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ASYMPTOTICALLY HOMOGENEOUS SOLUTIONS OF THE SUPERCRITICAL LANE-EMDEN SYSTEM

LOUIS DUPAIGNE, MARIUS GHERGU, AND HATEM HAJLAOUI

ABSTRACT. We consider the Lane-Emden system

$$-\Delta u = |v|^{p-1}v, \quad -\Delta v = |u|^{q-1}u \quad \text{in } \mathbb{R}^d.$$

When $p \geq q \geq 1$, it is known that there exists a positive radial stable solution $(u, v) \in C^2(\mathbb{R}^d)$ if and only if $d \geq 11$ and (p, q) lies on or above the so-called Joseph-Lundgren curve introduced in [4]. In this paper, we prove that for $d \leq 10$, there is no positive stable solution (or merely stable outside a compact set and (p, q) does not lie on the critical Sobolev hyperbola), while for $d \geq 11$, the Joseph-Lundgren curve is indeed the dividing line for the existence of such solutions, if one assumes in addition that they are asymptotically homogeneous (see Definition 1 below). Most of our results are optimal improvements of previous works in the litterature.

1. INTRODUCTION

The Lane-Emden system

$$-\Delta u = |v|^{p-1}v, \quad -\Delta v = |u|^{q-1}u \quad \text{in } \mathbb{R}^d, \tag{1.1}$$

where $d \geq 2$, $p \geq q > 0$, $pq > 1$ and $u, v \in C^2(\mathbb{R}^d)$, has been studied for the past three decades. Yet, not so much is known about its solutions. Thanks to the works of Mitidieri [27] and Serrin and Zou [33], there exists a radial positive solution to the system if and only if the exponents are supercritical i.e.

$$\frac{1}{p+1} + \frac{1}{q+1} \leq 1 - \frac{2}{d}. \tag{1.2}$$

The Lane-Emden conjecture states that, whether radial or not, no positive solution exists in the subcritical case i.e. when (1.2) does not hold. The current best known result is due to Souplet [34], who proved the conjecture for $d \leq 4$, while only partial results are available for $d \geq 5$.

We address here the supercritical case and focus on solutions which are stable outside a compact set, i.e. such that there exists a compact set $K \subset \mathbb{R}^d$ and two positive functions $\phi, \psi \in C^2(\mathbb{R}^d \setminus K)$ such that

$$-\Delta \phi \geq p|v|^{p-1}\psi, \quad -\Delta \psi \geq q|u|^{q-1}\phi \quad \text{in } \mathbb{R}^d \setminus K.$$

We will explain in a short moment why such an assumption is natural. Thanks to the scaling invariance of the equation, there exists a singular solution of the form

$$(u_s, v_s)(x) = (a|x|^{-\alpha}, b|x|^{-\beta}), \quad x \in \mathbb{R}^d \setminus \{0\} \tag{1.3}$$

where the scaling exponents are given by

$$\alpha = \frac{2(p+1)}{pq-1}, \quad \beta = \frac{2(q+1)}{pq-1}$$

and a, b are suitable (explicit) positive constants. In the supercritical regime, thanks to a suitable version of Hardy's inequality, such a solution is stable if and only if (p, q) lies on or above the Joseph-Lundgren curve i.e. when $d \geq 11$ and $H^2 \geq pq\lambda\mu$, where

$$\gamma = \alpha - \beta, \quad H = \frac{(d-2)^2 - \gamma^2}{4}, \quad \lambda = \alpha(d-2-\alpha) \quad \text{and} \quad \mu = \beta(d-2-\beta). \tag{1.4}$$

When $p \geq q \geq 1$, Chen and the first two named authors proved in [4] that there exists a positive radial stable solution $(u, v) \in C^2(\mathbb{R}^d)$ of (1.1) if and only if $d \geq 11$ and (p, q) lies on or above the Joseph-Lundgren curve. In the case $d \leq 10$, we obtain the following optimal improvement.

Theorem 1. *Let $d \leq 10$ and $p \geq q \geq 1$ be such that $pq > 1$. If $(u, v) \in C^2(\mathbb{R}^d)$ is a nonnegative solution which is stable, or merely stable outside a compact set but with*

$$\frac{1}{p+1} + \frac{1}{q+1} \neq 1 - \frac{2}{d}, \quad (1.5)$$

then $u = v = 0$.

In dimension $d \geq 11$, our results are sharp for a restricted class of solutions as we describe next.

Definition 1. *If it exists, a blow-down limit of a solution $(u, v) \in C^2(\mathbb{R}^d)$ is a cluster point⁽¹⁾ (u_∞, v_∞) of the family $(u_R, v_R)_{R \geq 1}$ of rescalings defined by*

$$(u_R, v_R)(x) = (R^\alpha u(Rx), R^\beta v(Rx)) \quad \text{for } x \in \mathbb{R}^d. \quad (1.6)$$

A solution (u, v) is said to be asymptotically homogeneous if all its blow-down limits are homogeneous i.e. if there exists $(f, g) \in C^2(S^{d-1})$ such that $(u_\infty, v_\infty)(r\theta) = (r^{-\alpha} f(\theta), r^{-\beta} g(\theta))$ for $r > 0$ and $\theta \in S^{d-1}$.

With this definition in mind, we obtain the following theorem.

Theorem 2. *Let $p \geq q \geq 1$ be such that $pq > 1$ and*

$$H^2 < pq\lambda\mu, \quad (1.7)$$

where H, λ, μ are given by (1.4). If $(u, v) \in C^2(\mathbb{R}^d)$ is a nonnegative solution which is stable (resp. stable outside a compact set and (1.5) holds) and asymptotically homogeneous, then $u = v = 0$.

Blow-down limits of nonnegative solutions of the Lane-Emden system which are stable outside a compact set are always well-defined:

Theorem 3. *Let $(u, v) \in C^2(\mathbb{R}^d)$ be a nonnegative solution of (1.1) which is stable outside a compact set. Then, the family $(u_R, v_R)_{R \geq 1}$ given by (1.6) is compact in $L_{loc}^{q+1}(\mathbb{R}^d \setminus \{0\}) \times L_{loc}^{p+1}(\mathbb{R}^d \setminus \{0\})$.*

In addition, in the cases $p = q$ and $p > q = 1$, all blow-down limits are homogeneous, thanks to the availability of a monotonicity formula (see [9]). Unfortunately, we do not know if this fact is also true for the Lane-Emden system for other choices of exponents. Still, we can prove that our results continue to hold for a possibly wider class of solutions.

Theorem 4. *Let $p \geq q \geq 1$ be such that $pq > 1$ and (1.7) holds. Let $K \subset \mathbb{R}^d$ be a compact set. If $(u, v) \in C^2(\mathbb{R}^d)$ is a nonnegative solution of (1.1) such that for all $\varphi \in C_c^1(\mathbb{R}^d)$ (resp. for all $\varphi \in C_c^1(\mathbb{R}^d \setminus K)$ and (1.5) holds),*

$$\int_{\mathbb{R}^d} |u|^{\frac{q-1}{2}} |v|^{\frac{p-1}{2}} \varphi^2 dx \leq \frac{1}{\sqrt{pq}} \left(\int_{\mathbb{R}^d} |\nabla \varphi|^2 dx - \frac{\gamma^2}{4} \int_{\mathbb{R}^d} \frac{\varphi^2}{|x|^2} dx \right) \quad (1.8)$$

then $u = v = 0$.

Assumption (1.8) is motivated by the following observation.

Theorem 5. *Let $(u, v) \in C^2(\mathbb{R}^d)$ be a nonnegative solution of (1.1) which is stable outside a compact set and asymptotically homogeneous. Let (u_∞, v_∞) denote a blow-down limit. Then, (u_∞, v_∞) satisfies (1.8) for any $\varphi \in C_c^1(\mathbb{R}^d)$.*

Finally, we obtain the following partial result for solutions which need not be asymptotically homogeneous.

⁽¹⁾for the topology of pointwise convergence

Theorem 6. *Let $p \geq q \geq 1$ be such that $pq > 1$ and*

$$d < 2 + 2x_0, \quad (1.9)$$

where x_0 is the largest root of the polynomial

$$H(x) = x^4 - pq\alpha\beta(4x^2 - 2(\alpha + \beta)x + \alpha\beta). \quad (1.10)$$

If $(u, v) \in C^2(\mathbb{R}^d)$ is a nonnegative solution of (1.1) which is stable (resp. stable outside a compact set and (1.5) holds), then $u = v = 0$.

Remark 1. *The sharp condition (1.7) can be reformulated as (1.9) where this time x_0 must be interpreted as the largest root of the polynomial*

$$H_{JL}(x) = (x^2 - \gamma^2/4)^2 - pq\alpha\beta(4x^2 - 2(\alpha + \beta)x + \alpha\beta).$$

The above theorems build upon a wealth of previously known results, as we describe next. Thanks to the comparison inequality

$$\frac{v^{p+1}}{p+1} \leq \frac{u^{q+1}}{q+1} \quad (1.11)$$

established for $p \geq q$ and for bounded positive solutions by Souplet in [34], such solutions of (1.1) must satisfy $u = v$ in the case $p = q$, so that the system becomes a single equation and the condition (1.7) reduces to $p < p_c(d)$, where $p_c(d)$ is the Joseph-Lundgren stability exponent discovered in [23]. Crandall and Rabinowitz proved in [8] that, given a smoothly bounded domain $\Omega \subset \mathbb{R}^d$ and the nonlinearity $f(u) = (1 + u)^p$, the equation

$$\begin{cases} -\Delta u = \lambda f(u) & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

admits a curve of solutions $\lambda \in [0, \lambda^*) \mapsto u_\lambda \in C^2(\overline{\Omega})$, which are stable. They applied Moser's method [29] with a twist: substituting Sobolev's inequality by the variational formulation of the stability property in the iteration process, they proved uniform boundedness of the family $(u_\lambda)_{\lambda \in [0, \lambda^*)}$ in $L^\infty(\Omega)$ for $p < p_c(d)$. This implies the existence of a smooth stable solution u^* also for the extremal parameter $\lambda = \lambda^*$. Conversely, for $p \geq p_c(d)$ and $\Omega = B_1$, the extremal solution u^* still exists but is singular (and $u^* - 1$ is given by (1.3) for $p = q$).

