



Anti-reflection subwavelength gratings for InP-based waveguide facets

Lukas Puts, Xaveer Leijtens, Pavel Cheben, Jens Schmid, Sander Reniers,
Daniele Melati

► To cite this version:

Lukas Puts, Xaveer Leijtens, Pavel Cheben, Jens Schmid, Sander Reniers, et al.. Anti-reflection subwavelength gratings for InP-based waveguide facets. *Optics Letters*, 2021, 46 (15), pp.3701-3704. 10.1364/OL.431353 . hal-04346315

HAL Id: hal-04346315

<https://hal.science/hal-04346315>

Submitted on 15 Dec 2023

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Anti-reflection subwavelength gratings for InP-based waveguide facets

LUKAS PUTS^{1,*}, XAVEER LEIJTENS¹, PAVEL CHEBEN², JENS SCHMID², SANDER RENIERS¹, AND DANIELE MELATI³

¹Photonic Integration Group, department of Electrical Engineering, Eindhoven University of Technology, P.O. Box 513, Eindhoven 5600MB, The Netherlands

²National Research Council Canada, Ottawa, Ontario K1A 0R6, Canada

³Centre for Nanoscience and Nanotechnologies, CNRS, Université Paris-Saclay, 10 Bv. Thomas Gobert, 91120 Palaiseau, France

*Corresponding author: l.puts@tue.nl

Compiled December 15, 2023

We demonstrate the anti-reflection properties of lithographically defined subwavelength gratings applied to the facets of InP-based integrated waveguides. The subwavelength gratings are based on the gradient index effect to create a smooth index transition between the core material and air, achieving reflections below -30 dB at a wavelength of 1550 nm for both TE and TM polarized modes, as shown by 3D finite-difference time-domain simulations. Characterizations performed on the experimental InP membrane-on-silicon platform using Mach-Zehnder interferometers test structures are in excellent agreement with simulation predictions. ©

2023 Optical Society of America

<http://dx.doi.org/10.1364/ao.XX.XXXXXX>

Controlling reflections at waveguide facets is crucial for optimal performance of photonic integrated devices and circuits. A notable example is the strong effect that small reflections can have on the dynamics of integrated lasers [1, 2]. For fiber-to-chip edge couplers low reflections are needed to achieve a high coupling efficiency. Waveguide facets represent one of the main sources of reflections, generated by the abrupt index transitions between the waveguide materials and air. For typical single mode waveguides in silicon and III-V semiconductor platforms, cleaved facets without adaptation structures have a strong reflectivity of about 30% [3]. Commonly, single or multilayer anti-reflection (AR) thin films with a precisely engineered thickness and refractive index are used to reduce these reflections. However, this comes with some important drawbacks. Complex vacuum deposition techniques are required to apply films on a cleaved facet, which can be done only as a post-processing step, increasing fabrication complexity. Furthermore, the films may become mechanically unstable under thermal cycling, resulting in loss of performance [4]. Another approach that is often applied is the use of angled facets. By changing the facet angle, reflections of about -40 dB for the fundamental mode has been demonstrated in III-V waveguides [5]. However, by increasing the facet angle, higher order mode reflections can be generated. This can be especially problematic in wide waveguides support-

ing higher order modes [6].

Another strategy to minimize reflections is the patterning of the waveguide facet with a subwavelength grating (SWG). This type of structure has been used in a wide range of applications to control reflections. For example, SWGs have been applied on bulk surfaces such as solar cells to optimize performance [7]. The same concept has been exploited to reduce reflections from the facets of quantum cascade lasers by patterning the facets using focused ion beam (FIB) milling after chip fabrication. Reflections as low as 1-3% have been achieved using this method [8]. However, FIB milling can only be applied during post-processing. Moreover, ion sputtering and a diverging ion beam resulted in significant differences between the designed and the fabricated grating geometry, indicating FIB might not be ideal to fabricate SWGs. The use of lithographically defined SWGs as anti-reflective coatings, first demonstrated by Schmid et al. [4], eliminates the need for chip-level post-processing. A grating with triangular teeth was fabricated using electron beam lithography (EBL) on a Silicon-on-Insulator (SOI) platform. Reflections as low as 2% and 2.4% for the fundamental TE and TM polarized modes were achieved, respectively [9]. While extensively studied for SOI waveguides [10–13], little investigations have been reported so far on the use of SWGs in InP-based platforms. The possibility to control reflections by subwavelength patterned geometries is of particular interest because of the wide range of active components available in III-V platforms, which often heavily rely on thin-film AR facet coatings.

