



Maximal reachable temperatures for Western Europe in current climate

Robin Noyelle, Yi Zhang, Pascal Yiou, Davide Faranda

► To cite this version:

Robin Noyelle, Yi Zhang, Pascal Yiou, Davide Faranda. Maximal reachable temperatures for Western Europe in current climate. *Environmental Research Letters*, 2023, 10.1088/1748-9326/acf679 . hal-04210152

HAL Id: hal-04210152

<https://hal.science/hal-04210152>

Submitted on 28 Sep 2023

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons Attribution 4.0 International License

LETTER • OPEN ACCESS

Maximal reachable temperatures for Western Europe in current climate

To cite this article: Robin Noyelle *et al* 2023 *Environ. Res. Lett.* **18** 094061

View the [article online](#) for updates and enhancements.

You may also like

- [Internal variability in European summer temperatures at 1.5 °C and 2 °C of global warming](#)
Laura Suarez-Gutierrez, Chao Li, Wolfgang A Müller et al.
- [Avoided population exposure to extreme heat under two scenarios of global carbon neutrality by 2050 and 2060](#)
Yadong Lei, Zhili Wang, Xiaoye Zhang et al.
- [Escalating global exposure to compound heat-humidity extremes with warming](#)
Dawei Li, Jiacan Yuan and Robert E Kopp

ENVIRONMENTAL RESEARCH
LETTERS

LETTER

OPEN ACCESS

RECEIVED
21 June 2023REVISED
4 August 2023ACCEPTED FOR PUBLICATION
4 September 2023PUBLISHED
18 September 2023

Original Content from
this work may be used
under the terms of the
[Creative Commons
Attribution 4.0 licence](#).

Any further distribution
of this work must
maintain attribution to
the author(s) and the title
of the work, journal
citation and DOI.



Maximal reachable temperatures for Western Europe in current climate

Robin Noyelle^{1,*}, Yi Zhang^{2,3}, Pascal Yiou¹ and Davide Faranda^{1,4,5}¹ Laboratoire des Sciences du Climat et de l'Environnement, UMR 8212 CEA-CNRS-UVSQ, Université Paris-Saclay & IPSL, Gif-sur-Yvette 91191, France² Department of Earth and Planetary Science, University of California, Berkeley, Berkeley, CA 94720, United States of America³ Miller Institute for Basic Research in Science, University of California, Berkeley, Berkeley, CA 94720, United States of America⁴ London Mathematical Laboratory, 8 Margravine Gardens London, W6 8RH London, United Kingdom⁵ Laboratoire de Météorologie Dynamique, Ecole Normale Supérieure, PSL Research University & IPSL, Paris, France

* Author to whom any correspondence should be addressed.

E-mail: robin.noyelle@lscce.ipsl.fr**Keywords:** heatwaves, climate change, maximum temperatures, flow analogues

Abstract

Human bodies, ecosystems and infrastructures display a non-linear sensibility to extreme temperatures occurring during heatwave events. Preparing for such events entails to know how high surface air temperatures can go. Here we examine the maximal reachable temperatures in Western Europe. Taking the July 2019 record-breaking heatwave as a case study and employing a flow analogues methodology, we find that temperatures exceeding 50 °C cannot be ruled out in most urban areas, even under current climate conditions. We analyze changes in the upper bound of surface air temperatures between the past (1940–1980) and present (1981–2021) periods. Our results show that the significant increase in daily maximum temperatures in the present period is only partially explained by the increase of the upper bound. Our results suggest that most of the warming of daily maximum surface temperatures result from strengthened diabatic surface fluxes rather than free troposphere warming.

1. Introduction

Extreme weather events have strong impacts on societies and ecosystems [1]. Among them, heatwaves have been the focus of extensive attention due to their increasing frequency and intensity with global warming [2]. Most biological and technical systems have a limiting capacity to cope with extreme temperatures and can respond non linearly once certain thresholds are reached. These include human bodies [3–6], plants [7, 8], ecosystems [9] and infrastructures [10]. Urban areas are particularly at risk due to the so-called heat island effect [11] which enhances maximum temperatures during the day and minimum temperatures during the night. The capacity to predict and anticipate future maximum intensities of heatwaves is therefore of primary importance for adaptation to climate change.

The physical mechanisms leading to heatwaves in the midlatitudes combine specific atmospheric synoptic conditions and anomalously low soil moisture, which can interact to further increase the intensity

of the event [11–14]. Heatwaves are associated with a slow moving, sometimes called quasi-stationary, high-amplitude Rossby wave [15] which is often embedded in a hemispheric pattern of wave number 5 or 7 [16–18]. Above the heatwave region, a blocking anticyclone builds up at mid- and upper-level troposphere in conjunction with a change in the jet stream's climatological path, towards a large poleward meridional meander. The anticyclone sustains the poleward advection of warm air along its western flank, adiabatic warming by subsidence and clear skies at its center. When linked to abnormally dry soils, these conditions promote warming through short-wave insolation, resulting in the allocation of incoming solar energy towards sensible rather than latent heat [19–26]. If the anticyclone stays stationary, extreme temperature can be reached. Although it is well established that global warming renders the temperature distribution hotter on average, how the coupled dynamics between the atmosphere, the oceans and the soils will evolve in response to this forcing during the hottest days is still debated [11, 14, 27–31].

One way to quantify the risk associated with heatwaves is to measure the tail probabilities of the temperature distribution, e.g. the risk of the maximum annual temperature reaching a certain level. A classical approach to deal with this assessment is extreme value theory (EVT). EVT has been used to determine statistical models for maxima (or minima) of climate variables of interest [32], such as temperatures. It is based on a convergence principle of the probability distribution of block maxima or peaks-over-threshold [33–35] towards a generalized extreme value (GEV) distribution. It allows to compute return values corresponding to very large return periods (i.e. longer than the period of observations), even in non stationary contexts [36]. This approach is for example extensively used by attribution methods to compare the probabilities of reaching extreme temperatures between a counterfactual world without climate change and a factual world with climate change [37–39]. This allows to estimate how climate change made a particular event more (or less) likely.

Fitting a GEV distribution to extreme temperatures usually gives estimates of the shape parameter that are robustly negative (e.g. [40, 41]). This implies that, in a stationary context, the distribution of annual maximum temperatures is bounded upwards. This upper bound is often the information asked to the scientific community by practitioners who want to prepare for the worst case scenario. However, recent intense heatwaves such as in 2010 in Russia [18], 2019 in Western Europe [42] and 2021 in the Pacific Northwest [43, 44] have challenged the reliability of this statistical upper bound by breaking it sometimes by a large margin. Another approach to estimate maximal reachable temperatures would be to run a long simulation with a climate model to sample more extreme events than those observed in the past. Apart from the inherent limitations of models to represent correctly the entire temperature distribution—especially at its tail [40, 45, 46]—this method is limited by the fact that estimating precisely small probabilities requires an extremely long simulation [47], which is too costly for most complex models. Various approaches have been suggested to address this problem, like rare events algorithms [48–50] or ensemble boosting [51].

Recently, [52] proposed a physics-based theory to address the question of the upper bound for extreme temperatures over midlatitude continental areas. They provide evidence for the hypothesis that moist convective instability sets a fundamental upper limit for surface temperatures. This allows for an estimate of maximal temperatures based on large scale circulation of the free-troposphere and surface specific humidity. In this paper, we make use of this theory to investigate the maximal reachable temperature in Western Europe in the current climate. We use

the flow analogues attribution method [53–55]. We condition the estimate of maximal temperatures on the free-troposphere large scale circulation observed during the most intense event of the recent years in Western Europe, namely the July 2019 heatwave. 2019 saw two exceptional heatwaves in June and July [56] with maximum temperatures very unlikely in a world without anthropogenic climate change [45, 57].

We focus on the July 2019 heatwave—when most records of Western Europe cities were broken—and we estimate the maximal temperatures this event could have reached in the present climate. We also investigate how these upper bounds have changed with respect to the past and detail the physical mechanisms related to the changes in daily maximum of surface temperatures.

This paper is organized as follows: in section 2, we present the data used, the computation of the heatwave maximal temperature and the flow analogs methodology. Section 3 presents and discusses the results. Finally, the conclusions drawn from our analysis are presented in section 4.

2. Data and methods

We use the ERA5 reanalysis data set of the European Center for Medium Range Weather Forecasting (ECMWF) over the period 1940–2021 [58]. The data have a horizontal resolution of $0.25^\circ \times 0.25^\circ$. For illustration purposes, we additionally use daily temperature data from the Paris Montsouris weather station taken from the Infoclimat website (<https://www.infoclimat.fr/climatologie/globale/paris-montsouris/07156.html>).

