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Do Stereochemical Effects Overcome a Charge-induced Perturbation in Isolated Protonated cyclo(Tyr-Tyr)?

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Abstract

The two diastereomers of the protonated diketopiperazine (DKP) dipeptide cyclo(Tyr-Tyr), namely, cyclo(LTyr-LTyr)H⁺ and cyclo(LTyr-DTyr)H⁺, are studied in a cryogenic ion trap by means of IR photodissociation spectroscopy combined with quantum chemical calculations. The two diastereomers have similar structures, in which one of the ring is folded over the DKP ring and the other one is extended in a *trans* geometry, allowing a strong OH⁺... π interaction to take place. This contrasts to the observation of a stacked geometry for neutral cyclo(LTyr-LTyr) only under supersonic expansion conditions that does not exist for cyclo(LTyr-DTyr). In the protonated form, the strength of the OH⁺... π interaction is different for the two diastereomers, resulting in a ~ 110 cm⁻¹ difference in the $\nu(\text{OH}^+)$ frequency and a smaller, but clearly identifiable difference in the protonated amide $\nu(\text{NH})$ frequency. Stereochemical effects are therefore still evidenced despite the strong perturbation due to the excess charge.

Introduction

2,5-Diketopiperazine (DKP) peptides are an important class of natural molecules synthesized mainly by bacteria in fungi or marine microorganisms. They result from the self-condensation of two α -amino acids.¹ Their rigid cyclic structure makes them a unique scaffold for the synthesis of innovative drugs.² Their cyclic nature also imposes the unusual *cis* conformation of the peptide bond, as shown in Figure 1. In contrast to the peptide backbone, which shows limited conformational mobility,³⁻⁴ the geometry of the substituents is quite flexible, in particular around the C _{α} -C _{β} bond. This flexibility results in different possible structures that

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3 depend on several factors, including the nature of the residue, the protonation or aggregation
4 state, the environment such as the solvent in solution or the crystal in the solid state.^{5-8 9-11}
5 Two different structures of peptides containing one aromatic residue coexist under jet-cooled
6 conditions: a folded structure, with the aromatic ring folded over the DKP ring, and an
7 extended one with the aromatic ring extended outwards the molecule.¹² The folded structure
8 is stabilised by solvation and is the only form observed in condensed-phase studies.⁹⁻¹¹ It
9 should be noted here that the cyclic nature of the DKP peptide deprives them from both C-
10 and N-termini. Therefore, no zwitterion is observed in the condensed phase. Dipeptides
11 containing two aromatic residues involve both repulsion between the bulky substituents and
12 dispersion interactions. Their structure is therefore dictated by the subtle balance between
13 these different forces. For cyclo(LPhe-LPhe), repulsion dominates and a single conformation
14 exists under jet-cooled conditions, with one of the phenyl substituents folded on the peptide
15 ring and the other extended outwards.¹³⁻¹⁴ The conformation does not change upon dimer
16 formation¹⁵ or in the solid state.^{5, 13} It is also very similar to that of cyclo(LPhe-DPhe).^{14, 16}
17 Cyclo(Tyr-Tyr) differs from cyclo(LPhe-LPhe) by the presence of the hydroxyl substituent in the
18 benzene ring. This substituent leads to additional *syn/anti* isomerism,¹⁷⁻¹⁸ and makes the
19 formation of an OH...O hydrogen bond possible. Conformer-selective IR-UV spectroscopy
20 showed that, besides the folded-extended conformation common to the two diastereomers,
21 a stacked conformation stabilised by an OH...O interaction exists for cyclo(LTyr-LTyr) only.¹⁹
22 Gas-phase cyclo(LTyr-LPhe) does not show this stacked geometry, although its existence was
23 suggested in polar solvents.²⁰ The OH... π interaction is therefore not sufficient to
24 counterbalance the strong repulsion between aromatic clouds.
25 Protonated cyclic dipeptides have been the subject of several studies, as possible
26 intermediates in the fragmentation of peptides.²¹⁻²⁴ Comparison between collision-induced
27 dissociation (CID) experiments and chemical dynamics simulations showed that the
28 fragmentation processes of the amide CO protonated systems does not depend much on
29 chirality, in contrast to those that are protonated on the side chain of a basic residue.²⁵⁻²⁸
30 Protonation modifies the balance between the different conformers of biomolecules.²⁹ For
31 cyclic dipeptides based on Tyr, Phe, or Pro residues, the proton is located on the amide CO
32 and does not influence much its conformation. The extended conformation of protonated
33 cyclo(Tyr-Pro) is dominant in a room-temperature ion trap, although minor amounts of folded
34 conformers are kinetically trapped.²⁶ The most stable form of protonated cyclo(Phe-Phe) is
35 similar for the two diastereomers and is a folded-extended geometry stabilized by an OH⁺... π
36 interaction.²⁸ This is the dominant species observed in protonated ions produced by
37 electrospray ionisation (ESI) at room temperature. However, correct assignment of the
38 vibrational spectrum requires the presence of higher-energy conformers resulting from kinetic
39 trapping. They involve an NH... π interaction and are suggested to be responsible for the
40 differences between the spectra of cyclo(LPhe-LPhe) and cyclo(LPhe-DPhe).²⁸ However, the
41 spectral broadening due to spectral congestion and temperature effects prevents drawing a
42 clear-cut conclusion. Low-temperature studies are therefore highly desirable for trapping the
43 low-energy species and evidence their contribution to the stereochemistry-dependent
44 spectroscopic results. Cryogenic ion traps are a powerful tool to stabilise and isolate small or
45 medium-size peptides and study their structure by diverse action spectroscopy such as Infra-
46 Red or Ultra-Violet Photodissociation (IRPD or UVPD), or a combination of both that results in
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3 conformer-specific spectroscopy.³⁰⁻³⁸ We apply here UVPD and IRPD to the study of
4 protonated cyclo(Tyr-Tyr) in a cryogenic ion trap. The aim of this study is twofold. It aims first
5 to test whether higher-energy conformers are necessary for observing a chirality effect in the
6 vibrational spectroscopy of this molecule or if an effect is already observed at low temperature
7 for the lower-energy conformer. On the other hand, the stereochemical effects observed in
8 the neutral system rest on the presence of an OH...O hydrogen bond that can be formed in
9 the cyclo(LTyr-LTyr) conformer only. The second aim of this work is thus to study the influence
10 of the perturbation due to the excess charge (proton). The question that should be answered
11 is whether the strong electrostatic interaction hides the chiral discrimination effect or if it is
12 still visible in the most stable conformer of the two diastereomers of protonated cyclo(Tyr-
13 Tyr), namely cyclo (LTyr-LTyr)H⁺ and cyclo (LTyr-DTyr)H⁺. They are denoted c-LLH⁺ and c-LDH⁺
14 in what follows and shown in Figure 1.
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20 **Material and methods**

