



HAL
open science

On slice measures of Green currents on $CP(2)$

Christophe Dupont, Virgile Tapiero

► **To cite this version:**

Christophe Dupont, Virgile Tapiero. On slice measures of Green currents on $CP(2)$. 2023. hal-04256303v2

HAL Id: hal-04256303

<https://hal.science/hal-04256303v2>

Preprint submitted on 23 Feb 2024

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

On slice measures of Green currents on $\mathbb{C}\mathbb{P}^2$

Christophe Dupont and Virgile Tapiero

19th december 2023

Abstract

Let f be a holomorphic map of $\mathbb{C}\mathbb{P}^2$ of degree $d \geq 2$, let T be its Green current and $\mu = T \wedge T$ be its equilibrium measure. We give a new proof of a theorem due to Dujardin asserting that $\mu \ll T \wedge \omega_{\mathbb{P}^2}$ implies $\lambda_2 = \frac{1}{2} \text{Log } d$, where $\lambda_1 \geq \lambda_2$ are the Lyapunov exponents of μ . Then, assuming $\mu \ll T \wedge \omega_{\mathbb{P}^2}$, we study slice measures $\nu := T \wedge dd^c|W|^2$, where W is a holomorphic local submersion. We give sufficient conditions on the Radon-Nikodym derivative of μ with respect to the trace measure $T \wedge \omega_{\mathbb{P}^2}$ ensuring $\mu = \nu$. The involved submersion W comes from normal coordinates for the inverse branches of the iterates of f .

Keywords— Holomorphic dynamics, Equilibrium measure, Green current, Lyapunov exponents, normal forms. *MSC 2020* : 32H50, 32U40, 37C40, 37D25

1 Introduction

Let f be a holomorphic map of \mathbb{P}^2 of degree $d \geq 2$. The Green current T and the equilibrium measure μ are invariant objects encoding the dynamical properties of f , we refer to the books [9], [19] by Dinh and Sibony. We recall that $T := \lim_n \frac{1}{d^n} f^{n*} \omega_{\mathbb{P}^2}$, where $\omega_{\mathbb{P}^2}$ is the normalized Fubini-Study $(1, 1)$ -form of \mathbb{P}^2 , this is a closed positive current, with local Hölder potentials. It satisfies $f^*T = dT$ and its auto-intersection $\mu := T \wedge T$ defines an invariant probability measure on \mathbb{P}^2 , which is mixing and satisfies $f^*\mu = d^2\mu$.

Berteloot-Loeb [5] proved that T is a smooth and non degenerate positive $(1, 1)$ -form on a non empty open subset of \mathbb{P}^2 if and only if f is a Lattès map. Berteloot-Dupont [1] established later that Berteloot-Loeb's condition on T characterizes the condition $\mu \ll \text{Leb}_{\mathbb{P}^2}$. Another characterization of Lattès maps involves the Lyapunov exponents $\lambda_1 \geq \lambda_2$ of μ (see Theorem 2.1 for their definitions). Briend-Duval [6] proved that $\lambda_2 \geq \frac{1}{2} \text{Log } d$. The equality $\lambda_1 = \lambda_2 = \frac{1}{2} \text{Log } d$ holds if and only if $\mu \ll \text{Leb}_{\mathbb{P}^2}$, see [12].

It is now natural to characterize the mappings satisfying $\lambda_1 > \lambda_2 = \frac{1}{2} \text{Log } d$. Hopefully, the presence of a minimal Lyapunov exponent should be equivalent to regular properties for T and μ , and perhaps (in some sense) to the existence of a one dimensional Lattès-like factor. The following result, due to Dujardin, nicely fits into this program.

Theorem 1.1 (Dujardin [10, Theorem 3.6]). *If $\mu \ll \sigma_T := T \wedge \omega_{\mathbb{P}^2}$ then $\lambda_2 = \frac{1}{2} \text{Log } d$.*

He also asked the question of the reverse implication, a partial answer is provided in [20]. The measure σ_T is called the trace of T and carries its mass. The proof of Dujardin's theorem is based on a construction of a df -invariant sub-bundle $\mathcal{T} \subset T\mathbb{P}^2$ (called Fatou directions) of rank ≥ 1 that satisfies for σ_T -almost every x and for any $\vec{v} \in \mathcal{T}_x \setminus \{0\}$:

$$\limsup_{n \rightarrow +\infty} \frac{1}{n} \text{Log} \|d_x f^n(\vec{v})\| \leq \frac{1}{2} \text{Log } d. \quad (1)$$

Indeed, using $\mu \ll \sigma_T$ and Oseledec's Theorem 2.1, for μ -almost every x , the lim sup in (1) tends to λ_1 or λ_2 , which implies in both cases $\lambda_2 \leq \frac{1}{2} \text{Log } d$. Briend-Duval inequality $\lambda_2 \geq \frac{1}{2} \text{Log } d$ then implies the equality. We note that these arguments only require that $\mu \ll \sigma_T$ on an open subset U charged by μ .

On a chart $U \subset \mathbb{P}^2$ equipped with holomorphic coordinates (Z, W) we have :

$$\sigma_T \asymp (T \wedge dd^c|Z|^2 + T \wedge dd^c|W|^2).$$

By construction $\text{Supp}(\mu) \subset \text{Supp}(\sigma_T)$, hence if $\mu(U) > 0$ then $(T \wedge dd^c|Z|^2)(U) > 0$ or $(T \wedge dd^c|W|^2)(U) > 0$. Actually in this case $T \wedge dd^c|Z|^2$ and $T \wedge dd^c|W|^2$ both charge U . The idea is that if $T \wedge dd^c|Z|^2 \equiv 0$ on U , then the potentials of T would be harmonic on almost every vertical disc contained in U , which implies that μ is null on U , see [13, §3.3].

The measures $T \wedge dd^c|Z|^2$ and $T \wedge dd^c|W|^2$ are called *slices* of T . Dujardin's theorem gives an information on λ_2 when $\mu \ll \sigma_T$, the proof is obtained by applying general results concerning Fatou directions [10]. Our purpose in this article is to analyse more deeply the relations between μ and slices of T when $\mu \ll \sigma_T$.

We first provide another proof of Dujardin's Theorem by using normal forms for the generic inverse branches of f^n and forward recurrent properties. We then provide another proof (of a weaker version of) Theorem 1.1 using backward recurrent properties (we indeed assume that the density of μ with respect to T is bounded). The interest of this second proof is to introduce a decomposition of μ using normal coordinates. This decomposition is the cornerstone of the proof of Theorem 1.2 that we now introduce.

Relations between μ and slices of T can be obtained when f preserves a pencil of lines, given for instance by the meromorphic function $\pi[z : w : t] = [z : w]$. Dupont-Taffin [14] proved that in this situation μ and T are related by the formula :

$$\mu = T \wedge \pi^* \mu_\theta \quad (\text{which implies } \mu \ll \sigma_T), \quad (2)$$

where θ is the rational map satisfying $\pi \circ f = \theta \circ \pi$. We note that Jonsson [15] previously established an analogous formula for polynomial skew products on \mathbb{C}^2 . On another hand, Berteloot-Loeb [4] proved that if θ is a Lattès map, then for every $a \in \mathbb{P}^1$ outside a finite subset, there exists a holomorphic coordinate ζ_a such that $\mu_\theta = dd^c|\zeta_a|^2$ on a neighborhood V_a of a (μ_θ is the equilibrium measure of θ). Combining this result with (2), we obtain on $\pi^{-1}(V_a)$:

$$\mu = T \wedge dd^c|W|^2, \text{ with } W := \zeta_a \circ \pi. \quad (3)$$

This gives an example where a slice of T is equal to μ . Theorem 1.2 shows a similar formula. We assume $\mu \ll \sigma_T$ and denote $\psi \in L^1(\sigma_T)$ the Radon-Nikodym derivative of μ with respect to σ_T . Let $J := \text{Supp}(\mu)$ and $\psi|_J$ be the restriction of ψ on J . Note that the measurable function ψ is defined on \mathbb{P}^2 , it is not unique since one can modify it outside J by still preserving $\mu = \psi \sigma_T$. However, one has $\mu = (\psi \mathbb{1}_J) \sigma_T$ on \mathbb{P}^2 for every version of ψ . We denote $\psi_\Omega := \psi \mathbb{1}_\Omega$ for every borel subset Ω of \mathbb{P}^2 . We also denote by $\hat{x} = (x_n)_{n \in \mathbb{Z}}$ the full orbits of the mapping f , by $\pi_0 : \hat{x} \mapsto x_0$ the projection map and by $\hat{\mu}$ the π_0 -pullback of μ on the set of full orbits, see Section 2.2. In the next statement, the local submersion $W_{\hat{x}} : B(x_0, \eta_\varepsilon(\hat{x})) \rightarrow \mathbb{C}$ is provided by the normal form Theorem 2.2.

Theorem 1.2. *Assume that $\mu = \psi_J \sigma_T$ for some measurable function ψ on \mathbb{P}^2 (in particular $\lambda_2 = \frac{1}{2} \text{Log } d$) and that λ_1 is strictly larger than λ_2 .*

1. If $\psi|_J$ is continuous, then for $\hat{\mu}$ -almost every \hat{x} there exist $r_1(\hat{x}) \leq \eta_\varepsilon(\hat{x})$ and $C_{\hat{x}} > 0$ such that :

$$\mu = C_{\hat{x}} (T \wedge dd^c |W_{\hat{x}}|^2) |_J \text{ on } B(x_0, r_1(\hat{x})).$$

2. If ψ is continuous on an open neighborhood V of J and if $\mu = \psi_V \sigma_T$, then there exist $r_2(\hat{x}) \leq \eta_\varepsilon(\hat{x})$ and $C_{\hat{x}} > 0$ such that :

$$\mu = C_{\hat{x}} T \wedge dd^c |W_{\hat{x}}|^2 \text{ on } B(x_0, r_2(\hat{x})).$$

The first Item implies that for $\hat{\mu}$ -almost every \hat{x}, \hat{y} satisfying $\pi_0(\hat{x}) = \pi_0(\hat{y})$, the measures $(T \wedge dd^c |W_{\hat{x}}|^2) |_J$ and $(T \wedge dd^c |W_{\hat{y}}|^2) |_J$ coincide on a small ball centered at $\pi_0(\hat{x})$. In particular, the germ of measure $(T \wedge dd^c |W_{\hat{x}}|^2) |_J$ does not depend on the full orbit \hat{x} but only on the π_0 -projection of \hat{x} . A similar remark holds for the second Item.

Let us also observe that the hypothesis $\lambda_1 > \lambda_2$ in Theorem 1.2 is not restrictive. Assume indeed that $\mu = \psi_J \sigma_T$ and that $\lambda_1 = \lambda_2$. Then the Lyapunov exponents are both minimal (equal to $\frac{1}{2} \text{Log } d$) and f is a Lattès mapping on \mathbb{P}^2 . In this case there exists a finite ramified covering $\sigma : \mathbb{C}^2/\Lambda \rightarrow \mathbb{P}^2$ such that σ^*T is equal to a positive definite hermitian form H with constant coefficients (Λ being a cocompact real lattice in \mathbb{C}^2), see [11]. This implies that, around every point p outside an algebraic subset of \mathbb{P}^2 (the critical values of σ), the current T coincides with H in local coordinates given by σ . It follows that, for every local submersion W defined near p , the measure $\mu = T \wedge T$ is equal to $T \wedge dd^c |W|^2$ multiplied by a smooth positive function.

