

Effect of water content and structural anisotropy on tensile mechanical properties of montmorillonite using molecular dynamics

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17 **Abstract:**

18 The mechanical behavior of montmorillonite is sensitive to water content and structural anisotropy. 19 In this study, the evolution of atomic structure and mechanical behavior of different hydrated 20 montmorillonite under tension has been investigated using Molecular Dynamics simulations. 21 External deformation has been applied to montmorillonite with a strain rate of 5×10^{-7} fs⁻¹ for 22 uniaxial tensile test, stretched along directions parallel (*x*- and *y*-direction) and perpendicular 23 (*z*-direction) to the clay sheets. The bond-breakage criterion based on the radial distribution function 24 of atom pairs and the evolution of broken bonds with tensile strain, were applied to explore the 25 deformation and failure mechanism of the montmorillonite-water system. Simulation results 26 indicated that the water content increased, causing a decrease in the ultimate tensile strength and 27 Young's modulus, as well as more and more plasticity for montmorillonite, especially in *z*-direction. 28 Moreover, the sequence of tensile mechanical strength was $y > x > z$. Stretching along the *x*- and 29 *y*-direction, fractures occurred in the silicon-oxygen tetrahedral and aluminum-oxygen octahedral 30 sheets, where the aluminum-oxygen bond was easier to be broken, and the normalized number of 31 total broken bonds in the *y*-direction was superior to the *x*-direction. Stretching along the *z*-direction, 32 only the hydrogen bonds in the interlaminar space were broken. In terms of three directions, the 33 anisotropy behavior was found in ultimate tensile strength, Young's modulus, residual tensile strength, 34 and the evolution of broken bonds.

35 **Keywords:** Montmorillonite; Molecular Dynamics; Water content; Tensile Mechanical properties; 36 Structural Anisotropy; Microstructure

37

38 **1. Introduction**

39 Clay minerals are widely distributed on the earth, and are the principal components of swelling 40 clay and claystone. The mechanical properties of clay minerals are always decisive to the 41 macroscopic properties of clay and claystone, which have been widely studied in geotechnical 42 engineering (Zheng et al., 2011; Liu et al., 2019; Yotsuji et al., 2021). As is well known, the 43 mechanical properties of clay minerals are very sensitive to water content and structural anisotropy 44 (Zhang et al., 2012; Liu et al., 2015; Zheng and Zaoui, 2018). So far, plenty of experimental 45 measurements (Mcneil and Grimsditch, 1993; Vanorio et al., 2003; Bobko and Ulm, 2008; Nazmul 46 et al., 2008; Hedan et al., 2015; Húlan et al., 2015) and theoretical derivation methods (Hueckel, 47 2002; Chen and Evans, 2006; Ortega et al., 2007) have been used to study the structural and 48 mechanical properties of claystone and clay minerals, such as the Brillouin scattering method 49 (Mcneil and Grimsditch, 1993), acoustic measurement (Vanorio et al., 2003), nano-indentation test 50 (Bobko and Ulm, 2008; Wu et al., 2020), in-situ test (Húlan et al., 2015), high energy synchrotron 51 X-ray, and the crystallographic Rietveld method (Wenk et al., 2007). These previous works showed 52 that the mechanical strength of clay minerals declined with increasing water content, and showed 53 obvious anisotropy in mechanical strength and failure form. On the other hand, the macroscopic 54 deformation of claystone is controlled by the local motion of clay particles at different scales (Zhu et 55 al., 2021). The macroscopic mechanical properties of claystone are strongly sensitive to water 56 saturation degree (Zhang et al., 2012). The microscopic mechanical behavior of clay particles also 57 depends on the hydration in interlayer spacing. Therefore, in terms of clays, especially swelling clay, 58 a fundamental molecular-level understanding of the clay-water interaction plays a significant role in 59 the evolution of its microstructure, hydromechanical behavior, and development of multiscale 60 approaches (Katti et al., 2007).

