wall thermal characterization, tabulated values for thermal properties are typically used as inputs of the energy models [Iglesias et al., 2018].

The in-situ measurement of the thermal resistance of walls should, for instance, contribute to decrease the gap between actual and predicted consumption to have more reliable energy-savings strategies or to have a more reliable assessment of the effectiveness of retrofit interventions [Simon et al., 2018].

The present study is part of the French National Research project named RESBATI 44 whose main objective is to develop a portable measurement device for evaluating the 45 thermal resistance of opaque building walls on site. The result of the thermal resistance 46 measurement should be provided with its uncertainty [Ha et al., 2020]. The active 47 method (a thermal gradient inside the wall is created by a step heating excitation applied 48 on a face) chosen in the project allows to estimate the thermal resistance of a building 49 wall in all seasons, for any type of building and use (occupied or otherwise) in quite a 50 short measurement time (less than three days) which improves on classical steady-state 51 or dynamic methods [François et al., 2020], [Ha et al., 2020]. In this paper, we consider 52 that the thermal resistance of a wall is measured indirectly from in-situ temperature 53 and flux measurements. This is an inverse problem classically encountered in building 54 physics [Rouchier, 2018]. 55

To carry out the identification method, a direct model (here a thermal model) 56 is embedded in a simulator $\eta(X,\theta)$ whose inputs are the calibration parameters θ 57 (called *primary unknowns* in [Kaipio and Fox, 2011]), and the experimental conditions 58 X, see Figure 1. In practice, [Kaipio and Fox, 2011] mention that it is very seldom that 59 the primary unknowns θ are the only *input unknowns*. The vector X gathers all the 60 uncertainty sources, called here *input parameters*, having an effect on the measurement 61 result of the thermal resistance and that can be modeled by a probability distribution. 62 In the following, the uppercase notation X denotes a random vector, whereas the 63 lowercase x denotes a particular observed value of X. In this paper, we have chosen 64 a one-dimensional transient thermal model as a good approximation of the physical 65 phenomenon with a low computational cost. Since the thermal model does not take 66 into account lateral flux in the building wall, the user has to make sure that the flux is 67 close to 1D, which makes the simulator fit for purpose. 68

The identification process, illustrated in Figure 1, also called calibration or inversion, consists to find the values $\hat{\theta}$ of the input calibration parameters of the thermal model so that the output $\eta(x, \hat{\theta})$ is as closed as possible to the experimental values y(x) obtained in the experimental conditions x. The classical techniques for solving the minimization problem pertaining to inversion are based on least squares methods (e.g. Levenberg-Marquardt algorihm), maximum-likelihood (ML) and



Figure 1. Schema of a calibration (inversion) procedure.

Bayesian analysis. A description and a comparison of these inversion techniques can 75 be found in [Rouchier, 2018] and [Kaipio and Fox, 2011]. In Bayesian statistics, the 76 inverse problem is regularized statistically through the modeling of prior distribution 77 [Kaipio and Fox, 2011]. Prior distributions represent prior degrees of belief on the 78 (usually) unique but unknown values of the inputs which are turned into posterior 79 degrees of belief (posterior distributions) via the Bayes formula when observations are 80 available. The reader interested in generalities about Bayesian statistics can refer to 81 [Gelman et al., 2013] and [Box and Tiao, 1992]. 82

Bayesian inversion requires advanced computational tools like Markov Chain 83 Monte Carlo methods (MCMC) to sample from the posterior distributions 84 [Metropolis et al., 1953], [Chib and Greenberg, 1995], [Robert and Casella, 2013]. For 85 a tutorial on the class of MCMC algorithms classically used for Bayesian inversion (the 86 Metropolis-Hastings algorithm) the reader is referred to [Klauenberg and Elster, 2016]. 87 Recently, Bayesian calibration using in-situ measurements has been widely used for 88 the inference of thermal properties and their associated uncertainties. To cite a few, 89 [Berger et al., 2016] tackled the estimation the thermal conductivity and the internal 90 convective heat transfer coefficient of a wall with 3 layers from one year experimental 91 data, [Iglesias et al., 2018] addressed the estimation of the thermal resistance and the 92 heat capacity of unit area of a homogeneous wall based on 7 days of measurements, 93 [Thébault and Bouchié, 2018] considered the estimation of the HLC with the ISABELE 94 method, [Simon et al., 2018] tackled the estimation of the thermal conductivity and the 95 heat capacity of the wall based on 8 days measurement campaign, [Rodler et al., 2019] 96 addressed the estimation of the thermal conductivity and the volumetric heat capacity 97 from a few days measurements. 98

This paper focuses on uncertainty propagation in inversion models. A widespread 99 approach is to consider input uncertainty sources as negligible or small enough to be 100 aggregated with uncertainties on the outputs [Perrin and Durantin, 2019]. Such an 101 approach is followed e.g. in [Simon et al., 2018] and [Berger et al., 2016]. However, 102 if the input uncertainties are not negligible or even if there is no information about 103 these input uncertainties, ignoring these uncertainties may result in biased estimated 104 calibration parameters [Perrin and Durantin, 2019]. The difficulty that arises for 105 uncertainty propagation is that there is usually no close functional form to describe 106 the relation between the uncertainty sources and the estimate of the calibration 107

Many approaches have been developped to parameters [Wang and Zabaras, 2004]. 108 create approximate relationships through which uncertainties are propagated. An 109 example of uncertainty propagation using the variance-covariance matrix of estimated 110 parameters with the ordinary least squares (OLS) method and a (thermal) model 111 reduction can be found in [François et al., 2020]. The analytical expression of 112 the outputs in function of the initial and the boundary conditions, e.g. using 113 discretized heat transfer equations as in [Iglesias et al., 2018] provides the "missing" 114 relationship between the uncertainty sources and the estimates of the calibration 115 parameters (obtained as Bayesian estimates of non linear regression parameters 116 in [Iglesias et al., 2018]). [Thébault and Bouchié, 2018] and [Gori and Elwell, 2018] 117 establish posterior relationships between the estimated parameters taken as the 118 maximum at posteriori (MAP) and the uncertain inputs (e.g. posterior profile at MAP). 119 Although the Bayesian approach conceptually allows to take into account measurement 120 uncertainty in both inputs and outputs, [Perrin and Durantin, 2019] point out that little 121 work has been done to uncertainties on inputs. The recent work by [Rodler et al., 2019] 122 combines Metropolis-Hastings with Monte Carlo sampling from the uncertainty sources 123 in an attempt of full Bayesian inversion. 124

The goal of the paper is to provide guidance on input uncertainty propagation in 125 Bayesian inversion for the in-situ estimation of thermal parameters from temperature 126 and flux measurements using a 1D thermal model. In this paper, we assume that 127 the steps of review of the uncertainty sources, selection of the most influential and 128 finally the prior quantification and modeling of uncertainty attached to all of them 129 has been done before performing the Bayesian uncertainty analysis of the inversion 130 problem. We show that uncertainty propagation of the experimental conditions through 131 the inversion model can be performed by a fully Bayesian analysis of the inversion 132 problem by posterior sampling from the joint prior distribution of the calibration 133 parameters and the uncertain experimental conditions, following an approach similar to 134 [Demeyer et al., 2021] and [Higdon et al., 2004]. The results are calculated with Markov 135 Chain Monte Carlo methods (MCMC) and an introduction to these methods is given. 136 Among the various simulation studies [Ha et al., 2020] and experimental works 137 carried out within the RESBATI project, the use case of this paper focuses on the 138 thermal measurements obtained on a IWI (Inner Wall Insulation) test wall built within 139 the project and circulated between LNE and CSTB. Here, we compare results from 140 the Bayesian analysis of the thermal measurements obtained in an energy room at 141

LNE with the results from the steady-state guarded hot box method at CSTB and we

perform a sensitivity analysis to study the effect of measurement duration and the level

144 of uncertainty.

