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# Strong Green Photoluminescence in $Zr_3N_4$ Thin Films with Orthorhombic Structure

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## ABSTRACT:

Green light causes the maximum luminous sensation in human eyes, which will bring many advantages for potential uses. However, the efficiency of light-emitting diodes (LEDs) in the green wavelength range (495-570 nm) is much lower than that in red or blue regions, which is known as the “green gap” problem. In this study, we report the strong green photoluminescence in metastable  $Zr_3N_4$  thin films with orthorhombic structure for the first time. Crystallized  $Zr_3N_4$  thin films have been successfully grown by high-power impulse magnetron sputtering. Notable green photoluminescence in the wavelength region of 490-535 nm has been identified in such orthorhombic  $Zr_3N_4$  thin films even at room temperature. In addition, a joint experimental and theoretical study has been performed to investigate its electronic structures and to reveal the origins of such green light emissions. The calculations of the transition dipole moments demonstrate that the direct transition between the conduction band minimum and the energy level II below the valence band maximum are responsible for such strong green photoluminescence. This work may pave a new way to solve to the “green gap” problem in LEDs with low cost.

**Keywords:** Nitride semiconductors; Thin films; Green photoluminescence; Electronic structure;

## 1. INTRODUCTION

Metal nitrides are a special class of materials, which have been widely studied in both traditionally mechanical parts and modern electronic or optoelectronic devices.<sup>1-4</sup> For an example, rocksalt nitrides (such as TiN, ZrN and HfN) with superior hardness and corrosion resistance have been extensively applied as wear-resistance coatings.<sup>5-8</sup> The high conductivity and stability of TiN and TaN also facilitate them to be the diffusion barrier layers in integrated circuits.<sup>9, 10</sup> The second case is the III-nitride-based light-emitting diodes (LEDs), which have become the most promising candidates for high efficiency solid state lighting (SSL), since the first breakthrough of high-brightness InGaN-based LEDs in 1990s.<sup>11, 12</sup> Currently, the external quantum efficiency (EQE) of InGaN-based blue LEDs already exceed 80%.<sup>13, 14</sup> However, for the wavelengths beyond 490 nm, the EQE of these InGaN-based LEDs decreases significantly when the wavelength increases. Coincidentally, the phosphide-based LEDs of AlGaInP, exhibit a dramatical drop of EQE for wavelengths shorter than 570 nm.<sup>15, 16</sup> Thus, the commercialized LEDs in the wavelength region of 490 - 570 nm to match the high spectral sensitivity of the human visual perception of brightness are still absent, which is known as the “green gap” problem.<sup>13, 17</sup> Therefore, it is of great value to investigate novel semiconductors with green emissions.

Among the various nitrides, binary Zr-N system has attracted much attention, where most of the research focused on the stable phase of ZrN with exceptional mechanical properties as mentioned above. In addition to this stable phase with cubic structure, theoretical calculations also predict the existence of metastable  $Zr_3N_4$  with four different structures:<sup>18-20</sup> (1) the orthorhombic phase (O- $Zr_3N_4$ ) with a space group of *Pnam*. (2) the cubic spinel phase (C- $Zr_3N_4$  (Spinel)) with a space group of *Fd3m*. (3) the defective phase (D- $Zr_3N_4$ ) formed by ordered Zr vacancies in simple cubic phase ZrN. (4) the cubic  $Th_3P_4$ -type phase (C- $Zr_3N_4$  ( $Th_3P_4$ )) with a space group of *I43d*. The stable ZrN exhibits the metallic conduction with the low resistivity, whereas, metastable phases of  $Zr_3N_4$  are calculated to be semiconductors. So far, only the metastable phase of C- $Zr_3N_4$ ( $Th_3P_4$ ) has been successfully synthesized by extreme conditions, which exhibits the hardness as large as 30-37 GPa.<sup>21, 22</sup> However, due to

the difficulty to fabricate the metastable phase of O-Zr<sub>3</sub>N<sub>4</sub>, its functional properties and electronic structures are still unknown.