For $p < p_S(d)$ where $p_S(d)$ is the Sobolev critical exponent and $f(u) = |u|^{p-1}u$, Bahri and Lions [1] proved that the L^∞ bound extends and is in fact equivalent to a bound on the Morse index of (possibly sign-changing) solutions. They argued by a blow-up argument and proved Theorem 4 for finite Morse index solutions of the subcritical equation. Farina extended their result to all $p < p_c(d)$, $p \neq p_S(d)$ and to solutions stable outside a compact set, observing that finite Morse index solutions belong to this class. In fact the two notions coincide, as proved by Devyver [11]. In other words, Theorem 2 generalizes Farina's result to the case of positive asymptotically homogeneous solutions of the system (1.1). We note that the restriction $p \neq p_S(d)$ is necessary in Farina's result, and so is the restriction (1.5) for the Lane-Emden system (1.1). Indeed, in the critical case, Lions proved in [25] that (1.1) has a ground state solution, which is unique up to scaling and translation. He observed that if (u, v) is a positive solution of (1.1), then

$$-\Delta((-\Delta v)^{\frac{1}{q}}) = v^p \quad \text{in } \mathbb{R}^d$$

and so a ground state solution can be sought by minimizing $\mathcal{E}(v) = \|\Delta v\|_{L^m(\mathbb{R}^d)}$, $m = \frac{q+1}{q}$, over the set of functions $v \in \mathcal{D}^{2,m}(\mathbb{R}^d)$ such that $\|v\|_{L^{p+1}(\mathbb{R}^d)} = 1$. Such a solution is positive, radial and stable outside a compact set by construction (or as follows from the asymptotics computed by Hulshof and Van der Vorst in [22]). As observed by Mtiri and Ye in [30], the second variation of the energy \mathcal{E} can be computed at positive stable solutions in the case $p > 1 \geq q$, leading to the following reformulation of stability:

$$\int_{\mathbb{R}^d} v^{p-1} \varphi^2 dx \leq \frac{1}{pq} \int_{\mathbb{R}^d} u^{1-q} |\Delta \varphi|^2 dx \quad \text{for all } \varphi \in C_c^2(\mathbb{R}^d). \quad (1.12)$$

Let us turn now to the well understood biharmonic case $q = 1$. In that case and for positive radial stable solutions, Theorem 4 follows from the works of Gazzola, Grunau [16], Guo, Wei [17] and Karageorgis [24]. Since the Moser iteration method is based on the chain rule, its adaptation to fourth order equations is nontrivial. Still, Wei and Ye [35] classified positive stable solutions when $d \leq 8$ thanks to the inequality (1.11) (a strategy inspired by earlier work of Cowan, Esposito and Ghoussoub [6]). Following the same Moser iteration strategy, Harrabi, Ye and the third named author classified positive stable solutions for $d \leq 12$ in [20]. They exploited the following interpolated version of the stability inequality:

$$\int_{\mathbb{R}^d} |u|^{\frac{q-1}{2}} |v|^{\frac{p-1}{2}} \varphi^2 dx \leq \frac{1}{\sqrt{pq}} \int_{\mathbb{R}^d} |\nabla \varphi|^2 dx \quad (1.13)$$

of which several proofs are available (based on Picone's identity in Cowan [5], Cowan and Ghoussoub [7] and Farina, Sirakov and the first author [12], interpolation theory in Goubet, Warnault and the first two named authors [13]). Finally, Theorem 4 was proved for the full range of exponents in the biharmonic case by Dávila, Wang, Wei and the first named author [9]. Their result crucially relies on a monotonicity formula which is unavailable for the system.

The complete system was studied by Cowan [5], who classified positive stable solutions for $d \leq 10$ and $p \geq q \geq 2$ thanks to (1.13). His result was improved by Harrabi, Mtiri and the third named author [19], who classified positive stable solutions for $d \leq 10$ and $p \geq q > 4/3$ and bounded positive stable solutions for $d \leq 6$ and $p \geq q > 1$. Also, Mtiri and Ye [30] completely classified positive solutions stable outside a compact set for (p, q) subcritical.

In order to prove Theorem 4, we propose a new strategy, where (near optimizers of) Hardy's inequality turn out to have a central role. Due to the lack of a monotonicity formula, any blow-down (resp. blow-up) limit of a given solution need not be homogeneous *a priori*. To bypass this difficulty, we assume that (1.8) holds. Now, rather than Moser iteration, we feed (1.8) in an iteration scheme which is closer to De Giorgi's original idea [10]. More precisely, we prove a reduction of the oscillation lemma for solutions satisfying (1.8). We capture the oscillations through the rescaled and renormalized Dirichlet energy

$$r^2 \int_{B_r} \left| \nabla \left(\frac{u}{u_s} \right)^a \right|^2 dx + r^2 \int_{B_r} \left| \nabla \left(\frac{v}{v_s} \right)^b \right|^2 dx \quad (1.14)$$

where $(a, b) = \left(\frac{q+1}{2}, \frac{p+1}{2} \right)$ and (u_s, v_s) is given by (1.3). Thanks to Campanato's characterization of Hölder spaces [3], this is enough to obtain the following universal Hölder estimate on local stable solutions of the system:

$$\|u/u_s\|_{C^\sigma(B_1)} + \|v/v_s\|_{C^\sigma(B_1)} \leq C \quad (1.15)$$

which holds for some $\sigma = \sigma(d, p, q) \in (0, 1)$ and $C = C(d, p, q) > 0$ under the assumption (1.8). Note that thanks to the translation invariance of the Lane-Emden system, the above estimate can be applied to any translation of a given solution and so Hölder regularity for the solution itself follows. Also, by a natural scaling argument, the classification follows for all solutions such that (1.8) holds in \mathbb{R}^d . In particular, blow-down limits are trivial under the assumptions of Theorem 4. This gives enough asymptotic information to conclude that solutions such that (1.8) holds only outside a compact set are also trivial. Estimate (1.15) bears resemblance with the work of Cabré, Figalli, Ros-Oton and Serra [2]. In their case, a version of Pohozaev's identity is central to the analysis. In our framework, we derive (1.15) through a much different road, where Hardy's inequality is used instead. It is also interesting to note that the scale invariance of the equation (resp. the nonlinear nature of the problem) appears explicitly in the definition of the Dirichlet energy (1.14) through the normalization by (u_s, v_s) (resp. through the exponents a, b reminiscent of Moser's iteration).

2. STABILITY REVISITED

In this section, we derive all the functional inequalities presented in the introduction. They are reformulations (or sometimes merely consequences of) stability and they are at the heart of the iteration

methods used in this paper. Recall that a solution $(u, v) \in C^2(\Omega)$ is said to be stable in an open set $\Omega \subset \mathbb{R}^d$ if there exist two positive functions $\phi, \psi \in C^2(\Omega)$ such that

$$-\Delta\phi \geq pv^{p-1}\psi, \quad -\Delta\psi \geq qu^{q-1}\phi \quad \text{in } \Omega. \quad (2.1)$$

In the next two lemmas, we give a variational reformulation of (2.1).

Lemma 1. *Let $\varphi \in C^2(\Omega)$ be a positive superharmonic function and $\eta \in C^2(\Omega)$. Then,*

$$\Delta \left(\frac{\eta^2}{\varphi} \right) \Delta\varphi \leq (\Delta\eta)^2.$$

Proof. Expand $\Delta \left(\frac{\eta^2}{\varphi} \right)$:

$$\begin{aligned} \Delta \left(\frac{\eta^2}{\varphi} \right) &= \frac{1}{\varphi} \Delta(\eta^2) + 2\nabla\eta^2 \cdot \nabla \frac{1}{\varphi} + \eta^2 \Delta \left(\frac{1}{\varphi} \right) \\ &= \frac{2}{\varphi} (\eta\Delta\eta + |\nabla\eta|^2) - 4\frac{\eta}{\varphi^2} \nabla\eta \cdot \nabla\varphi + \eta^2 \left(-\frac{1}{\varphi^2} \Delta\varphi + \frac{2}{\varphi^3} |\nabla\varphi|^2 \right) \\ &= \left[2\frac{\eta}{\varphi} \Delta\eta - \frac{\eta^2}{\varphi^2} \Delta\varphi \right] + \left[\frac{2}{\varphi} |\nabla\eta|^2 - 4\frac{\eta}{\varphi^2} \nabla\eta \cdot \nabla\varphi + 2\frac{\eta^2}{\varphi^3} |\nabla\varphi|^2 \right] \\ &= - \left[\frac{\eta^2}{\varphi^2} \Delta\varphi - 2\frac{\eta}{\varphi} \Delta\eta \right] + \frac{2}{\varphi} \left[|\nabla\eta|^2 - 2\frac{\eta}{\varphi} \nabla\eta \cdot \nabla\varphi + \frac{\eta^2}{\varphi^2} |\nabla\varphi|^2 \right]. \end{aligned}$$

Observe that the second bracket above is a perfect square. Multiply by $\Delta\varphi$ and complete the square in the first bracket to get

$$\Delta \left(\frac{\eta^2}{\varphi} \right) \Delta\varphi = \left[- \left(\frac{\eta}{\varphi} \Delta\varphi - \Delta\eta \right)^2 + (\Delta\eta)^2 \right] + 2\frac{\Delta\varphi}{\varphi} \left| \nabla\eta - \frac{\eta}{\varphi} \nabla\varphi \right|^2 \leq (\Delta\eta)^2, \quad (2.2)$$

since $\Delta\varphi \leq 0$. \square

Lemma 2. *Assume that $(u, v) \in C^2(\Omega)$ is positive and stable in Ω . Let $\Omega' \Subset \Omega$. Then, for every $\eta \in H_0^1(\Omega') \cap H^2(\Omega')$,*

$$pq \int v^{p-1} \eta^2 \leq \int u^{1-q} (\Delta\eta)^2$$

and

$$pq \int u^{q-1} \eta^2 \leq \int v^{1-p} (\Delta\eta)^2. \quad (2.3)$$

Proof. By symmetry, it suffices to prove the first inequality. We may also assume that $\eta \in C^2(\overline{\Omega'})$. Multiply the first inequality in (2.1) by $q\eta^2/\psi$ and integrate over Ω' to get

$$pq \int v^{p-1} \eta^2 \leq q \int (-\Delta\phi) \frac{\eta^2}{\psi} = q \int -\Delta \left(\frac{\eta^2}{\psi} \right) \phi \leq q \int_{[-\Delta(\frac{\eta^2}{\psi}) \geq 0]} -\Delta \left(\frac{\eta^2}{\psi} \right) \phi.$$

By the second inequality in (2.1), it follows that

$$pq \int v^{p-1} \eta^2 \leq \int_{[-\Delta(\frac{\eta^2}{\psi}) \geq 0]} u^{1-q} \Delta \left(\frac{\eta^2}{\psi} \right) \Delta\psi,$$

and the conclusion follows by Lemma 1. \square

In the next lemma, we characterize stability for homogeneous blow-down limits.

Lemma 3. *Let $(u, v) \in C^2(\mathbb{R}^d)$ be a solution of (1.1). Assume that $(u, v) \in C^2(\mathbb{R}^d)$ is positive, stable outside a compact set and asymptotically homogeneous. Let $(u_\infty, v_\infty)(r\theta) = (r^{-\alpha} f(\theta), r^{-\beta} g(\theta))$ denote a blow-down limit of u . Then, for every $\varphi \in C^2(S^{d-1})$,*

$$pq \int_{S^{d-1}} |g|^{p-1} \varphi^2 \leq \int_{S^{d-1}} f^{1-q} |\Delta_\theta \varphi - H\varphi|^2,$$

where H is given by (1.4).