61 The first-principles and molecular dynamics (MD) simulation methods have been used to study the 62 nanoscale mechanical behaviors of clay minerals. The first-principles simulation method has been 63 used to study the microstructure (Stixrude and Peacor, 2002), bulk modulus, and elastic constant 64 (Sato, 2005; Militzer et al., 2010) of clay minerals. Although the first-principles method is powerful 65 and reliable in the simulation results, the number of atoms in the studied system is limited due to the 66 important computational burden. MD simulation method is a well-established tool to model up to 67 millions of atoms and observe the system's deformation and failure details in nanoscale (Yang et al., 68 2019; Wei et al., 2022). The MD simulation method has been widely performed to investigate the 69 elastic and structural properties of montmorillonite (Manevitch and Rutledge, 2004; Suter et al., 2007; 70 Ebrahimi et al., 2012), illite (Hantal et al., 2014), kaolinite (Benazzouz and Zaoui, 2012), muscovite 71 (Teich-McGoldrick et al., 2012), etc. Schmidt et al. (Schmidt et al., 2005) explored the evolution of 72 the mechanical response of Na-montmorillonite interlayer with increasing hydration. Mazo et al. 73 (Mazo et al., 2008) found that the elastic modulus of montmorillonite containing monolayer water 74 molecules is more significant than that with bilayer water molecules. Recently, the mechanical 75 behavior of dehydrated kaolinite under tensile and compression has been studied in previous works 76 (Zhang et al., 2021), and Yang et al. (Han et al., 2019; Yang et al., 2019) also considered the effect of 77 temperature and hydrated state on that of kaolinite. Furthermore, the effect of loading path and 78 orientation on mechanical behaviors for dehydrated Na-montmorillonite has been explored, while 79 ignoring hydration (Zhu et al., 2021). Katti et al. (Katti et al., 2007) reported the mechanical response 80 of the interlayer of dehydrated and hydrated montmorillonite using steered MD simulation, and 81 Zhang et al. (Zhang et al., 2012) explored the effect of water content and structural anisotropy on the 82 mechanical property of claystone via micro-indentation and mini-compression tests. The above 83 research has well advanced the knowledge of the structural and mechanical properties of clay 84 minerals. However, some interesting questions arise: How many bonds would be broken in the 85 montmorillonite system in the deformation and failure process? And which bond breaks first? How 86 does the structural anisotropy of montmorillonite and its hydrated state affect the evolution of 87 bond-breakage and tensile mechanical behavior?

88 To the end, MD simulations have been used to investigate the structure and mechanical properties 89 of hydrated montmorillonite under tension. In the CLAYFF force field, a bond-breakage criterion 90 based on radial distribution function (RDF) of atom pairs would be performed to describe the 91 evolution of bond-breakage during the failure process of montmorillonite, considering Al…O, Si…O, 92 and O…H bonds. Furthermore, the effect of water content on stress-strain curve, ultimate tensile 93 strength, Young's modulus, and residual tensile strength of montmorillonite are evaluated.

94

95 **2. Simulation details**

96 **2.1. System setup and force field**

97 The Wyoming-type montmorillonite is a 2:1 type clay mineral, with one aluminum-oxygen 98 octahedral sheet sandwiched between two silicon-oxygen tetrahedral sheets. The simulation model 99 was derived from Viani et al. (Viani et al., 2002), with the chemical formula of unit cell as 100 Na0.75(Si7.75Al0.25)(Al3.5Mg0.5)O20(OH)4·*n*H2O. The lattice substitution of montmorillonite was 101 performed based on the following principles: (1) Adjacent atoms cannot be replaced at the same time; 102 (2) One out of every eight Al^{3+} is replaced by Mg^{2+} in the octahedral sheet, and one out of every 32 103 Si^{4+} in the tetrahedral sheet is replaced by Al^{3+} , where the replacement positions are randomly 104 distributed on the upper and lower layers of the clay sheet. The basal spacing of montmorillonite was 105 adjusted from an initial 15.0 Å to 9.6, 11.0, 12.5, and 14.0 Å, corresponding to the system with water 106 contents of 0%, 10%, 20%, and 30% weight percent, respectively. Water molecules were distributed 107 randomly in the interlayer space of montmorillonite. Thereafter, the supercells of montmorillonite are 108 composed of 128 unit cells $(8a \times 4b \times 4c)$, as shown in Figure 1. Moreover, the water molecules in 109 this system were represented by the rigid SPC/E water model (Berendsen et al., 1987), which was 110 adopted to describe interactions among hydrated montmorillonite.

