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## Raman spectroscopy and DFT modelling of Tl<sub>2</sub>S-GeS<sub>2</sub> crystals and glasses

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## Keywords

Thallium thiogermanate glasses, Raman spectroscopy, DFT modelling

## Abstract

Thallium thiogermanate glasses and crystals were studied by Raman spectroscopy accompanied by DFT modelling. The experimental Raman spectra of the known crystalline compounds in the Tl<sub>2</sub>S-GeS<sub>2</sub> system are consistent with the crystal structure and DFT modelling of the respective Tl-Ge-S clusters, representing structural motifs in c-Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub>, c-Tl<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub>, and c-Tl<sub>4</sub>GeS<sub>4</sub>. The Raman spectra of the (Tl<sub>2</sub>S)<sub>*x*</sub>(GeS<sub>2</sub>)<sub>1-*x*</sub> glasses can be divided into two groups: (1) semiconducting glasses at  $x \le 0.3$ , and (2) ion-conducting vitreous alloys at x > 0.3. At  $x \le 0.3$ , the continuous g-GeS<sub>2</sub> network is transformed into a fragmented structure consisting of adamantane-like units Ge<sub>4</sub>S<sub>10</sub> held together by thallium species. In the ion-conducting composition domain at x > 0.3 takes place further fragmentation of tetrahedral tetramers Ge<sub>4</sub>S<sub>10</sub> into edge-sharing ES-Ge<sub>2</sub>S<sub>6</sub> dimers and small corner-sharing CS-oligomers, i.e., CS-Ge<sub>2</sub>S<sub>7</sub> dimers or CS-Ge<sub>3</sub>S<sub>10</sub> trimers.

#### 1. Introduction

Heavy metal chalcogenide glasses appear to be promising optical materials transparent in far-IR [1-2], highly sensitive membranes for heavy metal ion detection [3-6] and interesting model materials exhibiting heavy ion transport in solids [7-13]. Structural studies of lead, cadmium, mercury or thallium chalcogenide glasses are rather rare, except for a few reports [14-20]. In the  $(Tl_2S)_x(GeS_2)_{1-x}$  system, glasses are formed over the entire composition region at  $x \le 0.5$  [13]. It should, however, be noted that thallium-poor glasses,  $x \le 0.15$ , show small Bragg peaks in both large neutron and small X-ray samples, which correspond to hightemperature  $\beta$ -GeS<sub>2</sub>. A spatial incompatibility between 2D network structure of glassy GeS<sub>2</sub> and isolated tetrahedral Ge<sub>4</sub>S<sub>10</sub> clusters (Fig. 1) seems to be the reason of enhanced crystallization ability of the Tl-poor glasses.

Three ternary compounds are known in the  $(Tl_2S)_x(GeS_2)_{1-x}$  system [21-23] in addition to two binary end-members GeS<sub>2</sub> and Tl<sub>2</sub>S (Fig. 1). In a remarkable contrast to low-temperature 3D lattice or high-temperature 2D lattice, Fig. 1(d), of the  $\alpha$ - and  $\beta$ -GeS<sub>2</sub> polymorphs [24-25], crystalline thallium thiogermanates exhibit isolated Ge-S anions held together by thallium species. Thallium-poor monoclinic Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub> ( $x = \frac{1}{3}$ ), space group C2/c, crystallizes in adamantane-type structure consisting of highly symmetric tetrahedral units Ge<sub>4</sub>S<sub>10</sub> ( $T_d$ symmetry), Fig. 1(c) [21]. The corner-sharing CS-GeSS<sub>3/2</sub> tetrahedra in the adamantane units share their corners and exhibit two types of Ge-S distances for bridging S<sub>b</sub> and terminal S<sub>t</sub> sulfur species:  $r(Ge - S_b) = 2.26\pm0.05$  Å and  $r(Ge - S_t) = 2.141\pm0.012$  Å. The average Ge-Ge distance between neighboring CS-GeSS<sub>3/2</sub> is 3.58±0.02 Å.

