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1 **Optical Properties of Organic Hazes in Water-rich Exoplanet Atmospheres:**  
2 **Implications for Observations with *JWST***

3 Chao He<sup>1\*</sup>, Michael Radke<sup>1</sup>, Sarah E. Moran<sup>1,2</sup>, Sarah M. Hörst<sup>1,3</sup>, Nikole K. Lewis<sup>4</sup>,  
4 Julianne I. Moses<sup>5</sup>, Mark S. Marley<sup>2</sup>, Eliza M.-R. Kempton<sup>6</sup>, Caroline V. Morley<sup>7</sup>, Jeff A.  
5 Valenti<sup>3</sup>, & Véronique Vuitton<sup>8</sup>

6 <sup>1</sup>Department of Earth and Planetary Sciences, Johns Hopkins University, Baltimore, MD,  
7 USA [che13@jhu.edu](mailto:che13@jhu.edu)

8 <sup>2</sup>Lunar and Planetary Laboratory, University of Arizona, Tucson, AZ, USA

9 <sup>3</sup>Space Telescope Science Institute, Baltimore, MD, USA

10 <sup>4</sup>Department of Astronomy and Carl Sagan Institute, Cornell University, Ithaca, NY, USA

11 <sup>5</sup>Space Science Institute, Boulder, CO, USA

12 <sup>6</sup>Department of Astronomy, University of Maryland, College Park, MD, USA

13 <sup>7</sup>Department of Astronomy, the University of Texas at Austin, Austin, TX, USA

14 <sup>8</sup>Institut de Planétologie et d'Astrophysique de Grenoble, Université Grenoble Alpes,  
15 CNRS, Grenoble, France

16 ***JWST* has begun its scientific mission, which includes the atmospheric**  
17 **characterization of transiting exoplanets. Some of the first exoplanets to be observed**  
18 **by *JWST* have equilibrium temperatures below 1000 K, which is a regime where**  
19 **photochemical hazes are expected to form. The optical properties of these hazes,**  
20 **which controls how they interact with light, are critical for interpreting exoplanet**  
21 **observations, but relevant data are not available. Here we measure the optical**  
22 **properties of organic haze analogues generated in water-rich exoplanet atmosphere**  
23 **experiments. We report optical constants (0.4 to 28.6  $\mu\text{m}$ ) of organic hazes for current**  
24 **and future observational and modeling efforts covering the entire wavelength range**  
25 **of *JWST* instrumentation and a large part of *Hubble*. We use these optical constants**  
26 **to generate hazy model atmospheric spectra. The synthetic spectra show that**  
27 **differences in haze optical constants have a detectable effect on the spectra, impacting**  
28 **our interpretation of exoplanet observations. This study emphasizes the need to**  
29 **investigate the optical properties of hazes formed in different exoplanet atmospheres,**  
30 **and establishes a practical procedure to determine such properties.**

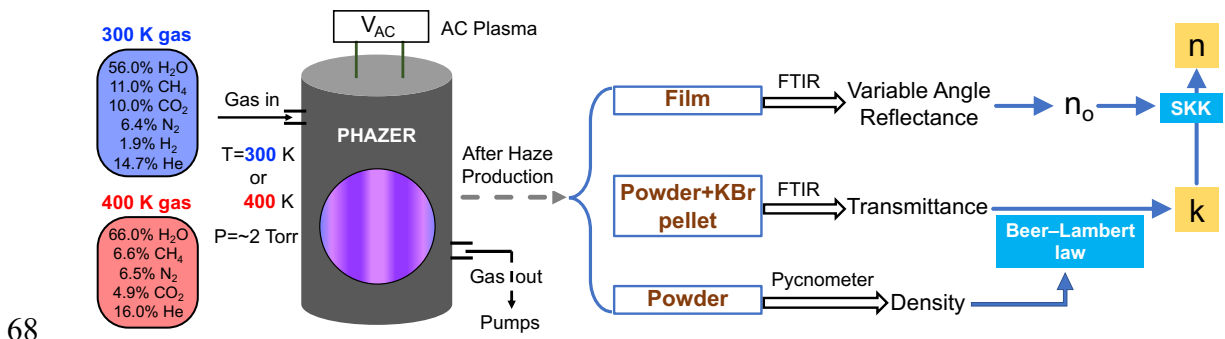
## 31 1. INTRODUCTION

32 Observations<sup>1-6</sup> indicate that many exoplanets could possess cloud and haze particles in  
33 their atmospheres and that these particles impact observed spectra. Recent modeling<sup>7,8</sup> and  
34 laboratory studies<sup>9-12</sup> suggest that organic haze particles are produced photochemically in  
35 temperate (<1000 K) exoplanet atmospheres, which are prime targets of observations for  
36 assessing habitability and searching for biosignatures beyond the Solar System. However,  
37 many exoplanets in this temperature range (including super-Earths and mini-Neptunes)  
38 have no Solar System analogs. The compositions and properties of organic hazes might be  
39 very distinct from what we know for Solar System bodies. These organic hazes can alter  
40 the transmission, emission, and reflected light spectra of exoplanets.<sup>8,13</sup> The optical  
41 properties of these hazes are essential to interpret exoplanet spectroscopic data and  
42 understand of their atmospheres.

43 Because the optical properties of exoplanet organic hazes are not yet known, the optical  
44 constants of Titan-like hazes (produced with 10% CH<sub>4</sub> in N<sub>2</sub>) from Khare et al. (1984)<sup>14</sup>  
45 and of soots (carbonaceous particles formed from incomplete combustion of  
46 hydrocarbons)<sup>15</sup> are widely used to generate models and interpret observations.<sup>8,16,17</sup> In  
47 reality, the organic hazes formed in diverse atmospheres have a variety of compositions<sup>18</sup>  
48 and therefore diverse optical properties. The optical properties of Titan-like hazes or soots  
49 simply cannot represent the wide variety of atmospheric hazes we anticipate on  
50 exoplanets.<sup>9-12</sup> It is therefore necessary to measure the optical properties of organic hazes  
51 formed over a broad range of atmospheric conditions, given that JWST has begun to deliver  
52 unprecedented observations of various exoplanets.<sup>6</sup>

53 Exoplanets are expected to exhibit a wide diversity of atmospheric compositions. The  
54 signature of water is particularly sought on potentially habitable exoplanets because water  
55 is a key element for life on Earth. For example, water vapor has been detected in the  
56 atmosphere of a non-terrestrial exoplanet (K2-18 b) in the habitable zone.<sup>19,20</sup> Modeling  
57 studies<sup>21-25</sup> suggest that water worlds could be very common among low mass exoplanets

58 and there might be many water-rich atmospheres. Previous laboratory studies<sup>9,10,18</sup> have  
 59 shown that water-rich atmospheres are likely to result in organic haze formation. Opacity  
 60 from organic hazes can mask spectral features from water and other gases.<sup>2,3,26,27</sup> Here we  
 61 measure the density and the optical properties of hazes analogous to those produced in  
 62 temperate water-rich atmospheres. Their optical constants (the real refractive indices,  $n$ ,  
 63 and the extinction coefficients,  $k$ ) are derived from 0.4 to 28.6  $\mu\text{m}$ , covering optical  
 64 wavelengths accessible with *Hubble* and ground-based facilities and the entire *JWST*  
 65 wavelength range. Fig. 1 summarizes our experimental setup, the initial gas compositions  
 66 and conditions, the measurements, and the analytical method of this study (for detailed  
 67 information, see Methods).



68  
 69 **Fig. 1.** Simplified schematic of the experimental setup, the simulated atmospheric  
 70 compositions and conditions, the measurements, and the analytical method for the current  
 71 study. Two haze analogues are produced with the PHAZER chamber by exposing the gas  
 72 mixture (300 K or 400 K) to AC plasma energy source. These two gas mixtures represent  
 73 the equilibrium compositions for atmospheres with 1000 times solar metallicity at 300 and  
 74 400 K, which are guided from chemical-equilibrium calculations (see Methods). The high  
 75 haze production rate in these two gas mixtures<sup>10</sup> allows further analysis of the properties  
 76 of the resulted haze particles. The density of the haze analogues is determined using a gas  
 77 pycnometer, and their transmittance and reflectance spectra are measured with a Fourier  
 78 transform infrared (FTIR) spectrometer. The extinction coefficients ( $k$ ) are calculated  
 79 based on the Beer-Lambert law, and the real refractive indices ( $n$ ) are derived using the  
 80 subtractive Kramers-Kronig (SKK) relation between  $n$  and  $k$ .

## 81 **2. RESULTS**

### 82 *2.1. Density of Exoplanet Haze Analogues*

83 The density of organic haze particles is an important property that impacts aerosol  
84 microphysics processes (coagulation, transport, and sedimentation) in exoplanet  
85 atmospheres. However, prior to this study the particle density was unknown due to lack of  
86 observational and experimental constraints. The haze mass density is often assumed to be  
87  $1 \text{ g cm}^{-3}$  in many microphysics models<sup>28-30</sup>, which is lower than the density we measured  
88 here. Our measurements here provide the first experimental constraints on the density of  
89 organic haze particles formed in water-rich exoplanet atmospheres, enabling realistic  
90 analysis and interpretation of observations of such exoplanets.