In this paper, we simulate, design and experimentally demonstrate an AR subwavelength waveguide facet grating on the InP membrane on silicon (IMOS) platform [14]. The IMOS platform comprises a 300 nm thick InP photonic membrane on a standard silicon wafer. The layer stack for a standard 400 nm wide single-mode IMOS waveguide is shown in Fig. 1(a). First, the InP waveguide pattern is transferred to a small InP wafer using EBL. The wafer is then flipped and bonded to a silicon wafer using the adhesive polymer divinylsiloxane-bis-benzocyclobutene (DVS-BCB) and a 200 nm thick SiO₂ bonding layer. The InP substrate and InGaAs etch-stop layers are removed by selective wet-chemical etching. The thin patterned InP membrane remains on the silicon wafer [14]. Due to the large index difference between InP ($n=3.17$), BCB ($n=1.54$), and SiO₂ ($n=1.45$), very

compact waveguides can be realized. The resulting electric field distribution for the fundamental TE and TM modes is shown in Fig. 1(c) and (d). Most of the electric field of the fundamental TE mode is concentrated in the waveguide core, whereas the TM mode is primarily localized in the oxide cladding below and air above the core. To realize the AR structure, the waveguide core is patterned with a gradient index (GRIN) SWG consisting of an array of triangles, as depicted in Fig. 1(b). When the period of the grating is smaller than the wavelength in the medium, the structure can be approximated as an equivalent material with a refractive index gradually changing from that of the waveguide core to the index of the cladding [15]. The gradual transition from core to cladding suppresses facets reflections by eliminating the abrupt termination of the waveguide. Since index grading does not rely on interference effects, it is inherently less wavelength sensitive compared to common thin film AR layers [4].

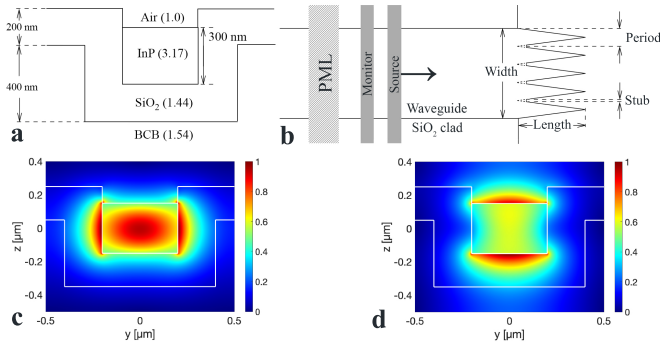


Fig. 1. (a) Schematic cross-section of a standard IMOS waveguide, indicating the materials with respective refractive indices and dimensions. (b) Top view of the SWG patterned on the waveguide facet and FDTD simulation layout. (c) and (d) Fundamental TE and TM electric field profiles in a 400 nm wide IMOS waveguide, respectively.

A combination of 3D finite-difference time-domain (FDTD) simulations for accurate reflection evaluation and 2D FDTD for tolerance analysis was used to design the AR structures. The model comprises a waveguide and a triangular grating with a fixed subwavelength period of 400 nm, as shown in Fig. 1(b). Grating length is varied from 0 to 2 μm , and the waveguide width is varied from 2 to 4 μm . The fundamental TE or TM mode is injected from the left at $\lambda_0 = 1550 \text{ nm}$ towards the grating. A power monitor left to the source measures the amount of light reflected from the SWG structure.