112 Figure 1. Modeling process of Na-montmorillonite system.

113

111

114 CLAYFF (Cygan et al., 2004) force field has been widely used for clay minerals in MD simulation 115 (Seppälä et al., 2016; Zheng and Zaoui, 2018; Wei et al., 2021; Al-Zaoari et al., 2022), which was 116 applied in this work. The non-bonded and bonded interaction parameters of the CLAYFF force field 117 were derived from the literature (Cygan et al., 2004), where Lennard-Jones parameters of 118 montmorillonite-water system are shown in Table 1. The total potential energy (*E*_{total}) of the system 119 could be calculated by equation (1).

120
$$
E_{\text{total}} = E_{\text{bond stretch}} + E_{\text{Coulomb}} + E_{\text{VDW}}
$$

121
$$
= k_1 (r_{ij} - r_0)^2 + \frac{e^2 q_i q_j}{4 \pi \epsilon_0 r_{ij}} + 4 \epsilon_{ij} \left[\left(\frac{\sigma_{ij}}{r_{ij}} \right)^{12} - \left(\frac{\sigma_{ij}}{r_{ij}} \right)^6 \right]
$$
(1)

122 where k_1 is the force constants, and r_0 represents the equilibrium bond length; r_{ij} is the distance between atoms *i* and *j*; q_i and q_j are the charges of atoms *i* and *j*, respectively; ε_0 is the dielectric constant; *σ* and *ε* are the size and energy parameters, respectively. Moreover, σ_{ij} and ϵ_{ij} can be obtained by Mixing Lorentz-Berthelot's law (Frenkel and Smit, 2001), as shown in equations (2) and 125 126 (3).

$$
\sigma_{ij} = \frac{\sigma_i + \sigma_j}{2} \tag{2}
$$

$$
\mathcal{E}_{ij} = \sqrt{\mathcal{E}_i \mathcal{E}_j} \tag{3}
$$

129 Table 1. Lennard-Jones parameters of montmorillonite-water system applied in this work.

130 a: (Berendsen et al., 1987); b: (Cygan et al., 2004); c: (Smith and Dang, 1994).

131

132 **2.2. Relaxation and simulation details**

133 MD simulations were carried out using LAMMPS (Plimpton, 1995) code. Three-dimensional 134 periodic boundary condition was applied in the simulations. The van der Waals forces were 135 calculated using the Lennard-Jones potential model with the cut-off radius set as 10 Å. Verlet 136 algorithm (Frenkel et al., 1997) was applied to integrate the motion equations. Long-range 137 electrostatic interactions with a cut-off radius of 8.5 Å were calculated using the Ewald sum method 138 (Darden et al., 1993). The conjugate gradient method was used to do the energy minimization for 139 geometry optimization, where the convergence tolerance was controlled at the energy of 1.0e-6 140 kcal/mol, the force of 1.0e-8 kcal/(mol•Å), and the maximum number of iteration steps of 1000.

141 To obtain a sufficient equilibrium configuration, the montmorillonite models were equilibrated 142 with the micro-canonical (NVE) ensemble for 200 ps, followed by the canonical (NVT) ensemble for 143 200 ps, and then with the isothermal-isobaric (NPT) ensemble for 600 ps, respectively. The 144 Nose-Hoover anisotropic thermostat and barostat were used to control temperature and pressure, with 145 a temperature of 298 K and pressure of 1.0 atm. The time step of the whole simulation was set to 1.0 146 fs. Finally, the equilibrium systems were carried out to uniaxial tensile test along the *x*-, *y*- and *z*-147 directions (corresponding to [1 0 0], [0 1 0], and [0 0 1] crystal orientation), respectively. The 148 uniaxial tensile test was achieved by changing the dimension of the simulation box of 149 montmorillonite supercell, called as strain-controlled loading approach. While stretching along the 150 *x*-direction, a certain strain rate was applied in this direction, with the temperature 298 K and the 151 pressure 1 atm in the other two directions during the whole simulation. Stretching along the *y*- and 152 *z*-directions was set in the same way.