The adamantane-like structural motif is widely adopted by analogous alkali thio- and selenogermanates and silicates (see, for example, [26-28] and references therein). In equimolar compound, triclinic Tl<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub> ( $x = \frac{1}{2}$ ), space group  $P\overline{1}$ , isolated edge-sharing ES-Ge<sub>2</sub>S<sub>6</sub> dimers are found, Fig. 1(b). The Ge-S distances also show a difference between S<sub>b</sub> and S<sub>t</sub> species:  $r(Ge - S_b) = 2.274\pm0.003$  Å and  $r(Ge - S_t) = 2.16\pm0.02$  Å. The Ge-Ge distance in the ES-Ge<sub>2</sub>S<sub>6</sub> dimer is 3.074\pm0.003 Å. Finally, in thallium-rich monoclinic Tl<sub>4</sub>GeS<sub>4</sub> ( $x = \frac{2}{3}$ ), space group Cc, isolated tetrahedra GeS<sub>4</sub> are held together by thallium ions, Fig. 1(a). The average Ge-S interatomic distance is  $r(Ge - S_t) = 2.216\pm0.025$  Å. In all three Tl<sub>2</sub>S-GeS<sub>2</sub> crystalline compounds the thallium coordination is high, 6 (Tl<sub>4</sub>GeS<sub>4</sub>)  $\leq N_{Tl-S} \leq 9$  (Tl<sub>4</sub>Ge4S<sub>10</sub>), and decreases with x. The average Tl-S interatomic distance also depends on thallium content x: 3.44\pm0.27 Å ( $x = \frac{1}{3}$ )  $\rightarrow 3.11\pm0.07$  Å ( $x = \frac{1}{2}$ )  $\rightarrow 3.02\pm0.10$  Å ( $x = \frac{2}{3}$ ).



Figure 1. Structural motifs in the Tl<sub>2</sub>S-GeS<sub>2</sub> system [21-23]: (a) thallium-rich monoclinic Tl<sub>4</sub>GeS<sub>4</sub> ( $x = \frac{2}{3}$ ), space group *Cc*, (b) equimolar triclinic Tl<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub> ( $x = \frac{1}{2}$ ), space group *P* $\overline{1}$ , and (c) thallium-poor monoclinic Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub> ( $x = \frac{1}{3}$ ), space group *C2/c*. (d) Characteristic layer structure in high-temperature monoclinic modification  $\beta$ -GeS<sub>2</sub>, space group *P2*<sub>1</sub>/*c* [24]. Thallium, germanium and sulfur atoms are presented in pink, green and yellow, respectively.

The goal of the present paper is to unveil the local environment in thallium thiogermanate glasses  $(Tl_2S)_x(GeS_2)_{1-x}$ ,  $0 \le x \le 0.5$ , compare it with known crystalline references and reveal a difference between semiconducting,  $x \le 0.3$ , and ion-conducting, x > 0.3, vitreous alloys [13]. As a first step, we report here Raman spectroscopy results for the  $Tl_2S$ -GeS<sub>2</sub> glasses and crystals accompanied by DFT modelling of the vibrational spectra. Further steps including pulsed neutron diffraction and high-energy X-ray scattering will be reported elsewhere.

## 2. Material and methods

## 2.1. Samples preparation

The glassy samples in the  $(Tl_2S)_x(GeS_2)_{100-x}$  system, where  $x = 0, 0.05, 0.1, 0.15, 0.2, 0.3, 04, 0.5 Tl_2S$  fraction were prepared from high purity elements by the standard melt-quenching technique. Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub>, Tl<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub> and Tl<sub>4</sub>GeS<sub>4</sub> crystals were prepared from its constituents by thermal synthesis in the evacuated silica tube according to the procedure described previously by Eulenberger [21-23]. The details for the glass and crystals synthesis as well as the glass forming region, macroscopic and transport properties of the solids are described in our previous paper [13].

## 2.2. Raman spectroscopy measurements, Raman data treatment and DFT modelling

The Raman spectra of the glasses were collected on a LABRAM Dilor spectrometer (Jobin Yvon Horiba Group) equipped with a liquid nitrogen cooled CCD detector and a microscope. He–Ne laser operating at an excitation wavelength of 632.8 nm was used. Each spectrum was recorded at room temperature in the 80-1200 cm<sup>-1</sup> spectral range using a power level of 0.1 mW (to avoid crystallisation of the glassy samples) and a resolution of 1 cm<sup>-1</sup>. The acquisition time varied between 60 and 150 s. In order to verify the sample homogeneity and the absence of photo-induced phenomena, two to four spectra were registered for each sample at different positions. The measured raw Raman spectra were analysed over the 100–600 cm<sup>-1</sup> range. The spectral background was approximated by a Voigt function and then subtracted from the experimental data with further normalization to the most intense feature. Typical raw Raman spectra treatment is shown in Fig. S1 (supplementary material).