The TE and TM reflections calculated by 3D FDTD as a function of the grating teeth length are shown in Fig. 2 and 3, respectively. For clarity, the simulation results for a 2 μm wide SWG are marked with the blue dots, the unmarked curves show the results for a 4 μm wide SWG. The dash-dotted lines show the total reflections (TE), the solid lines show the modal reflections (TE₀). The latter refers to the amount of reflected power coupled into the counter propagating fundamental mode. At a grating teeth length of 0 nm the structure is simply a flat facet with a reflection of $\sim 25\%$, as expected (Fig. 2). From these simulations it is clear that the facet reflection decreases rapidly with increasing grating teeth length. The TE₀ reflection drops to about -30 dB for a grating teeth length of approximately 0.5 μm . The difference between the modal and total reflection simulations is small, indicating that mode conversion is small. The oscillatory behavior

of the reflection results in three other minima at 1 μm , 1.4 μm and 1.85 μm . Simulated reflections for these minima are between -30 and -35 dB. The width of the waveguide has only a small influence on the reflection.

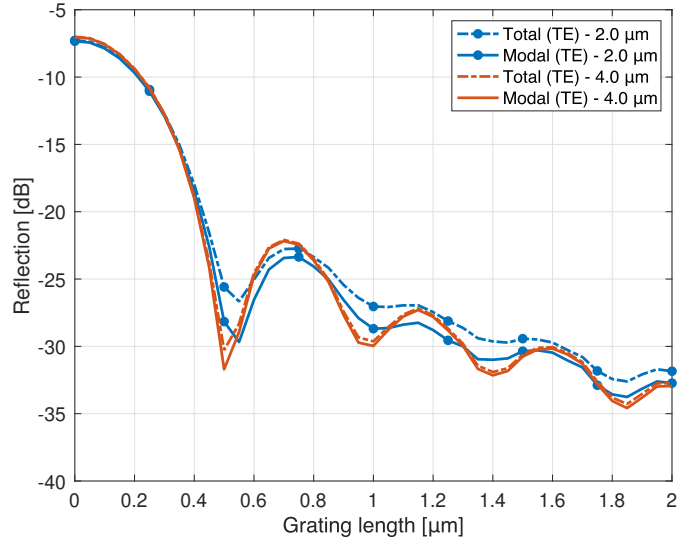


Fig. 2. 3D FDTD simulation results for the reflection of the TE polarized mode as a function of the grating teeth length.

From these simulations, the preferred design spot for a substantial reduction in TE reflection is the first minimum, at about 0.51 μm . This results in a compact device with the smallest possible aspect ratio. Even lower reflection could be achieved for longer grating teeth lengths, that however would require structures with a larger aspect ratio, making fabrication more challenging.

The results for TM simulations are shown in Fig. 3. The simulated structures are the same as described in the previous paragraph. Without any grating teeth (length of 0 nm) the reflections are around 30%, which decrease to values below -30 dB for grating teeth length larger than 1.4 μm . The SWG teeth are parallel to the oscillations of the electric field for the TM polarization and as a result the anti-reflection effect is less strong compared to TE polarization. This was previously observed also in AR SWGs on SOI waveguides [4]. The preferred design spot in this case is 1.55 μm , at which a reflection of about -35 dB is achieved.

Next, fabrication tolerances of the proposed designs are investigated. Due to the diffraction limit and electron backscattering in the photoresist, the resolution for EBL is limited to approximately 20 nm. There are two regions where the SWG feature size is below 20 nm and are hence difficult to define using EBL. The tip of each tooth will result in a rounded shape with a certain radius and the region between the bases of two adjacent teeth will not be defined accurately. Additionally, for an SWG with high aspect ratios, latter regions can be also difficult to etch, due to the etch lag effect [16]. In order to ensure a reliable fabrication, a 'stub' is placed between the bases of adjacent teeth, as shown in Fig. 1 (b). The size of this stub is made equal to the resolution of the EBL, decreasing the etch lag effect in this particular region. 2D FDTD simulations show that a stub of 20 nm adds approximately 3.2 and 2.1 dB to the total and modal TE reflections, respectively. Random and systematic process-induced fluctuations such as variations in focus, dose, and resist quality influence critical dimensions (CD). These fluctuations may result in uniform CD variations of several nanometers with respect to

