



HAL
open science

Global dissipative martingale solutions to the variational wave equation with stochastic forcing

Billel Guelmame, Julien Vovelle

► **To cite this version:**

Billel Guelmame, Julien Vovelle. Global dissipative martingale solutions to the variational wave equation with stochastic forcing. 2023. hal-04282928v1

HAL Id: hal-04282928

<https://hal.science/hal-04282928v1>

Preprint submitted on 24 Nov 2023 (v1), last revised 27 Sep 2024 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons Attribution - NonCommercial - NoDerivatives 4.0 International License

Global dissipative martingale solutions to the variational wave equation with stochastic forcing

Billel Guelmame* and Julien Vovelle†

November 24, 2023

Abstract

We consider the variational wave equation in one-dimensional space with stochastic forcing by an additive noise. Blow-up of local smooth solutions is established, and global existence is proved in the class of weak martingale solutions.

AMS Classification : 35R60, 60H15, 35L70, 35A01

Key words : Stochastic PDEs, Hyperbolic equations, Wave equation, blow-up, Young measure, Skorokhod–Jakubowski’s representation theorem.

Contents

1	Introduction	2
2	The equations and main results	3
2.1	The stochastic variational wave equation	3
2.2	Local and global existence	5
3	Local-in-time regular solutions	7
3.1	Generalized system	8
3.2	Local regular solutions	9
3.3	Conservative form, integration along the characteristics, Itô’s formula	15
3.4	Energy estimate	18
3.5	Solutions with finite explosion time	18
3.5.1	Some preliminary estimates	18
3.5.2	Singularity formation	20
4	Global-in-time regular solutions to the approximate system	22
4.1	Energy estimates	23
4.2	Global solutions	27
4.3	L^{3^-} estimates	28
4.4	One-sided entropy estimates	29

*UMPA, CNRS, ENS de Lyon, Université De Lyon, billel.guelmame@ens-lyon.fr

†UMPA, CNRS, ENS de Lyon, julien.vovelle@ens-lyon.fr

5	First limit elements	31
5.1	Young measures	32
5.2	Quasi-Polish spaces and the Skorokhod–Jakubowski theorem	33
5.3	Compactness	34
5.4	Limit equation with defect measure	36
5.4.1	Application of the Skorokhod–Jakubowski theorem	36
5.4.2	Convergence et identification of some deterministic elements	37
5.4.3	The stochastic integral	40
5.4.4	Limit equation with defect measure	41
6	Reduction of the Young measures	43
6.1	Compensated compactness	43
6.2	An evolution equation for the defect measure	45
6.3	Existence of weak martingale solutions	51

1 Introduction

We consider in this paper the variational wave equation with a stochastic forcing

$$du_t - c(u)(c(u)u_x)_x dt = \Phi dW, \quad (t, x) \in (0, T) \times \mathbb{T}, \quad (1.1)$$

where u_t denotes $\partial u / \partial t$ (*cf.* Remark 2.1 below), $\mathbb{T} = \mathbb{R}/\mathbb{Z}$ denotes the one-dimensional torus, and, over the filtered probability space $(\Omega, \mathcal{F}, \mathbb{P}, (\mathcal{F}_t)_{t \geq 0})$, W is a cylindrical Wiener process (see (2.1)). The equation (1.1) is the Euler–Lagrange equation associated to the Lagrangian

$$\mathcal{L} \stackrel{\text{def}}{=} \iint \left[\frac{1}{2} [u_t^2 - c(u)^2 u_x^2] dt + u \Phi dW(t) \right] dx.$$

Our aim is to study the well-posedness of (1.1) up to an explosion time. We prove then that regular solutions may blow-up in finite time. Finally we establish the existence of global-in-time weak martingale solutions to (1.1).

The deterministic variational wave equation is recovered taking $\Phi \equiv 0$. It appears in several physical contexts. For example, nematic liquid crystals [S89, HS91, GHZ97], long waves on a dipole chain [GHZ97, GI92, ZI92] and also in classical field theories and general relativity [GHZ97]. An asymptotic (and simpler) equation has been derived by Hunter and Saxton [HS91]

$$[u_t + u u_x]_x = \frac{1}{2} u_x^2. \quad (1.2)$$

Both the Hunter–Saxton equation (1.2) and the deterministic variational wave equation have been widely studied in the literature [HZ95a, HZ95b, ZZ98, ZZ00, ZZ01, ZZ03, ZZ05, BC05, BZ06, BZZ07, D11, BCZ15].

Considering the deterministic variational wave equation on the real line ($x \in \mathbb{R}$), the local (in time) well-posedness can be obtained using Kato’s theorem for the quasi-linear equations [K75]. Singularities may appear in finite time in the non-linear case $c'(\cdot) \not\equiv 0$ [GHZ96]. However, if the Riemann invariants $u_t \pm c(u)u_x$ are non-positive initially, then rarefactive solutions exist globally in time [ZZ98]. Global weak solutions to the variational wave equation are not unique, indeed, at least two types of solutions exist, dissipative and conservative ones. The conservative solutions were obtained in [BZ06] using an equivalent system in the Lagrangian coordinates. The total energy of the conservative solutions is equal to the initial energy for almost all $t > 0$. The uniqueness of the conservative solutions is established in [BCZ15]. Dissipative solutions are

obtained in [ZZ03, ZZ05] by studying an approximated system and then passing to the limit. The dissipative solutions satisfy an energy inequality and a one-sided entropy inequality. To the author’s knowledge, the uniqueness of the dissipative solutions remains an open problem. However, for the simpler Hunter–Saxton equation (1.2), the dissipative solutions are unique [D11].

Note that, in the stochastic case, the Hunter–Saxton equation (1.2) with noise has been studied recently in [HKP21].

We consider the problem (1.1) under study in the class of systems of stochastic non-linear hyperbolic equations. There are numerous works related to stochastic scalar conservation laws, see for instance [CP21] and references therein. In hyperbolic systems of conservation laws with stochastic forcing, which is our framework, the problematic and techniques are relatively different from the scalar case.

Let us first make some comments on our results and the way they are established. The problem (1.1) admits an equivalent formulation as a 2×2 system (see (2.6) below). The existence of local regular solutions to this 2×2 system is essentially a pathwise variation on the deterministic result and certainly classical in spirit, but we prefer to give the proof in full details. Global martingale solutions are obtained by the probabilistic compactness method. We adapt the method of P. Zhang and Y. Zheng, [ZZ01, ZZ03, ZZ05] (see also [Gue23] in the context of the Green–Naghdi equations with surface tension) for the deterministic treatment of the problem. If the stochastic aspects are not present in [ZZ03, ZZ05], our framework is also different, insofar as we consider periodic solutions and not solutions with certain localization properties. As a consequence, our approximated system involves some additional “correction terms” (see (2.11) below) to ensure the periodic character of the the solutions. One of the central technical tool to recover the desired equation at the limit are Young measures, as introduced by DiPerna in the context of systems of conservation laws, [DiP83]. Similar techniques have been used recently for the Hunter–Saxton equation (1.2), the Camassa–Holm equation and the Degasperis–Procesi equation with noise [CC24, HKP21, HKP23, GHKP22]. In addition, we exploit [BV19] (see also [FN08]), where Young measures for the study of the stochastic isentropic Euler equations play a central role too.

The paper is organized as follows. In Section 2 we introduce the kind of noise (white in time, colored in space) that we consider. We present the equivalent 2×2 system (2.6) and we state the main results of the paper. In Section 3, devoted to local-in-time solutions, we prove the existence of regular solutions up to an explosion time, and also give some criterion for blow-up (which amounts to the explosion of the Lipschitz norm of u) in finite time. Section 4 is devoted to the approximated system obtain by truncation and correction of (2.6) (see (2.11)). We prove that the approximate solutions exist globally in time, and that they satisfy some uniform estimates. In Section 5 we prove that the sequence of approximated solutions is tight (in a suitable space), and we use the Skorokhod–Jakubowski representation theorem to pass to the limit and to derive the limit equation with some defect measures. Finally, it is proved in Section 6 that the defect measures are trivial and that the limit is a global weak martingale solution to (1.1).

2 The equations and main results

2.1 The stochastic variational wave equation

Let \mathfrak{U} be an Hilbert space with an orthonormal basis $(g_k)_{k \geq 1}$ and let \mathfrak{U}_{-1} be another Hilbert space such that the injection $\mathfrak{U} \hookrightarrow \mathfrak{U}_{-1}$ is Hilbert–Schmidt. Let W be the cylindrical Wiener

process defined by

$$W(t) \stackrel{\text{def}}{=} \sum_{k \geq 1} g_k \beta_k(t), \quad t \geq 0, \quad (2.1)$$

where $(\beta_1(t), \beta_2(t), \dots)$ are independent one-dimensional Wiener processes, see Section 4.1.2 in [DPZ14]. Let $\Phi : \mathfrak{M} \rightarrow L^2(\mathbb{T})$ such that for any $k \geq 1$ we have $\sigma_k \stackrel{\text{def}}{=} \Phi g_k \in C(\mathbb{T})$ and

$$q_0 \stackrel{\text{def}}{=} \sum_k \|\sigma_k\|_{W^{1,\infty}(\mathbb{T})}^2 < \infty, \quad q(x) \stackrel{\text{def}}{=} \sum_k \sigma_k(x)^2. \quad (2.2)$$

By (2.2) and the injection $L^\infty(\mathbb{T}) \hookrightarrow L^2(\mathbb{T})$, the map Φ is Hilbert–Schmidt. Let us assume that $c \in C^\infty(\mathbb{R})$ satisfies

$$0 < c_1 \leq c(u) \leq c_2, \quad (2.3)$$

$$0 \leq c'(u) \leq c_3. \quad (2.4)$$

for some constants $c_1, c_2, c_3 \in (0, \infty)$. We consider the stochastic variational wave equation with additive noise in $(0, T) \times \mathbb{T}$

$$du_t - c(u) (c(u) u_x)_x dt = \Phi dW, \quad (2.5a)$$

$$u(0, \cdot) = u_0, \quad u_t(t=0, \cdot) = v_0. \quad (2.5b)$$

We can also consider the equivalent form

$$dR + c(u) R_x dt = \tilde{c}'(u) [R^2 - S^2] dt + \Phi dW, \quad (2.6a)$$

$$dS - c(u) S_x dt = \tilde{c}'(u) [S^2 - R^2] dt + \Phi dW, \quad (2.6b)$$

completed with the equation

$$u(t, x) = \mathcal{C}^{-1} \left\{ \mathcal{C} \left\{ \int_0^t \left(\frac{R+S}{2} \right) (s, 0) ds + u_0(0) \right\} + \int_0^x \frac{S-R}{2} (t, y) dy \right\}, \quad (2.7)$$

where

$$\mathcal{C}(r) = \int_0^r c(\sigma) d\sigma, \quad (2.8)$$

which expresses u as a non-local function of (R, S) . The system (2.6) is deduced from (2.5) by setting

$$R \stackrel{\text{def}}{=} u_t - c(u) u_x, \quad S \stackrel{\text{def}}{=} u_t + c(u) u_x, \quad \tilde{c}(u) \stackrel{\text{def}}{=} \frac{1}{4} \ln c(u). \quad (2.9)$$

The corresponding initial conditions for (2.6) are therefore

$$R(0, \cdot) = R_0 \stackrel{\text{def}}{=} v_0 - c(u_0) u_0', \quad S(0, \cdot) = S_0 \stackrel{\text{def}}{=} v_0 + c(u_0) u_0'. \quad (2.10)$$

Remark 2.1 (Notations). The subscript t , as in u_t , always denote the partial derivative with respect to t , and never the value at the given time t of a stochastic process X (the latter being simply denoted by $X(t)$).

We will establish in Theorem 2.2 the existence of local regular solutions to (2.6), which are strong in the probabilistic sense. It may happen that some of these solutions have a finite time of existence, cf. Theorem 2.3. Nevertheless, global-in-time solutions, which are weak in the probabilistic sense (martingale solutions), and defined in an L^2 -framework, are proved to

exist, see Theorem 2.5. These global-in-time solutions are obtained as limits of solutions to the following system

$$dR^\varepsilon + c(u^\varepsilon) R_x^\varepsilon dt = \check{c}'(u^\varepsilon) [(R^\varepsilon)^2 - (S^\varepsilon)^2 - \chi_\varepsilon(R^\varepsilon) + 2R^\varepsilon \Theta^\varepsilon] dt + \Phi^\varepsilon dW, \quad (2.11a)$$

$$dS^\varepsilon - c(u^\varepsilon) S_x^\varepsilon dt = \check{c}'(u^\varepsilon) [(S^\varepsilon)^2 - (R^\varepsilon)^2 - \chi_\varepsilon(S^\varepsilon) - 2S^\varepsilon \Theta^\varepsilon] dt + \Phi^\varepsilon dW, \quad (2.11b)$$

$$R^\varepsilon(0, \cdot) = R_0^\varepsilon \stackrel{\text{def}}{=} J_\varepsilon R_0, \quad S^\varepsilon(0, \cdot) = S_0^\varepsilon \stackrel{\text{def}}{=} J_\varepsilon S_0, \quad (2.11c)$$

coupled with the equation ($x \in [0, 1]$)

$$u^\varepsilon(t, x) = \mathcal{C}^{-1} \left\{ \mathcal{C} \left\{ \int_0^t \left(\frac{R^\varepsilon + S^\varepsilon}{2} \right) (s, 0) ds + u_0^\varepsilon(0) \right\} + \int_0^x \left[\frac{S^\varepsilon - R^\varepsilon}{2}(t, y) - \Theta^\varepsilon(t) \right] dy \right\}, \quad (2.12)$$

where J_ε is a Friedrichs mollifier, defined as the convolution operator $R \mapsto R * \rho_\varepsilon$, where (ρ_ε) is an approximation of the unit. In (2.11), the cut-off function χ_ε is defined by

$$\chi_\varepsilon(\xi) \stackrel{\text{def}}{=} \left(\xi - \frac{1}{\varepsilon} \right)^2 \mathbb{1}_{[\frac{1}{\varepsilon}, \infty)}(\xi) = \begin{cases} \left(\xi - \frac{1}{\varepsilon} \right)^2, & \xi \geq 1/\varepsilon, \\ 0, & \xi < 1/\varepsilon. \end{cases} \quad (2.13)$$

We have also introduced a ‘‘correction term’’

$$\Theta^\varepsilon(t) \stackrel{\text{def}}{=} \int_0^1 \frac{S^\varepsilon - R^\varepsilon}{2}(t, y) dy. \quad (2.14)$$

Note that this correction $\Theta^\varepsilon(t)$ is not necessary in the case of a problem set on the whole line \mathbb{R} , [ZZ03, ZZ05]. Finally, we define $\Phi^\varepsilon : \mathfrak{U} \rightarrow \cap_{s \geq 0} H^s(\mathbb{T})$ such that for any $k \geq 1$ and $\varepsilon > 0$ we have $\Phi^\varepsilon g_k \stackrel{\text{def}}{=} \sigma_k^\varepsilon \stackrel{\text{def}}{=} J_\varepsilon \sigma_k \in C^\infty(\mathbb{T})$. Clearly, we have the domination $|\sigma_k^\varepsilon| \leq |\sigma_k|$, and thus (see (2.2))

$$\sum_k \|\sigma_k^\varepsilon\|_{C(\mathbb{T})}^2 \leq \sum_k \|\sigma_k\|_{C(\mathbb{T})}^2 = q_0, \quad q^\varepsilon(x) \stackrel{\text{def}}{=} \sum_k \sigma_k(x)^2. \quad (2.15)$$

2.2 Local and global existence

In what follows, if E is a Banach space, the space $C([0, T]; E)$ of continuous functions $[0, T] \rightarrow E$ is endowed with the norm

$$\|u\|_{C([0, T]; E)} = \sup_{t \in [0, T]} \|u(t)\|_E.$$

The local existence theorem 2.2 below will be proved actually in a slightly more general version, Theorem 3.7. To solve (2.5) in the class $H^{s+1} \times H^s$, we need the noise to be a.s. H^{s+1} , which is ensured by the following hypothesis:

$$\sum_{k \geq 1} \|\sigma_k\|_{H^{s+1}(\mathbb{T})}^2 < \infty. \quad (2.16)$$

Indeed, a martingale inequality and (2.1) give

$$\mathbb{E} \|\Phi W\|_{C([0, T], H^{s+1}(\mathbb{T}))}^2 \leq CT \sum_{k \geq 1} \|\sigma_k\|_{H^{s+1}(\mathbb{T})}^2 < \infty. \quad (2.17)$$

Then (2.17) (considered for $T = 1, 2, \dots$) implies that there is a set Ω_s of probability one such that

$$\omega \in \Omega_s \Rightarrow \forall T > 0, \Phi W \in C([0, T], H^{s+1}(\mathbb{T})). \quad (2.18)$$

Definition 2.1 (Regular solution to (2.5), up to an explosion time). Let $(u_0, v_0) \in H^{s+1}(\mathbb{T}) \times H^s(\mathbb{T})$. Let $\tau > 0$ be a stopping time such that $\tau > 0$ a.s. Let Ω_s be defined in (2.18). A process $(u(t))_{0 \leq t < \tau}$ such that: a.s.,

$$u \in C([0, \tau), H^{s+1}(\mathbb{T})) \cap C^1([0, \tau), H^s(\mathbb{T})), \quad (2.19)$$

is said to be a regular (or H^{s+1}) solution to (2.5) up to the explosion time τ if,

1. $u(0, \cdot) = u_0$, \mathbb{P} -almost surely,
2. for all stopping time τ' such that $\tau' < \tau$ a.s., for all $x \in \mathbb{T}$, the process $(u(t \wedge \tau', x))_{t \geq 0}$ is predictable and, for all $\omega \in \Omega_s$ such that $\tau'(\omega) < \tau(\omega)$ (and thus \mathbb{P} -a.s.), we have the following identity in $H^{s-1}(\mathbb{T})$

$$u_t(\tau') = v_0 + \int_0^{\tau'} c(u) (c(u) u_x)_x ds + \Phi W(\tau'), \quad (2.20)$$

3. the stopping time τ is an explosion time, in the sense that

$$\tau = \sup_{n \geq 1} \tau_n, \quad \tau_n \stackrel{\text{def}}{=} \inf \{t \in [0, \tau); \|(u_t(t), u_x(t))\|_{L^\infty(\mathbb{T})} > n\}. \quad (2.21)$$

If $\tau = \infty$ a.s., the solution is said to be global.

In (2.21), we use the following convention.

Remark 2.2 (Convention for the infimum). If $T \geq 0$ and $(P_t)_{t \geq 0}$ are some properties depending on time, then the value of the infimum in the expressions

$$\inf \{t \in [0, T]; P_t \text{ is true}\}, \text{ or } \inf \{t \in [0, T]; P_t \text{ is true}\} \quad (2.22)$$

is set to the terminal value T if, respectively,

$$\forall t \in [0, T], P_t \text{ is false, or } \forall t \in [0, T], P_t \text{ is false.} \quad (2.23)$$

Theorem 2.2 (Local existence of regular solutions). Let $s > 3/2$. Let $c \in C^\infty(\mathbb{R})$ satisfies (2.3) and $|c'(u)| \leq c_3$. Let Φ satisfies (2.16) and let $(u_0, v_0) \in H^{s+1} \times H^s$. Then (2.5) admits a solution up to an explosion time τ^* and, for $i = 1, 2$, two solutions (u_i, τ_i^*) to (2.5) defined up to an explosion time τ_i^* coincide, in the sense that $\tau_1^* = \tau_2^*$ and $u_1 = u_2$ on $[0, \tau_1^*)$.

Remark 2.3. The local existence of regular solutions will be established for the equivalent system (2.6).

We then observe that singularity's formation can indeed occur in finite time.

Theorem 2.3 (Blow-up in finite time). Let $s > 3/2$. Assume that the conditions of Theorem 2.2 are satisfied, and that $c'(u^*) \neq 0$ for some $u^* \in \mathbb{R}$. For any $\gamma > 1/3$ and $\varepsilon > 0$, there is an initial datum $u_0^\varepsilon \in C^\infty(\mathbb{T})$ such that the regular solution (u, τ^*) to (2.5) defined up to the explosion time τ^* , with initial data $(u_0^\varepsilon, 0)$, satisfies

$$\mathbb{P}(\tau^* \leq \varepsilon^\gamma) \geq 1 - \varepsilon. \quad (2.24)$$

Eventually, we can state the existence of global-in-time weak solutions.

Definition 2.4 (Weak martingale solution). *Assume (2.2), (2.3) and (2.4). Let $u_0 \in H^1(\mathbb{T})$ and $v_0 \in L^2(\mathbb{T})$. We say that the problem (2.5) admits a weak martingale solution if there exists first a stochastic basis*

$$(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}}, (\tilde{\mathcal{F}}_t), (\tilde{W}(t))), \quad (2.25)$$

where $(\tilde{W}(t))$ is a cylindrical Wiener process on \mathfrak{U} , and, second, an adapted stochastic process $(u(t))$ with value in $H^1(\mathbb{T})$ such that, for all $T > 0$,

1. $u_t, u_x \in L^2(\tilde{\Omega}; L^\infty([0, T]; L^2(\mathbb{T})))$ and $u_t, u_x \in C([0, T]; L^2(\mathbb{T}) - \text{weak})$ $\tilde{\mathbb{P}}$ -a.s.,
2. $\tilde{\mathbb{P}}$ -a.s., $u(0, \cdot) = u_0$ and for all $\varphi \in C^1(\mathbb{T})$ and $t \in [0, T]$,

$$\int_{\mathbb{T}} u_t(t) \varphi \, dx - \int_{\mathbb{T}} v_0 \varphi \, dx + \int_0^t \int_{\mathbb{T}} (c(u(s)) \varphi)_x c(u(s)) u_x \, dx \, ds = \int_{\mathbb{T}} \varphi \Phi W(t) \, dx \quad (2.26)$$

3. $\tilde{\mathbb{P}}$ -a.s, the solution u satisfies the energy “dissipation” inequality

$$d \int_{\mathbb{T}} [u_t^2 + c(u)^2 u_x^2] \, dx \leq \|q\|_{L^1(\mathbb{T})} \, dt + 2 \int_{\mathbb{T}} u_t \Phi \, dx \, dW, \quad (2.27)$$

4. $u_t, u_x \in C_r([0, T]; L^2(\mathbb{T}))$ $\tilde{\mathbb{P}}$ -almost surely, i.e., for all $t_0 \in [0, T]$ we have

$$\lim_{t \downarrow t_0} \|(u_t(t) - u_t(t_0), u_x(t) - u_x(t_0))\|_{L^2} = 0. \quad (2.28)$$

Remark 2.4. The right-continuity condition (2.28) can be interpreted as a dissipation condition. Indeed, Dafermos [D11] proved that, in the case of the Hunter–Saxton equation (1.2), the right-continuity condition is equivalent to the dissipation of the energy. This plays a crucial role in the uniqueness of solutions [D11].

Theorem 2.5 (Global existence of weak martingale solutions). *Let $u_0 \in H^1(\mathbb{T})$ and $v_0 \in L^2(\mathbb{T})$. Assume (2.2), (2.3), (2.4) and suppose that*

$$c_0 \stackrel{\text{def}}{=} \inf_{x \in \mathbb{T}} c'(u_0(x)) > 0. \quad (2.29)$$

Then the problem (2.5) admits a weak martingale solution. Moreover, the solutions satisfies

- for all $p \in [1, 3)$ we have

$$\tilde{\mathbb{E}} \int_{[0, t] \times \mathbb{T}} c'(u) [|u_t|^p + |u_x|^p] \, dx \, dt \leq C(T, \alpha), \quad (2.30)$$

- for all $p \in [1, 2]$, there exists $C(p) > 0$ such that for all $t \in (0, T]$, we have the entropy inequality

$$\tilde{\mathbb{E}} \left\| [u_t \pm c(u) u_x]^- \right\|_{L^\infty}^p(t) \leq C(p, T) (1 + t^{-p}). \quad (2.31)$$

3 Local-in-time regular solutions

In this section we will consider the question of existence of local-in-time solutions in H^s to the system (3.1) below, which covers both systems (2.6) and (2.11).

3.1 Generalized system

For some given function $\theta \in C^\infty(\mathbb{R})$, we consider the system

$$dR + c(u) R_x dt = \tilde{c}'(u) [R^2 - S^2 - \theta(R) + 2R\Theta] dt + \Phi dW, \quad (3.1a)$$

$$dS - c(u) S_x dt = \tilde{c}'(u) [S^2 - R^2 - \theta(S) - 2S\Theta] dt + \Phi dW, \quad (3.1b)$$

$$R(0, \cdot) = R_0, \quad S(0, \cdot) = S_0, \quad (3.1c)$$

where u is defined by

$$u(t, x) = \mathcal{C}^{-1} \left\{ \mathcal{C} \left\{ \int_0^t \left(\frac{R+S}{2} \right) (s, 0) ds + u_0(0) \right\} + \int_0^x \left[\frac{S-R}{2}(t, y) - \Theta(t) \right] dy \right\}, \quad (3.2)$$

with

$$\Theta(t) \stackrel{\text{def}}{=} \int_0^1 \frac{S-R}{2}(t, y) dy. \quad (3.3)$$

The original system (2.6) simply corresponds to the case $\theta \equiv 0$ (one can prove that $\Theta(t) \equiv 0$ in that case¹), while the approximate system (2.11) corresponds to the choice $\theta = \chi_\varepsilon$ (note that χ_ε is not of class C^∞ though, but in the class $W_{\text{loc}}^{2,\infty}$, this is why Theorem 4.5 will be stated in the space $H^2(\mathbb{T})$). In essential, the system (3.1) will be solved pathwise, with the help of the following set of unknown functions:

$$P \stackrel{\text{def}}{=} R - \Phi W, \quad Q \stackrel{\text{def}}{=} S - \Phi W. \quad (3.4)$$

In terms of the unknown (P, Q) , the system (3.1) can be rewritten as the following system of PDEs with random coefficients:

$$P_t + c(u) P_x = \tilde{c}'(u) [P^2 - Q^2 - \theta(P + \Phi W) + 2(P + \Phi W)\Theta] - (c(u) \Phi W)_x, \quad (3.5a)$$

$$Q_t - c(u) Q_x = \tilde{c}'(u) [Q^2 - P^2 - \theta(Q + \Phi W) - 2(Q + \Phi W)\Theta] + (c(u) \Phi W)_x, \quad (3.5b)$$

$$P(0, \cdot) = R_0, \quad Q(0, \cdot) = S_0, \quad (3.5c)$$

where

$$u(t, x) = \mathcal{C}^{-1} \left\{ \mathcal{C} \left\{ \int_0^t \left(\frac{P+Q}{2} + \Phi W \right) (s, 0) ds + u_0(0) \right\} + \int_0^x \left[\frac{Q-P}{2}(t, y) - \Theta(t) \right] dy \right\}, \quad (3.6)$$

with

$$\Theta(t) \stackrel{\text{def}}{=} \int_0^1 \frac{Q-P}{2}(t, y) dy. \quad (3.7)$$

The system (3.5)-(3.6) has the form of the quasilinear system

$$V_t + A(V) V_x = F_\theta(V), \quad \text{in } (0, T) \times \mathbb{T}, \quad (3.8a)$$

$$V(0, \cdot) = V_0, \quad \text{in } \mathbb{T}. \quad (3.8b)$$

where

$$V \stackrel{\text{def}}{=} \begin{pmatrix} P \\ Q \end{pmatrix}, \quad (3.9)$$

¹indeed, Θ satisfies $\Theta' = \Theta \int_{\mathbb{T}} \tilde{c}'(u)(R+S) dx$ and $\Theta(0) = 0$.

and, (u, Θ) being given as a functions of V by (3.6)-(3.7), we set

$$A(V) \stackrel{\text{def}}{=} \begin{pmatrix} c(u) & 0 \\ 0 & -c(u) \end{pmatrix}, \quad (3.10)$$

and

$$F_\theta(V) \stackrel{\text{def}}{=} \begin{pmatrix} \tilde{c}'(u) [P^2 - Q^2 - \theta(P + \Phi W) + 2(P + \Phi W)\Theta] - (c(u)\Phi W)_x \\ \tilde{c}'(u) [Q^2 - P^2 - \theta(Q + \Phi W) - 2(Q + \Phi W)\Theta] + (c(u)\Phi W)_x \end{pmatrix}. \quad (3.11)$$

3.2 Local regular solutions

Let $s > 3/2$. We introduce the following definition.