21 **Experiment**

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24 Cyclo Tyr-Tyr(>99%) was purchased from Genecust (Luxembourg), and used without further
25 purification. The experimental set-up has already been described in detail.³⁹ Briefly, the
26 methanol solutions containing cyclo(Tyr-Tyr) at a concentration of 1.2×10^{-5} ML⁻¹ with 0.5% of
27 formic acid were electrosprayed from a nebuliser to generate fine droplets containing the
28 target molecule. The droplets were evaporated by travelling through a glass capillary heated
29 up to 70°C. The ions efficiently collected by an ion funnel were injected into a quadrupole
30 mass spectrometer (Q-MS: Extrel) via a hexapole ion guide. The target molecule was mass-
31 selected by the Q-MS and then introduced to a cryogenic quadrupole ion trap (QIT) via a
32 quadrupole ion bender and octopole ion guide. The QIT was kept at 4K by a closed-cycle two-
33 stage He cryostat (Sumitomo: RDK-408D2) during the whole experiment. A mixed buffer gas
34 of H₂ (20%) and helium was introduced to the QIT using a pulsed valve (Parker Hannifin:
35 General Valve Series 9). The mass-selected ions were allowed to collide with the cold buffer
36 gas so that they were trapped in the QIT and cooled down to ~10K. The collision-induced
37 dissociation (CID) experiments were performed in the QIT. The collision energy with He buffer
38 gas was increased by increasing the kinetic energy of the ions injected into the QIT. The CID
39 fragments were detected by a linear time-of-flight mass spectrometer (TOF-MS). For the IR
40 spectroscopy, H₂ molecules were allowed to weakly attach to the ions under such cryogenic
41 conditions.⁴⁰ The H₂-tagged ions were irradiated in the cold trap with a tunable IR laser
42 (LaserVision: OPO/OPA). Hydrogen molecules were detached from the cluster ions when the
43 wavenumber of the IR laser was resonant with a vibrational transition of the cluster ions.⁴¹⁻⁴²
44 The photofragments were extracted linearly and detected by the linear TOF-MS. The
45 photofragment ion signals from a dynode converter detector of TOFMS were amplified 10
46 times by a preamplifier and recorded on a fast digitizer. IR photodissociation (IRPD) spectra
47 were measured by monitoring the intensity of the photofragments (*m/z* 327) resulting from
48 H₂ loss as a function of the wavenumber of the IR laser. IR-UV double resonance spectroscopy
49 was performed to discuss the effect of H₂ tagging. UV photodissociation spectra were
50 measured by monitoring the photofragments resulting from photoexcitation by a UV dye laser
51 (Lumonics). IR-UV ion dip spectroscopy was not possible due to the broad nature of the UV
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bands. The IR-UV enhanced spectra was measured instead. In this method, the UV wavenumber was fixed to the absorption maximum of the broad UV band. The tunable IR laser illuminated the ions of interest ~50 ns prior to the UV irradiation. The photofragment yield was enhanced by opening an additional fragmentation channel from the states populated by the electronic transition from vibrationally-excited states populated by IR absorption. The IR-UV enhanced spectrum is totally free from tag molecules.

Theoretical method

The structures of the protonated species, cyclo(LTyr-LTyr)H⁺ and cyclo(LTyr-DTyr)H⁺ were optimised in the framework of the density functional theory (DFT) using the B3LYP-D3 functional and the 6-311++G(d,p) basis sets.⁴³⁻⁴⁵ All the calculations were performed with the Gaussian16 B01 software.⁴⁶ The frequencies were calculated at the same level of theory and scaled by 0.952 to account for anharmonicity and basis set incompleteness. The calculation level was chosen identical to that used for the study of the neutral molecule.¹⁹ The scaling factor was the same as for the neutral molecule.¹⁹ The potential energy surface was manually explored by building two types of initial structures in a systematic way. The first structures were built from the stable forms of protonated cyclo(PhePhe) modified by adding an hydroxyl substituent to the benzene ring, as it was successfully done before for finding the geometries of jet-cooled cyclo(Tyr-Phe) by removing an hydroxyl from neutral cyclo(Tyr-Tyr).^{28 47} To ensure that the exploration of the PES was complete, additional structures were obtained by adding a proton to the amide nitrogen or oxygen atom of the stable calculated structures of neutral cyclo(TyrTyr).¹⁹ The structures whose vibrational spectrum shows the best match with the experiment were re-optimised at the B3LYP-D3(BJ) level. D3(BJ) dispersion correction have indeed proved to yield better structures for systems containing two aromatic rings like cyclophane.⁴⁸ In what follows, the energetics and structures are those obtained at the B3LYP-D3 level. Those at the B3LYP-D3(BJ) level are very similar and we shall only emphasise the differences in calculated frequencies.