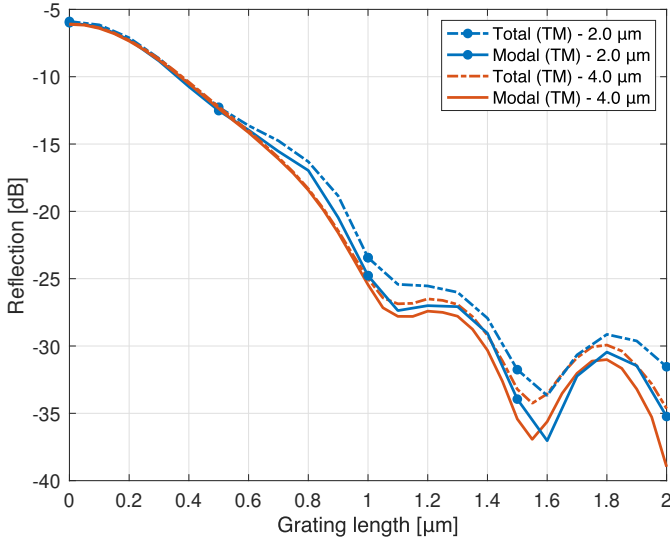


Fig. 3. 3D FDTD simulation of the reflection of the TM polarized mode as a function of the grating teeth length.

the nominal design. 2D FDTD simulations for the grating with a waveguide width of 4 μm , length of 0.51 μm , and a stub of 20 nm were carried out to investigate the effect of a uniform CD variation. Changing all structure dimensions by 10 nm resulted in an increase of the facet reflection of approximately 3 dB.

To experimentally characterize reflections induced by the facet grating, we designed the unbalanced Mach-Zehnder interferometer (MZI) schematically shown in Fig. 4. The MZI is based on three 2×2 multimode interference couplers (MMI), leveraging their low backscattering characteristics [17] and nominal waveguides with a width of 400 nm. One of the arms contains an MMI coupler and two 62 μm long linear adiabatic tapers terminated with two identical SWG gratings R_1 and R_2 . The phase shift $\Delta\phi$ in this arm is the result of the extra optical path length of approximately 1135 μm . Surface grating couplers are used to couple light into and out of the circuit (P_{in} , P_2 , P_3 , P_4).

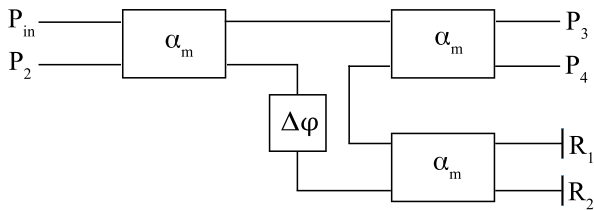


Fig. 4. Block diagram of the unbalanced MZI used for AR SWG characterization. Input and outputs ports P_{in} - P_4 and the SWGs grating under test R_1 and R_2 are marked.

In order to use this structure to measure the SWG reflectivity, the output power of the bar (P_3) and cross (P_4) ports as a function of the wavelength can be derived using the expression for the transfer function of a single MZI [18]. The contrast ratio C , defined as the ratio between the minima and maxima of the fringes in the MZI transmission spectrum for either of the two output ports P_3 and P_4 , can be calculated as follows:

$$C = \frac{P_{\text{out, max}}}{P_{\text{out, min}}} = \frac{\left((\sqrt{R_1} + \sqrt{R_2}) \alpha_m e^{-\alpha l} + 2 \right)^2}{\left((\sqrt{R_1} + \sqrt{R_2}) \alpha_m e^{-\alpha l} - 2 \right)^2} \quad (1)$$

