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Interferometric single-pixel imaging with a multicore fiber

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Abstract— Lensless illumination single-pixel imaging with a multicore fiber (MCF) is a computational imaging technique that enables potential endoscopic observations of biological samples at cellular scale. In this work, we show that this technique is tantamount to collecting multiple symmetric rank-one projections (SROP) of a Hermitian *interferometric matrix*—a matrix encoding the spectral content of the sample image. In this model, each SROP is induced by the complex *sketching* vector shaping the incident light wavefront with a spatial light modulator (SLM), while the projected interferometric matrix collects up to $O(Q^2)$ image frequencies for a Q -core MCF. While this scheme subsumes previous sensing modalities, such as raster scanning (RS) imaging with beamformed illumination, we demonstrate that collecting the measurements of M random SLM configurations—and thus acquiring M SROPs—allows us to estimate an image of interest if M and Q scale linearly (up to log factors) with the image sparsity level, hence requiring much fewer observations than RS imaging or a complete Nyquist sampling of the $Q \times Q$ interferometric matrix. This demonstration is achieved both theoretically, with a specific restricted isometry analysis of the sensing scheme, and with extensive Monte Carlo experiments. Experimental results made on an actual MCF system finally demonstrate the effectiveness of this imaging procedure on a benchmark image.

1 Introduction

Recently, the imaging community became more and more interested in lensless solutions, providing opportunities to design imaging systems free from the constraints imposed by traditional camera architectures. Cheaper, lighter, and enabling compressive imaging with large field-of-view (FOV), Lensless Imaging (LI) is convenient for medical applications such as microscopy [1, 2] and *in vivo* imaging [3, 4] where the extreme miniaturization of the imaging probe (diameter $\leq 200 \mu\text{m}$) offers a minimally invasive route to image at depths unreachable in microscopy [5]. Paving the way for deep biological tissues [6] and brain imaging with the capability to produce focal planes at various distances from the fiber tip, intensive research effort emerged for Lensless Endoscopy (LE) using multimode fibers [7, 8, 9, 10] or MultiCore Fibers (MCF) [11, 12, 13].

In the field of Computational Imaging (CI) to which LI belongs, a mathematical model is required to describe the observations as a function of the object to be imaged. Regarding the efficiency aspects, two categories of requirement must be considered when developing CI applications; (i) the model must be physically reliable but also computationally efficient to speed up the reconstruction algorithms (ii) the acquisition method must minimize the number of observations (also called *sample complexity*) needed to accurately estimate the object of interest while remaining fast. In single-pixel MCF-LI, *Speckle Imaging* (SI) consists in randomly shaping the wavefront of the light input to the (weakly coupled) cores entering the MCF to illuminate the entire object with a randomly distributed inten-

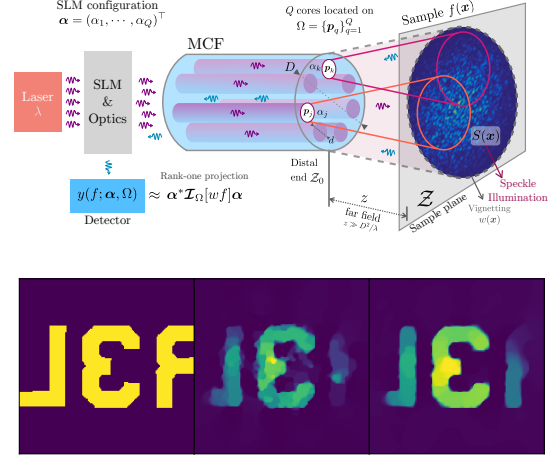


Figure 1: (top) Interferometric LI and its link with ROPs of the interferometric matrix. (bottom) left: ground truth, center: reconstruction with $(Q, M) = (110, 49)$, right: reconstruction with $(Q, M) = (110, 2 \cdot 10^4)$.

sity. The fraction of the light re-emitted (either at other wavelengths by fluorescence or by simple reflection) is integrated in a single-pixel sensor, playing the role of a complete projection of the speckle on the object. Compared to Raster Scanning (RS) the object with a translating focused spot [11], SI has been shown to reduce the sample complexity [14].