91 The measured densities are  $1.328$  and  $1.262 \text{ g cm}^{-3}$  (measurement uncertainties  $<1\%$ ) for  
92 the haze analogues from the simulated exoplanet atmospheres at  $300$  and  $400 \text{ K}$ ,  
93 respectively. Their densities are lower than the haze analogues produced with the PHAZER  
94 chamber for icy bodies ( $1.35$ - $1.43 \text{ g cm}^{-3}$  for Titan, Triton, and Pluto haze analogues) in  
95 our Solar System.<sup>31</sup> The density differences reflect their distinct chemical compositions.  
96 Elemental analysis has shown that the two haze analogues for exoplanets have higher  
97 oxygen ( $>15\%$  compared to  $<10\%$ ) but lower nitrogen ( $\sim 22\%$  compared to  $\sim 40\%$ ) contents  
98 relative to haze analogues for Titan, Triton and Pluto.<sup>18,32</sup> Besides the differences in  
99 elemental compositions, the chemical structures in each haze analogue, such as their  
100 molecular weight, polarity, and degree of unsaturation, also affect their density.

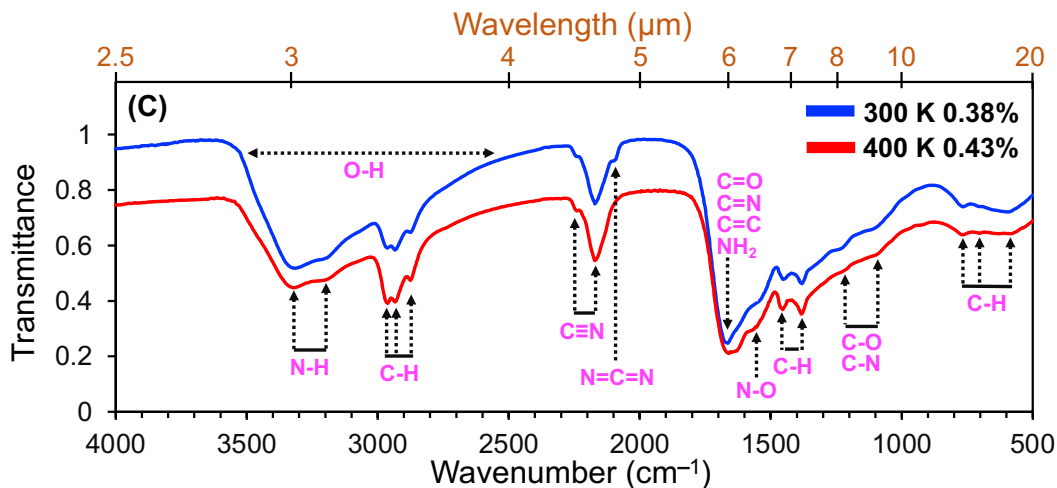
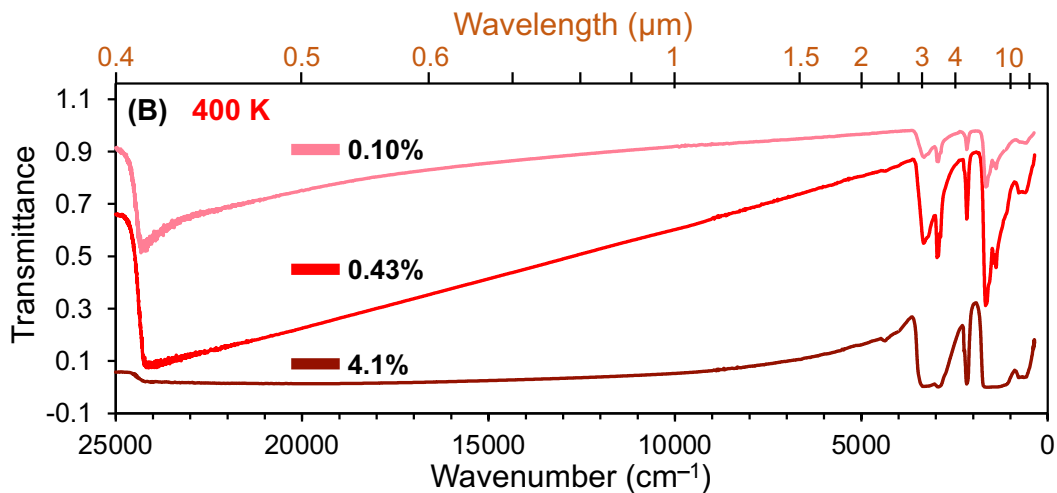
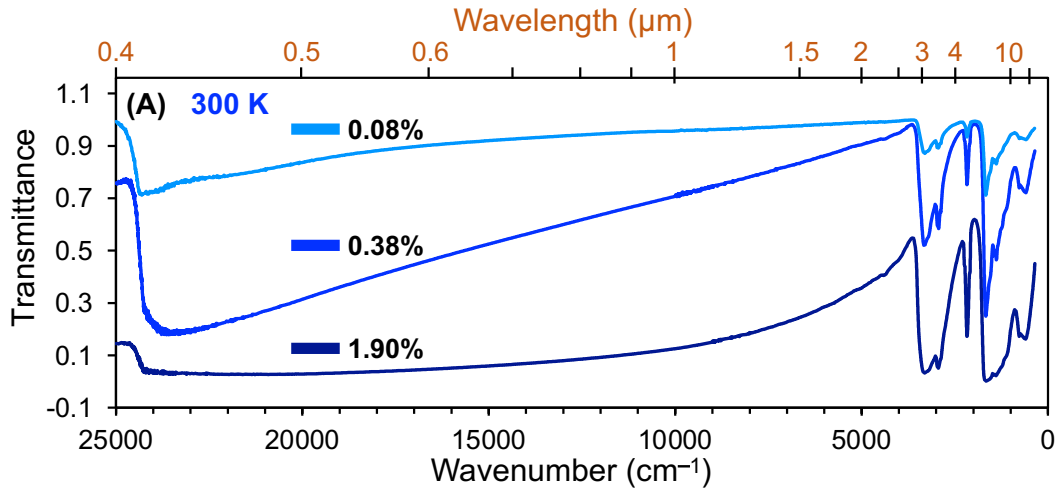
### 101 *2.2. Transmittance of Exoplanet Haze Analogues and their Functional Groups*

102 Exoplanet atmospheric transmission observations, obtained as the planet passes in front of  
103 host star, will be shaped by any hazes in the planet's atmosphere. Fig. 2 shows the  
104 transmittance of the two exoplanet haze analogues ( $300$  and  $400 \text{ K}$  samples). For each  
105 sample, the transmittance spectra are obtained at three different concentrations in

106 potassium bromide (KBr) to capture both strong and weak absorption features. KBr is a  
107 typical carrier (as pellet or disk) for spectroscopy measurements because it is optically  
108 transparent from ultraviolet (0.22  $\mu\text{m}$ ) to far IR ( $\sim 30 \mu\text{m}$ ). As shown in Fig. 2, the general  
109 spectral shape and features are similar at both temperatures. They both have an absorption  
110 feature at 0.42  $\mu\text{m}$ , suggesting that the samples contain aromatic compounds and/or  
111 unsaturated species with conjugated pi bonds.<sup>33,34</sup> The spectra are featureless from 0.5 to  
112 2.5  $\mu\text{m}$  but the absorption becomes weaker as the wavelength increases. Many absorption  
113 features appear from 2.5 to 20  $\mu\text{m}$  due to bond vibrations of various organic functional  
114 groups in the haze analogues.

115 We expand this region in Fig. 2C to show the characteristic frequencies of different bonds.  
116 Fig. 2C identifies the bonds responsible for each spectral feature, including the  
117 characteristic absorptions of O–H, N–H, C–H, C $\equiv$ N, –N=C=N–, C=O, C=N, C=C, N–O,  
118 C–O, and C–N bonds.<sup>35,36</sup> Table 1 summarizes the vibration modes of these bonds and their  
119 peak intensities. The absorption feature at 3500–2500  $\text{cm}^{-1}$  (2.86–4.00  $\mu\text{m}$ ) is due to O–H  
120 bond stretching. This strong, broad feature indicates the predominance of alcohols (3500–  
121 3200  $\text{cm}^{-1}$ ) and carbonic acids (3300–2500  $\text{cm}^{-1}$ ) in the haze analogues. The stretching of  
122 N–H (3350–3300 and 3215–3190  $\text{cm}^{-1}$ ) and C–H (2965, 2933, 2875  $\text{cm}^{-1}$ ) bonds is present  
123 in this range as well but appears as relatively sharp peaks. Other absorption features in the  
124 spectra indicate the presence of nitriles, aromatics, and unsaturated and saturated organics  
125 in both the 300 and 400 K haze analogues. A small amount of carbodiimide compounds (–  
126 N=C=N–,  $\sim 2100 \text{cm}^{-1}$ ) are only present in the 300 K haze sample, probably because lower  
127 temperature promotes carbodiimide formation from cyanamide.<sup>37</sup> Compared to Titan-like  
128 hazes<sup>38,39</sup>, the exoplanet haze analogues in this study have obvious compositional  
129 differences, which include various oxygen-containing groups. This is not unexpected  
130 because there are two major oxygen-containing gases ( $\text{H}_2\text{O}$  and  $\text{CO}_2$ ) in the initial  
131 atmospheres. The presence of oxygen-containing groups in the exoplanet haze analogues  
132 is supported by other analyses. Elemental analysis shows that these two haze analogues

133 have large oxygen contents (15% or 17% in mass), and high-resolution mass spectra reveal  
 134 many oxygen-containing molecules in these samples.<sup>18</sup> The compositional difference leads  
 135 to distinctive optical properties of the exoplanet haze analogues.