153 To evaluate the effect of strain rate on the tensile mechanical properties of montmorillonite, the 154 dehydrated montmorillonite was deformed at different strain rates $(10 \times 10^{-7} \text{ fs}^{-1}, 5 \times 10^{-7} \text{ fs}^{-1}, 1 \times 10^{-7} \text{ fs}^{-1})$ 155 fs⁻¹, and 5×10^{-8} fs⁻¹). The evolution of stress with tensile strain along the *x*- and *z*-directions is 156 illustrated in Figure S1 (it could be found in Supplementary Material). The decrease of critical strain 157 and maximum stress with decreasing strain rate is found especially for stretching along the 158 *x*-direction, but the difference between the cases of 5×10^{-8} fs⁻¹ and 5×10^{-7} fs⁻¹ was rather small. 159 Moreover, there is almost no difference in all cases for tensile along the *z-*direction. Therefore, the 160 tensile simulation ten times slower will not greatly change the simulation results, which agrees well 161 with the results of the literature (Yang et al., 2019; Zhang et al., 2021). To ensure reliability and save 162 the cost of computation time, the tensile strain rate of 5×10^{-7} fs⁻¹ was selected to carry out the 163 uniaxial tensile test for 600 ps to obtain a total strain of 0.3. In the tensile process, the relationship 164 between stress and strain, as well as the trajectory file of deformation and failure processes, were 165 obtained.

166 **3. Results and discussions**

167 **3.1. The atomic structure before loading**

168 The basal spacing of different hydrated montmorillonite increases with water content, as shown in 169 Table 2. The simulation results of montmorillonite in this work agree well with previous results. It 170 illustrates the accuracy of the models and the applied force field.

Table 2**.** The basal spacing of different hydrated montmorillonite. 171

	Water content $(\%)$	Results in this work (\AA)	Other simulation results (\AA)	Experiments (\AA)
		9.616	$9.51 \sim 10.01$ ^{c,d}	9.55 e; $9.60 - 9.80$ ^f
	10	11.147	$11.04 \sim 11.94$ ^d	$11.91 \sim 12.35$ ^f
	20	12.630	$12.08g$; 12.3 ^h ; 12.64 ⁱ ; 12.32 ^j ; 12.2 ^k ; 12.0 ¹ ; 12.16 ^a ; 12.55 ^d	$12.0 \sim 12.6$ e,m,n,b
	30	14.090	$13.09.14.75$ ^d	
172			a- (Zheng and Zaoui, 2011); b- (Keren R, 1975); c- (Kuang et al., 2013); d- (Chávez-Páez et al.,	

173 2001a); e- (Berend et al., 1995); f- (Fu et al., 1990); g- (Chang et al., 1995); h- (Marry et al., 2002); i-174 (Kosakowski et al., 2008); j- (Boek et al., 1995); k- (Chávez-Páez et al., 2001b); l- (Karaborni et al., 175 1996); m- (Porion et al., 2007); n- (Mooney et al., 1952).

176

177 The RDF and coordination number (CN) for atoms *j* around *i* were calculated respectively by:

 $g(r)_{ij} = \frac{n(r)_j}{\rho_i v}$ $rac{n(r)_j}{\rho_j v} \approx \frac{n(r)_j}{4 \pi \rho_j r^2}$ 178 $g(r)_{ij} = \frac{n(r)_j}{\rho_j v} \approx \frac{n(r)_j}{4\pi \rho_j r^2 \Delta r}$ (4)

179
$$
CN_{ij} = 4\pi\rho_j r^2 \int_0^{r_{ij}} g(r) r^2 dr
$$
 (5)

180 where $n(r)$ is the average number of atom *j* around *i* in a spherical shell of radius *r* to $r + \Delta r$. ρ_i is the 181 number density of atom *j*. *r* is the distance between atoms *i* and *j*.

182 To understand the effect of water content on the atomic structure of montmorillonite sheet, the 183 RDFs of Al-Ob, Ob-Hw (Figure 2), Al-Oh, and Si-Ob (Figure S2) of different hydrated 184 montmorillonite were illustrated, where the $g(r)$ is the number of the atom pairs between given 185 ranges of separation *r*. The intensity of first peak in these RDFs of Al-Ob, Al-Oh, and Si-Ob in 186 montmorillonite are positively correlated with water content, where the position of the first peak 187 remains the same, which are 1.95 Å for Al-Ob and Al-Oh, 1.55 Å for Si-Ob. It indicates that the 188 increasing water content could slightly rise the attraction between Al-Ob, Al-Oh, and Si-Ob atom 189 pairs in clay layer structure. Figure 2(b) shows the radius of first peak in RDF of Ob-Hw declined 190 with increasing water content, which may indicate that partial water molecules in the interlayer space 191 are closer to the surface oxygen (Ob) of montmorillonite at a higher hydrated state.