In this work, we push a leap forward towards real-time compressive LE, keeping the low sample complexity enabled by SI and introducing light propagation physics in the forward model of MCF imaging using a wavefront-shaping device. We point out that inserting the physics yields an interferometric sensing model similar to radio-interferometry applications [15, 16], where the interferences of the light emitted by the cores composing the MCF give specific access to the Fourier content of the object to be imaged. The sampling complexity of the underlying model is analyzed both theoretically and experimentally.

2 Sensing model

Considering an MCF with diameter D and Q fiber cores (see Fig. 1-top) whose locations on the MCF distal end Z_0 are in $\Omega := \{\mathbf{p}_q\}_{q=1}^Q \subset \mathbb{R}^2$, and assuming a planar sample in the plane Z at a distance z from Z_0 , by optically shaping the light wavefront with an SLM, we can set to $\alpha_q \in \mathbb{C}$ the complex amplitude of the electromagnetic field at each fiber core \mathbf{p}_q . Writing $\alpha = (\alpha_1, \dots, \alpha_Q)^T \in \mathbb{C}^Q$ and $\mathbf{x} \in \mathbb{R}^2$ a point on Z , under the far-field approximation $z \gg D^2/\lambda$ (with λ the laser wavelength) the illumination produced by the MCF on Z reads [14]

$$S(\mathbf{x}; \alpha) \approx w(\mathbf{x}) \left| \sum_{q=1}^Q \alpha_q e^{\frac{2\pi i}{\lambda z} \mathbf{p}_q^T \mathbf{x}} \right|^2,$$

where w is a smooth vignetting function—a Gaussian envelope with a diameter inversely proportional to the fiber cores

diameter—determining the FOV. In RS mode, a focused beam can be obtained in \mathcal{Z} when the Q fiber core locations in Ω are arranged in a Fermat’s golden spiral shape [17]; we will restrict our analysis to this configuration. The LE collects a fraction $c \in (0, 1)$ of the light y globally re-emitted by the sample—as modeled by the fluorophore density map $f(\mathbf{x})$ —under the illumination S . For short time exposure and low intensity illumination, fluorescence theory provides (in a noiseless regime)

$$\begin{aligned} y(f; \boldsymbol{\alpha}, \Omega) &= c \int_{\mathbb{R}^2} S(\mathbf{x}; \boldsymbol{\alpha}) f(\mathbf{x}) d\mathbf{x} \\ &= c \sum_{j,k=1}^Q \alpha_j^* \alpha_k \int_{\mathbb{R}^2} e^{\frac{2\pi i}{\lambda z} (\mathbf{p}_k - \mathbf{p}_j)^\top \mathbf{x}} w(\mathbf{x}) f(\mathbf{x}) d\mathbf{x} \in \mathbb{R}_+, \end{aligned}$$

with the number of collected photons y being Poisson distributed. Therefore, introducing the *interferometry* matrix $\mathcal{I}_\Omega[g] \in \mathbb{C}^{Q \times Q}$ such that, for a function $g : \mathbb{R}^2 \rightarrow \mathbb{R}$, $(\mathcal{I}_\Omega[g])_{jk} := \int_{\mathbb{R}^2} e^{\frac{2\pi i}{\lambda z} (\mathbf{p}_k - \mathbf{p}_j)^\top \mathbf{x}} g(\mathbf{x}) d\mathbf{x}$ with $\mathcal{I}_\Omega[g]$ Hermitian, assuming $c = 1$ and considering the scenario where we collect M LE observations $\mathbf{y} = (y_1, \dots, y_M)^\top$, each associated with a specific $\boldsymbol{\alpha}_m$ for $1 \leq m \leq M$,

$$y(f; \boldsymbol{\alpha}_m, \Omega) = \boldsymbol{\alpha}_m^* \mathcal{I}_\Omega[f^\circ] \boldsymbol{\alpha}_m = \langle \boldsymbol{\alpha}_m \boldsymbol{\alpha}_m^*, \mathcal{I}_\Omega[f^\circ] \rangle_F,$$

with $f^\circ := wf$ is the image f vignetted by w . Under a high photon counting regime, and gathering all possible noise sources in an additive term \mathbf{n} , we can thus compactly write the SI sensing as

$$\mathbf{y} = \mathcal{A} \circ \mathcal{I}_\Omega[f^\circ] + \mathbf{n}, \text{ with } (\mathcal{A}[\mathbf{H}])_m := \boldsymbol{\alpha}_m^* \mathbf{H} \boldsymbol{\alpha}_m, \quad (1)$$

for $\mathbf{H} = \mathbf{H}^*$ and $1 \leq m \leq M$.