Results

1) Theoretical results

Nomenclature: The conformational diversity of protonated cyclo(Tyr-Tyr) arises first from the orientation of the substituents, which in turn makes the protonation sites (CO or NH) non-equivalent. As described for the neutral molecule, each aromatic ring can take three possible orientations: two *gauche* and one *trans* forms, denoted as *g*⁺, *g*⁻ and *t*, respectively. They correspond to dihedral angles $\tau(\text{NCCC})$ of around 60°, -60° and 180°, respectively, for the L residue, with opposite values for the D residue. The *g*⁺ geometry corresponds the aromatic substituent folded over the DKP ring while the *g*⁻ geometry corresponds to an aromatic ring extended outwards; these geometries are frequently encountered for neutral DKP peptides.^{4, 13} The *trans* geometry corresponds to an aromatic ring extended outwards with an aromatic plane almost perpendicular to the peptide ring (see Figure 3). Moreover, the phenol OH can adopt two orientations that cannot be distinguished by their IR spectra.^{12, 19, 33, 39, 49-50} We will therefore think in terms of “families” including both orientations of the OH bond.

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3 The nomenclature used here rests on the geometry of the substituents and the nature of the
4 protonation site, in the following way: each structure will be called c-LLH⁺(AB_i) or c-LDH⁺(AB_i)
5 with A and B denoting the g⁺, g⁻ or t geometry and i the protonation site, O or N, denoted in
6 subscript of the corresponding residue.
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8 **Calculated structures:** Four protonation sites are identified for c-LLH⁺ and LDH⁺, namely the
9 CO and NH groups of each amide function. Protonation on NH results in conformers much
10 higher in energy than protonation on CO. As shown for the related cyclo(Phe-Phe) protonated
11 system, this difference is of the order of ~50 kJ mol⁻¹. We shall therefore limit the discussion
12 to the species protonated on the amide CO. The energetic data are collected in Table 1.
13 Because the distribution of conformers probably reflects a temperature between that of the
14 electrospray and that of the ion trap, Table 1 gives the Gibbs free energies ΔG at room
15 temperature (298 K) and at 350 K, together with the ZPE-corrected energies at 0 K. Unless
16 specified otherwise, the energetic values given in what follows are ΔG at 298 K. All the
17 calculated structures and energetic data, including those protonated on the amide NH, are
18 shown in Figures S1 and S2 of the Supporting Information (SI). Adding two hydroxyl
19 substituents to the most stable protonated structures of cyclo(LPhe-LPhe) or cyclo(LPhe-
20 DPhe) results to c-LLH⁺(g⁺t₀) and c-LDH⁺(g⁺t₀), which are the most stable structures (see Table
21 1 and Figure 3). In these structures, one of the phenol ring is folded over the peptide ring,
22 which allows a weak C_βH...π interaction, as observed in one of the stable neutral forms. The
23 other ring is extended in *trans* geometry, with the aromatic plane perpendicular to the DKP
24 ring, which results in a stabilising interaction between the OH⁺ and the folded aromatic ring.
25 The strength of the OH...π interaction is different in c-LLH⁺ and c-LDH⁺, as attested by the
26 length of the OH bond, which is slightly shorter in c-LL (0.985 Å) than in c-LD (0.990 Å),
27 reflecting a slightly longer distance between the proton and the centre of the aromatic ring
28 (2.51 Å for c-LLH⁺(g⁺t₀) vs. 2.48 Å for c-LDH⁺(g⁺t₀)). The second most stable structures derived
29 from c-LPheLPheH⁺ and c-LPheDPheH⁺ are c-LLH⁺(g⁻t₀) and c-LDH⁺(g⁻t₀) and are calculated at
30 9.5 and 8.1 kJ mol⁻¹ relative to the most stable form, respectively. Decreasing the temperature
31 disfavours the g⁻t₀ conformers. In terms of structure, they are fully extended and involve a
32 weak NH...π interaction. Both families of structures are stabilised by the interaction between
33 the OH⁺ and the aromatic ring, which is facilitated by the *trans* position of the latter. The last
34 two structures deduced from protonated c-LPheLPhe and c-LPheDPhe are c-LLH⁺(g⁻g⁺t₀) and c-
35 LLH⁺(g₀⁻g⁺) for c-LL and the equivalent diastereomer structures c-LD(g⁻g⁺t₀), c-LDH⁺(g₀⁻g⁺).
36 These structures have one ring folded and the other one extended. They are identical to those
37 obtained by adding a proton to the folded/extended conformer of the neutral form. They are
38 less stable by ~25 -28 kJ mol⁻¹ than the most stable form because the *gauche* geometry of the
39 extended ring prevent the OH...π to be formed. Last, adding a proton to the stacked conformer
40 of neutral c-LL g₀⁺g⁺ result in a structure energetically disfavoured by ~ 30 kJ mol⁻¹.
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54 2) Collision-induced dissociation