To study how more intense an observed heatwave event could have been, we use the criteria developed by [52] who provided evidence for the hypothesis that convective instability fundamentally limits maximal surface air temperatures over midlatitude land. This hypothesis states that, during heatwaves, surface air temperature T reaches its maximal value when the air column is neutrally stratified. Accordingly, the heatwave stops when precipitation is triggered by the convective instability of the boundary layer. This happens when the surface air parcel moist static energy

$$\text{MSE} = c_p T + L_v q + g z_s \quad (1)$$

is equal to the free-tropospheric saturation moist static energy MSE_a^* . Here c_p is the specific heat of air at constant pressure, L_v the latent heat of vaporization, q the specific humidity at the surface, g the gravitational acceleration and z_s the geopotential height of the surface. We choose to define the free-tropospheric situation as the state of the atmosphere at 500 hPa and MSE_a^* is therefore computed from equation (1) by

replacing T with the air temperature at 500 hPa T_{500} , q with the saturation specific humidity at T_{500} :

$$q_{\text{sat}}(T_{500}) \simeq \frac{\epsilon e_{\text{sat}}(T_{500})}{500 \text{ hPa}} \quad (2)$$

where ϵ is the molar ratio between water vapor and dry air and $e_{\text{sat}}(T_{500})$ is the saturation vapor pressure given by the Clausius–Clapeyron relation—and z_s with the geopotential height at 500 hPa z_{500} . The maximal surface temperature is therefore:

$$T_{s,\text{max}} = T_{500} + \frac{L_v}{c_p}(q_{\text{sat}}(T_{500}) - q) + \frac{g}{c_p}(z_{500} - z_s). \quad (3)$$

This temperature is maximized if q is null. However, such a low specific humidity is unrealistic for midlatitude land because of the presence of surface waters and vegetation. In the following, except specified explicitly, we therefore use the hourly June–July–August–September (JJAS) minimum q_{min} ever observed over the period 1940–2021 at each grid point. Using q at 2 m to compute $T_{s,\text{max}}$ is a conservative choice because when an air parcel rises, it entrains surrounding dryer air, which decreases its specific humidity. For T_{500} and z_{500} we use the daily mean values as the diurnal cycle does not influence strongly these quantities.

$T_{s,\text{max}}$ defines a heatwave potential in the sense that it gives the maximal temperature potentially reachable given the synoptic situation. Whether this temperature will be reached is however not certain because, for example, the circulation pattern may change too fast for the temperature to have time to increase under the forcing of incoming energy fluxes. Letting $TX1d$ be the observed daily temperature maximum at the surface, one can therefore define an index I_{TX1d} for how much an observed heatwave is close to its maximal potential:

$$I_{TX1d} := \frac{TX1d - T_{\text{clim}}}{T_{s,\text{max}} - T_{\text{clim}}}, \quad (4)$$

where T_{clim} is a climatological temperature and $TX1d$ is the observed daily maximum temperature. Here we use the July averaged daily mean temperature as a climatological value. When this index is close to 1, the event is close to its maximal potential intensity: it could not have been more intense.

To determine how global warming may have impacted the physical mechanisms leading to the extremely intense event of July 2019 in Western Europe, we employ an attribution methodology. The most common attribution method of a specific event [37] consists in comparing the probabilities of an observable reaching a certain level in a counterfactual world and in a factual world. The difference between the two periods usually lies in their global mean

temperature, i.e. worlds with or without anthropogenic global warming. The ratio of the probabilities gives how much more likely an observed event has been rendered because of global warming. This method however does not condition the result on the large scale circulation and therefore bears the risk of comparing atmospheric dynamics that are different even though they give rise to the same values for the observable of interest.

Here we make use of the flow analogues attribution methodology proposed by [55] and adapted from [53] (see also [59, 60]). The idea of the method is to compare the expectation of a variable of interest X in the counterfactual world ($F = 0$) and in the factual world ($F = 1$), conditional on the large scale circulation $\mathcal{C}(\zeta)$ of the observed event ζ :

$$\Delta_{\mathcal{C}(\zeta)}X = \mathbb{E}[X | F = 1, \mathcal{C}(\zeta)] - \mathbb{E}[X | F = 0, \mathcal{C}(\zeta)]. \quad (5)$$

In our analysis, ζ corresponds to the 25 July 2019 situation, the day where the heat peak was reached. The conditioning on the large scale circulation $\mathcal{C}(\zeta)$ should be understood as being in a reasonable vicinity of the circulation of the event rather than as a strict equality, which cannot be enforced. This framework allows to answer the question: how a similar large scale circulation pattern in the two worlds leads to different outcomes in an observable of interest? If the difference $\Delta_{\mathcal{C}}X$ is statistically significant, then one can say that in the factual world the event has been rendered more (or less) intense by $\Delta_{\mathcal{C}}X$. This method allows to disentangle the role of thermodynamics and dynamics (see [61] for the 2019 event). Whether the event has become more likely can in principle be obtained [53] if one can estimate the probabilities $\mathbb{P}[\mathcal{C}(\zeta) | F]$ and $\mathbb{P}[\mathcal{C}(\zeta) | X > x]$, where x is a given threshold, in the two worlds.

This method can be used with climate models output with or without anthropogenic forcing to define the factual and counterfactual worlds. However, models have known deficiencies, including biases [62] and incorrect dynamics of extremes under forcing over Western Europe [40, 45, 46]. As in [55] we therefore choose to rely on reanalysis outputs only, which, even though they also present limitations, are the closest we have to the actual past state of the atmosphere. Here, we define the counterfactual world as the past period (1940–1980) and the factual world as the present period (1981–2021).

To condition on the synoptic circulation, we select in both periods the best 40 analogs as the daily mean 500 hPa geopotential height (z_{500}) maps minimizing the pointwise Euclidean distance with the daily mean z_{500} map of the target day ζ . In the factual period we also impose that no analog can be found within a 20 calendar days window around the

25 July 2019 to avoid selecting the event itself. The reason for choosing to condition on z_{500} is that in the quasi-geostrophic approximation relevant for synoptic scales, z_{500} acts as an approximate streamfunction for the general circulation and has been shown to be a good conditioning field for temperatures [54]. Over the period 1940–2021, the troposphere has warmed and geopotential heights may therefore have inflated. To avoid this factor to bias the analogs selection, we remove a spatially uniform trend to the geopotential heights before computing the Euclidean distance. This ensures that the spatial gradients are preserved and therefore that analogs have a similar synoptic circulation. To compute the trend we average spatially the z_{500} maps and fit a second degree polynomial to this time series. The second order was chosen because of the so-called warming pause in the global temperature of the Earth in the period 1950–1970 [63]. The analogs are computed over the domain depicted with a dashed box in figure 1 panel (c). This domain is chosen because it encompasses the observed heatwave region at the surface and the relevant associated synoptic pattern, namely a large poleward meander of the jet stream associated with a strong anticyclonic anomaly over Western Europe. The analogs are chosen only in the months June, July, August and September to ensure closeness with the event. We verified that the results are not qualitatively sensitive to the choice of the number of analogs provided that we extract 20 or 30 analogs maps.

Forty-one years long factual and counterfactual periods ensure a large enough statistical sample of z_{500} maps to be able to select sufficiently representative analogs over such a limited geographical area. The counterfactual period can be assumed to satisfy the assumption of climate stationarity—for example with regards to the mean (not shown) or extreme (see e.g. figure 1 panel (a)) temperatures in summer—with respect to anthropogenic climate change. This period is also long enough for the interannual and interdecadal natural variability of the atmospheric circulation to be averaged over. The same is not necessarily true for the factual period for which there is a strong warming of Western Europe between the 1980s and the 2010s decades. A direct identification of the factual period and a stationary warmer world should therefore be made with care and it is more relevant to frame our results as showing a difference between a past world and a present world [64].

Once obtained the two sets of analogs for the factual and counterfactual periods, we average them out to search for significant differences in observables X of interest. To determine significant changes, we apply a two-tail Welch t-test [65], with different variances, at each grid point. We mark as significant only grid point changes for which the p -value of the test is below 0.05. In the following, as the reference event takes place in July we chose to display anomalies with respect to the July monthly mean and

standard deviation over the period 1990–2021. Apart when specified explicitly, this convention concerns all anomalies.