The DFT calculations have been carried out using GAUSSIAN<sup>®</sup> 16 [29] software associated with its graphical user interface GaussView<sup>®</sup> 6.1. Optimisation of the DFT frequency calculations was carried out with the Becke [30] three parameters hybrid exchange functional and the Lee–Yang–Parr correlation functional (B3LYP) [31] associated with 6-311++G(3df,2p) basis-set. In order to find a compromise between the cost of the calculations and the accuracy of the results, structural optimisations and harmonic vibrational frequency calculations were performed for isolated size-limited clusters: GeS<sub>4</sub>Tl<sub>4</sub>, Ge<sub>2</sub>S<sub>6</sub>Tl<sub>4</sub>, Ge<sub>4</sub>S<sub>10</sub>Tl<sub>4</sub>, Ge<sub>3</sub>S<sub>10</sub>Tl<sub>8</sub>, Ge<sub>3</sub>S<sub>9</sub>Tl<sub>6</sub>, corner-sharing CS-Ge<sub>3</sub>S<sub>10</sub>Tl<sub>8</sub> trimer and CS-Ge<sub>3</sub>S<sub>9</sub>Tl<sub>6</sub> ring. All the structures were optimized using the tight convergence option ensuring adequate convergence and reliability of computed wavenumbers. Further details of the DFT simulations are published elsewhere [20] [32].

## 3. Results and discussion

#### 3.1. Experimental and simulated Raman spectra in the Tl<sub>2</sub>S-GeS<sub>2</sub> crystals

The experimental Raman spectra of the synthesized crystalline compounds in the  $Tl_2S$ -GeS<sub>2</sub> system are shown in Fig. 2(a). We should note that neither IR nor Raman spectra of crystalline thallium thiogermanates were reported. Consequently, we can only compare the

obtained spectra with similar alkali thiogermanate crystals and carry out a DFT modelling of Tl-Ge-S clusters representing the structural motifs in the Tl<sub>2</sub>S-GeS<sub>2</sub> system (Fig. 3).





Figure 2. (a) Experimental Raman spectra of the crystalline reference compounds  $Tl_4Ge_4S_{10}$  ( $x = \frac{1}{3}$ ),  $Tl_4Ge_2S_6$  ( $x = \frac{1}{2}$ ), and  $Tl_4GeS_4$  ( $x = \frac{2}{3}$ ) within the Ge-S stretching spectral domain, and (b) DFT Raman spectra of the respective Tl-Ge-S clusters.

Figure 3. Optimized geometry of the Tl-Ge-S clusters: (a)  $GeS_4Tl_4$ , (b)  $Ge_2S_6Tl_4$ , and (c)  $Ge_4S_{10}Tl_4$ , used for DFT modelling of vibrational spectra of the crystalline reference compounds  $Tl_4GeS_4$ ,  $Tl_4Ge_2S_6$ , and  $Tl_4Ge_4S_{10}$ , respectively.

The optimized geometry of  $GeS_4Tl_4$ ,  $Ge_2S_6Tl_4$  and  $Ge_4S_{10}Tl_4$  clusters (interatomic distances and bond angles) are summarized in Table 1 and compared with crystallographic results [21-23]. The DFT simulations show a good agreement with the crystal structure.

Table 1. Optimized geometry parameters (Ge-S and Tl-S interatomic distances, S-Ge-S, Ge-S-Ge, S-Tl-S and Tl-S-Tl bond angles) of the Tl-Ge-S clusters used for DFT modelling of vibrational spectra in comparison with crystallographic results [21-23].

	Ge-S (Å)	S-Ge-S (deg)	Ge-S-Ge (deg)	Tl-S (Å)	S-Tl-S (deg)	Tl-S-Tl (deg)
GeS <sub>4</sub> Tl <sub>4</sub>	2.241	110(6)	-	2.956	72.6	90.9
c-Tl <sub>4</sub> GeS <sub>4</sub> [23]	2.22(3)	109(2)	-	3.02(10) <sup>a</sup>	73.4(7)	119(14)
$Ge_2S_6Tl_4$	2.25(2)	110(9)	84.4	2.938	72.0	85.4
c-Tl <sub>4</sub> Ge <sub>2</sub> S <sub>6</sub> [22]	2.22(5)	109(7)	85.0	3.11(7) <sup>a</sup>	75(2)	130(31)
$Ge_4S_{10}Tl_4$	2.25(7)	109(5)	108(3)	2.826	-	-
c-Tl <sub>4</sub> Ge <sub>4</sub> S <sub>10</sub> [21]	2.22(7)	109(3)	107(2)	3.07(4) <sup>a</sup>	72.0(10)	83.8(4)