55 The CID MS² spectra of c-LLH⁺ and c-LD H⁺ obtained in identical conditions are shown in Figure
56 2. The same fragmentation pattern is observed for the two diastereomers. CID is usually less
57 sensitive to the structure of epimer peptides than other dissociation methods such as
58 electron-capture⁵¹ or radical-induced dissociation.⁵² However, the result obtained here
59 contrasts to the stereoselectivity observed in histidine-containing cyclic dipeptides.^{25, 53}
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Besides the intense parent peak at m/z 327, two main fragments appear at m/z 136 and m/z 107. These fragments, together with the weak m/z 91 peak, are characteristic of tyrosine-containing peptides.⁵⁴ The m/z 136 fragment corresponds to the immonium ion, frequently observed in CID or UV-induced dissociation of charged peptides.⁵⁴⁻⁵⁶ The immonium ion of tyrosine was also observed in the CID spectrum of cyclo(Tyr-Pro)H⁺.²⁶ The fragments at m/z 107 arises from the immonium ion⁵⁴ and corresponds to the C_αC_β cleavage resulting in the hydroxybenzyl radical cation CH₂C₆H₄OH⁺. A minor peak at m/z 300 corresponds to HCN loss. A weak peak at m/z 221 is assigned to the protonated cyclo(Gly-Tyr) cation resulting from the C_αC_β cleavage (loss of p-quinomethane). The simultaneous observation of m/z 107 and m/z 221 suggests that the C_αC_β cleavage results in the possible location of the positive charge on either of the fragments. It should be noted here that the fragment resulting from CO loss (m/z 299) is indiscernible in the MS² spectrum, which contrasts to related systems like cyclo(Phe-Phe)H⁺ or cyclo(Tyr-Pro)H⁺ for which CO loss appears at low collision energy and is the dominant CID fragment in all the collision energy range studied. Intense peaks corresponding to sequential fragmentation after CO loss are also observed in the CID spectrum of cyclo(Tyr-Pro)H⁺, whose counterpart is completely absent here. The effect of stereochemistry is very limited in the system studied here. The CID efficiency is only slightly higher for c-LDH⁺, but with no modification of the fragmentation pattern.

3) Spectroscopic results and assignment

Experimental spectra: In what follows, we will discuss the IRPD spectra obtained by the messenger technique. The spectra of the two diastereomers are shown in Figure 4 and clearly differ from each other. They can be divided in three regions. The higher-energy region contains a single band at 3633 cm⁻¹ and 3628 cm⁻¹ for c-LLH⁺ and c-LDH⁺ respectively, assigned to the free phenol ν(OH), whose frequency is almost identical for the two diastereomers. The 3300-3500 cm⁻¹ region contains two bands characteristic of the amide ν(NH). They appear at different frequencies for c-LLH⁺ (3417 cm⁻¹ and 3348 cm⁻¹) and c-LDH⁺ (3412 cm⁻¹ and 3375 cm⁻¹). Lastly, the low-energy region shows a broad and intense band characteristic of the protonated amide ν(OH⁺), which appears a different frequencies for c-LLH⁺ and c-LDH⁺ (3190 cm⁻¹ and 3079 cm⁻¹, respectively).

Assignment: The spectra are well accounted for by that simulated for the most stable species, c-LLH⁺(g⁺t₀) and c-LDH⁺(g⁺t₀). The free phenol ν(OH) is mostly independent of the diastereomer. It shows a ~10 cm⁻¹ down shift relative to the neutral molecule, which is due to the effect of H₂ tagging (see below). A similar frequency was observed in protonated tyrosine or cyclo(Tyr-Pro).^{18, 26, 50} The ν(NH) region is well reproduced too by the simulations. Although the relative intensity of the two ν(NH) transitions is not quantitatively reproduced, the simulations account for the weaker intensity of the ν(NH) localised on the neutral amide and their similar frequencies for the two diastereomers. They also account for a larger down shift of the ν(NH) localised on the protonated amide in the case of c-LLH⁺. The same difference between c-LLH⁺ and c-LDH⁺ was observed in protonated cyclo(Phe-Phe). It should be mentioned here that both frequencies and absorption intensities are identical whether the D3 or D3(BJ) correction is used. Also the bound ν(OH⁺) frequency is calculated at different frequencies for c-LLH⁺ and c-LDH⁺, reproducing the lower frequency observed for c-LDH⁺,