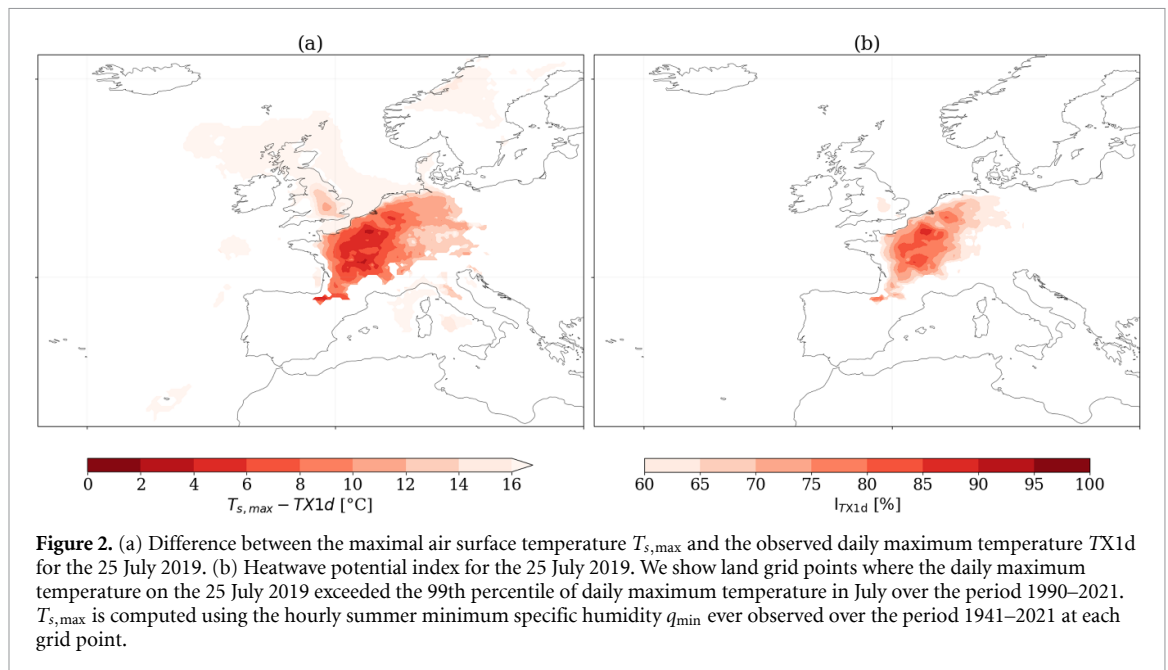
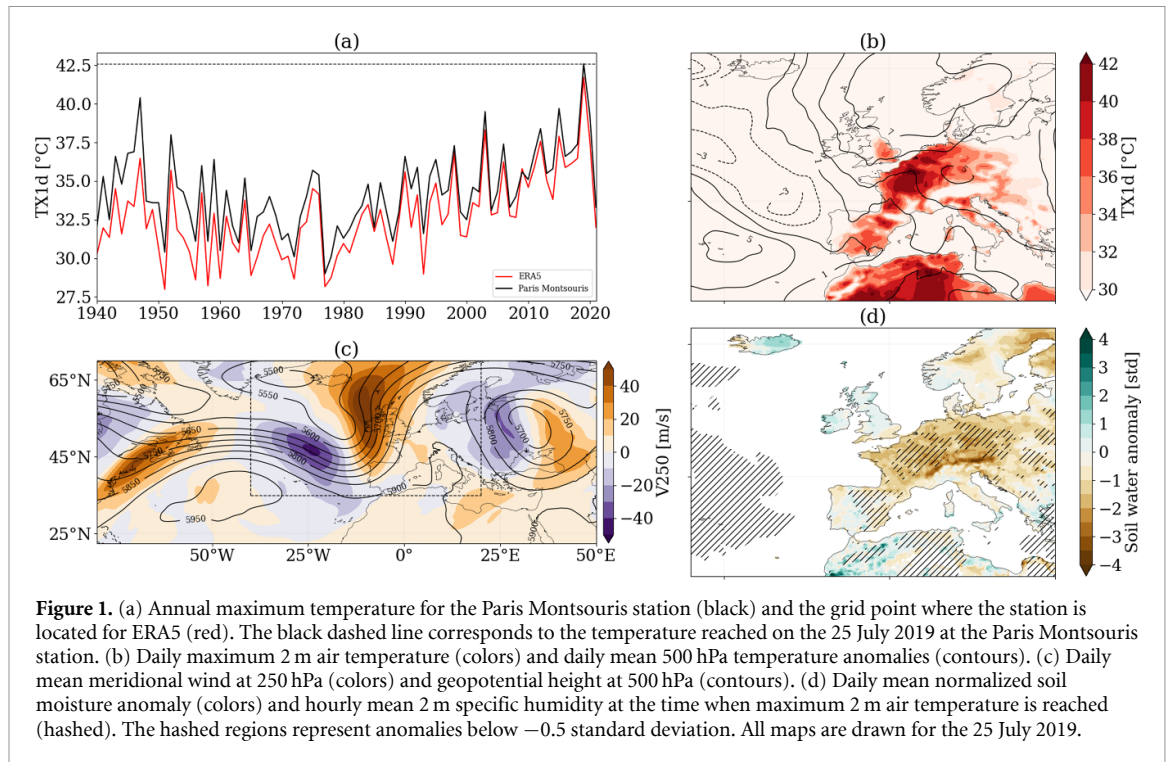
3. Results and discussion

3.1. Description of the event and intensity potential

In the end of July 2019, Western Europe was hit by a heatwave which broke many records in France, Belgium, the Netherlands and Western Germany ([45], figure 1 panel (b)). The heatwave peaked on the 25 July and, for instance, the Paris Montsouris station recorded its largest temperature since 1900 with a new record of 42.6 °C, breaking the previous 1940 record by a margin of 2.2 °C (figure 1 panel (a)). The synoptic situation was characterized by a strongly meridional meander of the high level jet (figure 1 panel (c)) sustaining advection of warm air and anomalies of up to 7 K at 500 hPa above the North Sea (figure 1 panel (b)). Extreme surface temperatures were reached after an anomalously dry and hot month (30% of climatological precipitations and +2.4 °C temperature anomaly for the Paris region) following the previous June heatwave [56], which made the event develop under desiccated soils (figure 1 panel (d)). The event however only lasted a few days and daily maximum temperatures decreased by more than 10 K on the 26 July.

Figure 2 panel (a) shows the difference between the maximal surface air temperature $T_{s,max}$ and the observed daily maximum temperature $TX1d$ for the 25 July 2019. To ensure realism in the value of $T_{s,max}$, we use the hourly JJAS minimum specific humidity q_{min} ever observed over the period 1940–2021 at each grid point. Over the core heatwave region (Northern France and Benelux), the daily maximum temperature provided by the reanalysis data set is close to the maximal potential intensity by a margin of 4–6 K, and consequently the heatwave potential index is close to 90%.

Table 1 shows the estimation of $T_{s,max}$ for several major urban areas in Western Europe which all set their 1940–2021 maximum temperature record during the July 2019 event. For all of them, we used a four grid point average encompassing most of the main city and its surroundings. We compare the estimated maximal surface temperature to the maximum value of a non-stationary GEV distribution fitted on the summer temperature maxima over the 1940–2021 period using the ERA5 reanalysis data set (see supplementary materials for more information). We emphasize that these values are conservative with respect to what can be actually measured by a weather station at the surface because they are averaged over a roughly 50 km by 50 km square. For example, the value for $TX1d$ for the Paris region with ERA5 gives 41.4 °C whereas the Paris Montsouris station measured 42.6 °C.



Strikingly, even in the current climate, it is not possible to rule out the possibility of reaching 50°C in all of these cities once we include the uncertainty on the minimum surface specific humidity. For example, for the Paris area the ERA5 dataset provides a daily maximum of 41.4°C , while $T_{s,max}$ is 46.6°C [45.2, 53.8]. Reaching such temperatures would imply either a very intense drought, with very low values of surface specific humidity, and/or massive advection of heat from neighbouring regions. The non-stationary GEV estimate for the upper bound is significantly lower than $T_{s,max}$ by 5 to