where R_1 and R_2 are the corresponding SWG reflectivities, α_m the MMI coupler loss, α the waveguide propagation loss, and l the MZI arm length difference. Assuming $R_1 = R_2 = R$, the reflectivity R can hence be determined from the measured contrast ratio. A total of 20 MZIs with 4 μm wide SWGs and grating teeth length ranging from 0 to 2 μm were designed on a $4 \times 6 \text{ mm}^2$ IMOS cell. The grating period was fixed at 400 nm. Two identical cells were patterned by EBL and proximity effect correction (PEC) and shape correction were applied to the GDS file to compensate for electron back scattering and optimize the dose locally.

TE polarized light was coupled in and out of the circuit using surface grating couplers optimized for TE coupling (Fig. 5). Due to variations in the fabrication, the optimal coupling efficiency of the surface grating couplers was shifted to 1535 nm from the designed 1550 nm. The wavelength of the input light was swept from 1465 to 1575 nm using a tuneable laser source. The chip temperature was maintained at 21°C using a Peltier element mounted on a copper chuck. The output light intensity was measured by a power meter.

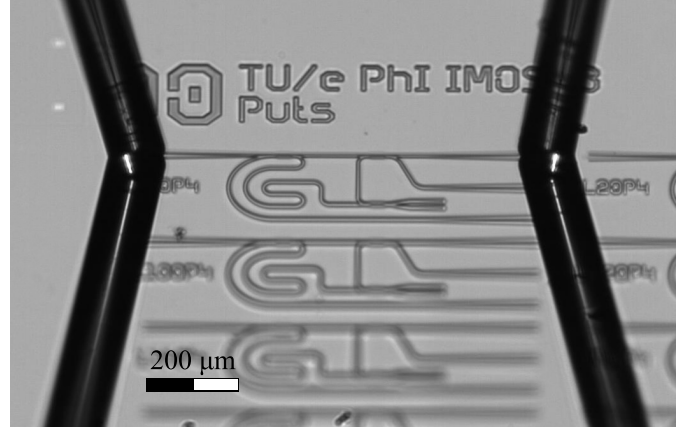


Fig. 5. Photograph of part of the IMOS chip showing three Mach-Zehnder interferometers and the input and output fibers at an angle of 10° for optimal coupling.

Two examples of the measured spectra at the output port P_3 for the test MZI with SWGs teeth lengths of $L=200$ and $L=500$ nm are shown in Fig. 6. The latter is the preferred design spot for the TE polarized mode and, as expected, its anti-reflection properties translate into a reduced contrast ratio of the spectral fringes compared to the SWG with a grating teeth length of 200 nm.

The mean and standard deviation of the contrast ratio C for each of the 20 MZIs of one design cell was determined from multiple fringes in the spectral range from 1520 to 1540 nm of the individual devices. The calculated relative reflections derived from Eq. (1) are then normalized using the 3D FDTD simulations of the SWG with a grating teeth length of 200 nm. The solid line in Fig. 7 shows the results of this analysis as a function of the grating teeth length. Squares mark the mean relative reflections and the error bars the standard deviations. The dashed line shows the modal TE 3D FDTD simulation results for reference (same data as in Fig. 2). Results of the simulations and measurements are in a very good agreement and show a similar trend of reflection as function of grating teeth length. The reflections drop rapidly with increasing grating teeth length and similar to the simulations the first minimum is located approximately at 500 nm at a value of -18 dB below the expected flat facet reflec-