Definition 3.1 (Regular solution to (3.8), up to an explosion time). *Let $V_0 \in H^s(\mathbb{T}; \mathbb{R}^2)$. Let $\tau > 0$ be a stopping time such that $\tau > 0$ a.s. Let Ω_s be defined in (2.18). A process $(V(t))_{0 \leq t < \tau} = (P(t), Q(t))_{0 \leq t < \tau}$ such that, we have: a.s.,*

$$(t \mapsto V(t)) \in C([0, \tau], H^s(\mathbb{T}; \mathbb{R}^2)) \cap C^1([0, \tau], H^{s-1}(\mathbb{T}; \mathbb{R}^2)), \quad (3.12)$$

is said to be a regular (or H^s) solution to (3.8) up to the explosion time τ if,

1. for all stopping time τ' such that $\tau' < \tau$ a.s., for all $x \in \mathbb{T}$, the process $(V(t \wedge \tau', x))_{t \geq 0}$ is predictable and, for all $\omega \in \Omega_s$ such that $\tau'(\omega) < \tau(\omega)$ (and thus \mathbb{P} -a.s.), we have the following identity in $H^{s-1}(\mathbb{T})$

$$V(\tau') = V_0 - \int_0^{\tau'} (A(V)V_x - F_\theta(V))(s) ds, \quad (3.13)$$

2. the stopping time τ is an explosion time, in the sense that

$$\tau = \sup_{n \geq 1} \tau_n, \quad \tau_n \stackrel{\text{def}}{=} \inf \{t \in [0, \tau); \|V(t)\|_{L^\infty(\mathbb{T}; \mathbb{R}^2)} > n\}. \quad (3.14)$$

If $\tau = \infty$ a.s., the solution is said to be global.

The integral identity (3.13) can be also expressed by integration along the characteristic curves. We will see several instances of this formulation: to solve the linear system (with frozen non-linear coefficients) associated to (3.8), and also to get some estimates on the solutions to (3.8), with either $\theta = 0$ or $\theta = \chi_\varepsilon$, so we give a rather general form to the statement.

Proposition 3.2 (Integration along characteristic curves). *Let $\tau > 0$ be a stopping time such that $\tau > 0$ a.s. Let $(v(t), \bar{c}(t), f(t))_{0 \leq t < \tau}$ be a stochastic process such that, we have: a.s.,*

$$[t \mapsto v(t)] \in C([0, \tau], H^s(\mathbb{T})) \cap C^1([0, \tau], H^{s-1}(\mathbb{T})), \quad (3.15)$$

and

$$[t \mapsto (\bar{c}, f)(t)] \in C^1([0, \tau], H^{s-1}(\mathbb{T})) \times C([0, \tau], H^{s-1}(\mathbb{T})). \quad (3.16)$$

Assume also: a.s., for all $t \geq 0$, $x \in \mathbb{T}$, $\bar{c}(t, x) \in [c_1, c_2]$ (where c_1, c_2 are the constants in (2.3)). Let $(X(t, x))$ denote the flow associated to \bar{c} , and let $y \mapsto Y(t, y)$ denote the inverse of the map $x \mapsto X(t, x)$. Then, there is equivalence between:

1. for all stopping time τ' such that $\tau' < \tau$ a.s.,

$$v(\tau') = v(0) - \int_0^{\tau'} (\bar{c}v_x - f)(s) ds, \quad (3.17)$$

and

2. a.s., for all $t \in [0, \tau)$ and all $x \in \mathbb{T}$,

$$v(t, x) = v(0, Y(t, x)) + \int_0^t f(s, X(s, Y(t, x))) ds. \quad (3.18)$$

Proof of Proposition 3.2. Note first that (3.17) is equivalent to: a.s., for all $t \in [0, \tau)$,

$$v(t) = v(0) - \int_0^t (\bar{c}v_x - f)(s) ds. \quad (3.19)$$

Indeed, (3.17) can be applied to the stopping time $t \wedge \tau'$, from which we deduce (3.19) for all $t \in [0, \tau'(\omega))$, and thus for all $t \in [0, \tau(\omega))$ since $\tau' < \tau$ is arbitrary. Conversely, (3.19) yields (3.17) when τ' is simple, and the general case follows by approximation. This being said, Proposition 3.2 is essentially a classical deterministic statement. It is not difficult to check that all the terms are sufficiently regular to justify the equivalence through differentiation. In particular, the equation

$$\partial_t v(t, x) + \bar{c}(t, x) \partial_x v(t, x) = f(t, x) \quad (3.20)$$

is satisfied a.s., for all $t \in [0, \tau)$ and all $x \in \mathbb{T}$. \square

Energy estimates in $H^s(\mathbb{T})$ are one of the main ingredient in the resolution of (3.8) locally in time. A basic issue in the justification of such estimates for solutions a priori in $H^s(\mathbb{T})$ is that the term $A(V)\partial_x$ involves an additional space derivative and necessitates the better H^{s+1} -regularity to be safely manipulated. There are several ways to circumvent this problem, for instance by working first on an approximate problem (by Friedrichs regularization of the term $A(V)\partial_x$, or by direct resolution of an iterative scheme in a class of sufficiently smooth solutions). We choose a different method, by a direct regularization of the equation and commutator estimates (based on Kato–Ponce’s estimates [KP88] and DiPerna–Lions estimates [DL89]). We give the result on the general equation (3.20).

Proposition 3.3 (H^s -estimate, linear equation). *Let $(v(t), \bar{c}(t), f(t))_{0 \leq t < \tau}$ be as in Proposition 3.2. There is a constant B_s depending on s only such that, a.s., for all $T \in [0, \tau)$ satisfying*

$$2TB_s \|\partial_x \bar{c}\|_{C([0, T]; L^\infty(\mathbb{T}))} \leq 1, \quad (3.21)$$

one has

$$\begin{aligned} \|v\|_{C([0, T]; H^s(\mathbb{T}))} &\leq 2\|v(0)\|_{H^s(\mathbb{T})} + 2T\|f\|_{C([0, T]; H^s(\mathbb{T}))} \\ &\quad + 2TB_s \|v\|_{C([0, T]; W^{1, \infty}(\mathbb{T}))} \|\bar{c}\|_{C([0, T]; H^s(\mathbb{T}))}. \end{aligned} \quad (3.22)$$

In particular, there is a constant C_s depending only on s such that, if T additionally satisfies

$$4TC_s \|\bar{c}\|_{C([0, T]; H^s(\mathbb{T}))} \leq 1, \quad (3.23)$$

then

$$\|v\|_{C([0, T]; H^s(\mathbb{T}))} \leq 4\|v(0)\|_{H^s(\mathbb{T})} + 4T\|f\|_{C([0, T]; H^s(\mathbb{T}))}. \quad (3.24)$$

We adopt the following notations: Λ^s denotes the operator $(I - \partial_x^2)^{s/2}$, $[A, B] = AB - BA$ the commutator of two operators A and B . In order to prove Proposition 3.3, we recall the following Kato–Ponce estimates [KP88]

$$\|fg\|_{H^r} \leq C (\|f\|_{L^\infty} \|g\|_{H^r} + \|f\|_{H^r} \|g\|_{L^\infty}), \quad (3.25)$$

$$\|[\Lambda^r, f]g\|_{L^2} \leq C (\|f_x\|_{L^\infty} \|g\|_{H^{r-1}} + \|f\|_{H^r} \|g\|_{L^\infty}), \quad (3.26)$$

where $r \geq 0$, and $C > 0$ is a constant depending only on r . We recall also the following estimate [CM02]

$$\|F(f) - F(0)\|_{H^s} \leq C (\|f\|_{L^\infty}) \|f\|_{H^s}, \quad (3.27)$$

for any $f \in H^s$ with $s > 1/2$ and $F \in C^\infty(\mathbb{R})$.

Proof of Proposition 3.3. The second estimate (3.24) follows from (3.22) and the continuous injection of $H^s(\mathbb{T})$ into $W^{1,\infty}(\mathbb{T})$ (the exponent s satisfying $s > 3/2$). To prove (3.22), we introduce (ρ_δ) , a standard approximation of the unit, and let J_δ denote the Friedrichs operator $f \mapsto f * \rho_\delta$, obtained by convolution by ρ_δ with respect to the variable x . Denote also by \bar{c} the operator of multiplication by \bar{c} . We apply J_δ to (3.20) to obtain

$$\partial_t v^\delta(t, x) + \bar{c}(t, x) \partial_x v^\delta(t, x) = f^\delta(t, x) + r^\delta(t, x), \quad (3.28)$$

where

$$v^\delta = J_\delta v, \quad f^\delta = J_\delta f, \quad r^\delta = \bar{c} \partial_x v^\delta - J_\delta(\bar{c} \partial_x v) = [\bar{c}, J_\delta] \partial_x v.$$

We have

$$v^\delta \in C^1([0, \tau]; H^s(\mathbb{T})), \quad \partial_x v^\delta, f^\delta, r^\delta \in C^0([0, \tau]; H^s(\mathbb{T})),$$

so we can apply Λ^s to (3.28) to get

$$\partial_t \Lambda^s v^\delta(t, x) + \bar{c}(t, x) \partial_x \Lambda^s v^\delta(t, x) = \Lambda^s f^\delta(t, x) + R^\delta(t, x), \quad (3.29)$$

where

$$R^\delta = \Lambda^s r^\delta + [\bar{c}, \Lambda^s] \partial_x v^\delta. \quad (3.30)$$

We can then multiply (3.29) by $\Lambda^s v^\delta$, integrate on \mathbb{T} and integrate by parts to get

$$\begin{aligned} \|\Lambda^s v^\delta(t)\|_{L^2(\mathbb{T})} &\leq \|\Lambda^s v^\delta(0)\|_{L^2(\mathbb{T})} \\ &+ \int_0^t \left\{ \frac{1}{2} \|\partial_x \bar{c}(\sigma)\|_{L^\infty(\mathbb{T})} \|\Lambda^s v^\delta(\sigma)\|_{L^2(\mathbb{T})} + \|\Lambda^s f^\delta(\sigma)\|_{L^2(\mathbb{T})} + \|R^\delta(\sigma)\|_{L^2(\mathbb{T})} \right\} d\sigma, \end{aligned} \quad (3.31)$$

which yields, for $T < \tau(\omega)$,

$$\begin{aligned} \|v^\delta\|_{C([0, T]; H^s(\mathbb{T}))} &\leq \|v^\delta(0)\|_{H^s(\mathbb{T})} \\ &+ T \left\{ \frac{1}{2} \|\partial_x \bar{c}\|_{C([0, T]; L^\infty(\mathbb{T}))} \|v^\delta\|_{C([0, T]; H^s(\mathbb{T}))} + \|f^\delta\|_{C([0, T]; H^s(\mathbb{T}))} + \|R^\delta\|_{C([0, T]; L^2(\mathbb{T}))} \right\}. \end{aligned} \quad (3.32)$$

Assume

$$\begin{aligned} \|R^\delta\|_{C([0, T]; L^2(\mathbb{T}))} &\leq B_s \left[\|\partial_x \bar{c}\|_{C([0, T]; L^\infty(\mathbb{T}))} \|v\|_{C([0, T]; H^s(\mathbb{T}))} + \|v\|_{C([0, T]; W^{1,\infty}(\mathbb{T}))} \|\bar{c}\|_{C([0, T]; H^s(\mathbb{T}))} \right], \end{aligned} \quad (3.33)$$

for a certain constant B_s . Then we can insert (3.33) in (3.32), pass to the limit $\delta \rightarrow 0$ and, after replacing B_s by $B_s + 1/2$, conclude to (3.22). To establish (3.33), we write

$$R^\delta = \Lambda^s [\bar{c}, J_\delta] \partial_x v + [\bar{c}, \Lambda^s] \partial_x v^\delta = [\Lambda^s, [\bar{c}, J_\delta]] \partial_x v + [\bar{c}, J_\delta] \Lambda^s \partial_x v + [\bar{c}, \Lambda^s] \partial_x v^\delta, \quad (3.34)$$

and use the Jacobi identity

$$[\Lambda^s, [\bar{c}, J_\delta]] + [[\bar{c}, [J_\delta, \Lambda^s]] + [J_\delta, [\Lambda^s, \bar{c}]] = 0,$$

and the fact that $[J_\delta, \Lambda^s] = 0$ to get

$$R^\delta = [J_\delta, [\bar{c}, \Lambda^s]]\partial_x v + [\bar{c}, J_\delta]\Lambda^s\partial_x v + [\bar{c}, \Lambda^s]\partial_x v^\delta = J_\delta[\bar{c}, \Lambda^s]\partial_x v + [\bar{c}, J_\delta]\partial_x(\Lambda^s v). \quad (3.35)$$

The intermediate identities between (3.30) and (3.35) can be understood in the sense of distribution, but we must check that each term in (3.35) is in $L^2(\mathbb{T})$ and satisfies the expected estimate. We have $[\bar{c}, \Lambda^s]\partial_x v \in L^2(\mathbb{T})$ by Kato-Ponce's commutator estimate (3.26), with

$$\begin{aligned} & \|[\bar{c}(t), \Lambda^s]\partial_x v(t)\|_{L^2(\mathbb{T})} \\ & \leq C_{\text{KP}} \left[\|\partial_x \bar{c}(t)\|_{L^\infty(\mathbb{T})} \|\partial_x v(t)\|_{H^{s-1}(\mathbb{T})} + \|\bar{c}(t)\|_{H^s(\mathbb{T})} \|\partial_x v(t)\|_{L^\infty(\mathbb{T})} \right], \end{aligned}$$

and thus

$$\begin{aligned} & \|[\bar{c}(t), \Lambda^s]\partial_x v(t)\|_{L^2(\mathbb{T})} \\ & \leq C_{\text{KP}} \left[\|\partial_x \bar{c}\|_{C([0, T]; L^\infty(\mathbb{T}))} \|v\|_{C([0, T]; H^s(\mathbb{T}))} + \|v\|_{C([0, T]; W^{1, \infty}(\mathbb{T}))} \|\bar{c}\|_{C([0, T]; H^s(\mathbb{T}))} \right], \end{aligned}$$

for $t \in [0, T]$. By the commutator estimate Lemma II-1. in [DL89] (and more exactly, by the proof of Lemma II-1), we have $[\bar{c}, J_\delta]\partial_x(\Lambda^s v) \in L^2(\mathbb{T})$ with

$$\|[\bar{c}(t), J_\delta]\partial_x(\Lambda^s v(t))\|_{L^2(\mathbb{T})} \leq C_{\text{DPL}} \|\partial_x \bar{c}(t)\|_{L^2(\mathbb{T})} \|\Lambda^s v(t)\|_{L^2(\mathbb{T})},$$

which yields $R^\delta(t) \in L^2(\mathbb{T})$ for $t \in [0, T]$, and the estimate (3.33). \square

To go on our analysis, we consider now the deterministic quasilinear (but non-local) equation

$$\partial_t v + \bar{c}[v]\partial_x v = F[v], \quad (3.36)$$

on $(0, T) \times \mathbb{T}$, where $c, F: C([0, T]; H^s(\mathbb{T})) \rightarrow C([0, T]; H^s(\mathbb{T}))$ satisfy

$$\|\partial_x \bar{c}[v]\|_{C([0, T]; L^\infty(\mathbb{T}))} \leq \mathfrak{C}(\|v\|_{C([0, T]; L^\infty(\mathbb{T}))}), \quad (3.37)$$

and

$$\|\partial_x F[v]\|_{C([0, T]; L^\infty(\mathbb{T}))} \leq \mathfrak{C}(\|v\|_{C([0, T]; L^\infty(\mathbb{T}))}) (1 + \|\partial_x v\|_{C([0, T]; L^\infty(\mathbb{T}))}), \quad (3.38)$$

as well as the bounds

$$\|\bar{c}[v]\|_{C([0, T]; H^s(\mathbb{T}))} + \|F[v]\|_{C([0, T]; H^s(\mathbb{T}))} \leq \mathfrak{C}(\|v\|_{C([0, T]; L^\infty(\mathbb{T}))}) (1 + \|v\|_{C([0, T]; H^s(\mathbb{T}))}), \quad (3.39)$$

for all $v \in C([0, T]; H^s(\mathbb{T}))$, for a given function $\mathfrak{C}: \mathbb{R}_+ \rightarrow \mathbb{R}_+$. We will also consider the Lipschitz condition

$$\begin{aligned} & \|\bar{c}[v_1] - \bar{c}[v_2]\|_{C([0, T]; L^2(\mathbb{T}))} + \|F[v_1] - F[v_2]\|_{C([0, T]; L^2(\mathbb{T}))} \\ & \leq \mathfrak{C}(\|v_1\|_{C([0, T]; W^{1, \infty}(\mathbb{T}))} + \|v_2\|_{C([0, T]; W^{1, \infty}(\mathbb{T}))}) \|v_1 - v_2\|_{C([0, T]; L^2(\mathbb{T}))}. \end{aligned} \quad (3.40)$$

Proposition 3.4 (H^s -estimate, linear equation). *Assume \bar{c}, F in (3.36) satisfy (3.37) and (3.39). Then, given $R > 0$, there is a constant $\nu_R > 0$ such that, if $T\nu_R \leq 1$, then all solutions $v \in C([0, T]; H^s(\mathbb{T}))$ to the equation*

$$\partial_t v + \bar{c}[w]\partial_x v = F[w], \quad (3.41)$$

with initial datum $v_0 \in H^s(\mathbb{T})$ and entry $w \in C([0, T]; H^s(\mathbb{T}))$ of size

$$\|v_0\|_{H^s(\mathbb{T})} \leq R/3, \quad \|w\|_{C([0, T]; H^s(\mathbb{T}))} \leq R,$$

satisfy the bound

$$\|v\|_{C([0, T]; H^s(\mathbb{T}))} \leq R. \quad (3.42)$$

Proof of Proposition 3.4. This is a straightforward consequence of the estimate (3.22), using (3.37)-(3.39) and the injection $H^s(\mathbb{T}) \hookrightarrow L^\infty(\mathbb{T})$, which yields

$$\|v\|_{C([0, T]; H^s(\mathbb{T}))} \leq 2\|v_0\|_{H^s(\mathbb{T})} + C_R T \leq \frac{2}{3}R + C_R T, \quad (3.43)$$

for $T\nu_R \leq 1$, where $\nu_R > 0$, $C_R \geq 0$. We obtain then (3.42) by adjusting the value of ν_R if necessary. \square

Next, we state the contraction property in the low norm L^2 .

Proposition 3.5 (L^2 -contraction, linear equation). *Assume \bar{c}, F in (3.36) satisfy (3.37), (3.39) and also (3.40). Given $R > 0$, there is a constant $\nu_R > 0$ such that, if $T\nu_R \leq 1$, and $v_1, v_2 \in C([0, T]; H^s(\mathbb{T}))$ are solutions of the respective equations*

$$\partial_t v_i + \bar{c}[w_i] \partial_x v_i = F[w_i], \quad i \in \{1, 2\}, \quad (3.44)$$

with respective initial datum $v_{1,0}, v_{2,0} \in H^s(\mathbb{T})$ and entries $w_1, w_2 \in C([0, T]; H^s(\mathbb{T}))$ of size

$$\|v_0\|_{H^s(\mathbb{T})} \leq R/3, \quad \|w_i\|_{C([0, T]; L^\infty(\mathbb{T}))} \leq R, \quad i \in \{1, 2\},$$

then

$$\|v_1 - v_2\|_{C([0, T]; L^2(\mathbb{T}))} \leq 4\|v_{1,0} - v_{2,0}\|_{L^2(\mathbb{T})} + \frac{1}{2}\|w_1 - w_2\|_{C([0, T]; L^2(\mathbb{T}))}. \quad (3.45)$$

Proof of Proposition 3.5. The difference $\bar{v} := v_1 - v_2$ satisfies the equation

$$\partial_t \bar{v} + \bar{c}[w_1] \partial_x \bar{v} = \bar{f} := F[w_1] - F[w_2] + (\bar{c}[w_2] - \bar{c}[w_1]) \partial_x v_2, \quad (3.46)$$

starting from $\bar{v}_0 := v_{1,0} - v_{2,0}$, and thus the L^2 -estimate

$$\begin{aligned} \|\bar{v}\|_{C([0, T]; L^2(\mathbb{T}))} &\leq \|v_{1,0} - v_{2,0}\|_{L^2(\mathbb{T})} \\ &\quad + T \|\partial_x \bar{c}[w_1]\|_{C([0, T]; L^\infty(\mathbb{T}))} \|\bar{v}\|_{C([0, T]; L^2(\mathbb{T}))} + 2T \|\bar{f}\|_{C([0, T]; L^2(\mathbb{T}))}. \end{aligned} \quad (3.47)$$

We apply Proposition 3.4: if T is sufficiently small, then the bound (3.42) is satisfied by v_1 and v_2 . As a consequence of (3.37) and (3.40) (and of the injection $H^s(\mathbb{T}) \hookrightarrow W^{1, \infty}(\mathbb{T})$ again), we obtain

$$\|\bar{v}\|_{C([0, T]; L^2(\mathbb{T}))} \leq 2\|v_{1,0} - v_{2,0}\|_{L^2(\mathbb{T})} + C'_R T \|\bar{v}\|_{C([0, T]; L^2(\mathbb{T}))} + C''_R T \|\bar{w}\|_{C([0, T]; L^2(\mathbb{T}))}, \quad (3.48)$$

with $\bar{w} := w_1 - w_2$, from which (3.45) follows for T sufficiently small. \square

We can then give a result of existence of solutions to (3.36) up to a time of explosion of the Lipschitz norm of the solution.

Proposition 3.6 (Existence up to explosion of the sup-norm). *Assume \bar{c}, F in (3.36) satisfy (3.37)-(3.38)-(3.39) and also (3.40). Let $v_0 \in H^s(\mathbb{T})$ and let $T > 0$. There exists a solution $v \in C([0, T^*]; H^s(\mathbb{T})) \cap C^1([0, T^*]; H^{s-1}(\mathbb{T}))$ to (3.36) with initial datum v_0 defined up to a time T^* which satisfies*

$$T^* = \sup_{n \geq 1} T^n, \quad T^n \stackrel{\text{def}}{=} \inf \{t \in [0, T^*]; \|v(t)\|_{L^\infty(\mathbb{T})} > n\}. \quad (3.49)$$

Proof of Proposition 3.6. We consider the iterative scheme $v^k \mapsto v^{k+1}$, where v^{k+1} is the solution (given by the method of characteristics, cf. Proposition 3.2) to the equation

$$\partial_t v^{k+1} + \bar{c}[v^k] \partial_x v^{k+1} = F[v^k], \quad (3.50)$$

with initial datum v_0 . By Proposition 3.4 and Proposition 3.5, there is a time T_ℓ depending on $\|v_0\|_{H^s(\mathbb{T})}$ only such that the sequence (v^k) converges in $C([0, T_\ell]; H^s(\mathbb{T}))$ to a solution v to (3.36) with initial datum v_0 defined on $[0, T_\ell]$. This gives the local existence. Local uniqueness follows from (3.45) (possibly applied repeatedly). So the problem (3.36) with initial datum v_0 admits a maximal solution defined up to the time S^* defined by

$$S^* \stackrel{\text{def}}{=} \sup_{n \geq 1} S^n, \quad S^n \stackrel{\text{def}}{=} \inf \{t \in [0, T^*]; \|v(t)\|_{H^s(\mathbb{T})} > n\}. \quad (3.51)$$

Let us first introduce the explosion time of the Lipschitz norm:

$$\hat{T}^* \stackrel{\text{def}}{=} \sup_{n \geq 1} \hat{T}^n, \quad \hat{T}^n \stackrel{\text{def}}{=} \inf \{t \in [0, T^*]; \|v(t)\|_{W^{1,\infty}(\mathbb{T})} > n\}. \quad (3.52)$$

The norm in $H^s(\mathbb{T})$ dominates the norm in $W^{1,\infty}(\mathbb{T})$ so $S^* \leq \hat{T}^*$. On the other hand, if $t < \hat{T}_n \wedge S^*$, so that v solves (3.36) on $[0, t]$ and

$$\|v\|_{C([0,t]; W^{1,\infty}(\mathbb{T}))} \leq n,$$

we observe that (3.22) gives (for arbitrary times $0 \leq r < r + \sigma \leq t$)

$$\|v\|_{C([r,r+\sigma]; H^s(\mathbb{T}))} \leq 2\|v(r)\|_{H^s(\mathbb{T})} + 2\sigma \left(\|F[v]\|_{C([r,r+\sigma]; H^s(\mathbb{T}))} + B_s n \|\bar{c}[v]\|_{C([r,r+\sigma]; H^s(\mathbb{T}))} \right). \quad (3.53)$$

We deduce then from (3.39) that

$$\|v\|_{C([r,r+\sigma]; H^s(\mathbb{T}))} \leq 2\|v(r)\|_{H^s(\mathbb{T})} + 2\sigma \mathfrak{C}(n) (1 + B_s n) (\|v\|_{C([r,r+\sigma]; H^s(\mathbb{T}))} + 1), \quad (3.54)$$

and thus

$$\|v\|_{C([r,r+\sigma]; H^s(\mathbb{T}))} \leq K(n) (\|v(r)\|_{H^s(\mathbb{T})} + 1), \quad (3.55)$$

if $\sigma < \sigma(n)$ for a $\sigma(n)$ sufficiently small depending on n (and s) only. By iterating (3.55), we obtain the bound $\|v\|_{C([0,t]; H^s(\mathbb{T}))} \leq H(n)$ for a given function $H: \mathbb{N} \rightarrow \mathbb{N}$, and thus $t < S_{H(n)}$. Letting $t \uparrow \hat{T}_n \wedge S^*$ gives $\hat{T}_n \wedge S^* \leq S_{H(n)} \leq S^*$ and thus $\hat{T}_n \leq S^*$. We conclude that $\hat{T}^* \leq S^*$ and finally that $\hat{T}^* = S^*$.

Let us now prove that $T^* = \hat{T}^*$. Clearly $\hat{T}^* \leq T^*$. We consider a time $t < T_n \wedge \hat{T}^*$. By differentiation in x in (3.36) on $[0, t]$, we obtain the equation

$$\partial_t w + c[v] \partial_x w = \partial_x F[v] - w \partial_x c[v], \quad (3.56)$$

where $w := \partial_x v$. By the comparison principle, and for $s \in [0, t]$,

$$\|w(s)\|_{L^\infty(\mathbb{T})} \leq \|\partial_x v(0)\|_{L^\infty(\mathbb{T})} + \int_0^s \|\partial_x F[v] - w \partial_x c[v]\|_{L^\infty(\mathbb{T})}(\sigma) d\sigma. \quad (3.57)$$

We use (3.37), (3.38) and the Grönwall lemma, to deduce from (3.57) that

$$\sup_{0 \leq s \leq t} \|w(s)\|_{L^\infty(\mathbb{T})} \leq (\|\partial_x v(0)\|_{L^\infty(\mathbb{T})} + C(n)t) \exp(C(n)t), \quad (3.58)$$

for a certain constant $C(n)$. This implies $t < \hat{T}_{\tilde{H}(n)}$ for a certain function $\tilde{H}: \mathbb{N} \rightarrow \mathbb{N}$, which, as above, gives $T^* \leq \hat{T}^*$, and so finally $T^* = \hat{T}^*$. \square

Theorem 3.7 (Local existence of regular solutions). *Let $s > 3/2$. Assume (2.3), (2.16) and $|c'(u)| \leq c_3$. Let $V_0 \in H^s(\mathbb{T}; \mathbb{R}^2)$ such that $\int_{\mathbb{T}} (Q_0 - P_0) dx = 0$. Then (3.8) admits a solution up to an explosion time τ and two solutions (V_i, τ_i) , $i = 1, 2$ to (3.8) coincide, in the sense that $\tau_1 = \tau_2$ and $V_1 = V_2$ on $[0, \tau_1)$.*

Proof of Theorem 3.7. We use Proposition 3.6 (the results can be extended to the case of two equations in a straightforward way) to solve (3.8) pathwise. Therefore it will be sufficient to prove that, a.s., $u = u[V]$ being defined by (3.6), and F_θ given by (3.11), the functions

$$V \mapsto c[V] := c(u), \quad V \mapsto F_\theta(V)$$

satisfy (3.37)-(3.38)-(3.39)-(3.40). If we take into account the regularity (2.18) and the definition of (P, Q) as the shift of (R, S) by ΦW , then the property (3.37) follows from the identity

$$c(u)u_x = \frac{S - R}{2} - \Theta, \quad (3.59)$$

and the bounds in (2.3). Similarly, we deduce (3.39) and (3.40) from (2.3), (2.18), (3.6), (3.59) and (3.27). The property (3.38) is a straightforward consequence of (2.18) and the expression (3.11) of F_θ . The local uniqueness is a consequence of (3.45), which gives $V_1 = V_2$ on $[0, \tau_1^* \wedge \tau_2^*)$, but then also yields $\tau_1^* = \tau_2^*$ since both are explosion times. Finally, the predictable character of the stochastic processes can be deduced from the predictable character of each element of the iteration (3.50). \square

3.3 Conservative form, integration along the characteristics, Itô's formula

The identity (3.13) should be understood in $H^{s-1}(\mathbb{T})$. In particular, (3.13) is satisfied point-wise for every $x \in \mathbb{T}$ and so is the system (3.1) (once interpreted in the standard integral form of SDEs). The Itô formula can therefore be applied for every $x \in \mathbb{T}$ to the real-valued processes $(R(t, x))_{t < \tau}$ and $(S(t, x))_{t < \tau}$ and gives, for all smooth function $h \in C^\infty(\mathbb{R})$,

$$\begin{aligned} dh(R) + c(u)h(R)_x dt &= \tilde{c}'(u) [h'(R)(R^2 - S^2 - \theta(R) + 2R\Theta) + \frac{1}{2}qh''(R)] dt \\ &\quad + h'(R)\Phi dW, \end{aligned} \quad (3.60)$$

and

$$\begin{aligned} dh(S) - c(u)h(S)_x dt &= \tilde{c}'(u) [h'(S)(S^2 - R^2 - \theta(S) - 2S\Theta) + \frac{1}{2}qh''(S)] dt \\ &\quad + h'(S)\Phi dW. \end{aligned} \quad (3.61)$$

We can then use the identity (3.59) to obtain the conservative form

$$\begin{aligned} dh(R) + (c(u)h(R))_x dt &= h'(R)\Phi dW \\ + \tilde{c}'(u) [(S - R)B_h(R, S) - h'(R)\theta(R) + 2\Theta(Rh'(R) - 2h(R)) + \frac{1}{2}qh''(R)] dt, \end{aligned} \quad (3.62)$$

and

$$\begin{aligned} dh(S) - (c(u)h(S))_x dt &= h'(S)\Phi dW \\ = \tilde{c}'(u) [(R - S)B_h(S, R) - h'(S)\theta(S) - 2\Theta(S h'(S) - 2h(S)) + \frac{1}{2}qh''(S)] dt, \end{aligned} \quad (3.63)$$

where

$$B_h(R, S) \stackrel{\text{def}}{=} 2h(R) - (R + S)h'(R). \quad (3.64)$$

By taking $h(R) = R^2$, we obtain

$$B_h(R, S) = -2RS = B_h(S, R), \quad R h'(R) - 2h(R) = 0, \quad (3.65)$$

and, by adding (3.62) to (3.63),

$$\begin{aligned} (R^2 + S^2)(\tau') &= R^2(0) + S^2(0) + \int_0^{\tau'} [(c(u)(S^2 - R^2))_x - 2R\theta(R) - 2S\theta(S) + 2q] dt \\ &\quad + \int_0^{\tau'} 2(R + S) \Phi dW, \end{aligned} \quad (3.66)$$

for all stopping time $\tau' < \tau$.