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3 which in turns reflect a stronger $\text{OH}^+\dots\pi$ interaction. The downshift of the bound $\nu(\text{OH}^+)$
4 frequency is underestimated for both diastereomers, which is a deficiency of the currently
5 used method for modes involved in strong hydrogen bonds.⁵⁷⁻⁵⁸ However, the D3(BJ)
6 correction performs slightly better in this respect by shifting the calculated value further down
7 in energy, thus closer to the experimental value. With both D3 and D3(BJ) corrections, the
8 smaller red shift of $\nu(\text{OH})$ in $c\text{-LLH}^+$ relative to $c\text{-LDH}^+$ correlates well with the strength of the
9 $\text{OH}^+\dots\pi$ interaction, as reflected by the distance between the proton and the centre of the
10 aromatic ring mentioned above.
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14 **Effect of H_2 tagging:** To assess the effect of the H_2 tag on the vibrational spectrum, the IRPD
15 spectra are compared to the corresponding IR-UV enhanced spectra. To this end, the UV
16 spectrum of $c\text{-LLH}^+$ and $c\text{-LDH}^+$ were recorded and are shown in Figure S3 of the SI. Despite
17 the low temperature of the trap, they are broad and do not show any structure. This strongly
18 contrasts to the behaviour of protonated tyrosine, the UV spectrum of which is well structured
19 under ion trap conditions.^{37, 39, 59} This result points towards fast deactivation processes in the
20 cyclic dipeptide relative to the protonated amino acid. Figure S4 shows the IR-UV enhanced
21 spectrum recorded by setting the UV wavelength at the maximum of the broad absorption.
22 The band shifts induced by H_2 tagging are smaller than 10 cm^{-1} for the amide $\nu(\text{NH})$ stretch,
23 but the shifts are slightly larger for the phenol $\nu(\text{OH})$ stretch ($11\text{-}13\text{ cm}^{-1}$). The largest shift
24 ($\sim 30\text{ cm}^{-1}$) is observed for the protonated $\nu(\text{OH}^+)$. However, the spectral patterns themselves
25 did not change significantly, in agreement with recent results suggesting that H_2 tagging only
26 perturbs the spectrum for highly symmetrical systems such as CH_5^+ .⁶⁰ It should be noted that
27 the $\nu(\text{OH}^+)$ is up-shifted by the H_2 tag, in contrast to the other modes that are down shifted.
28 Moreover, the two phenolic $\nu(\text{OH})$ undergo similar upshift upon tagging, as do the two amide
29 $\nu(\text{CH})$. These observations suggest the presence of an H_2 molecule on each aromatic ring in
30 the tagged species. This hypothesis has been checked by measuring the mass spectrum of $c\text{-LLH}^+$
31 in the absence of the IR laser. It is shown in Figure S5 of the Supplementary Information.
32 A strong peak appears at m/z 331, assigned to $c\text{-LLH}^+(\text{H}_2)_2$.
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35 A very weak band is experimentally observed at 3551 cm^{-1} for both diastereomers. In addition,
36 a very weak band appears for $c\text{-LDH}^+$ at $\sim 3340\text{ cm}^{-1}$. These bands could be due to an additional
37 conformer kinetically trapped in our experimental conditions. The calculated spectrum of $c\text{-LDH}^+(\text{g}^-\text{g}^+\text{o})$
38 accounts well for the additional bands observed in the spectrum of $c\text{-LDH}^+$, with a
39 band at 3551 cm^{-1} assigned to the free $\nu(\text{OH}^+)$.²⁸ The band at 3340 cm^{-1} would correspond to
40 $\nu(\text{NH})$, slightly shifted down in frequency due to the $\text{NH}\dots\pi$ interaction. However, $c\text{-LD}(\text{g}^-\text{g}^+\text{o})$
41 is not the most stable $c\text{-LD}(\text{g}^-\text{g}^+)$ structure, and is 31.7 kJ mol^{-1} higher than $c\text{-LD}(\text{g}^-\text{o}\text{g}^+)$. The
42 spectrum of $c\text{-LLH}^+$ only shows the additional band at 3551 cm^{-1} , which could be due to $c\text{-LL}(\text{g}^-\text{g}^+\text{o})$
43 or $c\text{-LL}(\text{g}^-\text{o}\text{g}^+)$. These two structures have relative Gibbs energy of $+25.3\text{ kJ mol}^{-1}$ for $c\text{-LL}(\text{g}^-\text{o}\text{g}^+)$
44 and $+29.5\text{ kJ mol}^{-1}$ for $c\text{-LL}(\text{g}^-\text{g}^+\text{o})$. Although the corresponding structures are among the
45 most stable for cyclo (Phe-Phe) H^+ and contribute to its IRMPD spectrum at room temperature,
46 their contribution here is expected to be very limited due to the low temperature of the trap.
47 We therefore conclude that these bands are due to species complexed by more than two H_2
48 tags. The weak band at 3551 cm^{-1} could be due to a species in which an additional H_2 molecule
49 attaches to one of the phenol OH. The band at 3340 cm^{-1} could be due to a species in which
50 an additional H_2 molecule attaches to one of the amide NH.
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3 The most stable structures of c-LLH⁺ and c-LDH⁺ are similar to those calculated for protonated
4 cyclo(Phe-Phe), and so are the spectra simulated at the same level of calculations. The $\nu(\text{OH}^+)$
5 frequency is shifted down by 110 cm⁻¹ in c-LDH⁺ and appears at 3079 cm⁻¹ for LD vs. 3190 cm⁻¹
6 for LL. This difference in frequency was calculated also in cyclo(Phe-Phe)H⁺ but was difficult to
7 observe experimentally due to the experimental conditions at room temperature resulting in
8 broad transitions.
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11 **Conclusion**

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13 The structure responsible for the experimentally observed spectrum is the most stable
14 calculated form, for c-LLH⁺ and c-LDH⁺ alike. The most stable c-LLH⁺ structure is more stable
15 than the most stable c-LDH⁺ by 1.77 kJ mol⁻¹. Although very small, probably within the error
16 of the calculation, this energy difference parallels the CID efficiency, which is slightly larger for
17 c-LDH⁺. The effect of the charge is dominant in these systems and protonation induces a
18 dramatic change in the structure, preventing the formation of the stacked geometry observed
19 in the neutral molecule for c-LL. However, the chiral selection overcomes the strong
20 perturbation induced by the excess charge and a clear difference is observed in the IRPD
21 spectra of the two diastereomers. The cold temperature achieved here minimises the
22 presence of higher-energy conformers in contrast to the room temperature experiments⁶¹
23 and makes it possible to evidence the differences in the $\nu(\text{OH}^+)$ and protonated amide $\nu(\text{NH})$
24 stretch frequencies between the two diastereomers. These results confirm that the chirality
25 effects are visible in the most stable structure of protonated cyclic peptides.
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32 **Supporting Information.**