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¹⁴⁵ 2. Bayesian calibration under uncertainty

146 2.1. General formulation

Denoting $y = (y_1, ..., y_N)$ the vector of output measurements, $X = (X_1, ..., X_N)$ where $X_i = (X_{i1}, ..., X_{id})^T$, the $N \times d$ matrix of uncertain input variables and θ the vector of calibration parameters, the fully Bayesian inversion model writes

$$y = \eta(X, \theta) + \varepsilon \tag{1}$$

$$X \sim \pi(X) \tag{2}$$

$$\theta \sim \pi(\theta) \tag{3}$$

where $\eta(X,\theta) = (\eta(X_1,\theta),...\eta(X_N,\theta))$ denotes the vector of simulation outputs, $\varepsilon = (\varepsilon_1,...,\varepsilon_N)$ denotes the vector of output measurement error, $\pi(X)$ and $\pi(\theta)$ denote respectively the prior distributions of X and θ .

The unknown vector X is classically represented with an error in variable model (see [Perrin and Durantin, 2019]) as a deviation from a known central value x (see discussion section 2.5)

$$X = x + \zeta \tag{4}$$

where $\zeta = (\zeta_1, ..., \zeta_N)$ denotes the matrix of input measurement error and $x = (x_1, ..., x_N)$ denotes the matrix of input observations.

The random variables ε_i and ζ_i , for i = 1, ..., N, are commonly modeled as centered Gaussian variables with known covariance matrix denoted Σ_{ε} and Σ_{ζ} respectively [Perrin and Durantin, 2019]. Furthermore θ , ε and ζ are assumed statistically independent.

159 2.2. Modeling excess variance

Model (1) assumes that output observations are equivalent to the calibrated simulations, their difference being explained by the reported measurement uncertainties. In practice, an excess variability of the y_i with respect to (w.r.t.) the variability explained by the measurement uncertainty may be observed. Under the assumption that the simulator is fit for purpose, not taking into account excess variance may lead to biased estimates of the calibration parameters.

Under the hypothesis that the excess variability is due to a missing term in the uncertainty budget \ddagger , we propose to model the excess variability (also called heterogeneity of data) with an adjustment factor $\sigma > 0$ as in [Bodnar and Elster, 2014] and [Mana et al., 2012]

$$y = \eta(X, \theta) + \sigma\varepsilon \tag{5}$$

‡ In this paper, we assume that the excess variability of the time-varying thermal measurements y_i where *i* indexes time, comes from a missing influential source of uncertainty or from varying, not controllable, conditions during the recording period (noise).

The parameter σ is related to the Birge ratio, that is often used in metrology to enlarge quoted uncertainties when combining inconsistent measurement results on the same measurand [Bodnar and Elster, 2014]. As shown in [Mana et al., 2012], the Birge ratio can be interpreted as the most probable value of such an adjustment factor.

According to [Bodnar and Elster, 2014], it can be assumed that σ^2 follows a priori an inverse chi-squared distribution $\sigma^2 \sim \text{InvChi}^2(\nu_0, s_0^2)$, in which case the marginal distribution of observations w.r.t σ^2 is multivariate t-distributed

$$y \sim t_{\nu_0} \left(\eta(X, \theta), s_0^2 \Sigma_{\varepsilon} \right) \tag{6}$$

A poorly informative prior for σ^2 , centered on 1, can be obtained with $\nu_0 = 2$ and $s_0^2 = 2$.

172 Remarks

• The prior degree of freedom ν_0 can be viewed as allowing "uncertainty on the uncertainty" contained in Σ_{ε} which comes to considering reported uncertainties as point estimates of the unknown standard deviations.§ The assumption that the missing source of uncertainty is common to all measurements translates into ν_0 being common to all measurements.

• Additive random effects \parallel could be used to model excess variance if measurement results y_i were obtained from uncertainty analyses performed independently, for instance by various teams or with different measurement devices. They have become common practice for instance when building consensus estimates in meta-analysis and interlaboratory comparisons [Bodnar et al., 2017]. For such applications, individual degrees of freedom ν_i are taken into account when they are available, see [Toman and Possolo, 2009].

185 2.3. Bayesian inference

In order to perform a full Bayesian analysis of the inversion problem, calibration parameters θ are augmented with the latent variables X and the heterogeneity variance parameter σ^2 . The Bayes formula provides the joint posterior distribution for (θ, X, σ^2)

$$\pi(\theta, X, \sigma^2 | y) \propto l(y | \eta(X, \theta), \sigma^2) \pi(\theta) \pi(\sigma^2) \pi(X)$$
(8)

 \S For instance, the GUM uses the so-called Welch-Satterthwaite formula to estimate the degrees of liberty associated with a measurement result assuming it is t-distributed.

|| The conventional random effects model considers the following model for the output observations

$$y = \eta(X, \theta) + \lambda + \varepsilon \tag{7}$$

where $\lambda = (\lambda_1, ..., \lambda_N)$ gathers for instance the team effects under the assumptions that $\lambda_i \sim_{iid} N(0, \sigma^2)$, $\varepsilon_i \sim N(0, u_i^2)$ and that the λ_i and the ε_i are independent. where $l(y|\eta(X,\theta),\sigma^2)$ is the likelihood, $\pi(\theta)$, $\pi(\sigma^2)$, $\pi(X)$ are the prior distributions of θ , σ^2 and X respectively, assumed independent. The prior distribution $\pi(\theta)$ is usually based on expert knowledge or tabulated values. The prior distribution for σ^2 can be chosen as poorly informative when the dataset is large. In order to account for the input uncertainty sources X, $\pi(X)$ is a data-driven prior distribution based on (4).