Proof. Fix an open set $\Omega \Subset \mathbb{R}^d \setminus \{0\}$. Since (u, v) is stable outside a compact set K , its rescaling $(u_R, v_R)(x) = (R^\alpha u(Rx), R^\beta v(Rx))$ is stable outside K/R and so it is stable in an open neighborhood of $\bar{\Omega}$ for R large enough. Let

$$\lambda_1 = \inf \left\{ \int v_R^{1-p} (\Delta \eta)^2 - pq \int u_R^{q-1} \eta^2 : \eta \in H_0^1(\Omega) \cap H^2(\Omega) \text{ s.t. } \|\eta\|_{L^2(\Omega)} = 1 \right\}.$$

By Lemma 2, we have $\lambda_1 \geq 0$. Since (u_R, v_R) is bounded and bounded away from zero in the compact set $\bar{\Omega}$, λ_1 is achieved by some function ϕ_R weakly solving the linearized equation

$$-\Delta \phi_R = p v_R^{p-1} \psi_R, \quad -\Delta \psi_R = (q u_R^{q-1} + p^{-1} \lambda_1) \phi_R \quad \text{in } \Omega,$$

where ψ_R is defined through the first equation above. By standard elliptic regularity, $(\phi_R, \psi_R) \in C^2(\Omega)$ and the equation is satisfied in the classical sense. Note that the solution $\tilde{\phi}_R \in H_0^1(\Omega) \cap H^2(\Omega)$ to $-\Delta \tilde{\phi}_R = |\Delta \phi_R|$ also solves the minimization problem, so that we can assume $(\phi_R, \psi_R) > 0$. In fact, $\lambda_1 > 0$. If not, inequality (2.3) (applied to the solution (u_R, v_R)) would be an equality for the choice $\eta = \phi_R$. This would imply in turn that (2.2) is an equality for the choice $(\eta, \varphi) = (\phi_R, \phi)$, leading to $\phi_R = c\phi$ in $\bar{\Omega}$, for some $c \in \mathbb{R}^*$. The latter equality cannot hold on $\partial\Omega$. Hence, $\lambda_1 > 0$ and repeating the proof of (2.3), we obtain the slightly finer inequality

$$pq \int v_R^{p-1} \eta^2 \leq \int \left(u_R^{q-1} + \frac{\lambda_1}{pq} \right)^{-1} (\Delta \eta)^2, \quad \eta \in C_c^2(\Omega)$$

Thanks to Fatou's lemma and the dominated convergence theorem (applied to the family (u_R, v_R) in the limit $R \rightarrow +\infty$), we arrive at

$$pq \int v_\infty^{p-1} \eta^2 \leq \int \left(u_\infty^{q-1} + \frac{\lambda_1}{pq} \right)^{-1} (\Delta \eta)^2 \leq \int u_\infty^{1-q} (\Delta \eta)^2.$$

Since (u_∞, v_∞) is homogeneous, we deduce that for $\eta = h(r)\varphi(\theta)$, $h \in C_c^2(0, +\infty)$, $\varphi \in C^2(S^{d-1})$,

$$pq \left(\int_{\mathbb{R}_+} h^2 r^{-\beta(p-1)+(d-1)} dr \right) \left(\int_{S^{d-1}} |g|^{p-1} \varphi^2 \right) \leq \int_{\mathbb{R}_+ \times S^{d-1}} r^{\alpha(q-1)} f^{1-q} (\Delta \eta)^2. \quad (2.4)$$

We have

$$\Delta \eta = (\Delta h)\varphi + \frac{h}{r^2} \Delta_\theta \varphi$$

and so

$$(\Delta \eta)^2 = (\Delta h)^2 \varphi^2 + 2 \left(\frac{h \Delta h}{r^2} \right) \varphi \Delta_\theta \varphi + \frac{h^2}{r^4} (\Delta_\theta \varphi)^2.$$

Thus, (2.4) becomes

$$\begin{aligned} pq \left(\int_{\mathbb{R}_+} h^2 r^{-\beta(p-1)+d-1} dr \right) \left(\int_{S^{d-1}} |g|^{p-1} \varphi^2 \right) &\leq \left(\int_{\mathbb{R}_+} h^2 r^{d-5+\alpha(q-1)} dr \right) \left(\int_{S^{d-1}} f^{1-q} (\Delta_\theta \varphi)^2 \right) + \\ &\quad \left(\int_{\mathbb{R}_+} 2h \Delta h r^{d-3+\alpha(q-1)} dr \right) \left(\int_{S^{d-1}} f^{1-q} \varphi \Delta_\theta \varphi \right) + \\ &\quad \left(\int_{\mathbb{R}_+} (\Delta h)^2 r^{d-1+\alpha(q-1)} dr \right) \left(\int_{S^{d-1}} f^{1-q} \varphi^2 \right). \end{aligned}$$

Recalling the definitions of the scaling exponents (α, β) , the first integral on the left-hand side is equal to the first integral on the right-hand side in the above inequality. So,

$$pq \left(\int_{S^{d-1}} |g|^{p-1} \varphi^2 \right) \leq \left(\int_{S^{d-1}} f^{1-q} (\Delta_\theta \varphi)^2 \right) + A \left(\int_{S^{d-1}} f^{1-q} \varphi \Delta_\theta \varphi \right) + B \left(\int_{S^{d-1}} f^{1-q} \varphi^2 \right),$$

where

$$A = \frac{\int_{\mathbb{R}_+} 2h\Delta h r^{d-3+\alpha(q-1)} dr}{\int_{\mathbb{R}_+} h^2 r^{d-5+\alpha(q-1)} dr} \quad \text{and} \quad B = \frac{\int_{\mathbb{R}_+} (\Delta h)^2 r^{d-1+\alpha(q-1)} dr}{\int_{\mathbb{R}_+} h^2 r^{d-5+\alpha(q-1)} dr}.$$

Choosing $h(r) = r^{-\frac{d-2-\gamma}{2}} k_n(r) \in C_c^2(\mathbb{R}_+)$, where k_n is a standard cut-off function (vanishing near the origin and near infinity) such that $k_n(r) \rightarrow 1$ for every $r \in \mathbb{R}_+^*$, we find that $A \rightarrow -2H$ and $B \rightarrow H^2$ as $n \rightarrow +\infty$, where H is given by (1.4). Hence,

$$pq \left(\int_{S^{d-1}} |g|^{p-1} \varphi^2 \right) \leq \int_{S^{d-1}} f^{1-q} ((\Delta_\theta \varphi)^2 - 2H\varphi\Delta_\theta \varphi + H^2\varphi^2) = \int_{S^{d-1}} f^{1-q} (\Delta_\theta \varphi - H\varphi)^2.$$

□

In the last two lemmas of this section, we prove that stability implies the weaker but more handy inequality (1.8).

Lemma 4. *Assume that $(u, v) \in C^2(\mathbb{R}^d)$ is positive and stable outside a compact set. Assume also that (u, v) is asymptotically homogeneous and let $(u_\infty, v_\infty)(r\theta) = (r^{-\alpha} f(\theta), r^{-\beta} g(\theta))$ denote a blow-down limit. Then, for every $\eta \in C^1(S^{d-1})$,*

$$\sqrt{pq} \int_{S^{d-1}} f^{\frac{q-1}{2}} g^{\frac{p-1}{2}} \eta^2 \leq \int_{S^{d-1}} (|\nabla_\theta \eta|^2 + H\eta^2),$$

where H is given by (1.4).

Proof. Let

$$\lambda_1 = \inf \left\{ \int_{S^{d-1}} f^{1-q} |\Delta_\theta \psi - H\psi|^2 - pq \int_{S^{d-1}} |g|^{p-1} \psi^2 : \psi \in H^2(S^{d-1}) \text{ s.t. } \|\psi\|_{L^2(S^{d-1})} = 1 \right\}.$$

Then, $\lambda_1 \geq 0$ by Lemma 3. Since f is bounded away from zero and g is bounded, λ_1 is attained by some function $\psi \in H^2(S^{d-1})$ such that $\|\psi\|_{L^2(S^{d-1})} = 1$. Replacing ψ by $\tilde{\psi}$ the weak solution to

$$-\Delta_\theta \tilde{\psi} + H\tilde{\psi} = |\Delta_\theta \psi - H\psi| \quad \text{in } S^{d-1},$$

we may assume that $\psi \geq 0$ and $\varphi := q^{-1} f^{1-q} (-\Delta_\theta \psi + H\psi) \geq 0$. By the strong maximum principle, $(\varphi, \psi) > 0$ a.e. Also, (φ, ψ) is a weak solution to

$$-\Delta_\theta \varphi + H\varphi = (p|g|^{p-1} + \lambda_1)\psi, \quad -\Delta_\theta \psi + H\psi = q|f|^{q-1}\varphi \quad \text{in } S^{d-1}$$

and by elliptic regularity, $(\varphi, \psi) \in C^2(S^{d-1})$. Take $\eta \in C^1(S^{d-1})$, multiply the first equation above by η^2/φ and integrate by parts. Since $\lambda_1 \geq 0$,

$$p \int |g|^{p-1} \frac{\psi}{\varphi} \eta^2 \leq \int \nabla_\theta \varphi \cdot \nabla_\theta \left(\frac{\eta^2}{\varphi} \right) + H \int \eta^2 \leq \int (|\nabla_\theta \eta|^2 + H\eta^2).$$

Similarly,

$$q \int |f|^{q-1} \frac{\varphi}{\psi} \eta^2 \leq \int (|\nabla_\theta \eta|^2 + H\eta^2).$$

Multiplying the above two inequalities and applying the Cauchy-Schwarz inequality, the result follows. □

Now, we are ready to provide the proof of Theorem 5.

Proof of Theorem 5. Fix $\varphi \in C_c^1(\mathbb{R}^d)$ and $r > 0$. By the previous lemma,

$$\frac{1}{\sqrt{pq}} \int_{S^{d-1}} (|\nabla_\theta \varphi(r\theta)|^2 + H\varphi(r\theta)^2) d\sigma(\theta) \geq \int_{S^{d-1}} f^{\frac{q-1}{2}} g^{\frac{p-1}{2}} \varphi(r\theta)^2 d\sigma(\theta),$$

Multiply by r^{d-3} and integrate the above inequality in the r -variable. We get

$$\frac{1}{\sqrt{pq}} \int_{S^{d-1} \times \mathbb{R}_+} (|\nabla_\theta \varphi(r\theta)|^2 + H\varphi(r\theta)^2) r^{d-3} dr d\sigma(\theta) \geq \int_{\mathbb{R}^d} u_\infty^{\frac{q-1}{2}} v_\infty^{\frac{p-1}{2}} \varphi^2 dx.$$

In addition, by Hardy's inequality, for every $h \in C_c^1(\mathbb{R}_+)$,

$$\int_{\mathbb{R}_+} \left(\frac{\partial h}{\partial r} \right)^2 r^{d-1} dr \geq \left(\frac{d-2}{2} \right)^2 \int_{\mathbb{R}_+} \frac{h^2}{r^2} r^{d-1} dr.$$

Applying Hardy's inequality to $h(r) = \varphi(r\theta)$ for fixed $\theta \in S^{d-1}$, recalling that $|\nabla\varphi|^2 = \left(\frac{\partial\varphi}{\partial r} \right)^2 + \frac{1}{r^2} |\nabla_\theta\varphi|^2$ and summing the above two inequalities, (1.8) follows. \square

3. THE CASE $\alpha + \beta \geq d - 4$

In this section, we derive integral estimates on solutions thanks to their stability properties. These estimates are the central tool to prove Theorems 1 and 6. They also imply Theorem 3. To begin with, recall the following classical lemma, which holds for all positive solutions of (1.1), see [28, 33] and e.g. Lemma 2 in [19] for a proof.