Figure 2. Radial distribution function (RDF) of (a) Al-Ob; (b) Ob-Hw, in montmorillonite system. 194

195

 To understand the coordination environments of montmorillonite-water system, the CNs of 196 197 different atom pairs were calculated by integrating the RDF curves from origin to the first minimum based on equation (5). As shown in Table S1 and S2, the CNs of total Si-O and Al-O atom pairs for 198 all hydrated cases are always equal to 4.0 and 6.0, respectively. This is consistent with the 199 coordination number of Si-O and Al-O in silicon-oxygen tetrahedral and aluminum-oxygen 200 201 octahedral sheets of clay. As shown in Table S3, at $0\% \sim 20\%$ hydrated cases, the CNs of Ob-Hw and 202 Na-Ow rise with increasing water content, but that of Na-Ob decreases, which agrees well with the previous simulations for NH4-montmorillonite (Peng et al., 2020) and Na-smectite (Zhang et al., 203 204 2014). With increasing water content, the Clay-Water $(C...W)$ hydrogen bonds (HBonds) could be 205 constantly generated by the hydrogen atoms (Hw) of interlayer water molecules and oxygen atoms

206 (Ob) of clay surface; and the interaction between interlayer cation $Na⁺$ and clay surface could 207 constantly reduce, and the partial Na⁺ cations left clay surface and diffused into water. On the other hand, compared the 20% hydrated case with the 30% hydrated one, the CNs of Ob-Hw, Na-Ob, and 208 209 Na-Ow were almost the same, indicating that the coordination of cation Na⁺ by atoms Ow and Ob, as well as Ob and Hw were saturated and stable at 20% hydrated case (see Figure S3). 210

211 A bond-breakage criterion based on the RDFs of atom pairs was proposed by Yang et al. (Yang et 212 al., 2019), that the first local minimum after the first peak of RDF curves means that the chance of 213 finding atoms pairs with the corresponding separation distance is approximately zero. Thus, it could 214 assume that the bonds would be broken once the increasing distance between atom pairs under strain 215 is greater than this distance. As defined in CLAYFF force field, there are no bond between atoms 216 except hydrogen and oxygen in hydroxyl and water molecules, in which exist the covalent bonds. As 217 a difference, virtual bonds were defined here with the symbol "…" in Al…O, Al…Oh, Si…O, and 218 O…H bonds in clay system, with their bond energy equal to the non-bonded interactions. The RDFs 219 of Al-Ob, Al-Oh, and Si-Ob bonds for montmorillonite are shown in Figure 2 and Figure S2, and the 220 cut-off distances of each virtual bond breakage are listed in Table 3. Additionally, the hydrated 221 montmorillonite contains two types of HBonds: (1) the C…W HBonds (Clay-Water) is formed 222 between water molecules and the acceptor oxygen of montmorillonite surface; (2) the W…W 223 HBonds (Water-Water) is formed between water molecules. It is worth noting that the C…W 224 HBonds play a more important role in the mechanical properties of montmorillonite than W…W 225 HBonds, so the C…W HBonds would be considered in this study. The criterion for forming the 226 HBonds is the cut-off distance O-H, as 2.4 Å , as proposed by Martí (Martí, 1999).

227

 228 Table 3. The cut-off distance of individual bond types for Na-montmorillonite, compared with that of

kaolinite in previous work (Yang et al., 2019). 229

Bond types Cut-off distance (A)

230

231 **3.2. Influences of structural anisotropy and water content**

232 **3.2.1. Stress-strain relationship**

233 The maximum tensile stress of montmorillonite along the *x*-, *y*-, and *z*-directions decline as the 234 water content increases, as shown in Figure 3. The most critical decrease of the maximum tensile 235 stress was found from the dehydrated case to 10% hydrated cases of montmorillonite. The existence 236 of water greatly reduced the tensile mechanical strength of montmorillonite. Moreover, to further 237 analyze the deformation process of each stage in the stress-strain curve, point O referred to the initial 238 system without strain, point D for the final system with a strain of 0.3; and points A, B, and C were 239 the tensile strain states at the maximum stress, under failure, and residue respectively (Figure 3(a)). 240 Due to the fluctuations of stress while tensile along the *z*-direction, Savitzky-Golay method 241 (Gallagher, 2020) was performed to smooth the curves, as displayed in Figure 3(c).