In this paper, the quantity of interest is the joint posterior distribution $\pi(\theta, X|y)$ obtained by integrating the joint posterior distribution out σ^2

$$\pi(\theta, X|y) = \int_0^{+\infty} \pi(\theta, X, \sigma^2|y) \mathrm{d}\sigma^2$$
(9)

For a well-chosen prior distribution $\pi(\sigma^2)$, a closed form expression can be computed for (9). For instance, using $\sigma^2 \sim \text{Inv} - \text{chi2}(\nu_0, s_0^2)$ yields

$$\pi(\theta, X|y) \propto \pi(\theta)\pi(X)l_{int}(y|X, \theta) \tag{10}$$

where $l_{int}(y|X,\theta) = t_{\nu_0}(\eta(X,\theta), s_0^2 \Sigma_{\varepsilon})$

199 2.4. Posterior simulation with Markov Chain Monte Carlo algorithms

Integrals (8), (9) or (10) are usually intractable and require simulation methods like 200 Markov Chain Monte Carlo (MCMC) methods. MCMC methods provide a flexible and 201 powerful tool for sampling from an arbitrary distribution. These methods construct a 202 sequence of dependent values which form a Markov chain with stationary distribution 203 equal to the sought posterior distribution. The Metropolis-Hastings (MH) algorithm 204 [Metropolis et al., 1953], [Chib and Greenberg, 1995] constitutes a popular class of 205 MCMC methods. The interesting feature of the MH algorithm is that it only requires to 206 know the posterior distribution up to a proportionality constant, i.e. the right-hand part 207 of Eq.8 or Eq.10, which makes it particularly suited for calibration problems. The MH 208 algorithm can be used in conjunction with the Gibbs algorithm if conditional posterior 209 distributions are tractable for a subset of parameters (the so-called *Metropolis-within-*210 *Gibbs* algorithm). The sequence of values is considered after a burn-in period and often 211 the chains are thinned (i.e. only each 10th value is used) to decrease the autocorrelation 212 of the chains. 213

The Metropolis Hastings algorithm for sampling from the posterior distribution of (θ, X) given y in (10) is given in algorithm (1). The *initialization* step (line 1) provides a starting point for the iterative algorithm which can be sampled from the prior distributions or near a best physical point if available or obtained from computational techniques based e.g. on the gradient (Lagrange,...). At each iteration l, a (multidimensional) candidate $(\theta^{(c)}, X^{(c)})$ is sampled (line 5) and the corresponding output of the thermal simulation $\eta(\theta^{(c)}, X^{(c)})$ is computed (line 6). The *acceptance/rejection* step lines 8, 9 and 10 is based on the ratio of posterior distributions (known up to the same constant, that disappears in the ratio)

$$\alpha = \frac{\pi(\theta^{(c)}, X^{(c)}|y)}{\pi(\theta^{(l-1)}, X^{(l-1)}|y)}$$
(11)

A candidate point with ratio $\alpha < 1$ has probability α to be selected which corresponds to $P(u < \alpha)$ where $u \sim Unif(0, 1)$. If a candidate is accepted, the candidate is the new point of the chain, otherwise the previous point is added to the chain. It is important to note that at each iteration an element is added to the chain, being either the candidate or the previous point. For computational reasons, the log-posteriors are used (line 2 for the log-posterior of the starting point and line 7 for the log-posterior of the candidate point).

Algorithm 1 Metropolis-Hastings algorithm for Bayesian calibration under uncertainty Input: measurements x_i , y_i , t_i ;

Output: M samples from the posterior distribution $\pi(\theta, X|y)$ according to (10); 1: initialize: $(\theta^{(1)}, X^{(1)}, \eta^{(1)} = \eta(X^{(1)}, \theta^{(1)}));$ 2: compute $a = \ln(\pi(\theta^{(1)}, X^{(1)}|y));$ 3: repeat $l \leftarrow l + 1;$ 4: sample $\theta^{(c)} \sim N_d(\theta^{(l-1)}, \Sigma_{\theta}^{(l-1)})$ and $X^{(c)} \sim N_N\left(X^{(l-1)}, \Sigma_X^{(l-1)}\right)$; 5:generate $\eta^{(c)} = \eta(X^{(c)}, \theta^{(c)})$; 6: compute $b = \ln(\pi(\theta^{(c)}, X^{(c)} | \eta^{(c)}));$ 7: let $\alpha = \min(\exp(b-a), 1)$ and $u \sim \text{Unif}(0, 1)$; 8: if $\alpha \geq u$ then $\theta^{(l)} = \theta^{(c)}$ and $X^{(l)} = X^{(c)}$; 9: else $\theta^{(l)} = \theta^{(l-1)}$ and $X^{(l)} = X^{(l-1)}$; 10: $a \leftarrow b$ 11: 12: **until** l = M

221 2.5. Discussion

The Bayesian approach allows a flexible modelling of the prior distributions $\pi(X)$ 222 and $\pi(\theta)$ due to the use of MCMC simulation methods, among which the Gaussian, 223 Student, rectangular, triangular and trapezoïdal distributions that are commonly 224 used to represent uncertainty on input quantities e.g. in metrology (see GUM-225 For instance, the Gaussian distribution can be used S1 [BIPM et al., 2008b]). 226 to model inputs resulting from an uncertainty propagation (e.g. following the 227 GUM propagation of variance [BIPM et al., 2008a]) like the solar aperture and the 228 volumetric flowrate. Generally, the prior on θ consists in a rectangular distribution 229 defining bounds inside which the unknown value is supposed to lie, to improve 230 the computational efficiency of the inversion method. For example, the bounds 231 associated with thermophysical properties of material can usually be found in literature 232

²³³ [Heo et al., 2012], [ASHRAE, 2017]. More generally, guidelines for assessing the ²³⁴ uncertainty of inputs used in building energy models can be found in [Macdonald, 2002] ²³⁵ and [Heo et al., 2012].

To account for complex uncertain inputs like those related to time (e.g. 236 occupancy, weather, usage...), a discrepancy term $\delta(x)$, function of these inputs, is 237 usually added to (1) [Heo et al., 2012] and $\pi(X)$ is modeled as a Gaussian Process 238 [Rasmussen and Williams, 2006], [Santner and Notz, 2003]. A complete Bayesian 239 framework for modelling an additional discrepancy term in inversion is proposed in 240 [Higdon et al., 2004], [Kennedy and O'Hagan, 2001]. In this paper, due to the active 241 method, temperature and flux measurements are likely to be non-stationary time 242 processes that would require modeling using e.g. Gaussian processes functions of both 243 x and y. For simplicity, we choose to model the uncertain time varying inputs and 244 outputs following (4) and (1) respectively and to leave the thorough modeling of input 245 and output time series and their associated uncertainties in Bayesian inversion as a 246 perspective of this work. Indeed, for the analysis of the real case study, such further 247 work did not prove necessary. 248