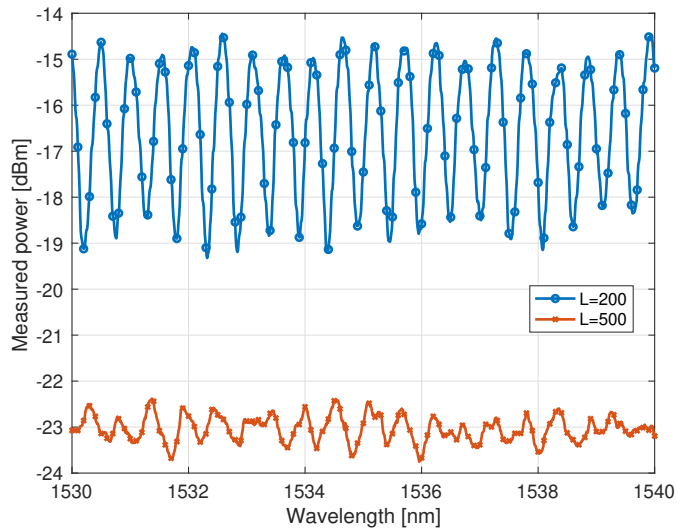


Fig. 6. Output power at P_3 of two MZIs with facet grating lengths of 200 nm and 500 nm. Lower fringe contrast of the SWG with longer grating teeth length is due to the AR properties. Contrast ratio for these test structures is 2.484 ± 0.332 and 1.176 ± 0.0815 for lengths of 200 and 500 nm, respectively. To prevent cluttering of data points, only a quarter of the measurements are marked.

tion of -6.8 dB. Relative reflectivity between 600 and 1600 nm is substantially higher compared to the simulations. This is likely related to the challenges in etching SWGs with large aspect ratios. Moreover, the reflection of a flat facet is substantially lower than what is expected from simulations. A possible explanation for this is that non-perfect vertical etching may have resulted in a slightly slanted facet which would suppress coupling of the reflected light to the waveguide mode.

In conclusion, we have designed and experimentally demonstrated an effective anti-reflection strategy for InP IMOS waveguides by exploiting subwavelength gratings. These lithographically defined structures are particularly interesting for InP-based technology platforms since the performance of active components often relies heavily on anti-reflection coatings deposited during post-processing. Simulated TE_0 mode reflections are as low as -30 dB for an SWG with a grating teeth length of 0.51 μm , width of 4 μm and a fixed period of 400 nm. This results in a compact design with a relatively low aspect ratio, which is beneficial for fabrication. Reflections as low as -35 dB are achieved for the TM_0 mode for SWGs with a grating teeth length of 1.55 μm . The anti-reflection effect was experimentally demonstrated on waveguide facets using Mach-Zehnder interferometers as test structures. Reflections for a 4 μm wide SWG and grating teeth lengths varied from 0 to 2 μm yield excellent agreement with 3D FDTD simulations.

Acknowledgments. The authors would like to thank Marc Spiegelberg and Jeroen Bolk for their contribution to the design and characterization. Nazca Design was used to generate the mask layout in this work.

Disclosures. The authors declare no conflicts of interest.

REFERENCES

1. P. Zorabedian, *IEEE J. Quantum Electron.* **30**, 1542 (1994).

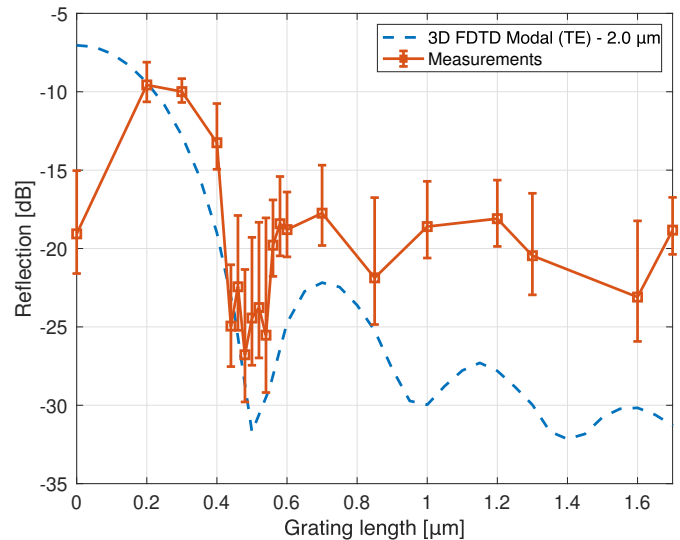


Fig. 7. Overlay of 3D FDTD simulations and experimental results in the IMOS platform. The error bars show the standard deviation in the measurement.