242 As shown in Figure 3(a), during the stretching process along the *x*- and *y*-direction for all systems, 243 the stress increases almost linearly with strain from the initial state (point O) to maximum (point A). 244 Thereafter, a sharp drop occurs from maximum (point A) to relative lower value (point B), which 245 means a sudden and brittle failure. However, while stretching along the *z*-direction for hydrated 246 montmorillonite, the stress decreased relatively flat, which signified a good plastic (Figure 3(c)). 247 This indicates that the increase of water content for montmorillonite improves its plastic, which 248 agrees well with the experimental results of Zhang et al. (Zhang et al., 2012). Besides, the evolutions 249 of stress with strain were similar to the cases of dehydrated montmorillonite for stretching along the 250 *x*-, *y*, and *z*-directions, but the maximum stress along *z*-direction are much smaller than the others.

Figure 3. Tensile stress-strain curves of Na-montmorillonite with different water contents at stretched 255 256 along different directions: (a) *x*-direction; (b) *y*-direction; (c) *z*-direction, where fitting curves are obtained by Savitzky-Golay method to perform smooth. 257

258

259 **3.2.2. Broken bonds during the failure process**

260 The normalization number of broken bonds (*r*bond) of hydrated montmorillonite could be 261 calculated by the following equation.

$$
r_{\text{bond}} = \frac{N - N_0}{N_0} \times 100\% = \frac{\Delta N}{N_0} \times 100\% \tag{6}
$$

263 where *N* and *N*0 are the numbers of the certain bond at the state corresponding to a certain strain and 264 the initial strain 0, respectively.

265 Since the deformation behavior of different hydrated montmorillonite were very similar, the 266 montmorillonite with a water content of 10% was taken thereafter for analysis. The deformation of 267 hydrated montmorillonite with water content of 10% under tensile along the *x*-direction (Figure 4(a)) 268 could be divided into four stages: (1) at OA stage, no bond in the crystal structure was broken when 269 the tensile strain was less than 0.05. Thereafter, just very few bonds were broken as the tensile strain 270 increased to ultimate tensile strain (Point A). (2) At AB stage, the largest number of bonds were 271 broken, and the crack grew rapidly. The clay sheets were seriously destroyed. (3) At BC stage, the 272 number of broken bonds for Al…Ob and Al…Oh rose slowly with tensile strain, accompanied by 273 further propagation of the crack and hole. However, the Si…Ob bonds were not further damaged. (4) 274 At CD stage, the number of broken bonds for Al…Ob and Al…Oh fluctuated around a certain value, 275 and the one of Si…Ob remained unchanged.

276 As shown in Figure 4(a) and (b), the number of broken bonds for Al…Ob is greater than that of 277 Al…Oh as stretching along the *x*-direction, but it is just the opposite as stretching along the 278 *y*-direction. In both cases, the broken bonds of Si…Ob remained the weakest, compared with the 279 other two bond types. The normalized number of broken bonds of Al…Ob and Si…Ob was 280 approximately the same when stretching along the *x*- and *y*-direction. However, the normalized 281 number of broken bonds of Al…Oh in the *y*-direction was much greater than that of *x*-direction, 282 which was due to the different arrange for Al…Oh bond along the *x*- and *y*-directions (Figure S4), 283 shown as the more total broken bonds in the *y*-direction. Therefore, more energy was required to 284 overcome the interaction of these bonds as stretching along the *y*-direction, so the tensile mechanical 285 strength of *y*-direction for montmorillonite was superior to that of *x*-direction. This finding also could 286 explain why the tensile mechanical strength of *y*-direction is greater than the one of *x*-direction for 287 kaolinite (Han et al., 2019; Yang et al., 2019; Zhang et al., 2021).