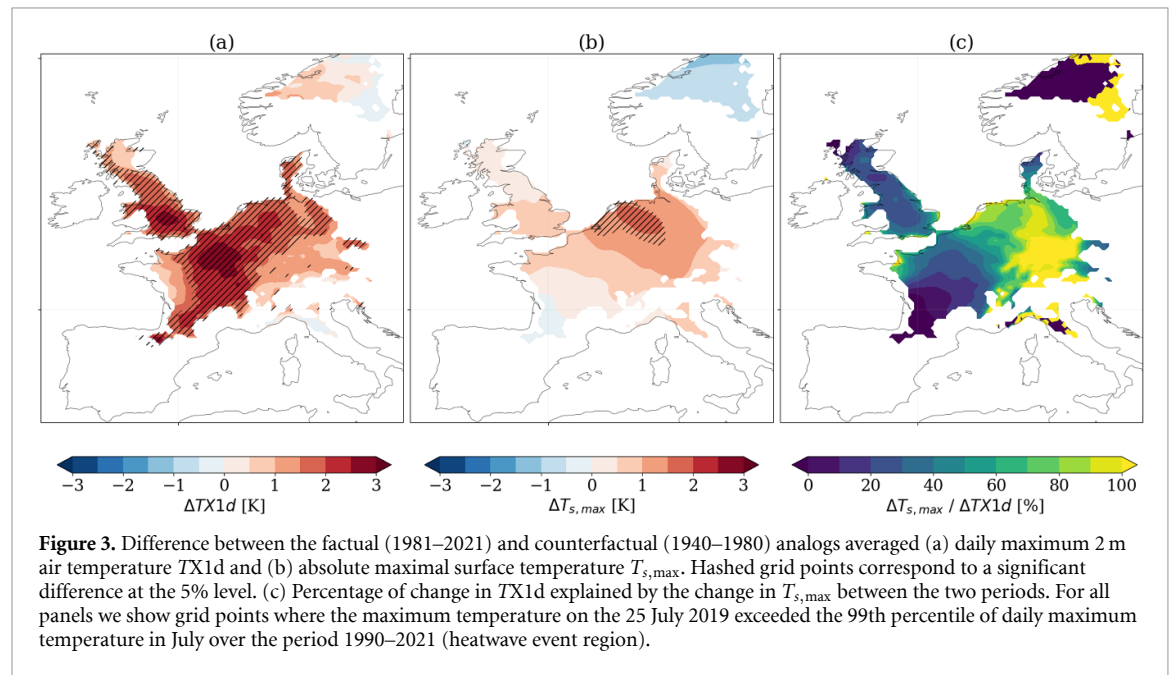
10 K. The median estimate is even lower than the 2019 value for Paris and Lille. For most cities, the estimated $T_{s,max}$ is inside the uncertainty range of the GEV estimate, but much closer to the upper bound than to the lower one. This suggests that the statistical analysis of past time series of maximum temperatures underestimates the risk of reaching extreme temperatures.

3.2. Flow analogs analysis

Figure 3 panel (a) shows the difference between the factual and counterfactual periods of the analogs averaged daily maximum 2 m air temperature $TX1d$

Table 1. Maximal temperatures in °C for several urban areas in Western Europe. All of them reached their 1940–2021 maximum temperature record during the July 2019 event. See supplementary materials for the detail of the computation of the uncertainty range.

City	TX1d (ERA5)	$T_{s,max}$ (realistic)	$T_{s,max}$ (absolute)	GEV estimate
Paris	41.4	46.6 [45.2, 54.6]	54.2 [53.8, 54.6]	40.8 [40.0, 47.2]
Lille	39.7	47.1 [45.0, 55.0]	54.9 [54.6, 55.0]	39.2 [38.3, 48.5]
Bruxelles	38.9	47.1 [45.4, 54.9]	54.8 [54.5, 54.9]	42.8 [37.6, 49.7]
Amsterdam	35.6	47.8 [44.7, 56.2]	55.9 [55.7, 56.2]	37.3 [33.7, 41.5]
Köln	40.2	48.4 [45.9, 55.5]	55.3 [55.0, 55.5]	42.1 [39.6, 48.4]
Franckfurt	39.3	50.1 [46.4, 56.8]	56.6 [56.3, 56.8]	42.2 [39.3, 49.2]
London	35.6	48.6 [46.9, 56.0]	55.8 [55.6, 56.0]	42.9 [36.3, 51.0]
Luxembourg	37.8	45.4 [45.2, 54.1]	53.8 [53.5, 54.1]	41.1 [37.6, 47.5]



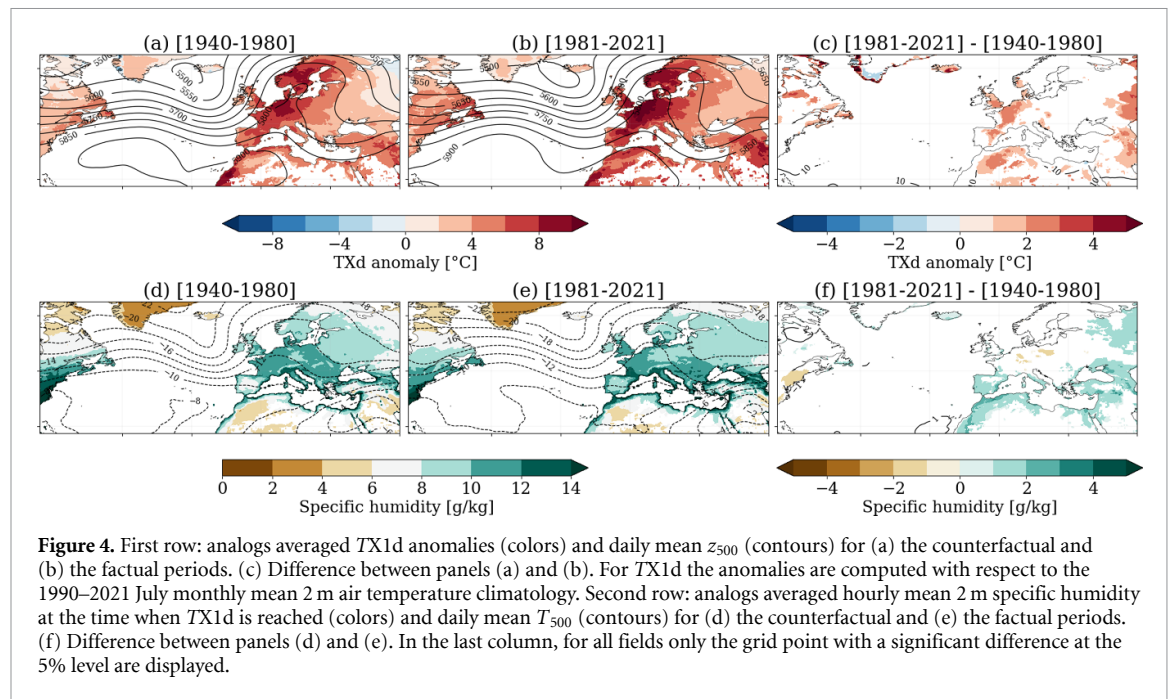
over Western Europe. In the heatwave region (France, England, Benelux and Western Germany) $TX1d$ is between 2 and 3 K more intense in the recent period. It should also be noted that this analogs difference is stronger by approximately 1 K than the difference in the daily mean temperature (see figure C1 in supplementary materials).

Figure 3 panel (b) displays the difference of the absolute maximal surface temperature $T_{s,max}$. This quantity has not changed significantly over most of Western Europe—except a small region above the Netherlands which coincides with an increase of T_{500} (not shown). Figure 3 panel (c) shows the percentage of change in $TX1d$ explained by the increase in $T_{s,max}$. The latter corresponds to no more than 30% of the actual change between the two period over France (Paris region) but as much as 70% over Benelux and Western Germany.

The change in T_{500} is the dominant factor for the stronger increase in absolute $T_{s,max}$ over the Netherlands, with an increase of 0.7 K, whereas only 0.4 K over Paris (see vertical profiles in figure C2). There is also a more significant drying of surface air over the Netherlands (figure C2 panels (b) and (d)). These two elements lead to a much stronger

increase in the analogs averaged lifting condensation level (LCL) between the two regions: 21 hPa over the Netherlands vs 14 hPa over Paris, although the difference between the two periods is non significant at the 5% level. Drier air parcels in a hotter troposphere follow the dry adiabatic temperature profile longer, which leads to increased surface temperature as can be seen when comparing white and magenta lines in figure C2 panels (a) and (c). These elements suggest that under the synoptic atmospheric conditions associated with the 25 July 2019 event, most of the observed change in 2 m air temperature maximum in France between the two periods is caused by surface processes—either local diabatic heating or advection from neighboring regions.

Figure 4 shows the analogs analysis for $TX1d$ and z_{500} (first row), surface specific humidity q and T_{500} (second row). In both periods the averaged z_{500} pattern matches the event z_{500} pattern (figure 1 panel (c)) and the difference between the periods is minimal as can be seen by the fact that there are almost no significant differences in the z_{500} field (figure 4 panel (c)). We are therefore confident on the quality of the analogs and the relevance of the comparison



between mean observables conditional on the synoptic circulation.

Figure 4 panel (f) shows no significant difference in the T_{500} fields between the two periods. This should however be taken with care because in equation (3) this variable is used to compute $e_{\text{sat}}(T_{500})$ which evolves strongly non-linearly with T_{500} . A small difference in T_{500} , even non significant, can therefore lead to significant $T_{s,\text{max}}$ differences at the surface as shown in figure 3 panel (b). Finally, over the heatwave region the difference in specific humidity is not significant (figure 4 panel (f)), except over a small area in Western Germany.

In figure C3 in supplementary materials we present the same plot for upper layer soil moisture, which is a variable closely related to 2 m specific humidity (low soil moisture reduces the evapotranspiration potential). There is a general drying pattern over most of the heatwave region in the factual period, but it is significant only over the same region as specific humidity in figure 4 panel (f).