Let us finally note that the assumption of Item 2 morally replaces the function $\psi_J = \psi \mathbb{1}_J$ of the first Item (which is singular due to the multiplication by $\mathbb{1}_J$) by a continuous function on a neighborhood V of J . More precisely, it requires the existence of an open neighborhood V of J such that $\psi|_V$ is continuous and such that $\text{Supp}(\psi \sigma_T)|_V = J$. This assumption is fulfilled if one assumes that the restriction $\psi|_J$ is continuous and that $\text{Supp } T = J$, by applying Tietze-Urysohn theorem to $\psi|_J$.

Question 1.3. *Is it sufficient to assume $\psi|_J$ continuous to get $\mu = T \wedge dd^c |W_{\hat{x}}|^2$ locally?*

We remark that the continuity of $\psi|_J$ is satisfied by suspensions of Lattès maps on \mathbb{P}^1 , we provide explicit computations of T, μ and ψ in Section 8.

Acknowledgements : The first author thanks the Simons Foundation, Laura DeMarco and Mattias Jonsson for the invitation at the Simons Symposium on Algebraic, Complex and Arithmetic Dynamics, in Schloss Elmau in August 2022. His research is partially funded by the European Research Council (ERC GOAT 101053021). The two authors benefited from the support of the French government "Investissements d'Avenir" program integrated to France 2030 (ANR-11-LABX-0020-01). We thank the referee, whose advices and careful reading enabled us to improve the exposition of the article.

2 Classical results

2.1 Lyapunov exponents

Let f be a holomorphic mapping of \mathbb{P}^2 of degree $d \geq 2$. Let T be its Green current and $\mu = T \wedge T$ be its equilibrium measure, it is f -invariant and ergodic. Since T has local continuous psh potentials, μ integrates local psh functions. In particular μ does not charge analytic subsets of \mathbb{P}^2 and $\mu(\mathcal{C}) = 0$, where $\mathcal{C} := \bigcup_{n \in \mathbb{Z}} f^n(\text{Crit } f)$. It also allows to define the Lyapunov exponents $\lambda_1 \geq \lambda_2$:

$$\lambda_1 = \lim_{n \rightarrow +\infty} \frac{1}{n} \int_{\mathbb{P}^2} \text{Log} \|df^n\| d\mu \text{ and } \lambda_1 + \lambda_2 = \int_{\mathbb{P}^2} \text{Log} |\det_{\mathbb{C}} df| d\mu.$$

Briend-Duval [6] proved that $\lambda_2 \geq \frac{1}{2} \text{Log } d$. The Oseledec theorem states as follows.

Theorem 2.1 (Oseledec).

1. If $\lambda_1 = \lambda_2 = \lambda$, then there exists an invariant borel subset A_{os} of full μ -measure and disjoint from \mathcal{C} such that for every $x \in A_{os}$:

$$\forall \vec{v} \in T_x \mathbb{P}^2 \setminus \{0\}, \quad \lim_{n \rightarrow +\infty} \frac{1}{n} \text{Log} \|d_x f^n(\vec{v})\| = \lambda.$$

2. If $\lambda_1 > \lambda_2$, then there exists an invariant borel subset A_{os} of full μ -measure and disjoint from \mathcal{C} such that for every $x \in A_{os}$, there exists $v_s(x) \in \mathbb{P}(T_x \mathbb{P}^2)$ satisfying :

$$\forall \vec{v} \in T_x \mathbb{P}^2 \setminus v_s(x), \quad \lim_{n \rightarrow +\infty} \frac{1}{n} \text{Log} \|d_x f^n(\vec{v})\| = \lambda_1.$$

$$\forall \vec{v} \in v_s(x) \setminus \{0\}, \quad \lim_{n \rightarrow +\infty} \frac{1}{n} \text{Log} \|d_x f^n(\vec{v})\| = \lambda_2.$$

Moreover v_s is measurable and satisfies $[d_x f](v_s(x)) = v_s(f(x))$ for every $x \in A_{os}$, where $[d_x f]$ is the projectivization of the tangent map $d_x f$.

2.2 Normal forms for inverse branches

We denote $X := \mathbb{P}^2 \setminus \mathcal{C}$ which is totally invariant by f , and we denote \widehat{X} the set of all sequences $\widehat{x} = (x_n)_{n \in \mathbb{Z}}$ of elements of X that are orbits under the action of f i.e. $x_{n+1} = f(x_n)$. Equipped with the natural projection $\pi_0 : \widehat{X} \rightarrow X$ defined by $\pi_0(\widehat{x}) = x_0$, the dynamical system (X, f, μ) admits a natural extension $(\widehat{X}, \widehat{f}, \widehat{\mu})$, where \widehat{f} is the left shift on \widehat{X} and $\widehat{\mu}$ is the unique probability measure on \widehat{X} invariant by \widehat{f} such that $(\pi_0)_* \widehat{\mu} = \mu$. The measure $\widehat{\mu}$ is ergodic as μ , see [8, Chapter 10].

Given an orbit $\widehat{x} \in \widehat{X}$, there exists a family of inverse branches $(f_{\widehat{x}}^{-n})_{n \in \mathbb{N}}$ such that $f_{\widehat{x}}^{-n}$ is defined on $B(x_0, \eta_\varepsilon(\widehat{x}))$ for every $n \geq 0$, and such that $\text{Lip}(f_{\widehat{x}}^{-n}) \leq \beta_\varepsilon(\widehat{x}) e^{-n(\lambda_2 - \varepsilon)}$. The functions $\widehat{x} \mapsto \eta_\varepsilon(\widehat{x})$ and $\widehat{x} \mapsto \beta_\varepsilon(\widehat{x})$ are ε -tempered, i.e. they satisfy :

$$\forall n \in \mathbb{Z}, \quad e^{-|n|\varepsilon} \varphi(\widehat{x}) \leq \varphi(\widehat{f}^n(\widehat{x})) \leq e^{+|n|\varepsilon} \varphi(\widehat{x}).$$

We refer to Briend-Duval [6] (see also [7, 12]) for details about the construction of the inverse branches. The following result provides normal forms for those mappings, see [2, 3, 16].

Theorem 2.2. *For every $\varepsilon > 0$ small enough with respect to the Lyapunov exponents $\lambda_1 \geq \lambda_2$, there exists an invariant borel subset Λ_ε of full $\widehat{\mu}$ -measure, ε -tempered functions $\rho_\varepsilon : \Lambda_\varepsilon \rightarrow]0, 1]$, $\beta_\varepsilon, L_\varepsilon, M_\varepsilon : \Lambda_\varepsilon \rightarrow [1, +\infty[$ and a function $N : \Lambda_\varepsilon \rightarrow \mathbb{N}$ satisfying the following properties. For every $\widehat{x} \in \Lambda_\varepsilon$, there exist injective holomorphic mappings*

$$\xi_{\widehat{x}} : B(x_0, \eta_\varepsilon(\widehat{x})) \longrightarrow \mathbb{D}^2(\rho_\varepsilon(\widehat{x}))$$

satisfying :

1. $\xi_{\widehat{x}}(x_0) = 0$ and $d_{x_0} \xi_{\widehat{x}}(v_s(x_0))$ is the vertical axis in \mathbb{C}^2 ,
2. $\forall p, q \in B(x_0, \eta_\varepsilon(\widehat{x})), \quad \frac{1}{2} \text{dist}(p, q) \leq \|\xi_{\widehat{x}}(p) - \xi_{\widehat{x}}(q)\| \leq \beta_\varepsilon(\widehat{x}) \text{dist}(p, q)$,
3. the following diagram commutes for every $n \geq N(\widehat{x})$:

$$\begin{array}{ccc} B(x_{-n}, \eta_\varepsilon(\widehat{x}_{-n})) & \xleftarrow{f_{\widehat{x}}^{-n}} & B(x_0, \eta_\varepsilon(\widehat{x})) \\ \xi_{\widehat{x}_{-n}} \downarrow & & \downarrow \xi_{\widehat{x}} \\ \mathbb{D}^2(\rho_\varepsilon(\widehat{x}_{-n})) & \xleftarrow{R_{n, \widehat{x}}} & \mathbb{D}^2(\rho_\varepsilon(\widehat{x})) \end{array} \quad (4)$$

The mappings $R_{n,\hat{x}}$ have the following form depending on (λ_1, λ_2) :

(i) If $\lambda_1 = \lambda_2 = \lambda$, then $R_{n,\hat{x}}$ is a linear and satisfies

$$e^{-n(\lambda+\varepsilon)}\|(z, w)\| \leq \|R_{n,\hat{x}}(z, w)\| \leq e^{-n(\lambda-\varepsilon)}\|(z, w)\|.$$

(ii) If $\lambda_1 = k\lambda_2$ for some $k \geq 2$ (we say that λ_1 and λ_2 are resonant), then $R_{n,\hat{x}}(z, w) = (\alpha_{n,\hat{x}}z, \beta_{n,\hat{x}}w) + (\gamma_{n,\hat{x}}w^k, 0)$.

(iii) If $\lambda_1 \notin \{k\lambda_2, k \geq 1\}$, then $R_{n,\hat{x}}(z, w) = (\alpha_{n,\hat{x}}z, \beta_{n,\hat{x}}w)$.

Moreover, in the cases (ii) and (iii), we have

$$e^{-n(\lambda_1+\varepsilon)} \leq |\alpha_{n,\hat{x}}| \leq e^{-n(\lambda_1-\varepsilon)} \quad , \quad |\gamma_{n,\hat{x}}| \leq M_\varepsilon(\hat{x})e^{-n(\lambda_1-\varepsilon)}$$

and

$$e^{-n(\lambda_2+\varepsilon)} \leq |\beta_{n,\hat{x}}| \leq e^{-n(\lambda_2-\varepsilon)}.$$

Remark 2.3. The diagram (4) commutes for every fixed $n \geq 0$ by reducing the radius of the balls. The property $n \geq N(\hat{x})$ actually ensures that $f_{\hat{x}}^{-n}(B(x_0, \eta_\varepsilon(\hat{x}))) \subset B(x_{-n}, \eta_\varepsilon(\hat{x}_{-n}))$.

We shall need the following control of $(f^n)^*\omega_{\mathbb{P}^2}$ in Section 3.

Lemma 2.4. For every $\hat{x} \in \Lambda_\varepsilon$ and $n \geq 0$, there exists $0 < r_n \leq \eta_\varepsilon(\hat{x})$ such that

$$(f^n)^*\omega_{\mathbb{P}^2} \geq 4\beta_\varepsilon(\hat{x})^{-2}e^{2n(\lambda_2-2\varepsilon)}\omega_{\mathbb{P}^2} \text{ on } B(x_0, r_n).$$

Proof : We assume to simplify that $\lambda_1 > \lambda_2$ and that λ_1, λ_2 are not resonant, the case of equality and the resonant case can be treated similarly (up to modifying the function β_ε). Let $\varepsilon > 0$ be small enough such that $\lambda_1 - \lambda_2 > 2\varepsilon$, ensuring $|\alpha_{n,\hat{x}_n}|^{-1} \geq |\beta_{n,\hat{x}_n}|^{-1}$. According to Remark 2.3, there exists $r_n > 0$ and $0 < r'_n \leq \eta_\varepsilon(\hat{x}_n)$ such that $f_{\hat{x}_n}^{-n} = \xi_{\hat{x}}^{-1} \circ R_{n,\hat{x}_n} \circ \xi_{\hat{x}_n}$ on $B(x_n, r'_n)$ and $B(x_0, r_n) \subset f_{\hat{x}_n}^{-n}(B(x_n, r'_n))$. Since we have on $B(x_0, r_n)$:

$$(f^n)^*\omega_{\mathbb{P}^2} \geq \|(df^n)^{-1}\|^{-2}\omega_{\mathbb{P}^2}$$

and

$$\|(df^n)^{-1}\|^{-2} = \|df_{\hat{x}_n}^{-n}\|^{-2} \geq \|d\xi_{\hat{x}}^{-1}\|^{-2}\|dR_{n,\hat{x}_n}\|^{-2}\|d\xi_{\hat{x}_n}\|^{-2},$$

Theorem 2.2 implies

$$(f^n)^*\omega_{\mathbb{P}^2} \geq 4\|dR_{n,\hat{x}_n}\|^{-2}\beta_\varepsilon(\hat{x}_n)^{-2}\omega_{\mathbb{P}^2} \text{ on } B(x_0, r_n).$$

Since $|\alpha_{n,\hat{x}_n}|^{-1} \geq |\beta_{n,\hat{x}_n}|^{-1}$ and β_ε is ε -tempered, we deduce on $B(x_0, r_n)$:

$$(f^n)^*\omega_{\mathbb{P}^2} \geq 4|\beta_{n,\hat{x}_n}|^{-2}\beta_\varepsilon(\hat{x})^{-2}e^{-2n\varepsilon}\omega_{\mathbb{P}^2}.$$

The conclusion follows from $|\beta_{n,\hat{x}_n}|^{-2} \geq e^{2n(\lambda_2-\varepsilon)}$. □

We shall also need the following lemma in Section 5.