We will also need the evolution equation of $\frac{h(R)}{c(u)}$ (and similarly for S , see below). To that purpose, we compute the derivative of u .

Proposition 3.8 (Computation of u_t). *Let (V, τ) be a regular solution to (3.8) up to the explosion time τ and let (u, Θ) being given as a functions of V by (3.6)-(3.7). We have then*

$$\Theta'(t) = \alpha(t) + \beta(t)\Theta(t), \quad (3.67)$$

where

$$\alpha(t) = \frac{1}{2} \int_0^1 \tilde{c}'(u) (\theta(R) - \theta(S)) dy, \quad \beta(t) = \int_0^1 \tilde{c}'(u) (R + S) dy, \quad (3.68)$$

and

$$u_t = \frac{R + S}{2} - \frac{1}{c(u)} \int_0^x [\zeta(t, y) - \bar{\zeta}(t)] dy, \quad (3.69)$$

where

$$\zeta = \tilde{c}'(u) \left[\frac{\theta(R) - \theta(S)}{2} + (R + S)\Theta \right], \quad \bar{\zeta}(t) = \int_0^1 \zeta(t, y) dy. \quad (3.70)$$

Proof of Proposition 3.8. We consider the equations (3.62)-(3.63) with $h(R) = R$, and subtract the first equation from the second one to obtain

$$\partial_t(S - R) + (c(u)(R + S))_x = \tilde{c}'(u) [(\theta(R) - \theta(S)) + 2(R + S)\Theta]. \quad (3.71)$$

By integration on $(0, 1)$ in (3.71), we obtain (3.67). Then we observe that, taking $x = 0$ in (3.2) gives

$$u(t, 0) = \int_0^t \left(\frac{R + S}{2} \right) (s, 0) ds + u_0(0), \quad (3.72)$$

so that

$$\mathcal{C}(u(t, x)) = \mathcal{C}(u(t, 0)) - \int_0^x \left[\frac{S - R}{2}(t, y) - \Theta(t) \right] dy. \quad (3.73)$$

By differentiation in (3.72) and (3.73) with respect to t , we obtain

$$c(u)u_t = c(u(t, 0)) \left(\frac{R + S}{2} \right) (t, 0) - \int_0^x \left[\frac{\partial_t(S - R)}{2}(t, y) - \Theta'(t) \right] dy. \quad (3.74)$$

We use (3.71) and (3.67) to get

$$c(u)u_t = c(u) \frac{R + S}{2} - \int_0^x [\zeta(t, y) - \bar{\zeta}(t)] dy, \quad (3.75)$$

where ζ and $\bar{\zeta}$ are defined in (3.70). Dividing by $c(u)$ yields (3.69). \square

Set

$$\Xi \stackrel{\text{def}}{=} u_t - \frac{R+S}{2} = -\frac{1}{c(u)} \int_0^x [\zeta(t, y) - \bar{\zeta}(t)] dy. \quad (3.76)$$

It follows from (3.60)-(3.61) and (3.69) that

$$\begin{aligned} & d \left[\frac{h(R)}{c(u)} \right] + [h(R)]_x dt = h'(R) \Phi dW \\ & + \frac{\tilde{c}'(u)}{c(u)} [4h(R)\Xi + (S+R)F_h(R, S) - h'(R)\theta(R) + 2\Theta R h'(R) + \frac{1}{2}q h''(R)] dt, \end{aligned} \quad (3.77)$$

and

$$\begin{aligned} & d \left[\frac{h(S)}{c(u)} \right] - [h(S)]_x dt = h'(S) \Phi dW \\ & + \frac{\tilde{c}'(u)}{c(u)} [4h(S)\Xi + (R+S)F_h(S, R) - h'(S)\theta(S) - 2\Theta S h'(S) + \frac{1}{2}q h''(S)] dt, \end{aligned} \quad (3.78)$$

where

$$F_h(R, S) := 2h(R) - (R-S)h'(R). \quad (3.79)$$

We complete this section with the the expression of (3.60)-(3.61) after integration along the characteristic curves (cf. Proposition 3.2).

Proposition 3.9 (Integration along the characteristic curves). *Let $(V(t))_{t < \tau}$ be a regular solution to (3.8) up to the explosion time τ . Let $h \in C^\infty(\mathbb{R})$. Let X^\pm denote the flow associated to the vector field $\pm c(u(t, \cdot))$:*

$$\frac{d}{dt} X^\pm(t, x) = \pm c(u(t, X^\pm(t, x))), \quad X^\pm(0, x) = x, \quad x \in \mathbb{T}, \quad t \in [0, \tau). \quad (3.80)$$

Let also $Y^\pm(t, \cdot)$ denote the inverse of the map $x \mapsto X^\pm(t, x)$. We have then

1. X^\pm and Y^\pm are well-defined, C^1 -maps on $[0, \tau) \times \mathbb{T}$; for all $x \in \mathbb{T}$, $(X^\pm(t, x))$ and $(Y^\pm(t, x))$ are predictable processes, and the stochastic integrals

$$\mathcal{M}_h(t, x) = \sum_k \int_0^t \{h'(R)\sigma_k\}(s, X^+(s, x)) d\beta_k(s), \quad (3.81)$$

and

$$\mathcal{N}_h(t, x) = \sum_k \int_0^t \{h'(S)\sigma_k\}(s, X^-(s, x)) d\beta_k(s), \quad (3.82)$$

define martingales,

2. we have

$$\begin{aligned} & h(R)(t, x) = h(R_0)(Y^+(t, x)) + \mathcal{M}_h(t, Y^+(t, x)) \\ & + \int_0^t \left\{ \tilde{c}'(u) [h'(R)(R^2 - S^2 - \theta(R) + 2R\Theta) + \frac{1}{2}q h''(R)] \right\} (s, X^+(s, Y^+(t, x))) ds, \end{aligned} \quad (3.83)$$

and

$$\begin{aligned} & h(S)(t, x) = h(S_0)(Y^-(t, x)) + \mathcal{N}_h(t, Y^-(t, x)) \\ & + \int_0^t \left\{ \tilde{c}'(u) [h'(S)(S^2 - R^2 - \theta(S) - 2S\Theta) + \frac{1}{2}q h''(S)] \right\} (s, X^-(s, Y^-(t, x))) ds, \end{aligned} \quad (3.84)$$

for all stopping time $t \in [0, \tau)$ and all $x \in \mathbb{T}$.

3.4 Energy estimate

We give here an energy estimate for the system (2.6). The corresponding estimate for (2.11) will be proved later in Section 4.1.

Proposition 3.10 (Energy estimate). *Let $(R(t), S(t))_{t < \tau}$ be a regular solution to (2.6) up to a stopping time τ . Let*

$$\mathcal{E}(t) \stackrel{\text{def}}{=} \|(R, S)(t)\|_{L^2(\mathbb{T})}^2 = \int_{\mathbb{T}} (R^2 + S^2)(t, x) dx \quad (3.85)$$

denote the total energy of the system. Then, for all stopping time $t < \tau$ a.s.,

$$\mathcal{E}(t) = \mathcal{E}(0) + 2 \|q\|_{L^1} t + 2 \mathcal{M}(t), \quad (3.86)$$

where $\mathcal{M}(t)$ is a martingale satisfying the bound

$$\mathbb{E} [\mathcal{M}(t \wedge T)^2] \leq 2 q_0 \mathcal{E}(0) T + 2 q_0^2 T^2, \quad (3.87)$$

for all $T > 0$.

Proof of Proposition 3.10. We apply (3.66) (where $\theta \equiv 0$). By integration with respect to $x \in \mathbb{T}$ in (and by the stochastic Fubini Theorem), we obtain (3.86) where $\mathcal{M}(t)$ is the stochastic integral

$$\begin{aligned} \mathcal{M}(t) &\stackrel{\text{def}}{=} \int_0^t \int_{\mathbb{T}} \sum_{k \geq 1} (R + S)(s, x) \sigma_k(x) dx d\beta_k(s) \\ &= \int_0^\infty \mathbf{1}_{s < t} \int_{\mathbb{T}} \sum_{k \geq 1} (R + S)(s, x) \sigma_k(x) dx d\beta_k(s), \end{aligned} \quad (3.88)$$

with variance

$$\mathbb{E} [\mathcal{M}(t)^2] = \mathbb{E} \int_0^t \sum_{k \geq 1} \left| \int_{\mathbb{T}} (R + S)(s, x) \sigma_k(x) dx \right|^2 ds. \quad (3.89)$$

From (3.86) follows the global control on the energy

$$\mathbb{E} [\mathcal{E}(t \wedge T)] \leq \mathcal{E}(0) + 2 q_0 T, \quad (3.90)$$

which in turns (using the Cauchy–Schwarz inequality and (2.2)) gives (3.87). \square

3.5 Solutions with finite explosion time

3.5.1 Some preliminary estimates

In this section, we consider solely the system (2.6). We establish a bound on the square R^2 (resp. S^2) integrated along the characteristic curve X^- (resp. X^+). If applied to the first equation (2.6a) for instance, the bound of S^2 along X^+ can be used to control the negative term $-S^2$ in the right-hand side of (2.6a), which then cannot prevent the blow-up due to the term R^2 . This is exploited in the following section 3.5.2, where solutions with a finite blow-up time are exhibited.

Proposition 3.11 (Control along the opposite characteristic curve). *Let $(R(t), S(t))_{t < \tau}$ be a regular solution to (2.6) up to a stopping time τ . Let X^\pm denote the characteristic curves (3.80).*

There is a constant C depending on the constants c_1 and c_2 in (2.3) only, such that: for all $x_1 \in \mathbb{T}$,

$$\int_0^{t \wedge T} [R^2(s, X^-(s, x_1 + 2c_2 T)) + S^2(s, X^+(s, x_1))] ds \leq C(T+1) \|(R_0, S_0)\|_{L^2(\mathbb{T})}^2 + C q_0 T^2 + M_{x_1}(t \wedge T), \quad (3.91)$$

for all stopping time t such that $t < \tau$ a.s., and all $T > 0$, where $M_{x_1}(t)$ satisfies

$$\mathbb{E} [M_{x_1}(t \wedge T)^2] \leq C q_0 \|(R_0, S_0)\|_{L^2(\mathbb{T})}^2 T + C q_0^2 T^2, \quad \mathbb{E} M_{x_1}(t \wedge T) = 0, \quad (3.92)$$

for all $T > 0$ and all stopping time $t < \tau$ a.s.

Proof of Proposition 3.11. The letter C will denote any constant (possibly varying from line to line) depending solely on c_1 and c_2 . We consider the solution (R, S) as a periodic function defined on the real line. To establish (3.91), we define

$$x_2 \stackrel{\text{def}}{=} x_1 + 2c_2 T. \quad (3.93)$$

The definition (3.93) and the bound of $c(u)$ by c_2 in (2.3) ensure that

$$X^+(t \wedge T, x_1) \leq X^-(t \wedge T, x_2). \quad (3.94)$$

Since $c(u)$ is bounded from below by c_1 , the map $s \mapsto X^+(s, x_1)$ is a C^1 -diffeomorphism from $[0, t]$ onto $[x_1, X^+(t, x_1)]$, we denote by $\tau_{x_1}^+$ its inverse and extend the domain of definition of $\tau_{x_1}^+$ by setting

$$\tau_{x_1}^+(y) \stackrel{\text{def}}{=} \inf \{s \in [0, t \wedge T]; X^+(s, x_1) \geq y\}. \quad (3.95)$$

Similarly, we set

$$\tau_{x_2}^-(y) \stackrel{\text{def}}{=} \inf \{s \in [0, t \wedge T]; X^-(s, x_2) \leq y\}. \quad (3.96)$$

We define $\bar{x} \stackrel{\text{def}}{=} (x_1 + x_2)/2 \in [X^+(t \wedge T, x_1), X^-(t \wedge T, x_2)]$ (see Figure 1), and we use the two

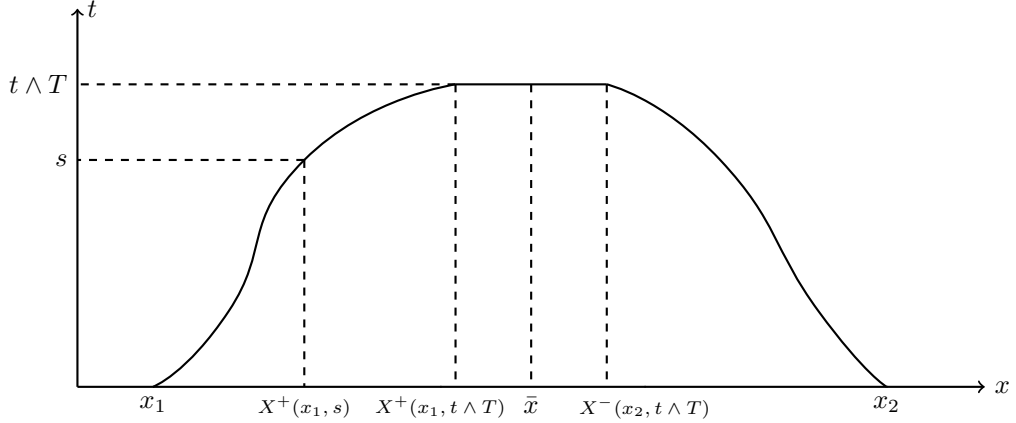


Figure 1: Characteristics.

parametrizations

$$A_1 \stackrel{\text{def}}{=} \{(s, y), x_1 \leq y \leq \bar{x}, 0 \leq s \leq \tau_{x_1}^+(y)\} = \{(s, y), 0 \leq s \leq t \wedge T, X^+(x_1, s) \leq y \leq \bar{x}\},$$

$$A_2 \stackrel{\text{def}}{=} \{(s, y), \bar{x} \leq y \leq x_2, 0 \leq s \leq \tau_{x_2}^-(y)\} = \{(s, y), 0 \leq s \leq t \wedge T, \bar{x} \leq y \leq X^-(x_2, s)\}.$$

Define

$$G \stackrel{\text{def}}{=} \left(\frac{R^2 + S^2 - \tilde{\mathcal{M}}}{c(u) [R^2 - S^2]} \right), \quad \tilde{\mathcal{M}}(t, x) \stackrel{\text{def}}{=} \int_0^t 2 [R(s, x) + S(s, x)] \Phi(x) dW(s). \quad (3.97)$$

Then, using (3.66) (with $\theta \equiv 0$), we obtain

$$\operatorname{div}_{t,x} G \leq 2q_0. \quad (3.98)$$

Integrating now (3.98) on $A_1 \cup A_2$, using the divergence theorem and some changes of variables we obtain

$$\begin{aligned} & 2 \int_0^{t \wedge T} [c(u) R^2(s, X^-(s, x_2)) + c(u) S^2(s, X^+(s, x_1))] ds \\ & \leq 2q_0 T (x_2 - x_1) + \int_{x_1}^{x_2} (R_0^2 + S_0^2) dx - \int_{X^+(t \wedge T, x_1)}^{X^-(t \wedge T, x_2)} (R^2 + S^2) dx \\ & \quad + \int_{x_1}^{\bar{x}} \tilde{\mathcal{M}}(\tau_{x_1}^+(y), y) dy + \int_{\bar{x}}^{x_2} \tilde{\mathcal{M}}(\tau_{x_2}^-(y), y) dy. \end{aligned} \quad (3.99)$$

This inequality (3.99) and the lower bound $c(u) \geq c_1$ imply (3.91), where

$$M_{x_1}(t \wedge T) \stackrel{\text{def}}{=} (2c_1)^{-1} \left[\int_{x_1}^{\bar{x}} \tilde{\mathcal{M}}(\tau_{x_1}^+(y), y) dy + \int_{\bar{x}}^{x_2} \tilde{\mathcal{M}}(\tau_{x_2}^-(y), y) dy \right].$$

Since $\tilde{\mathcal{M}}$ is a martingale and $\tau_{x_1}^+(y)$ and $\tau_{x_2}^-(y)$ are stopping times, then $\mathbb{E}M_{x_1}(t \wedge T) = 0$. Using

$$|M_{x_1}(t \wedge T)|^2 \leq C(T) \int_{\mathbb{T}} \sup_{s \in [0, t \wedge T]} |\tilde{\mathcal{M}}(s, y)|^2 dy \quad (3.100)$$

and the energy inequality (3.90) we obtain (3.92). \square

3.5.2 Singularity formation

Let $u^* \in \mathbb{R}$ be such that $c'(u^*) > 0$. We introduce (using the continuity of $c'(\cdot)$) a length $L > 0$ such that

$$|u - u^*| \leq L \implies c'(u) \geq c'(u^*)/2 > 0. \quad (3.101)$$

Let $\varphi \in C_c^\infty(\mathbb{R})$ be supported in the interval $[1/4, 3/4]$ and non trivial: we assume that $\varphi'(x_0) < 0$, for a given $x_0 \in (1/4, 3/4)$. For $\varepsilon \in (0, 1)$ and α, ν, γ such that

$$\alpha > 1, \quad \nu \in (0, \alpha - 1), \quad \gamma > 1/3, \quad (3.102)$$

we consider the initial data

$$u_0^\varepsilon(x) = u^* + \varepsilon^\alpha \varphi\left(\frac{x}{\varepsilon^{\alpha+\nu+\gamma}}\right), \quad v_0^\varepsilon(x) = 0, \quad (3.103)$$

defined for $x \in [0, 1)$ and extended by periodicity.

Theorem 3.12 (Blow-up in finite time). *Let $(u(t))_{t < \tau}$ be a regular solution up to the explosion time τ of (2.5) with initial data given by (3.103). There exists $\varepsilon_0 > 0$ depending on the constants c_1, c_2 in (2.3) and c_3 , on L , on q_0 and on the profile φ only, such that*

$$\mathbb{P}(\tau \leq \varepsilon^\gamma) \geq 1 - \varepsilon, \quad (3.104)$$

for all $\varepsilon < \varepsilon_0$.

Proof of Theorem 3.12. Set $x_\varepsilon = \varepsilon^{\alpha+\nu+\gamma}x_0$ and

$$R_{\chi^+}(t) = R(t, X^+(t, x_\varepsilon)), \quad S_{\chi^+}(t) = S(t, X^+(t, x_\varepsilon)), \quad (3.105)$$

where X^+ is the characteristic curve defined by (3.80). We will estimate the possible blow-up of R_{χ^+} before a time ε^γ . The letters ε_0 and C will denote any positive constant depending on the constants c_1, c_2 in (2.3), on c_3 , on L , on q_0 and on the profile φ only, with a value possibly changing from line to line. We will however use the specific notation

$$\delta \stackrel{\text{def}}{=} \frac{1}{3} c_1 |\varphi'(x_0)|. \quad (3.106)$$

By (3.83), we have

$$R_{\chi^+}(t) = R_0(x_\varepsilon) + \int_0^t \tilde{c}'(u)(s, X^+(s, x_\varepsilon)) R_{\chi^+}^2(s) ds \quad (3.107a)$$

$$- \int_0^t (\tilde{c}'(u) S^2)(s, X^+(s, x_\varepsilon)) ds + \int_0^t \sum_{k \geq 1} \sigma_k(X^+(s, x_\varepsilon)) d\beta_k(s), \quad (3.107b)$$

for all stopping time $t < \tau$ a.s. Using (2.3), the first term in the right-hand side (3.107a) is bounded as follows:

$$R_0(x_\varepsilon) = -c(u_0(x_\varepsilon)) \varepsilon^{-\nu-\gamma} \varphi'(x_0) \geq 3\delta \varepsilon^{-\nu-\gamma}. \quad (3.108)$$

To bound from below the coefficient $c'(u)$ in factor of $R_{\chi^+}^2(s)$ in (3.107a), we will use (3.101). By (2.9) we have,

$$u(s, X^+(s, x_\varepsilon)) - u_0(x_\varepsilon) = \int_0^s S_{\chi^+}(r) dr, \quad (3.109)$$

and so

$$\sup_{s \leq \varepsilon^\gamma \wedge t} |u(s, X^+(s, x_\varepsilon)) - u_0(x_\varepsilon)|^2 \leq \varepsilon^\gamma \int_0^{\varepsilon^\gamma \wedge t} S_{\chi^+}^2(r) dr. \quad (3.110)$$

We use the estimate (3.91) to get

$$\mathbb{E} \int_0^{\varepsilon^\gamma \wedge t} S_{\chi^+}^2(r) dr \leq C \mathcal{E}(0) + C \varepsilon^{2\gamma}. \quad (3.111)$$

The initial energy satisfies

$$\mathcal{E}(0) = \int_{\mathbb{T}} (R_0^2 + S_0^2)(x) dx \leq 2c_2^2 \varepsilon^{-2\nu-2\gamma} \int_{\mathbb{T}} |\varphi'(\varepsilon^{-\alpha-\nu-\gamma}x)|^2 dx \leq C \varepsilon^{\alpha-\nu-\gamma}, \quad (3.112)$$

so

$$\varepsilon^\gamma \mathbb{E} \int_0^{\varepsilon^\gamma \wedge t} S_{\chi^+}^2(r) dr \leq C (\varepsilon^{\alpha-\nu} + \varepsilon^{3\gamma}), \quad (3.113)$$

and we infer from (3.110) and the Markov inequality that

$$\mathbb{P} \left(\sup_{s \leq \varepsilon^\gamma \wedge t} |u(s, X^+(s, x_\varepsilon)) - u_0(x_\varepsilon)| \geq L/2 \right) \leq C (\varepsilon^{\alpha-\nu} + \varepsilon^{3\gamma}). \quad (3.114)$$

For $\varepsilon < \varepsilon_0$ we have $\|u_0 - u^*\|_{L^\infty} \leq \varepsilon^\alpha \|\varphi\|_{L^\infty} \leq L/2$, and then (3.101) and (3.114) imply

$$\mathbb{P} \left(\inf_{s \leq \varepsilon^\gamma \wedge t} c'(u(s, X^+(s, x_\varepsilon))) \geq c'(u^*)/2 \right) \geq 1 - C (\varepsilon^{\alpha-\nu} + \varepsilon^{3\gamma}). \quad (3.115)$$

Let us now estimate the terms in (3.107b). Using the Doob maximal inequality, Ito's isometry and the definition (2.2) of the variance q , we have

$$\begin{aligned} & \mathbb{P} \left(\sup_{s \leq \varepsilon^\gamma \wedge t} \left| \int_0^s \sum_{k \geq 1} \sigma_k(X^+(s, x_\varepsilon)) d\beta_k(s) \right| \geq \delta \varepsilon^{-\nu-\gamma} \right) \\ & \leq C \delta^{-2} \varepsilon^{2\nu+2\gamma} \mathbb{E} \left| \int_0^{\varepsilon^\gamma \wedge t} \sum_{k \geq 1} \sigma_k(X^+(s, x_\varepsilon)) d\beta_k(s) \right|^2 \leq C \varepsilon^{2\nu+3\gamma}. \end{aligned} \quad (3.116)$$

By (2.3) and $|c'(u)| \leq c_3$, we also have

$$\left| \int_0^{\varepsilon^\gamma \wedge t} (\tilde{c}'(u) S^2)(s, X^+(s, x_\varepsilon)) ds \right| \leq C \int_0^{\varepsilon^\gamma \wedge t} S_{\chi^+}^2(r) dr, \quad (3.117)$$

so (3.113) and the Markov inequality give

$$\mathbb{P} \left(\sup_{s \leq \varepsilon^\gamma \wedge t} \left| \int_0^s (\tilde{c}'(u) S^2)(r, X^+(r, x_\varepsilon)) dr \right| \geq \delta \varepsilon^{-\nu-\gamma} \right) \leq C (\varepsilon^\alpha + \varepsilon^{\nu+3\gamma}). \quad (3.118)$$

We gather the estimates (3.108), (3.115), (3.116), (3.118) to obtain, based on (3.107), and for $\varepsilon < \varepsilon_0$, the inequality

$$R_{\chi^+}(\varepsilon^\gamma \wedge t) \geq \delta \varepsilon^{-\nu-\gamma} + \frac{c'(u^*)}{2} \int_0^{\varepsilon^\gamma \wedge t} R_{\chi^+}^2(s) ds, \quad (3.119)$$

with a probability greater than $1 - C\varepsilon^{\min\{\alpha-\nu, 3\gamma\}}$. Since $\min\{\alpha - \nu, 3\gamma\} > 1$ and since the blow-up in (3.119) occurs before the time $\frac{2\varepsilon^{\nu+\gamma}}{c'(u^*)\delta}$, we can in particular conclude to (3.104). \square

4 Global-in-time regular solutions to the approximate system

This section is devoted to the analysis of the approximate system (2.11). As already explained when introducing the generalized system (3.1), the cut-off function χ_ε is in the class $W_{loc}^{2,\infty}$, but not in C^∞ . Nevertheless, the non-linear estimate (3.27) holds true for $s = 2$ and $F \in W_{loc}^{2,\infty}$, and so the results of Section 3.2 as well, if we limit ourselves to $s = 2$. By the change of unknown

$$P^\varepsilon \stackrel{\text{def}}{=} R^\varepsilon - \Phi^\varepsilon W, \quad Q^\varepsilon \stackrel{\text{def}}{=} S^\varepsilon - \Phi^\varepsilon W, \quad (4.1)$$

we transform the stochastic system (3.1) into a system with random coefficients

$$\begin{aligned} P_t^\varepsilon + c(u^\varepsilon) P_x^\varepsilon &= \tilde{c}'(u^\varepsilon) [|P^\varepsilon|^2 - |Q^\varepsilon|^2 - \chi_\varepsilon(P^\varepsilon + \Phi^\varepsilon W) + 2(P^\varepsilon + \Phi^\varepsilon W) \Theta^\varepsilon] \\ &\quad - (c(u^\varepsilon) \Phi^\varepsilon W)_x, \end{aligned} \quad (4.2a)$$

$$\begin{aligned} Q_t^\varepsilon - c(u^\varepsilon) Q_x^\varepsilon &= \tilde{c}'(u^\varepsilon) [|Q^\varepsilon|^2 - |P^\varepsilon|^2 - \chi_\varepsilon(Q^\varepsilon + \Phi^\varepsilon W) - 2(Q^\varepsilon + \Phi^\varepsilon W) \Theta^\varepsilon] \\ &\quad + (c(u^\varepsilon) \Phi^\varepsilon W)_x, \end{aligned} \quad (4.2b)$$

$$P^\varepsilon(0, \cdot) = R_0 * \rho_\varepsilon, \quad Q^\varepsilon(0, \cdot) = S_0 * \rho_\varepsilon, \quad (4.2c)$$

where

$$\begin{aligned}
& u^\varepsilon(t, x) \\
&= \mathcal{C}^{-1} \left\{ \mathcal{C} \left\{ \int_0^t \left(\frac{P^\varepsilon + Q^\varepsilon}{2} + \Phi^\varepsilon W \right) (s, 0) \, ds + u_0^\varepsilon(0) \right\} + \int_0^x \left[\frac{Q^\varepsilon - P^\varepsilon}{2}(t, y) - \Theta^\varepsilon(t) \right] dy \right\},
\end{aligned} \tag{4.3}$$

with (for $x \in [0, 1)$)

$$\Theta^\varepsilon(t) \stackrel{\text{def}}{=} \int_0^1 \frac{Q^\varepsilon - P^\varepsilon}{2}(t, y) \, dy, \quad u_0^\varepsilon(x) = \mathcal{C}^{-1} \left(\int_0^x \frac{S_0^\varepsilon - R_0^\varepsilon}{2} \, dx \right). \tag{4.4}$$

Then in the definition 3.1 of the solution V stands for $V^\varepsilon = (P^\varepsilon, Q^\varepsilon)$.