In this paper, we propose a fully Bayesian approach of inversion problems to 249 propagate uncertainties of input parameters during the inversion procedure, which comes 250 to treat all parameters as *calibration parameters* in classical Bayesian inversion. Such 251 an approach is even recommended as a baseline of any Bayesian inversion procedure in 252 [Higdon et al., 2004]. Indeed, according to [Higdon et al., 2004], even if we believe that 253 the true physical value is known, allowing a slight deviation from the true physical value 254 (similar to representation (4)) may produce an empirically better model of the reality. 255 Furthermore, estimations of the calibration parameters may not even be interpreted as 256 estimates of the true physical values of these parameters [Higdon et al., 2004]. Rather, 257 they are by nature a 'best fit' estimate of θ depending on the many hypotheses (among 258 which that the thermal model is a perfect fit of the physical phenomenon) and the 259 experimental conditions involved in the whole process. 260

In this paper, the use of an active method on a short period (24 hours) allows to limit 261 bias due to the sensitivity of the inversion procedure to the environmental conditions. 262 Indeed, the effect of occupancy, weather and usage are potentially the most influential 263 on the measure of thermal parameters (or more generally the energy performance) but 264 also the most uncertain [Goffart et al., 2017]. If any, model approximation (assumed 265 small) is taken into account in the free variance parameter σ^2 used to capture the excess 266 variance with respect to the reported uncertainties, with the aim of limiting bias of the 267 estimated calibration parameters due to modeling. 268

²⁶⁹ 3. Application to the in-situ estimation of the thermal resistance of a wall

270 3.1. Motivation and context

The thermal resistance $R \,(\text{m}^2 \,\text{K} \,\text{W}^{-1})$ of the enveloppe of a building is an indicator of the thermal performance of the building. Within the RESBATI project, a measurement method of the thermal resistance of opaque walls has been developed to produce in-situ results with an associated uncertainty in less than three days.

In order to evaluate the performance of the method, a test wall of dimensions 2 275 meters by 2 meters with inner wall insulation (IWI) was built. The insulation layer 276 is made of EPS (expanded polystyrene). In this application, we compare the results 277 obtained by the active method on the test wall placed in an energy room at LNE, France 278 (see sections 3.2 and 3.5.3), with results obtained with the steady-state guarded hot box 279 method [ISO 8990:1994, 1994] at CSTB, France (see section 3.5.2). For this application, 280 the steady-state guarded hot plate method [ISO 8302:1991, 1991] (see section 3.5.1) is 281 used to estimate the thermal resistance of the insulation layer (EPS). The description 282 of the experimental conditions of each method is given. 283

²⁸⁴ 3.2. Description of the energy room REBECCA at LNE

The energy room at LNE is a climatic chamber inside another climatic chamber. It 285 is composed of an internal cell with the dimensions of a dwelling (3.42 m of side on)286 $2.29 \,\mathrm{m}$ high, surface of the test specimen $7.83 \,\mathrm{m}^2$). The latter is surrounded by 4 287 climatic chambers in which it is possible to modify the temperature between -7 °C 288 and 35 °C and to generate any type of transient regime (hot cold, cycle, etc.). The 289 four surrounding boxes (front cell, guard, floor and ceiling) can be controlled separately 290 (Figure 2). The REBECCA (Research and Testing of Buildings and Heat Emitters 291 under Artificial Climate) cell is built to reproduce the principle of a guarded hot box 292 test but with more room inside to install additional devices. The front panel is equipped 293 with the specimen to be studied. In order to ensure that the entire heat flow passes 294 through the test specimen, the 5 walls of the inner cell in contact with the thermal guard 295 are insulated with Vacuum Insulated Panels having a thermal resistance greater than 296 $10 \,\mathrm{m}^2 \,\mathrm{K} \,\mathrm{W}^{-1}$. The test can be carried out under steady state condition, which occurs 297 after a long period of time, or in a transient way like temperature ramp. 298

²⁹⁹ 3.3. Description of the measurement process with the active method

The test wall is installed in the REBECCA cell and an active method is used to produce faster in-situ results using a heating module, temperature and flux sensors placed on the wall, as displayed in Figure 3.

The active method uses a prototype built by the laboratory CERTES (Centre d'Études et de Recherche en Thermique, Environnement et Systèmes, Université Paris Est Créteil, France) within the RESBATI project. The prototype is a cube of 60 cm of side. One side is open; in the opposite side 24 DC halogen spots of 20 W are



Figure 2. Schematic diagram of the energy room called 'REBECCA' at LNE, France, from [Koenen et al., 2019].

positioned. The maximum electrical power is 480 W that can be adjusted depending on 307 the test scenario. The four remained faces are insulated with reflective material with 308 a low emissivity. In the front side in contact with the IWI, a heat flow meter with 309 thermocouples is positioned. The IWI is then heated by the halogen spots. The heat 310 flow that goes in the wall through the heat flow sensor and the surface temperatures of 311 the wall are measured every 2s. A step heating excitation which can reach several days 312 is used for all experiments. A schematic description and a view of the experimental set-313 up are presented in Figure 3 and Figure 4. A number of 30 thermocouples are positioned 314 along the heat flow meter to validate the one dimension hypothesis. 315

The heat flow sensors were calibrated in the LNE guarded hot plate at different mean temperatures T (mean of external and internal temperatures) and temperature difference ΔT between the internal and external temperatures. All thermocouples were calibrated at LNE. A test lasts 24 h after steady state condition is achieved. This can continue several days depending on the thermal properties of the walls.

The RESBATI device induces a constant power of 105 W. The external cell was cooled at 0 °C and the internal room was heated at 20 °C in order to obtain a temperature difference ΔT of 20 °C between the two sides of the wall. The heat flow going through the wall in the RESBATI area is measured continuously and reaches around 10 W m⁻² after 10 h and remains constant afterwards.

In this application, the standard uncertainty (evaluated at LNE) associated with the inner and outer surface temperatures (Tsi and Tse respectively) is u(Tsi) = u(Tse) =0.5 °C and the standard uncertainty associated with the absorbed flux q_{int} is $u(q_{int}) =$ Bayesian uncertainty analysis of inversion models applied to the inference of thermal properties of walls12 329 $3\% q_{int} \text{ W m}^{-2}$.



Figure 3. Schematic view of the experimental setup (side view, face view), from [Koenen et al., 2019].



Figure 4. View of the active method prototype, from [Koenen et al., 2019].

The thermal resistance of a multi-layer wall is measured indirectly from temperature 330 and flux measurements after the identification of the parameters θ of a thermal model 331 $\eta(X,\theta)$ (here with $X = (Tse, q_{int})$, see section 3.4) using the Bayesian approach 332 described section 2. A 1D transient thermal model solved with a finite element 333 method in space and a Euler implicit time integration scheme was chosen in this 334 study that takes the heat capacity of unit area cw_i (J m⁻² K⁻¹) and the thermal 335 conductivity k_i (W m⁻¹ K⁻¹) of each layer *i* as calibration inputs gathered in the vector 336 $\theta = \{cw_i, k_i\}_{i=1,\dots,I}$ where I = 4 is the number of layers. 337

The thermal resistance of the wall is expressed as the sum of the thermal resistance

 R_i of each layer

$$R = \sum_{i=1}^{I} R_i, \text{ with } R_i = \frac{l_i}{k_i}$$
(12)

where l_i (m) is the thickness of layer *i* (in this study, l_i is considered known, see values in Table 7). The estimate of the thermal resistance and its associated uncertainty are obtained with the Monte Carlo method applied to (12) using the samples from the posterior distributions of the k_i . In particular, R_2 denotes the thermal resistance of the insulation layer.