Lemma 5. *Let $p \geq q \geq 1$ and let $(u, v) \in C^2(\mathbb{R}^d)$ be a positive solution of (1.1). Then, there exists a constant $C > 0$ depending on p, q, d only such that for any $R \geq 1$, there holds*

$$\int_{B_R} v^p dx \leq CR^{d-\beta p}, \quad \int_{B_R} u^q dx \leq CR^{d-\alpha q}.$$

We will also use the following comparison inequality, due to Souplet [34]. See e.g. Q. H. Phan [31] for a proof.

Lemma 6. *Let $p \geq q \geq 1$ be such that $pq > 1$. Then, any positive solution $(u, v) \in C^2(\mathbb{R}^d)$ of (1.1) verifies*

$$v^{p+1} \leq \frac{p+1}{q+1} u^{q+1} \quad \text{in } \mathbb{R}^d. \quad (3.1)$$

The next lemma pertains to stable solutions and uses ideas from [19, 20].

Lemma 7. *Assume that $(u, v) \in C^2(\Omega)$ is a positive stable solution of (1.1) in an open set $\Omega \subset \mathbb{R}^d$. Let $a > \frac{q+1}{4}$ and $b = \frac{p+1}{q+1}a$. Assume that $AB > 1$, where $A = \sqrt{pq} \frac{(2a-1)}{a^2}$ and $B = \sqrt{pq} \frac{(2b-1)}{b^2}$. Then, there exists $C > 0$ depending on p, q, d, a only such that*

$$\int_{\Omega} v^p u^{2a-1} \eta^2 dx \leq C \int_{\Omega} u^{2a} [|\nabla\eta|^2 + \Delta(\eta^2)] dx, \quad \forall \eta \in C_c^2(\Omega).$$

Proof. Take $\eta \in C_c^2(\Omega)$. Let (u, v) be a stable solution of (1.1) in Ω and $a > \frac{1}{2}$. Integrating by parts,

$$\begin{aligned} \int_{\Omega} |\nabla u^a|^2 \eta^2 dx &= a^2 \int_{\Omega} u^{2a-2} |\nabla u|^2 \eta^2 dx \\ &= \frac{a^2}{2a-1} \int_{\Omega} \eta^2 \nabla(u^{2a-1}) \nabla u dx \\ &= \frac{a^2}{2a-1} \int_{\Omega} u^{2a-1} v^p \eta^2 dx + \frac{a}{2(2a-1)} \int_{\Omega} u^{2a} \Delta(\eta^2) dx, \end{aligned} \quad (3.2)$$

and

$$2a \int_{\Omega} u^{2a-1} \eta \nabla u \nabla \eta dx = \frac{1}{2} \int_{\Omega} \nabla(u^{2a}) \nabla(\eta^2) dx = -\frac{1}{2} \int_{\Omega} u^{2a} \Delta(\eta^2) dx. \quad (3.3)$$

Take $\varphi = u^a \eta$ in inequality (1.13). Using (3.2)-(3.3), we obtain

$$\sqrt{pq} \int_{\Omega} u^{\frac{q-1}{2}} v^{\frac{p-1}{2}} u^{2a} \eta^2 dx \leq \int_{\Omega} |\nabla\varphi|^2 dx \leq \frac{a^2}{2a-1} \int_{\Omega} u^{2a-1} v^p \eta^2 dx + C \int_{\Omega} u^{2a} [|\nabla\eta|^2 + \Delta(\eta^2)] dx.$$

So,

$$\int_{\Omega} u^{\frac{q-1}{2}} v^{\frac{p-1}{2}} u^{2a} \eta^2 dx \leq A^{-1} \int_{\Omega} u^{2a-1} v^p \eta^2 dx + C \int_{\Omega} u^{2a} [|\nabla\eta|^2 + \Delta(\eta^2)] dx. \quad (3.4)$$

Similarly, applying inequality (1.13) with $\varphi = v^b \eta$, $b \geq 1$, we obtain

$$\int_{\Omega} u^{\frac{q-1}{2}} v^{\frac{p-1}{2}} v^{2b} \eta^2 dx \leq B^{-1} \int_{\Omega} u^q v^{2b-1} \eta^2 dx + C \int_{\Omega} v^{2b} [|\nabla \eta|^2 + \Delta(\eta^2)] dx. \quad (3.5)$$

Let I_1 denote the left-hand side of (3.4) (resp. I_2 that of (3.5)). Then,

$$\begin{aligned} A^{\frac{4a}{q+1}} I_1 + I_2 &= A^{\frac{4a}{q+1}} \int_{\Omega} u^{\frac{q-1}{2}} v^{\frac{p-1}{2}} u^{2a} \eta^2 dx + \int_{\Omega} u^{\frac{q-1}{2}} v^{\frac{p-1}{2}} v^{2b} \eta^2 dx \\ &\leq A^{\frac{4a}{q+1}-1} \int_{\Omega} u^{2a-1} v^p \eta^2 dx + B^{-1} \int_{\Omega} u^q v^{2b-1} \eta^2 dx \\ &\quad + C \int_{\Omega} (u^{2a} + v^{2b}) [|\nabla \eta|^2 + \Delta(\eta^2)] dx. \end{aligned} \quad (3.6)$$

Fix now

$$b = \frac{p+1}{q+1} a. \quad (3.7)$$

Assume that $a > \frac{q+1}{4}$ so that $m = \frac{q+1}{4a} = \frac{p+1}{4b} \in (0, 1)$. By Young's inequality, there holds

$$\begin{aligned} B^{-1} \int_{\Omega} u^q v^{2b-1} \eta^2 dx &= B^{-1} \int_{\Omega} u^{\frac{q-1}{2}} v^{\frac{p-1}{2}} u^{2am} v^{2b(1-m)} \eta^2 dx \\ &\leq m B^{-\frac{1}{m}} I_1 + (1-m) I_2 \end{aligned}$$

and similarly we have

$$A^{\frac{4a}{q+1}-1} \int_{\Omega} u^{2a-1} v^p \eta^2 dx \leq (1-m) A^{\frac{1}{m}} I_1 + m I_2.$$

Combining the above two estimates with (3.6), we derive that

$$A^{\frac{1}{m}} I_1 \leq \left(m B^{-\frac{1}{m}} + (1-m) A^{\frac{1}{m}} \right) I_1 + C \int_{\Omega} (u^{2a} + v^{2b}) [|\nabla \eta|^2 + \Delta(\eta^2)] dx,$$

hence

$$m \left((AB)^{\frac{1}{m}} - 1 \right) I_1 \leq C \int_{\Omega} (u^{2a} + v^{2b}) [|\nabla \eta|^2 + \Delta(\eta^2)] dx.$$

Thus, if $AB > 1$, there holds

$$\int_{\Omega} u^{\frac{q-1}{2}} v^{\frac{p-1}{2}} u^{2a} \eta^2 dx \leq C \int_{\Omega} (u^{2a} + v^{2b}) [|\nabla \eta|^2 + \Delta(\eta^2)] dx.$$

Using (3.7) and (3.1), there holds $u^{2a} \geq C v^{2b}$ and $u^{\frac{q-1}{2}} v^{\frac{p-1}{2}} u^{2a} \geq C u^{2a-1} v^p$. So if $AB > 1$ and $a > \frac{q+1}{4}$,

$$\int_{\Omega} v^p u^{2a-1} \eta^2 dx \leq C \int_{\Omega} u^{2a} [|\nabla \eta|^2 + \Delta(\eta^2)] dx,$$

the proof is completed. \square

Using the above lemma, we obtain the following estimate.

Lemma 8. *Assume that $(u, v) \in C^2(B_2)$ is a positive solution of (1.1) which is stable in B_2 . Let $\theta \in \left[1, \frac{d}{d-2}\right)$. Then for any $a > \frac{q+1}{4}$ satisfying $AB > 1$, there exists $C > 0$ depending on p, q, d, θ, a only such that*

$$\int_{B_1} u^{2a\theta} dx \leq C. \quad (3.8)$$

Proof of Lemma 8. Let (u, v) be a solution of (1.1) which is stable in B_2 and take $\theta \in \left[1, \frac{d}{d-2}\right)$. Take $\eta \in C_c^\infty(B_2)$. By L^1 elliptic regularity theory, there exists $C > 0$ such that

$$\left(\int_{B_2} u^{2a\theta} \eta^{2\theta} \right)^{\frac{1}{\theta}} \leq C \int_{B_2} |\Delta(u^{2a} \eta^2)|.$$

Expanding $\Delta(u^{2a}\eta^2)$ in the right-hand side of the above inequality, we find

$$|\Delta(u^{2a}\eta^2)| \leq (2a)u^{2a-1}v^p\eta^2 + 2a(2a-1)u^{2a-2}|\nabla u|^2\eta^2 + u^{2a}|\Delta(\eta^2)| + 4au^{2a-1}|\nabla u||\nabla\eta^2|.$$

Then using Young's inequality, there holds

$$\left(\int_{B_2} u^{2a\theta}\eta^{2\theta}\right)^{\frac{1}{\theta}} \leq C \left[\int_{B_2} u^{2a-1}v^p\eta^2 + \int_{B_2} u^{2a-2}|\nabla u|^2\eta^2 + \int_{B_2} u^{2a}(|\Delta(\eta^2)| + |\nabla\eta|^2) \right]. \quad (3.9)$$

Multiplying $-\Delta u = v^p$ by $u^{2a-1}\eta^2$, integrating by parts and applying Young's inequality, we get

$$\int_{B_2} u^{2a-2}|\nabla u|^2\eta^2 dx \leq \frac{1}{2a-1} \int_{B_2} u^{2a-1}v^p\eta^2 dx + C \int_{B_2} u^{2a}|\Delta(\eta^2)| dx. \quad (3.10)$$

Combining the last two estimates, we obtain

$$\left(\int_{B_2} u^{2a\theta}\eta^{2\theta}\right)^{\frac{1}{\theta}} \leq C \left[\int_{B_2} u^{2a-1}v^p\eta^2 + \int_{B_2} u^{2a}(|\Delta(\eta^2)| + |\nabla\eta|^2) \right].$$

By Lemma 7 with $\Omega = B_2$, we conclude that

$$\left(\int_{B_2} u^{2a\theta}\eta^{2\theta}\right)^{\frac{1}{\theta}} \leq C \int_{B_2} u^{2a}(|\Delta(\eta^2)| + |\nabla\eta|^2).$$

Take φ a cut-off function in $C_c^\infty(B_2)$ verifying $0 \leq \varphi \leq 1$, $\varphi = 1$ in B_1 . Letting $\eta = \varphi^m$, $m \geq 1$, we arrive at

$$\left(\int_{B_2} u^{2a\theta}\varphi^{2\theta m} dx\right)^{\frac{1}{\theta}} \leq C \int_{B_2} u^{2a}\varphi^{2m-2} dx. \quad (3.11)$$

Let

$$J_1 := \int_{B_2} u^{2a\theta}\varphi^{2\theta m} dx, \quad J_2 := \int_{B_2} u^{2a}\varphi^{2m-2} dx.$$

As $a > \frac{q+1}{4}$, we have $q < 2a < 2a\theta$. A direct calculation yields

$$2a = q\lambda + 2a\theta(1-\lambda) \quad \text{with } \lambda = \frac{2a(\theta-1)}{2a\theta-q} \in (0,1).$$

Take m large such that $m(1-\theta(1-\lambda)) = m\frac{(2a-1)(\theta-1)}{2a\theta-q} > 1$. By Hölder's inequality and (3.11), we get

$$\begin{aligned} J_2 &\leq J_1^{1-\lambda} \left(\int_{B_2} u^q \varphi^{\frac{2m(1-\theta(1-\lambda))-2}{\lambda}} dx \right)^\lambda \\ &\leq (CJ_2)^{\theta(1-\lambda)} \left(\int_{B_2} u^q dx \right)^\lambda \\ &\leq C' J_2^{\theta(1-\lambda)}, \end{aligned}$$

where in the last inequality we have applied Lemma 5. This implies, using again (3.11),

$$J_1 \leq C^\theta (J_2)^\theta \leq C.$$

Since $\varphi = 1$ in B_1 , (3.8) follows. \square

Thanks to the energy estimate (3.8), Theorem 3 easily follows.