Theorem 4.1 (Regular solutions to the approximate system). *Let $\varepsilon > 0$, $R_0, S_0 \in H^2(\mathbb{T})$ such that $\int_{\mathbb{T}} (S_0 - R_0) \, dx = 0$. Then (2.11) admits a regular solution $(V^\varepsilon, \tau^\varepsilon)$ in $H^2(\mathbb{T})$, associated to the initial data $R_0^\varepsilon, S_0^\varepsilon$, and defined up to the explosion time τ^ε .*

Our purpose is to prove that $\tau^\varepsilon = \infty$ a.s. We will use various estimates on the solution $(R^\varepsilon, S^\varepsilon)$ to show that blow-up cannot happen in finite time. Although we work at fixed ε here, it will be useful, when possible, to prove some estimates uniform in ε to prepare the limit $\varepsilon \rightarrow 0$, established in the next Section 5 and Section 6.

Notations. We use the following convention: the letter \mathbf{C} will denote any constant independent on ε , that may otherwise depend on the constant c_0 in (2.29), on the constant c_1 and c_2 in (2.3), on c_3 , on the value of the covariance q in (2.2), and on the L^2 -norms $\|R(0)\|_{L^2(\mathbb{T})}$ and $\|S(0)\|_{L^2(\mathbb{T})}$ only. So

$$\mathbf{C} = \mathbf{C}(c_0, c_1, c_2, c_3, q_0, \|R(0)\|_{L^2(\mathbb{T})}, \|S(0)\|_{L^2(\mathbb{T})}). \tag{4.5}$$

The precise value of the constant may vary from lines to lines. An additional dependence on other parameters will be indicated, as $\mathbf{C}(T)$ in particular to indicate the additional dependence on a given final time T .

4.1 Energy estimates

The aim of this section is to obtain energy estimates of the approximated system as in Proposition 3.10 and Proposition 3.11.

Proposition 4.2 (Energy estimate). *Let $(R^\varepsilon(t), S^\varepsilon(t))_{t < \tau^\varepsilon}$ be a regular solution to (2.11) in $H^2(\mathbb{T})$, up to a stopping time τ^ε . Let*

$$\mathcal{E}^\varepsilon(t) \stackrel{\text{def}}{=} \|(R^\varepsilon, S^\varepsilon)(t)\|_{L^2(\mathbb{T})}^2 = \int_{\mathbb{T}} ((R^\varepsilon)^2 + (S^\varepsilon)^2)(t, x) \, dx \tag{4.6}$$

denote the total energy of the system. Then, for all stopping time $t < \tau^\varepsilon$ a.s.,

$$\begin{aligned}
\mathcal{E}^\varepsilon(t) + 2 \int_0^t \int_{\mathbb{T}} \tilde{c}'(u^\varepsilon) [R^\varepsilon \chi_\varepsilon(R^\varepsilon) + S^\varepsilon \chi_\varepsilon(S^\varepsilon)] \, dx \, dt &= \mathcal{E}^\varepsilon(0) + 2 \|q^\varepsilon\|_{L^1} t + \mathcal{M}^\varepsilon(t) \\
&\leq \mathbf{C} + 2q_0 t + \mathcal{M}^\varepsilon(t),
\end{aligned} \tag{4.7}$$

where $\mathcal{M}^\varepsilon(t)$ is a martingale satisfying the bound

$$\mathbb{E} \left[\left| \sup_{0 \leq s \leq \tau^\varepsilon \wedge T} \mathcal{M}^\varepsilon(s) \right|^p \right] \leq \mathbf{C}(p, T), \tag{4.8}$$

for all $p \geq 1$. Moreover, we have

$$\mathbb{E} \left[\exp \left(\eta \sup_{0 \leq s \leq \tau^\varepsilon \wedge T} \mathcal{E}^\varepsilon(s) \right) \right] \leq \mathfrak{C}(T), \quad (4.9)$$

for all η satisfying the smallness condition $4q_0\eta \leq 1$.

A direct consequence of (4.7)-(4.8) is the bound

$$\mathbb{E} \left[\left| \sup_{0 \leq s \leq \tau^\varepsilon \wedge T} \mathcal{E}^\varepsilon(s) \right|^p \right] \leq \mathfrak{C}(p, T), \quad (4.10)$$

for all $p \geq 1$. We also state the following corollary to Proposition 4.2.

Corollary 4.3 (Bound on the correction terms). *Let $(R^\varepsilon(t), S^\varepsilon(t))_{t < \tau^\varepsilon}$ be a regular solution to (2.11) in $H^2(\mathbb{T})$. The correction Θ^ε defined in (2.14) satisfies the bound*

$$\mathbb{E} \left[\|\Theta^\varepsilon\|_{L^\infty(0, \tau^\varepsilon \wedge T)}^p \right] \leq \mathfrak{C}(p, T) \varepsilon^{1/2}, \quad (4.11)$$

for all $p \geq 1$, while

$$\mathbb{E} \left[\|u_t^\varepsilon - (R^\varepsilon + S^\varepsilon)/2\|_{L^1((0, \tau^\varepsilon \wedge T); L^\infty(\mathbb{T}))}^p \right] \leq \mathfrak{C}(T, p) \varepsilon^{1/2}. \quad (4.12)$$

and

$$\mathbb{E} \left[\|u_t^\varepsilon - (R^\varepsilon + S^\varepsilon)/2\|_{L^\infty((0, \tau^\varepsilon \wedge T) \times \mathbb{T})}^p \right] \leq \mathfrak{C}(T, p), \quad (4.13)$$

for all $p \geq 1$.

Proof of Proposition 4.2. We use the fact that $\xi\chi_\varepsilon(\xi) \geq 0$ and integrate over $x \in \mathbb{T}$ in (3.66) (where $\theta = \chi_\varepsilon$), to obtain (4.7) with

$$\mathcal{M}^\varepsilon(t) \stackrel{\text{def}}{=} 2 \int_0^t \int_{\mathbb{T}} \sum_{k \geq 1} (R^\varepsilon + S^\varepsilon)(s, x) \sigma_k^\varepsilon(x) dx d\beta_k(s). \quad (4.14)$$

Using Itô's isometry, Jensen's inequality and the bound (2.15), we then have

$$\mathbb{E}[\mathcal{M}^\varepsilon(t \wedge T)]^2 \leq 4q_0 \int_0^T \mathbb{E} \int_{\mathbb{T}} [(R^\varepsilon + S^\varepsilon)^2] dx ds \leq 8q_0 \mathcal{E}(0)T + 8q_0^2 T^2 \leq \mathfrak{C}(T).$$

There remains to prove (4.9) and (4.8). Let $\eta \in (0, 1]$. By (4.7), we have

$$\eta \sup_{0 \leq s \leq t \wedge T} \mathcal{E}^\varepsilon(s) \leq \mathfrak{C} + \eta \sup_{0 \leq s \leq t \wedge T} \mathcal{M}^\varepsilon(s). \quad (4.15)$$

Let

$$\langle \mathcal{M}^\varepsilon, \mathcal{M}^\varepsilon \rangle(t) \stackrel{\text{def}}{=} 4 \int_0^t \sum_{k \geq 1} \left| \int_0^1 \sigma_k(R^\varepsilon + S^\varepsilon) dx \right|^2 ds \quad (4.16)$$

denote the quadratic variation of \mathcal{M}^ε . We use an exponential martingale inequality:

$$\mathbb{P} \left(\sup_{0 \leq s \leq t \wedge T} \mathcal{M}^\varepsilon(s) \geq (a + b \langle \mathcal{M}^\varepsilon, \mathcal{M}^\varepsilon \rangle(t \wedge T)) \lambda \right) \leq e^{-2ab\lambda^2}, \quad (4.17)$$

with $\lambda = 1$ to obtain

$$\mathbb{E}[\exp(2\eta Z)] \leq \mathbf{C}(b), \quad Z \stackrel{\text{def}}{=} \sup_{0 \leq s \leq t \wedge T} \mathcal{M}^\varepsilon(s) - b \langle \mathcal{M}^\varepsilon, \mathcal{M}^\varepsilon \rangle(t \wedge T), \quad (4.18)$$

if $\eta \leq b/2$. By (4.15) and the Cauchy–Schwarz inequality, we have then

$$\mathbb{E} \left[\exp \left(\eta \sup_{0 \leq s \leq t \wedge T} \mathcal{E}^\varepsilon(s) \right) \right] \leq \mathbf{C} \left\{ \mathbb{E}[\exp(2\eta Z)] \right\}^{1/2} \left\{ \mathbb{E}[\exp(2b\eta \langle \mathcal{M}^\varepsilon, \mathcal{M}^\varepsilon \rangle(t \wedge T))] \right\}^{1/2}. \quad (4.19)$$

The bound (4.18) therefore gives

$$\mathbb{E} \left[\exp \left(\eta \sup_{0 \leq s \leq t \wedge T} \mathcal{E}^\varepsilon(s) \right) \right] \leq \mathbf{C}(b) \left\{ \mathbb{E}[\exp(2b\eta \langle \mathcal{M}^\varepsilon, \mathcal{M}^\varepsilon \rangle(t \wedge T))] \right\}^{1/2}. \quad (4.20)$$

By the Cauchy–Schwarz inequality, we also have

$$\langle \mathcal{M}^\varepsilon, \mathcal{M}^\varepsilon \rangle(t \wedge T) \leq q_0 \mathcal{E}^\varepsilon(t \wedge T) \leq q_0 \sup_{0 \leq s \leq t \wedge T} \mathcal{E}^\varepsilon(s), \quad (4.21)$$

so

$$\mathbb{E} \left[\exp \left(\eta \sup_{0 \leq s \leq t \wedge T} \mathcal{E}^\varepsilon(s) \right) \right] \leq \mathbf{C}(b) \left\{ \mathbb{E} \left[\exp \left(2bq_0\eta \sup_{0 \leq s \leq t \wedge T} \mathcal{E}^\varepsilon(s) \right) \right] \right\}^{1/2}. \quad (4.22)$$

We can let $t \uparrow \tau^\varepsilon$ and choose $b = 1/(2q_0)$ to conclude to (4.9), under the condition $4q_0\eta \leq 1$. Let us now establish (4.8). By the Burkholder–Davis–Gundy inequality, we have

$$\mathbb{E} \left[\left| \sup_{0 \leq s \leq t \wedge T} \mathcal{M}^\varepsilon(s) \right|^p \right] \leq \mathbf{C}(p) \left\{ \mathbb{E} \left[\langle \mathcal{M}^\varepsilon, \mathcal{M}^\varepsilon \rangle(t \wedge T)^{p/2} \right] \right\}^{1/2}. \quad (4.23)$$

We use (4.21), the bound $|x|^{p/2} \leq \mathbf{C}(p)(1 + e^{|x|/(8q_0)})$ and (4.9) to conclude. \square

Proof of Corollary 4.3. We assume $2 \leq p$. We use the equations (3.67)–(3.68), where $\theta = \chi_\varepsilon$ (this is why we will add an exponent ε to the quantities α and β below). By integration in (3.67), we obtain

$$\Theta^\varepsilon(t) = \int_0^t \alpha^\varepsilon(s) e^{-\int_s^t \beta^\varepsilon(\sigma) d\sigma} ds, \quad (4.24)$$

which gives

$$\|\Theta^\varepsilon\|_{L^\infty(0, t \wedge T)}^p \leq \int_0^{t \wedge T} |\alpha^\varepsilon(s)|^p ds \left(\int_0^{t \wedge T} e^{p' \int_s^t |\beta^\varepsilon(\sigma)| d\sigma} ds \right)^{p/p'}, \quad (4.25)$$

where p' is the conjugate exponent to p , and then

$$\begin{aligned} & \mathbb{E} \left[\|\Theta^\varepsilon\|_{L^\infty(0, t \wedge T)}^p \right] \\ & \leq \left\{ \mathbb{E} \left[\int_0^{t \wedge T} |\alpha^\varepsilon(s)|^{2p} ds \right] \right\}^{1/2} \left\{ \mathbb{E} \left[\left(\int_0^{t \wedge T} \exp \left(p' \int_s^t |\beta^\varepsilon(\sigma)| d\sigma \right) ds \right)^{(2p)/p'} \right] \right\}^{1/2}. \end{aligned} \quad (4.26)$$

We estimate α^ε in two manners. Let us introduce the notation

$$\bar{\mathcal{M}}^\varepsilon(t) = \sup_{0 \leq s \leq t} \mathcal{M}^\varepsilon(s), \quad (4.27)$$

where \mathcal{M}^ε is the martingale (4.14). First we have

$$|\alpha^\varepsilon(s)| \leq c_3 \int_0^1 |\chi_\varepsilon(R^\varepsilon) + \chi_\varepsilon(S^\varepsilon)|(s, x) dx \leq c_3 \mathcal{E}(s) \leq \mathbf{C}(T) + c_3 \bar{\mathcal{M}}^\varepsilon(s). \quad (4.28)$$

We can also use the bound $\chi_\varepsilon(R) \leq R\chi_\varepsilon(R)\varepsilon$ and the second term in the left-hand side of (4.7) to obtain

$$\int_0^{t \wedge T} |\alpha^\varepsilon(s)| ds \leq \varepsilon \int_0^{t \wedge T} \int_0^1 \tilde{c}'(u^\varepsilon) [R^\varepsilon \chi_\varepsilon(R^\varepsilon) + S^\varepsilon \chi_\varepsilon(S^\varepsilon)] dx ds \leq \varepsilon [\mathbf{C}(T) + \bar{\mathcal{M}}^\varepsilon(t \wedge T)]. \quad (4.29)$$

We combine (4.28) with (4.29) to obtain

$$\int_0^{t \wedge T} |\alpha^\varepsilon(s)|^{2p} ds \leq \mathbf{C}(T) \left[1 + |\bar{\mathcal{M}}^\varepsilon(s)|^{2p} \right] \varepsilon, \quad (4.30)$$

and the bound (4.8) gives

$$\mathbb{E} \left[\int_0^{t \wedge T} |\alpha^\varepsilon(s)|^{2p} ds \right] \leq \mathbf{C}(p, T) \varepsilon. \quad (4.31)$$

To estimate the remaining term in (4.26), we first note that

$$\left(\int_0^{t \wedge T} \exp \left(p' \int_s^{t \wedge T} |\beta^\varepsilon(\sigma)| d\sigma \right) ds \right)^{(2p)/p'} \leq \mathbf{C}(p, T) \exp \left(\mathbf{C}(p, T) \sup_{0 \leq s \leq t \wedge T} \mathcal{E}(\sigma)^{1/2} \right), \quad (4.32)$$

and

$$\mathbf{C}(p, T) \sup_{0 \leq s \leq t \wedge T} \mathcal{E}(\sigma)^{1/2} \leq \mathbf{C}(p, T) + (8q_0)^{-1} \sup_{0 \leq s \leq t \wedge T} \mathcal{E}(s). \quad (4.33)$$

Using (4.9), (4.32) and (4.33) yields,

$$\mathbb{E} \left[\left(\int_0^{t \wedge T} \exp \left(p' \int_s^{t \wedge T} |\beta^\varepsilon(\sigma)| d\sigma \right) ds \right)^{(2p)/p'} \right] \leq \mathbf{C}(p, T) \quad (4.34)$$

We combine (4.31) and (4.34) with (4.26) to conclude to (4.11) with t instead of τ^ε . Then we let $t \uparrow \tau^\varepsilon$ to get (4.11). The estimates (4.12) and (4.13) are a consequence of (3.69) (where all quantities are indexed by ε) and the bounds

$$\mathbb{E} \left[\|\zeta^\varepsilon\|_{L^1((0, \tau^\varepsilon \wedge T) \times \mathbb{T})}^p \right] \leq \mathbf{C}(T, p) \varepsilon^{1/2}, \quad (4.35)$$

and

$$\mathbb{E} \left[\|\zeta^\varepsilon\|_{L^\infty((0, \tau^\varepsilon \wedge T); L^1(\mathbb{T}))}^p \right] \leq \mathbf{C}(T, p), \quad (4.36)$$

where

$$\zeta^\varepsilon = \tilde{c}'(u^\varepsilon) \left[\frac{\chi_\varepsilon(R^\varepsilon) - \chi_\varepsilon(S^\varepsilon)}{2} + (R^\varepsilon + S^\varepsilon)\Theta^\varepsilon \right]. \quad (4.37)$$

This follows easily from the bounds (4.28)-(4.29) and (4.7)-(4.8)-(4.11). \square

Let $X^{\varepsilon, \pm}$ denote the flow associated to the vector field $\pm c(u^\varepsilon(t, \cdot))$ (characteristic curves):

$$\frac{d}{dt} X^{\varepsilon, \pm}(t, x) = \pm c(u^\varepsilon(t, X^{\varepsilon, \pm}(t, x))), \quad X^{\varepsilon, \pm}(0, x) = x, \quad x \in \mathbb{T}, \quad t \in [0, \tau^\varepsilon]. \quad (4.38)$$

In parallel with Proposition 3.11, we have the following result.

Proposition 4.4 (Control along the opposite characteristic curve). *Consider a regular solution $(R^\varepsilon(t), S^\varepsilon(t))_{t < \tau^\varepsilon}$ to (2.11) up to a stopping time τ^ε . Let $X^{\varepsilon, \pm}$ denote the characteristic curves (4.38). There is a constant \mathbf{C} as in (4.5), such that: for all $x_1 \in \mathbb{T}$ and $x_2 = x_1 + 2c_2 T$,*

$$\int_0^{t \wedge T} [(R^\varepsilon)^2(s, X^-(s, x_2)) + (S^\varepsilon)^2(s, X^+(s, x_1))] ds \leq \mathbf{C} + \mathbf{C} M_{x_1}^\varepsilon(t \wedge T), \quad (4.39)$$

for all stopping time t such that $t < \tau^\varepsilon$ a.s., where $M_{x_1}^\varepsilon(t)$ satisfies the bound

$$\mathbb{E} \left[\sup_{s \in [0, \tau^\varepsilon \wedge T]} \sup_{x_1} |M_{x_1}^\varepsilon(s)|^2 \right] \leq \mathbf{C}(T). \quad (4.40)$$

Proof. The proof of Proposition 4.4 is similar to the one of Proposition 3.11. The constant in (3.100) is independent on x_1 , then we obtain (4.40) by following the proof of (3.92) and using (4.10). \square

4.2 Global solutions

We can now use the estimates established in Section 4.1 to show that the regular solutions of (2.11) given by Theorem 4.1 exist globally in time almost surely.

Theorem 4.5. *Let $\varepsilon > 0$, $R_0, S_0 \in L^2(\mathbb{T})$ such that $\int_{\mathbb{T}} (S_0 - R_0) dx = 0$. Let τ^ε be the explosion time of the solution of (2.11) given by Theorem 4.1. Then $\mathbb{P}(\tau^\varepsilon = \infty) = 1$.*

Proof of Theorem 4.5. By (2.16), (2.18) (satisfied by $\Phi^\varepsilon W$) and the bounds (4.40) and (4.11), there exists $\bar{\Omega}^\varepsilon$ such that $\mathbb{P}(\bar{\Omega}^\varepsilon) = 1$ and, for all $\omega \in \bar{\Omega}^\varepsilon$, for all $T > 0$,

$$\Phi^\varepsilon W \in C([0, T], H^3(\mathbb{T})), \quad M^\varepsilon \in L^\infty([0, \tau^\varepsilon \wedge T] \times \mathbb{R}), \quad \Theta^\varepsilon \in L^\infty(0, \tau^\varepsilon \wedge T). \quad (4.41)$$

Let $\omega \in \bar{\Omega}^\varepsilon$ be fixed. We have to establish the bound

$$\sup_{t \in [0, \tau^\varepsilon \wedge T]} \|V^\varepsilon(t)\|_{L^\infty(\mathbb{T})} < \infty, \quad (4.42)$$

for all $T > 0$. We remark that $\xi^2 - \chi_\varepsilon(\xi) \geq 0$, so the equations (4.2) imply

$$P_t^\varepsilon + c(u^\varepsilon) P_x^\varepsilon \geq -\tilde{c}'(u^\varepsilon) [(S^\varepsilon)^2 + 2([P^\varepsilon]^- + |\Phi^\varepsilon W|) |\Theta^\varepsilon|] - c(u^\varepsilon) (\Phi^\varepsilon W)_x, \quad (4.43)$$

where $s^- = \max(-s, 0)$ is the negative part of a real $s \in \mathbb{R}$. We integrate along the characteristics in (4.43) (cf. (3.83) for instance) and use (4.39) to get

$$\|[P^\varepsilon]^- (t)\|_{L^\infty(\mathbb{T})} \leq \mathbf{C}(\omega, \varepsilon, T) \left(1 + \int_0^t \|[P^\varepsilon]^- (s)\|_{L^\infty(\mathbb{T})} ds \right), \quad (4.44)$$

for $t \in [0, \tau^\varepsilon \wedge T]$, which gives

$$\sup_{t \in [0, \tau^\varepsilon \wedge T]} \|[P^\varepsilon]^- (t)\|_{L^\infty(\mathbb{T})} \leq \mathbf{C}(\omega, \varepsilon, T), \quad (4.45)$$

by the Grönwall lemma. The bounds on the positive part of P^ε is established in a similar manner: it follows from (4.2) and the bound $R^2 - \chi_\varepsilon(R) \leq 2\varepsilon^{-1}R^+$ that

$$P_t^\varepsilon + c(u^\varepsilon) P_x^\varepsilon \leq \mathbf{C}(\omega, \varepsilon, T) [1 + [P^\varepsilon]^+]. \quad (4.46)$$

By the comparison principle (or integration along the characteristics) and the Grönwall lemma, (4.46) implies

$$\sup_{t \in [0, \tau^\varepsilon \wedge T]} \|[P^\varepsilon]^+(t)\|_{L^\infty(\mathbb{T})} \leq \mathbf{C}(\omega, \varepsilon, T). \quad (4.47)$$

We have of course the bounds similar to (4.45)-(4.47) on Q^ε , so (4.42) is indeed satisfied. \square

4.3 L^{3^-} estimates

In this section we will establish the following almost L^3 estimates with weights on R^ε and S^ε .

Proposition 4.6 (Estimates in $L_{t,x}^{3^-}$). *Let $\alpha \in [0, 1)$. We have the bound*

$$\sup_\varepsilon \mathbb{E} \int_0^T \int_{\mathbb{T}} c'(u^\varepsilon) [|R^\varepsilon|^{2+\alpha} + |S^\varepsilon|^{2+\alpha}] dx dt \leq \mathbf{C}(T, \alpha), \quad (4.48)$$

for all $T \geq 0$.

Proof of Proposition 4.6. Let $h: \mathbb{R} \rightarrow \mathbb{R}$ be the non-decreasing function of class $W^{2,\infty}$ defined as

$$h(R) \stackrel{\text{def}}{=} \int_0^R (r^2 + 1)^{\frac{\alpha}{2}} dr.$$

The function h satisfies the growth conditions

$$|h(R)| \leq C(1 + |R|^2)^{\frac{\alpha+1}{2}}, \quad |h'(R)| \leq C(1 + |R|)^\alpha, \quad |h''(R)| \leq C, \quad R \in \mathbb{R}. \quad (4.49)$$

We sum up the two equations (3.62) and (3.63) with $\theta = \chi_\varepsilon$ and integrate on \mathbb{T} to obtain (using $h'(R)\chi_\varepsilon(R) \geq 0$ and the energy estimate (4.10))

$$\mathbb{E} \int_0^T \int_{\mathbb{T}} c'(u^\varepsilon) \Delta_1(R^\varepsilon, S^\varepsilon) dx dt \leq \mathbf{C}(T) + \mathbf{C}(T) \mathbb{E} \int_0^T |\Theta^\varepsilon(t)| \int_{\mathbb{T}} (|R^\varepsilon|^2 + |S^\varepsilon|^2)^{\frac{1+\alpha}{2}} dx dt, \quad (4.50)$$

where

$$\Delta_1(R, S) := (R - S)(B_h(R, S) - B_h(S, R)), \quad (4.51)$$

and $B_h(R, S)$ is defined in (3.64). By Hölder's inequality, (4.11) and (4.10), we can bound the last term in (4.50) and get

$$\mathbb{E} \int_0^T \int_{\mathbb{T}} c'(u^\varepsilon) \Delta_1(R^\varepsilon, S^\varepsilon) dx dt \leq \mathbf{C}(T, \alpha). \quad (4.52)$$

Similarly, starting from (3.77)-(3.78) and using the estimates (4.13) and (4.10) we can establish the bound

$$\mathbb{E} \int_0^T \int_{\mathbb{T}} \frac{c'(u^\varepsilon)}{c(u^\varepsilon)} \Delta_2(R^\varepsilon, S^\varepsilon) dx dt \leq \mathbf{C}(T, \alpha), \quad (4.53)$$

where

$$\Delta_2(R, S) := (R + S)(F_h(R, S) + F_h(S, R)), \quad (4.54)$$

and $F_h(R, S)$ is defined in (3.79).

We claim that

$$(R - S)^2 (|R|^\alpha + |S|^\alpha) \leq \mathfrak{C}(\alpha) \Delta_1(R, S), \quad (R + S)^2 (|R|^\alpha + |S|^\alpha) \leq \mathfrak{C}(\alpha) \Delta_2(R, S). \quad (4.55)$$

Summing up the two estimates (4.52) and (4.53) and using (2.3) and (4.55) one then gets (4.48). Let us now give a proof of the first inequality in (4.55) (the second inequality is obtained by similar arguments). Using the definition of $B_h(R, S)$ and the monotonicity of $\xi \mapsto \xi/h'(\xi)$ we obtain

$$\begin{aligned} \Delta_1(R, S) &\geq \frac{1-\alpha}{1+\alpha} (R - S) [R h'(S) - S h'(R)] + (R - S) [2h(R) - R h'(R) - 2h(S) + S h'(S)] \\ &= \frac{1-\alpha}{1+\alpha} (R - S)^2 [h'(R) + h'(S)] + 2(R - S) \left[h(R) - \frac{1}{1+\alpha} R h'(R) - h(S) + \frac{1}{1+\alpha} S h'(S) \right]. \end{aligned}$$

The result follows from the monotonicity of $\xi \mapsto h(\xi) - \xi h'(\xi)/(\alpha + 1)$ and from the bound from below $h'(\xi) \geq |\xi|^\alpha$. \square

4.4 One-sided entropy estimates

The central result in this section is Proposition 4.8. We will need first a control on $c'(u(t))$ for small times. More precisely, recall that

$$c_0 = \inf_{x \in \mathbb{T}} c'(u_0(x)) > 0 \quad (4.56)$$

by (2.29). For ε sufficiently small, we have then

$$\inf_{x \in \mathbb{T}} c'(u^\varepsilon(0, x)) \geq c_0/2. \quad (4.57)$$

Consider the stopping time

$$\bar{t}^\varepsilon \stackrel{\text{def}}{=} \inf \left\{ t \geq 0, \inf_{x \in \mathbb{T}} c'(u^\varepsilon(t, x)) \leq c_0/4 \right\}. \quad (4.58)$$

We have the following estimate on \bar{t}^ε .

Lemma 4.7. *We have*

$$\mathbb{E} [(\bar{t}^\varepsilon)^{-p}] \leq \mathfrak{C}(p), \quad (4.59)$$

for all $p \geq 1$.