288 As shown in Figure 4(c), the layers of montmorillonite system are delaminated at the tensile 289 process along the *z*-direction, and the tetrahedra and octahedra sheets remain unchanged, which 290 agrees well with the results of the literature (Zhang et al., 2021). The maximum stress in the 291 stress-strain curve of *z*-direction was much smaller than that of *x*-direction and *y*-direction (Figure 3), 292 due to the weaker interactions between layers. Moreover, (1) at OA stage, with increasing tensile 293 strain, the number of broken bonds for C…W HBonds rose significantly. It was due to the primary 294 interaction between layers by C…W HBonds. With the increasing tensile strain, the distance between 295 the surface of montmorillonite and interlayer water increased, causing the broken bonds for C…W 296 Hbonds. (2) At AB and BC stage, the number of broken bonds for C…W HBonds declined rapidly. 297 The layers of montmorillonite system was delaminated as two portions, the water molecules around 298 the division layer of clay sheet moved gradually close to the surface of montmorillonite, resulting in 299 the forming bonds for C…W HBonds. (3) At CD stage, the number of broken bonds for C…W 300 HBonds fluctuated around a certain value, indicating that the montmorillonite system was separated 301 into two portions, and then into an equalized state.

302

303 Figure 4. Evolution of structure deformation and the normalization number of broken bonds of 304 hydrated montmorillonite with water content of 10% under different tensile strains: (a), (b) The 305 Al…Ob, Al…Oh, and Si…Ob bonds are considered, stretching along *x*-direction and *y*-direction, 306 respectively; (c) The C…W HBonds are analyzed, stretching along *z*-direction. The color code marks 307 the bonds with less deformation as blue and those with more deformation as red.

308

309 To study the effect of water content and structural anisotropy on the evolution of broken bonds, the 310 relationship between the total normalization number of broken bonds and tensile strains for different 311 hydrated montmorillonite are shown in Figure 5. As for dehydrated montmorillonite (water content 312 0%), no C…W HBond was contained, so the 0% curve was not shown in Figure 5(c). As shown in 313 Figure 5(b), as stretching along the *y*-direction, the normalized number of total broken bonds under 314 different hydrated states are relatively closed. However, in Figure 5(a) and (c), the lower the water 315 content, the more the broken bonds. On the other hand, the normalized number of total broken bonds 316 in the *y*-direction was superior to the *x*-direction in all cases. This evolution of broken bonds further 317 shows the structural anisotropy of montmorillonite, especially in the parallel (*x*- and *y*-) and 318 perpendicular (*z*-) directions.

montmorillonite with tensile strains: (a) total bonds for *x*-direction, and (b) *y*-direction; (c) C…W 324 HBonds for *z*-direction. 325

326

327 **3.2.3. Tensile mechanical parameters**

328 The ultimate tensile strength and ultimate tensile strain are equal to the maximum stress and its 329 corresponding strain on the stress-strain curve, and the residual tensile strength is the stress at a strain 330 of 0.3 (Figure 3). Young's modulus could be obtained by taking the slope of the linear elastic regime 331 with a small strain range, where the selected strain range is $0 \sim 0.01$. Based on the ultimate tensile 332 strength of three directions, the overall ultimate tensile strength (σ_{toveral}) is proposed to illustrate the 333 overall tensile mechanical, computed by equation (7).

334
$$
\sigma_{\text{toverall}} = \sqrt{\frac{\sigma_{ix}^2 + \sigma_{iy}^2 + \sigma_{iz}^2}{3}}
$$
 (7)

335 Where σ_{tx} , σ_{ty} , and σ_{tz} correspond to the ultimate tensile strength of the *x*-, *y*-, and *z*-directions, 336 respectively. Similarly, the overall Young's modulus (*E*overall) could be determined.

337 Young's modulus of dehydrated montmorillonite and previous results, have been collected in Table 338 4, reflecting the binding strength among atoms or molecules in nanoscale. The simulation results in 339 this study agreed well with the results of the previous literature (Ledbetter et al., 2001; Chen and 340 Evans, 2006; Nazmul et al., 2008; Zheng and Zaoui, 2018), which could validate the accuracy of 341 applied force field and established model. There is a certain difference in the value of Young's 342 modulus between this work and previous studies, likely due to the difference in ion substitution site, 343 the applied force field, and selecting ensemble for montmorillonite during whole simulations. 344 Furthermore, the sequence of Young's modulus for dehydrated montmorillonite was $y > x > z$, which 345 is consistent with results found by Zheng and Zaoui (Zheng and Zaoui, 2018), for the cases before 346 swelling and after shrinking of montmorillonite.