136

137 **Fig. 2.** The transmittance spectra of two exoplanet haze analogues formed in water-rich  
 138 atmospheres at 300 K (A) and 400 K (B). The lines with different shades of color show the  
 139 spectra of the samples with different concentrations in KBr (lighter color indicates lower  
 140 concentration and darker indicates higher). The sample is mixed with KBr, and the  
 141 percentage concentration is by mass. The transmittance spectra are obtained at different  
 142 concentrations to capture both strong and weak absorption features in the samples. Panel  
 143 C is the expansion of two spectra from panels A (0.38% of the 300 K sample) and B (0.43%  
 144 of the 400 K sample) from 2.5 to 20  $\mu\text{m}$  to show the absorption features of various  
 145 functional groups. Different bonds are labeled in the figure near their absorption features.  
 146

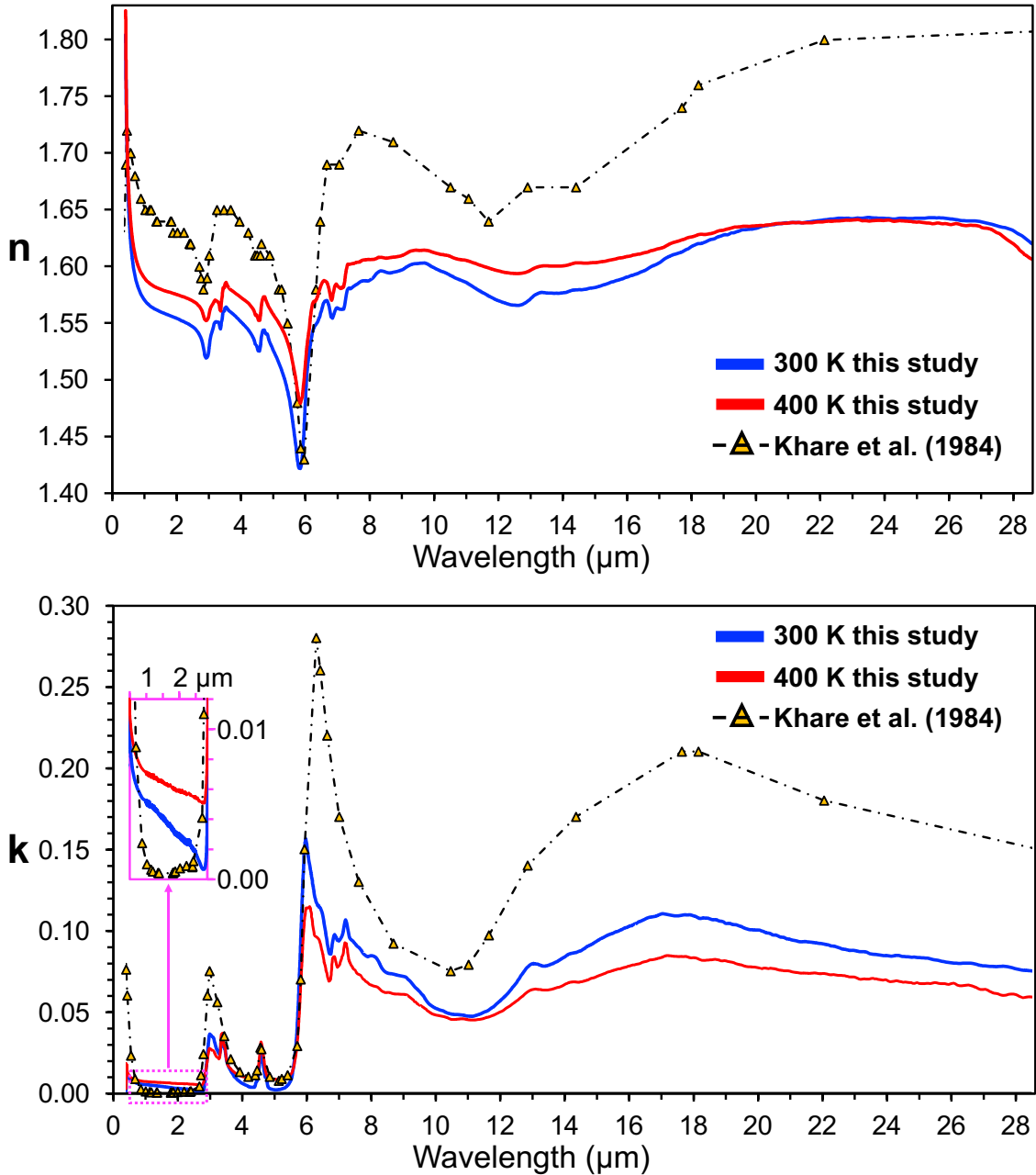
147 **Table 1.** Functional group assignments to characteristic absorption frequencies of two  
 148 exoplanet haze analogues as observed in the transmittance spectra (Fig. 2)

Frequency ( $\nu$ ) $\text{cm}^{-1}$	Wavelength ( $\lambda$ ) $\mu\text{m}$	Functional group	Intensity
3500-2500	2.86-4.00	-O-H stretching, in alcohols and carbonic acids	strong, broad
3350-3300	2.98-3.03	-NH <sub>2</sub> or -NH- stretching	strong
3215-3190	3.11-3.13	-NH <sub>2</sub> stretching or overtone of NH <sub>2</sub> bending	strong
2965	3.37	-CH <sub>3</sub> asymmetric stretching	strong, sharp
2933	3.41	-CH <sub>2</sub> - asymmetric stretching	strong, sharp
2875	3.48	-CH <sub>3</sub> symmetric stretching	strong, sharp
2244-2232	4.46-4.48	-C $\equiv$ N or R-C=C-R stretching	medium, shoulder
2172	4.60	Conjugated -C $\equiv$ N stretching	strong, sharp
2108-2094	4.74-4.77	-N=C=N- stretching	medium, shoulder
1680-1626	5.95-6.15	C=O, C=N, C=C stretching or -NH <sub>2</sub> scissors bending	very strong
1580-1540	6.33-6.49	C=C stretching (aromatic) or N-O stretching (nitro)	medium, shoulder
1462-1450	6.84-6.90	sp <sup>3</sup> C-H bending	strong, sharp
1386-1378	7.21-7.26	sp <sup>3</sup> C-H bending or aldehydic C-H bending	strong, sharp
1247-1224	8.02-8.17	C-O or C-N stretching (aromatic)	weak, shoulder
1126-1104	8.88-9.06	C-O or C-N stretching (aliphatic)	weak, shoulder
775-762	12.9-13.12	sp <sup>2</sup> C-H bending	medium
714-701	14.0-14.26	sp <sup>2</sup> C-H bending	medium
612-584	16.34-17.12	sp <sup>2</sup> C-H bending	medium



149 *2.3. Optical Constants of Exoplanet Haze Analogues*

150 Optical constants for hazes can be used incorporate their features into theoretical spectra  
151 and for comparisons with retrieved values from observations. Fig. 3 shows the optical  
152 constants of two exoplanet haze analogues (300 and 400 K). The real refractive indices ( $n$ )  
153 are in the range between 1.43 and 1.82 for both samples, and the extinction coefficients ( $k$ )  
154 vary in a broad range from  $\sim 10^{-3}$  to  $10^{-1}$ . The detailed data and their uncertainties can be  
155 found in Supplementary Information (Table S1). The uncertainties of the  $k$  values are  
156 estimated by considering the uncertainties of the transmittance measurements, the  
157 thickness calculations, and the average relative standard deviation of  $k$  values obtained for  
158 the samples at three different concentrations. For the  $n$  values, the uncertainties are  
159 determined from the error propagations of the uncertainty of  $n_0$  at the anchor point and the  
160 integration of the  $k$  values. The uncertainties for the  $k$  values are about 1-4% at longer  
161 wavelengths ( $>2.9 \mu\text{m}$ ) where the  $k$  values are larger than 0.01, but increase up to 9% at  
162 shorter wavelengths (0.4 to  $2.9 \mu\text{m}$ ) where the  $k$  values are smaller than 0.01. The  
163 determined  $n$  values have small uncertainties (3-4%) across the measured wavelength range  
164 because of the accuracy of the  $n_0$  (uncertainty less than 1%) and the integration of the  $k$   
165 values over the entire wavelength range. The  $n$  and  $k$  values of the two haze analogues (300  
166 and 400 K) are similar in general and share many fine-scale features. Over most of the  
167 measured wavelength range, the 300 K sample has slightly lower  $n$  values but slightly  
168 higher  $k$  values relative to the 400 K sample. This could be caused by differences in their  
169 chemical compositions. For example, the 300 K sample has higher extinction coefficient  
170 at  $\sim 3 \mu\text{m}$ , indicating that it contains more N-H bonds. This is consistent with the elemental  
171 analysis that shows the 300 K sample has higher nitrogen content than the 400 K sample  
172 (27% vs 21%).<sup>18</sup> In addition, the higher extinction coefficients at 5.9-6.5  $\mu\text{m}$  suggest that  
173 there are more double bonds in the 300 K sample.