In this work, orthorhombic Zr<sub>3</sub>N<sub>4</sub> thin films have been successfully deposited on glass and silicon (100) substrates by high-power impulse magnetron sputtering (HiPIMS). Strong green photoluminescence has been identified in such thin films even at room temperature for the first time. In addition, the electronic structure, optical and mechanical properties of orthorhombic Zr<sub>3</sub>N<sub>4</sub> have been investigated.

## 2. EXPERIMENTAL SECTION

**2.1. Experimental details.** Metastable Zr<sub>3</sub>N<sub>4</sub> thin films were deposited on glass and Si(100) substrates by HIPIMS in Ar-N<sub>2</sub> reactive mixture. It was shown that the phase structure and preferred orientation of thin films are independent of the nature of non-matched substrate (glass or silicon), as the native amorphous SiO<sub>2</sub> layer on single crystal Si (100) substrate was not removed. No intentional heating was applied to the substrate holder. A Zr target (99.9% purity) with rectangular shape (360 x 80 x 8 mm) was connected to a HIPIMS power supplier (Analoge und Digitale Leistungselektronik GmbH, GX 150/1000). The Ar and N<sub>2</sub> flow rate were fixed at 40 and 6 sccm, respectively. The total sputtering pressure was 0.9 Pa. The distance between target and substrate was 115 mm. The power supply worked on the constant power mode with the average power of 300 W, the frequency of 1000 Hz and the duty cycle of 2%. In addition, ZrN thin films with gold color were also deposited by HIPIMS for a comparison. For this purpose, the Ar and N<sub>2</sub> flow rate were fixed at 160 and 4 sccm, respectively, and the total sputtering pressure was 0.5 Pa. The power supply worked on the constant power mode with the average power of 1000 W, the frequency of 1000 Hz and the duty cycle of 5%.

Thin film thickness was measured by a surface profilometer (XP-1). The resistance was measured by a four-point probing system (RTS-8). X-ray diffraction (XRD, Rigaku Ultima IV) with Cu K $\alpha$  radiation ( $\lambda = 0.15406$  nm) in Bragg Brentano configuration was used to analyze the phase structures. X-ray photoelectron spectroscopy (XPS) analyses were performed in an ultrahigh vacuum (UHV)

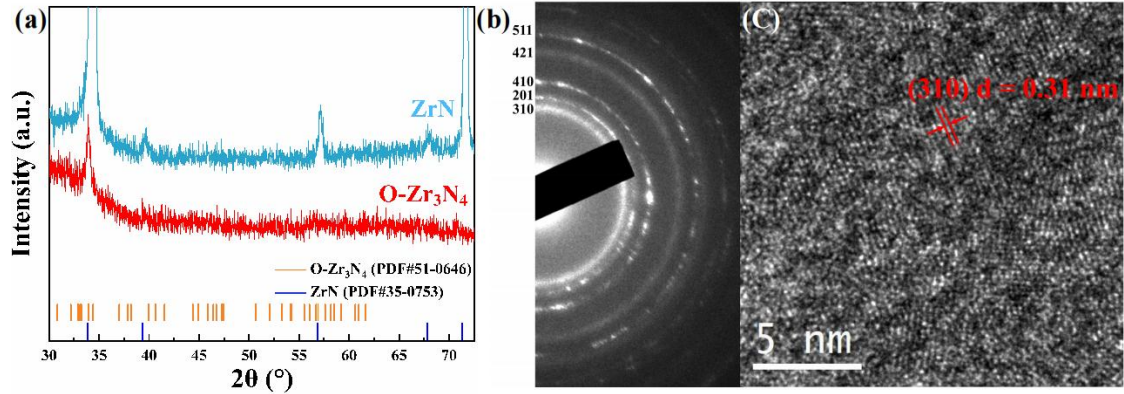
experimental setup equipped with a photoemission analyzer and an Al K $\alpha$  (1486.7 eV) x-ray source (Thermo Scientific Escalab 250Xi+). The Ar<sup>+</sup> ion etching was performed to clean the surface until there is no evolution in the C-1s XPS core level spectra. Scanning electron microscope (SEM, FEI Nova NanoSEM 450) and transmission electron microscope (TEM, JEOL ARM 200F) were used to study the surface morphology and microstructure of thin films, respectively. The optical absorption coefficient was determined by an ultraviolet-visible-near-infrared (UV-Vis-NIR) spectrophotometry (Lambda 1050). Photoluminescence (PL, FLS980) spectra were recorded at different temperatures (80 – 300 K) using a laser with the wavelength of 365 nm. The hardness of thin films was measured by a Nanoindenter (TTX-NHT2).