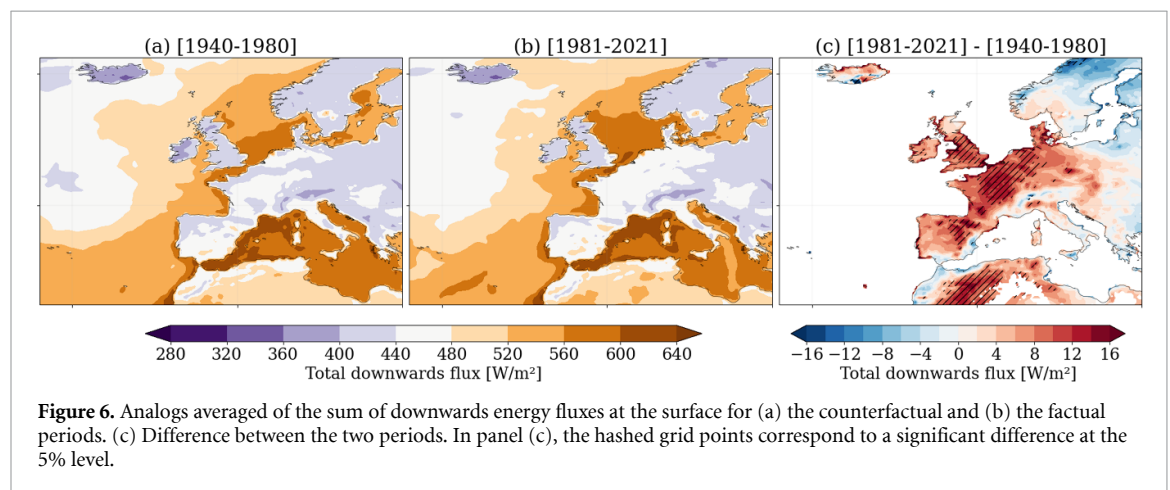
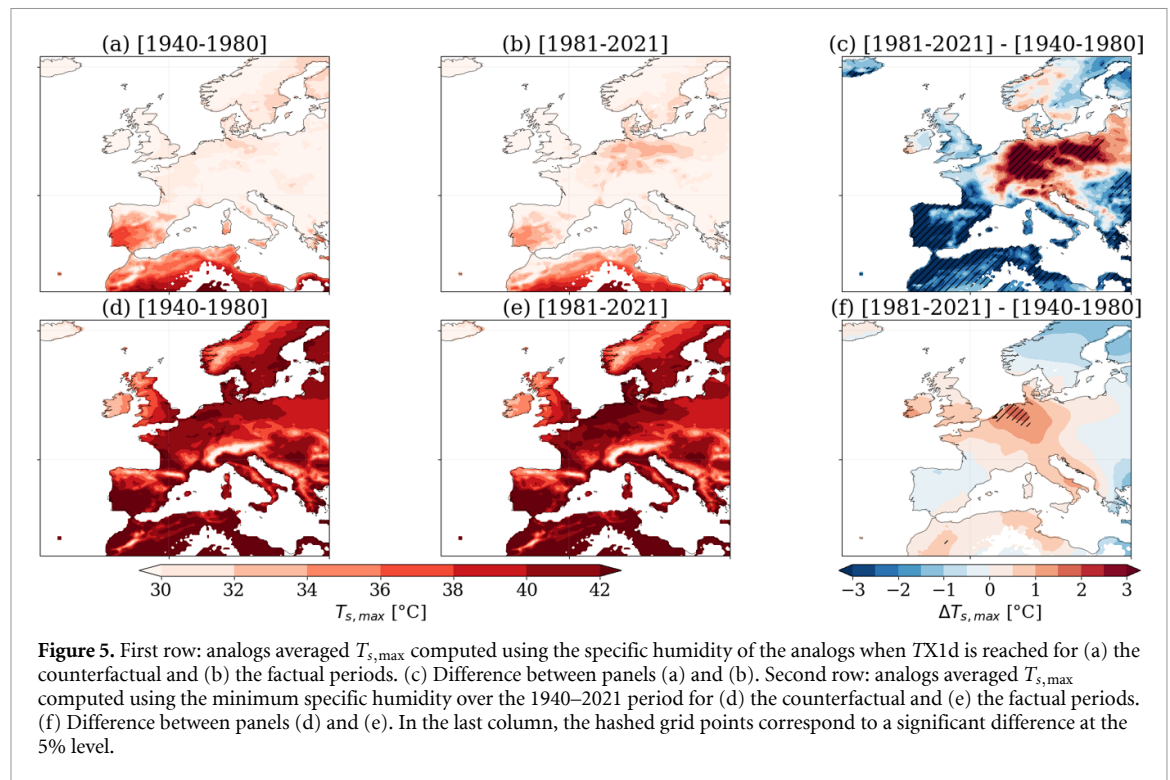
The flow analogs method for these two variables is however less robust because soil moisture may show a much slower dynamics than z_{500} . The distribution of soil moisture (and therefore specific humidity) is likely to depend on the full trajectory followed by the atmospheric dynamics rather than its last state. The association between the synoptic circulation $\mathcal{C}(\zeta)$ of the 25 July 2019 and dry soils/low surface specific humidity may therefore be a coincidence, especially for such a short extreme. As a consequence, it is possible that the framework presented in equation (5) may not be relevant for these variables.

Figure C4 in supplementary materials indeed shows that the JJAS daily minimum of hourly mean specific humidity at 2 m displays a significant

increase between the two periods for Western Europe (figure C4 panel (c)), probably as a result of moisture advection from nearby oceans [66]. This is apparently in contradiction with the analogs analysis presented here only if one dismisses the conditioning on the z_{500} pattern. Nonetheless, whether the soil desiccation pattern observed between the two periods is a relevant feature of the change in the physical mechanisms leading to extreme temperatures in the two periods can be questioned.

Previous literature showed that there is a decrease of surface specific humidity during the hottest days [56, 67–69], but this may arise as both a cause and a consequence of intensifying surface temperature extremes [70]. We note also that although soil moisture is well represented in ERA5, the temporal drying trend may be underestimated [71].

Figure 5 shows the $T_{s,\text{max}}$ indicator computed in both periods using either the actual specific humidity at the surface q or the 1940–2021 minimum specific humidity q_{min} . As already shown above, figure 5 panel (f) shows a significant change in absolute $T_{s,\text{max}}$ only over a limited region, which cannot explain the intensity change in TX1d between the two periods. Moreover, if one takes into account the moistening of minimum surface specific humidity over Western Europe in the factual period, this small difference may disappear completely as it would decrease $T_{s,\text{max}}$ in the factual period. When one uses the actual q of the analogs, the difference is much stronger between the two periods: there is an increase of more than 3 K of $T_{s,\text{max}}$ over a large region encompassing most of Germany. This region however does not coincide perfectly with the region with a significant change in TX1d in figure 4 panel (c).



To investigate this question, we analyze incoming energy fluxes. Even though the heatwave potential intensity has not changed much between the two periods, the actual capacity of a synoptic circulation to reach this maximal intensity may be greater in the factual period. Figure 6 displays the analogues average of the sum of daily mean incoming energy fluxes at the surface (downward positive). It includes latent and sensible heat fluxes, in addition to net shortwave radiation and downward longwave radiation. Figure 6 panel (c) shows that there is a significant and strong difference in incoming energy fluxes between the two periods which coincide with the heatwave region where the difference in $TX1d$ is significant. These fluxes are strongly coupled to the boundary layer and soils states—especially with regards to the partitioning between sensible and latent heat. Therefore one

cannot claim that energy fluxes are a fully external cause of the increase of $TX1d$ between the periods. However, they are an indicator of the strong increase in local diabatic heating.

The observed increase in $TX1d$ could also be the result of increased sensible heat advection from neighboring regions with increasingly desiccated land surface [24, 72]. It is possible that this phenomenon of non-local self-propagation of dry and hot conditions played an important role for the 25 July 2019 event as shown by [56], especially with regards to the aftermath of the June heatwave which desiccated soils in the south of France and the Iberian peninsula. However, it is unlikely to be the case for all of the 40 analogues in both periods: as shown by [73], in Europe, hot events are generally associated with weaker horizontal transport, but strong adiabatic

warming and local temperature increase caused by enhanced surface fluxes.