Proof of Lemma 4.7. By continuity of the map $u \mapsto c'(u)$ and by (4.57), there exists $\eta > 0$ such that

$$\|v - u^\varepsilon(0, \cdot)\|_{L^\infty(\mathbb{T})} \leq \eta \quad \implies \quad \inf_{x \in \mathbb{T}} c'(v(x)) > c_0/4, \quad (4.60)$$

for any $v \in C(\mathbb{T})$. In particular,

$$\eta \leq \sup_{s \leq \bar{t}^\varepsilon} \|u^\varepsilon(s) - u_0^\varepsilon\|_{L^\infty(\mathbb{T})}. \quad (4.61)$$

We apply the estimate

$$\|v\|_{L^\infty(\mathbb{T})}^2 \leq \|v\|_{L^2(\mathbb{T})} \|\partial_x v\|_{L^2(\mathbb{T})},$$

valid for any $v \in H^1(\mathbb{T})$, to $v = u^\varepsilon(t) - u_0^\varepsilon$. Writing

$$u^\varepsilon(t) - u_0^\varepsilon = \int_0^t u_t^\varepsilon(s) \, ds \quad (4.62)$$

yields

$$\|u^\varepsilon(t) - u_0^\varepsilon\|_{L^\infty(\mathbb{T})}^2 \leq 2t \|u_t^\varepsilon\|_{L^\infty((0,t),L^2(\mathbb{T}))} \|\partial_x u^\varepsilon\|_{L^\infty((0,t);L^2(\mathbb{T}))}, \quad (4.63)$$

and, by (4.61)

$$(\bar{t}^\varepsilon)^{-1} \leq (\bar{t}^\varepsilon \wedge T)^{-1} \leq 2\eta^{-1} \|u_t^\varepsilon\|_{L^\infty((0,T),L^2(\mathbb{T}))} \|\partial_x u^\varepsilon\|_{L^\infty((0,T);L^2(\mathbb{T}))}, \quad (4.64)$$

where T is an arbitrary positive time (we may take $T = 1$). We use the identities (3.59)-(3.69) and the bounds (4.11), (4.13), (4.10) to conclude to (4.59). \square

We can now state the following one-sided estimate, where $s^- = -\min\{s, 0\}$ denotes the negative part of a real s .

Proposition 4.8 (Oleinik's estimate). *We have*

$$\mathbb{E} \left[\left\| (R^\varepsilon)^- (\tau \wedge T, \cdot) \right\|_{L^\infty(\mathbb{T})}^p \right] + \mathbb{E} \left[\left\| (S^\varepsilon)^- (\tau \wedge T, \cdot) \right\|_{L^\infty(\mathbb{T})}^p \right] \leq \mathfrak{C}(p, T) \left(1 + \mathbb{E} \left[\frac{1}{(\tau \wedge T)^p} \right] \right), \quad (4.65)$$

for all $p \in [1, 2]$, for any time $T > 0$ and any stopping time $\tau > 0$ a.s.

Proof of Proposition 4.8. For simplicity, we will drop the superscript ε on the quantities R, S , etc. in the proof. Let us consider the equation for $P = R - \Phi W$: use

$$R^2 = (P + \Phi W)^2 \geq \frac{1}{2} P^2 - |\Phi W|^2, \quad (4.66)$$

and

$$2R\Theta = 2(P + \Phi W)\Theta \geq -\frac{1}{4} P^2 - C(|\Theta|^2 + |\Phi W|^2), \quad (4.67)$$

to obtain

$$P_t + c(u) P_x \geq \frac{c'(u)}{4} P^2 - CS^2 - C\chi_\varepsilon(P + \Phi W) - CK_T^2, \quad (4.68)$$

where

$$K_T := \left\{ \sup_{t \in [0, T]} \left(\|\Phi W(t)\|_{W^{1,\infty}(\mathbb{T})}^2 + |\Theta(t)|^2 \right) \right\}^{1/2}. \quad (4.69)$$

We deduce from (4.68) that

$$H_t + c(u) H_x \leq -\frac{c'(u)}{8} H^2 + CS^2 + CK_T^2, \quad (4.70)$$

where $H \stackrel{\text{def}}{=} (P + K_T)^-$. Integrate along the characteristics: for $t \leq \tau \wedge T$,

$$H(t, X^+(t, x)) \leq H(s, X^+(s, x)) - \int_s^t \frac{c'(u)}{8} H^2(r, X^+(r, x)) \, dr + A_T. \quad (4.71)$$

where

$$A_T \stackrel{\text{def}}{=} C \int_0^{\tau \wedge T} S^2(s, X^+(s, x)) \, ds + CT K_T^2. \quad (4.72)$$

On the interval $[0, \bar{t}^\varepsilon]$, assuming $\tau \wedge T \geq \bar{t}^\varepsilon$, we have

$$H(\sigma, X^+(\sigma, x)) \leq H(s, X^+(s, x)) - \frac{c_0}{32} \int_s^\sigma H^2(r, X^+(r, x)) dr + A_T, \quad [s, \sigma] \subset [0, \bar{t}^\varepsilon]. \quad (4.73)$$

We deduce from (4.73) that

$$H(t, X^+(t, x)) \leq \frac{32}{c_0 t} + A_T, \quad \forall t \in (0, \bar{t}^\varepsilon]. \quad (4.74)$$

Indeed, assume that (4.74) is not realized, and let $t_1 \in (0, \bar{t}^\varepsilon]$ be such that

$$H(t_1, X^+(t_1, x)) > \frac{32}{c_0 t_1} + A_T. \quad (4.75)$$

Define

$$t_0 \stackrel{\text{def}}{=} \sup \left\{ t \in (0, t_1]; H(t, X^+(t, x)) \leq \frac{32}{c_0 t} \right\}. \quad (4.76)$$

We have $t_0 > 0$ since $H(0, x) < \infty$ and, by continuity,

$$H(t_0, X^+(t_0, x)) = \frac{32}{c_0 t_0}, \quad H(t, X^+(t, x)) > \frac{32}{c_0 t}, \quad \forall t \in (t_0, t_1]. \quad (4.77)$$

Taking $s = t_0$ and $\sigma = t_1$ in (4.73) therefore gives

$$H(t_1, X^+(t_1, x)) \leq \frac{32}{c_0 t_0} + \left[\frac{32}{c_0 r} \right]_{t_0}^{t_1} + A_T = \frac{32}{c_0 t_1} + A_T, \quad (4.78)$$

which is in contradiction with (4.75). Whatever the value of τ may be, infer from (4.74) that, for all $x \in \mathbb{T}$,

$$H(\tau \wedge T \wedge \bar{t}^\varepsilon, x) \leq \frac{32}{c_0 \tau \wedge T \wedge \bar{t}^\varepsilon} + A_T. \quad (4.79)$$

If $\tau \wedge T > \bar{t}^\varepsilon$, we complete the argument as follows: we use (4.71) between the times \bar{t}^ε and $\tau \wedge T$ and exploit (4.74) with $t = \bar{t}^\varepsilon$ to get

$$H(\tau \wedge T, x) \leq \frac{32}{c_0 \bar{t}^\varepsilon} + 2A_T. \quad (4.80)$$

Finally, we deduce from (4.79) and (4.80) that

$$H(\tau \wedge T \wedge \bar{t}^\varepsilon, x) \leq \frac{32}{c_0 \tau \wedge T \wedge \bar{t}^\varepsilon} + 2A_T, \quad (4.81)$$

and thus, using (2.17), (4.11) and Proposition 4.4 to estimate $\mathbb{E}[A_T^p]$,

$$\mathbb{E} \left[\|H(\tau \wedge T, \cdot)\|_{L^\infty(\mathbb{T})}^p \right] \leq \mathfrak{C}(p, T) \left(1 + \mathbb{E} \left[\frac{1}{(\tau \wedge T)^p} \right] \right). \quad (4.82)$$

Then we conclude with the inequality $R^- \leq H + 2K_T$. \square

5 First limit elements

It follows from (3.62)-(3.63) that $(R^\varepsilon, S^\varepsilon)$ satisfies the following set of equations

$$\begin{aligned} dh(R^\varepsilon) + (c(u^\varepsilon) h(R^\varepsilon))_x dt &= h'(R^\varepsilon) \Phi^\varepsilon dW + 2\tilde{c}'(u^\varepsilon) \Theta^\varepsilon (R^\varepsilon h'(R^\varepsilon) - 2h(R^\varepsilon)) dt \\ &= \tilde{c}'(u^\varepsilon) [(S^\varepsilon - R^\varepsilon) B_h(R^\varepsilon, S^\varepsilon) - h'(R^\varepsilon) \chi_\varepsilon(R^\varepsilon) + \frac{1}{2} q^\varepsilon h''(R^\varepsilon)] dt, \end{aligned} \quad (5.1)$$

and

$$\begin{aligned} dh(S^\varepsilon) - (c(u^\varepsilon)h(S^\varepsilon))_x dt + h'(S^\varepsilon)\Phi^\varepsilon dW - 2\check{c}'(u^\varepsilon)\Theta^\varepsilon(S^\varepsilon h'(S^\varepsilon) - 2h(S^\varepsilon)) dt \\ = \check{c}'(u^\varepsilon) \left[(R^\varepsilon - S^\varepsilon) B_h(S^\varepsilon, R^\varepsilon) - h'(S^\varepsilon) \chi_\varepsilon(S^\varepsilon) + \frac{1}{2} q^\varepsilon h''(S^\varepsilon) \right] dt. \end{aligned} \quad (5.2)$$

In this section, we use the estimates derived in the previous part to obtain a first set of limit elements, including random Young measures (see Proposition 5.6). Then we pass to the limit in (5.1)-(5.2) to obtain a first set of limit equations with defect measures. In the next section 6, we will complete the analysis of the limit elements, prove that the Young measures reduce to Dirac masses, and establish the limit equation for good.

The plan of this section is the following one: some elements on random Young measures are given in Section 5.1. Taking the limit of the stochastic integral in (5.1)-(5.2) is done as usual by strengthening the probabilistic mode of convergence, passing from convergence in law to a.s.-convergence, at the expense of a modification of the probability spaces and measures. We use the Skorokhod–Jakubowski theorem to that effect, explained in Section 5.2. Compactness results are established in Section 5.3 and the limit equation with defect measure derived in Section 5.4.

5.1 Young measures

Let $T > 0$ and let $\mathcal{P}_T([0, T] \times \mathbb{T} \times \mathbb{R}^2)$ denote the space of (non-negative) measures μ on $[0, T] \times \mathbb{T} \times \mathbb{R}^2$ having total measure

$$\mu([0, T] \times \mathbb{T} \times \mathbb{R}^2) = T. \quad (5.3)$$

We start this section by defining Young measures on $[0, T] \times \mathbb{T}$.

Definition 5.1. *Let $\mu \in \mathcal{P}_T([0, T] \times \mathbb{T} \times \mathbb{R}^2)$, we say that μ is a Young measure on $[0, T] \times \mathbb{T}$ (with state space \mathbb{R}^2) if for almost all $(t, x) \in [0, T] \times \mathbb{T}$ there exists a probability measure $\mu_{t,x}$ on \mathbb{R}^2 such that*

- *the map $(t, x) \mapsto \langle \mu_{t,x}, \varphi \rangle$ is measurable for any $\varphi \in C_b(\mathbb{R}^2)$*
- *$\langle \mu, \psi \rangle = \int_{[0, T] \times \mathbb{T}} \int_{\mathbb{R}^2} \psi(t, x, \xi) d\mu_{t,x}(\xi, \eta) dx dt$ for any $\psi \in C_b([0, T] \times \mathbb{T} \times \mathbb{R}^2)$.*

We denote by \mathcal{Y} the space of Young measures on $[0, T] \times \mathbb{T}$ with state space \mathbb{R}^2 .

The space \mathcal{Y} is equipped with the topology of weak convergence in $\mathcal{P}_T([0, T] \times \mathbb{T} \times \mathbb{R}^2)$, i.e.

$$\mu^\varepsilon \rightarrow \mu \quad \text{iff} \quad \langle \mu^\varepsilon, \psi \rangle \rightarrow \langle \mu, \psi \rangle \quad \forall \psi \in C_b([0, T] \times \mathbb{T} \times \mathbb{R}^2). \quad (5.4)$$

The weak convergence in the space $\mathcal{P}_T([0, T] \times \mathbb{T} \times \mathbb{R}^2)$ (and also \mathcal{Y}) coincides with the convergence with respect to the Prokhorov metric [Bil99, p.72]. Then, \mathcal{Y} is a Polish (separable and completely metrizable) space, [Bil99, Theorem 6.8]. In order to obtain tightness in the space \mathcal{Y} , one needs to identify some compact subsets of \mathcal{Y} . For that purpose, we recall the following result (see Proposition 4.1 in [BV19])

Proposition 5.2. *Let $f \in C(\mathbb{R}^2)$ be a non-negative function satisfying*

$$\lim_{|(\xi, \eta)| \rightarrow \infty} f(\xi, \eta) = \infty. \quad (5.5)$$

Let also $C > 0$, then the set

$$\left\{ \mu \in \mathcal{Y}, \int_{[0, T] \times \mathbb{T} \times \mathbb{R}^2} f(\xi, \eta) d\mu(t, x, \xi, \eta) \leq C \right\} \quad (5.6)$$

is a compact subset of \mathcal{Y} .

Remark 5.1 (Convergence in \mathcal{Y}). If μ and μ_1, μ_2, \dots are Young measures in the set (5.6), and if (5.5) is satisfied, then to check the weak convergence $\mu_n \rightarrow \mu$ (which corresponds to convergence against continuous and bounded functions on $\mathbb{T} \times [0, T] \times \mathbb{R}^2$, cf. (5.4)), it is sufficient to check convergence against functions in $C_0(\mathbb{T} \times [0, T] \times \mathbb{R}^2)$ (functions which tend to 0 at infinity). Actually, we can even conclude that

$$\int_0^T \int_{\mathbb{T}} \int_{\mathbb{R}^2} \varphi(t, x) g(\xi, \eta) d\mu_n \rightarrow \int_0^T \int_{\mathbb{T}} \int_{\mathbb{R}^2} \varphi(t, x) g(\xi, \eta) d\mu, \quad (5.7)$$

for all $\varphi \in C([0, T] \times \mathbb{T})$ and $g \in C(\mathbb{R}^2)$ satisfying $|g(\xi, \eta)| \leq C f(\xi, \eta)^\delta$, where $\delta \in [0, 1)$, since for any μ in the set (5.6), and for any cut-off function $\bar{\chi}_R \in C([0, \infty))$ such that $\chi_R \equiv 1$ on the interval $[0, R]$, we have the uniform bound

$$\left| \int_0^T \int_{\mathbb{T}} \int_{\mathbb{R}^2} \varphi(t, x) [g(\xi, \eta) - g(\xi, \eta) \bar{\chi}_R(f(\xi, \eta))] d\mu \right| \leq \|\varphi\|_{C([0, T] \times \mathbb{T})} C R^{\delta-1}. \quad (5.8)$$

In the same spirit as Remark 5.1, and following the proof of Proposition 4.4 in [BV19] we can establish the following result.

Proposition 5.3. *Let $[t_1, t_2] \subset [0, T]$ and let $(a_n)_n$ be a sequence of non-negative $C([t_1, t_2] \times \mathbb{T})$ -valued random variables. Let also $(\mu_n)_n$ be a sequence of random Young measures (\mathcal{Y} -valued random variables) satisfying*

$$C(f) \stackrel{\text{def}}{=} \sup_n \mathbb{E} \left(\int_{[t_1, t_2] \times \mathbb{T} \times \mathbb{R}^2} a_n(t, x) f(\xi, \eta) d\mu_n(t, x, \xi, \eta) \right)^\gamma < \infty, \quad (5.9)$$

where $f \in C(\mathbb{R}^2)$ is a non-negative function and $\gamma > 0$. Assume that almost-surely $\mu_n \rightarrow \mu$ in \mathcal{Y} , and $a_n \rightarrow a$ in $C([t_1, t_2] \times \mathbb{T})$, then

$$\mathbb{E} \left(\int_{[t_1, t_2] \times \mathbb{T} \times \mathbb{R}^2} a(t, x) f(\xi, \eta) d\mu(t, x, \xi, \eta) \right)^\gamma \leq C(f). \quad (5.10)$$

Moreover, for all $g \in C(\mathbb{R}^2)$ satisfying $|g(\xi, \eta)| \leq f(\xi, \eta)^\delta$ with $\delta \in (0, 1)$ and all $\varphi \in L^{\frac{1}{1-\delta}}([t_1, t_2] \times \mathbb{T})$ we have

$$\lim_n \mathbb{E} \left| \int_{[t_1, t_2] \times \mathbb{T} \times \mathbb{R}^2} a_n(t, x) \varphi(t, x) g(\xi, \eta) d\mu_n(t, x, \xi, \eta) - \int_{[t_1, t_2] \times \mathbb{T} \times \mathbb{R}^2} a(t, x) \varphi(t, x) g(\xi, \eta) d\mu(t, x, \xi, \eta) \right|^\gamma = 0. \quad (5.11)$$

Proof of Proposition 5.3. The proof follows from [BV19, Proposition 4.4] applied to the measure $\mu'_n := a_n \mu_n$. Although μ'_n may not have the required total mass (5.3), the arguments hold true. Note also that [BV19, Proposition 4.4] addresses the case $\gamma = 1$ only, but the arguments can be adapted easily if one uses Egorov's theorem directly, instead of Vitali's theorem. \square

5.2 Quasi-Polish spaces and the Skorokhod–Jakubowski theorem

In order to study the limit $\varepsilon \rightarrow 0$, we need to prove the compactness of $(R^\varepsilon, S^\varepsilon)$ in the non-metrizable space $C([0, T], L_w^2(\mathbb{T}))$, where L_w^2 is the L^2 space equipped with its weak topology. Therefore, the classical Skorokhod theorem cannot be used. We use instead the generalized Skorokhod–Jakubowski theorem [J97].

Definition 5.4. A topological space (\mathcal{X}, τ) is said to be *quasi-Polish* if there is a sequence $(f_n)_n$ of continuous functions $f_n : \mathcal{X} \rightarrow [-1, 1]$ separating points in \mathcal{X} .

We refer to [BOS16, Section 3] (where the terminology “quasi-Polish” was introduced by the way), [GHKP22, Appendix B], and the references therein for more details on quasi-Polish spaces.

Remark 5.2. (i) Any Polish space is a quasi-Polish space; (ii) A countable product of quasi-Polish spaces is a quasi-Polish space (if, for each k , $(f_{k,n})$ is a separating sequence of continuous functions on \mathcal{X}_k , and if $\pi^k : \mathcal{X} \rightarrow \mathcal{X}_k$ is the projection from the product space \mathcal{X} onto the factor \mathcal{X}_k , then $\{f_{k,n} \circ \pi^k; k, n \in \mathbb{N}\}$ is a separating sequence of continuous functions on \mathcal{X}); (iii) The spaces $C([0, T], H_w^1(\mathbb{T}))$, $C([0, T], L_w^2(\mathbb{T}))$ are quasi-Polish spaces.

We recall now the Skorokhod–Jakubowski theorem [J97, Theorem 2]

Theorem 5.5. Let (\mathcal{X}, τ) be a quasi-Polish space and let $(X_n)_n$ be a sequence of \mathcal{X} -valued random variable. Suppose that for all $\delta > 0$ there exists a compact subset $K_\delta \subset \mathcal{X}$ such that

$$\inf_n \mathbb{P}(X_n \in K_\delta) > 1 - \delta.$$

Then, up to a subsequence of $(X_n)_n$ (noted also $(X_n)_n$), there exists a probability space $(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})$, an \mathcal{X} -valued random variable \tilde{X} and an \mathcal{X} -valued sequence of random variables $(\tilde{X}_n)_n$ defined on $(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})$ such that

$$X_n \sim \tilde{X}_n, \quad n = 1, 2, \dots \quad (5.12)$$

$$\tilde{X}_n \rightarrow_\tau \tilde{X} \text{ for almost all } \omega \in \tilde{\Omega}. \quad (5.13)$$

Remark 5.3. The probability space $(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})$ can be chosen to be $([0, 1], \mathcal{B}_{[0,1]}, \mathcal{L})$.

5.3 Compactness

Let $R^\varepsilon, S^\varepsilon$ and u^ε be the global solutions of (2.11)–(2.12) given by Theorems 4.1 and 4.5. Let \mathcal{L}^2 denote the Lebesgue measure on $[0, T] \times \mathbb{T}$. We define for any $(t, x) \in [0, T] \times \mathbb{T}$ and any $\varepsilon > 0$ the measures

$$\mu_{t,x}^\varepsilon \stackrel{\text{def}}{=} \delta_{R^\varepsilon(t,x)} \otimes \delta_{S^\varepsilon(t,x)}, \quad \mu^\varepsilon \stackrel{\text{def}}{=} \mu_{t,x}^\varepsilon \times \mathcal{L}^2. \quad (5.14)$$

That is to say

$$\langle \mu^\varepsilon, \psi \rangle = \int_{[0,T] \times \mathbb{T}} \int_{\mathbb{R}^2} \psi(t, x, \xi, \eta) d\mu_{t,x}^\varepsilon(\xi, \eta) dx dt = \int_{[0,T] \times \mathbb{T}} \psi(t, x, R^\varepsilon(t, x), S^\varepsilon(t, x)) dx dt,$$

for all $\psi \in C_b([0, T] \times \mathbb{T} \times \mathbb{R}^2)$.

Let \mathcal{S} be a countable set of functions in $C_c^\infty(\mathbb{R})$ which is dense in $C_0(\mathbb{R})$. Set

$$X_R^\varepsilon \stackrel{\text{def}}{=} (f(R^\varepsilon))_{f \in \mathcal{S}}, \quad X_S^\varepsilon \stackrel{\text{def}}{=} (f(S^\varepsilon))_{f \in \mathcal{S}}. \quad (5.15)$$

Let also L_w^2 denote the L^2 space equipped with its weak topology and define $C([0, T], L_w^2(\mathbb{T}))^{\aleph_0}$ as the countable product

$$C([0, T], L_w^2(\mathbb{T}))^{\aleph_0} \stackrel{\text{def}}{=} \prod_{\kappa \in \mathbb{N}} C([0, T], L_w^2(\mathbb{T})), \quad (5.16)$$

with the product topology. Let \mathcal{X} denote the product (which is a quasi-Polish space) of the following spaces

$$\mathcal{Y}, (C([0, T], L_w^2(\mathbb{T}))^{\aleph_0})^2, C([0, T] \times \mathbb{T}), C([0, T], \mathfrak{M}_{-1}(\mathbb{T})), C(0, T), L^1((0, T); C(\mathbb{T})). \quad (5.17)$$

Our aim is to show that the (law of the) sequence of \mathcal{X} -valued Random variables $(X^\varepsilon)_\varepsilon$ defined by

$$X^\varepsilon \stackrel{\text{def}}{=} (\mu^\varepsilon, X_R^\varepsilon, X_S^\varepsilon, u^\varepsilon, W, \Theta^\varepsilon, \Xi^\varepsilon) \quad (5.18)$$

is tight in \mathcal{X} .

Proposition 5.6. *For any $\delta > 0$, there exists a compact $K_\delta \in \mathcal{X}$ such that*

$$\sup_\varepsilon \mathbb{P}(X^\varepsilon \notin K_\delta) \leq \delta.$$

Proof of Proposition 5.6. The proof will be done in several steps. We show that every element in X^ε is tight in its proper space. We use the same convention as in (4.5) for the use of the constant \mathbf{C} .

Step 1. (Young measure). By the energy estimate (4.10) (with $p = 1$), we obtain the bound

$$\mathbb{E} \int_{[0, T] \times \mathbb{T} \times \mathbb{R}^2} [\xi^2 + \eta^2] d\mu^\varepsilon(t, x, \xi, \eta) = \mathbb{E} \int_{[0, T] \times \mathbb{T}} \int_{\mathbb{R}^2} [\xi^2 + \eta^2] d\mu_{t,x}^\varepsilon(\xi, \eta) dx dt \leq \mathbf{C}(T), \quad (5.19)$$

and then, by the Markov inequality,

$$\mathbb{P} \left(\int_{[0, T] \times \mathbb{T} \times \mathbb{R}^2} [\xi^2 + \eta^2] d\mu^\varepsilon(t, x, \xi, \eta) \geq \mathbf{C}(T)/\delta \right) \leq \delta,$$

for all $\delta > 0$: Proposition 5.2 shows that $(\mu^\varepsilon)_\varepsilon$ is tight.

Step 2. (Factor $(X_R^\varepsilon, X_S^\varepsilon)_\varepsilon$). It follows from the energy estimate (4.10) (still with $p = 1$) that, for $f \in \mathcal{S}$,

$$\mathbb{E} \left[\|f(R^\varepsilon)\|_{C([0, T]; L^2(\mathbb{T}))}^2 \right] \leq \mathbf{C}(T). \quad (5.20)$$

Let now $\varphi \in H^1(\mathbb{T})$. By (5.1) we have

$$\begin{aligned} \int_{\mathbb{T}} [f(R^\varepsilon)(t) - f(R^\varepsilon)(s)] \varphi dx &= \int_s^t \int_{\mathbb{T}} c(u^\varepsilon) f(R^\varepsilon) \varphi_x dx d\sigma + \int_{\mathbb{T}} \left[\int_s^t f'(R^\varepsilon) \Phi^\varepsilon dW(\sigma) \right] \varphi dx \\ &+ \int_s^t \int_{\mathbb{T}} \left[\frac{1}{2} q^\varepsilon f''(R^\varepsilon) + \tilde{c}'(u^\varepsilon) [(R^\varepsilon)^2 - (S^\varepsilon)^2] f'(R^\varepsilon) + 2(S^\varepsilon - R^\varepsilon) f(R^\varepsilon) \right] \varphi dx d\sigma \\ &+ \int_s^t \int_{\mathbb{T}} \tilde{c}'(u^\varepsilon) [(2\Theta^\varepsilon R^\varepsilon - \chi_\varepsilon(R^\varepsilon)) f'(R^\varepsilon) - 4\Theta^\varepsilon f(R^\varepsilon)] \varphi dx d\sigma. \end{aligned}$$

Using the energy estimate (4.10) and straightforward computations gives us

$$\mathbb{E} \left| \int_{\mathbb{T}} [f(R^\varepsilon)(t) - f(R^\varepsilon)(s)] \varphi dx \right|^4 \leq \mathbf{C}(T, \|\varphi\|_{H^1(\mathbb{T})}) |t - s|^2. \quad (5.21)$$

By [Bas11, Theorem 8.2] we deduce from (5.21) that, for all $R > 0$,

$$\mathbb{P} \left(\sup_{0 \leq s < t \leq T} \frac{\int_{\mathbb{T}} f(R^\varepsilon)(t) - f(R^\varepsilon)(s) \varphi dx}{|t - s|^{1/4}} \geq R \right) \leq \mathbf{C}(T, \|\varphi\|_{H^1(\mathbb{T})}) R^{-4}. \quad (5.22)$$

Combining (5.20) with (5.22) (and the Markov inequality) gives us

$$\mathbb{P} \left(\|f(R^\varepsilon)\|_{C([0,T];L^2(\mathbb{T}))}^2 + \sup_{0 \leq s < t \leq T} \frac{\int_{\mathbb{T}} f(R^\varepsilon)(t) - f(R^\varepsilon)(s) \varphi dx}{|t - s|^{1/4}} \geq R \right) \leq \mathfrak{C}(T, \|\varphi\|_{H^1(\mathbb{T})}) R^{-1}. \quad (5.23)$$

Then, using [L96, Lemma C.1] we obtain that $(f(R^\varepsilon))_\varepsilon$ is tight in $C([0, T], L_w^2(\mathbb{T}))$. Now if $\delta > 0$ is fixed and if $\{f_1, f_2, \dots\}$ is an enumeration of \mathcal{S} , then, for any $n \geq 1$, there exists a compact K_n such that $\mathbb{P}(f_n(R^\varepsilon) \in K_n^c) \leq \delta/2^n$. From the union bound

$$\mathbb{P}(\exists n \geq 1, f_n(R^\varepsilon) \in K_n^c) \leq \sum_{n \geq 1} \mathbb{P}(f_n(R^\varepsilon) \in K_n^c) \leq \delta,$$

and Tychonoff's theorem, we can then deduce that (X_R^ε) is tight. The proof for (X_S^ε) is similar.