Lemma 2.5. Assume that $\mu = \psi_J\sigma_T$ for some measurable function ψ on \mathbb{P}^2 . Let $\hat{x} \in \Lambda_\varepsilon$ and $B := B(x_0, \eta_\varepsilon(\hat{x}))$. Let $A_n := \{p \in B \cap J, \psi \circ f_{\hat{x}}^{-n}(p) > 0\}$.

1. For every $n \geq 0$, the measures σ_T and $(f_{\hat{x}}^{-n})^*\sigma_T$ are equivalent on B , and $A_n = A_0$ modulo a borel subset of zero σ_T -measure.
2. If $A \subset B \cap J$ satisfies $\mu(A) = 0$, then $\sigma_T(A \cap A_0) = 0$. In particular, for $a, b \geq 0$,
 - (i) If $\psi \circ f_{\hat{x}}^{-n} \leq b$ μ -a.e. on $B \cap J$, then $\psi \circ f_{\hat{x}}^{-n} \leq b$ σ_T -a.e. on $B \cap J$.

(ii) If $\psi \circ f_{\hat{x}}^{-n} \geq a$ μ -a.e. on $B \cap J$, then $\psi \circ f_{\hat{x}}^{-n} \geq a$ σ_T -a.e. on $B \cap J \cap A_0$.

Proof : Since $f_{\hat{x}}^{-n} : B \rightarrow f_{\hat{x}}^{-n}(B)$ is a biholomorphism (and up to shrinking B), there exists $0 < \alpha < \beta$ such that

$$\alpha \omega_{\mathbb{P}^2} \leq (f_{\hat{x}}^{-n})^* \omega_{\mathbb{P}^2} \leq \beta \omega_{\mathbb{P}^2} \text{ on } B.$$

Using that T is f -invariant, we deduce that σ_T and $(f_{\hat{x}}^{-n})^* \sigma_T$ are equivalent on B . Now for every $p \in B \cap J$ and every r small enough such that $B(p, r) \subset B$,

$$\mu(B(p, r)) = \int_{B(p, r)} \psi_J \sigma_T.$$

But $\mu(B(p, r)) = d^{2n} \mu(f_{\hat{x}}^{-n}(B(p, r)))$ and

$$\mu(f_{\hat{x}}^{-n}(B(p, r))) = \int_{B(p, r)} \psi_J \circ f_{\hat{x}}^{-n} (f_{\hat{x}}^{-n})^* \sigma_T \simeq_{\alpha, \beta} d^{-n} \int_{B(p, r)} \psi_J \circ f_{\hat{x}}^{-n} \sigma_T,$$

where $\simeq_{\alpha, \beta}$ means that the equality holds up to a multiplicative constant between α and β . We deduce by dividing by $\sigma_T(B(p, r))$ (which is positive since $p \in J \subset \text{Supp } T$) :

$$\int_{B(p, r)} \psi_J \sigma_T \simeq_{\alpha, \beta} d^{-n} \int_{B(p, r)} \psi_J \circ f_{\hat{x}}^{-n} \sigma_T.$$

By Lebesgue Theorem [17, Theorem 2.12] (taking limits when r tends to zero), we get for σ_T -a.e. $p \in B \cap J$ that $\psi(p) > 0$ if and only if $\psi \circ f_{\hat{x}}^{-n}(p) > 0$. The second item comes from $0 = \mu(A) = \int_{A \cap A_0} \psi_J \sigma_T$, which gives $\sigma_T(A \cap A_0) = 0$ since ψ_J is positive on $A \cap A_0$. To obtain (i) and (ii), we apply this fact to $\{p \in B \cap J, \psi \circ f_{\hat{x}}^{-n}(p) > b\}$ and to $\{p \in B \cap J, \psi \circ f_{\hat{x}}^{-n}(p) < a\}$ (which both have zero μ -measure by assumption) and the fact that $A_n = A_0$ modulo a borel subset of zero σ_T -measure (provided by the first item). \square

3 A new proof of Dujardin's theorem

We provide another proof of Theorem 1.1. We shall use Theorem 2.2 and forward dynamics. We assume that $\mu \ll \sigma_T = T \wedge \omega_{\mathbb{P}^2}$. This implies that there exists $\psi \in L^1(\sigma_T)$ such that $\mu = \psi_J \sigma_T$. We recall that $\lambda_2 \geq \frac{1}{2} \text{Log } d$ for every endomorphism of \mathbb{P}^2 , it remains to prove the reverse inequality $\lambda_2 \leq \frac{1}{2} \text{Log } d$ in our case of absolute continuity. For every $n \geq 0$, the set

$$E_n := \left\{ x \in \mathbb{P}^2 : \lim_{\rho \rightarrow 0^+} \int_{B(x, \rho)} \psi_J \circ f^n \sigma_T = \psi_J \circ f^n(x) \right\}$$

is a borel subset of full σ_T -measure (and so of full μ -measure) by Lebesgue Theorem. We define the borel set

$$E := \left(\bigcap_{n \in \mathbb{N}} E_n \right) \setminus \left(\bigcup_{n \geq 1} \text{Crit}(f^n) \right),$$

which is also of full μ -measure (μ does not charge analytic subsets). Let us define

$$F := E \cap \left\{ \frac{1}{\tau} \leq \psi_J \leq \tau \right\} \subset J,$$

where $\tau > 0$ is chosen large enough to have $\mu(F) > 0$. According to Poincaré recurrence theorem, for $\hat{\mu}$ -almost every $\hat{x} \in \pi_0^{-1}(F)$, there exists a sub-sequence \hat{x}_{n_k} which satisfies

$\widehat{x}_{n_k} \in \pi_0^{-1}(F)$. Let us fix $\varepsilon > 0$ and apply the normal form Theorem 2.2, it yields a borel subset Λ_ε of full $\widehat{\mu}$ -measure of good orbits. Let us fix $\widehat{x} \in \pi_0^{-1}(F) \cap \Lambda_\varepsilon$ and a sub-sequence $(n_k)_k$ as before. Let us also fix an integer k . Since $x_0 \in E$ avoids the critical set of f^{n_k} , there exists $\rho_k > 0$ small enough such that f^{n_k} is injective on the ball $B(x_0, \rho_k)$. Then we have for any $\rho \in]0, \rho_k]$:

$$\mu(A_\rho) = d^{-2n_k} \mu(f^{n_k}(A_\rho)), \quad (5)$$

where $A_\rho := B(x_0, \rho)$. Up to taking a smaller ρ_k , we can assume that

$$f^{n_k}(A_\rho) \subset B(x_{n_k}, \eta_\varepsilon(\widehat{x}_{n_k})) \text{ and } A_\rho \subset B(x_0, r_{n_k}),$$

where r_{n_k} comes from Lemma 2.4. From $\mu = \psi_J \sigma_T$, Formula (5) yields for every $\rho \in]0, \rho_k]$:

$$\begin{aligned} \int_{A_\rho} \psi_J \sigma_T &= d^{-2n_k} \int_{f^{n_k}(A_\rho)} \psi_J \sigma_T \\ &= d^{-2n_k} \int_{A_\rho} (\psi_J \circ f^{n_k}) d^{n_k} T \wedge (f^{n_k})^* \omega_{\mathbb{P}^2} \end{aligned}$$

where the second equality uses the invariance of the Green current $(f^{n_k})^* T = d^{n_k} T$ and the fact that f^{n_k} is injective on A_ρ . Since $A_\rho \subset B(x_0, r_{n_k})$, we can use the estimate given by Lemma 2.4 with $n = n_k$ to deduce

$$\begin{aligned} \int_{A_\rho} \psi_J \sigma_T &\geq d^{-n_k} \int_{A_\rho} (\psi_J \circ f^{n_k}) 4\beta_\varepsilon(\widehat{x})^{-2} e^{2n_k(\lambda_2 - 2\varepsilon)} T \wedge \omega_{\mathbb{P}^2} \\ &= 4d^{-n_k} \beta_\varepsilon(\widehat{x})^{-2} e^{2n_k(\lambda_2 - 2\varepsilon)} \int_{A_\rho} (\psi_J \circ f^{n_k}) \sigma_T. \end{aligned}$$

We chose $x_0 \in J$, thus $\sigma_T(A_\rho) > 0$ and we have :

$$\int_{A_\rho} \psi_J \sigma_T \geq 4d^{-n_k} \beta_\varepsilon(\widehat{x})^{-2} e^{2n_k(\lambda_2 - 2\varepsilon)} \int_{A_\rho} (\psi_J \circ f^{n_k}) \sigma_T.$$

But we also have $x_0 \in E \subset E_0 \cap E_{n_k}$, thus we can take limits when $\rho \rightarrow 0^+$ to get

$$\psi_J(x_0) \geq 4d^{-n_k} \beta_\varepsilon(\widehat{x})^{-2} e^{2n_k(\lambda_2 - 2\varepsilon)} \times \psi_J(f^{n_k}(x_0)).$$

Since x_0 and $x_{n_k} = f^{n_k}(x_0)$ belongs to $F \subset \{\frac{1}{\tau} \leq \psi_J \leq \tau\}$, we deduce that :

$$\tau \geq \frac{4}{\tau} \beta_\varepsilon(\widehat{x})^{-2} e^{-2n_k(\frac{1}{2} \text{Log } d - \lambda_2 + 2\varepsilon)}.$$

Since this is true for every integer k , we get $\frac{1}{2} \text{Log } d - \lambda_2 + 2\varepsilon \geq 0$. We obtain $\lambda_2 \leq \frac{1}{2} \text{Log } d$ as desired, since $\varepsilon > 0$ is arbitrary small.