We emphasize that the results presented here are very sensitive to changes in surface specific humidity. This variable has significantly increased in the factual period, which according to equation (3), should compensate the increase in T_{500} due to mean tropospheric warming with climate change. If we take the example of the Paris grid point, the change between the two periods in $T_{s,max}$ due to the change in analogs averaged T_{500} is equal to $+0.58$ K. On the contrary, if we assume that minimum specific humidity is expected to shift by the same amount as the mean difference between the distribution of daily minimum of hourly mean specific humidity ($+0.52$ g kg⁻¹), then one could expect a change of $T_{s,max}$ by -1.29 K, which would largely compensate the tropospheric warming. This simple reasoning however assumes that the minimum of specific humidity during the hottest days would increase as the rest of the distribution. This is in opposition to the analogs analysis presented here which shows in contrast drying during the days associated with a synoptic circulation similar to the 2019 event. This analysis therefore supports the idea that the increase in the variability in temperatures, especially at the tail, is the result of increased variability in soil moisture and surface specific humidity. As a consequence, under the appropriate atmospheric conditions—namely a large anticyclonic anomaly—the response in daily maxima of surface temperature is stronger in the recent period.

4. Conclusion

We investigated the maximal reachable temperatures in Western Europe in the current climate. We took the July 2019 record-shattering heatwave event as a paradigmatic example and investigated how the maximum temperatures reached during this event could have been higher. We found that even in the current climate, we cannot rule out the possibility that 50°C cannot be reached in most urban areas of Western Europe. With a similar atmospheric pattern as the 2019 event, this would need a massive and large scale drought, with levels of surface specific humidity lower than the lowest observed in the last 80 years.

We then employed a flow analogues method to investigate how the upper bound of surface temperatures has changed between the past (1940–1980) and the present (1981–2021). We demonstrated that in the present period, the conditional expectation of daily maximum 2 m air temperature has increased significantly. This increase is stronger by around 1 K than the increase in the daily mean temperature. The change in the expected heatwave intensity potential $T_{s,max}$ explains part of this increase over the Netherlands and Western Germany where a small increase in T_{500} combined with dry soils to increase $T_{s,max}$ by more

than 3 K in the recent period. The same does not apply over France and Belgium, where the change in $T_{s,max}$ is small and insignificant with limited and non significant surface drying and no T_{500} difference. However, over these regions, incoming energy fluxes at the surface have increased strongly in the present period which may explain why higher $TX1d$ are reached compared to the past period even though the intensity potential has not changed. Our results suggest that most of the increase in observed daily maximum surface temperature may result from strengthened surface diabatic warming processes rather than free troposphere warming.

What would be the impact of a similar weather pattern as the July 2019 event in the future has been investigated by [74] under various warming scenarios and using a storyline approach. They found that an analogous event would entail peak temperatures around 50°C in Central Europe under a high emission scenario. Combined with our results, these elements show that the past extreme temperatures are not a good guide of the worst case scenario under global warming. The fast increase of maximum surface temperatures already observed demands a massive update of adaptation strategies.

Data availability statement

The data that support the findings of this study are available upon request from the authors. ERA5 reanalysis data are available on the Copernicus website (<https://cds.climate.copernicus.eu/cdsapp#!/home>).

Acknowledgments

The authors would like to thank Robert Vautard and Fabio D'Andrea for fruitful discussions, and the two anonymous reviewers for their valuable comments.

Funding

The authors acknowledge the support of the Grant ANR-20-CE01-0008-01 (SAMPRACE: PY). This work has received support from the European Union's Horizon 2020 research and innovation Programme under Grant Agreement No. 101003469 (XAIDA) and by the Marie Skłodowska-Curie Grant Agreement No. 956396 (EDIPI). R N was supported by a doctoral Grant from CEA.

Conflict of interest

The authors declare no conflict of interest.

Code availability

The main results of this work were obtained using Python. The scripts are available upon request.

Authors' contributions

R N did the data collection and analysis. All authors contributed to the design of the analysis and the writing of the paper.

Appendix A. Non-stationary GEV estimation

If we consider the maximum $M_n = \max\{X_1, \dots, X_n\}$ of n i.i.d. random variables X_1, \dots, X_n , the extremal theorems [33, 34] state that if there exist sequences $a_n > 0$ and b_n so that the distribution $\mathbb{P}(\frac{M_n - b_n}{a_n} \leq z)$ converges towards $G(z)$ when n goes to infinity, then $G(z)$ is given by:

$$G(z) = \begin{cases} \exp\left\{-\left[1 + \xi\left(\frac{z - \mu}{\sigma}\right)\right]^{-1/\xi}\right\} & \text{if } \xi \neq 0 \\ \exp\left\{-\exp\left[-\left(\frac{z - \mu}{\sigma}\right)\right]\right\} & \text{if } \xi = 0 \end{cases} \quad (\text{A.1})$$

where μ is called the location parameter, σ the scale parameter and ξ the shape parameter. This distribution is usually called the GEV distribution. Importantly, it is bounded upwards for $\xi < 0$ and the upper bound is given by:

$$z_{\max} = \mu - \frac{\sigma}{\xi}. \quad (\text{A.2})$$

Here we estimate a non-stationary GEV by letting the parameter μ depend on another covariate T encompassing the warming of the Earth: $\mu = \mu_1 + \mu_2 T$. We chose T to be the annual average of 2 m air temperature of North-Hemisphere over land as described in [75]. The GEV distribution is then estimated using the maximum likelihood method and we compute the upper bound using the formula (A.2).

We then report the 'upper' upper bound for T in 2021:

$$z_{\max}(2021) = \mu_0 + \mu_1 T_{2021} - \frac{\sigma}{\xi}. \quad (\text{A.3})$$

We compute the uncertainty range over this upper bound using $n = 1000$ bootstrap samples of

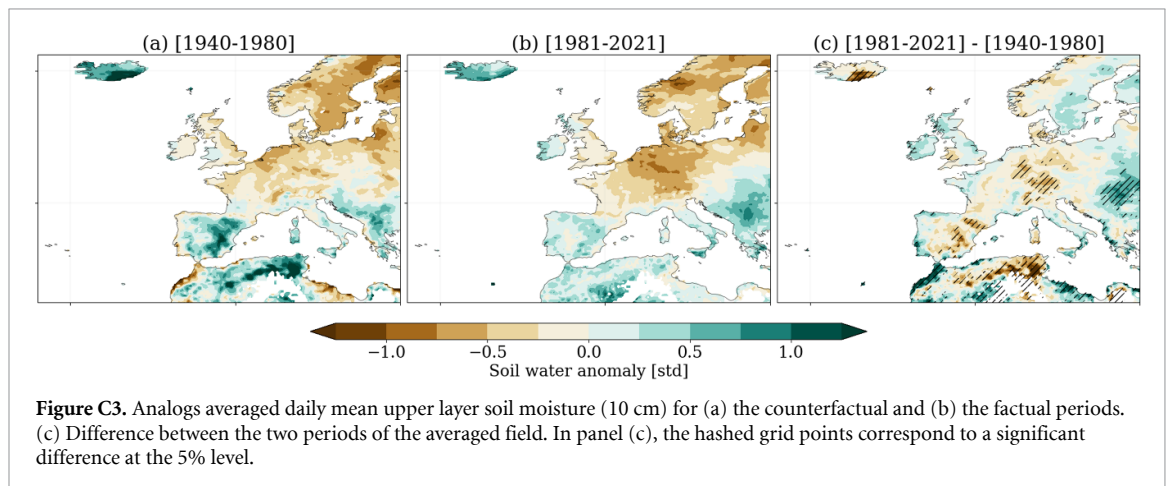
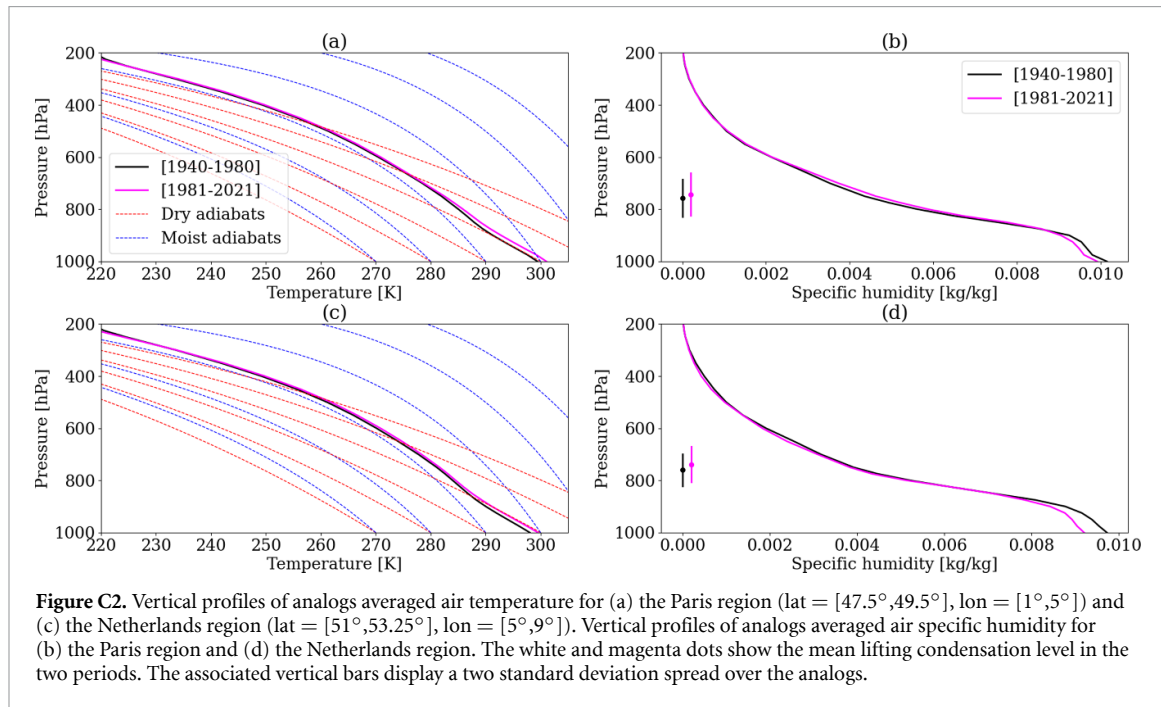
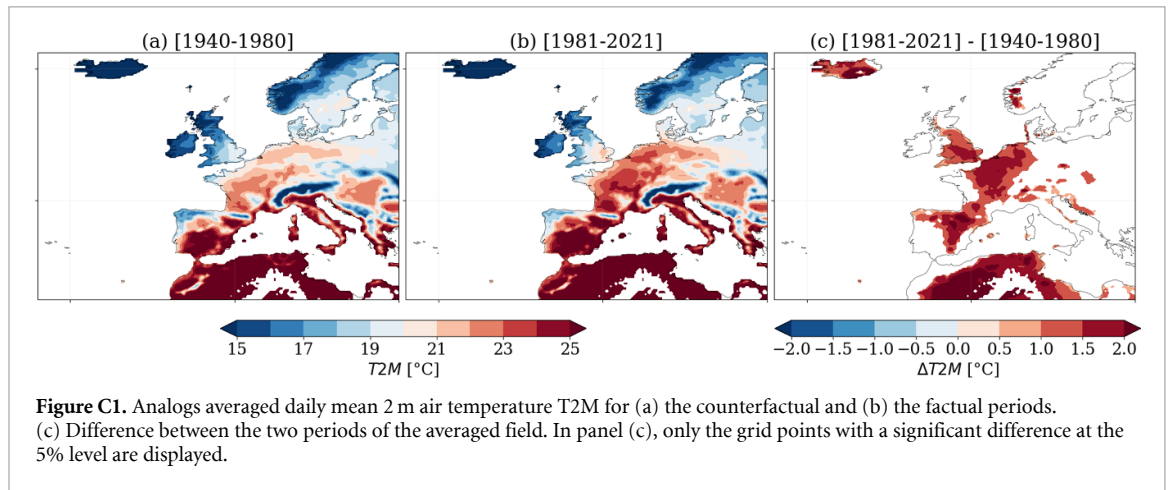
the initial series. For each bootstrap series, we fit a non-stationary GEV law (with T as a covariate on μ) and compute the upper bound. We take the median values of these bootstrapped estimates as the best estimate and we define the lower (resp. higher) bound of the estimate as the 97.5% and 2.5% quantiles of the bootstrapped z_{\max} .