We thus observe that (1) is tantamount to first sampling the Fourier transform of f° over frequencies selected in the difference set $\mathcal{V} := \frac{2\pi}{\lambda z}(\Omega - \Omega) = \left\{ \frac{2\pi}{\lambda z}(\mathbf{p}_j - \mathbf{p}_k) \right\}_{j,k=1}^Q$, and next performing M symmetric rank-one projections (SROP [18, 19]) of $\mathcal{I}_\Omega[f^\circ]$ as determined by \mathcal{A} and the complex amplitude vectors $\{\boldsymbol{\alpha}_m\}_{m=1}^M$. Thus, the SI sensing corresponds to a specific interferometric system: assuming we collect enough SROP observations, we can potentially estimate the interferometry matrix $\mathcal{I}_\Omega[f^\circ]$. The system is thus equivalent to the radio-interferometry principles [16]—each fiber core playing somehow the role of a radio telescope and each entry of $(\mathcal{I}_\Omega[f^\circ])_{jk}$ probing the frequency content of f° on the “visibility” $\nu_{jk} := \frac{2\pi}{\lambda z}(\mathbf{p}_j - \mathbf{p}_k)$.

3 Interferometric structural models

If one aims to image a vignetted sample that is K -sparse in the spatial domain, *i.e.*, it is composed of a few spikes as $f^\circ(\mathbf{x}) = \sum_{i=1}^K \rho_i \delta(\mathbf{x} - \mathbf{x}_i)$ for $K \ll Q$, the entry “ jk ” of the interferometric matrix reads $\mathcal{I}_\Omega[f^\circ]_{jk} = \sum_{i=1}^K \rho_i e^{i2\pi(\mathbf{p}_k - \mathbf{p}_j)^\top \mathbf{x}_i}$. We can therefore write the interferometric matrix as

$$\mathcal{I}_\Omega[f^\circ] = \sum_i \rho_i \mathbf{u}(\mathbf{x}_i) \mathbf{u}^*(\mathbf{x}_i), \quad \mathbf{u}(\mathbf{x})_j := e^{-i2\pi \mathbf{p}_j^\top \mathbf{x}},$$

which shows that it is low-rank with rank K . More generally, for a sample f° that can be assumed sparsely represented in a collection of functions $\{\psi_k\}_{k=1}^d$ (*e.g.*, a wavelet basis for bandlimited function supported inside Ω), *i.e.*, $f^\circ(\mathbf{x}) = \sum_{k=1}^d \rho_k \psi_k(\mathbf{x})$ with $\|\boldsymbol{\rho}\|_0 = K \ll d$, then the interferometric matrix belongs to a subspace of dimension K and writes as $\mathcal{I}_\Omega[f^\circ] = \sum_{k|\rho_k \neq 0} \rho_k \mathcal{I}_\Omega[\psi_k]$.

4 Image reconstruction

Assuming the sample $f^\circ \in \Omega$ is band-limited, we are interested in accurately estimating a discretisation $\mathbf{f} \in \mathbb{R}^N$ of f° . We consider a discretisation of (1) that reads

$$\mathbf{y} = \mathcal{A} \circ \tilde{\mathcal{I}}_\Omega[\mathbf{f}] + \mathbf{n}, \text{ with } \tilde{\mathcal{I}}_\Omega[\mathbf{f}] = \mathbf{R}_{\tilde{\mathcal{V}}} \mathbf{F} \mathbf{f},$$