174

175 **Fig. 3.** Optical constants of two exoplanet haze analogues along with those of Titan-like  
 176 hazes from Khare et al (1984): (A) the real refractive indices ( $n$ ), and (B) the extinction  
 177 coefficients ( $k$ ) as function of wavelength (0.4 to 28.6  $\mu\text{m}$ ). The inset of (B) shows the  
 178 difference of  $k$  values between this study and the one from Khare et al (1984), when  $k$   
 179 is smaller than 0.01 from 0.5 to 2.9  $\mu\text{m}$ .

180

181 The optical constants of Titan-like hazes from Khare et al. (1984)<sup>14</sup> are also plotted in Fig.

182 3 for comparison. Due to its extensive spectral coverage (from soft X-ray, 0.025  $\mu\text{m}$ , to

183 microwave frequencies, 1000  $\mu\text{m}$ ), the optical constants from Khare et al. (1984)<sup>14</sup> are  
184 widely used in atmospheric modeling and observation interpretations for both Solar System  
185 bodies (Titan<sup>40</sup>, Triton<sup>41</sup>, and Pluto<sup>42</sup>) and many exoplanets.<sup>30,31,43</sup> However, the haze  
186 analogues used in Khare et al. (1984)<sup>14</sup> for the optical measurements were produced with  
187 DC plasma discharge in  $\text{N}_2/\text{CH}_4 = 90/10$  gas mixture at 0.2 mbar, with the purpose of  
188 simulating haze particles formed in Titan's upper atmosphere.

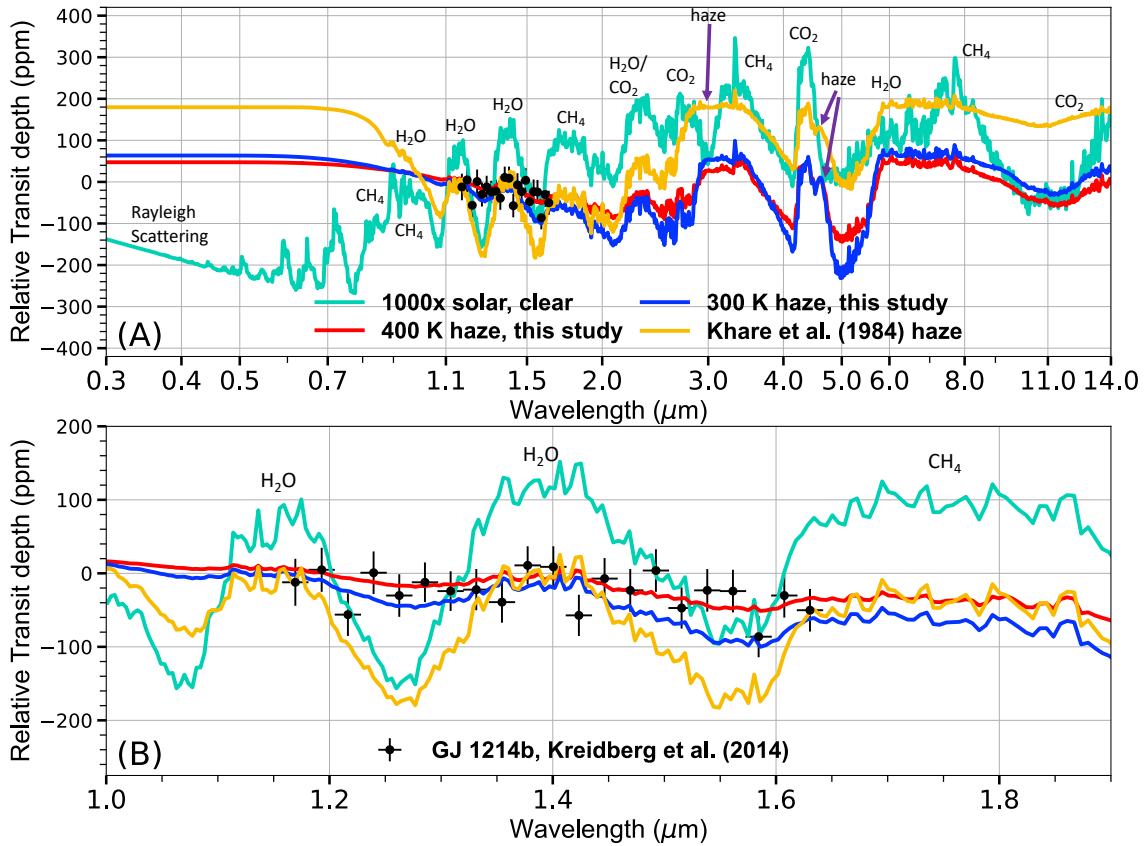
189 As shown in Fig. 3, the general trend of the optical constants of two exoplanet haze  
190 analogues in this study is similar to that of the Titan-like hazes from Khare et al. (1984)<sup>14</sup>  
191 because of the organic nature of these materials. However, our  $n$  and  $k$  values are  
192 substantially smaller than those of Khare et al. (1984)<sup>14</sup> over most of the measured  
193 wavelength range. Both of our two samples have smaller  $n$  values than the Titan-like hazes  
194 from Khare et al. (1984)<sup>14</sup> except at 5.8-6.2  $\mu\text{m}$  where their  $n$  values are similar to those of  
195 our 300 K sample. The  $k$  values from Khare et al. (1984)<sup>14</sup> at 0.4-0.6  $\mu\text{m}$ ,  $\sim 3$   $\mu\text{m}$ , and above  
196 6  $\mu\text{m}$  are about two times greater than our  $k$  values; they are comparable to the  $k$  values of  
197 our 400 K sample from 3.4 to 5.8  $\mu\text{m}$ ; but they are smaller than our  $k$  values in the short  
198 wavelength region (0.7-2.6  $\mu\text{m}$ ). Besides the differences of the absolute  $n$  and  $k$  values, the  
199 data from Khare et al. (1984)<sup>14</sup> has fewer features than our values mainly due to two reasons.  
200 First, the spectral measurements in Khare et al. (1984)<sup>14</sup> were done at very low resolution,  
201 especially over the longer wavelengths (6-28.6  $\mu\text{m}$ ); second, the compositions are distinct  
202 between our samples and those in Khare et al. (1984). The low-resolution can lead to fine  
203 features not being resolvable (such as those from 6-10  $\mu\text{m}$ ), and the compositional  
204 difference may cause some features to be weaker or absent (such as those at 3-3.5  $\mu\text{m}$ ).  
205 The higher spectral resolution of the current study reveals new spectral features that may  
206 enable measurements of haze composition, rather than just haze detection. Nonetheless, the  
207 differences in the absolute values and the features can have consequential impacts on  
208 atmospheric modeling and observational interpretation of exoplanets.

209 *2.4. Effect of New Optical Properties of Exoplanetary Hazes on Atmospheric Transmission*  
210 *Spectra*

211 We next briefly explore the effects of these newly derived exoplanet haze optical properties  
212 on transmission spectra for a representative exoplanet atmosphere. Since the experiments  
213 were run for hazy, water-rich atmospheres, we choose the parameters of the well-studied,  
214 likely aerosol-laden exoplanet GJ 1214 b as a reasonable example planet to demonstrate  
215 how our new optical properties would affect observations. GJ 1214 b has a mass of  $8.17 \pm$   
216  $0.43 M_{\oplus}$ , a radius of  $2.742 \pm 0.05 R_{\oplus}$ , and an equilibrium temperature of 596 K.<sup>44</sup> The  
217 planet’s mass and radius suggest its bulk density could be consistent with a hydrogen-rich  
218 atmosphere, or a steam atmosphere with a significant hydrogen/helium fraction<sup>44</sup> as in our  
219 experiments. The temperature is sufficiently warm that water should not condense out of  
220 the upper atmosphere. Observations with *Hubble* show its atmosphere has a featureless  
221 transmission spectrum from 1.1 to 1.7 microns, which could be indicative of a high altitude  
222 ( $\sim$ mbar) haze layer.<sup>1</sup> The incident UV flux is expected to drive significant photochemistry  
223 and potentially haze formation.<sup>45,46</sup>

224 Fig. 4 shows the synthetic spectra of a water-rich atmosphere reflecting the initial gas  
225 composition from the experiments at 400 K around a GJ 1214 b-like planet with the effects  
226 of a haze layer. We use the open-source aerosol modeling code *Virga*<sup>47,48</sup> and the open-  
227 source radiative transfer suite *PICASO*<sup>49</sup> to compute Mie properties for our hazes and  
228 implement them into a synthetic transmission spectrum model of an exoplanetary  
229 atmosphere (for detailed information about the haze modeling and generating the synthetic  
230 spectra, see Methods). We show spectra based on our newly derived optical properties for  
231 300 and 400 K water-rich atmospheric hazes and for the commonly used Titan-like haze  
232 of Khare et al. (1984)<sup>14</sup>. Our aim in this work was to isolate the effect of the different haze  
233 optical properties on the atmospheric transmission spectra rather than perform a full model  
234 parameter study, and so we held all other haze parameters (e.g., particle size distribution,  
235 number density, vertical mixing) fixed. Interpretation of actual atmospheric spectral data

236 would require a full investigation into all these variables but is beyond the scope of this  
 237 work.