Step 3. (Factor $(u^\varepsilon)_\varepsilon$). Using the identity $u_t^\varepsilon = (R^\varepsilon + S^\varepsilon)/2 + \Xi^\varepsilon$ we obtain

$$\|u^\varepsilon(t) - u^\varepsilon(s)\|_{L^2(\mathbb{T})} \leq \mathfrak{C}(T) |t - s| \sup_{t \in [0, T]} \|(R^\varepsilon, S^\varepsilon, \Xi^\varepsilon)\|_{L^2(\mathbb{T})}.$$

Then, using the identity $u_x^\varepsilon = (S^\varepsilon - R^\varepsilon)/(2c(u^\varepsilon)) - \Theta^\varepsilon/c(u^\varepsilon)$ and the bounds (4.10), (4.11), (4.13), we obtain

$$\mathbb{E} \sup_{t \in [0, T]} \|u^\varepsilon\|_{H^1(\mathbb{T})}^2 + \mathbb{E} \sup_{\substack{t, s \in [0, T] \\ s \neq t}} \frac{\|u^\varepsilon(t) - u^\varepsilon(s)\|_{L^2(\mathbb{T})}^2}{|t - s|^2} \leq \mathfrak{C}(T).$$

This implies that

$$\mathbb{P} \left(\sup_{t \in [0, T]} \|u^\varepsilon\|_{H^1(\mathbb{T})}^2 + \sup_{\substack{t, s \in [0, T] \\ s \neq t}} \frac{\|u^\varepsilon(t) - u^\varepsilon(s)\|_{L^2(\mathbb{T})}^2}{|t - s|^2} \geq \mathfrak{C}(T)/\delta \right) \leq \delta. \quad (5.24)$$

By [Sim87, Theorem 5], (5.24) implies that $(u^\varepsilon)_\varepsilon$ is tight in the space $C([0, T] \times \mathbb{T})$.

Step 4. (Factor $(W)_\varepsilon$). The process W belongs to the Polish space $C([0, T], \mathfrak{M}_{-1})$ almost surely [DPZ14, Theorem 4.5]. It is tight by [Bil99, Theorem 1.3], but also for the more elementary reason that $\mathbb{E}\|W\|_{C([0, T], \mathfrak{M}_{-1})}^2 < \infty$.

Step 5. (Factor $(\Theta^\varepsilon, \Xi^\varepsilon)_\varepsilon$). The bounds (4.11) and (4.13) show that $(\Theta^\varepsilon, \Xi^\varepsilon)$ is tight in $C(0, T) \times L^1((0, T); C(\mathbb{T}))$. □

5.4 Limit equation with defect measure

5.4.1 Application of the Skorokhod–Jakubowski theorem

From Proposition 5.6 and Theorem 5.5 we deduce the following result.

Theorem 5.7 (Convergence up to a change of the probability space). *There exists a sequence $(\varepsilon_n)_n$, a probability space $(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})$, a \mathcal{X} -valued sequence of random variables $(\tilde{X}^n)_n$ and another \mathcal{X} -valued random variable \tilde{X} defined on $(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})$ such that for all $n \in \mathbb{N}$ we have*

$$X^{\varepsilon_n} \sim \tilde{X}_n, \quad n = 1, 2, \dots \quad (5.25)$$

$$\tilde{X}_n \rightarrow \tilde{X} \text{ in } \mathcal{X} \text{ } \tilde{\mathbb{P}}\text{-a.s. in } \tilde{\Omega}. \quad (5.26)$$

Remark 5.4. For the sake of simplicity, we omit the subsequence in the remainder of this paper and we write

$$X^\varepsilon \sim \tilde{X}^\varepsilon, \quad \forall \varepsilon \text{ (in fact for a countable values of } \varepsilon) \quad (5.27)$$

$$\tilde{X}^\varepsilon \rightarrow \tilde{X} \text{ in } \mathcal{X} \text{ for almost all } \omega \in \tilde{\Omega}. \quad (5.28)$$

The $\tilde{\mathbb{P}}$ -a.s. convergence (which also gives convergence in probability) in Theorem 5.7 will be exploited to pass to the limit in the equations (5.1) and (5.2). We examine first all terms but the stochastic integral in Section 5.4.2 below, and then devote Section 5.4.3 to the analysis of the stochastic integral.

5.4.2 Convergence et identification of some deterministic elements

As in (5.18) we have the decomposition

$$\tilde{X}^\varepsilon \stackrel{\text{def}}{=} (\tilde{\mu}^\varepsilon, \tilde{X}_R^\varepsilon, \tilde{X}_S^\varepsilon, \tilde{u}^\varepsilon, \tilde{W}^\varepsilon, \tilde{\Theta}^\varepsilon, \tilde{\Xi}^\varepsilon), \quad (5.29)$$

$$\tilde{X} \stackrel{\text{def}}{=} (\tilde{\mu}, \tilde{X}_R, \tilde{X}_S, \tilde{u}, \tilde{W}, \tilde{\Theta}, \tilde{\Xi}). \quad (5.30)$$

If any of the original factors $\mu^\varepsilon, R^\varepsilon, \dots$ satisfy \mathbb{P} -a.s. a relation which can be characterized in terms of a Borel subset A of \mathcal{X} , then this relation is still satisfied $\tilde{\mathbb{P}}$ -a.s. by the factors with tildas. If moreover A is closed for the topology considered on \mathcal{X} , then the relation is also satisfied $\tilde{\mathbb{P}}$ -a.s. by the limit factors since

$$\tilde{\mathbb{P}}(\tilde{X} \in A) \geq \limsup_{\varepsilon \rightarrow 0} \tilde{\mathbb{P}}(\tilde{X}^\varepsilon \in A) = 1. \quad (5.31)$$

As a consequence, we have the following identities.

Proposition 5.8. *The following identities hold $\tilde{\mathbb{P}}$ -almost surely:*

$$\tilde{\Theta} = 0, \quad \tilde{\Xi} = 0, \quad (5.32)$$

as well as (2.12) with tildas and

$$\tilde{u}_x^\varepsilon = \frac{\tilde{S}^\varepsilon - \tilde{R}^\varepsilon}{2c(\tilde{u}^\varepsilon)} - \frac{\tilde{\Theta}^\varepsilon}{c(\tilde{u}^\varepsilon)}, \quad \tilde{u}_t^\varepsilon = \frac{\tilde{S}^\varepsilon + \tilde{R}^\varepsilon}{2} + \tilde{\Xi}^\varepsilon, \quad \tilde{u}_x = \frac{\tilde{S} - \tilde{R}}{2c(\tilde{u})}, \quad \tilde{u}_t = \frac{\tilde{S} + \tilde{R}}{2}, \quad (5.33)$$

and

$$\tilde{\mu}^\varepsilon = \delta_{\tilde{R}^\varepsilon(t,x)} \otimes \delta_{\tilde{S}^\varepsilon(t,x)} \rtimes \mathcal{L}^2, \quad \tilde{\mu} = \tilde{\mu}_{t,x} \rtimes \mathcal{L}^2, \quad (5.34)$$

where, for \mathcal{L}^2 -almost all $(t,x) \in [0,T] \times \mathbb{T}$, $\tilde{\mu}_{t,x}$ is probability measures on \mathbb{R}^2 . In particular, we have $\tilde{u} \in H^1((0,T) \times \mathbb{T})$ with

$$\|\tilde{u}\|_{H^1((0,T) \times \mathbb{T})} \leq \mathfrak{C} (\|\tilde{S}\|_{L^2((0,T) \times \mathbb{T})} + \|\tilde{R}\|_{L^2((0,T) \times \mathbb{T})}). \quad (5.35)$$

Finally, we have the following convergence results: for all function $g: \mathbb{R} \rightarrow \mathbb{R}_+$ such that $g(\xi) = o(\xi)$ when $\xi \rightarrow \infty$, $\tilde{\mathbb{P}}$ -a.s.

$$\lim_{\varepsilon \rightarrow 0} \int_0^T \int_{\mathbb{T}} \tilde{c}'(\tilde{u}^\varepsilon) [g(\tilde{R}^\varepsilon) \chi_\varepsilon(\tilde{R}^\varepsilon) + g(\tilde{S}^\varepsilon) \chi_\varepsilon(\tilde{S}^\varepsilon)] dx dt = 0. \quad (5.36)$$

Due the equality of laws (5.27), all the estimates obtained in Section 4 are valid with the tildas.

Proof of Proposition 5.8. The identities in (5.32) follow from (4.11)-(4.13) (apply (5.31) with A a closed ball around the origin of arbitrary small radius). We also deduce (5.36) from (4.7) and the fact that $\chi_\varepsilon(R) \neq 0 \Rightarrow \varepsilon R \geq 1$. Let us prove also (5.34) for instance: the first identity between the components $\mu^\varepsilon, R^\varepsilon, S^\varepsilon$ of X^ε can be characterized by the four relations

$$\begin{aligned} & \int_{[0,T] \times \mathbb{T} \times \mathbb{R} \times \mathbb{R}} \varphi(t,x)(\alpha_1 \xi + \alpha_2 \eta + \alpha_3 \xi^2 + \alpha_4 \eta^2) d\mu(t,x,\xi,\eta) \\ &= \int_{[0,T] \times \mathbb{T}} \varphi(t,x)(\alpha_1 R(t,x) + \alpha_2 S(t,x) + \alpha_3 R(t,x)^2 + \alpha_4 S(t,x)^2) dx dt, \end{aligned} \quad (5.37)$$

where $\alpha_i \in \{0,1\}$, for φ ranging in a dense countable subset of $C([0,T] \times \mathbb{T})$, and is therefore a Borel set since each term in (5.37) defines a Borel function (in fact continuous or semi-continuous function) of

$$(\mu, R, S) \in \mathcal{Y} \times C([0,T], L_w^2(\mathbb{T})) \times C([0,T], L_w^2(\mathbb{T})). \quad (5.38)$$

The last equation $\mu = \mu_{t,x} \times \mathcal{L}^2$ can be characterized by the relations

$$\int_{[0,T] \times \mathbb{T} \times \mathbb{R} \times \mathbb{R}} \varphi(t,x) d\mu(t,x,\xi,\eta) = \int_{[0,T] \times \mathbb{T}} \varphi(t,x) dx dt, \quad (5.39)$$

for φ ranging in a dense countable subset of $C([0,T] \times \mathbb{T})$, [BV19, Section 4.1], so is a Borel set as well. \square

Proposition 5.9. *After redefinition of $(t,x) \mapsto \tilde{\mu}_{t,x}$ on a \mathcal{L}^2 -negligible set, we have the following convergence properties: $\tilde{\mathbb{P}}$ -a.s., for all functions $f_1, f_2 \in C(\mathbb{R})$ satisfying the growth condition*

$$|f(\xi)| \leq C_f(1 + |\xi|^r), \quad |f'(\xi)| \leq C_f(1 + |\xi|^{r-1}), \quad |f''(\xi)| \leq C_f, \quad r \in [1,2), \quad (5.40)$$

for all $\varphi \in C(\mathbb{T})$, we have

$$\int_{\mathbb{T}} f_1(\tilde{R}^\varepsilon) + f_2(\tilde{S}^\varepsilon) \varphi dx \rightarrow \left[t \mapsto \int_{\mathbb{T} \times \mathbb{R}^2} \varphi(x) [f_1(\xi) + f_2(\zeta)] d\tilde{\mu}_{t,x}(\xi, \zeta) dx \right] \quad (5.41)$$

in $C([0,T])$. Moreover, at the initial time $t = 0$, we have

$$\tilde{\mu}_{0,x} = \delta_{R_0(x)} \otimes \delta_{S_0(x)}. \quad (5.42)$$

Proof of Proposition 5.9. We first observe that if a function $g \in L^1(0,T)$ coincide a.e. with a function $h \in C([0,T])$, then the set of right-Lebesgue points of g is the full set $[0,T)$ and

$$h(t) = \lim_{\delta \rightarrow 0} \frac{1}{\delta} \int_t^{t+\delta} g(s) ds \quad (5.43)$$

for all $t \in [0,T)$. We consider therefore the representative $\tilde{\mu}_{t,x}^*$ of $\tilde{\mu}_{t,x}$ given by

$$\int_{\mathbb{T} \times \mathbb{R}^2} \varphi(x) F(\xi, \zeta) d\tilde{\mu}_{t,x}^*(\xi, \zeta) dx \stackrel{\text{def}}{=} \lim_{\delta \rightarrow 0} \frac{1}{\delta} \int_t^{t+\delta} \int_{\mathbb{T} \times \mathbb{R}^2} \varphi(x) F(\xi, \zeta) d\tilde{\mu}_{s,x}(\xi, \zeta) dx ds, \quad (5.44)$$

where $\varphi \in C(\mathbb{T})$, $F \in C_0(\mathbb{R}^2)$. By differentiation's theory, (5.44) is satisfied for all $t \in I_{\varphi,F}$, where $I_{\varphi,F}$ is the set of right-Lebesgue points of the function

$$t \mapsto \int_{\mathbb{T} \times \mathbb{R}^2} \varphi(x) F(\xi, \zeta) d\tilde{\mu}_{s,x}(\xi, \zeta) dx \quad (5.45)$$

and is of full measure in $[0, T)$. We denote by I the intersection of the sets $I_{\varphi, f}$ over φ in a dense countable subset of $C(\mathbb{T})$ and F in a dense countable subset of $C_0(\mathbb{R}^2)$. Without loss of generality, we can assume that the functions F of the form

$$F_{1,2}(\xi, \zeta) = f_1(\xi) + f_2(\zeta), \quad f_1, f_2 \in \mathcal{S}, \quad (5.46)$$

are here taken into account. If $f_1, f_2 \in \mathcal{S}$, then, by (5.28), the left-hand side of (5.41) is converging in $C([0, T])$ to a continuous function Λ . It is also converging in $\mathcal{D}'(0, T)$ to the function

$$\Gamma(t) \stackrel{\text{def}}{=} \int_{\mathbb{T} \times \mathbb{R}^2} \varphi(x) F_{1,2}(\xi, \zeta) d\tilde{\mu}_{t,x}(\xi, \zeta) dx, \quad F_{1,2}(\xi, \zeta) \stackrel{\text{def}}{=} f_1(\xi) + f_2(\zeta), \quad (5.47)$$

since $\tilde{\mu}^\varepsilon \rightarrow \tilde{\mu}$ in \mathcal{Y} . Consequently Γ and Λ coincide a.e. and by our initial observation, (5.41) is satisfied with the representative $\tilde{\mu}_{t,x}^*$ of $\tilde{\mu}_{t,x}$. A simple limiting argument shows that the result holds true for f_1 and f_2 satisfying (5.40). The last assertion (5.42) is a simple consequence of the fact that R_0^ε and S_0^ε are deterministic so \mathbb{P} -a.s.,

$$(\tilde{X}^\varepsilon(0), \tilde{X}_S^\varepsilon(0)) = (f(R_0^\varepsilon), f(S_0^\varepsilon))_{f \in \mathcal{S}}. \quad (5.48)$$

This is equivalent to $(\tilde{R}^\varepsilon(0), \tilde{S}^\varepsilon(0)) = (R_0^\varepsilon, S_0^\varepsilon)$, \mathbb{P} -a.s. Since $(R_0^\varepsilon, S_0^\varepsilon)$ converges strongly to (R_0, S_0) in $L^2(\mathbb{T})$, $\tilde{\mu}_{0,x}$ reduces to the Dirac mass $\delta_{(R_0(x), S_0(x))} = \delta_{R_0(x)} \otimes \delta_{S_0(x)}$. \square

Using the estimates (4.10), (4.48), (4.65) with Proposition 5.3 we obtain the following result.

Proposition 5.10. *Let $t_0 \in (0, T)$ and $r \in (1, 2)$, $p \in (1, \infty)$, $q \in (1, 3)$. Then the measure $\tilde{\mu}$ satisfies*

$$\tilde{\mathbb{E}} \int_{[0, T] \times \mathbb{T} \times \mathbb{R}^2} [|\xi|^2 + |\eta|^2] d\tilde{\mu}(t, x, \xi, \eta) < \infty, \quad (5.49)$$

and

$$\tilde{\mathbb{E}} \left(\int_{[t_0, T] \times \mathbb{T} \times \mathbb{R}^2} [|\xi^-|^p + |\eta^-|^p] d\tilde{\mu}(t, x, \xi, \eta) \right)^{2/p} < \infty, \quad (5.50)$$

$$\tilde{\mathbb{E}} \int_{[0, T] \times \mathbb{T} \times \mathbb{R}^2} c'(\tilde{u}(t, x)) [|\xi|^q + |\eta|^q] d\tilde{\mu}(t, x, \xi, \eta) < \infty. \quad (5.51)$$

Moreover, for any functions $f, g \in C(\mathbb{R}^2)$ satisfying

$$|f(\xi, \eta)| \leq C [1 + |\xi|^r + |\eta|^r], \quad |g(\xi, \eta)| \leq C [1 + |\xi^-|^p + |\eta^-|^p + |\xi^+|^q + |\eta^+|^q], \quad (5.52)$$

we have

$$\lim_{\varepsilon \rightarrow 0} \tilde{\mathbb{E}} \left| \int_{[0, T] \times \mathbb{T} \times \mathbb{R}^2} \phi(t, x) f(\xi, \eta) d\tilde{\mu}^\varepsilon(t, x, \xi, \eta) - \int_{[0, T] \times \mathbb{T} \times \mathbb{R}^2} \phi(t, x) f(\xi, \eta) d\tilde{\mu}(t, x, \xi, \eta) \right| = 0, \quad (5.53)$$

and

$$\lim_{\varepsilon \rightarrow 0} \tilde{\mathbb{E}} \left| \int_{[t_0, T] \times \mathbb{T} \times \mathbb{R}^2} c'(\tilde{u}^\varepsilon(t, x)) \varphi(t, x) g(\xi, \eta) d\tilde{\mu}^\varepsilon(t, x, \xi, \eta) - \int_{[t_0, T] \times \mathbb{T} \times \mathbb{R}^2} c'(\tilde{u}(t, x)) \varphi(t, x) g(\xi, \eta) d\tilde{\mu}(t, x, \xi, \eta) \right|^{2/p} = 0, \quad (5.54)$$

for all $\phi \in L^{\frac{2}{2-r}}([0, T] \times \mathbb{T})$, $\varphi \in \cup_{q' \in (q, 3)} L^{\frac{q'}{q'-q}}([0, T] \times \mathbb{T})$. Finally, the convergence (5.41) holds true when $f_1, f_2 \in C(\mathbb{R})$ satisfy the growth condition $|f_1(\xi)| + |f_2(\xi)| \leq C(1 + |\xi|^r)$ and $\varphi \in L^{\frac{2}{2-r}}(\mathbb{T})$.

5.4.3 The stochastic integral

We would like to show that the equations (5.1)-(5.2) are satisfied $\tilde{\mathbb{P}}$ -a.s. by the quantities with tildas and then pass to the limit as $\varepsilon \rightarrow 0$. So we would like to establish, for all $\varphi \in C^1(\mathbb{T})$ and $h \in \mathcal{S}$, the identity

$$\tilde{M}_{\varphi, h}(t) = \int_{\mathbb{T}} \int_{t_0}^t h'(\tilde{R}^\varepsilon) \varphi \Phi^\varepsilon d\tilde{W}^\varepsilon(s) dx, \quad (5.55)$$

where

$$\begin{aligned} \tilde{M}_{\varphi, h}(t) &\stackrel{\text{def}}{=} \int_{\mathbb{T}} h(\tilde{R}^\varepsilon)(t, x) \varphi(x) dx - \int_{\mathbb{T}} h(\tilde{R}^\varepsilon)(t_0, x) \varphi(x) dx - \int_{t_0}^t \int_{\mathbb{T}} (c(\tilde{u}^\varepsilon) h(\tilde{R}^\varepsilon)) \varphi_x dx ds \\ &\quad - \int_{t_0}^t \int_{\mathbb{T}} \tilde{c}'(\tilde{u}^\varepsilon) [(\tilde{S}^\varepsilon - \tilde{R}^\varepsilon) B_h(\tilde{R}^\varepsilon, \tilde{S}^\varepsilon) - h'(\tilde{R}^\varepsilon) \chi_\varepsilon(\tilde{R}^\varepsilon) - \frac{1}{2} q^\varepsilon h''(\tilde{R}^\varepsilon)] \varphi dx ds \\ &\quad - 2 \int_{t_0}^t \int_{\mathbb{T}} \tilde{c}'(\tilde{u}^\varepsilon) \tilde{\Theta}^\varepsilon (\tilde{R}^\varepsilon h'(\tilde{R}^\varepsilon) - 2h(\tilde{R}^\varepsilon)) dx ds, \end{aligned} \quad (5.56)$$

(the term B_h being defined by (3.64)). The identification (5.55) is not trivial. We can either use a characterization in law of the stochastic integral in terms of martingales as in [Ond10, BO11, HS12, Hof13, DHV16] for instance, or resort to the trick of Bensoussan, [Ben95, p. 282], which uses a regularization by convolution of the integrand of the stochastic integral and the stochastic Fubini theorem to “fully integrate” the $dW(s)$. This last approach will indeed give (5.55)-(5.56), the details are left to the reader. Let us however specify what is the filtration involved here and what are the properties of the Wiener process \tilde{W}^ε . Let $(\tilde{\mathcal{F}}_t^\varepsilon)$ denote the filtration generated by the process \tilde{X}^ε and let $(\tilde{\mathcal{F}}_t^\varepsilon)$ be the augmented filtration, obtained by the completion of the right-continuous filtration $(\tilde{\mathcal{F}}_{t+}^{0, \varepsilon})$.

Lemma 5.11 (Equivalent cylindrical Wiener process). *The process \tilde{W}^ε is a $(\tilde{\mathcal{F}}_t^\varepsilon)$ -adapted cylindrical Wiener process, the increment $\tilde{W}^\varepsilon(t) - \tilde{W}^\varepsilon(s)$ is independent on $\tilde{\mathcal{F}}_s^\varepsilon$ for all $t \geq s \geq 0$ and $\tilde{\mathbb{P}}$ -a.s*

$$\tilde{W}^\varepsilon = \sum_k \tilde{\beta}_k^\varepsilon g_k \quad (5.57)$$

in $C([0, T], \mathfrak{U}_{-1})$, where $(\tilde{\beta}_1^\varepsilon(t), \tilde{\beta}_2^\varepsilon(t), \dots)$ are independent one-dimensional Wiener processes.

Proof. We obtain directly from the definition of the filtration $(\tilde{\mathcal{F}}_t^\varepsilon)$ that \tilde{W}^ε is a $(\tilde{\mathcal{F}}_t^\varepsilon)$ -adapted Wiener process. If $t \geq s \geq 0$, that $\tilde{W}^\varepsilon(t) - \tilde{W}^\varepsilon(s)$ is independent on $\tilde{\mathcal{F}}_s^{0, \varepsilon}$ means

$$\tilde{\mathbb{E}} [G(\tilde{W}^\varepsilon(t) - \tilde{W}^\varepsilon(s)) H(\tilde{X}^\varepsilon(s_1), \dots, \tilde{X}^\varepsilon(s_m))] = 0, \quad (5.58)$$

for all $m \geq 1$, all times $0 \leq s_1 \leq \dots \leq s_m \leq s$, and all continuous functions $G: \mathfrak{U}_{-1} \rightarrow \mathbb{R}$, $H: \mathcal{X}^m \rightarrow \mathbb{R}$. The condition (5.58) is clearly invariant by change of the probability space and so true, since satisfied by the original variables when the tildas are removed. We have then

$$\tilde{\mathbb{E}} [G(\tilde{W}^\varepsilon(t) - \tilde{W}^\varepsilon(s')) \tilde{Y}_s] = 0, \quad (5.59)$$

for all $s < s' \leq t$ and all bounded $\tilde{\mathcal{F}}_{s+}^{0,\varepsilon}$ -measurable random variable \tilde{Y}_s . By continuity of $r \mapsto \tilde{W}^\varepsilon(r)$, we can let $s' \downarrow s$ in (5.59) and obtain the independence of $\tilde{W}^\varepsilon(t) - \tilde{W}^\varepsilon(s)$ and $\tilde{\mathcal{F}}_{s+}^\varepsilon$, and thus also independence with the completed σ -algebra. Finally, we have the expansion (5.57) with

$$\tilde{\beta}_k^\varepsilon(t) \stackrel{\text{def}}{=} \langle \tilde{W}^\varepsilon(t), g_k \rangle_{\mathfrak{U}_{-1}}. \quad (5.60)$$

The map $w \mapsto \langle w, g_k \rangle_{\mathfrak{U}_{-1}}$ is continuous $C([0, T]; \mathfrak{U}_{-1}) \rightarrow C([0, T])$, so, for each J finite, $(\tilde{\beta}_k^\varepsilon(t))_{k \in J}$ and $(\beta_k(t))_{k \in J}$ have the same laws in $C([0, T])^J$. It follows that $(\tilde{\beta}_1^\varepsilon(t), \tilde{\beta}_2^\varepsilon(t), \dots)$ are independent one-dimensional Wiener processes. \square

Once (5.55)-(5.56) is established we can use [DGHT11, Lemma 2.1] to pass to the limit in the stochastic integral. This is explained in the next section 5.4.4.

5.4.4 Limit equation with defect measure

We introduce the following notations

$$\langle f(R, S) \rangle(t, x) \stackrel{\text{def}}{=} \int_{\mathbb{R}^2} f(\xi, \eta) d\tilde{\mu}_{t,x}(\xi, \eta), \quad (5.61)$$

for any f such that (5.61) is well defined. As a consequence of Proposition 5.9 and Proposition 5.10, if $t_0 \in (0, T)$ and if f and g satisfy (5.52) while f_1, f_2 satisfy (5.40), then, up to a subsequence, we have, $\tilde{\mathbb{P}}$ -a.s.,

$$\begin{aligned} f(\tilde{R}^\varepsilon, \tilde{S}^\varepsilon) &\rightarrow \langle f(R, S) \rangle && \text{in } L^{2/r}([0, T] \times \mathbb{T}), \\ c'(\tilde{u}^\varepsilon) g(\tilde{R}^\varepsilon, \tilde{S}^\varepsilon) &\rightarrow c'(\tilde{u}) \langle g(R, S) \rangle && \text{in } L^{q'/q}([t_0, T] \times \mathbb{T}) \quad \forall q' \in (q, 3), \\ \int_{\mathbb{T}} f_1(\tilde{R}^\varepsilon) + f_2(\tilde{S}^\varepsilon) \varphi dx &\rightarrow \int_{\mathbb{T}} \langle f_1(R) + f_2(S) \rangle \varphi dx && \text{in } C([0, T]). \end{aligned}$$

Using the dominated convergence theorem, we also have

$$\lim_{\varepsilon} \|\Phi - \Phi^\varepsilon\|_{L_2(\mathfrak{U}, L^2(\mathbb{T}))}^2 = \sum_{k \geq 1} \lim_{\varepsilon} \|\sigma_k - \sigma_k^\varepsilon\|_{L^2(\mathbb{T})}^2 = 0. \quad (5.62)$$

Let $(\tilde{\mathcal{F}}_t)$ be defined as the augmented filtration generated by \tilde{X} . Lemma 5.11 has a completely equivalent version for $((\tilde{\mathcal{F}}_t), (\tilde{W}(t)))$ so we can claim that

$$(\tilde{\Omega}, \tilde{\mathbb{P}}, \tilde{\mathcal{F}}, (\tilde{\mathcal{F}}_t), (\tilde{W}(t))) \text{ is a stochastic basis.} \quad (5.63)$$

We can now pass to the limit $[\varepsilon \rightarrow 0]$ in (5.55)-(5.56), using [DGHT11, Lemma 2.1] and a similar argument as (5.62) for the stochastic integral, (5.36), and Proposition 5.10, to obtain the limit equation

$$\begin{aligned} &\int_{\mathbb{T}} \langle h(R) \rangle(t, x) \varphi(x) dx - \int_{\mathbb{T}} \langle h(R) \rangle(t_0, x) \varphi(x) dx - \int_{t_0}^t \int_{\mathbb{T}} c(\tilde{u}) \langle h(R) \rangle \varphi_x dx ds \\ &= \int_{t_0}^t \int_{\mathbb{T}} \tilde{c}'(\tilde{u}) [\langle 2(S - R) h(R) + h'(R) (R^2 - S^2) \rangle + \frac{1}{2} q \langle h''(R) \rangle] \varphi dx ds \\ &\quad + \int_{\mathbb{T}} \int_{t_0}^t \langle h'(R) \rangle \varphi \Phi d\tilde{W}(s) dx, \end{aligned} \quad (5.64)$$

and, similarly,

$$\begin{aligned}
& \int_{\mathbb{T}} \langle h(S) \rangle(t, x) \varphi(x) dx - \int_{\mathbb{T}} \langle h(S) \rangle(t_0, x) \varphi(x) dx + \int_{t_0}^t \int_{\mathbb{T}} c(\tilde{u}) \langle h(S) \rangle \varphi_x dx ds \\
&= \int_{t_0}^t \int_{\mathbb{T}} \tilde{c}'(\tilde{u}) [\langle 2(R-S)h(S) + h'(S)(S^2 - R^2) \rangle + \frac{1}{2}q \langle h''(S) \rangle] \varphi dx ds \\
&\quad + \int_{\mathbb{T}} \int_{t_0}^t \langle h'(S) \rangle \varphi \Phi d\tilde{W}(s) dx, \tag{5.65}
\end{aligned}$$

those equations (5.64) and (5.65) being satisfied when $h \in W_{\text{loc}}^{2,\infty}(\mathbb{R})$ is such that

$$|h(\xi)| \leq C(1 + |\xi|^r), \quad |h'(\xi)| \leq C(1 + |\xi|^{r-1}), \quad |h''(\xi)| \leq C, \tag{5.66}$$

for some given constant $C \geq 0$ and some exponent $r \in [1, 2)$. In particular, for $h(\xi) = \xi$, we obtain

$$\begin{aligned}
& \int_{\mathbb{T}} \langle R \rangle(t, x) \varphi(x) dx - \int_{\mathbb{T}} \langle R \rangle(t_0, x) \varphi(x) dx - \int_{t_0}^t \int_{\mathbb{T}} c(\tilde{u}) \langle R \rangle \varphi_x dx ds \\
&= - \int_{t_0}^t \int_{\mathbb{T}} \tilde{c}'(\tilde{u}) [\langle R^2 \rangle - 2 \langle RS \rangle + \langle S^2 \rangle] \varphi dx ds + \int_{\mathbb{T}} \int_{t_0}^t \varphi \Phi d\tilde{W}(s) dx, \tag{5.67}
\end{aligned}$$

and, similarly,

$$\begin{aligned}
& \int_{\mathbb{T}} \langle S \rangle(t, x) \varphi(x) dx - \int_{\mathbb{T}} \langle S \rangle(t_0, x) \varphi(x) dx + \int_{t_0}^t \int_{\mathbb{T}} c(\tilde{u}) \langle S \rangle \varphi_x dx ds \\
&= - \int_{t_0}^t \int_{\mathbb{T}} \tilde{c}'(\tilde{u}) [\langle R^2 \rangle - 2 \langle RS \rangle + \langle S^2 \rangle] \varphi dx ds + \int_{\mathbb{T}} \int_{t_0}^t \varphi \Phi d\tilde{W}(s) dx, \tag{5.68}
\end{aligned}$$

Assume $\langle RS \rangle = \langle R \rangle \langle S \rangle$ (actually, this is established as a first step in the next section 6). Then (5.67)-(5.68) is the expected limit equation (2.6) in conservative form, up to the defect (non-negative) terms $\langle R^2 \rangle - \langle R \rangle^2$ and $\langle S^2 \rangle - \langle S \rangle^2$. We will show in Section 6 that these defect measures are trivial. To that effect, we will need an evolution equation for some non-linear functions of $\langle R \rangle$, as stated in the following proposition.