**2.2. Theoretical calculations.** The electronic structure calculations of O-Zr<sub>3</sub>N<sub>4</sub> were performed by the projector-augmented-wave (PAW) pseudopotentials, as implemented in Vienna *ab initio* Simulation Package (VASP) code. The plane-wave cut off was set to 400 eV. The HSE06 hybrid functional was used, which can accurately predict the electronic structures of nitride semiconductors.<sup>23-25</sup> The HF mixing constant  $\alpha$  of 0.25 was employed.<sup>26</sup> The generalized gradient approximation (GGA) of the Perdew-Burke-Ernzerh (PBE)<sup>27</sup> functional was used to calculate sum of the squares of the transition dipole moments ( $P^2$ ) for the analysis of the probability of optical transition between different energy levels.<sup>28</sup>

### 3. RESULTS AND DISCUSSIONS

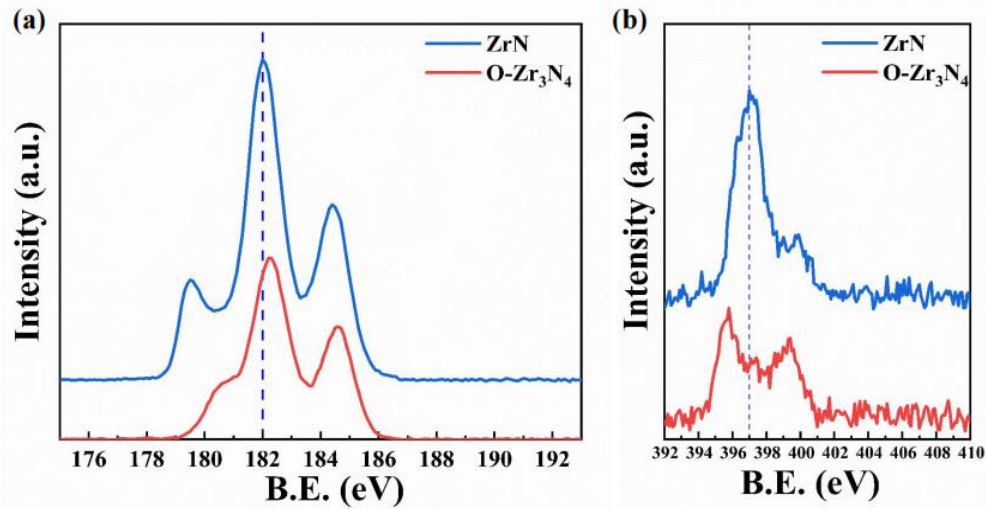
**3.1. Identifying the structure of O- Zr<sub>3</sub>N<sub>4</sub> thin films.** The XRD diffractogram of nitrogen-rich zirconium nitride thin film with the thickness of 156 nm is shown in **Figure 1a**, which exhibits a diffraction peak at  $2\theta$  of about 34.0°. Referring to the standard diffraction patterns of ZrN (PDF#35-0753) and O-Zr<sub>3</sub>N<sub>4</sub> (PDF#51-0646) in **Figure 1a**, it is difficult to identify the phase structure of such nitrogen-rich zirconium nitride thin film. But the electron diffraction pattern of TEM provides additional and valuable information (Supplementary Table S1, the lattice constants calculated by electron diffraction pattern are consistent with the reported result<sup>29, 30</sup>). As seen in **Figure 1b**, the electron diffraction pattern can be well indexed by the O-Zr<sub>3</sub>N<sub>4</sub>, rather