238  
 239 **Fig. 4.** Model spectra of a water-rich atmosphere around a GJ 1214 b -like planet, showing  
 240 the effect of our newly measured haze optical properties. We also show the existing *Hubble*  
 241 data of GJ 1214 b (Kreidberg et al. 2014)<sup>1</sup>. All spectra have been normalized by subtracting  
 242 the mean of each transit depth from the spectrum, to better show the relative sizes of  
 243 absorption features. The top panel (A) shows the spectra of our modeled atmospheres from  
 244 0.4 to 14 microns; bottom panel (B) shows our model atmospheres focused on *Hubble*  
 245 Wide Field Camera 3/G141 instrument wavelengths. In teal, we show models of clear  
 246 atmospheres with the laboratory atmospheric composition. Orange lines show hazy models  
 247 generated using the optical properties for Titan-like hazes as measured by Khare et al.  
 248 (1984). Red and blue lines show models generated using our newly derived optical  
 249 properties for hazes made from 400 and 300 K water-rich atmospheres, respectively. The  
 250 molecules and/or hazes responsible for the atmospheric features are indicated in each plot.  
 251 **Given the same other atmospheric assumptions, the presence of our newly measured**  
 252 **hazes would be differentiable from a Titan-like haze with observations from existing**  
 253 **space-based observatories like *Hubble* and *JWST*, as with the GJ 1214 b data shown**  
 254 **here.**  
 255

256 Throughout the spectrum, the values of  $n$  and  $k$  between our newly measured water-rich  
257 derived hazes and those of the Titan-like hazes<sup>14</sup> show similar behavior but at different  
258 intensities, as shown in Fig. 3. As such, all the hazy synthetic spectra we show in Fig. 4  
259 mute the gaseous spectral features of the atmosphere, but to differing extents depending on  
260 the optical properties assumed for the haze particles. Notably, the Khare Titan-like hazes  
261 more heavily mute the atmospheric spectral features longward of  $3.0\ \mu\text{m}$ , while our water-  
262 rich hazes more heavily mute the atmospheric spectrum between  $0.7$  and  $2.5\ \mu\text{m}$ — that is,  
263 the wavelength region of focus for most *Hubble* exoplanet studies. In particular, the Titan-  
264 like hazes more strongly absorb at  $\sim 3\ \mu\text{m}$  and  $\sim 6\ \mu\text{m}$ , which probe functional groups due  
265 to amine (N-H) as well as carbon- and nitrogen-bearing (C=C, C=N, NH<sub>2</sub>) bonds. The  
266 transit depth differences between haze models range from  $\sim 100$  ppm in the mid-IR to  $\sim 150$   
267 ppm in the visible and near-IR, all of which are well within the sensitivity of *JWST* and  
268 *Hubble* measurements,<sup>1,51-55</sup> though this result relies on the same particle size distributions  
269 for the different kinds of hazes.

270 Additionally, the strength of the Rayleigh scattering slope from the visible towards the  
271 near-UV is enhanced by all hazes, though the Titan-like hazes absorb more strongly  
272 because of their larger  $k$  over the shortest wavelengths ( $0.4$  to  $0.7\ \mu\text{m}$ ). The difference  
273 between the visible-UV slopes of Titan-like and water-rich derived hazes is  $\sim 150$  ppm for  
274 our modeled atmospheres. Smaller particles would enhance the variation between the  
275 visible-UV slopes (see Extended Data Fig. 3). Such differences are within the precision of  
276 *Hubble*'s WFC3/UVIS instrument, which covers  $0.2$  to  $0.8\ \mu\text{m}$ . Since we have explicitly  
277 modeled the same haze particle size distributions and number densities for each set of  
278 optical properties, these differences here stem purely from the stronger absorption by the  
279 Titan-like haze over these wavelengths. In other words, using the optical properties from  
280 this study or from Khare et al. (1984)<sup>14</sup> to fit spectra from real observations would result in  
281 different inferred particle sizes and number densities for the same atmosphere to achieve  
282 the same scattering slope.

283 Most interestingly, our water-rich derived hazes almost completely mute the atmospheric  
284 spectrum over the HST WFC3/G141 bandpass, but the Titan-like haze of Khare et al.  
285 (1984)<sup>14</sup> allows the water and methane atmospheric features to peek through from 1 to 2  
286  $\mu\text{m}$  because its  $k$  values are an order of magnitude lower in the NIR. This results in our  
287 water-rich hazes providing a better match to the existing GJ 1214 b data from *Hubble*<sup>1</sup>  
288 compared to the Khare et al. (1984)<sup>14</sup> Titan-like hazes for the same abundance of haze mass  
289 loading. Both sets of haze optical properties allow atmospheric features to peek through at  
290 redder wavelengths which are covered by *JWST*. The size of these atmospheric features at  
291 longer wavelengths, even muted by haze, is still well within the precision of *JWST*  
292 measurement. Future model development and modeling studies, beyond the scope of this  
293 work, are needed to fully assess how vertical mixing and sedimentation of haze layers with  
294 these optical properties influence atmospheric spectra.

### 295 3. DISCUSSION

296 So far, the optical constants of the Titan-like hazes by Khare et al. (1984)<sup>14</sup> and various  
297 soots<sup>15</sup> have served as the only source for optical properties of organic haze in exoplanet  
298 atmospheres. Our study provides optical constants (0.4 to 28.6  $\mu\text{m}$ ) of organic hazes  
299 photochemically produced in water-rich exoplanet atmospheres. The optical constants of  
300 our water-rich derived exoplanet hazes differ from those of the Titan-like hazes<sup>14</sup>, therefore  
301 affecting transmission, thermal emission, and reflected light spectra of exoplanets to  
302 different extents. When interpreting actual observations of exoplanets, using different  
303 optical constants of hazes can lead to different conclusions; in some cases, using  
304 inapplicable optical constants may lead to misinterpretations. Such misinterpretations  
305 could include incorrect atmospheric abundances of key species such as water and  
306 methane<sup>56,57</sup>, or incorrect aerosol particle sizes and extents<sup>30</sup>, which could then further  
307 affect the temperature structure<sup>17</sup> and thus the inferred dynamics and climate feedbacks  
308 within an atmosphere.<sup>58</sup>

309 JWST has an unprecedented capability to detect faint chemical signatures in exoplanet  
310 atmospheres and is set to observe a wide variety of exoplanet types. The haze particles  
311 formed in those exoplanet atmospheres are likely to have different compositions, which  
312 would then have distinct optical properties. The optical properties of different hazes are  
313 essential to interpreting observations, but we should be cautious when applying the existing  
314 haze optical constants to different types of exoplanets. In addition, the optical constant data  
315 of different hazes are necessary input parameters for atmospheric modeling of various  
316 exoplanets to understand their physical and chemical processes (i.e., their temperature  
317 structures, climates, etc.). The two sets of optical constants reported here are applicable for  
318 temperate water-rich exoplanet atmospheres and can be used for current and future  
319 observational and modeling efforts of such atmospheres. More laboratory work is needed  
320 to determine optical constants of haze analogues formed in various exoplanet atmospheric  
321 regimes. This study demonstrates a feasible avenue to determine such properties in the  
322 future.

## 323 **4. METHODS**

### 324 *4.1 Production of Haze Analogues*

325 The exoplanet haze analogues are produced using the Planetary Haze Research (PHAZER)  
326 experimental setup at Johns Hopkins University.<sup>31</sup> The detailed experimental procedure  
327 has been described previously.<sup>9-12</sup> This study uses haze analogues produced in the two  
328 water-rich atmospheres from our previous investigations.<sup>9,10,18</sup> The initial gas mixtures for  
329 our experiments are guided by the calculations from the chemical-equilibrium models of  
330 Moses et al. (2013).<sup>59</sup> Photochemical hazes are expected to play an important role in the  
331 atmospheres of exoplanets with equilibrium temperature below 1000 K, especially those  
332 planets with enhanced metallicities or enhanced C/O ratios.<sup>7,9,10</sup> The detailed compositions  
333 of actual exoplanet atmospheres have not yet been measured well, so we resort to chemical  
334 equilibrium calculations<sup>59</sup> to guide our experiments. The equilibrium atmospheric



335 compositions provide a reasonable starting point to investigate the photochemical  
336 processes in the atmospheres of super-Earths and mini-Neptunes with higher metallicity.  
337 In our previous studies<sup>9-12</sup>, we explored atmospheres with a range of metallicities (100×,  
338 1000×, and 10,000× solar metallicity) and equilibrium temperatures (300, 400, 600, and  
339 800 K). The two water-rich atmospheres represent the equilibrium compositions for  
340 atmospheres with 1000× solar metallicity at 300 and 400 K. These two atmospheres  
341 produce haze at a higher rate than the other cases<sup>10</sup>, which allows further analysis of the  
342 properties of the resulting haze particles.