4 Decomposition of μ

4.1 Normal coordinates and absolute continuity

According to Theorem 2.2, there exists an invariant subset of full $\widehat{\mu}$ -measure Λ_ε such that for any $\widehat{x} \in \Lambda_\varepsilon$, there exist inverse branches

$$f_{\widehat{x}}^{-n} : B(x_0, \eta_\varepsilon(\widehat{x})) \longrightarrow B(x_{-n}, \eta_\varepsilon(\widehat{x}_{-n})), \quad n \geq N(\widehat{x}),$$

and holomorphic charts

$$\xi_{\hat{x}} = (Z_{\hat{x}}, W_{\hat{x}}) : B(x_0, \eta_\varepsilon(\hat{x})) \longrightarrow \mathbb{D}^2(\rho_\varepsilon(\hat{x})) \quad (6)$$

such that $\frac{1}{2} \leq \|d\xi_{\hat{x}}\| \leq \beta_\varepsilon(\hat{x})$ on $B(x_0, \eta_\varepsilon(\hat{x}))$ and

$$f_{\hat{x}}^{-n} = \xi_{\hat{x}_{-n}}^{-1} \circ R_{n, \hat{x}} \circ \xi_{\hat{x}}$$

for any $n \geq N(\hat{x})$, where $R_{n, \hat{x}}$ is a polynomial map $(\mathbb{C}^2, (Z_{\hat{x}}, W_{\hat{x}})) \rightarrow (\mathbb{C}^2, (Z_{\hat{x}_{-n}}, W_{\hat{x}_{-n}}))$. When $\lambda_1 > \lambda_2$ are not resonant, $R_{n, \hat{x}}$ is a linear diagonal map with eigenvalues $|\alpha_{n, \hat{x}}| \simeq e^{-n(\lambda_1 \pm \varepsilon)}$ and $|\beta_{n, \hat{x}}| \simeq e^{-n(\lambda_2 \pm \varepsilon)}$. Let $(\omega_{\hat{x}, AB})_{A, B \in \{Z, W\}}$ be the smooth functions defined by the formula

$$\omega_{\mathbb{P}^2} = \sum_{A, B \in \{Z, W\}} \omega_{\hat{x}, AB} \frac{i}{2} (dA_{\hat{x}} \wedge d\bar{B}_{\hat{x}}) \quad \text{on } B(x_0, \eta_\varepsilon(\hat{x})). \quad (7)$$

For every A, B in $\{Z, W\}$ and for every $n \geq N(\hat{x})$ we define the following functions :

$$F_{\hat{x}, n}^{AB} := \omega_{\hat{x}_{-n}, AB} \circ f_{\hat{x}}^{-n} : B(x_0, \eta_\varepsilon(\hat{x})) \longrightarrow \mathbb{C}.$$

Proposition 4.1. *Assume that $\mu \ll \sigma_T$ and let ψ_J be the Radon-Nikodym derivative of μ with respect to σ_T . We assume that the Lyapunov exponents $\lambda_1 > \lambda_2$ are not resonant. For every $\hat{x} \in \Lambda_\varepsilon$, for every borel set $U \subset B(x_0, \eta_\varepsilon(\hat{x}))$ and for every $n \geq N(\hat{x})$ we have :*

$$\mu(U) = d^n \int_U (\psi_J \circ f_{\hat{x}}^{-n}) \sum_{A, B \in \{Z, W\}} F_{\hat{x}, n}^{AB} \alpha_{n, \hat{x}}^A \overline{\alpha_{n, \hat{x}}^B} T \wedge \frac{i}{2} dA_{\hat{x}} \wedge d\bar{B}_{\hat{x}},$$

where $\alpha_{n, \hat{x}}^Z := \alpha_{n, \hat{x}}$ and $\alpha_{n, \hat{x}}^W := \beta_{n, \hat{x}}$.

Proof : From the equalities $\mu(U) = d^{2n} \mu(f_{\hat{x}}^{-n}(U))$ and $\mu = \psi_J(T \wedge \omega_{\mathbb{P}^2})$, we obtain

$$\mu(U) = d^{2n} \int_U (\psi_J \circ f_{\hat{x}}^{-n}) (f_{\hat{x}}^{-n})^* T \wedge (f_{\hat{x}}^{-n})^* \omega_{\mathbb{P}^2}.$$

By invariance of the Green current, we get $(f_{\hat{x}}^{-n})^* T = d^{-n} T$ on $B(x_0, \eta_\varepsilon(\hat{x}))$ and

$$\mu(U) = d^n \int_U (\psi_J \circ f_{\hat{x}}^{-n}) T \wedge (f_{\hat{x}}^{-n})^* \omega_{\mathbb{P}^2}.$$

We conclude by pulling back Formula (7) (\hat{x} being replaced by \hat{x}_{-n}) by $f_{\hat{x}}^{-n}$ and using the fact that $R_{n, \hat{x}}$ is linear and diagonal with eigenvalues $\alpha_{n, \hat{x}}$ and $\beta_{n, \hat{x}}$. \square

Definition 4.2. *Assume that $\mu \ll \sigma_T$ and let $\psi_J = \psi \mathbb{1}_J$ be the Radon-Nikodym derivative of μ with respect to σ_T . We assume that the Lyapunov exponents $\lambda_1 > \lambda_2$ are not resonant. For every $U \subset B(x_0, \eta_\varepsilon(\hat{x}))$ and $n \geq 0$, we have according to Proposition 4.1 :*

$$\mu(U) = I_n(U \cap J) + J_n(U \cap J),$$

where for every $D \subset B(x_0, \eta_\varepsilon(\hat{x}))$:

$$I_n(D) := d^n |\beta_{n, \hat{x}}|^2 \int_D (\psi \circ f_{\hat{x}}^{-n}) F_{\hat{x}, n}^{WW} T \wedge \frac{i}{2} (dW_{\hat{x}} \wedge d\bar{W}_{\hat{x}})$$

and

$$J_n(D) := d^n \sum_{(A, B) \neq (W, W)} \alpha_{n, \hat{x}}^A \overline{\alpha_{n, \hat{x}}^B} \int_D (\psi \circ f_{\hat{x}}^{-n}) F_{\hat{x}, n}^{AB} T \wedge \frac{i}{2} (dA_{\hat{x}} \wedge d\bar{B}_{\hat{x}})$$

where $\alpha_{n, \hat{x}}^Z := \alpha_{n, \hat{x}}$ and $\alpha_{n, \hat{x}}^W := \beta_{n, \hat{x}}$.

Remark 4.3. *We observe that if $\mu = \psi_V \sigma_T$ as in the second item of Theorem 1.2, then*

$$\mu(U) = I_n(U \cap V) + J_n(U \cap V)$$

for every $U \subset B(x_0, \eta_\varepsilon(\hat{x}))$.

4.2 Recurrence in \mathcal{A}_τ

Definition 4.4. For $\tau > 0$ let $\mathcal{A}_\tau \subset \Lambda_\varepsilon$ be the subset consisting of elements \hat{y} satisfying

1. $|\omega_{\hat{y}, AB}| \leq \tau$,
2. $\frac{1}{\tau} \leq \omega_{\hat{y}, AA} \leq \tau$,
3. $S_{AB}(\hat{y}) := \sup_{B(y_0, \eta_\varepsilon(\hat{y}))} \|d\omega_{\hat{y}, AB}\| \leq \tau$,
4. $\psi_J \circ \pi_0(\hat{y}) \geq \frac{1}{\tau}$,

the two first items holding on $B(y_0, \eta_\varepsilon(\hat{y}))$, the three first ones for every $A, B \in \{Z, W\}$.

Since the distortion of the coordinates $\xi_{\hat{x}} = (Z_{\hat{x}}, W_{\hat{x}})$ is controlled by Theorem 2.2, there exists τ large enough such that \mathcal{A}_τ has positive $\hat{\mu}$ -measure. By Poincaré recurrence theorem, for $\hat{\mu}$ -almost every \hat{x} , there exists $(n_k)_{k \in \mathbb{N}}$ such that $n_k \geq N(\hat{x})$ and

$$\hat{x}_{-n_k} \in \mathcal{A}_\tau.$$

We have specified $n_k \geq N(\hat{x})$ to have $f_{\hat{x}}^{-n_k}(B(x_0, \eta_\varepsilon(\hat{x}))) \subset B(x_{-n_k}, \eta_\varepsilon(\hat{x}_{-n_k}))$. In this way the functions $F_{\hat{x}, n_k}^{AB}$ are well defined on $B(x_0, \eta_\varepsilon(\hat{x}))$. By definition of $F_{\hat{x}, n}^{AB}$, we get from the first item of Definition 4.4 :

$$|F_{\hat{x}, n_k}^{AB}| \leq \tau \text{ on } B(x_0, \eta_\varepsilon(\hat{x})). \quad (8)$$

Moreover, using the third item of Definition 4.4 and $\text{Lip}(f_{\hat{x}}^{-n_k}) \leq \beta_\varepsilon(\hat{x})e^{-n_k(\lambda_2 - \varepsilon)}$ (given by Theorem 2.2), we get

$$\|dF_{\hat{x}, n_k}^{AB}\| \leq \tau \beta_\varepsilon(\hat{x})e^{-n_k(\lambda_2 - \varepsilon)} \text{ on } B(x_0, \eta_\varepsilon(\hat{x})).$$

We deduce that

$$F_{\hat{x}, n_k}^{AB} \text{ uniformly converges on } B(x_0, \eta_\varepsilon(\hat{x})) \text{ to a constant } F_{\hat{x}}^{AB} \text{ up to a subsequence} \quad (9)$$

and that (according to the second item of Definition 4.4) :

$$F_{\hat{x}, n_k}^{AA} \geq \frac{1}{\tau} \text{ on } B(x_0, \eta_\varepsilon(\hat{x})). \quad (10)$$

4.3 A weaker version of Dujardin's theorem

We give another proof of Theorem 1.1 assuming that the Radon-Nikodym derivative $\psi_J \in L^\infty(\sigma_T)$. Our motivation is to use backward iterates, namely the inverse branches $f_{\hat{x}}^{-n_k}$, in order to prepare the proof of Theorem 1.2 in Section 6.1. As in Section 3, it suffices to prove that $\lambda_2 \leq \frac{1}{2} \text{Log } d$.

For sake of simplicity, we assume that the Lyapunov exponents $\lambda_1 > \lambda_2$ are not resonant. This implies that the polynomial map $R_{n, \hat{x}}$ appearing in the diagram (4) is linear and diagonal. The proof can be adapted in the resonant case, see Section 7.

We apply the normal form Theorem 2.2 with $\varepsilon > 0$, it yields a borel subset Λ_ε of full $\hat{\mu}$ -measure of good orbits. We use the set \mathcal{A}_τ defined in Section 4.2. Let us fix a $\hat{\mu}$ -generic element $\hat{x} \in \Lambda_\varepsilon$, in particular $x_0 \in J$. Let $(n_k)_{k \in \mathbb{N}}$ be such that $\hat{x}_{-n_k} \in \mathcal{A}_\tau$.

1) Construction of a set of continuity C .

Let $\delta > 0$ be such that $\mu(B(x_0, \eta_\varepsilon(\hat{x}))) > 2\delta$. By Lusin's theorem there exists a compact set $C_0 \subset B(x_0, \eta_\varepsilon(\hat{x}))$ such that $\mu(B(x_0, \eta_\varepsilon(\hat{x})) \setminus C_0) < \delta$ and ψ_J is continuous on C_0 . Using

again Lusin's theorem we construct by induction a sequence $(C_n)_n$ of compact subsets of $B(x_0, \eta_\varepsilon(\hat{x}))$ such that

$$C_{n+1} \subset C_n \text{ and } \mu(C_n \setminus C_{n+1}) < 2^{-(n+1)}\delta$$

and

$$\psi_J \circ f_{\hat{x}}^{-(n+1)} \in C^0(C_{n+1}).$$

Then $\mu(C_n) > \mu(C_{n-1}) - 2^{-n}\delta > \dots > \mu(C_0) - \delta > \mu(B(x_0, \eta_\varepsilon(\hat{x}))) - 2\delta > 0$. Thus

$$C := \bigcap_{n \in \mathbb{N}} C_n \subset B(x_0, \eta_\varepsilon(\hat{x}))$$

satisfies $\mu(C) > 0$. Let μ_C be the measure on $B(x_0, \eta_\varepsilon(\hat{x}))$ defined by

$$\mu_C(\cdot) := \mu(\cdot \cap C).$$

Its support $\text{supp } \mu_C$ is included in $C \cap J$ because C is a closed subset of $B(x_0, \eta_\varepsilon(\hat{x}))$.

2) Construction of x'_0 in $\text{supp } \mu_C$, with boundedness properties.