than cubic phases of ZrN, C-Zr<sub>3</sub>N<sub>4</sub> (Spinel), D-Zr<sub>3</sub>N<sub>4</sub> and C-Zr<sub>3</sub>N<sub>4</sub> (Th<sub>3</sub>P<sub>4</sub>). The high resolution TEM (HRTEM) image is shown in **Figure 1c**, where the crystal planes with a distance of 0.31 nm are determined as the O-Zr<sub>3</sub>N<sub>4</sub> (310). These evidences demonstrate the successful growth of O-Zr<sub>3</sub>N<sub>4</sub> thin films. The resistivity of O-Zr<sub>3</sub>N<sub>4</sub> thin films at 300 K is as high as 76000 Ω cm, whereas the ZrN thin film possesses the resistivity of 1.1 x 10<sup>-3</sup> Ω cm.



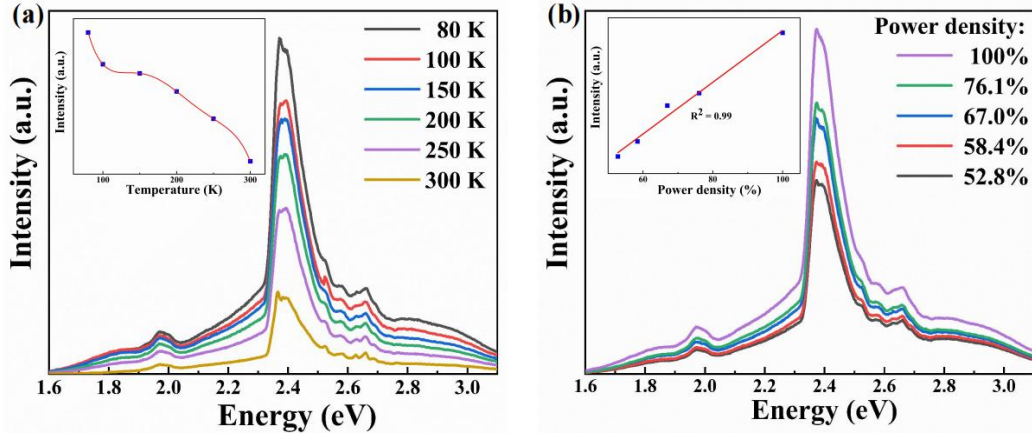
**Figure 1.** (a) The XRD diffractograms of ZrN and O-Zr<sub>3</sub>N<sub>4</sub> thin films. The standard diffraction patterns of ZrN (PDF#35-0753) and O-Zr<sub>3</sub>N<sub>4</sub> (PDF#51-0646) are also presented in the bottom. (b) The electron diffraction pattern and (c) the HRTEM image of O-Zr<sub>3</sub>N<sub>4</sub> thin film.

The Zr-3d and N-1s core level XPS spectra of O-Zr<sub>3</sub>N<sub>4</sub> are compared with those of ZrN (see **Figure 2**). As seen in **Figure 2a**, the Zr-3d peaks at the binding energy of 182 eV exhibit slight shift to higher energy, when changing the phases from ZrN to O-Zr<sub>3</sub>N<sub>4</sub>. On the contrary, the N-1s peaks move towards to lower binding energy, as shown in **Figure 2b**. These opposite directions of Zr-3d and N-1s are in agreement with the reported results,<sup>30</sup> where metallic ZrN was compared with insulator Zr<sub>3</sub>N<sub>4</sub> with unknown phase structure. Such evolution tendency of binding energy is an indicator of electronic changes and charge transfer from Zr to N, which occurs when the phases change from ZrN to O-Zr<sub>3</sub>N<sub>4</sub>. Thus, the ionic bonding becomes stronger and the charge is more localized, giving rise to an insulator.<sup>30</sup>