343 Here, we briefly recount the production procedure. Gas mixtures, excluding water vapor,  
344 are premixed in a stainless-steel cylinder with high-purity gases purchased from Airgas  
345 (H<sub>2</sub>-99.9999%, He-99.9995%, N<sub>2</sub>-99.9997%, CH<sub>4</sub>-99.999%, CO<sub>2</sub>-99.999%). The total  
346 flow rate is 10 standard cubic centimeters per minute (sccm), and the premixed gas mixture  
347 flows at a proportional rate of 10 sccm based on the mixing ratio. Water vapor is introduced  
348 to the system at a pressure (corresponding to the mixing ratio) from HPLC water (Fisher  
349 Chemical) at the desired temperature maintained by a dry ice/methanol/water cold bath. A  
350 heating coil is employed to heat the gas mixture (including water vapor) to the experimental  
351 temperature (400 or 300 K). The gas mixture is then exposed to an AC glow discharge in  
352 a reaction chamber, which initiates complex chemical processes and leads to the formation  
353 of new gas-phase products and solid particles. After 72 hours of continuous discharge flow,  
354 we collect the haze analogues as powders and films on quartz substrate disks in a dry (<0.1  
355 ppm H<sub>2</sub>O), oxygen free (<0.1 ppm O<sub>2</sub>), N<sub>2</sub> glove box (Inert Technology Inc., I-lab 2GB).  
356 The haze analogues are kept in the glovebox until further analysis to avoid contamination  
357 from Earth's atmosphere and light sources.

#### 358 ***4.2 Density Measurements***

359 The density of exoplanet haze analogues has not been reported previously. We determine  
360 the density of our haze analogues by measuring the mass and volume of a certain amount

361 of sample by using an analytical balance ( $\pm 0.0001$  gram) and a gas pycnometer (AccPyc II  
362 1340, Micrometrics), respectively. We first fill the powder sample to  $\sim 70\%$  of a  $0.1 \text{ cm}^3$   
363 cup and measure the mass of the sample in the cup on the balance (where the mass of the  
364 cup is known). Next, we measure the volume of the sample at ambient temperature with  
365 the gas pycnometer that uses the gas (helium) displacement method to determine the  
366 sample volume based on Boyle's law of volume–pressure relationships. We set the  
367 pycnometer to measure 20 cycles, and the average volume from the 20 measurements is  
368 reported with a standard deviation less than  $0.0001 \text{ cm}^3$ . The density is then calculated by  
369 dividing the mass by the volume.

### 370 ***4.3 Vacuum Fourier-transform Infrared Spectroscopy (FTIR) Measurement***

371 We employ a Vertex 70v FTIR spectrometer (Bruker Optics) to characterize the spectral  
372 properties of the haze analogues. The Vertex 70v is a vacuum spectrometer, which can  
373 eliminate the spectral features ( $\text{H}_2\text{O}$  or  $\text{CO}_2$  absorption) from Earth's atmosphere and  
374 increase the peak sensitivity without masking very weak spectral features. The wavelength  
375 range of the spectrometer is  $0.4$  to  $28.6 \text{ }\mu\text{m}$  ( $25,000 \text{ cm}^{-1}$  to  $350 \text{ cm}^{-1}$ ) with a maximum  
376 resolution of  $0.4 \text{ cm}^{-1}$ , covering the whole observable wavelength range of *JWST* and a  
377 large part of *Hubble*. This spectrometer is configured with a Reflection/Transmission  
378 accessory (A510 Q/T, Bruker) and Seagull Variable Angle Reflection Accessory (Harrick  
379 Scientific), allowing transmittance measurements and reflectance measurements at  
380 different angles.

381 To derive the optical constants of the haze analogues, we measure the reflectance of the  
382 film samples on quartz discs and transmittance of the powder samples using the potassium  
383 bromide (KBr) pellet method. The films deposited on optical-grade quartz substrates in the  
384 reaction chamber are directly used for the reflectance measurements (4.3.1). The procedure  
385 for making KBr pellets is described in 4.3.2 and the transmittance measurements of the  
386 pellets are described in 4.3.3.

### 387 **4.3.1 Reflectance Measurements of the films**

388 Using the Seagull Variable Angle Reflection Accessory, we measure the reflectance of  
389 each haze analogue film at two different angles of incidence, 15° and 45°. The reflectance  
390 of the films on quartz substrates is measured under vacuum (below 0.2 mbar) at room  
391 temperature (294 K). We measure the reflectance in the wavelength range from 0.4 to 1.1  
392  $\mu\text{m}$  using a quartz beamsplitter and silicon diode detector. We acquire 1000 scans and  
393 average them to obtain spectra with a resolution of  $\sim 5\text{ cm}^{-1}$ . The reflectance of an  
394 aluminum standard mirror is measured as reference. The reflectance spectrum of the  
395 sample is the ratio of the sample measurement to the reference measurement.

### 396 **4.3.2 Preparation of the KBr Pellets**

397 A quality KBr pellet is critical for quantitative measurements. To ensure that the pellet is  
398 dry, KBr powder is dried in an oven overnight at 110 °C and immediately transferred to  
399 the dry N<sub>2</sub> glove box after heating. Subsequent pellet-making steps are all performed in the  
400 glove box. First, the KBr and the sample powders are separately ground to fine particles (<  
401 2  $\mu\text{m}$ ) with ShakIR (a ball-mill grinder from Pike Technologies). Dry and fine KBr is used  
402 to make a pure KBr pellet as a spectral reference. The fine particles of the sample and KBr  
403 are mixed using the ShakIR to create a homogenous mixture of sample/KBr before making  
404 pellets.

405 The concentration of the sample in KBr needs to be optimized for transmittance  
406 measurements, usually between 0.5-3% depending on the absorption coefficients. Due to  
407 the complex nature of the haze analogues, we prepare the haze/KBr mixtures with three  
408 different concentrations to reveal both strong and weak absorption features in the samples.  
409 The actual concentrations of our mixtures range from 0.08 to 4.1%. Because the quantities  
410 of the haze analogues in the KBr pellet are relatively small, we use a two-step dilution  
411 method to minimize errors in the measurement process. As reported in Myers et al. (2019)<sup>60</sup>,  
412 this method can reduce gravimetric error for the mass measurements from nearly 10% to

413 less than 1%. For example, to prepare a 0.1% mixture, we weigh 25 mg of the haze  
414 analogue using the analytical balance ( $\pm 0.0001$  gram) and mix it with 975 mg of KBr in  
415 the ShakIR to create a 2.5% mixture; next, we take 40 mg of the 2.5% mixture and 960 mg  
416 of pure KBr, and then mix them thoroughly using the ShakIR. After two dilutions, a 1:1000  
417 (or 0.1%) homogenous sample/KBr mixture is ready to be pressed into a pellet.

418 Approximately 200 mg of finely ground KBr powder or homogenous mixture is placed  
419 into a 13 mm pellet die (Pike Technologies), and then is pressed on a CrushIR 15 Ton  
420 Digital Press (Pike Technologies). We apply 3 tons of force for 2 seconds, 7 tons for 30  
421 seconds, and 10 tons for 120 seconds. The pressed pellet (diameter: 13 mm, thickness:  $\sim 0.5$   
422 mm) is removed from the die, mounted on a standard sample holder, and loaded into the  
423 FTIR spectrometer for transmittance measurement.

#### 424 **4.3.3 Transmittance Measurements of the Pellets**

425 We measure the transmittance of the pressed pellets over a wavelength range from 0.4 to  
426 28.6  $\mu\text{m}$  under vacuum (below 0.2 mbar) at room temperature (294 K). Two detectors  
427 (silicon diode and DLaTGS detector) and two beamsplitters (quartz beamsplitter and KBr  
428 beamsplitter) are used to cover the whole wavelength range. The silicon diode detector and  
429 quartz beamsplitter are used from 0.4 to 1.11  $\mu\text{m}$  (25000 to 9000  $\text{cm}^{-1}$ ); the DLaTGS  
430 detector and quartz beamsplitter are used from 0.83 to 1.25  $\mu\text{m}$  (12000 to 8000  $\text{cm}^{-1}$ ); and  
431 the DLaTGS detector and KBr beamsplitter are used from 1.11 to 28.6  $\mu\text{m}$  (9000 to 350  
432  $\text{cm}^{-1}$ ). Overlapping data confirm that the spectrometer is calibrated properly across  
433 different wavelength ranges. For each measurement, 500 scans are acquired with a  
434 resolution of  $\sim 2$   $\text{cm}^{-1}$ . The transmittance of a pure KBr pellet is measured as a reference.  
435 The transmittance spectrum of the sample is the ratio of the sample measurement to the  
436 reference measurement.