2. K. Petermann, *IEEE J. Sel. Top. Quantum Electron.* **1**, 480 (1995).
3. D. Melati, F. Morichetti, and A. Melloni, *Opt. Quantum Electron.* **45**, 309 (2013).
4. J. Schmid, P. Cheben, S. Janz, J. Lapointe, A. Del  ge, A. Densmore, B. Lamontagne, P. Waldron, and D. Xu, *Opt. Express* **6796** (2007).
5. Y. Matsuoaka, S. Mathonni  re, M. P. Semtsiv, and W. T. Masselink, "Angled facet waveguide quantum cascade laser for external cavity system," *International Society for Optics and Photonics (SPIE)*, (2018), pp. 176 – 181.
6. P. Zorabedian, "Tunable external-cavity semiconductor lasers," (Academic Press, San Diego, 1995), *Optics and Photonics*, pp. 349–442.
7. N. Das and S. I. Sen, "Optimization of nano-grating structure to reduce the reflection losses in GaAs solar cells," (2012), pp. 1–5.
8. A. O. Dirisu, G. Silva, Z. Liu, C. F. Gmachl, F. J. Townner, J. Bruno, and D. L. Sivco, *IEEE Photonics Technol. Lett.* **19**, 221 (2007).
9. J. Schmid, P. Cheben, S. Janz, J. Lapointe, E. Post, and D.-X. Xu, *Opt. Lett.* **32**, 1794 (2007).
10. P. Cheben, R. Halir, J. Schmid, H. Atwater, and D. Smith, *Nature* **560** (2018).
11. R. Halir, A. Ortega-Mo  ux, D. Benedikovic, G. Z. Mashanovich, J. G. Wang  emert-P  rez, J. H. Schmid, I. Molina-Fern  ndez, and P. Cheben, *Proc. IEEE* **106**, 2144 (2018).
12. R. Halir, P. J. Bock, P. Cheben, A. Ortega-Mo  ux, C. Alonso-Ramos, J. H. Schmid, J. Lapointe, D.-X. Xu, J. G. Wang  emert-P  rez, I. Molina-Fern  ndez, and S. Janz, *Laser & Photonics Rev.* **9**, 25 (2015).
13. J. Schmid, P. Cheben, S. Janz, J. Lapointe, E. Post, A. Del  ge, A. Densmore, B. Lamontagne, P. Waldron, and D.-X. Xu, *Adv. Opt. Technol.* **2008** (2008).
14. J. J. G. M. van der Tol, Y. Jiao, J. P. Van Engelen, V. Pogoretskiy, A. A. Kashi, and K. Williams, *IEEE J. Quantum Electron.* **56**, 1 (2020).
15. S. M. Rytov, *Sov. Phys. JETP* **2**, 466 (1956).
16. J. Bolk, "ArF scanner lithography for InP photonic integrated circuit fabrication," Ph.D. thesis, Eindhoven University of Technology (2020).
17. E. Kleijn, M. K. Smit, and X. J. M. Leijtens, "Analysis of parasitic effects in PICs using circuit simulation," *International Society for Optics and Photonics (SPIE)*, (2013), pp. 12 – 20.
18. C. K. Madsen and J. H. Zhao, *Digital Filter Concepts for Optical Filters* (John Wiley & Sons, Ltd, 2001), chap. 3, pp. 95–164.