Comments on the identifiability of the thermal parameters With this choice of thermal 343 model and the low number of used sensors, we numerically observe that only the 344 joint distribution of the input thermal parameters is identifiable but not its individual 345 constituents (various combinations of values for cw_i and k_i could produce close thermal 346 resistance values). This is a common situation in inversion problems, see the point 347 of [Higdon et al., 2004] discussed in section 2.5, which is overcome here by considering 348 that the thermal resistance (here, the quantity of interest) is identifiable but a priori 349 not the cw_i and k_i taken individually. It is thus important to take samples from the joint 350 posterior distribution to perform the Monte Carlo estimation of the thermal resistance. 351 Still, from virtual testing and model-based sensitivity analysis, the thermal 352 conductivity of the insulation layer k_2 should be identifiable for this application. Thus, 353 the comparison of the posterior distribution of k_2 with an experimental result (e.g. 354 obtained with GHP) could be used as an indication of the ability of the method to 355 estimate the thermal resistance of the insulation layer (calculated as $R_2 = l_2/k_2$), see 356 section 3.5.3. Besides, in this study, the thickness of the layers of the wall is assumed to 357 be known, but in general the thickness should be estimated as well with prior knowledge 358 e.g. based on approximate measurements or from the design of the wall. More generally, 359 identifiability issues are tackled with the use of informative prior distributions. 360

$_{361}$ 3.4. Review of uncertainty sources to specify X

In order to perform a sound uncertainty analysis, the first step is to deeply analyse the measurement process by performing the most exhaustive review of the uncertainty sources pertaining to it (even those that are not quantifiable). This is a demanding brainstorming step involving both experts of the field and statisticians. The classical tool for such a task is the Ishikawa diagram, which is used in the GUM [BIPM et al., 2008a] to review the sources of uncertainty involved in a measurement process.

Figure 5 displays an instance of an Ishikawa diagram pertaining to the measurement process of the thermal resistance of a wall using the RESBATI prototype. Uncertainty sources are divided into five categories shortly described thereafter.

Means: refers to all the uncertainty sources involved with the measurements (sensors, reference materials, standards), here limited to sensors;

³⁷³ Method: refers to all uncertainty sources related to the measurement procedure and ³⁷⁴ the data analysis which consists here in an identification method;

Wall/Enveloppe: refers to the object of the measurement, here the wall; for this application, the thermo-physical properties are the calibration parameters of the

thermal model (see Figure 1);

Environment: refers to the variations of the experimental conditions during the process, here in the energy room;

³⁸⁰ **Operator:** refers to the uncertainty arising from the interpretation and the ³⁸¹ implementation of the procedure by operators involved in the whole process.



Figure 5. Ishikawa diagram for the thermal resistance.

Figure 6 describes the uncertainty sources pertaining to the identification model i.e to the thermal model and the inversion procedure.

In practice, only the most influential quantifiable uncertainty sources X are kept for the uncertainty analysis. Such a selection can be performed with sensitivity analysis methods [Saltelli et al., 2004] and/or by experts based on prior knowledge or information. It is important to note that the combined effect of unquantifiable uncertainty sources, assumed to vary randomly when repeating measurements, can be taken into account when processing measurements.

Here, we choose to focus only on $X = (Tse, q_{int})$ as an example of uncertainty propagation associated with temporal measurements.

392 3.5. Results for IWI test wall

³⁹³ In this section, we compare the estimates of the thermal resistance of the IWI wall ³⁹⁴ obtained with the active method using the Bayesian inversion under uncertainty with

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Figure 6. Ishikawa diagram for the identification model.

those obtained with the steady-state guarded hot box at CSTB, France in comparable 395 experimental conditions ($T = 10 \,^{\circ}\text{C}$ and $\Delta T = 20 \,^{\circ}\text{C}$). As discussed section 3.3, many 396 calibration parameters are not identifiable but the thermal conductivity of the insulation 397 material k_2 should be identifiable. The comparison of the posterior distribution of k_2 398 with the experimental result obtained by LNE using the guarded hot plate method is 399 also used as an indication of the performance of the Bayesian inversion method. Indeed, 400 the thermal resistance of the insulation layer (calculated as $R_2 = l_2/k_2$) is far the most 401 important contribution to the global thermal resistance of the IWI wall. A sensitivity 402 study is performed to show the effect of duration and input uncertainties on the results 403 of the Bayesian inversion. 404

3.5.1. Guarded hot plate results The guarded hot plate apparatus [ISO 8302:1991, 1991]
determines steady-state thermal transmission properties of flat slab specimens having a
low thermal conductivity. It uses the one-dimensional steady-state thermal conductivity
equation. The relative standard uncertainty associated with GHP results is evaluated
around 0.5%.

To conduct the measurement, we have cut two specimen of 600×600 mm from IWI coming from the same batch than those used to build the test wall. The measurements in the guarded hot plate were done at T = 10 °C mean temperature and $\Delta T = 15$ °C. The result is computed as a mean over 8 hours after the steady state is reached. Results are displayed in Table 1. Tabulated values for the plasterboard and the cinderblock are given in Table 2.

⁴¹⁶ 3.5.2. Guarded hot box results The guarded hot box [ISO 8990:1994, 1994] assesses the ⁴¹⁷ thermal performance of walls at full scale. The wall to be tested is positioned between ⁴¹⁸ two ambient conditions: one hot and one cold. Here, the complete wall 2×2 m was ⁴¹⁹ measured in a GHB at CSTB, France at T = 10 °C and $\Delta T = 20$ °C. The result ⁴²⁰ is computed as a mean over 10 hours after the steady state is reached. Under these ⁴²¹ conditions, the thermal resistance of the specimen wall using the GHB is given by the

| Layer | Heat capacity of unit area $cw \ /\mathrm{Jm^{-3}K}$ | Thermal conductivity $k / W m^{-1} K^{-1}$ | Thermal resistance R $/m^2 \mathrm{K} \mathrm{W}^{-1}$ |
|--------------------|--|--|--|
| EPS | $cw_2 = 1.35 \times 10^4$ | $k_2 = 0.031$ | $R_2 = 3.85$ |
| EPS + Plasterboard | - | - | 3.9 |

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 Table 1. Measured thermal performance on GHP.

| Layer | Heat capacity of unit area $cw \ / \mathrm{J m^{-3} K}$ | Thermal conductivity $k / W m^{-1} K^{-1}$ | Thermal resistance R $/m^2 \mathrm{K} \mathrm{W}^{-1}$ |
|--------------|---|--|--|
| Plasterboard | $cw_1 = 7.30 \times 10^4$ | $k_1 = 0.250$ | $\begin{array}{c} 0.05 \\ 0.2 \end{array}$ |
| Cinderblock | $cw_3 = 9.25 \times 10^4$ | $k_3 = 0.850$ | |

Table 2. Tabulated thermal performance values.