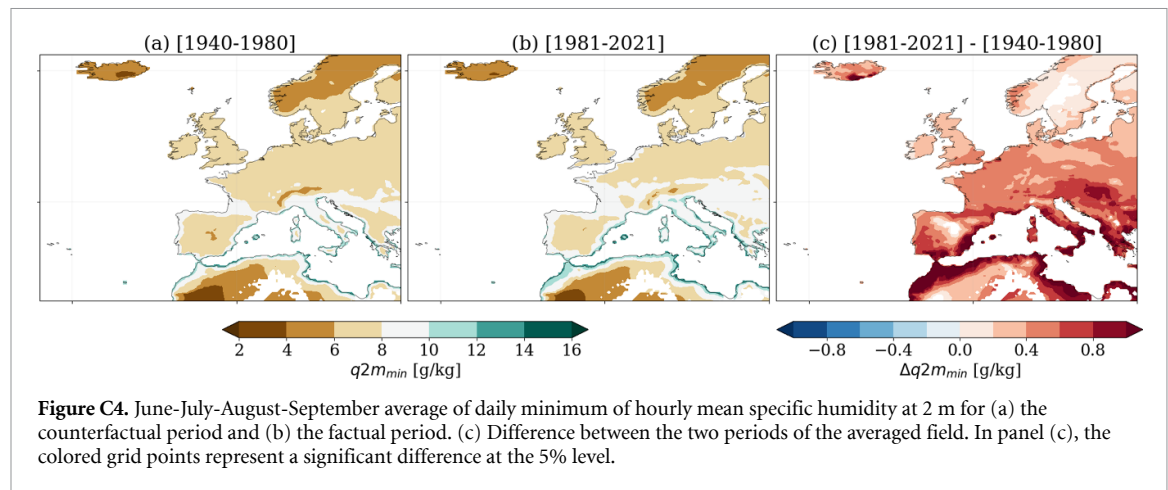
Appendix B. Uncertainty in $T_{s,\max}$ estimation

For computing the uncertainty range of the absolute $T_{s,\max}$ for the urban areas, we use for T_{500} and z_{500} the spread of the 10 ensemble members of the ERA5 reanalysis simulation. The lowest (resp. highest) bound of the range uses the lowest (resp. highest) T_{500} and z_{500} over the members. The ensemble spread is small therefore the uncertainty range for the absolute $T_{s,\max}$ is around 1 K.

For the computation of the more realistic $T_{s,\max}$, we also need to take into account the uncertainty in the q_{\min} value. To do so, we fit a stationary GEV law over the JJAS annual minimum of hourly 2 m air specific humidity over the 82 years between 1940–2021. This allows to give an approximate value to $P(q) := \mathbb{P}(q_{\min} \geq q)$ where q is a fixed threshold. We compute the uncertainty range over this estimated probability using $n = 1000$ bootstrap samples of the initial series. For each bootstrap series, we fit a GEV law and compute the return level q_{82} associated with a probability $P(q_{82}) = \frac{1}{82}$, similar to the probability of the minimum observed in the ERA5 dataset. We then take the 97.5% and 2.5% quantiles of the bootstrapped q_{82} as the uncertainty range for q_{\min} . Due to the resampling procedure, the 2.5% quantile is negative for all cities. As this is not physical, we take this lower bound to be null. We then combine the lowest (resp. highest) T_{500} and z_{500} over the members to the highest (resp. lowest) q_{\min} to get the uncertainty range. Due to the large uncertainty in q_{\min} , the uncertainty range for the realistic $T_{s,\max}$ is also large.