Let G be a borel subset of full σ_T -measure (and so of full μ -measure) satisfying

$$\forall p \in G, 0 \leq \psi_J(p) \leq \|\psi_J\|_{L^\infty(\sigma_T)},$$

and let G_k be the borel set

$$G_k := f_{\hat{x}}^{-n_k}(B(x_0, \eta_\varepsilon(\hat{x}))) \cap G.$$

Let us verify that $\mu(f^{n_k}(G_k)) = \mu(B(x_0, \eta_\varepsilon(\hat{x})))$. Since $\mu(G) = 1$, we have $\mu(G_k) = \mu(f_{\hat{x}}^{-n_k}(B(x_0, \eta_\varepsilon(\hat{x})))) = d^{-2n_k} \mu(B(x_0, \eta_\varepsilon(\hat{x})))$. But f^{n_k} is injective on $f_{\hat{x}}^{-n_k}(B(x_0, \eta_\varepsilon(\hat{x})))$, which contains G_k , hence $\mu(f^{n_k}(G_k)) = d^{2n_k} \mu(G_k)$, completing our verification.

Therefore

$$\mu\left(\bigcap_{k \in \mathbb{N}} f^{n_k}(G_k)\right) = \mu(B(x_0, \eta_\varepsilon(\hat{x}))),$$

which is positive since $x_0 \in J$. Let

$$x'_0 \in \bigcap_{k \in \mathbb{N}} f^{n_k}(G_k) \cap \left\{ \frac{1}{\tau'} \leq \psi_J \right\} \cap \text{supp } \mu_C, \quad (11)$$

where $\tau' > 1$ is large enough so that the set defined in (11) is not empty. Using the fact that $f_{\hat{x}}^{-n_k} \circ (f^{n_k})|_{G_k} = \text{Id}_{G_k}$, we observe that

$$\forall k \in \mathbb{N}, f_{\hat{x}}^{-n_k}(x'_0) \in G_k \subset G. \quad (12)$$

3) Conclusion.

For a given k , let $\rho_k > 0$ be small enough so that for any $0 < \rho < \rho_k$:

$$A_\rho := (B(x'_0, \rho) \cap C) \subset B(x_0, \eta_\varepsilon(\hat{x})).$$

Since $\|d\xi_{\hat{x}}\| \leq \beta_\varepsilon(\hat{x})$ we can assume that $\sum_{A,B} |T \wedge \frac{i}{2} dA_{\hat{x}} \wedge d\bar{B}_{\hat{x}}| \leq \beta_\varepsilon(\hat{x})^2 \sigma_T$ on $B(x_0, \eta_\varepsilon(\hat{x}))$. According to Proposition 4.1, $|F_{\hat{x}, n_k}^{AB}| \leq \tau$ (see Equation (8)) and $|\alpha_{n, \hat{x}}| \leq |\beta_{n, \hat{x}}|$, we deduce

$$\int_{A_\rho} \psi_J d\sigma_T = \mu(A_\rho) \leq d^{n_k} |\beta_{n_k, \hat{x}}|^2 \tau \beta_\varepsilon(\hat{x})^2 \int_{A_\rho} (\psi_J \circ f_{\hat{x}}^{-n_k}) d\sigma_T.$$

According to Equation (11), we have $x'_0 \in \text{supp } \mu_C$, thus $\sigma_T(A_\rho) > 0$ and

$$\int_{A_\rho} \psi_J \, d\sigma_T \leq d^{n_k} |\beta_{n_k, \hat{x}}|^2 \tau \beta_\varepsilon(\hat{x})^2 \int_{A_\rho} (\psi_J \circ f_{\hat{x}}^{-n_k}) \, d\sigma_T.$$

Equation (11) also gives $\frac{1}{\tau'} \leq \psi_J(x'_0)$ and $x'_0 \in \text{supp } \mu_C \subset C \subset C_0 \cap C_{n_k}$. We moreover have $A_\rho \subset C$ by definition. Hence, if ρ tends to 0, we get by continuity of ψ_J (resp. $\psi_J \circ f_{\hat{x}}^{-n_k}$) on C_0 (resp. on C_{n_k}):

$$\frac{1}{\tau'} \leq \psi_J(x'_0) \leq d^{n_k} |\beta_{n_k, \hat{x}}|^2 \tau \beta_\varepsilon(\hat{x})^2 (\psi_J \circ f_{\hat{x}}^{-n_k})(x'_0) \leq d^{n_k} |\beta_{n_k, \hat{x}}|^2 \tau \beta_\varepsilon(\hat{x})^2 \|\psi_J\|_{L^\infty(\sigma_T)},$$

the last inequality coming from the choice of x'_0 , see Equation (12). We obtain for every k :

$$(\tau' \tau \beta_\varepsilon(\hat{x})^2 \|\psi_J\|_{L^\infty(\sigma_T)})^{-1} \leq d^{n_k} |\beta_{n_k, \hat{x}}|^2 \leq e^{2n_k(\frac{1}{2} \text{Log } d - \lambda_2 + \varepsilon)}.$$

Hence $\lambda_2 \leq \frac{1}{2} \text{Log } d + \varepsilon$, which gives $\lambda_2 \leq \frac{1}{2} \text{Log } d$ when ε tends to zero, as desired. \square

5 Estimates on I_n and J_n in the decomposition of μ

5.1 Study of I_n and J_n

We recall that I_n and J_n were introduced in Definition 4.2.

Proposition 5.1. *Assume that $\mu \ll \sigma_T$ and let $\psi \mathbb{1}_J$ be the Radon-Nikodym derivative of μ with respect to σ_T . Assume that $\lambda_1 > \lambda_2$ are not resonant and let $\varepsilon > 0$ be small enough such that*

$$\lambda_1 + \lambda_2 - 2\varepsilon > \text{Log } d \quad , \quad \lambda_1 > \frac{1}{2} \text{Log } d + \varepsilon.$$

Let \hat{x} be a $\hat{\mu}$ -generic orbit and let $(n_k)_k$ be a sub-sequence such that $\hat{x}_{-n_k} \in \mathcal{A}_T$. Let

$$\sigma_{\hat{x}} := T \wedge dd^c |W_{\hat{x}}|^2.$$

Let $B := B(x_0, \eta_\varepsilon(\hat{x}))$, $A_0 := \{p \in B \cap J, \psi(p) > 0\}$ and $U \subset B$ be a borel set.

1. If $\psi \circ f_{\hat{x}}^{-n_k} \leq b \mu - a.e.$ on B for every $k \geq 0$, then there exists $C_{\hat{x}} > 0$ which does not depend on $U \subset B$ such that

$$J_{n_k}(U \cap J) \leq C_{\hat{x}} e^{-n_k(\lambda_1 + \lambda_2 - \text{Log } d)} e^{2n_k \varepsilon} \xrightarrow[k \rightarrow \infty]{} 0.$$

2. If $\psi \circ f_{\hat{x}}^{-n_k} \leq b \mu - a.e.$ on B for every $k \geq 0$ and if $\mu(U) > 0$, then

$$\sigma_{\hat{x}}(U \cap J \cap A_0) > 0.$$

3. If $0 < a \leq \psi \circ f_{\hat{x}}^{-n_k} \leq b \mu - a.e.$ on B for every $k \geq 0$, then

$$d^{n_k} |\beta_{n_k, \hat{x}}|^2 a \tau^{-1} \sigma_{\hat{x}}(U \cap J \cap A_0) \leq I_{n_k}(U \cap J) \leq d^{n_k} |\beta_{n_k, \hat{x}}|^2 b \tau \sigma_{\hat{x}}(U \cap J).$$

In particular $(d^{n_k} |\beta_{n_k, \hat{x}}|^2)_k$ converges, up to a sub-sequence, to some $u > 0$.

Remark 5.2. *We assume in Proposition 5.1 that λ_1, λ_2 are not resonant. The same statement actually holds in the resonant case, this is explained in Section 7.*

Proof of Proposition 5.1 : Let us prove the first item. By Definition 4.2,

$$J_{n_k}(U \cap J) = \sum_{(A,B) \neq (W,W)} d^{n_k} \alpha_{n_k, \hat{x}}^A \overline{\alpha_{n_k, \hat{x}}^B} \int_{U \cap J} (\psi \circ f_{\hat{x}}^{-n_k}) F_{\hat{x}, n_k}^{AB} T \wedge dA_{\hat{x}} \wedge d\overline{B}_{\hat{x}}.$$

Let us introduce the measures (which are $\ll \sigma_T$) :

$$\sigma_{\hat{x}} := T \wedge dd^c |W_{\hat{x}}|^2, \quad \lambda_{\hat{x}} := T \wedge dd^c |Z_{\hat{x}}|^2.$$

Using $|\alpha_{n, \hat{x}}| \leq |\beta_{n, \hat{x}}|$, the (assumed) upper bound on $\psi \circ f_{\hat{x}}^{-n_k}$, the upper bound on $F_{\hat{x}, n_k}^{AB}$ provided by (8), the item 2.(i) of Lemma 2.5 and Cauchy-Schwarz inequality, we get

$$|J_{n_k}(U \cap J)| \leq d^{n_k} |\alpha_{n_k, \hat{x}} \beta_{n_k, \hat{x}}| b \tau (\sigma_{\hat{x}} + \lambda_{\hat{x}})(U \cap J).$$

The first item then follows from $|\alpha_{n, \hat{x}}| \leq e^{-n(\lambda_1 - \varepsilon)}$ and $|\beta_{n, \hat{x}}| \leq e^{-n(\lambda_2 - \varepsilon)}$.

We show the second item. By using Item 1 of Lemma 2.5 (for the last equality), we get

$$\mu(U) = \mu(U \cap J) = d^{2n} \mu(f_{\hat{x}}^{-n}(U \cap J)) = d^{2n} \int_{U \cap J \cap A_0} (\psi \circ f_{\hat{x}}^{-n}) (f_{\hat{x}}^{-n})^*(T \wedge \omega_{\mathbb{P}^2}).$$

Cauchy-Schwarz inequality implies (up to multiplication by a constant) :

$$T \wedge \omega_{\mathbb{P}^2} \leq \omega_{\hat{x}_{-n}, WW} \sigma_{\hat{x}_{-n}} + \omega_{\hat{x}_{-n}, ZZ} \lambda_{\hat{x}_{-n}} \text{ on } B(\hat{x}_{-n}, \eta_\varepsilon(\hat{x}_{-n})).$$

We now fix $n = n_k$. Using the (assumed) upper bound on $\psi \circ f_{\hat{x}}^{-n_k}$ on B combined with item 2.(i) of Lemma 2.5, and the upper bounds on $F_{\hat{x}, n_k}^{WW} = \omega_{\hat{x}_{-n}, WW} \circ f_{\hat{x}}^{-n_k}$, $F_{\hat{x}, n_k}^{ZZ} = \omega_{\hat{x}_{-n}, ZZ} \circ f_{\hat{x}}^{-n_k}$ provided by Equation (8), we get

$$\mu(U) \leq b d^{2n_k} \tau \left(\int_{U \cap J \cap A_0} (f_{\hat{x}}^{-n_k})^*(\sigma_{\hat{x}_{-n_k}}) + \int_{U \cap J \cap A_0} (f_{\hat{x}}^{-n_k})^*(\lambda_{\hat{x}_{-n_k}}) \right).$$

We deduce

$$\mu(U) \leq b d^{2n_k} \tau \left(\int_{U \cap J \cap A_0} d^{-n_k} |\beta_{n_k, \hat{x}}|^2 \sigma_{\hat{x}} + \int_{U \cap J \cap A_0} d^{-n_k} |\alpha_{n_k, \hat{x}}|^2 \lambda_{\hat{x}} \right),$$

hence

$$0 < \mu(U) \leq b d^{n_k} \tau (|\beta_{n_k, \hat{x}}|^2 \sigma_{\hat{x}}(U \cap J \cap A_0) + |\alpha_{n_k, \hat{x}}|^2 \lambda_{\hat{x}}(U \cap J \cap A_0)).$$

To conclude we observe that

$$d^{n_k} |\alpha_{n_k, \hat{x}}|^2 \leq e^{n_k(\text{Log } d - 2\lambda_1 + 2\varepsilon)},$$

which tends to 0 when k tends to infinity. Hence $\sigma_{\hat{x}}(U \cap J \cap A_0) > 0$.