 $_{422}$ 95% coverage interval defined by $4.08 \pm 0.86 \text{ m}^2 \text{ K W}^{-1}$, see details in Appendix A.

⁴²³ 3.5.3. Active method coupled with Bayesian inversion Measurements recorded every ⁴²⁴ two seconds for 24 hours are displayed in Figure 7 for the surface temperatures y = Tsi⁴²⁵ and Tse and in Figure 8 for the absorbed flux q_{int} ($X = (Tse, q_{int})$). It can be observed ⁴²⁶ that the absorbed flux becomes constant after 10 hours. For computational issues, ⁴²⁷ measurements are sampled every 300 seconds for a total of N = 289 measurements.

⁴²⁸ A sensitivity analysis is performed to show the effects of duration and experimental ⁴²⁹ uncertainties on the estimation of both the thermal conductivity k_2 of the insulation ⁴³⁰ layer (EPS) and the estimation of the global thermal resistance R.

Table 3 describes the three uncertainty configurations ("no input uncertainty", 431 "low level", "medium level") considered in the study. It is important to note that 432 u(Tsi) = u(Tse) in all configurations due to the experimental setting. The case "no 433 input uncertainty" corresponds to the classical modelling where the input uncertainty 434 is not taken into account but the output uncertainty is set to its experimental value 435 u(Tsi) = 0.5 °C. The case "low level input uncertainty" corresponds to the case where 436 all experimental uncertainties would have been underestimated. The case "medium 437 level uncertainty" corresponds to taking into account all the experimental values of 438 uncertainties provided by LNE. This configuration is recommended a priori. 439

The effect of duration is studied for the recommended "medium level of uncertainty" and the effect of uncertainty is studied for the duration 24 hours. The resulting labels used in the legends of Figures 9, 10 and 11, are displayed in Table 4.

To conduct the Bayesian analysis, uniform prior distributions are assigned to the calibration parameters, whose bounds are displayed in Table 5. The resulting prior for the thermal resistance is obtained with Monte Carlo simulations from (12) and displayed as the grey histogram in Figure 10. The prior distribution for the excess



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Figure 7. Surface temperature measurements for the IWI wall placed in the energy room REBECCA at LNE, France.



Figure 8. Absorbed flux measurements for the IWI wall placed in the energy room REBECCA at LNE, France.

variance parameter is chosen as $\pi(\sigma^2) \sim \text{InvChi}^2(\nu_0, s_0^2)$ with $\nu_0 = 2$ and $s_0^2 = 0$ to make it poorly informative as explained in section 2.2. Prior distributions of uncertain inputs in the vector $X = (Tse^T, q_{int}^T)^T$ are modeled following (4) with

$$\Sigma_{\zeta} = \begin{pmatrix} \Sigma_{Tse} & 0\\ 0 & \Sigma_{q_{int}} \end{pmatrix}$$
(13)

where $\Sigma_{Tse} = \text{diag}(\{u^2(Tse)\}_N)$ and $\Sigma_{q_{int}} = \text{diag}(\{u^2(q_{int})\}_N)$ are the diagonal matrices 450 of dimension $N \times N$ of reported variances of Tse and q_{int} respectively. Similarly, the 451 covariance matrix of output observations y is chosen diagonal as $\Sigma_{\epsilon} = \text{diag}(\{u^2(Tsi)\}_N)$. 452 The MCMC procedure described in Algorithm 1 was run for each case from Table 453 4 with 50000 iterations, a burn-in of 20000 iterations and a thinning of 30 iterations. 454 For each case, we controlled that the acceptance rate falls between 15% and 30%. For 455 the "no input uncertainty" case the number of parameters is 8 (number of calibration 456 parameters), whereas in the two other cases the propagation of input uncertainties 457 involves 2N additional parameters corresponding to the elements of vector X, for a 458 total of 586 parameters. 459

| Level of uncertainty | u(Tsi) /°C | u(Tse) /°C | $\frac{u(q_{int})}{/\mathrm{Wm^{-2}}}$ |
|----------------------|---------------|---------------|--|
| no input uncertainty | 0.5 | - | - |
| low level | 0.1 | 0.1 | $1\% q_{int}$ |
| medium level | 0.5 | 0.5 | $3\% q_{int}$ |

 Table 3. Description of uncertainty configurations.

| | Duration | | | | |
|----------------------|--------------------|------------------------|--|--|--|
| Level of uncertainty | 12h | 24h | | | |
| no input uncertainty | _ | 24h_no_uncertainty | | | |
| low level | - | $24h_T001_Q1pct$ | | | |
| medium level | $12h_T005_Q3pct$ | $24h_{-}T005_{-}Q3pct$ | | | |

Table 4. Labels used in the legends of plots resulting from the choice of duration and uncertainty configurations.

Results for the thermal conductivity of the insulation layer (EPS, Expanded Polystyrene) Figure 9 shows the effect of duration on the estimation of the thermal conductivity k_2 for the medium uncertainty level. The comparison with the guarded hot plate result (GHP, see section 3.5.1) displayed as the black dashed Gaussian distribution shows that the posterior estimate of k_2 is biased for 12 hours whereas the posterior estimate for 24

| Layer | Thickness /m | Thermal conductivity k /W m ⁻¹ K ⁻¹ | Heat capacity of unit area cw $/{\rm Jm^{-2}K^{-1}}$ |
|------------------|-----------------|--|--|
| Plasterboard | $l_1 = 0.013$ | $0.2 \le k_1 \le 0.4$ | $7\times 10^5 \le cw_1 \le 8\times 10^5$ |
| EPS | $l_2 = 0.120$ | $0.02 \le k_2 \le 0.04$ | $1 \times 10^4 \le cw_2 \le 3 \times 10^4$ |
| Cinderblock | $l_3 = 0.150$ | $0.7 \le k_3 \le 1.2$ | $8.5 \times 10^5 \le cw_3 \le 2 \times 10^6$ |
| Exterior coating | $l_4 = 0.015$ | $0.5 \le k_4 \le 1.2$ | $1 \times 10^6 \le cw_4 \le 2 \times 10^6$ |

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Table 5. Caracterisation of each layer: thickness (assumed known in this application) and bounds for the uniform prior distributions of the thermal parameters, found in the literature.

hours is unbiased (the posterior distribution of k_2 is centered on the GHP best estimate).

⁴⁶⁶ In other words, for 24 hours, the estimate of the thermal conductivity of the insulation

467 material (k_2) is consistent with the physical value assessed with the guarded hot plate

468 method.