**Figure 2.** (a) Zr-3d and (b) N-1s core level XPS spectra of ZrN and O-Zr<sub>3</sub>N<sub>4</sub> thin films.

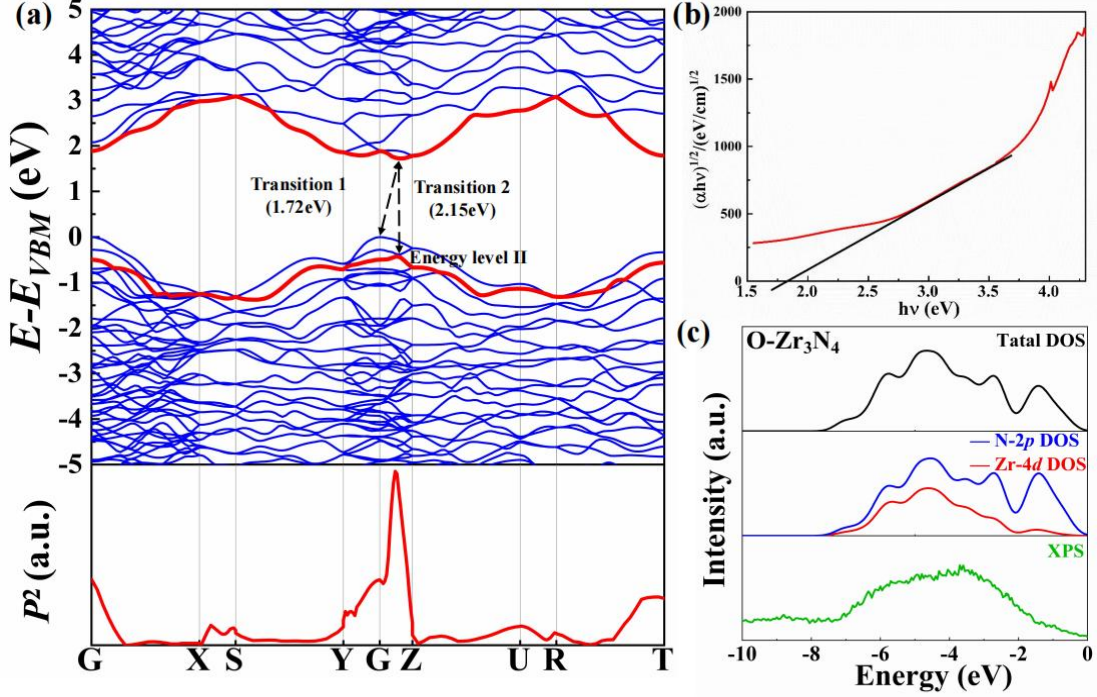
**3.2. Strong green photoluminescence identified in O-Zr<sub>3</sub>N<sub>4</sub> thin film.** The temperature-dependent photoluminescence (PL) properties of O-Zr<sub>3</sub>N<sub>4</sub> thin film have been investigated from 80 to 300 K, as shown in **Figure 3**. Two luminescence peaks at about 2.4 and 1.97 eV are clearly evident in all PL spectra acquired at different temperature, as shown in **Figure 3a**. Especially, the luminescence peak of 2.4 eV located on the green wavelength range exhibit quite strong intensity, even at 300 K (see **Figure 3a**). It is worth noting that the both PL peaks do not shift with increasing temperature, which means that these light-emission are the intrinsic characters of O-Zr<sub>3</sub>N<sub>4</sub>, rather than from the defects. Meanwhile, the radiative lifetime of 1.6  $\mu$ s obtained the time-resolved PL (see **Figure S1**) by is much larger than the PL lifetime of exciton emission in GaN and ZnO.<sup>31,32</sup> Such luminescence peaks are asymmetrical broadening at high energy part, which is the characteristic of the PL spectra of indirect bandgap semiconductors.<sup>33</sup> We also investigate the power-dependent PL properties of O-Zr<sub>3</sub>N<sub>4</sub> thin films (see **Figure 3b**). With the increase of laser power, the intensity of both peaks increases, while their positions do not shift. Specially, the peak intensities at 2.4 eV exhibit a linear relationship with the laser power. This also confirms that these light-emission are the intrinsic characters of O-Zr<sub>3</sub>N<sub>4</sub>.