Let us prove the third item. According to item 2.(ii) of Lemma 2.5, we have

$$\psi \circ f_{\hat{x}}^{-n_k} \geq a, \quad \sigma_{\hat{x}} - a.e. \text{ on } B \cap J \cap A_0.$$

Hence, according to Definition 4.2, we get

$$I_{n_k}(U \cap J \cap A_0) \geq d^{n_k} |\beta_{n_k, \hat{x}}|^2 a \int_{U \cap J \cap A_0} F_{\hat{x}, n_k}^{WW} \sigma_{\hat{x}} \geq d^{n_k} |\beta_{n_k, \hat{x}}|^2 a \frac{1}{\tau} \sigma_{\hat{x}}(U \cap J \cap A_0),$$

where the last inequality comes from Equation (10). Similarly, by using item 2.(i) of Lemma 2.5 and Equation (8), we get

$$I_{n_k}(U \cap J) \leq d^{n_k} |\beta_{n_k, \hat{x}}|^2 b \tau \sigma_{\hat{x}}(U \cap J)$$

This proves the two stated inequalities on $I_{n_k}(U \cap J)$. Let us verify the last fact. Since $0 < \mu(B) = I_{n_k}(B \cap J) + J_{n_k}(B \cap J)$ and since $J_{n_k}(B)$ tends to 0 by the first item, we get $\frac{1}{2}\mu(B) \leq I_{n_k}(B \cap J) \leq 2\mu(B)$ for k large enough, which implies

$$0 < \frac{1}{2b\tau} \frac{\mu(B)}{\sigma_{\hat{x}}(B \cap J)} \leq d^{n_k} |\beta_{n_k, \hat{x}}|^2 \leq \frac{2\tau}{a} \frac{\mu(B)}{\sigma_{\hat{x}}(B \cap J \cap A_0)} < +\infty,$$

since $\sigma_{\hat{x}}(B \cap J \cap A_0) > 0$ by the second item. \square

5.2 Application : an equivalence between μ and a slice of T

We prove that μ can be approximated by a slice of T under absolute continuity assumptions. The proof is a direct application of the decomposition of μ fixed in Definition 4.2 and of Proposition 5.1.

Proposition 5.3. *Assume that $\mu \ll \sigma_T$ and $\lambda_1 > \lambda_2 = \frac{1}{2} \text{Log } d$. Let $\psi \mathbb{1}_J$ be the Radon-Nikodym derivative of μ with respect to σ_T . Let \hat{x} be a $\hat{\mu}$ -generic element and let $(n_k)_k$ be a sub-sequence such that $\hat{x}_{-n_k} \in \mathcal{A}_\tau$. Let $B := B(x_0, \eta_\varepsilon(\hat{x}))$ and $A_0 := \{p \in B \cap J, \psi(p) > 0\}$. If*

$$0 < a \leq \psi \circ f_{\hat{x}}^{-n_k} \leq b \mu - \text{a.e. on } B \text{ for every } k \geq 0,$$

then there exists $C_{\hat{x}} > 0$ such that on B :

$$C_{\hat{x}}^{-1}(T \wedge dd^c |W_{\hat{x}}|^2)|_{J \cap A_0} \leq \mu \leq C_{\hat{x}}(T \wedge dd^c |W_{\hat{x}}|^2)|_J.$$

Proof : Let $U \subset B$ be a borel set and let $I_k := I_{n_k}(U \cap J)$, $J_k := J_{n_k}(U \cap J)$. According to Definition 4.2, we have $\mu(U) = I_k + J_k$. The first item of Proposition 5.1 yields $J_k \rightarrow 0$. Using the third item of Proposition 5.1 (which in particular provides $\lim_k d^{n_k} |\beta_{n_k, \hat{x}}|^2 = u > 0$), Equation (8) and item 2.(i) of Lemma 2.5, we get

$$\mu(U) \leq u b \tau \sigma_{\hat{x}}(U \cap J).$$

Similarly, by using Equation (10) and item 2.(ii) of Lemma 2.5, we get

$$\mu(U) \geq \frac{u a}{\tau} \sigma_{\hat{x}}(U \cap J \cap A_0).$$

We conclude by setting $C_{\hat{x}} := \max\{\tau/ua, ub\tau\}$. \square

6 Proof of Theorem 1.2

6.1 Proof of the first item

Lemma 6.1. *Assume that $\mu \ll \sigma_T$ and $\lambda_1 > \lambda_2 = \frac{1}{2} \text{Log } d$. Let $\psi \mathbb{1}_J$ be the Radon-Nikodym derivative of μ with respect to σ_T , and assume that $\psi|_J$ is continuous. Let \hat{x} be a $\hat{\mu}$ -generic element and let $(n_k)_k$ be a sub-sequence such that $\hat{x}_{-n_k} \in \mathcal{A}_\tau$ with $n_k \geq N(\hat{x})$. There exists $r_1(\hat{x}) \leq \eta_\varepsilon(\hat{x})$ and $x_{-\infty} \in J$ such that, up to a sub-sequence :*

1. $\lim_k d^{n_k} |\beta_{n_k, \hat{x}}|^2 = u > 0$ and $\lim_k J_{n_k}(U \cap J) = 0$ for any borel set $U \subset B(x_0, r_1(\hat{x}))$,
2. $\psi \circ f_{\hat{x}}^{-n_k}$ uniformly converges to $\psi(x_{-\infty})$ on $B(x_0, r_1(\hat{x})) \cap J$.

Proof : To prove the first item, in view of Proposition 5.1, it suffices to verify

$$\frac{1}{2\tau} \leq \psi \circ f_{\hat{x}}^{-n_k}$$

on some $B(x_0, r_1(\hat{x})) \cap J$. By uniform continuity, there exists $\delta(\tau) > 0$ such that

$$\forall a, b \in J, \text{dist}(a, b) \leq \delta(\tau) \implies |\psi(a) - \psi(b)| \leq \frac{1}{2\tau}.$$

Defining $r_1(\hat{x}) := \min \left\{ \eta_\varepsilon(\hat{x}), \frac{\delta(\tau)}{\beta_\varepsilon(\hat{x})} \right\}$, we have for every $p \in B(x_0, r_1(\hat{x}))$:

$$\text{dist}(f_{\hat{x}}^{-n_k}(p), f_{\hat{x}}^{-n_k}(x_0)) \leq \text{Lip}(f_{\hat{x}}^{-n_k})r_1(\hat{x}) \leq \beta_\varepsilon(\hat{x})r_1(\hat{x}) \leq \delta(\tau).$$

Using $\psi(x_{-n_k}) \geq \frac{1}{\tau}$ (coming from $\hat{x}_{-n_k} \in \mathcal{A}_\tau$), we get $\psi \circ f_{\hat{x}}^{-n_k} \geq \frac{1}{2\tau}$ on $B(x_0, r_1(\hat{x})) \cap J$.

For the second item, the sequence $(x_{-n_k})_k \in J$ converges by compactness to some $x_{-\infty} \in J$. Since $\text{Lip} f_{\hat{x}}^{-n_k}$ tends to 0, $(f_{\hat{x}}^{-n_k})_k$ converges uniformly on $B(x_0, r_1(\hat{x}))$ to the constant mapping $x_{-\infty}$. The conclusion follows from the continuity of ψ on J . \square

We now prove the first item of Theorem 1.2. Let $U \subset B(x_0, r_1(\hat{x}))$ be a borel set. We recall that $\sigma_{\hat{x}} = T \wedge dd^c |W_{\hat{x}}|^2$ and that by Definition 4.2 :

$$\mu(U) = \mu(U \cap J) = d^{n_k} |\beta_{n_k, \hat{x}}|^2 \int_{U \cap J} (\psi \circ f_{\hat{x}}^{-n_k}) F_{\hat{x}, n_k}^{WW} \sigma_{\hat{x}} + J_{n_k}(U \cap J).$$

According to Equation (9), $F_{\hat{x}, n_k}^{WW}$ uniformly converges on $B(x_0, r_1(\hat{x}))$ to some constant $F_{\hat{x}}$, up to a sub-sequence. We conclude by using the two items of Lemma 6.1, which yield $\mu(U) = C_{\hat{x}} \sigma_{\hat{x}}(U \cap J)$, where $C_{\hat{x}} := u F_{\hat{x}} \psi(x_{-\infty})$. Note that $C_{\hat{x}}$ does not depend on U and is positive by taking $U = B(x_0, r_1(\hat{x}))$.

6.2 Proof of the second item

The following lemma is the counterpart of Lemma 6.1, the function ψ being now continuous on a neighborhood V of J . The two items hold without any restriction to J . The proof is similar using Remark 4.3 (decomposition of μ on V), the uniform continuity of ψ on V , and by introducing a radius $r_2(\hat{x})$ satisfying $B(x_0, r_2(\hat{x})) \subset V$. Note that $f_{\hat{x}}^{-n_k}(B(x_0, r_2(\hat{x}))) \subset V$ for k large enough.

Lemma 6.2. *Assume that $\mu \ll \sigma_T$ and $\lambda_1 > \lambda_2 = \frac{1}{2} \text{Log } d$. We assume that there exists a neighborhood V of J and a function $\psi \in L^1(\sigma_T)$ such that ψ is continuous on V and satisfies $\mu = \psi_V \sigma_T$ (we recall that $\psi_V = \psi \mathbb{1}_V$). Let \hat{x} be a $\hat{\mu}$ -generic element and let $(n_k)_k$ be a sub-sequence such that $\hat{x}_{-n_k} \in \mathcal{A}_\tau$ with $n \geq N(\hat{x})$. There exists $r_2(\hat{x}) \leq \eta_\varepsilon(\hat{x})$ (small enough to have $B(x_0, r_2(\hat{x})) \subset V$) and $x_{-\infty} \in J$ such that, up to a sub-sequence :*

1. $\lim_k d^{n_k} |\beta_{n_k, \hat{x}}|^2 = u > 0$ and $\lim_k J_{n_k}(U) = 0$ for any borel set $U \subset B(x_0, r_2(\hat{x}))$,
2. $\psi \circ f_{\hat{x}}^{-n_k}$ uniformly converges to $\psi(x_{-\infty})$ on $B(x_0, r_2(\hat{x}))$.

The second item of Theorem 1.2 follows as in the end of Section 6.1, by using

$$\mu(U) = d^{n_k} |\beta_{n_k, \hat{x}}|^2 \int_U (\psi \circ f_{\hat{x}}^{-n_k}) F_{\hat{x}, n_k}^{WW} \sigma_{\hat{x}} + J_{n_k}(U)$$

for every $U \subset B(x_0, r_2(\hat{x})) \subset V$, this formula is provided by Definition 4.2 and Remark 4.3.