Proof of Theorem 3. Indeed, applying (3.8) with $a = \frac{q+1}{2}$ and some $\theta \in \left(1, \frac{d}{d-2}\right)$, we deduce that the family of rescaled functions $(u_R)_{R \geq 1}$ is bounded in $L_{loc}^{\tilde{q}+1}(\mathbb{R}^d)$ for some $\tilde{q} > q$. By the comparison inequality (3.1), we deduce that so must be $(v_R)_{R \geq 1}$ in $L_{loc}^{\tilde{p}+1}(\mathbb{R}^d)$ for some $\tilde{p} > p$. Thanks to the equation (1.1), it also follows that $(u_R, v_R)_{R \geq 1}$ is bounded in $W_{loc}^{1,1}(\mathbb{R}^d)$. This fact and the $L_{loc}^{\tilde{q}+1}(\mathbb{R}^d) \times L_{loc}^{\tilde{p}+1}(\mathbb{R}^d)$ bound imply compactness of $(u_R, v_R)_{R \geq 1}$ in $L_{loc}^{q+1}(\mathbb{R}^d) \times L_{loc}^{p+1}(\mathbb{R}^d)$ as claimed. \square

The energy estimate (3.8) also implies the following Liouville type result for positive stable solutions of (1.1).

Theorem 7. *Suppose $p \geq q > 1$. Then, (1.1) has no stable classical solution if $d < 2 + 2x_0$, where x_0 is the largest root of the polynomial H given by (6).*

Proof. First, for $a > \frac{q+1}{4}$, noting $x = a\alpha$, we can easily check that $AB > 1 \Leftrightarrow H(x) < 0$. Furthermore, by [18, Lemma 3.1] (corresponds to $p = \theta$, $q = p$ and $H(x) = (\frac{\alpha}{2})^4 L(\frac{2}{\alpha}x)$ according to the notations in this paper), we have $H(a\alpha) = H(\frac{q+1}{2}\alpha) < 0$ and x_0 is the unique root of H in $(\frac{q+1}{2}\alpha, +\infty)$. Suppose that $d < 2 + 2x_0$. Then, there exists $\frac{q+1}{2} < a < \frac{x_0}{\alpha}$ such that $d < 2 + 2a\alpha$. Let $R > 1$. The function defined for $x \in \mathbb{R}^d$ by

$$(u_R(x), v_R(x)) = (R^\alpha u(Rx), R^\beta v(Rx))$$

is a positive stable solution of (1.1). Replacing u in (3.8) by u_R , we deduce that

$$\int_{B_R} u^{2a\theta} dx \leq CR^{d-2a\alpha\theta}, \quad (3.12)$$

for any $1 < \theta < \frac{d}{d-2}$. Now, as $d < 2 + 2a\alpha$, we can take θ close to $\frac{d}{d-2}$ such that $d < 2a\alpha\theta$ and then, letting $R \rightarrow \infty$ in (3.12), we deduce that $u = 0$. \square

Turning to solutions merely stable outside a compact set, we first make the following observation.

Remark 2. *Let (u, v) be a positive classical solution of (1.1) which is stable outside a compact set K of \mathbb{R}^d . Let $R_0 > 0$ so that $K \subset B_{R_0}$. Suppose that $d < 2 + 2x_0$, where x_0 is the largest root of the polynomial H given by (6). Let $\gamma < \frac{2x_0}{\alpha} \frac{d}{d-2}$.*

(1) *Using (3.12), we have*

$$\int_{B_R} u^\gamma dx \leq CR^{d-\gamma\alpha}, \quad \text{for any } R > 0 \text{ and } B_R \subset \mathbb{R}^d \setminus B_{R_0}. \quad (3.13)$$

(2) *In addition, all the computations in the proof of Lemma 8 hold true by considering a cut-off function with support in $\Omega = B_{2R} \setminus \overline{B_{R_0}}$, for $R < R_0 + 3$ (see for example [14]), to find*

$$\int_{R_0+2 < |x| < R} u^\gamma dx \leq C_1 + C_2 R^{d-\gamma\alpha}, \quad \text{for any } R > R_0 + 3. \quad (3.14)$$

Following a strategy used in [14], we deduce the following decay estimate.

Lemma 9. *Let $p \geq q \geq 1$ be such that $pq > 1$ and $d < 2 + 2x_0$, where x_0 is the largest root of the polynomial H given by (6). Suppose that $(u, v) \in C^2(\mathbb{R}^d)$ is a nonnegative solution of (1.1) which is stable outside a compact set and (1.5) holds. Then,*

$$u(x) = o(|x|^{-\alpha}), \quad v(x) = o(|x|^{-\beta}) \quad \text{as } |x| \rightarrow \infty.$$

Proof. Fix $\varepsilon > 0$. As $d < 2 + 2x_0$ and (1.5) holds, we have $2 < \frac{d}{\alpha} < \frac{2x_0}{\alpha} \frac{d}{d-2}$. So, Applying 2) of Remark 2 with $\gamma = \frac{d}{\alpha}$, there exists $R_1 > R_0$ such that

$$\left(\int_{|x| \geq R_1} u^{\frac{d}{\alpha}} \right)^{\frac{\alpha}{d}} < \varepsilon. \quad (3.15)$$

Consider $y \in \mathbb{R}^d$ such that $|y| > 2R_1$ and set $\rho = \frac{|y|}{4}$. With this choice, we have

$$B_{2\rho}(y) \subset \{x \in \mathbb{R}^d; |x| > R_1\} \subset \{x \in \mathbb{R}^d; |x| > R_0\}.$$

Remark that u is a solution of the linear equation

$$\Delta u + \ell(x)u = 0 \text{ in } B_{2\rho}(y),$$

where $\ell(x) = \frac{v^p}{u}$. By the comparison inequality (3.1), we have $\ell \leq C u^{\frac{pq-1}{p+1}} = C u^{\frac{2}{\alpha}}$. In addition, since $\frac{d}{\alpha} < \frac{2x_0}{\alpha} \frac{d}{d-2}$, there exists δ sufficiently small such that $\frac{2d}{(2-\delta)\alpha} < \frac{2x_0}{\alpha} \frac{d}{d-2}$. So, by Remark 2, we derive that $\ell \in L^{\frac{d}{2-\delta}}(B_{2\rho}(y))$. Therefore, according to [33, Theorem 1], we obtain

$$\|u\|_{L^\infty(B_\rho(y))} \leq C_S \rho^{-\frac{d}{2}} \|u\|_{L^2(B_{2\rho}(y))}, \quad (3.16)$$

where C_S is a constant depending only on p, d, R_0 . $\rho^\delta \|\ell\|_{L^{\frac{d}{2-\delta}}(B_{2\rho}(y))}$. Now, we have

$$\rho^\delta \|\ell\|_{L^{\frac{d}{2-\delta}}(B_{2\rho}(y))} \leq C \rho^\delta \left(\int_{B_{2\rho}(y)} u^{\frac{2d}{(2-\delta)\alpha}} dx \right)^{\frac{2-\delta}{d}}.$$

Applying Remark 2 with $\gamma = \frac{2d}{(2-\delta)\alpha}$, it follows that

$$\rho^\delta \|\ell\|_{L^{\frac{d}{2-\delta}}(B_{2\rho}(y))} \leq C \rho^\delta \rho^{(d-\gamma\alpha)\frac{2-\delta}{d}} = C. \quad (3.17)$$

Hence, the constant C_S depends only on p, d and R_0 . Using (3.16) and applying Hölder's inequality, we obtain

$$\begin{aligned} \|u\|_{L^\infty(B_\rho(y))} &\leq C_S \rho^{-\frac{d}{2}} \|u\|_{L^2(B_{2\rho}(y))} \\ &\leq C_S C \rho^{-\frac{d}{2}} \rho^{d(\frac{1}{2}-\frac{\alpha}{d})} \|u\|_{L^{\frac{d}{\alpha}}(B_{2\rho}(y))} \\ &= C_1 \rho^{-\alpha} \|u\|_{L^{\frac{d}{\alpha}}(B_{2\rho}(y))}. \end{aligned}$$

Now, using (3.15), we get

$$\begin{aligned} |u(y)| &\leq \|u\|_{L^\infty(B_\rho(y))} \\ &\leq C_1 \rho^{-\alpha} \varepsilon \\ &= C_2 |y|^{-\alpha} \varepsilon. \end{aligned}$$

In summary, we have proven that for any $\varepsilon > 0$ there exists $M = 2R_1$ such that

$$y \in \mathbb{R}^d, |y| > M \Rightarrow |y|^\alpha |u(y)| \leq C_2 \varepsilon.$$

Therefore, $u(x) = o(|x|^{-\alpha})$ as $|x| \rightarrow \infty$ and by the comparison inequality (3.1), we deduce that $v(x) = o(|x|^{-\beta})$, which completes the proof. \square

4. FAST DECAY SOLUTIONS

In order to go further, we proceed by proving that the solution has in fact faster decay rate. As we shall see, this will be enough to conclude (in the range of parameters studied in the previous section), thanks to Pohozaev's identity.

Theorem 8. *Let $d \geq 3$, $p \geq q > 1$ supercritical i.e. $\frac{1}{p+1} + \frac{1}{q+1} < 1 - \frac{2}{d}$. Then, the system (1.1) has no classical positive solution (u, v) satisfying*

$$u(x) = o(|x|^{-\alpha}), \quad v(x) = o(|x|^{-\beta}), \quad \text{as } |x| \rightarrow \infty. \quad (4.1)$$

To prove Theorem 8, we first establish the following decay estimates.