33 Calculated structures and simulated spectra for the structures not discussed in the main
34 text. UVPD spectra. Comparison between H₂ tagging and IR-UV enhanced spectroscopy
35 results.
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38 **Acknowledgments**

39
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Structure name	$\Delta G/ \text{KJ mol}^{-1}$ @298K	$\Delta G/ \text{KJ mol}^{-1}$ @350K	ZPE-corrected $\Delta G/ \text{KJ mol}^{-1}$	Structure name	$\Delta G/ \text{KJ mol}^{-1}$ @298K	$\Delta G/ \text{KJ mol}^{-1}$ @350K	ZPE-corrected $\Delta G/ \text{KJ mol}^{-1}$
c-LL(g ⁺ t ₀)	0.0	0.0	0.0	c-LD(g ⁺ t ₀)	0.0	0.0	0.0
c-LL(g ⁻ t ₀)	9.5	8.5	14.4	c-LD(g ⁻ t ₀)	8.1	6.0	18.7
c-LL(g ⁻ ₀ g ⁺)	25.3	25.2	25.7	c-LD(g ⁻ g ⁺ ₀)	20.3	19.3	25.4
c-LL(g ⁻ t ₀)	28.0	26.5	35.6	c-LD(g ⁻ ₀ g ⁻)	23.3	21.0	34.5
c-LL(g ⁻ ₀ t)	28.0	26.5	35.6	c-LD(g ⁻ g ⁻ ₀)	24.9	22.6	35.8
c-LL(g ⁺ g ⁺ ₀)	29.5	30.3	25.5	c-LD(g ⁻ ₀ g ⁻)bis	28.2	26.0	38.5
c-LL(g ⁻ g ₀ ⁺)	31.4	31.4	30.9	c-LD(g ⁻ g ⁻ ₀)bis	29.3	27.0	40.0
c-LL(g ⁻ ₀ t)	33.1	31.6	40.7	c-LD(g ⁻ g ⁺ ₀)	31.7	30.7	36.5
c-LL(g ⁻ ₀ g ⁻)	40.5	39.2	46.7	c-LD(g ⁻ ₀ t)	34.9	32.7	45.5
c-LL(g ⁺ g ⁺ _N)	52.1	52.7	48.7	c-LD(g ⁻ t) bis	46.7	44.0	60.0
c-LL(g ⁻ g ⁺ _N)	69.1	69.0	69.3	c-LD(g ⁻ _N g ⁺)	51.4	49.6	59.8
c-LL(g ⁻ g ⁺ _N)	73.9	73.5	76.1	c-LD(g ⁻ t) ter	55.7	53.2	67.3
				c-LD(g ⁻ g ⁺ _N)	68.9	67.0	78.7

Table 1: Energetic data for the most stable calculated structures of c-LLH⁺ and c-LDH⁺

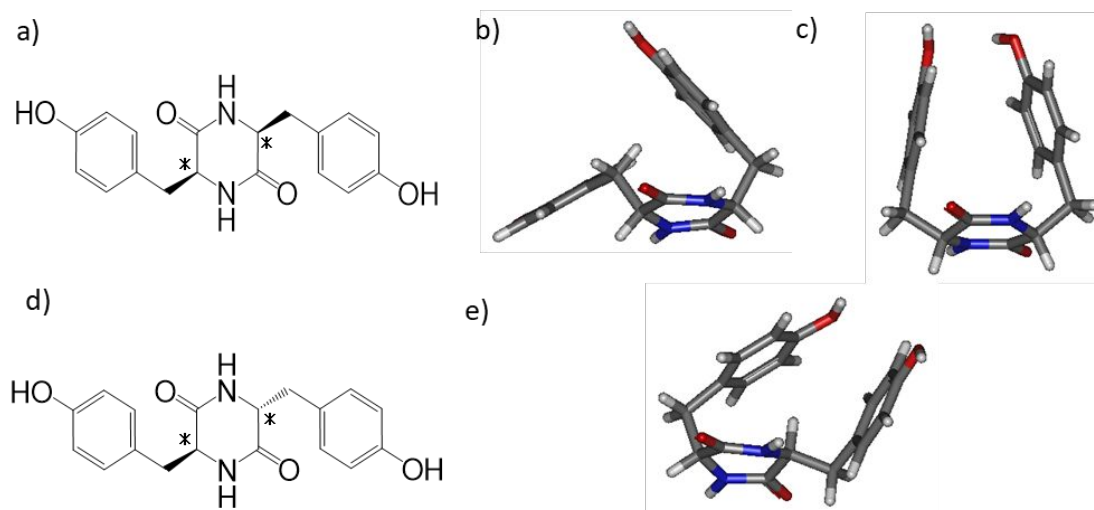


Figure 1: Scheme of the studied molecules and structures observed in the gas phase. The chiral centres are indicated by *. a) Schematic view of the homochiral neutral dipeptide c-LL. b) and c) observed structures for neutral c-LL. d) Schematic view of the heterochiral neutral dipeptide c-LD. e) observed structure for neutral c-LD. Note that the stacked structure (c) is only observed for c-LL.