Figure 9. Histograms of the posterior samples for k_2 obtained after 12 hours and 24 hours and plots of the Gaussian distribution associated with the guarded hot plate (GHP) result and the prior distribution.

Table 6 gives the posterior estimates for all the uncertainty configurations. It can be observed that for 24 hours, all posterior results are consistent with the GHP result (the best estimate of the GHP 0.031 is included in the 95% Bayesian credible intervals).

It is important to note that the number of significant digits provided on the GHP result does not allow to compare the posterior results in terms of bias because all rounded posterior mean results are unbiased equal to 0.031. So the comparison with the GHP result does not allow to study the effect of uncertainty but globally says that the Bayesian methodology provides unbiased estimates of k_2 (w.r.t GHP result) and then of R_2 for a 24 h observation.

| Duration /h | u(Tsi) /°C | u(Tse) /°C | $u(q_{int})$ /W m ⁻² | \hat{k}_2 /W m ⁻¹ K ⁻¹ | $u(\hat{k}_2)$ /W m ⁻¹ K ⁻¹ | Shortest 95% CI $/W m^{-1} K^{-1}$ |
|----------------|---------------|---------------|------------------------------------|---|--|------------------------------------|
| 24 | 0.1 | - | _ | 0.0309 | 0.0001 | [0.0307, 0.0311] |
| 24 | 0.5 | - | - | 0.0309 | 0.0001 | [0.0307, 0.0311] |
| 24 | 0.1 | 0.1 | $1\% q_{int}$ | 0.0309 | < 0.0001 | [0.0307, 0.0310] |
| 12 | 0.5 | 0.5 | $3\% q_{int}$ | 0.0337 | 0.0003 | [0.0333, 0.0342] |
| 24 | 0.5 | 0.5 | $3\% q_{int}$ | 0.0310 | < 0.0001 | [0.0308, 0.0312] |

Table 6. Results of the sensitivity study for the thermal conductivity k_2 .

Results for the thermal resistance of the wall Figure 10 displays the results for the 478 global resistance of the wall obtained with the guarded hot box method (GHB, see 479 section 3.5.2) represented by a Gaussian approximation with 95% coverage interval 480 (dashed vertical bars), the estimate of R_2 from the GHP result (see Table 1) and its 481 95% coverage interval (in black) and the posterior estimates of the thermal resistance 482 and their associated 95% credible intervals (represented with horizontal turquoise and 483 yellow dashes) obtained with the active method for the medium level of uncertainty for 484 12 hours and 24 hours respectively. Note that the histogram plot of the prior distribution 485 shows that the prior for the global thermal resistance of the wall is consistent with the 486 GHB result. 487

The 95% coverage interval for the GHB result covers all results (even those incomplete for only 12 hours) which means that the comparison which the GHB result cannot be used alone to evaluate the performance of the Bayesian method in the various configurations. We can underline that active solicitation coupled with Bayesian technique, which was developed for in-situ application, can lead to identification results as accurate as those obtained in GHP.

The ability of the Bayesian method to provide an unbiased estimate of R_2 was demontrated previously. Since R_2 represents 95% of the expected global thermal resistance R, the fact that the posterior estimate of the global resistance for 24 hours is closed to the best estimate of the GHB makes highly credible the ability of the Bayesian method to also provide an unbiased estimate of R for 24 hours.

⁴⁹⁹ The comparison with the *thermal resistance of the EPS* R_2 obtained from the GHP ⁵⁰⁰ estimate of k_2 (black line) also shows that the estimate of the thermal resistance obtained ⁵⁰¹ after 12 hours (blue line) nearly allows to retrieve the thermal resistance of the EPS.

It is important to note that the comparison with the GHB result does not allow to study finely the effect of uncertainty in terms of resulting bias on the global thermal resistance.

Results of the sensitivity study in terms of the global thermal resistance are displayed in Table 7. We observe that the posterior uncertainty is similar (with two significant digits) for the configurations involving 24 hours records.

Finally, the plot of the posterior distribution of the thermal resistance of the wall for 24 hours and medium input uncertainty is displayed in Figure 11 with its 95% credible interval.

⁵¹¹ 3.5.4. Remarks For both experiments conducted in a climatic chamber with similar ⁵¹² controlled environmental conditions, the active method developed in the RESBATI ⁵¹³ project allows to estimate the thermal resistance with a relative uncertainty of less ⁵¹⁴ than 2% instead of 10% for the guarded hot box.

More generally, for in-situ measurements, the use of an active method should lessen the influence of uncontrolled environmental conditions so that, in practice, the same modelling should be used for in-situ measurements than for measurements obtained in the climatic chamber. Precisely, a sharp modelling is not required for the complex input quantities like weather, occupation,...

| Duration /h | u(Tsi) /°C | u(Tse) /°C | $\begin{array}{c} u(q_{int}) \\ / \mathrm{W}\mathrm{m}^{-2} \end{array}$ | $\hat{R} \\ /\mathrm{m}^2\mathrm{K}\mathrm{W}^{-1}$ | $\begin{array}{c} u(\hat{R}) \\ /\mathrm{m}^2\mathrm{K}\mathrm{W}^{-1} \end{array}$ | Shortest 95% CI $/m^2 K W^{-1}$ |
|----------------|---------------|---------------|--|---|---|---------------------------------|
| 24 | 0.1 | - | - | 4.064 | 0.014 | [4.039, 4.091] |
| 24 | 0.5 | - | - | 4.072 | 0.014 | [4.045, 4.101] |
| 24 | 0.1 | 0.1 | $1\% q_{int}$ | 4.083 | 0.014 | [4.054, 4.107] |
| 12 | 0.5 | 0.5 | $3\% q_{int}$ | 3.759 | 0.025 | [3.712, 3.808] |
| 24 | 0.5 | 0.5 | $3\% q_{int}$ | 4.058 | 0.014 | [4.031, 4.087] |

Table 7. Results of the sensitivity study for the global thermal resistance.

520 4. Conclusion and discussion

This paper presents a Bayesian approach for the indirect measurement of thermal 521 parameters of a wall from thermal measurements obtained with thermocouples and 522 fluxmeters using an active method. This work is part of the ANR RESBATI project 523 whose main objective is to develop a portable measurement device for evaluating the 524 thermal resistance of opaque building walls on site. The case study chosen in this paper 525 concerns the analysis of measurements performed on an internal wall insulation (IWI) 526 built within the project. The work takes advantage of the experimental work performed 527 on the wall for its global thermal caracterisation in a guarded hot box and for the 528

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Figure 10. In grey: prior distribution of R, in chocolate: Gaussian distribution associated with the GHB estimate of R and its 95% coverage interval, in black: GHP estimate of R_2 and its 95% coverage interval, in turquoise: Bayesian estimate of R after 12 h with the medium uncertainty level and its 95% credible interval, in yellow: Bayesian estimate of R after 24 h with the medium uncertainty level and its 95% credible interval, is 95% credible interval.

thermal caracterisation of the insulation layer performed with a guarded hot plate, for assessing the performance of the Bayesian uncertainty analysis.