7 About the resonant case

We assume that λ_1, λ_2 are resonant : $\lambda_1 = q\lambda_2$ for some $q \geq 2$. We assume that $\mu \ll \sigma_T$, let ψ denote the Radon-Nikodym derivative of μ with respect to σ_T . Let us first explain how Proposition 4.1 is affected. We shall compute the corresponding $J_n(D)$ appearing in Definition 4.2. According to Theorem 2.2, the map $R_{n,\hat{x}}$ is equal to :

$$R_{n,\hat{x}}(z, w) = (\alpha_{n,\hat{x}}z, \beta_{n,\hat{x}}w) + (\gamma_{n,\hat{x}}w^q, 0),$$

where $|\gamma_{n,\hat{x}}| \leq M_\varepsilon(\hat{x})e^{-n(\lambda_1-\varepsilon)}$. Let $U \subset B(x_0, \eta_\varepsilon(\hat{x}))$ be a borel set. Denoting $\psi_n := \psi \circ f_{\hat{x}}^{-n}$, we have $\mu(U) = I_n(U \cap J) + J_n(U \cap J)$, where $I_n(D)$ has the same expression as in the non resonant case :

$$I_n(D) = d^n |\beta_{n,\hat{x}}|^2 \int_D \psi_n F_{\hat{x},n}^{WW} T \wedge \frac{i}{2} (dW_{\hat{x}} \wedge d\bar{W}_{\hat{x}}).$$

In the resonant case, the formula giving $J_n(D)$ is modified as follows :

$$\begin{aligned} J_n(D) &= d^n |\alpha_{n,\hat{x}}|^2 \int_D \psi_n F_{\hat{x},n}^{ZZ} T \wedge \frac{i}{2} (dZ_{\hat{x}} \wedge d\bar{Z}_{\hat{x}}) \\ &\quad + 2d^n \int_D \psi_n \operatorname{Re} \left[\left((qF_{\hat{x},n}^{ZZ} \alpha_{n,\hat{x}} \bar{\gamma}_{n,\hat{x}} \bar{W}_{\hat{x}}^{q-1} + F_{\hat{x},n}^{ZW} \alpha_{n,\hat{x}} \bar{\beta}_{n,\hat{x}}) T \wedge \frac{i}{2} (dZ_{\hat{x}} \wedge d\bar{W}_{\hat{x}}) \right) \right] \\ &\quad + d^n \int_D \psi_n q^2 F_{\hat{x},n}^{ZZ} |\gamma_{n,\hat{x}}|^2 |W_{\hat{x}}|^{2(q-1)} T \wedge \frac{i}{2} (dW_{\hat{x}} \wedge d\bar{W}_{\hat{x}}) \\ &\quad + d^n \int_D 2\operatorname{Re} \left(qF_{\hat{x},n}^{ZW} \gamma_{n,\hat{x}} \bar{\beta}_{n,\hat{x}} W_{\hat{x}}^{q-1} \right) T \wedge \frac{i}{2} (dW_{\hat{x}} \wedge d\bar{W}_{\hat{x}}). \end{aligned}$$

Let us verify that the first item of Proposition 5.1 remains valid in the resonant case. Let \hat{x} be a $\hat{\mu}$ -generic element and let $(n_k)_k$ be a sub-sequence such that $\hat{x}_{-n_k} \in \mathcal{A}_\tau$. Assume

$$\psi_n = \psi \circ f_{\hat{x}}^{-n_k} \leq b, \quad \mu - a.e. \text{ on } B$$

and let us verify that there exists $C'(\hat{x})$ such that for every borel set $U \subset B$:

$$J_{n_k}(U) \leq C'(\hat{x}) e^{-n_k(\lambda_1 + \lambda_2 - \operatorname{Log} d)} e^{2n_k \varepsilon}.$$

We use the formula on $J_n(D)$ given above and Cauchy-Schwarz inequality. We also use $|F_{\hat{x},n}^{AB}| \leq \tau$ given by Equation (8), item (i) of Lemma 2.5 and $|W_{\hat{x}}| \leq \rho_\varepsilon := \rho_\varepsilon(\hat{x})$ on B provided by Equation (6). If $\sigma_{\hat{x}} := T \wedge dd^c |W_{\hat{x}}|^2$, $\lambda_{\hat{x}} := T \wedge dd^c |Z_{\hat{x}}|^2$, we indeed obtain

$$\begin{aligned} J_{n_k}(U) &\leq d^{n_k} |\alpha_{n_k,\hat{x}}|^2 b\tau \lambda_{\hat{x}}(U) \\ &\quad + 2d^{n_k} |\alpha_{n_k,\hat{x}}| b\tau \left(q|\gamma_{n_k,\hat{x}}| \rho_\varepsilon^{q-1} + |\beta_{n_k,\hat{x}}| \right) (\lambda_{\hat{x}} + \sigma_{\hat{x}})(U) \\ &\quad + qd^{n_k} |\gamma_{n_k,\hat{x}}| b\tau \left(q|\gamma_{n_k,\hat{x}}| \rho_\varepsilon^{2(q-1)} + 2|\beta_{n_k,\hat{x}}| \rho_\varepsilon^{q-1} \right) \sigma_{\hat{x}}(U). \end{aligned}$$

The existence of $C'(\hat{x})$ then comes from the upper bounds $|\alpha_{n,\hat{x}}| \leq e^{-n(\lambda_1-\varepsilon)}$, $|\beta_{n,\hat{x}}| \leq e^{-n(\lambda_2-\varepsilon)}$, $|\gamma_{n,\hat{x}}| \leq M_\varepsilon(\hat{x})e^{-n(\lambda_1-\varepsilon)}$ and from $\lambda_1 > \lambda_2$.

8 Suspensions of one-dimensional Lattès maps

Let $\theta = [P : Q]$ be Lattès map on \mathbb{P}^1 of degree $d \geq 2$ and let $f := [P : Q : t^d]$ be its suspension on \mathbb{P}^2 . These maps were studied by Berteloot-Loeb [4]. We identify the first affine chart with \mathbb{C}^2 . Let

$$G_\theta := \lim_{n \rightarrow +\infty} \frac{1}{d^n} \operatorname{Log} \|(P^n, Q^n)\| : \mathbb{C}^2 \rightarrow \mathbb{R} \cup \{-\infty\}$$

be the Green function of the polynomial mapping (P, Q) and let $G := \max\{G_\theta, 0\}$. The attracting basin A of $(0, 0)$ is bounded and equal to $\{G_\theta < 0\}$. The Green current of f satisfies $T = dd^c G$ on \mathbb{C}^2 , the support of μ coincides with the boundary ∂A of A .

By [4, Proposition 3.1] (see also [11] in higher dimensions), for every p outside a finite number of circles drawn on ∂A , there exists a biholomorphism $\mathbf{p} : (\mathbb{D}^2, 0) \longrightarrow (\mathbb{D}^2, p)$ such that

$$G_0(z, w) := G_\theta \circ \mathbf{p}(z, w) = \operatorname{Re}(z) + |w|^2 \text{ on } (\mathbb{D}^2, 0).$$

We denote

$$T_0 := \mathbf{p}^* T = dd^c(\max\{G_0, 0\}) \text{ and } \mu_0 := \mathbf{p}^* \mu.$$

The closed positive current T_0 has the following matrix representation

$$T_0 = \begin{bmatrix} T_{11} & 0 \\ 0 & T_{22} \end{bmatrix},$$

where $T_{11} := T_0 \wedge dd^c |w|^2$ and $T_{22} := T \wedge dd^c |z|^2$. We describe below T_0 and μ_0 on $(\mathbb{D}^2, 0)$. We assume to simplify that the germ $(\mathbb{D}^2, 0)$ is the open set $D :=]-1, 1[\times \mathbb{D}$. Let

$$\Omega := \{\operatorname{Re}(z) + |w|^2 > 0\} \cap D \text{ and } M_0 := \{\operatorname{Re}(z) + |w|^2 = 0\} \cap D.$$

Observe that M_0 is the image (intersected with D) of the euclidian 3-sphere of \mathbb{C}^2 by the classical Cayley transformation, see [18, Chapter 2.3]. Let us parametrize M_0 by

$$\Phi :]-1, 0[\times]-1, 1[\times]0, 2\pi[\rightarrow M_0, \quad (u, v, \theta) \mapsto (u + iv, \sqrt{-u}e^{i\theta}) \quad (13)$$

and let $\operatorname{Leb}_{M_0} := \Phi_* \operatorname{Leb}_{\mathbb{R}^3}$. Let also ω_0 be the standard hermitian $(1, 1)$ -form of \mathbb{C}^2 .

Proposition 8.1. *With the preceding notations,*

$$T_0 = \begin{bmatrix} \frac{1}{8} \operatorname{Leb}_{M_0} & 0 \\ 0 & \operatorname{Leb}_\Omega + \frac{|w|^2}{2} \operatorname{Leb}_{M_0} \end{bmatrix},$$

and

$$\mu_0 = \frac{1}{8} \operatorname{Leb}_{M_0} = T_0 \wedge dd^c |w|^2.$$

In particular,

$$\mu_0 = \psi_0(T_0 \wedge \omega_0), \text{ where } \psi_0(z, w) := \frac{1}{1 + 4|w|^2} \mathbb{1}_{M_0}.$$

The remainder of this Section is devoted to the proof of Proposition 8.1.

Computation of T_{11} .— Let φ be a test function with compact support in D . We denote $\varphi_{z\bar{z}} := \frac{\partial}{\partial z} \frac{\partial}{\partial \bar{z}} \varphi$, similarly for $\varphi_{w\bar{w}}$. By definition we have :

$$\langle T_{11}, \varphi \rangle = \int_\Omega G_0(z, w) \varphi_{z\bar{z}}(z, w) \, d\operatorname{Leb}(z, w).$$

Let Δ_w denote the subset $] -1, 1[\times] -1, 1[\times \{w\}$ intersected with Ω , so that

$$\langle T_{11}, \varphi \rangle = \int_{\mathbb{D}} \left[\int_{\Delta_w} G_0(z, w) \varphi_{z\bar{z}}(z, w) \, d\operatorname{Leb}(z) \right] d\operatorname{Leb}(w) =: \int_{\mathbb{D}} I_w(\varphi) \, d\operatorname{Leb}(w). \quad (14)$$

We write $z = u + iv$, hence $\varphi_{z\bar{z}} = \frac{1}{4} \Delta_{u,v} \varphi$. Observe also that G_0 only depends on (u, w) . The number $4I_w(\varphi)$ is then equal to

$$\int_{-1}^1 \left[\int_{-|w|^2}^1 G_0(u, w) \frac{\partial^2 \varphi}{\partial u^2}(u + iv, w) \, du \right] dv + \int_{-|w|^2}^1 G_0(u, w) \left[\int_{-1}^1 \frac{\partial^2 \varphi}{\partial v^2}(u + iv, w) \, dv \right] du.$$

The second term vanishes since $\text{Supp}(\varphi(\cdot, w)) \subset]-1, 1[\times]-1, 1[$. Integrating by parts and using $\partial_u G_0 = 1$, $\partial_v G_0 = 0$ for the first one, we get

$$4I_w(\varphi) = - \int_{-1}^1 \left[\int_{-|w|^2}^1 \frac{\partial \varphi}{\partial u}(u + iv, w) du \right] dv .$$

Using again $\text{Supp}(\varphi(\cdot, w)) \subset]-1, 1[\times]-1, 1[$, we get

$$I_w(\varphi) = \frac{1}{4} \int_{-1}^1 \varphi(-|w|^2 + iv, w) dv .$$

We therefore obtain by (14) :

$$\langle T_{11}, \varphi \rangle = \frac{1}{4} \int_{\mathbb{D}} \int_{-1}^1 \varphi(-|w|^2 + iv, w) dv d\text{Leb}(w)$$

Observe that $(w, v) \mapsto (-|w|^2 + iv, w)$ is another parametrization of M_0 different from Φ defined in (13). Let us write this integral with polar coordinates :

$$\langle T_{11}, \varphi \rangle = \frac{1}{8} \int_{-1}^1 \int_0^{2\pi} \int_0^1 \varphi(-\rho^2 + iv, \rho e^{i\theta}) 2\rho d\rho d\theta dv .$$

Applying the change of variable $u = -\rho^2$ one gets back to the parametrization Φ :

$$\langle T_{11}, \varphi \rangle = \frac{1}{8} \int_{]-1, 0] \times]-1, 1[\times]0, 2\pi[} \varphi(u + iv, \sqrt{-u}e^{i\theta}) d\text{Leb}(u, v, \theta) = \frac{1}{8} \int_{M_0} \varphi d\text{Leb}_{M_0} .$$

Computation of T_{22} . - Let Δ_z be the vertical disc $\{z\} \times \mathbb{D}$ intersected with Ω . One has :

$$\langle T_{22}, \varphi \rangle = \int_{]-1, 1[\times]-1, 1[} \left[\int_{\Delta_z} G_0(z, w) \varphi_{w\bar{w}}(z, w) d\text{Leb}(w) \right] d\text{Leb}(z) .$$

Writing $z = u + iv$, we get

$$\langle T_{22}, \varphi \rangle = \int_{-1}^1 \left[\int_{-1}^0 J_{u,v}(\varphi) du \right] dv + \int_{-1}^1 \left[\int_0^1 J_{u,v}(\varphi) du \right] dv ,$$

where

$$J_{u,v}(\varphi) := \int_{\mathbb{D} \setminus \mathbb{D}(\sqrt{-u})} G_0(u + iv, w) \varphi_{w\bar{w}}(u + iv, w) d\text{Leb}(w) ,$$

with the convention $\mathbb{D}(\sqrt{-u}) = \emptyset$ if $0 \leq u \leq 1$.