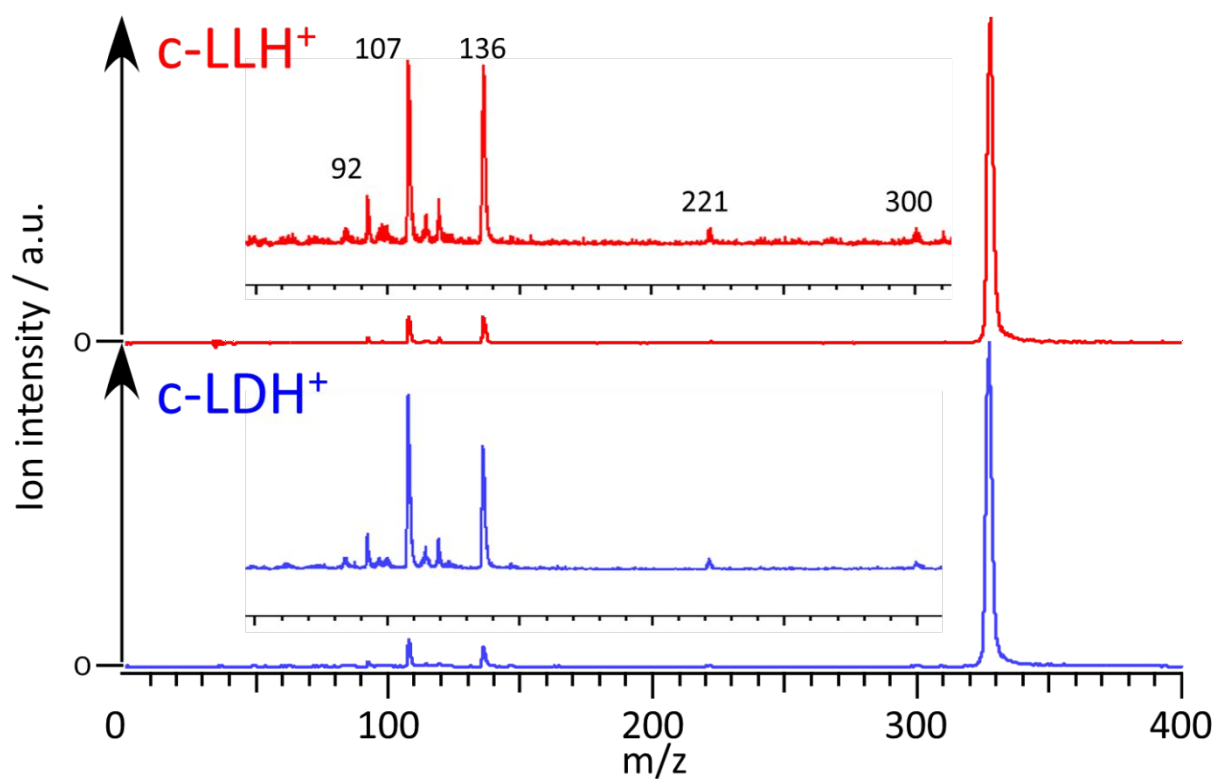


Figure 2: MS² spectrum obtained by collision-induced dissociation of the parent m/z 327 for the protonated diastereoisomers of cyclo(Tyr-Tyr) cLLH⁺ (top) and c-LDH⁺ (bottom).