<sup>a</sup> The shortest Tl-S interatomic distances. <sup>b</sup> The last digit(s) in parentheses corresponds to the mean-square deviation (MSD) of the average experimental or calculated values. The missing MSD for the calculated geometry parameters means either a single calculated value or a negligible difference between several nearly identical geometry parameters.

#### Tl4Ge4S10

The Raman spectrum of the thallium-poor crystal within the Ge-S stretching domain ( $\approx 300 \le v \le \approx 450 \text{ cm}^{-1}$ ) exhibits two intense vibrational features at 339 and 447 cm<sup>-1</sup>. Similar spectra were observed for alkali thiogermanate crystals M<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub> (M = K, Rb, Cs) [28]; the intense features were attributed to the  $A_1$  vibration modes representing Ge-S stretching of bridging and terminal sulfur species, respectively. Our DFT modelling reproduces rather well the Raman features of Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub>, Fig. 2(b). The positions of the two most intense  $A_1$  modes are consistent with experimental results. The 339 cm<sup>-1</sup> mode corresponds to the in-phase breathing of the Ge<sub>4</sub>S<sub>10</sub> tetrahedral cluster. The high-frequency feature at 447 cm<sup>-1</sup> is related to the in-phase symmetric stretching of the terminal sulfur species. The apparent discrepancy is observed for the amplitudes of multiple asymmetric stretching modes while their positions reflect the experimental features. It is worth mentioning that the calculated DFT intensities are similar to those in heavy alkali analogues of the thallium compound, i.e., Cs<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub> or Rb<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub> [28]. As expected, the TI-related vibrations appear at low frequencies, v < 100 cm<sup>-1</sup>.

#### Tl4Ge2S6

The agreement between DFT and experimental Raman spectra is better for equimolar  $Tl_4Ge_2S_6$  in both peak positions and amplitudes. The most intense feature at 370 cm<sup>-1</sup> corresponds to the in-phase breathing of the edge-sharing ES-Ge<sub>2</sub>S<sub>6</sub> dimer of the  $A_1$  symmetry. The high-frequency mode at 409 cm<sup>-1</sup> is related to the  $B_1$  in-phase asymmetric stretching of bridging and terminal sulfur atoms. The bending and deformation Ge-S vibrations are observed below 200 cm<sup>-1</sup>, similar to those in adamantane units. The Tl-related features appear below 100 cm<sup>-1</sup>, as in Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub>. There is no available Raman data for similar alkali thiogermanates. The only exception appears to be monoclinic Na<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub> with different crystal structure consisting of infinite (CS-GeS<sub>2</sub>S<sub>2/2</sub>)<sub>∞</sub> chains instead of isolated ES-Ge<sub>2</sub>S<sub>6</sub> dimers [33]. Consequently, its Raman spectrum [34] looks differently compared to Tl<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub>.

## Tl4GeS4

Finally, the thallium-rich crystal Tl<sub>4</sub>GeS<sub>4</sub> also exhibits a good agreement with DFT modelling. The most intense spectroscopic feature at 349 cm<sup>-1</sup> corresponds to the  $A_1$  breathing of isolated GeS<sub>4</sub> tetrahedra while asymmetric stretching and bending/deformation modes are observed either at higher frequencies or below 230 cm<sup>-1</sup>, respectively. As expected, the Tl-related modes appear below 100 cm<sup>-1</sup>.

Resuming, we observe an excellent consistency between the crystal structure of  $Tl_2S$ -GeS<sub>2</sub> compounds, their experimental Raman spectra and DFT modelling.

#### 3.2. Experimental Raman spectra in the Tl<sub>2</sub>S-GeS<sub>2</sub> glasses

Before discussing the evolution of Raman features in the quasi-binary glassy system, it seems to be reasonable to remind the main structural motifs in glassy g-GeS<sub>2</sub>, reflected by their characteristic vibrational modes and DFT replicas [35], shown in Fig. 4.