Lemma 8.2.

1. $J_{u,v}(\varphi) = \int_{\mathbb{D}} \varphi(u + iv, w) d\text{Leb}(w)$ for $0 \leq u \leq 1$.
2. $J_{u,v}(\varphi) = \int_{\mathbb{D} \setminus \mathbb{D}(\sqrt{-u})} \varphi(z, w) d\text{Leb}(w) - \frac{u}{2} \int_0^{2\pi} \varphi(z, \sqrt{-u}e^{i\theta}) d\theta$ for $-1 \leq u \leq 0$.

Proof : We write $w = s + it$ and use $\varphi_{w\bar{w}} = \frac{1}{4} \Delta_{s,t} \varphi$. Since φ has compact support in D , one gets by Green-Stokes formula :

$$4J_{u,v}(\varphi) := \int_{\mathbb{D}} G_0(u + iv, w) \Delta_{s,t} \varphi d\text{Leb}(w) = \int_{\mathbb{D}} \Delta_{s,t} G_0(u + iv, w) \varphi d\text{Leb}(w)$$

The first item follows since $\partial_s^2 G_0 = \partial_t^2 G_0 = 2$. In the case $-1 \leq u \leq 0$, we have :

$$J_{u,v}(\varphi) = \int_{\mathbb{D} \setminus \mathbb{D}(\sqrt{-u})} G_0(u + iv, w) \varphi_{w\bar{w}}(u + iv, w) d\text{Leb}(w) . \quad (15)$$

Let us denote $z = u + iv$ and $w = s + it \in \mathbb{D} \setminus \mathbb{D}(\sqrt{-u})$. Using $\text{Supp}(\varphi(z, \cdot)) \subset \mathbb{D}$ and $G_0(z, \cdot)|_{\partial\mathbb{D}(\sqrt{-u})} = 0$, Green-Stokes formula gives :

$$\begin{aligned} & \int_{\mathbb{D} \setminus \mathbb{D}(\sqrt{-u})} G_0(z, s + it) \Delta_{s,t} \varphi(z, s + it) - \Delta_{s,t} G_0(z, s + it) \varphi(z, s + it) \, d\text{Leb}(s, t) \\ &= \int_{\partial\mathbb{D}(\sqrt{-u})} \varphi(z, s + it) (\nabla_{(s,t)} G_0(z, s + it)) \cdot (s, t) \, d\sigma(s, t), \end{aligned}$$

where σ is the Lebesgue measure on the circle $\partial\mathbb{D}(\sqrt{-u})$. Since $\Delta_{s,t} G_0(z, s + it) = 4$ and $(\nabla_{(s,t)} G_0(z, s + it)) \cdot (s, t) = 2|w|^2 = -2u$, we infer

$$4 \int_{\mathbb{D} \setminus \mathbb{D}(\sqrt{-u})} \varphi(z, w) - G_0(z, w) \varphi_{w\bar{w}}(z, w) \, d\text{Leb}(w) = \int_{\partial\mathbb{D}(\sqrt{-u})} 2u \varphi(z, w) \, d\sigma(w).$$

The second item follows by combining this formula with (15). \square

Let us now return to the computation of $\langle T_{22}, \varphi \rangle$. Let $\Omega^+ := \Omega \cap \{u > 0\}$ and $\Omega^- := \Omega \cap \{u < 0\}$. According to Lemma 8.2, we get

$$\int_{-1}^1 \int_0^1 J_{u,v}(\varphi) \, du \, dv = \int_{\Omega^+} \varphi(z, w) \, d\text{Leb}(z, w),$$

and that $\int_{-1}^1 \int_{-1}^0 J_{u,v}(\varphi) \, du \, dv$ is equal to

$$\int_{\Omega^-} \varphi(z, w) \, d\text{Leb}(z, w) + \int_{-1}^1 \left[\int_{-1}^0 \left[\int_0^{2\pi} -\frac{u}{2} \varphi(u + iv, \sqrt{-u} e^{-i\theta}) \, d\theta \right] du \right] dv.$$

If $h(z, w) := \frac{|w|^2}{2}$ the second integral is equal to :

$$\int_{]-1, 0[\times]-1, 1[\times]0, 2\pi[} (h\varphi) \circ \Phi(u, v, \theta) \, d(u, v, \theta) = \int_{M_0} h\varphi \, d\text{Leb}_{M_0}.$$

We deduce as desired

$$\langle T_{22}, \varphi \rangle = \int_{-1}^1 \int_{-1}^1 J_{u,v}(\varphi) \, du \, dv = \int_{\Omega} \varphi(z, w) \, d\text{Leb}(z, w) + \int_{M_0} h\varphi \, d\text{Leb}_{M_0}.$$

Computation of μ_0 .— Writing $G_0^+ := \max\{G_0, 0\}$, we have $T_0 = dd^c G_0^+$ and $\mu_0 = dd^c(G_0^+ dd^c G_0^+)$. Hence

$$\begin{aligned} \langle \mu_0, \varphi \rangle &= \langle T_0, G_0^+ dd^c \varphi \rangle = \langle T_{11} dd^c |z|^2 + T_{22} dd^c |w|^2, G_0^+ dd^c \varphi \rangle \\ &= \frac{1}{8} \int_{M_0} G_0 \varphi_{w\bar{w}} \, d\text{Leb}_{M_0} + \int_{\Omega} G_0 \varphi_{z\bar{z}} \, d\text{Leb} + \int_{M_0} G_0 \varphi_{z\bar{z}} h \, d\text{Leb}_{M_0}. \end{aligned}$$

Since $G_0 = 0$ on M_0 , we get $\langle \mu_0, \varphi \rangle = \int_{\Omega} G_0 \varphi_{z\bar{z}} \, d\text{Leb}_{\Omega} = \langle T_{11}, \varphi \rangle$ as desired (observe that this is the Lebesgue part on Ω in T_{22} which provides $\mu = T_{11}$). The formula $\mu_0 = \frac{1}{1+8h} \mathbb{1}_{M_0} T_0 \wedge \omega_0$ immediately follows.

References

- [1] F. Berteloot and C. Dupont. Une caractérisation des endomorphismes de Lattès par leur mesure de Green. *Comment. Math. Helv.*, 80(2):433–454, 2005.

- [2] F. Berteloot and C. Dupont. A distortion theorem for the iterated inverse branches of a holomorphic endomorphism of $\mathbb{C}\mathbb{P}^k$. *Journal of the London Mathematical Society*, 99(1):153–172, 2019.
- [3] F. Berteloot, C. Dupont, and L. Molino. Normalization of bundle holomorphic contractions and applications to dynamics. *Ann. Inst. Fourier (Grenoble)*, 58(6):2137–2168, 2008.
- [4] F. Berteloot and J.-J. Loeb. Spherical hypersurfaces and Lattès rational maps. *J. Math. Pures Appl. (9)*, 77(7):655–666, 1998.
- [5] F. Berteloot and J.-J. Loeb. Une caractérisation géométrique des exemples de Lattès de \mathbb{P}^k . *Bull. Soc. Math. France*, 129(2):175–188, 2001.
- [6] J.-Y. Briend and J. Duval. Exposants de Liapounoff et distribution des points périodiques d’un endomorphisme de $\mathbb{C}\mathbb{P}^k$. *Acta Math.*, 182(2):143–157, 1999.
- [7] J. Buzzzi. The coding of non-uniformly expanding maps with application to endomorphisms of $\mathbb{C}\mathbb{P}^k$. *Ergodic Theory and Dynamical Systems*, 23(4):1015–1024, 2003.
- [8] I. Cornfeld, A. Sossinskii, S. Fomin, and Y. Sinai. *Ergodic Theory*. Grundlehren der mathematischen Wissenschaften. Springer New York, 2012.
- [9] T.-C. Dinh and N. Sibony. *Dynamics in several complex variables: endomorphisms of projective spaces and polynomial-like mappings*, volume 1998 of *Lecture Notes in Math.*, pages 165–294. Springer, Berlin, 2010.
- [10] R. Dujardin. Fatou directions along the Julia set for endomorphisms of $\mathbb{C}\mathbb{P}^k$. *J. Math. Pures Appl. (9)*, 98(6):591–615, 2012.
- [11] C. Dupont. Exemples de Lattès et domaines faiblement sphériques de \mathbb{C}^n . *Manuscripta Math.*, 111(3):357–378, 2003.
- [12] C. Dupont. Formule de Pesin et applications méromorphes. *Bull. Braz. Math. Soc. (N.S.)*, 37(3):393–418, 2006.
- [13] C. Dupont and A. Rogue. Dimension of ergodic measures and currents on $\mathbb{C}\mathbb{P}^2$. *Ergodic Theory and Dynamical Systems*, 40(8):2131–2155, 2020.
- [14] C. Dupont and J. Taflin. Dynamics of fibered endomorphisms of \mathbb{P}^k . *Ann. Sc. Norm. Super. Pisa Cl. Sci. (5)*, 22(1):53–78, 2021.
- [15] M. Jonsson. Dynamics of polynomial skew products on \mathbb{C}^2 . *Math. Ann.*, 314(3):403–447, 1999.
- [16] M. Jonsson and D. Varolin. Stable manifolds of holomorphic diffeomorphisms. *Invent. Math.*, 149:409–430, Aug. 2002.
- [17] P. Mattila. *Geometry of sets and measures in Euclidean spaces*, volume 44 of *Cambridge Studies in Advanced Mathematics*. Cambridge University Press, Cambridge, 1995. Fractals and rectifiability.
- [18] W. Rudin. *Function theory in the unit ball of \mathbb{C}^n* . Classics in Mathematics. Springer-Verlag, Berlin, 2008. Reprint of the 1980 edition.
- [19] N. Sibony. Dynamique des applications rationnelles de \mathbb{P}^k . In *Dynamique et géométrie complexes (Lyon, 1997)*, volume 8 of *Panor. Synthèses*, pages ix–x, xi–xii, 97–185. Soc. Math. France, Paris, 1999.
- [20] V. Tapiero. Absolute continuity for semi-extremal holomorphic mappings, arxiv:2211.17006, 2022.

Université de Rennes
CNRS, IRMAR - UMR 6625
35042 Campus de Beaulieu, Rennes, France
christophe.dupont@univ-rennes.fr
virgile.tapiero@univ-rennes.fr