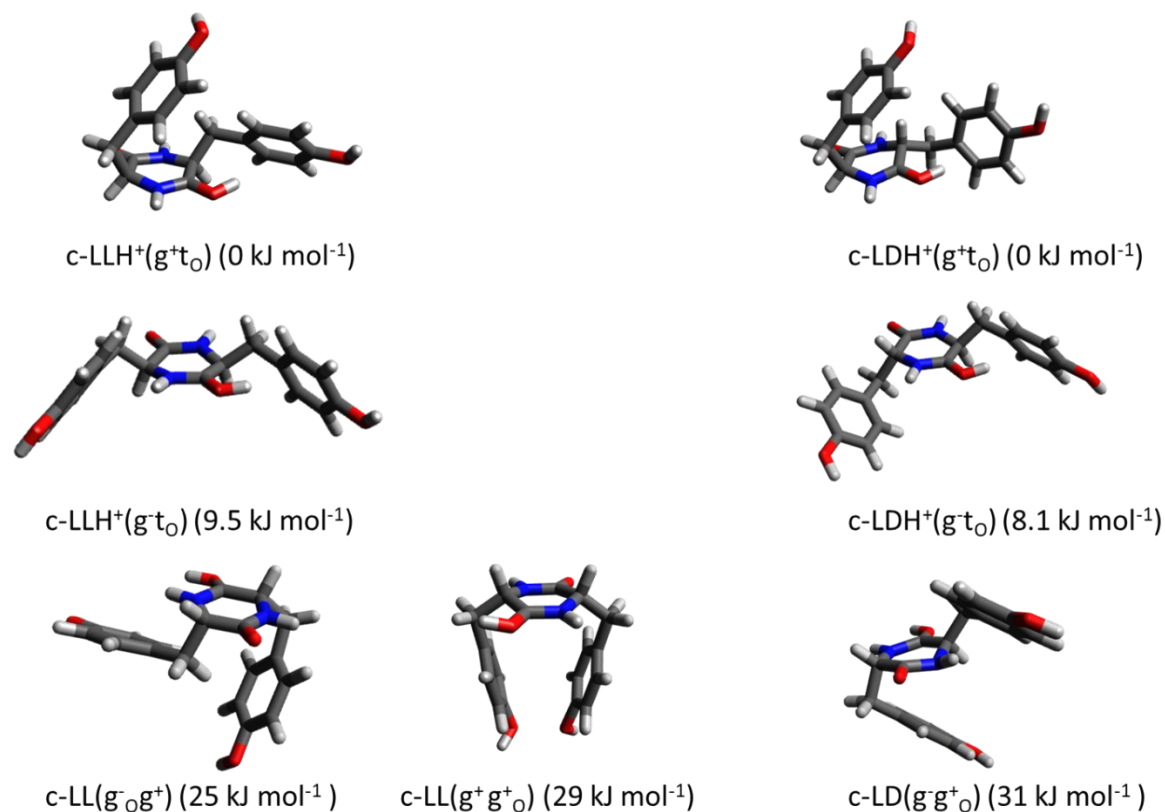


Figure 3: Most stable calculated structures of the protonated diastereoisomers of cyclo(Tyr-Tyr) $c\text{-LLH}^+$ (left) and $c\text{-LDH}^+$ (right).

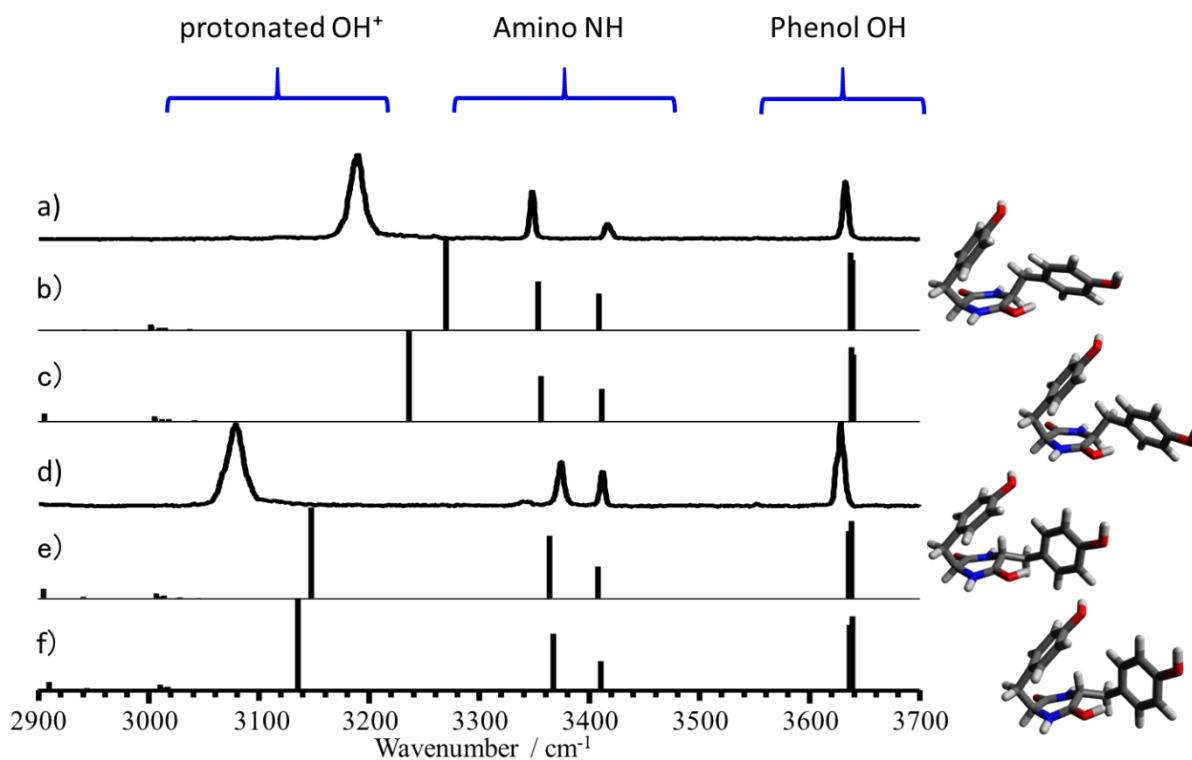
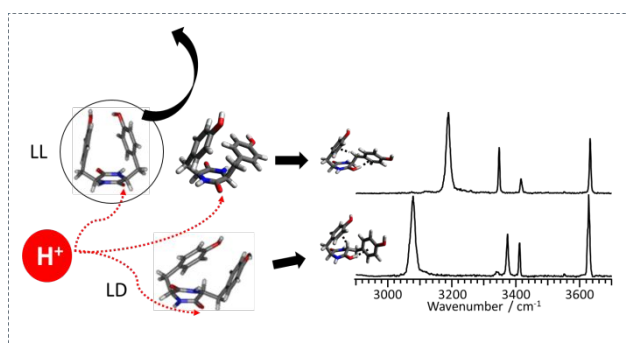


Figure 4: Comparison between the experimental spectra of (a) c-LLH⁺ and (d) c-LDH⁺ and those calculated for the most stable structures: (b) c-LLH⁺(g⁺t₀) at the B3LYP-D3/6-311++g(d,p) level of theory. (c) Same with D3BJ dispersion correction. (e) c-LDH⁺(g⁺t₀) at the B3LYP-D3/6-311++g(d,p) level of theory. (f) Same with D3BJ dispersion correction. The frequencies have been scaled by 0.952 (see text).



Graphical Abstract for TOC

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