Figure 4. (a) Experimental Raman spectrum of glassy  $GeS_2$  and DFT Raman spectra of corner-sharing CS- $GeS_{4/2}$  and edge-sharing ES- $GeS_{4/2}$  tetrahedra, as well as of ethane-like ETH- $Ge_2S_{6/2}$  units [35]; (b) schematic representation of (1) CS- $Ge_2S_7$ , (2) ES- $Ge_2S_6$ , and (3) ETH- $Ge_2S_6$  clusters used for DFT modelling. The terminal hydrogen species are not shown and H-related vibrational modes are removed from the DFT spectra.

The GeS<sub>2</sub> stoichiometric glass exhibits well-known spectroscopic features consistent with neutron and X-ray diffraction results (see for example [36-38] and references therein). The most intense line at 344 cm<sup>-1</sup> correspond to the  $A_1$  symmetric in-phase breathing of cornersharing CS-GeS<sub>4/2</sub> tetrahedra [39-41], accompanied by a  $A_{1c}$  companion line at 372 cm<sup>-1</sup> [42] related to the symmetric breathing of edge-sharing ES-GeS<sub>4/2</sub> units. The broad asymmetric feature at 259 cm<sup>-1</sup> originates from Ge-Ge stretching in ethane-like ETH-Ge<sub>2</sub>S<sub>6/2</sub> units [42-44]. Recent *ab initio* calculations [35,44,45] have shown that the broad poorly resolved feature around 400 cm<sup>-1</sup> contains contributions from both CS- and ES-tetrahedra as well as from ETH-units, but the peak at 436 cm<sup>-1</sup> is due entirely to the highest  $F_2$  mode in the ES-GeS<sub>4/2</sub> tetrahedra [35,45].

As it was mentioned previously [13], the Tl-poor glasses ( $0.05 \le x \le 0.15$ ) are hardly to be obtained without crystalline impurities of high-temperature  $\beta$ -GeS<sub>2</sub>. The Raman spectroscopy measurements, Fig. 5, confirm the diffraction results. In particular, for these glass compositions it was possible to find the  $\beta$ -GeS<sub>2</sub> crystallites at the glass surface and collect good quality spectra, Fig. 6(b). The reason of such behavior seems to the difference in crystal structure of  $\beta$ -GeS<sub>2</sub> and Tl-poor Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub> ( $x = \frac{1}{3}$ ), Fig. 1. The high-temperature  $\beta$ -GeS<sub>2</sub> polymorph consists of 2D layers formed by CS- and ES-GeS<sub>4/2</sub> tetrahedra forming two-, three- and eight-membered rings. The crystal structure of  $\beta$ -GeS<sub>2</sub> appears to be the main structural motif of glassy GeS<sub>2</sub> [21,35-44]. In contrast, monoclinic Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub> consists of isolated adamantane units Ge<sub>4</sub>S<sub>10</sub> connected by thallium species, i.e., its local and intermediate-range order is drastically different from that in  $\beta$ - and glassy GeS<sub>2</sub>. Consequently, the spatial incompatibility of the structural motifs leads to enhanced crystallization ability of the Tl-poor glasses.



Figure 5. Raman spectra of (a)  $(Tl_2S)_{0.05}(GeS_2)_{0.95}$  glass and (b) crystalline impurities of high-temperature  $\beta$ -GeS<sub>2</sub>, found on the surface of the *x* = 0.05 sample.

Experimental Raman spectra of the  $(Tl_2S)_x(GeS_2)_{1-x}$  glasses are shown in Fig. 6. Two contrasting composition trends are observed for semiconducting ( $x \le 0.3$ ) and ion-conducting ( $x \ge 0.3$ ) glasses, consistent with composition limits of the  $(Tl_2S)_x(GeS_2)_{1-x}$  compounds [13].



Figure 6. Experimental Raman spectra of the  $(Tl_2S)_x(GeS_2)_{1-x}$  glasses: (a)  $0 \le x \le 0.15$  and (b)  $0.2 \le x \le 0.5$ . Characteristic spectral features related to adamantane-like units, corner-sharing CS-Ge<sub>m</sub>S<sub>n</sub> oligomers and ES-Ge<sub>2</sub>S<sub>6</sub> dimers are highlighted in yellow, red and blue, respectively. The  $A_1$  mode of high-temperature  $\beta$ -GeS<sub>2</sub> crystalline impurities at  $\approx 361$  cm<sup>-1</sup>, observed in thallium-poor glasses (x = 0.05 and 0.10), is indicated by a magenta star.