Lemma 10. *Let $d \geq 3$, $p \geq q > 1$ and (u, v) a classical positive solution to the system (1.1) such that*

$$u(x) = o(|x|^{-\alpha}), \quad \text{as } |x| \rightarrow \infty. \quad (4.2)$$

Then, for any small $\epsilon > 0$, there exists two positive constants c and C such that

$$c|x|^{2-d} \leq u(x) \leq C|x|^{2-d+\epsilon} \quad \text{and} \quad |\nabla u(x)| \leq C|x|^{1-d+\epsilon} \quad \text{for } |x| \geq 1. \quad (4.3)$$

In addition,

$$|\nabla v(x)| = o(|x|^{-\beta-1}) \quad \text{as } |x| \rightarrow \infty. \quad (4.4)$$

Proof. Since u is a non-trivial, positive and superharmonic function, the first inequality in (4.3) is a standard comparison result. For the proof see e.g. [33, Lemma 2.1].

Now, by Lemma 6, we have

$$-\Delta u = v^p \leq \left(\frac{p+1}{q+1}\right)^{\frac{1}{p+1}} u^{\frac{2}{\alpha}} u.$$

Let $0 < \epsilon < \frac{1}{2}$. By the decay estimate (4.2), there exists $R_\epsilon > 1$ such that $\left(\frac{p+1}{q+1}\right)^{\frac{1}{p+1}} u^{\frac{2}{\alpha}} \leq \frac{\epsilon^2}{|x|^2}$, for any $|x| > R_\epsilon$. Therefore,

$$-\Delta u - \frac{\epsilon^2}{|x|^2} u \leq 0, \text{ for any } |x| > R_\epsilon.$$

Now, we can easily check that the functions $f_a := |x|^{-a}$, $f_b := |x|^{-b}$ are solutions of

$$-\Delta f - \frac{\epsilon^2}{|x|^2} f = 0 \quad \text{in } \mathbb{R}^d \setminus \{0\}, \quad (4.5)$$

where $a = \frac{d-2}{2} + \sqrt{(\frac{d-2}{2})^2 - \epsilon^2}$ and $b = \frac{d-2}{2} - \sqrt{(\frac{d-2}{2})^2 - \epsilon^2}$. Let $R > R_\epsilon$ and consider $w(x) = R_\epsilon^\alpha \|u\|_\infty f_a(x) + \epsilon^2 R^{b-\alpha} f_b(x)$. Then, we have

$$\begin{cases} \left(-\Delta - \frac{\epsilon^2}{|x|^2}\right)(u - w) \leq 0, & \text{in } A_{R_\epsilon}^R := B_R \setminus \overline{B_{R_\epsilon}}, \\ u - w \leq 0, & \text{on } \partial A_{R_\epsilon}^R. \end{cases} \quad (4.6)$$

By the maximum principle, we deduce that $u \leq w$ in $A_{R_\epsilon}^R$ for any $R > A_{R_\epsilon}^R$. Choosing $\epsilon > 0$ so small that $b - \alpha < 0$ and letting $R \rightarrow \infty$, we conclude that for $|x| \geq R_\epsilon$,

$$u(x) \leq C|x|^{-a} \leq C|x|^{2-d+\epsilon}.$$

To obtain the gradient estimates, we scale and apply standard elliptic regularity. Precisely, given $z \in \mathbb{R}^d \setminus \{0\}$ and $\rho = |z|/2$, let $(\tilde{u}(x), \tilde{v}(x)) = (\rho^\alpha u(z + \rho x), \rho^\beta v(z + \rho x))$ for $x \in \mathbb{R}^d$. Then, (\tilde{u}, \tilde{v}) still solves the Lane-Emden system and by standard elliptic regularity

$$|\nabla \tilde{u}(0)| \leq C (\|\tilde{v}^p\|_{L^\infty(B_1)} + \|\tilde{u}\|_{L^\infty(B_1)}) \leq C' \|\tilde{u}\|_{L^\infty(B_1)},$$

where the last inequality follows from the comparison inequality Lemma 6 and the fact that $pq \geq 1$. Hence, for $|z| \geq 1$,

$$|\nabla u(z)| \leq C|z|^{1-d+\epsilon},$$

as claimed. The gradient estimate on v follows similarly. \square

The last crucial ingredient in the proof of Theorem 8 is the following identity of Pohozaev-type (see [26, 32, 34]).

Lemma 11. *Let $p, q > 0$ and (u, v) a positive classical solution of (1.1). Then, for any $R > 0$, there holds*

$$\begin{aligned} & \left(\frac{d}{p+1} - a_1\right) \int_{B_R} v^{p+1} + \left(\frac{d}{q+1} - a_2\right) \int_{B_R} u^{q+1} \\ &= R^d \int_{S^{d-1}} \left(\frac{u^{q+1}}{q+1} + \frac{v^{p+1}}{p+1}\right) (R\theta) d\sigma(\theta) + R^{d-1} \int_{S^{d-1}} (a_1 u \partial_r v + a_2 v \partial_r u) (R\theta) d\sigma(\theta) \\ & \quad + R^d \int_{S^{d-1}} \left(\partial_r u \partial_r v - \frac{\nabla_\theta u \cdot \nabla_\theta v}{R^2}\right) (R\theta) d\sigma(\theta), \end{aligned} \quad (4.7)$$

where $a_2, a_2 \in \mathbb{R}$ satisfy $a_1 + a_2 = d - 2$.

Proof of Theorem 8. We argue by contradiction and suppose that there exists a classical positive solution (u, v) to the system (1.1) which decays faster than the homogeneous solution. Thanks to (4.3), since $q+1 > \frac{d}{d-2}$, $u \in L^{q+1}(\mathbb{R}^d)$. By the comparison inequality (3.1), $v \in L^{p+1}(\mathbb{R}^d)$. We deduce that

the left-hand side of (4.7) converges as $R \rightarrow +\infty$. Thanks to (4.3) and (4.4), we also deduce that all integrals on the right-hand side of (4.7) converge to 0 as $R \rightarrow +\infty$. Hence,

$$\left(\frac{d}{p+1} - a_1\right) \int_{\mathbb{R}^d} v^{p+1} dx + \left(\frac{d}{q+1} - a_2\right) \int_{\mathbb{R}^d} u^{q+1} dx = 0,$$

for any a_1, a_2 satisfying $a_1 + a_2 = d - 2$. Take $a_1 = \frac{d}{p+1}$, it follows that

$$\left(\frac{d}{q+1} - a_2\right) \int_{\mathbb{R}^d} u^{q+1} dx = 0.$$

Since $\frac{1}{p+1} + \frac{1}{q+1} < \frac{d-2}{2}$, we have

$$\left(\frac{d}{q+1} - a_2\right) = \frac{d}{q+1} + \frac{d}{q+1} - (d-2) < 0.$$

Therefore, $u \equiv 0$ in \mathbb{R}^d , a contradiction. \square

From the above two sections, we conclude the proof of Theorem 6. Moreover, as shown in page 277 of [19], we have $x_0 > 4, \forall p \geq q > 1$. Hence, Theorem 1 is a consequence of Theorem 6 and Theorem 7.

5. STABLE HOMOGENEOUS SOLUTIONS

This section is not needed for the proofs of the theorems stated in the introduction. It is written in order to classify homogeneous solutions and to prepare the reader for the more delicate iteration method of the next section. Let us restrict to the class of stable homogeneous solutions to the Lane-Emden system (1.1), in the supercritical case (1.2). Let $u(r\theta) = r^{-\alpha}f(\theta)$, $v(r\theta) = r^{-\beta}g(\theta)$, $r > 0, \theta \in S^{d-1}$ be a stable homogeneous solution of (1.1). Then (f, g) satisfies

$$-\Delta_\theta f = -\lambda f + |g|^{p-1}g, \quad -\Delta_\theta g = -\mu g + |f|^{q-1}f \quad \text{in } S^{d-1}. \quad (5.1)$$

Moreover, from Lemma 4, the stability of the homogeneous solution (u, v) , it follows that

$$\int_{S^{d-1}} |\nabla_\theta \phi|^2 d\sigma + H \int_{S^{d-1}} \phi^2 d\sigma \geq \sqrt{pq} \int_{S^{d-1}} |f|^{\frac{q-1}{2}} |g|^{\frac{p-1}{2}} \phi^2 d\sigma, \quad \forall \phi \in C^\infty(S^{d-1}). \quad (5.2)$$

We prove the following theorem.

Theorem 9. *Let $p \geq q \geq 1$ such that $pq > 1$. There exists a stable homogeneous solution to (1.1) if and only if (1.7) holds.*

Proof of Theorem 9. By the discussion in the introduction, it suffices to prove the "only if" part of Theorem 9. For this, we proceed by contradiction. Suppose that $H^2 < pq\lambda\mu$ and $(u, v) = (r^{-\alpha}f(\theta), r^{-\beta}g(\theta))$ is a stable homogeneous solution of (1.1). As in Section 3, we denote $A = \sqrt{pq} \frac{2a-1}{a^2}$ and $B = \sqrt{pq} \frac{2b-1}{b^2}$, with $b = \frac{p+1}{q+1}a$. Multiplying the equations of the system (5.1) by $\sqrt{pq}|f|^{2a-2}f$ and $\sqrt{pq}|g|^{2b-2}g$ respectively and integrating over S^{d-1} , there holds

$$A \int_{S^{d-1}} |\nabla_\theta(|f|^{a-1}f)|^2 d\sigma + \sqrt{pq}\lambda \int_{S^{d-1}} |f|^{2a} d\sigma = \sqrt{pq} \int_{S^{d-1}} |g|^{p-1}|f|^{2a-2}f d\sigma, \quad (5.3)$$

and

$$B \int_{S^{d-1}} |\nabla_\theta(|g|^{b-1}g)|^2 d\sigma + \sqrt{pq}\mu \int_{S^{d-1}} |g|^{2b} d\sigma = \sqrt{pq} \int_{S^{d-1}} |f|^{q-1}|g|^{2b-2}g d\sigma. \quad (5.4)$$

Take $a > \frac{q+1}{4}, b = \frac{p+1}{q+1}a > \frac{p+1}{4}$ and let $m := \frac{q+1}{4a} = \frac{p+1}{4b} \in (0, 1)$. By Hölder's inequality, we get

$$\sqrt{pq} \int_{S^{d-1}} |g|^{p-1}|f|^{2a-2}f d\sigma \leq \sqrt{pq} \int_{S^{d-1}} |f|^{\frac{q-1}{2}} |g|^{\frac{p-1}{2}} |f|^{2a(1-m)} |g|^{\frac{p-1}{2}} \leq \sqrt{pq} I_1^{1-m} I_2^m \quad (5.5)$$

and

$$\sqrt{pq} \int_{S^{d-1}} |f|^{q-1}|g|^{2b-2}g d\sigma \leq \sqrt{pq} \int_{S^{d-1}} |f|^{\frac{q-1}{2}} |g|^{\frac{p-1}{2}} g^{2b(1-m)} |f|^{\frac{q-1}{2}} \leq \sqrt{pq} I_1^m I_2^{1-m}, \quad (5.6)$$

where $I_1 := \int_{S^{d-1}} |f|^{\frac{q-1}{2}+2a} |g|^{\frac{p-1}{2}} d\sigma$ and $I_2 := \int_{S^{d-1}} |f|^{\frac{q-1}{2}} |g|^{\frac{p-1}{2}+2b} d\sigma$. Combining (5.3)–(5.6), we obtain

$$\begin{aligned} & \left(A \int_{S^{d-1}} |\nabla_\theta(|f|^{a-1}f)|^2 d\sigma + \sqrt{pq}\lambda \int_{S^{d-1}} |f|^{2a} d\sigma \right) \left(B \int_{S^{d-1}} |\nabla_\theta(|g|^{b-1}g)|^2 d\sigma + \sqrt{pq}\mu \int_{S^{d-1}} |g|^{2b} d\sigma \right) \\ & \leq pqI_1I_2. \end{aligned}$$