**Figure 3.** (a) Temperature-dependent and (b) power-dependent PL emission spectra of O-Zr<sub>3</sub>N<sub>4</sub> deposited on glass substrate. The inset of (a) shows the relationship between the peak intensity at 2.4 eV and temperature, while the inset of (b) exhibits the linear relationship between the peak intensity at 2.4 eV and power density.

**3.3. Electronic structures of O-Zr<sub>3</sub>N<sub>4</sub>.** The electronic structures of O-Zr<sub>3</sub>N<sub>4</sub> (Supplementary **Figure S2** for the crystal structure) have been calculated by first-principles calculations using the HSE06 functionals. The computed crystal constants ( $a = 0.976$  nm;  $b = 1.081$  nm;  $c = 0.328$  nm) are consistent with the experimental ones ( $a = 0.973$  nm;  $b = 1.082$  nm;  $c = 0.328$  nm) achieved above. **Figure 4a** shows the electronic band dispersion of O-Zr<sub>3</sub>N<sub>4</sub>, demonstrating an indirect band gap of 1.72 eV. Such a value is much larger than the previous reported 0.68 eV calculated by GGA-PBE.<sup>34</sup> In addition, the experimental band gap of 1.8 eV is identified in O-Zr<sub>3</sub>N<sub>4</sub> thin film by Tauc fitting of  $(\alpha E)^{1/2}$  vs.  $E$  ( $\alpha$  and  $E$  are the absorption coefficient and the photon energy, respectively), as shown in **Figure 4b**. This means the theoretical and experimental band gaps are well consistent.





**Figure 4.** (a) Electronic band dispersion and corresponding  $P^2$  of O-Zr<sub>3</sub>N<sub>4</sub>. (b) Tauc fitting of  $(\alpha E)^{1/2}$  vs.  $E$  for O-Zr<sub>3</sub>N<sub>4</sub> thin film. (c) Theoretically total DOS, Zr-4d and N-2p PDOS, are compared with the experimental photoemission valence band spectrum of O-Zr<sub>3</sub>N<sub>4</sub>.

The XPS valence band spectrum of O-Zr<sub>3</sub>N<sub>4</sub> thin film is compared with the calculated density of states (DOS). Before discussing the results in detail, we briefly comment on the relative sensitivities of the Al  $K\alpha$  (1486.7 eV) photoemission source on the N-2p and Zr-4d spectral weights. The cross-section ratio of  $\sigma(\text{N-}2p) / \sigma(\text{Zr-}4d) \approx 0.07$  for Al  $K\alpha$ , is determined utilizing the known energy dependence of the photoionization cross-section.<sup>35</sup> This indicates that XPS primarily probes the Zr- $d$  states. The XPS valence band spectrum and the DOS obtained from first-principles calculations of O-Zr<sub>3</sub>N<sub>4</sub> are shown in **Figure 4c**, where the valence band maximum (VBM) has been calibrated to 0 eV. The valence band structure of O-Zr<sub>3</sub>N<sub>4</sub> can be divided into two main parts, as can be seen from the total DOS in **Figure 4c**. In the energy range of 0 to -1.8 eV, it is dominated by N-2p character. When the energy decreases to -1.9 to -6.7 eV, the contribution of Zr-4d gradually increases, and the electronic state is contributed from both Zr-4d and N-2p. Comparing Zr-4d partial density of states (PDOS) with XPS valence spectra, their spectral shape and peak

positions look quite close, as XPS primarily probes the Zr-*d* states as mentioned above.