Proposition 5.12 (Renormalization in (5.67)). *Let $h \in C^\infty(\mathbb{R})$ satisfy the growth condition (5.66). We have then*

$$\begin{aligned}
& \int_{\mathbb{T}} h(\langle R \rangle)(t, x) \varphi(x) dx - \int_{\mathbb{T}} h(\langle R \rangle)(t_0, x) \varphi(x) dx - \int_{t_0}^t \int_{\mathbb{T}} c(\tilde{u}) h(\langle R \rangle) \varphi_x dx ds \\
&= \int_{t_0}^t \int_{\mathbb{T}} \tilde{c}'(\tilde{u}) \left[2(\langle R \rangle - \langle S \rangle) [\langle R \rangle h'(\langle R \rangle) - h(\langle R \rangle)] - h'(\langle R \rangle) \langle (R-S)^2 \rangle + \frac{1}{2}q h''(\langle R \rangle) \right] \varphi dx ds \\
&\quad + \int_{\mathbb{T}} \int_{t_0}^t h'(\langle R \rangle) \varphi \Phi d\tilde{W}(s) dx, \tag{5.69}
\end{aligned}$$

for all $0 \leq t_0 \leq t$.

Proof of Proposition 5.12. The proof is the same as in [DL89]: by regularization and study of the remainder terms. Since we have to treat additional stochastic terms here, and since renormalization will again be invoked later (in the derivation of (6.32) precisely), we give few

details. Let (ρ_δ) be an approximation of the unit constituted of even functions and let J_δ be the Friedrichs regularization operator $F \mapsto F * \rho_\delta$. We use $\varphi := J_\delta \varphi$ in (5.67) to obtain the equation

$$dR^\delta(t, x) + [c(\tilde{u}(t, x)) R^\delta(t, x)]_x dt = F^\delta(t, x) dt + \gamma^\delta(t, x) dt + \Phi^\delta(x) d\tilde{W}(t), \quad (5.70)$$

where $R^\delta = J_\delta \langle R \rangle$, $\Phi^\delta = J_\delta \Phi$ and

$$F^\delta(t) = J_\delta F(t), \quad F \stackrel{\text{def}}{=} -\tilde{c}'(\tilde{u}) [\langle R^2 \rangle - 2 \langle RS \rangle + \langle S^2 \rangle], \quad (5.71)$$

and γ^δ is the commutator

$$\gamma^\delta = [c(\tilde{u}) R^\delta]_x - J_\delta [c(\tilde{u}) \langle R \rangle]_x. \quad (5.72)$$

The equation (5.70) is an equation satisfied for all $x \in \mathbb{T}$ by the real-valued stochastic process $(R^\delta(t, x))$. Let $h \in C_c^\infty(\mathbb{R})$. The Itô formula and (5.70) give

$$\begin{aligned} dh(R^\delta(t, x)) + h'(R^\delta(t, x)) [c(\tilde{u}(t, x)) R^\delta(t, x)]_x dt &= h'(R^\delta(t, x)) F^\delta(t, x) dt \\ + h'(R^\delta(t, x)) \gamma^\delta(t, x) dt + \frac{1}{2} h''(R^\delta(t, x)) q^\delta(x) dt &+ h'(R^\delta(t, x)) \Phi^\delta(x) d\tilde{W}(t), \end{aligned} \quad (5.73)$$

where q^δ is defined in (2.15). By the chain-rule for H^1 -functions and (5.33) we have

$$h'(R^\delta) [c(\tilde{u}) R^\delta]_x = [c(\tilde{u}) h(R^\delta)]_x + 2 \tilde{c}'(\tilde{u}) (\tilde{S} - \tilde{R}) [h'(R^\delta) R^\delta - h(R^\delta)]. \quad (5.74)$$

We combine (5.74) with (5.33), integrate (5.73) against a function $\varphi \in C^1(\mathbb{T})$ and pass to the limit $\delta \rightarrow 0$ (on the formulation integrated in time) to obtain (5.69). The term involving the commutator γ^δ is

$$\int_{t_0}^t \int_{\mathbb{T}} h'(R^\delta) \gamma^\delta \varphi dx ds, \quad (5.75)$$

which converges to 0 since $h'(R^\delta)$ and φ are bounded, while $\gamma^\delta \rightarrow 0$ in $L^1((t_0, t) \times \mathbb{T})$ by [DL89, Lemma II.1.ii)] applied with $\alpha = 2$, $p = 2$. The limit of the other terms is obtained by standard considerations. \square

6 Reduction of the Young measures

Our aim is to show that $\tilde{\mu}_{t,x} = \delta_{\tilde{R}(t,x)} \otimes \delta_{\tilde{S}(t,x)}$. We prove in a first step that $\tilde{\mu}_{t,x} = \tilde{\nu}_{t,x}^1 \otimes \tilde{\nu}_{t,x}^2$ where $\tilde{\nu}_{t,x}^1$ and $\tilde{\nu}_{t,x}^2$ are two probability measures on \mathbb{R} that will be identified later.

6.1 Compensated compactness

For almost all $(t, x) \in [0, T] \times \mathbb{T}$ we can define the marginals $\tilde{\nu}_{t,x}^1$ and $\tilde{\nu}_{t,x}^2$ of $\tilde{\mu}_{t,x}$:

$$\tilde{\nu}_{t,x}^1(A) \stackrel{\text{def}}{=} \tilde{\mu}_{t,x}(A \times \mathbb{R}), \quad \tilde{\nu}_{t,x}^2(A) \stackrel{\text{def}}{=} \tilde{\mu}_{t,x}(\mathbb{R} \times A) \quad \forall A \in \mathcal{B}(\mathbb{R}).$$

The aim of this section is to prove the following lemma

Lemma 6.1. *For almost all $(\omega, t, x) \in \tilde{\Omega} \times [0, T] \times \mathbb{T}$ we have*

$$\tilde{\mu}_{t,x} = \tilde{\nu}_{t,x}^1 \otimes \tilde{\nu}_{t,x}^2. \quad (6.1)$$

Proof of Lemma 6.1. We will use the div-curl lemma. Our aim is to prove that, given $f \in C_c^\infty(\mathbb{R})$, the sequence

$$\{f(\tilde{R}^\varepsilon)_t + [c(\tilde{u})f(\tilde{R}^\varepsilon)]_x\}_\varepsilon \quad (6.2)$$

is tight in $H^{-1}((0, T) \times \mathbb{T})$. By (5.1), we can write (6.2) as the sum

$$\partial_t M_{\tilde{R}}^\varepsilon + T^\varepsilon + [(c(\tilde{u}) - c(\tilde{u}^\varepsilon))f(\tilde{R}^\varepsilon)]_x, \quad (6.3)$$

where

$$\begin{aligned} T^\varepsilon &\stackrel{\text{def}}{=} \frac{1}{2} q^\varepsilon f''(\tilde{R}^\varepsilon) + \tilde{c}'(\tilde{u}^\varepsilon) [((\tilde{R}^\varepsilon)^2 - (\tilde{S}^\varepsilon)^2) f'(\tilde{R}^\varepsilon) + 2(\tilde{S}^\varepsilon - \tilde{R}^\varepsilon) f(\tilde{R}^\varepsilon)] \\ &\quad + \tilde{c}'(\tilde{u}^\varepsilon) [(2\tilde{\Theta}^\varepsilon \tilde{R}^\varepsilon - \chi_\varepsilon(\tilde{R}^\varepsilon)) f'(\tilde{R}^\varepsilon) - 4\tilde{\Theta}^\varepsilon f(\tilde{R}^\varepsilon)], \end{aligned}$$

and

$$M_{\tilde{R}}^\varepsilon(t, x) \stackrel{\text{def}}{=} \sum_k \int_0^t f'(\tilde{R}^\varepsilon(s, x)) \sigma_k^\varepsilon(x) d\tilde{\beta}_k^\varepsilon(s).$$

Clearly, the quantity

$$[c(\tilde{u}) - c(\tilde{u}^\varepsilon)] f(\tilde{R}^\varepsilon) \quad (6.4)$$

tends to 0 $\tilde{\mathbb{P}}$ -almost surely in $L^2((0, T) \times \mathbb{T})$ and so the sequence

$$\{[(c(\tilde{u}) - c(\tilde{u}^\varepsilon))f(\tilde{R}^\varepsilon)]_x\}_\varepsilon \quad (6.5)$$

is tight in $H^{-1}((0, T) \times \mathbb{T})$. There remains to examine the terms $M_{\tilde{R}}^\varepsilon$ and T^ε , which is done in the following two steps, before we conclude the argument in Step 3.

Step 1. Using the Itô isometry with the embedding $L^2 \hookrightarrow H^{-1}$ gives us

$$\tilde{\mathbb{E}} \|M_{\tilde{R}}^\varepsilon(t+h, \cdot) - M_{\tilde{R}}^\varepsilon(t, \cdot)\|_{H^{-1}(\mathbb{T})}^2 \leq \tilde{\mathbb{E}} \|M_{\tilde{R}}^\varepsilon(t+h, \cdot) - M_{\tilde{R}}^\varepsilon(t, \cdot)\|_{L^2(\mathbb{T})}^2 \leq CT|h|.$$

Integrating with respect to t , we obtain

$$\tilde{\mathbb{E}} \|M_{\tilde{R}}^\varepsilon(\cdot + h, \cdot) - M_{\tilde{R}}^\varepsilon\|_{L^2((0, T-h), H^{-1}(\mathbb{T}))}^2 \leq C|h|.$$

Using the Itô isometry again we also have

$$\tilde{\mathbb{E}} \|M_{\tilde{R}}^\varepsilon(t, \cdot)\|_{L^2(\mathbb{T})}^2 \leq C T t,$$

and integrating with respect to t yields $\tilde{\mathbb{E}} \|M_{\tilde{R}}^\varepsilon(t, \cdot)\|_{L^2((0, T) \times \mathbb{T})}^2 \leq C$. By [Sim87, Theorem 3], the set

$$\left\{ M \in L^2((0, T) \times \mathbb{T}); \|M\|_{L^2((0, T) \times \mathbb{T})} + \sup_{h \in [0, T/2]} \left\| \frac{M(\cdot + h) - M}{\sqrt{h}} \right\|_{L^2((0, T-h), H^{-1}(\mathbb{T}))} \leq R \right\}$$

is compact in $L^2((0, T), H^{-1}(\mathbb{T}))$. It follows then that $(M_{\tilde{R}}^\varepsilon)_\varepsilon$ is tight in $L^2([0, T], H^{-1}(\mathbb{T}))$. Finally, we can define $\partial_t M_{\tilde{R}}^\varepsilon$ as a $H^{-1}((0, T) \times \mathbb{T})$ -valued random variable and the sequence $(\partial_t M_{\tilde{R}}^\varepsilon)_\varepsilon$ is tight in $H^{-1}((0, T) \times \mathbb{T})$.

Step 2. Since $f \in C_c^\infty(\mathbb{R})$, $\chi_\varepsilon(\xi) \leq \xi^2$ and by the inequality (4.48) we obtain that $(T^\varepsilon)_\varepsilon$ is bounded in $L^p(\tilde{\Omega} \times (0, T) \times \mathbb{T})$ for all $p \in (1, 3/2)$. By the compact embedding $L^p((0, T) \times \mathbb{T}) \hookrightarrow$

$H^{-1}([0, T] \times \mathbb{T})$, it follows that $(T^\varepsilon)_\varepsilon$ is tight in $H^{-1}((0, T) \times \mathbb{T})$. We can therefore conclude that the sequence (6.2) is tight in $H^{-1}((0, T) \times \mathbb{T})$.

Step 3. The counterpart of the analysis above for the variable S^ε is that, for all $g \in C_c^\infty(\mathbb{R})$, the sequence

$$\{g(\tilde{S}^\varepsilon)_t - [c(\tilde{u})g(\tilde{S}^\varepsilon)]_x\}_\varepsilon$$

is tight in $H^{-1}((0, T) \times \mathbb{T})$. Set

$$\tilde{Y}^\varepsilon \stackrel{\text{def}}{=} \begin{pmatrix} f(\tilde{R}^\varepsilon) \\ c(\tilde{u})f(\tilde{R}^\varepsilon) \end{pmatrix}, \quad \tilde{Z}^\varepsilon \stackrel{\text{def}}{=} \begin{pmatrix} c(\tilde{u})g(\tilde{S}^\varepsilon) \\ g(\tilde{S}^\varepsilon) \end{pmatrix},$$

and

$$\tilde{Y} \stackrel{\text{def}}{=} \begin{pmatrix} \langle f(R) \rangle \\ c(\tilde{u})\langle f(R) \rangle \end{pmatrix}, \quad \tilde{Z} \stackrel{\text{def}}{=} \begin{pmatrix} c(\tilde{u})\langle g(S) \rangle \\ \langle g(S) \rangle \end{pmatrix}.$$

We know that $(\tilde{Y}^\varepsilon, \tilde{Z}^\varepsilon)$ converges weakly to (\tilde{Y}, \tilde{Z}) as $\varepsilon \rightarrow 0$ and from the previous steps, we have that $\{\text{div}_{t,x}\tilde{Y}^\varepsilon\}_\varepsilon$ and $\{\text{curl}_{t,x}\tilde{Z}^\varepsilon\}_\varepsilon$ are tight in $H^{-1}((0, T) \times \mathbb{T})$. Then, for any $\alpha > 0$ there exists a compact set $K_\alpha \subset H^{-1}((0, T) \times \mathbb{T})$ such that

$$\text{div}_{t,x}\tilde{Y}^\varepsilon \in K_\alpha \text{ and } \text{curl}_{t,x}\tilde{Z}^\varepsilon \in K_\alpha \quad (6.6)$$

with a probability larger than $1 - \alpha$. If (6.6) is realized, the div-curl lemma ensures that $\tilde{Y}^\varepsilon \cdot \tilde{Z}^\varepsilon$ converges weakly to $\tilde{Y} \cdot \tilde{Z}$ as $\varepsilon \rightarrow 0$. We have therefore

$$\langle f(R)g(S) \rangle = \langle f(R) \rangle \langle g(S) \rangle \quad (6.7)$$

for almost all $(\omega, t, x) \in A_\alpha$ where $\tilde{\mathbb{P}} \times \mathcal{L}^2(A_\alpha) \geq (1 - \alpha)T$. We choose a decreasing sequence $(\alpha_n)_n$ that converges to 0 and assume without loss of generality that $(K_{\alpha_n})_n$ and $(A_{\alpha_n})_n$ to be increasing sequences of sets. Then $A := \cup_n A_{\alpha_n}$ is of full measure, $\tilde{\mathbb{P}} \times \mathcal{L}^2(A) = T$, and (6.7) on A . \square

Remark 6.1. As a consequence of (6.7) and (5.50), and provided $t_0 > 0$, we can partially relax the growth condition (5.66) on the function h in (5.64) and (5.69), to admit the following growth

$$|h(\xi)| \leq C(1 + |\xi^-|^p + |\xi^+|^r), \quad p \in [1, \infty), r \in [1, 2). \quad (6.8)$$

6.2 An evolution equation for the defect measure

Our goal here is to show that the measures $\tilde{\nu}_{t,x}^1$ and $\tilde{\nu}_{t,x}^2$ are Dirac measures, which we will characterize by the identities $\langle R^2 \rangle = \langle R \rangle^2$ and $\langle S^2 \rangle = \langle S \rangle^2$. We introduce the non-negative quantity $\Delta \stackrel{\text{def}}{=} \frac{1}{2} (\langle R^2 \rangle - \langle R \rangle^2)$ and derive an evolution transport equation (inequality, see (6.20)) satisfied by Δ . Since $\Delta(t, \cdot) \rightarrow 0$ when $t \downarrow 0$ (see Proposition 6.4), this will imply $\Delta(t, \cdot) = 0$ at all positive times $t \in [0, T]$.

Since the map $\xi \mapsto \xi^2$ does not satisfy (6.8), we have to work with suitable truncate functions. For $\kappa \geq 0$, we define the truncate function

$$Q_\kappa(\xi) \stackrel{\text{def}}{=} \frac{1}{2} \xi^2 - \frac{1}{2} ((\xi - \kappa)^+)^2. \quad (6.9)$$

Lemma 6.2. For any $T > 0$ we have, $\tilde{\mathbb{P}}$ -a.s., and in $L^2(\Omega)$,

$$\lim_{\kappa \rightarrow \infty} \|\langle Q'_\kappa(R) \rangle - Q'_\kappa(\langle R \rangle)\|_{L^2([0,T] \times \mathbb{T})} = \lim_{\kappa \rightarrow \infty} \|\langle Q'_\kappa(R) \rangle - \langle R \rangle\|_{L^2([0,T] \times \mathbb{T})} = 0. \quad (6.10)$$

Also, for all $\kappa \geq 0$, we have for almost all $(\omega, t, x) \in \tilde{\Omega} \times [0, T] \times \mathbb{T}$,

$$\frac{1}{2} [\langle Q'_\kappa(R) \rangle - Q'_\kappa(\langle R \rangle)]^2 \leq \langle Q_\kappa(R) \rangle - Q_\kappa(\langle R \rangle), \quad (6.11)$$

and

$$\Delta_\kappa \stackrel{\text{def}}{=} \langle Q_\kappa(R) \rangle - Q_\kappa(\langle R \rangle) \leq \Delta \stackrel{\text{def}}{=} \frac{1}{2} [\langle R^2 \rangle - \langle R \rangle^2], \quad (6.12)$$

and, $\tilde{\mathbb{P}}$ -a.s.,

$$\lim_{\kappa \rightarrow \infty} \|\Delta_\kappa - \Delta\|_{L^1([0,T] \times \mathbb{T})} = 0. \quad (6.13)$$

Proof of Lemma 6.2. By differentiation in (6.9), we obtain

$$\xi - Q'_\kappa(\xi) = (\xi - \kappa)^+. \quad (6.14)$$

By the triangular inequality, we have then

$$\begin{aligned} |\langle Q'_\kappa(R) \rangle - Q'_\kappa(\langle R \rangle)| &\leq |\langle Q'_\kappa(R) \rangle - \langle R \rangle| + |Q'_\kappa(\langle R \rangle) - \langle R \rangle| \\ &= \langle R - Q'_\kappa(R) \rangle + \langle R \rangle - Q'_\kappa(\langle R \rangle). \end{aligned}$$

Since $\xi \mapsto (\xi - \kappa)^+$ is convex, we obtain

$$|\langle Q'_\kappa(R) \rangle - Q'_\kappa(\langle R \rangle)| \leq 2\langle (R - \kappa)^+ \rangle,$$

and by Jensen's inequality again,

$$|\langle Q'_\kappa(R) \rangle - Q'_\kappa(\langle R \rangle)|^2 \leq 4\langle [(R - \kappa)^+]^2 \rangle.$$

Since $\langle R^2 \rangle \in L^1(\Omega \times (0, T) \times \mathbb{T})$, (6.10) follows. Let us now establish (6.11). We want to prove that the function

$$\varphi(\kappa) \stackrel{\text{def}}{=} \langle Q_\kappa(R) \rangle - Q_\kappa(\langle R \rangle) - \frac{1}{2} [\langle Q'_\kappa(R) \rangle - Q'_\kappa(\langle R \rangle)]^2$$

is non-negative. We have $\partial_\kappa Q_\kappa(\xi) = (\xi - \kappa)^+ = \xi \vee \kappa - \kappa$ and $Q'_\kappa(\xi) = \xi - \xi \vee \kappa$, so

$$\varphi'(\kappa) = (\langle R \vee \kappa \rangle - \langle R \rangle \vee \kappa) [1 - (\langle \mathbf{1}_{\{R \geq \kappa\}} \rangle - \mathbf{1}_{\{\langle R \rangle \geq \kappa\}})]. \quad (6.15)$$

Since $\xi \mapsto \xi \vee \kappa$ is convex, $\langle R \vee \kappa \rangle - \langle R \rangle \vee \kappa$ is non-negative by Jensen's inequality, while

$$1 - (\langle \mathbf{1}_{\{R \geq \kappa\}} \rangle - \mathbf{1}_{\{\langle R \rangle \geq \kappa\}}) \geq 0. \quad (6.16)$$

We have then $\varphi'(\kappa) \geq 0$, so it will be sufficient to study the case $\kappa = 0$. Set $\xi^- = -\min(\xi, 0)$. We have $Q_0(\xi) = (\xi^-)^2/2$, $Q'_0(\xi) = -\xi^-$, so

$$2\varphi(0) = \langle (R^-)^2 \rangle - (\langle R^- \rangle)^2 - (\langle R^- \rangle - \langle R \rangle^-)^2. \quad (6.17)$$

If $\langle R \rangle \geq 0$, then

$$2\varphi(0) = \langle (R^-)^2 \rangle - (\langle R^- \rangle)^2 \quad (6.18)$$

is non-negative, by Jensen's inequality. If $\langle R \rangle < 0$, then, using Jensen's inequality again, we get

$$2\varphi(0) = \langle (R^-)^2 \rangle - (\langle R \rangle)^2 - (\langle R^- \rangle + \langle R \rangle)^2 \geq -2\langle R \rangle (\langle R^- \rangle + \langle R \rangle) = -2\langle R \rangle \langle R^+ \rangle \geq 0,$$

and this achieves the proof of (6.11). The inequality (6.12) follows from Jensen's inequality since

$$\frac{1}{2}\xi^2 - Q_\kappa(\xi) = \frac{1}{2}((\xi - \kappa)^+)^2 \quad (6.19)$$

is a convex function of ξ . By (6.19) and (6.12), we also have

$$\tilde{\mathbb{E}} \|\Delta_\kappa - \Delta\|_{L^1([0,T] \times \mathbb{T})} = \frac{1}{2} \tilde{\mathbb{E}} \int_0^T \int_{\mathbb{T}} \int_{\mathbb{R}} ((\xi - \kappa)^+)^2 d\tilde{\mu}_{t,x} dx dt - \frac{1}{2} \tilde{\mathbb{E}} \int_0^T \int_{\mathbb{T}} ((\tilde{R} - \kappa)^+)^2 dx dt,$$

and (6.13) follows by dominated convergence. \square

Here is the central proposition of this section.

Proposition 6.3 (Evolution equation for the defect measure). *Let Δ be defined by (6.12). Let $t_0 \in (0, T)$. Then $\tilde{\mathbb{P}}$ -almost surely, we have*

$$\left[\sqrt{\Delta/c(\tilde{u})} \right]_t + \left[\sqrt{c(\tilde{u})\Delta} \right]_x \leq 0 \quad \text{in } \mathcal{D}'((t_0, T) \times \mathbb{T}). \quad (6.20)$$

Proof of Proposition 6.3. The proof breaks into several steps.