On the other hand, testing $|f|^{a-1}f$ and $|g|^{b-1}g$ in the stability inequality (5.2), there holds

$$\sqrt{pq}I_1 \leq \int_{S^{d-1}} |\nabla_\theta(|f|^{a-1}f)|^2 d\sigma + H \int_{S^{d-1}} |f|^{2a} d\sigma, \quad \sqrt{pq}I_2 \leq \int_{S^{d-1}} |\nabla_\theta(|g|^{b-1}g)|^2 d\sigma + H \int_{S^{d-1}} |g|^{2b} d\sigma.$$

From the last two inequalities, it follows that

$$\begin{aligned} & \left(A \int_{S^{d-1}} |\nabla_\theta(|f|^{a-1}f)|^2 d\sigma + \sqrt{pq}\lambda \int_{S^{d-1}} |f|^{2a} d\sigma \right) \left(B \int_{S^{d-1}} |\nabla_\theta(|g|^{b-1}g)|^2 d\sigma + \sqrt{pq}\mu \int_{S^{d-1}} |g|^{2b} d\sigma \right) \\ & \leq \left(\int_{S^{d-1}} |\nabla_\theta(|f|^{a-1}f)|^2 d\sigma + H \int_{S^{d-1}} |f|^{2a} d\sigma \right) \left(\int_{S^{d-1}} |\nabla_\theta(|g|^{b-1}g)|^2 d\sigma + H \int_{S^{d-1}} |g|^{2b} d\sigma \right), \end{aligned}$$

or equivalently,

$$\begin{aligned} & (AB - 1) \int_{S^{d-1}} |\nabla_\theta(|f|^{a-1}f)|^2 d\sigma \int_{S^{d-1}} |\nabla_\theta(|g|^{b-1}g)|^2 d\sigma \\ & + (\sqrt{pq}A\mu - H) \int_{S^{d-1}} |\nabla_\theta(|f|^{a-1}f)|^2 d\sigma \int_{S^{d-1}} |g|^{2b} d\sigma \\ & + (\sqrt{pq}B\lambda - H) \int_{S^{d-1}} |\nabla_\theta(|g|^{b-1}g)|^2 d\sigma \int_{S^{d-1}} |f|^{2a} d\sigma \\ & + (pq\lambda\mu - H^2) \int_{S^{d-1}} |f|^{2a} d\sigma \int_{S^{d-1}} |g|^{2b} d\sigma \leq 0. \end{aligned} \tag{5.7}$$

Choose $a = \frac{q+1}{2}$ and so $b = \frac{p+1}{2}$. Since $p \geq q > 1$, we have

$$AB - 1 = \frac{1}{(q+1)^2(p+1)^2} (16p^2q^2 - (q+1)^2(p+1)^2) > 0.$$

Hence, applying Theorem 7, it follows that $f = g = 0$ if $d - 4 \leq \alpha + \beta$. So, we can assume that $d - 4 > \alpha + \beta$. In addition, since $p \geq q > 1$, we have $\frac{\beta}{\alpha}A^2 = \frac{16pq^3}{(p+1)(q+1)^3} = \frac{2p}{p+1} \left(\frac{2q}{q+1} \right)^3 \geq 1$. Recall that $\alpha \geq \beta$, so that $\frac{\lambda}{\mu} = \frac{\alpha}{\beta} \frac{d-2-\alpha}{d-2-\beta} \leq \frac{\alpha}{\beta} \leq A^2$. Since $H^2 < pq\lambda\mu$, we find $H < \sqrt{pq}\sqrt{\lambda}\sqrt{\mu} \leq \sqrt{pq}A\mu$. Using $d - 4 > \alpha + \beta$, we can easily check that $A\mu \leq B\lambda$ and so $H < \sqrt{pq}A\mu \leq \sqrt{pq}B\lambda$. We have just proved that all the constants appearing in (5.7) are positive, whence $f = g = 0$. So, we are done. \square

6. THE CASE $d - 4 > \alpha + \beta$: AN ITERATION METHOD IN THE SPIRIT OF DE GIORGI

Lemma 12. *Let (u, v) be a positive solution of (1.1) satisfying (1.8) in B_2 . Assume that $d - 4 > \alpha + \beta$ and (1.7) holds. Then, there exist constants $C > 0$ and $\sigma \in (0, 1)$ depending on d, p, q only such that letting $U = u/u_s$, $V = v/v_s$, $a = (q + 1)/2$ and $b = (p + 1)/2$, there holds*

$$\| |U|^{a-1}U \|_{C^\sigma(B_1)} + \| |V|^{b-1}V \|_{C^\sigma(B_1)} \leq C. \tag{6.1}$$

Proof. We divide our proof into three steps.

Step 1: Basic identities. Let $a \geq 1$, $\varphi \in C_c^1(B_2)$, multiply the first equation of the system by $|u|^{2a-2}u\varphi^2$ and integrate. The left-hand side is equal to

$$\begin{aligned} \int \nabla u \cdot \nabla(|u|^{2a-2}u\varphi^2) &= (2a - 1) \int |u|^{2a-2} |\nabla u|^2 \varphi^2 + \int |u|^{2a-2} u \nabla u \nabla \varphi^2 \\ &= \frac{2a - 1}{a^2} \int |\nabla(|u|^{a-1}u)|^2 \varphi^2 + \frac{1}{2a} \int \nabla |u|^{2a} \nabla \varphi^2 \end{aligned}$$

$$\begin{aligned}
&= \frac{2a-1}{a^2} \int |\nabla(|u|^{a-1}u\varphi)|^2 - \frac{2a-1}{a^2} \int \nabla\varphi \nabla(|u|^{2a}\varphi) + \frac{1}{2a} \int \nabla|u|^{2a} \nabla\varphi^2 \\
&= \frac{2a-1}{a^2} \int |\nabla(|u|^{a-1}u\varphi)|^2 + \int |u|^{2a} \left(\frac{2a-1}{a^2} \varphi \Delta\varphi - \frac{1}{2a} \Delta\varphi^2 \right).
\end{aligned}$$

Next, choose $\varphi = \varphi_0\psi$, where $\varphi_0(x) = |x|^{\alpha a - \frac{d-2}{2}}$ and $\psi \in C_c^1(\mathbb{R}^d \setminus \{0\})$. Then,

$$\Delta\varphi = (\Delta\varphi_0)\psi + 2\nabla\varphi_0 \cdot \nabla\psi + \varphi_0\Delta\psi,$$

and

$$\Delta\varphi^2 = (\Delta\varphi_0^2)\psi^2 + 2\nabla\varphi_0^2 \cdot \nabla\psi^2 + \varphi_0^2\Delta\psi^2.$$

Let R (for rest) be defined by

$$\frac{2a-1}{a^2} \varphi \Delta\varphi - \frac{1}{2a} \Delta\varphi^2 = \left(\frac{2a-1}{a^2} \varphi_0 \Delta\varphi_0 - \frac{1}{2a} \Delta\varphi_0^2 \right) \psi^2 - R,$$

where

$$R = \frac{1}{2a^2} \nabla\varphi_0^2 \cdot \nabla\psi^2 + \left(\frac{1}{a} |\nabla\psi|^2 - \frac{a-1}{a^2} \psi \Delta\psi \right) \varphi_0^2. \quad (6.2)$$

By direct computation, we have

$$\begin{aligned}
\frac{2a-1}{a^2} \varphi_0 \Delta\varphi_0 - \frac{1}{2a} \Delta\varphi_0^2 &= \left[\frac{2a-1}{a^2} \left((\alpha a)^2 - \left(\frac{d-2}{2} \right)^2 \right) - \frac{1}{2a} (2\alpha a(2\alpha a - (d-2))) \right] \frac{\varphi_0^2}{|x|^2} \\
&= \left[(2a-1)\alpha^2 - \frac{2a-1}{a^2} \left(\frac{d-2}{2} \right)^2 - 2\alpha^2 a + \alpha(d-2) \right] \frac{\varphi_0^2}{|x|^2} \\
&= \left[\lambda - \frac{2a-1}{a^2} \left(\frac{d-2}{2} \right)^2 \right] \frac{\varphi_0^2}{|x|^2}.
\end{aligned}$$

So, we just proved that for $\varphi = \varphi_0\psi$,

$$\int \nabla u \cdot \nabla(|u|^{2a-2}u\varphi^2) = \frac{2a-1}{a^2} \int |\nabla(|u|^{a-1}u\varphi)|^2 + \left(\lambda - \frac{2a-1}{a^2} \left(\frac{d-2}{2} \right)^2 \right) \int \frac{(|u|^a\varphi)^2}{|x|^2} - \int R|u|^{2a}.$$

The right hand-side is equal to

$$\int v^{p-1}v|u|^{2a-2}u\varphi^2.$$

Introducing the quadratic forms

$$Q_1(f) = \int \left(|\nabla f|^2 - \left(\frac{d-2}{2} \right)^2 \frac{f^2}{|x|^2} \right), \quad Q_2(f) = \int \frac{f^2}{|x|^2}$$

and the numbers $Q_1 = Q_1(|u|^{a-1}u\varphi)$, $Q_2 = Q_2(|u|^{a-1}u\varphi)$, $S = \int v^{p-1}v|u|^{2a-2}u\varphi^2$, $T = \int R|u|^{2a}$, we conclude that

$$\frac{2a-1}{a^2} Q_1 + \lambda Q_2 = S + T. \quad (6.3)$$

The choice $a = \frac{q+1}{2}$ makes the expression of $\varphi_0(x) = |x|^{\frac{(p+1)(q+1)}{pq-1} - \frac{d-2}{2}}$ symmetric in the variables (p, q) . Hence, choosing $b = \frac{p+1}{2}$, we may assert that for the *same* functions $\varphi = \varphi_0\psi$ and R , there holds

$$\frac{2b-1}{b^2} Q'_1 + \mu Q'_2 = S' + T', \quad (6.4)$$

where $Q'_1 = Q(|v|^{b-1}v\varphi)$, $Q'_2 = Q_2(|v|^{b-1}v\varphi)$, $S' = \int |u|^{q-1}u|v|^{2b-2}v\varphi^2$ and $T' = \int R|v|^{2b}$. For this choice of a and b we have $S = S'$.

Step 2: $Q_1 + Q'_1 \leq |T| + |T'|$ and $Q_2 + Q'_2 \leq |T| + |T'|$.