#### 437 **4.4 Optical Constants Derivation**

438 With the data from the measurements described above, we can derive the optical constants  
 439 (complex refractive index,  $n+ik$ , where  $n$  is the real refractive index and  $k$  is the imaginary  
 440 refractive index or the extinction coefficient) of the haze analogues. Based on the Beer-  
 441 Lambert law, the absorbance ( $A$ ) can be expressed as Equation 1 (Eq. 1):

$$442 \quad A = \alpha(\nu) * d \quad \text{Eq. 1}$$

443 where  $\alpha$  is absorption coefficient,  $\nu$  is the wavenumber, and  $d$  is the effective thickness of  
 444 the sample. The absorbance ( $A$ ) can be determined from the FTIR measurements (Eq. 2):

$$445 \quad A = -\ln T = -\ln \frac{I}{I_0} = \ln \frac{I_0}{I} \quad \text{Eq. 2}$$

446 where  $T$  is the transmittance,  $I$  is the light intensity passing through the sample (haze/KBr  
 447 pellet), and  $I_0$  is the light intensity of the reference (pure KBr pellet). The effective  
 448 thickness ( $d$ ) of the sample in the pellet is equal to:

$$449 \quad d = \frac{m}{\pi r^2 * \rho} \quad \text{Eq. 3}$$

450 in which  $m$  is the mass of the haze analogues in the pellet,  $r$  is the pellet radius (6.5 mm),  
 451 and  $\rho$  is the measured density of the haze analogues.

452 Therefore, the absorption coefficient ( $\alpha$ ) can be calculated:

$$453 \quad \alpha(\nu) = \ln \frac{I_0}{I} / d \quad \text{Eq. 4}$$

454 Then, the extinction coefficient ( $k$ ) can be determined:

$$455 \quad k(\nu) = \frac{\alpha(\nu)}{4\pi\nu} = \frac{1}{4\pi\nu d} \ln \frac{I_0}{I} \quad \text{Eq. 5}$$

456 For each sample, we obtain three sets of the extinction coefficients ( $k$ ) independently by  
 457 measuring the transmittance of three pellets with different effective thicknesses. The final  
 458  $k$  values are determined by averaging the three sets of  $k$  values.

459 With the determined extinction coefficient ( $k$ ), we can calculate the real refractive index  
 460 ( $n$ ) based on the subtractive Kramers-Kronig (SKK) relation<sup>61-63</sup> between  $n$  and  $k$  as in Eq.  
 461 6:

$$462 \quad n(\nu) = n_0 + \frac{2(\nu^2 - \nu_0^2)}{\pi} P \int_0^\infty \frac{\nu' k(\nu')}{(\nu'^2 - \nu^2)(\nu'^2 - \nu_0^2)} d\nu' \quad \text{Eq. 6}$$

463 where  $\nu$  is wavenumber ( $\text{cm}^{-1}$ ),  $n_0$  is the real refraction index at  $\nu_0$ , and  $P$  indicates the  
 464 Cauchy principal value. The principal  $P$  value is integrated for the entire wavelength range;  
 465 we use the  $k$  values calculated (Eq. 5) from 350 to 25000  $\text{cm}^{-1}$  in the integrand and assume  
 466 constant  $k$  values for wavelengths beyond the measured range. The assumption is generally  
 467 valid unless there are large local absorption peaks outside the measurement range. Using  
 468 an anchor point ( $n_0$ ) can reduce the uncertainty for the numerical integration as  
 469 demonstrated in previous studies.<sup>62,63</sup> However, the real refraction index of the haze  
 470 analogues is initially unknown. From the interference fringes on the reflectance spectra of  
 471 the film samples at two angles ( $15^\circ$  and  $45^\circ$ ) of incidence (see Extended Data Fig. 1), we  
 472 determine  $n$  values from 0.4 to 1.1  $\mu\text{m}$  using Eq. 7.

$$473 \quad n_0 = \sqrt{\frac{\sin^2\theta_1 * (\Delta\nu_1)^2 - \sin^2\theta_2 * (\Delta\nu_2)^2}{(\Delta\nu_1)^2 - (\Delta\nu_2)^2}} \quad \text{Eq. 7}$$

474 where  $\theta_1$  and  $\theta_2$  are the two different angles of incidence in our reflectance measurements  
 475 of the film samples ( $15^\circ$  and  $45^\circ$ ), and  $\Delta\nu_1$  and  $\Delta\nu_2$  are the average fringe spacing in the  
 476 reflectance spectrum at respective incidence angle ( $\theta_1$  and  $\theta_2$ ). The  $n_0$  value determined at  
 477 0.625  $\mu\text{m}$  ( $\nu_0=16000 \text{ cm}^{-1}$ ) is used as the anchor point for the numerical integration in Eq.  
 478 6 to calculate the real refractive index ( $n$ ) from 0.4 to 28.6  $\mu\text{m}$  using the SKK relation. We  
 479 used different anchor points ( $n_0$  at different wavelengths) for the calculation and found that  
 480 the yielded  $n$  values are very close (the difference is less than 2%).

#### 481 *4.5 Atmospheric Transmission Spectra Simulation*

482 We use the open-source aerosol modeling code *Virga*<sup>48</sup> and the open-source radiative  
 483 transfer suite *PICASO*<sup>49</sup> to compute haze Mie properties and implement them into a  
 484 synthetic transmission spectrum model of an exoplanetary atmosphere. To generate hazy  
 485 atmospheric model spectra, we require a stellar radius, planetary mass and radius, pressure-  
 486 temperature profile, an atmospheric gas mixture, and haze profiles.

487 For the atmospheric composition of our synthetic spectra, we input the system parameters  
 488 of GJ 1214 b for the stellar and planetary mass and radius<sup>44</sup> and the same gas mixing ratios  
 489 as in the experimental run of the 400 K experiment shown in Fig. 1 into *PICASO*, using a

490 simple parametrized isothermal pressure-temperature profile at GJ 1214 b's  $T_{eq}$  of 600 K.  
491 Into *Virga*, we input the optical properties of our experimental hazes as their real and  
492 imaginary refractive indices  $n$  and  $k$ , binned down from the native measured resolution to  
493 a spacing of 10 nm, which we verify is still a fine enough sampling to capture the spectral  
494 features of the haze. We then used *Virga*'s optics functionality to generate wavelength-  
495 dependent Mie coefficients using haze particle size distributions as measured in the  
496 laboratory by He et al. 2018<sup>64</sup>. In addition to the optical properties and particle size of the  
497 experimental haze, we also need the vertical extent and mass loading of the haze layer  
498 within *Virga* in order to compute a synthetic atmospheric transmission spectrum with  
499 *PICASO*. *Virga* currently does not have the capability to self-consistently compute  
500 photochemical haze profiles. Therefore, for simplicity, we implement an isobaric haze  
501 layer from a pressure of 0.1 bar to 0.1  $\mu$ bar with haze mass loading of  $\sim 25$  particle/cm<sup>3</sup>.  
502 These values are consistent with the extent of haze in Titan's mesosphere<sup>65</sup> as well as self-  
503 consistent haze formation modeling that has been performed for exoplanets.<sup>66</sup> The resulting  
504 optical depth of the wavelength-dependent haze layers as a function of pressure and  
505 wavelength from our model set-up is shown in Extended Data Fig. 2. For the Titan-like  
506 haze cases, we use the same settings as described above for our 400 K and 300 K haze  
507 layers, but substitute the refractive indices of Khare et al. (1984)<sup>14</sup>.  
508 The purpose of our spectra simulation is to show the effect of the different haze optical  
509 properties on the atmospheric transmission spectra rather than perform a full model  
510 parameter study. Simplified assumptions were used in our simulation, such as atmospheric  
511 pressure-temperature profiles, the particle radius and mass distribution. Future model  
512 development and modeling studies, employing self-consistent pressure-temperature  
513 profiles including the radiative effects of the haze<sup>17</sup>, and full microphysics treatment of  
514 particle parameters<sup>67-71</sup>, are required to fully explore haze impacts on observations of  
515 exoplanet atmospheres.