At  $x \le 0.3$ , the Raman spectrum of g-GeS<sub>2</sub> exhibits characteristic changes. (i) One observes the appearance and growth of a new 450 cm<sup>-1</sup> mode, similar to the A<sub>1</sub> mode in Ge<sub>4</sub>S<sub>10</sub> adamantane units corresponding to in-phase Ge-S<sub>t</sub> stretching. (ii) The A<sub>1</sub> symmetric Ge-S stretching mode in CS-GeS<sub>4/2</sub> tetrahedra at 344 cm<sup>-1</sup> shifts to lower frequencies. (iii) The A<sub>1</sub> symmetric (372 cm<sup>-1</sup>) and  $F_2$  asymmetric (435 cm<sup>-1</sup>) modes in ES-GeS<sub>4/2</sub> tetrahedra disappear with increasing x. The above changes indicate a gradual transformation of the continuous g-GeS<sub>2</sub> network into a fragmented disordered structure consisting of isolated adamantane-like units and connecting Tl species. Similar trend was reported earlier for alkali thiogermanate glasses [28,34].



Figure 7. Experimental Raman spectra of  $(Tl_2S)_x$ (GeS<sub>2</sub>)<sub>1-x</sub> glasses in comparison with respective crystalline compounds: (a) x = 0.3 and c-Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub> ( $x = \frac{1}{3}$ ); (b) x = 0.5 and c-Tl<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub> ( $x = \frac{1}{2}$ ).

Figure 8. Experimental Raman spectra of the x = 0.5 glass and (a) crystalline references c-Na<sub>2</sub>GeS<sub>3</sub> and c-Na<sub>6</sub>Ge<sub>2</sub>S<sub>7</sub> [34]; (b) DFT Raman spectra of Ge<sub>3</sub>S<sub>10</sub>Tl<sub>8</sub> and Ge<sub>3</sub>S<sub>9</sub>Tl<sub>6</sub> clusters.

The x = 0.3 glass has a nearly identical composition to monoclinic Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub> ( $x = \frac{1}{3}$ ) and shows a similar Raman spectrum with one important difference, a significant population of Ge-S stretching modes between 350 and 420 cm<sup>-1</sup>, Fig. 7(a). Changes in this spectral region become predominant for ion-conducting (Tl<sub>2</sub>S)<sub>x</sub>(GeS<sub>2</sub>)<sub>1-x</sub> glasses at x > 0.3.

The Tl-rich x = 0.5 glass has the same composition as triclinic Tl<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub> ( $x = \frac{1}{2}$ ). Nevertheless, the two materials exhibit quite different Raman spectra, Fig. 7(b). The  $A_1$  breathing mode of ES-Ge<sub>2</sub>S<sub>6</sub> dimers at 370 cm<sup>-1</sup> is present in the ion-conducting glasses, Fig. 6(b). However, this mode is not the most intense even for the x = 0.5 vitreous alloy, in contrast to its crystalline counterpart. The  $A_1$  in-phase breathing of corner-sharing tetrahedral units at 325 cm<sup>-1</sup> remains the predominant spectral feature for the ion-conducting  $(Tl_2S)_x(GeS_2)_{1-x}$  glasses. In other words, the structural motif of triclinic Tl<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub>, ES-Ge<sub>2</sub>S<sub>6</sub> dimers, does not determine the local and intermediate-range order of Tl-rich glasses at x > 0.3. In addition, the second most intense feature at 400 cm<sup>-1</sup> (x = 0.5) does not have the crystalline analogue in the Tl<sub>2</sub>S-GeS<sub>2</sub> system, Fig. 2(a), suggesting a new bonding pattern for the ion-conducting glasses.

The fast ion-conducting Na<sub>2</sub>S-GeS<sub>2</sub> system [34] allows the structural puzzle to be solved. Two crystalline compounds in this system, monoclinic Na<sub>2</sub>GeS<sub>3</sub>, space group  $P2_1/c$  [33], and monoclinic Na<sub>6</sub>Ge<sub>2</sub>S<sub>7</sub>, space group C2/c [46], reveal different crystal structures compared to heavy alkali and thallium thiogermanates, discussed previously. The Raman spectra of these two monoclinic crystals [34] show some resemblances with unidentified spectral features in the Tl-rich thiogermanate glasses, Fig. 8(a).