Step 1. (Truncation) We derive a first equation for the version Δ_κ of Δ with truncation defined in (6.12). We use the equations (5.64) and (5.69) with $h = Q_\kappa$ (cf. Remark 6.1). Subtracting the second equation from the first one gives us

$$[\Delta_\kappa]_t + [c(\tilde{u})\Delta_\kappa]_x = \frac{1}{2}q[\langle Q_\kappa''(R) \rangle - Q_\kappa''(\langle R \rangle)] + \partial_t M_\kappa \quad (6.21)$$

$$+ \tilde{c}'(\tilde{u})[\langle \Psi_\kappa(R) \rangle - \Psi_\kappa(\langle R \rangle)] + \tilde{c}'(\tilde{u})Q_\kappa'(\langle R \rangle)(\langle R^2 \rangle - \langle R \rangle^2) \quad (6.22)$$

$$+ \tilde{c}'(\tilde{u})[2\langle S \rangle \Delta_\kappa + \langle S^2 \rangle [Q_\kappa'(\langle R \rangle) - \langle Q_\kappa'(R) \rangle]] \quad (6.23)$$

in $\mathcal{D}'((t_0, T) \times \mathbb{T})$, where

$$\Psi_\kappa(\xi) \stackrel{\text{def}}{=} \xi^2 Q_\kappa'(\xi) - 2\xi Q_\kappa(\xi), \quad (6.24)$$

and

$$M_\kappa(t) \stackrel{\text{def}}{=} \int_{t_0}^t [\langle Q_\kappa'(R) \rangle - Q_\kappa'(\langle R \rangle)] \Phi d\tilde{W}(s). \quad (6.25)$$

Step 2. (Preparation for renormalization) Different obstructions prevent us from passing to the limit $[\kappa \rightarrow \infty]$ on the parameter of truncation κ in (6.21)-(6.22)-(6.23). Indeed, we can make no sense of a term $\tilde{c}'(\tilde{u})L(\langle R \rangle, \langle S \rangle)$ (or $\tilde{c}'(\tilde{u})\langle L(R, S) \rangle$) unless L satisfies the growth condition $|L(\xi, \eta)| \leq C(L)(1 + |\xi|^{3-\gamma} + |\eta|^{3-\gamma})$ with $\gamma > 0$ for positive ξ, η . It is therefore not possible to define the limit as $\kappa \rightarrow \infty$ of the three last terms of (6.21)-(6.22)-(6.23). Renormalization will help to solve this problem. Indeed, the first term in (6.23) has a factor Δ_κ , the second term in (6.23) is controlled by $\Delta_\kappa^{1/2}$ as a consequence of (6.11). However, the second term in (6.22) has not the desired form, since this is $2\tilde{c}'(\tilde{u})Q_\kappa'(\langle R \rangle)\Delta$ and not $2\tilde{c}'(\tilde{u})Q_\kappa'(\langle R \rangle)\Delta_\kappa$. We exploit the fact that $Q_\kappa'(\xi) \leq \kappa$ and the inequality (6.12) to write

$$(6.22) \leq \tilde{c}'(\tilde{u})[\langle \Phi_\kappa(R) \rangle - \Phi_\kappa(\langle R \rangle)] + 2\tilde{c}'(\tilde{u})Q_\kappa'(\langle R \rangle)\Delta_\kappa, \quad (6.26)$$

where

$$\Phi_\kappa(\xi) = \Psi_\kappa(\xi) + \kappa(\xi^2 - 2Q_\kappa(\xi)) = -\kappa^2(\xi - \kappa)^+. \quad (6.27)$$

Since Φ_κ is concave, we obtain

$$(6.22) \leq 2 \tilde{c}'(\tilde{u}) Q'_\kappa(\langle R \rangle) \Delta_\kappa, \quad (6.28)$$

and thus

$$[\Delta_\kappa]_t + [c(\tilde{u}) \Delta_\kappa]_x \leq \frac{1}{2} q [\langle Q''_\kappa(R) \rangle - Q''_\kappa(\langle R \rangle)] + \partial_t M_\kappa \quad (6.29)$$

$$+ \tilde{c}'(\tilde{u}) [2 \langle Q'_\kappa(\langle R \rangle) \rangle + \langle S \rangle] \Delta_\kappa + \langle S^2 \rangle [Q'_\kappa(\langle R \rangle) - \langle Q'_\kappa(R) \rangle]. \quad (6.30)$$

Step 3. (Renormalization) Let $H(c, \Delta)$ be a smooth function of its argument. We can do the formal computations

$$\begin{aligned} dH(c(\tilde{u}), \Delta_\kappa) + \partial_x [c(\tilde{u}) H(c(\tilde{u}), \Delta_\kappa)] dt &= \frac{\partial H}{\partial c}(c(\tilde{u}), \Delta_\kappa) [\partial_t c(\tilde{u}) + c(\tilde{u}) \partial_x c(\tilde{u})] dt \\ &+ \frac{\partial H}{\partial \Delta}(c(\tilde{u}), \Delta_\kappa) [d\Delta_\kappa + \partial_x (c(\tilde{u}) \Delta_\kappa)] dt + \frac{1}{2} \frac{\partial^2 H}{\partial \Delta^2}(c(\tilde{u}), \Delta_\kappa) d\langle M_\kappa, M_\kappa \rangle(t) \\ &+ \partial_x c(\tilde{u}) (H - \Delta \frac{\partial H}{\partial \Delta})(c(\tilde{u}), \Delta_\kappa) dt, \end{aligned}$$

where $\langle M_\kappa, M_\kappa \rangle_t$ is the quadratic variation of the martingale M_κ defined in (6.25), *i.e.*

$$\langle M_\kappa, M_\kappa \rangle(t) \stackrel{\text{def}}{=} q \int_{t_0}^t |\langle Q'_\kappa(R) \rangle - Q'_\kappa(\langle R \rangle)|^2 ds. \quad (6.31)$$

Using the formulas

$$\partial_t c(\tilde{u}) = c'(\tilde{u}) \tilde{u}_t = 2c(\tilde{u}) \tilde{c}'(\tilde{u}) (\langle R \rangle + \langle S \rangle), \quad \partial_x c(\tilde{u}) = c'(\tilde{u}) \tilde{u}_x = 2\tilde{c}'(\tilde{u}) (\langle S \rangle - \langle R \rangle),$$

we obtain

$$\begin{aligned} dH_\kappa + \partial_x (c(\tilde{u}) H_\kappa) dt &= 4c(\tilde{u}) \tilde{c}'(\tilde{u}) \frac{\partial H_\kappa}{\partial c} \langle S \rangle dt + \frac{\partial H_\kappa}{\partial \Delta} [d\Delta_\kappa + \partial_x (c(\tilde{u}) \Delta_\kappa)] dt \\ &+ \frac{1}{2} \frac{\partial^2 H_\kappa}{\partial \Delta^2} d\langle M_\kappa, M_\kappa \rangle(t) + 2\tilde{c}'(\tilde{u}) (H_\kappa - \Delta_\kappa \frac{\partial H_\kappa}{\partial \Delta}) (\langle S \rangle - \langle R \rangle) dt, \end{aligned}$$

where, for a function $F = F(c, \Delta)$, F_κ stands for $F(c(\tilde{u}), \Delta_\kappa)$. By (6.29)-(6.30), we then deduce the following inequality in $\mathcal{D}'((t_0, T) \times \mathbb{T})$

$$\begin{aligned} \partial_t H_\kappa + \partial_x (c(\tilde{u}) H_\kappa) &\leq 4c(\tilde{u}) \tilde{c}'(\tilde{u}) \frac{\partial H_\kappa}{\partial c} \langle S \rangle + \frac{1}{2} q \frac{\partial H_\kappa}{\partial \Delta} [\langle Q''_\kappa(R) \rangle - Q''_\kappa(\langle R \rangle)] + \partial_t N_\kappa \\ &+ \frac{\partial H_\kappa}{\partial \Delta} \tilde{c}'(\tilde{u}) [2 \langle Q'_\kappa(\langle R \rangle) \rangle + \langle S \rangle] \Delta_\kappa + \langle S^2 \rangle [Q'_\kappa(\langle R \rangle) - \langle Q'_\kappa(R) \rangle] \\ &+ \frac{1}{2} \frac{\partial^2 H_\kappa}{\partial \Delta^2} \partial_t \langle M_\kappa, M_\kappa \rangle(t) + 2\tilde{c}'(\tilde{u}) (H_\kappa - \Delta_\kappa \frac{\partial H_\kappa}{\partial \Delta}) (\langle S \rangle - \langle R \rangle), \end{aligned} \quad (6.32)$$

where

$$N_\kappa(t) \stackrel{\text{def}}{=} \int_{t_0}^t \frac{\partial H}{\partial \Delta} [\langle Q'_\kappa(R) \rangle - Q'_\kappa(\langle R \rangle)] \Phi d\tilde{W}(s). \quad (6.33)$$

The rigorous derivation of (6.32) uses a preliminary step of convolution and DiPerna-Lions' commutator lemma and the chain-rule for the H^1 -function $c(\tilde{u})$. This is very similar to the proof of Proposition 5.12, so we will not give the details here.

Step 4. (Limit $[\kappa \rightarrow \infty]$) Assume that the function H in (6.32) satisfies the following bounds: there is a constant $C_H \geq 0$ such that, for all $c \in [c_1, c_2]$, for all $\Delta, \Delta' \geq 0$,

$$|\Delta \frac{\partial H}{\partial \Delta}(c, \Delta)| \leq C_H (1 + \Delta^{1/2}), \quad (1 + \Delta^{1/2}) |\frac{\partial H}{\partial \Delta}(c, \Delta)| \leq C_H, \quad (6.34)$$

and

$$|F(c, \Delta) - F(c, \Delta')|^2 \leq C_H |\Delta - \Delta'|, \quad F \in \{H, \frac{\partial H}{\partial c}, \Delta \frac{\partial H}{\partial \Delta}\}. \quad (6.35)$$

Moreover, we assume that for all $c \in [c_1, c_2]$, the map $\Delta \mapsto H(c, \Delta)$ is concave. We consider the limit of (6.32) when $\kappa \rightarrow \infty$ along a given sequence. By taking a subsequence if necessary, we can assume that the convergences

$$\langle Q'_\kappa(R) \rangle - Q'_\kappa(\langle R \rangle) \rightarrow 0, \quad \langle Q'_\kappa(R) \rangle - \langle R \rangle \rightarrow 0 \quad (6.36)$$

are satisfied $\tilde{\mathbb{P}}$ -a.s. in $L^2((0, T) \times \mathbb{R})$ as in (6.10), but also a.s. in (ω, t, x) . From (6.35), we can then deduce the following convergence results in $\mathcal{D}'((t_0, T) \times \mathbb{T})$:

$$\partial_t H_\kappa + \partial_x (c(\tilde{u})H_\kappa) \rightarrow \partial_t H + \partial_x (c(\tilde{u})H), \quad (6.37)$$

$$4c(\tilde{u})\tilde{c}'(\tilde{u})\frac{\partial H_\kappa}{\partial c}\langle S \rangle \rightarrow 4c(\tilde{u})\tilde{c}'(\tilde{u})\frac{\partial H}{\partial c}\langle S \rangle, \quad (6.38)$$

$$2\tilde{c}'(\tilde{u})(H_\kappa - \Delta_\kappa\frac{\partial H_\kappa}{\partial \Delta})(\langle S \rangle - \langle R \rangle) \rightarrow 2\tilde{c}'(\tilde{u})(H - \Delta\frac{\partial H}{\partial \Delta})(\langle S \rangle - \langle R \rangle), \quad (6.39)$$

$$\frac{\partial H_\kappa}{\partial \Delta}\tilde{c}'(\tilde{u})\langle S \rangle \Delta_\kappa \rightarrow \frac{\partial H}{\partial \Delta}\tilde{c}'(\tilde{u})\langle S \rangle \Delta, \quad (6.40)$$

where, for a function $F = F(c, \Delta)$, F stands for $F(c(\tilde{u})(t, x), \Delta(t, x))$. Next, the bounds (6.11), (6.34) and the convergence (6.36) show that

$$\partial_t N_\kappa, \quad \frac{\partial H_\kappa}{\partial \Delta}[\langle Q''_\kappa(R) \rangle - Q''_\kappa(\langle R \rangle)], \quad \tilde{c}'(\tilde{u})\frac{\partial H}{\partial \Delta}\langle S^2 \rangle [Q'_\kappa(\langle R \rangle) - \langle Q'_\kappa(R) \rangle] \quad (6.41)$$

all converge to 0 in $\mathcal{D}'((t_0, T) \times \mathbb{T})$. Let us give few details about the convergence of the last term in (6.41): by (6.11) and (6.34), we have the bound

$$|\tilde{c}'(\tilde{u})\frac{\partial H}{\partial \Delta}\langle S^2 \rangle [Q'_\kappa(\langle R \rangle) - \langle Q'_\kappa(R) \rangle]| \leq C_H \tilde{c}'(\tilde{u})\langle S^2 \rangle, \quad (6.42)$$

while the term converges to 0 $\tilde{\mathbb{P}}$ -a.s., a.e. in (t, x) . The $\tilde{\mathbb{P}}$ -a.s. convergence to 0 follows by dominated convergence. We also have

$$\frac{1}{2}\frac{\partial^2 H_\kappa}{\partial \Delta^2}\partial_t \langle M_\kappa, M_\kappa \rangle(t) \leq 0. \quad (6.43)$$

There remains to examine the term

$$\frac{\partial H_\kappa}{\partial \Delta}\tilde{c}'(\tilde{u})Q'_\kappa(\langle R \rangle)\Delta_\kappa, \quad (6.44)$$

which can be split as

$$F\tilde{c}'(\tilde{u})\langle R \rangle + [F_\kappa - F]\tilde{c}'(\tilde{u})\langle R \rangle - F_\kappa\tilde{c}'(\tilde{u})[\langle R \rangle - Q'_\kappa(\langle R \rangle)], \quad (6.45)$$

with $F = \Delta\frac{\partial H}{\partial \Delta}(c, \Delta)$ and $F_\kappa = \Delta_\kappa\frac{\partial H}{\partial \Delta}(c, \Delta_\kappa)$. The same arguments as above then show that the two last terms in (6.45) converge to 0 and we finally deduce from (6.32) the inequality

$$\partial_t H + \partial_x (c(\tilde{u})H) \leq 2\tilde{c}'(\tilde{u})\left\{2c(\tilde{u})\frac{\partial H}{\partial c}\langle S \rangle + \Delta\frac{\partial H}{\partial \Delta}(\langle R \rangle + \langle S \rangle) + (H - \Delta\frac{\partial H}{\partial \Delta})(\langle S \rangle - \langle R \rangle)\right\},$$

that is to say

$$\partial_t H + \partial_x (c(\tilde{u})H) \leq 2\tilde{c}'(\tilde{u})\left\{[2c(\tilde{u})\frac{\partial H}{\partial c} + H]\langle S \rangle + [2\Delta\frac{\partial H}{\partial \Delta} - H]\langle R \rangle\right\}. \quad (6.46)$$

Step 5. (Conclusion) We take now

$$H(c, \Delta) = H_\delta(c, \Delta) \stackrel{\text{def}}{=} \sqrt{\Delta/c + \delta}, \quad (6.47)$$

where δ is a positive parameter. It is easy to check that the map $\Delta \mapsto H(c, \Delta)$ is concave and that (6.34) and (6.35) are satisfied. In (6.46), we obtain

$$\partial_t H_\delta + \partial_x (c(\tilde{u})H_\delta) \leq 2\delta\tilde{c}'(\tilde{u})\frac{\langle S \rangle - \langle R \rangle}{\sqrt{\Delta/c + \delta}} \leq 2\delta^{\frac{1}{2}}\tilde{c}'(\tilde{u})|\langle S \rangle - \langle R \rangle|. \quad (6.48)$$

Taking the limit $[\delta \rightarrow 0]$ in (6.48) yields (6.20). \square

To fully exploit (6.20), we need to show that $\Delta(t_0) \rightarrow 0$ when $t_0 \rightarrow 0$. This is the content of the following proposition.

Proposition 6.4 (No initial boundary layer). *Recall that Δ is defined by (6.12). Set also*

$$\check{\Delta} \stackrel{\text{def}}{=} \frac{1}{2} (\langle S^2 \rangle - \langle S \rangle^2). \quad (6.49)$$

Then, as $t \rightarrow 0$, we have, $\tilde{\mathbb{P}}$ -almost surely, the convergences

$$\langle R \rangle(t) \rightarrow R_0, \quad \langle S \rangle(t) \rightarrow S_0, \quad \text{in } L^2(\mathbb{T}), \quad (6.50)$$

and

$$\Delta(t), \check{\Delta}(t) \rightarrow 0, \quad \text{in } L^1(\mathbb{T}). \quad (6.51)$$

Proof of Proposition 6.4. Let \bar{Q}_κ denote the truncation of $\xi \mapsto \xi^2/2$ from above and from below defined by

$$\bar{Q}_\kappa(\xi) = \frac{1}{2} \left[\xi^2 - ((\xi - \kappa)^+)^2 - ((\xi + \kappa)^-)^2 \right]. \quad (6.52)$$

We apply (5.41) with $f_1 = f_2 = \bar{Q}_\kappa$ and $\varphi \equiv 1$ to obtain $\tilde{\mathbb{P}}$ -almost surely, for $t > 0$,

$$\int_{\mathbb{T}} [\langle \bar{Q}_\kappa(R) \rangle + \langle \bar{Q}_\kappa(S) \rangle] (t, x) dx = \lim_{\varepsilon \rightarrow 0} \int_{\mathbb{T}} [\bar{Q}_\kappa(\tilde{R}^\varepsilon) + \bar{Q}_\kappa(\tilde{S}^\varepsilon)] (t, x) dx. \quad (6.53)$$

Since $\bar{Q}_\kappa(\xi) \leq \xi^2/2$, the estimate

$$\int_{\mathbb{T}} [\langle \bar{Q}_\kappa(R) \rangle + \langle \bar{Q}_\kappa(S) \rangle] (t, x) dx \leq \frac{1}{2} \liminf_{\varepsilon \rightarrow 0} \|(\tilde{R}^\varepsilon, \tilde{S}^\varepsilon)(t)\|_{L^2(\mathbb{T})}^2 \quad (6.54)$$

follows. We then use the energy inequality (4.7) to get a bound on the right-hand side of (6.54). Indeed, we can pass to the limit in the energy inequality (4.7) since the explicit expression of the martingale \mathcal{M}^ε is given by (4.14) and

$$2 \int_{\mathbb{T}} \int_{t_0}^{\cdot} (\tilde{R}^\varepsilon + \tilde{S}^\varepsilon) \Phi^\varepsilon d\tilde{W}^\varepsilon dx \rightarrow M_e \stackrel{\text{def}}{=} 2 \int_{\mathbb{T}} \int_{t_0}^{\cdot} (\langle R \rangle + \langle S \rangle) \Phi dx \quad \text{in } C([t_0, T]).$$

Using also (5.42) to treat the initial terms, we obtain

$$\int_{\mathbb{T}} [\langle \bar{Q}_\kappa(R) \rangle + \langle \bar{Q}_\kappa(S) \rangle] (t, x) dx \leq \frac{1}{2} \|(R_0, S_0)\|_{L^2(\mathbb{T})}^2 + \|q\|_{L^1(\mathbb{T})} t + \frac{1}{2} M_e(t). \quad (6.55)$$

Taking $\kappa \rightarrow \infty$ in (6.55) then gives us

$$\int_{\mathbb{T}} [\langle R^2 \rangle + \langle S^2 \rangle] (t, x) dx \leq \|(R_0, S_0)\|_{L^2(\mathbb{T})}^2 + 2\|q\|_{L^1(\mathbb{T})} t + M_e(t), \quad (6.56)$$

and so

$$\|(\langle R \rangle, \langle S \rangle)\|_{L^2(\mathbb{T})}^2(t) + 2 \int_{\mathbb{T}} (\Delta(t) + \check{\Delta}(t)) dx \leq \|(R_0, S_0)\|_{L^2(\mathbb{T})}^2 + 2q_0 t + M_e(t). \quad (6.57)$$

The convergence of $(\tilde{R}^\varepsilon, \tilde{S}^\varepsilon)$ to $(\langle R \rangle, \langle S \rangle)$ in $C([0, T], L_w^2(\mathbb{T}))$ has the consequence that

$$\|(R_0, S_0)\|_{L^2(\mathbb{T})}^2 \leq \liminf_{t \rightarrow 0} \|(\langle R \rangle, \langle S \rangle)\|_{L^2(\mathbb{T})}^2. \quad (6.58)$$

From (6.57) and (6.58), we deduce (6.51) and the convergence of the norms

$$\lim_{t \rightarrow 0} \|(\langle R \rangle, \langle S \rangle)\|_{L^2(\mathbb{T})}^2 = \|(R_0, S_0)\|_{L^2(\mathbb{T})}^2, \quad (6.59)$$

which, combined with the weak convergence, gives (6.50). \square

We can now conclude the proof of reduction (to Dirac masses) of the Young measures.

Proposition 6.5 (Reduction of the Yong measures). *For almost all $(\omega, t, x) \in \tilde{\Omega} \times [0, T] \times \mathbb{T}$, we have*

$$\tilde{\mu}_{t,x} = \delta_{(\langle R \rangle(t,x), \langle S \rangle(t,x))}. \quad (6.60)$$

Proof of Proposition 6.5. By (6.20) and (6.51), we have $\Delta(t) = 0$ for all t , which is equivalent to $\nu_{t,x}^1 = \delta_{\langle R \rangle(t,x)}$ $\tilde{\mathbb{P}}$ -a.s., for a.e. (t, x) . We can also establish in a completely similar manner the identity $\nu_{t,x}^2 = \delta_{\langle S \rangle(t,x)}$ $\tilde{\mathbb{P}}$ -a.s., for a.e. (t, x) . Then (6.60) follows from the decomposition (6.1). \square

6.3 Existence of weak martingale solutions

Proof of Theorem 2.5. Once the reduction (6.60) of the Young measure is established, the equations (5.67)-(5.68) take the desired form

$$\begin{aligned} \int_{\mathbb{T}} \langle R \rangle(t, x) \varphi(x) dx &= \int_{\mathbb{T}} R_0(x) \varphi(x) dx + \int_0^t \int_{\mathbb{T}} c(\tilde{u}) \langle R \rangle \varphi_x dx ds \\ &\quad - \int_0^t \int_{\mathbb{T}} \tilde{c}'(\tilde{u}) [\langle R \rangle - \langle S \rangle]^2 \varphi dx ds + \int_{\mathbb{T}} \int_0^t \varphi \Phi d\tilde{W}(s) dx, \end{aligned} \quad (6.61)$$

$$\begin{aligned} \int_{\mathbb{T}} \langle S \rangle(t, x) \varphi(x) dx &= \int_{\mathbb{T}} S_0(x) \varphi(x) dx - \int_0^t \int_{\mathbb{T}} c(\tilde{u}) \langle S \rangle \varphi_x dx ds \\ &\quad - \int_0^t \int_{\mathbb{T}} \tilde{c}'(\tilde{u}) [\langle R \rangle - \langle S \rangle]^2 \varphi dx ds + \int_{\mathbb{T}} \int_0^t \varphi \Phi d\tilde{W}(s) dx. \end{aligned} \quad (6.62)$$

where R_0 and S_0 are defined by (2.10). Summing up (6.61) and (6.62) and using the identities (5.33) with the chain-rule for H^1 -functions we obtain (2.26). We have already establish that $\tilde{R}, \tilde{S} \in L^2(\Omega; L^\infty([0, T]; L^2(\mathbb{T})))$ and $\tilde{R}, \tilde{S} \in C([0, T]; L^2(\mathbb{T}) - \text{weak})$ $\tilde{\mathbb{P}}$ -a.s. Doing the same proof of Proposition 6.4 for any t_0 we obtain the right continuity (2.28). The energy dissipation (2.27) follows from (6.56). Taking into account (5.63), we can conclude to the existence of a weak martingale solution to (2.5). Finally, the estimates (2.30)-(2.31) follow from (5.51)-(4.65). \square

Acknowledgement. This work was supported by the LABEX MILYON (ANR-10-LABX-0070) of Université de Lyon, within the program "Investissements d'Avenir" (ANR-11-IDEX-0007) operated by the French National Research Agency (ANR). It was also supported by the Unité de mathématiques pure et appliquées, UMPA (CNRS and ENS de Lyon).

References

- [Bas11] R. F. Bass. *Stochastic processes*, volume 33 of *Cambridge Series in Statistical and Probabilistic Mathematics*. Cambridge University Press, Cambridge, 2011.
- [Ben95] A. Bensoussan. Stochastic Navier-Stokes equations. *Acta Appl. Math.*, 38(3):267–304, 1995.
- [BV19] F. Berthelin and J. Vovelle. Stochastic isentropic Euler equations. *Ann. Sci. Éc. Norm. Supér. (4)*, 52(1):181–254, 2019.

- [Bil99] P. Billingsley. *Convergence of probability measures*. Wiley Series in Probability and Statistics: Probability and Statistics. John Wiley & Sons Inc., New York, second edition, 1999. A Wiley-Interscience Publication.
- [BCZ15] A. Bressan, G. Chen, and Q. Zhang. Unique conservative solutions to a variational wave equation. *Arch. Ration. Mech. Anal.*, 217(3):1069–1101, 2015.
- [BC05] A. Bressan and A. Constantin. Global solutions of the hunter–saxton equation. *SIAM Journal on Mathematical Analysis.*, 37(3): 996–1026, 2005.
- [BZ06] A. Bressan and Y. Zheng. Conservative solutions to a nonlinear variational wave equation. *Comm. Math. Phys.*, 266(2):471–497, 2006.
- [BZZ07] A. Bressan, P. Zhang and Y. Zheng. Asymptotic variational wave equations. *Arch. Rational Mech. Anal.*, 183, 1, 163–185, 2007.
- [BO11] Z. Brzeźniak and M. Ondreját. Weak solutions to stochastic wave equations with values in Riemannian manifolds. *Comm. Partial Differential Equations*, 36(9):1624–1653, 2011.
- [BOS16] Z. Brzeźniak, M. Ondreját and J. Seidler. Invariant measures for stochastic nonlinear beam and wave equations. *Journal of Differential Equations*, 260(5):4157–4179, 2016.
- [CC24] N. V. Chemetov and F. Cipriano. Weak solution for stochastic Degasperis-Procesi equation. *Journal of Differential Equations*, 382:1–49, 2024.
- [CP21] G. Q. G. Chen and P. H. C. Pang. Nonlinear anisotropic degenerate parabolic-hyperbolic equations with stochastic forcing. *J. Funct. Anal.*, 281(12):Paper No. 109222, 48, 2021.
- [CM02] A. Constantin and L. Molinet. The initial value problem for a generalized Boussinesq equation. *Differential Integral Equations*, 15(9):1061–1072, 2002.
- [D11] C.M. Dafermos. Generalized characteristics and the Hunter–Saxton equation. *Journal of Hyperbolic Differential Equations*, 8:01 159–168, 2011.
- [DPZ14] G. Da Prato and J. Zabczyk. *Stochastic equations in infinite dimensions*, volume 152 of *Encyclopedia of Mathematics and its Applications*. Cambridge University Press, Cambridge, second edition, 2014.
- [DGHT11] A. Debussche, N. Glatt-Holtz, and R. Temam. Local martingale and pathwise solutions for an abstract fluids model. *Phys. D*, 240(14-15):1123–1144, 2011.
- [DHV16] A. Debussche, M. Hofmanová, and J. Vovelle. Degenerate parabolic stochastic partial differential equations: quasilinear case. *Ann. Probab.*, 44(3):1916–1955, 2016.
- [DiP83] R. J. DiPerna. Convergence of approximate solutions to conservation laws. *Arch. Rational Mech. Anal.*, 82(1):27–70, 1983.
- [DL89] R. J. DiPerna and P.-L. Lions. Ordinary differential equations, transport theory and Sobolev spaces. *Invent. Math.*, 98(3):511–547, 1989.
- [FN08] J. Feng and D. Nualart. Stochastic scalar conservation laws. *J. Funct. Anal.*, 255(2):313–373, 2008.

- [GHKP22] L. Galimberti, H. Holden, K. H. Karlsen and P. H. C. Pang. Global existence of dissipative solutions to the Camassa–Holm equation with transport noise. *arXiv:2211.07046*, 2022.
- [GHZ96] R. T. Glassey, J. K. Hunter, and Y. Zheng. Singularities of a variational wave equation. *J. Differential Equations*, 129(1):49–78, 1996.
- [GHZ97] R. T. Glassey, J. K. Hunter, and Y. Zheng. Singularities and oscillations in a nonlinear variational wave equation. *Singularities and Oscillations*, Springer, 37–60, 1997.
- [GI92] A. M. Grundland and E. Infeld. A family of nonlinear Klein–Gordon equations and their solutions. *Journal of Mathematical Physics*, 33(7):2498–2503, 1992.
- [Gue23] B. Guelmame. Global weak solutions of the Serre–Green–Naghdi equations with surface tension. *Annales de l’Institut Henri Poincaré C, Analyse non linéaire*, 2023.
- [Hof13] M. Hofmanová. Degenerate parabolic stochastic partial differential equations. *Stochastic Process. Appl.*, 123(12):4294–4336, 2013.
- [HS12] M. Hofmanová and J. Seidler. On weak solutions of stochastic differential equations. *Stoch. Anal. Appl.*, 30(1):100–121, 2012.
- [HKP21] H. Holden, K. H. Karlsen, and P. H. C. Pang. The Hunter–Saxton equation with noise. *J. Differential Equations*, 270:725–786, 2021.
- [HKP23] H. Holden, K. H. Karlsen, and P. H. C. Pang. Global well-posedness of the viscous Camassa–Holm equation with gradient noise. *Discrete Contin. Dyn. Syst.*, 43(2):568–618, 2023.
- [HS91] J. K. Hunter and R. Saxton. Dynamics of director fields. *SIAM J. Appl. Math.*, 51(6):1498–1521, 1991.
- [HZ95a] J. K. Hunter and Y. Zheng. On a nonlinear hyperbolic variational equation: I. Global existence of weak solutions *Arch. Rat. Mech. Anal*, 129:305–353, 1995.
- [HZ95b] J. K. Hunter and Y. Zheng. On a nonlinear hyperbolic variational equation: II. The zero-viscosity and dispersion limits. *Arch. Rat. Mech. Anal*, 129:355–383, 1995.
- [J97] A. Jakubowski. The almost sure skorokhod representation for subsequences in non-metric spaces. *Theory Probab. Appl.*, 42:167–174, 1997.
- [K75] T. Kato. Quasi-linear equations of evolution, with applications to partial differential equations. *Spectral theory and differential equations.*, Lecture Notes in Math., vol. 448, Springer-Verlag, Berlin, 25–70, 1975
- [KP88] T. Kato and G. Ponce. Commutator estimates and the Euler and Navier–Stokes equations. *Comm. Pure Appl. Math.*, 41(7):891–907, 1988.
- [L96] P.L. Lions. Mathematical topics in fluid mechanics. Vol. 1. Incompressible Models. *Oxford Lecture Series in Mathematics and its Applications*, 3., 1996.
- [Ond10] M. Ondreját. Stochastic nonlinear wave equations in local Sobolev spaces. *Electron. J. Probab.*, 15:no. 33, 1041–1091, 2010.
- [S89] R. A. Saxton. Dynamic instability of the liquid crystal director. *Contemporary Mathematics*, 100: 325–330, 1989.

- [Sim87] J. Simon. Compact sets in the space $L^p(0, T; B)$. *Ann. Mat. Pura Appl. (4)*, 146:65–96, 1987.
- [ZZ98] P. Zhang and Y. Zheng. On oscillations of an asymptotic equation of a nonlinear variational wave equation. *Asymptot. Anal.*, 18(3-4):307–327, 1998.
- [ZZ00] P. Zhang and Y. Zheng. Existence and uniqueness of solutions to an asymptotic equation of a variational wave equation with general data. *Arch. Ration. Mech. Anal.*, 155:49–83 2000.
- [ZZ01] P. Zhang and Y. Zheng. Rarefactive solutions to a nonlinear variational wave equation of liquid crystals. *Comm. Partial Differential Equations*, 26(3-4):381–419, 2001.
- [ZZ03] P. Zhang and Y. Zheng. Weak solutions to a nonlinear variational wave equation. *Arch. Ration. Mech. Anal.*, 166(4):303–319, 2003.
- [ZZ05] P. Zhang and Y. Zheng. Weak solutions to a nonlinear variational wave equation with general data. *Ann. Inst. H. Poincaré C Anal. Non Linéaire*, 22(2):207–226, 2005.
- [ZI92] H. Zorski and E. Infeld. New soliton equations for dipole chains. *A Phys. Rev. Lett.*, 68: 1180–1183 , 1992.