Appendix C. Additional figures





References

- Pörtner H O et al 2022 IPCC sixth assessment report
- Seneviratne S I, Zhang X, Adnan M, Badi W, Dereczynski C, Di Luca A, Vicente-Serrano S M, Wehner M and Zhou B 2021 Weather and climate extreme events in a changing climate *Climate Change 2021: The Physical Science Basis. Contribution of Working Group I to the Sixth Assessment Report of the Intergovernmental Panel on Climate Change* (Cambridge University Press)
- Campbell S, Remenyi T A, White C J and Johnston F H 2018 *Health Place* **53** 210–8
- Breshears D D, Fontaine J B, Ruthrof K X, Field J P, Feng X, Burger J R, Law D J, Kala J and Hardy G E S J 2021 *New Phytol.* **231** 32–39
- Huang W T K, Braithwaite I, Charlton-Perez A, Sarran C and Sun T 2022 *Environ. Res. Lett.* **17** 034017
- Masselot P et al 2023 *Lancet Planet. Health* **7** e271–81
- Hatfield J L and Prueger J H 2015 *Weather Clim. Extremes* **10** 4–10
- Brás T A, Seixas J, Carvalhais N and Jägermeyr J 2021 *Environ. Res. Lett.* **16** 065012
- Bastos A et al 2021 *Earth Syst. Dyn.* **12** 1015–35
- Zuo J, Pullen S, Palmer J, Bennetts H, Chileshe N and Ma T 2015 *J. Cleaner Prod.* **92** 1–12
- Barriopedro D, García-Herrera R, Ordóñez C, Miralles D and Salcedo-Sanz S 2023 *Rev. Geophys.* **61** e2022RG000780
- Perkins S E 2015 *Atmos. Res.* **164** 242–67
- Horton R M, Mankin J S, Lesk C, Coffel E and Raymond C 2016 *Curr. Clim. Change Rep.* **2** 242–59
- Domeisen D I, Eltahir E A, Fischer E M, Knutti R, Perkins-Kirkpatrick S E, Schär C, Seneviratne S I, Weisheimer A and Wernli H 2022 *Nat. Rev. Earth Environ.* **4** 1–15
- Petoukhov V, Rahmstorf S, Petri S and Schellnhuber H J 2013 *Proc. Natl Acad. Sci.* **110** 5336–41
- Coumou D, Petoukhov V, Rahmstorf S, Petri S and Schellnhuber H J 2014 *Proc. Natl Acad. Sci.* **111** 12331–6
- Kornhuber K, Coumou D, Vogel E, Lesk C, Donges J F, Lehmann J and Horton R M 2020 *Nat. Clim. Change* **10** 48–53
- Di Capua G, Sparrow S, Kornhuber K, Rousi E, Osprey S, Wallom D, van den Hurk B and Coumou D 2021 *npj Clim. Atmos. Sci.* **4** 1–14
- Seneviratne S I, Lüthi D, Litschi M and Schär C 2006 *Nature* **443** 205–9
- Fischer E M, Seneviratne S I, Lüthi D and Schär C 2007 *Geophys. Res. Lett.* **34** L05707
- Seneviratne S I, Corti T, Davin E L, Hirschi M, Jaeger E B, Lehner I, Orlowsky B and Teuling A J 2010 *Earth-Sci. Rev.* **99** 125–61
- Hirschi M, Seneviratne S I, Alexandrov V, Boberg E, Boroneant C, Christensen O B, Formayer H, Orlowsky B and Stepanek P 2011 *Nat. Geosci.* **4** 17–21
- Miralles D G, Van Den Berg M, Teuling A and De Jeu R 2012 *Geophys. Res. Lett.* **39** L21707
- Miralles D G, Teuling A J, Van Heerwaarden C C and Vilà-Guerau de Arellano J 2014 *Nat. Geosci.* **7** 345–9
- Rasmijn L, Van der Schrier G, Bintanja R, Barkmeijer J, Sterl A and Hazeleger W 2018 *Nat. Clim. Change* **8** 381–5
- Dirmeyer P A, Balsamo G, Blyth E M, Morrison R and Cooper H M 2021 *AGU Adv.* **2** e2020AV000283
- Petoukhov V, Petri S, Rahmstorf S, Coumou D, Kornhuber K and Schellnhuber H J 2016 *Proc. Natl Acad. Sci.* **113** 6862–7
- Kornhuber K, Petoukhov V, Karoly D, Petri S, Rahmstorf S and Coumou D 2017 *J. Clim.* **30** 6133–50
- Mann M E, Rahmstorf S, Kornhuber K, Steinman B A, Miller S K and Coumou D 2017 *Sci. Rep.* **7** 1–12
- Mann M E, Rahmstorf S, Kornhuber K, Steinman B A, Miller S K, Petri S and Coumou D 2018 *Sci. Adv.* **4** eaat3272
- Kornhuber K and Tamarin-Brodsky T 2021 *Geophys. Res. Lett.* **48** e2020GL091603
- Ghil M et al 2011 *Nonlinear Process. Geophys.* **18** 295–350
- Fisher R A and Tippett L H C 1928 Limiting forms of the frequency distribution of the largest or smallest member of a sample *Mathematical Proceedings of the Cambridge Philosophical Society* vol 24 (Cambridge University Press) pp 180–90
- Gnedenko B 1943 *Ann. Math.* **44** 423–53
- Coles S, Bawa J, Trenner L and Dorazio P 2001 *An Introduction to Statistical Modeling of Extreme Values* vol 208 (Springer)
- Cheng L, AghaKouchak A, Gilleland E and Katz R W 2014 *Clim. Change* **127** 353–69
- National Academies of Sciences Engineering and Medicine 2016 *Attribution of Extreme Weather Events in the Context of Climate Change* (The National Academies Press)
- Philip S et al 2020 A protocol for probabilistic extreme event attribution analyses *Adv. Stat. Climatol. Meteorol. Oceanogr.* **6** 177–203
- Van Oldenborgh G J et al 2021 *Clim. Change* **166** 13
- Van Oldenborgh G J, Wehner M F, Vautard R, Otto F E, Seneviratne S I, Stott P A, Hegerl G C, Philip S Y and Kew S F 2022 *Earth's Future* **10** e2021EF002271
- Auld G, Hegerl G and Papastathopoulos I 2021 *Adv. Stat. Clim. Meteorol. Oceanogr.* **9** 45–66
- Mitchell D, Kornhuber K, Huntingford C and Uhe P 2019 *Lancet Planet. Health* **3** e290–2
- Philip S Y et al 2022 *Earth Syst. Dyn.* **13** 1689–713
- Thompson V, Kennedy-Asser A T, Vosper E, Lo Y E, Huntingford C, Andrews O, Collins M, Hegerl G C and Mitchell D 2022 *Sci. Adv.* **8** eabm6860
- Vautard R et al 2020 *Environ. Res. Lett.* **15** 094077

- [46] Patterson M 2023 *Geophys. Res. Lett.* **50** e2023GL102757
- [47] Wouters J and Bouchet F 2016 *J. Phys. A: Math. Theor.* **49** 374002
- [48] Ragone F, Wouters J and Bouchet F 2018 *Proc. Natl Acad. Sci.* **115** 24–29
- [49] Yiou P and Jézéquel A 2020 *Geosci. Model Dev.* **13** 763–81
- [50] Ragone F and Bouchet F 2021 *Geophys. Res. Lett.* **48** e2020GL091197
- [51] Gessner C, Fischer E M, Beyerle U and Knutti R 2021 *J. Clim.* **34** 6619–34
- [52] Zhang Y and Boos W R 2023 *Proc. Natl Acad. Sci.* **120** e2215278120
- [53] Yiou P, Jézéquel A, Naveau P, Otto F E, Vautard R and Vrac M 2017 *Adv. Stat. Climatol. Meteorol. Oceanogr.* **3** 17–31
- [54] Jézéquel A, Yiou P and Radanovics S 2018 *Clim. Dyn.* **50** 1145–59
- [55] Faranda D, Bourdin S, Ginesta M, Krouma M, Messori G, Noyelle R, Pons F and Yiou P 2022 *Weather Clim. Dyn.* **3** 1311–40
- [56] Sousa P M, Barriopedro D, García-Herrera R, Ordóñez C, Soares P M and Trigo R M 2020 *Commun. Earth Environ.* **1** 48
- [57] Ma F, Yuan X, Jiao Y and Ji P 2020 *Geophys. Res. Lett.* **47** e2020GL087809
- [58] Hersbach H et al 2020 *Q. J. R. Meteorol. Soc.* **146** 1999–2049
- [59] Cadiou C, Noyelle R, Malhomme N and Faranda D 2023 *Asia Pac. J. Atmos. Sci.* **59** 83–94
- [60] Faranda D, Pascale S and Bulut B 2023 *Environ. Res. Lett.* **18** 034030
- [61] Xu P, Wang L, Huang P and Chen W 2021 *Atmos. Res.* **252** 105446
- [62] Maraun D 2016 *Curr. Clim. Change Rep.* **2** 211–20
- [63] Zhang L 2016 *Clim. Dyn.* **47** 3157–69
- [64] van Haren R, van Oldenborgh G J, Lenderink G and Hazeleger W 2013 *Environ. Res. Lett.* **8** 014053
- [65] Welch B L 1947 *Biometrika* **34** 28–35
- [66] Chadwick R, Good P and Willett K 2016 *J. Clim.* **29** 7613–32
- [67] Coffel E D, Horton R M, Winter J M and Mankin J S 2019 *Environ. Res. Lett.* **14** 084003
- [68] Padrón R S, Gudmundsson L, Decharme B, Ducharne A, Lawrence D M, Mao J, Peano D, Krinner G, Kim H and Seneviratne S I 2020 *Nat. Geosci.* **13** 477–81
- [69] McKinnon K A, Poppick A and Simpson I R 2021 *Nat. Clim. Change* **11** 598–604
- [70] Merrifield A L, Simpson I R, McKinnon K A, Sippel S, Xie S P and Deser C 2019 *Geophys. Res. Lett.* **46** 14082–92
- [71] Li M, Wu P and Ma Z 2020 *Int. J. Climatol.* **40** 5744–66
- [72] Schumacher D L, Keune J, Van Heerwaarden C C, Vilà-Guerau de Arellano J, Teuling A J and Miralles D G 2019 *Nat. Geosci.* **12** 712–7
- [73] Bieli M, Pfahl S and Wernli H 2015 *Q. J. R. Meteorol. Soc.* **141** 98–108
- [74] Sánchez-Benítez A, Goessling H, Pithan F, Semmler T and Jung T 2022 *J. Clim.* **35** 2373–90
- [75] Osborn T J, Jones P D, Lister D H, Morice C P, Simpson I R, Winn J, Hogan E and Harris I C 2021 *J. Geophys. Res.: Atmos.* **126** e2019JD032352