As discussed above, the theoretically indirect bandgap of O-Zr<sub>3</sub>N<sub>4</sub> is estimated to be about 1.72 eV (see **Figure 4a**). Thus, the peak at 1.97 eV in the PL spectra may be attributed to the indirect band gap transition between the conduction band maximum (CBM) and the valence band maximum (VBM), noted as transition 1 in **Figure 4a**. In order to investigate the origin of the notable green photoluminescence at about 2.4 eV in **Figure 4a**, the squares of the transition dipole moments ( $P^2$ ) at various k points between several energy levels near the CBM and VBM were calculated (see Supplementary **Figure S3**), demonstrating that the direct transition between the CBM and the energy level II possesses the maximum transition probability (see **Figure 4a**). Looking at the energy band dispersion of O-Zr<sub>3</sub>N<sub>4</sub> in **Figure 4a**, the direct transition between the CBM and the energy level II in the valence band exhibits a gap of 2.15 eV, which may correspond to the PL peak at 2.4 eV in **Figure 3a**, considering the slight dispersion between experiments and calculations. Therefore, it is believed that the direct transition between the CBM and the energy level II produces the strong green PL speak in O-Zr<sub>3</sub>N<sub>4</sub>. In addition, the potential epitaxial relationships between O-Zr<sub>3</sub>N<sub>4</sub> and other nitride semiconductors to form heterojunction LEDs have also been investigated, which shows a mismatch of less than 1.5% of Zr<sub>3</sub>N<sub>4</sub>(001)||AlN(110) and Zr<sub>3</sub>N<sub>4</sub>(100)||GaN(010) (see Supplementary **Figure S4**). Therefore, it's believe that O-Zr<sub>3</sub>N<sub>4</sub> may be able to fabricate high-efficiency green heterojunction LEDs.

**3.4. Large hardness of O-Zr<sub>3</sub>N<sub>4</sub> thin film.** The hardness of O-Zr<sub>3</sub>N<sub>4</sub> thin film has been measured by the Nanoindenter, yielding the hardness of 8.8 GPa (see Supplementary **Figure S5**). Such a value is even larger than traditional semiconductors in the forms of single crystal, such as Ge (111) (7.12 GPa), GaP (7.73 GPa), GaAs (6.8 GPa) and ZnO (4.7 GPa).<sup>36</sup> The hardness of O-Zr<sub>3</sub>N<sub>4</sub> thin film is even close to that of GaN (0001) single crystal (10.2 GPa).<sup>36</sup> Such a large hardness of the O-Zr<sub>3</sub>N<sub>4</sub> may be an additional benefit for its semiconductor applications.

## 4. CONCLUSIONS

In conclusion, the strong green photoluminescence in crystallized O-Zr<sub>3</sub>N<sub>4</sub> thin films grown by high-power impulse magnetron sputtering was reported for the first time. Two emission peaks at approx.1.9 eV and 2.4 eV were observed in the photoluminescence spectra of O-Zr<sub>3</sub>N<sub>4</sub> thin films, and these light-emissions are considered to be the intrinsic characters of O-Zr<sub>3</sub>N<sub>4</sub>. By combining experiments and first-principle calculations, the bandgap of O-Zr<sub>3</sub>N<sub>4</sub> was determined to be about 1.8 eV, and the origins of two luminescence peaks were identified. Considering that O-Zr<sub>3</sub>N<sub>4</sub> has the perfect lattice match with AlN and GaN (lattice mismatch less than 1.5%), O-Zr<sub>3</sub>N<sub>4</sub> may be a promising candidate to fabricate the green LEDs with high efficiency.

## **SUPPORTING INFORMATION**

See the supplementary materials for the additional crystal structure characteristics, electronic band dispersions, and the potential epitaxial relationships between O-Zr<sub>3</sub>N<sub>4</sub> and other nitride semiconductors.

## **DATA AVAILABILITY**

All data supporting the findings of this article and its Supplementary Information will be made available upon reasonable request to the authors. Source data are provided with this paper.

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### **Author contributions**

Y.L. and Y.W. conceived the idea and designed the experiments. Y.L. carried out the growth, structural and optical characterizations, as well as the first-principle calculations. N.G. and B.D. contributed to first principle calculations and data analysis. K.W. participated the growth of thin film and performed the measurements of the hardness. N.L., Y.X., and H.S. performed the structural analysis. X.S. performed the TEM experiments and analysis. J.P. contributed the analysis and discussion for the results. Y.L. and Y.W. wrote the manuscript with the contributions from all the authors.

### **Notes**

The authors declare no competing financial interests.

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