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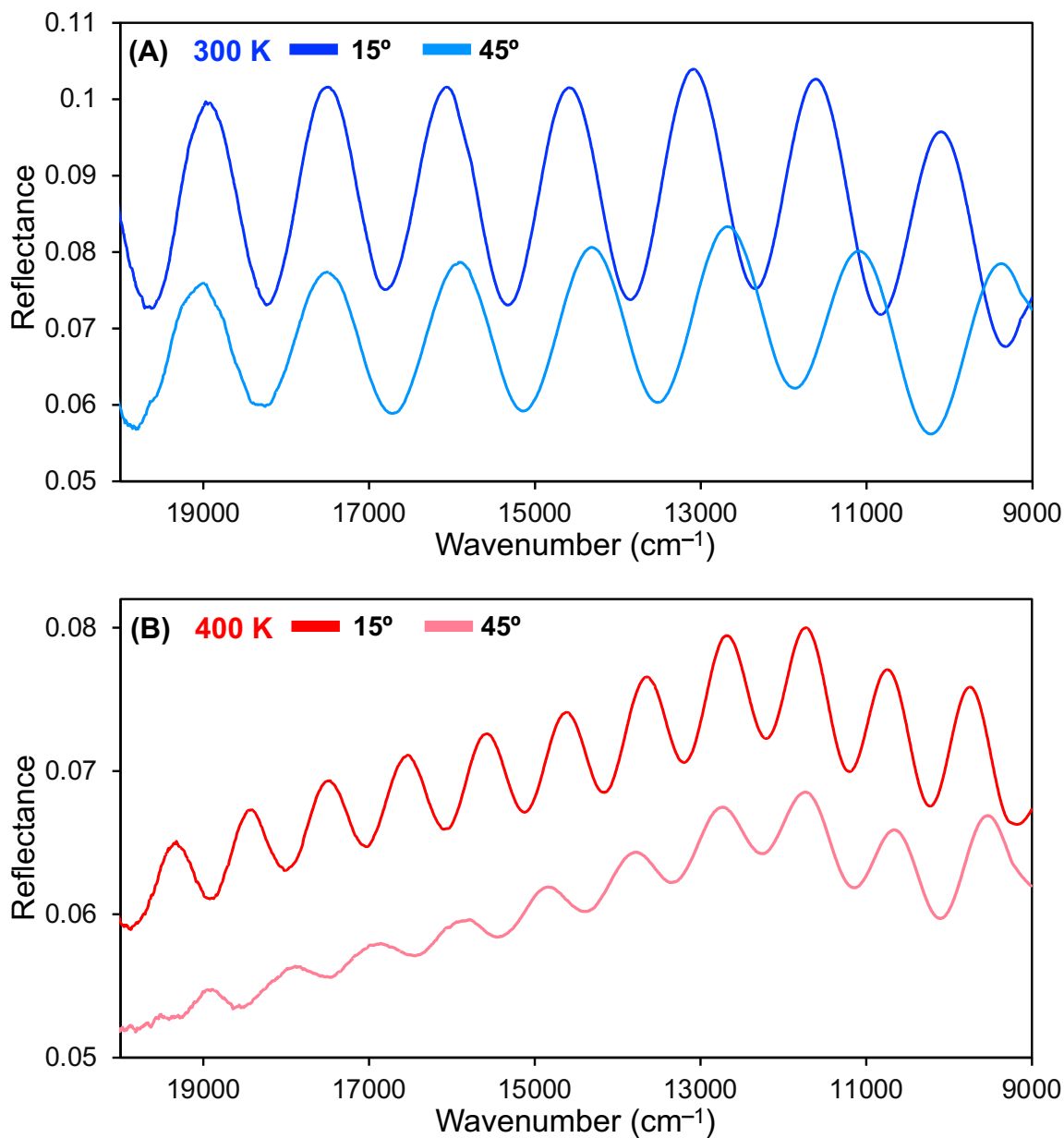
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## 681 **Acknowledgements**

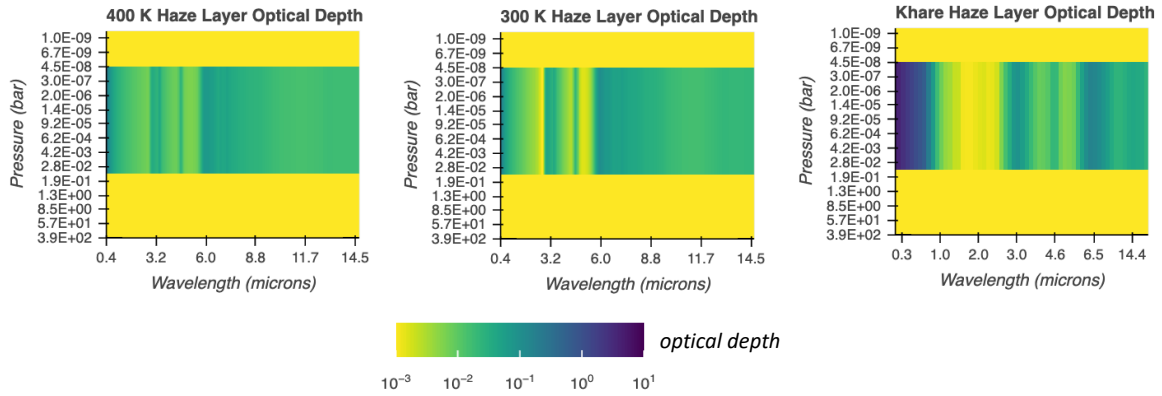
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## 685 **Author contributions**

686 C.H., M.R., S.E.M., S.M.H., N.K.L., M.S.M., and J.I.M. conceived the study. J.I.M.  
687 calculated the starting gas mixtures. C.H. carried out the experiments. C.H. and M.R.  
688 performed the optical measurements. S.E.M. and C.H. simulated the synthetic transmission  
689 spectra. C.H. conducted the data analysis and prepared the manuscript. All authors  
690 participated in discussions regarding interpretation of the results and edited the manuscript.

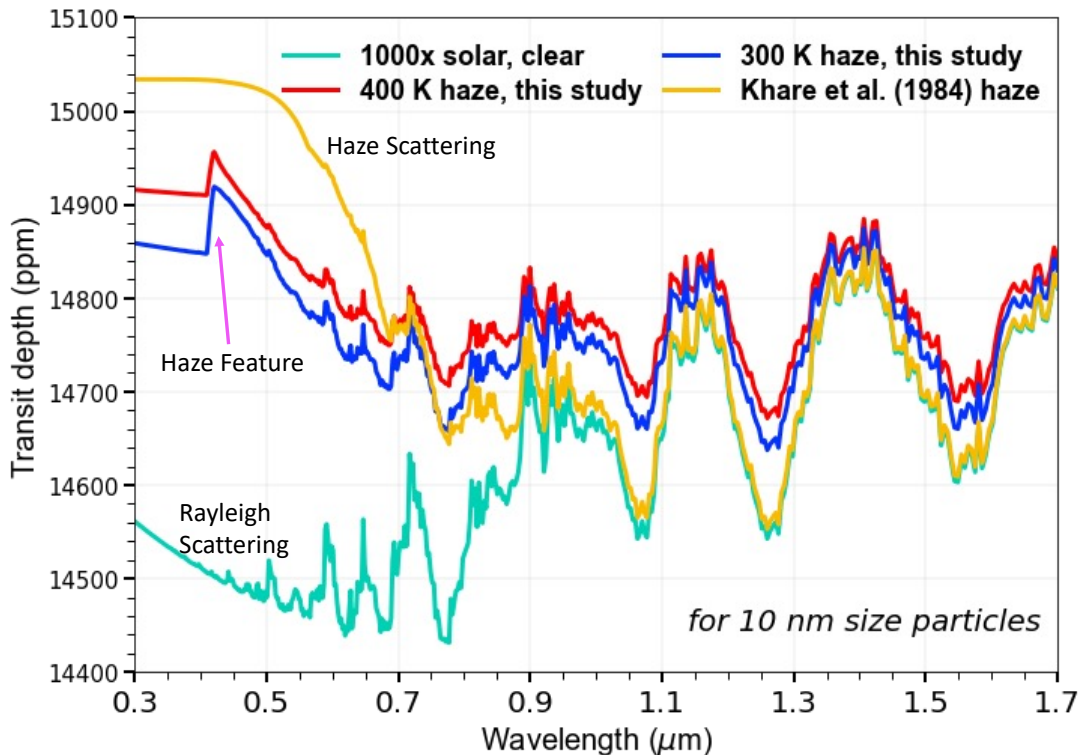


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 692 **Extended Data Fig. 1.** Reflectance spectra of two exoplanet haze analogues formed in  
 693 water-rich atmospheres at 300 K (A) and 400 K (B). Using Seagull Variable Angle  
 694 Reflection Accessory, we measure the reflectance of each haze analogue film at two  
 695 different angles of incidence, 15° (darker shade) and 45° (lighter shade). The spectra from  
 696 20000 to 9000 cm<sup>-1</sup> (0.5 to 1.1 μm) are shown here. From the interference fringes on the  
 697 reflectance spectra at two different angles, the  $n$  values of the samples at corresponding  
 698 wavelengths can be determined using Eq. 7.



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**Extended Data Fig. 2.** Haze slab profiles as a function of pressure and wavelength as implemented with the *Virga*-derived Mie coefficients. The color bar indicates the optical depth at each wavelength, with darker shading corresponding to higher optical depths. The haze layer extends from 0.1 bar to 0.1  $\mu$ bar with a haze particle radii distribution centered around 25 nm. Note that the Khare et al. (1984) data has wider wavelength coverage than shown, but at lower resolution than our measured optical properties.



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**Extended Data Fig. 3.** Model spectra of a water-rich atmosphere around a GJ 1214 b -like planet. We show the effect of our newly measured haze optical properties using small radii (10 nm) haze particles, focusing on the wavelength range accessible to *Hubble*. The method

711 and settings for generating the spectra here are the same as described in 4.5, except the  
712 haze particle radii (10 nm) and haze mass loading (4 particles/cm<sup>3</sup>). With sufficiently small  
713 particles, the large scattering slopes between different haze compositions are differentiable  
714 with *Hubble*'s ultraviolet/visible capabilities even if such hazes less strongly impact the  
715 NIR.