where \mathbf{F} is the Fourier matrix and $\mathbf{R}_{\tilde{\mathcal{V}}} : \mathbb{C}^N \mapsto \mathbb{C}^{Q \times Q}$ is the restriction to the set $\tilde{\mathcal{V}}$ obtained as a Cartesian gridding of the (off-grid) difference set \mathcal{V} (reached by nearest neighbors). From the factorization of this model, we first conclude that the set $\tilde{\mathcal{V}}$ should ideally be composed of as many distinct frequencies as possible (except for the zero frequency that has multiplicity Q) to improve our knowledge of \mathbf{f} . Interestingly, we can show numerically that Fermat’s gold spiral arrangement ensures the unicity of the visibilities $\tilde{\nu}_{jk}$ when $j \neq k$, *i.e.*, $\tilde{\mathcal{V}}$ is composed of $Q(Q-1) + 1$ distinct frequencies. In a noiseless scenario, we have shown that there exists a combination of $M_0 = O(Q^2)$ deterministic ROP observations that exactly recovers $\tilde{\mathcal{I}}_\Omega[f^\circ]$. Therefore, in a compressive setting, we could first leverage the low-complexity structure of $\tilde{\mathcal{I}}_\Omega[f^\circ]$ —as induced from that of \mathbf{f} —to recover $\tilde{\mathcal{I}}_\Omega[\mathbf{f}]$ from $M < M_0$ random complex ROPs, and then infer \mathbf{f} from its $Q(Q-1) + 1$ frequencies encoded in $\tilde{\mathcal{I}}_\Omega[\mathbf{f}]$, this second step being similar to the inverse problem posed in radio-interferometry [16, 15]. In a simpler case where the cores are placed at all integer positions, the interferometric matrix is shown to be circulant and low-rank, *i.e.*, $\tilde{\mathcal{I}}_\Omega[\mathbf{f}] = \mathcal{T} \circ \mathbf{F} \mathbf{f}$ where $\mathcal{T} : \mathbb{C}^N \mapsto \mathbb{C}^{N \times N}$ is the operator that turns a vector into a circulant matrix. In this situation, the sensing operator \mathcal{B} respects a specific RIP- ℓ_2/ℓ_1 property over the set of sparse images—thus extending former approaches restricted to real sparse and low-rank matrices [18], and the RIP of random partial Fourier sensing characterized by $\tilde{\mathcal{V}}$ [20]. Proposition 4.1 shows that proving this RIP- ℓ_2/ℓ_1 for \mathcal{B} implies we can reliably estimate it in a single basis pursuit denoising program (BPDN) with an ℓ_1 fidelity term. This happens with high probability if the vectors $\{\boldsymbol{\alpha}_m\}_{m=1}^M$ are random and sub-Gaussian, and both M and Q^2 are large compared to the sparsity level of \mathbf{f} .

Proposition 4.1 (ℓ_2/ℓ_1 instance optimality of BPDN $_{\ell_1}$). *Let $\mathcal{B} := \mathcal{A} \circ \mathcal{T} \circ \mathbf{F}$ be an operator that respects the RIP $_{\ell_2/\ell_1}(k, \alpha_k, \beta_k)$ for $k \in \{K, K + K'\}$ with $K' > 2K$, and $\frac{1}{\sqrt{2}} m_{K+K'} - M_{K'} \frac{\sqrt{K}}{\sqrt{K'}} \geq \gamma > 0$, for some $M_{K'} > 0$. Then, $\forall \mathbf{f} \in \mathbb{R}^N$, the estimate*

$$\hat{\mathbf{f}} \in \arg \min_{\mathbf{u}} \|\mathbf{u}\|_1 \text{ s.t. } \left\| \underbrace{\mathcal{B}(\mathbf{f}) + \mathbf{n} - \mathcal{B}(\mathbf{u})}_{\mathbf{y}} \right\|_1 \leq \epsilon$$

satisfies

$$\|\mathbf{f} - \hat{\mathbf{f}}\|_2 \leq C \frac{\|\mathbf{f} - \mathbf{f}_K\|_1}{\sqrt{K}} + D \frac{\epsilon}{m}$$

The current theoretical derivations are accompanied by numerical (not shown in this abstract) and experimental reconstruction results (see lead-in in Fig. 1-bottom) that suggest Proposition 4.1 may also hold when relaxing the cited assumptions. Additionally, this may extend to other optimisation problems like LASSO [21] or lagrangian formulations with various regularization terms (ℓ_1 in the identity or orthonormal basis, total variation).

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