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The coupling phase in cavity magnonics

 $Alan$ $Gardin^{1,2}, Jeremy Bourhill^{2,3}, Vincent Vlaminck², Christian Person², Christophe Fumeaux¹,$ Vincent Castel² & Giuseppe Tettamanzi¹

Abstract

- •Cavity magnonics relevant for quantum transducers, quantum memories (using dark mode physics), non-reciprocal devices [1]: such applications justify studying cavity magnonics in multimode systems.
- Microwave photon/magnon coupling due to Zeeman coupling involves a ω pling phase factor which is usually rightfully neglected.
- \Rightarrow unless the system is "loop-coupled" (see next section for definition and examples), coupling phases are unobservables.
- (Experimentally) Simplest system in which the coupling phases manifest: two magnon modes coupling simultaneously to two cavity modes. The various coupling phases accumulate in a single quantity θ : the physical phase. Physical phase found to influence the presence of dark modes, and the strength of the cavity-mediated coupling.

Figure 1: Representation of the magnon/photon coupling. Coupling phase φ measures the difference in direction of the (local) yellow arrows.

with coupling strength $g/2\pi \propto \eta$ √ $\overline{\omega_c}$ and

Relevance of the coupling phase

- Origin of the coupling phase φ : Zeeman coupling between macrospin $S =$ $S_x\hat{\mathbf{x}} + S_y\hat{\mathbf{y}} + S_z\hat{\mathbf{z}}$ and the cavity mode's RF magnetic field $\mathbf{h} = h_x\hat{\mathbf{x}} + h_y\hat{\mathbf{y}}$ (h $\perp H_0$ \hat{z} static applied magnetic field)
- Interaction Hamiltonian using the rotating wave approximation [2]:

Figure 2: Non loop-coupled systems $(k < n)$

$$
H_I = \int_{V_m} d^3 r \, \mathbf{S} \cdot \mu_0 \mathbf{h}
$$

= $\hbar g \left(e^{i \varphi} cm^\dagger + e^{-i \varphi} c^\dagger m \right)$ (2)

Figure 4: Reduction of the coupling phases to a single physical phase θ through unitary transformations.

$$
\eta = \sqrt{\frac{\left(\int_{V_m} \mathrm{d}^3 r h_x\right)^2 + \left(\int_{V_m} \mathrm{d}^3 r h_y\right)^2}{V_m \int_{V_c} \mathrm{d}^3 r |\mathbf{h}|^2}}, \ \varphi = \arg \left\{ \int_{V_m} \mathrm{d}^3 r h_x + i \int_{V_m} \mathrm{d}^3 r h_y \right\} \tag{3}
$$

•A unitary transformation can gauge the phase away by rotating one of the modes \Rightarrow if φ can be eliminated, it cannot have physical consequences •Coupling phase can be "rotated away" when the system is not "loopIndirect magnon/magnon coupling at $\omega_m = \omega_c$

- •Coupling phase can have physical consequences in loop-coupled multimode systems
- Resulting discrete gauge-invariant phase θ akin to a synthetic gauge field, which can break time-reversal symmetry for specific values [3]

An example with a physical phase

System definition

- System under study: two magnon modes (annihilation operators m_0 and m_1 , angular frequencies $\omega_m - \delta_m$ and $\omega_m + \delta_m$) coupling to two cavity modes (c₀ and c₁, angular frequencies $\omega_c - \delta_c$ and $\omega_c + \delta_c$), with coupling strength $g_0/2\pi$ to c_0 and $g_1/2\pi$ to c_1 .
- Loop-coupled system: the various coupling phases combine to form a $physi$ cal phase θ characterising the physics:

• Synthetic gauge field+dissipation engineering \Rightarrow non-reciprocal behaviour [4] However, the physical phase is fixed by the cavity geometry...

Figure 6: Spectra for $\theta = 0$ in solid lines and $\theta = \pi$ in dots. Other parameters same as in Fig 5. G_{θ} is given by equation (4) below.

with θ -dependent indirect magnon/magnon coupling strength $G_{\theta}/2\pi$:

$$
G_{\theta} = \frac{\delta_c}{\delta_c^2 - \delta_m^2} \left(g_0^2 - e^{i\theta} g_1^2 \right).
$$
 (4)

Dark mode physics

coupled", i.e. we have n modes but $k < n$ couplings between modes: we have *n* degrees of freedom but only $k < n$ constraints

• For $\delta_m = 0$ and $\theta = 0$, M_π does not couple: dark mode

Outlook

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