The fully Bayesian approach is particularly suited for uncertainty quantification, 531 especially when prior knowledge is available, and is used here for its ability to propagate 532 input uncertainties in an inversion problem, which is in practice an issue rarely addressed 533 with no consensus methodology so far up to our knowledge. The advocated uncertainty 534 propagation approach is easy to apply since it consists in treating all uncertain inputs 535 as calibration parameters and thus requires only a small adaptation of existing MCMC 536 algorithms, for a slightly higher computational cost. For instance, in the case study, we 537 experienced that the computational cost comes primarily from the duration. 538

Under the experimental conditions of the climatic chamber equipped with calibrated thermocouples and fluxmeters, this study demonstrates the efficiency of the Bayesian analysis of the measurements obtained with an active method to produce unbiased (w.r.t the guarded hot plate method) estimates of the thermal conductivity of the insulation layer (EPS) for 24 hours with a low associated standard uncertainty (less than 2%). Nevertheless, smaller observation time, e.g. 12h, has lead to a biased identification.

In this application, it appears that the impact of uncertainties is quite small and cannot be interpreted in terms of resulting bias w.r.t. a reference value. In particular,

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Figure 11. Plot of the posterior distribution of the global thermal resistance of the IWI wall obtained with the active method, and its associated 95% credible interval for 24 hours observations (vertical dashed lines).

⁵⁴⁷ due to the experimental settings we were not able to quantify the impact of input ⁵⁴⁸ uncertainties on the estimated thermal parameters.

For setting input uncertainties, we recommend to perform calibrations and/or verifications of measurement means. If no input uncertainties are available, this study confirms the recommendation from the literature to set a small uncertainty on input parameters to help reducing bias.

The general formulation of the methodology makes it easily applicable to other types of walls and insulation. In this respect, one of the perspectives of the RESBATI project is to address sustainable materials like bio-sourced materials and raw earth considering a hygro-thermal model.

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⁶⁶³ Appendix A. Details on GHB uncertainty evaluation at CSTB, France

The estimated U value for the IWI wall for a mean temperature $T = 10 \,^{\circ}\text{C}$ is

$$U = 0.245 \,\mathrm{W}\,\mathrm{m}^{-2}\,\mathrm{K} \tag{A.1}$$

The global thermal resistance is

$$R = \frac{1}{U} \tag{A.2}$$

Denoting u(U) the standard uncertainty associated with U, GUM [BIPM et al., 2008a] uncertainty propagation gives

$$u(R) = \frac{u(U)}{U^2} = 0.43 \,\mathrm{m}^2 \,\mathrm{K} \,\mathrm{W}^{-1}$$
 (A.3)

Equivalently, the 95% coverage interval associated with R is

$$4.08 \pm 0.86 \,\mathrm{m}^2 \,\mathrm{K} \,\mathrm{W}^{-1}(k=2) \tag{A.4}$$

The following describes the uncertainty evaluation of U. The specimen is installed on a surrounding wall, the measurement zone is larger than the sample surface so that, in steady state conditions, a part of the heating power injected in the measurement zone ϕ_{in} is split between :

- A thermal flow through the metering box, ϕ_{out} , controlled to be neglectable by controlling the guarded zone temperature equal to the measuring zone temperature. The residual thermal flow through the metering box is then controlled to be neglectable by measuring it using a thermopile ($\phi_{out} = 0$ W);
- A thermal flow through the specimen ϕ_{sp} ;
- A thermal flow through the surrounding wall ϕ_{sur} ;
- A thermal flow through thermal bridge between the sample and surrounding wall ϕ_{edge} ;

The thermal flow through the specimen is then obtained by

$$\phi_{sp} = \phi_{in} - \phi_{sur} - \phi_{edge} \tag{A.5}$$

The U value is then obtained by

$$U = \frac{\phi_{sp}}{A_{sp} \left(\theta_{n,i} - \theta_{n,e}\right)} \tag{A.6}$$

where A_{sp} is the surface of the specimen, $\theta_{n,i}$ and $\theta_{n,e}$ are the environment internal/external temperatures included by air temperature ($\theta_{c,i}, \theta_{c,e}$) and radiant temperature viewed by the specimen ($\theta_{r,i}, \theta_{r,e}$) temperatures:

$$\theta_{n,i} = F_{c,i}\theta_{c,i} + (1 - F_{c,i})\theta_{r,i} \tag{A.7}$$

$$\theta_{n,e} = F_{c,e}\theta_{c,e} + (1 - F_{c,e})\theta_{r,e} \tag{A.8}$$

where $F_{c,i}$ and $F_{c,e}$ are convective fractions.

Radiant temperatures $(\theta_{r,i}, \theta_{r,e})$ are described in Annex A of NF EN ISO 12567-1 [NF EN ISO 12567-1:2013, 2013] as an explicit function of surface temperatures of the specimen $(\theta_{s,i}, \theta_{s,e})$ and the baffle $(\theta_{b,i}, \theta_{b,e})$ panel situated in front of each side.

The thermal flow through the surrounding wall ϕ_{sur} and convective fractions $F_{c,i}$ 681 and $F_{c,e}$ are obtained using a calibration process as described in NF EN ISO 12567-682 1 [NF EN ISO 12567-1:2013, 2013]. This calibration consists in using homogeneous 683 insulated calibration panels as specimen. The thermal resistance of these calibration 684 panels have been measured using guarded hot plate method so that their R-values are 685 well known. Measuring surface temperatures during the calibration test in the guarded 686 hot box make possible determining both R_{sur} (and then ϕ_{sur}) as a function of average 687 support wall temperature, and convective fractions as a function of thermal flow through 688 the specimen, using linear regression techniques. The linear regression coefficients are 689 then used to predict ϕ_{sur} , $F_{c,i}$ and $F_{c,e}$ during all experiment. The uncertainty on this 690 prediction is estimated using less squares methods [XP ISO/TS 28037:2013, 2013]. 691

The uncertainty on the thermal flow through thermal bridge between the sample and surrounding wall ϕ_{edge} is estimated by simulating each junction [NF EN ISO 10211:2017, 2017]. A sensibility analysis has been provided on each unknown input: thermal conductivity of wood element for instance may vary between 0.13 to 0.18 W/(m.K). The possible variation interval of the calculated psi-values, estimated with this sensibility analysis, is then used to define a uniform law with standard deviation:

$$u_{\psi_{edge}} = \frac{\psi_{edge_max} - \psi_{edge_min}}{\sqrt{3}} \tag{A.9}$$

⁶⁹² Finally, all errors are propagated using GUM recommendations [BIPM et al., 2008a].

⁶⁹³ Appendix B. R code

The R code associated with the low and medium uncertainty levels implementing 1 is given in the Supplementary material.



Figure A1. Notations for uncertainty propagation in GHB at CSTB, France.