Figure 9. The crystal structure of (a) monoclinic Na<sub>2</sub>GeS<sub>3</sub> = Na<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub>, space group  $P2_1/c$  [33], and (b) monoclinic Na<sub>6</sub>Ge<sub>2</sub>S<sub>7</sub>, space group C2/c [46].

Figure 10. The optimized structure of corner-sharing oligometric clusters used for DFT modelling: (a) CS-Ge<sub>3</sub>S<sub>10</sub>Tl<sub>8</sub> trimer, and (b) CS-Ge<sub>3</sub>S<sub>9</sub>Tl<sub>6</sub> ring.

The two Na<sub>2</sub>S-GeS<sub>2</sub> crystals consist of CS-GeS<sub>4/2</sub> tetrahedra (Fig. 9). In monoclinic Na<sub>2</sub>GeS<sub>3</sub>, the tetrahedra form infinite chains, (CS-GeS<sub>2</sub>S<sub>2/2</sub>)<sub>∞</sub>, while Na<sub>6</sub>Ge<sub>2</sub>S<sub>7</sub> is built up by CS-Ge<sub>2</sub>S<sub>7</sub> dimers. The two Ge-S structural motifs are connected by sodium species. The stretching frequencies of the  $A_1$  in-phase breathing modes of CS-units for the two crystals show a blue shift, +6 cm<sup>-1</sup> (Na<sub>2</sub>GeS<sub>3</sub>) or +17 cm<sup>-1</sup> (Na<sub>6</sub>Ge<sub>2</sub>S<sub>7</sub>), consistent with a smaller average oscillator mass in sodium crystals compared to thallium thiogermanate glasses. The symmetric and asymmetric Ge-S<sub>t</sub> stretching frequencies involving terminal sulfur are matching better a poorly resolved bimodal high-frequency feature at 400 and 435 cm<sup>-1</sup> for the x = 0.5 glass, Fig. 8(a). A preliminary DFT modelling of several oligometric CS-Ge<sub>m</sub>S<sub>n</sub>Tl<sub>k</sub> clusters, Fig. 10, shows a general agreement with the experimental Raman spectra, Fig. 8(b), but needs to be continue to precise the exact stoichiometry of oligometric CS-units.

## 4. Conclusions

The experimental Raman spectra of the known crystalline compounds in the  $Tl_2S$ -GeS<sub>2</sub> system are consistent with the crystal structure and DFT modelling of the respective Tl-Ge-S

clusters, representing structural motifs in c-Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub>, c-Tl<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub>, and c-Tl<sub>4</sub>GeS<sub>4</sub>. Structural motifs in (Tl<sub>2</sub>S)<sub>*x*</sub>(GeS<sub>2</sub>)<sub>1-*x*</sub> glasses, where  $0 \le x \le 0.5$ , are reminiscent of respective Tl<sub>2</sub>S-GeS<sub>2</sub> crystals and, more specifically, of thallium-poor monoclinic Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub> (x = 1/3) and equimolar triclinic Tl<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub> (x = 1/2). The Raman spectra of the (Tl<sub>2</sub>S)<sub>*x*</sub>(GeS<sub>2</sub>)<sub>1-*x*</sub> glasses can be divided into two groups: (1) semiconducting glasses at  $x \le 0.3$ , and (2) ion-conducting vitreous alloys at x > 0.3. At  $x \le 0.3$ , the continuous g-GeS<sub>2</sub> network is transformed into a fragmented structure consisting of adamantane-like units Ge<sub>4</sub>S<sub>10</sub>, related to Tl<sub>4</sub>Ge<sub>4</sub>S<sub>10</sub>, held together by thallium species. In the ion-conducting composition domain at x > 0.3 takes place further fragmentation of tetrahedral tetramers Ge<sub>4</sub>S<sub>10</sub> into ES-Ge<sub>2</sub>S<sub>6</sub> dimers, as in Tl<sub>4</sub>Ge<sub>2</sub>S<sub>6</sub>. However, the glasses at x > 0.3 consists also of small CS-oligomers, i.e., CS-Ge<sub>2</sub>S<sub>7</sub> dimers or CS-Ge<sub>3</sub>S<sub>10</sub> trimers, absent in the Tl<sub>2</sub>S-GeS<sub>2</sub> crystals but present in light alkali analogues, i.e., monoclinic Na<sub>6</sub>Ge<sub>2</sub>S<sub>7</sub>.

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