347

Table 4. Young's modulus (GPa) of dehydrated Na-montmorillonite compared to previous 348

experimental and simulation results. 349

351 al., 2005); e- (Zheng and Zaoui, 2018); f- (Zartman et al., 2010); g- (Carrier et al., 2014); h-352 (Ledbetter et al., 2001); i- (Prasad, 2001).

353

354 The intensity attenuation coefficient (*D*s) of the sample is used to illustrate its post-peak strength 355 decay behavior, proposed by Peng et al. (Peng et al., 2015), which can be expressed as:

$$
D_s = \frac{\sigma_t - \sigma_r}{\sigma_t} \tag{8}
$$

357 where σ_t represents the ultimate tensile strength (corresponding to point A in stress-strain curve of 358 Figure 3(a)), and σ_r represents the residual tensile strength (point D in Figure 3(a)).

359 As shown in Figure 6 and Table 5, the ultimate tensile strength and Young's modulus of 360 montmorillonite decrease almost linearly with increasing water content. It indicates that the existence 361 of water reduced the capability to resist deformation and failure of montmorillonite. For all cases, the 362 ultimate tensile strength, Young's modulus, and residual tensile strength along the *y*-direction are 363 slightly superior to those along the *x*-direction, and significantly superior to those along the 364 *z*-direction. At 0% ~ 20% hydrated cases, the intensity attenuation coefficient along *z*-direction were 365 close to 1.0, and greater than that of *x*- and *y*-direction (Table 5), illustrating that the tensile

366 mechanical strength of *z*-direction decayed the most. Moreover, there was still residual tensile 367 strength for *y*- and *x*- directions at the strain of 0.3. These findings indicated that the weakest surface 368 of montmorillonite is (0 0 1) crystal plane, which is consistent with some previous work on 369 montmorillonite (Zheng and Zaoui, 2018) and kaolinite (Zhang et al., 2021). It further provides 370 evidence for the anisotropy of mechanical properties of montmorillonite.

373 Figure 6. Evolution of tensile mechanical parameters with water content for montmorillonite, 374 corresponding to *x*-, *y*-, *z*-, and overall directions: (a) ultimate tensile strength; (b) Young's modulus. 375

Table 5. Residual tensile strength and intensity attenuation coefficient of different hydrated 376

montmorillonite. 377

378

379 **4. Conclusions**

380 Molecular dynamics (MD) simulations have been used to study the effect of water content and 381 structural anisotropy on atomic structure and tensile mechanical properties of montmorillonite in 382 three directions. It is noted that the CLAYFF force field cannot present the bond breaking and 383 rearrangement, so the bond-breakage criterion based on radial distribution function (RDF) of atom 384 pairs was proposed to describe the deformation and failure process of montmorillonite.

385 The Young's modulus of dehydrated montmorillonite and basal spacing of all hydrated 386 montmorillonite obtained in this work were consistent with previous simulation and experiment 387 results. The increasing water content augmented the first peak in RDFs of Al-Ob, Al-Oh, and Si-Ob, 388 but without effect on the position of these peaks, illustrating that the higher water content could 389 slightly rise the attraction between Al-Ob, Al-Oh, and Si-Ob atom pairs in clay layer structure. 390 Moreover, stretching along the *x*- and *y*-directions, the Al…Ob and Al…Oh bonds were broken 391 earlier and significant higher quantities than the Si…Ob bonds, indicating that the silicon-oxygen 392 tetrahedral sheets was stronger than aluminum-oxygen octahedral sheet. The normalized number of 393 total broken bonds in the *y*-direction was superior to the *x*-direction. On the other hand, stretching 394 along *z*-direction, the layers of the system were delaminated, where Clay-Water hydrogen bonds 395 (HBonds) was broken for the hydrated system, while the tetrahedral and octahedral sheets remained 396 intact. Furthermore, the higher the water content, the greater the plasticity of the montmorillonite 397 system. Additionally, based on Young's modulus, ultimate tensile strength, and residual tensile 398 strength of three directions, the sequence of tensile mechanical strength for montmorillonite was *y* > 399 $x > z$, showing obvious anisotropy of montmorillonite. The increase of water molecules for 400 montmorillonite gradually reduced its tensile mechanical properties. This work could lay the 401 foundation for further research on the compression and shear mechanical properties of clay minerals. 402 Especially, the mechanical properties of clay are also sensitive to water under compression, which 403 will be further discussed in the future.

404

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Graphical Abstract