Multiplying (6.3) and (6.4) we find

$$\left(\frac{4q}{(q+1)^2}Q_1 + \lambda Q_2\right) \left(\frac{4p}{(p+1)^2}Q'_1 + \mu Q'_2\right) \leq (S+T)(S+T') = S^2 + S(T+T') + TT'.$$

If $\varepsilon > 0$ is small enough, the above inequality yields

$$\left(\frac{4q}{(q+1)^2}Q_1 + \lambda Q_2\right) \left(\frac{4p}{(p+1)^2}Q'_1 + \mu Q'_2\right) \leq (1+\varepsilon)S^2 + C(\varepsilon)(T^2 + T'^2). \quad (6.5)$$

Now, recalling the choice $a = \frac{q+1}{2}$, $b = \frac{p+1}{2}$, the Cauchy-Schwarz inequality yields

$$\left|\int |u|^{q-1}u|v|^{2b-2}v\varphi^2\right| \leq \int |u|^q|v|^p\varphi^2 \leq \left(\int |u|^{\frac{q-1}{2}}|v|^{\frac{3p+1}{2}}\varphi^2\right)^{\frac{1}{2}} \left(\int |u|^{\frac{3q+1}{2}}|v|^{\frac{p-1}{2}}\varphi^2\right)^{\frac{1}{2}}.$$

Applying (1.8) with test functions $|v|^{b-1}v\varphi$ and $|u|^{a-1}u\varphi$, we deduce that

$$\left|\int |u|^{q-1}u|v|^{2b-2}v\varphi^2\right| \leq \frac{1}{\sqrt{pq}} \left(\int |\nabla(|v|^{b-1}v\varphi)|^2 - \frac{\gamma^2}{4} \int \frac{(|v|^b\varphi)^2}{|x|^2}\right)^{1/2} \left(\int |\nabla(|u|^{a-1}u\varphi)|^2 - \frac{\gamma^2}{4} \int \frac{(|u|^a\varphi)^2}{|x|^2}\right)^{1/2}.$$

Similarly,

$$\left|\int v^{p-1}v|u|^{2a-2}u\varphi^2\right| \leq \frac{1}{\sqrt{pq}} \left(\int |\nabla(|v|^{b-1}v\varphi)|^2 - \frac{\gamma^2}{4} \int \frac{(|v|^b\varphi)^2}{|x|^2}\right)^{1/2} \left(\int |\nabla(|u|^{a-1}u\varphi)|^2 - \frac{\gamma^2}{4} \int \frac{(|u|^a\varphi)^2}{|x|^2}\right)^{1/2}.$$

Multiplying both inequalities, we find

$$S^2 \leq \frac{1}{pq}(Q_1 + HQ_2)(Q'_1 + HQ'_2). \quad (6.6)$$

We plug this last estimate into (6.5) to find

$$(AB - (1+\varepsilon))Q_1Q'_1 + (pq\lambda\mu - (1+\varepsilon)H^2)Q_2Q'_2 + (\sqrt{pq}A\mu - (1+\varepsilon)H)Q_1Q'_2 \\ + (\sqrt{pq}B\lambda - (1+\varepsilon)H)Q'_1Q_2 \leq C(T^2 + T'^2),$$

where, as above, we denote $A := \sqrt{pq}\frac{2a-1}{a^2} = \sqrt{pq}\frac{4q}{(q+1)^2}$ and $B := \sqrt{pq}\frac{2b-1}{b^2} = \sqrt{pq}\frac{4p}{(p+1)^2}$. By the computations made at the end of the previous section, we have $AB > 1$, $\sqrt{pq}A\mu > H$ and $\sqrt{pq}B\lambda > H$. Then, using (1.7) and choosing $\varepsilon > 0$ small enough we find that all coefficients of the above inequality are positive. This yields

$$\max\{Q_1Q'_1, Q_2Q'_2, Q_1Q'_2, Q'_1Q_2\} \leq C(T^2 + T'^2).$$

Combing back to (6.6) we find $S^2 \leq C(T^2 + T'^2)$, so $|S| \leq C(|T| + |T'|)$.

Step 3: Conclusion.

Let us start with the inequality $Q_2 + Q'_2 \leq C(|T| + |T'|)$ which we established in Step 2 above. Setting

$$U = \frac{u}{u_s} \quad \text{and} \quad V = \frac{v}{v_s},$$

we find

$$\int |x|^{-d}(|U|^{q+1} + |V|^{p+1})\psi^2 \leq C \int R(|U|^{q+1} + |V|^{p+1}).$$

Take first $\psi(x) = \psi_1(x)$ a standard cut-off function which is identical to 1 in B_1 and vanishes outside of B_2 . Then, the expression of R in (6.2) and the above estimate yield

$$\int_{B_1} |x|^{-d}(|U|^{q+1} + |V|^{p+1}) \leq C \int_{B_2 \setminus B_1} |x|^{-d}(|U|^{q+1} + |V|^{p+1}),$$

that is,

$$\int_{B_1} |x|^{-d}(|U|^{q+1} + |V|^{p+1}) \leq c \int_{B_2} |x|^{-d}(|U|^{q+1} + |V|^{p+1}),$$

for some $c \in (0, 1)$ which depends only on p, q and d . By rescaling and iterating, we deduce that there exist $C, C' > 0$ and $\sigma \in (0, 1)$ depend only on p, q and d such that

$$\int_{B_r} |x|^{-d} (|U|^{q+1} + |V|^{p+1}) \leq C r^{2\sigma} \int_{B_2 \setminus B_1} |x|^{-d} (|U|^{q+1} + |V|^{p+1}) \leq C' r^{2\sigma} \quad \text{for all } r \in (0, 1), \quad (6.7)$$

where we used the energy estimate (3.8) with $(a, \theta) = (\frac{q+1}{2}, 1)$ and the comparison between components (3.1) in the last inequality.

We now turn to the inequality $Q_1 + Q'_1 \leq C(|T| + |T'|)$ in which we take $\psi(x) = \psi_1(x/r)$. Using (6.7) we deduce

$$\int_{B_r} \left| \nabla (|U|^{a-1} U |x|^{-\frac{d-2}{2}}) \right|^2 + \left| \nabla (|V|^{b-1} V |x|^{-\frac{d-2}{2}}) \right|^2 \leq \tilde{C} r^{2\sigma}, \quad (6.8)$$

where, as before \tilde{C} depends only on p, q and d . The identity

$$\int \left| \nabla (h |x|^{-\frac{d-2}{2}}) \right|^2 = \left(\frac{d-2}{2} \right)^2 \int \frac{h^2}{|x|^2} + \int |x|^{-d+2} |\nabla h|^2$$

applied to (6.8) leads us to

$$r^2 \int_{B_r} \left(|\nabla |U|^{a-1} U|^2 + |\nabla |V|^{b-1} V|^2 \right) \leq C_1 r^{2\sigma}. \quad (6.9)$$

By the invariance to translation, the above estimate holds for any ball $B_r(x) \subset B_1$, $r \in (0, 1)$. If we set $\zeta = |U|^{a-1} U$ and $\zeta_{x,r}$ the spherical average of ζ over $B_r(x)$, from Poincaré-Wirtinger and (6.9) we find

$$\int_{B_r(x)} |\zeta - \zeta_{x,r}| \leq \left(\int_{B_r(x)} |\zeta - \zeta_{x,r}|^2 \right)^{1/2} \leq C_1 \left(r^2 \int_{B_r(x)} |\nabla \zeta|^2 \right)^{1/2} \leq C_2 r^\sigma.$$

The characterization of Hölder functions due to Campanato yields

$$|\zeta(x) - \zeta(y)| \leq C_3 |x - y|^\sigma \quad \text{for all } x, y \in B_{1/2}. \quad (6.10)$$

A similar inequality holds for $\xi = |V|^{b-1} V$ and this concludes our proof. \square

7. THE CASE $d - 4 > \alpha + \beta$: ASYMPTOTICS OF SOLUTIONS STABLE OUTSIDE A COMPACT SET

Lemma 13. *Let (u, v) be a positive solution of (1.1) satisfying (1.8). Assume that (1.7) holds. Then,*

$$|x|^\alpha u(x) \rightarrow 0 \quad \text{and} \quad |x|^\beta v(x) \rightarrow 0 \quad \text{as } |x| \rightarrow \infty. \quad (7.1)$$

Proof. Without loosing any generality, we may assume that (u, v) satisfies (1.8) with $K = B_1$. Let $z \in \mathbb{R}^d \setminus B_2$ and $\rho = |z|/4$. Then the pair (\tilde{u}, \tilde{v}) defined as

$$\tilde{u}(x) = \rho^\alpha u(z + \rho x), \quad \tilde{v}(x) = \rho^\beta v(z + \rho x) \quad (7.2)$$

is a solution of (1.1) satisfying (1.8) in B_2 . Letting $\tilde{U} = \tilde{u}/u_s$ and $\tilde{V} = \tilde{v}/v_s$, we have by Lemma 12 that \tilde{U}^a and \tilde{V}^b are Hölder continuous in $B_{1/2}$. In particular (6.10) yields

$$|\tilde{U}^a(x) - \tilde{U}^a(y)| \leq C |x - y|^\sigma \quad \text{for all } x, y \in B_{1/2}.$$

Letting $X = z + \rho x$ and $Y = z + \rho y$, the above estimate implies

$$|U^a(X) - U^a(Y)| \leq C \rho^{-\alpha\sigma} |X - Y|^\sigma \quad \text{for all } X, Y \in B_{\rho/2}(z),$$

where, as in Lemma 12, we denote $U = u/u_s$ and $V = v/v_s$. Averaging the above inequality in the Y variable over a ball of small radius $1/\rho$ leads to

$$\left| U^a(X) - \int_{B_{1/\rho}(X)} U^a(Y) dY \right| \leq C \rho^{-\alpha\sigma} \int_{B_{1/\rho}(X)} |X - Y|^\sigma dY = C \rho^{-\alpha\sigma - 2\sigma} \quad \text{for all } X \in B_{\rho/4}(z).$$

We next let $\rho = |z|/4 \rightarrow \infty$. Since $X \in B_{\rho/4}(z)$, this also implies $|X| \geq |z| - \rho/4 \rightarrow \infty$. Hence, for $\rho = |z|/2$ large enough the above estimate yields

$$U^a(X) - \int_{B_{1/\rho}(X)} U^a \rightarrow 0 \quad \text{as } \rho \rightarrow \infty.$$

By (6.7) (applied to a translation of U), one has

$$\int_{B_{1/\rho}(X)} U^a \rightarrow 0 \quad \text{as } \rho \rightarrow \infty.$$

Thus, $U^a(X) \rightarrow 0$ as $X \rightarrow \infty$. A similar conclusion holds for V^b and we conclude. \square

We can now apply Section 4 to conclude that Theorems 4 and 2 hold.

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INSTITUT CAMILLE JORDAN, UMR CNRS 5208, UNIVERSITÉ CLAUDE BERNARD LYON 1, VILLEURBANNE, FRANCE
Email address: louis.dupaigne@math.cnrs.fr

SCHOOL OF MATHEMATICS AND STATISTICS, UNIVERSITY COLLEGE DUBLIN, BELFIELD CAMPUS, DUBLIN 4, IRELAND

INSTITUTE OF MATHEMATICS OF THE ROMANIAN ACADEMY, 21 CALEA GRIVITEI ST., 010702 BUCHAREST, ROMANIA
Email address: marius.ghergu@ucd.ie

HIGHER INSTITUTE OF APPLIED MATHEMATICS AND COMPUTER SCIENCES OF KAIROUAN, TUNISIA
Email address: hajlaouihatem@gmail.com