FULL REFERENCES

1. P. Zorabedian, "Axial-mode instability in tunable external-cavity semiconductor lasers," *IEEE J. Quantum Electron.* **30**, 1542–1552 (1994).
2. K. Petermann, "External optical feedback phenomena in semiconductor lasers," *IEEE J. Sel. Top. Quantum Electron.* **1**, 480–489 (1995).
3. D. Melati, F. Morichetti, and A. Melloni, "Modeling reflections induced by waveguide transitions," *Opt. Quantum Electron.* **45**, 309–316 (2013).
4. J. Schmid, P. Cheben, S. Janz, J. Lapointe, A. Delâge, A. Densmore, B. Lamontagne, P. Waldron, and D. Xu, "Subwavelength grating structures in planar waveguide facets for modified reflectivity," *Opt. Express* **6796** (2007).
5. Y. Matsuoka, S. Mathonnière, M. P. Semtsiv, and W. T. Masselink, "Angled facet waveguide quantum cascade laser for external cavity system," *International Society for Optics and Photonics (SPIE)*, (2018), pp. 176 – 181.
6. P. Zorabedian, "Tunable external-cavity semiconductor lasers," (Academic Press, San Diego, 1995), *Optics and Photonics*, pp. 349–442.
7. N. Das and S. I. Sen, "Optimization of nano-grating structure to reduce the reflection losses in GaAs solar cells," (2012), pp. 1–5.
8. A. O. Dirisu, G. Silva, Z. Liu, C. F. Gmachl, F. J. Towner, J. Bruno, and D. L. Sivco, "Reduction of Facet Reflectivity of Quantum-Cascade Lasers With Subwavelength Gratings," *IEEE Photonics Technol. Lett.* **19**, 221–223 (2007).
9. J. Schmid, P. Cheben, S. Janz, J. Lapointe, E. Post, and D.-X. Xu, "Gradient-index antireflective subwavelength structures for planar waveguide facets," *Opt. Lett.* **32**, 1794–1796 (2007).
10. P. Cheben, R. Halir, J. Schmid, H. Atwater, and D. Smith, "Subwavelength integrated photonics," *Nature* **560** (2018).
11. R. Halir, A. Ortega-Moñux, D. Benedikovic, G. Z. Mashanovich, J. G. Wangüemert-Pérez, J. H. Schmid, I. Molina-Fernández, and P. Cheben, "Subwavelength-Grating Metamaterial Structures for Silicon Photonic Devices," *Proc. IEEE* **106**, 2144–2157 (2018).
12. R. Halir, P. J. Bock, P. Cheben, A. Ortega-Moñux, C. Alonso-Ramos, J. H. Schmid, J. Lapointe, D.-X. Xu, J. G. Wangüemert-Pérez, I. Molina-Fernández, and S. Janz, "Waveguide Sub-Wavelength Structures: a Review of Principles and Applications," *Laser & Photonics Rev.* **9**, 25–49 (2015).
13. J. Schmid, P. Cheben, S. Janz, J. Lapointe, E. Post, A. Delâge, A. Densmore, B. Lamontagne, P. Waldron, and D.-X. Xu, "Subwavelength grating structures in silicon-on-insulator waveguides," *Adv. Opt. Technol.* **2008** (2008).
14. J. J. G. M. van der Tol, Y. Jiao, J. P. Van Engelen, V. Pogoretskiy, A. A. Kashi, and K. Williams, "InP Membrane on Silicon (IMOS) Photonics," *IEEE J. Quantum Electron.* **56**, 1–7 (2020).
15. S. M. Rytov, "Electromagnetic Properties of a Finely Stratified Medium," *Sov. Phys. JETP* **2**, 466–475 (1956).
16. J. Bolk, "ArF scanner lithography for InP photonic integrated circuit fabrication," Ph.D. thesis, Eindhoven University of Technology (2020).
17. E. Kleijn, M. K. Smit, and X. J. M. Leijtens, "Analysis of parasitic effects in PICs using circuit simulation," *International Society for Optics and Photonics (SPIE)*, (2013), pp. 12 – 20.
18. C. K. Madsen and J. H. Zhao, *Digital Filter Concepts for Optical Filters* (John Wiley & Sons, Ltd, 2001), chap. 3, pp. 95–164.